## Study of $\pi^{-} p \rightarrow \pi^{-} \eta(\eta) p$ at 190 GeV with the COMPASS experiment



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## Agenda

- Exotic $\pi(1400)$ observations.
- Lightest scalar nonet and beyond.


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- PWA description and comparison with standard formalisms.
- Conclusion and outlook.

E852


CBAR


Seen by E852 exp. in $\pi^{-} p \rightarrow \eta \pi^{-} p$ at $18 \mathrm{GeV} / \mathrm{c}$ (publ. in 1997) and by CBAR exp. in $\bar{p} d \rightarrow \pi^{-} \pi^{0} \eta p_{\text {spectator }}$ (publ. in 1998).


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Confirmed again by E852 in 2007 in $\pi^{-} p \rightarrow \eta \pi^{0} n$ at $18 \mathrm{GeV} / \mathrm{c}$, but with a lower mass ( $M=1257 \pm 20 \pm 25 \mathrm{MeV}$ ).

## Hypothetical lightest scalar nonets configurations and beyond

$$
\text { Hypo } 1 \text { for } 0^{++} \text {Nonet }
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Hypo 2 for $0^{++}$Nonet


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Hypo 5: $f_{0}(1370), f_{0}(1500), f_{0}(1710)$ are the result of the mixing of the glueball (and a tetraquark) with ordinary mesons.

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COMPASS goal in centrally produced data is to confirm and improve the observation of WA102: measure the decay branching widths in $K \bar{K}, \pi \pi, \eta \eta, \eta \eta^{\prime}, 4 \pi, \eta^{\prime} \eta^{\prime}, \ldots$

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$\pi^{-} p \rightarrow \pi^{-} \eta \eta p$ very selective: $X \rightarrow \eta \eta$ has $I\left(J^{P C}\right)=0\left(0^{++}, 2^{++}, 4^{++}, \ldots\right)$

## COMPASS setup and detector description



- Two arm spectrometer
- Tracking: Straw, Drift chambers, MicroMegas, PixelGEM, Recoil Proton Detector
- Calorimetry: ECAL1 (2006), ECAL2, HCAL1, HCAL2, Sandwich Veto
- Cherenkov: CEDAR, RICH


## Electromagnetic Calorimeters

## ECAL1



- 11.1 m downstream, low energetic photon detection, $L \times H: 3.97 \times 2.86 \mathrm{~m}^{2}$
- 1500 channels:
- OLGA: 302 cells, $14.3 \times 14.3 \mathrm{~cm}^{2}$
- MAINZ: 572 cells, $7.5 \times 7.5 \mathrm{~cm}^{2}$
- GAMS: 608 cells, $3.8 \times 3.8 \mathrm{~cm}^{2}$


## ECAL2

- 33.2 downstream, high energetic photon detection, $L \times H: 2.45 \times 1.94 \mathrm{~m}^{2}$
- 3068 channels:
- peripheral area: GAMS lead glass blocks $3.8 \times 3.8 \mathrm{~cm}^{2}$
- central area: new $\sim 900$ radiation hard SHASHLYK modules $3.8 \times 3.8 \mathrm{~cm}^{2}$
- New ADC (2008) with 32 sample converters


## Pre-selection of exclusive events

- Trigger dedicated to diffractive and "central" reactions.
- Loop to all primary vertexes.
- Interaction in the target: $-69<z_{\text {vertex }}<-29 \mathrm{~cm}$ and $r_{\text {vertex }}<1.5 \mathrm{~cm}$.
- 1 outgoing negative track with $E_{\text {track }}<180 \mathrm{GeV}$.
- 2 and 4 good clusters in ECAL1 and in ECAL2 for the $2 \gamma$ and $4 \gamma$ channels, respectively:


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- not pointed by a track.
- noise suppression.
$>E_{\text {clus }_{\text {min }}}>1 \mathrm{GeV}$ in ECAL1 and $E_{\text {clus }_{\text {min }}}>4 \mathrm{GeV}$ in ECAL2.
- in time with the beam: $-3<t_{\text {cluster }}-t_{\text {beam }}<5 n s$.


