

Cold nuclear matter effects from charmonium measurements

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LABORATÓRIO DE INSTRUMENTAÇÃO E FÍSICA EXPERIMENTAL DE PARTÍCULAS



Motivation

Ultra-relativistic heavy ion collisions: Quark Gluon Plasma (QGP)

One of the signatures for the QGP:

 J/ψ suppression in central collisions (Matsui & Satz, PLB 178(1986)416)

But: J/ψ suppression also occurs in pA collisions, which provide a needed baseline to interpret the suppression in QGP.

- J/ψ production in light systems allows to measure cold nuclear matter effects:
- parton energy loss
- nuclear absorption
- Cronin effect (p_{τ} broadening)
- nPDF

A useful observable to study this is the **Nuclear modification factor:**



baryon versus QGP (from P. Preuss, Berkely Lab)

$$R_{hA} = \frac{dN_{hA}^{\mathrm{J}/\psi}}{\langle N_{coll} \rangle \ dN_{hh}^{\mathrm{J}/\psi}}$$

If no nuclear effects: $R_{hA} = 1$

Nuclear shadowing and parton energy loss

Different phenomena observed. At low x, driven by partons multiple scattering.

Try to encode it all in process-independent **nPDF**s

Partons may also lose energy via soft gluon emissions when crossing the cold nuclear matter

Different hard processes allow to study the **energy loss effect**:

- Drell-Yan \rightarrow initial state radiation
- J/ψ production \rightarrow initial and final state radiation, interference





Some J/ ψ measurements from past experiments



Nuclear modification factor in COMPASS

COMPASS: 190 GeV π^- beam on fixed target (W or Al). Sqrt(s)=18.9 GeV

In COMPASS we measure inclusive J/ψ production:

 $\pi^{-} \ A \ \rightarrow \ J/\psi \ X \ \rightarrow \ \mu^{+}\mu^{-} \ X$

 $R_{\pi A}$ (x_F, p_T)|_{W/AI} is defined as the ratio of J/ψ production cross sections (per nucleon) between W and AI targets, in a given (x_F, p_T) bin.

In the center-of-mass of the hadrons collision (*),

$$x_F = \frac{2 \ p_L^*}{\sqrt{s}}$$

 p_L and p_T : longitudinal and transverse momentum of the dimuon. 5

COMPASS experiment @ CERN



Measurements done with the "**Drell-Yan set-up**" In 2015 and 2018.

Two dimuon triggers based on hodoscope pairs:

- "2 muons in LAS" (LL);
- "1 muon in LAS and 1 muon in SAS" (LO).

• NH3 polarized target

NH3

- Aluminium target
- Tungsten target

Data selection

- All data recorded in 2018
- Muon pairs of opposite charge
- Dimuon trigger fired (LL or LO)
- 1.5 < M < 8.5 GeV/c²
- 0< x_c < 0.9
- p₋ < 4 GeV/c
- Vertices inside one of targets:
 - → -73.5 < Zvtx < -66.5 cm (Al)</p>
 - \rightarrow -30 < Zvtx < -20 cm (W)
 - ~80K J/ ψ in Al and ~600K J/ ψ in W



Full Monte Carlo simulation of all relevant physics processes, and propagation in the spectrometer (pythia 8 + GEANT4):

- J/ψ
- ψ(2S)
- Drell-Yan
- Open charm

 and combinatorial background estimated from real data like-sign muon pairs.

Analysis done in bins (x_{r}, p_{T}) : dimuon mass



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Analysis done in bins (x_{r}, p_{T}) : dimuon mass



J/ψ signal extraction using "cocktail fit"

Reconstructed MC physics components and combinatorial background from real data are used to fit the dimuon mass spectra, in each of the (x_F, p_T) bins, separately per trigger and per target.



Example: W target, LL trigger $0.1 < x_{F} < 0.2$, in bins of p_{T}

Cocktail fits done in the mass range $2.1 < M < 8.5 \text{ GeV/c}^2$

Monte Carlo to real data comparison



Example: W target, LL trigger, μ^+

Good description, once all MC physics components (in the proportions given by cocktail fit) and combinatorial background from real data (RD) are taken into account.

J/ψ acceptance

The acceptance – including smearing effects and detector and trigger efficiencies – is obtained in $(x_{_F}, p_{_T})$ bins, separately per trigger and per target



LO trigger

$R_{\pi A}$ as a function of x_F , in bins of p_T



$R_{\pi A}$ as a function of p_{τ} , in bins of x_{F}





Error bars show the statistical uncertainty. Systematic uncertainty estimated: <10%

Comparing COMPASS result



To a model of energy loss, with transport coefficient q_0 by F. Arléo and S. Peigné, JHEP 03 (2013) 122. $\hat{q}_0 = 0.075 \text{ GeV}^2/\text{fm}$



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Comparing COMPASS result



 M_{\perp} : transverse mass



Still to come from COMPASS

- ✓ These nuclear modification factor results from COMPASS are very recent and still preliminary.
- Comparison to models of coherent parton energy loss is ongoing.
- ✓ An analysis of the J/ ψ <p₁²> from W and Al targets is ongoing, which should allow comparison to models of nuclear p₁-broadening.
- \checkmark J/ ψ absolute cross section in AI and W targets should follow soon.
- \checkmark An analysis on the J/ ψ polarization from W and NH3 targets is also ongoing.

The nuclear modification factor from the Drell-Yan process, and Drell-Yan cross-section should also be released soon.

ADDBER Apparatus for Meson and Baryon

Experimental Research

What else can we learn from quarkonium measurements?

