



2024  
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Joint XX-th International Workshop on Hadron Structure and Spectroscopy and 5-th Workshop on Correlations in Partonic and Hadronic Interactions

Yerevan, Armenia  
30 September – 4 October, 2024

IWHS COMPASS  
Yerevan Armenia

# GPDs study at COMPASS

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For the COMPASS Collaboration

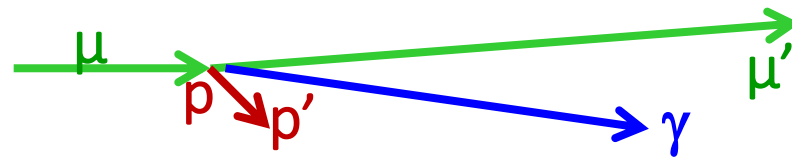
# Hard Exclusive Reactions at COMPASS at CERN

Exclusive photon (DVCS) and meson (HEMP) production at small transfer for GPD studies



Deeply Virtual Compton Scattering

$$\text{DVCS: } \mu \ p \rightarrow \mu' \ p' \ \gamma$$



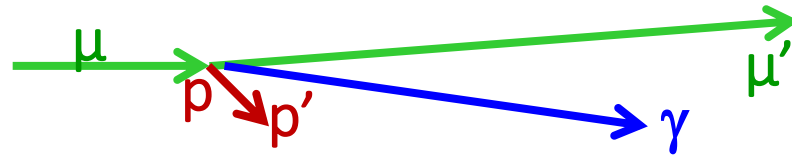
$$\text{Pseudo-Scalar Meson : } \mu \ p \rightarrow \mu' \ p' \ \pi^0$$

$$\text{Vector Meson : } \mu \ p \rightarrow \mu' \ p' \ \rho \text{ or } \omega$$



# Measurement of exclusive cross sections at COMPASS

DVCS :  $\mu p \rightarrow \mu' p' \gamma$  at small transfer

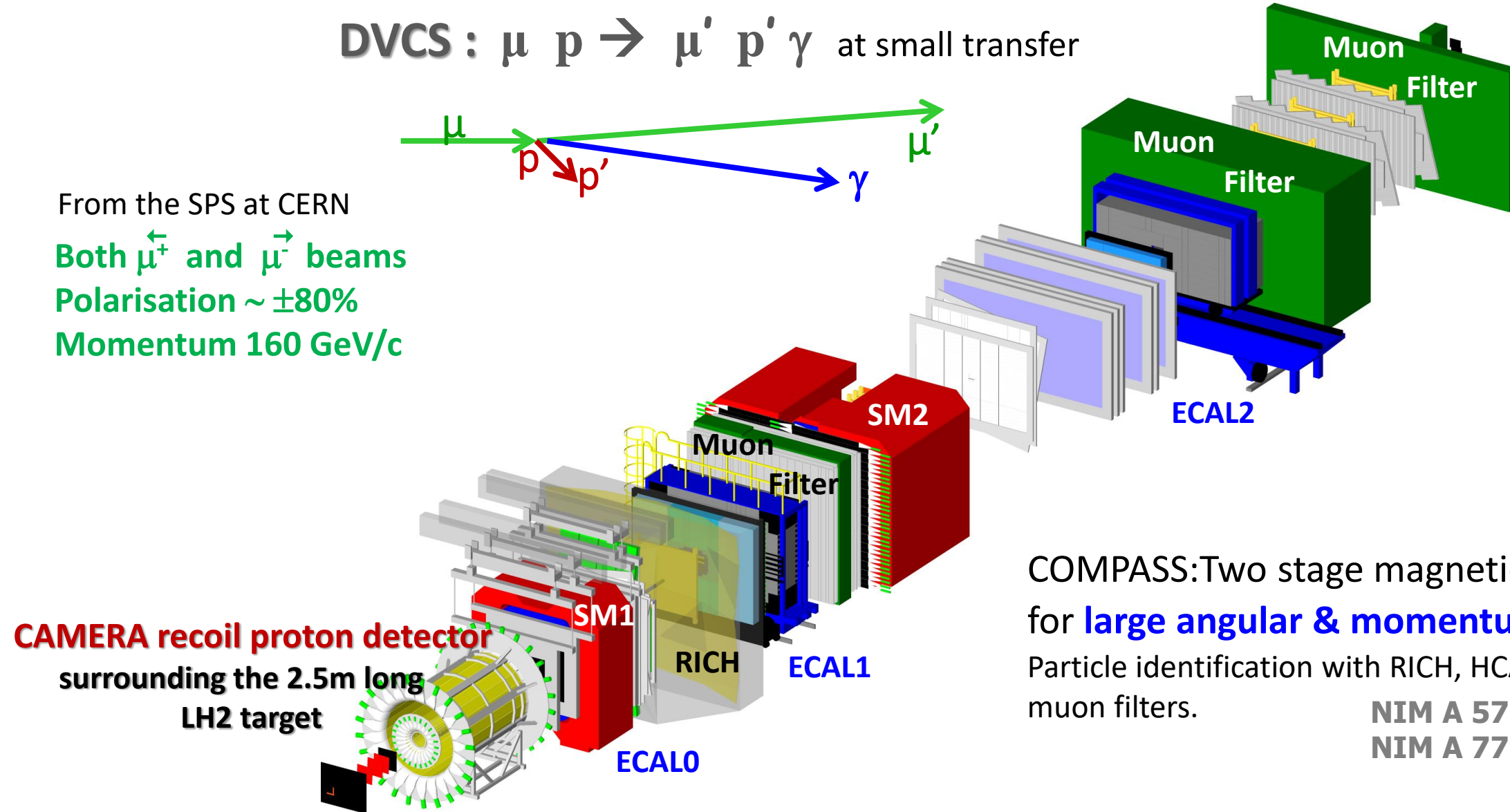


From the SPS at CERN

Both  $\mu^+$  and  $\mu^-$  beams

Polarisation  $\sim \pm 80\%$

Momentum 160 GeV/c



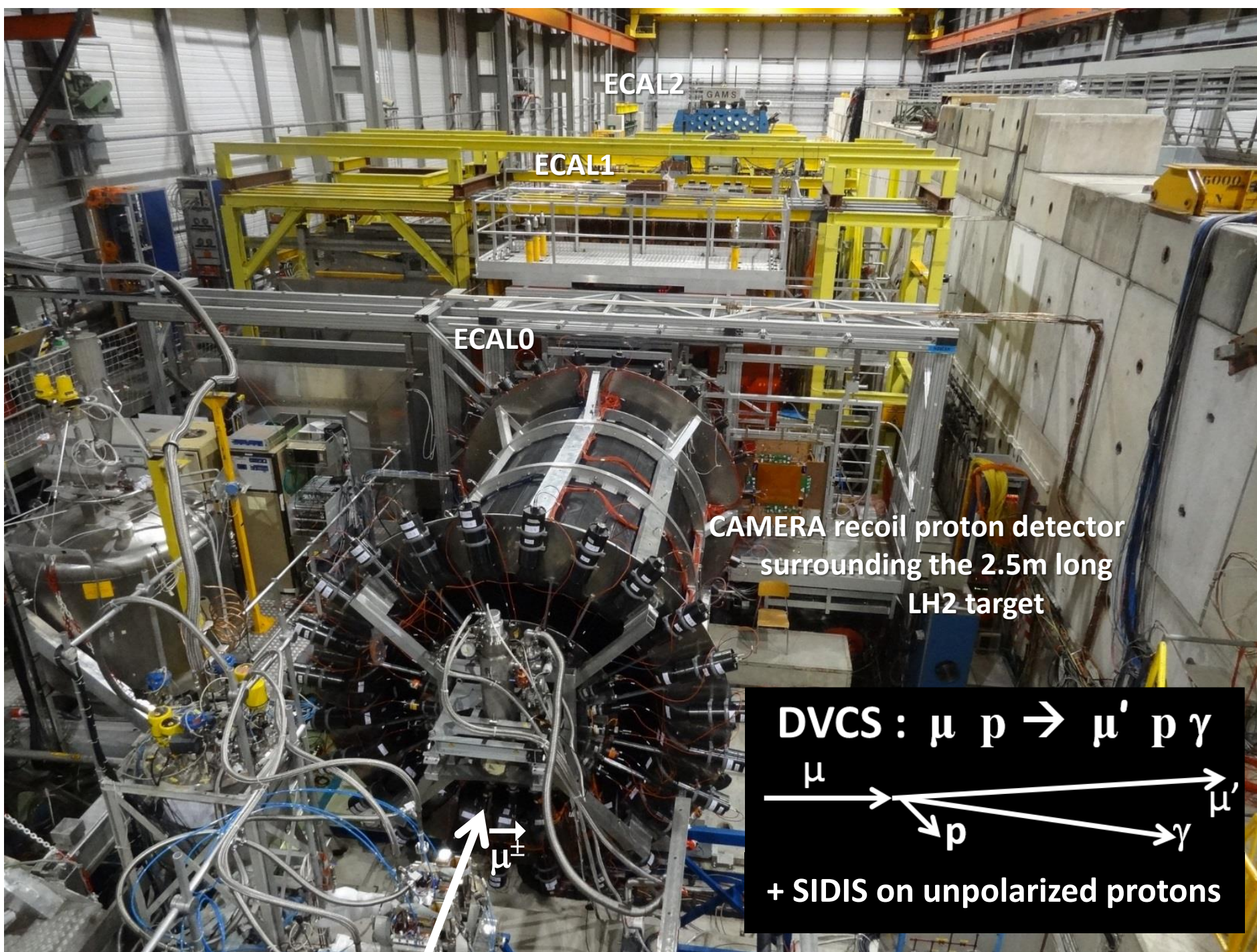
COMPASS: Two stage magnetic spectrometer for **large angular & momentum acceptance**

Particle identification with RICH, HCALs, ECALs and muon filters.

NIM A 577 (2007) 455

NIM A 779 (2015) 69



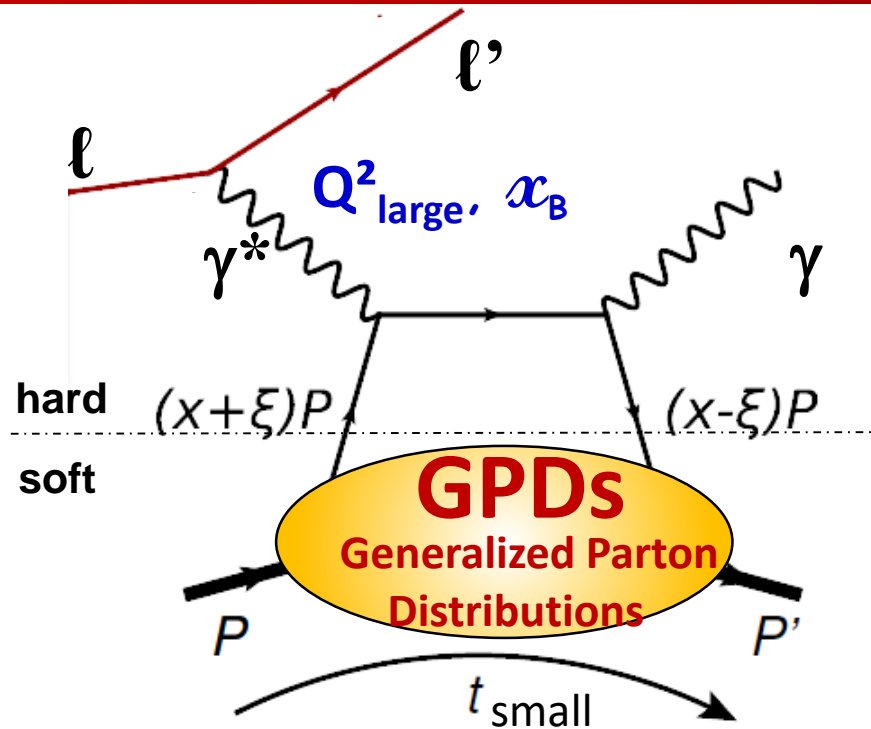


**2012:**  
1 month pilot run

**2016-17:**  
2 x 6 month  
data taking



# Deeply virtual Compton scattering (DVCS)



$x$  inside the loop: average longitudinal momentum fraction

$\xi \approx x_B/2$ : transferred longitudinal momentum fraction

$t$ : total proton momentum transfer squared related to  $b_\perp$  via Fourier transform (when  $\xi = 0$ )

The amplitude DVCS at LT & LO in  $\alpha_s$  (GPD  $\mathcal{H}$ ):

$$\mathcal{H} = \int_{-1}^{+1} dx \frac{\mathcal{H}(x, \xi, t)}{x - \xi + i\epsilon} = \mathcal{P} \int_{-1}^{+1} dx \frac{\mathcal{H}(x, \xi, t)}{x - \xi} - i \pi \mathcal{H}(x = \pm \xi, \xi, t)$$

Real part                      Imaginary part

In an experiment we measure Compton Form Factor  $\mathcal{H}$

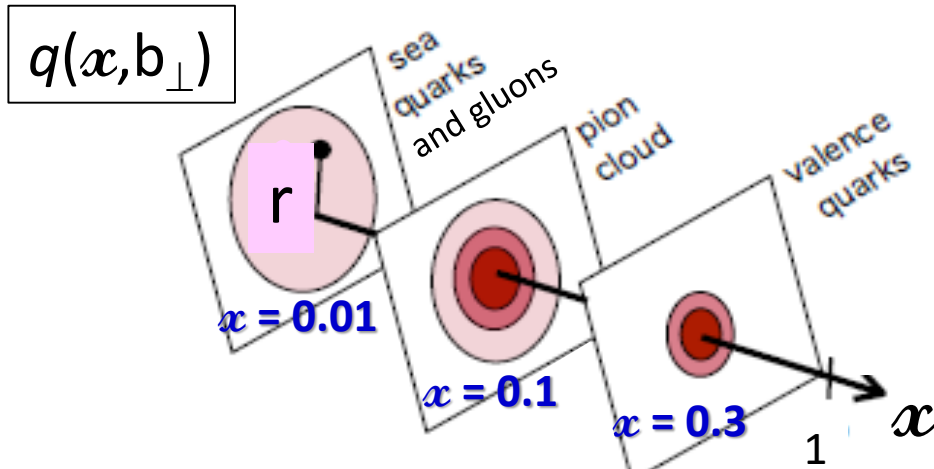
$$\text{Re}\mathcal{H}(\xi, t) = \pi^{-1} \int dx \frac{\text{Im}\mathcal{H}(x, t)}{x - \xi} + \Delta(t)$$

# Deeply virtual Compton scattering (DVCS)

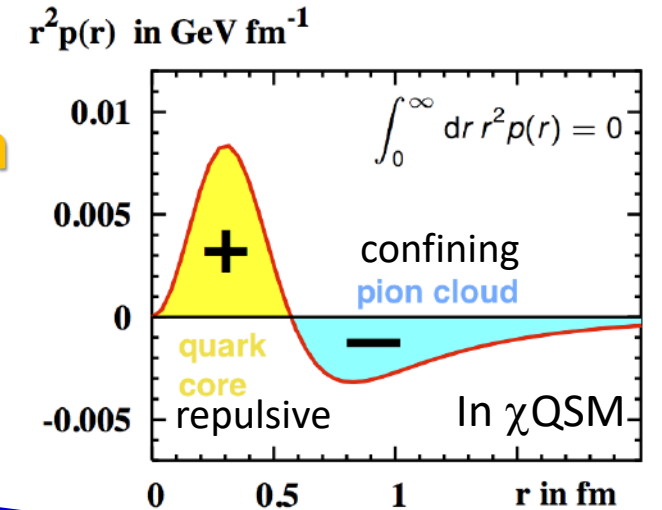
M. Burkardt, PRD66(2002)

M. Polyakov, P. Schweitzer, Int.J.Mod.Phys. A33 (2018)

## Mapping in the transverse plane



## Pressure Distribution



FT of  $H(x, \xi=0, t)$

The amplitude DVCS at LT & LO in  $\alpha_s$  (GPD  $\mathbf{H}$ ):

$$\mathcal{H} = \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi + i\epsilon} = \mathcal{P} \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi} - i \pi H(x = \pm \xi, \xi, t)$$

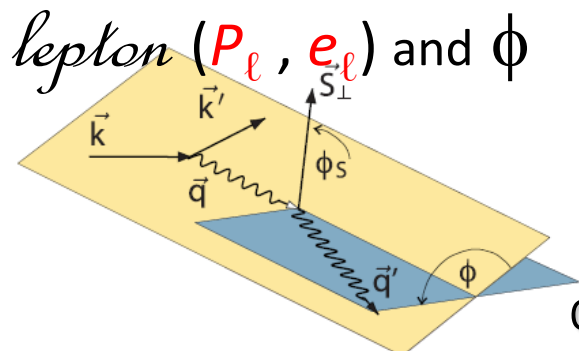
In an experiment we measure Compton Form Factor  $\mathcal{H}$

$$\text{Re}\mathcal{H}(\xi, t) = \pi^{-1} \int dx \frac{\text{Im}\mathcal{H}(x, t)}{x - \xi} + \Delta(t)$$

$d_1(t)$   
D-term

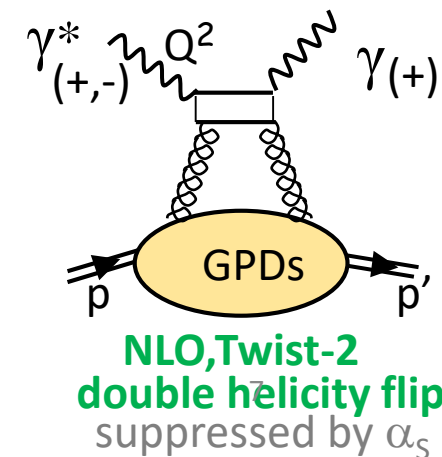
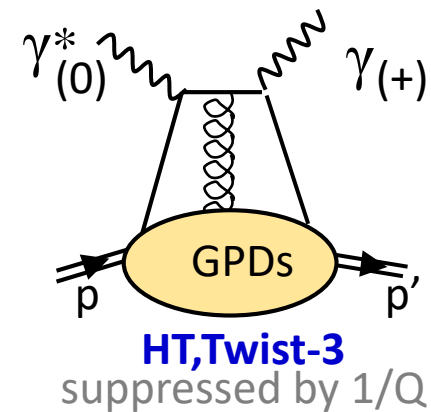
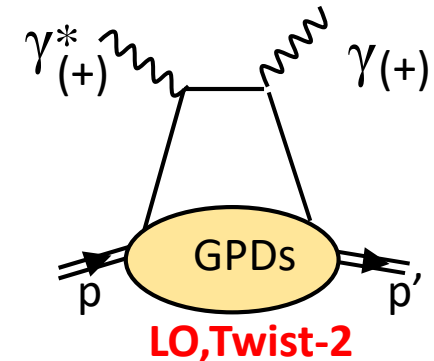


# Exclusive single photon production (BH + DVCS)



$$d\sigma = |T^{BH}|^2 + |T^{DVCS}|^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left( d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) - (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$



