Study of $\pi^- p \to \pi^- \eta(\eta) p$ at 190 GeV with the COMPASS experiment



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on behalf of the COMPASS collaboration

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- Lightest scalar nonet and beyond.





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- A first glimpse of $f_0(1500) \rightarrow \eta \eta$





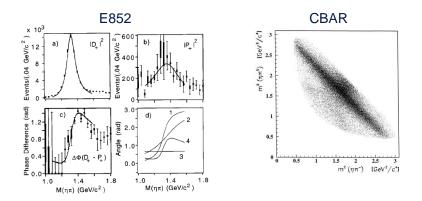
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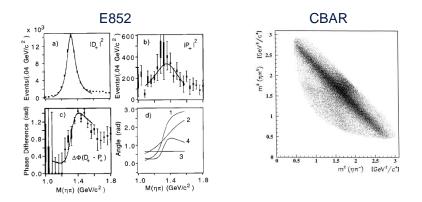
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- Conclusion and outlook.



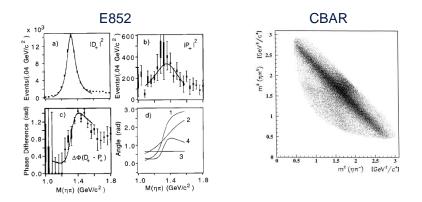


Seen by E852 exp. in $\pi^- p \rightarrow \eta \pi^- p$ at 18 GeV/c (publ. in 1997) and by CBAR exp. in $\bar{p}d \rightarrow \pi^- \pi^0 \eta p_{spectator}$ (publ. in 1998).





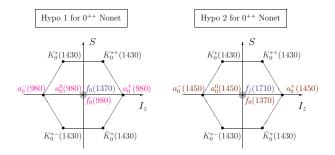
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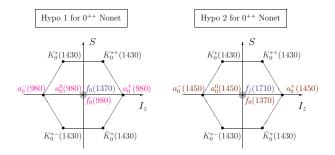
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Confirmed again by E852 in 2007 in $\pi^- p \rightarrow \eta \pi^0 n$ at 18 GeV/c, but with a lower mass ($M = 1257 \pm 20 \pm 25$ MeV).



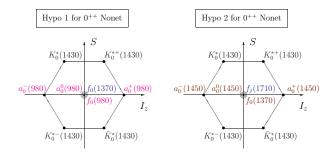






Hypo 3: $f_0(1500)$ supernumerary and therefore may be a glueball

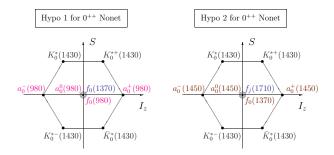




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Hypo 4: $a_0(980)$, $f_0(980)$ cusps or members of a tetraquark nonet.



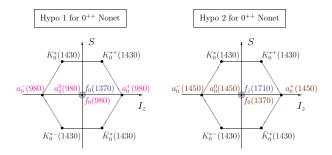


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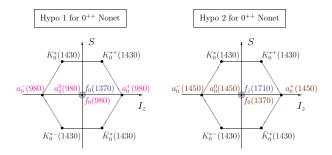
Hypo 5: $f_0(1370), f_0(1500), f_0(1710)$ are the result of the mixing of the glueball (and a tetraquark) with ordinary mesons.

Mixing scheme is based mainly on the results of WA102 experiment.

COMPASS goal in centrally produced data is to confirm and improve the observation of WA102:

measure the decay branching widths in $K\bar{K}, \pi\pi, \eta\eta, \eta\eta', 4\pi, \eta'\eta', ...$





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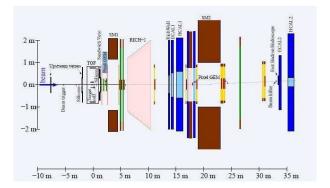
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 $\pi^- p \rightarrow \pi^- \eta \eta p$ very selective: $X \rightarrow \eta \eta$ has $I(J^{PC}) = 0(0^{++}, 2^{++}, 4^{++}, ...)$