Pion-induced J/ ψ production: access to the **pion structure...**

... provided we understand and manage to distinguish experimentally its production mechanisms (namely $q\bar{q}$ annihilation and gg fusion)



AMBER/NA66 experiment

The AMBER/NA66 experiment had phase-I measurements approved at CERN in December 2020 (\rightarrow proposal)

Several topics related to the emergence of hadron mass:

- Hadron charge radii
- Hadron spectroscopy
- Hadron structure

 π^{\pm} and K[±] beam particles

190 GeV/c

Setup similar to COMPASS with added/upgraded detectors and target region

Beam



Pion structure studies at AMBER





J/ Ψ : access to gluon content in the pion

- Large statistics on J/ψ production at dimuon channel: $\pi^{\pm} C \rightarrow J/\psi X \rightarrow \mu^{+}\mu^{-} X$
- Inclusive: due to the hadron absorber, we cannot distinguish prompt production from the rest
- Expected significant feed-down: $\psi(2S)$, χ_{c1} , χ_{c2}
- In the low-p_r regime
- Dominant contribution from $2 \rightarrow 1$ processes
- Use the polarization, $x_{_{\rm F}}$ and $p_{_{\rm T}}$ dependences to distinguish production mechanisms...



AMBER phase-II: RF-separated kaon beams for access to kaon structure

The RF-separation technique allows to improve the kaon purity in a hadron beam. But:

– kaon intensity limited to 5x10⁵/second (as opposed to 2x10⁷ assumed in AMBER LoI)



- Drell-Yan process: not enough statistics
- **Charmonium measurements**: possible! The lower Intensity might allow for open spectrometer

AMBER phase-II: RF-separated kaon beams J/ψ and access to valence and glue

- Beam intensity 10⁵ kaons/second
- ~ 10 000 J/ ψ events for each beam
- Model-dependent access to the gluon distribution in kaons
- J/ψ production cross section (LO):



• Cross-section difference isolates val-val term:

 $\sigma(K^{-}) - \sigma(K^{+}) \propto \overline{U}_{v}^{K} U_{v}^{p}$

CEM, Int.J.Mod.Phys. A 10 (1995) 3043 JAM18 "pion" PDFs, PRL 121, 152001 (2018)



Higher charmonium states at AMBER

With RF-separated kaon beams:

- much lower beam intensity
 - open spectrometer measurements become possible
 - Much better mass resolution, allowing to distinguish J/ψ and ψ(2S)

ψ(2S):

Cleaner than J/ ψ , since no feed-down from χ_c states.

Its polarization is a stronger discriminant than J/ψ between qq and gg mechanisms

But: relatively low statistics

E771 Collaboration / Physics Letters B 374 (1996) 271–276



Summary



First measurements of the J/ ψ nuclear modification factor from COMPASS.

 $R_{\pi A}(x_{F}, p_{T})|_{W/AI}$ is seen to decrease for higher x_{F} and lower p_{T} .

The measured J/ψ suppression is in qualitative agreement with past experiments, and with the coherent parton energy loss assumption.

Many more results expected soon.

The AMBER new experiment at CERN focuses on the studies of meson structure: pion and kaon.

A new generation of charmonium measurements at intermediate energies are proposed:

- understand production mechanisms;
- access to valence and gluon content in the pion and kaon. \neg

starting!

SPARES

The double differential J/ψ cross section:

$$\frac{\mathrm{d}^2 \sigma^{\pi^- \mathrm{A}}}{\mathrm{d} \mathrm{x}_\mathrm{F} \mathrm{d} \mathrm{p}_\perp} J/\psi = \frac{\mathrm{N}_\mathrm{events}^{\mathrm{J}/\psi}(\mathrm{x}_\mathrm{F},\mathrm{p}_\perp)}{\epsilon^{\mathcal{A}}.\mathrm{BR}.\Delta\mathrm{x}_\mathrm{F}.\Delta\mathrm{p}_\perp.\mathcal{L}}$$

The integrated luminosity:

$$\mathcal{L} = \Phi^0 \; \alpha^A \; L^A_{eff} \; \rho^A \; \mathcal{N}_A \; / M^A$$

The effective target length, taking into account the beam attenuation inside the target:

$$L_{\text{eff}}^{i} = \frac{\lambda_{\text{int}}}{\rho} \left[1 - \exp\left(\frac{-\rho \ L}{\lambda_{\text{int}}}\right) \right]$$

 ϵ^{A} : acceptance for target A BR: branching ratio J/ $\psi \rightarrow \mu\mu$ \pounds : integrated luminosity

Φ^0 : integrated beam flux

- α^A : beam attenuation at entry of target A
- L^A_{eff}: effective target length
- ρ^A : density of target A
- $\mathcal{N}_{\!\scriptscriptstyle A}$: Avogadro number
- M^A: molar mass of target A

$$R_{\pi^{-}A}^{J/\psi}(W/AI) = \frac{N_{W}^{J/\psi}(x_{F}, p_{\perp})}{\epsilon_{W}.\alpha^{W}.L_{eff}^{W}.\rho^{W}} / \frac{N_{Al}^{J/\psi}(x_{F}, p_{\perp})}{\epsilon^{AI}.\alpha^{AI}.L_{eff}^{AI}.\rho^{AI}}$$

J/ψ production and model interpretation on the mechanisms of production



CEM model, Int.J.Mod.Phys. A 10 (1995) 3043

Using global fit extracted pion and nucleon PDFs (GRS and CT10)



Daniele Binosi, at this workshop

Using pion structure predictions from 29 the continuum approach