

TMD effects in unpolarised processes – experiment overview

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in hard processes and the transverse structure of the proton
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



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and physics

Supported by 'FORTE' project, CZ.02.01.01/00/22-008/0004632, from Czech MEYS, co-funded by the EU.



		Parent hadron polarization		
		Unpolarised	Longitudinal	Transverse
Parton polarisation	U	 $f_1(x, k_T^2)$ (unpolarised)		 $f_{1T}^\perp(x, k_T^2)$ (Sivers)
	L/C		 $g_1(x, k_T^2)$ (helicity)	 $g_{1T}(x, k_T^2)$ (Kotzinian–Mulders)
	T/L	 $h_1^\perp(x, k_T^2)$ (Boer–Mulders)	 $h_{1L}^\perp(x, k_T^2)$ (worm-gear)	 $h_1(x, k_T^2)$ (transversity)
				 $h_{1T}^\perp(x, k_T^2)$ (pretzelosity)

Parton polarisation:



L/C – longitudinal (quarks) or circular (gluons)

T/L – transverse (quarks) or linear (gluons)

8 leading twist TMD PDFs, 2 are relevant for unpolarised hadrons.

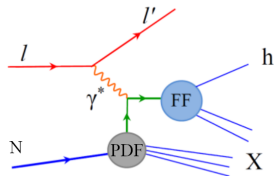
		Unpolarised hadron
Parton polarisation	U	 $f_1(x, k_T^2)$ (unpolarised)
	L/C	
	T/L	 $h_1^\perp(x, k_T^2)$ (Boer–Mulders)

Parton distribution functions

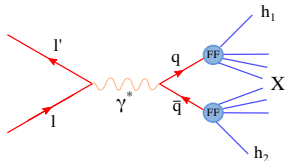
		Unpolarised hadron
Parton polarisation	U	 $D_1(z, P_\perp^2)$ (unpolarised)
	L/C	
	T/L	 $H_1^\perp(z, P_\perp^2)$ (Collins)

Fragmentation functions

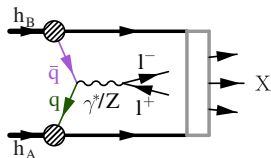
- Parton polarisation:
 - L/C: longitudinal (quarks) or circular (gluons).
 - T/L: transverse (quarks) or linear (gluons).
- Boer–Mulders function
 - Quarks: chiral-odd,
 - Gluons: chiral-even.



Hadron production in DIS (SIDIS)



Hadron production in e^+e^- annihilation.



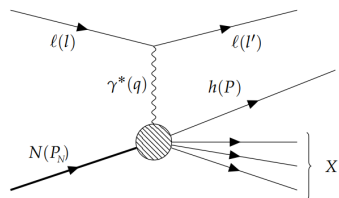
Drell-Yan process or Z boson production.

$$l(l) + N(P_N) \rightarrow l'(l') + h(P) + X$$

The cross section is [A. Bacchetta *et al.*, JHEP 02 (2007) 093]

$$\begin{aligned} \frac{d\sigma}{dx dy dz d\phi_h dP_T^2} &= \frac{2\pi\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{2xM^2}{Q^2}\right) \left(F_{UU,T} + \varepsilon F_{UU,L} \right. \\ &\quad \left. + \sqrt{2\varepsilon(1+\varepsilon)} F_{UU}^{\cos\phi_h} \cos\phi_h + \varepsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h + \lambda \sqrt{2\varepsilon(1-\varepsilon)} F_{LU}^{\sin\phi_h} \sin\phi_h \right) \\ &= \sigma_0 \left(1 + \varepsilon_1 A_{UU}^{\cos\phi_h} \cos\phi_h + \varepsilon_2 A_{UU}^{\cos 2\phi_h} \cos 2\phi_h + \lambda \varepsilon_3 A_{LU}^{\sin\phi_h} \sin\phi_h \right) \end{aligned}$$

- x, y, Q^2 : usual DIS variables,
- λ : beam polarisation (≈ 0.8 at COMPASS),
- ε : ratio of longitudinal and transverse photon flux,
- z : fraction of γ^* energy carried by h .
- P_T : transverse momentum of h in the γN frame, ϕ_h is its azimuthal angle.
- $F_{XU}^{f(\phi_h)}(x, z, P_T^2, Q^2)$ are structure functions.
- $A_{XU}^{f(\phi_h)}(x, z, P_T^2, Q^2)$ are commonly called azimuthal asymmetries.



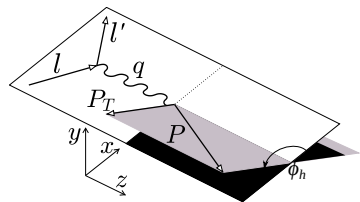
Semi-inclusive DIS.

$$l(l) + N(P_N) \rightarrow l'(l') + h(P) + X$$

The cross section is [A. Bacchetta *et al.*, JHEP 02 (2007) 093]

$$\begin{aligned} \frac{d\sigma}{dx dy dz d\phi_h dP_T^2} &= \frac{2\pi\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{2xM^2}{Q^2}\right) \left(F_{UU,T} + \varepsilon F_{UU,L} \right. \\ &\quad \left. + \sqrt{2\varepsilon(1+\varepsilon)} F_{UU}^{\cos\phi_h} \cos\phi_h + \varepsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h + \lambda\sqrt{2\varepsilon(1-\varepsilon)} F_{LU}^{\sin\phi_h} \sin\phi_h \right) \\ &= \sigma_0 \left(1 + \varepsilon_1 A_{UU}^{\cos\phi_h} \cos\phi_h + \varepsilon_2 A_{UU}^{\cos 2\phi_h} \cos 2\phi_h + \lambda\varepsilon_3 A_{LU}^{\sin\phi_h} \sin\phi_h \right) \end{aligned}$$

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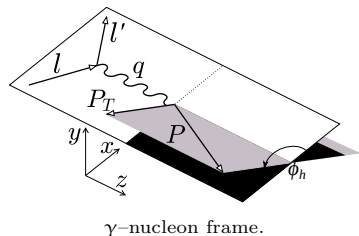
γ -nucleon frame.

$$l(l) + N(P_N) \rightarrow l'(l') + h(P) + X$$

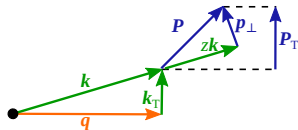
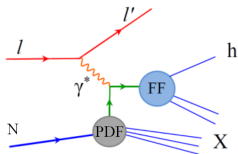
The cross section is [A. Bacchetta *et al.*, JHEP 02 (2007) 093]

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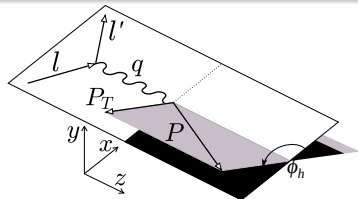
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TMD factorisation ($P_T \ll Q$): $\sigma \propto$ TMD PDF \otimes hard part \otimes TMD FF.



$$P_T \approx z k_T + P_{\perp}$$



γ -nucleon system.

$$C[wfD] = x \sum_q e_q^2 \int d^2 \mathbf{k}_T d^2 \mathbf{P}_{\perp} \delta^{(2)}(z \mathbf{k}_T + \mathbf{P}_{\perp} - \mathbf{P}_T) w(k_T, P_{\perp}) f^q(x, k_T, Q^2) D^{q \rightarrow h}(z, P_{\perp}, Q^2)$$

Up to order $\frac{1}{Q}$ and in Wandzura–Wilczek-type approximation:

TMD PDFs:

f_1 : unpolarised

h_1^{\perp} : Boer–Mulders

TMD FFs:

D_1 : unpolarised

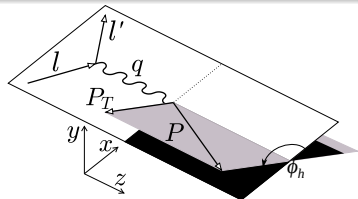
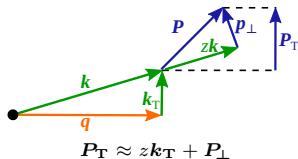
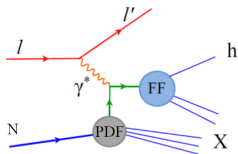
H_1^{\perp} : Collins

$$F_{UU,T} = C[f_1 D_1] \quad F_{UU,L} = 0 \quad F_{LU}^{\sin \phi_h} = 0 + \dots \leftarrow \text{pure } qgq$$

$$F_{UU}^{\cos 2\phi_h} = C \left[\frac{2(\hat{h} \cdot \mathbf{k}_T)(\hat{h} \cdot \mathbf{P}_{\perp}) - (\mathbf{k}_T \cdot \mathbf{P}_{\perp})}{z M M_h} h_1^{\perp} H_1^{\perp} \right] \leftarrow +\text{Cahn as } 1/Q^2$$

$$F_{UU}^{\cos \phi_h} = -\frac{2M}{Q} C \left[\underbrace{\frac{(\hat{h} \cdot \mathbf{k}_T)}{M} f_1 D_1}_{\text{Cahn effect}} + \underbrace{\frac{k_T^2 (\hat{h} \cdot \mathbf{P}_{\perp})}{z M^2 M_h} h_1^{\perp} H_1^{\perp}}_{\text{Boer–Mulders effect}} + \dots \right] \quad \hat{h} = \frac{\mathbf{P}_T}{|\mathbf{P}_T|}$$

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Up to order $\frac{1}{Q}$ and in Wandzura–Wilczek-type approximation:

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TMD FFs:

D_1 : unpolarised

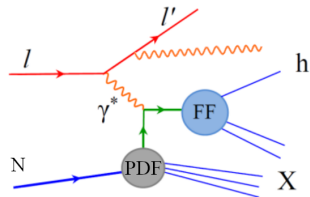
H_1^\perp : Collins

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 \end{aligned}$$

TMD interpretation defined at tree level

→ QED radiative effects need to be accounted for:

- renormalisation of the vertices,
- radiation of photons along the ℓ , ℓ' and γ^* ,
- changes in x , Q^2 , tail from elastic scattering,
- orientation of γ -nucleon system distorted.



