

Partial-Wave Analysis of the $K_S^0 K^-$ Final State: Ambiguities and Physics Results

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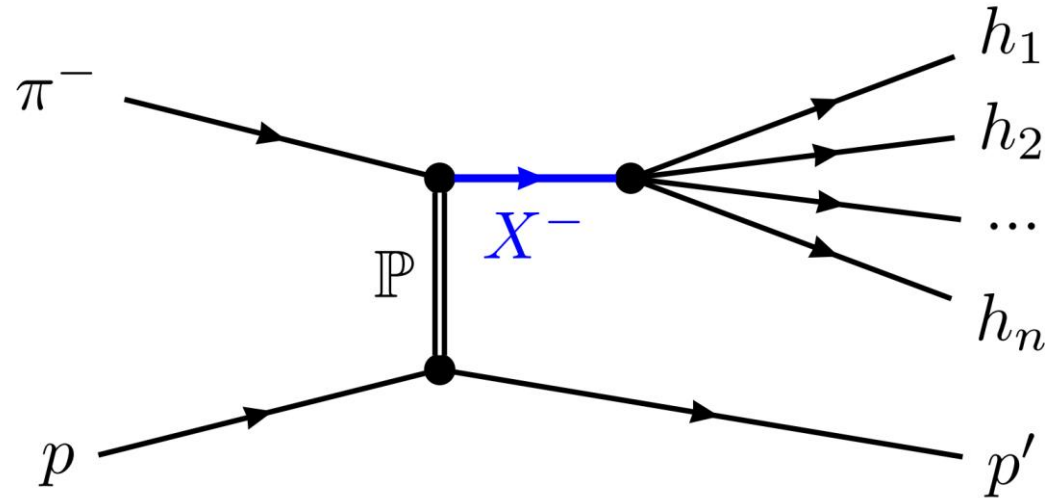
PWA/ATHOS 2024

May 28th, 2024



Excited Light Mesons at COMPASS

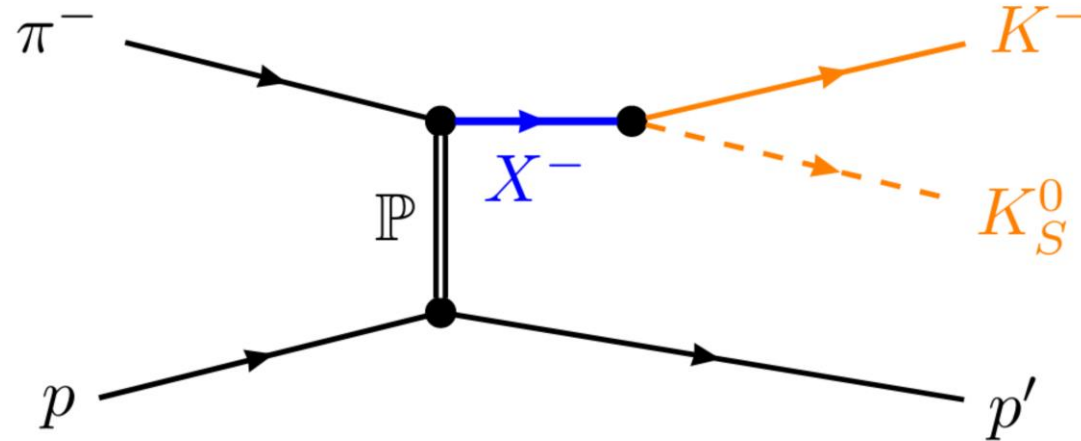
- Inelastic scattering reactions of high-energetic meson beam



- Strong interaction (Pomeron exchange) between beam meson and target proton
- Intermediate hadronic resonances X^- are created, then decay into n -body final state
→ wide range of allowed (spin) quantum numbers
- Final-state particles measured in the COMPASS spectrometer

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The COMPASS Experiment



Large-acceptance magnetic spectrometer @ CERN-SPS

Beam:

- Secondary hadrons (π^- , K^-) at 190 GeV/c
- produced via primary proton beam from SPS



The COMPASS Experiment



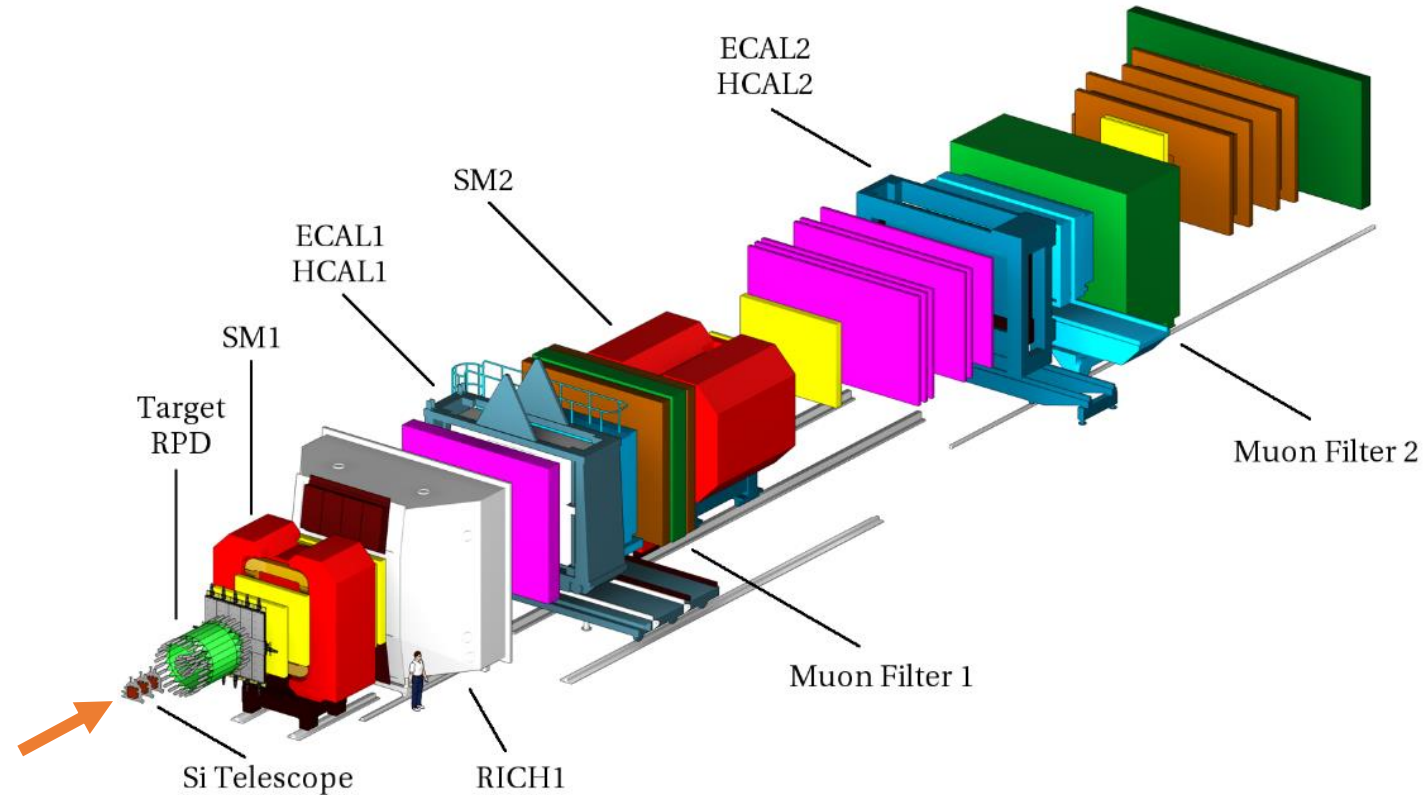
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Spectrometer:

- Liquid-hydrogen target
- Two-stage spectrometer setup around two dipole magnets SM1/2



From COMPASS Collab., The COMPASS Setup for Physics with Hadron Beams (Nucl. Instrum. Methods Phys. Res. A 779 (2014), pp. 69–115)

Quantum Numbers of the $K_S^0 K^-$ Final State

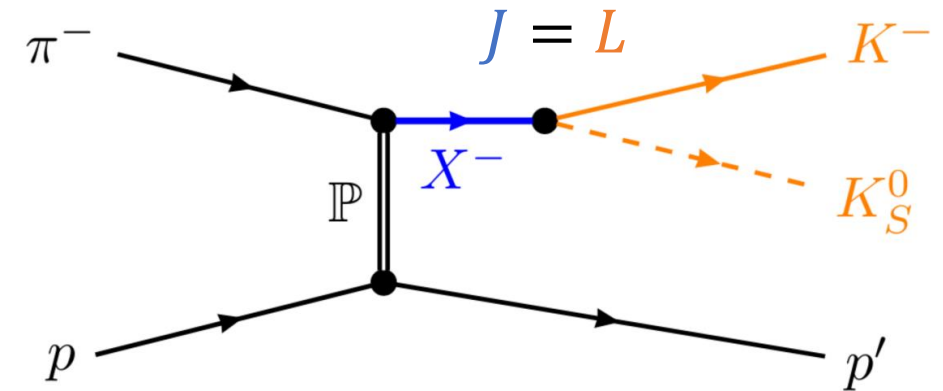
- Quantum numbers of a $K_S^0 K^-$ system

$$I = 1$$

$$G = (-1)^{L+1}$$

$$P = (-1)^L \quad C = G(-1)^I = (-1)^L$$

all given
by $L!$



$$I^G J^{PC} = 1^+ 1^{--}, 1^- 2^{++}, 1^+ 3^{--}, 1^- 4^{++}, \dots$$

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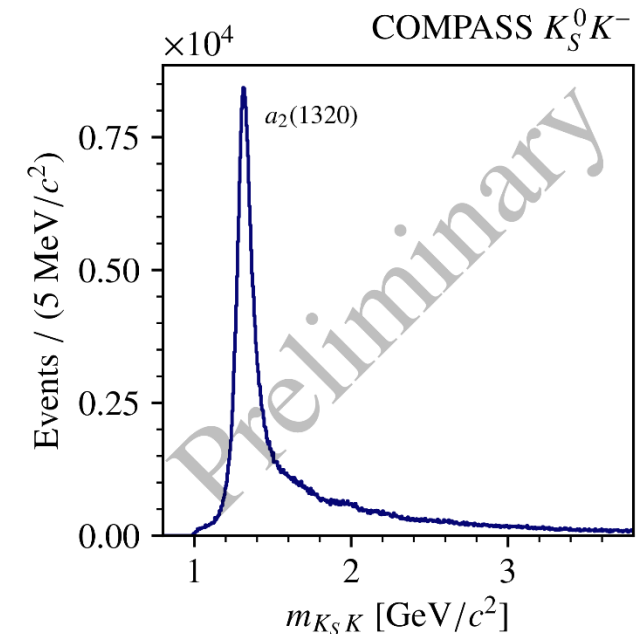
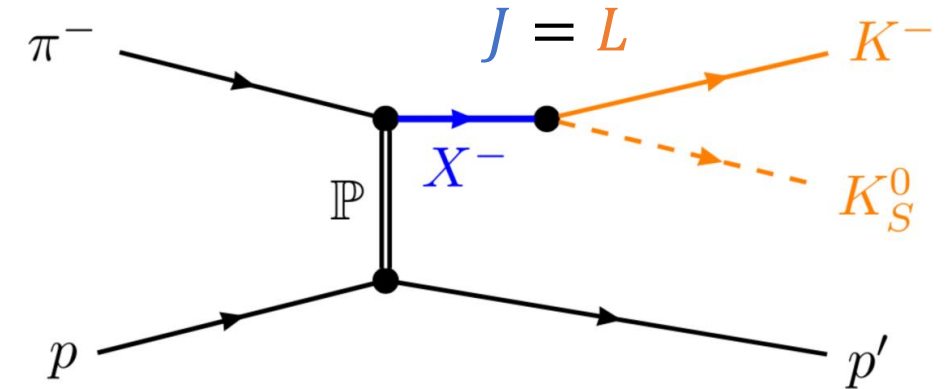
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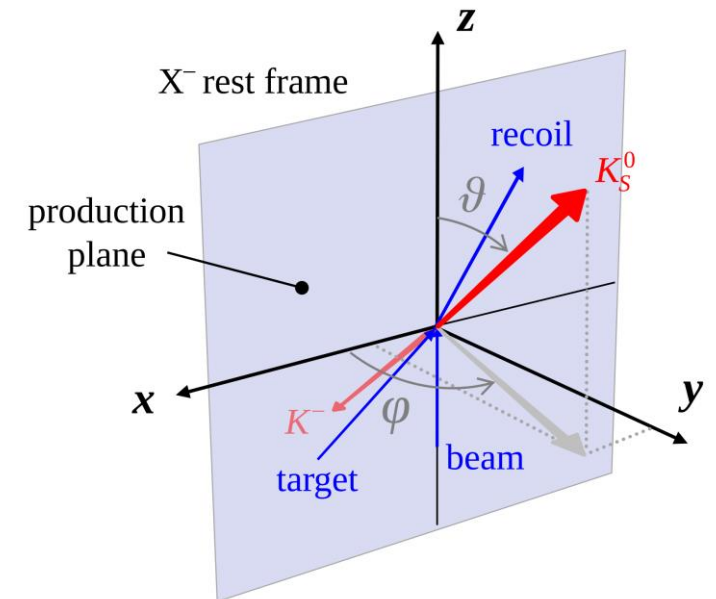
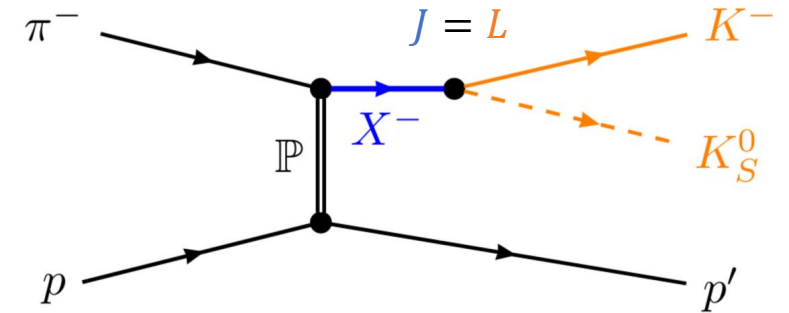


$$I^G J^{PC} = 1^+ 1^{--}, 1^- 2^{++}, 1^+ 3^{--}, 1^- 4^{++}, \dots$$

Partial-Wave Analysis Procedure

$$A(m_{KK}, t'; \theta, \phi) = \sum_{J,M} T_{J,M}(m_{KK}, t') \psi_{J,M}(\theta, \phi)$$

