



GPD Studies at the COMPASS Experiment

***3D Structure of the Nucleon
Via Generalized Parton Distributions***

Howard Johnson Incheon Airport Hotel

June 27, 2024

Po-Ju Lin

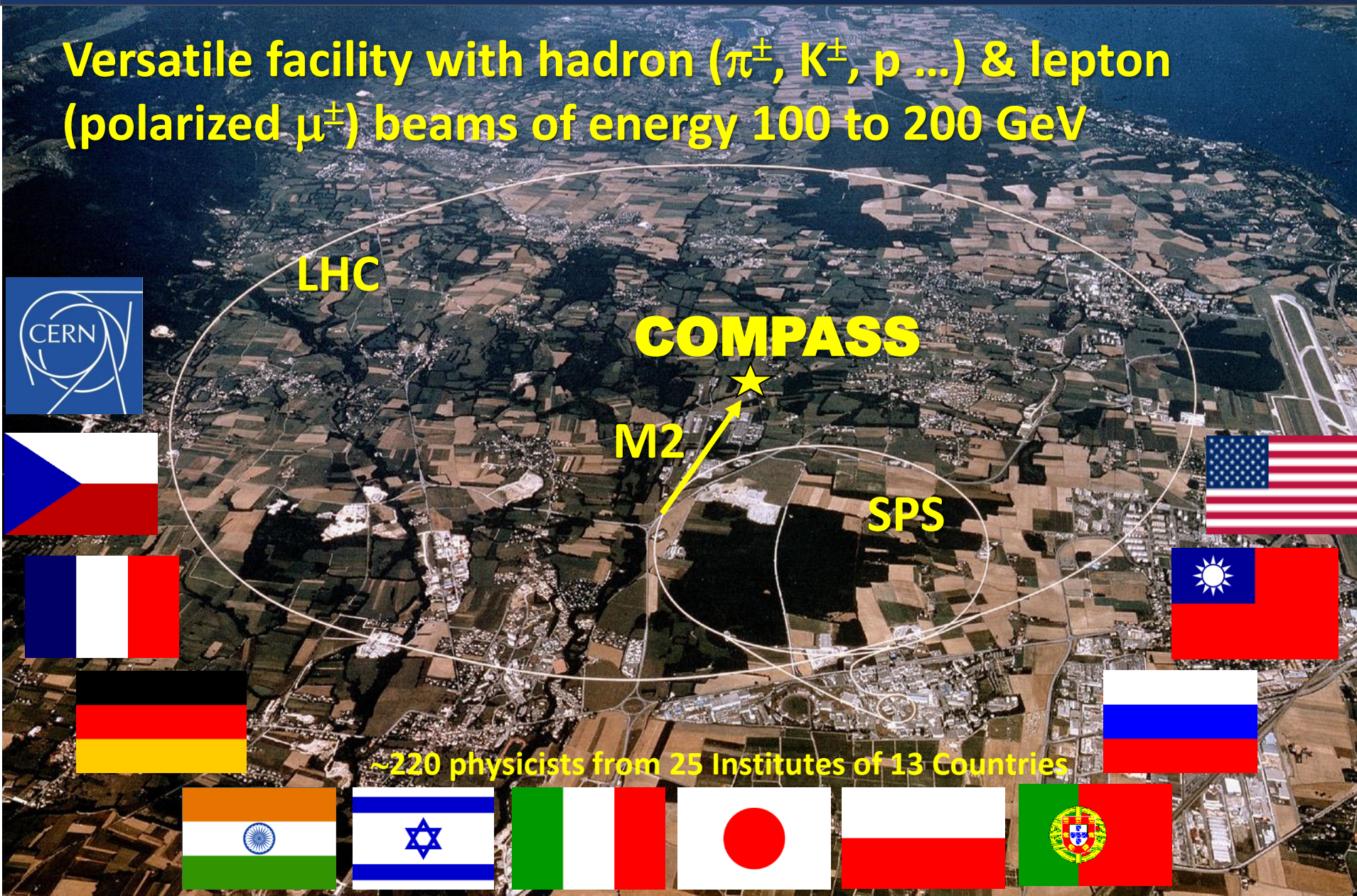
Department of Physics, National Central University

- The COMPASS Experiment
- Deeply Virtual Compton Scattering (DVCS)
- Hard Exclusive Meson Production (HEMP)
- Summary

COMPASS Experiment



Versatile facility with hadron (π^\pm , K^\pm , p ...) & lepton (polarized μ^\pm) beams of energy 100 to 200 GeV

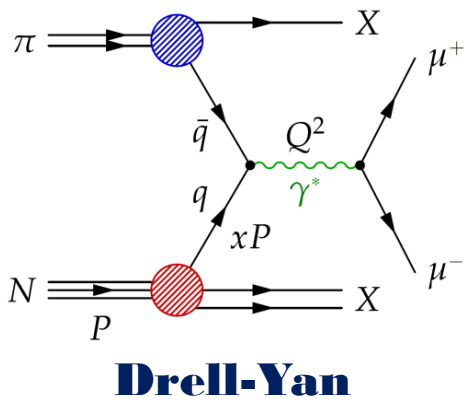
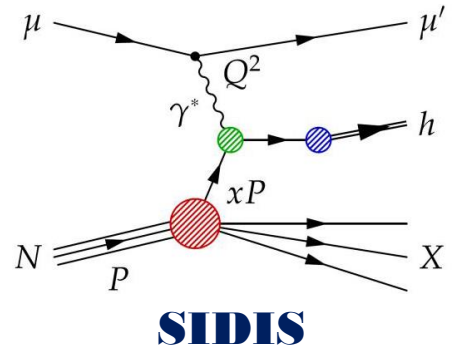


COmmun
Muon and
Proton
Apparatus for
Structure and
Spectroscopy

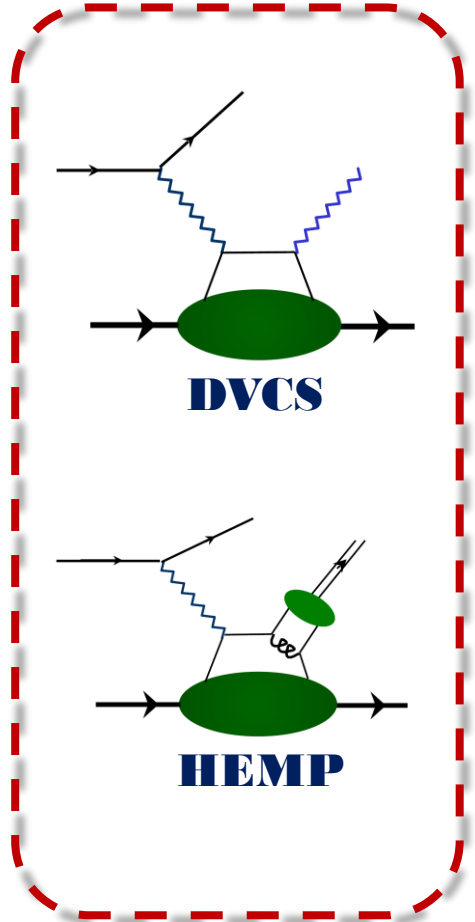
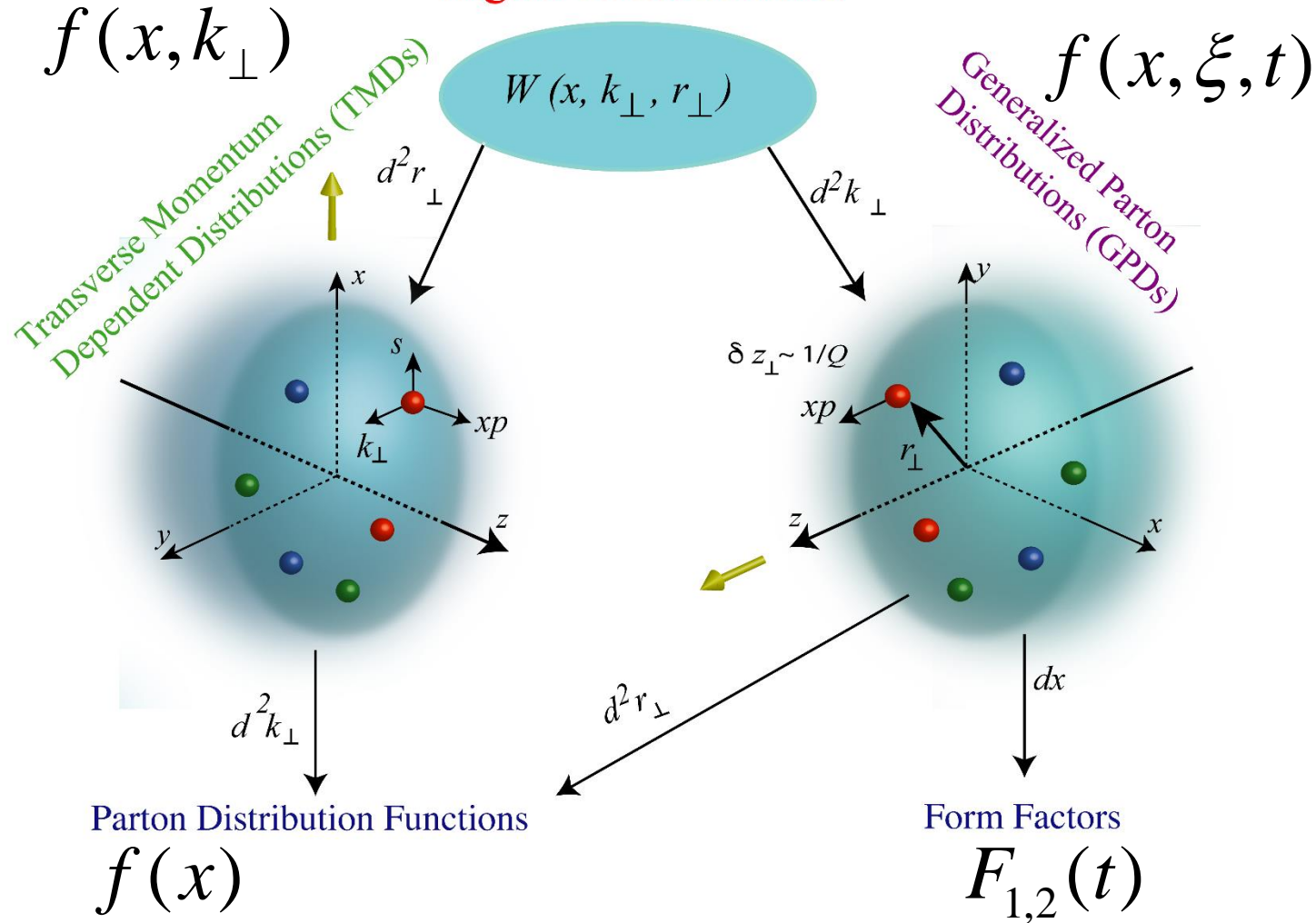
~220 physicists from 25 Institutes of 13 Countries

Multi-dimensional Partonic Structures

<http://www.int.washington.edu/PROGRAMS/17-3/>

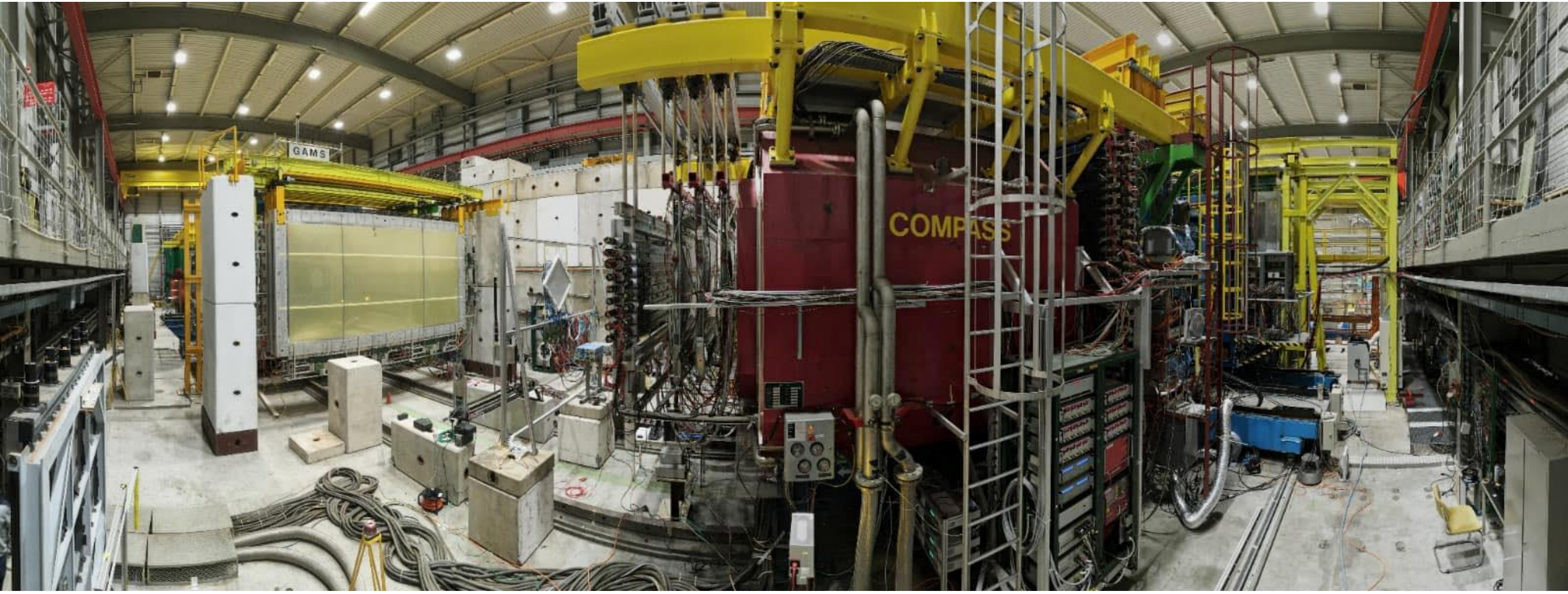


Wigner Distributions

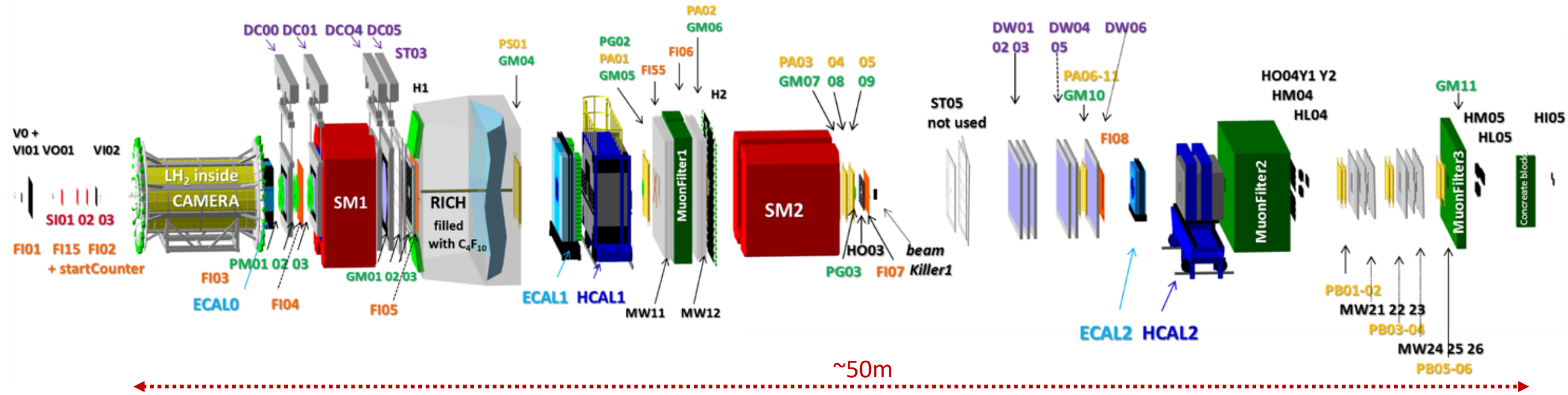


COMPASS investigates the multi-dimensional structure of nucleon via various processes

COMPASS Experimental Setup



COMPASS Experimental Setup



- Primary beam – 400 GeV p from SPS
 - Impinging on Be production target
- 190 GeV secondary hadron beams
 - h^- beam: 97% π^- , 2% K^- , 1% p
 - h^+ beam: 75% π^+ , 24% p, 1% K^+
- 160 GeV tertiary muon beams
 - μ^\pm longitudinally polarized

Large-acceptance forward spectrometer

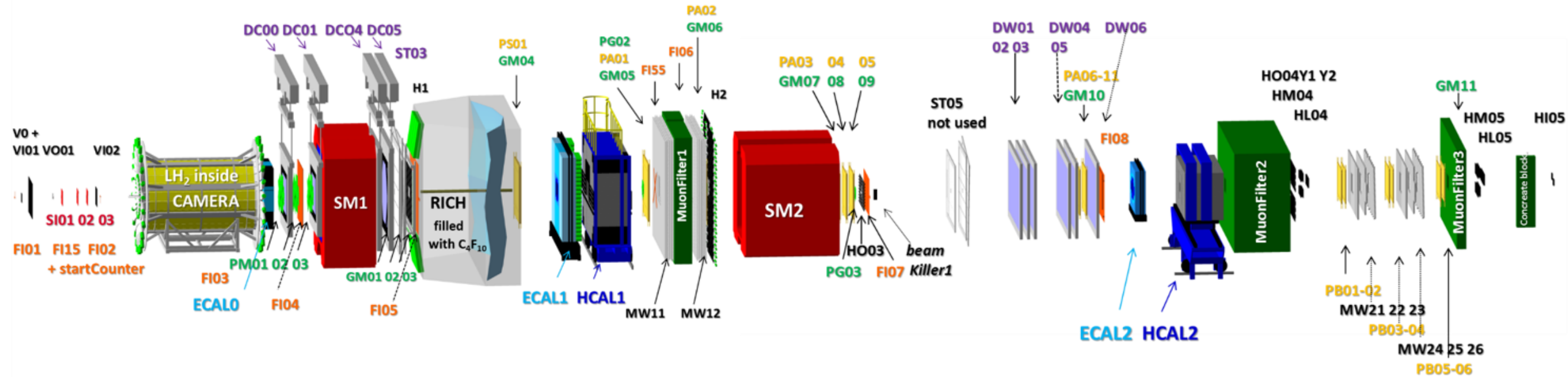
- Precise tracking (350 planes)
SciFi, Silicon, MicroMegas, GEM, MWPC, DC, straw
- PID – CEDARs, RICH, calorimeters, Muon Walls

Various targets:

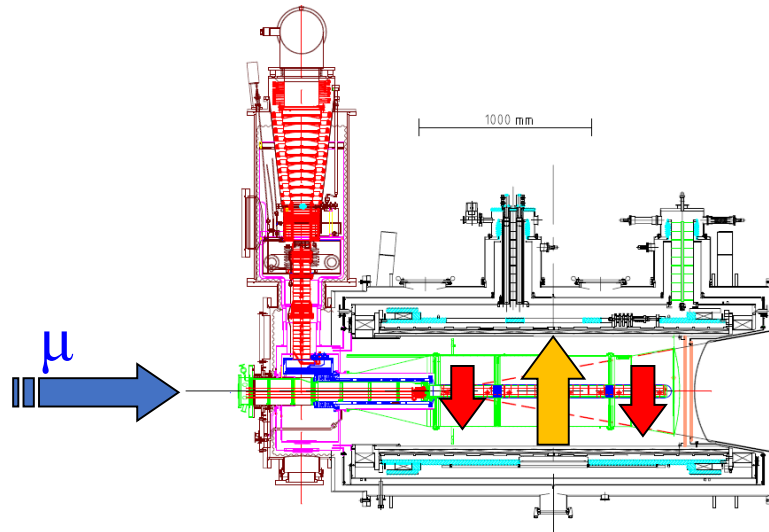
- Polarized solid-state NH_3 or ^6LiD
- Liquid H_2
- Solid-state nuclear targets