With unpolarized target:

Belitsky, Müller, Kirner, NPB629 (2002)

$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$$

$$d\sigma_{unpol}^{DVCS} \propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$$

$$d\sigma_{pol}^{DVCS} \propto s_1^{DVCS} \sin \phi$$

$$\text{Re } I \propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$$

$$\text{Im } I \propto s_1^I \sin \phi + s_2^I \sin 2\phi$$

# Exclusive single photon production (BH + DVCS)

With both  $\mu^+$  and  $\mu^-$  beams we can build:

## ① beam charge-spin sum

$$\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-} =$$

$$\begin{aligned} d\sigma^{BH} &\propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi \\ + d\sigma_{unpol}^{DVCS} &\propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi \\ + \text{Im } I &\propto s_1^I \sin \phi + s_2^I \sin 2\phi \end{aligned}$$

## ② difference

$$\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-} =$$

$$\begin{aligned} d\sigma_{pol}^{DVCS} &\propto s_1^{DVCS} \sin \phi \\ + \text{Re } I &\propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi \end{aligned}$$

$$\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-} \rightarrow s_1^I \propto \text{Im } \mathcal{F}$$

$$\text{and } c_0^{DVCS} \propto (\text{Im } \mathcal{H})^2$$

$$\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-} \rightarrow c_1^I \propto \text{Re } \mathcal{F}$$

$$\mathcal{F} = F_1 \mathcal{H} + \xi (F_1 + F_2) \tilde{\mathcal{H}} - t/4m^2 F_2 \mathcal{E}$$

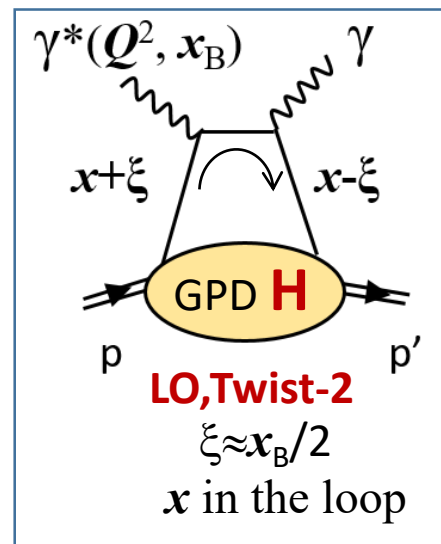
for proton target



at small  $x_B$   
COMPASS domain

$$F_1 \mathcal{H}$$

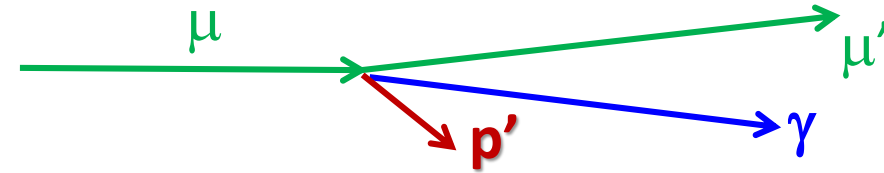
Compton Form Factor  
linked to the GPD  $\mathcal{H}$





# COMPASS 2016 data Selection of exclusive single photon production

Comparison between the observables given by the spectro or by CAMERA

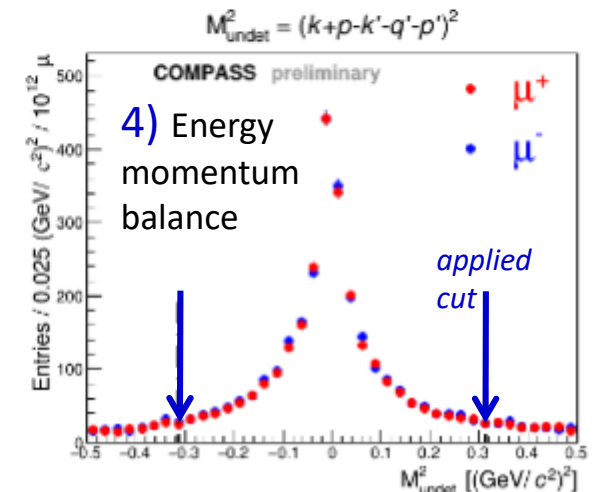
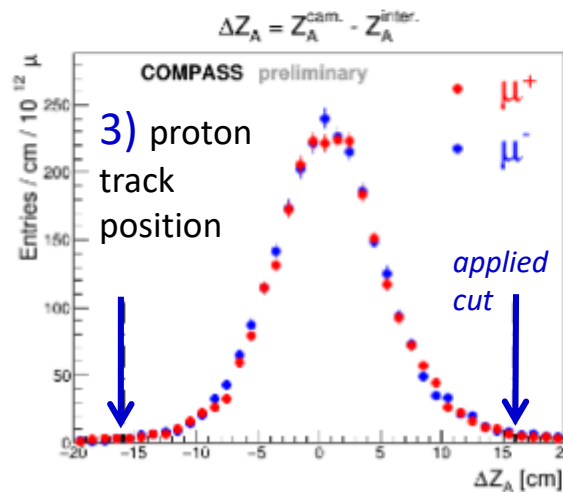
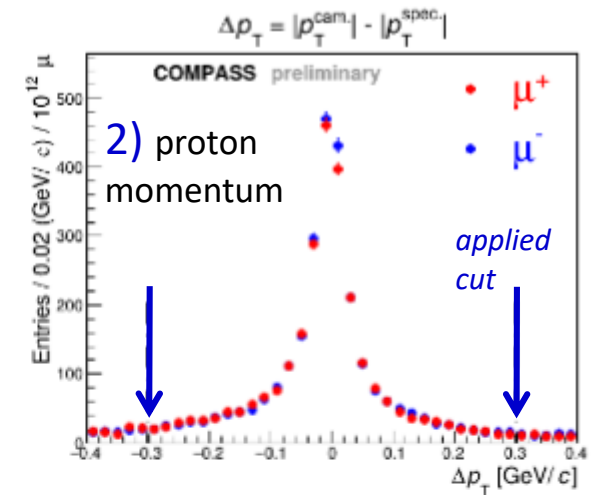
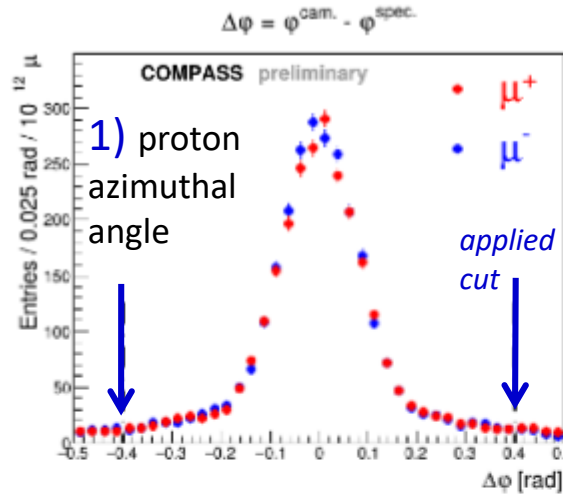


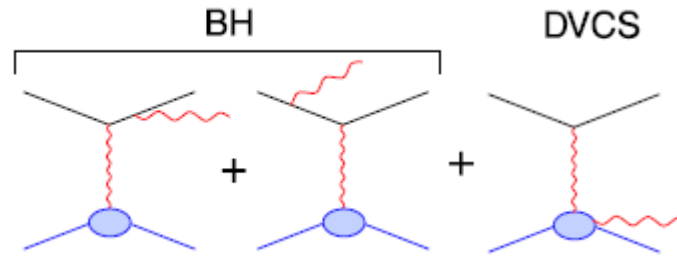
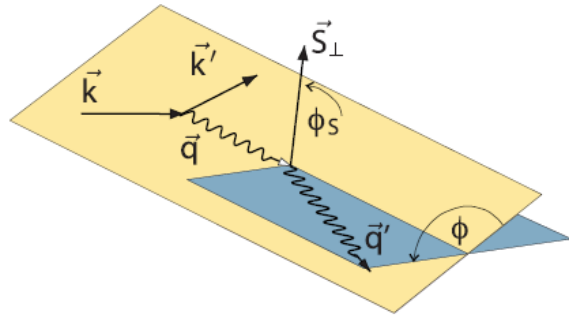
## DVCS: $\mu p \rightarrow \mu' p \gamma$

- 1)  $\Delta\varphi = \varphi^{\text{cam}} - \varphi^{\text{spec}}$
- 2)  $\Delta p_T = p_T^{\text{cam}} - p_T^{\text{spec}}$
- 3)  $\Delta Z_A = Z_A^{\text{cam}} - Z_A^{\text{inter}}$  and vertex
- 4)  $M_{X=0}^2 = (p_{\mu_{\text{in}}} + p_{p_{\text{in}}} - p_{\mu_{\text{out}}} - p_{p_{\text{out}}} - p_{\gamma})^2$

Good agreement between  $\vec{\mu}^+$  and  $\vec{\mu}^-$  yields important achievement for:

- ①  $\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$  **Easier, done first**  
**Mapping in Transverse plane**
- ②  $\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$  **Challenging, but promising**  
**Related to EMT and pressure**





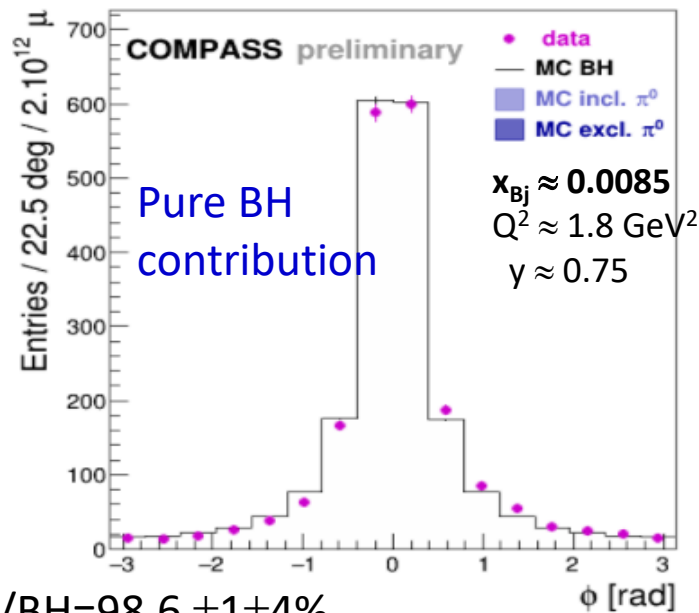
$$\Sigma = d\sigma(\mu^+) + d\sigma(\mu^-)$$

$$d\sigma \propto |T^{\text{BH}}|^2 + \text{Interference Term} + |T^{\text{DVCS}}|^2$$

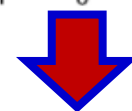
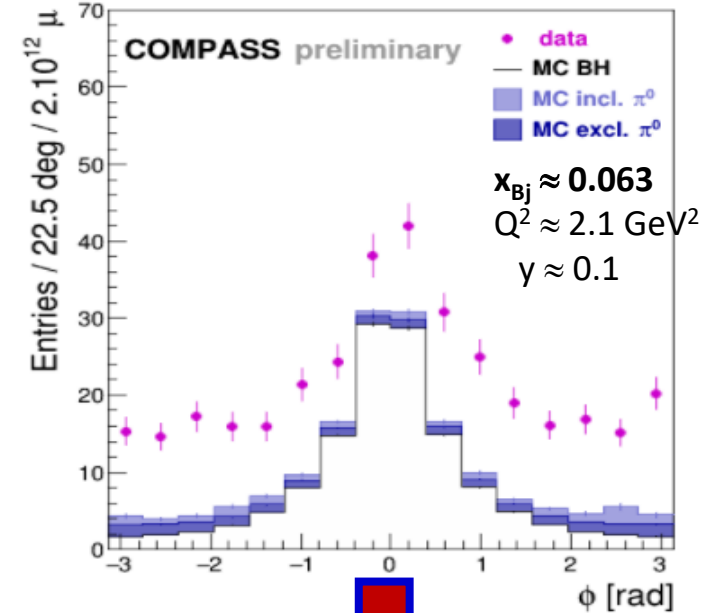
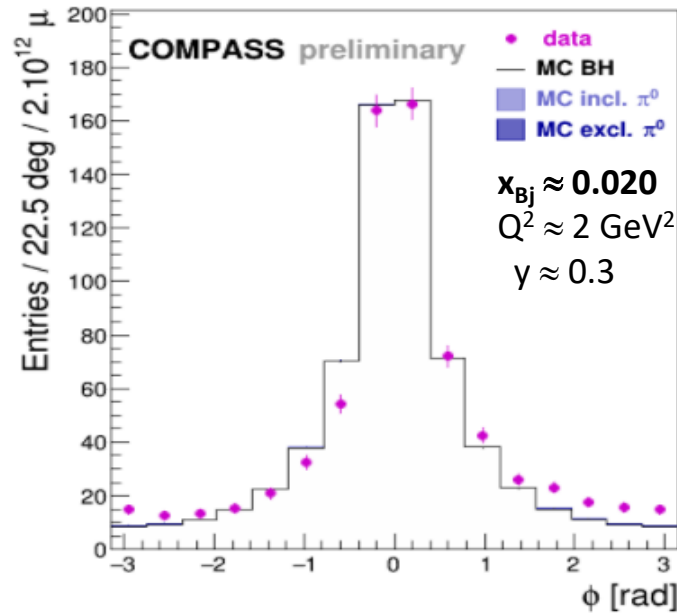
$80 < v \text{ [GeV]} < 144$

$32 < v \text{ [GeV]} < 80$

$10 < v \text{ [GeV]} < 32$



Data/BH =  $98.6 \pm 1 \pm 4\%$



**DVCS** above the **BH** contrib.

MC: BH contribution evaluated for the integrated luminosity  
 $\pi^0$  background contribution from SIDIS (LEPTO) + exclusive production (HEPGEN)



At COMPASS using polarized positive and negative muon beams:

$$\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-} = 2[d\sigma^{BH} + d\sigma_{unpol}^{DVCS} + \text{Im } I]$$

$$= 2[d\sigma^{BH} + c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi + s_1^I \sin \phi + s_2^I \sin 2\phi]$$

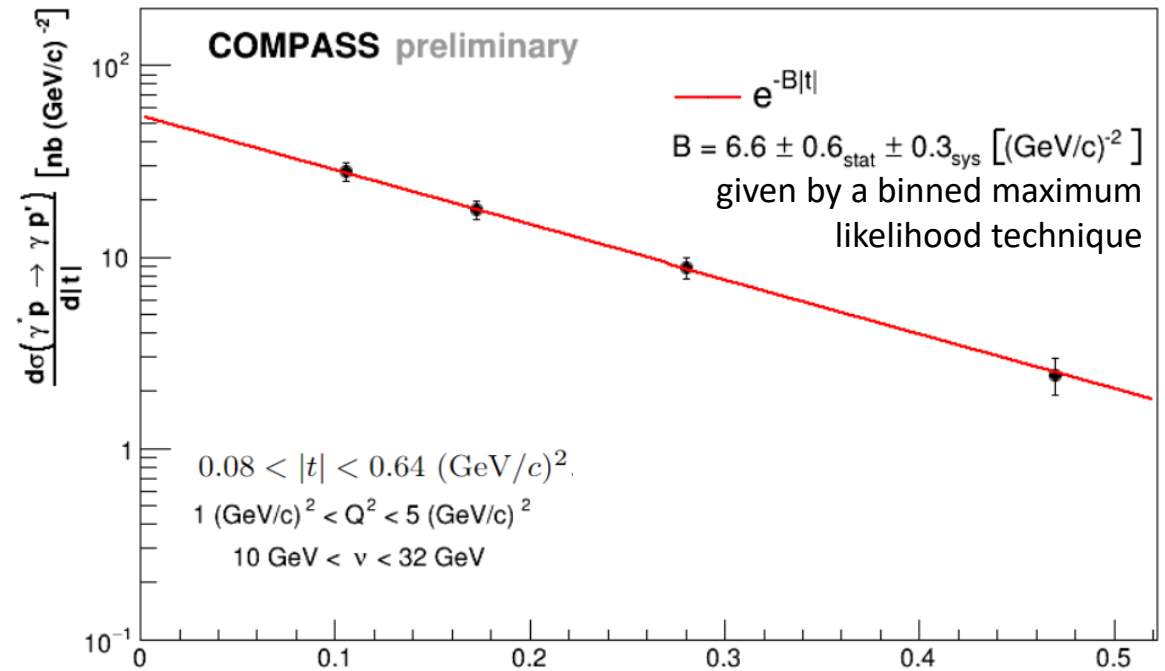
calculable  
can be subtracted

All the other terms are cancelled in the integration over  $\phi$

$$\frac{d^3\sigma_T^{\mu p}}{dQ^2 d\nu dt} = \int_{-\pi}^{\pi} d\phi (d\sigma - d\sigma^{BH}) \propto c_0^{DVCS}$$

$$\frac{d\sigma^{\gamma^* p}}{dt} = \frac{1}{\Gamma(Q^2, \nu, E_\mu)} \frac{d^3\sigma_T^{\mu p}}{dQ^2 d\nu dt}$$