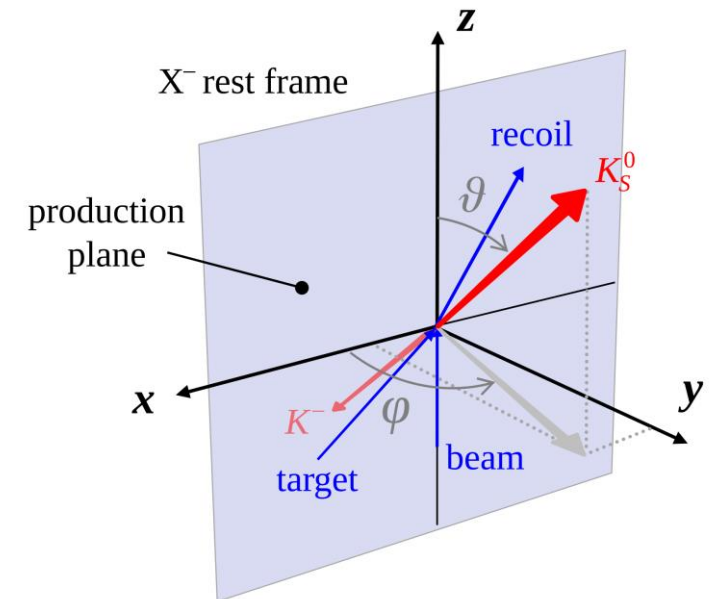
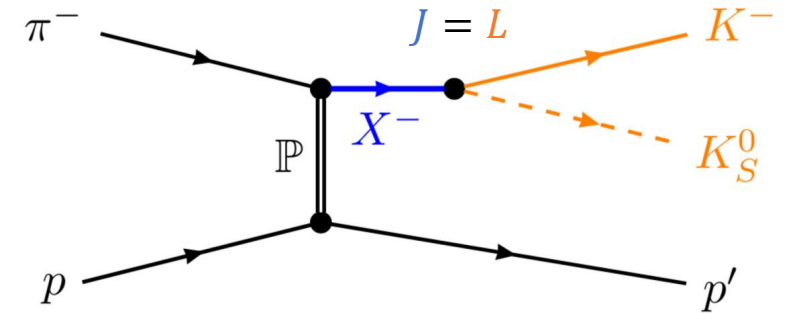
- Decompose total process amplitude into **partial waves**
 - Depend on spin J and spin-projection M
 - Other quantum numbers fixed by J, M



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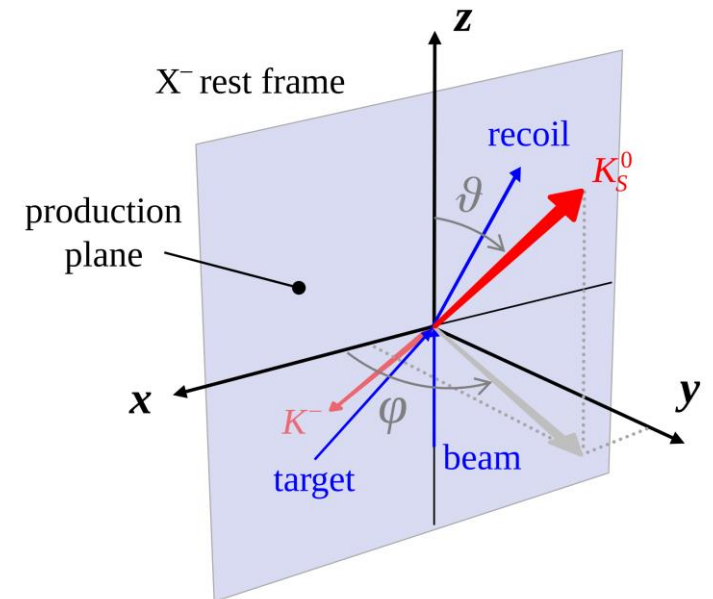
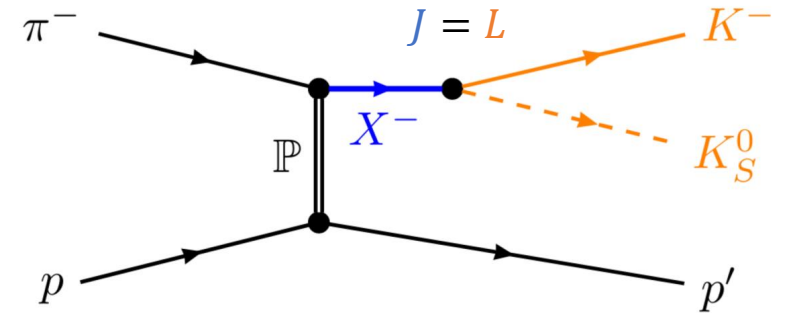
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- Partial-wave amplitudes split into
 - production and propagation $\rightarrow T_{J,M}(m_{KK}, t')$
- and
 - decay of $X^- \rightarrow \psi_{J,M}(\theta, \phi) = Y_J^M(\theta, \phi)$



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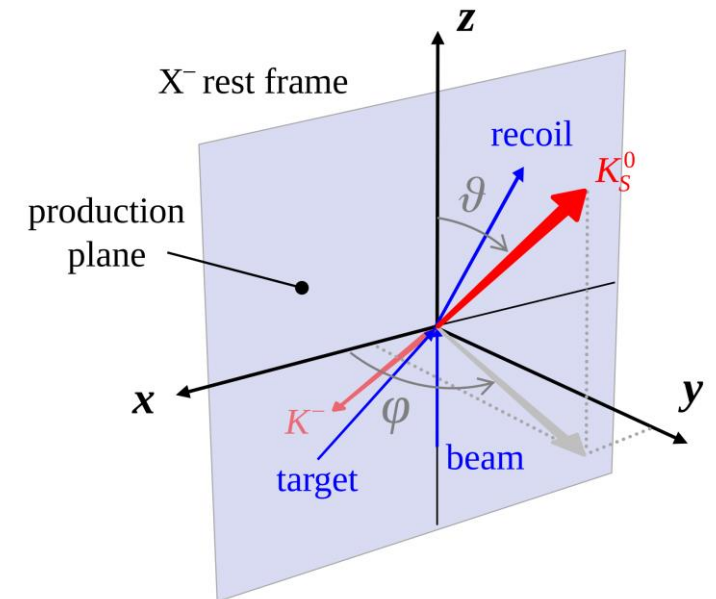
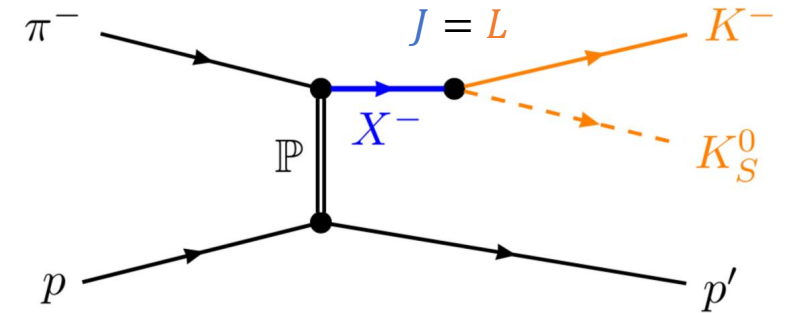
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Partial-Wave Analysis Procedure

$$\frac{dN}{d\Phi(\theta, \phi)} \sim I(m_{KK}, t'; \theta, \phi) = \left| \sum_{J,M} T_{J,M}(m_{KK}, t') Y_J^M(\theta, \phi) \right|^2$$

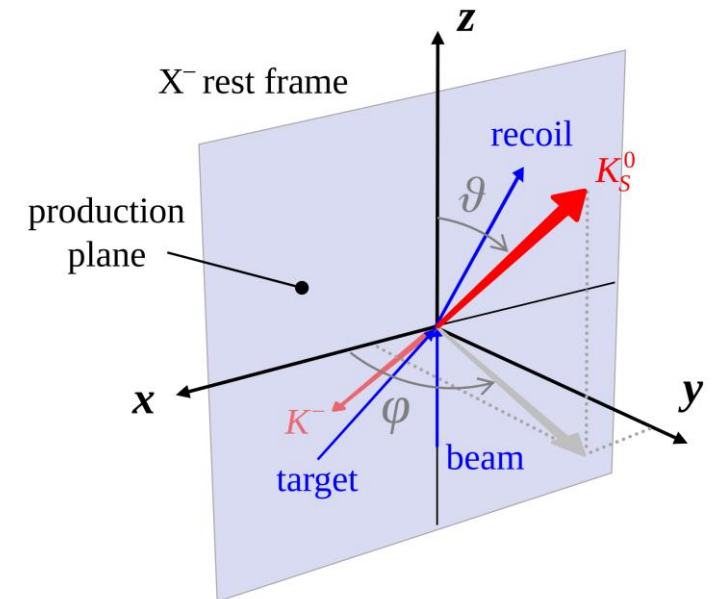
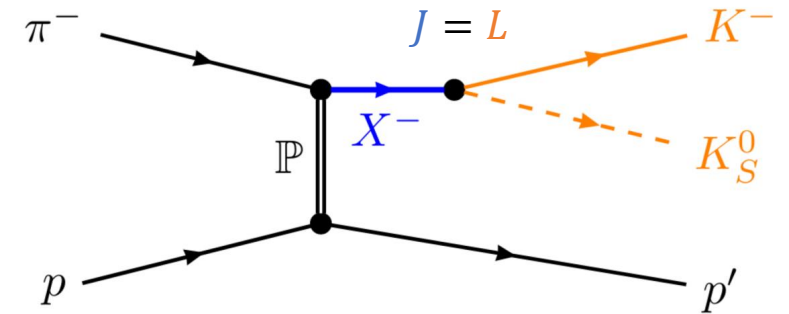
- Decompose total process amplitude into **partial waves**
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Ambiguities in the Partial-Wave Decomposition

For any final state with **two spinless** particles ($\pi\pi, KK, \eta\pi, \dots$):

- Decomposition of intensity into $\{T_{J,M}\}$ is not **unique**

→ **Several sets of $\{T_{J,M}\}$** lead to the **same $I(\theta, \phi)$** in each (m_{KK}, t') bin

$$I(\theta, \phi) = \left| \sum_{J,M} T_{J,M}^{(1)} Y_J^M(\theta, \phi) \right|^2 = \left| \sum_{J,M} T_{J,M}^{(2)} Y_J^M(\theta, \phi) \right|^2$$

- The fit cannot distinguish between the **mathematically equivalent** solutions!

Ambiguities in the Partial-Wave Decomposition

$$I(\theta, \phi) = \left| \sum_{J,M} T_{J,M} Y_J^M(\theta, \phi) \right|^2$$

Assume strong dominance of $|M| = 1$ *

- Pomeron exchange dominant $\rightarrow M \neq 0$
- Higher $|M|$ suppressed

*using reflectivity basis for ψ_{JM} :
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Ambiguities in the Partial-Wave Decomposition

$$I(\theta, \phi) = \left| \sum_J T_J Y_J^1(\theta, \phi) \right|^2$$

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$$a(\theta) = \sum_{j=0}^{J_{\max}-1} c_j(\{T_J\}) \tan^{2j}(\theta) = c(\{T_J\}) \prod_{k=1}^{J_{\max}-1} (\tan^2(\theta) - r_k(\{T_J\}))$$

root decomposition
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$$= c^2 \prod_{k=1}^{J_{\max}-1} |\tan^2(\theta) - r_k|^2 |\sin \phi|^2 = c^2 \prod_{k=1}^{J_{\max}-1} |\tan^2(\theta) - r_k^*|^2 |\sin \phi|^2$$

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$$\{T_J'\} \neq \{T_J\}$$

$$\{c_j'\}$$

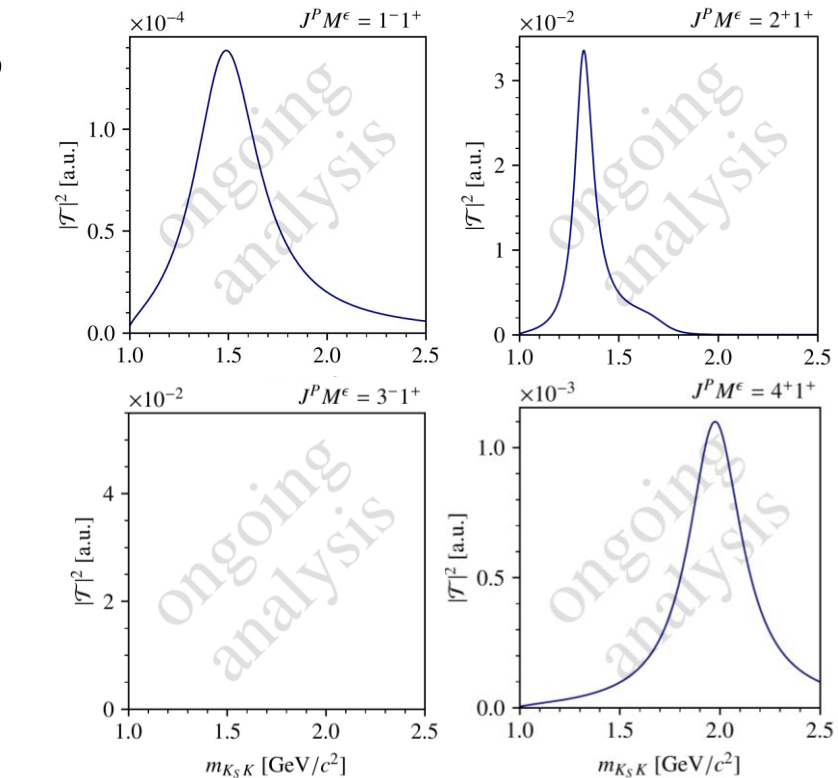
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Study of the Ambiguities

I. Continuous amplitude model

How do the ambiguous solutions look like (**continuity, signals, ...**)?

