



Transverse spin-dependent asymmetries at COMPASS experiment

International Conference on High Energy Physics

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20 VII 2024, Prague

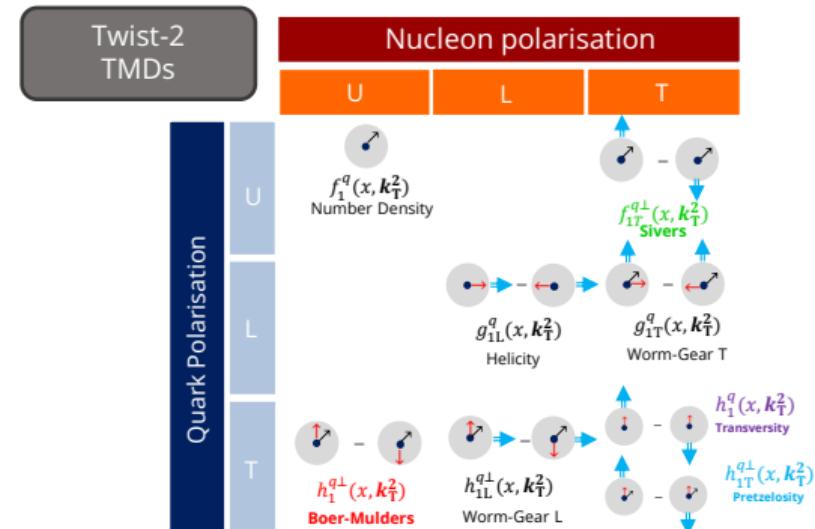
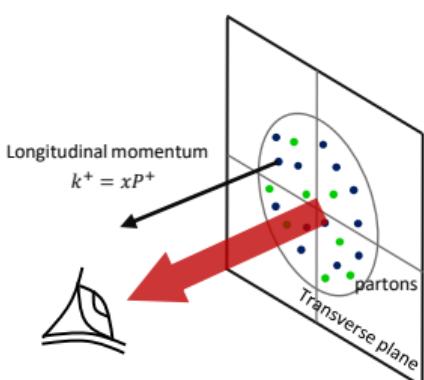
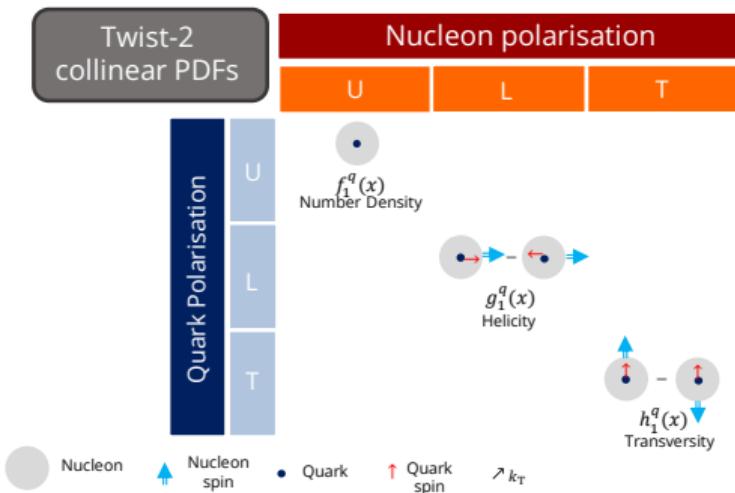
Transverse Momentum Dependent Parton Distribution Functions

"Well begun is half done."

Old Proverb



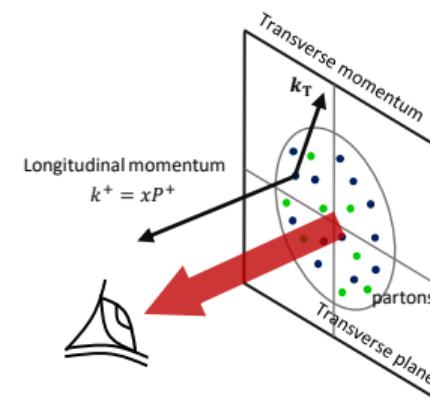
Nucleon spin structure



TMD PDFs accessed through target spin dependent azimuthal asymmetries both in SIDIS and Drell-Yan processes.

$$h_1^{q\perp}(\text{SIDIS}) = -h_1^{q\perp}(\text{DY})$$

$$f_{1T}^{q\perp}(\text{SIDIS}) = -f_{1T}^{q\perp}(\text{DY})$$



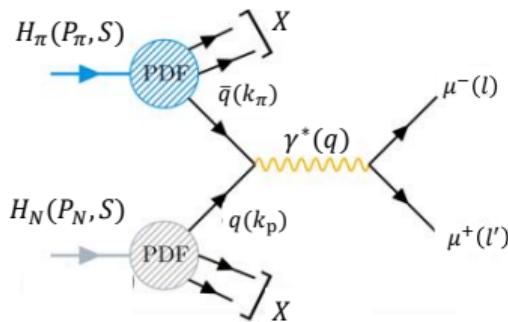


Single polarised Drell-Yan process

*"Polarisation data has often been the graveyard of fashionable theories.
If theorists had their way, they might well ban such measurements
altogether out of self-protection."*

James Bjorken

Single polarised Drell-Yan process



Boer-Mulders

$$A_U^{\cos(2\varphi_{CS})} \propto h_{1,\pi}^{q \perp} \otimes h_{1,N}^{q \perp}$$

Sivers

$$A_T^{\sin(\varphi_S)} \propto f_{1,\pi}^q \otimes f_{1T,N}^{q \perp}$$

Pretzelosity

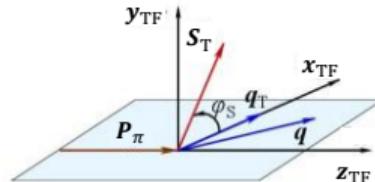
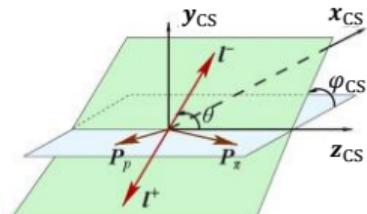
$$A_T^{\sin(2\varphi_{CS} + \varphi_S)} \propto h_{1,\pi}^{q \perp} \otimes h_{1T,N}^{q \perp}$$

Transversity

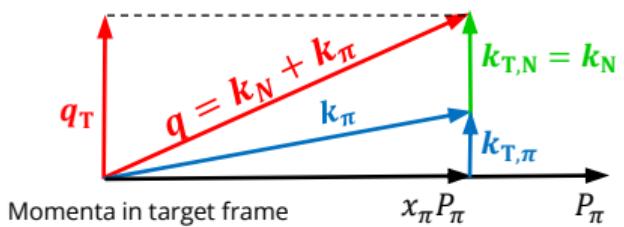
$$A_T^{\sin(2\varphi_{CS} - \varphi_S)} \propto h_{1,\pi}^{q \perp} \otimes h_{1,N}^q$$

Cross-section, LO TMD approach for transversely polarised target:

$$\frac{d\sigma^{LO}}{dx_p dx_\pi d^2 q_T d\varphi_{CS} d(\cos\theta) d\varphi_S} \propto \left\{ \begin{array}{l} 1 + D \sin^2 \theta \cos(2\varphi_{CS}) A_U^{\cos(2\varphi_{CS})} \\ \quad \sin(\varphi_S) A_T^{\sin(\varphi_S)} + \\ + |S_T| \left[D \sin^2 \theta \left(\begin{array}{l} \sin(2\varphi_{CS} + \varphi_S) A_T^{\sin(2\varphi_{CS} + \varphi_S)} \\ + \sin(2\varphi_{CS} - \varphi_S) A_T^{\sin(2\varphi_{CS} - \varphi_S)} \end{array} \right) \right] \end{array} \right\}$$



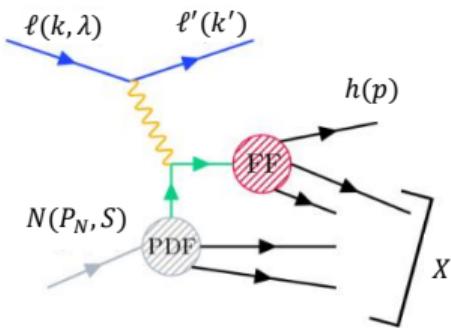
The convolution of TMD PDFs runs over the intrinsic transverse momenta k_T .