Flux for transverse  
virtual photons



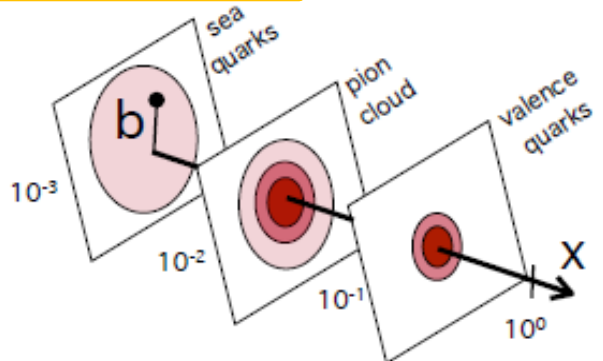
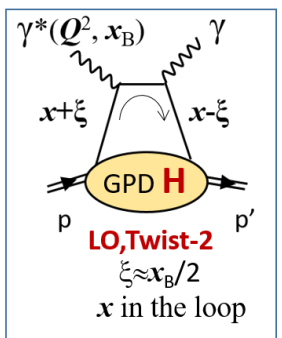
# COMPASS 12-16 Transverse extension of partons in the sea quark range

$$d\sigma^{DVCS}/dt = e^{-B|t|} = c_0^{DVCS} \propto (\text{Im}\mathcal{H})^2$$

$$\text{Im}\mathcal{H} = H(x=\xi, \xi, t)$$

$$x = \xi \approx x_B/2 \text{ close to } 0$$

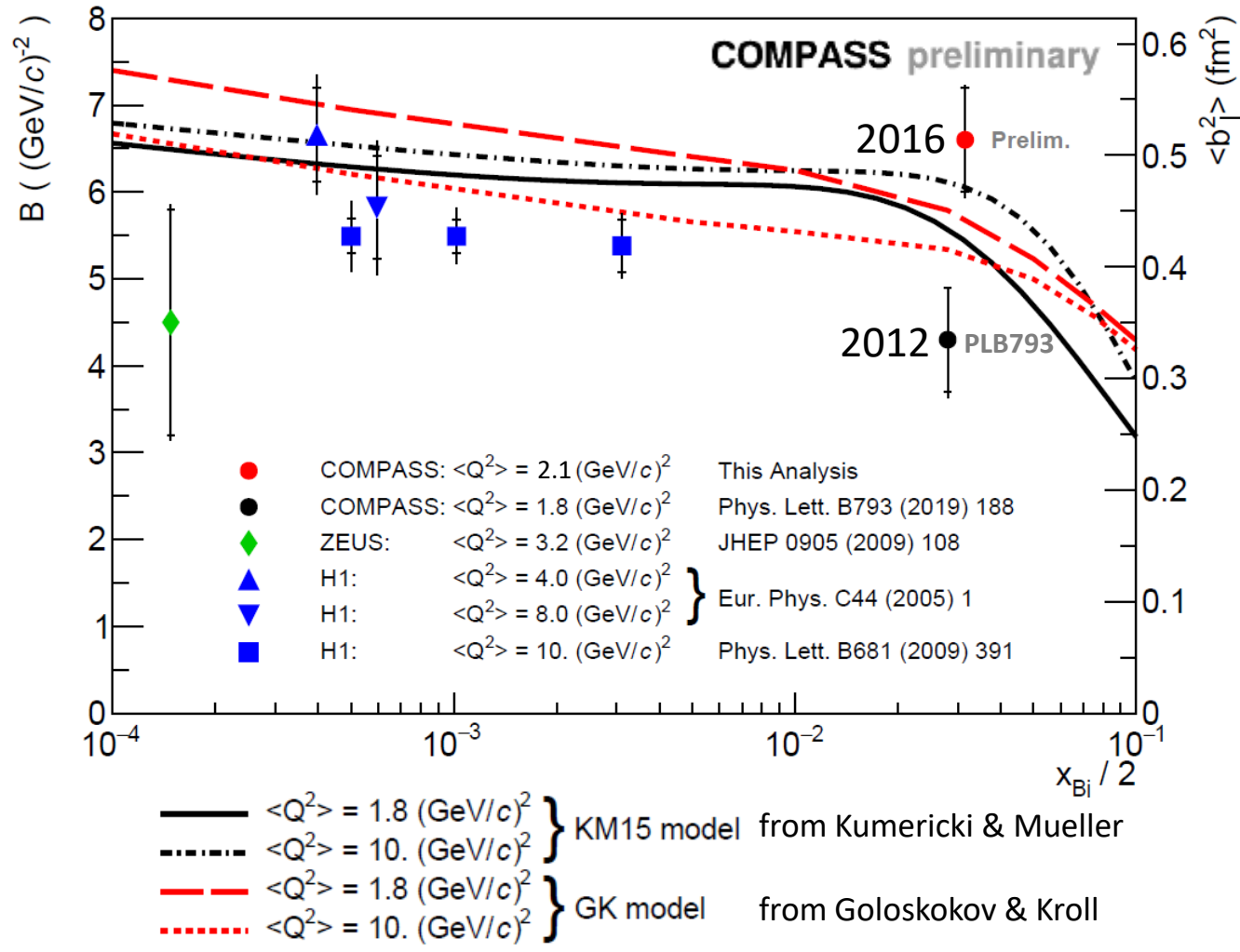
$$\langle b_{\perp}^2(x) \rangle \approx 2B(\xi)$$



## Improvements in 2016 analysis compared to 2012

- same intensity with  $\mu^+$  and  $\mu^-$  beam in 2016
- more advanced analysis with 2016 data, still ongoing
- $\pi^0$  contamination with different thresholds
- better MC description of the evolution in  $v$
- binning with 3 variables ( $t, Q^2, v$ ) or 4 variables ( $t, \phi, Q^2, v$ )
- different binning in  $t$

2012 statistics = Ref  
 2016 analysed statistics =  $2.3 \times \text{Ref}$   
 2016+2017 expected statistics =  $10 \times \text{Ref}$

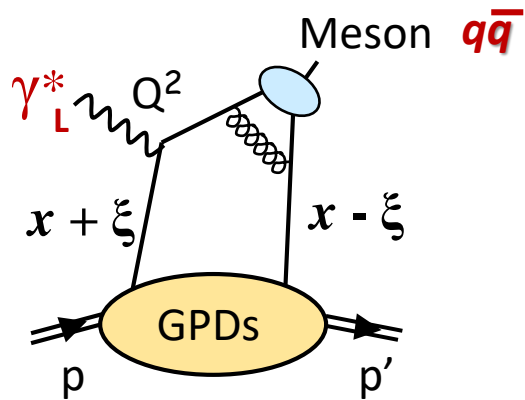


# GPDs and Hard Exclusive Meson Production

Factorisation proven only for  $\sigma_L$

The meson wave function is an additional non-perturbative term

Quark contribution



## For Pseudo-Scalar Meson, as $\pi^0$

chiral-even GPDs: helicity of parton unchanged

$$\tilde{H}^q(x, \xi, t) \quad \tilde{E}^q(x, \xi, t)$$

+ chiral-odd or transversity GPDs: helicity of parton changed

$$H_T^q(x, \xi, t) \quad (\text{as the transversity TMD})$$

related in the forward limit to transversity and the tensor charge

$$\bar{E}_T^q = 2 \tilde{H}_T^q + E_T^q \quad (\text{as the Boer-Mulders TMD})$$

related to the distortion of the polarized quark distribution in the transverse plane for the unpolarized proton and to its transverse anomalous magnetic moment

$\sigma_T$  should be asymptotically suppressed by  $1/Q^2$  but large contribution observed

GK model:  $k_T$  of  $q$  and  $\bar{q}$  and Sudakov suppression factor are considered

Chiral-odd GPDs with a twist-3 meson wave function

$\mu^\pm p \rightarrow \mu^\pm \pi^0 p$   
 $\mu^\pm$  beams with opposite polarization

$$\frac{1}{2} \left( \frac{d^2\sigma^+}{dt d\phi_\pi} + \frac{d^2\sigma^-}{dt d\phi_\pi} \right) = \frac{1}{2\pi} \left[ \left( \epsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} \right) + \epsilon \cos 2\phi_\pi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_\pi \frac{d\sigma_{LT}}{dt} \right]$$

COMPASS  
 $\langle x_B \rangle = 0.13$   
 $\epsilon$  close to 1

$$\frac{d\sigma_L}{dt} \propto |\langle \tilde{H} \rangle|^2 - \frac{t'}{4m^2} |\langle \tilde{E} \rangle|^2$$

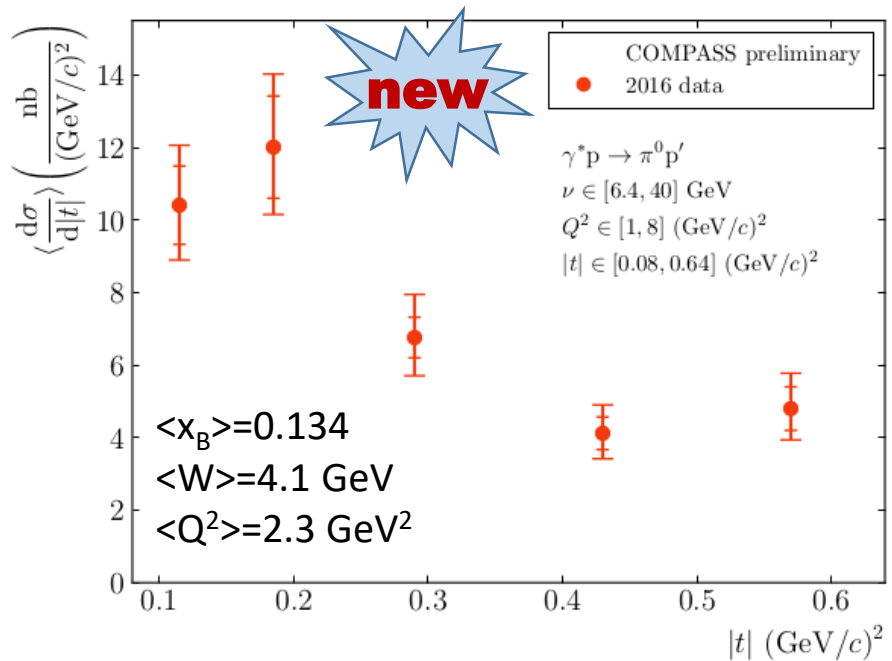
$$\frac{d\sigma_T}{dt} \propto |\langle H_T \rangle|^2 - \frac{t'}{8m^2} |\langle \bar{E}_T \rangle|^2$$

$$\frac{\sigma_{TT}}{dt} \propto \frac{t'}{16m^2} |\langle \bar{E}_T \rangle|^2$$

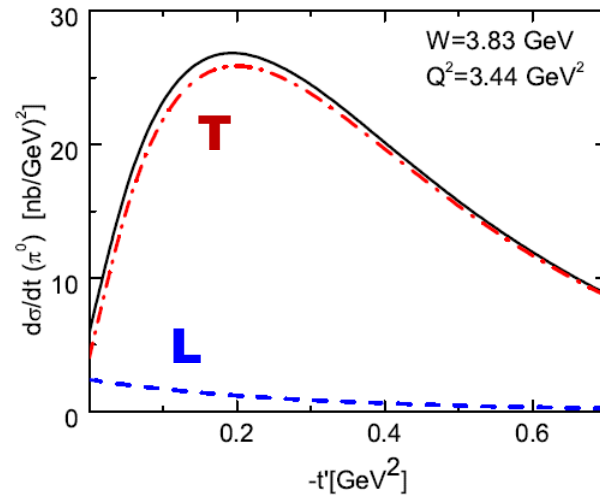
$$\frac{\sigma_{LT}}{dt} \propto \frac{\sqrt{-t'}}{2m} \text{Re} \left[ \langle H_T \rangle^* \langle \tilde{E} \rangle + \langle \bar{E}_T \rangle^* \langle \tilde{H} \rangle \right]$$

$F\pi^0 = 2/3 F^u + 1/3 F^d$  ( $\tilde{H}^u \tilde{H}^d$ ) ( $\tilde{E}^u \tilde{E}^d$ ) ( $H_T^u H_T^d$ ) of opposite sign

( $\bar{E}_T^u \bar{E}_T^d$ ) of same sign  $\rightarrow$  **clearly enhanced contribution**



S. Goloskokov, P. Kroll, EPJC47 (2011)



Typical dip of the cross section as a function of  $-t' = -(t-t_0) \approx |t|$   
 $|t_0| \approx 10^{-2} \text{ GeV}^2$



$\mu^\pm p \rightarrow \mu^\pm \pi^0 p$   
 $\mu^\pm$  beams with opposite polarization

$$\frac{1}{2} \left( \frac{d^2\sigma^+}{dt d\phi_\pi} + \frac{d^2\sigma^-}{dt d\phi_\pi} \right) = \frac{1}{2\pi} \left[ \left( \epsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} \right) + \epsilon \cos 2\phi_\pi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_\pi \frac{d\sigma_{LT}}{dt} \right]$$

COMPASS  
 $\langle x_B \rangle = 0.13$   
 $\epsilon$  close to 1

$$\frac{d\sigma_L}{dt} \propto |\langle \tilde{H} \rangle|^2 - \frac{t'}{4m^2} |\langle \tilde{E} \rangle|^2$$

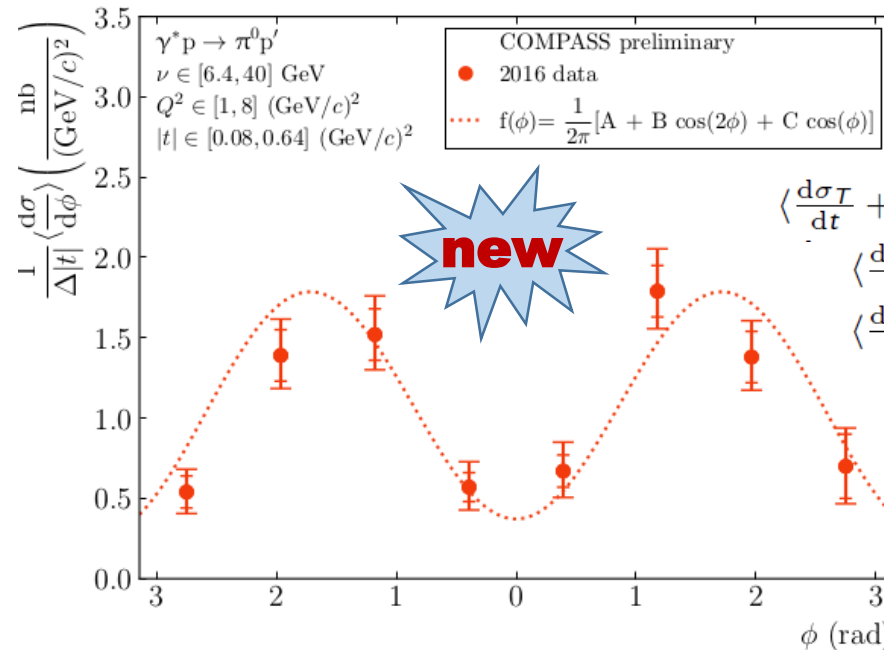
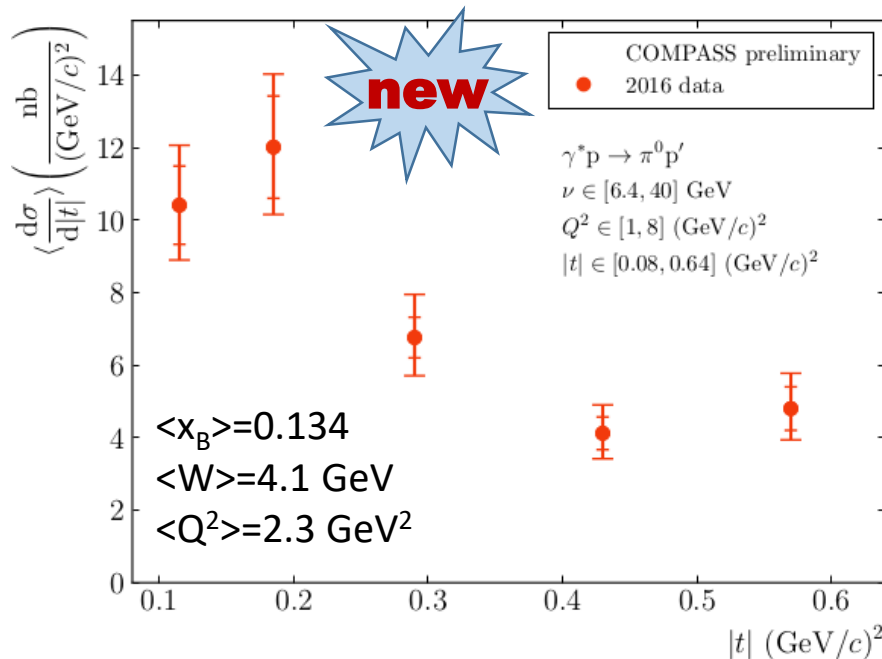
$$\frac{d\sigma_T}{dt} \propto |\langle H_T \rangle|^2 - \frac{t'}{8m^2} |\langle \bar{E}_T \rangle|^2$$

$$\frac{\sigma_{TT}}{dt} \propto \frac{t'}{16m^2} |\langle \bar{E}_T \rangle|^2$$

$$\frac{\sigma_{LT}}{dt} \propto \frac{\sqrt{-t'}}{2m} \text{Re} \left[ \langle H_T \rangle^* \langle \tilde{E} \rangle + \langle \bar{E}_T \rangle^* \langle \tilde{H} \rangle \right]$$

$F\pi^0 = 2/3 F^u + 1/3 F^d$  ( $\tilde{H}^u \tilde{H}^d$ ) ( $\tilde{E}^u \tilde{E}^d$ ) ( $H_T^u H_T^d$ ) of opposite sign

$(\bar{E}_T^u \bar{E}_T^d)$  of same sign  $\rightarrow$  **clearly enhanced contribution**



$$\left\langle \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right\rangle = (6.6 \pm 0.3_{\text{stat}} + 0.9_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{TT}}{dt} \right\rangle = (-4.6 \pm 0.5_{\text{stat}} + 0.3_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{LT}}{dt} \right\rangle = (0.2 \pm 0.2_{\text{stat}} + 0.2_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

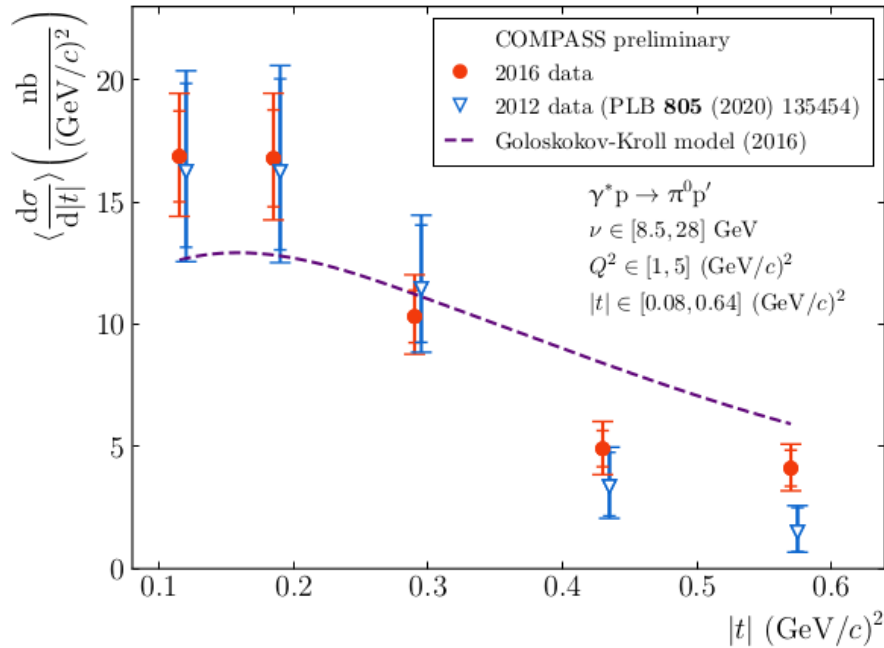
$$\langle \epsilon \rangle = 0.997$$

$|\sigma_{TT}|$  almost as large as  $\sigma_T + \epsilon \sigma_L$   
 $\rightarrow$  **impact of  $\bar{E}_T$**

Cross section for  $\nu \in [6.4, 40]$  GeV and  $Q^2 \in [1, 8]$  GeV<sup>2</sup>  $\langle x_B \rangle = 0.13$

$\sigma_{LT}$  rather small 15/30

2016 kinematic domain: Cross section for  $\nu \in [6.4, 40]$  GeV and  $Q^2 \in [1, 8]$  GeV<sup>2</sup>  $\langle x_B \rangle = 0.13$   
 2012 kinematic domain for comparison:  $\nu \in [8.5, 28]$  GeV and  $Q^2 \in [1, 5]$  GeV<sup>2</sup>  $\langle x_B \rangle = 0.10$