❖ NIM A 577 (2007) & NIM A 779 (2015) 69

COMPASS Experimental Setup

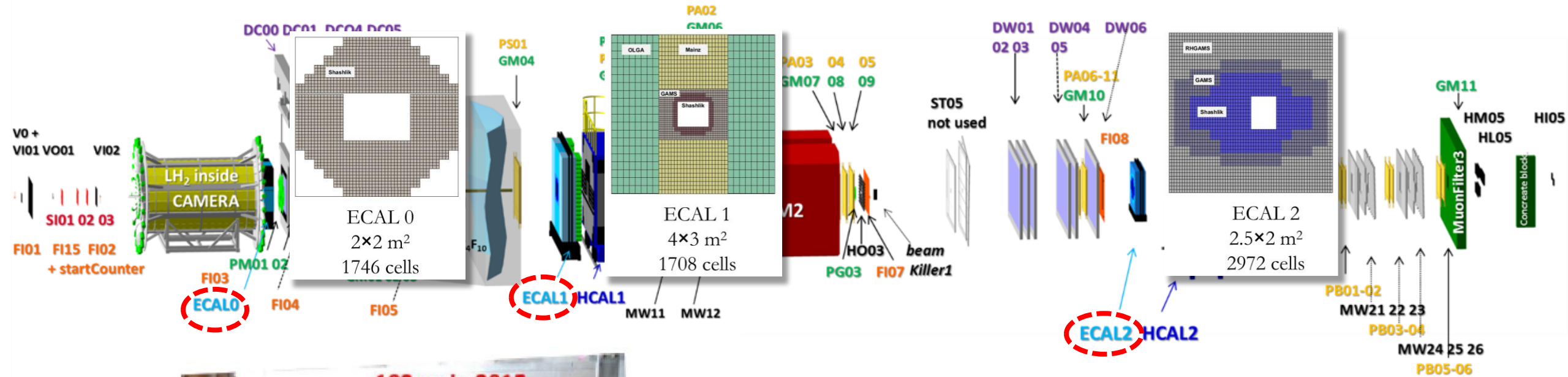


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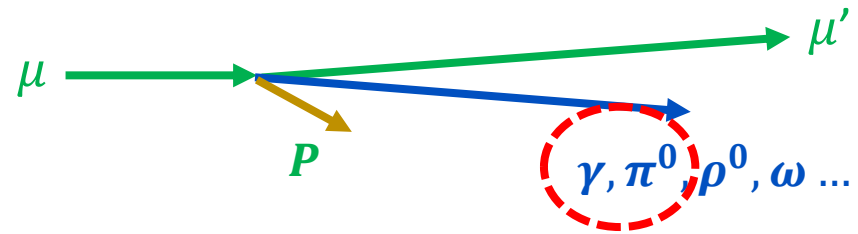


- In early GPD studies, transversely polarized target was used.
- Polarization reversal by magnetic field rotation
- 2.5m unpolarized LH₂ target used in GPD dedicated runs

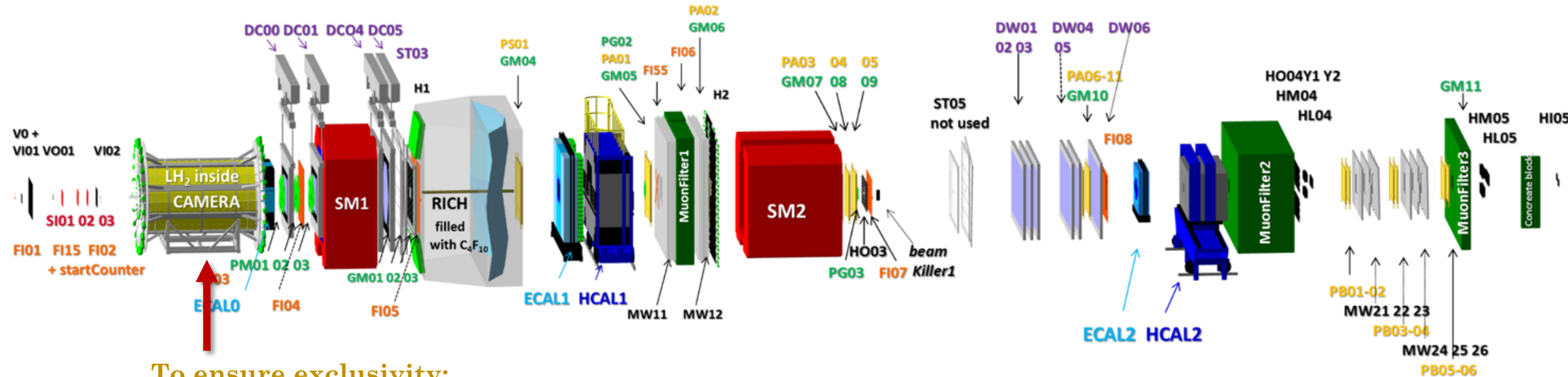
COMPASS Setup for Exclusive Processes



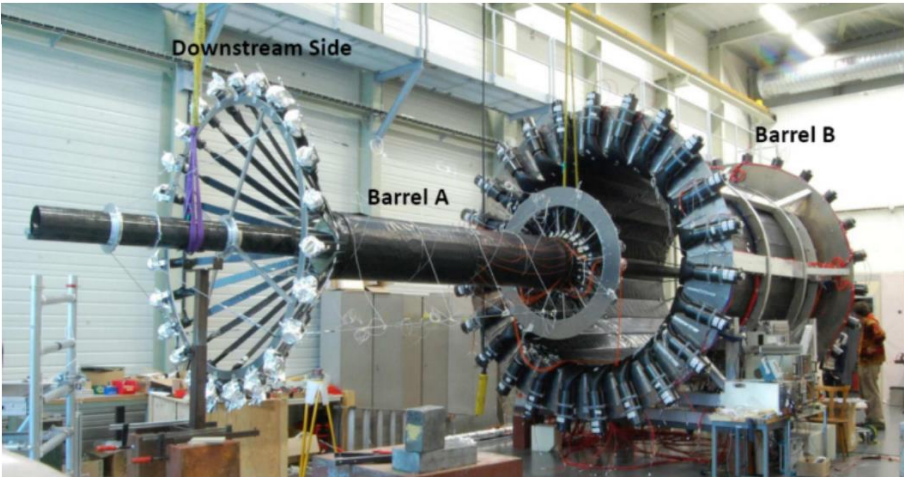
Exclusive Muoproduction



COMPASS Setup for Exclusive Processes

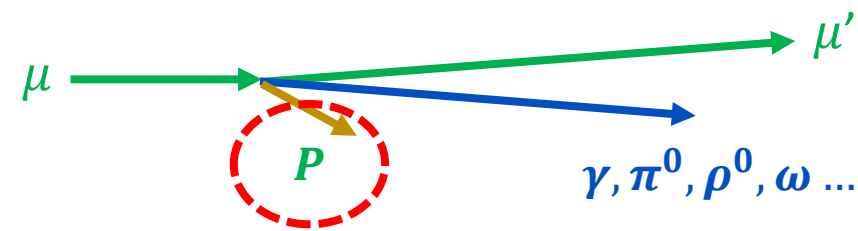


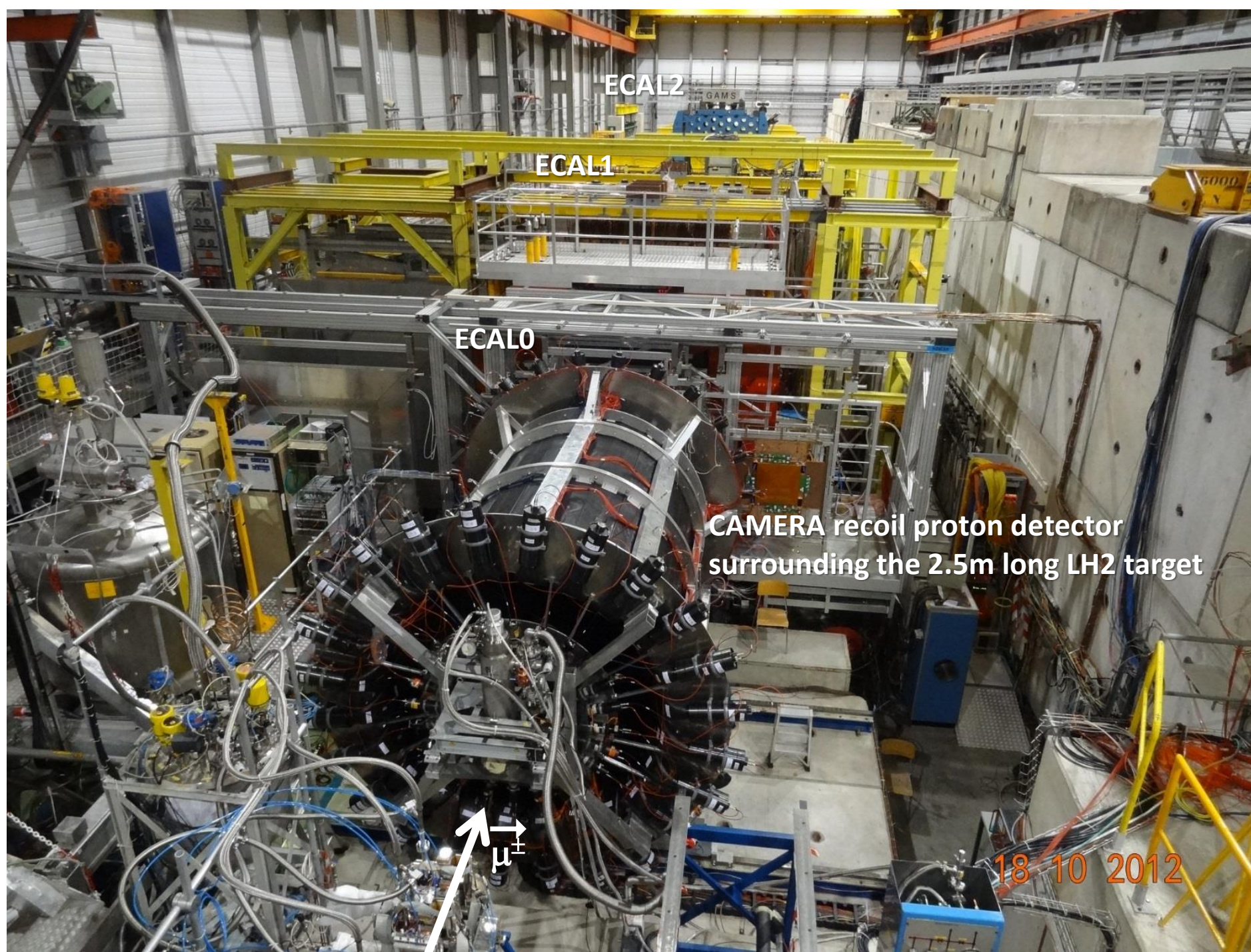
To ensure exclusivity:



CAMERA recoil proton detector

Exclusive Muoproduction





COMPASS Experiment



2002-2022 COMPASS data taking

2002-2004	DIS & SIDIS, μ^+ - d , 160 GeV, L & T polarized target
2005	<i>CERN accelerator shutdown, increase of COMPASS acceptance</i>
2006	DIS & SIDIS, μ^+ - d , 160 GeV, L polarized target
2007	DIS & SIDIS, μ^+ - p , 160 GeV, L & T polarized target
2008-2009	Hadron spectroscopy & Primakoff reaction, $\pi/K/p$ beam
2010	SIDIS, μ^+ - p , 160 GeV, T polarized target
2011	DIS & SIDIS, μ^+ - p , 200 GeV, L polarized target
2012	Primakoff reaction, $\pi/K/p$ beam
2012 pilot run	DVCS/HEMP/SIDIS, μ^+ & μ^- - p , 160 GeV, unpolarized target
2013	<i>CERN accelerator shutdown, LS1</i>
2014-2015	Drell-Yan, π^- - p , T polarized target
2016-2017	DVCS/HEMP/SIDIS, μ^+ & μ^- - p , 160 GeV, unpolarized target
2018	Drell-Yan, π^- - p , T polarized target
2019-2020	<i>CERN accelerator shutdown, LS2</i>
2021-2022	SIDIS, μ^+ - d , 160 GeV, T polarized target

Study hadron structure with complementary tools:

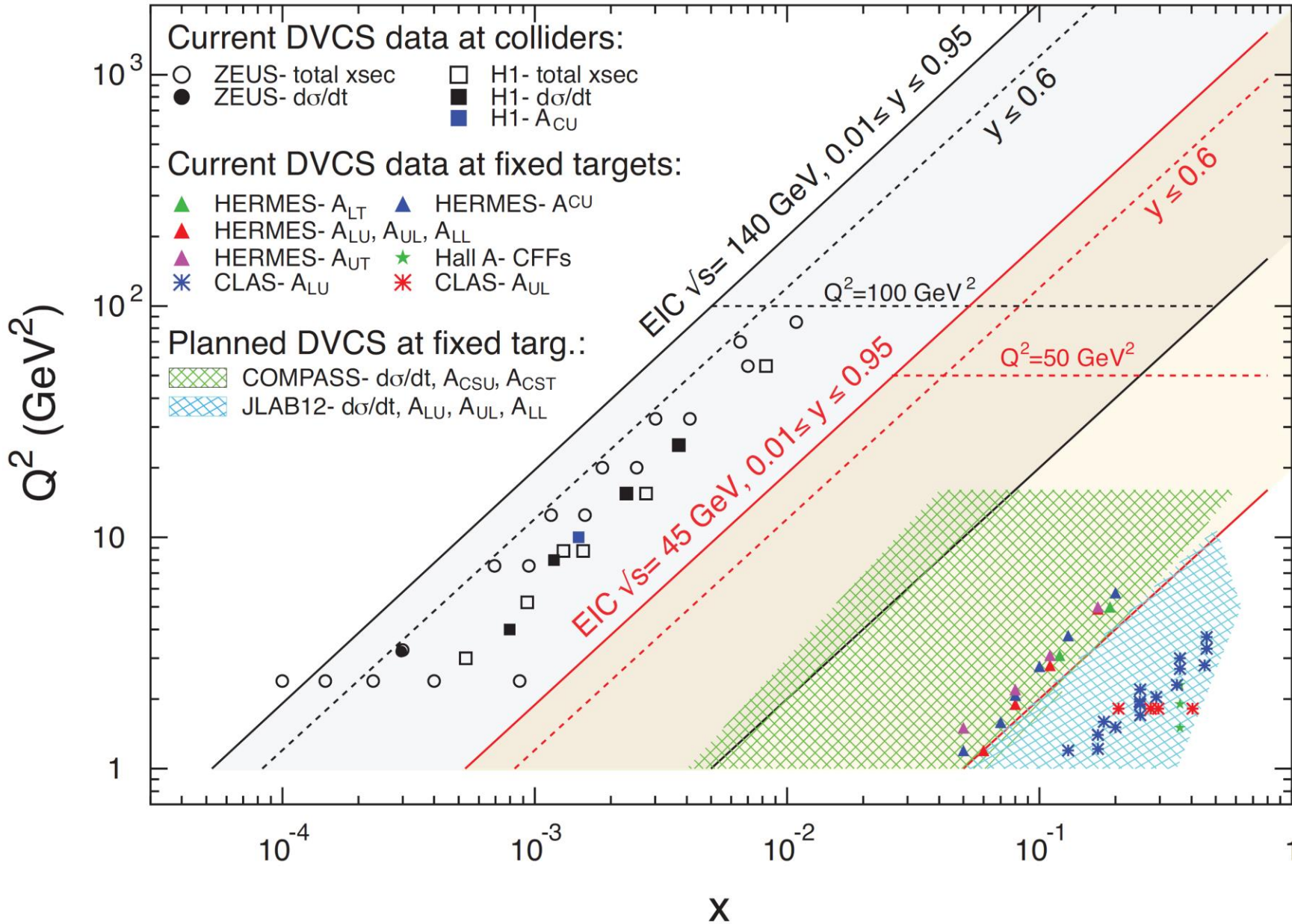
➤ **COMPASS holds the record for the longest-running CERN experiment**

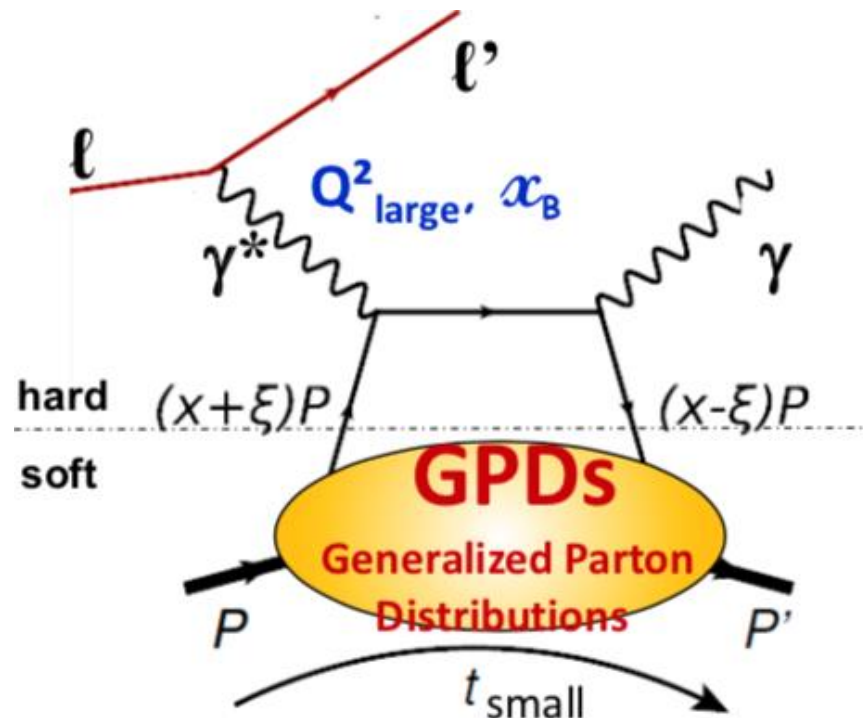


- 2012 pilot run with 4-week data taking
- 2016-17 dedicated run. 2 x 6 months.

Deeply Virtual Compton Scattering @ COMPASS

Landscape – Global Programs of DVCS





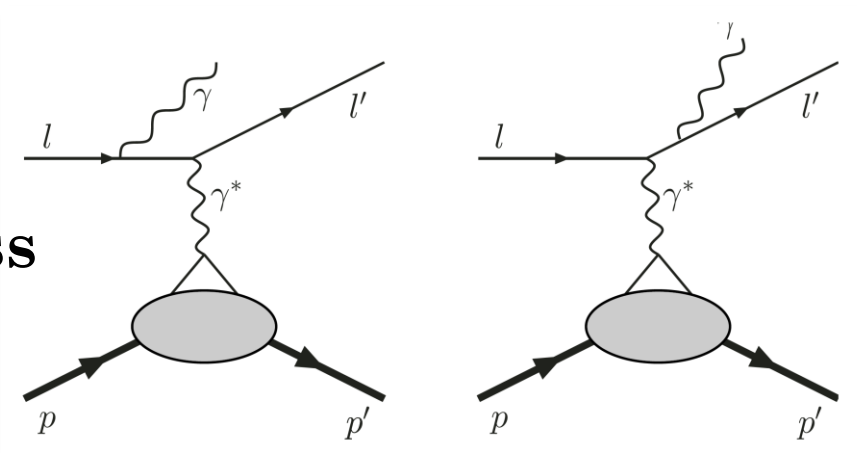
$$\text{DVCS: } l + p \rightarrow l' + p' + \gamma$$

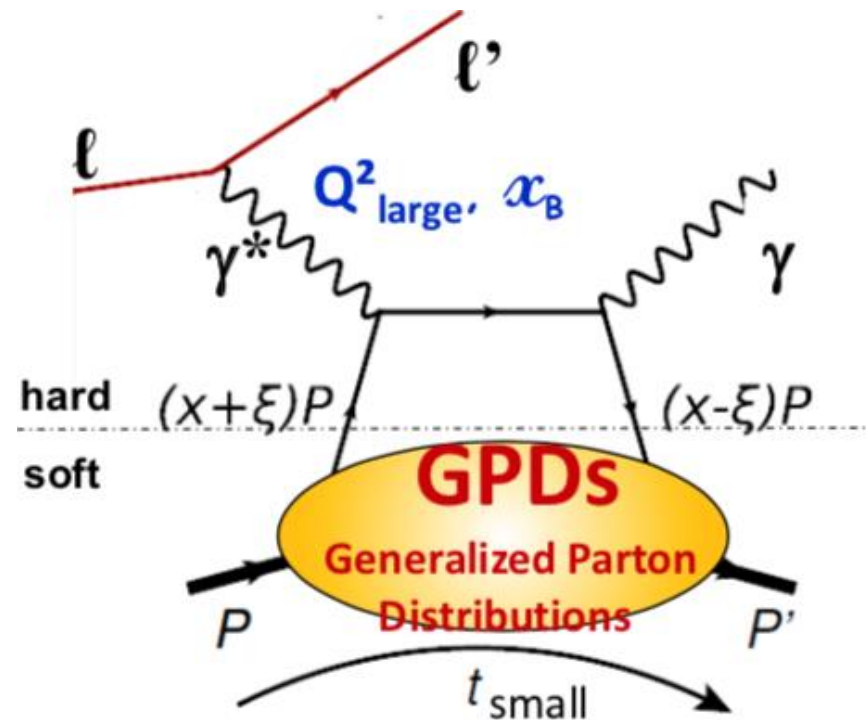
- As the golden channel to access GPDs, DVCS has been the workhorse for GPD Extraction.
- Its interference with the well-understood Bethe-Heitler process gives access to more info.