Inclusive corrections – mostly used in the past.

- [A.A. Akhundov *et al.*, *Fortschr. Phys.* 44 (1996) 373]
 - Semi-analytical approach, parametrised in TERAD program.
- MC generators RADGEN, POLRAD (for longitudinally polarised e^\pm)
[I. Akushevich, Böttcher, Ryekbosh, *hep-ph/9906408* (1998)]

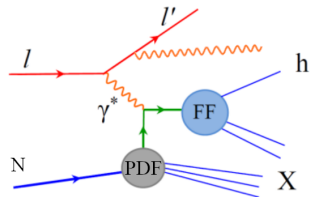
Taking into account hadron phase space, z -, P_T - and ϕ_h -dependences is needed for TMDs.

- HAPRAD [I. Akushevich, N. Shumeiko, A. Soroko, *Eur.Phys.J.C*10 (1999) 681–687]:
 - 5D (x, y, z, P_T^2, ϕ_h) corrections for inclusive h leptoproduction,
 - [I. Akushevich, A. Ilyichev, M. Osipenko, *Phys. Lett. B*672 (2009) 35-44.]: tail from excl. $ep \rightarrow e\pi\gamma$
- DJANGO (DJANGO6)
 - radiative effects from HERACLES.
 - hadronic final state using LEPTO (JETSET).
 - [K. Charchula, G.A. Schuler, H. Spiesberger, *Comput.Phys.Commun.* 81 (1994) 381-402]
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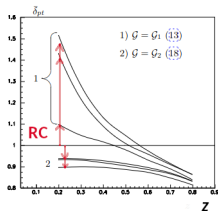


Fig. 3. Radiative correction to $\langle p_T^2 \rangle$ defined in eq. (23) for HERMES kinematics, $\sqrt{S}=7.19$ GeV, $y=0.4$. Curves from top to bottom corresponds to $x=0.15, 0.05$ and 0.45 .

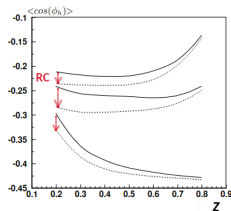
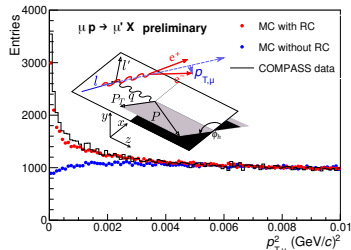


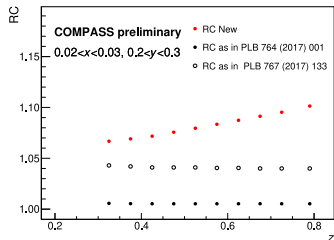
Fig. 4. Azimuthal asymmetry $\langle \cos(\phi_h) \rangle$ vs z for $y=0.2$ within HERMES kinematics; $\sqrt{S}=7.19$ GeV. Dashed (solid) lines correspond to born(observable) asymmetries. Curves from top to bottom correspond to $x=0.7, 0.45$ and 0.05 .

[I. Akushevich, N. Shumeiko, A. Soroko,
Eur.Phys.J.C10 (1999) 681–687]

- Sizeable corrections in both approaches.
- MC–data checks are very important.
- Model-dependence \rightarrow iterative approach (e.g. the P_T -slope, possibly ϕ_h -modulations).
- Useful to publish the corrections used.

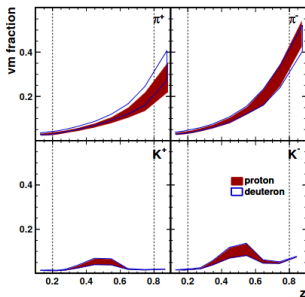
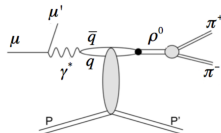


DJANGO H reproduces e^\pm from radiative γ in COMPASS data.

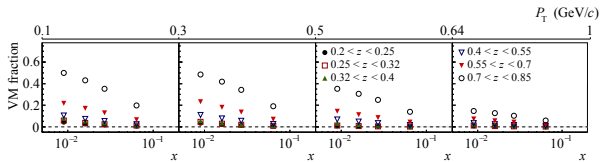


DJANGO H RC for h multiplicity
[M. Stolarski (COMPASS), MENU2023]

- Exclusive vector mesons (EVM) inherit polarization from γ^*
→ large amplitudes of azimuthal modulations for decay products.
- Significant contributors: $\rho \rightarrow \pi^+\pi^-$ and $\phi \rightarrow K^+K^-$



[HERMES, Phys.Rev.D87 (2013) 074029]



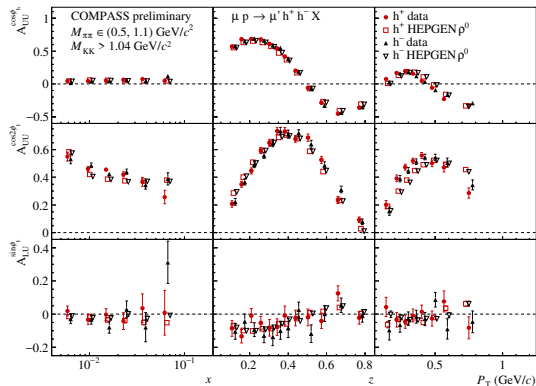
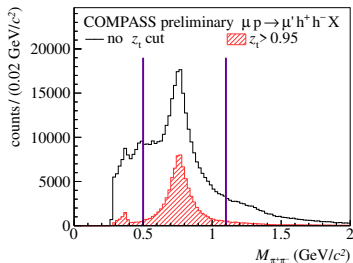
[COMPASS, Nucl.Phys.B 956 (2020) 115039]

- Strong x -, z - and P_T -dependence.
- Up to 50% of observed hadrons at low x , low P_T and high z are from EVMs!
- Important for multiplicities, cross sections.

Azimuthal dependence studied:

[COMPASS, Nucl.Phys.B 956 (2020) 115039] [V. Benešová (COMPASS), DIS2024]

- Events with only $\mu'h^+h^-$ reconstructed in the spectrometer.
- $z_t = z_1 + z_2 > 0.95$ and ρ^0 mass selected.



→ large modulations seen for the h^+ and h^- , here reasonably reproduced by HEPGEN MC

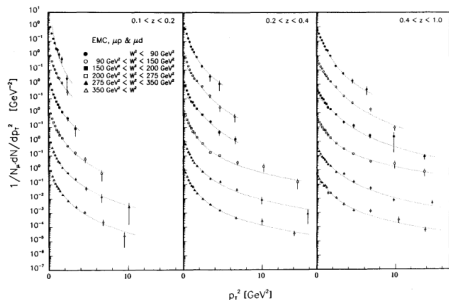
[A. Sandacz, P. Sznajder, arXiv:1207.0333].

First measurements:

[SLAC, Phys.Rev.Lett. 31 (1973) 786].

[Cornell, Phys.Rev.Lett.37 (1976) 651].

EMC

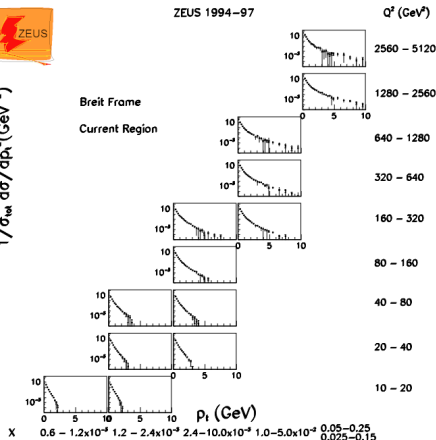


Hadron multiplicity differential in P_T^2
in bins of z : W^2 [EMC, Z.Phys.C52 (1991) 361-388]

- $\mu p/d \rightarrow \mu' h X$, $\sqrt{s} = 14-24$ GeV.
- Inclusive RC.
- no EVM corr.