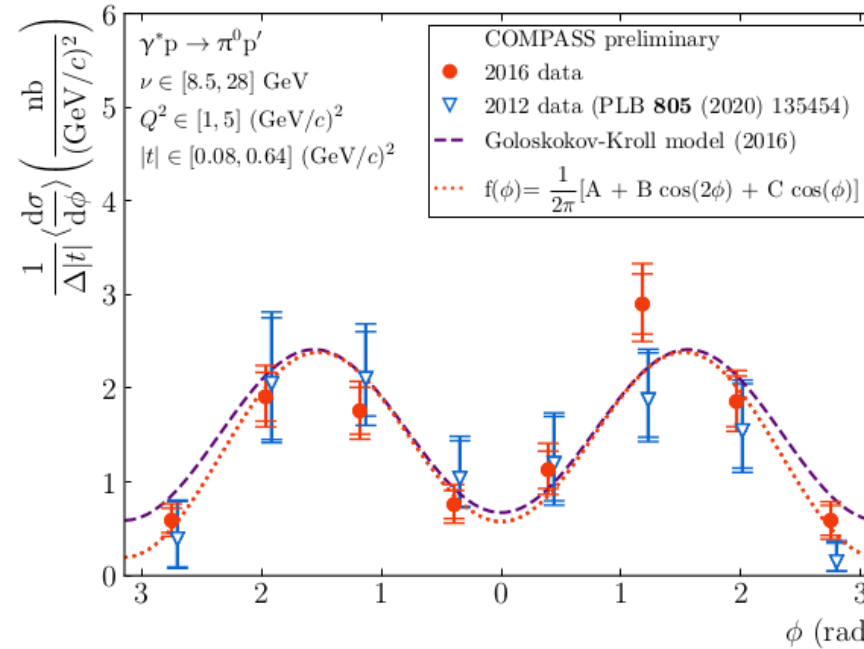
2012 data:

$$\left\langle \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right\rangle = (8.1 \pm 0.9_{\text{stat}} + 1.1 |_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{TT}}{dt} \right\rangle = (-6.0 \pm 1.3_{\text{stat}} + 0.7 |_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{LT}}{dt} \right\rangle = (1.4 \pm 0.5_{\text{stat}} + 0.3 |_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\langle \epsilon \rangle = 0.996$$



2016 data:

$$\left\langle \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right\rangle = (8.7 \pm 0.5_{\text{stat}} + 1.0 |_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{TT}}{dt} \right\rangle = (-6.3 \pm 0.8_{\text{stat}} + 0.4 |_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{LT}}{dt} \right\rangle = (0.6 \pm 0.3_{\text{stat}} + 0.3 |_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\langle \epsilon \rangle = 0.996$$



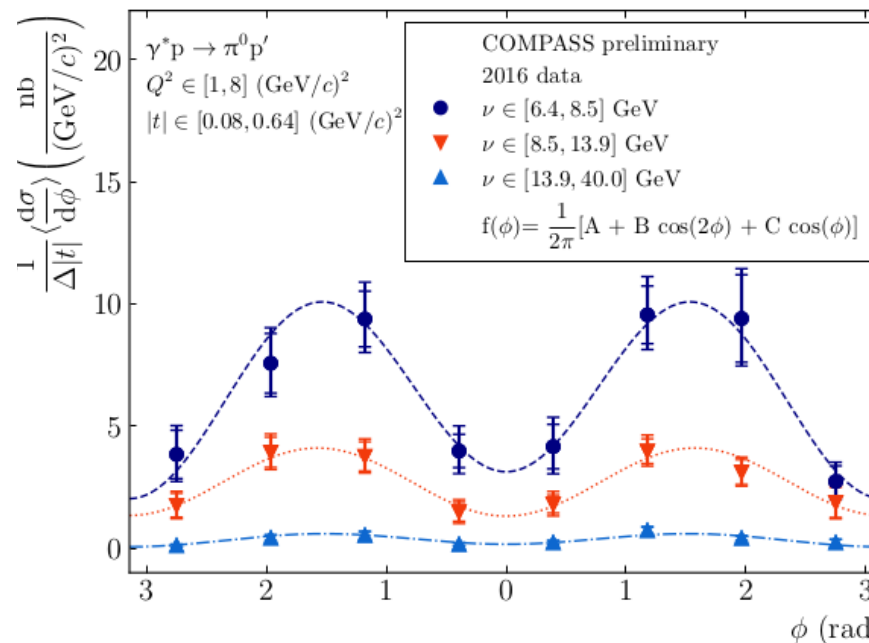
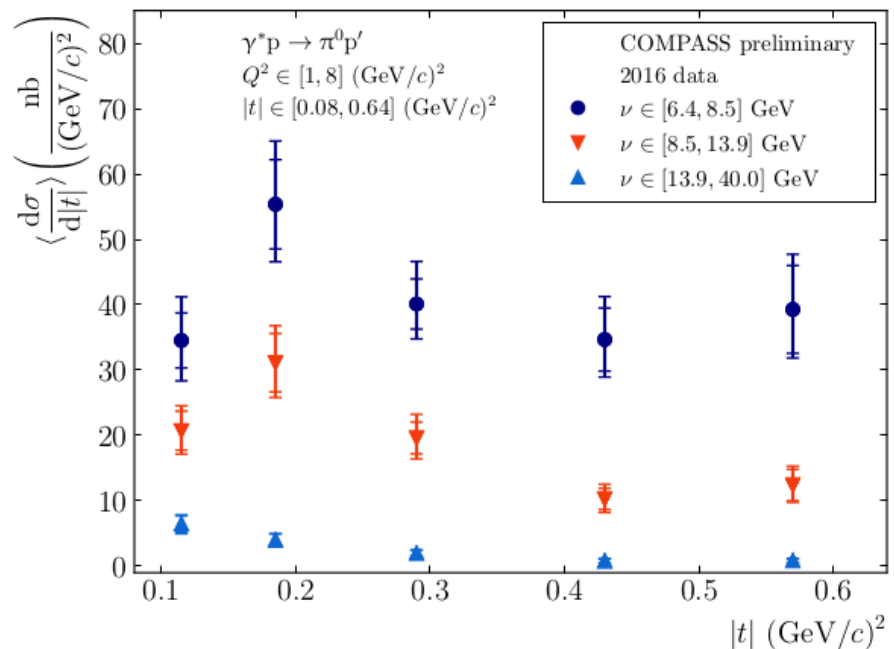
**$|\sigma_{TT}|$  almost as large as**

$$\sigma_T + \epsilon \sigma_L$$

**→ impact of  $\bar{E}_T$**

$\sigma_{LT}$  rather small

Evolution of the cross section with  $\nu$ :  $\sigma \searrow$  when  $\nu \nearrow$



→ Extraction of

$$\sigma_T + \epsilon \sigma_L$$

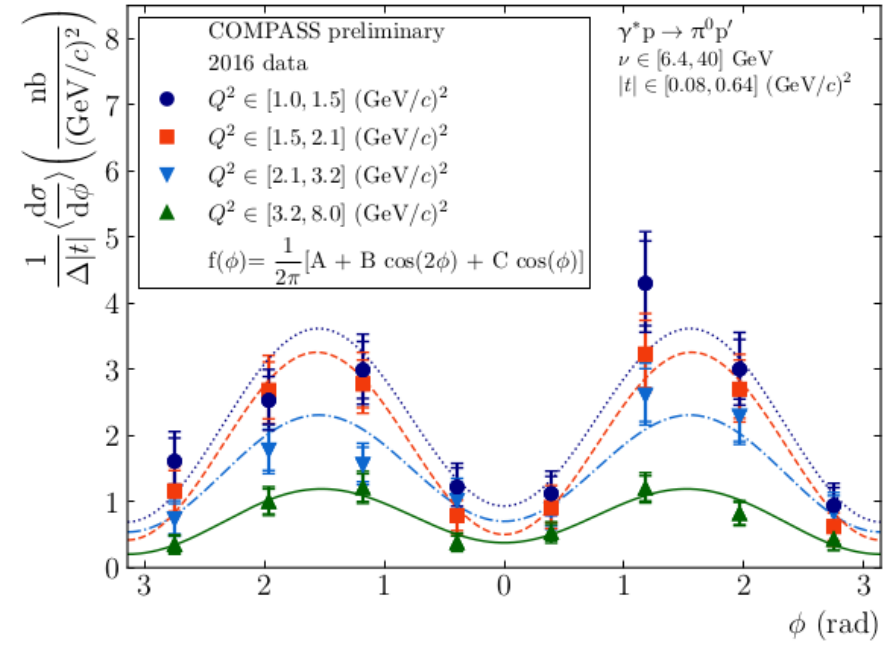
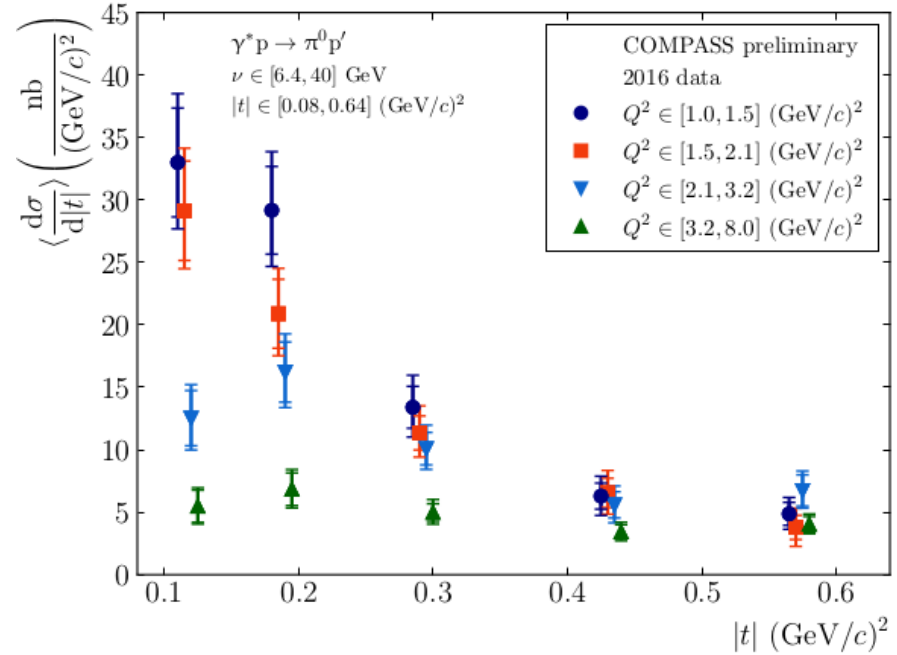
$$\sigma_{TT}$$

$$\sigma_{LT}$$

in 3  $\nu$  bins

	$\langle \nu \rangle$ [GeV]	$\langle Q^2 \rangle$ [GeV <sup>2</sup> /c <sup>2</sup> ]	$\langle x_B \rangle$	$\langle \epsilon \rangle$
$\nu \in [6.4, 8.5]$	7.35	2.15	0.156	0.999
$\nu \in [8.5, 13.9]$	10.32	2.50	0.131	0.998
$\nu \in [13.9, 40.0]$	21.08	2.09	0.057	0.989

Evolution of the cross section with  $Q^2$ :  $\sigma \searrow$  when  $Q^2 \nearrow$

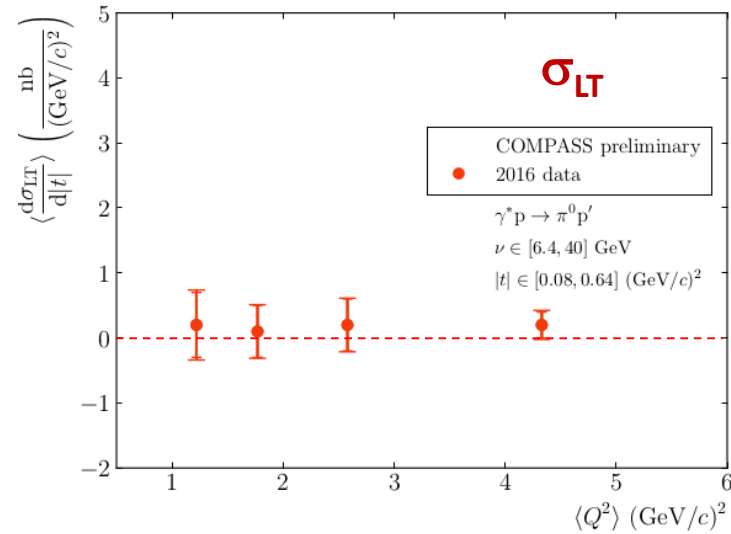
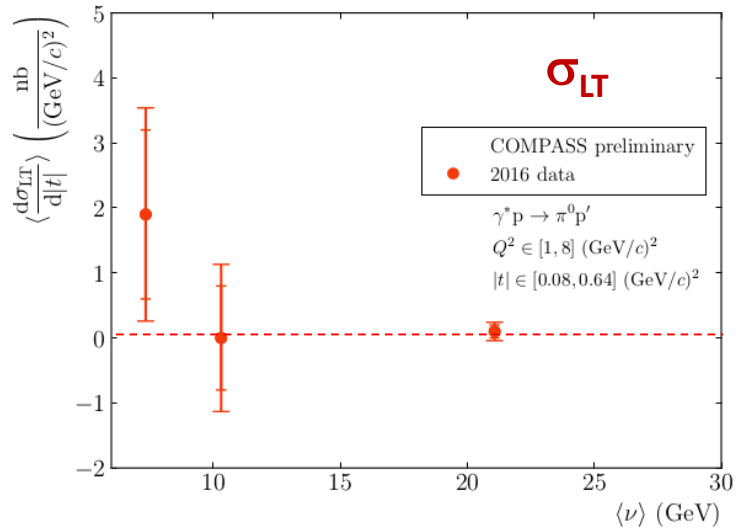


→ Extraction of  
 $\sigma_T + \epsilon \sigma_L$   
 $\sigma_{TT}$   
 $\sigma_{LT}$   
 in 4  $Q^2$  bins

	$\langle Q^2 \rangle \text{ [GeV}^2/c^2]$	$\langle \nu \rangle \text{ [GeV]}$	$\langle x_B \rangle$	$\langle \epsilon \rangle$
$Q^2 \in [1.0, 1.5]$	1.22	10.54	0.072	0.997
$Q^2 \in [1.5, 2.1]$	1.77	9.81	0.109	0.997
$Q^2 \in [2.1, 3.2]$	2.58	9.82	0.157	0.997
$Q^2 \in [3.2, 8.0]$	4.33	10.39	0.247	0.997

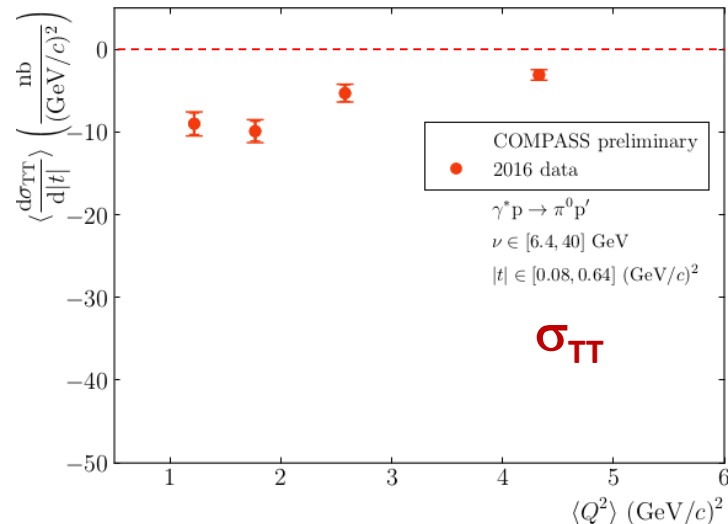
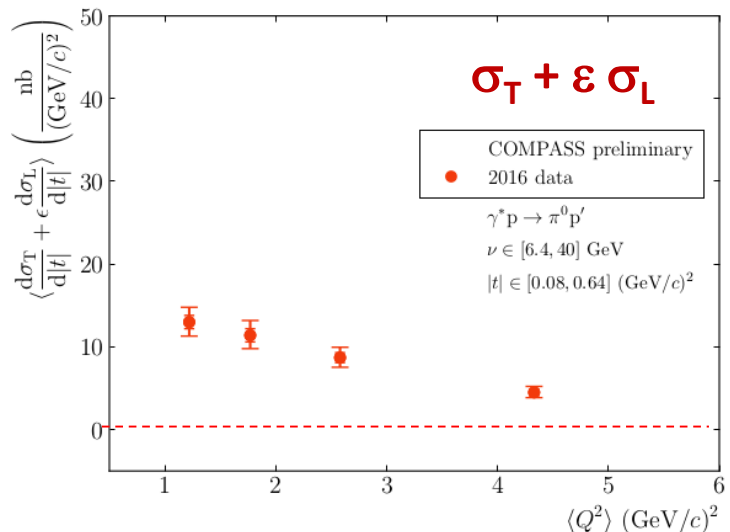
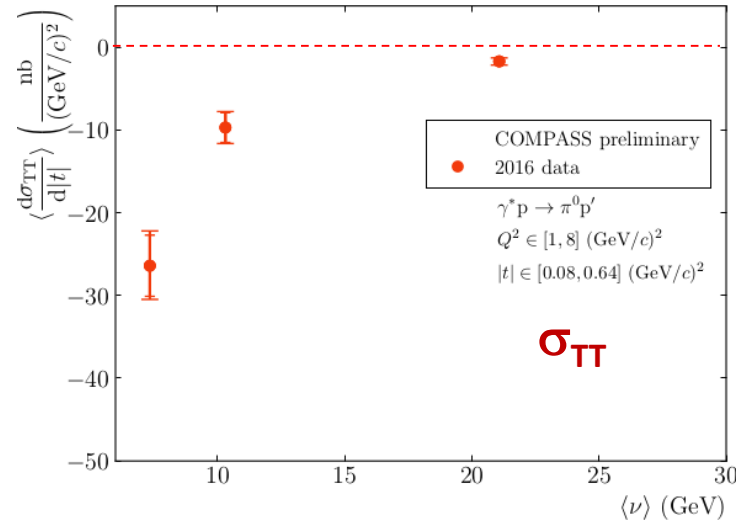
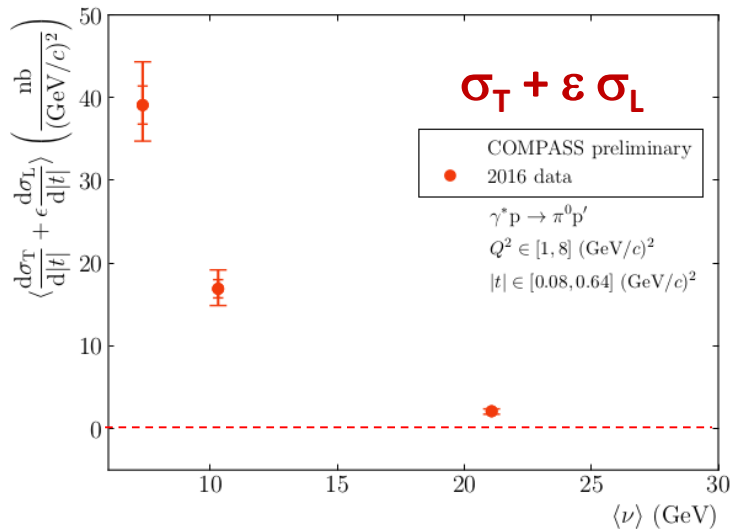


Evolution of the structure functions with  $\nu$  and  $Q^2$



$\sigma_{LT}$  close to 0

## Evolution of the structure functions with $\nu$ and $Q^2$



Both  $\sigma_T + \epsilon \sigma_L$  and  $\sigma_{TT}$  large evolution with  $\nu$   
 small evolution with  $Q^2$

Impact of these data for modeling  $\bar{E}_T$  (and other GPDs) contributions at twist-3 and NLO

Recent work on twist-3 contribution  
 G. Duplančić, P. Kroll and  
 K. Passek-Kumerički, PRD109 (2024)

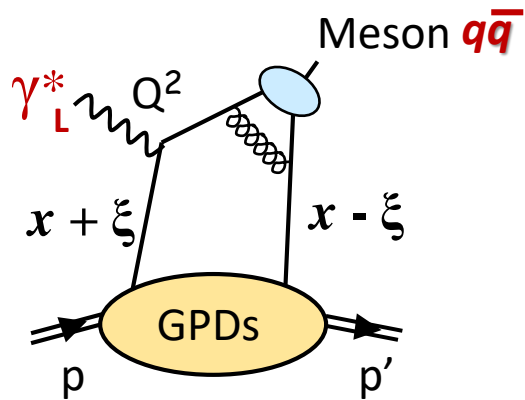
Also S. Golosgokov et al.  
 S. Liuti et al.