COMPASS setup and detector description

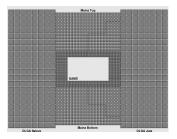


- Two arm spectrometer
- Tracking: Straw, Drift chambers, MicroMegas, PixelGEM, Recoil Proton Detector
- Calorimetry: ECAL1 (2006), ECAL2, HCAL1, HCAL2, Sandwich Veto
- Cherenkov: CEDAR, RICH



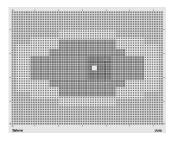
Electromagnetic Calorimeters

ECAL1



- 11.1 m downstream, low energetic photon detection, L × H: 3.97 × 2.86 m²
- 1500 channels:
- OLGA: 302 cells, 14.3 × 14.3 cm²
- MAINZ: 572 cells, 7.5 × 7.5 cm²
- GAMS: 608 cells, 3.8 × 3.8 cm²

ECAL2



- 33.2 downstream, high energetic photon detection, L × H: 2.45 × 1.94 m²
- 3068 channels:
- peripheral area: GAMS lead glass blocks 3.8 × 3.8 cm²
- central area: new ~ 900 radiation hard SHASHLYK modules 3.8 × 3.8 cm²
- New ADC (2008) with 32 sample converters



Pre-selection of exclusive events

- Trigger dedicated to diffractive and "central" reactions.
- Loop to all primary vertexes.
- Interaction in the target: $-69 < z_{vertex} < -29$ cm and $r_{vertex} < 1.5$ cm.
- 1 outgoing negative track with $E_{track} < 180$ GeV.
- 2 and 4 good clusters in ECAL1 and in ECAL2 for the 2γ and 4γ channels, respectively:



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- 2 and 4 good clusters in ECAL1 and in ECAL2 for the 2γ and 4γ channels, respectively:
 - not pointed by a track.
 - noise suppression.
 - E_{clusmin} > 1 GeV in ECAL1 and E_{clusmin} > 4 GeV in ECAL2.
 - in time with the beam: $-3 < t_{cluster} t_{beam} < 5ns$.

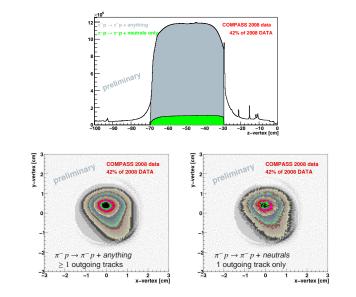


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 - ▶ in time with the beam: -3 < t_{cluster} t_{beam} < 5ns.</p>
- Correction of the photons momenta assuming they originate from the primary vertex.
- Correlation with RPD: $-0.3 < \phi_{\pi^- n\gamma} \phi_p < 0.3$ rad.
- Energy balance: $180 < E_{\pi^- n\gamma} < 200 \text{ GeV}$ assuming the track to be a pion.
- $\pi^0, \eta \to \gamma_1 \gamma_2$: 1 combination. $\pi^0_1, \eta_1 \to \gamma_i \gamma_j, \pi^0_2, \eta_2 \to \gamma_k \gamma_m$: 3 combinations.

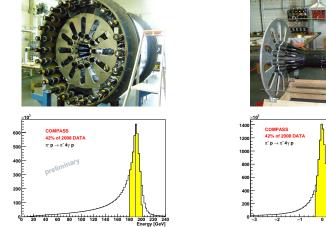


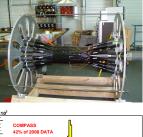
Vertex distributions

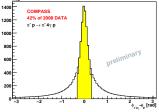




Recoil Proton Detector and exclusivity cuts

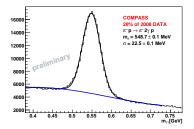


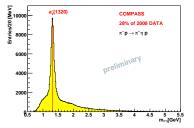






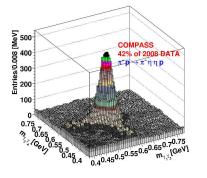
η and two-body $\eta\pi^-$ invariant masses in the 2γ channel