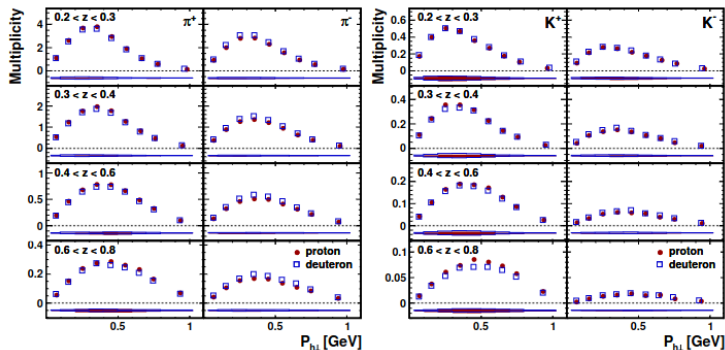
$1/\sigma_{tot} d\sigma/dp_T^2$ (GeV⁻²)



Hadron multiplicity differential in P_T^2
in bins of x : Q^2 [ZEUS, Eur.Phys.J.C11 (1999) 251]

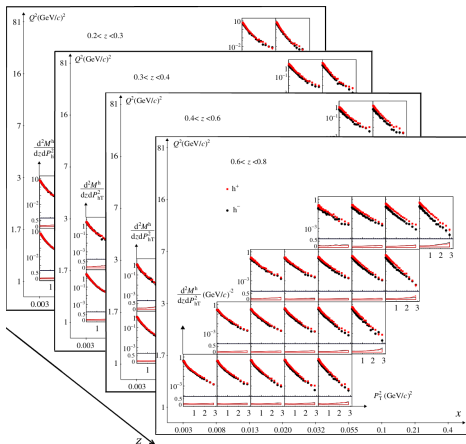
- $e p \rightarrow e' h X$, $\sqrt{s} = 14-300$ GeV.
- RC from DJANGO
- no EVM corr. (likely small large Q^2).

HERMES



π^\pm and K^\pm multiplicity in $z : P_T$
[HERMES, Phys.Rev.D87 (2013) 074029]

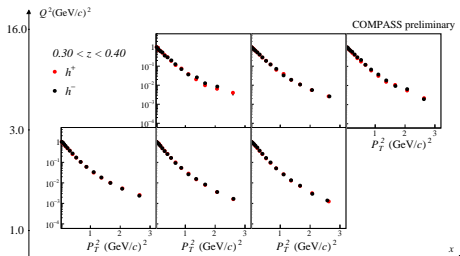
- $e p/d \rightarrow e' \pi/K^\pm X$, $\sqrt{s} = 18$ GeV.
- 3D binning in $Q^2 : z : P_T^2$ and $x : z : P_T^2$.
- Inclusive RC from RADGEN.
- EVM correction: PYTHIA (tuned to EVM prod. at HERMES).



h^\pm multiplicity in $x : Q^2 : z : P_T^2$
[COMPASS, Phys.Rev.D97 (2018) 032006]

- $\mu \text{ } ^6\text{LiD} \rightarrow \mu' h^\pm X, \sqrt{s} = 18 \text{ GeV}$.
- Inclusive RC from RADGEN.
- EVM correction: HEPGEN (no SDMEs).

COMPASS

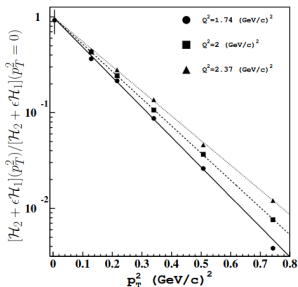


h^\pm distributions in $x : Q^2 : z : P_T^2$ (2nd z bin)
[A. Moretti (COMPASS), Proc. of ICNFP 2020]

- $\mu p \rightarrow \mu' h^\pm X, \sqrt{s} = 18 \text{ GeV}$.
- Normalised to the lowest- P_T point.
- Kinematic domain to be expanded.
- No RC, coming soon (DJANGO)
- Visible EVM decays excluded.
- Remaining EVMs subtracted (HEPGEN with SDMEs).
- More in Andrea Bressan's talk.

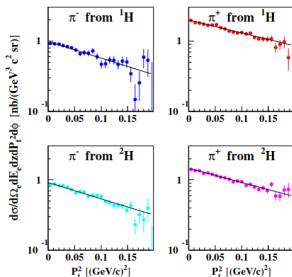
JLab6 ($\sqrt{s} = 3.5$ GeV): CLAS Hall C

JLab E06-010



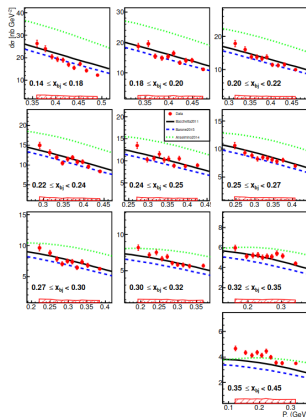
[CLAS, Phys.Rev.D 80 (2009) 032004]

- $e p \rightarrow e' \pi^+ X$.
- Binning in $Q^2 : x : z : P_T$.
- RC from HAPRAD.
- No EVM correction.



[Hall C, Phys.Rev.C85 (2012) 015202]

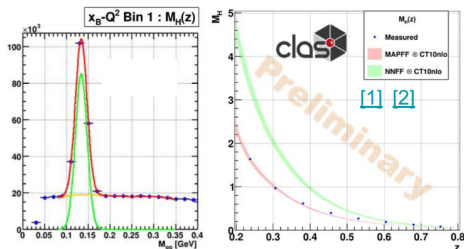
- $e p/d \rightarrow e' \pi^\pm X$.
- Binning in P_T^2 .
- Semi-incl. RC (SIMC, HAPRAD).
- EVM correction: SIMC+PYTHIA (tuned).



[Phys.Rev.C 95 (2017) 035209]

- $e {}^3\text{He} \rightarrow e' \pi^\pm X$.
- Binning in $x : P_T^2$.
- RC from HAPRAD.
- No EVM correction.

CLAS12 – ongoing work



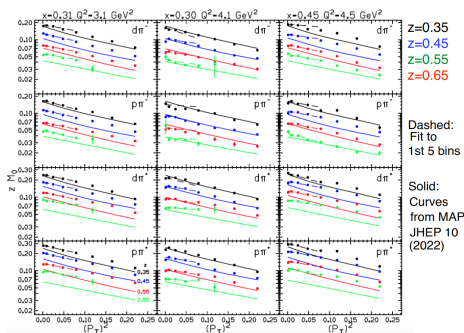
[M. Diehl (CLAS), Transversity 2024]

- $e p \rightarrow e' \pi^0 X$.
- No EVM contribution to π^0 .

More to come...

- $e p/d \rightarrow e' \pi^\pm X$.
- High-statistics, fully differential.
- RC from HAPRAD.
- EVM contribution being studied.
- Separation of $F_{UU,T}, F_{UU,L}$.

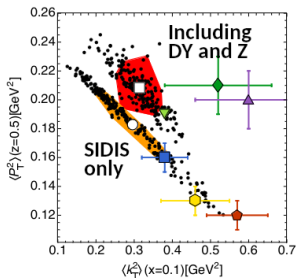
Hall C 11 GeV – ongoing work



[E.R. Kinney, DIS2024]

- $e p/d \rightarrow e' \pi^\pm / K^\pm X$.
- Multi-D binning.
- RC?
- Missing mass cut to avoid exclusive proc.
- Soon also π^0 production.

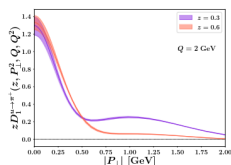
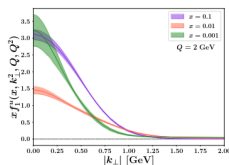
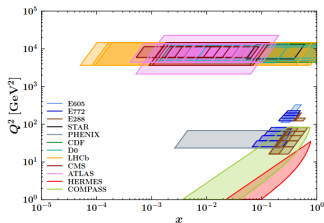
Fitting the unpolarised TMD PDFs



$\langle k_T^2 \rangle$ and $\langle P_T^2 \rangle$ from global fits
 → strong anti-correlation with SIDIS-only.

[A. Bacchetta *et al.*, JHEP 06 (2017) 081]

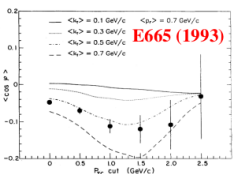
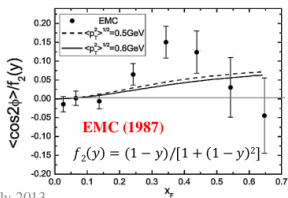
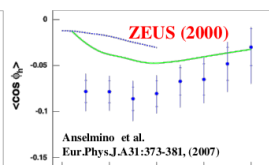
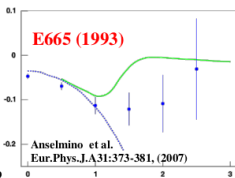
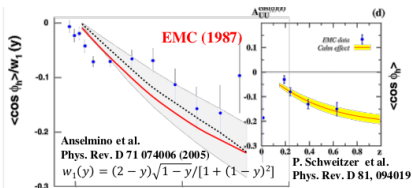
- Precision era: N³LL, global analysis
 - fixed-target Drell–Yan,
 - collider Drell–Yan and Z production,
 - e^+e^- (?)
- Similar approach (SIDIS and DY):
 [I. Scimemi, A. Vladimirov, JHEP 06 (2020) 137]
- More on Thursday morning...