- The GPDs depend on the following variables:

- x : average longitudinal momentum frac.
- ξ : longitudinal momentum diff.
- t : four momentum transfer
(correlated to b_{\perp} via Fourier transform)
- Q^2 : virtuality of γ^*

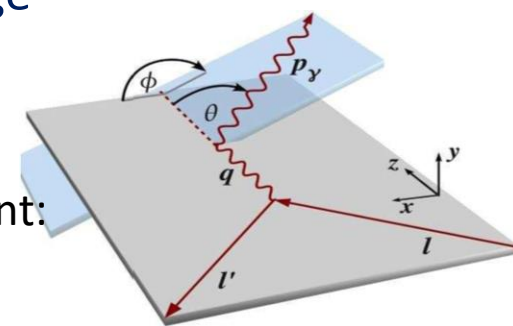
**BH
Process**





DVCS: $l + p \rightarrow l' + p' + \gamma$

- With LH₂ target and small x_B coverage
→ focuses on H at COMPASS



- The variables measured in the experiment:

$$E_\ell, Q^2, x_{Bj} \sim 2\xi / (1+\xi),$$

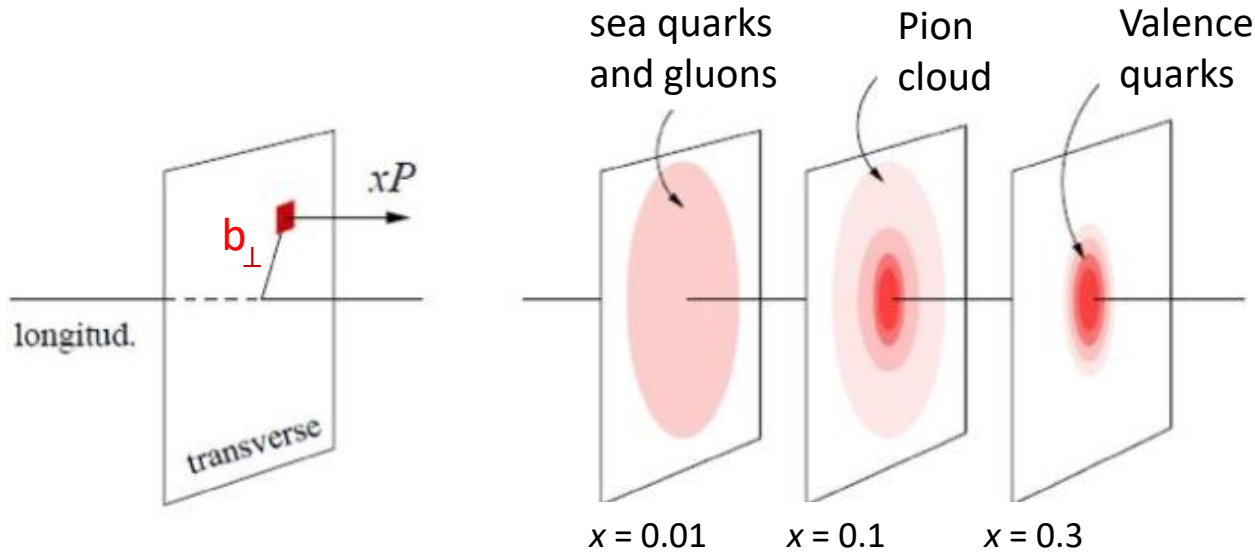
$$t \text{ (or } \theta_{\gamma^*\gamma} \text{) and } \phi \text{ (} \ell\ell' \text{ plane}/\gamma\gamma^* \text{ plane)}$$

$$\overset{GPD}{\mathcal{H}(\xi, t)} = \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi + i\epsilon} + \dots = \overset{REAL \text{ part}}{\mathcal{P} \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi}} - \overset{Imaginary \text{ part}}{i\pi \mathbf{H}(x = \pm \xi, \xi, t)} + \dots$$

$$\text{Re } \mathcal{H}(\xi, t) = \mathcal{P} \int dx \frac{\text{Im } \mathcal{H}(x, t)}{x - \xi} + \Delta(t)$$

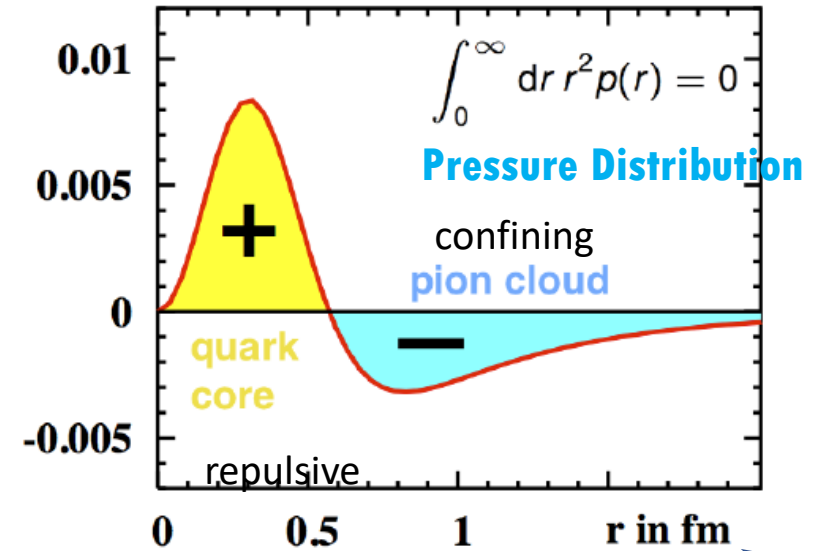
Transverse Imaging and Pressure Distribution

Mapping in the transverse plane



M. Polyakov, P. Schweitzer, *Int.J.Mod.Phys. A33* (2018)

$r^2 p(r)$ in GeV fm^{-1}



$\overset{\text{GPD}}{\mathcal{H}(\xi, t)} = \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi + i\epsilon} + \dots = \mathcal{P} \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi} - i\pi \mathbf{H}(x = \pm \xi, \xi, t) + \dots$

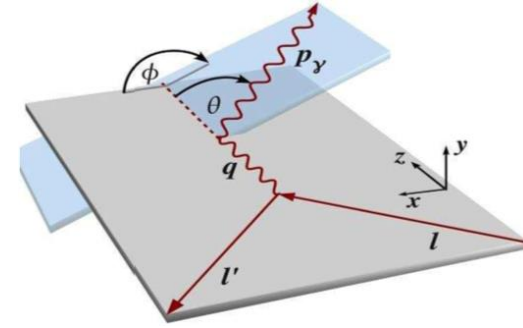
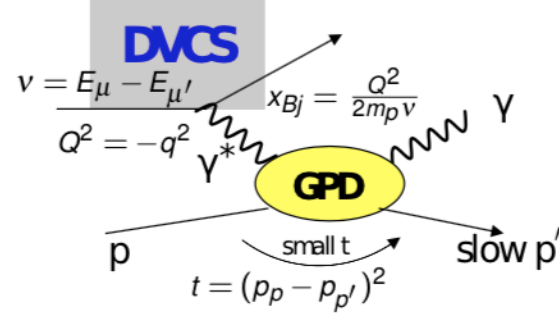
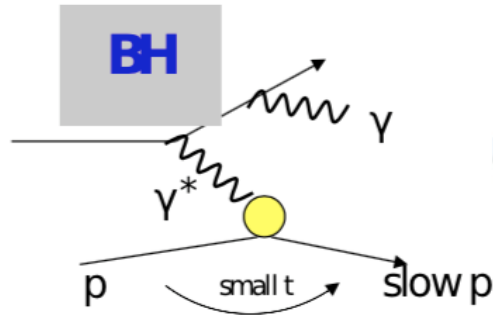
$\text{REAL part} \quad \text{Imaginary part}$

$\text{Re } \mathcal{H}(\xi, t) = \mathcal{P} \int dx \frac{\text{Im } \mathcal{H}(x, t)}{x - \xi} + \Delta(t)$

$\Delta(t) \equiv d_1(t)$ (D-term)

$\leftarrow \text{FT of } H(x, \xi=0, t)$

Azimuthal Dependence of BH & DVCS



$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = d\sigma^{BH} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$

Well known

Beam Charge-spin difference & sum

$$\mathcal{D}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow})$$

$$\mathcal{S}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow})$$

$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$$

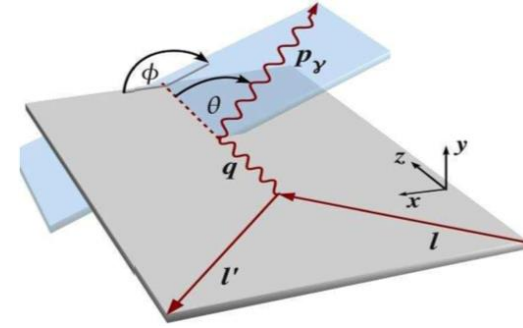
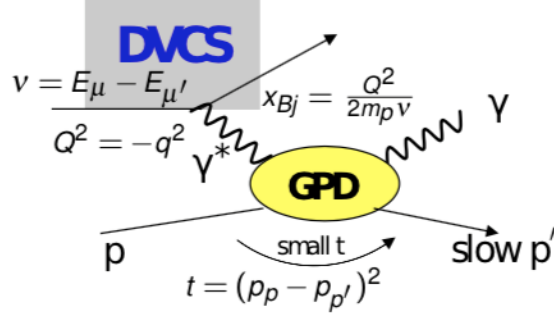
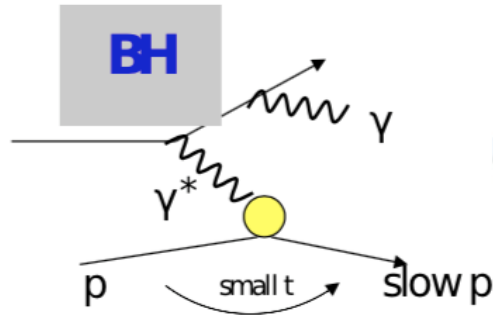
$$d\sigma_{unpol}^{DVCS} \propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$$

$$d\sigma_{pol}^{DVCS} \propto s_1^{DVCS} \sin \phi$$

$$\text{Re } I \propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$$

$$\text{Im } I \propto s_1^I \sin \phi + s_2^I \sin 2\phi$$

Azimuthal Dependence of BH & DVCS



$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = d\sigma^{BH} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$

Well known

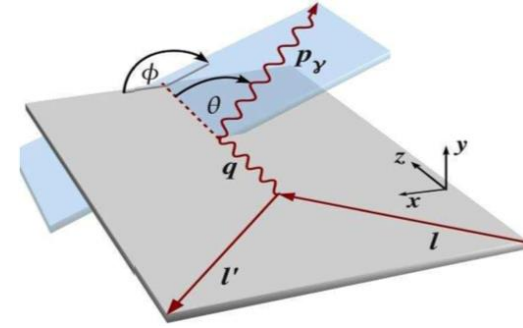
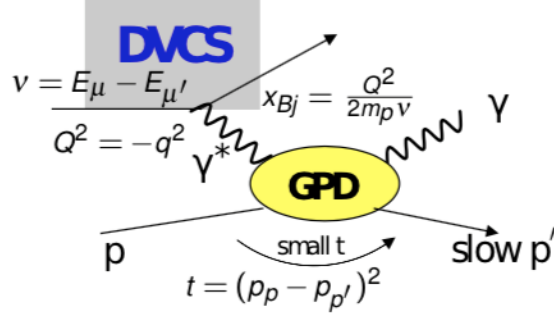
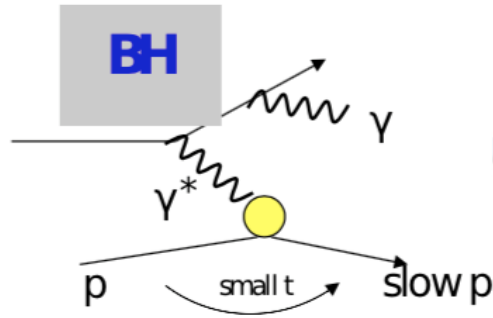
Beam Charge-spin difference & sum

$$\mathcal{D}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow})$$

$$\mathcal{S}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow})$$

$d\sigma^{BH}$	\propto	$c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$	$\mathcal{D}_{CS,U}(\phi)$
$d\sigma_{unpol}^{DVCS}$	\propto	$c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$	
$d\sigma_{pol}^{DVCS}$	\propto	$s_1^{DVCS} \sin \phi$	
$\text{Re } I$	\propto	$c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$	
$\text{Im } I$	\propto	$s_1^I \sin \phi + s_2^I \sin 2\phi$	

Azimuthal Dependence of BH & DVCS



$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = d\sigma^{BH} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$

Well known

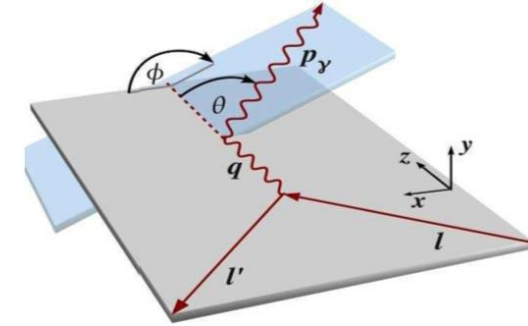
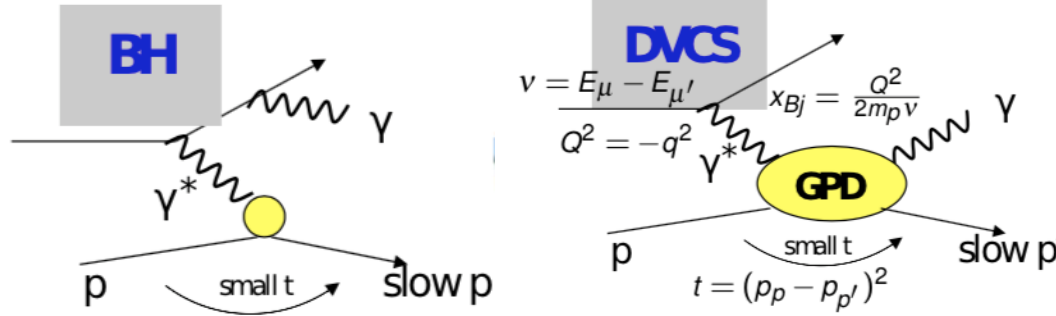
Beam Charge-spin difference & sum

$$\mathcal{D}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow})$$

$$\mathcal{S}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow})$$

$d\sigma^{BH}$	$\propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$	$\mathcal{S}_{CS,U}(\phi)$
$d\sigma_{unpol}^{DVCS}$	$\propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$	
$d\sigma_{pol}^{DVCS}$	$\propto s_1^{DVCS} \sin \phi$	
$\text{Re } I$	$\propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$	
$\text{Im } I$	$\propto s_1^I \sin \phi + s_2^I \sin 2\phi$	$\mathcal{S}_{CS,U}(\phi)$

Azimuthal Dependence of BH & DVCS



$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = d\sigma^{BH}_{\text{Well known}} + \left(d\sigma_{\text{unpol}}^{DVCS} + P_\ell d\sigma_{\text{pol}}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$

Beam Charge-spin difference & sum

$$\mathcal{D}_{CS,u}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow}) \rightarrow \underline{C_1^I \propto \text{Re } \mathcal{F}}$$

More challenging

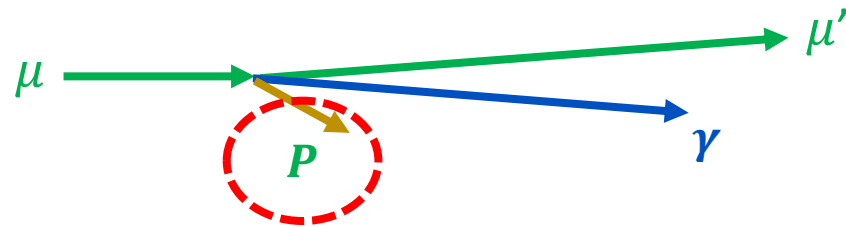
$$\mathcal{S}_{CS,u}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow}) \rightarrow \underline{C_0^{DVCS} \propto (\text{Im } \mathcal{H})^2} \text{ and } \underline{S_1^I \propto \text{Im } \mathcal{F}}$$

Easier to measure

$$\mathcal{F} = F_1 \mathcal{H} + \xi(F_1 + F_2) \mathcal{H} + t/4m^2 F_2 \mathcal{E} \xrightarrow[\text{Small } x_B \text{ at COMPASS}]{\text{Proton Target}} F_1 \mathcal{H}$$

Compton Form factor linked to GPD \mathcal{H}

COMPASS 2016 Preliminary Results



$$\Delta\phi = \phi^{\text{cam.}} - \phi^{\text{spec.}}$$

Proton azimuthal angle

$$\Delta p_T = |p_T^{\text{cam.}}| - |p_T^{\text{spec.}}|$$

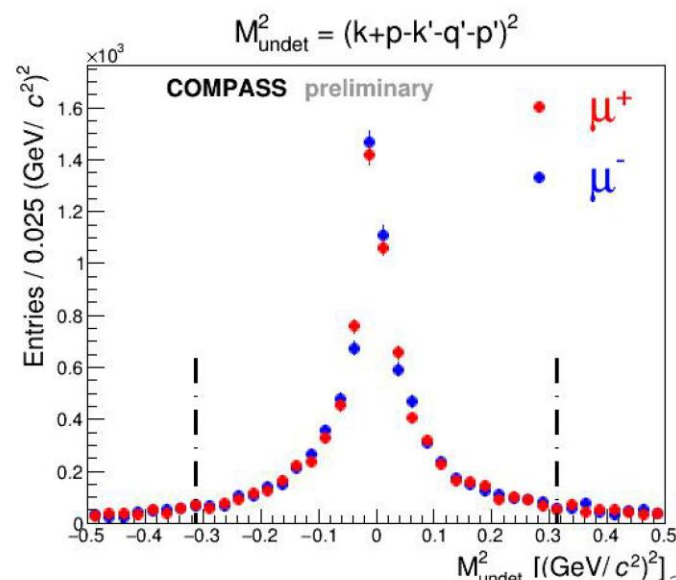
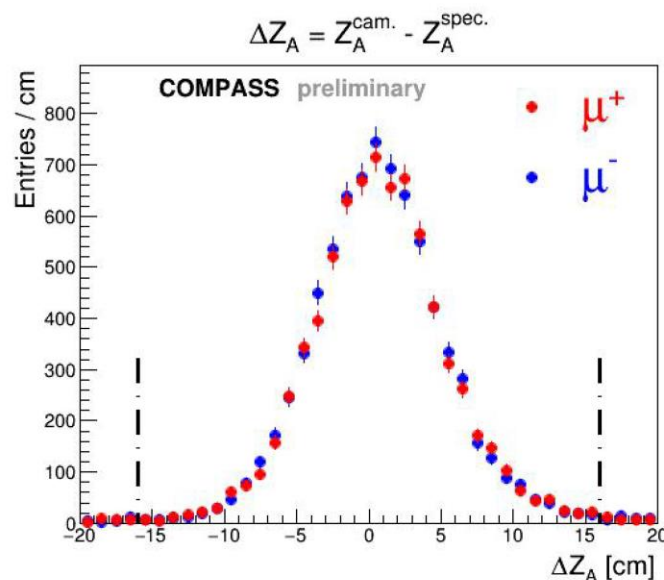
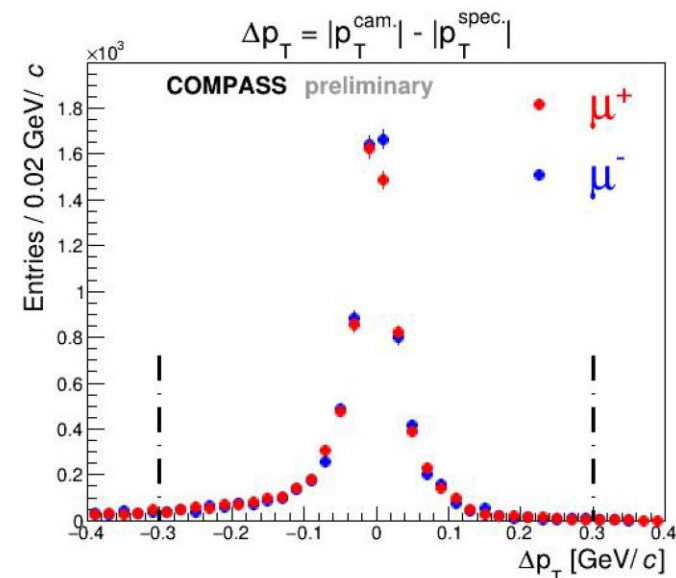
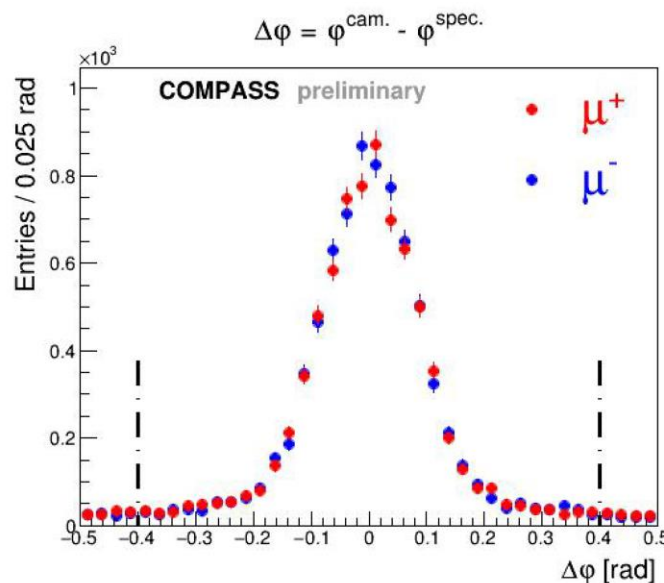
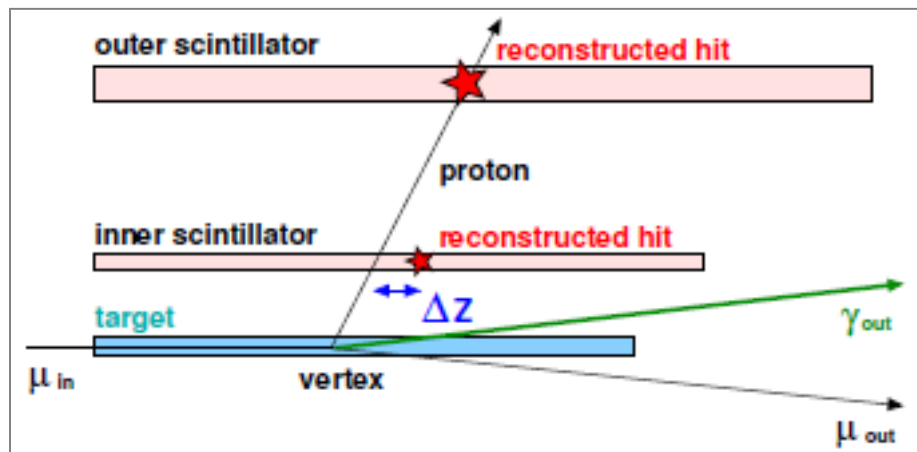
Proton momentum