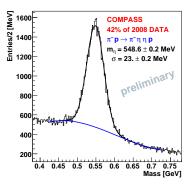






η masses in the 4γ channel







Preliminary statistics of $\pi^- p \rightarrow \pi^- \pi^0(\eta) p$

Amount of processed data (28% of 2008 data)	100.00%
DT0 trigger	73.17%
Majority < 6 for CEDAR1 and CEDAR2	71.75%
Primary vertex	66.15%
$-69 < z_{vertex} < -29$ cm	54.44%
$r_{vertex} < 1.5 \text{ cm}$	52.92%
1 negative track	4.89%
Two golden clusters	0.93%
$-0.3 < \phi_{\pi^- 2\gamma} - \phi_p < 0.3$	0.25%
Exclusivity (180 < $E_{\pi^- 2\gamma}$ < 200 GeV)	0.06%
$100 < m_{\pi 0} < 170 \ MeV$	31.2% of excl. events
$\pi^0 \text{ 1C } CL > 10\% \ (\pi^0 \text{ mass})$	14.6% of excl. events
$450 < m_{\eta} < 650 \ MeV$	21.6% of excl. events
η 1C $CL > 10\%$ (η mass)	8.1% of excl. events

- A preliminary sample of about 150K fitted $\pi^- p \rightarrow \pi^- \eta p$ events is used for the amplitude analysis.
- Better statistics will be achieved with improved calorimeter calibration (2008 data) and additional LASER and LED calorimeter monitoring system (2009 DATA).



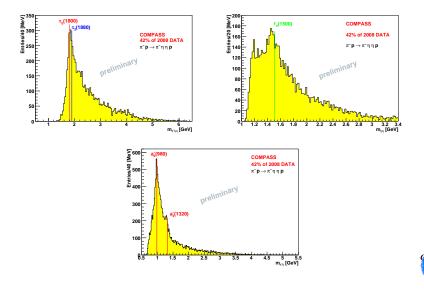
Preliminary statistics of $\pi^- p \rightarrow \pi^- \eta \eta p$

Amount of processed data (42% of 2008 data)	100.00%
DT0 trigger	73.78%
Majority < 6 for CEDAR1 and CEDAR2	72.49%
Primary vertex	66.91%
$-69 < z_{vertex} < -29$ cm	54.81%
$r_{vertex} < 1.5 \text{ cm}$	53.36%
1 negative track	4.94%
Four good clusters	0.61%
$-0.3 < \phi_{\pi^- 4\gamma} - \phi_p < 0.3$	0.21%
Exclusivity (180 < $E_{\pi^- 4\gamma}$ < 200 GeV)	0.10%
$\sqrt{(m_{\gamma_1\gamma_2} - m_{\pi^0})^2 + (m_{\gamma_3\gamma_4} - m_{\pi^0})^2} < 25 \text{ MeV}$	69.78% of excl. events
$2\pi^0 \text{ 2C } CL > 10\% \ (\pi^0 \text{ mass})$	27.39% of excl. events
$\sqrt{(m_{\gamma_1\gamma_2} - m_{\eta})^2 + (m_{\gamma_3\gamma_4} - m_{\eta})^2} < 25 \text{ MeV}$	0.17% of excl. events
$2\eta \text{ 2C } CL > 10\% \ (\eta \text{ mass})$	0.13% of excl. events

- A preliminary sample of about 5K fitted $\pi^- p \rightarrow \eta \eta p$ events is used for the amplitude analysis.
- Comparable amound of data is available in $\pi^- p$ and in pp at 190 GeV in the 2009 run.
- Better statistics will be achieved with improved calorimeter calibration (2008 data) and additional LASER and LED calorimeter monitoring system (2009 DATA).
- The statistics will be further increased by using the mixed decay mode of one of both ηs in π⁺π⁻π⁰.