- Create an amplitude model for four partial waves
- Sample points in m_{KK} and calculate ambiguous solutions



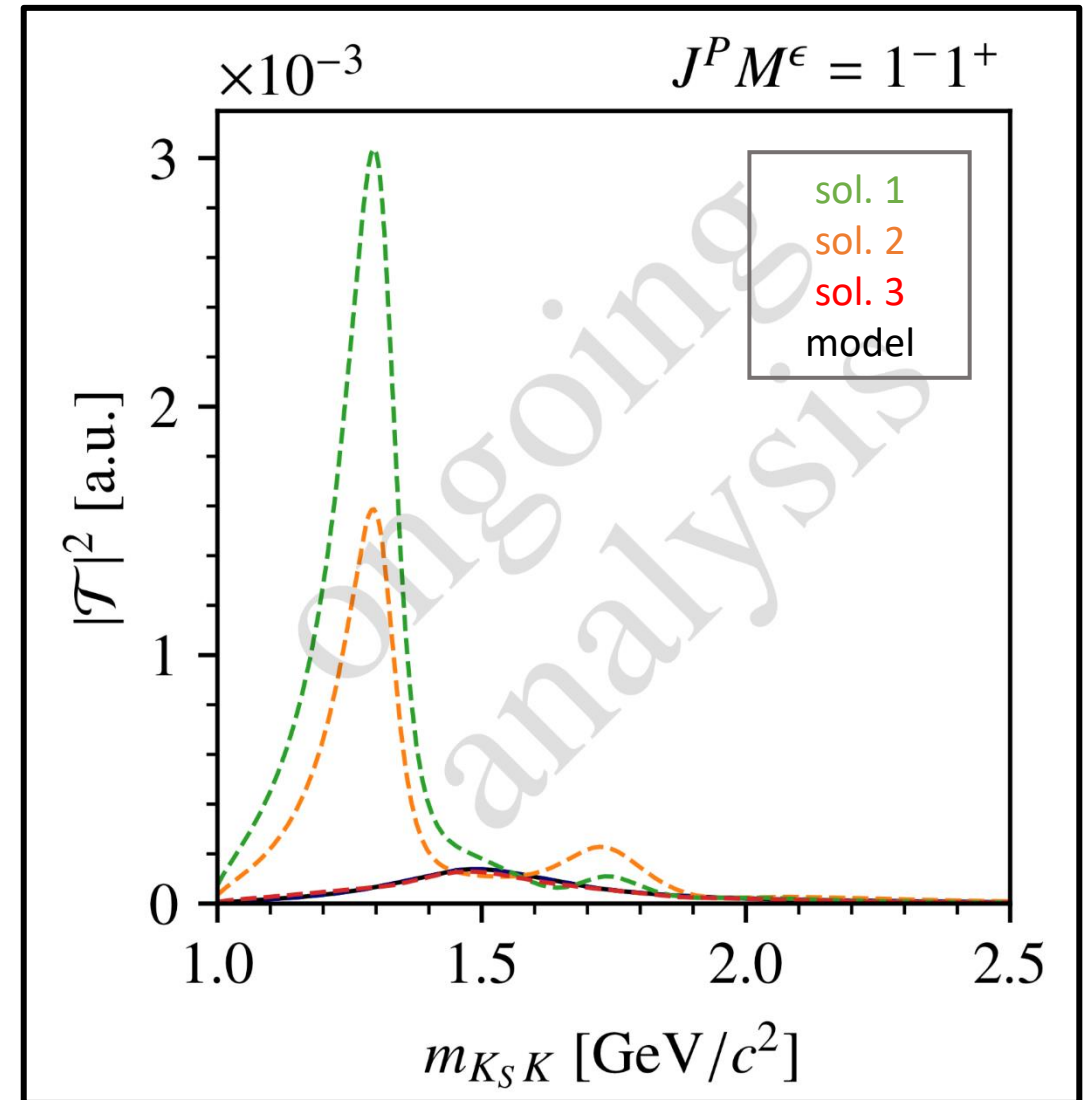
Continuous Amplitude Model

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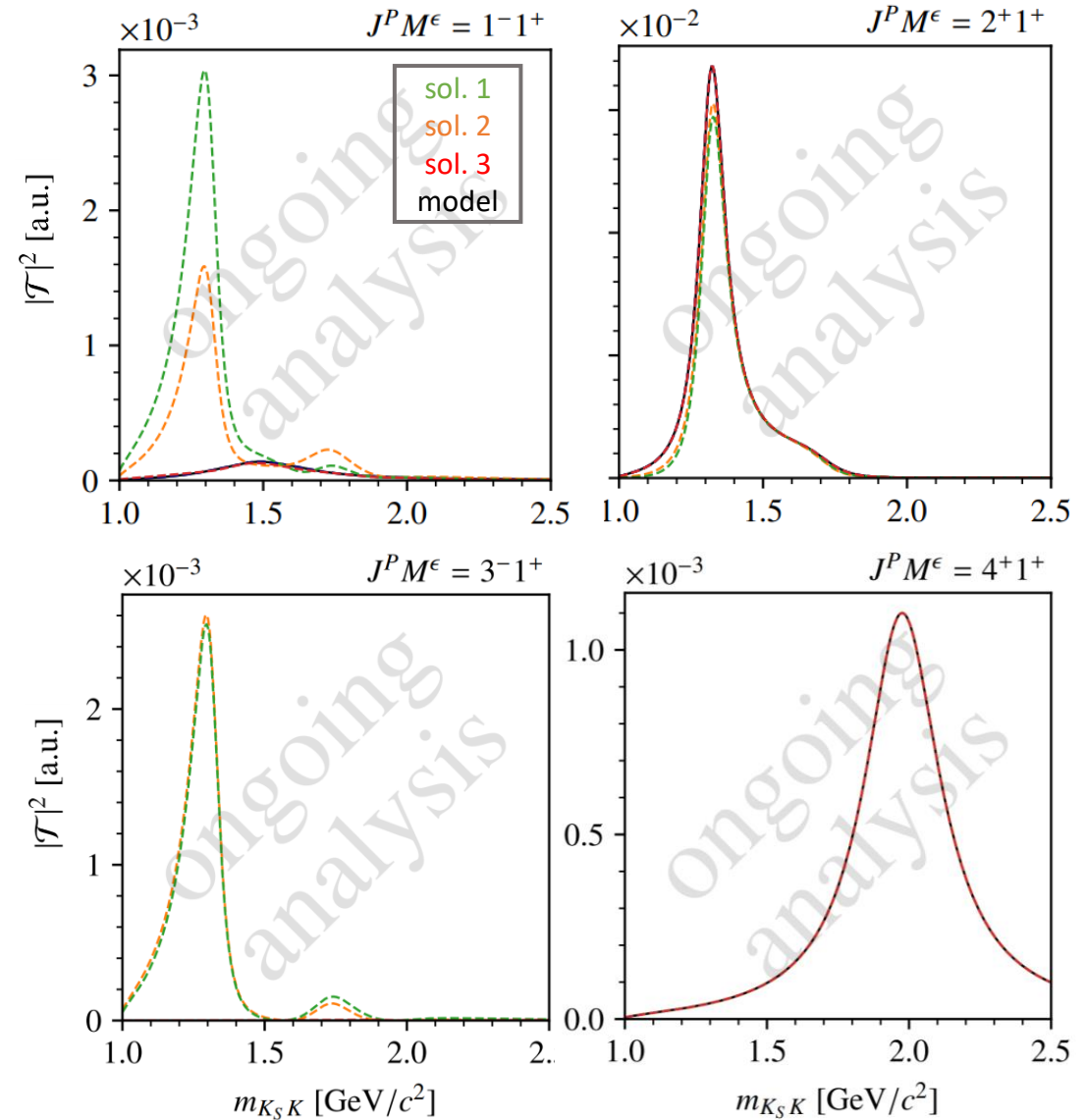
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Partial-Wave Decomposition Fits on Pseudodata

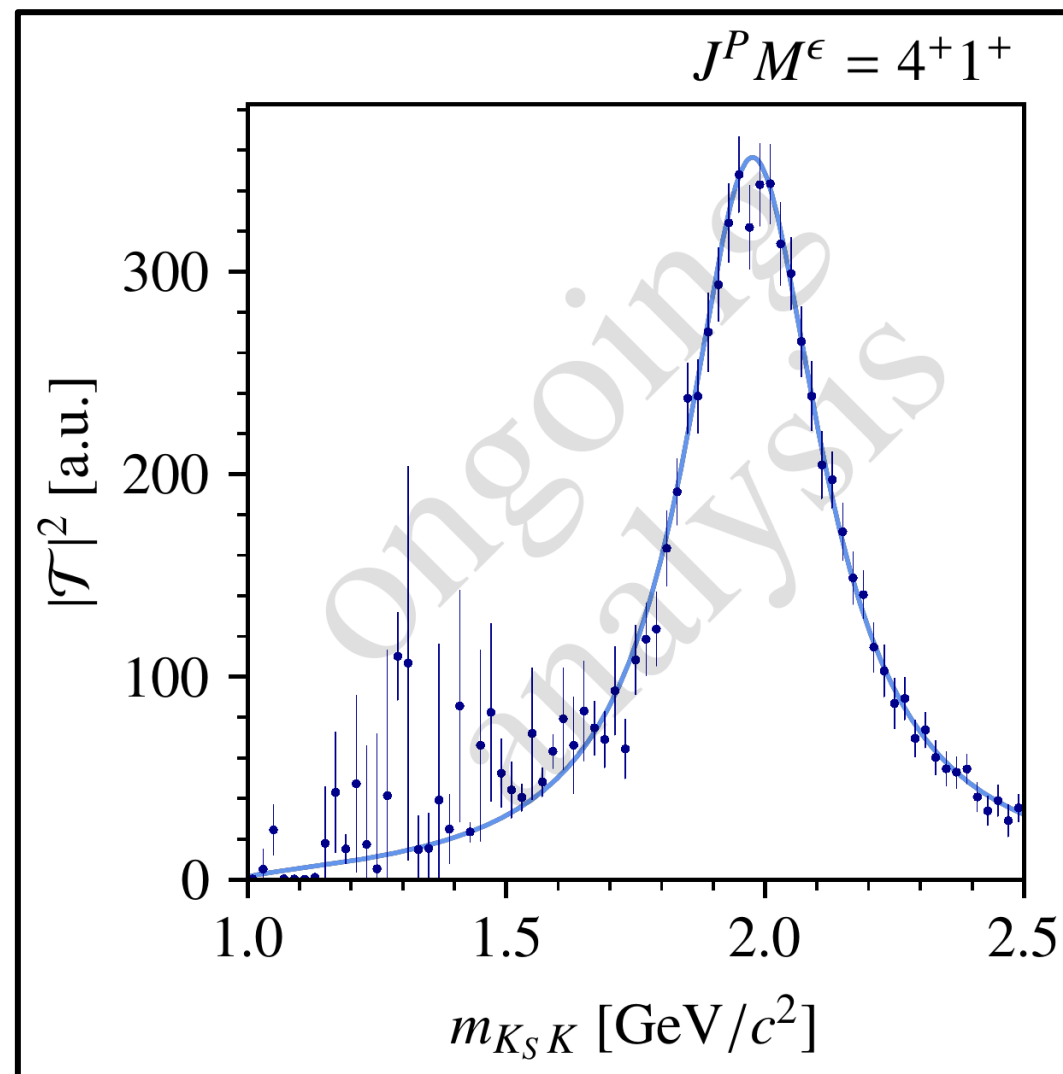
II. Finite pseudo-data

- reality: **finite data** and **amplitudes unknown**
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- 3000 attempts with **random** start values

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- Overall, amplitude values found by the fit follow the calculated distributions
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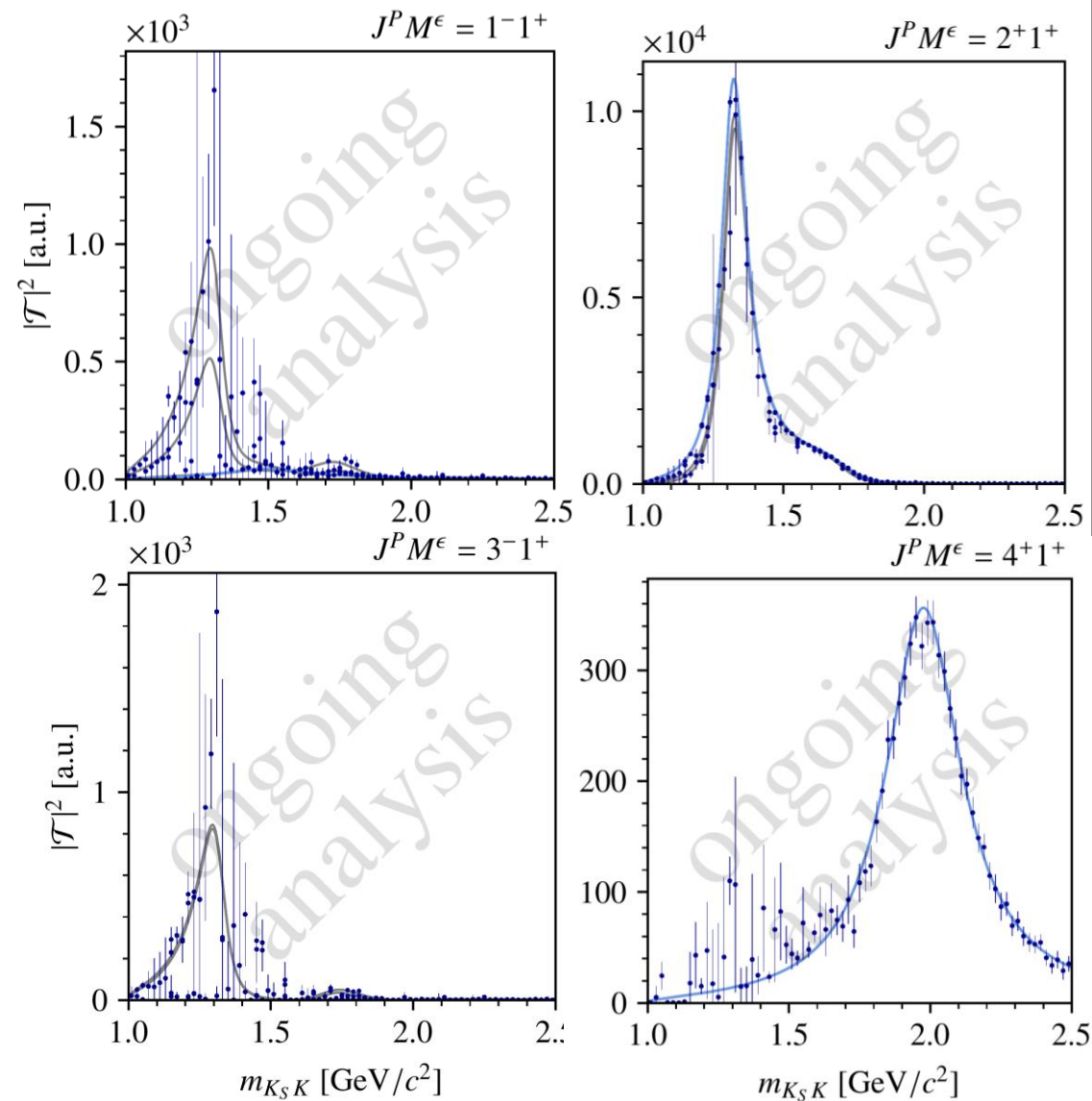


