

**XVII International Workshop  
on Hadron Structure  
and Spectroscopy**

**Trieste**

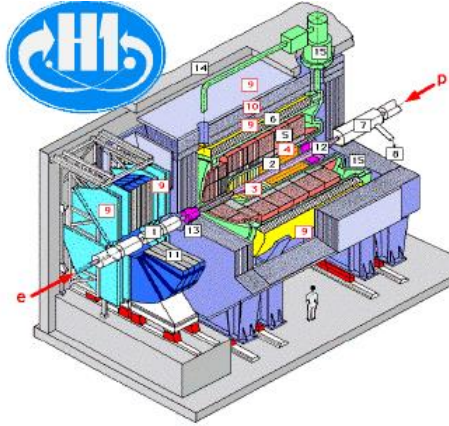
**16-18 November 2020**



**Experimental overview of exclusive reactions  
(related to GPDs)**

**Nicole d'Hose, CEA Université Paris-Saclay**

# Past & present for exclusive experiments at small $t$ : $\ell p \rightarrow \ell' p' \gamma$ or meson



## Collider mode e-p: forward fast proton

**HERA: H1 and ZEUS**

Polarised **27 GeV** e-/e+

Unpolarized **920 GeV** proton

~ Full event reconstruction (proton in Roman Pots)

## Fixed target mode: slow recoiling proton

**HERMES:** Polarised **27 GeV** e-/e+

Long, Trans polarised p, d target

*Missing mass technique*

*2006-07 with recoil detector*

**Jlab: Hall A, C, CLAS, CLAS12** High Luminosity Polar. **6 & 12 GeV** e-

Long, (Trans) polarised p, d target

*Missing mass technique (A,C) and complete detection (CLAS)*

**COMPASS @ CERN:** Polarised **160 GeV**  $\mu^+/\mu^-$

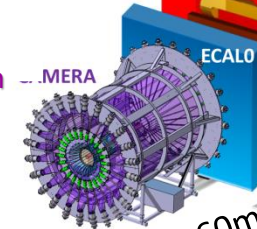
p target, (Trans) polarised target

*with recoil detection*

*Rejection of background: SIDIS, exclusive  $\pi^0$ /DVCS, dissociation of the proton*



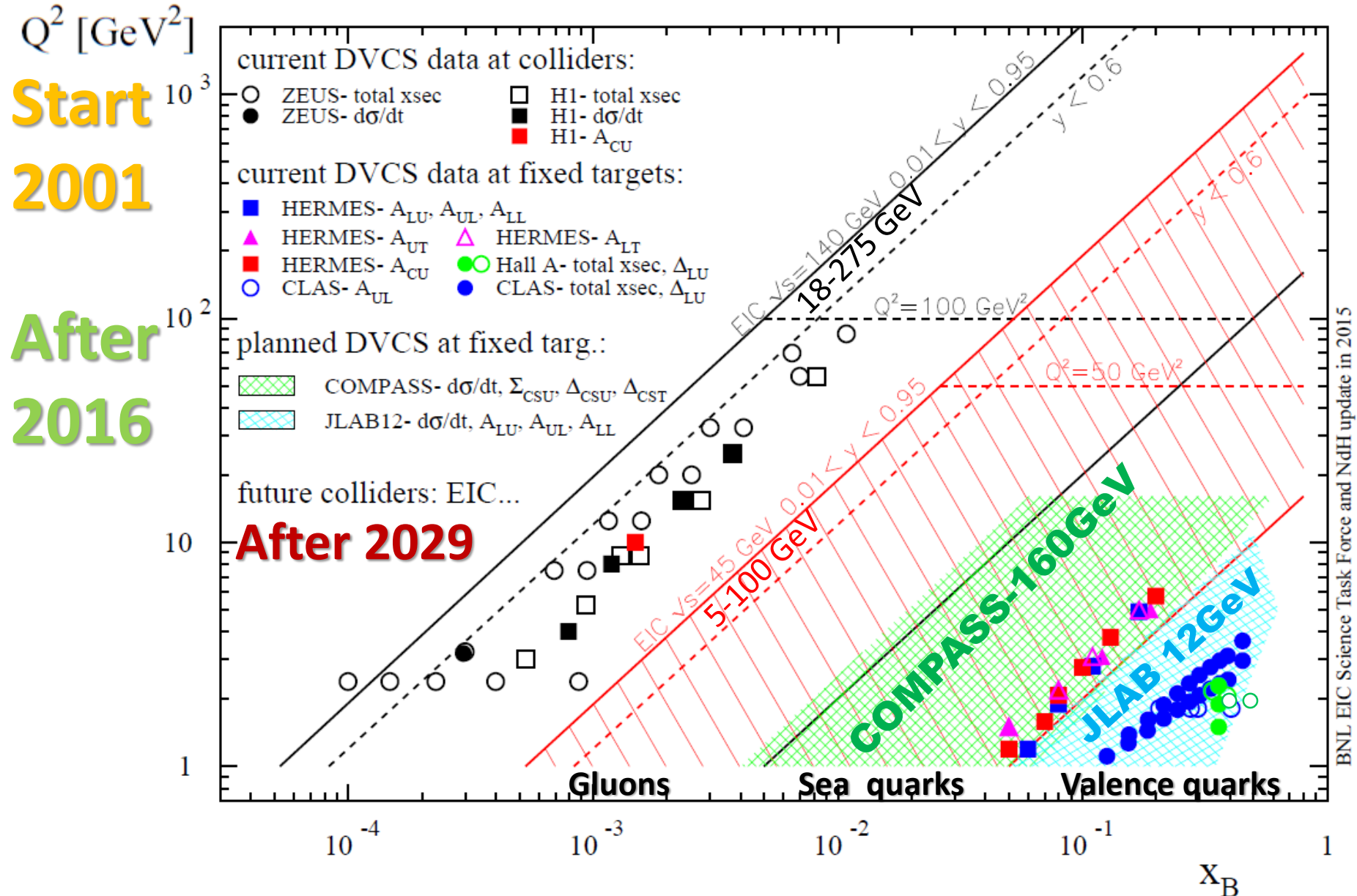
recoil proton detector CAMERA



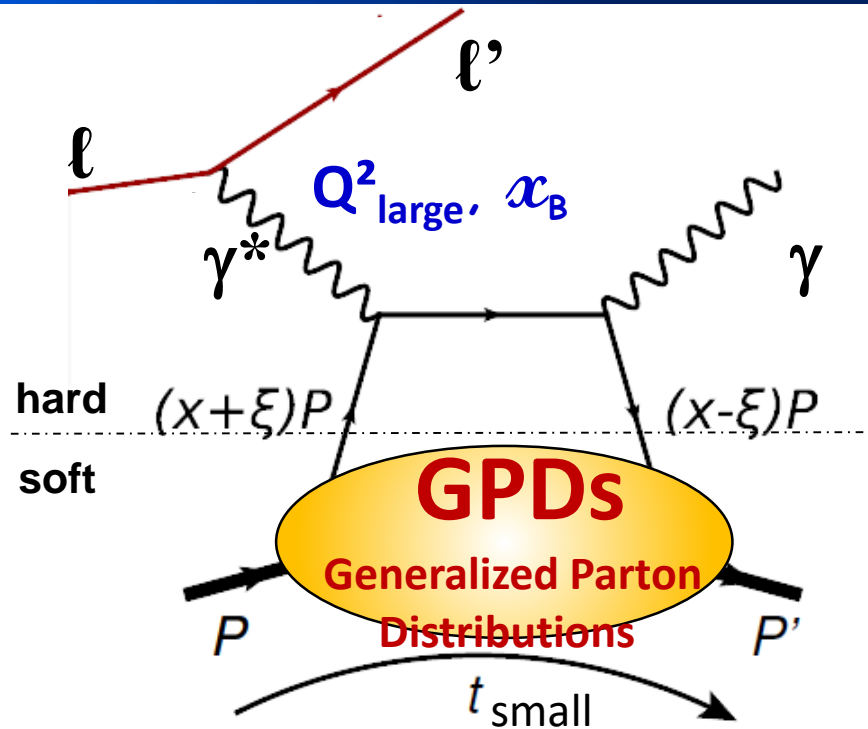
COMPASS

+ 60m long magnetic spectrometer of large acceptance with 3 EM Calos

# Past and future experiments for DVCS $\ell p \rightarrow \ell' p' \gamma$



# Deeply virtual Compton scattering (DVCS)



D. Mueller *et al*, Fortsch. Phys. 42 (1994)

X.D. Ji, PRL 78 (1997), PRD 55 (1997)

A. V. Radyushkin, PLB 385 (1996), PRD 56 (1997)

DVCS:  $l p \rightarrow l' p' \gamma$

the golden channel

because it interferes with  
the Bethe-Heitler process

also meson production

$l p \rightarrow l' p' \pi, \rho, \omega$  or  $\phi$  or  $J/\psi \dots$

The GPDs depend on the following variables:

$x$ : average long. momentum

$\xi$ : long. mom. difference

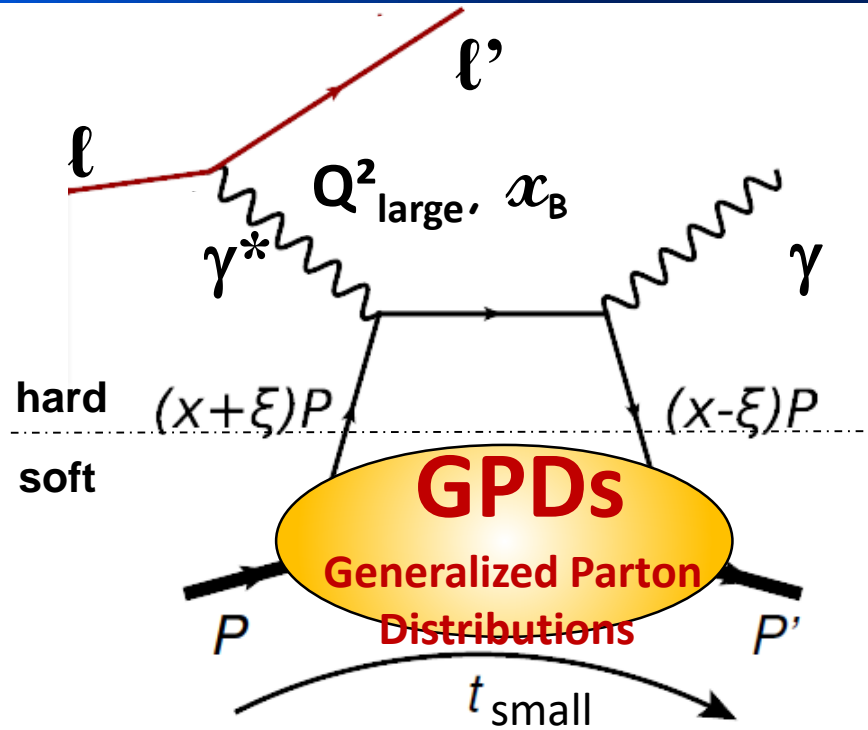
$t$ : four-momentum transfer  
related to  $b_{\perp}$  via Fourier transform

The variables measured in the experiment:

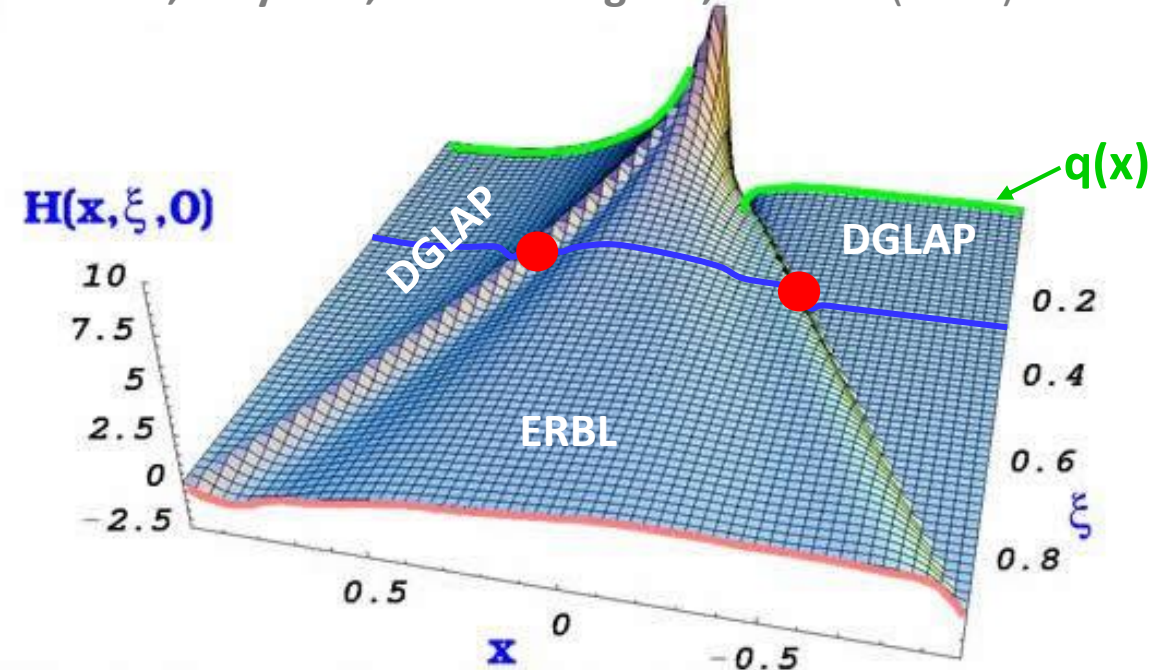
$E_{\ell}, Q^2, x_B \sim 2\xi / (1+\xi),$

$t$  (or  $\theta_{\gamma^* \gamma}$ ) and  $\phi$  ( $l l'$  plane /  $\gamma \gamma^*$  plane)

# Deeply virtual Compton scattering (DVCS)



Goeke, Polyakov, Vanderhaeghen, PPNP47 (2001)



The amplitude DVCS at LT & LO in  $\alpha_s$  (GPD  $\mathcal{H}$ ):

$$\mathcal{H} = \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi + i\epsilon} = \mathcal{P} \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi} - i\pi H(x = \pm \xi, \xi, t)$$