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- in time with the beam: $-3<t_{\text {cluster }}-t_{\text {beam }}<5 n s$.
- Correction of the photons momenta assuming they originate from the primary vertex.
- Correlation with RPD: $-0.3<\phi_{\pi^{-n \gamma}}-\phi_{p}<0.3 \mathrm{rad}$.
- Energy balance: $180<E_{\pi^{-} n \gamma}<200 \mathrm{GeV}$ assuming the track to be a pion.
- $\pi^{0}, \eta \rightarrow \gamma_{1} \gamma_{2}$ : 1 combination.
$\pi_{1}^{0}, \eta_{1} \rightarrow \gamma_{i} \gamma_{j}, \pi_{2}^{0}, \eta_{2} \rightarrow \gamma_{k} \gamma_{m}: 3$ combinations.


## Vertex distributions





## Recoil Proton Detector and exclusivity cuts





## $\eta$ and two-body $\eta \pi^{-}$invariant masses in the $2 \gamma$ channel




## $\eta$ masses in the $4 \gamma$ channel




## Preliminary statistics of $\pi^{-} p \rightarrow \pi^{-} \pi^{0}(\eta) p$

| Amount of processed data (28\% of 2008 data) | $100.00 \%$ |
| :--- | ---: |
| DT0 trigger | $73.17 \%$ |
| Majority $<6$ for CEDAR1 and CEDAR2 | $71.75 \%$ |
| Primary vertex | $66.15 \%$ |
| $-69<z_{\text {vertex }}<-29 \mathrm{~cm}$ | $54.44 \%$ |
| $r_{\text {vertex }}<1.5 \mathrm{~cm}$ | $52.92 \%$ |
| 1 negative track | $4.89 \%$ |
| Two golden clusters | $0.93 \%$ |
| $-0.3<\phi_{\pi^{-}-2 \gamma}-\phi_{p}<0.3$ | $0.25 \%$ |
| Exclusivity $\left(180<E_{\pi}-2 \gamma<200 \mathrm{GeV}\right)$ | $0.06 \%$ |
| $100<m_{\pi^{0}}<170 \mathrm{MeV}$ | $31.2 \%$ of excl. events |
| $\pi^{0} 1 \mathrm{CL}>10 \%\left(\pi^{0}\right.$ mass) | $14.6 \%$ of excl. events |
| $450<m_{\eta}<650 \mathrm{MeV}$ | $21.6 \%$ of excl. events |
| $\eta 1 \mathrm{C} \mathrm{CL}>10 \%(\eta$ mass $)$ | $8.1 \%$ of excl. events |

- A preliminary sample of about 150 K fitted $\pi^{-} p \rightarrow \pi^{-} \eta p$ events is used for the amplitude analysis.
- Better statistics will be achieved with improved calorimeter calibration (2008 data) and additional LASER and LED calorimeter monitoring system (2009 DATA).


## Preliminary statistics of $\pi^{-} p \rightarrow \pi^{-} \eta \eta p$

| Amount of processed data (42\% of 2008 data) | $100.00 \%$ |
| :--- | ---: |
| DT0 trigger | $73.78 \%$ |
| Majority $<6$ for CEDAR1 and CEDAR2 | $72.49 \%$ |
| Primary vertex | $66.91 \%$ |
| $-69<z_{\text {vertex }}<-29 \mathrm{~cm}$ | $54.81 \%$ |
| $r_{\text {vertex }}<1.5 \mathrm{~cm}$ | $53.36 \%$ |
| 1 negative track | $4.94 \%$ |
| Four good clusters | $0.61 \%$ |
| $-0.3<\phi_{\pi^{-}} 4 \gamma-\phi_{p}<0.3$ | $0.21 \%$ |
| Exclusivity $\left(180<E_{\pi^{-} 4 \gamma}<200 \mathrm{GeV}\right)$ | $0.10 \%$ |
| $\sqrt{\left(m_{\gamma_{1} \gamma_{2}}-m_{\pi^{0}}\right)^{2}+\left(m_{\gamma_{3} \gamma_{4}}-m_{\pi^{0}}\right)^{2}}<25 \mathrm{MeV}$ | $69.78 \%$ of excl. events |
| $2 \pi^{0} 2 \mathrm{C} C L>10 \%\left(\pi^{0}\right.$ mass) | $27.39 \%$ of excl. events |
| $\sqrt{\left(m_{\gamma_{1} \gamma_{2}}-m_{\eta}\right)^{2}+\left(m_{\gamma_{3} \gamma_{4}}-m_{\eta}\right)^{2}}<25 \mathrm{MeV}$ | $0.17 \%$ of excl. events |
| $2 \eta 2 \mathrm{C} C L>10 \%(\eta$ mass $)$ | $0.13 \%$ of excl. events |