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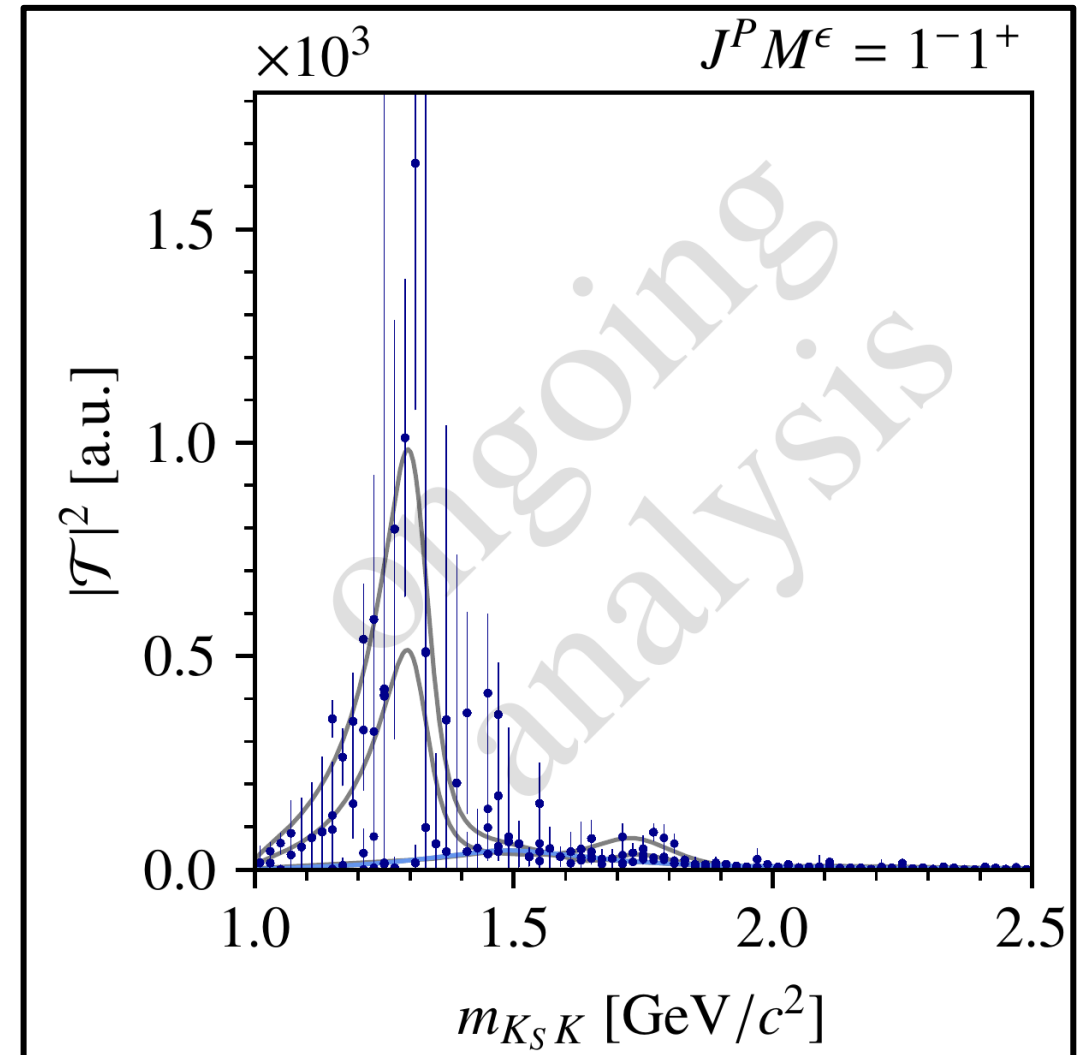
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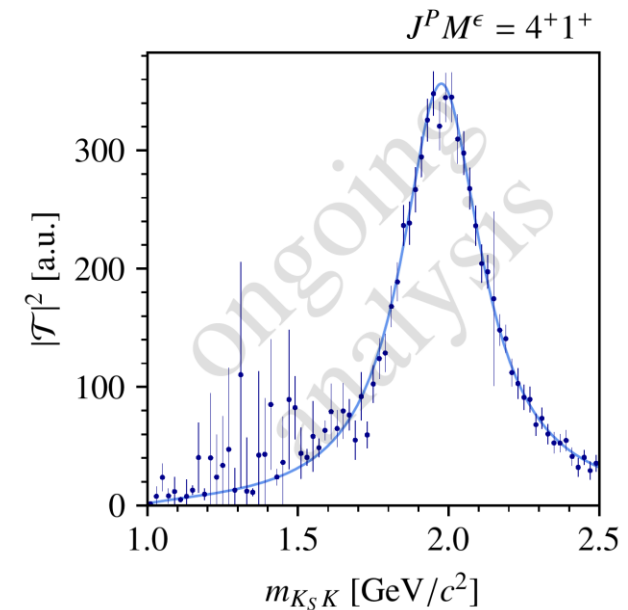
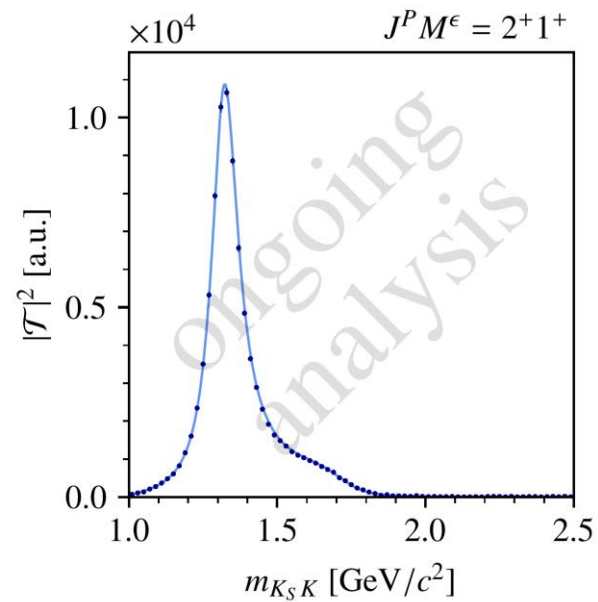
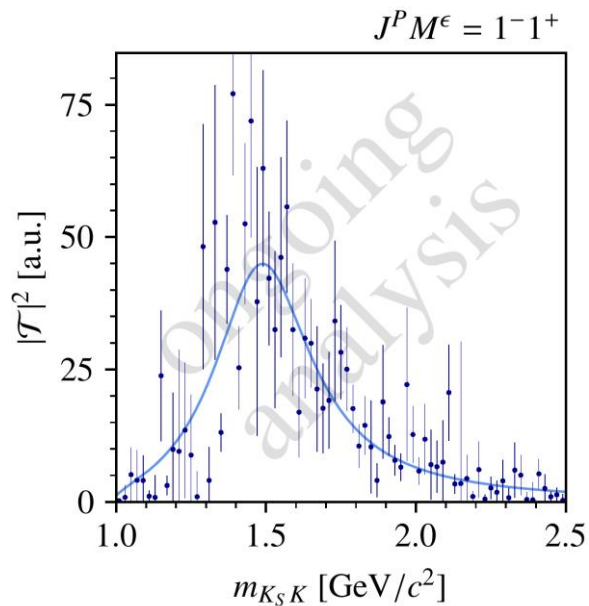


Reducing the Ambiguities

- Intensity of highest-spin wave is unaffected by ambiguities

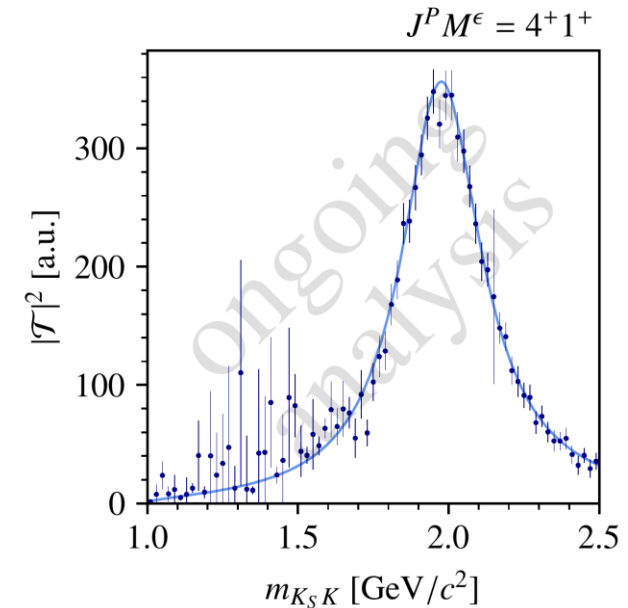
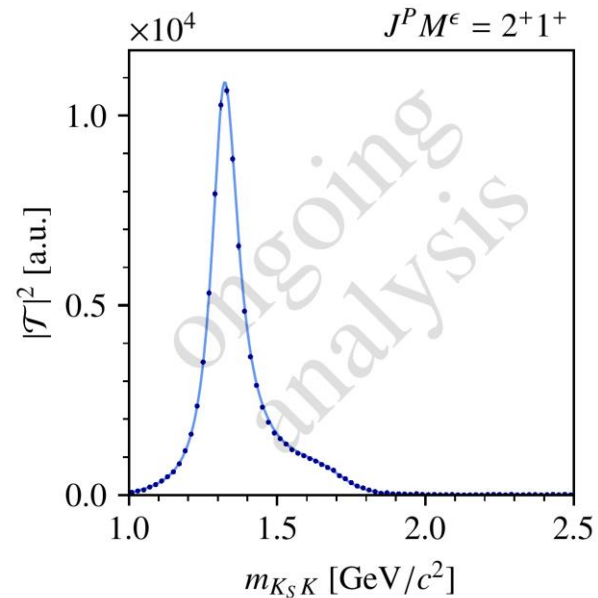
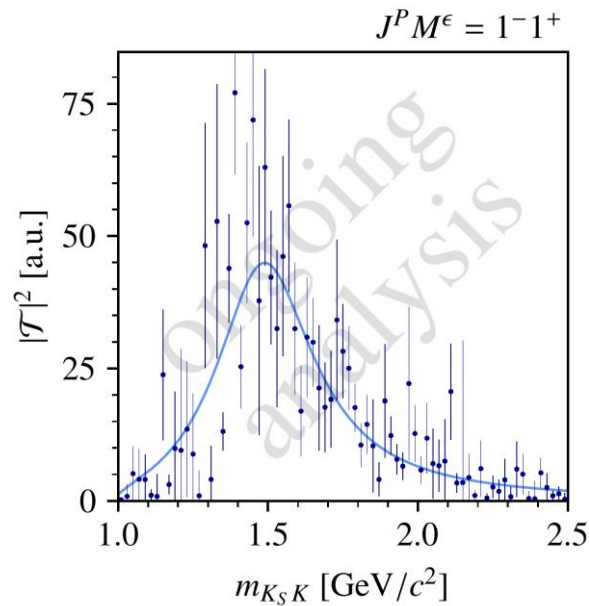
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- Remove one wave with $J < J_{\max}$ → **resolves ambiguities**



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At COMPASS: **odd-spin waves strongly suppressed!**

→ assume no ambiguities in the decomposition!

- Except: Two solutions in the phases of the partial waves

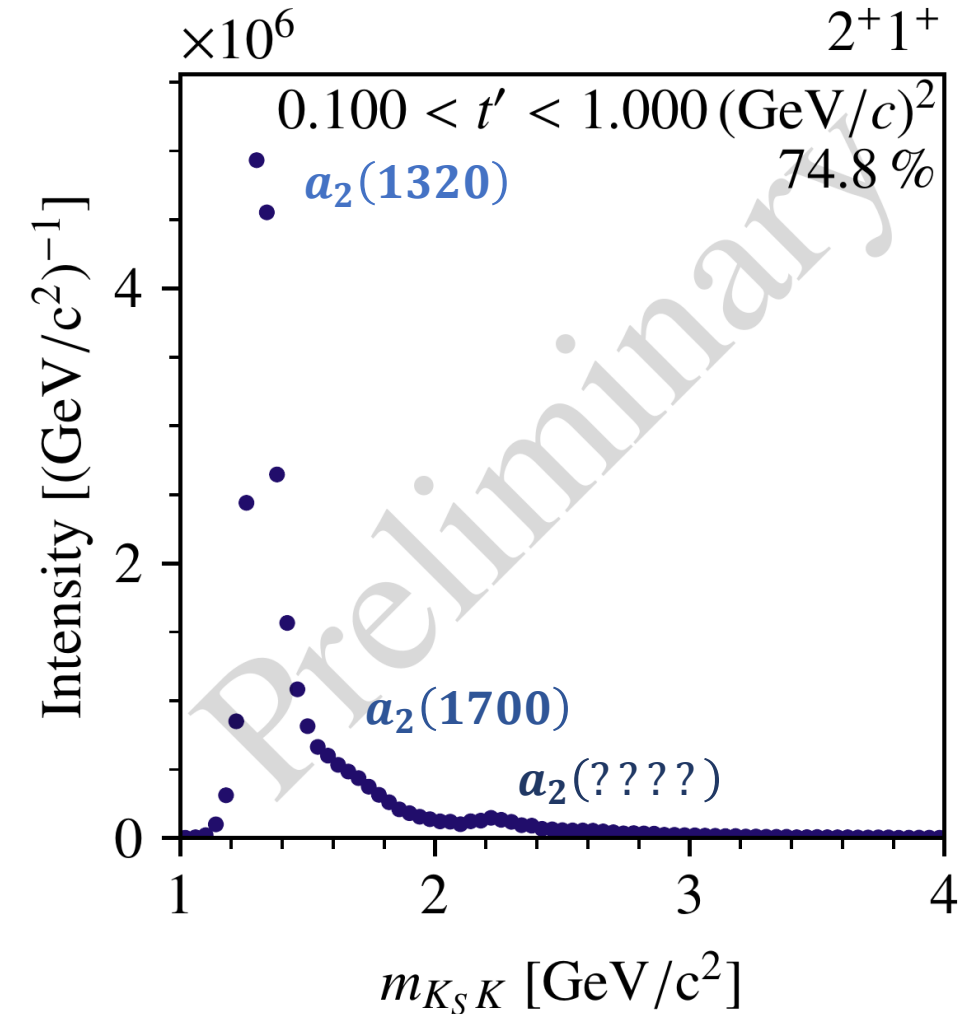
$$I(\theta, \phi) = \left| \sum_J T_J Y_J^1(\theta, \phi) \right|^2 = \left| \sum_J T_J^* Y_J^1(\theta, \phi) \right|^2$$

Partial-Wave Analysis of the $K_S^0 K^-$ Final State at COMPASS: Physics Results

J^{PC}	M^ϵ	Resonances
2^{++}	1^+	a_2 -like
	2^+	
4^{++}	1^+	a_4 -like
6^{++}	1^+	a_6 -like
(1^{--})	1^+	ρ -like) → later

Previous study by . E. Cleland et al. "Resonance Production in the Reaction $\pi^\pm p \rightarrow K0K^\pm p$ at 30 GeV/c and 50 GeV/c". In: Nucl. Phys. B 208 (1982)