In an experiment we measure  
Compton Form Factor  $\mathcal{H}$

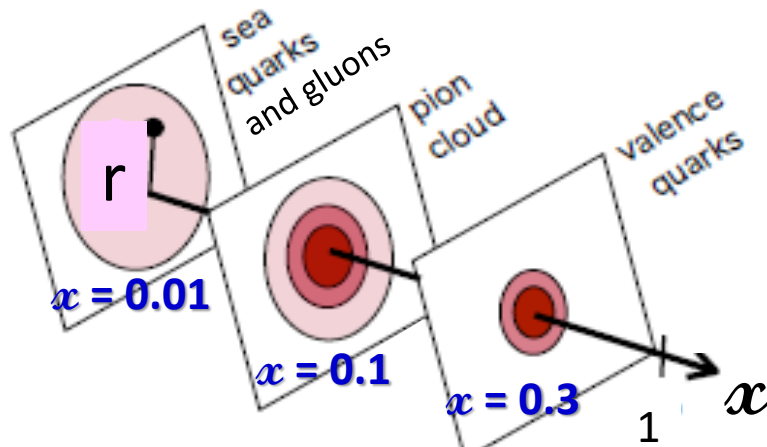
$$\text{Re}\mathcal{H}(\xi, t) = \pi^{-1} \int dx \frac{\text{Im}\mathcal{H}(x, t)}{x - \xi} + \Delta(t)$$

# Deeply virtual Compton scattering (DVCS)

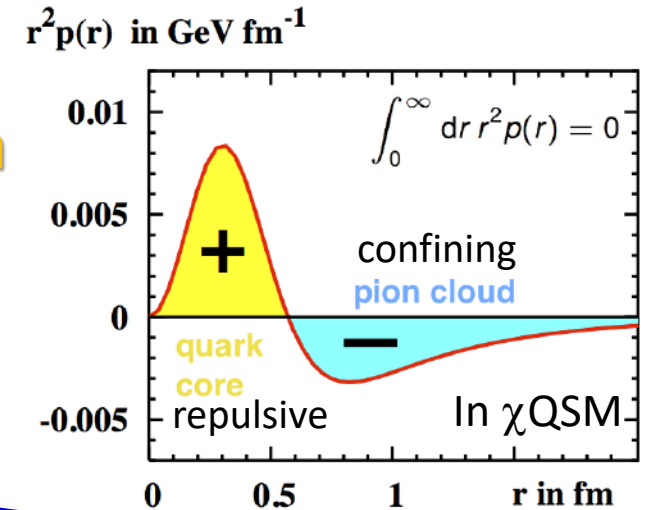
M. Burkardt, PRD66(2002)

M. Polyakov, P. Schweitzer, Int.J.Mod.Phys. A33 (2018)

## Mapping in the transverse plane



## Pressure Distribution



FT of  $H(x, \xi=0, t)$

The amplitude DVCS at LT & LO in  $\alpha_s$  (GPD  $\mathbf{H}$ ):

$$\mathcal{H} = \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi + i\epsilon} = \mathcal{P} \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi} - i \pi H(x = \pm \xi, \xi, t)$$

In an experiment we measure Compton Form Factor  $\mathcal{H}$

$$\text{Re}\mathcal{H}(\xi, t) = \pi^{-1} \int dx \frac{\text{Im}\mathcal{H}(x, t)}{x - \xi} + \Delta(t)$$

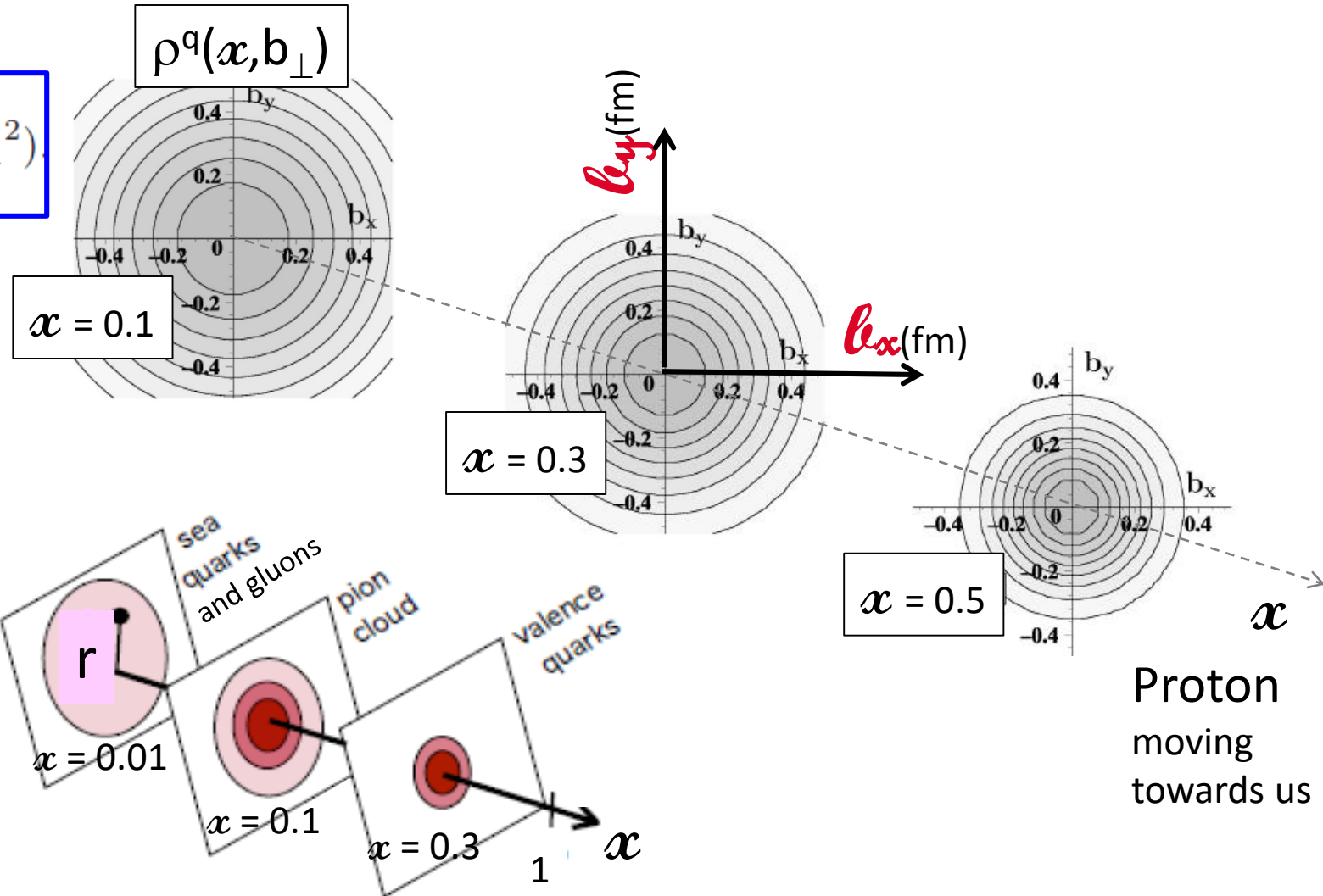
$d_1(t)$   
D-term

# GPDs and 3D imaging

M. Burkardt, PRD66(2002)

$$\rho^q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{-i\mathbf{b}_\perp \cdot \Delta_\perp} H_-^q(x, 0, -\Delta_\perp^2)$$

mapping in the transverse plane  
Impact parameter distribution



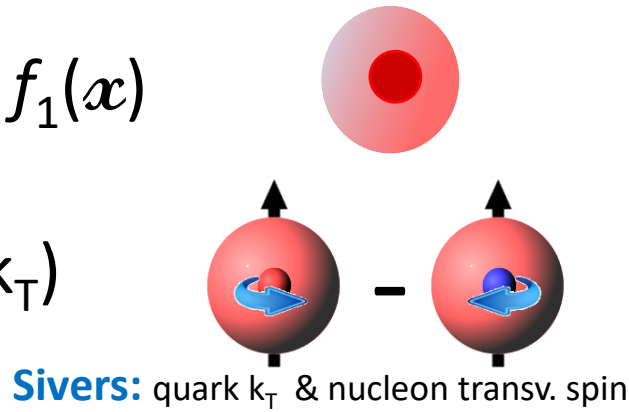
Correlation between the spatial distribution of partons and the longitudinal momentum fraction

# GPDs and Energy-Momentum Tensor and Confinement

$$\mathbf{H}^q(x, \xi, t) \xrightarrow{t \rightarrow 0} q(x) \text{ or } f_1(x)$$

“Elusive”

$$\mathbf{E}^q(x, \xi, t) \leftrightarrow f_{1T}(x, k_T)$$



mass & energy distribution

$$\int_{-1}^1 dx x H^a(x, \xi, t) = A^a(t) + \xi^2 d_1^a(t)$$

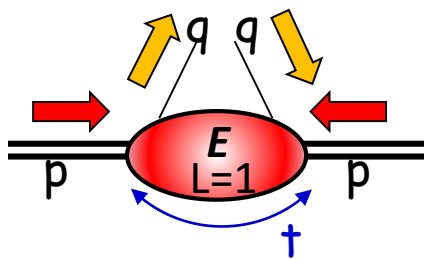
Angular momentum distribution

$$\int_{-1}^1 dx x E^a(x, \xi, t) = 2J^a(t) - A^a(t) - \xi^2 d_1^a(t)$$

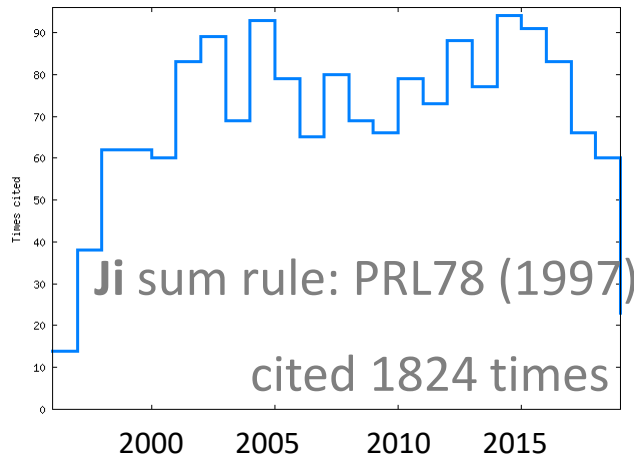
Force & Pressure distribution

$$2J^q = \lim_{t \rightarrow 0} \int x (\mathbf{H}^q(x, \xi, t) + \mathbf{E}^q(x, \xi, t)) dx$$

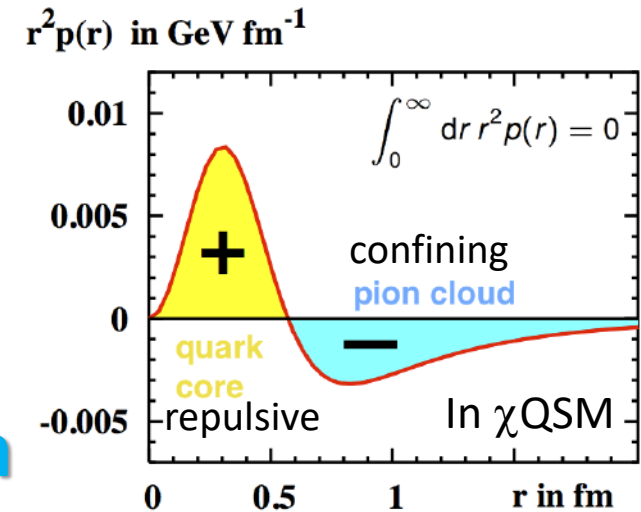
M. Polyakov, P. Schweitzer, Int.J.Mod.Phys. A33 (2018)



Relation to OAM



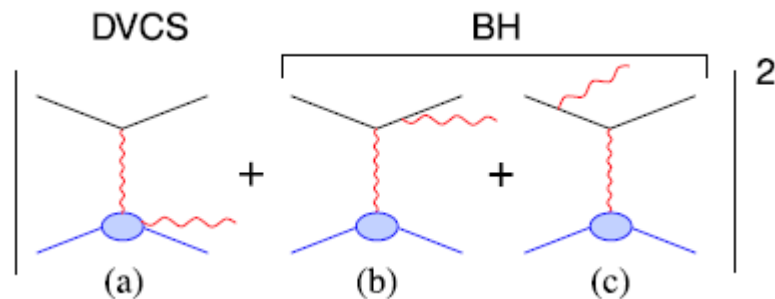
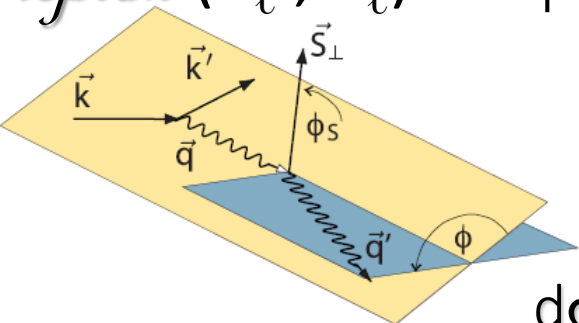
Pressure Distribution