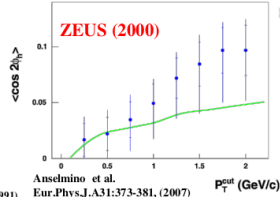
[MAP, JHEP 10 (2022) 127]

Flavour separation

- Relies mostly on HERMES data.
- [A. Signori *et al.*, JHEP 11 (2013) 194]
 - $\langle P_{T, \text{unfav.}}^2 \rangle = 1.2 \langle P_{T, \text{fav.}}^2 \rangle \approx \langle P_{T, u \rightarrow K}^2 \rangle$
- Soon: $\pi^{\pm, 0}$, K^{\pm} from JLab, COMPASS



J. Chay, S. D. Ellis, and J. W. Stirling,
Phys. Rev. D 45, 46 (1992), Phys. Lett. B 269, 175 (1991).
Bakur Parsamyan

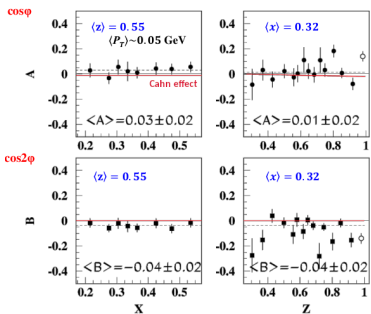


23 July 2013

29

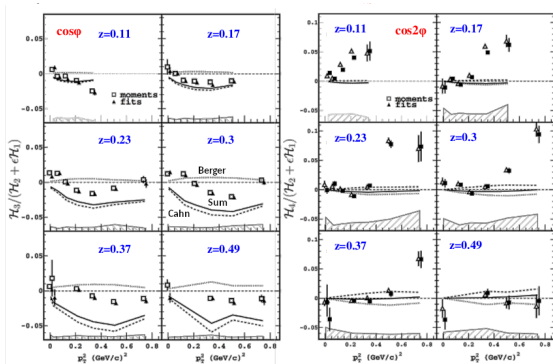
First measurements (compilation [B. Parsamyan, IWHSS2013])

- [EMC, Phys.Lett.B130 (1983) 118–122]
- [EMC, Z.Phys.C34 277 (1987)]
- [E665, Phys.Rev.D48, 5057 (1993)]
- [ZEUS, Phys.Lett.B481 (2000) 199]
 - $e p \rightarrow e' h X$, $\sqrt{s} = 300 \text{ GeV}$,
 - RC from DJANGO.
 - no EVM corr. (likely small large Q^2).



[Hall C, Phys.Lett.B665 (2008) 20]

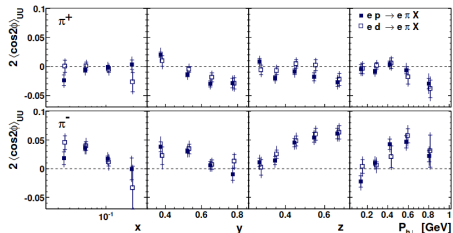
- $e p/d \rightarrow e' \pi^\pm X$.
- Binning in P_T^2 .
- Semi-incl. RC.
- EVM correction.
- Similar: [Hall C, Phys.Rev.C85 (2012) 015202]



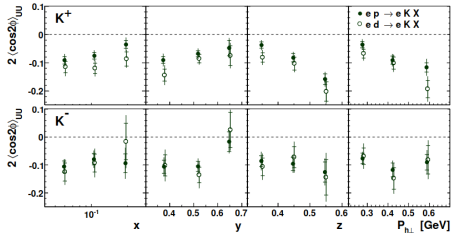
Data: [CLAS, Phys.Rev.D 80 (2009) 032004]

Predictions: [R.N. Cahn, Phys.Rev.D40 (1989) 3107.],
[M. Anselmino *et al.*, Phys.Rev.D71 (2005) 074006],
[A. Brandenburg *et al.*, Phys.Lett.B347 (1995) 413].

- $e p \rightarrow e' \pi^\pm X$.
- Binning in $Q^2 : x : z : P_T$.
- RC from HAPRAD.
- No EVM correction.



$\varepsilon_2 A_{UU}^{\cos 2\phi_h}$ for π^\pm .



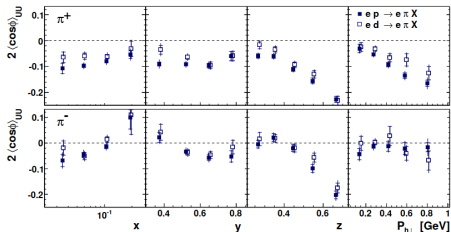
$\varepsilon_2 A_{UU}^{\cos 2\phi_h}$ for K^\pm .

[HERMES, Phys.Rev.D87 (2013) 012010]

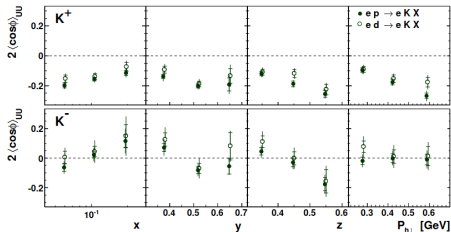
- $e \text{ p/d} \rightarrow e' \pi^\pm / K^\pm / h^\pm X$.
($\sqrt{s} = 7.5 \text{ GeV}$)
- Binning in $x : y : z : P_T$.
- Inclusive (?) RC from RADGEN.
- No EVM correction.

- Similar on p and d $\rightarrow h_1^{\perp,u} \approx h_1^{\perp,d}$.
- Known: $H_{1,\text{fav.}}^\perp \approx -H_{1,\text{unfav.}}^\perp$
- π behaviour \approx expected [V. Barone, A. Prokudin, B.Q. Ma, Phys.Rev.D78 (2008) 045022]
- K very different from π .
- **Interesting flavour dependence...**

$$F_{UU}^{\cos 2\phi_h} = C \left[\frac{2(\hat{h} \cdot \mathbf{k}_T)(\hat{h} \cdot \mathbf{P}_\perp) - (\mathbf{k}_T \cdot \mathbf{P}_\perp)}{zMM_h} h_1^\perp H_1^\perp \right] + \frac{M^2}{Q^2} C \left[-\frac{2(\hat{h} \cdot \mathbf{k}_T)^2 - k_T^2}{M^2} f_1 D_1 + \dots \right]$$



$\varepsilon_1 A_{UU}^{\cos \phi_h}$ for π^\pm .



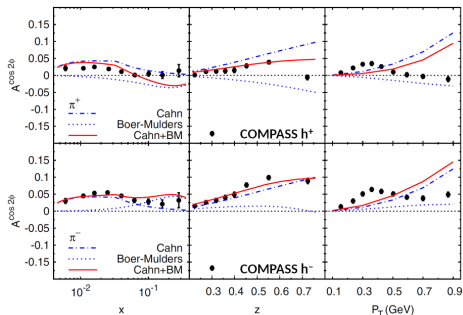
$\varepsilon_1 A_{UU}^{\cos \phi_h}$ for K^\pm .

[HERMES, Phys.Rev.D87 (2013) 012010]

- $e \text{ p/d} \rightarrow e' \pi^\pm / K^\pm / h^\pm X$.
($\sqrt{s} = 7.5 \text{ GeV}$)
- Binning in $x : y : z : P_T$.
- Inclusive (?) RC from RADGEN.
- No EVM correction.

- Small difference between p and d targets.
- Difference between π^+ , π^-
→ Boer–Mulders effect, different $\langle k_T^2 \rangle$?
- K similar to π
(while being different for $\cos 2\phi_h$).

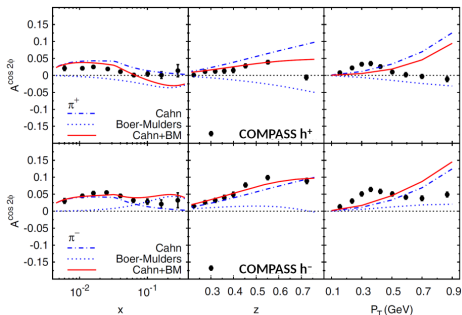
$$F_{UU}^{\cos \phi_h} = -\frac{2M}{Q} C \left[\frac{(\hat{h} \cdot \mathbf{k}_T)}{M} f_1 D_1 + \frac{k_T^2 (\hat{h} \cdot \mathbf{P}_\perp)}{z M^2 M_h} h_1^\perp H_1^\perp + \dots \right]$$



COMPASS $A_{UU}^{\cos 2\phi_h}$ for h^\pm on ${}^6\text{LiD}$ (d) target
 [COMPASS, Nucl.Phys.B886 (2014) 1046–1077],
 fitted [V. Barone *et al.*, Phys.Rev.D91 (2015) 074019].