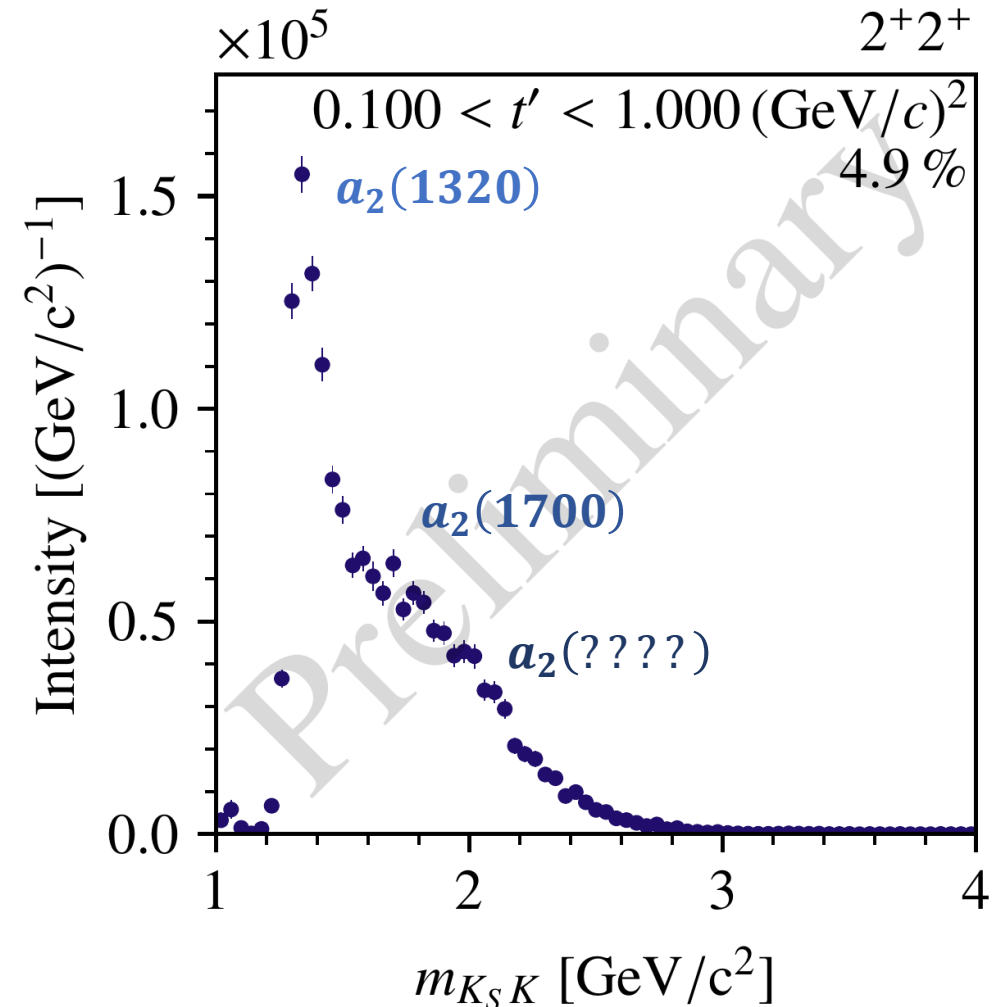
Results: Intensities and Phases



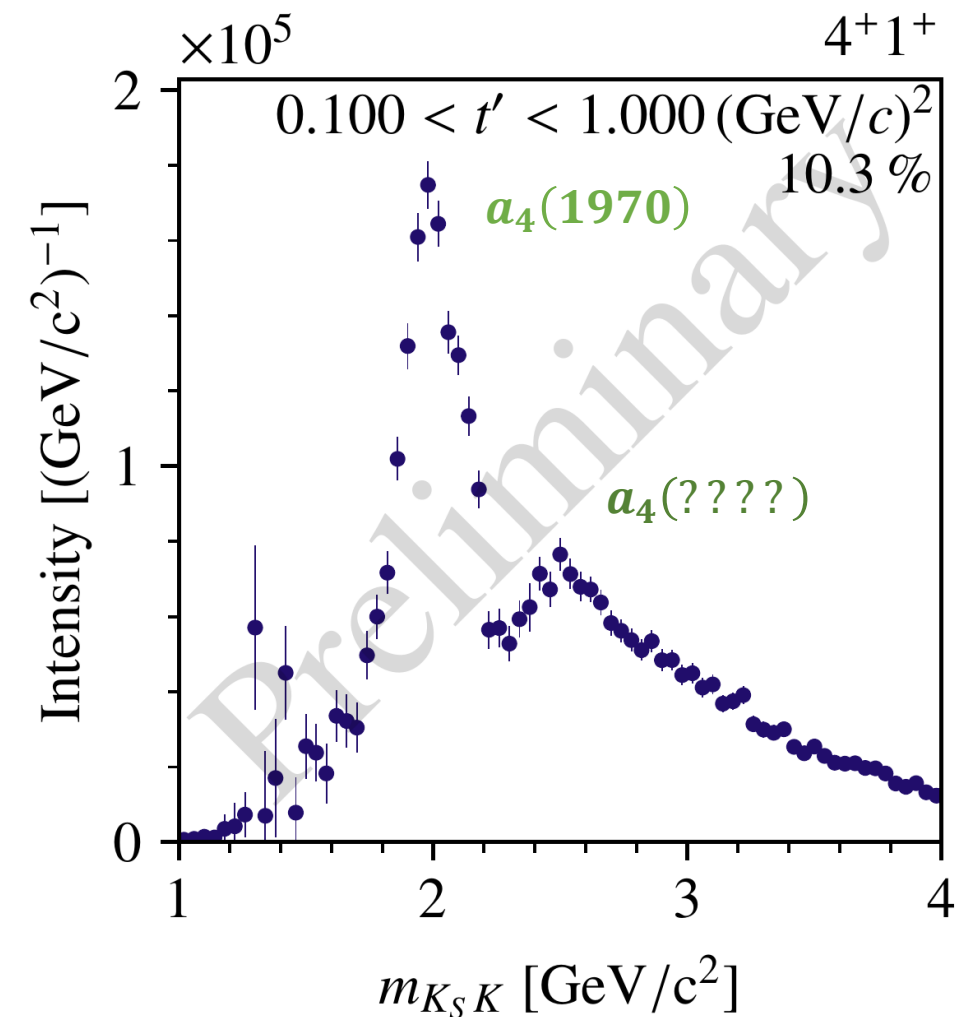
$a_2(1320)$
 $m_{\text{PDG}} = 1.32 \text{ GeV}$
 $\Gamma_{\text{PDG}} = 0.10 \text{ GeV}$

$a_2(1700)$
 $m_{\text{PDG}} = 1.71 \text{ GeV}$
 $\Gamma_{\text{PDG}} = 0.38 \text{ GeV}$

$a_2(????)$

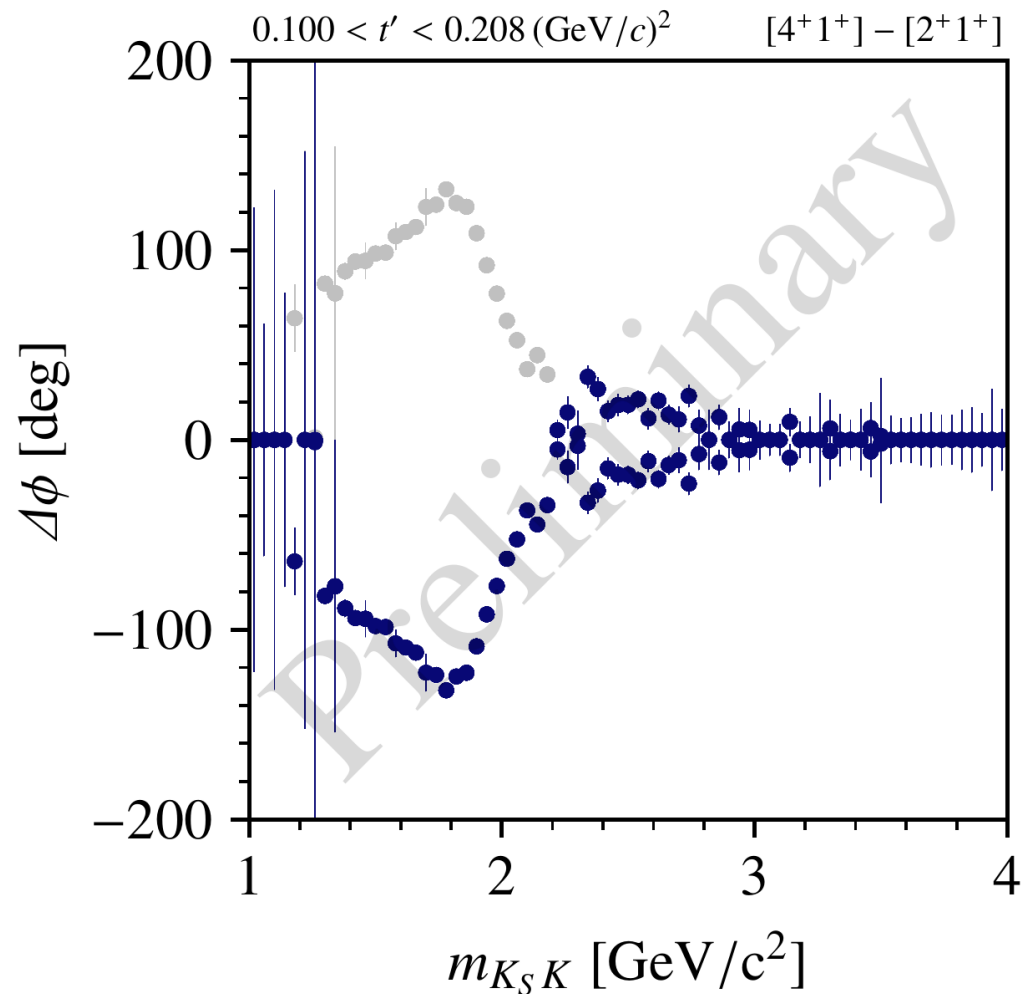


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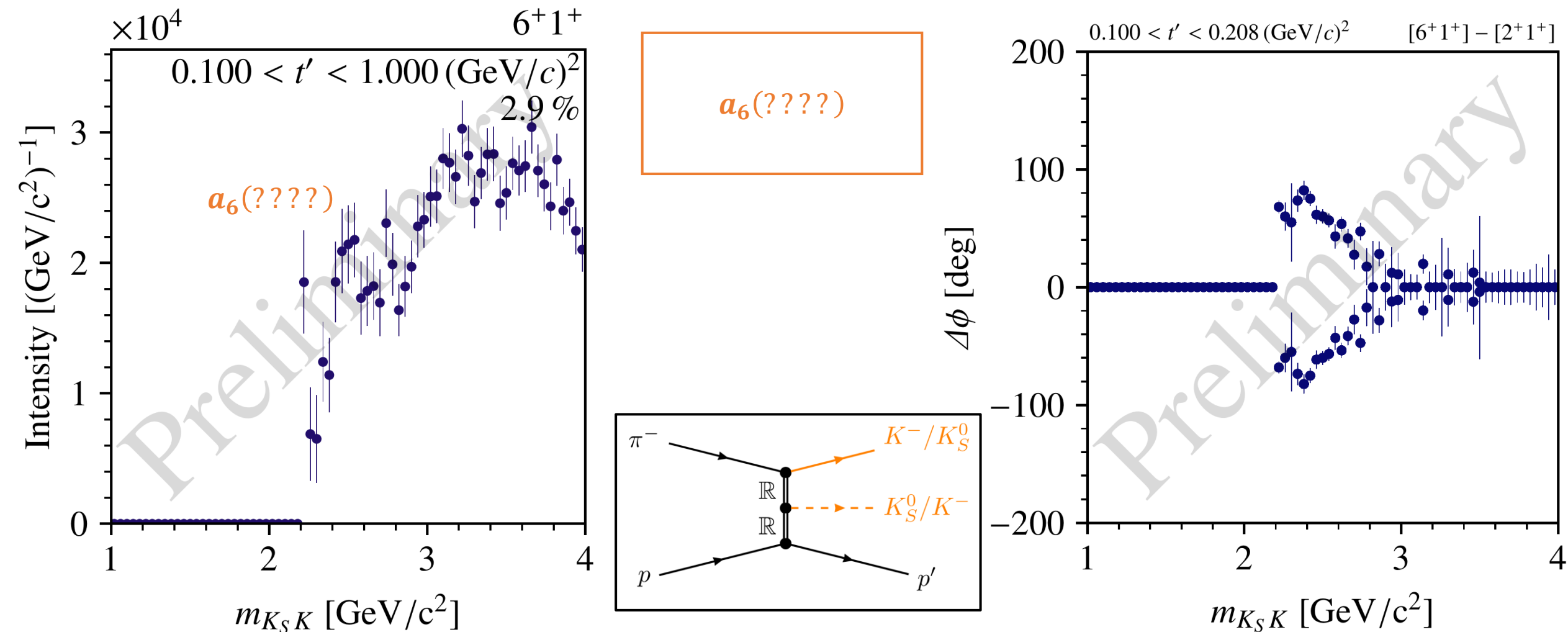


$a_4(1970)$
 $m_{\text{PDG}} = 1.97 \text{ GeV}$
 $\Gamma_{\text{PDG}} = 0.32 \text{ GeV}$

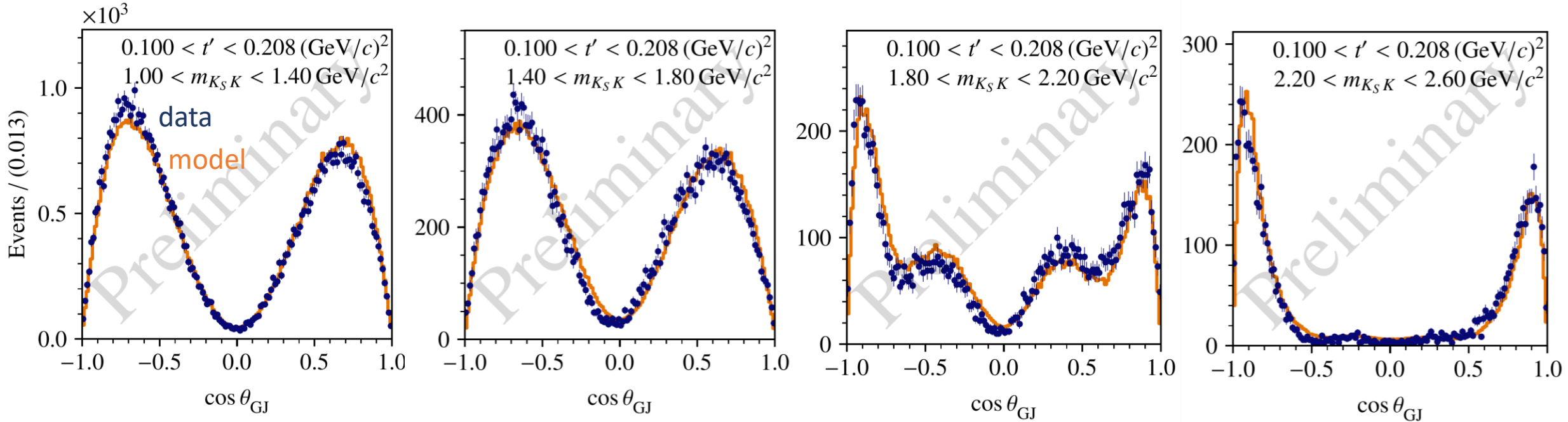
$a_4(????)$



Results: Intensities and Phases



Agreement Between Model and Data



- Angular distributions as **predicted by the PWA model** after the fit vs **real data**
- Good agreement, but data exhibits larger asymmetry than predicted in the model

Adding an Odd-Spin Wave

How can we introduce asymmetry in the model?