# Deeply virtual Compton scattering (DVCS)

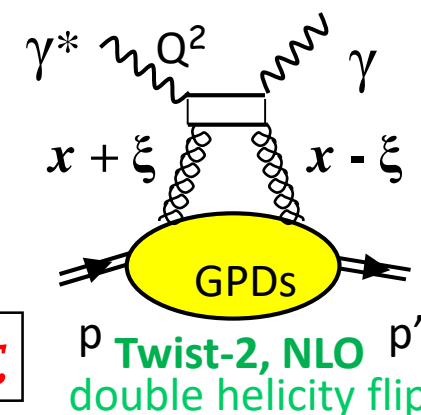
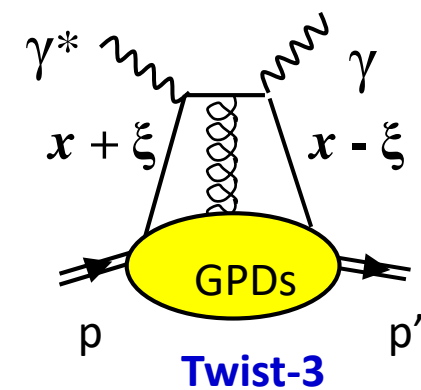
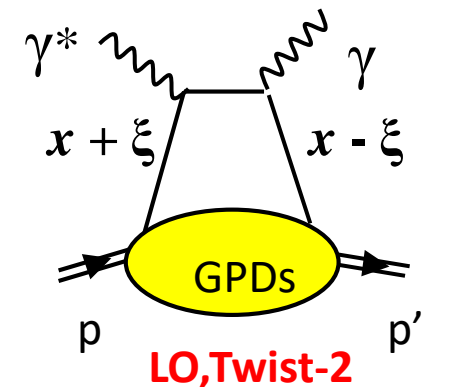
lepton ( $P_\ell, e_\ell$ ) and  $\phi$



$$d\sigma = |T^{BH}|^2 + |T^{DVCS}|^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 dt |d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \underbrace{\left( d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right)}_{|T^{DVCS}|^2} + \underbrace{\left( e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I \right)}_{\text{Interference Term}}$$

$$\begin{aligned} d\sigma^{BH} &\propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi \\ d\sigma_{unpol}^{DVCS} &\propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi \\ d\sigma_{pol}^{DVCS} &\propto s_1^{DVCS} \sin \phi \\ \text{Re } I &\propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi \\ \text{Im } I &\propto s_1^I \sin \phi + s_2^I \sin 2\phi \end{aligned}$$



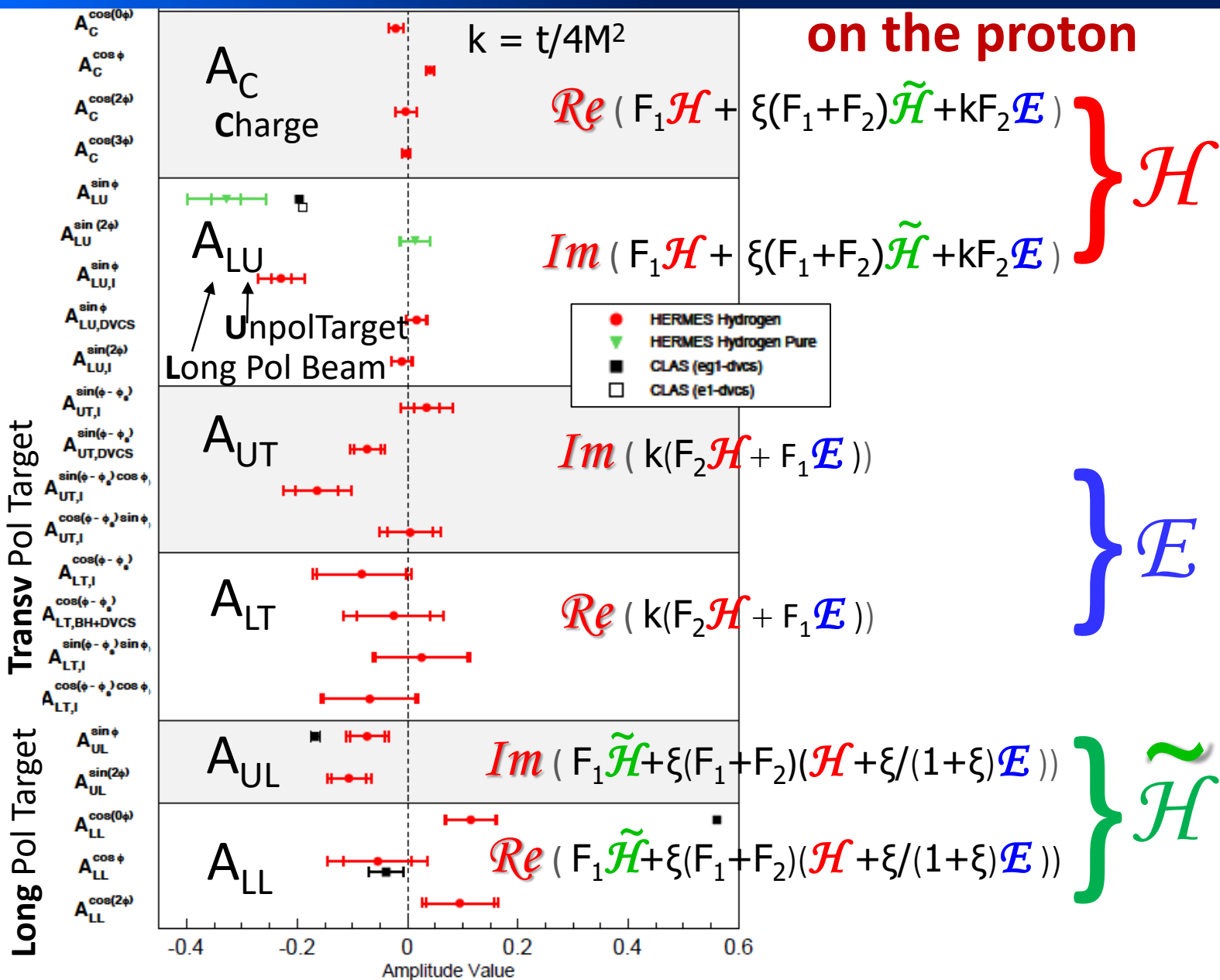
Changing  $P_\ell \rightarrow s_1^I = \text{Im } \mathcal{F}$

Changing  $e_\ell P_\ell \rightarrow c_1^I = \text{Re } \mathcal{F}$

$\mathcal{F} = F_1 \mathcal{H} + \xi(F_1 + F_2) \tilde{\mathcal{H}} + t/4m^2 F_2 \mathcal{E}$

**Twist-2, NLO double helicity flip**

# 2001-2012: A complete set of DVCS asymmetries at Hermes



HERMES 27 GeV provided a complete set of observables

2001: 1<sup>st</sup> DVCS publication as CLAS & H1

2007: end of data taking

**2012:** still important publications

JHEP 07 (2012) 032  $A_C A_{LU}$

JHEP10(2012) 042  $A_{LU}$

with recoil detection (2006-7)

Note: **the neutron** allows

✓ flavor decomposition

✓ access to  $\mathcal{E}$

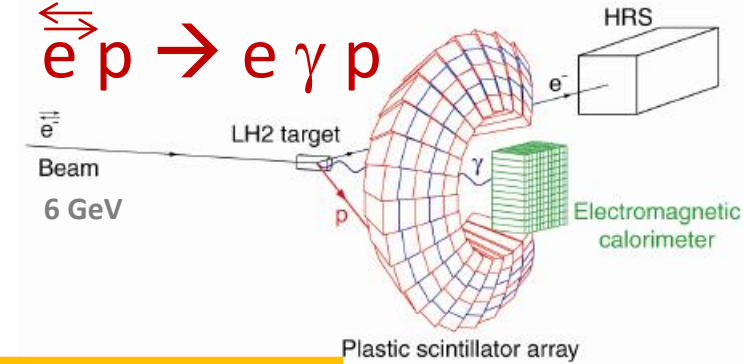
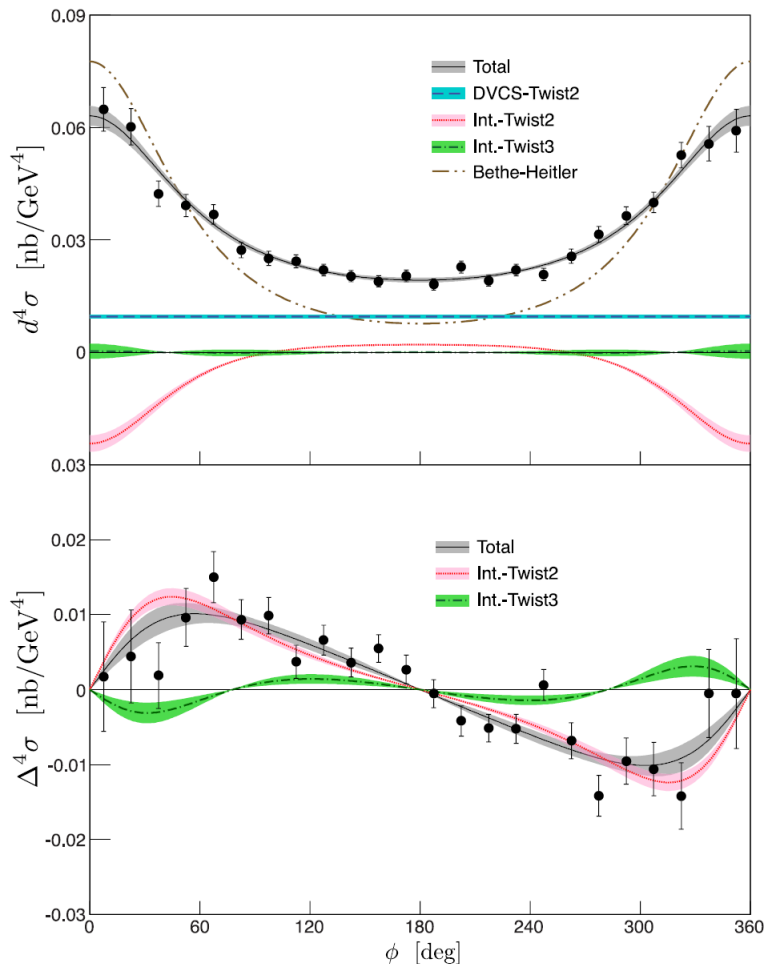
# 2004-2015: Beam Spin Sum and Diff for DVCS - HallA

E00-110 pioneer experiment in 2004 with magnetic spectrometer

$x_B=0.36$   $Q^2= 1.5, 1.9, 2.3 \text{ GeV}^2$  First analysis: Munoz et al. PRL97, 262002 (2006)

$x_B=0.34, x_B=0.39$   $Q^2= 2.1 \text{ GeV}^2$  Final analysis: Defurne et al., PRC92, 055202 (2015)

$x_B=0.36, Q^2= 2.3 \text{ GeV}^2, -t= 0.32 \text{ GeV}^2$



## Unpolarized cross section

$$\begin{aligned}
 d\sigma^{\leftarrow} + d\sigma^{\rightarrow} &\propto d\sigma^{BH} + d\sigma_{unpol}^{DVCS} + \text{Re } I \\
 &\rightarrow d\sigma^{BH} + \underbrace{c_0^{DVCS}}_{\text{blue}} + \underbrace{c_1^{DVCS}}_{\text{blue}} \cos \phi + \underbrace{c_2^{DVCS}}_{\text{green}} \cos 2\phi \\
 &\quad + \underbrace{c_0^I}_{\text{pink}} + \underbrace{c_1^I}_{\text{pink}} \cos \phi + \underbrace{c_2^I}_{\text{green}} \cos 2\phi + \underbrace{c_3^I}_{\text{green}} \cos 3\phi
 \end{aligned}$$

## Helicity Dependent cross section

$$\begin{aligned}
 d\sigma^{\leftarrow} - d\sigma^{\rightarrow} &\propto d\sigma_{pol}^{DVCS} + \text{Im } I \\
 &\rightarrow \underbrace{s_1^{DVCS}}_{\text{blue}} \sin \phi + \underbrace{s_1^I}_{\text{pink}} \sin \phi + \underbrace{s_2^I}_{\text{green}} \sin 2\phi
 \end{aligned}$$

→ Further DVCS/Interference separation with different beam energies with 2010 data

# 2010-2017: Beam Spin Sum and Diff for DVCS - HallA

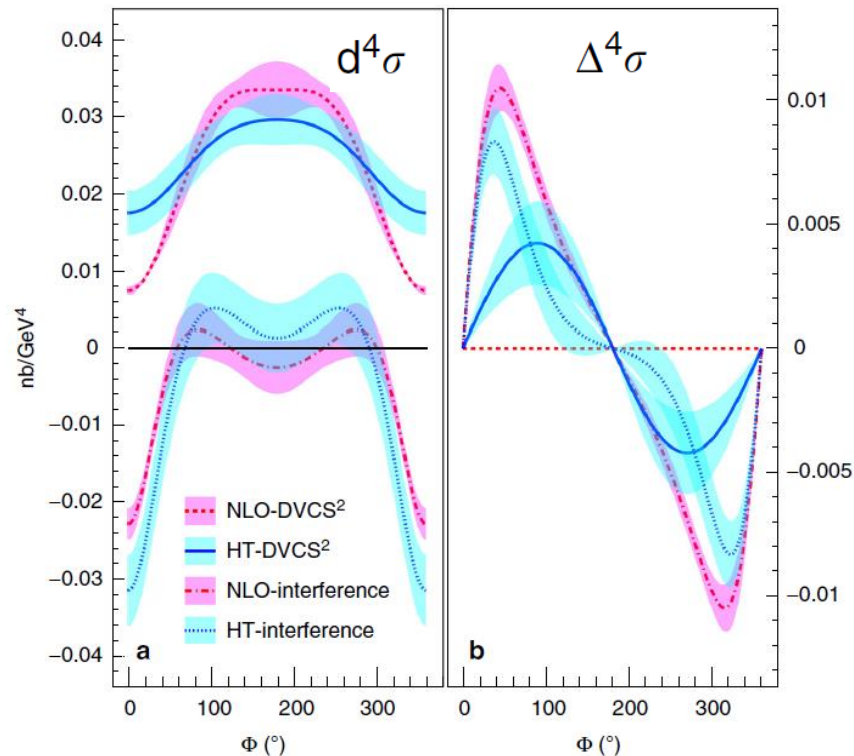
E07-007 Hall-A experiment in 2010 with magnetic spectrometer