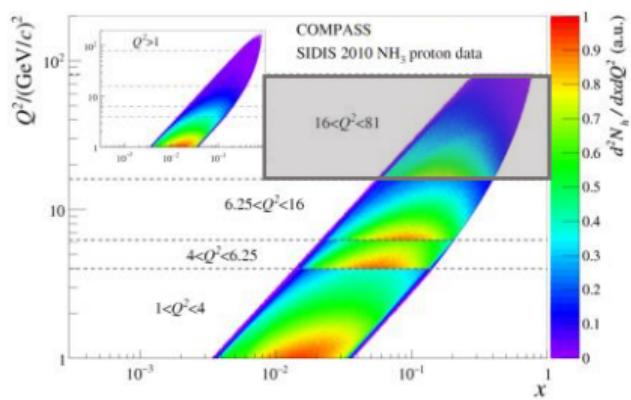
Semi Inclusive Deep Inelastic Scattering

Drell-Yan process

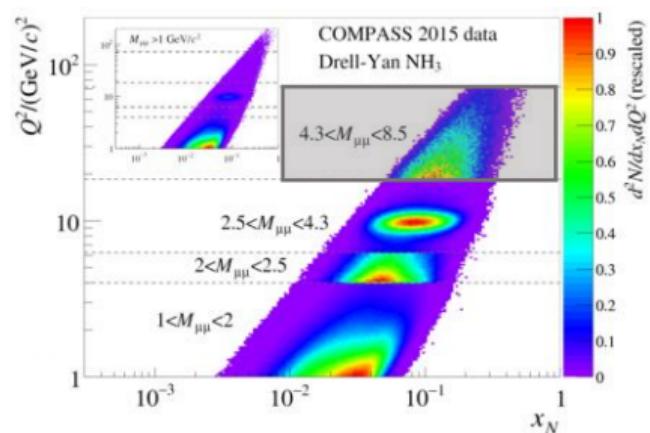
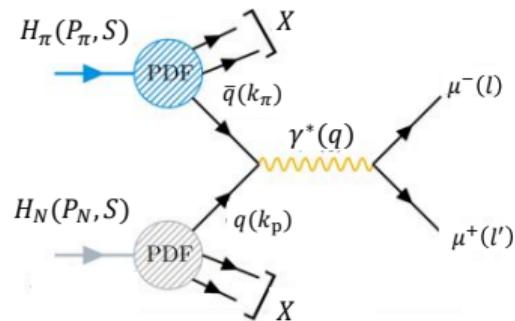
$$A_{SIDIS} \propto PDF_N \otimes FF$$



SIDIS vs. DY



COMPASS achieves comparable $Q^2 - x$ kinematic coverage
Minimizing possible Q^2 evolution effects



COMPASS Experiment

"Knowledge is of no value unless you put it into practice."

Anton Chekhov



COMPASS Collaboration



Common Muon and Proton Apparatus for Structure and Spectroscopy



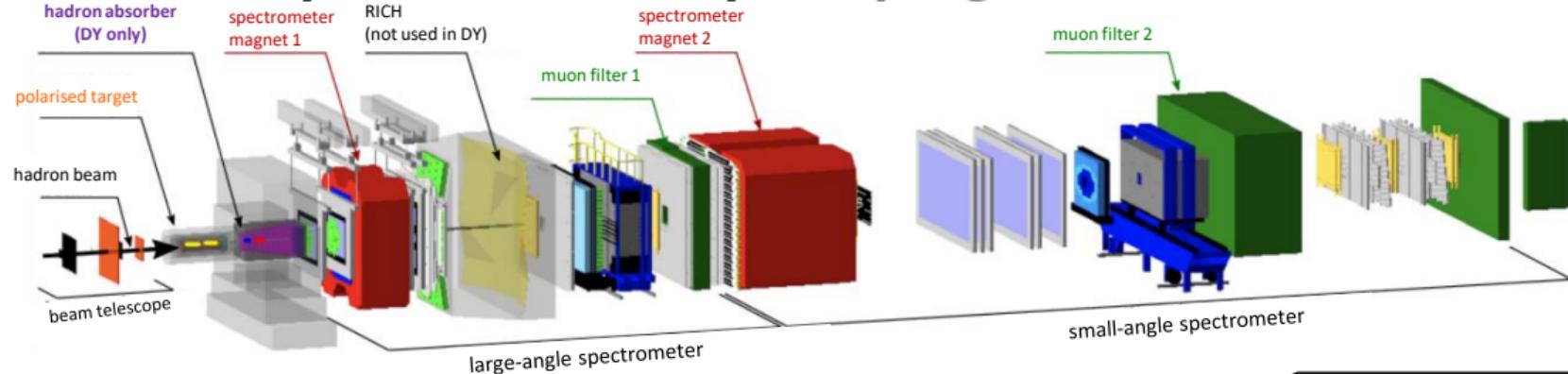
An extensive research programme on the structure of nucleons, including spin and on hadron spectroscopy

Drell Yan data taking
2015 + 2018

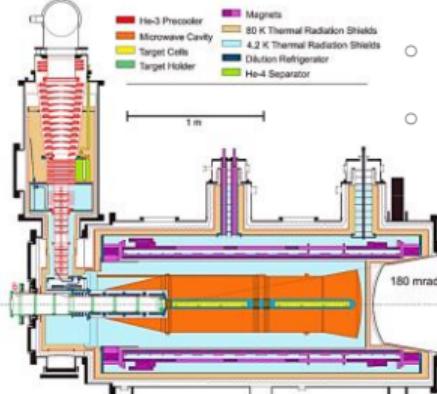
- 24 institutions from 14 countries (approximately 220 physicists)
- CERN SPS North Area
- Fixed target experiment



COMPASS experimental setup: DY programme

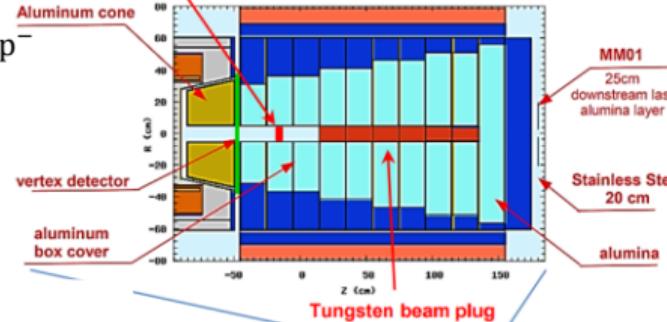


Polarised target



- High energy negative hadron beam: 190 GeV Nuclear target (Al)

Hadron absorber



- Beam composition: **97% π^- , 2% K^- , 1% p^-**

- Targets:

- polarised NH_3
 - Al, W

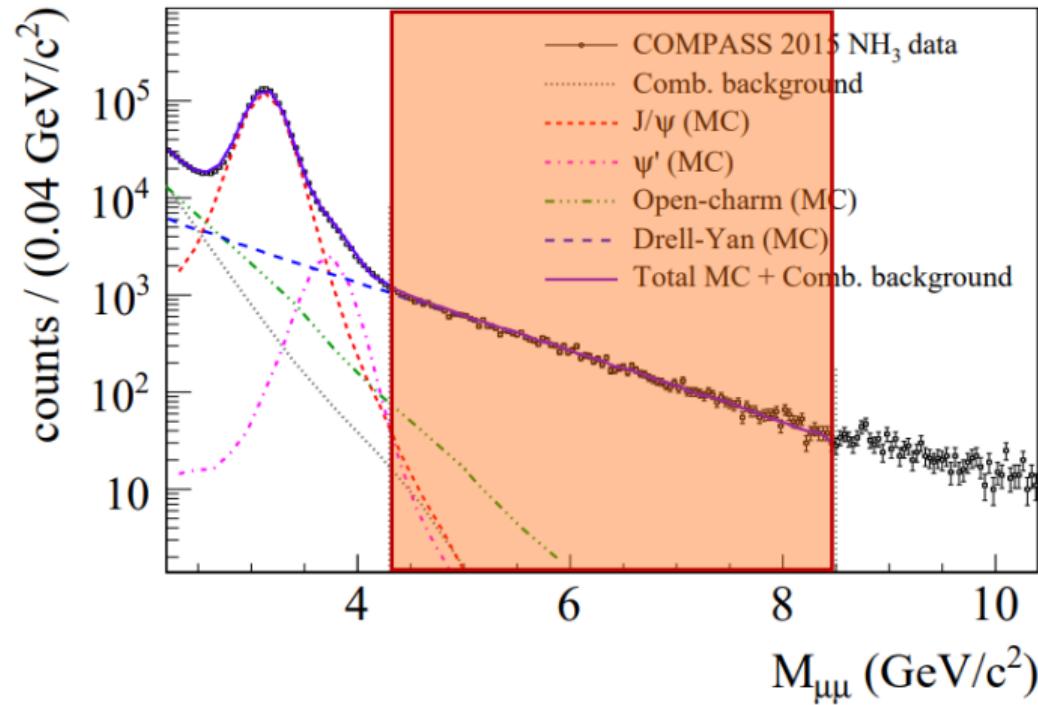
- Hadron absorber
- Muon identification system
- 2 spectrometer stages for a wide phase space coverage