- A preliminary sample of about 5 K fitted $\pi^{-} p \rightarrow \eta \eta p$ events is used for the amplitude analysis.
- Comparable amound of data is available in $\pi^{-} p$ and in $p p$ at 190 GeV in the 2009 run.
- Better statistics will be achieved with improved calorimeter calibration (2008 data) and additional LASER and LED calorimeter monitoring system (2009 DATA).
- The statistics will be further increased by using the mixed decay mode of one of both $\eta \mathrm{s}$ in $\pi^{+} \pi^{-} \pi^{0}$.


## Two and three-body inv. masses in the $4 \gamma$ channel




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## Production mechanisms

At 190 GeV incoming beam energy two compelling mechanisms for the production process of a state X are possible:

- as a product of the decay of a diffractively produced state $\mathrm{Y}: \pi^{-} p \rightarrow Y p, Y \rightarrow \pi^{-} X, X \rightarrow \eta \eta$
- centrally produced via Double Pomeron Exchange: $\pi^{-} p \rightarrow \pi_{\text {fast }}^{-} X p, X \rightarrow \eta \eta$

Diffractive dissociation


MC of diffractive dissociation


Central production


MC of central production

$x_{f}$ and rapidities overlap: both processes have to be fitted simultaneously!
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## Amplitude Ansatz for the decay process

- Amplitude (isobar model):

$$
A_{J}^{\lambda}=G_{\lambda} e^{i \delta_{\lambda}} F_{J}(q) \frac{Y_{J}^{\lambda}(\alpha, \beta)}{m_{0}^{2}-s-i m_{0} \Gamma(m)}
$$

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- Blatt-Weisskopf barrier factors

- Angular part: spherical harmonics, decay angles $\alpha, \beta$ after "Wick rotations" (no D-functions needed).


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- Resonance mass dependent width

$$
\downarrow_{\Gamma(m)}^{\downarrow}=\Gamma_{0}\left(\frac{m_{0}}{m} \frac{q}{q_{0}} \frac{F_{J}^{2}(q)}{F_{J}^{2}\left(q_{0}\right)}\right)
$$

Mass of the two-body system
Break-up momentum

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- Break-up momentum
- Intensity with two resonances with masses $m_{0}$ and $m_{1}$, spin $J$ and $J^{\prime}$ : $w\left(m, m_{0}, m_{1}\right)=\sum_{\lambda}\left[\left|A_{X_{J}}^{\lambda}\left(m, m_{0}\right)\right|^{2}+\left|A_{Y_{J^{\prime}}}^{\lambda}\left(m, m_{1}\right)\right|^{2}+2 c_{\lambda} \Re\left(A_{X_{J}}^{\lambda}\left(m, m_{0}\right) A_{Y_{J^{\prime}}}^{\lambda *}\left(m, m_{1}\right)\right]\right.$


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- $\lambda$ spin component along $z,-1 \leq c_{\lambda} \leq 1$ degree of coherence


## Minimization and comparison with standard PWA

- Minimization of total intensity of the negative log-likelihood:

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-\ln \mathcal{L}=\left(-\sum_{j=1}^{N} \ln w_{j}\right)+N \ln \left(\sum_{i=1}^{M} w_{i}\right)
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- $N$ : number of data events
$M$ : number of MC events
Well established resonance parameters fixed at PDG values.
$G_{\lambda}, \delta_{\lambda}, c_{\lambda}$ : free parameters of the fit
With this definition, and for a fixed set of parameters, a reduction of $\ln \mathcal{L}$ by 0.5 is statistically significant and corresponds to one standard deviation in mass and width optimizations.
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Main differences between this formalism and the standard PWA formalism used in the BNL E852 and WA102 experiments.
- Resonance rest frame after Wick rotation vs. Gottfried-Jackson reference frame
- Partial coherence vs. reflectivity basis with natural and unnatural-parity exchange
- Fitting procedure: unbinned log-likelihood fit vs. mass independent (binned) log-likelihood angular fit + mass fit


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- Disadvantages:
- Computing limitations (presently fast convergence only for $<100 \mathrm{~K}$ events).
- Additional mass and width scans for all possible spin combination of all unknown resonances needed.