Defurne et al., Nature Communications 8 (2017) 1408



$x_B=0.36$ ,  $Q^2=1.75 \text{ GeV}^2$ ,  $-t=0.30 \text{ GeV}^2$

Ebeam=5.55 GeV



## Unpolarized cross section

$$d^4\sigma \quad d\sigma^{\leftarrow} + d\sigma^{\rightarrow} \propto d\sigma^{BH} + d\sigma_{unpol}^{DVCS} + \text{Re } I$$

$$\longrightarrow d\sigma^{BH} + \underline{c_0^{DVCS}} + \underline{c_1^{DVCS} \cos \phi} + \underline{c_2^{DVCS} \cos 2\phi}$$

$$+ \underline{c_0^I + c_1^I \cos \phi} + \underline{c_2^I \cos 2\phi} + \underline{c_3^I \cos 3\phi}$$

## Helicity Dependent cross section

$$\Delta^4\sigma \quad d\sigma^{\leftarrow} - d\sigma^{\rightarrow} \propto d\sigma_{vol}^{DVCS} + \text{Im } I$$

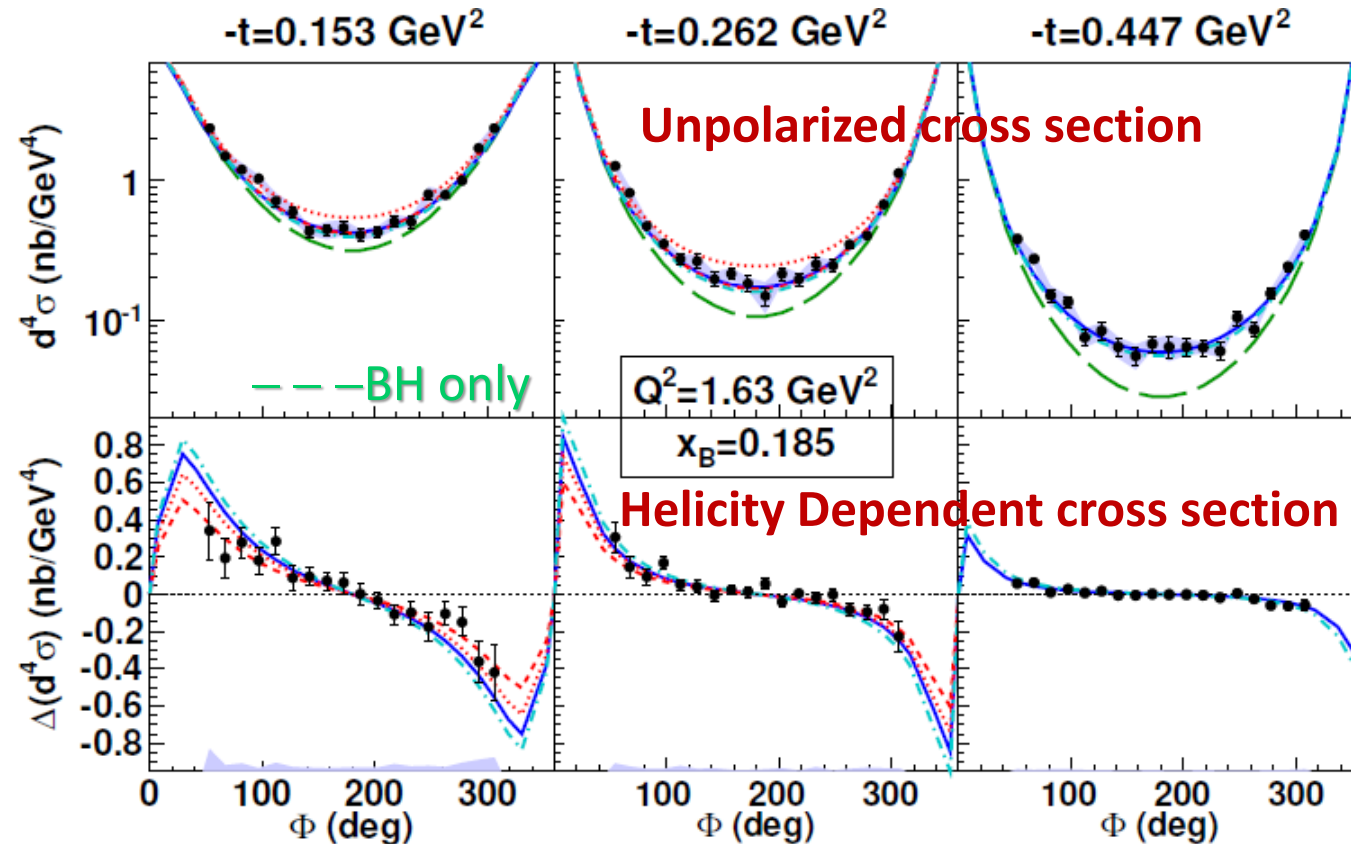
$$\longrightarrow \underline{s_1^{DVCS} \sin \phi} + \underline{s_1^I \sin \phi} + \underline{s_2^I \sin 2\phi}$$

2 solutions: higher-twist OR next-to-leading order

# 2005-2018: Beam Spin Sum and Diff for DVCS - CLAS

21 bins in  $(x_B, Q^2)$  or 110 bins  $(x_B, Q^2, t)$  3 months data taken in 2005

Girod et al. PRL100 (2008), Jo et al. PRL115 (2015), Hirlinger Saylor et al. PRC98 (2018)



models:

**VGG** Vanderhaeghen, Guichon, Guidal  
PRL80(1998), PRD60(1999), PPNP47(2001), PRD72(2005)  
1st model of GPDs improved regularly

**KMS12** Kroll, Moutarde, Sabatié, EPJC73 (2013)

using the **GK** model

Goloskokov, Kroll, EPJC42,50,53,59,65,74

for GPD adjusted on the hard exclusive meson production at small  $x_B$

**“universality”** of GPDs

**KM10a** --- (KM10 .....) Kumericki, Mueller, NPB (2010) 841

Flexible parametrization of the GPDs based on both a Mellin-Barnes representation and dispersion integral which entangle skewness and  $t$  dependences

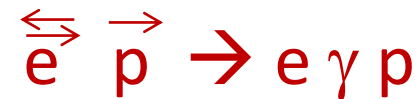
**Global fit** on the world data ranging from H1, ZEUS to HERMES, JLab

# 2009-2015: Single Spin and Double Spin - CLAS

$$A_{LU(UL)} = \frac{\alpha_{LU(UL)} \sin \phi}{1 + \beta \cos \phi}$$

$$A_{LL} = \frac{\kappa_{LL} + \lambda_{LL} \cos \phi}{1 + \beta \cos \phi}$$

2009: Longitudinally polarized NH3 target



Seder et al. PRL114, 032001 (2015)

Pisano et al. PRD91, 052014 (2015)

CLAS  $A_{LU}$  on  $H_2$

data

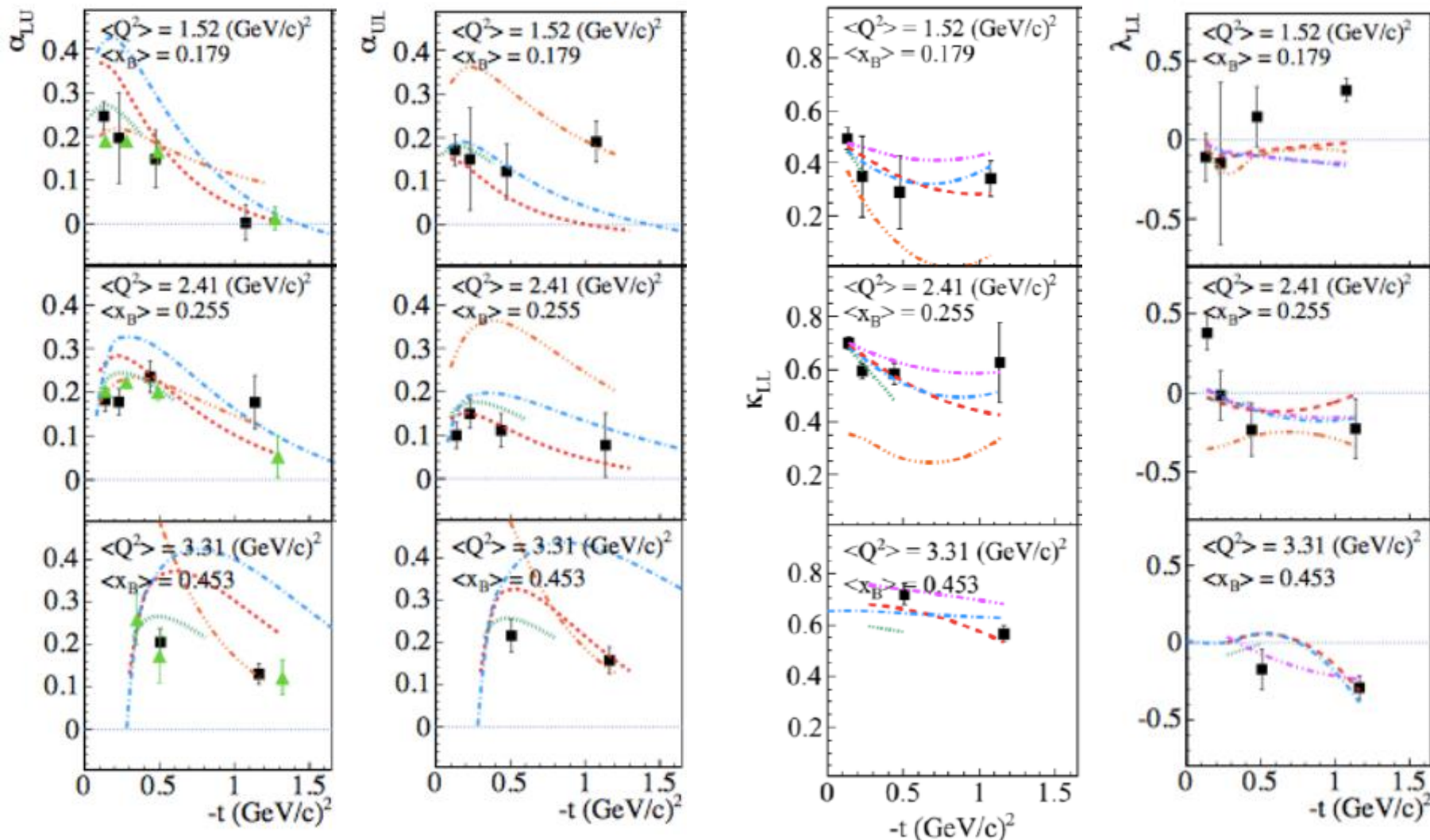
VGG Vanderhaeghen, Guichon, Guidal

KMM Kumericki, Mueller, Murray

GK Goloskokov, Kroll

GGL Goldstein, Gonzalez, Luiti

BH



# → nucleon tomography in the valence domain

Fit of 8 CFFs at L.O and L.T.

$\underline{\text{Im}\mathcal{H}}, \text{Re}\mathcal{H}, \text{Im}\mathcal{E}, \text{Re}\mathcal{E}, \underline{\text{Im}\tilde{\mathcal{H}}}, \text{Re}\tilde{\mathcal{H}}, \text{Im}\tilde{\mathcal{E}}, \text{Re}\tilde{\mathcal{E}}$

$$H(x, 0, 0) = q(x)$$

$$\tilde{H}(x, 0, 0) = \Delta q(x)$$

$$\int_{-1}^{+1} H dx = F_1$$

$$\int_{-1}^{+1} \tilde{H} dx = G_A$$

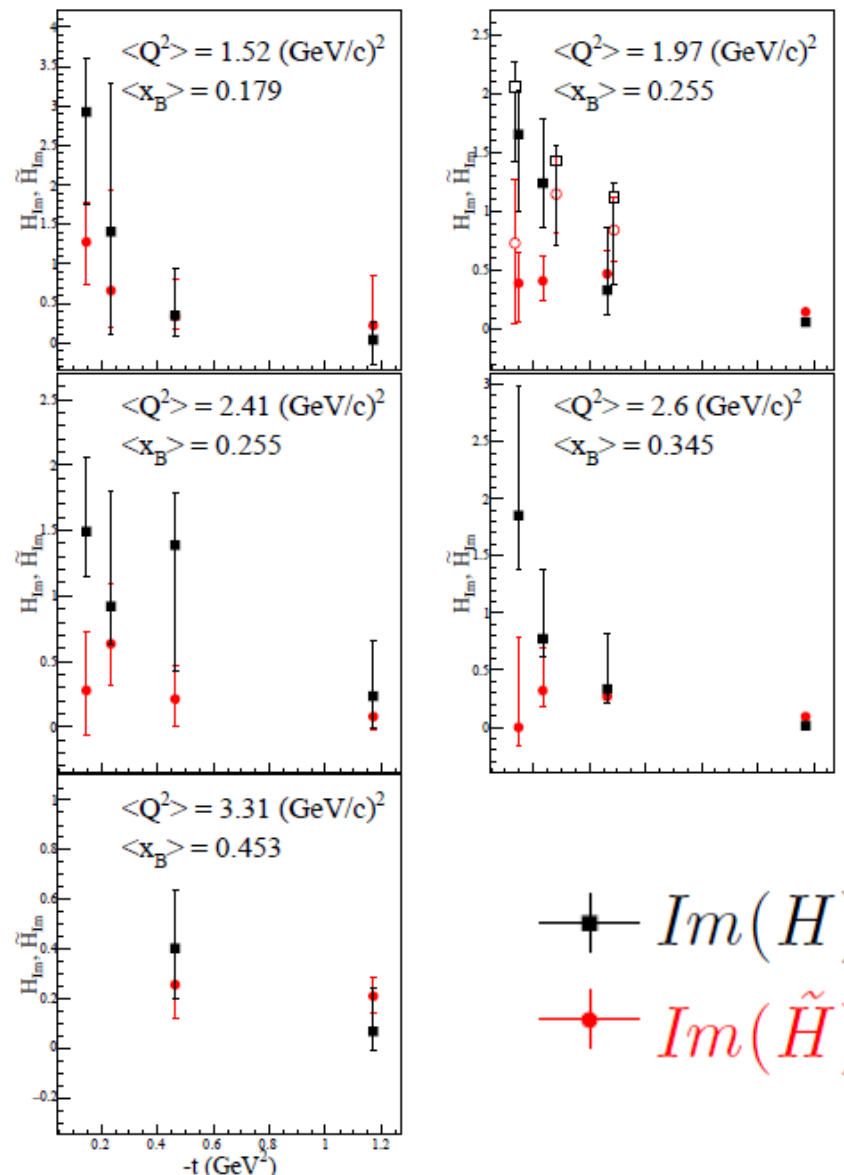
$\text{Im}(H) \rightarrow$  electromagnetic charge distribution