- Effects of the detector acceptance
- Introduce partial wave(s) with odd spin J

$$I^{GJPC} = 1^+ 1^{--}$$

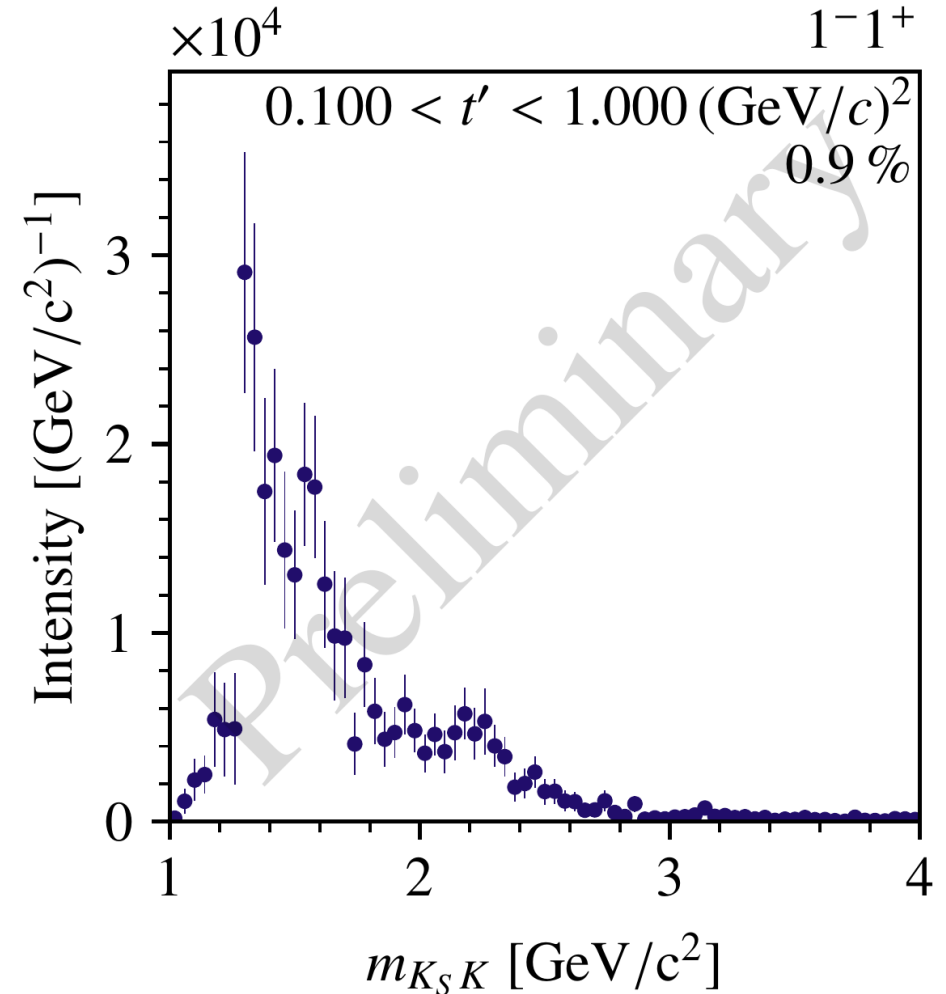
- Cannot be produced via Pomeron exchange!
(but e.g. by ω exchange)

$\rho(1450)$

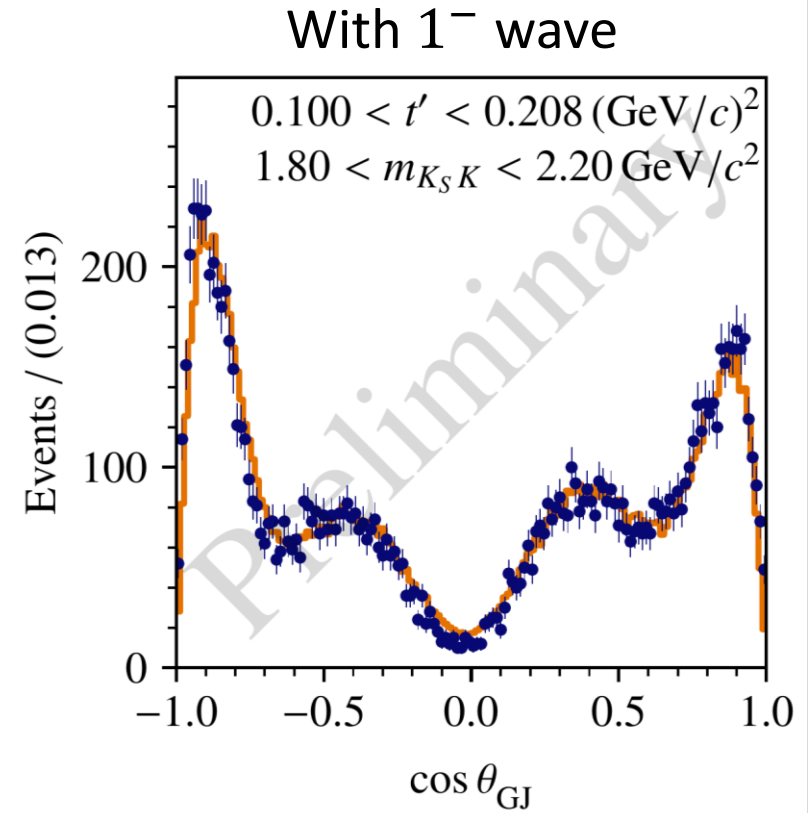
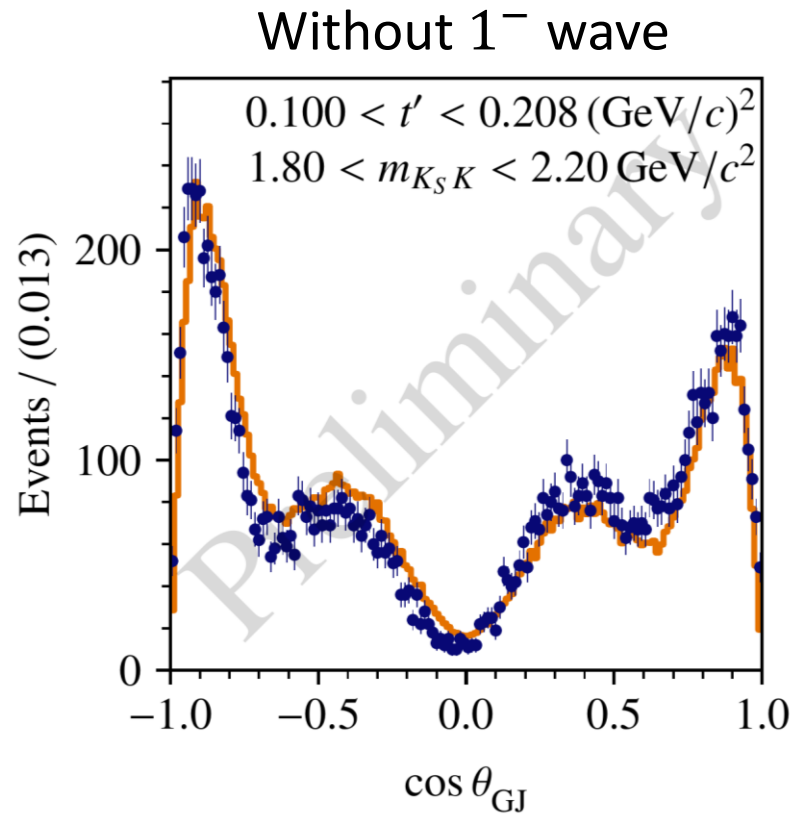
- $m = 1.47$ GeV
- $\Gamma = 0.40$ GeV

$\rho(1700)$

- $m = 1.73$ GeV
- $\Gamma = 0.25$ GeV



Agreement Between Model and Data II



- Interference of $J = 1$ and even J could explain the forward/backward asymmetry

Conclusion and Outlook

- Resonances appearing in the final state:

$J^{PC} = 2^{++}, 4^{++}, \dots \rightarrow a_J$ states (Pomeron exchange)

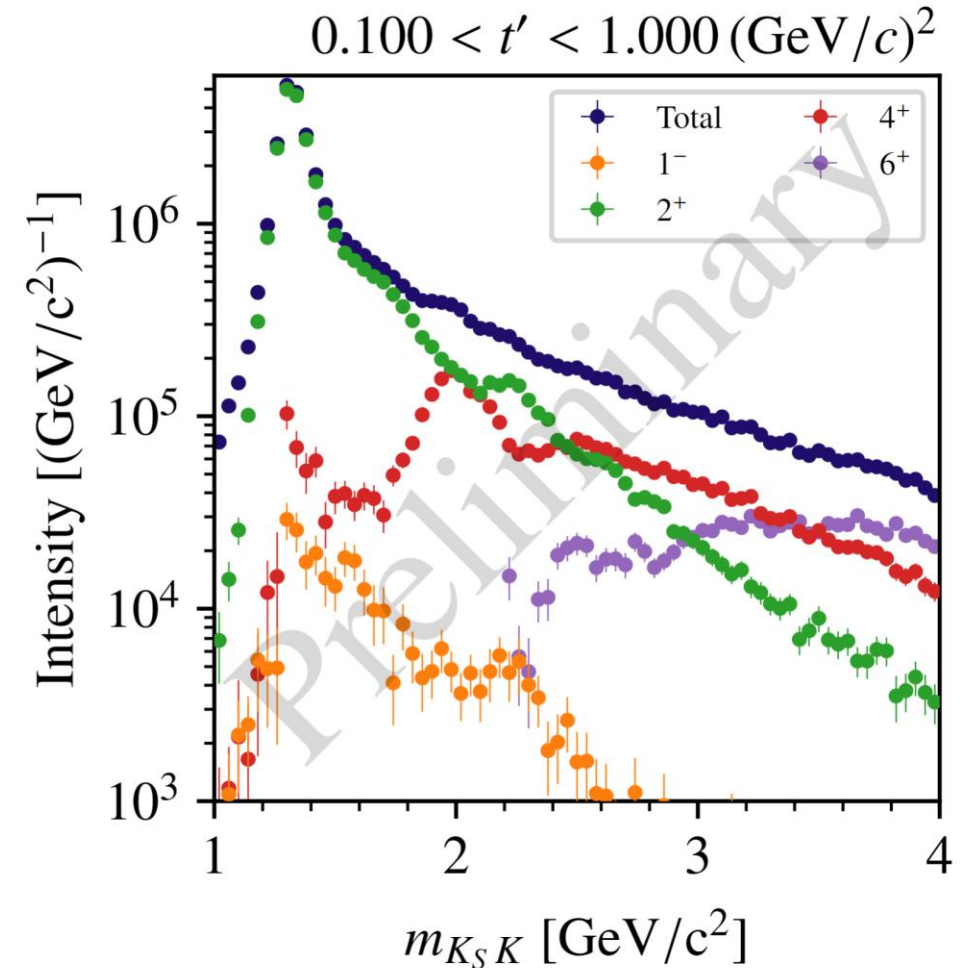
other exchanges may yield $J^{PC} = 1^{--}, 3^{--}, \dots$

- Performed the **partial-wave decomposition**:

- Clear signals of $a_2(1320)$ and $a_4(1970)$
- Possible higher-lying a_2, a'_4 and a_6
- Indications of small intensity in 1^{--} partial wave

- Next step: **Resonance-model fit**

Extract resonance parameters by modelling m_{KK}, t' dependences



Thank you for your attention!

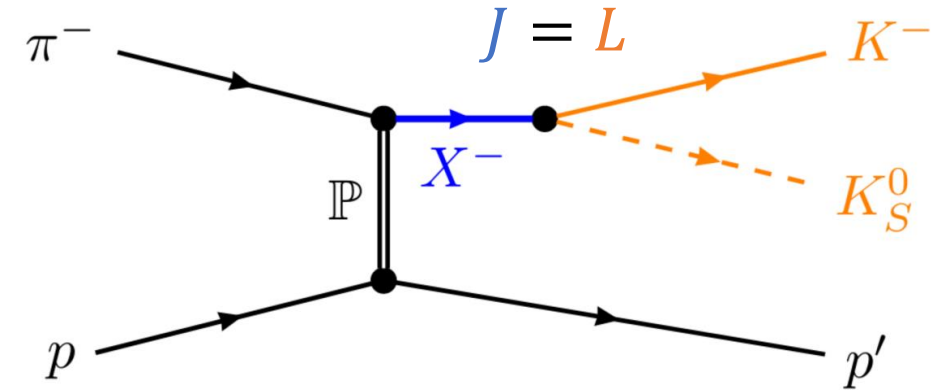
BACKUP

Quantum Numbers of the $K_S^0 K^-$ Final State

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$$\begin{aligned}
 I &= 1 \\
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 P &= (-1)^L \quad C = G(-1)^I = (-1)^L
 \end{aligned}$$

all given by $L!$



- Quantum numbers of the produced X^- :

$$I^G = 1^- \text{ for Pomeron } (I^G = 0^+) \text{ exchange}$$

→ odd spins J are suppressed!

- Reflectivity ($\epsilon = \pm 1$) basis:

$$\text{eigenvalue of } \hat{\epsilon} = \hat{P}R_{\perp}(\pi)$$

at high \sqrt{s} , $\epsilon \approx \eta$ η : naturality of the exchange particle

→ $\epsilon = +1$ for Pomeron exchange

$$I^G J^{PC} = 1^+ 1^{--}, 1^- 2^{++}, 1^+ 3^{--}, 1^- 4^{++}, \dots$$

S. U. Chung, doi.org/10.1103/PhysRevD.11.633

S. U. Chung, C- and G-parity: A New Definition and Applications. <https://suchung.web.cern.ch/Cparity7b.pdf>

Ambiguities in Incoherent Sectors

$$\varepsilon = \pm 1: \quad I(m_X, t'; \theta, \phi) = \left| \sum_{JM} T_{JM}^+(m_X, t') \psi_{JM}^+(\theta, \phi) \right|^2 + \left| \sum_{JM} T_{JM}^-(m_X, t') \psi_{JM}^-(\theta, \phi) \right|^2$$

$$a_1^+ = \sum_{J=1}^{J_{\max}^+} T_{J1}^+ Y_J^1(\theta, 0) \quad \varepsilon = +1, M = 1$$

$$a_0^- = \sum_{J=0}^{J_{\max}^-} T_{J0}^- Y_J^0(\theta, 0) \quad \varepsilon = -1, M = 0$$

$$a_1^- = \sum_{J=1}^{J_{\max}^-} T_{J1}^- Y_J^1(\theta, 0) \quad \varepsilon = -1, M = 1$$

- Define $a_s^- = a_0^- + a_1^-$, then same procedure as for only positive-reflectivity sector
- New amplitudes for $\varepsilon = +1$: $|a_1^+|^2 = |a_1^-|^2 - \text{const.}$ → **positivity requirement!**

Study of the Ambiguities

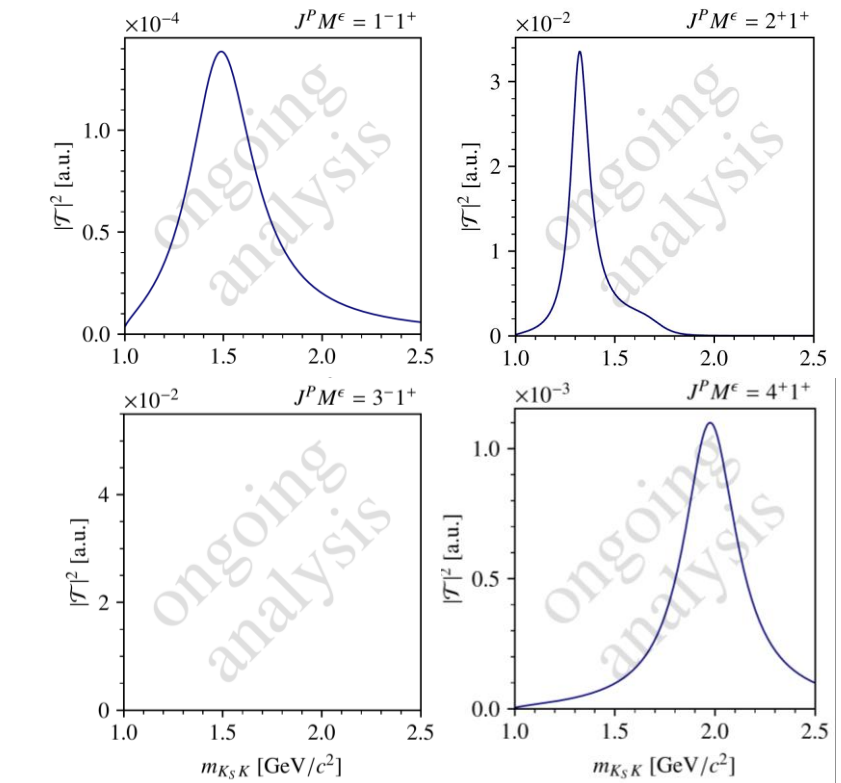
- How do the ambiguous solutions look like (**continuity, signals, ...**)?
- What are the effects of the **partial-wave decomposition fit on finite data** on the ambiguities?