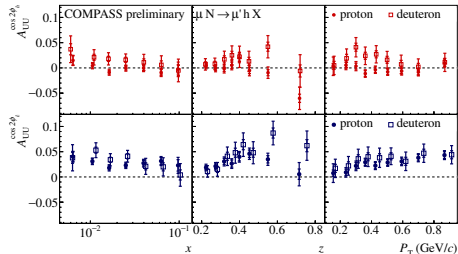
- μ ${}^6\text{LiD} \rightarrow \mu' h^\pm X$. ($\sqrt{s} = 18$ GeV)
- Binning in $x : z : P_T$.
- No RC.
- No EVM correction.

$$F_{UU}^{\cos 2\phi_h} = C \left[\frac{2(\hat{h} \cdot \mathbf{k}_T)(\hat{h} \cdot \mathbf{P}_\perp) - (\mathbf{k}_T \cdot \mathbf{P}_\perp)}{zMM_h} h_1^\perp H_1^\perp \right] + \frac{M^2}{Q^2} C \left[-\frac{2(\hat{h} \cdot \mathbf{k}_T)^2 - k_T^2}{M^2} f_1 D_1 + \dots \right]$$



COMPASS $A_{UU}^{\cos 2\phi_h}$ for h^\pm on ${}^6\text{LiD}$ (d) target
 [COMPASS, Nucl.Phys.B886 (2014) 1046–1077],
 fitted [V. Barone *et al.*, Phys.Rev.D91 (2015) 074019].

- $\mu {}^6\text{LiD} \rightarrow \mu' h^\pm X$. ($\sqrt{s} = 18$ GeV)
- Binning in $x : z : P_T$.
- No RC.
- No EVM correction.

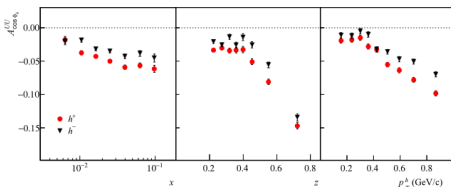


COMPASS $A_{UU}^{\cos 2\phi_h}$ for h^\pm on d and p targets
 with EVM contribution subtracted (no RC).

Updates:

- EVM contribution subtracted ($x : z : P_T$)
 [COMPASS, Nucl.Phys.B956 (2020) 115039]
- Data on **p target** being analysed
 - Visible EVM decays excluded.
 - Remaining EVMs subtracted (HEPGEN with SDMEs).
 - RC from DJANGO.
 - **More in Andrea Bressan's talk.**

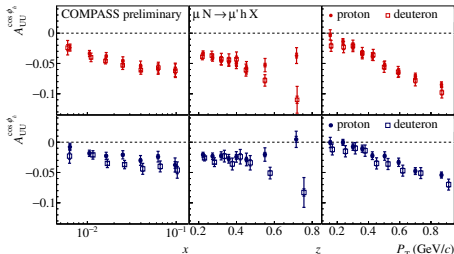
$$F_{UU}^{\cos 2\phi_h} = C \left[\frac{2(\hat{h} \cdot \mathbf{k}_T)(\hat{h} \cdot \mathbf{P}_\perp) - (\mathbf{k}_T \cdot \mathbf{P}_\perp)}{zMM_h} h_1^\perp H_1^\perp \right] + \frac{M^2}{Q^2} C \left[-\frac{2(\hat{h} \cdot \mathbf{k}_T)^2 - k_T^2}{M^2} f_1 D_1 + \dots \right]$$



COMPASS $A_{UU}^{\cos\phi_h}$ for h^\pm on ${}^6\text{LiD}$ (d) target
 [COMPASS, Nucl.Phys.B886 (2014) 1046–1077]
 (no EVM correction, no RC).

- μ ${}^6\text{LiD} \rightarrow \mu' h^\pm X$. ($\sqrt{s} = 18$ GeV)
- Binning in $x : z : P_T$.
- No RC.
- No EVM correction.

$$F_{UU}^{\cos\phi_h} = -\frac{2M}{Q} C \left[\frac{(\hat{h} \cdot \mathbf{k}_T)}{M} f_{1D1} + \frac{k_T^2 (\hat{h} \cdot \mathbf{P}_\perp)}{z M^2 M_h} h_1^\perp H_1^\perp + \dots \right]$$

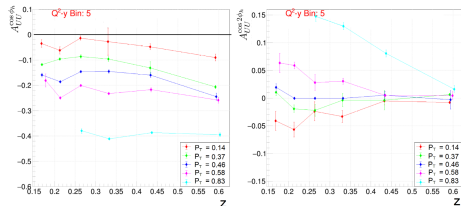


COMPASS $A_{UU}^{\cos\phi_h}$ for h^\pm on d and p targets
 with EVM contribution subtracted (no RC).

Updates:

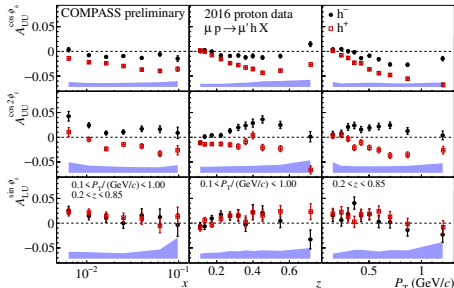
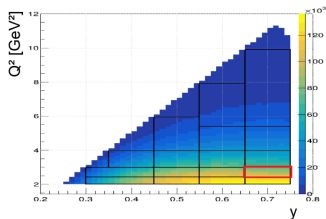
- EVM contribution subtracted ($x : z : P_T$)
 [COMPASS, Nucl.Phys.B956 (2020) 115039]
- Data on **p target** being analysed
 - Visible EVM decays excluded.
 - Remaining EVMs subtracted (HEPGEN with SDMEs).
 - RC from DJANGO.
 - **More in Andrea Bressan's talk.**

More interesting results on the horizon...



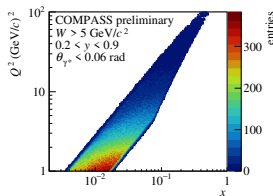
Stefan Diehl's talk.

- Ongoing CLAS12 analysis.
- $e p/d \rightarrow e' \pi^+ X$.
- 4D binning $Q^2 : y : z : P_T$.



See Andrea Bressan's talk...

- $\mu p \rightarrow \mu' h^\pm X$. ($\sqrt{s} = 18$ GeV)
- New RC from DJANGOH applied.



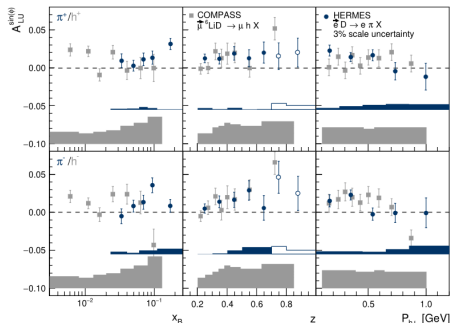
$A_{LU}^{\sin \phi_h}$ is a pure twist-3 object (qgq-correlations):

$$F_{LU}^{\sin \phi} = \frac{2M}{Q} \mathcal{C} \left(-\frac{\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M_h} \left(x e H_1^\perp + \frac{M_h}{M} f_1 \frac{\tilde{G}^\perp}{z} \right) + \frac{\hat{\mathbf{h}} \cdot \mathbf{p}_T}{M} \left(x g^\perp D_1 + \frac{M_h}{M} h_1^\perp \frac{\tilde{E}}{z} \right) \right)$$

twist-3 pdf
Collins FF
unpolarized dist. function
twist-3 FF
twist-3 t-odd dist. function
Boer-Mulders
twist-3 FF
→ TMDs and FFs

(borrowed from Stefan Diehl's slides)

- HERMES ($Q \approx 1.6$ GeV)
 - [HERMES, Phys.Lett.B 797 (2019) 134886]
 - $e \text{ p/d} \rightarrow e' \pi^\pm / K^\pm / p^\pm \text{ X}$.
- COMPASS ($Q \approx 1.7$ GeV)
 - [COMPASS, Nucl.Phys.B886 (2014) 1046–1077]
 - $\mu \text{ }^6\text{LiD} \rightarrow \mu' \text{ h}^\pm \text{ X}$.
- CLAS ($Q \approx 1.4$ GeV)
 - [CLAS, Phys. Rev. D89 (2014) 072011]
 - $e \text{ p} \rightarrow e' \pi^{\pm,0} \text{ X}$.
- RC are negligible for $\sin \phi_h$.
- EVM correction not applied.
- Limited statistics (1D/2D binning).