$\text{Im}(\tilde{H}) \rightarrow$  axial charge distribution

**Axial charge is more concentrated than electromagnetic charge**

Seder et al. PRL114, 032001 (2015)

Pisano et al. PRD91, 052014 (2015)



# → nucleon tomography in the valence domain

Fit of 8 CFFs at L.O and L.T. Dupré, Guidal, Nicolai, Vanderhaeghen, PRD95, 011501(R)(2017) Eur.Phys.J. A53 (2017)

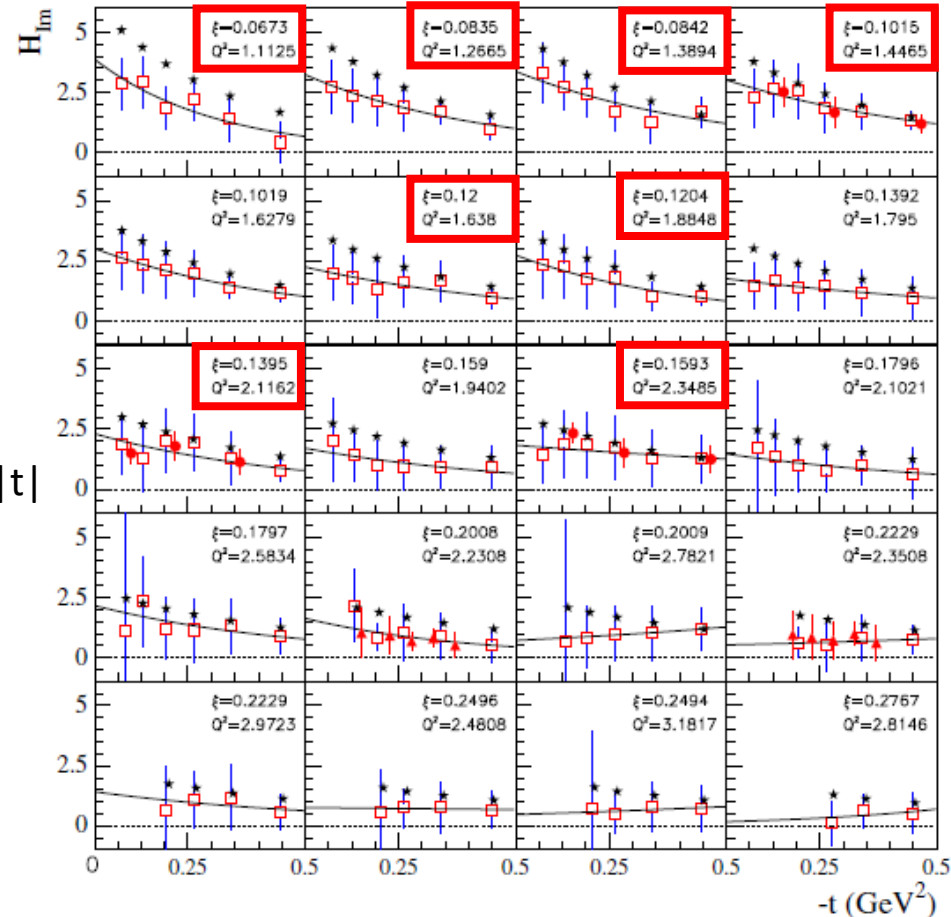
$s_1^I = \text{Im } F_1 \mathcal{H}$  is the best constrained

$$\rho^q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{-i\mathbf{b}_\perp \cdot \Delta_\perp} H_-^q(x, 0, -\Delta_\perp^2).$$

$$\langle b_\perp^2 \rangle^q(x) = -4 \frac{\partial}{\partial \Delta_\perp^2} \ln H_-^q(x, 0, -\Delta_\perp^2) \Big|_{\Delta_\perp=0}.$$

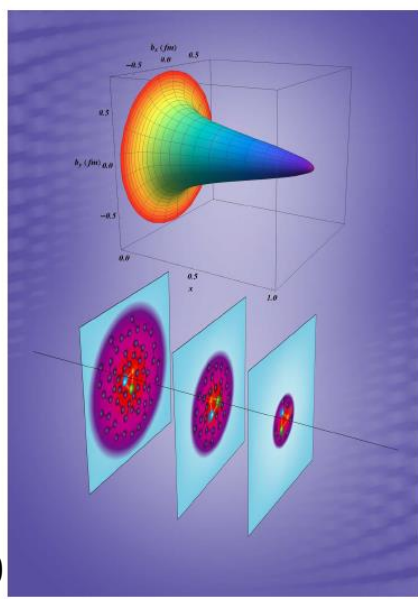
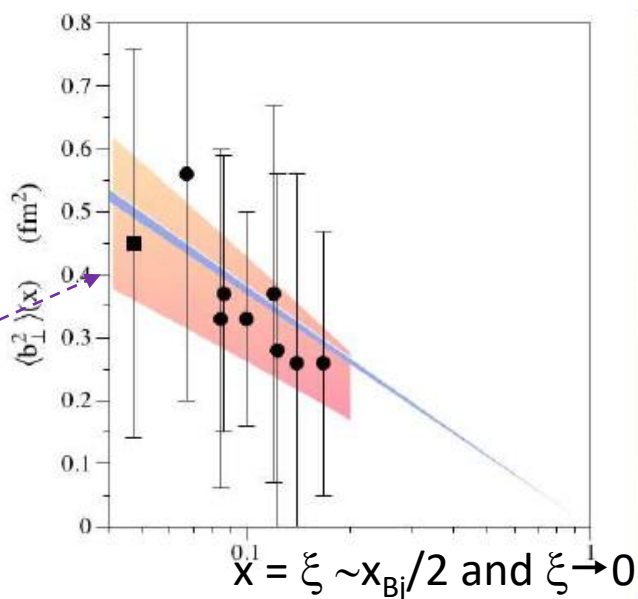
$$\langle b_\perp^2 \rangle \approx 4 B$$

— Fit  
 $\text{Im } \mathcal{H} = A e^{-B|t|}$



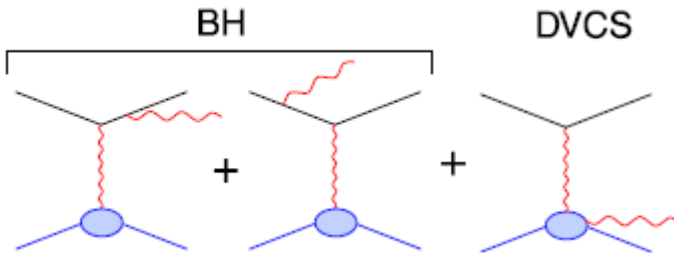
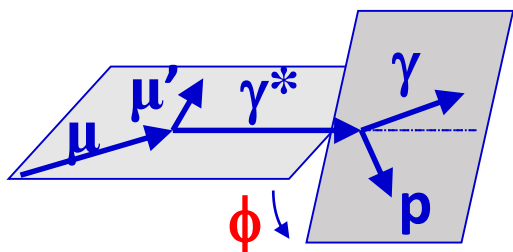
- CLAS  $\sigma$  and  $\Delta\sigma$
- ▲ HallA  $\sigma$  and  $\Delta\sigma$
- CLAS  $A_{UL}$  and  $A_{LL}$
- ★ VGG model

HERMES  
 + 8 data  
 from JLab





# 2012-2019: Beam Charge & Spin Sum for DVCS - COMPASS



$\mu^{\leftarrow}$  and  $\mu^{\rightarrow}$  beams of 160 GeV

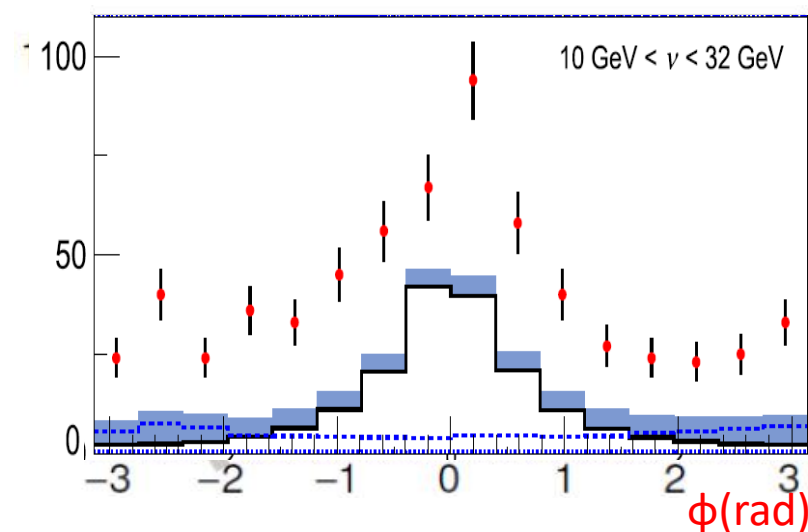
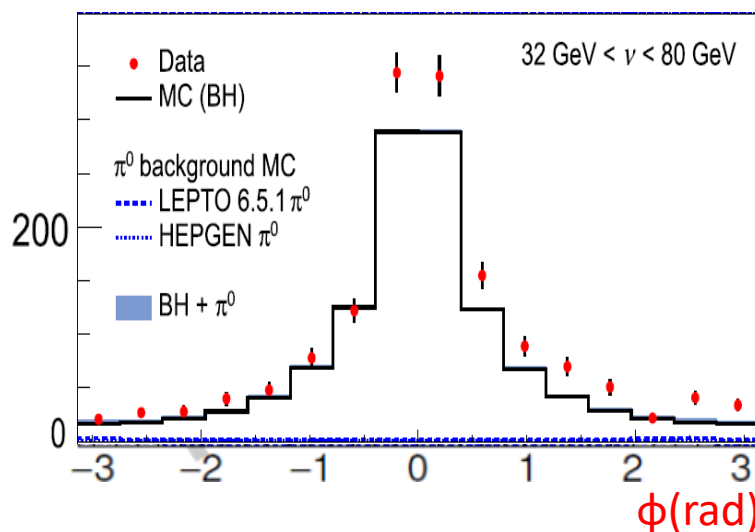
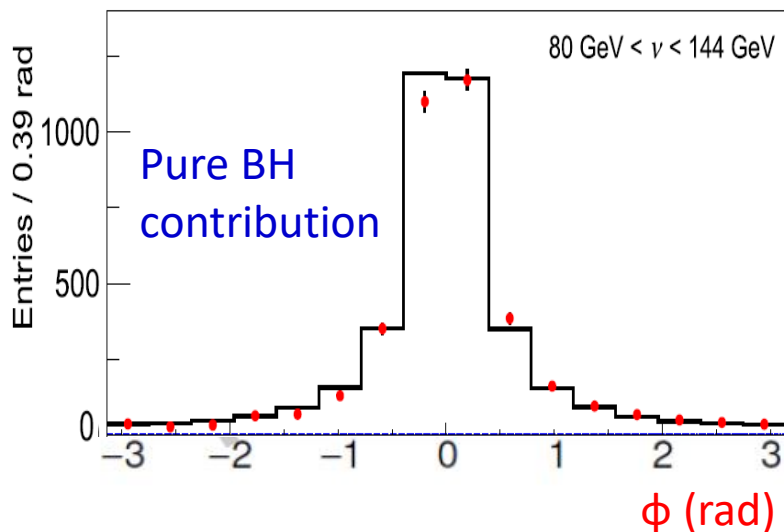
$$S_{CS,U} \equiv d\sigma^{\leftarrow} + d\sigma^{\rightarrow}$$

$$d\sigma \propto |T^{BH}|^2 + \text{Interference Term} + |T^{DVCS}|^2$$

$0.005 < x_{Bi} < 0.01$

$0.01 < x_{Bj} < 0.03$

$x_{Bj} > 0.03$



MC: — BH normalisation based on integrated luminosity

■  $\pi^0$  background contribution from SIDIS (LEPTO) + exclusive production (HEPGEN)

**DVCS > BH**

# 2012-2019: Beam Charge & Spin Sum for DVCS - COMPASS

At COMPASS using polarized positive and negative muon beams:

when **DVCS** > **BH**

$$S_{CS,U} \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-} = 2[d\sigma^{BH} + d\sigma_{unpol}^{DVCS} + \text{Im } I]$$

$$= 2[d\sigma^{BH} + c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi + s_1^I \sin \phi + s_2^I \sin 2\phi]$$

calculable  
can be subtracted

All the other terms are cancelled in the integration over  $\phi$

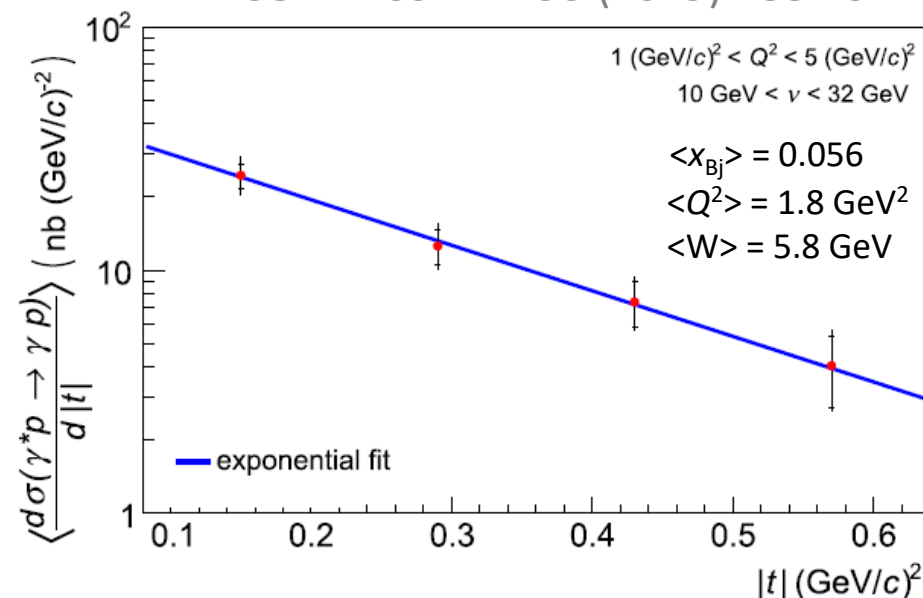
$$d\sigma^{DVCS}/dt = e^{-B|t|} = c_0^{DVCS}$$

$$B = (4.3 \pm 0.6_{\text{stat}} \pm 0.1_{\text{sys}}) (\text{GeV}/c)^{-2}$$

$$\langle r_{\perp}^2(x_B) \rangle \approx 2B'(x_B)$$

$$\sqrt{\langle r_{\perp}^2 \rangle} = (0.58 \pm 0.04_{\text{stat}} \pm 0.01_{\text{sys}} \pm 0.04_{\text{model}}) \text{ fm}$$

