Investigating cold nuclear matter effects in charmonia at the fixed-target COMPASS experiment

Anisa Khatun For the COMPASS Collaboration 31<sup>st</sup> August, CERN

#### International Workshop on Hadron Structure and Spectroscopy







# Introduction

- Heavy quarkonia suppression is one of the most distinctive signatures of QGP in heavy-ion collisions.
- Suppression in hh and hA collisions Cold Nuclear Matter (CNM) effects.
- Disentangle the CNM effects from QGP to interpret AA collisions.
- Quarkonia and Drell-Yan cross-sections crucial tool to study CNM effects.

### CNM effects ?

- Initial state effects:
  - Nuclear modification of the PDFs Anti-shadowing, EMC effects are dominant at SPS energies.
  - Initial state energy loss.
- Final state effects:
  - Final state energy loss.
- These effects can be quantified by measuring nuclear modification factor ( $R_{\pi A}$ ).



### Parton energy loss effects

A high energy parton travelling in a medium can radiate gluons induced by the elastic scatterings with the constituents of the medium



Parton energy loss effects in different hard processes:

- Drell-Yan process:  $hA \rightarrow \ell^+ \ell^- + X$ 
  - Initial state radiation
- Hadron production:  $hA \rightarrow q/g(\rightarrow h') + X$ 
  - Initial state radiation
  - Final state radiation
  - Interference of both

Anisa Khatun, IWHSS-2022



CNM effects are studied at COMPASS via J/ $\psi$  production at  $\sqrt{s}$ =18.9 GeV.

### Observations from previous fixed target experiments

- J/ $\psi$  is more suppressed than DY as a function of  $x_{F}$  and  $p_{T}$
- Different CNM effects for J/ψ and DY [PRL 84 (2000) 3256]



# Observations from previous experiments

- No scaling as a function of x<sub>2</sub>. [Arleo, Naïm, Platchkov, JHEP01(2019)129]
- $J/\psi$  suppression depends on center of mass energy.



- Coherent energy loss regime explains alone E866 J/ $\psi$  data at  $\sqrt{s}$ =38.7 GeV. [Arleo, Peigne, JHEP03(2013)122]
- Energy loss model explains the strong suppression at large x<sub>F</sub> for DY. [Arleo, Naïm, Platchkov, JHEP01(2019)129]

# COMPASS experimental set up

#### $J/\psi = \mu^{+} + \mu^{-} (2 < M < 4.3 \text{ GeV/c}^{2})$ DY = $\mu^{+} + \mu^{-} (4.3 < M < 8.5 \text{ GeV/c}^{2})$



 $\pi^{-}$ 

NH<sub>3</sub><sup>1</sup>

-300

He

NH<sub>3</sub><sup>2</sup>

-100

-200

- Al 27 nucleons
- W 184 nucleons

### Anisa Khatun, IWHSS-2022

w

0

100 (cm)

### Main Observable

The double differential cross-section

$$rac{\mathrm{d}^2 \sigma^{\pi^- \mathrm{A}}}{\mathrm{d} \mathrm{x}_\mathrm{F} \mathrm{d} \mathrm{p}_\perp} J/\psi = rac{\mathrm{N}_\mathrm{events}^{\mathrm{J}/\psi}(\mathrm{x}_\mathrm{F},\mathrm{p}_\perp)}{\epsilon^{\mathcal{A}}.\mathrm{BR}.\Delta\mathrm{x}_\mathrm{F}.\Delta\mathrm{p}_\perp.\mathcal{L}}$$

With integrated luminosity,

$$\mathcal{L} = lpha^i \Phi^0 imes L^i_{ ext{eff}} imes 
ho^i imes rac{\mathcal{N}_A}{M^i}$$

With effective length, taking into account the beam attenuation inside the target

$$L_{\text{eff}}^{i} = \frac{\lambda_{\text{int}}}{\rho} \left[ 1 - \exp\left(\frac{-\rho L}{\lambda_{\text{int}}}\right) \right]$$

All the targets have the same initial beam flux measured by beam telescopes and the Avogadro's number is constant, therefore

The double ratio cross-section

$$\mathcal{R}_{\pi^-\mathcal{A}}^{J/\psi}(W/\mathcal{A}I) = rac{\mathrm{N}_{\mathrm{W}}^{J/\psi}(\mathrm{x}_{\mathrm{F}},\mathrm{p}_{\perp})}{\epsilon_{W}.lpha^W.L_{ ext{eff}}^W.
ho^W} / rac{\mathrm{N}_{\mathrm{Al}}^{J/\psi}(\mathrm{x}_{\mathrm{F}},\mathrm{p}_{\perp})}{\epsilon^{\mathcal{A}I}.lpha^{\mathcal{A}I}.L_{ ext{eff}}^{\mathcal{A}I}.
ho^{\mathcal{A}I}}$$

# **Analysis Ingredients**

- Kinematic variables: Exact definition (see back-up slide)
- Targets range:
  - $\circ$  W : -30 < Z<sub>vtx</sub> < -20 (cm)
  - AI : -73.5 <  $Z_{vtx}$  < -66.5 (cm)
- 0 < xF < 0.9
- 0 < pT < 4 GeV/c
- Migration correction:
   W ----> AI



### **Dimuon invariant mass**





Double differential analysis using data taken in 2018 per trigger basis.