## Conclusion and outlook

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- The rest of 2008 and all of the 2009 data will added to the final sample.


## Backup Slides

## Simulation of the production of a diffractive X in $\pi^{-} p \rightarrow X p$ with $X \rightarrow \pi^{-} \eta$



- $M_{X}$ uniformly from $m_{\pi}+m_{\eta}$ to 3.5 GeV
- $t_{X}$ as $e^{-b t}$ with $b=6 \mathrm{GeV}^{-2}$ with $0<t<1 \mathrm{GeV}^{2}$. To take into account a resonance dependent production mechanism the shape of the t-distribution will be optimized from the data in different mass ranges around the resonance masses.
- $\phi_{X}\left(\phi_{p}\right)$ uniformly from 0 to $2 \pi$

$$
\begin{gathered}
1-x_{X}=\frac{M_{X}^{2}-m_{\pi^{-}}^{2}}{s} \\
p_{T, X}^{2}=-t_{X}
\end{gathered}
$$




## Simulation of the production of a central X in $\pi^{-} p \rightarrow X p$ with $X \rightarrow \eta \eta$



$$
x_{\mathbb{P}_{1}}-\frac{M^{2}}{s} \frac{1}{x_{\mathrm{P}_{1}}}=\frac{M_{T}}{\sqrt{s}}\left(e^{y}-e^{-y}\right)
$$

Solution:

$$
x_{\mathrm{P}_{1}}=\frac{M_{T}}{\sqrt{s}}\left[ \pm \sqrt{\left(\frac{M_{X}}{M_{T}}\right)^{2}+(\sinh y)^{2}}+\sinh y\right]
$$

- $M_{X}$ uniformly from $2 m_{\eta}$ to 3.5 GeV
- $t_{X}$ as $e^{-b t}$ with an average $b=6 \mathrm{GeV}^{-2}$ with $0<t<1 \mathrm{GeV}^{2}$. The optimization of a resonance dependent $t$-distribution will be obtained from the data.
- Flat rapidity distribution $-1<y(X)<1$
- $\phi_{X}\left(\phi_{p}\right)$ uniformly from 0 to $2 \pi$

$$
M_{X}^{2}=-x_{\mathbb{P}_{1}} x_{\mathbb{P}_{2}} s
$$

$x_{\mathrm{P}_{2}}=1-x_{\pi}$ on the $\pi$ side, $x_{\mathrm{P}_{1}}=x_{p}-1$ on the $p$ side In the center of mass


$$
x_{p}+x_{\pi}+x_{X}=0 \quad x_{X}=M_{T} \frac{e^{y}-e^{-y}}{\sqrt{s}}=\frac{2 M_{T} \sinh y_{c m}}{\sqrt{s}}
$$

## Wick rotation and amplitude Ansatz



Definition of angles for a diffractive $X$ in $\pi^{-} p \rightarrow \pi^{-} X p, X \rightarrow \pi^{-} \eta p$ :

- The $z$ axis is defined in the $\pi p$ c.m. frame. The $x, y$ axes are defined by the angle formed by the production plane and the decay plane.
- The Wick rotation by angles $-\phi$ and $\theta$ to the direction of flight of the diffractive $X$ are followed by a Lorentz boost to the its rest frame $\left(x^{\prime}, y^{\prime}, z^{\prime}\right)$ and by another rotation by $-\theta$ and $-\phi$ so that the direction of the new reference frame $x^{\prime \prime}, y^{\prime \prime}, z^{\prime \prime}$ correspond to one of $x, y, z$.
- $\alpha, \beta$ define the direction of one $\eta$ in the rest frame of X after the Wick rotations. The effect of the Lorentz boost is to leave the $\eta$ with final momenta different from those in the overall $\pi p$ rest frame.
The angles $\alpha, \beta$ obtained in this reference frame after Wick rotations enter in the decay amplitude definition:
- $A_{J}^{\lambda}=G_{\lambda} e^{i \delta_{\lambda}} F_{J}(q) \frac{Y_{J}^{\lambda}(\alpha, \beta)}{m_{0}^{2}-s-i m_{0} \Gamma(m)}$