COMPASS PLB793 (2019) 188-194



# → nucleon tomography in the gluon and sea quark domains

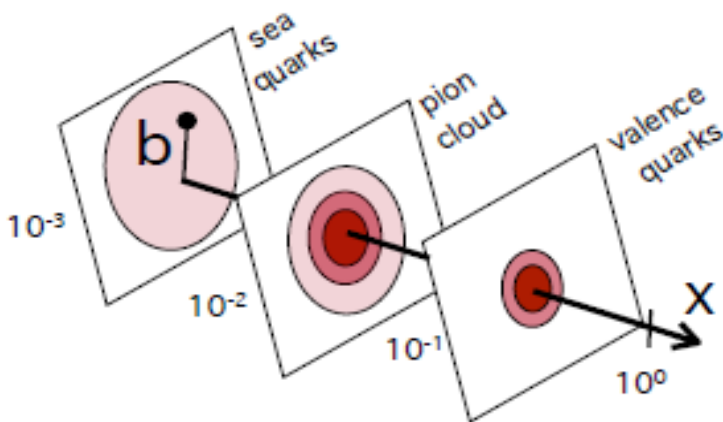
$$d\sigma^{DVCS}/dt = e^{-B'} |t| = c_0^{DVCS}$$

At COMPASS:  $\langle x_{Bj} \rangle = 0.056$ ;  
 $t$  varies from 0.08 to 0.64 GeV<sup>2</sup>

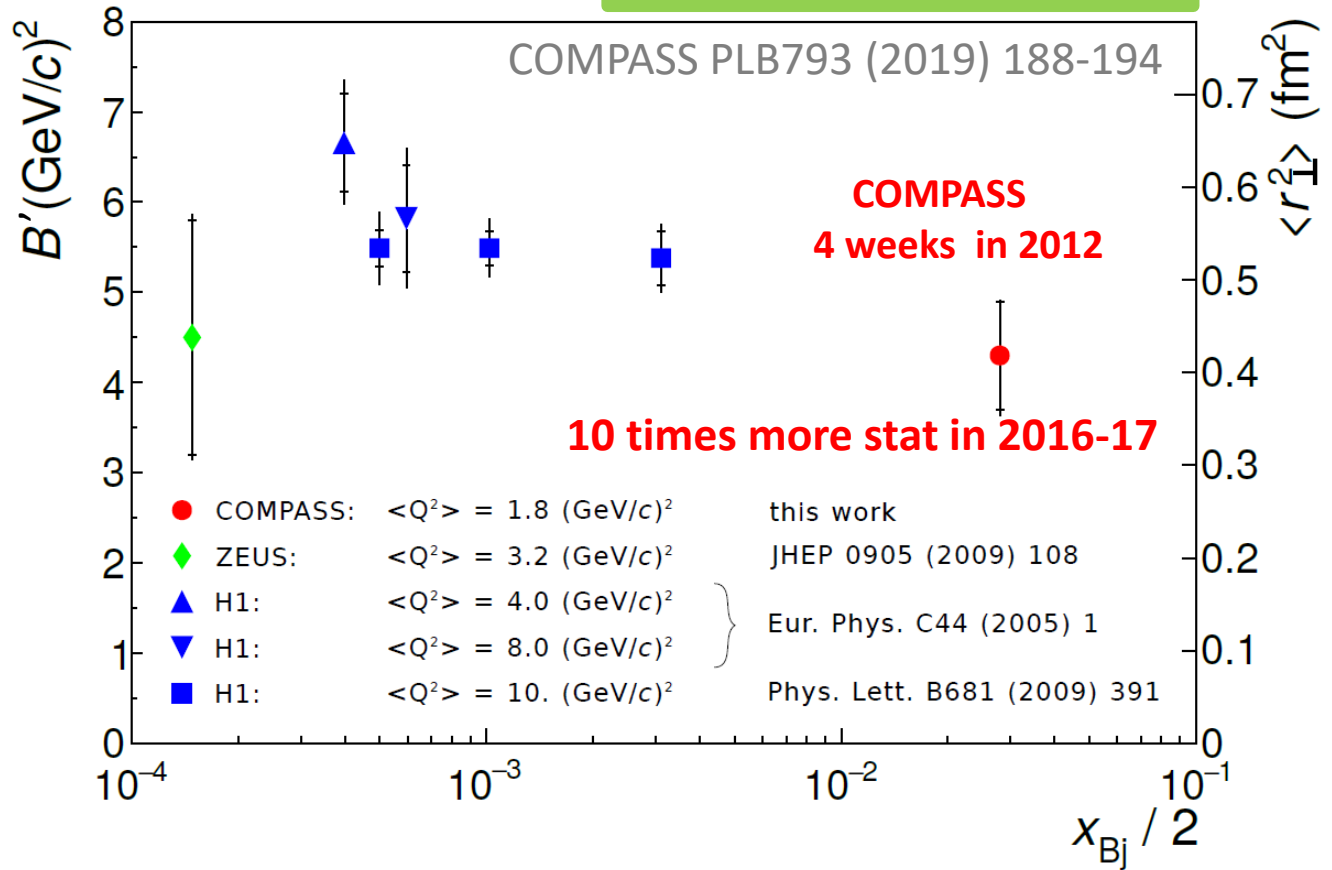
At small  $x_{Bj}$  and small  $t$ :

$$c_0^{DVCS} \propto 4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) - \frac{t}{M^2} \mathcal{E}\mathcal{E}^*$$

Dominance of  $Im\mathcal{H}$   
 (with respect of  $Re\mathcal{H}$  and other  $CFFs$ )



$$\langle r_{\perp}^2(x_B) \rangle \approx 2B'(x_B)$$



# → nucleon tomography in the gluon and sea quark domains

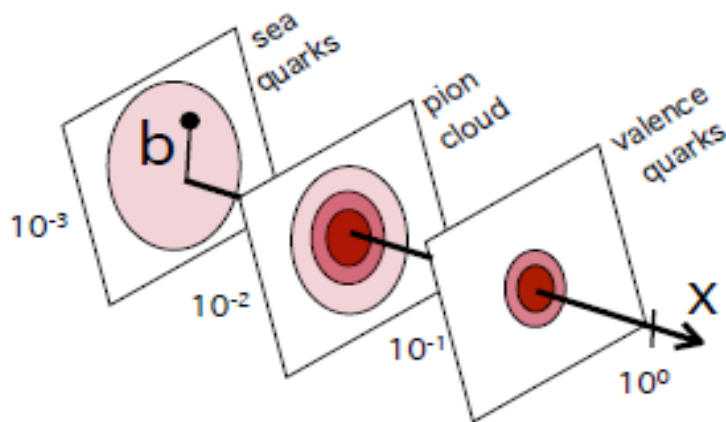
$$d\sigma^{DVCS}/dt = e^{-B'} |t| = c_0^{DVCS}$$

At COMPASS:  $\langle x_{Bj} \rangle = 0.056$ ;  
 $t$  varies from 0.08 to 0.64 GeV<sup>2</sup>

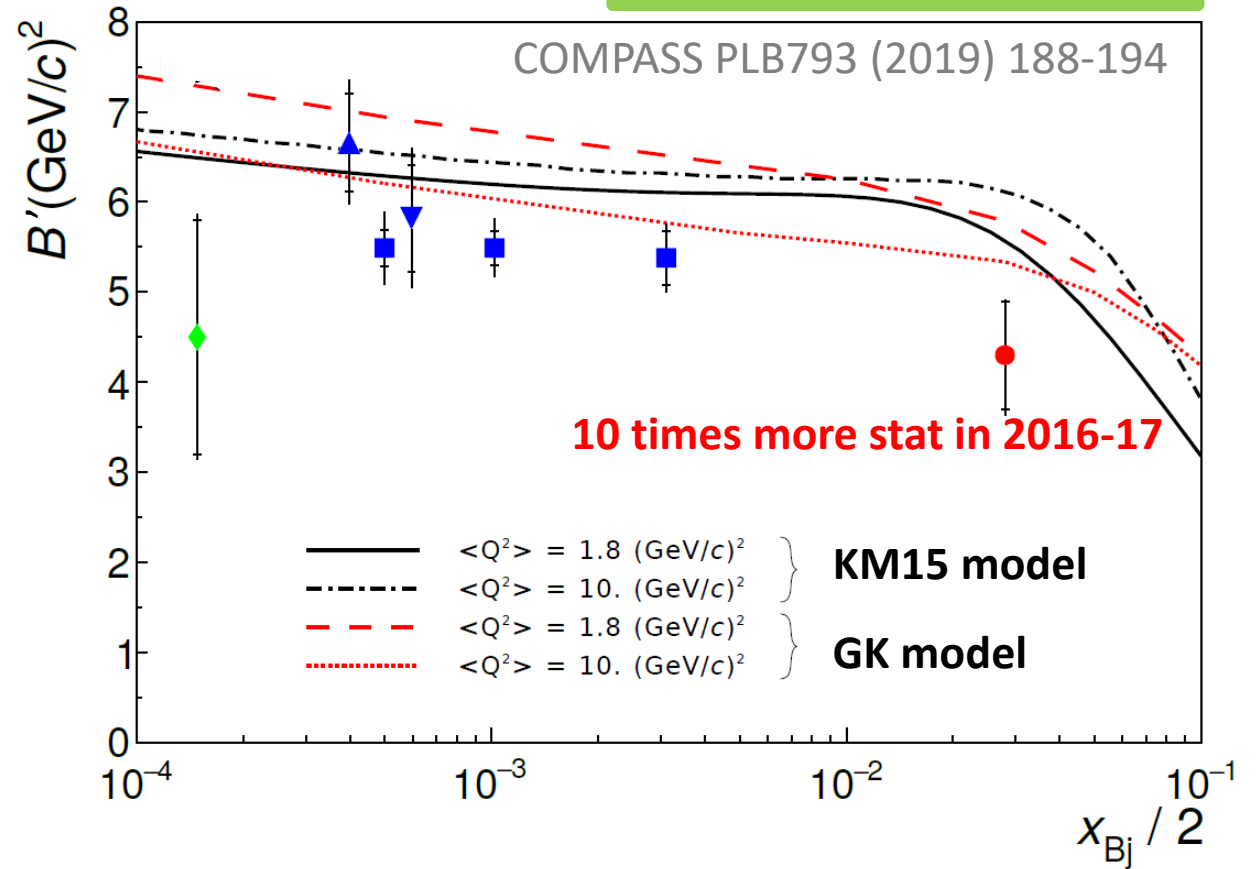
At small  $x_{Bj}$  and small  $t$ :

$$c_0^{DVCS} \propto 4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) - \frac{t}{M^2} \mathcal{E}\mathcal{E}^*$$

Dominance of  $Im\mathcal{H}$   
 (with respect of  $Re\mathcal{H}$  and other  $CFFs$ )



$$\langle r_{\perp}^2(x_B) \rangle \approx 2B'(x_B)$$



# → D-term and Pressure distribution in the proton

$\Delta(t)$  subtraction constant of the DVCS dispersion relation:

$$\text{Re}\mathcal{H}(\xi, t) = \Delta(t) + \frac{1}{\pi} \text{P.V.} \int_0^1 dx \left( \frac{1}{\xi - x} - \frac{1}{\xi + x} \right) \text{Im}\mathcal{H}(x, t)$$

Relation with  $D(z,t)$ , the **D-term** of the GPD

$$-1 < z = \frac{x}{\xi} < 1$$

& with  $d_1^q$ , the **proton gravitational FF** (the spherical Bessel transform of the pressure):

$$\Delta(t) = 2 \sum_q Q_q^2 \int_{-1}^1 dz \frac{D_q(z,t)}{1-z} = 4 \sum_q Q_q^2 (d_1^q(t) + \dots), q = u, d, \dots$$

next order terms  $\ll$

Q is the quark charge

With assumptions, considering only u and d quarks:

$$d_1^Q(t) = \frac{9}{10} \Delta(t)$$

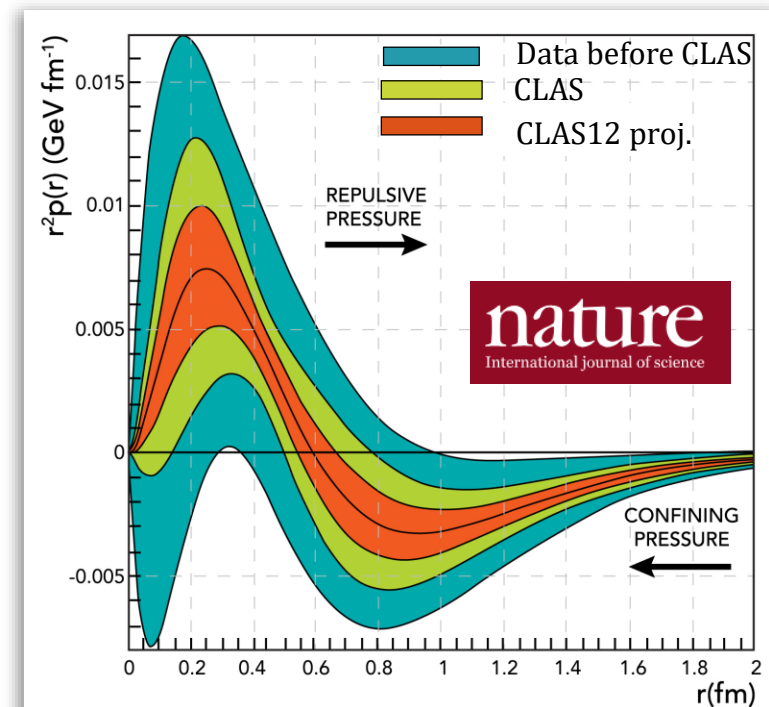
The spherical Bessel transform of the pressure :