$$\Delta z_A = z_A^{\text{cam.}} - z_A^{\text{spec.}}$$

Proton track position

$$M_{\text{undet}}^2 = (k + p - k' - q' - p')^2$$

Energy momentum balance



➤ Main background of exclusive single photon events: π^0 decay

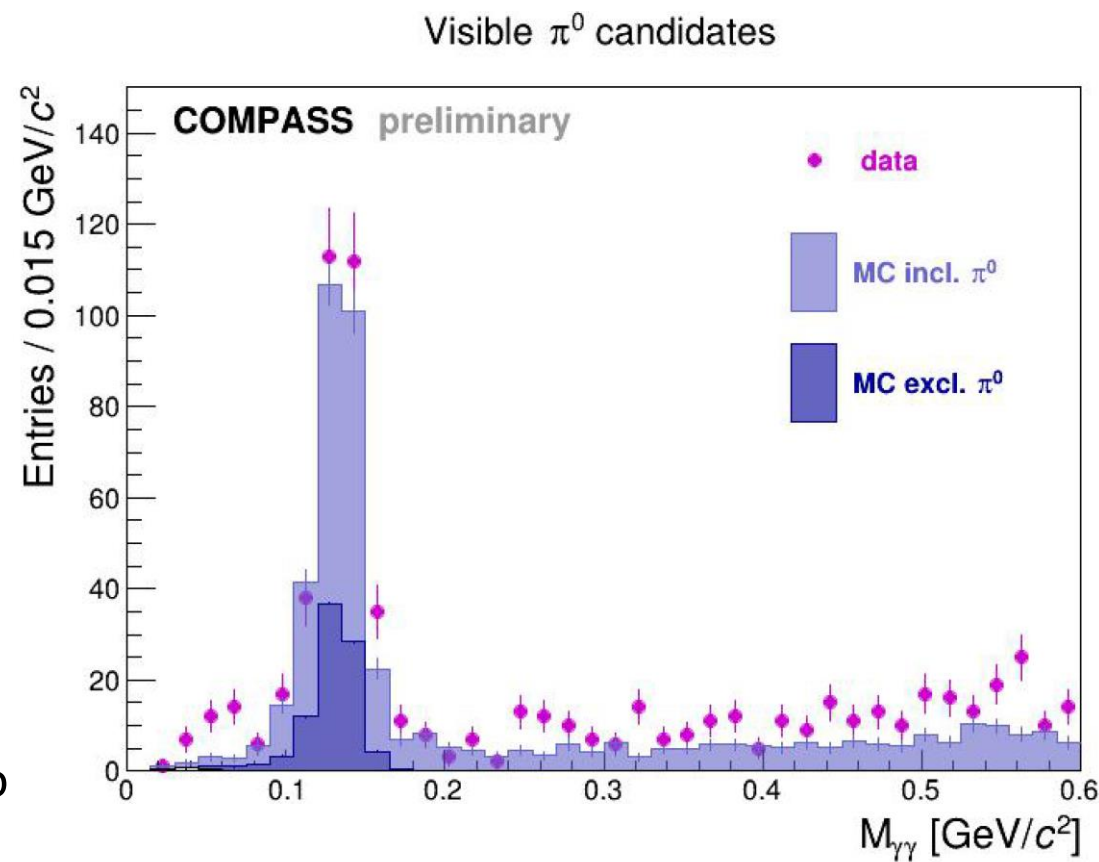
➤ Visible (both γ detected) – subtracted

A high-energy DVCS photon candidate is combined with all detected photons with energies lower than the DVCS threshold: (4,5) GeV in Ecal (0,1) respectively

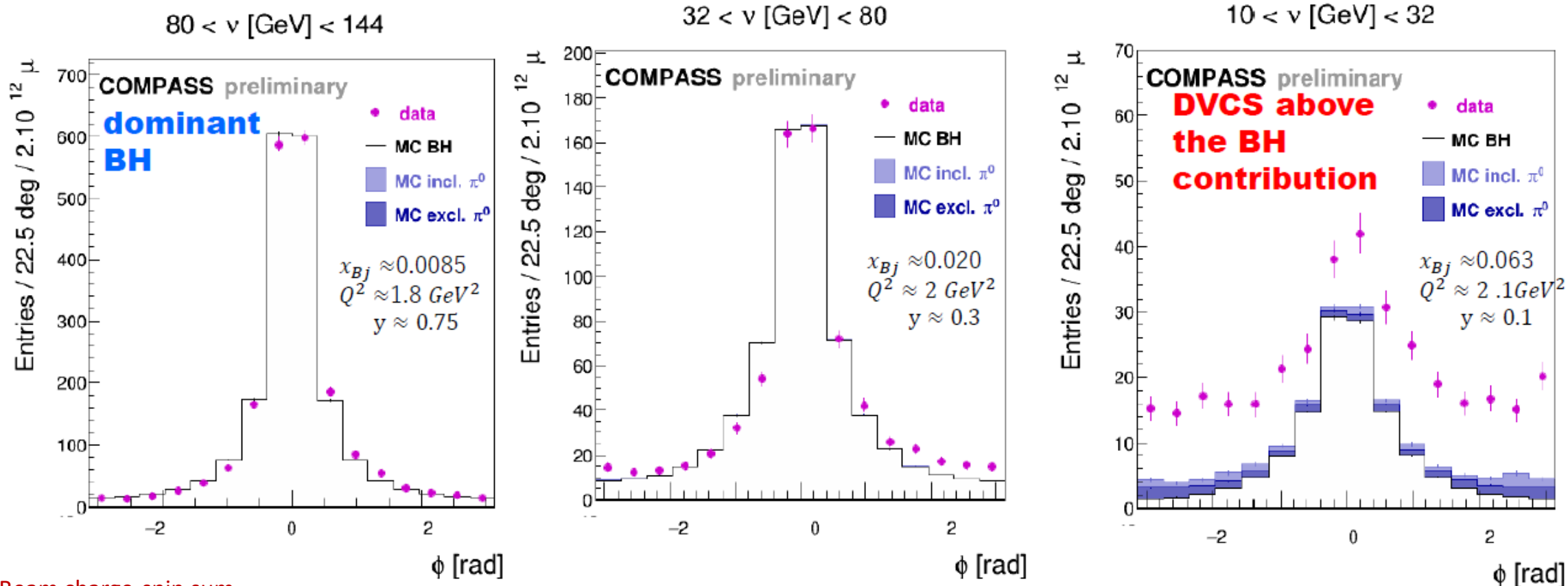
➤ Invisible (one γ lost) – estimated by MC

- Semi-inclusive LEPTO 6.1
- Exclusive HEPGEN π^0 (GK model)

The sum of LEPTO and HEPGEN contributions is normalized to the π^0 peak in $M_{\gamma\gamma}$ of the real data



COMPASS 2016 Preliminary Results



➤ Beam charge-spin sum

$$S_{CS, U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) + d\sigma(\mu^{-\rightarrow}) = 2[d\sigma^{BH} + d\sigma_{unpol}^{DVCS} + \text{Im } I]$$

$$= 2[d\sigma^{BH} + c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi + s_1^I \sin \phi + s_2^I \sin 2\phi]$$

$$c_0^{DVCS} \underset{\text{small } x_{Bj}}{\propto} 4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) + \frac{t}{M^2} \mathcal{E}\mathcal{E}^* \rightarrow 4 (\text{Im } \mathcal{H})^2$$

model dependent

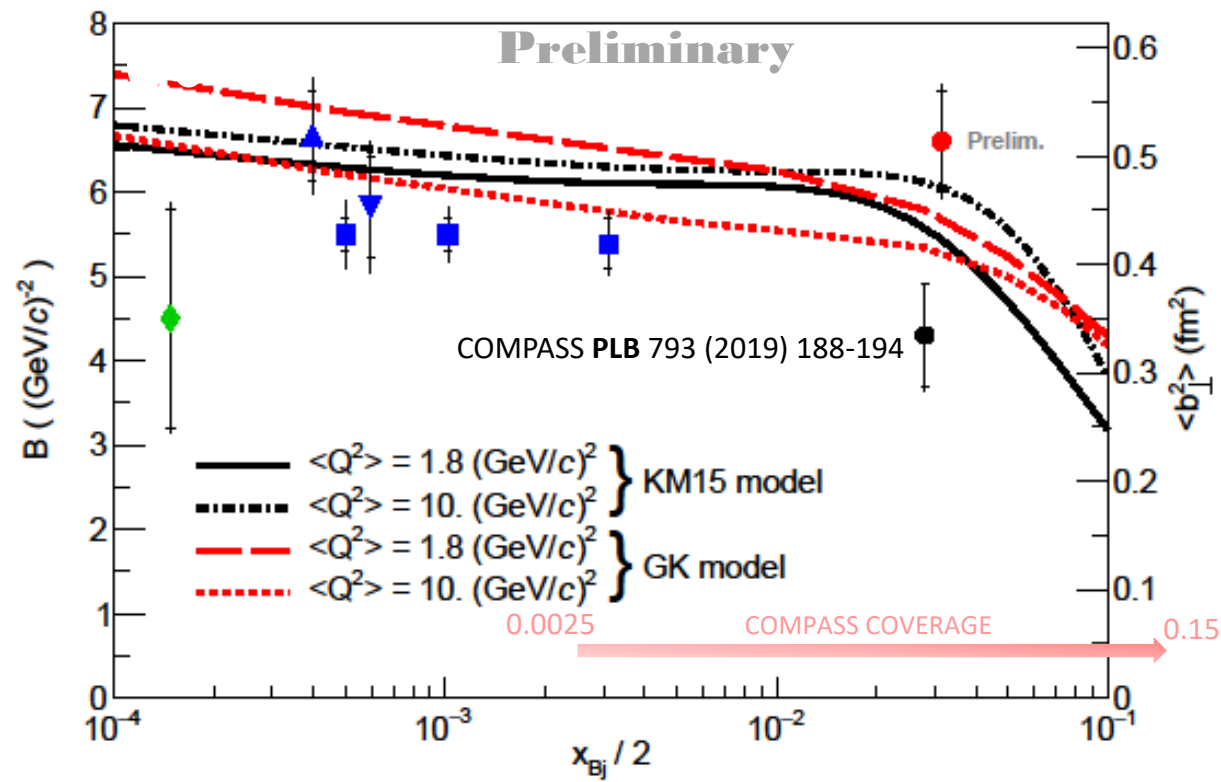
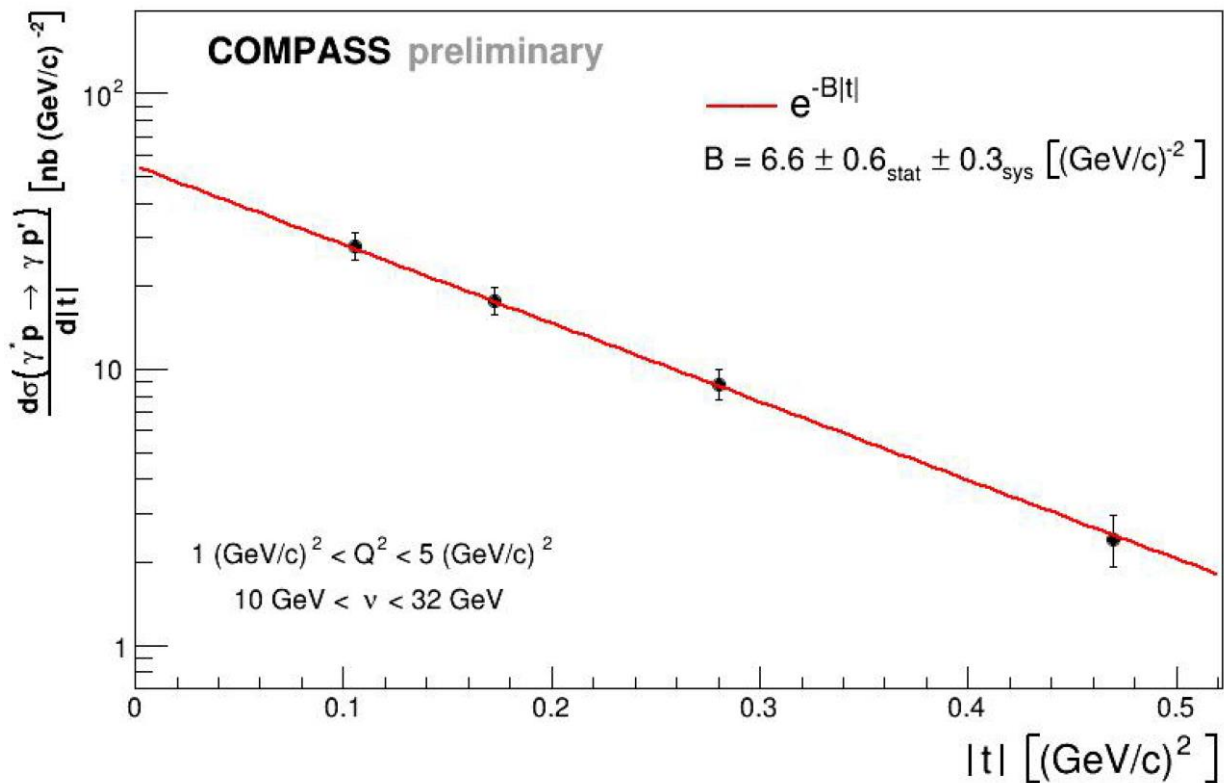
Transverse extension of partons – 2016 data



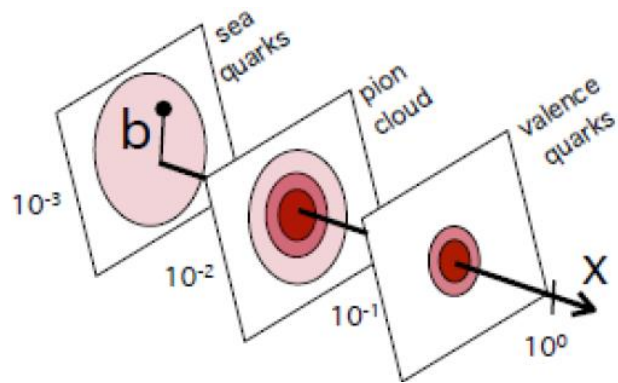
$$d\sigma^{DVCS}/d|t| \propto e^{-B|t|}$$

$$\langle r_{\perp}^2(x_B) \rangle \approx 2B(x_B) \quad \text{At small } x_B$$

- COMPASS: $\langle Q^2 \rangle = 1.8 \text{ (GeV/c)}^2$
- ◆ ZEUS: $\langle Q^2 \rangle = 3.2 \text{ (GeV/c)}^2$
- ▲ H1: $\langle Q^2 \rangle = 4.0 \text{ (GeV/c)}^2$
- ▼ H1: $\langle Q^2 \rangle = 8.0 \text{ (GeV/c)}^2$
- H1: $\langle Q^2 \rangle = 10. \text{ (GeV/c)}^2$



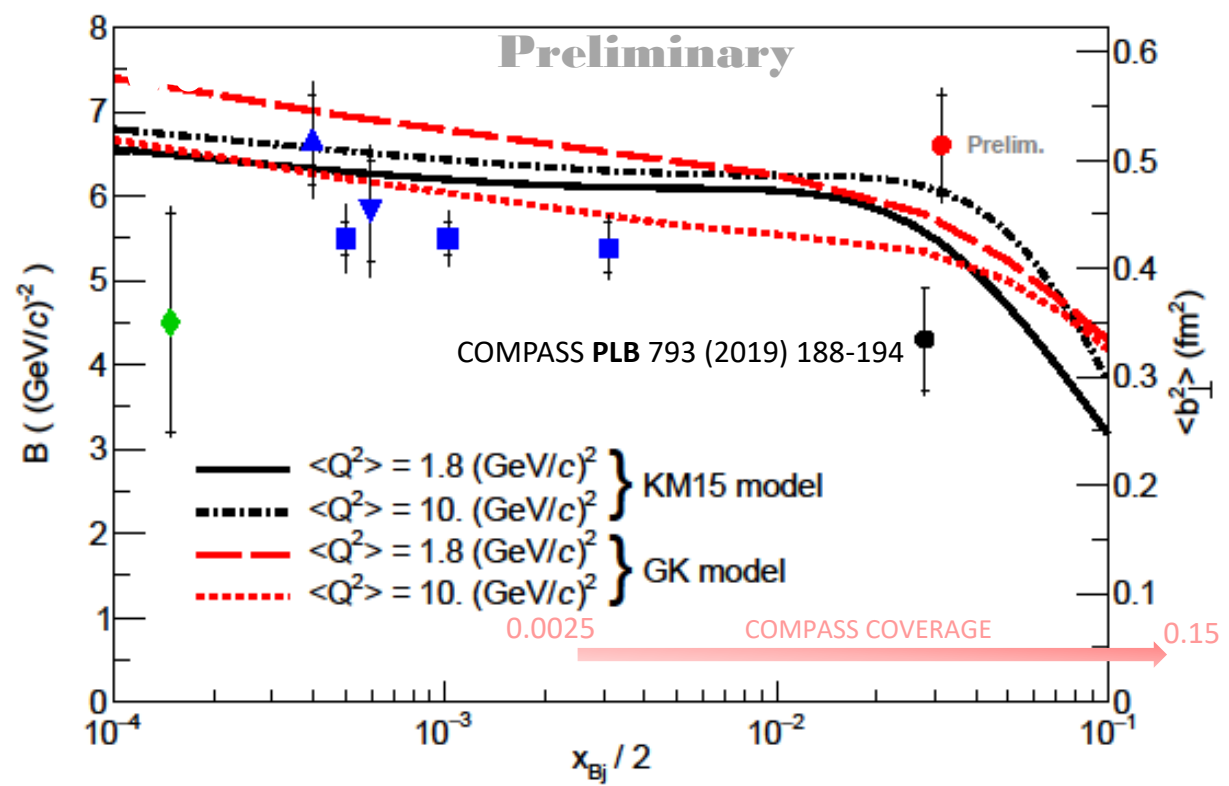
Transverse extension of partons – 2016 data



- Improvements in the 2016 analysis, relative to 2012
 - μ^+ and μ^- beams at same intensity
 - More advanced analysis with 2016 data, ongoing
 - Improved π^0 contamination estimation
 - Better MC description in ν

$$\langle r_{\perp}^2(x_B) \rangle \approx 2B(x_B) \quad \text{At small } x_B$$

- COMPASS: $\langle Q^2 \rangle = 1.8 \text{ (GeV/c)}^2$
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- ▼ H1: $\langle Q^2 \rangle = 8.0 \text{ (GeV/c)}^2$
- H1: $\langle Q^2 \rangle = 10. \text{ (GeV/c)}^2$



➤ The transverse-size evolution as a function of x_{Bj} → Expect at least 3 x_{Bj} bins from 2016-17 data