# GPDs and Hard Exclusive Meson Production

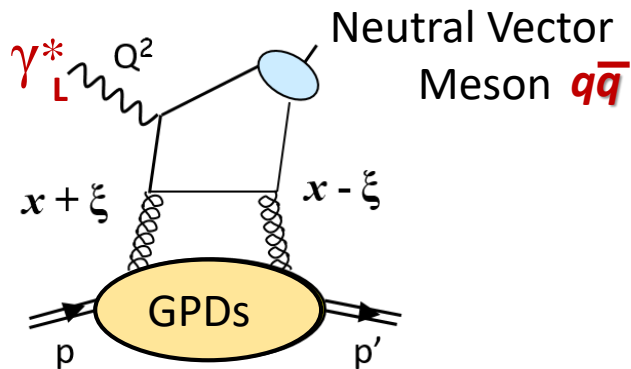
Factorisation proven only for  $\sigma_L$

The meson wave function is an additional non-perturbative term

Quark contribution



Gluon contribution at the same order in  $\alpha_s$



For Vector Meson, as  $\rho, \omega, \phi...$

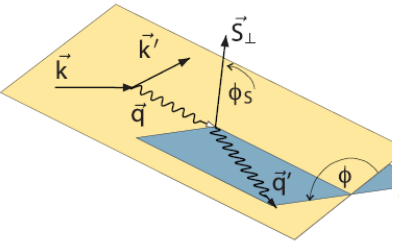
chiral-even GPDs: helicity of parton unchanged

$$H^q(x, \xi, t) \quad E^q(x, \xi, t)$$

+ chiral-odd or transversity GPDs: helicity of parton changed

$$H_T^q(x, \xi, t) \quad (\text{as the transversity TMD})$$

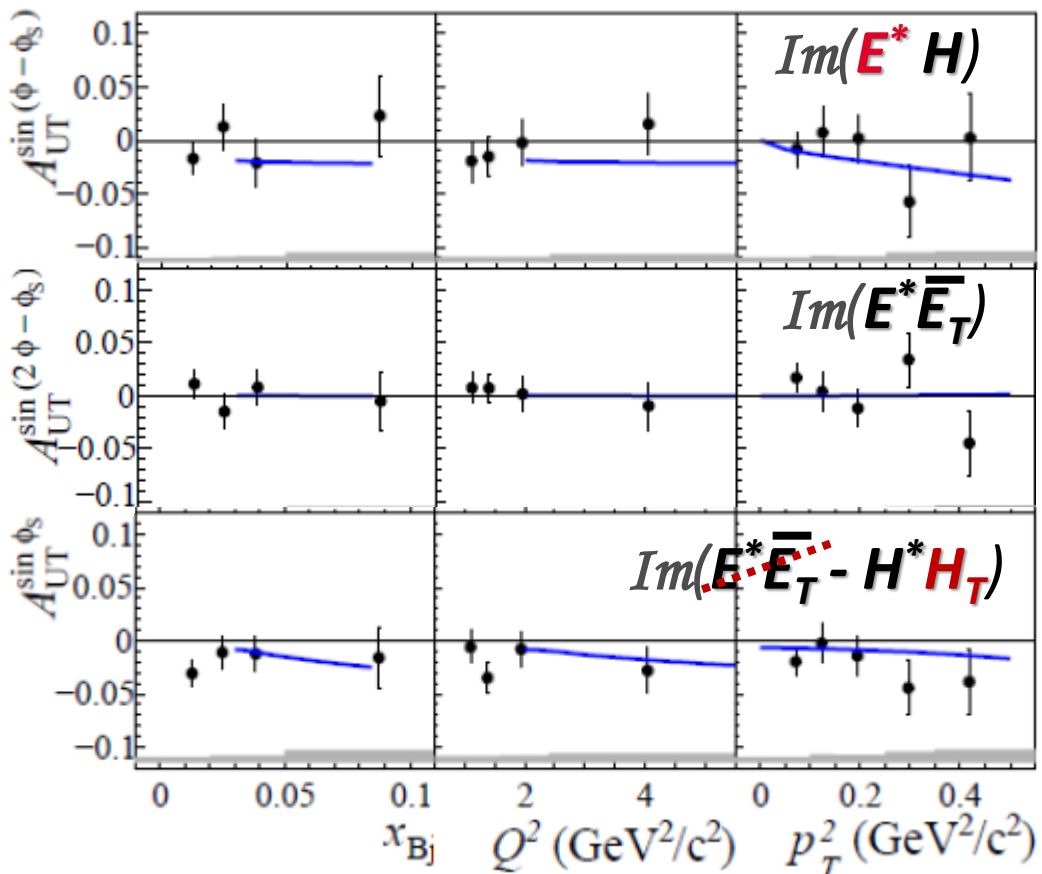
$$\bar{E}_T^q = 2 \tilde{H}_T^q + E_T^q \quad (\text{as the Boer-Mulders TMD})$$



$\rho^0 \rightarrow \pi^+ \pi^-$

$$E_{\rho^0} = \frac{1}{\sqrt{2}} \left( \frac{2}{3} E_u \oplus \frac{1}{3} E_d + \frac{3}{4} \frac{E_g}{x} \right)$$

COMPASS, NPB 865 (2012) 1-20, PLB731 (2014) 19

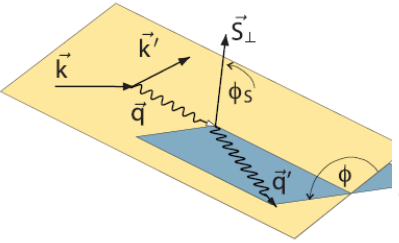


Sensitivity to **E** and **H<sub>T</sub>**

GK model EPJC42,50,53,59,65,74



# HEMP with Transversely Polarized Target without RD



$\rho^0 \rightarrow \pi^+ \pi^-$

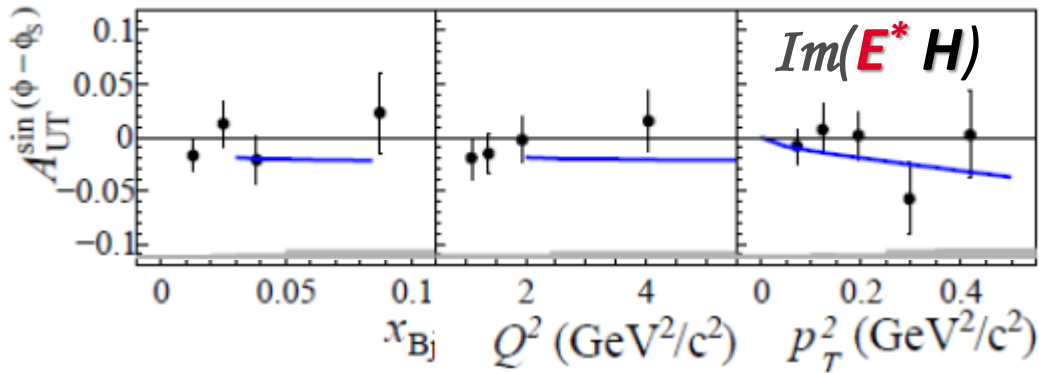
$$E_{\rho^0} = \frac{1}{\sqrt{2}} \left( \frac{2}{3} E^u \oplus \frac{1}{3} E^d + \frac{3}{4} \frac{E^g}{x} \right)$$

$\omega \rightarrow \pi^+ \pi^- \pi^0$

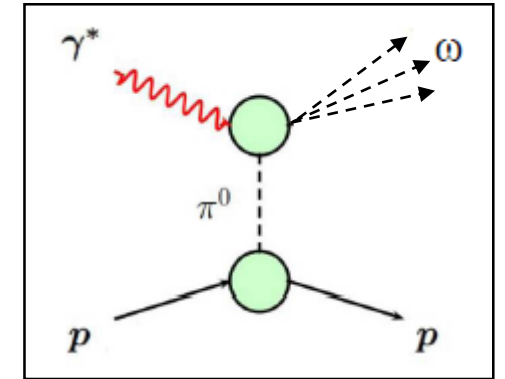
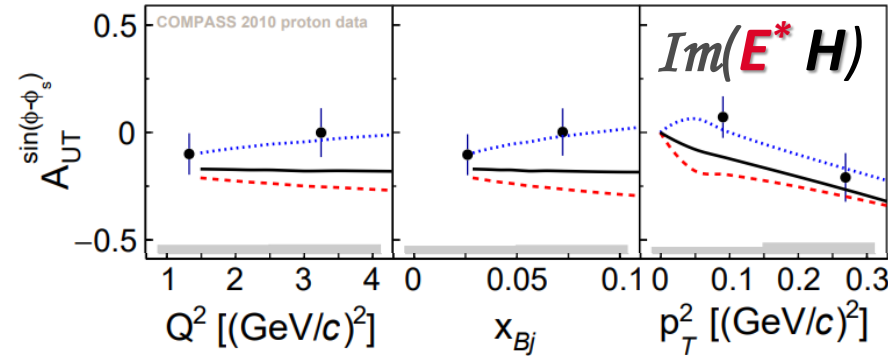
$$E_{\omega} = \frac{1}{\sqrt{2}} \left( \frac{2}{3} E^u \ominus \frac{1}{3} E^d + \frac{1}{4} \frac{E^g}{x} \right)$$

$E^u$  and  $E^d$  of opposite sign

COMPASS, NPB 865 (2012) 1-20, PLB731 (2014) 19



COMPASS, NPB 915 (2017)



$\Gamma(\omega \rightarrow \pi^0 \gamma) = 9 \times \Gamma(\rho^0 \rightarrow \pi^0 \gamma)$

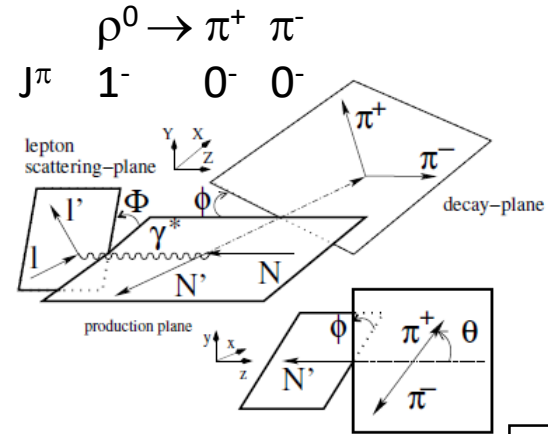
Same for  $\pi\omega$  FF but sign unknown

$\omega$  is more promising (see the larger scale)  
but there is the inherent pion pole contribution

- ▶ positive  $\pi\omega$  form factor
- ▶ no pion pole
- ▶ negative  $\pi\omega$  form factor

GK model EPJC42,50,53,59,65,74

# COMPASS 2012-16 exclusive VM production with Unpolarised Target and SDME



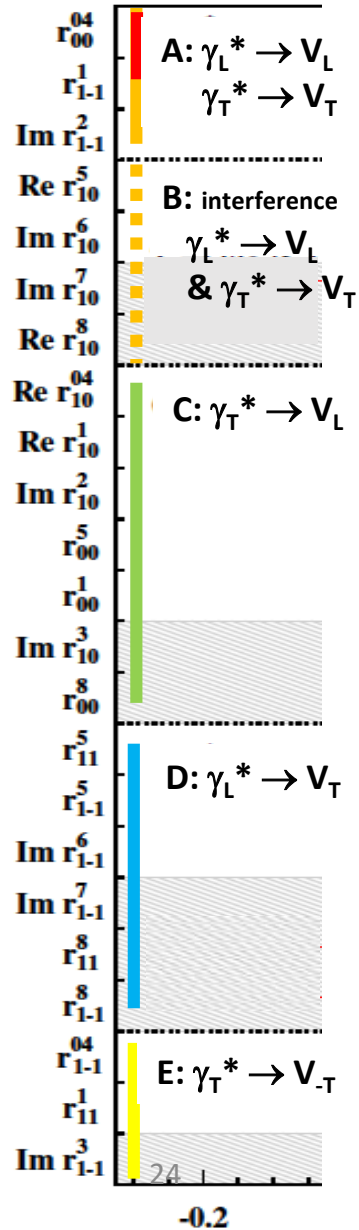
experimental angular distributions:

$$\mathcal{W}^{U+L}(\Phi, \phi, \cos \Theta) = \mathcal{W}^U(\Phi, \phi, \cos \Theta) + P_b \mathcal{W}^L(\Phi, \phi, \cos \Theta)$$

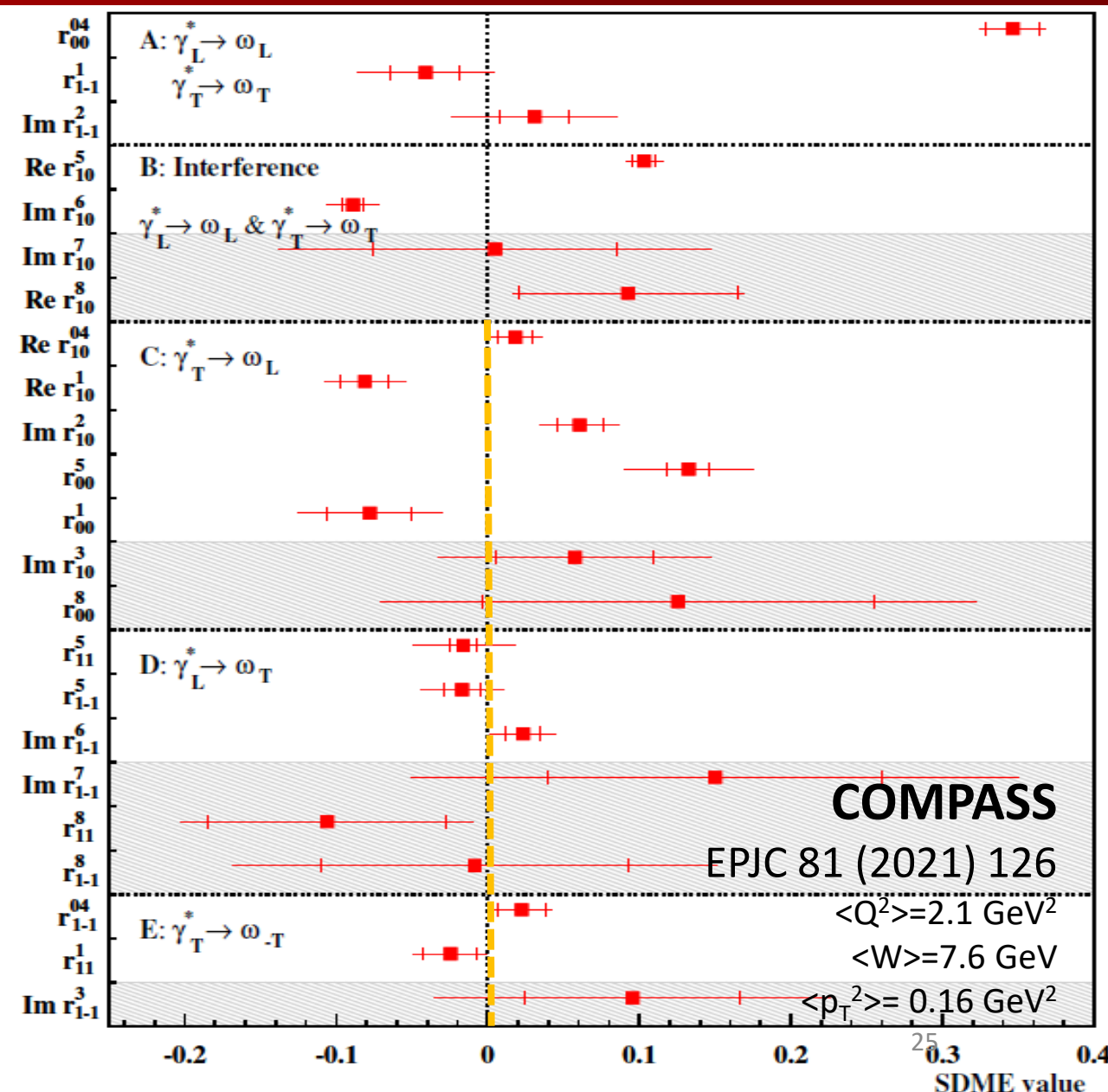
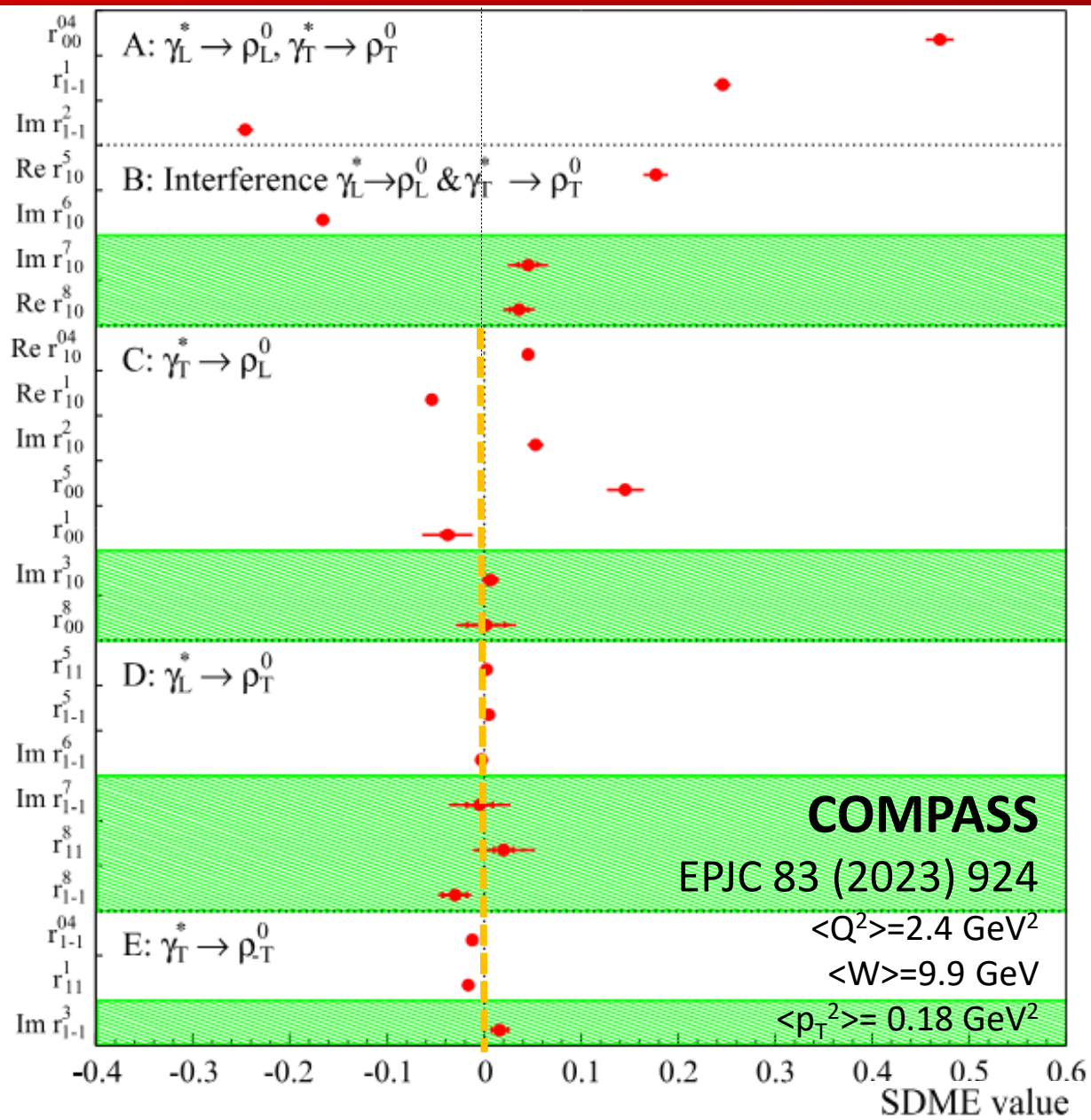
15 'unpolarized' and 8 'polarized' SDMEs