Drell-Yan measurement at COMPASS

96% pure Drell-Yan in dimuon mass range: $4.3 < M_{\mu\mu}/(\text{GeV}/c^2) < 8.5$

WTSA

$4.3 < M_{\mu\mu}/(\text{GeV}/c^2) < 8.5$



Drell-Yan measurement at COMPASS

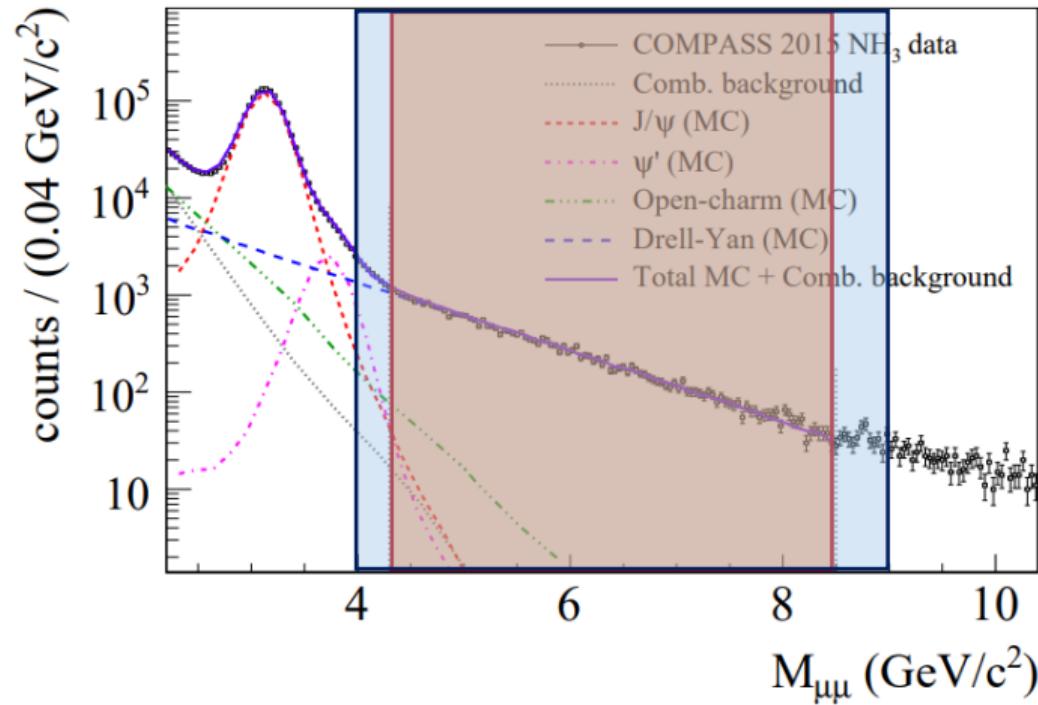
With background correction, achieve mass range: $4.0 < M_{\mu\mu}/(\text{GeV}/c^2) < 9.0$

WTSA

$$4.3 < M_{\mu\mu}/(\text{GeV}/c^2) < 8.5$$

TSA

$$4.0 < M_{\mu\mu}/(\text{GeV}/c^2) < 9.0$$





TSA

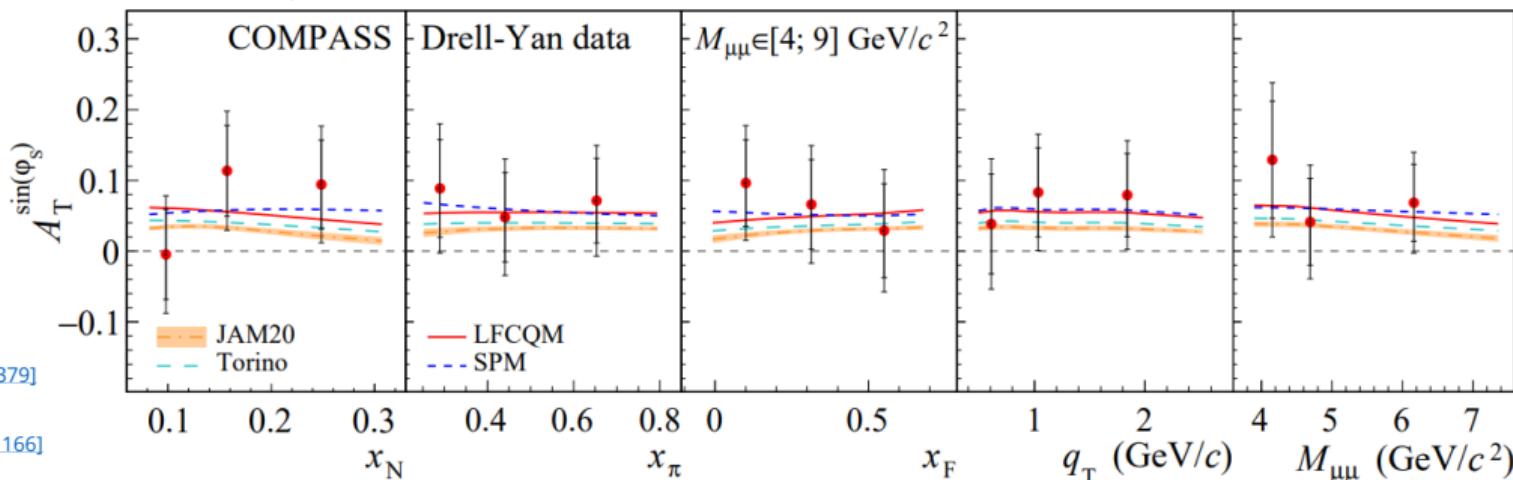
“Nothing in life is to be feared, it is only to be understood. Now is the time to understand more, so that we may fear less.”

Maria Skłodowska-Curie

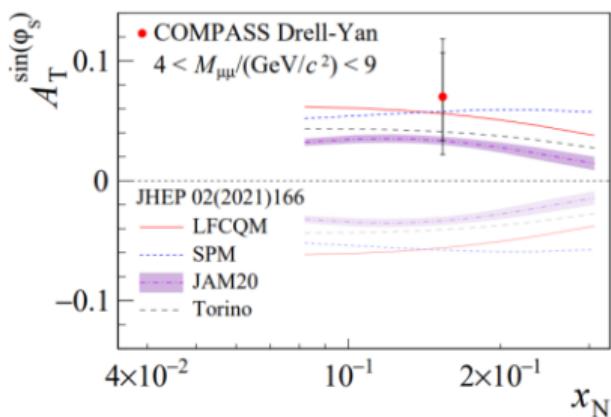
$$A_T^{\sin(\phi_S)} \propto f_{1,\pi}^q \otimes f_{1T,N}^{q\perp}$$

TSAs Results: Sivers (SIDIS and DY)

**Sivers
DY TSA**

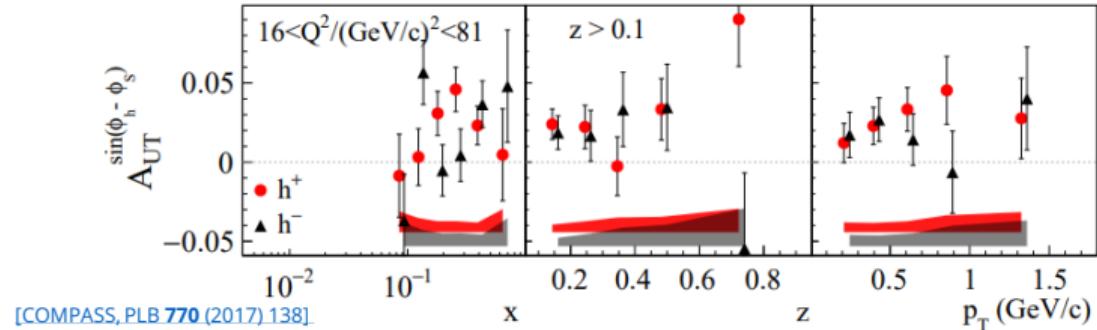


Theo. Pred.: [JHEP 02 (2021) 166]