## Comparison between Real Data and Monte Carlo

- MC production at COMPASS: (PYTHIA8 + GEANT4)
- Real Data (RD) → J/ψ + ψ(2S) + DY + OC + CB

#### "Cocktail Fit"

- CB: correlated muon pairs from RD.
- MC-RD comparison: The momenta and angle of single muon after adding all the components from MC and CB.
- Good agreement between MC-RD after including all the components.

Anisa Khatun, IWHSS-2022



W, LO trigger, *θ*μ-



#### Al, LL trigger, Pµ-



#### W, LO trigger, Pµ+



10

## Signal extraction



### Acceptance



The total acceptance -> Geometrical acceptance + detector & trigger efficiency.

No contamination correction applied for Aluminium. Acceptance goes up to 12% .

# Results: $p_T$ differential $R_{\pi-A}$ as a function of $x_F$



# Results: $x_F$ differential $R_{\pi-A}$ as a function of $p_T$



# Results: $\textbf{R}_{\pi\text{-}A}$ integrated over $\textbf{x}_{F}$ and $\textbf{p}_{T}$



- Combined for two trigger by taking average over common kinematic range.
- Suppression towards high  $x_{F}$  and low  $p_{T}$  observed similar to 2D results.

## Comparison with other fixed-target experiments



- Qualitative comparison with previous fixed target experiments E866, NA3 results.
- Data from E866, NA3, E537,NA60 are well described by the energy loss model [Arleo, Peigne, JHEP03(2013)122].
- Hint of energy loss effect for COMPASS data.

## Summary and Outlook

- Preliminary results of nuclear dependence of  $J/\psi$  data have been presented.
- Suppression towards large  $x_F$  and low  $p_T$ .
- Maximum suppression at the highest  $x_{F}$  and the lowest  $p_{T}$  interval ~40%
- Qualitative comparison with previous fixed target experiments shows similar trend.
- Suppression towards large  $x_{F}$ , indicating possible energy loss effect.

### Thank you for your attention!



#### Kinematic variable definition

x-Feynman definition used by previous fixed target experiments NA3 [Z. Phys. C20 (1983) 101]

$$\begin{aligned} x_F &= \frac{p_z}{p_z^{max}} = \frac{p_z}{\sqrt{s/2}} \\ x_1 &= 0.5 \times (\sqrt{x_F^2 + 4Q^2/s} + x_F) \\ x_2 &= 0.5 \times (\sqrt{x_F^2 + 4Q^2/s} - x_F) \end{aligned}$$

# **Physics motivation**

• The cold nuclear effects in hA collisions characterized by nuclear modification factor:

 $R_{hA} = \frac{1}{A} \frac{d\sigma_{hA}}{dx_F} \frac{d\sigma_{hp}}{dx_F}$ = (1 or no nuclear effects)

- Nuclear modification factor depends on nPDF  $f_j^{p/A} \neq f_j^p$
- nPDF depends on bjorken-*x*, distinguishes among Nuclear Shadowing, Anti-Shadowing, EMC effects.
- At COMPASS access to wider positive x-Feynman range, possible to study nPDF.
- The anti-shadowing (0.01 ≤ x ≤ 0.3) and EMC region (0.3 ≤ x ≤ 0.7) are covered by COMPASS.



### Parton energy loss regimes

• Landau Pomeranchuk Migdal or the LPM effect (small formation time  $t_f \leq L$ )

 $\langle \epsilon 
angle_{LPM} \propto lpha_{s} \hat{q} L^{2}$ 

• Drell-Yan process:  $hA \rightarrow \ell^+ \ell^- + X$ 

• Full coherent parton energy loss effect (large formation time  $t_f \gg L$ )

 $\langle\epsilon
angle_{coherent}\propto\sqrt{\hat{q}L}/M.E\gg\langle\epsilon
angle_{LPM}$ 

• Quarkonium production:  $hA 
ightarrow [Q\hat{Q}(g)]_8 + X$ 

**Transport coefficient :** The scattering properties of the medium, depends on  $x_F$  and  $p_{\perp}$  distribution

$$\hat{\boldsymbol{q}} = rac{4\pi^2 lpha_s C_R}{N_c^2 - 1} \rho imes x G(x, Q^2), \hat{\boldsymbol{q}} \equiv rac{\mu^2}{\lambda} = rac{\mathrm{d}\Delta \mathrm{p}_\perp^2}{\mathrm{dL}}$$

Single  $\hat{q}$  to study both energy loss effects and  $p_{\perp}$  broadening. These nuclear effects are worth investigating with COMPASS DY data.