$$\begin{aligned} \mathcal{W}^U(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[ \frac{1}{2}(1 - r_{00}^{04}) + \frac{1}{2}(3r_{00}^{04} - 1) \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^{04}\} \sin 2\Theta \cos \phi - r_{1-1}^{04} \sin^2 \Theta \cos 2\phi \right. \\ & - \epsilon \cos 2\Phi \left( r_{11}^1 \sin^2 \Theta + r_{00}^1 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^1\} \sin 2\Theta \cos \phi - r_{1-1}^1 \sin^2 \Theta \cos 2\phi \right) \\ & \left. - \epsilon \sin 2\Phi \left( \sqrt{2}\text{Im}\{r_{10}^2\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^2\} \sin^2 \Theta \sin 2\phi \right) \right. \\ & + \sqrt{2\epsilon(1+\epsilon)} \cos \Phi \left( r_{11}^5 \sin^2 \Theta + r_{00}^5 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^5\} \sin 2\Theta \cos \phi - r_{1-1}^5 \sin^2 \Theta \cos 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \sin \Phi \left( \sqrt{2}\text{Im}\{r_{10}^6\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^6\} \sin^2 \Theta \sin 2\phi \right) \right], \\ \mathcal{W}^L(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[ \sqrt{1-\epsilon^2} \left( \sqrt{2}\text{Im}\{r_{10}^3\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^3\} \sin^2 \Theta \sin 2\phi \right) \right. \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos \Phi \left( \sqrt{2}\text{Im}\{r_{10}^7\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^7\} \sin^2 \Theta \sin 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \sin \Phi \left( r_{11}^8 \sin^2 \Theta + r_{00}^8 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^8\} \sin 2\Theta \cos \phi - r_{1-1}^8 \sin^2 \Theta \cos 2\phi \right) \right] \end{aligned}$$

$\epsilon$  close to 1,  
small  $\mathcal{W}^L$   
no L/T separation

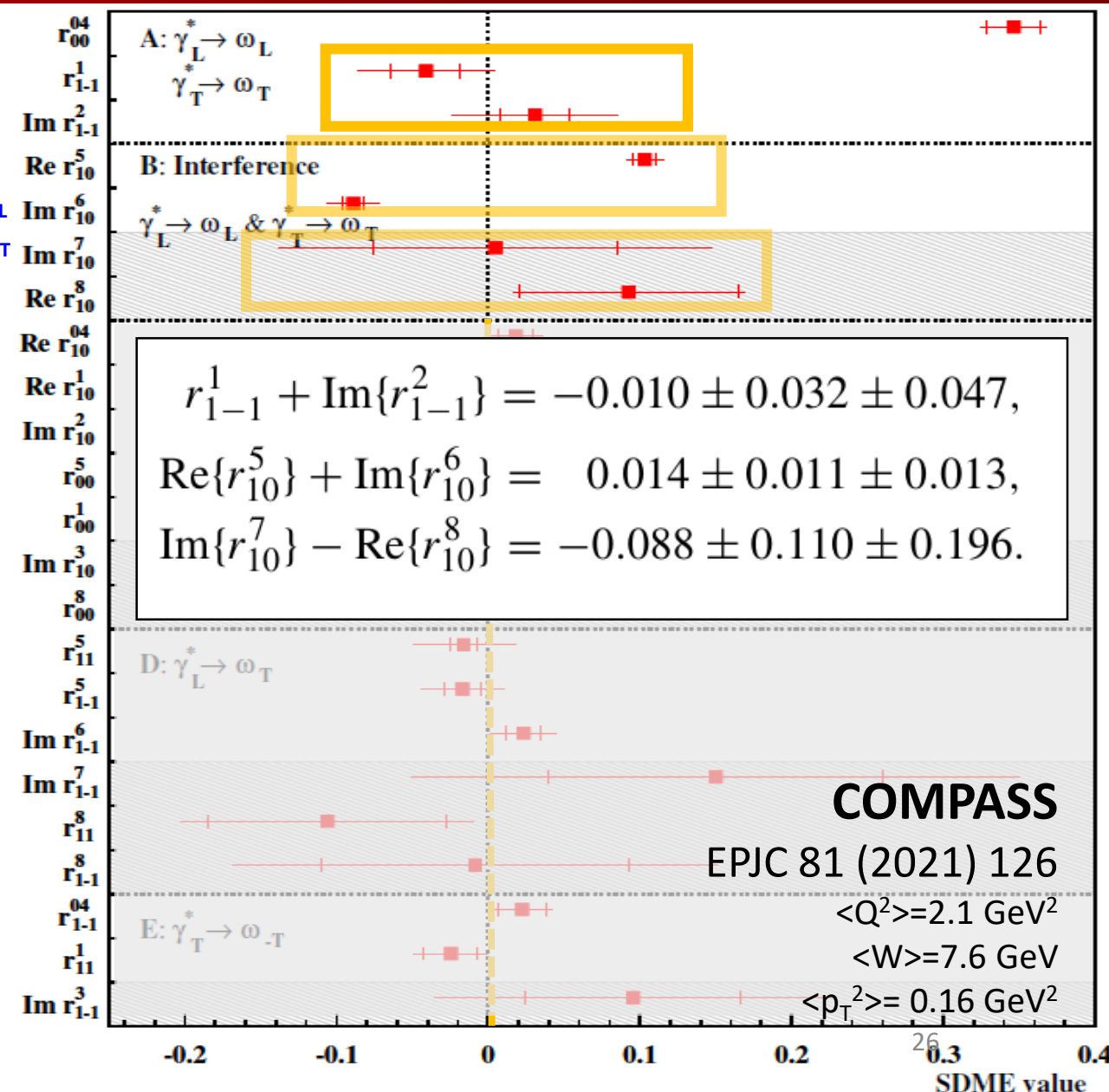
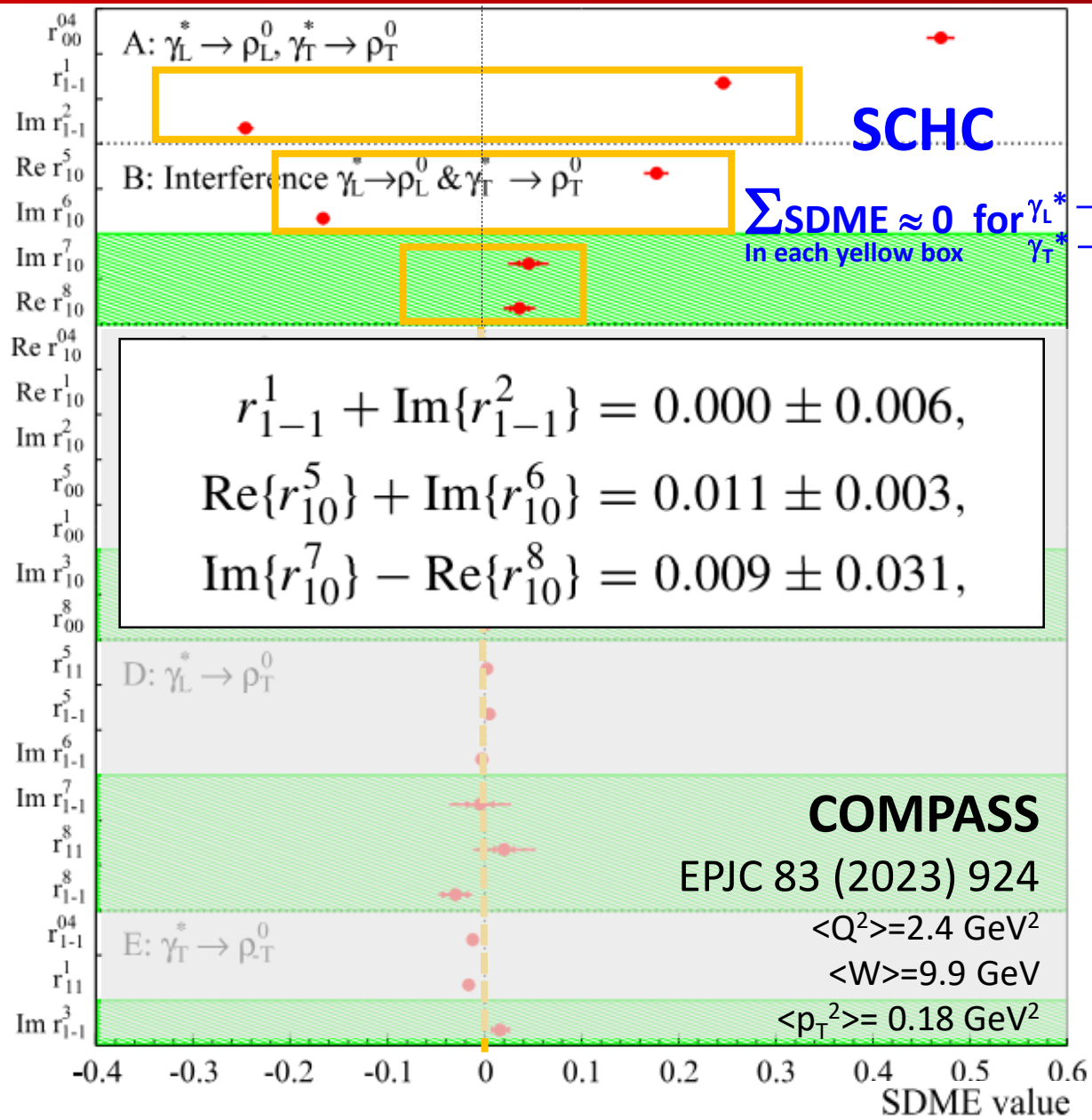


# COMPASS 2012 Exclusive $\rho^0$ and $\omega$ production on unpolarized proton



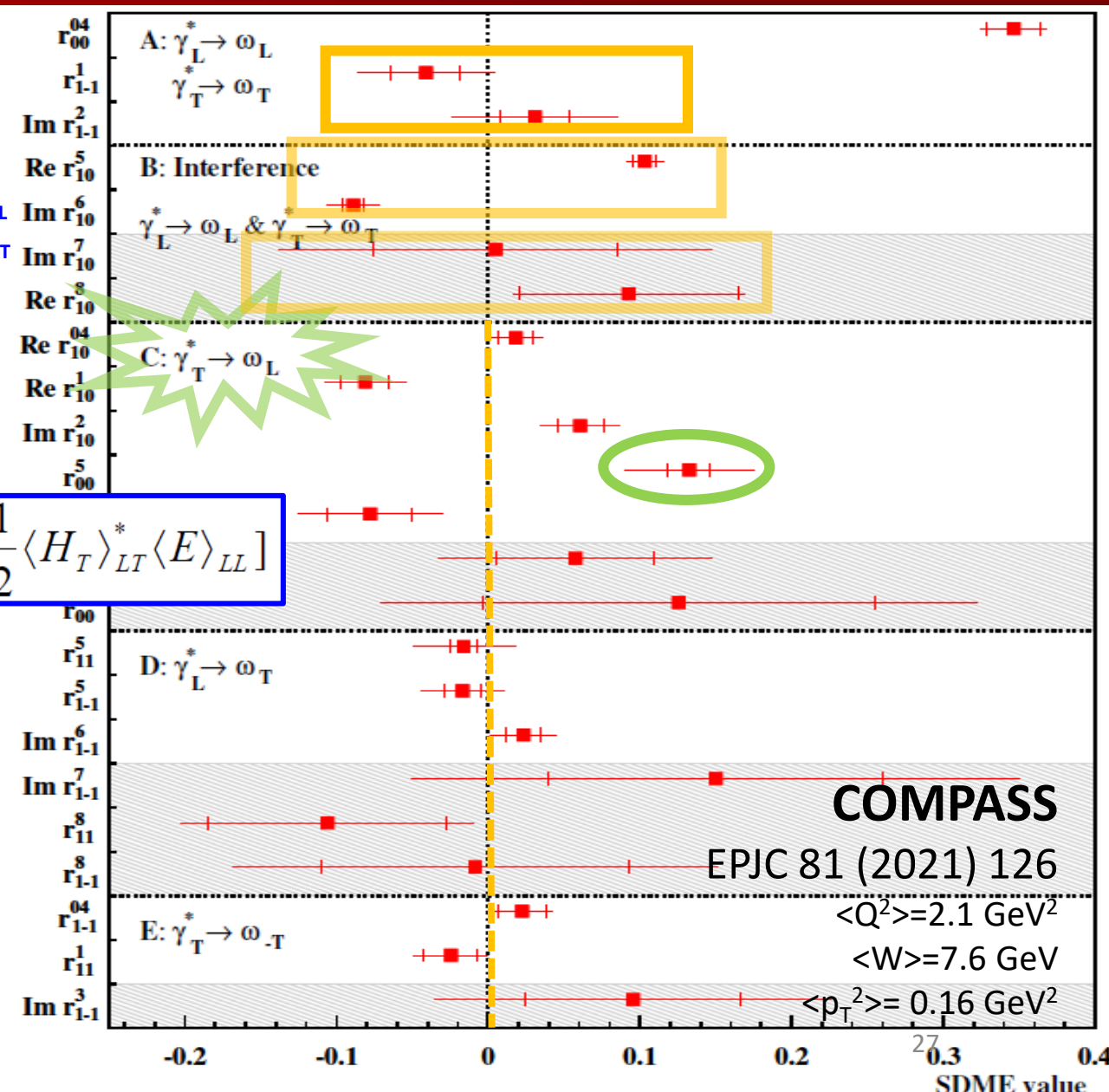
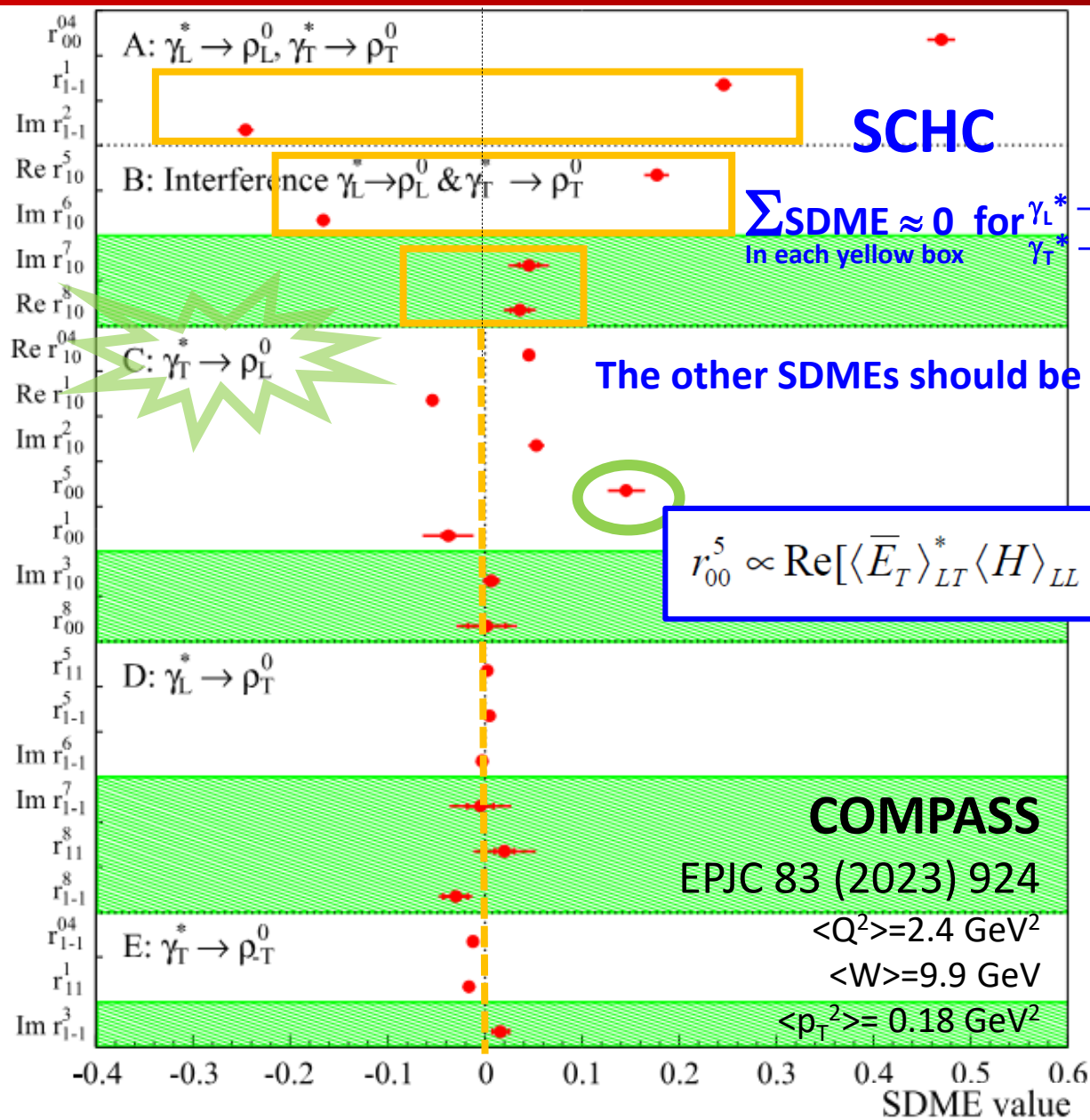