SIDIS Sivers TSA

$$A_{UT}^{\sin(\phi_h - \phi_S)} \propto f_{1T,N}^{q\perp} \otimes D_{1q}^h$$

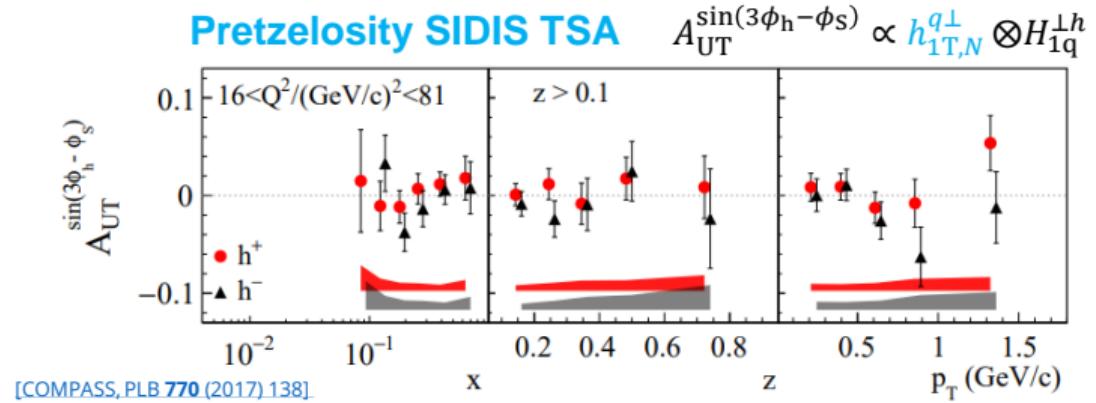
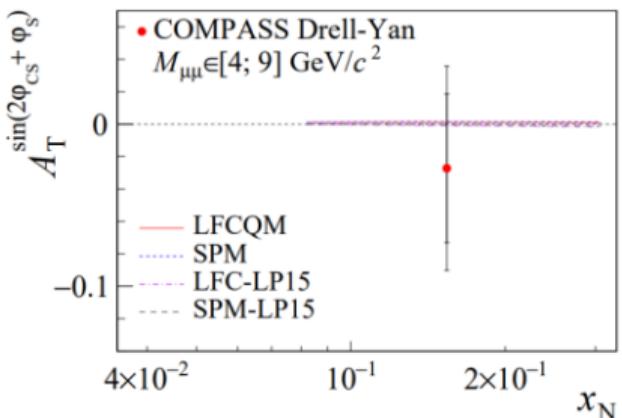
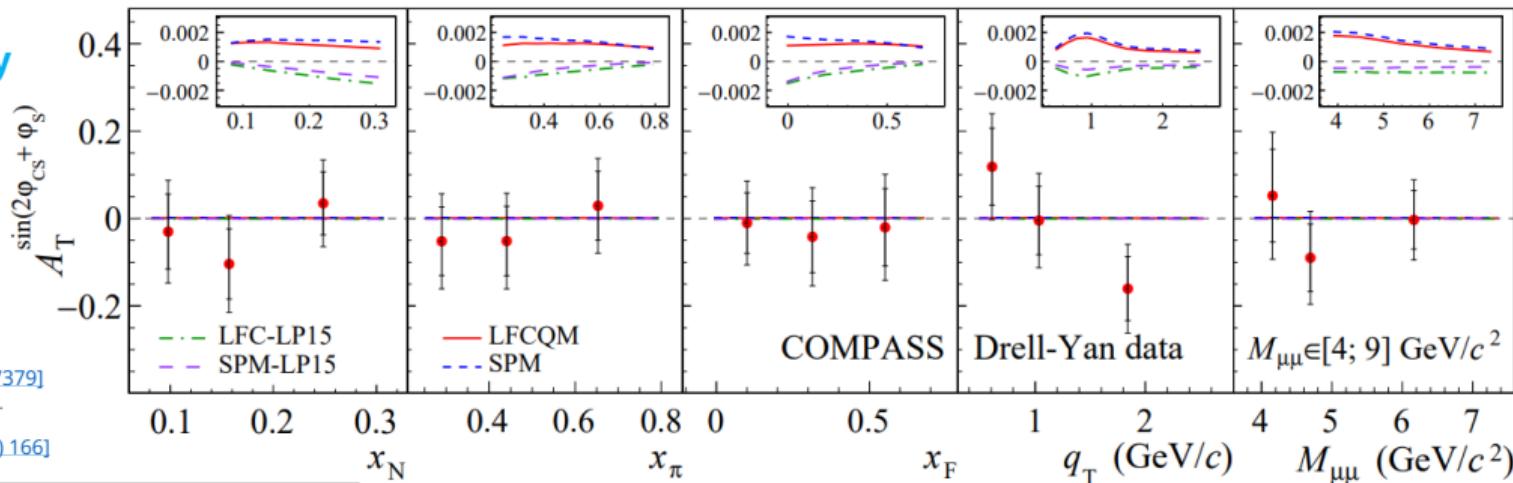


Transverse spin-dependent asymmetries at COMPASS experiment

$$A_T^{\sin(2\phi_{CS} + \phi_S)} \propto h_{1,\pi}^{q\perp} \otimes h_{1T,N}^{q\perp}$$

TSA Results: Pretzelosity (SIDIS and DY)

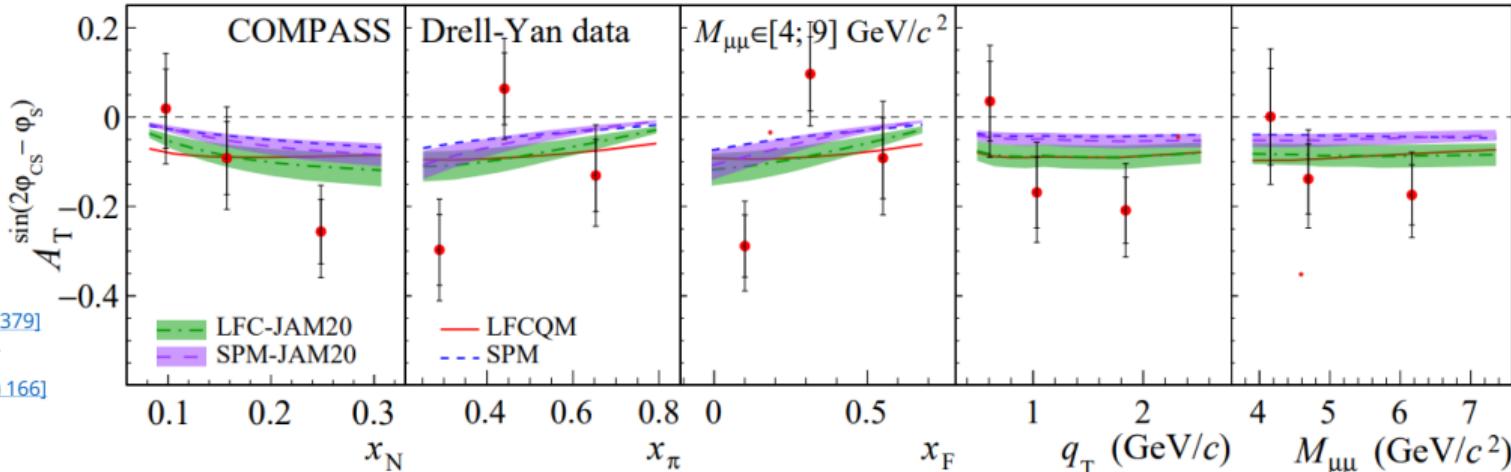
Pretzelosity DY TSA



$$A_T^{\sin(2\phi_{CS} - \phi_S)} \propto h_{1,\pi}^{q_\perp} \otimes h_{1,N}^q$$