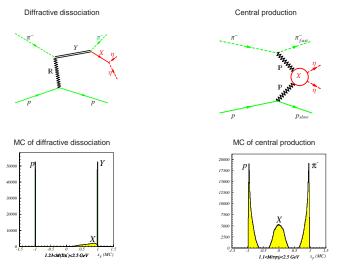
Two and three-body inv. masses in the 4γ channel



Production mechanisms

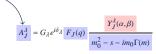
At 190 GeV incoming beam energy two compelling mechanisms for the production process of a state X are possible:

- as a product of the decay of a diffractively produced state Y: $\pi^- p \to Y p$, $Y \to \pi^- X$, $X \to \eta \eta$
- centrally produced via Double Pomeron Exchange: $\pi^- p \to \pi^-_{fast} X p, X \to \eta \eta$



xf and rapidities overlap: both processes have to be fitted simultaneously!

Amplitude (isobar model): -



- Blatt-Weisskopf barrier factors
- Angular part: spherical harmonics, decay angles α,β after "Wick rotations" (no D-functions needed).
- Relativistic Breit Wigner

$$\Gamma(m) = \Gamma_0 \left(\frac{m_0}{m} \frac{q}{q_0} \frac{F_J^2(q)}{F_J^2(q_0)} \right)$$

$$w(m,m_0,m_1) = \sum_{\lambda} [|A_{X_J}^{\lambda}(m,m_0)|^2 + |A_{Y_{J'}}^{\lambda}(m,m_1)|^2 + 2c_{\lambda} \Re(A_{X_J}^{\lambda}(m,m_0)A_{Y_{J'}}^{\lambda *}(m,m_1)]$$



Amplitude (isobar model):

$$A_{J}^{2} = G_{\lambda}e^{i\delta_{\lambda}}F_{J}(q) \frac{Y_{J}^{\lambda}(\alpha,\beta)}{m_{0}^{2} - s - im_{0}\Gamma(m)}$$

Weisskoof barrier factors

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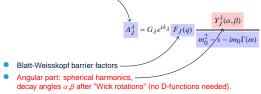
Blatt

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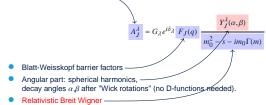
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Amplitude (isobar model): -



Resonance mass dependent width

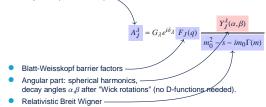
$$\Gamma(m) = \Gamma_0 \left(\frac{m_0}{m} \frac{q}{q_0} \frac{F_J^2(q)}{F_J^2(q_0)} \right)$$

- Mass of the two-body system
- Breek-up momentum

$$w(m,m_0,m_1) = \sum_{A} \left[|A_{XJ}^{A}(m,m_0)|^2 + |A_{YJ}^{A}(m,m_1)|^2 + 2 c_A \Re(A_{XJ}^{A}(m,m_0)A_{YJ'}^{A*}(m,m_1)) + |A_{YJ'}^{A*}(m,m_1)|^2 + 2 c_A \Re(A_{XJ}^{A*}(m,m_0)A_{YJ'}^{A*}(m,m_1)) + |A_{XJ}^{A*}(m,m_1)|^2 + 2 c_A \Re(A_{XJ}^{A*}(m,m_1)A_{YJ'}^{A*}(m,m_1)) + |A_{XJ}^{A*}(m,m_1)|^2 + 2 c_A \Re(A_{XJ}^{A*}(m,m_1)A_{YJ'}^{A*}(m,m_1)) + |A_{XJ}^{A*}(m,m_1)|^2 + 2 c_A \Re(A_{XJ}^{A*}(m,m_1)A_{YJ'}^$$



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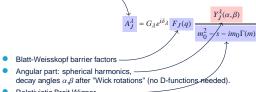