# COMPASS 2012 Exclusive $\rho^0$ and $\omega$ production on unpolarized proton





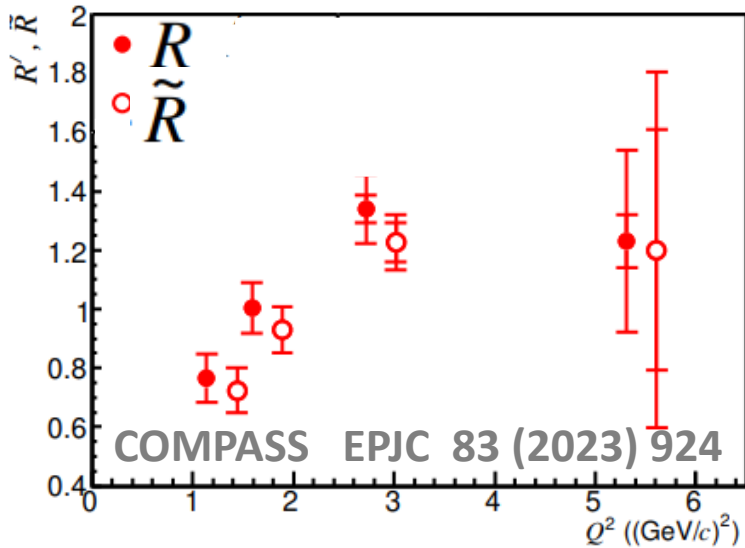
# COMPASS 2012 Exclusive $\rho^0$ and $\omega$ production on unpolarized proton



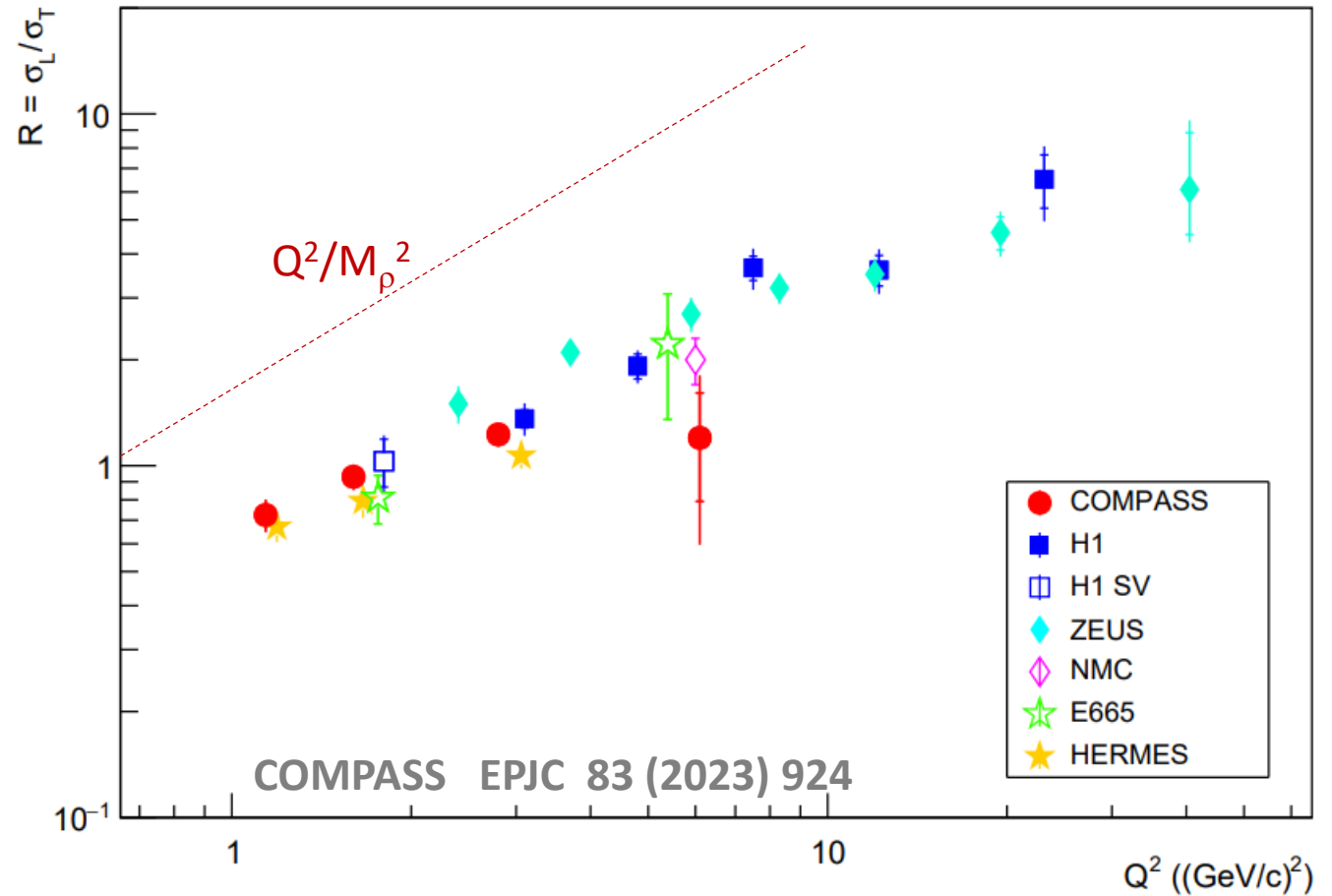
$$R = \frac{\sigma_L(\gamma_L^* \rightarrow V)}{\sigma_T(\gamma_T^* \rightarrow V)}$$

$$R = \frac{1}{\epsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}} \quad \text{only if SCHC}$$

In COMPASS domain evaluation of  $R$  and  $\tilde{R}$  considering violation of SCHC (but only NPE)



for all the experiments with  $Q^2 > 1 \text{ GeV}^2$



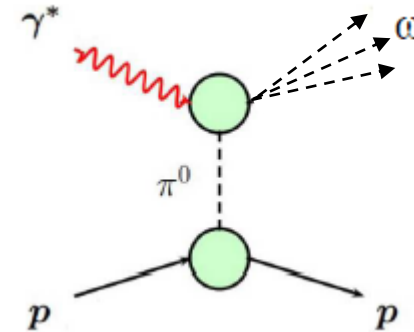
Deviations from the pQCD LO prediction in  $Q^2/M_\rho^2$  due to QCD evolution and  $q_T$  Transverse size effects of the meson smaller for  $\sigma_L$  than for  $\sigma_T$

## Natural (N) to Unnatural (U) Parity Exchange for $\gamma_T^* \rightarrow V_T$

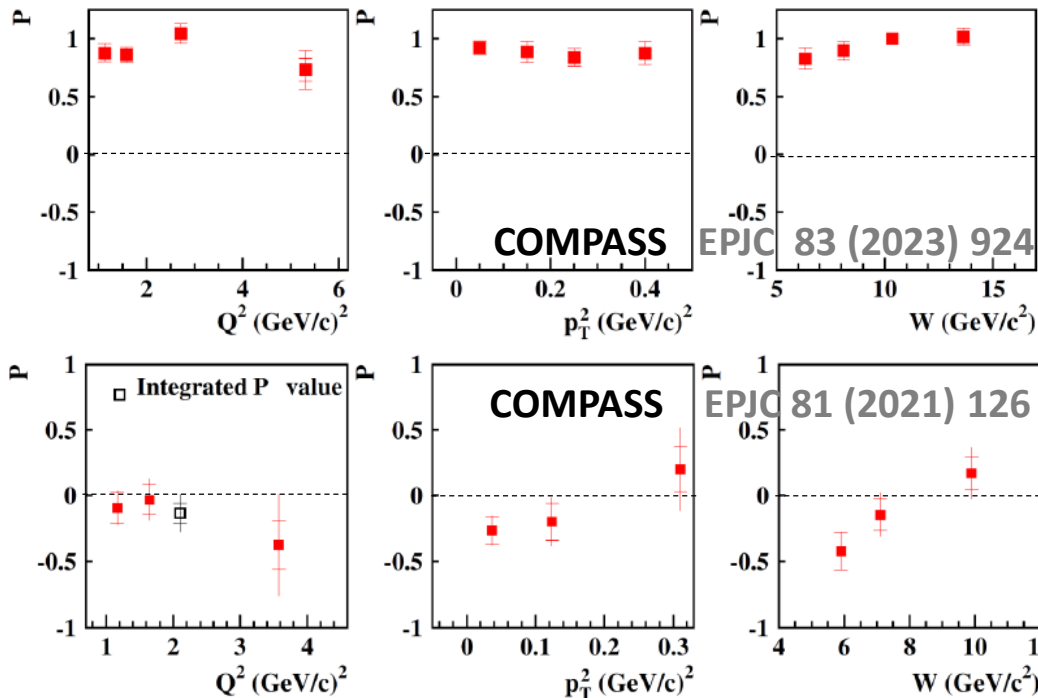
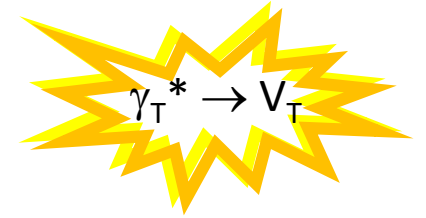
$$P = \frac{d\sigma_T^N(\gamma_T^* \rightarrow V_T) - d\sigma_T^U(\gamma_T^* \rightarrow V_T)}{d\sigma_T^N(\gamma_T^* \rightarrow V_T) + d\sigma_T^U(\gamma_T^* \rightarrow V_T)} \approx \frac{2r_{1-1}^1}{1 - r_{00}^{04} - 2r_{1-1}^{04}}$$

The pion pole exchange (UPE) is large for  $\omega$  compared to  $\rho^0$

$\Gamma(\omega \rightarrow \pi^0 \gamma) = 9 \times \Gamma(\rho^0 \rightarrow \pi^0 \gamma)$  as for  $\pi^0$  Vector Meson FF



It plays an important role in  $\omega$  production for:



$\rho^0$ :  $P \sim 1 \rightarrow$  NPE dominance  $P \sim 1$   
NPE with GPDs  $H, E$

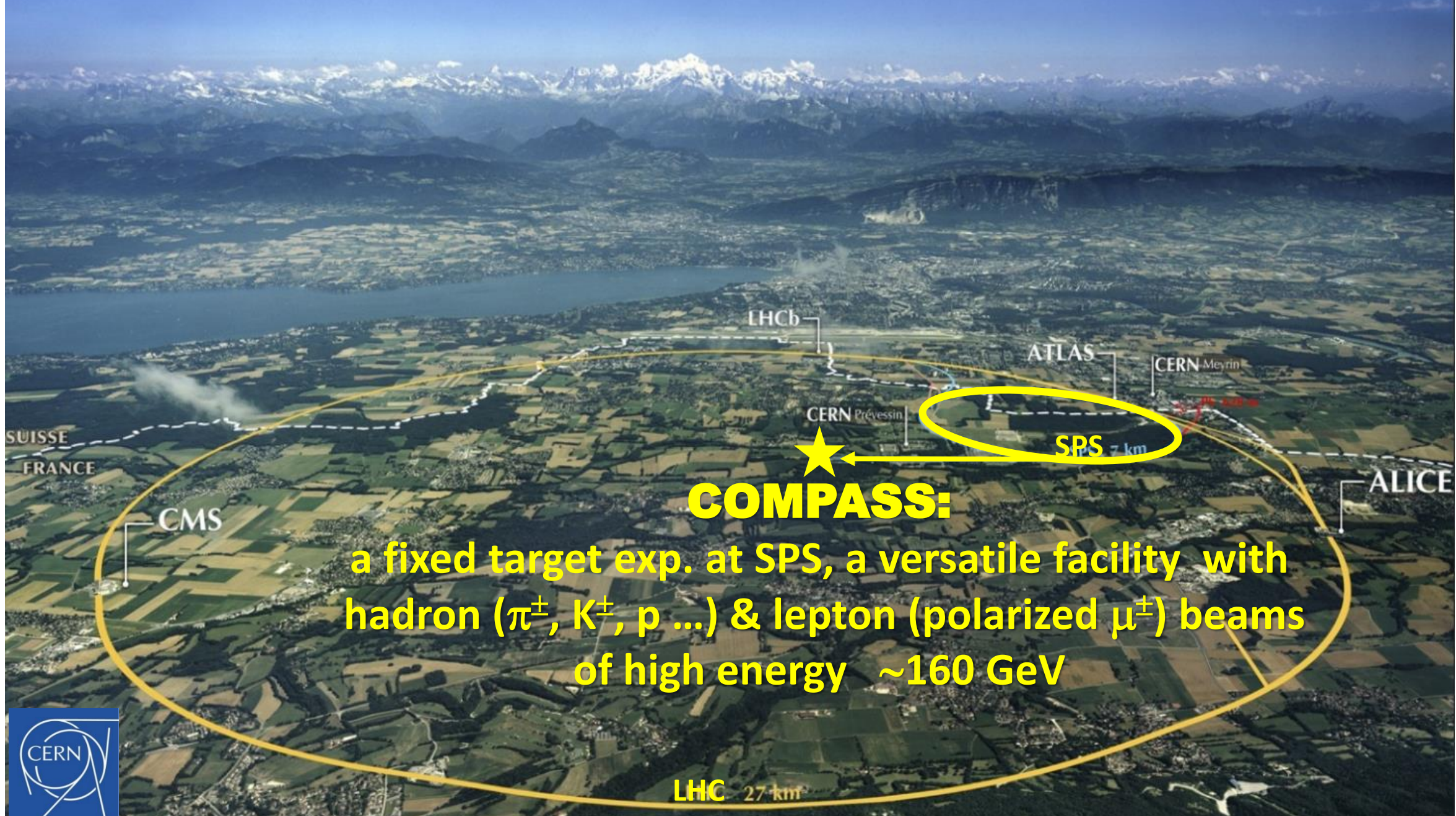
$\omega$ :  $P \sim 0 \rightarrow$  NPE  $\sim$  UPE  
UPE dominance at small  $W$  and  $p_T^2$   
UPE with GPDs  $\tilde{H}, \tilde{E}$  and the dominant pion pole

# Summary and perspective using 2016 + 2017 data

- ✓ **DVCS** and the **sum**  $\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$ 
  - $c_0$  and  $s_1$  and constrain on  $\text{Im}\mathcal{H}$  and Transverse extension of partons
- ✓ **DVCS** and the **difference**  $\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$ 
  - $c_1$  and constrain on  $\text{Re}\mathcal{H}$  (>0 as H1 or <0 as HERMES)  
for D-term and pressure distribution
- ✓ **Cross section** for  $\pi^0$
- ✓ **Cross section** and **SDME** for  $\rho^0, \omega, \phi, J/\psi$ 
  - ✓ Transversity GPDs
  - ✓ Gluon GPDs
  - ✓ Flavor decomposition

THANK YOU FOR YOUR ATTENTION



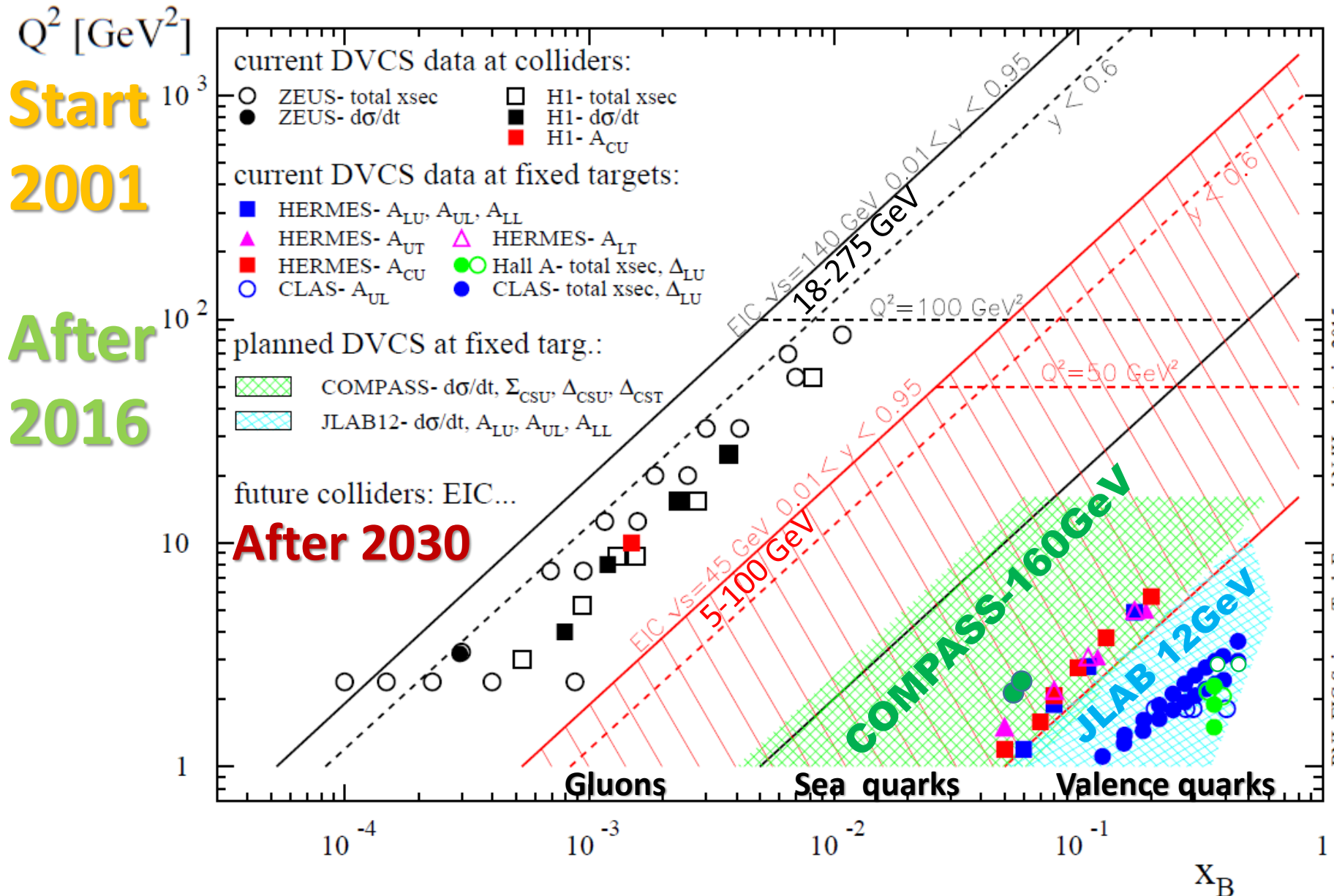


**COMPASS:**  
a fixed target exp. at SPS, a versatile facility with  
hadron ( $\pi^\pm$ ,  $K^\pm$ ,  $p$  ...) & lepton (polarized  $\mu^\pm$ ) beams  
of high energy  $\sim 160$  GeV