I. Continuous intensity model

- create an amplitude model for selected partial waves
- calculate exact ambiguities

II. Finite pseudo-data

- generate pseudo-data according to model
- perform partial-wave decomposition

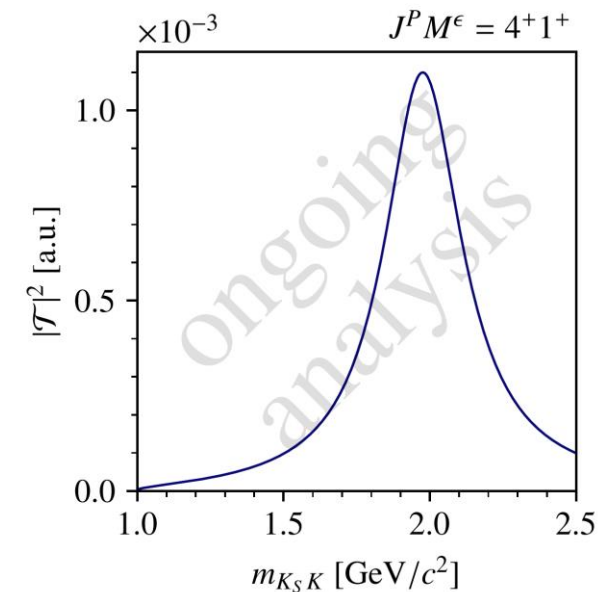
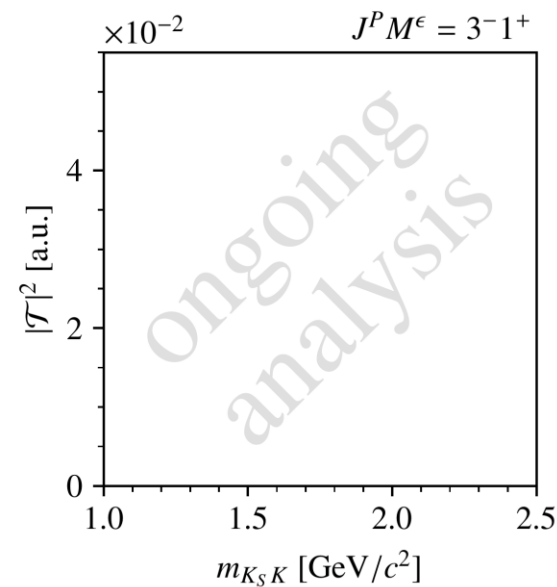
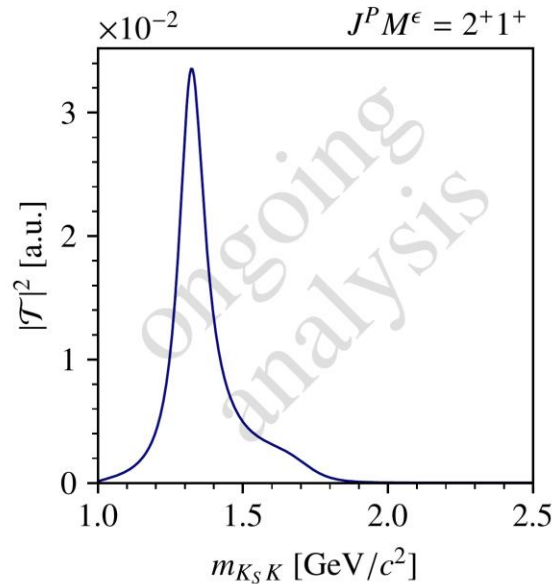
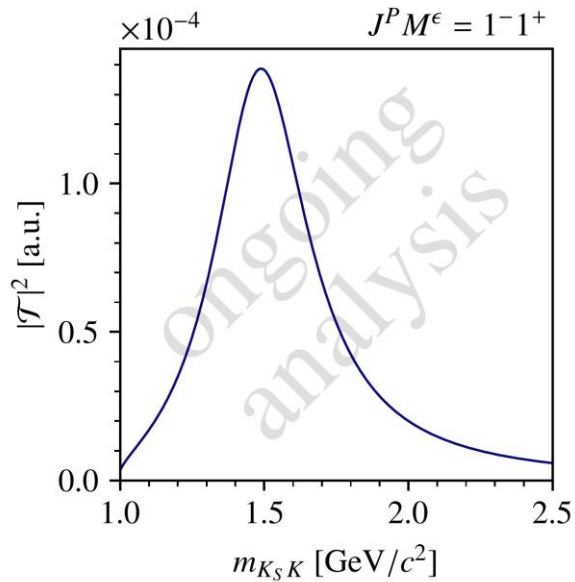


Continuous Amplitude Model

I. Continuous intensity model

- create an amplitude model for **four** selected partial waves
- In $1.0 < m_X < 2.5 \text{ GeV}/c^2$
- m_X -dependence by Breit-Wigner amplitudes (PDG parameters)

J^{PC}	Resonances
1^{--}	$\rho(1450)$
2^{++}	$a_2(1320), a'_2(1700)$
3^{--}	None
4^{++}	$a_4(1970)$

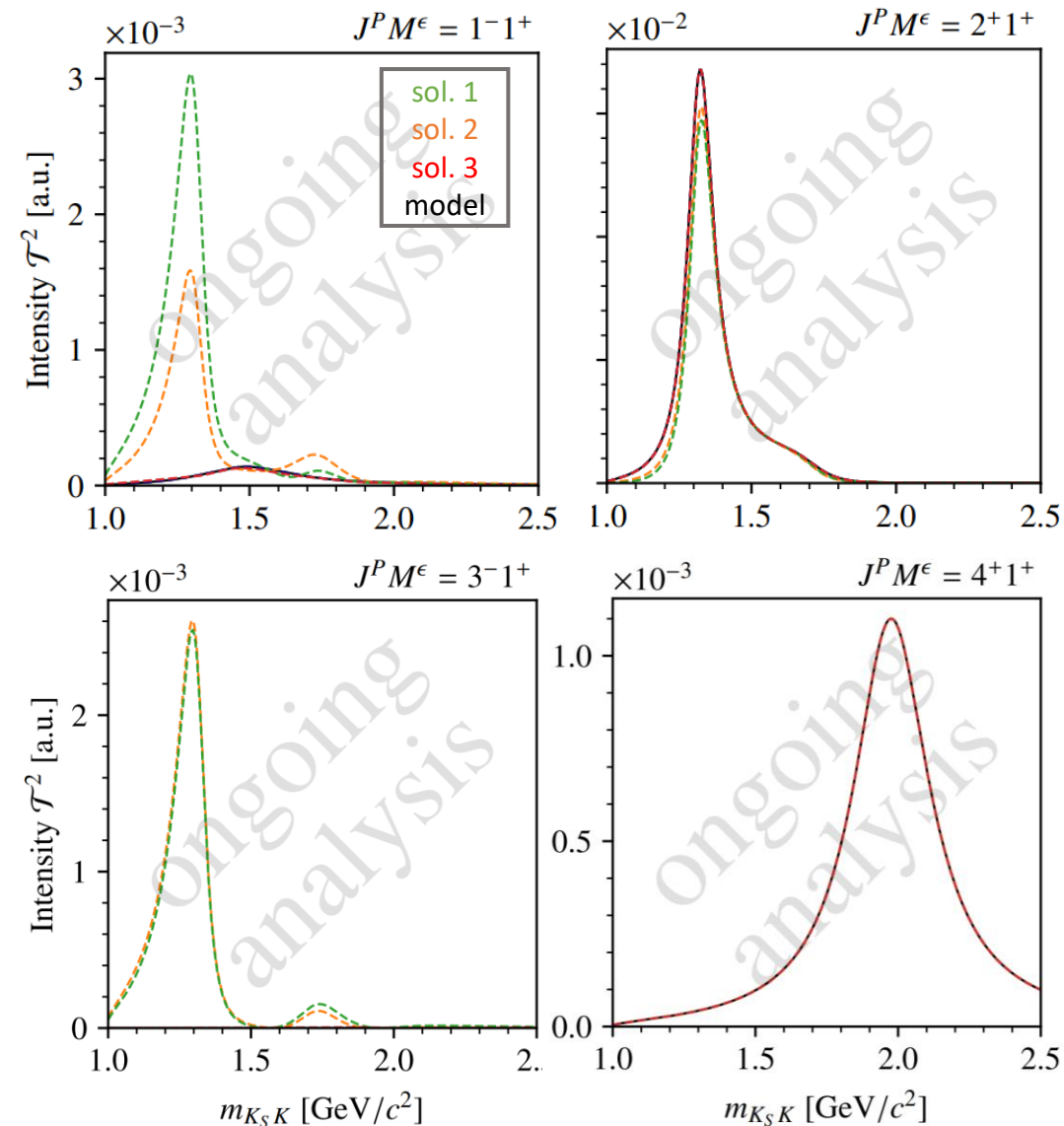


Continuous Amplitude Model

I. Continuous intensity model

$$N_a = 3$$

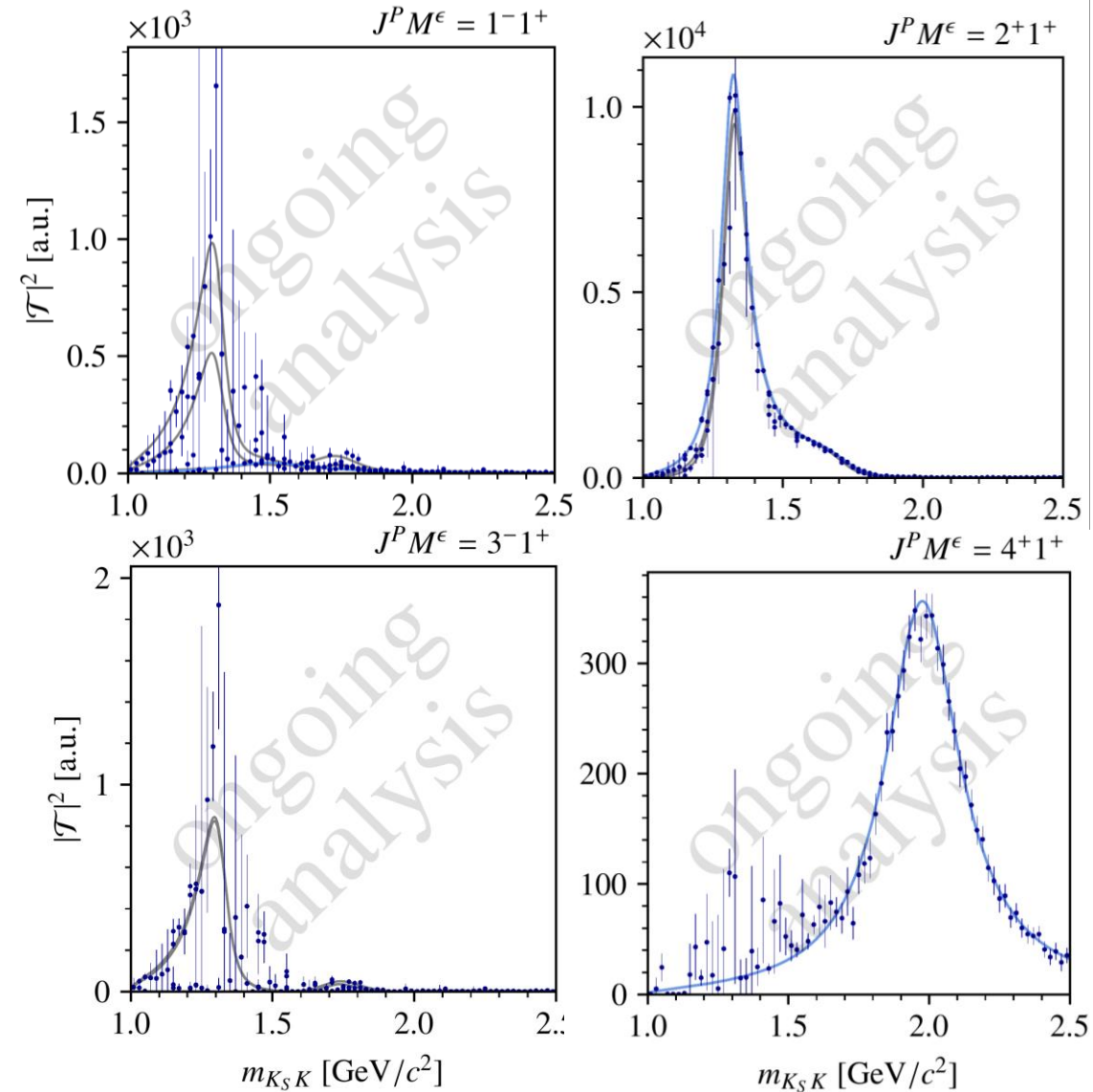
- Sample points in m_X and calculate ambiguous solutions
- Ambiguous intensities are also continuous
- Not all solutions are different from each other!
- Highest-spin (4^{++}) intensity is invariant!