M.V. Polyakov, Phys. Lett. B555 (2003) 57

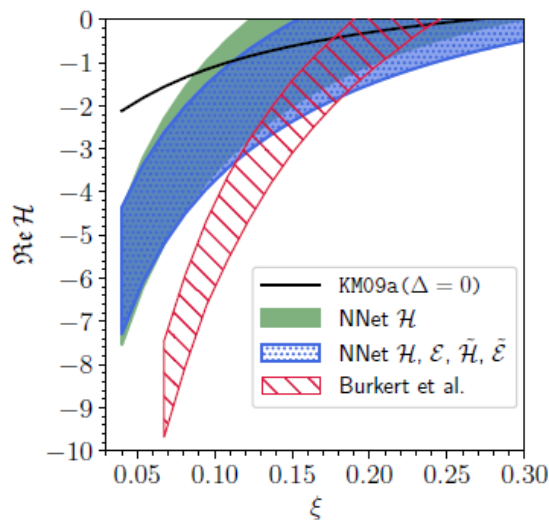
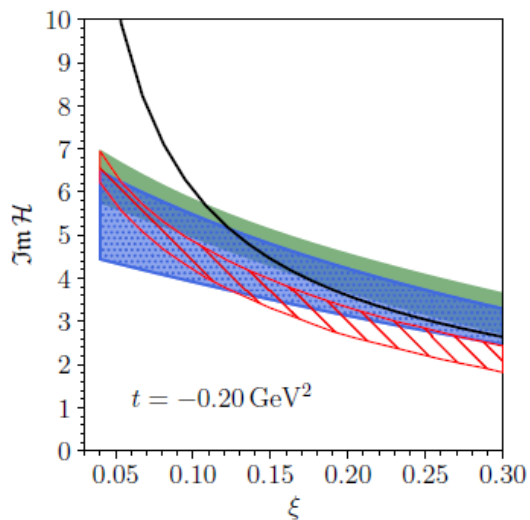
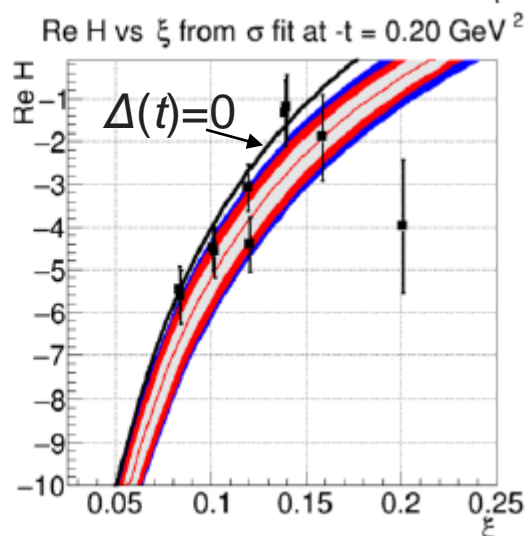
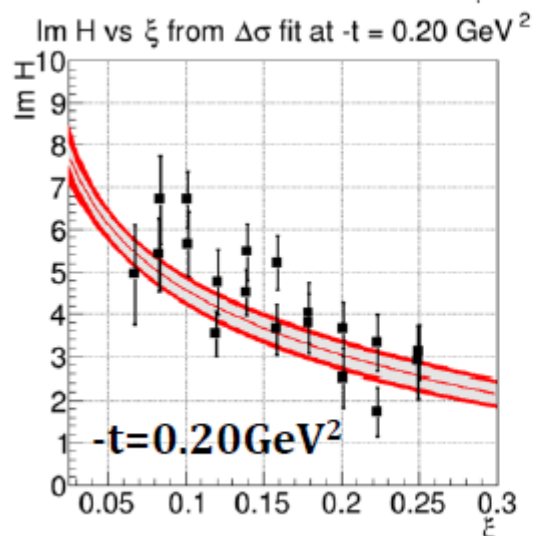
M.V. Polyakov, P. Schweitzer,

Int.J.Mod.Phys. A33 (2018)

$$d_1(t) \propto \int d^3\mathbf{r} \frac{j_0(r\sqrt{-t})}{2t} p(r)$$



# → D-term and Pressure distribution in the proton



V. Burkert et al., Nature 557, 396-399 (2018)

accurate  $\Delta(t)$  to determine D-term and pressure within some assumptions

$$\Delta(t=0) = -1.63 \pm 0.11 \pm 0.24$$

This is a critical result, required for dynamical stability of the proton. Deeply rooted in chiral symmetry breaking.

*however improvement of uncertainties*

*Using flexible parametrization by neural networks*

K. Kumericki, Nature 570, E1–E2 (2019)

$$\Delta(t) = 0.78 \pm 1.5 \text{ (statistical uncertainty)}$$

with almost no dependence on  $t$

→ D-term and pressure consistent with 0

→ waiting for more data sensitive to Re H

(importance of  $\mu^\pm$  at COMPASS and  $e^+$  at JLab)

# next future: Beam Charge and Spin Diff @ COMPASS

$$\mathcal{D}_{CS,U} \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-} \xrightarrow{L.T.} c_0^I + c_1^I \cos \phi$$

$$c_1^I = \text{Re } F_1 \mathcal{H}$$

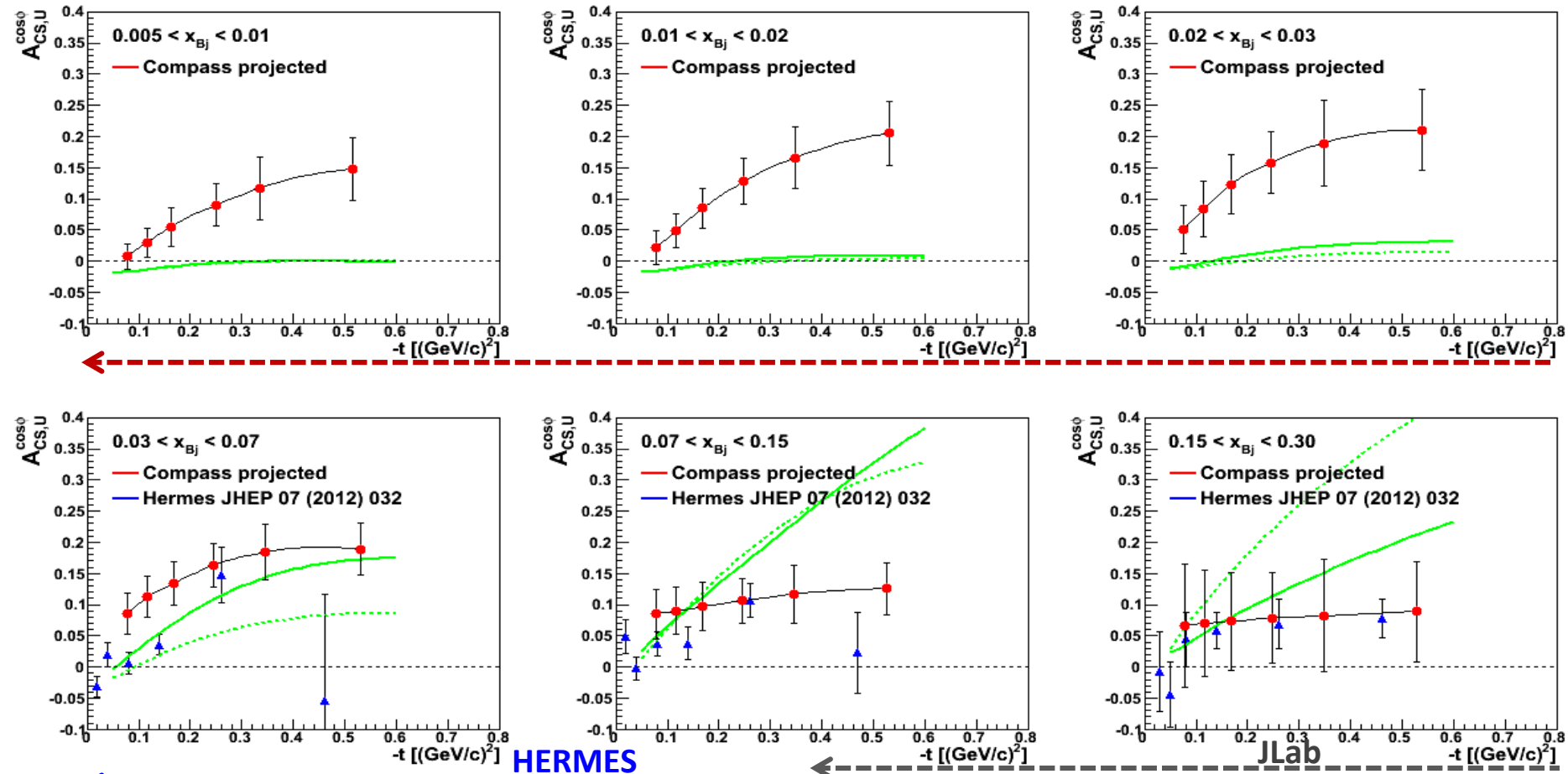
$\text{Re } \mathcal{H} > 0$  at H1  
 $< 0$  at HERMES  
 Value of  $x_B$  for the node?

Predictions with **VGG** and **KM10**

2016-17: 2x 6 months of data taking

Analysis on going

Impact on the D-term



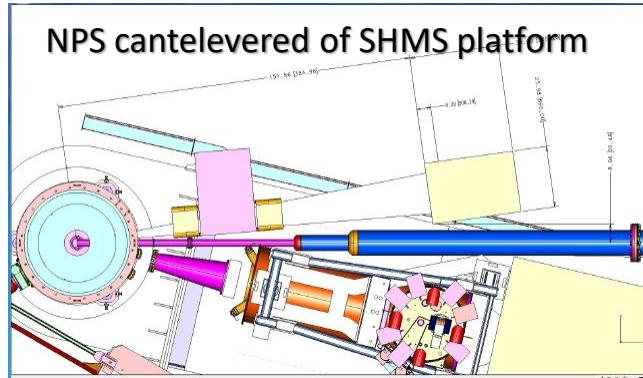
**COMPASS** 2 years of data  $E_\mu = 160 \text{ GeV}$   $1 < Q^2 < 8 \text{ GeV}^2$

# next future: Beam Spin Sum and Diff @ JLab12

with high resolution magnetic spectrometer  
+ Calorimeter in Halls A and C

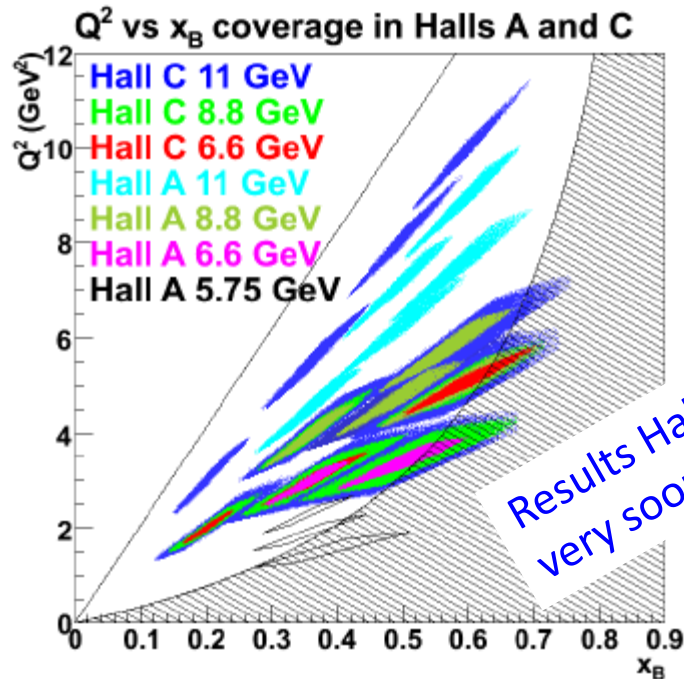
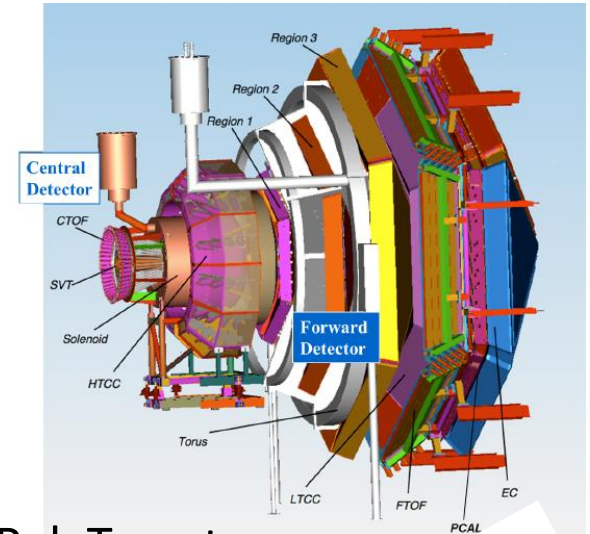
done in 2016-17: Hall A: E12-06-114  
>2021: Hall C: E12-13-010

Different beam energies for a  
Rosenbluth-like DVCS<sup>2</sup>/Interf.  
separation