- $\Gamma(m) = \Gamma_0 \left(\frac{m_0}{m} \frac{q}{q_0} \frac{F_J^2(q)}{F_J^2(q_0)} \right)$
- Mass of the two-body system
- Break-up momentum

$$w(m,m_0,m_1) = \sum_{j} \left[|A_{X_J}^{\dagger}(m,m_0)|^2 + |A_{Y_{J'}}^{\dagger}(m,m_1)|^2 + 2 c_J \Re(A_{X_J}^{\dagger}(m,m_0)A_{Y_{J'}}^{\dagger *}(m,m_1)) + 2 c_J \Re(A_{X_J}^{\dagger *}(m,m_1)A_{Y_{J'}}^{\dagger *}(m,m_1)) \right]$$



Amplitude (isobar model): -



- Relativistic Breit Wigner –
- Resonance mass dependent width <</p>

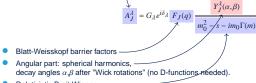
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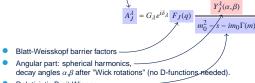


Break-up momentum

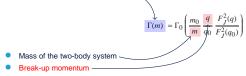
$$\mathbf{w}(m,m_0,m_1) = \sum_{j} \left[|A_{Xj}^{d}(m,m_0)|^2 + |A_{Yj'}^{d}(m,m_1)|^2 + 2c_{A} \Re(A_{Xj}^{d}(m,m_0)A_{Yj'}^{d*}(m,m_1)) + |A_{Yj'}^{d*}(m,m_1)|^2 + 2c_{A} \Re(A_{Xj}^{d}(m,m_0)A_{Yj'}^{d*}(m,m_1)) \right]$$



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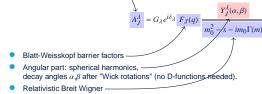


Intensity with two resonances with masses m₀ and m₁, spin J and J²:

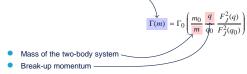
$$w(m,m_0,m_1) = \sum_{\lambda} [|A_{X_f}^{\lambda}(m,m_0)|^2 + |A_{Y_{f'}}^{\lambda}(m,m_1)|^2 + 2 c_{\lambda} \Re(A_{X_f}^{\lambda}(m,m_0)A_{Y_{f'}}^{\lambda^*}(m,m_1))^2 + 2 c_{\lambda} \Re(A_{X_f}^{\lambda}(m,m_0)A_{Y_{f'}}^{\lambda^*}(m,m_1))^2 + 2 c_{\lambda} \Re(A_{X_f}^{\lambda}(m,m_0)A_{Y_{f'}}^{\lambda^*}(m,m_1))^2 + 2 c_{\lambda} \Re(A_{X_f}^{\lambda}(m,m_0)A_{Y_{f'}}^{\lambda^*}(m,m_1))^2 + 2 c_{\lambda} \Re(A_{X_f}^{\lambda^*}(m,m_0)A_{Y_{f'}}^{\lambda^*}(m,m_1))^2 + 2 c_{\lambda} \Re(A_{X_f}^{\lambda^*}(m,m_0)A_{Y_{f'}}^{\lambda^*}(m,m_0)A_{Y_{f'$$



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- Resonance mass dependent width



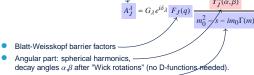
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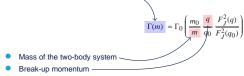


I spin component along $z_i - 1 \le c_i \le 1$ degree of coherence

Amplitude (isobar model):



- Relativistic Breit Wigner -
- Resonance mass dependent width

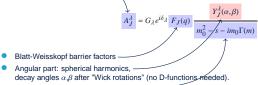


Intensity — with two resonances with masses m_0 and m_1 , spin J and J': $w(m,m_0,m_1) = \sum_{i} [|A_{X_J}^{\lambda}(m,m_0)|^2 + |A_{Y_{J'}}^{\lambda}(m,m_1)|^2 + 2 \frac{c_{\lambda}}{c_{\lambda}} \Re(A_{X_J}^{\lambda}(m,m_0)A_{Y_{J'}}^{\lambda*}(m,m_1)]$

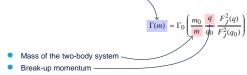


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- λ spin component along z, −1 ≤ c_λ ≤ 1 degree of coherence -



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With this definition, and for a fixed set of parameters, a reduction of $ln\mathcal{L}$ by 0.5 is statistically significant and corresponds to one standard deviation in mass and width optimizations.