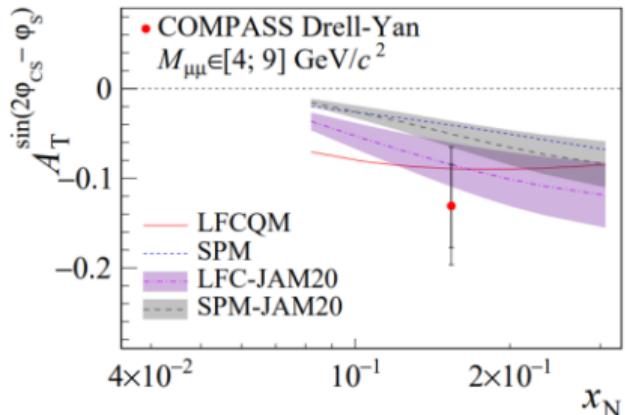
TSA Results: Transversity (SIDIS and DY)

Transversity DY TSA

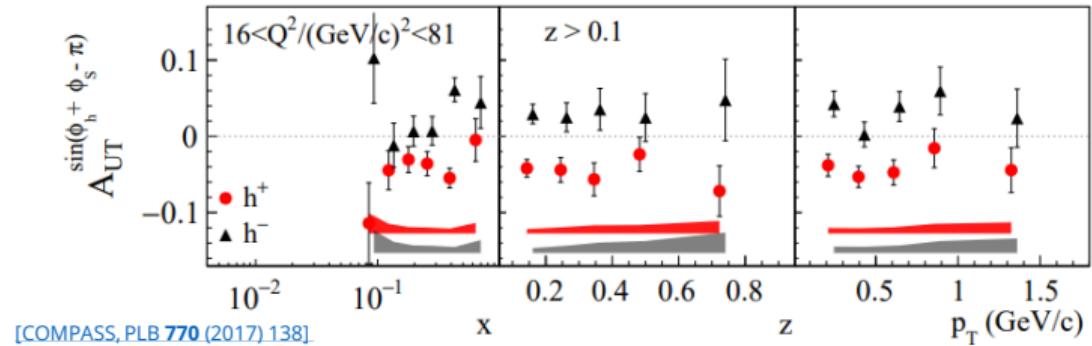


[COMPASS, hep-ex/2312.17379]
to appear in Phys. Rev. Lett.

Theo. Pred.: [JHEP 02 (2021) 166]



Transversity SIDIS TSA

$$A_{UT}^{\sin(\phi_h + \phi_s)} \propto h_{1,N}^q \otimes H_{1q}^{\perp h}$$


[COMPASS, PLB 770 (2017) 138]

Weighted TSAs

*"With four parameters I can fit an elephant,
and with five I can make him wiggle his
trunk."*

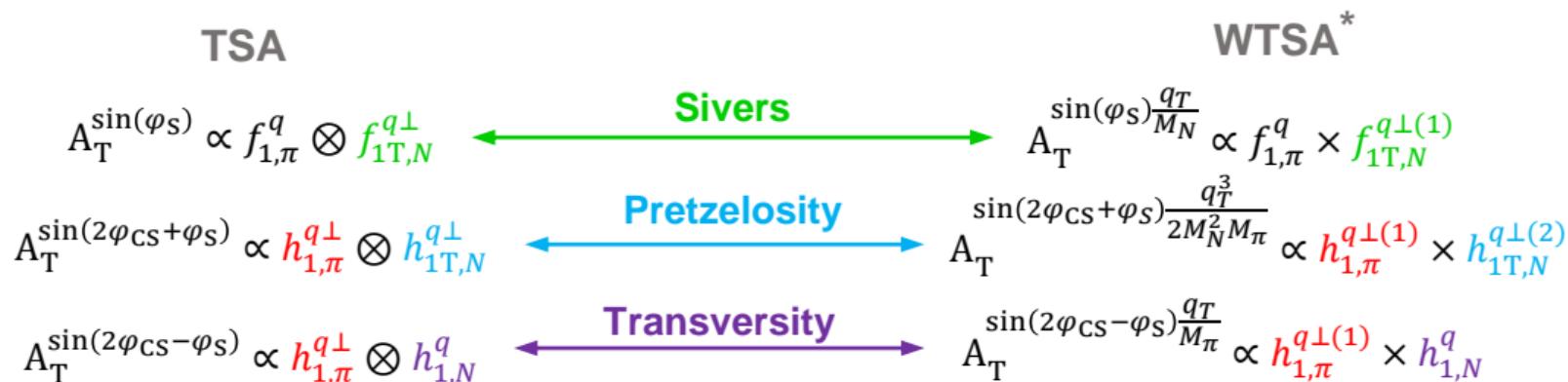
John von Neumann



Weighted TSAs in Drell-Yan

The convolution cannot be resolved without assumptions about the dependence of the TMD PDF on the intrinsic transverse momentum.

Weighting with powers of the transverse momentum allows to avoid assumptions on k_T .



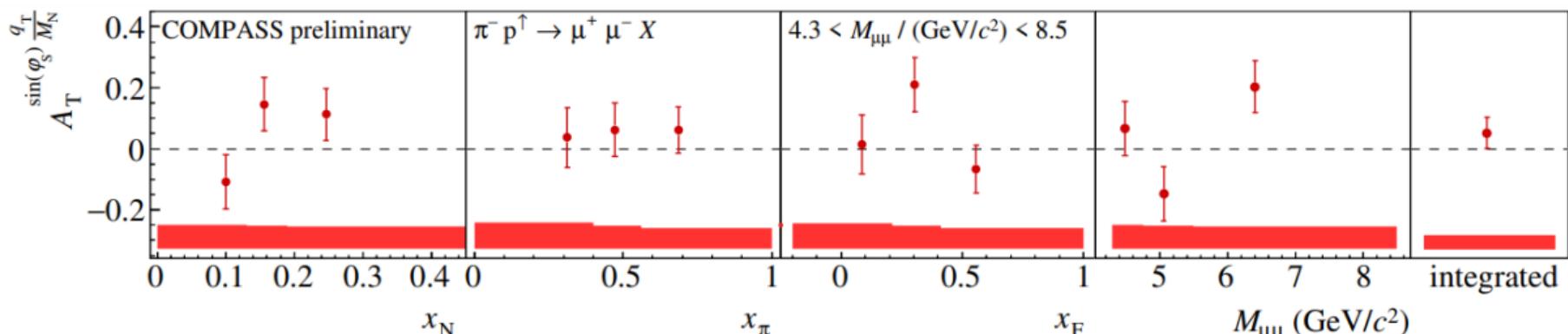
* The n-th moment of a TMD PDF of a pion or proton:

$$f^{(n)}(x) = \int d^2 k_T \left(\frac{k_T^2}{2M^2} \right)^n f(x, k_T^2)$$

WTSA Results: Sivers (SIDIS and DY)

Sivers DY WTSAs

$$A_T^{\sin(\phi_S) \frac{q_T}{M_N}} \propto f_{1,\pi}^q \times f_{1T,N}^{q\perp(1)}$$

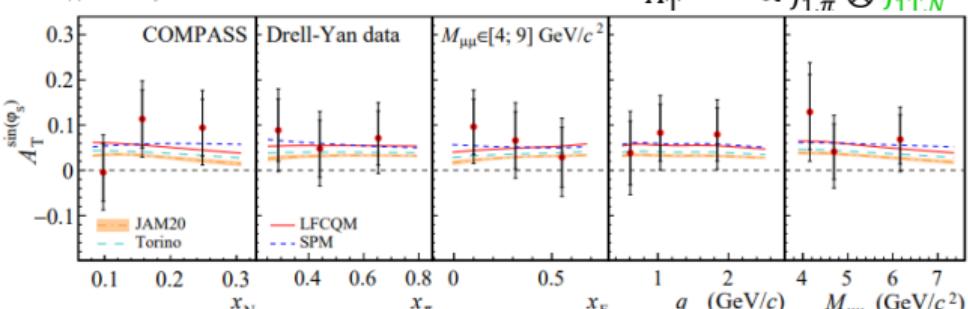


1 σ positive Sivers WTSAs compatible with Sivers TSA

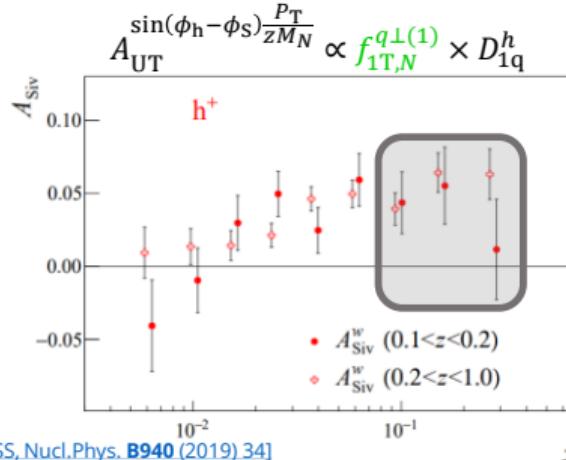
[COMPASS, hep-ex/2312.17379]
to appear in Phys. Rev. Lett.