[HERMES, Phys.Lett.B 797 (2019) 134886]

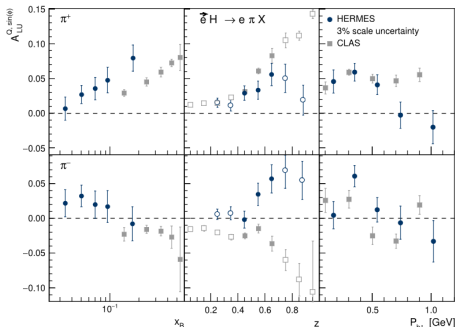
$A_{LU}^{\sin \phi_h}$ is a pure twist-3 object (qgq-correlations):

$$F_{LU}^{\sin \phi} = \frac{2M}{Q} \mathcal{C} \left(-\frac{\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M_h} \left(x e H_1^\perp + \frac{M_h}{M} f_1 \frac{\tilde{G}^\perp}{z} \right) + \frac{\hat{\mathbf{h}} \cdot \mathbf{p}_T}{M} \left(x g^\perp D_1 + \frac{M_h}{M} h_1^\perp \frac{\tilde{E}}{z} \right) \right)$$

twist-3 pdf
Collins FF
unpolarized dist. function
twist-3 FF
twist-3 t-odd dist. function
Boer-Mulders
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→ TMDs and FFs

(borrowed from Stefan Diehl's slides)

- HERMES ($Q \approx 1.6$ GeV)
 - [HERMES, Phys.Lett.B 797 (2019) 134886]
 - $e p/d \rightarrow e' \pi^\pm / K^\pm / p^\pm X$.
- COMPASS ($Q \approx 1.7$ GeV)
 - [COMPASS, Nucl.Phys.B886 (2014) 1046–1077]
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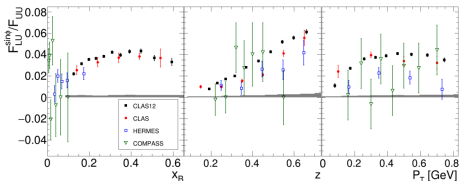
[HERMES, Phys.Lett.B 797 (2019) 134886]

$$F_{LU}^{\sin \phi} = \frac{2M}{Q} \mathcal{C} \left(-\frac{\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M_h} \left(x e H_1^+ + \frac{M_h}{M} f_1 \frac{\tilde{G}^\perp}{z} \right) + \frac{\hat{\mathbf{h}} \cdot \mathbf{p}_T}{M} \left(x g^\perp D_1 + \frac{M_h}{M} h_1^\perp \frac{\tilde{E}}{z} \right) \right)$$

twist-3 pdf
Collins FF
unpolarized dist. function
twist-3 FF
twist-3 t-odd dist. function
Boer-Mulders
twist-3 FF
→ TMDs and FFs

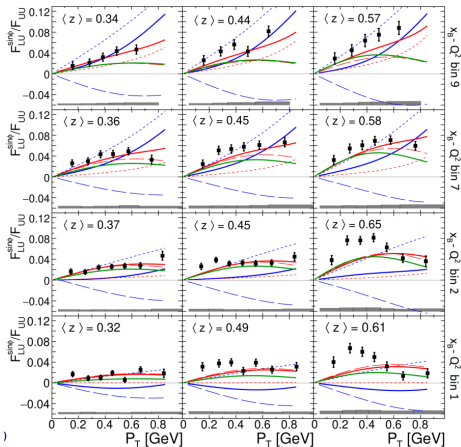
(borrowed from Stefan Diehl's slides)

NEW from CLAS12:



π^+ (CLAS, HERMES) and h^+ (COMPASS).

- $e p \rightarrow e' \pi^+ X$
[CLAS, Phys.Rev.Lett. 128 (2022) 062005]
- $e p \rightarrow e' \pi^{\pm,0} / K^+ X$
(preliminary, Stefan Diehl's talk)
- 4D binning ($Q^2 : x : z : P_T$).
- EVM contribution being studied.

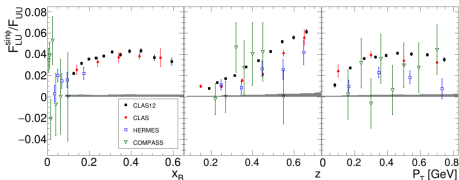


$$F_{LU}^{\sin \phi} = \frac{2M}{Q} C \left(-\frac{\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M_h} \left(x e H_1^+ + \frac{M_h}{M} f_1 \tilde{G}^+ \right) + \frac{\hat{\mathbf{h}} \cdot \mathbf{p}_T}{M} \left(x g^+ D_1 + \frac{M_h}{M} h_1^+ \tilde{E} \right) \right)$$

twist-3 pdf
Collins FF
unpolarized dist. function
twist-3 FF
twist-3 t-odd dist. function
Boer-Mulders
twist-3 FF
→ TMDs and FFs

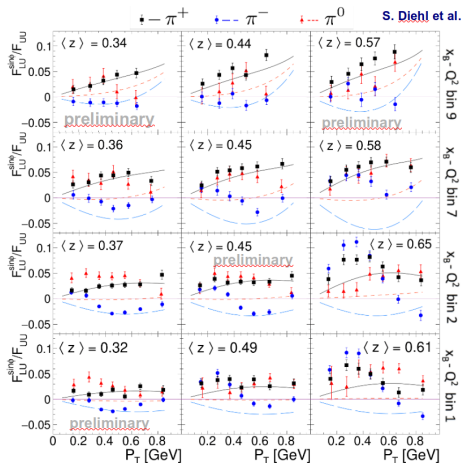
(borrowed from Stefan Diehl's slides)

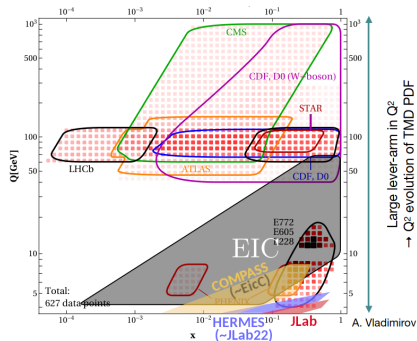
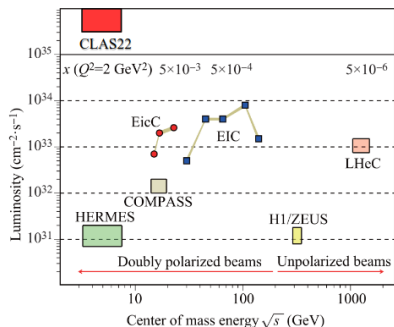
NEW from CLAS12:



π^+ (CLAS, HERMES) and h^+ (COMPASS).

- $e p \rightarrow e' \pi^+ X$
[CLAS, Phys.Rev.Lett. 128 (2022) 062005]
- $e p \rightarrow e' \pi^{\pm,0} / K^+ X$
(preliminary, Stefan Diehl's talk)
- 4D binning ($Q^2 : x : z : P_T$).
- EVM contribution being studied.





- **SoLID @ JLab12** (Haylin Gao's talk)

- Large luminosity \rightarrow multi-D.
- High x region.

- **JLab22**

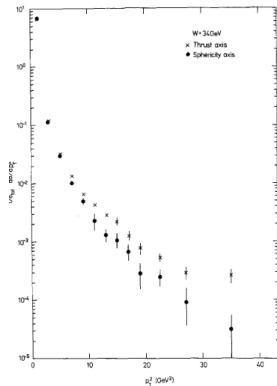
- HERMES-like kinematic region.
- Really large luminosity.

- **EicC** (Yuxiang Zhao's talk)

- COMPASS-like kinematic region.

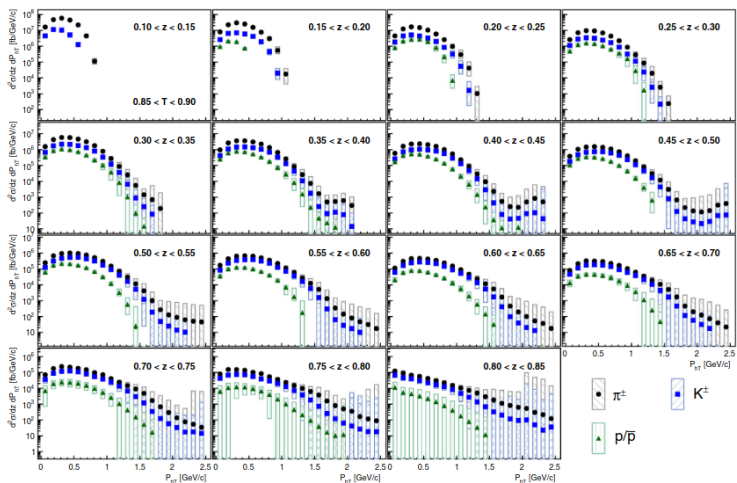
- **EIC** (Elke Aschenauer's talk)

- [Snowmass 2021, 2203.13199 [hep-ph]]
- Large lever arm in Q^2
 \rightarrow evolution,
 \rightarrow wider region of $P_T \ll Q$
 \rightarrow closer to high-energy DY, Z data.
- **Gluon TMDs.**
 - E.g. $1 N \rightarrow 1'$ jet jet X, jets back-to-back
 - Quarkonium + γ , quarkonium + jet [A. Mukherjee, IWHSS2023]
 - Daniel Boer's talk.