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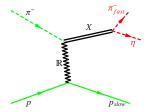
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- The rest of 2008 and all of the 2009 data will added to the final sample.



Backup Slides



Simulation of the production of a diffractive X in $\pi^- p \to X p$ with $X \to \pi^- \eta$



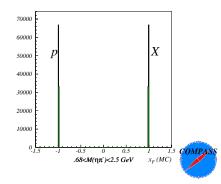


•
$$M_X$$
 uniformly from $m_{\pi} + m_{\eta}$ to 3.5 GeV

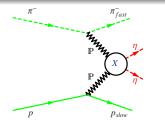
- t_X as e^{-bt} with b = 6 GeV⁻² with 0 < t < 1 GeV². To take into account a resonance dependent production mechanism the shape of the t-distribution will be optimized from the data in different mass ranges around the resonance masses.
- $\phi_X(\phi_p)$ uniformly from 0 to 2π

$$1 - x_X = \frac{M_X^2 - m_\pi^2}{s}$$

$$p_{T,X}^2 = -t_X$$



Simulation of the production of a central X in $\pi^- p \to X p$ with $X \to \eta \eta$



$$x_{\mathbb{P}_1} - \frac{M^2}{s} \frac{1}{x_{\mathbb{P}_1}} = \frac{M_T}{\sqrt{s}} (e^y - e^{-y})$$

Solution:

20000

$$x_{\mathbb{P}_1} = \frac{M_T}{\sqrt{s}} \left[\pm \sqrt{\left(\frac{M_X}{M_T}\right)^2 + (\sinh y)^2} + \sinh y \right].$$

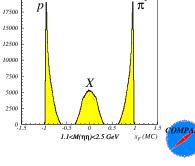
- M_X uniformly from $2m_p$ to 3.5 GeV
- t_X as e^{-bt} with an average $b = 6 \ GeV^{-2}$ with $0 < t < 1 \ GeV^2$. The optimization of a resonance dependent t-distribution will be obtained from the data.
- Flat rapidity distribution -1 < y(X) < 1
- $\phi_X(\phi_p)$ uniformly from 0 to 2π

$$M_X^2 = -x_{\mathbb{P}_1} x_{\mathbb{P}_2} s$$

 $x_{\mathbb{P}_2} = 1 - x_{\pi}$ on the π side, $x_{\mathbb{P}_1} = x_p - 1$ on the p side In the center of mass

$$x_p + x_\pi + x_X = 0$$
 $x_X = M_T \frac{e^y - e^{-y}}{\sqrt{s}} = \frac{2M_T \sinh y_{cm}}{\sqrt{s}}$





Wick rotation and amplitude Ansatz



Definition of angles for a diffractive X in $\pi^- p \rightarrow \pi^- X p$, $X \rightarrow \pi^- \eta p$:

- The z axis is defined in the πp c.m. frame. The x, y axes are defined by the angle formed by the production plane and the decay plane.
- The Wick rotation by angles −φ and θ to the direction of flight of the diffractive X are followed by a Lorentz boost to the its rest frame (x', y', z') and by another rotation by −θ and −φ so that the direction of the new reference frame x', y', z'' correspond to one of x, y, z.
- α, β define the direction of one η in the rest frame of X after the Wick rotations. The effect of the Lorentz boost is to leave the η with final momenta different from those in the overall πp rest frame.

The angles α, β obtained in this reference frame after Wick rotations enter in the decay amplitude definition:

•
$$A_J^{\lambda} = G_{\lambda} e^{i\delta_{\lambda}} F_J(q) \frac{Y_J^{\lambda}(\alpha, \beta)}{m_0^2 - s - im_0 \Gamma(m)}$$