Beam Charge-spin Difference



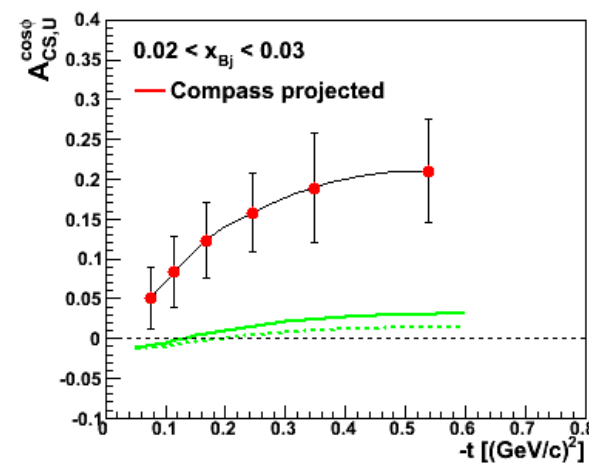
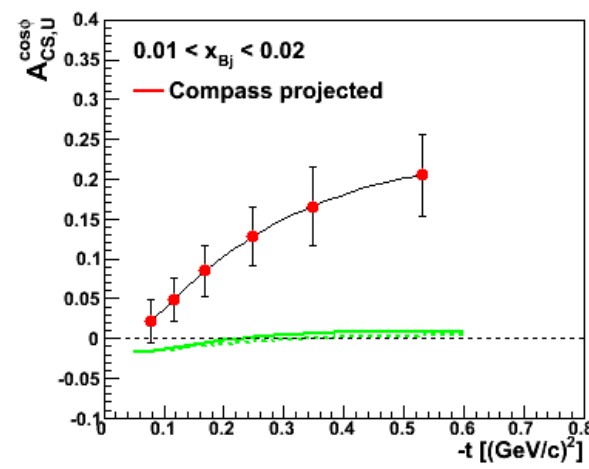
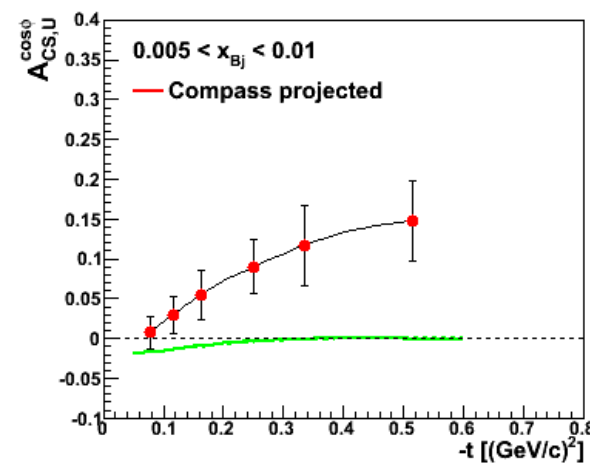
$$\mathcal{D}_{CS,U}(\phi) \equiv d\sigma(\mu^{+\leftarrow}) - d\sigma(\mu^{-\rightarrow}) \rightarrow c_0^I + c_1^I \cos \phi$$

$$BCSA = \mathcal{D}_{CS,U} / S_{CS,U} = A_0 + A_{CS,U}^{\cos\phi} \cos\phi + A_2 \cos 2\phi$$

$$c_1^I \rightarrow \text{Re } F_1 \mathcal{H}$$

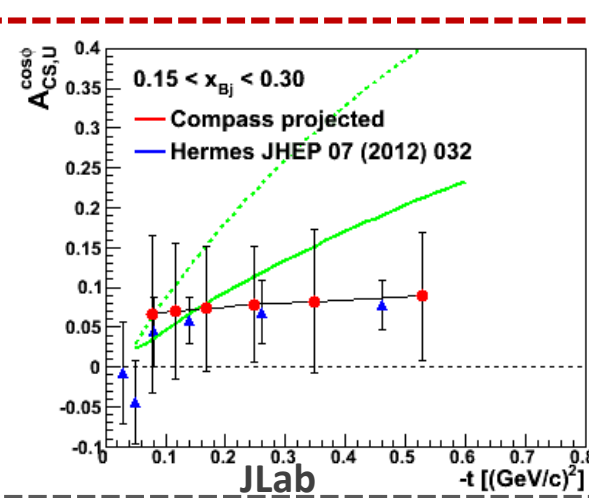
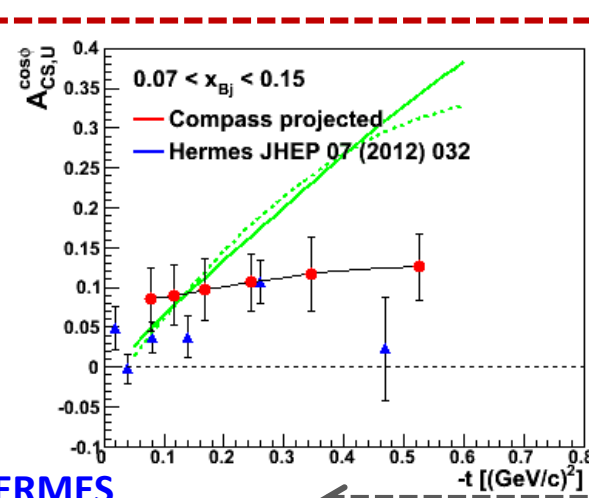
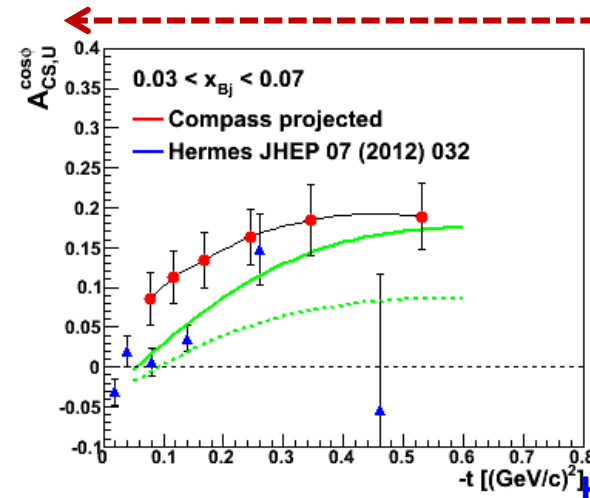
VGG

KM10 – fit to data



➤ With $\text{Re } F_1 \mathcal{H}$ and $\text{Im } F_1 \mathcal{H}$
 → Extraction of **D-term**

$\text{Re } \mathcal{H} > 0$ at H1
 < 0 at HERMES
 Value of x_{Bj} for the node?



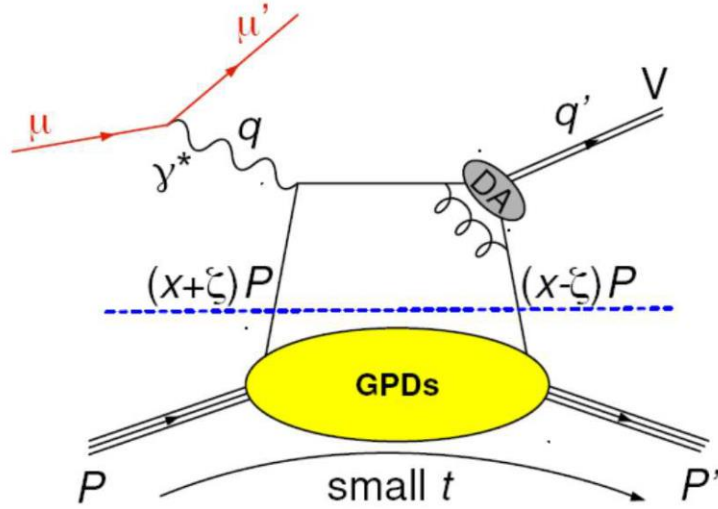
COMPASS 2 years of data $E_\mu = 160 \text{ GeV}$ $1 < Q^2 < 8 \text{ GeV}^2$

**Hard Exclusive Meson Production
@ COMPASS**

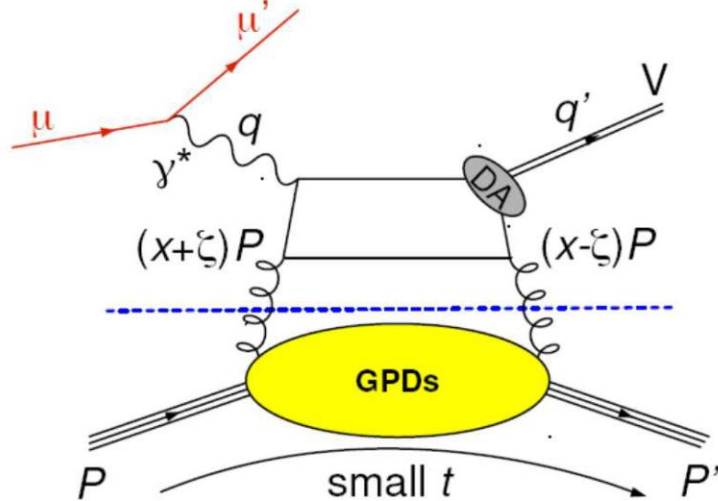
GPDs in Hard Exclusive Meson Production



quark contribution



gluon contribution



4 chiral-even GPDs: helicity of parton unchanged

$$\mathbf{H}^q(x, \xi, t) \quad \mathbf{E}^q(x, \xi, t) \quad \rightarrow \text{Vector Meson}$$

$$\tilde{\mathbf{H}}^q(x, \xi, t) \quad \tilde{\mathbf{E}}^q(x, \xi, t) \quad \rightarrow \text{Pseudo-Scalar Meson}$$

+ 4 chiral-odd (transversity) GPDs: helicity of parton changed
(not possible in DVCS)

$$\mathbf{H}_T^q(x, \xi, t)$$

$$\mathbf{E}_T^q(x, \xi, t)$$

$$\tilde{\mathbf{H}}_T^q(x, \xi, t)$$

$$\tilde{\mathbf{E}}_T^q(x, \xi, t)$$

$$\bar{\mathbf{E}}_T^q = 2 \tilde{\mathbf{H}}_T^q + \mathbf{E}_T^q$$

- Ability to probe the chiral-odd GPDs.
- Universality of GPDs, quark flavor filter
- In addition to nuclear structure, provide insights into reaction mechanism.
- Additional non-perturbative term from meson wave function.

Exclusive π^0 Production on Unpolarized Proton



$\mu p \rightarrow \mu \pi^0 p$

$$\frac{d^2\sigma}{dt d\phi_\pi} = \frac{1}{2\pi} \left[\left(\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right) + \epsilon \cos 2\phi_\pi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_\pi \frac{d\sigma_{LT}}{dt} \right]$$

ϵ : degree of longitudinal polarization

$$\frac{d\sigma_L}{dt} = \frac{4\pi\alpha}{k'} \frac{1}{Q^6} \left\{ (1-\xi^2) |\langle \tilde{H} \rangle|^2 - 2\xi^2 \text{Re} [\langle \tilde{H} \rangle^* \langle \tilde{E} \rangle] - \frac{t'}{4m^2} \xi^2 |\langle \tilde{E} \rangle|^2 \right\}$$

Leading twist expected be dominant
But measured as \approx only a few % of $\frac{d\sigma_T}{dt}$

The other contributions arise from coupling between chiral-odd (quark helicity flip) GPDs to the **twist-3** pion amplitude

$$\frac{d\sigma_T}{dt} = \frac{4\pi\alpha}{2k'} \frac{\mu_\pi^2}{Q^8} \left[(1-\xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8m^2} |\langle \bar{E}_T \rangle|^2 \right]$$

$$\frac{d\sigma_{LT}}{dt} = \frac{4\pi\alpha}{\sqrt{2}k'} \frac{\mu_\pi}{Q^7} \xi \sqrt{1-\xi^2} \frac{\sqrt{-t'}}{2m} \text{Re} [\langle H_T \rangle^* \langle \tilde{E} \rangle]$$

$$\frac{d\sigma_{TT}}{dt} = \frac{4\pi\alpha}{k'} \frac{\mu_\pi^2}{Q^8} \frac{t'}{16m^2} |\langle \bar{E}_T \rangle|^2$$

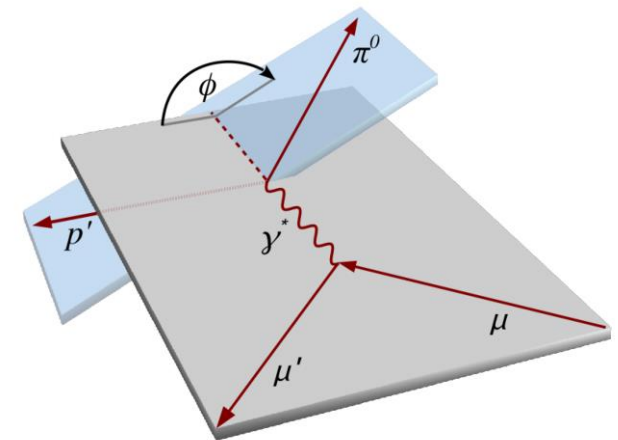


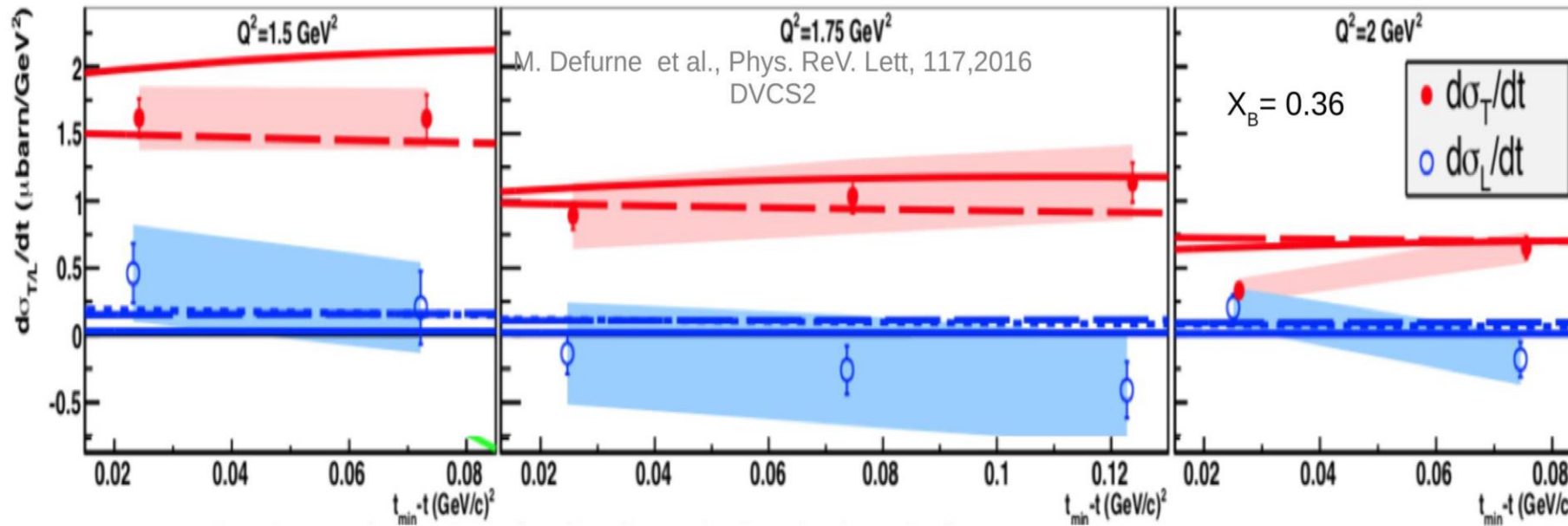
Fig: M.G. Alexeev et al. *Phys.Lett.B* 805 (2020)

Exclusive π^0 Production on Unpolarized Proton



$\mu p \rightarrow \mu \pi^0 p$

$d^2\sigma$ $1 / [(d\sigma_T \quad d\sigma_L) \quad d\sigma_{TT} \quad d\sigma_{LT}]$



— S. V. Goloskokov and P. Kroll, Eur. Phys. J. C65:137 (2010)

- - - G. R. Goldstein, J. O. Hernandez, S. Liuti, Phys. Rev. D84 (2011)

$$\frac{\sigma_{LT}}{dt} = \frac{4\pi\alpha}{\sqrt{2}k'} \frac{\mu_\pi}{Q^7} \xi \sqrt{1-\xi^2} \frac{\sqrt{-t'}}{2m} \operatorname{Re} \left[\langle H_T \rangle \langle \tilde{E} \rangle \right]$$

$$\frac{\sigma_{TT}}{dt} = \frac{4\pi\alpha}{k'} \frac{\mu_\pi^2}{Q^8} \frac{t'}{16m^2} |\langle \bar{E}_T \rangle|^2$$

tion

Exclusive π^0 Production on Unpolarized Proton



$\mu p \rightarrow \mu \pi^0 p$

$$\frac{d^2\sigma}{dt d\phi_\pi} = \frac{1}{2\pi} \left[\left(\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right) + \epsilon \cos 2\phi_\pi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_\pi \frac{d\sigma_{LT}}{dt} \right]$$

ϵ : degree of longitudinal polarization

$$\frac{d\sigma_L}{dt} = \frac{4\pi\alpha}{k'} \frac{1}{Q^6} \left\{ (1-\xi^2) |\langle \tilde{H} \rangle|^2 - 2\xi^2 \text{Re} [\langle \tilde{H} \rangle^* \langle \tilde{E} \rangle] - \frac{t'}{4m^2} \xi^2 |\langle \tilde{E} \rangle|^2 \right\}$$

Leading twist expected to be dominant
But measured as \approx only a few % of $\frac{d\sigma_T}{dt}$

The other contributions arise from the coupling between chiral-odd (quark helicity flip) GPDs to the **twist-3** pion amplitude

$$\frac{d\sigma_T}{dt} = \frac{4\pi\alpha}{2k'} \frac{\mu_\pi^2}{Q^8} \left[(1-\xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8m^2} |\langle \bar{E}_T \rangle|^2 \right]$$

$$\frac{d\sigma_{LT}}{dt} = \frac{4\pi\alpha}{\sqrt{2}k'} \frac{\mu_\pi}{Q^7} \xi \sqrt{1-\xi^2} \frac{\sqrt{-t'}}{2m} \text{Re} [\langle H_T \rangle^* \langle \tilde{E} \rangle]$$

$$\frac{d\sigma_{TT}}{dt} = \frac{4\pi\alpha}{k'} \frac{\mu_\pi^2}{Q^8} \frac{t'}{16m^2} |\langle \bar{E}_T \rangle|^2$$

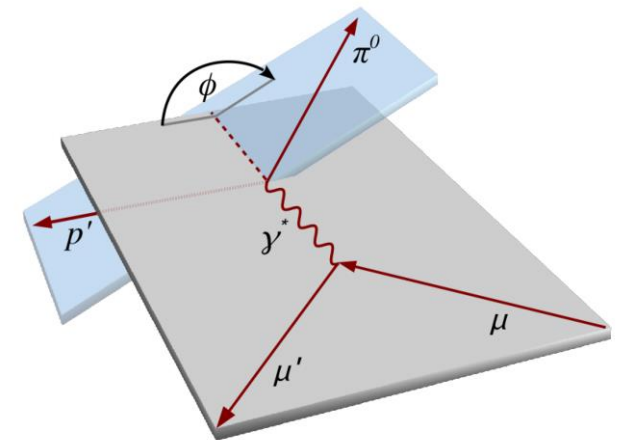
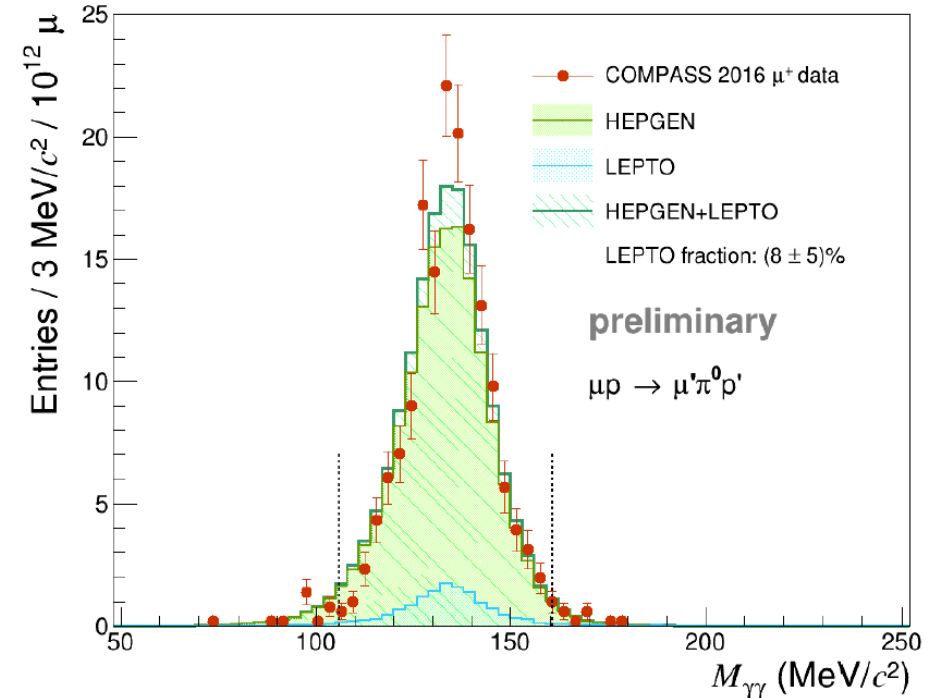
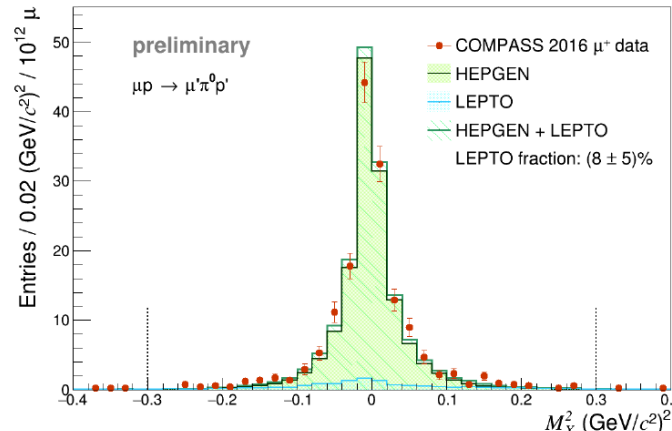
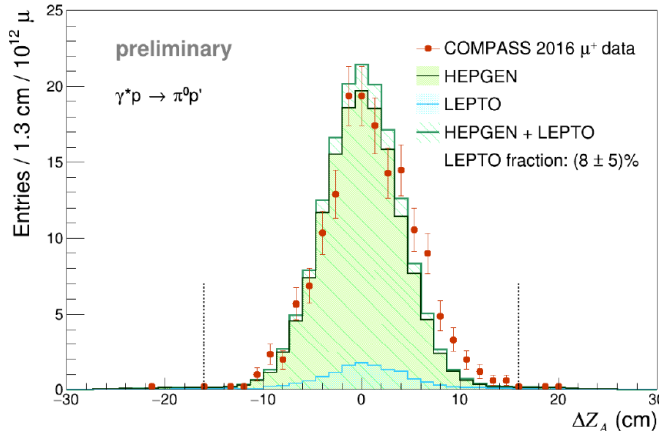
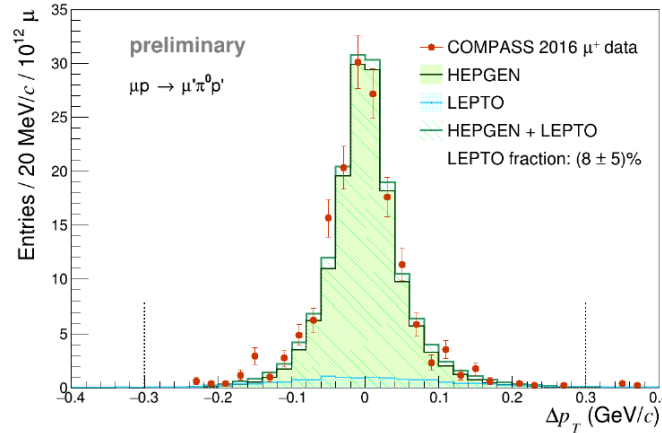
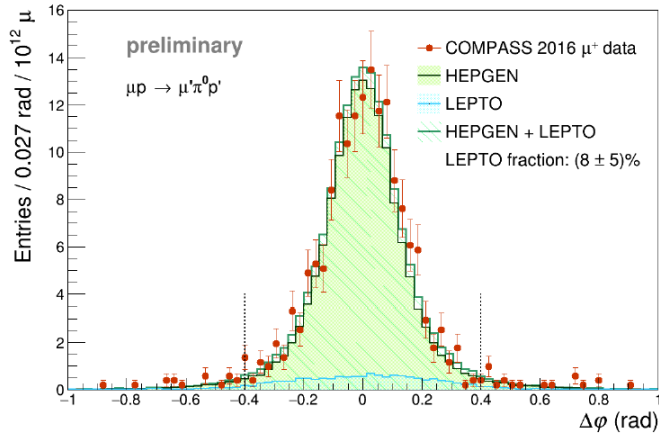