# Past and future experiments for DVCS $\ell p \rightarrow \ell' p' \gamma$



# COMPASS 12-16 Transverse extension of partons in the sea quark range

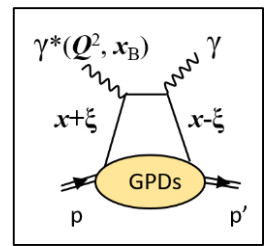
$$d\sigma^{DVCS}/dt = e^{-B|t|} = c_0^{DVCS} \propto (\text{Im}\mathcal{H})^2$$

$$c_0^{DVCS} \propto 4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) + \frac{t}{M^2}\mathcal{E}\mathcal{E}^*$$

In the COMPASS kinematics,  $x_B \approx 0.06$ , dominance of  $\text{Im}\mathcal{H}$   
 97% (GK model) 94% (KM model)

$$\text{Im}\mathcal{H} = H(x=\xi, \xi, t)$$

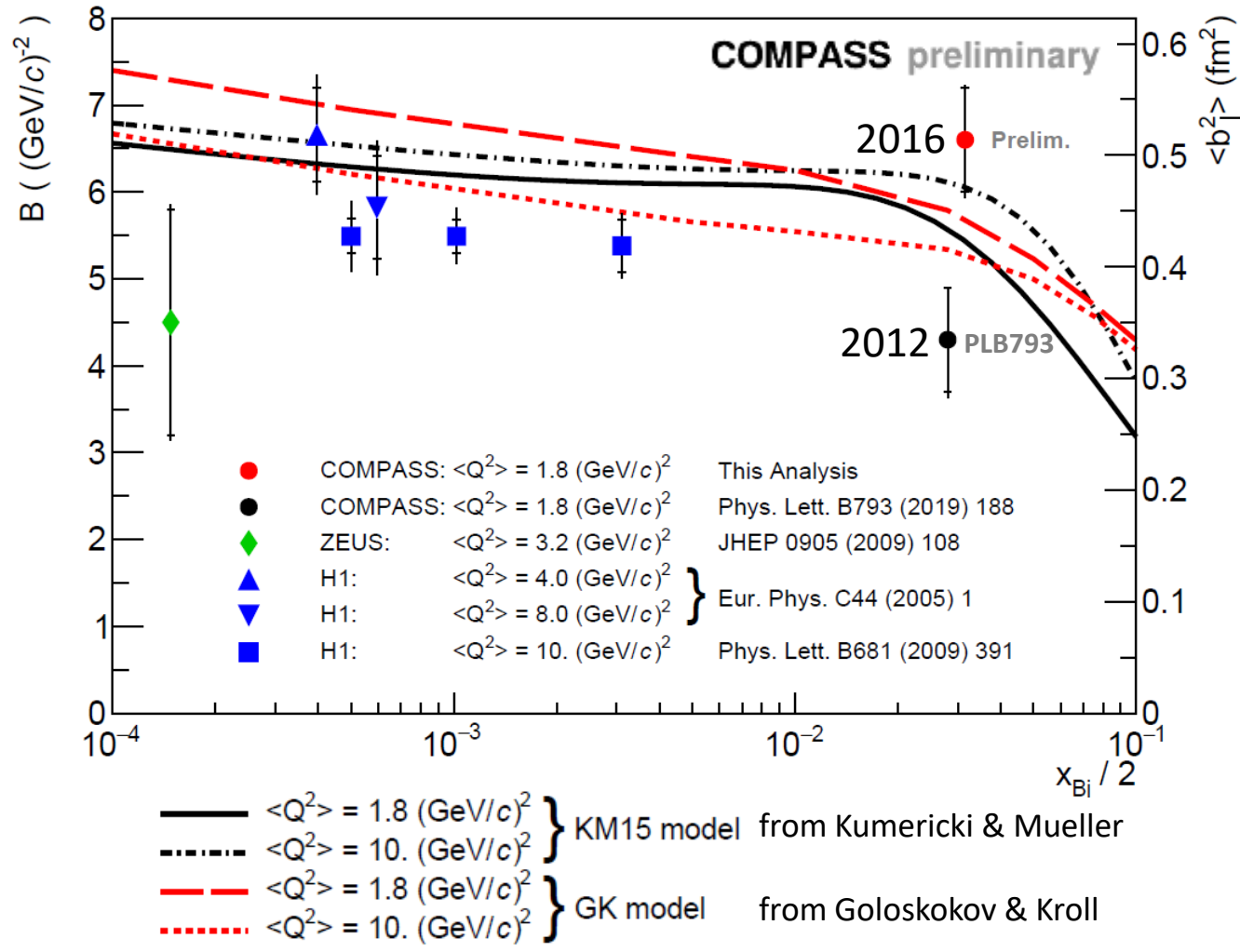
$$x = \xi \approx x_B/2 \text{ close to } 0$$



$$q(x, \mathbf{b}_\perp) = \int \frac{d^2\Delta_\perp}{(2\pi)^2} e^{-i\mathbf{b}_\perp \cdot \Delta_\perp} H^q(x, 0, -\Delta_\perp^2)$$

$$\langle b_\perp^2 \rangle_x^f = \frac{\int d^2b_\perp b_\perp^2 q_f(x, b_\perp)}{\int d^2b_\perp q_f(x, b_\perp)} = -4 \frac{\partial}{\partial t} \log H^f(x, \xi=0, t) \Big|_{t=0}$$

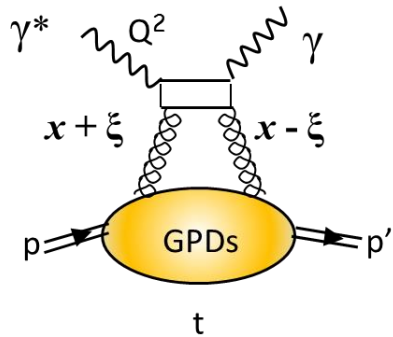
$$\langle b_\perp^2(x) \rangle \approx 2B(\xi)$$



# → nucleon tomography in the gluon domain at HERA

$$d\sigma^{\text{DVCS}}/dt = e^{-B|t|}$$

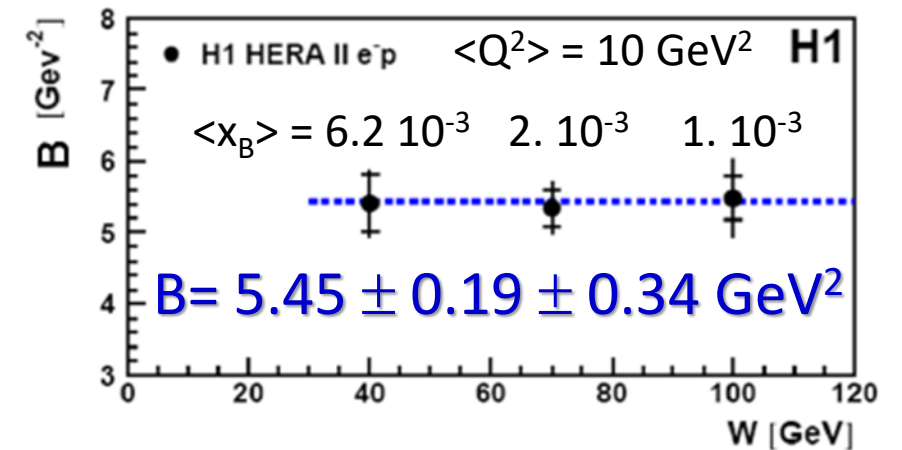
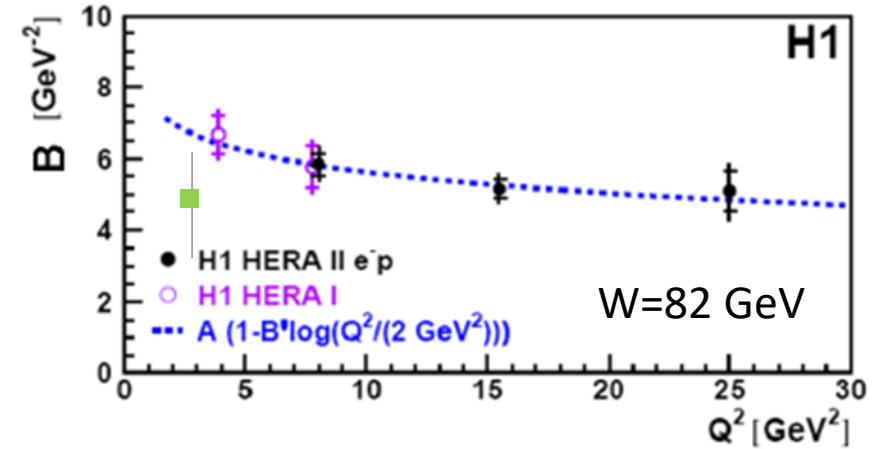
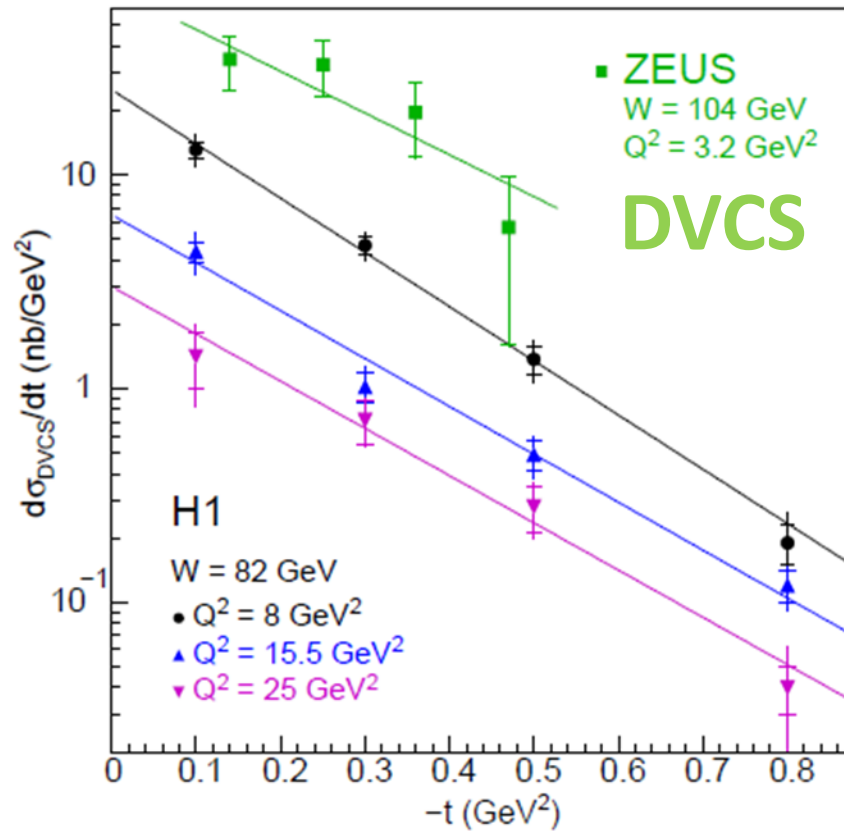
Dominance of  $\text{Im}\mathcal{H}$



ZEUS-H1  
Data collected  
1995-2007

B is related to the transversed size of the scattering object

Aaron et al., H1 Coll, PLB659 (2008)



$$\langle b_{\perp}^2(x) \rangle \approx 2B(\xi)$$

$$\sqrt{\langle b_{\perp}^2 \rangle} = 0.65 \pm 0.02 \text{ fm} \quad \text{to be compared to} \quad \sqrt{4 \frac{d}{dt} F_1^p} \Big|_{t=0} = 0.67 \pm 0.01 \text{ fm}$$

$$\frac{d^2\sigma_{\gamma^*p}^{\leftrightarrow}}{dtd\phi} = \frac{1}{2\pi} \left[ \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \epsilon \cos(2\phi) \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos\phi \frac{d\sigma_{LT}}{dt} \mp |P_l| \sqrt{2\epsilon(1-\epsilon)} \sin\phi \frac{d\sigma'_{LT}}{dt} \right]$$

$$\frac{d\sigma_L}{dt} \propto \left[ (1-\xi^2) |\langle \tilde{H} \rangle|^2 - 2\xi^2 \operatorname{Re} [\langle \tilde{H} \rangle^* \langle \tilde{E} \rangle] - \frac{t'}{4M^2} \xi^2 |\langle \tilde{E} \rangle|^2 \right],$$

$$\frac{d\sigma_T}{dt} \propto \left[ (1-\xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8M^2} |\langle \bar{E}_T \rangle|^2 \right],$$

$$\frac{d\sigma_{TT}}{dt} \propto t' |\langle \bar{E}_T \rangle|^2,$$

$$\frac{d\sigma_{LT}}{dt} \propto \xi \sqrt{1-\xi^2} \sqrt{-t'} \operatorname{Re} [\langle H_T \rangle^* \langle \tilde{E} \rangle + \langle \bar{E}_T \rangle^* \langle \tilde{H} \rangle],$$

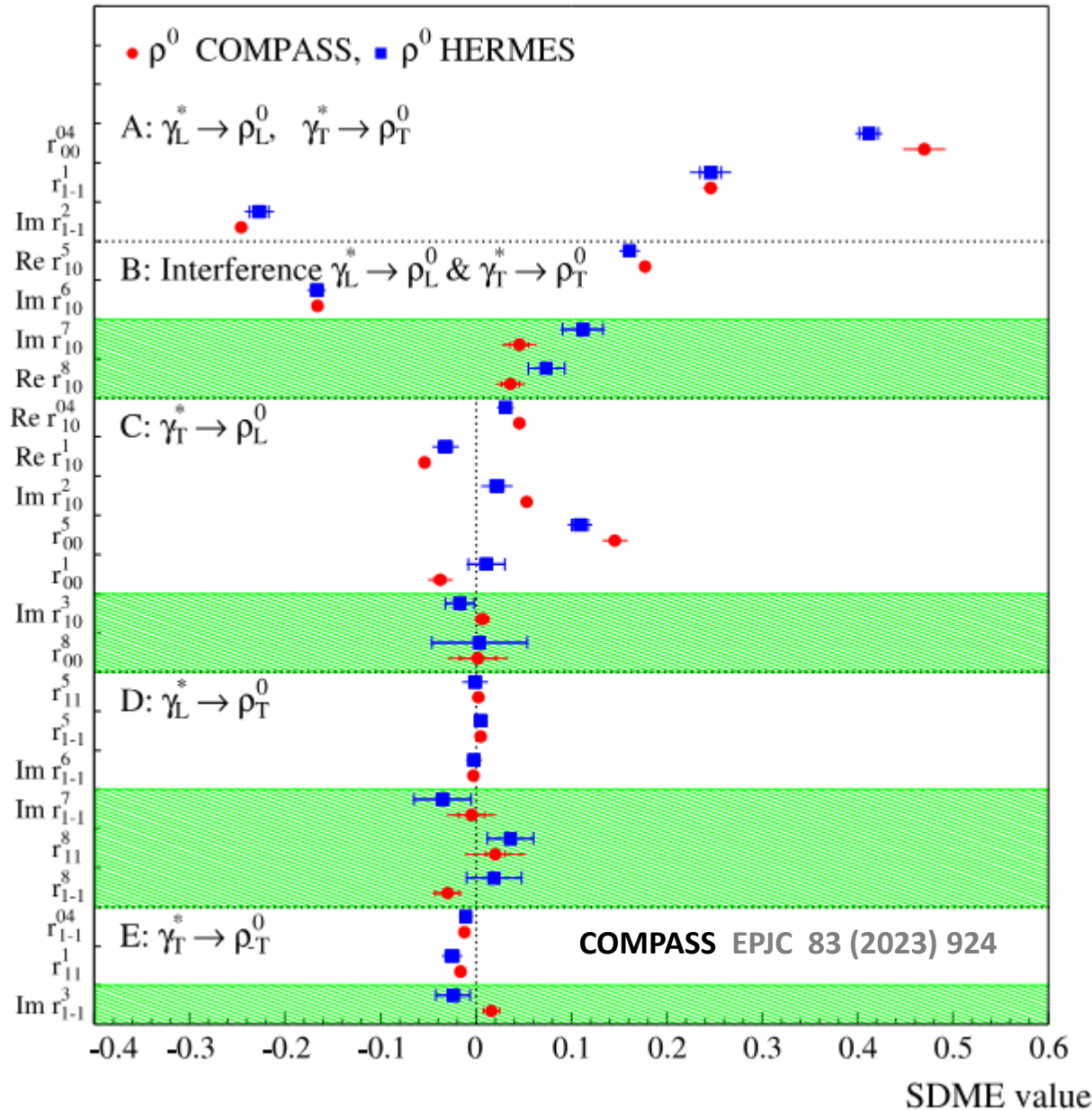
$$\frac{d\sigma_{LT'}}{dt} \propto \xi \sqrt{1-\xi^2} \sqrt{-t'} \operatorname{Im} [\langle H_T \rangle^* \langle \tilde{E} \rangle + \langle \bar{E}_T \rangle^* \langle \tilde{H} \rangle].$$

At COMPASS

$$|P_l| \sqrt{2\epsilon(1-\epsilon)} \simeq 0.06$$



# Comparison $\rho^0$ SDMEs at COMPASS and HERMES



**Fig. 12** Comparison of the 23 SDMEs for exclusive  $\rho^0$  lepton production on the proton extracted in the entire kinematic regions of the HERMES and COMPASS experiments. For HERMES the average kinematic values are  $\langle Q^2 \rangle = 1.96$  (GeV/c)<sup>2</sup>,  $\langle W \rangle = 4.8$  GeV/c<sup>2</sup>,  $\langle |t'| \rangle = 0.13$ , while those for COMPASS are  $\langle Q^2 \rangle = 2.40$  (GeV/c)<sup>2</sup>,  $\langle W \rangle = 9.9$  GeV/c<sup>2</sup>,  $\langle p_T^2 \rangle = 0.18$  (GeV/c)<sup>2</sup>. Inner error bars represent statistical uncertainties and outer ones statistical and systematic uncertainties added in quadrature. Unpolarised (polarised) SDMEs are displayed in unshaded (shaded) areas