TASSO, 1984

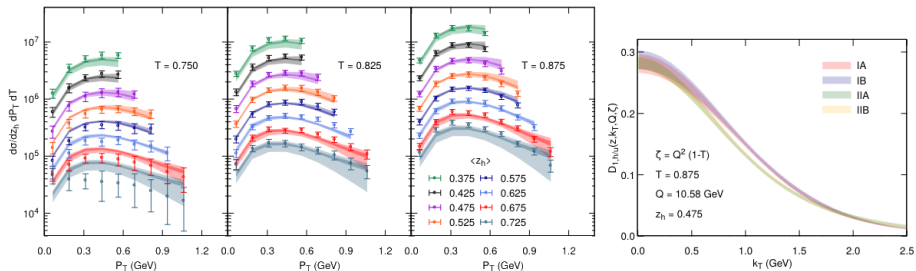
- [TASSO, Z.Phys.C22 (1984) 307–340]
- [TASSO, Z.Phys. C47 (1990) 187–198]
- [PLUTO, Z.Phys. C22 (1984) 103]
- [MARK-II, Phys.Rev.D37 (1988) 1]



[Belle, Phys.Rev.D99 (2019) 11, 112006]

- P_{\perp} defined with respect to the thrust axis.
- Factorisation formalism developed, possibility to enter global fits?

[Z.B. Kang, D.Y. Shao, F. Zhao, JHEP12 (2020) 127] [M. Boglione, A. Simonelli, JHEP 02 (2021) 076]



[M. Boglione, J.O. Gonzalez-Hernandez, A. Simonelli, *Phys.Rev.D*106 (2022) 074024]

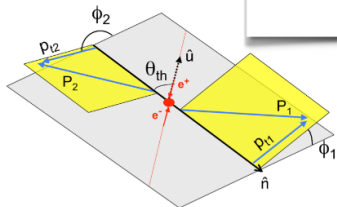
See D. Boer, NPB 806, 23 (2009)

RF12

$$\frac{d\sigma(e^+e^- \rightarrow h_1 h_2 X)}{d\Omega dz_1 dz_2 d\phi_1 d\phi_2} = \sum_{q,\bar{q}} \frac{3\alpha^2 e_q^2}{Q^2} \frac{z_1^2 z_2^2}{4} \left[(1 + \cos^2\theta) D_1^{q,(0)}(z_1) \bar{D}_1^{q,(0)}(z_2) + \sin^2(\theta) \cos(\phi_1 + \phi_2) H_1^{\perp,(1),q}(z_1) \bar{H}_1^{\perp,(1),q}(z_2) \right]$$

All quantities in e^+e^- center of mass

θ : angle between the e^+e^- axis and the thrust axis;
 $\phi_{1,2}$: azimuthal angles between $P_{h_1(h_2)}$ and the scattering plane



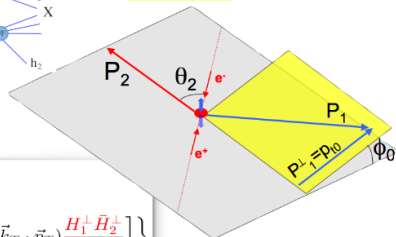
θ_2 : angle between the e^+e^- axis and P_{h_2} ;
 ϕ_0 : angle between the plane spanned by P_{h_2} and the e^+e^- axis, and the direction of P_{h_1} perpendicular to P_{h_2} .

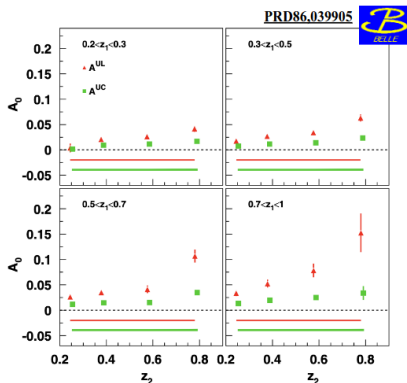
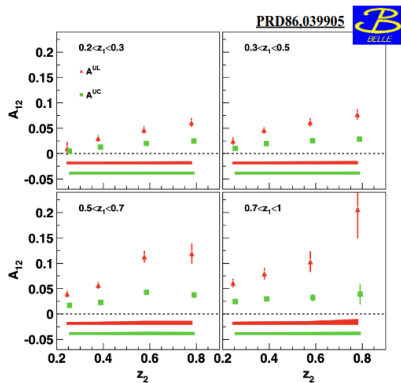
All quantities in e^+e^- center of mass

$$\frac{d\sigma(e^+e^- \rightarrow h_1 h_2 X)}{d\Omega dz_1 dz_2 d^2\vec{q}_T} = \frac{3\alpha^2}{Q^2} z_1^2 z_2^2 \left\{ A(y) \mathcal{F}[D_1 \bar{D}_2] + B(y) \cos(2\phi_0) \mathcal{F} \left[(2\hat{h} \cdot \vec{k}_T \hat{h} \cdot \vec{p}_T - \vec{k}_T \cdot \vec{p}_T) \frac{H_1^\perp \bar{H}_2^\perp}{M_1 M_2} \right] \right\}$$

[I. Garzia, IWHSS2023.]

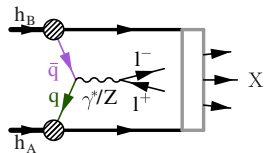
RF0





• $\mathcal{L} \sim 547 \text{ fb}^{-1}$ at $\sim 10.58 \text{ GeV}$

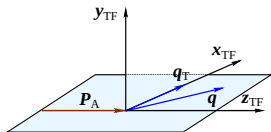
- 1st measurement: [Belle, Phys.Rev.Lett.96 (2006) 232002] [Belle, Phys.Rev.D 78 (2008) 032011]
- Belle $\pi^\pm\pi^\pm, \pi^\pm\eta$ [Phys.Rev.D 100 (2019) 9, 092008]
- BaBar $\pi\pi$ [Phys.Rev.D 90 (2014) 5, 052003]
- BaBar $KK, K\pi$ [Phys.Rev.D 92 (2015) 11, 111101]
 $\rightarrow D_1^u \rightarrow K$ [M. Anselmino *et al.*, Phys.Rev.D93 (2016) 3, 034025]
- BESIII $\pi\pi$ [Phys.Rev.Lett. 116 (2016) 4, 042001]



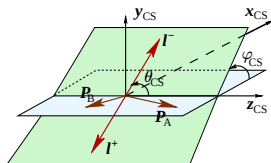
Cross-section with unpolarised target:

$$\frac{d\sigma}{dQ^2 dq_T^2 d\eta} = \frac{4\pi^2 \alpha_{em}^2}{3Q^2 s} \left(F_{UU}^1 + F_{UU}^{\cos 2\varphi_{CS}} \right).$$

$$\frac{d\sigma}{d\Omega} = \frac{3}{4\pi} \frac{1}{\lambda + 3} \left(\begin{array}{l} + \lambda \cos^2 \theta_{CS} \\ + \mu \sin 2\theta \cos \varphi_{CS} \\ + \frac{\nu}{2} \sin^2 \theta \cos 2\varphi_{CS} \end{array} \right).$$



Target frame.



Collins–Soper frame

$$\bullet \lambda = \frac{F_{UU}^1 - 2F_{UU}^{\cos 2\varphi_{CS}}}{F_{UU}^1 + 2F_{UU}^{\cos 2\varphi_{CS}}},$$

$$\bullet \mu = \frac{F_{UU}^{\cos \varphi_{CS}}}{F_{UU}^1 + 2F_{UU}^{\cos 2\varphi_{CS}}},$$

$$\bullet \nu = \frac{2F_{UU}^{\cos 2\varphi_{CS}}}{F_{UU}^1 + 2F_{UU}^{\cos 2\varphi_{CS}}}.$$

• At leading twist:

$$F_{UU}^{\cos \varphi_{CS}} = 0,$$

$$F_{UU}^1 \propto f_{1,A} \otimes f_{1,B},$$

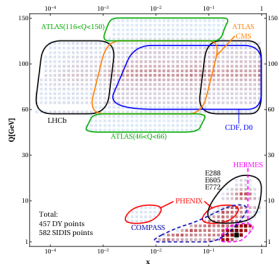
$$F_{UU}^{\cos 2\varphi_{CS}} \propto h_{1,A}^\perp \otimes h_{1,B}^\perp.$$

• $h h \rightarrow Z X$: cross section also $\propto F_{UU}^1 + F_{UU}^{\cos 2\varphi_{CS}}$

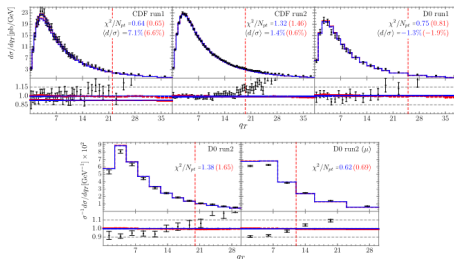
Numerous pp experiments, e.g.:

- A.S. Ito et al., *Measurement of the continuum of dimuons produced in high-energy proton-nucleus collisions*, *Phys. Rev. D* **23** (1981) 604 [inSPIRE].
- G. Moreno et al., *Dimuon production in proton-copper collisions at $\sqrt{s} = 38.8$ GeV*, *Phys. Rev. D* **43** (1991) 2815 [inSPIRE].
- E772 collaboration, *Cross-sections for the production of high mass muon pairs from 800 GeV proton bombardment of H-2*, *Phys. Rev. D* **50** (1994) 3038 [Erratum *ibid.* **D 60** (1999) 119903] [inSPIRE].
- PHENIX collaboration, *Measurements of $\mu\mu$ pairs from open heavy flavor and Drell-Yan in $p + p$ collisions at $\sqrt{s} = 200$ GeV*, *Phys. Rev. D* **99** (2019) 072003 [arXiv:1805.02448] [inSPIRE].
- CDF collaboration, *The transverse momentum and total cross section of e^+e^- pairs in the Z-boson region from pp collisions at $\sqrt{s} = 1.8$ TeV*, *Phys. Rev. Lett.* **84** (2000) 845 [hep-ex/0001021] [inSPIRE].
- CDF collaboration, *Transverse momentum cross section of e^+e^- pairs in the Z-boson region from pp collisions at $\sqrt{s} = 1.96$ TeV*, *Phys. Rev. D* **86** (2012) 052010 [arXiv:1207.7138] [inSPIRE].
- D0 collaboration, *Measurement of the inclusive differential cross section for Z bosons as a function of transverse momentum in pp collisions at $\sqrt{s} = 1.8$ TeV*, *Phys. Rev. D* **61** (2000) 032004 [hep-ex/9907009] [inSPIRE].
- D0 collaboration, *Measurement of the shape of the boson transverse momentum distribution in pp $\rightarrow Z/\gamma^* \rightarrow e^+e^- + X$ events produced at $\sqrt{s} = 1.96$ TeV*, *Phys. Rev. Lett.* **100** (2008) 102002 [arXiv:0712.0803] [inSPIRE].
- D0 collaboration, *Measurement of the normalized $Z/\gamma^* \rightarrow \mu^+\mu^-$ transverse momentum distribution in pp collisions at $\sqrt{s} = 1.96$ TeV*, *Phys. Lett. B* **693** (2010) 522 [arXiv:1006.0618] [inSPIRE].
- LHCb collaboration, *Measurement of the forward Z boson production cross-section in pp collisions at $\sqrt{s} = 7$ TeV*, *JHEP* **08** (2015) 039 [arXiv:1505.07024] [inSPIRE].
- LHCb collaboration, *Measurement of forward W and Z boson production in pp collisions at $\sqrt{s} = 8$ TeV*, *JHEP* **01** (2016) 155 [arXiv:1511.08039] [inSPIRE].
- LHCb collaboration, *Measurement of the forward Z boson production cross-section in pp collisions at $\sqrt{s} = 13$ TeV*, *JHEP* **09** (2016) 136 [arXiv:1607.06495] [inSPIRE].
- ATLAS collaboration, *Measurement of the Z/γ^* boson transverse momentum distribution in pp collisions at $\sqrt{s} = 7$ TeV with the ATLAS detector*, *JHEP* **09** (2014) 145 [arXiv:1406.3660] [inSPIRE].
- ATLAS collaboration, *Measurement of the transverse momentum and ϕ_a^* distributions of Drell-Yan lepton pairs in proton-proton collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector*, *Eur. Phys. J. C* **76** (2016) 291 [arXiv:1512.02192] [inSPIRE].
- CMS collaboration, *Measurement of the rapidity and transverse momentum distributions of Z bosons in pp collisions at $\sqrt{s} = 7$ TeV*, *Phys. Rev. D* **85** (2012) 032002 [arXiv:1110.4973] [inSPIRE].
- CMS collaboration, *Measurement of the transverse momentum spectra of weak vector bosons produced in proton-proton collisions at $\sqrt{s} = 8$ TeV*, *JHEP* **02** (2017) 096 [arXiv:1606.05864] [inSPIRE].

Used in global TMD fits, e.g.

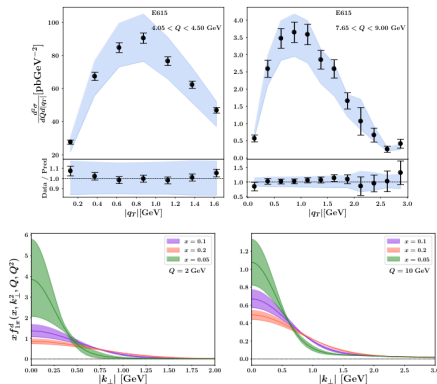


[I. Scimemi, A. Vladimirov, JHEP 06 (2020) 137]



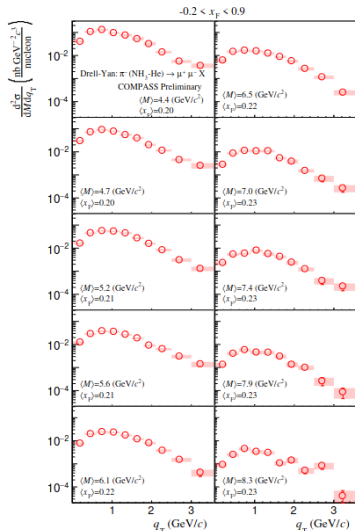
Example: fit of CDF and D0 data.

π -p Drell-Yan \rightarrow TMD distribution of the pion

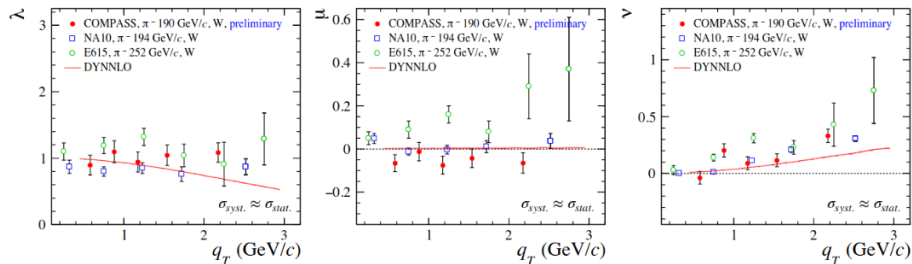


[MAP, Phys.Rev.D 107 (2023) 1, 014014]

- E615, NA10.
- **NEW: COMPASS π -N**
- Future: AMBER $\pi^\pm C \rightarrow$ access to the sea.
(both in Catarina's talk)



[V. Andrieux (COMPASS), DIS2024]



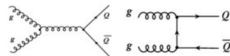
Azimuthal modulations from π -N DY (NA10, E615 and preliminary COMPASS).
 → Catarina's talk.

- Perturbative and possible non-perturbative origin of the modulations.
- There seems to be room for Boer–Mulders...

$$\nu \propto h_{1,A}^\perp \otimes h_{1,B}^\perp.$$

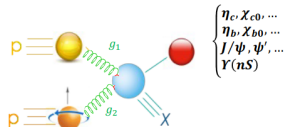
- pp, pd Drell–Yan:
 - E866 NuSea, Phys. Rev. Lett. 99 (2007) 082301
 - E866 NuSea, Phys. Rev. Lett. 102 (2009) 182001
 - **New:** [SeaQuest, SPIN2023].
 - Small asymmetry (role of the sea).

In high-energy hadron collisions, heavy quarks are dominantly produced through gg fusion:

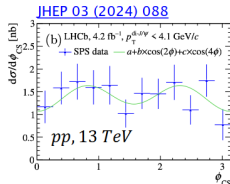


The most efficient way to access the gluon dynamics inside the proton at LHC is to **measure heavy-quark observables**

- **Inclusive quarkonia production in (un)polarized pp interaction** ($pp^{(T)} \rightarrow [Q\bar{Q}]X$) turns out to be an ideal observable to access gTMDs (assuming TMD factorization)
- TMD factorization requires $q_T(Q) \ll M_Q$. Can look at **associate quarkonia production**, where only the relative q_T needs to be small, e.g.: $pp^{(T)} \rightarrow J/\psi + J/\psi + X$



- η_c, χ_{c0}, \dots
- η_b, χ_{b0}, \dots
- $J/\psi, \psi', \dots$
- $Y(nS)$



$$d\sigma_{J/\psi+J/\psi} = a + b \times \cos(2\phi_{CS}) + c \times \cos(4\phi_{CS})$$

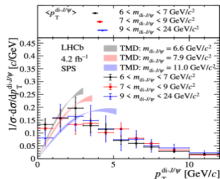
$$a = F_1 \mathcal{C}[f_1^g f_1^g] + F_2 \mathcal{C}[w_2 h_1^{1g} h_1^{1g}],$$

$$b = F_3 \mathcal{C}[w_3 f_1^g h_1^{1g}] + F_3' \mathcal{C}[w_3' h_1^{1g} f_1^g],$$

$$c = F_4 \mathcal{C}[w_4 h_1^{1g} h_1^{1g}],$$

$$\langle \cos 2\phi_{CS} \rangle = -0.029 \pm 0.050 \text{ (stat)} \pm 0.009 \text{ (syst)},$$

$$\langle \cos 4\phi_{CS} \rangle = -0.087 \pm 0.052 \text{ (stat)} \pm 0.013 \text{ (syst)},$$



...but very challenging at fixed-target kinematics!

[LHCb, JHEP 03 (2024)]

- More collider results,
- AMBER,
- SpinQuest
- LHCSpin (**Charlotte's talk**)