Fig: M.G. Alexeev et al. *Phys.Lett.B* 805 (2020)

Exclusive π^0 Selection and Background Estimation

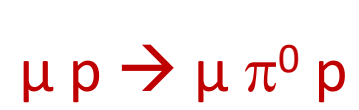


- Exclusivity ensured by cuts on *exclusivity variables*, similar to DVCS.
- Background fraction determined by fitting the exclusivity variables with Monte Carlo simulations.
 - *LEPTO* for non-exclusive background
 - *HEPGEN* of exclusive π^0 for signal

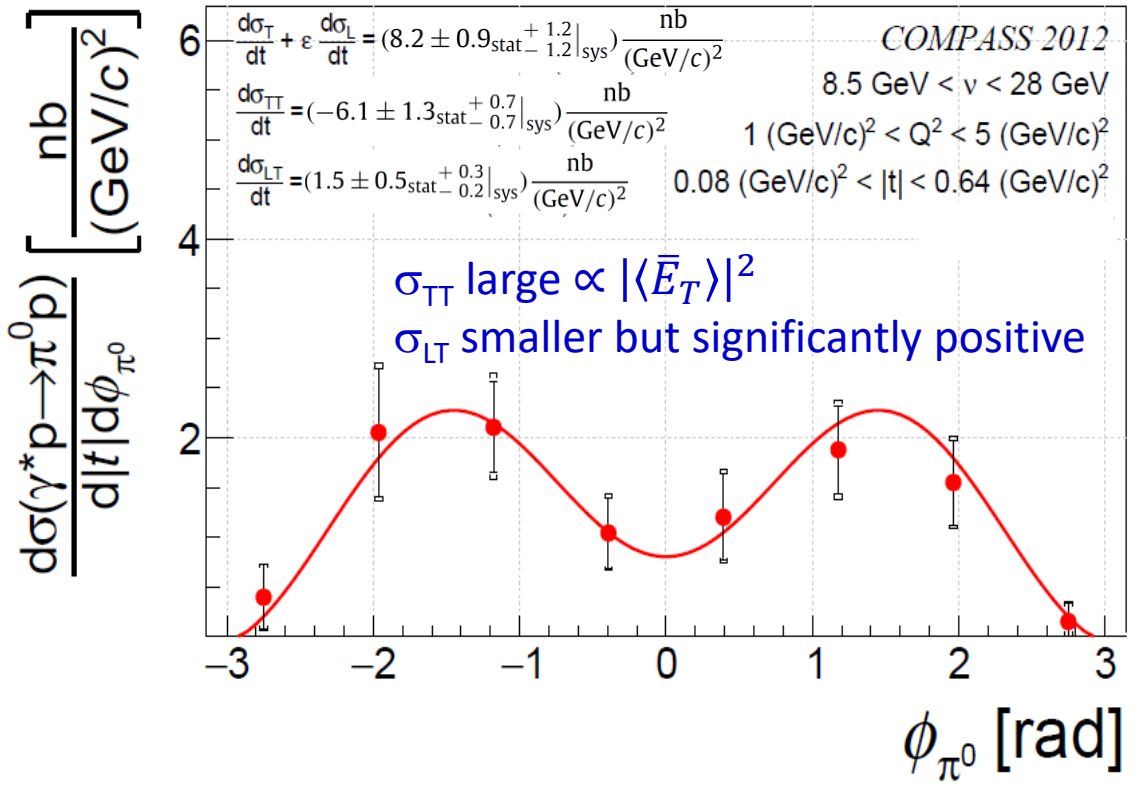


- In 2016 data, non-exclusive background fraction in data $\rightarrow 8 \pm 5\%$

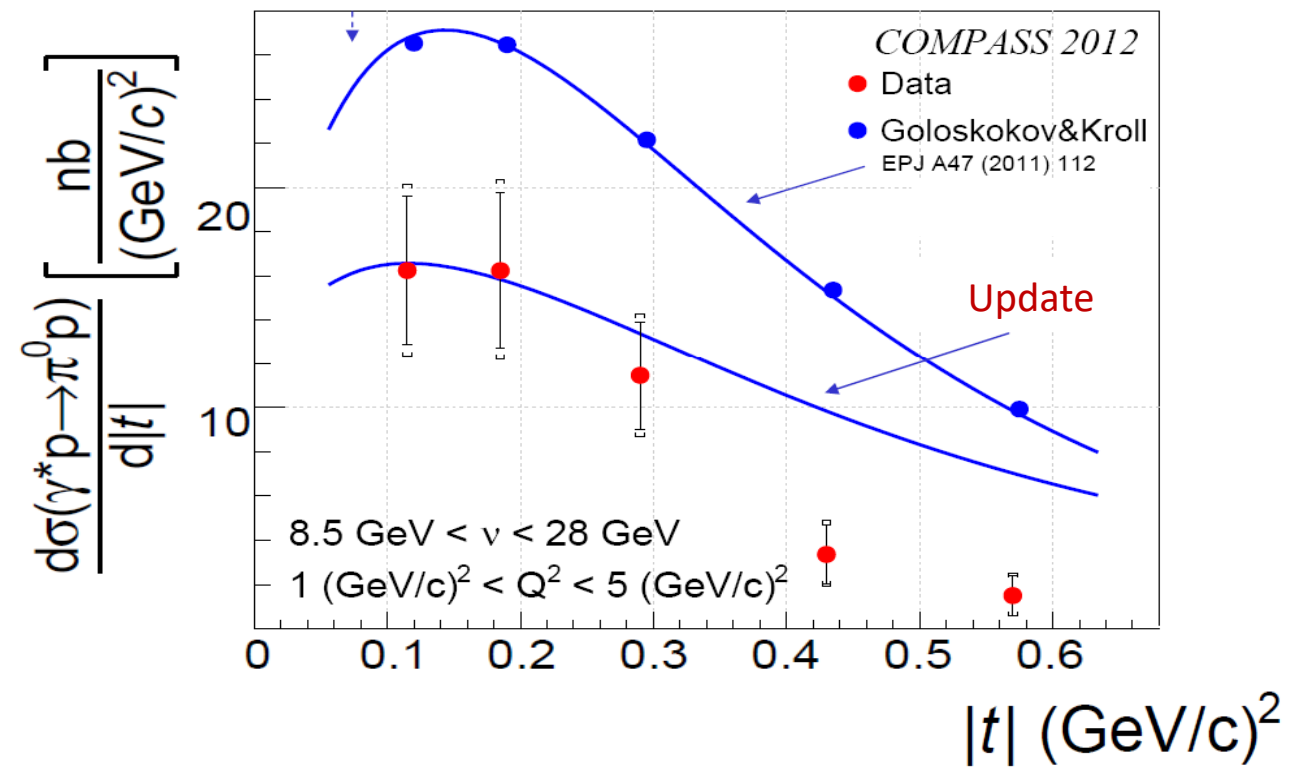
2012 Exclusive π^0 Prod. on Unpolarized Proton



$$\frac{d^2\sigma}{dt d\phi_\pi} = \frac{1}{2\pi} \left[\left(\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right) + \epsilon \cos 2\phi_\pi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_\pi \frac{d\sigma_{LT}}{dt} \right]$$



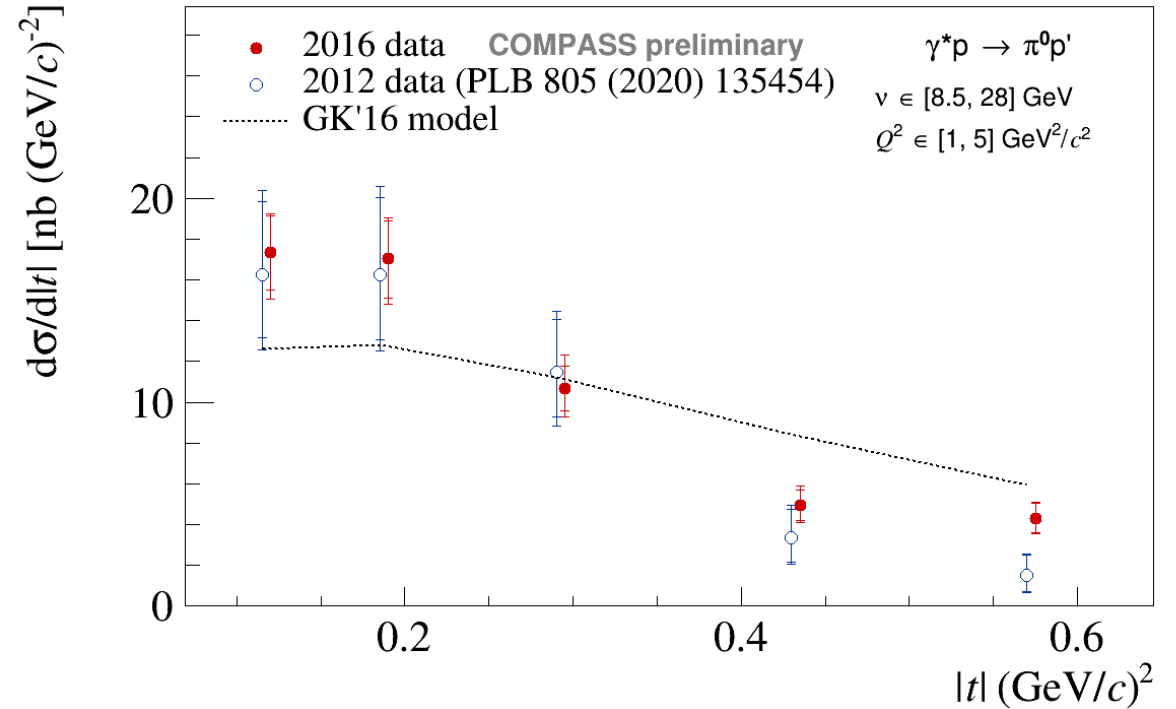
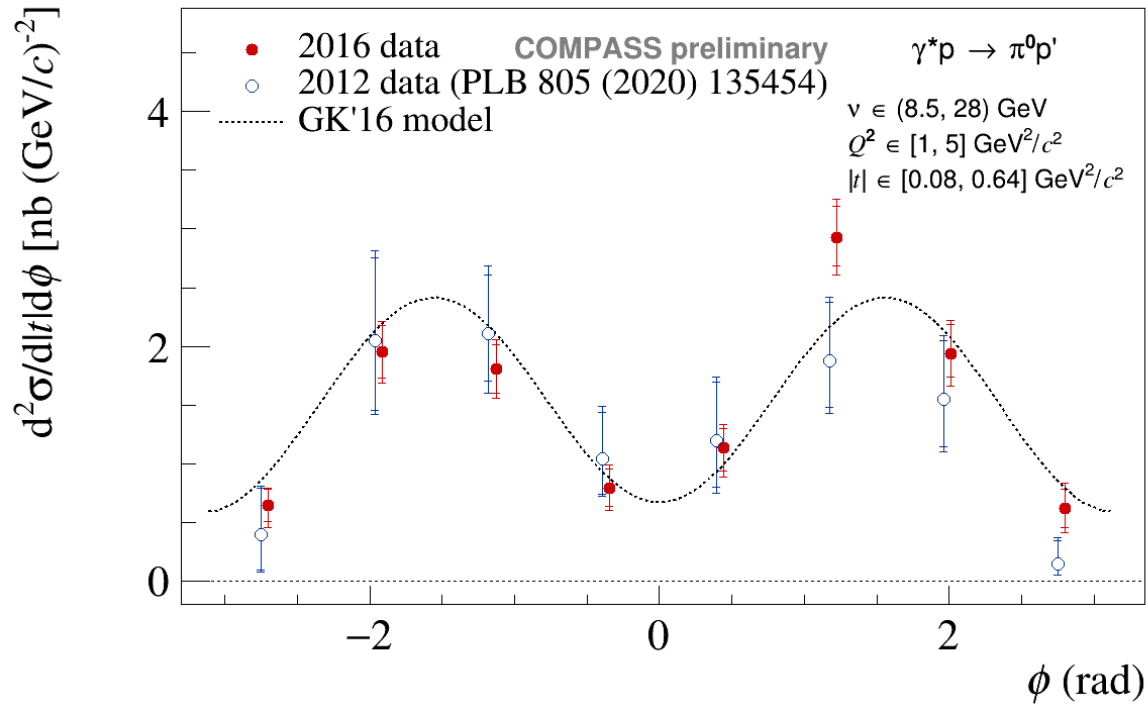
A dip at small t would indicate the significance of \bar{E}_T



2016 Exclusive π^0 Prod. on Unpolarized Proton



➤ **New 2016 data release:** statistics about 2.3 times larger than the published 2012 pilot run.

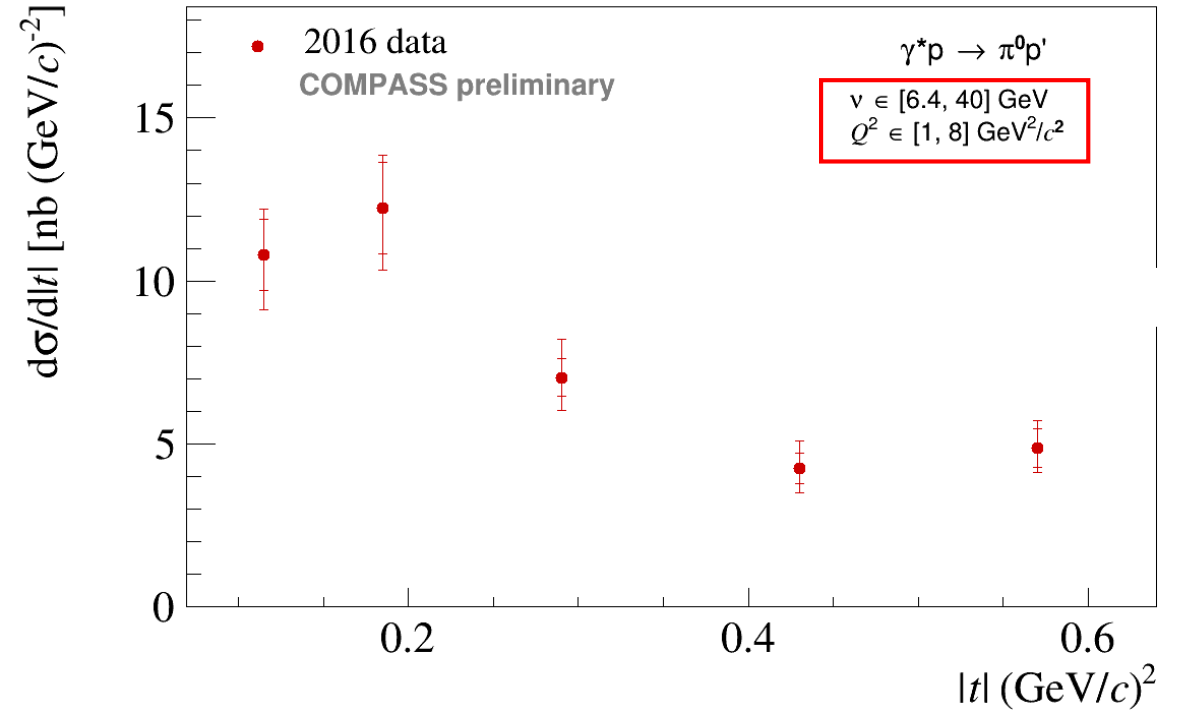
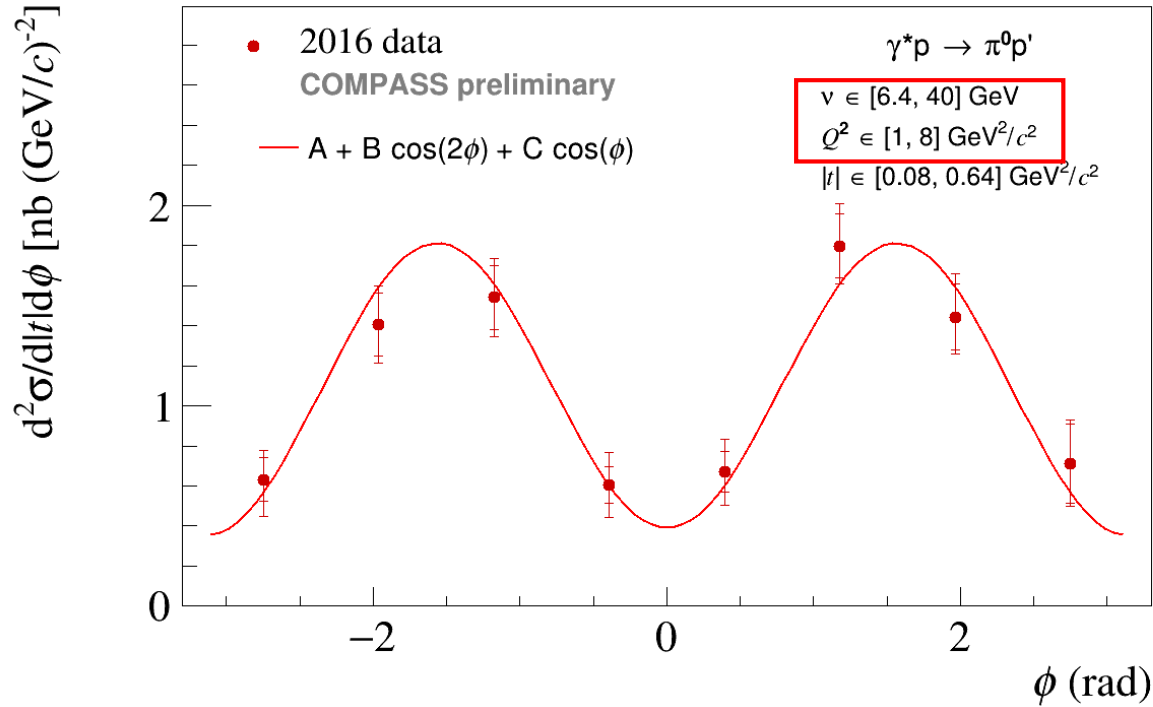


➤ **Agree with previous measurements, with better uncertainty**

2016 Exclusive π^0 Prod. on Unpolarized Proton



➤ **New 2016 data release:** statistics about 2.3 times larger than the published 2012 pilot run.