Sivers DY TSA

$$A_T^{\sin(\phi_S) \propto f_{1,\pi}^q \otimes f_{1T,N}^{q\perp}}$$



Sivers
SIDIS
WTSAs



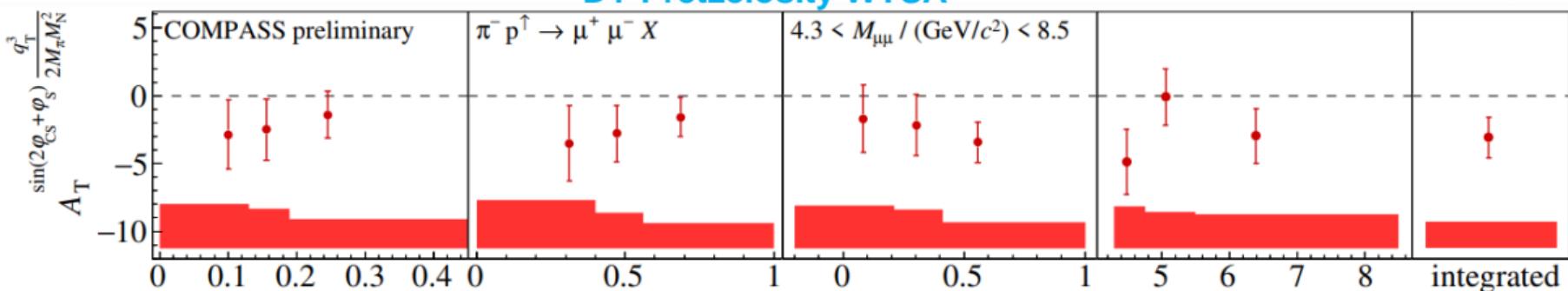
[COMPASS, Nucl.Phys. B940 (2019) 34]

WTSA Results: Pretzelosity (DY)

Comparison of WTSAs and TSA for **Pretzelosity** asymmetry

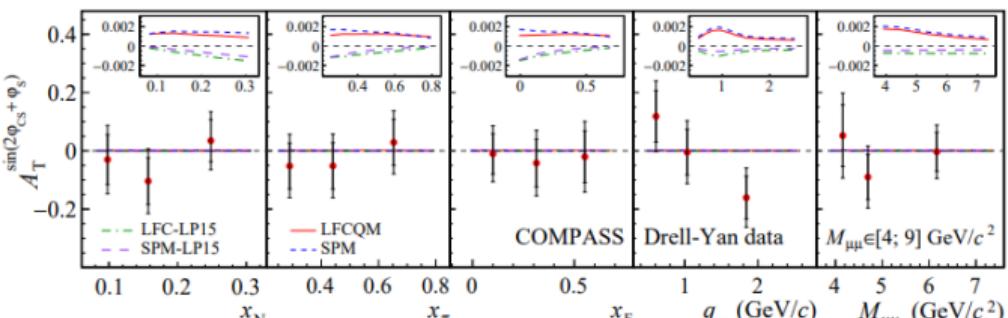
$$A_T \propto h_{1,\pi}^{q\perp(1)} \times h_{1T,N}^{q\perp(2)}$$

DY Pretzelosity WTSAs



$$A_T \propto h_{1,\pi}^{\perp} \otimes h_{1T,N}^{\perp}$$

DY Pretzelosity TSA



Pretzelosity is expected to be zero
2 σ negative Pretzelosity WTSAs

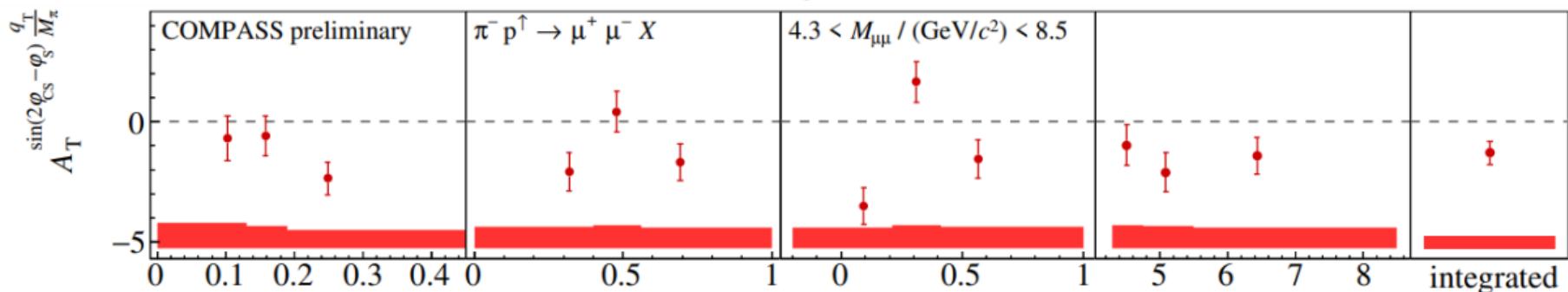
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WTSA Results: Transversity (DY)

Comparison of WTSAs and TSA for **Transversity** asymmetry

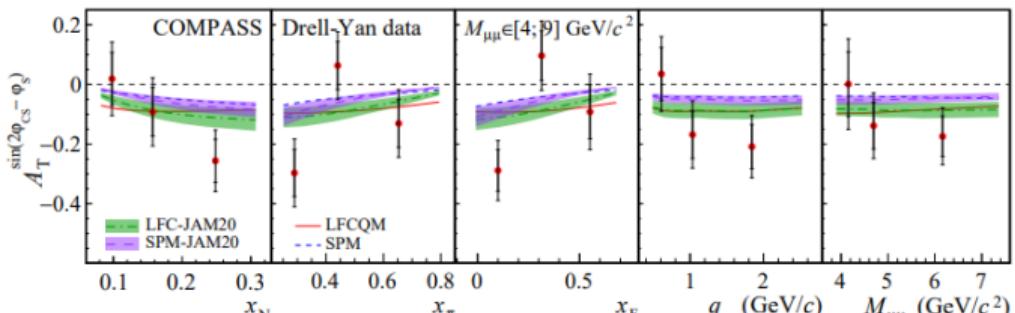
DY Transversity WTSAs

$$A_T^{\sin(2\varphi_{CS} - \varphi_S) \frac{q_T}{M_\pi}} \propto h_{1,\pi}^{q \perp(1)} \times h_{1,N}^q$$



$$A_T^{\sin(2\varphi_{CS} - \varphi_S) \propto h_{1,\pi}^\perp \otimes h_{1,N}} \quad x_N \quad x_\pi \quad x_F \quad M_{\mu\mu} (\text{GeV}/c^2)$$

DY Transversity TSA



2 σ negative **Transversity**
Results compatible with
Transversity TSA

[COMPASS, hep-ex/2312.17379]
to appear in Phys. Rev. Lett.

Conclusions

- **COMPASS probes 3-dimensional structure of nucleon**
- COMPASS SIDIS and Drell-Yan TSAs measurements represent a unique experimental input to study the universality of TMD PDFs

Drell-Yan TSAs

- 1 σ positive **Sivers TSA**
- **Pretzelosity TSA** found to be small and compatible with zero
- 2 σ negative **Transversity TSA**
- Results agree with theoretical predictions and consistent with analogous measurements for SIDIS

$$A_{DY} \propto PDF_N \otimes PDF_{\pi^-}$$

Transverse momentum weighted Drell-Yan TSA

- A way to overcome the convolution over intrinsic k_T
- A direct access to the k_T^2 -moments of TMD PDFs
- ~ 1 σ positive **Sivers WTSAs** compatible with DY TSA and SIDIS P_T -weighted TSA
- ~ 2 σ negative **Pretzelosity WTSAs** effect
- ~ 2 σ negative **Transversity WTSAs** consistent with TSAs

$$A_{DY}^W \propto PDF_N \times PDF_{\pi^-}$$

Prospects

- Analysis of a WTSAs ongoing, paper in preparation



Thank you for attention!

"All of physics is either impossible or trivial. It is impossible until you understand it, and then it becomes trivial."