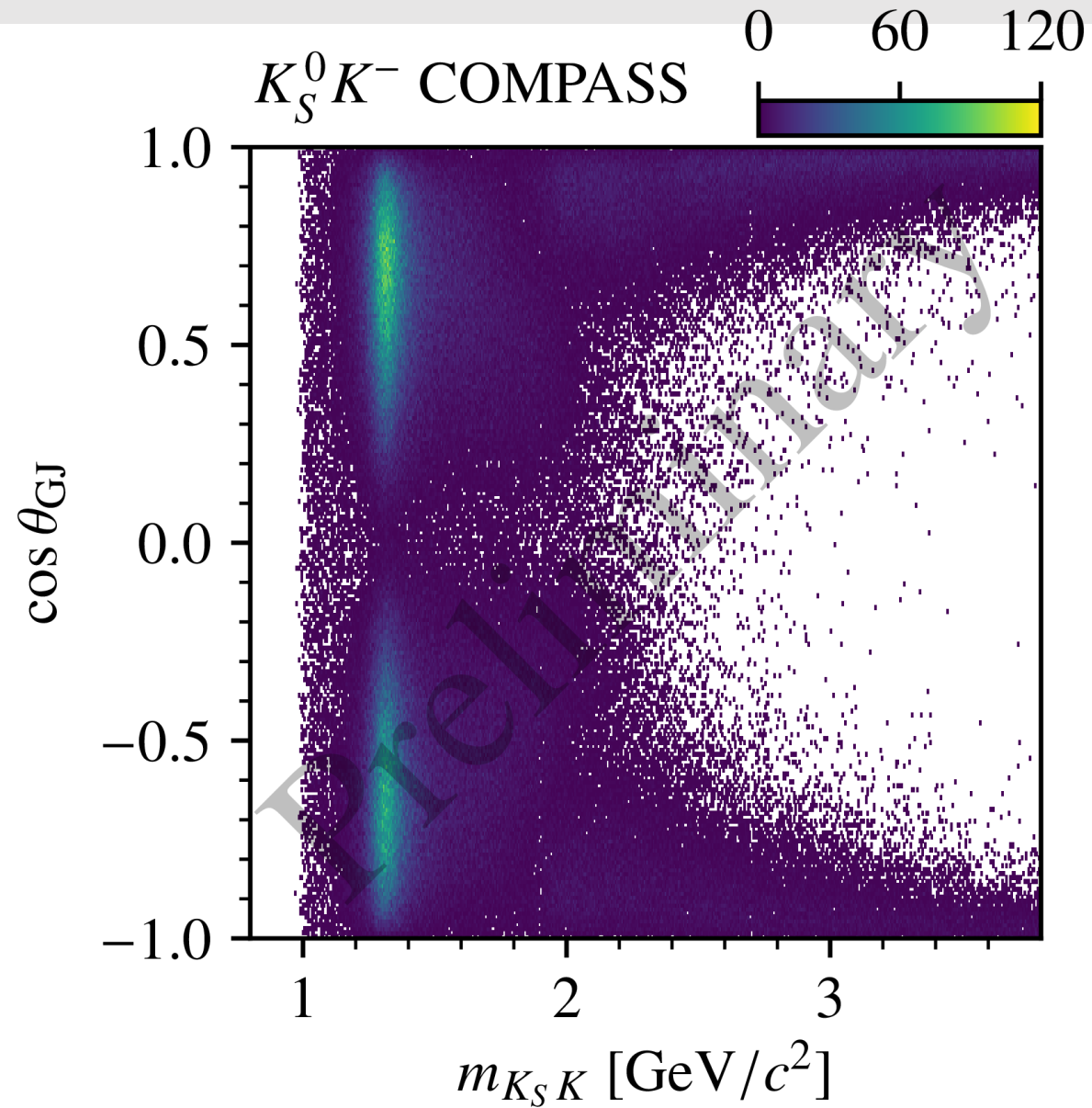


Partial-Wave Decomposition Fits on Pseudodata

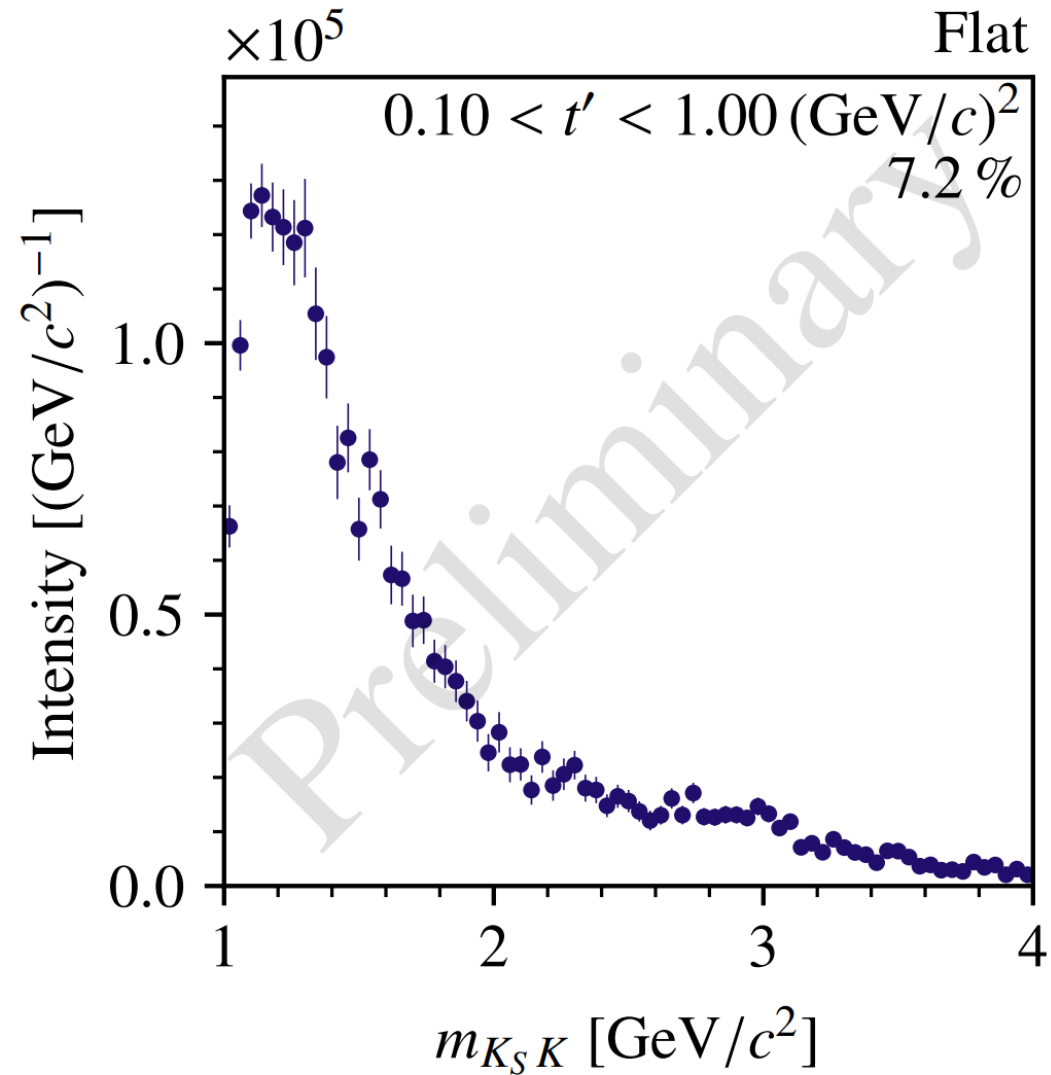
II. Finite pseudo-data

- 4^{++} intensity is still invariant!
- Overall, amplitude values found by the fit follow the calculated distributions
- Not all solutions are found in each m_X bin
→ PWD fit distorts the intensity distribution!





Results: Flat-Wave Intensity



Kinematic Binning

Analyzed invariant mass range:

$$1.0 < m_{K_S K} < 4.0 \text{ GeV}/c^2$$

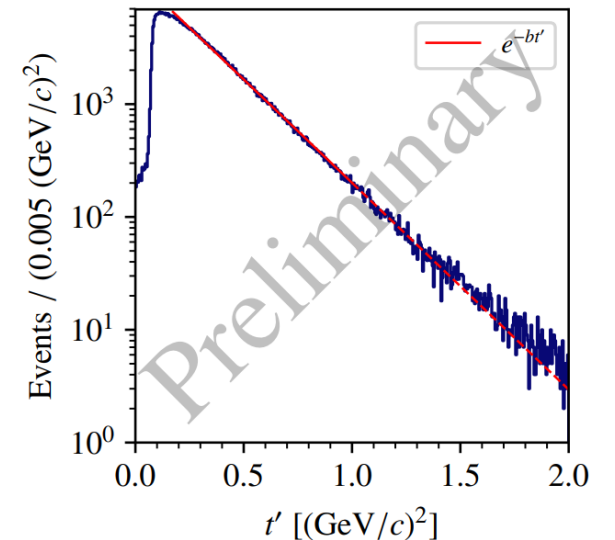
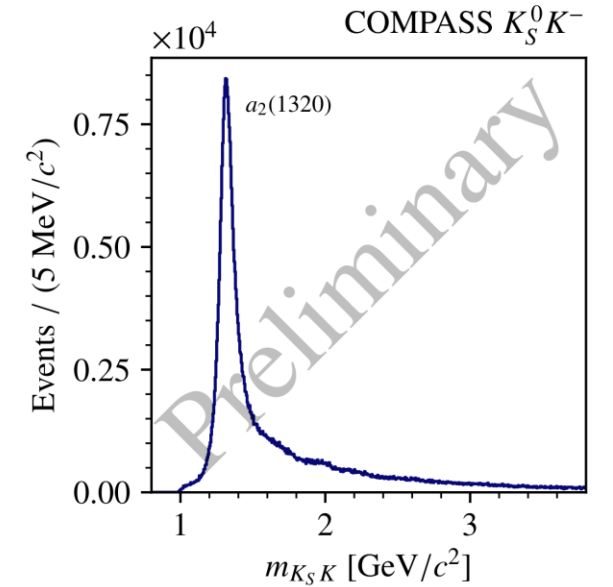
divided into 75 bins of $40 \text{ MeV}/c^2$ width

Analyzed t' range

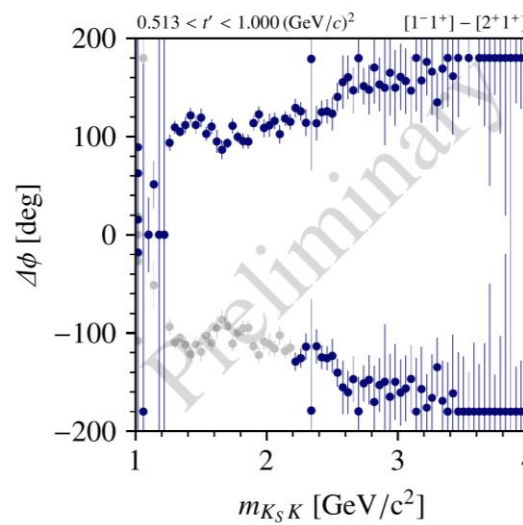
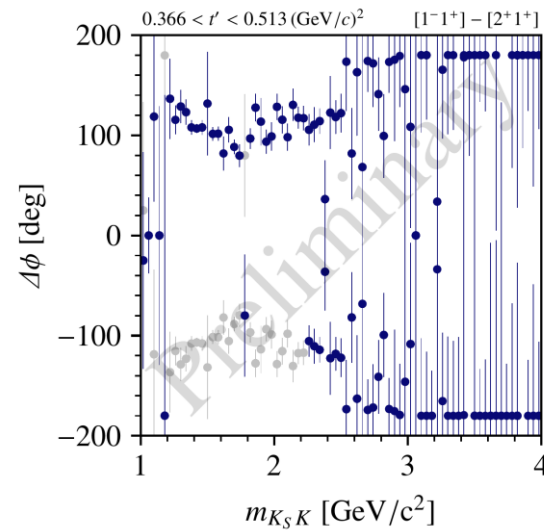
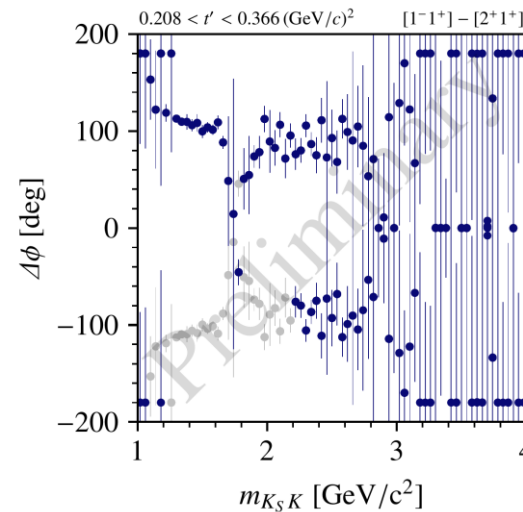
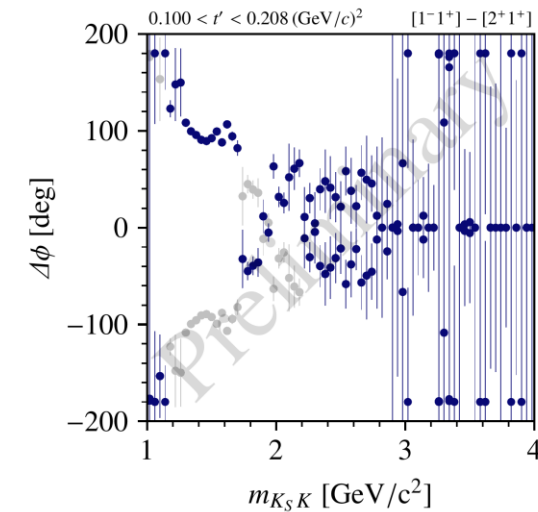
$$0.1 < t' < 1.0 \text{ GeV}^2/c^2$$

divided into four bins:

t' bin borders in $(\text{GeV}/c)^2$				
0.100	0.208	0.366	0.513	1.000



$J^{PC} = 1^{--}$ Partial Wave

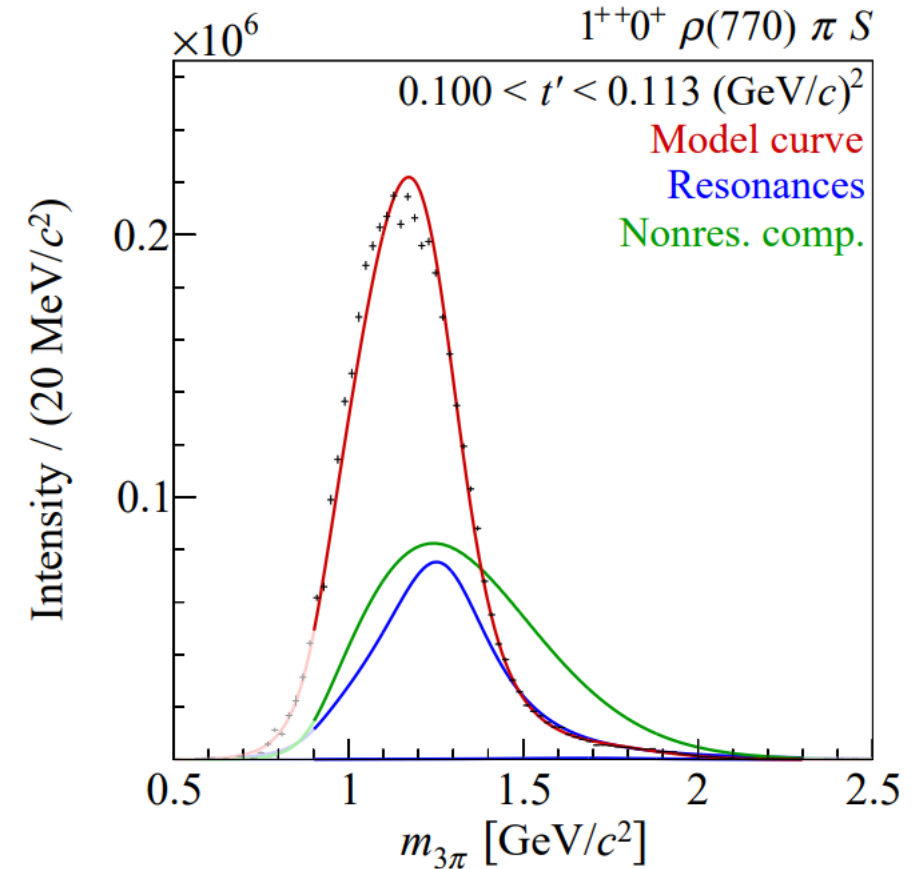


Resonance-Model Fit

Second step: **extract resonance parameters**

- Build **model** for mass dep. of partial-wave amplitudes:
resonant (e.g. Breit-Wigner distribution)
+ **non-resonant background** components
- χ^2 fit to output of partial-wave decomposition

→ get **masses and widths** of parameterized resonances



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