➤ **Larger (ν, Q^2) domain achievable.**

$$\nu \in [8.5, 28] \rightarrow [6.4, 40] \text{ GeV}$$

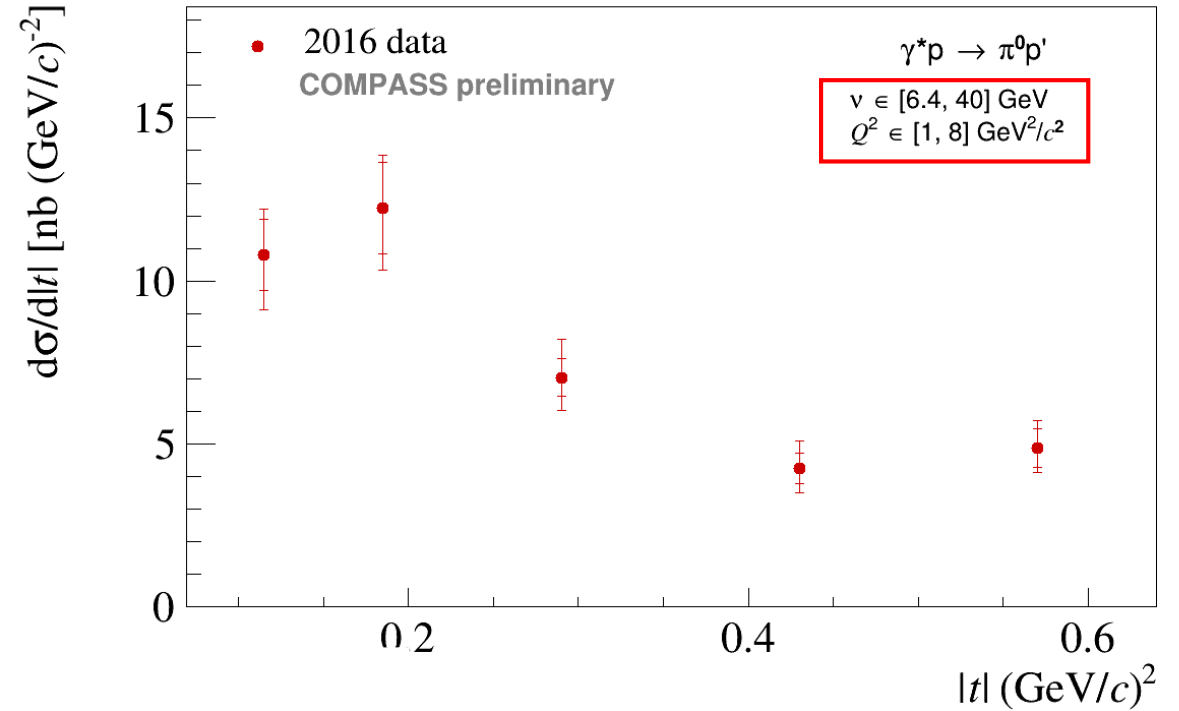
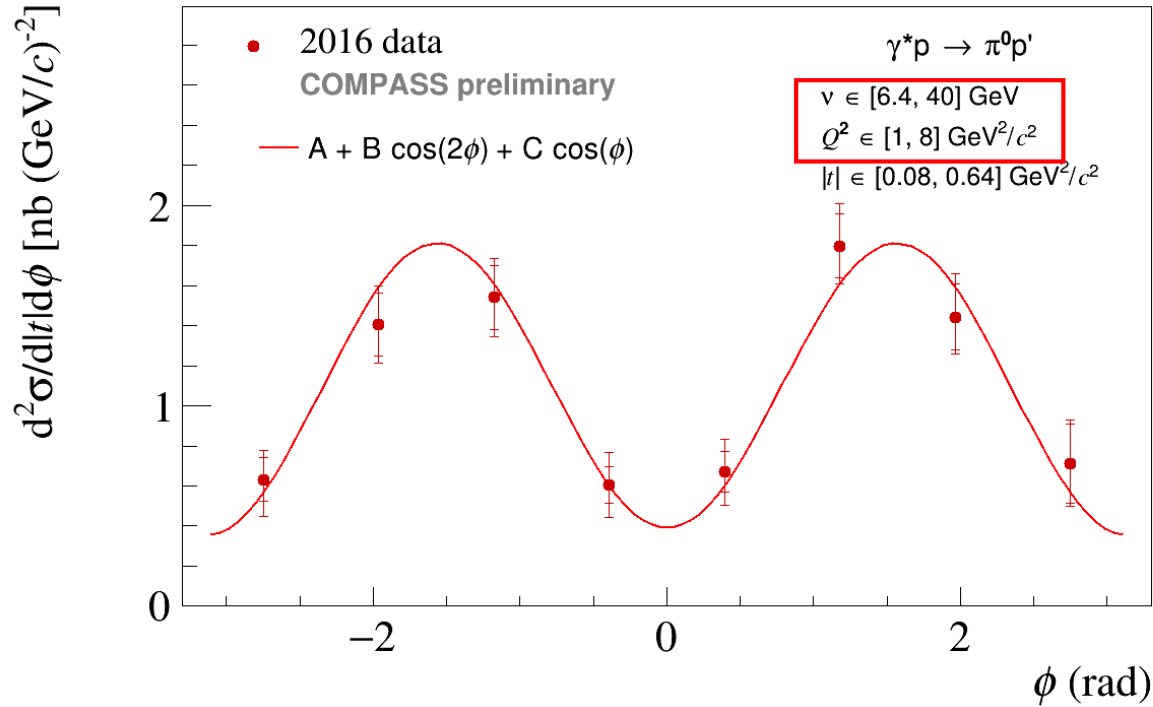
$$Q^2 \in [1, 5] \rightarrow [1, 8] \text{ GeV}^2/c^2$$

$$|t| \in [0.08, 0.64] \text{ GeV}^2/c^2$$

2016 Exclusive π^0 Prod. on Unpolarized Proton



➤ **New 2016 data release:** statistics about 2.3 times larger than the published 2012 pilot run.



$$\left\langle \frac{\sigma_T}{|t|} + \epsilon \frac{\sigma_L}{|t|} \right\rangle = (6.9 \pm 0.3_{\text{stat}} \pm 0.8_{\text{syst}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{\sigma_{TT}}{|t|} \right\rangle = (-4.5 \pm 0.5_{\text{stat}} \pm 0.2_{\text{syst}}) \frac{\text{nb}}{(\text{GeV}/c)^2} \rightarrow \bar{E}_T$$

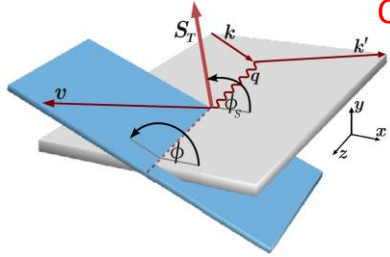
$$\left\langle \frac{\sigma_{LT}}{|t|} \right\rangle = (0.06 \pm 0.2_{\text{stat}} \pm 0.1_{\text{syst}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

- Main systematic uncertainty comes from the evaluation of the π^0 background from SIDIS
- We will provide the evolution with 3 bins in ν and 4 bins in Q^2

2007 & 2010 HEMP with Transversely Polarized Target



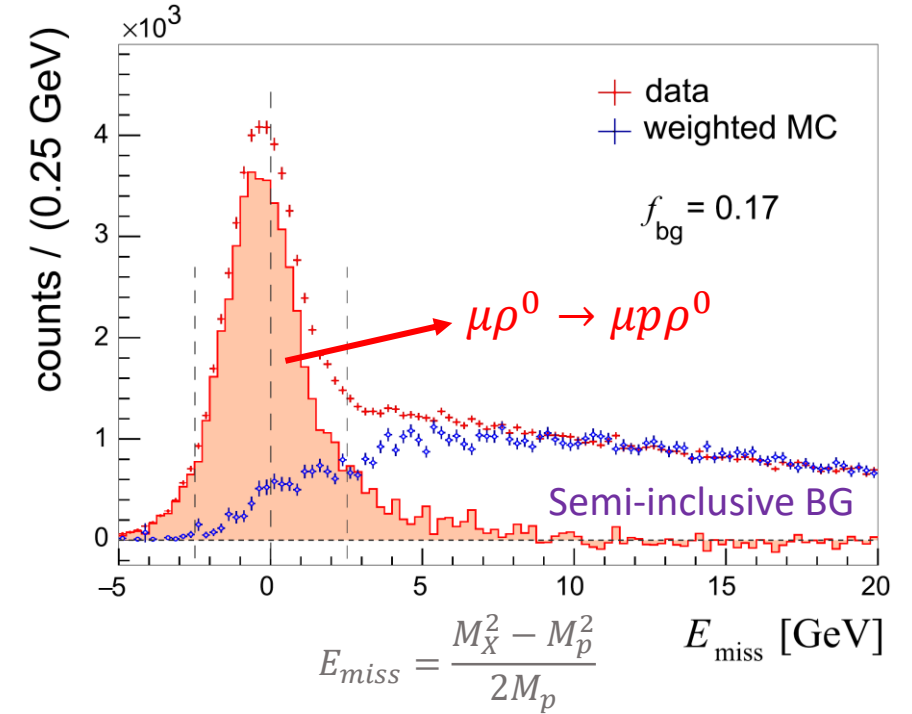
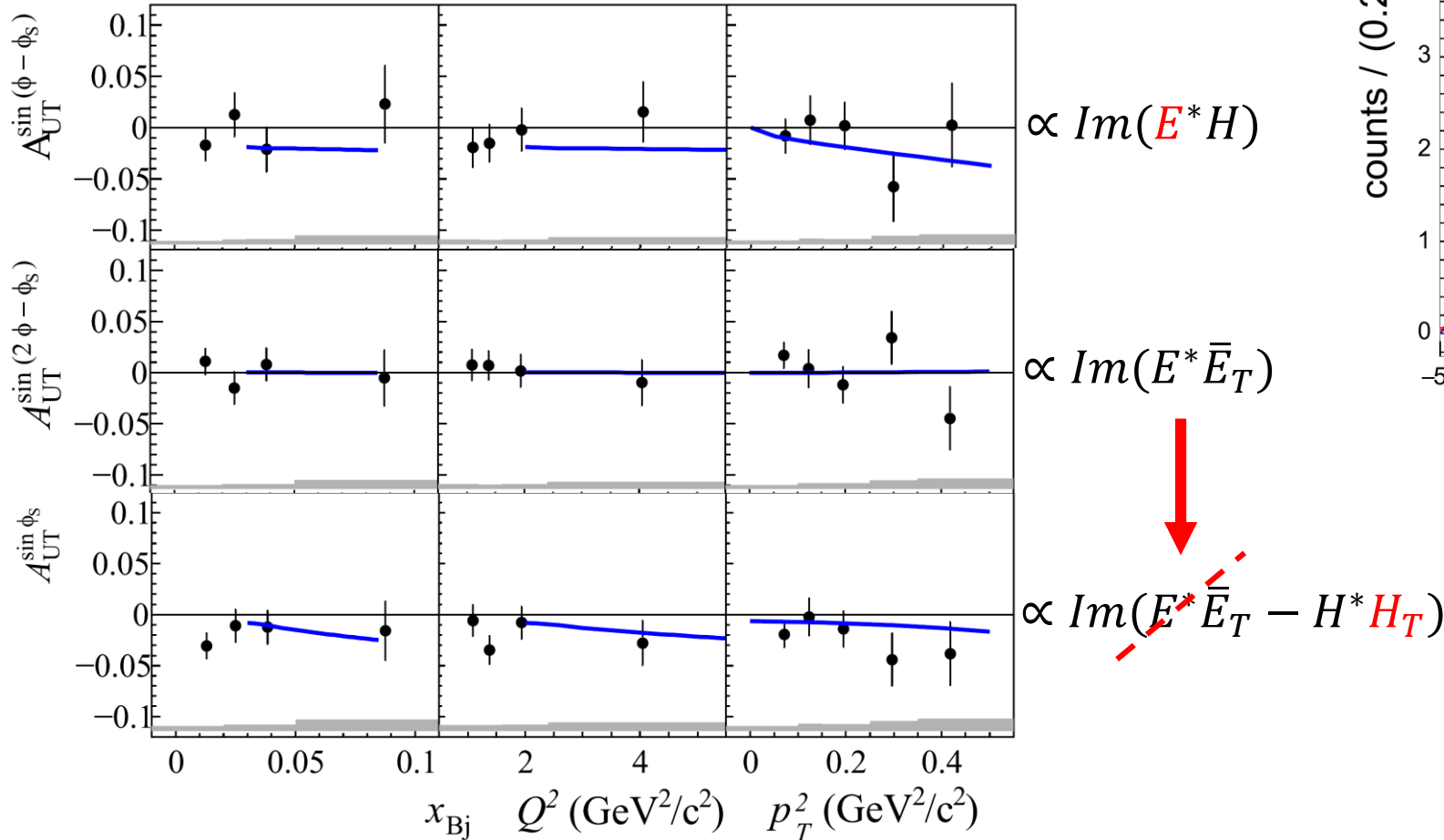
COMPASS, **NPB 865** (2012) 1-20, **PLB 731** (2014) 19



$$\rho^0 \rightarrow \pi^+ \pi^-$$

$$E_{\rho^0} = \frac{1}{\sqrt{2}} \left(\frac{2}{3} E^u + \frac{1}{3} E^d + \frac{3}{4} E^g / x \right)$$

➤ Exclusivity ensured by “missing mass technique”

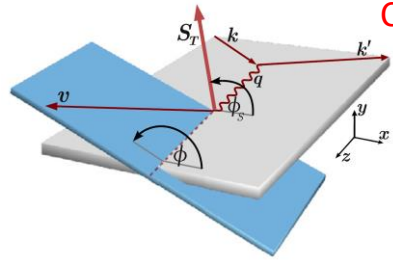


➤ Sensibility to E and H_T

2007 & 2010 HEMP with Transversely Polarized Target

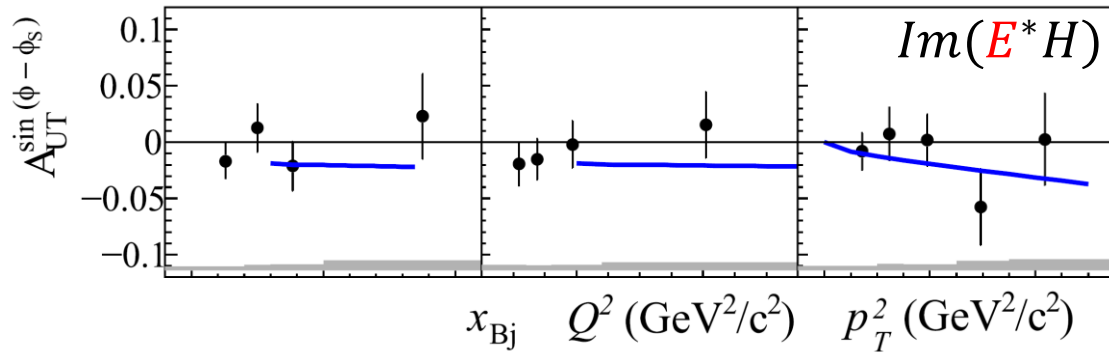


COMPASS, NPB 865 (2012) 1-20, PLB 731 (2014) 19



$$\rho^0 \rightarrow \pi^+ \pi^-$$

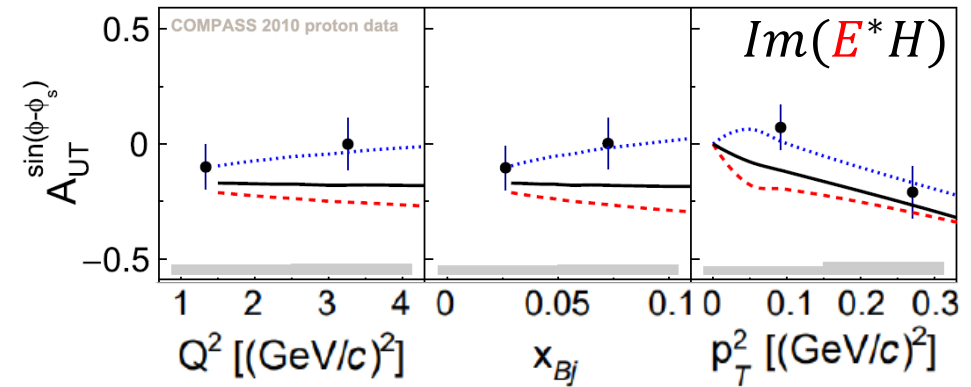
$$E_{\rho^0} = \frac{1}{\sqrt{2}} \left(\frac{2}{3} E^u + \frac{1}{3} E^d + \frac{3}{4} E^g / x \right)$$



COMPASS, NPB 865 (2012) 1-20, PLB 731 (2014) 19

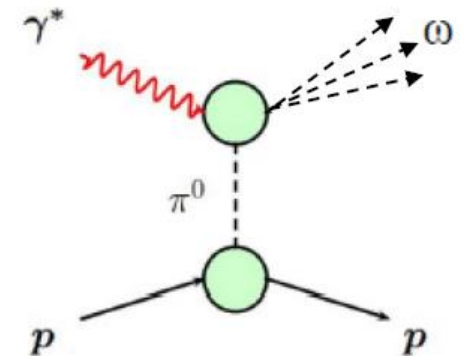
$$\omega \rightarrow \pi^+ \pi^- \pi^0$$

$$E_{\omega} = \frac{1}{\sqrt{2}} \left(\frac{2}{3} E^u - \frac{1}{3} E^d + \frac{1}{4} E^g / x \right)$$

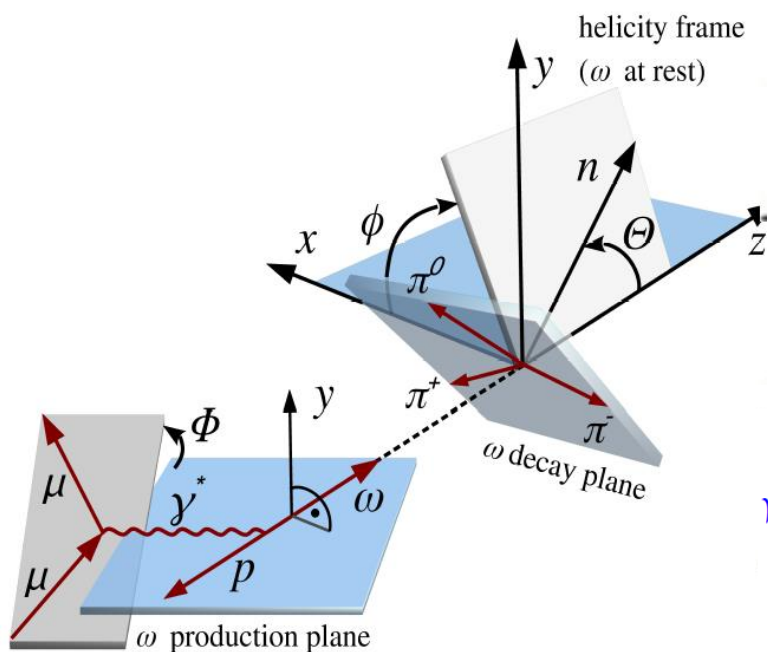


— No pion pole
 ⋯ Negative $\pi\omega$ form factor
- - - Positive $\pi\omega$ form factor

- E^u and E^d are of opposite sign $\rightarrow \omega$ is more promising for GPD study
- Nevertheless, obscured by the inherent pion pole contribution



Exclusive ω Production on Unpolarized Proton



Experimental angular distributions

$$\mathcal{W}^{U+L}(\Phi, \phi, \cos \Theta) = \mathcal{W}^U(\Phi, \phi, \cos \Theta) + P_b \mathcal{W}^L(\Phi, \phi, \cos \Theta)$$

15 unpolarized SDMEs in \mathcal{W}^U and 8 polarized in \mathcal{W}^L

$$\begin{aligned} \mathcal{W}^U(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[\frac{1}{2}(1 - r_{00}^{04}) + \frac{1}{2}(3r_{00}^{04} - 1) \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^{04}\} \sin 2\Theta \cos \phi - r_{1-1}^{04} \sin^2 \Theta \cos 2\phi \right. \\ & - \epsilon \cos 2\Phi \left(r_{11}^1 \sin^2 \Theta + r_{00}^1 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^1\} \sin 2\Theta \cos \phi - r_{1-1}^1 \sin^2 \Theta \cos 2\phi \right) \\ & - \epsilon \sin 2\Phi \left(\sqrt{2}\text{Im}\{r_{10}^2\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^2\} \sin^2 \Theta \sin 2\phi \right) \\ & + \sqrt{2\epsilon(1+\epsilon)} \cos \Phi \left(r_{11}^5 \sin^2 \Theta + r_{00}^5 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^5\} \sin 2\Theta \cos \phi - r_{1-1}^5 \sin^2 \Theta \cos 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \sin \Phi \left(\sqrt{2}\text{Im}\{r_{10}^6\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^6\} \sin^2 \Theta \sin 2\phi \right) \right], \\ \mathcal{W}^L(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[\sqrt{1-\epsilon^2} \left(\sqrt{2}\text{Im}\{r_{10}^3\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^3\} \sin^2 \Theta \sin 2\phi \right) \right. \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos \Phi \left(\sqrt{2}\text{Im}\{r_{10}^7\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^7\} \sin^2 \Theta \sin 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \sin \Phi \left(r_{11}^8 \sin^2 \Theta + r_{00}^8 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^8\} \sin 2\Theta \cos \phi - r_{1-1}^8 \sin^2 \Theta \cos 2\phi \right) \right] \end{aligned}$$

➤ $\epsilon \rightarrow 1$, small \mathcal{W}^L

2012 Exclusive ω Prod. on Unpolarized Proton



SCHC ($\lambda_\gamma = \lambda_V$)