Ernest Rutherford

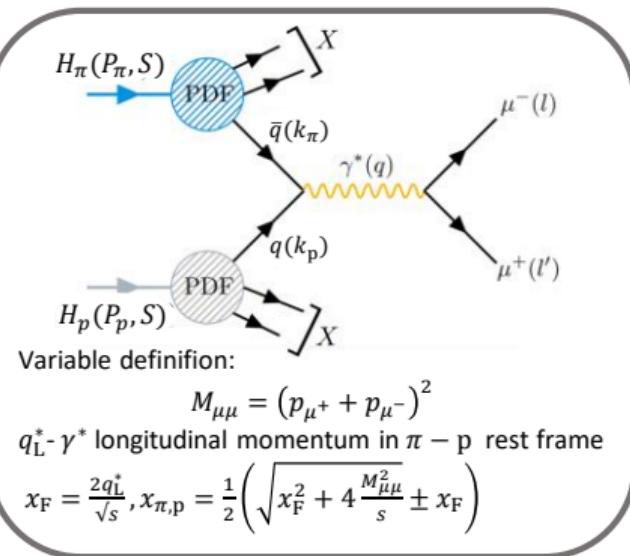
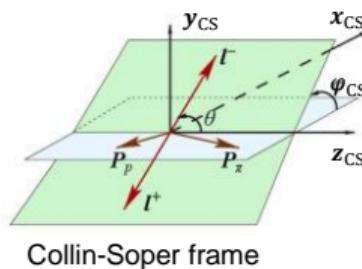
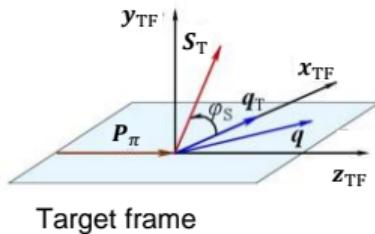
Backup slides



Single polarised Drell-Yan process

Cross-section, LO TMD approach for transversely polarised target in terms of structure functions:

$$\frac{d\sigma^{LO}}{dx_p dx_\pi d^2 q_T d\varphi_{CS} d(\cos\theta) d\varphi_S} = C_0 \left\{ \begin{array}{l} (1 + \cos^2 \theta) F_U^1 + \sin^2 \theta \cos 2\varphi F_U^{\cos 2\varphi_{CS}} \\ + |S_T| \left[\begin{array}{l} (1 + \cos^2 \theta) \sin(\varphi_S) F_T^{\sin \varphi_S} + \\ \sin^2 \theta \left(\begin{array}{l} \sin(2\varphi + \varphi_S) F_T^{\sin(2\varphi_{CS} + \varphi_S)} \\ + \sin(2\varphi - \varphi_S) F_T^{\sin(2\varphi_{CS} - \varphi_S)} \end{array} \right) \end{array} \right] \end{array} \right\}$$



Boer-Mulders

$$F_U^{\cos 2\varphi_{CS}} \propto h_{1,\pi}^{q\perp} \otimes h_{1,N}^{q\perp}$$

Sivers

$$F_T^{\sin \varphi_S} \propto f_{1,\pi}^q \otimes f_{1T,N}^{q\perp}$$

Pretzelosity

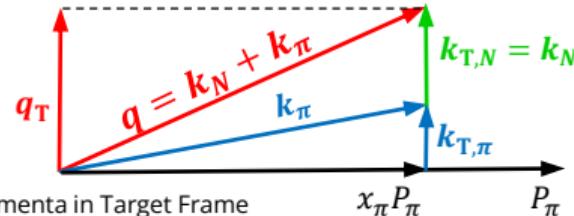
$$F_T^{\sin(2\varphi_{CS} + \varphi_S)} \propto h_{1,\pi}^{q\perp} \otimes h_{1T,N}^{q\perp}$$

Transversity

$$F_T^{\sin(2\varphi_{CS} - \varphi_S)} \propto h_{1,\pi}^{q\perp} \otimes h_{1,N}^{q\perp}$$

Single polarised Drell-Yan process

The convolution of TMD PDFs runs over the intrinsic transverse momenta k_T .



TMD PDFs are accessed through measurement of target spin dependent azimuthal asymmetries TSA.

Sivers Asymmetries $A_T^{\sin(\varphi_S)} = \frac{F_T^{\sin(\varphi_S)}}{F_U^1}$ **Sivers** for nucleon, number density for π^-

Pretzelosity Asymmetries $A_T^{\sin(2\varphi_{CS} + \varphi_S)} = \frac{F_T^{\sin(2\varphi_{CS} + \varphi_S)}}{2F_U^1}$ **Pretzelosity** for nucleon, Boer-Mulders for π^-

Transversity Asymmetries $A_T^{\sin(2\varphi_{CS} - \varphi_S)} = \frac{F_T^{\sin(2\varphi_{CS} - \varphi_S)}}{2F_U^1}$ **Transversity** for nucleon, Boer-Mulders for π^-

Backup: Single Polarised Drell-Yan Process

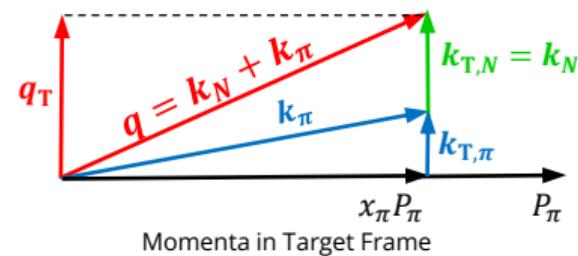
Each structure function can be written as a TMD PDF convolution over the intrinsic transverse momenta.

$$F_T^{\sin(\varphi_S)} = \mathcal{C} \left[\frac{\mathbf{q}_T \mathbf{k}_{T,N}}{q_T M_N} f_{1,\pi} \mathbf{f}_{1T,p}^\perp \right] \quad \text{Sivers for proton, number density for } \pi^-$$

$$F_T^{\sin(2\varphi_{CS} + \varphi_S)} = -\mathcal{C} \left[\frac{2(\mathbf{q}_T \mathbf{k}_{T,N})[2(\mathbf{q}_T \mathbf{k}_{T,\pi})(\mathbf{q}_T \mathbf{k}_{T,\pi}) - q_T^2 (\mathbf{k}_{T,N} \mathbf{k}_{T,\pi})] - q_T^2 k_{T,N}^2 (\mathbf{q}_T \mathbf{k}_{T,\pi})}{2q_T^3 M_N^2 M_\pi} \mathbf{h}_{1,\pi}^\perp \mathbf{h}_{1T,N}^\perp \right] \quad \text{Pretzelosity for proton, Boer–Mulders for } \pi^-$$

$$F_T^{\sin(2\varphi_{CS} - \varphi_S)} = -\mathcal{C} \left[\frac{\mathbf{q}_T \mathbf{k}_{T,\pi}}{q_T M_\pi} \mathbf{h}_{1,\pi}^\perp \mathbf{h}_{1N}^\perp \right] \quad \text{Transversity for proton, Boer–Mulders for } \pi^-$$

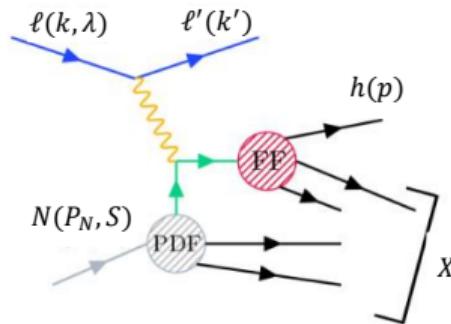
$$\mathcal{C}[w(\mathbf{k}_{T,\pi}; \mathbf{k}_{T,N}; \mathbf{q}_T) f_\pi f_N] = \frac{1}{N_c} \sum_q \left\{ e_q^2 \int d^2 \mathbf{k}_{T,\pi} d^2 \mathbf{k}_{T,N} \delta^{(2)}(\mathbf{q}_T - \mathbf{k}_{T,\pi} - \mathbf{k}_{T,N}) \right. \\ \times w(\mathbf{k}_{T,\pi}; \mathbf{k}_{T,N}; \mathbf{q}_T) \left[f_\pi^{\bar{q}}(x_\pi, k_{T,\pi}^2) f_N^q(x_N, k_{T,N}^2) + f_\pi^q(x_\pi, k_{T,\pi}^2) f_p^{\bar{q}}(x_p, k_{T,p}^2) \right] \left. \right\}$$