(S-Channel Helicity Conservation)

SCHC implies:

• $r_{1-1}^1 + \text{Im} r_{1-1}^2 = 0$

= $-0.010 \pm 0.032 \pm 0.047$ OK

• $\text{Re} r_{10}^5 + \text{Im} r_{10}^6 = 0$

= $0.014 \pm 0.011 \pm 0.013$ OK

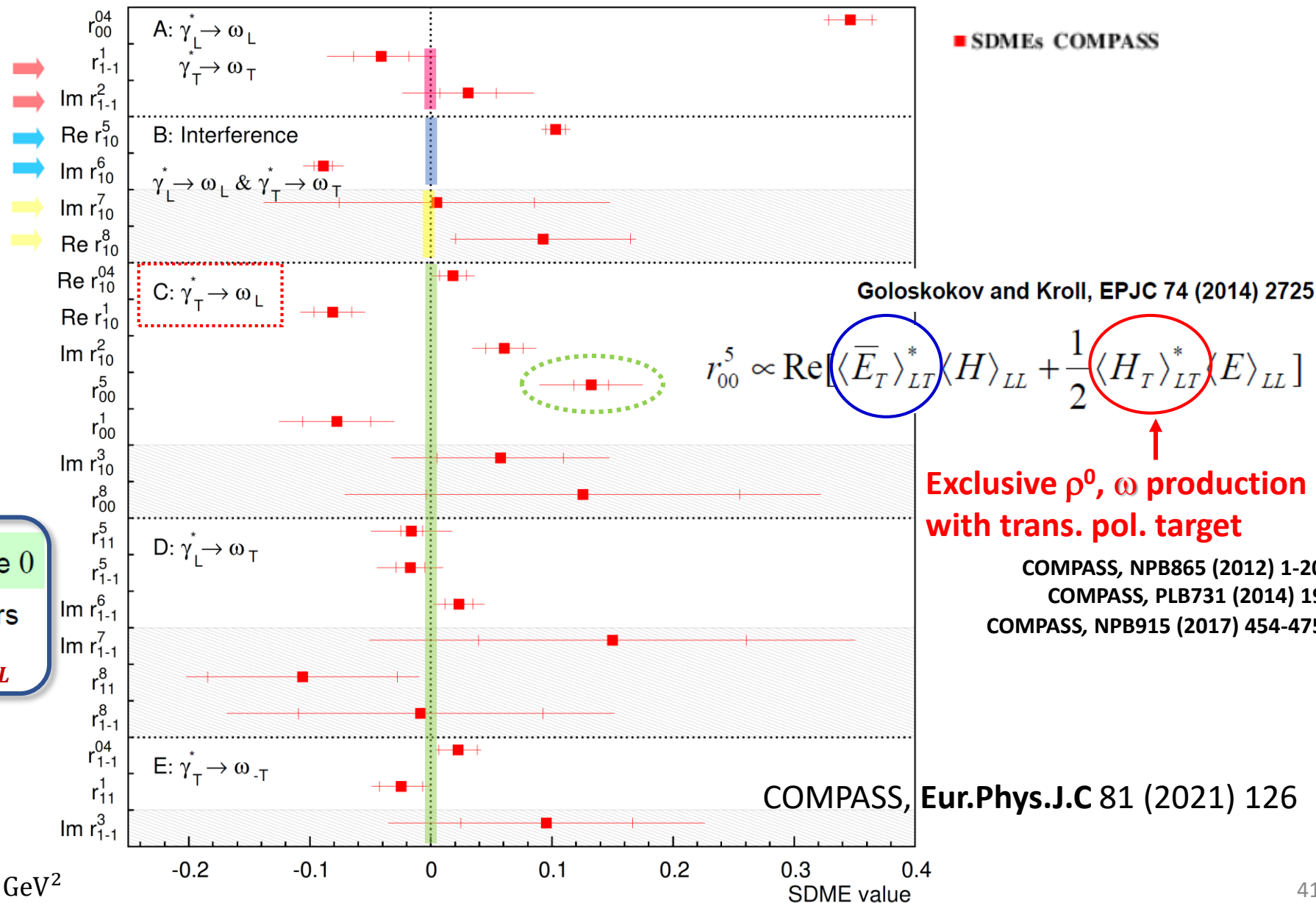
• $\text{Im} r_{10}^7 - \text{Re} r_{10}^8 = 0$

= $-0.088 \pm 0.110 \pm 0.196$ OK

• all elements of classes C, D, E should be 0

for $\gamma_L^* \rightarrow \omega_T$ and $\gamma_T^* \rightarrow \omega_T$ OK within errors

NOT OBSERVED for transitions $\gamma_T^* \rightarrow \omega_L$



$\langle Q^2 \rangle = 2.1 \text{ GeV}^2, \langle W \rangle = 7.6 \text{ GeV}, \langle P_T^2 \rangle = 0.16 \text{ GeV}^2$

2012 Exclusive ρ^0 Prod. on Unpolarized Proton



SCHC ($\lambda_\gamma = \lambda_V$)

(S-Channel Helicity Conservation)

SCHC implies:

• $r_{1-1}^1 + \text{Im} r_{1-1}^2 = 0$

= -0.000 ± 0.006 OK

• $\text{Re} r_{10}^5 + \text{Im} r_{10}^6 = 0$

= -0.011 ± 0.003 Violation

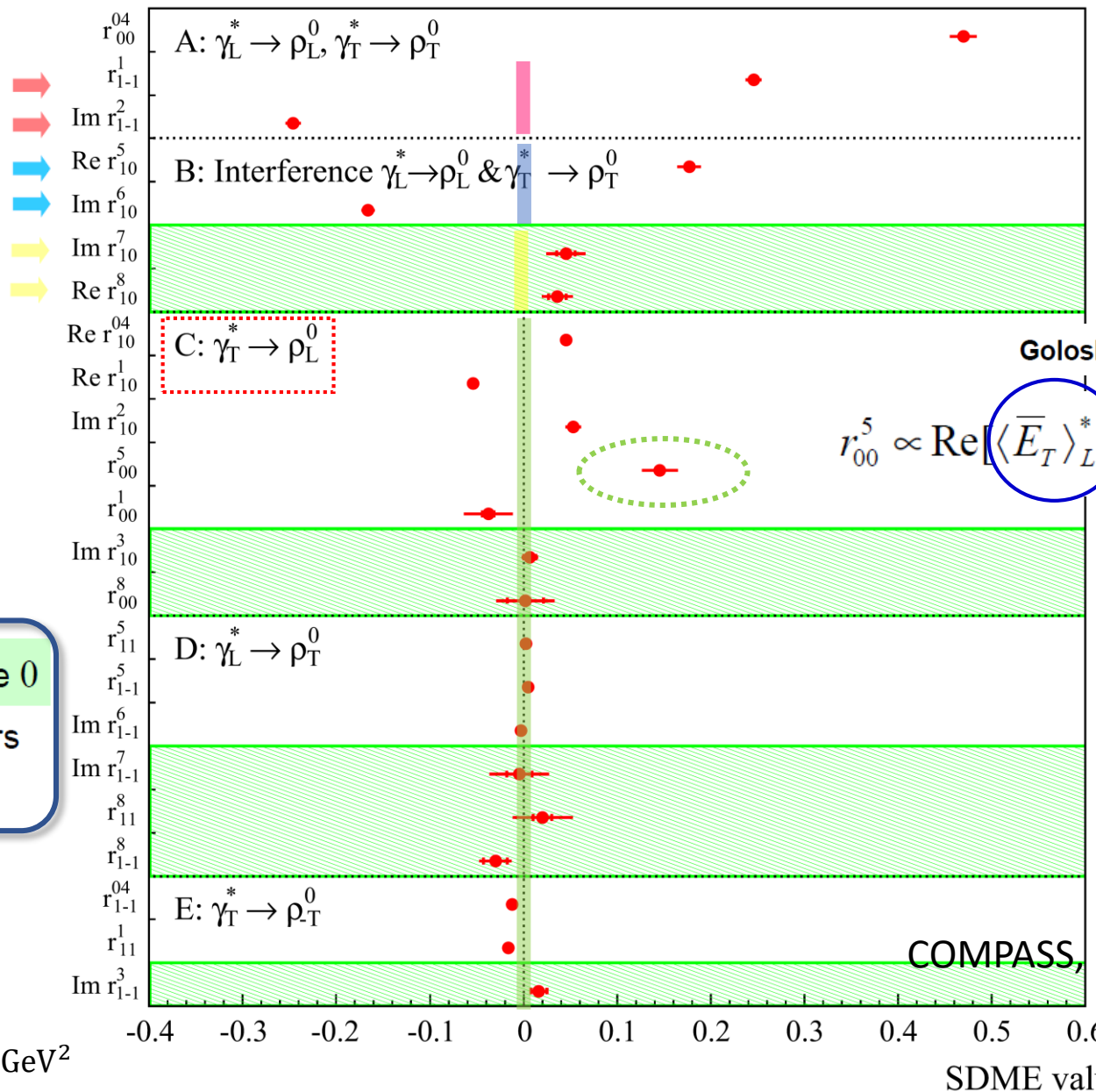
• $\text{Im} r_{10}^7 - \text{Re} r_{10}^8 = 0$

= -0.009 ± 0.031 OK

• all elements of classes C, D, E should be 0

for $\gamma_L^* \rightarrow \omega_T$ and $\gamma_T^* \rightarrow \omega_T$ OK within errors

NOT OBSERVED for transitions $\gamma_T^* \rightarrow \rho_L^0$



■ SDMEs COMPASS

Goloskokov and Kroll, EPJC 74 (2014) 2725

$$r_{00}^5 \propto \text{Re} \left[\langle \bar{E}_T \rangle_{LT}^* \langle H \rangle_{LL} + \frac{1}{2} \langle H_T \rangle_{LT}^* \langle E \rangle_{LL} \right]$$

First term dominates
→ Probes \bar{E}_T

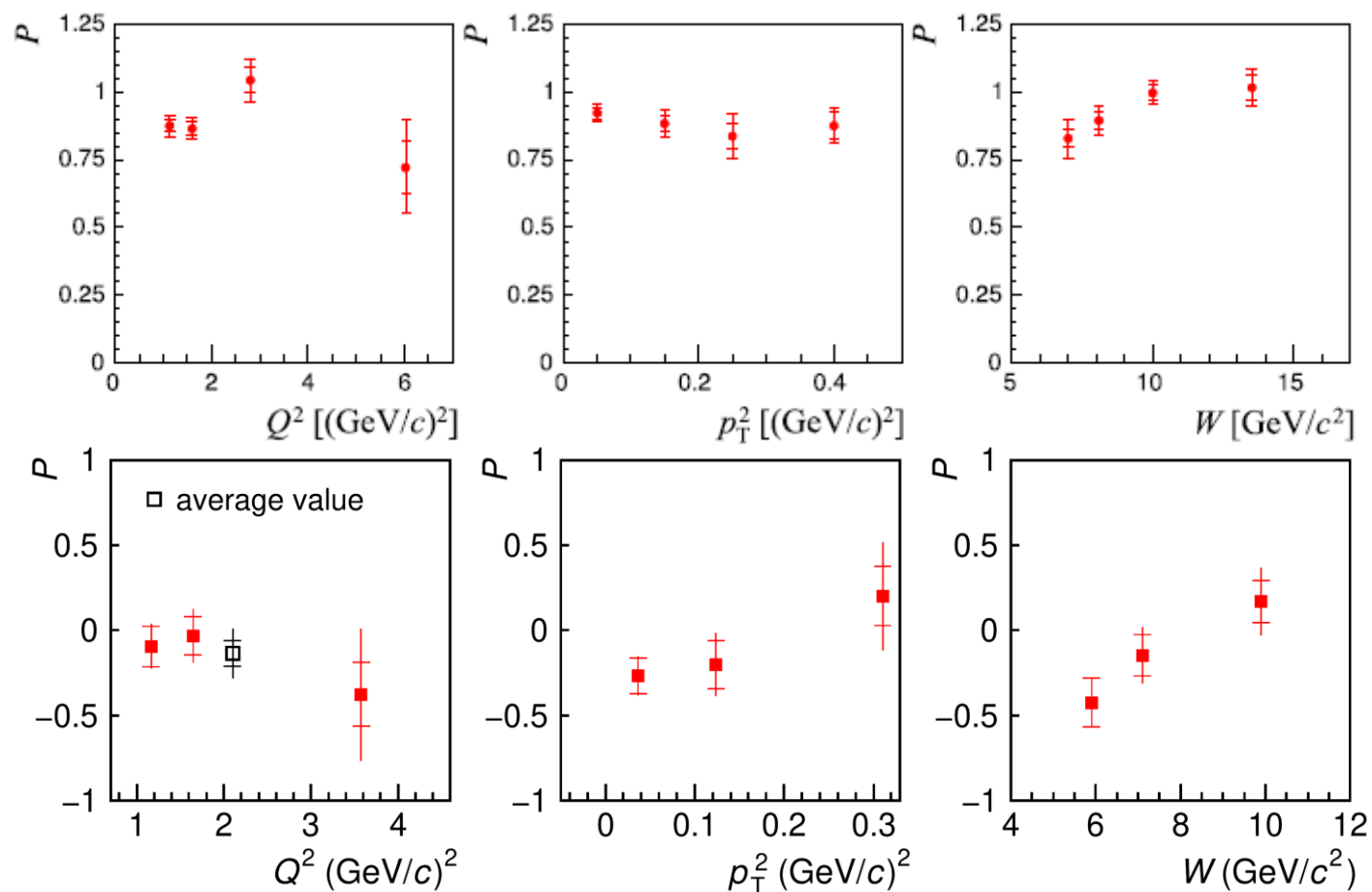
COMPASS, Eur.Phys.J.C 83 (2023) 924

2012 NPE-to-UPE Asymmetry

$$P = \frac{2r_{1-1}^1}{1 - r_{00}^{04} - 2r_{1-1}^{04}} \approx \frac{d\sigma_T^N(\gamma_T^* \rightarrow V_T) - d\sigma_T^U(\gamma_T^* \rightarrow V_T)}{d\sigma_T^N(\gamma_T^* \rightarrow V_T) + d\sigma_T^U(\gamma_T^* \rightarrow V_T)}$$

NPE-to-UPE asymmetry of cross sections for transitions $\gamma_T^* \rightarrow V_T$

- NPE: Natural Parity Exchange
- UPE: Unnatural Parity Exchange



ρ^0 COMPASS, Eur.Phys.J.C 83 (2023) 924

- NPE Dominance
- NPE \rightarrow GPDs E, H

ω COMPASS, Eur.Phys.J.C 81 (2021) 126

- NPE \approx UPE on average
- UPE Dominance at small W and p_T^2
- UPE \rightarrow GPDs \tilde{E}, \tilde{H}
+ Pion pole (dominant)

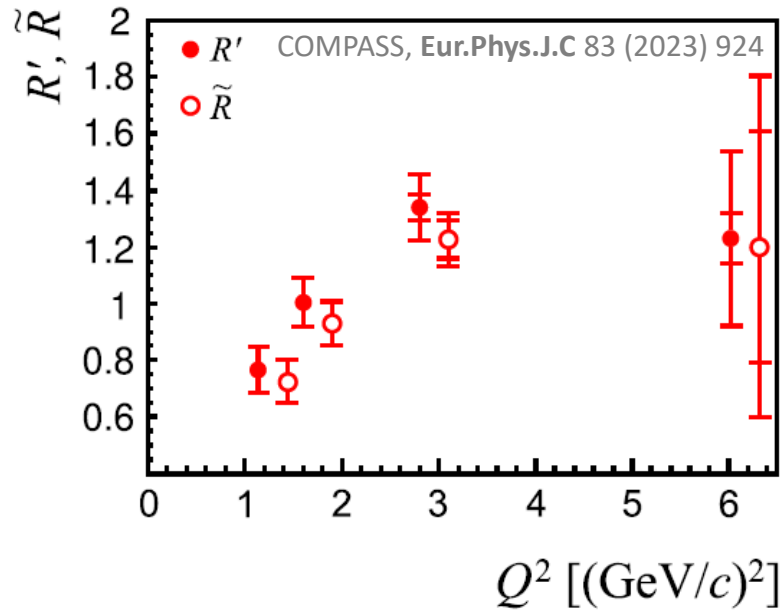
2012 $R = \sigma_L/\sigma_T$ for Exclusive ρ^0 Production



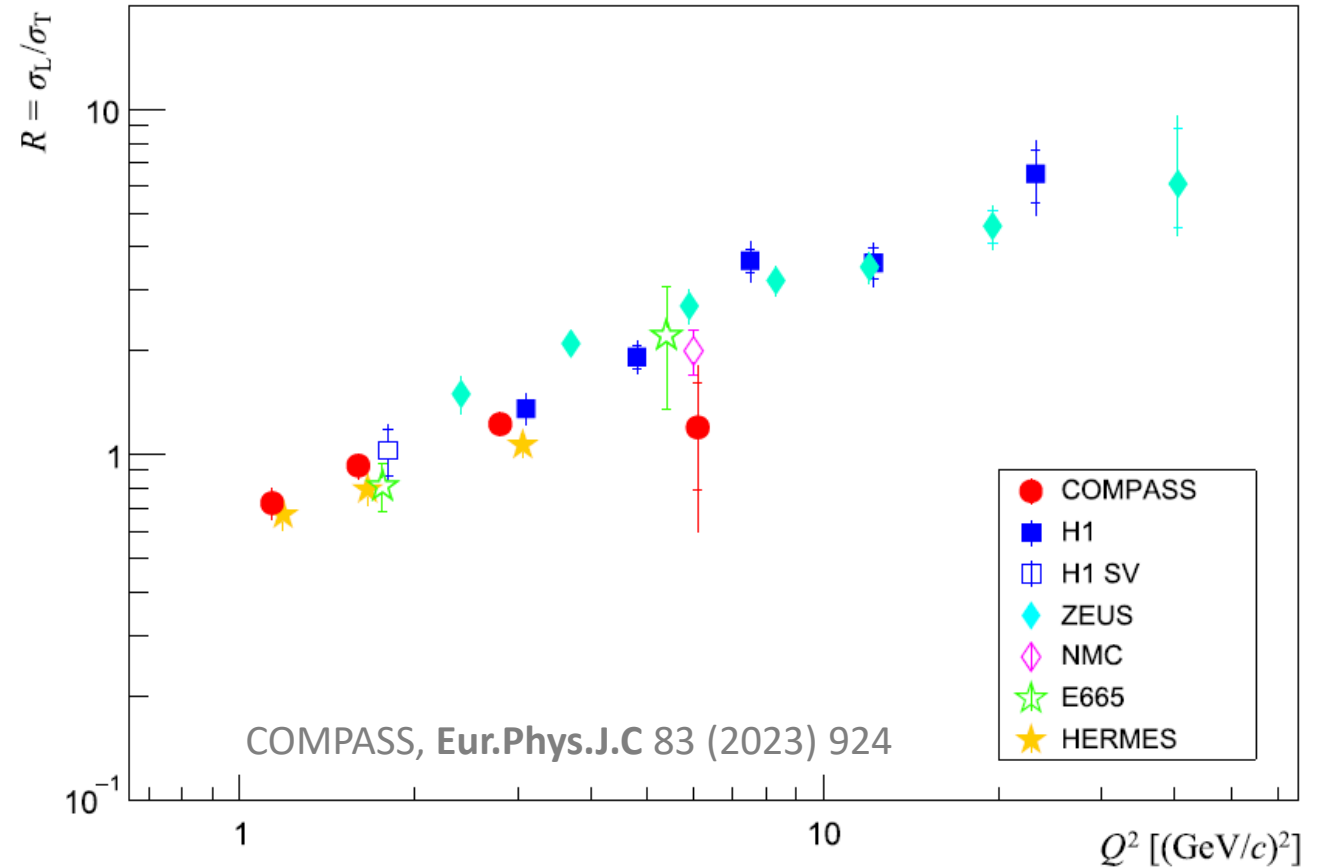
➤ Longitudinal-to-transverse γ^* cross section ratio: $R = \frac{\sigma_L(\gamma_L^* \rightarrow V)}{\sigma_T(\gamma_T^* \rightarrow V)}$

• Commonly used “effective” ratio ($R' = R$ only if SCHC): $R' = \frac{1}{\epsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}}$

• Use of \tilde{R} , which takes SCHC violation into consideration, is preferred.



Results of all experiments with $Q^2 > 1$ (GeV/c)²



➤ Leading-order pQCD prediction: $Q^2/M_\rho^2 \rightarrow$ deviation due to effect of QCD evolution and q_T

DVCS cross sections with polarized μ^+ and μ^-

- Beam charge-spin sum $\rightarrow \text{Im}\mathcal{H}(\xi,t) \rightarrow$ Transverse extension of partons as a function of x_{Bj}
- Beam charge-spin difference $\rightarrow \text{Re}\mathcal{H}(\xi,t) \rightarrow$ D-term, pressure distribution

HEMP of π^0 , ρ , ω , ϕ , J/ψ

- Cross section of π^0 , SDME of ρ & $\omega \rightarrow$ Transversity GPDs & Flavor Decomposition
- ϕ , $J/\psi \rightarrow$ underway

➤ **COMPASS has entered its analysis phase, expect more results soon!**