TMD PDFs can be accessed through measurement of target spin (in)dependent azimuthal asymmetries

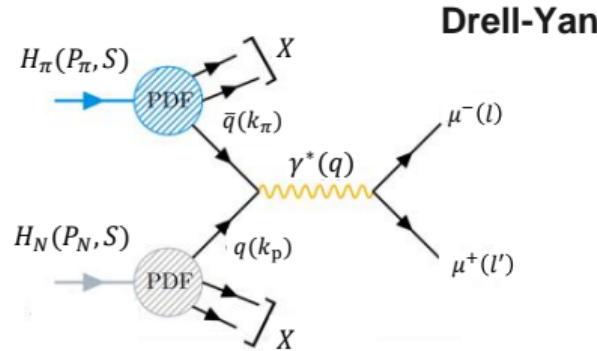
$$A_U^{\cos(\varphi_{CS})} = \frac{F_U^{\cos(\varphi_{CS})}}{F_U^1} \quad A_T^{\sin(\varphi_S)} = \frac{F_T^{\sin(\varphi_S)}}{F_U^1} \quad A_T^{\sin(2\varphi_{CS} + \varphi_S)} = \frac{F_T^{\sin(2\varphi_{CS} + \varphi_S)}}{2F_U^1} \quad A_T^{\sin(2\varphi_{CS} - \varphi_S)} = \frac{F_T^{\sin(2\varphi_{CS} - \varphi_S)}}{2F_U^1}$$

SIDIS



$$A_{SIDIS} \propto PDF_N \otimes FF$$

SIDIS vs. DY



$$A_{DY} \propto PDF_N \otimes PDF_{\pi^-}$$

$$\frac{d\sigma_{SIDIS}^{LO}}{dx dy dz d^2 p_T d\phi_h d\phi_S} \propto \left\{ + |S_T| \begin{bmatrix} 1 + \cos(2\phi_h) \varepsilon A_{UU}^{\cos(2\phi_h)} \\ \sin(\phi_h - \phi_S) A_{UT}^{\sin(\phi_h - \phi_S)} + \\ \sin(\phi_h + \phi_S) \varepsilon A_{UT}^{\sin(\phi_h + \phi_S)} + \\ \sin(3\phi_h - \phi_S) \varepsilon A_{UT}^{\sin(3\phi_h - \phi_S)} + \end{bmatrix} \right\}$$

$$\frac{d\sigma_{DY}^{LO}}{d\Omega d^4 q} \propto \left\{ + |S_T| \begin{bmatrix} 1 + D_{\sin^2 \theta} \cos(2\varphi_{CS}) A_U^{\cos(2\varphi_{CS})} \\ \sin(\varphi_S) A_T^{\sin(\varphi_S)} + \\ D_{\sin^2 \theta} \begin{pmatrix} \sin(2\varphi_{CS} + \varphi_S) A_T^{\sin(2\varphi_{CS} + \varphi_S)} \\ \sin(2\varphi_{CS} - \varphi_S) A_T^{\sin(2\varphi_{CS} - \varphi_S)} \end{pmatrix} \end{bmatrix} \right\}$$

$$\begin{aligned} A_{UU}^{\cos(2\phi_h)} &\propto h_{1,N}^{q\perp} \otimes H_{1q}^{\perp h} \\ A_{UT}^{\sin(\phi_h - \phi_S)} &\propto f_{1T,N}^{q\perp} \otimes D_{1q}^h \\ A_{UT}^{\sin(3\phi_h - \phi_S)} &\propto h_{1T,N}^{q\perp} \otimes H_{1q}^{\perp h} \\ A_{UT}^{\sin(\phi_h + \phi_S)} &\propto h_{1T,N}^q \otimes H_{1q}^{\perp h} \end{aligned}$$

Boer-Mulders

Sivers

Pretzelosity

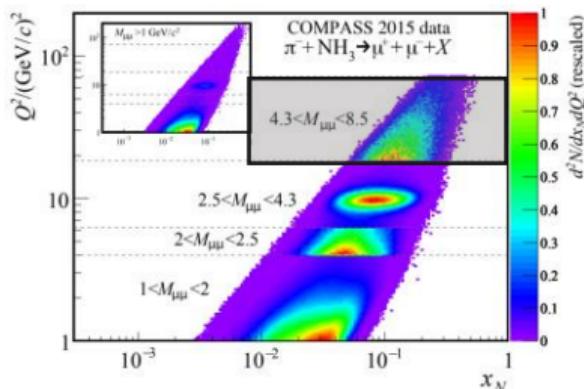
Transversity

$$\begin{aligned} A_U^{\cos(2\varphi_{CS})} &\propto h_{1,\pi}^{q\perp} \otimes h_{1,N}^{q\perp} \\ A_T^{\sin(\varphi_S)} &\propto f_{1,\pi}^q \otimes f_{1T,N}^{q\perp} \\ A_T^{\sin(2\varphi_{CS} + \varphi_S)} &\propto h_{1,\pi}^{q\perp} \otimes h_{1T,N}^{q\perp} \\ A_T^{\sin(2\varphi_{CS} - \varphi_S)} &\propto h_{1,\pi}^{q\perp} \otimes h_{1,N}^q \end{aligned}$$

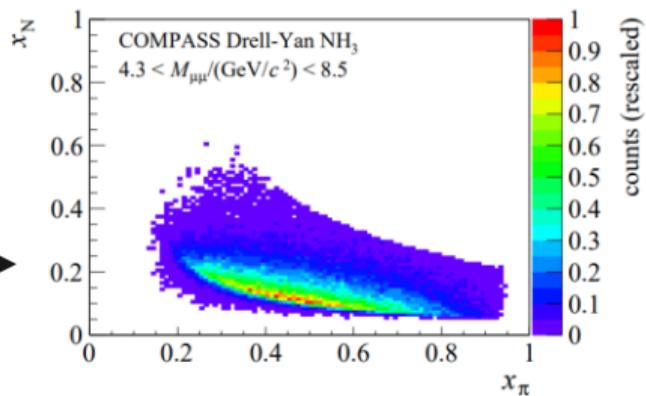
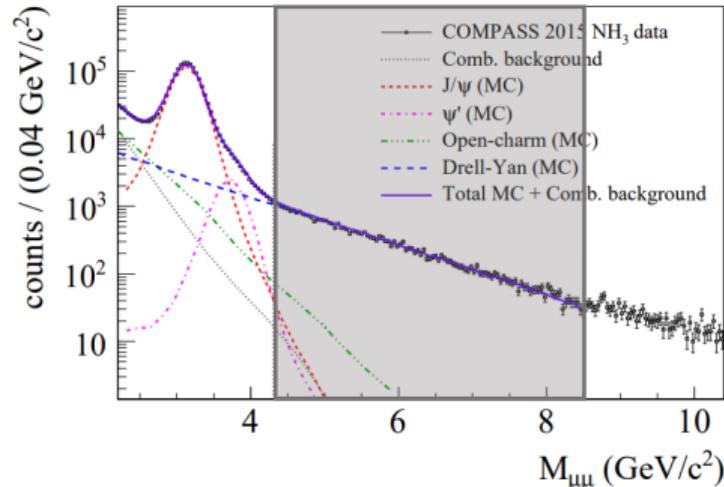
Drell-Yan measurement at COMPASS

The dimuon mass range $4.3 < M_{\mu\mu}/(\text{GeV}/c^2) < 8.5$ is 96% pure Drell-Yan.

- Low background
- Valence region



Valence region of x_N and x_π in a given $Q^2 - x$ kinematic coverage



Weighted TSAs in Drell-Yan

The convolution cannot be resolved without assumptions about the dependence of the TMD PDF on the intrinsic transverse momentum.

Weighting with powers of the transverse momentum allows to avoid assumptions on k_T .

Asymmetries in terms of structure functions:

$$A_T^{\sin\Phi W_\Phi} = \frac{\int d^2\mathbf{q}_T W_\Phi F_T^{\sin\Phi}}{\int d^2\mathbf{q}_T F_U^1},$$

where W_Φ is weight for $\Phi = \varphi_S, 2\varphi_{CS} + \varphi_S, 2\varphi_{CS} - \varphi_S$.

The n-th moment of a TMD PDF of a pion or proton:

$$f^{(n)}(x) = \int d^2\mathbf{k}_T \left(\frac{k_T^2}{2M^2}\right)^n f(x, k_T^2).$$