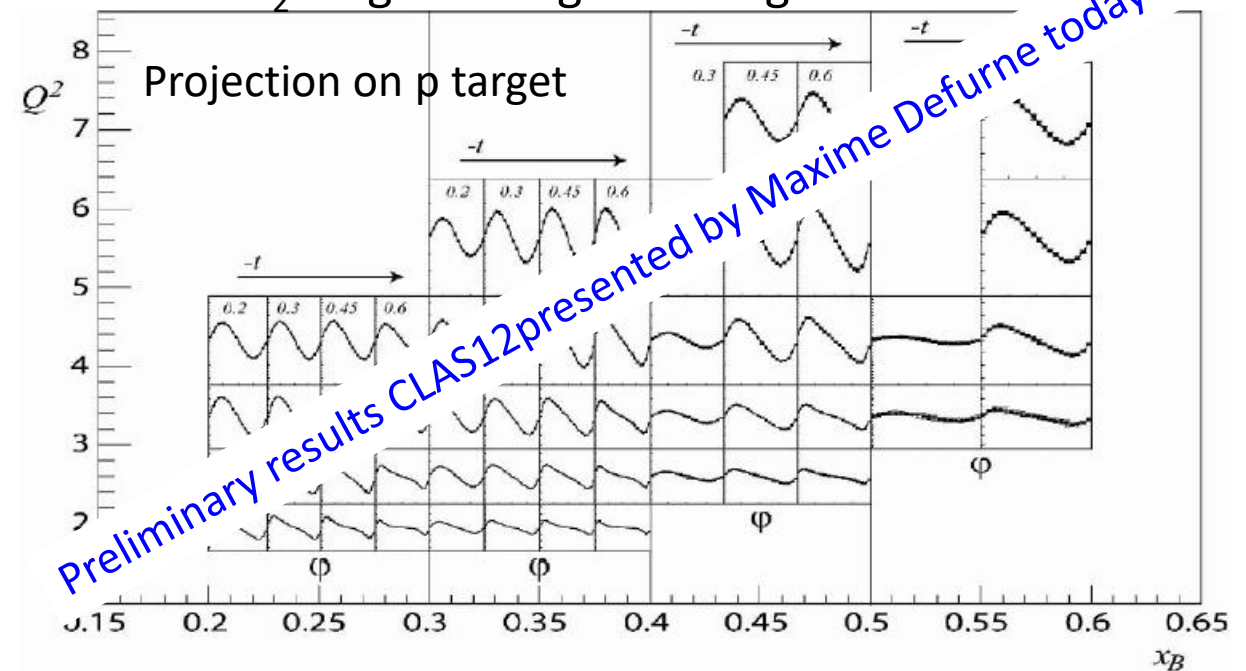


with CLAS12

E12-06-119  
2018-19: LH2  
2020: Long Pol Target



LH<sub>2</sub> Target & Long. Pol. Target





# 2007-8: GPD $\mathcal{E}$ from Jlab 6 GeV and HERMES

$$\ell d \rightarrow \ell n \gamma (p)$$

$$\Delta\sigma_{LU}^{\sin\phi} = \text{Im} (F_{1n} \mathcal{H} + \xi (F_{1n} + F_{2n}) \tilde{\mathcal{H}} + t/4m^2 F_{2n} \mathcal{E})$$

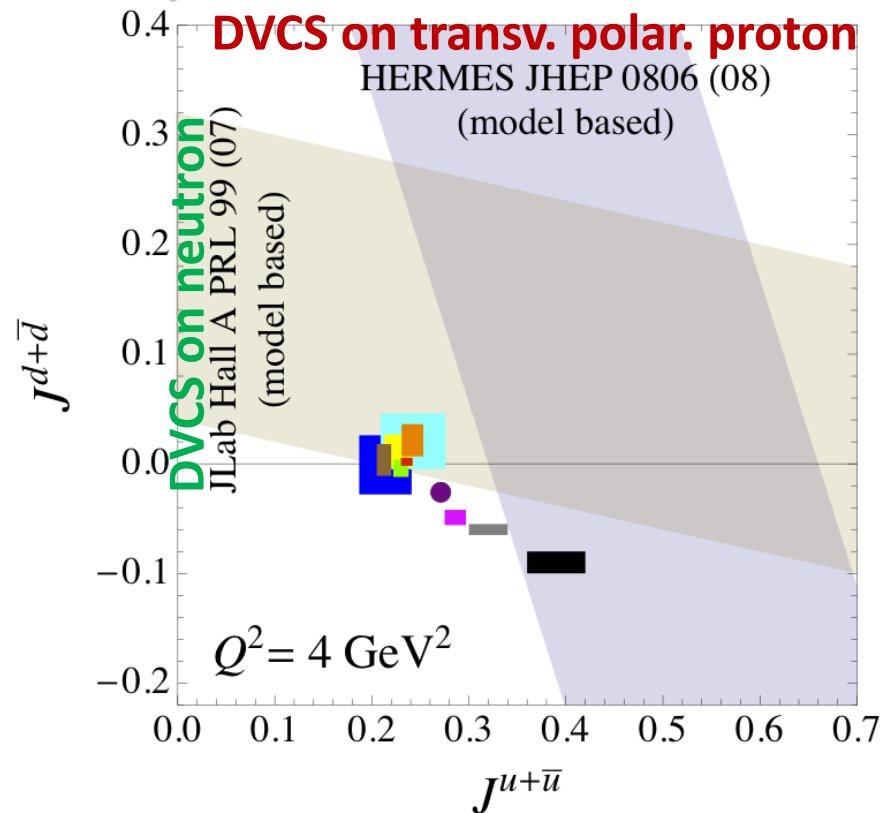
analysis still on going for a Hall-A experiment done in 2010  
 Note: neutron target  $\rightarrow$  quark flavor separation

$$\vec{\ell} p^\uparrow \rightarrow \ell p \gamma$$

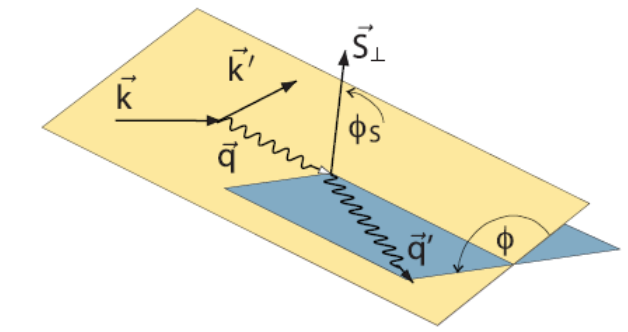
$$\Delta\sigma_{UT}^{\sin(\phi - \phi_s) \cos\phi} = -t/4m^2 \text{Im} (F_{2p} \mathcal{H} - F_{1p} \mathcal{E})$$

$$\Delta\sigma_{LT}^{\sin(\phi - \phi_s) \cos\phi} = -t/4m^2 \text{Re} (F_{2p} \mathcal{H} - F_{1p} \mathcal{E})$$

## Model dependent extraction of $J^u$ and $J^d$



- Goloskokov & Kroll, EPJ C59 (09) 809
- Diehl et al., EPJ C39 (05) 1
- Guidal et al., PR D72 (05) 054013
- Liuti et al., PRD 84 (11) 034007
- Bacchetta & Radici, PRL 107 (11) 212001
- LHPC-1, PR D77 (08) 094502
- LHPC-2, PR D82 (10) 094502
- QCDSF, arXiv:0710.1534
- Wakamatsu, EPJ A44 (10) 297
- Thomas, PRL 101 (08) 102003
- Thomas, INT 2012 workshop



**LATTICE QCD**

Dudek et al., EPJA48 (2012)

# next future: GPD $\mathcal{E}$ @ JLab12 with CLAS12

$\ell d \rightarrow \ell n \gamma (p)$

$\vec{\ell} p \uparrow \rightarrow \ell p \gamma$

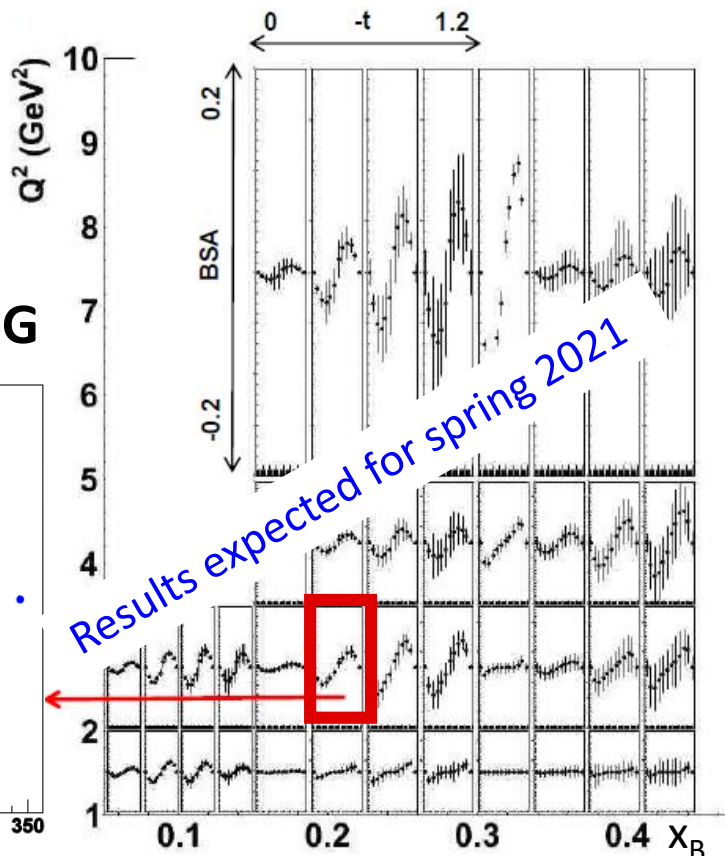
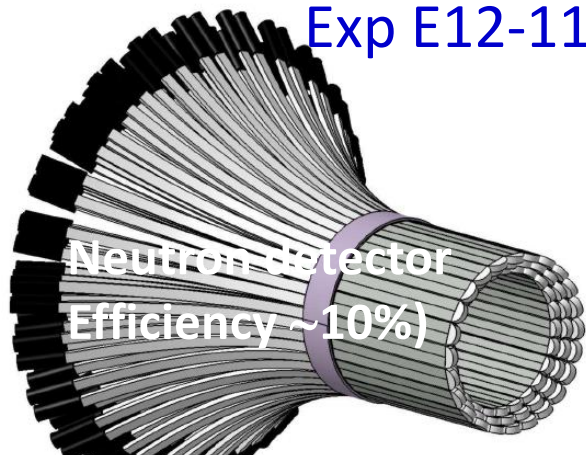
$$\Delta\sigma_{LU}^{\sin\phi} = \text{Im} (F_{1n} \mathcal{H} + \xi (F_{1n} + F_{2n}) \tilde{\mathcal{H}} + t/4m^2 F_{2n} \mathcal{E})$$

$$\Delta\sigma_{UT}^{\sin(\phi - \phi_s) \cos\phi} = -t/4m^2 \text{Im} (F_{2p} \mathcal{H} - F_{1p} \mathcal{E})$$

$$\Delta\sigma_{LT}^{\sin(\phi - \phi_s) \cos\phi} = -t/4m^2 \text{Re} (F_{2p} \mathcal{H} - F_{1p} \mathcal{E})$$

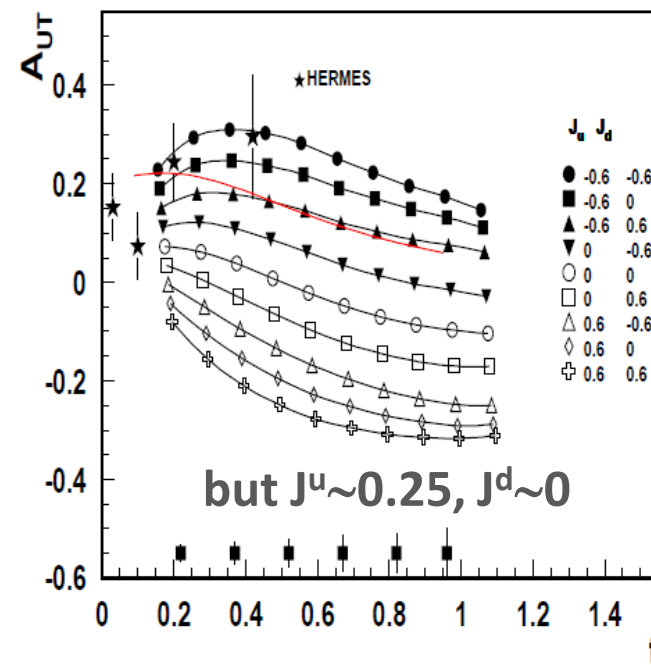
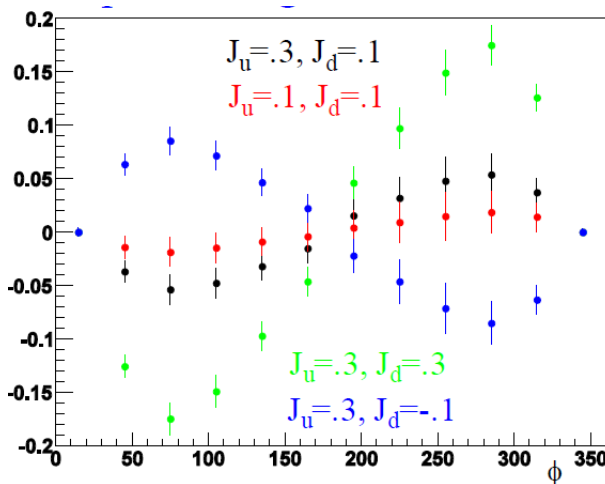
Exp E12-11-003: DVCS on the neutron  
2019-20: 90 days on LD2 target  
Lumi =  $10^{35} \text{ cm}^{-2} \text{ s}^{-1}/\text{nucleon}$

Exp E12-12-010: DVCS on a transversely polarized HD-Ice target Pol H = 60% Pol D = 35%  
2021: 110 days on HD-Ice target  
Lumi =  $5 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}/\text{nucleon}$



Results expected for spring 2021

Model prediction using VGG



# 2010-2020 : DVCS off the neutron in Hall A @ 6 GeV



$$\ell d \rightarrow \ell n \gamma (p)$$

Benali et al., Nature Physics 16, 191-198 (2020)

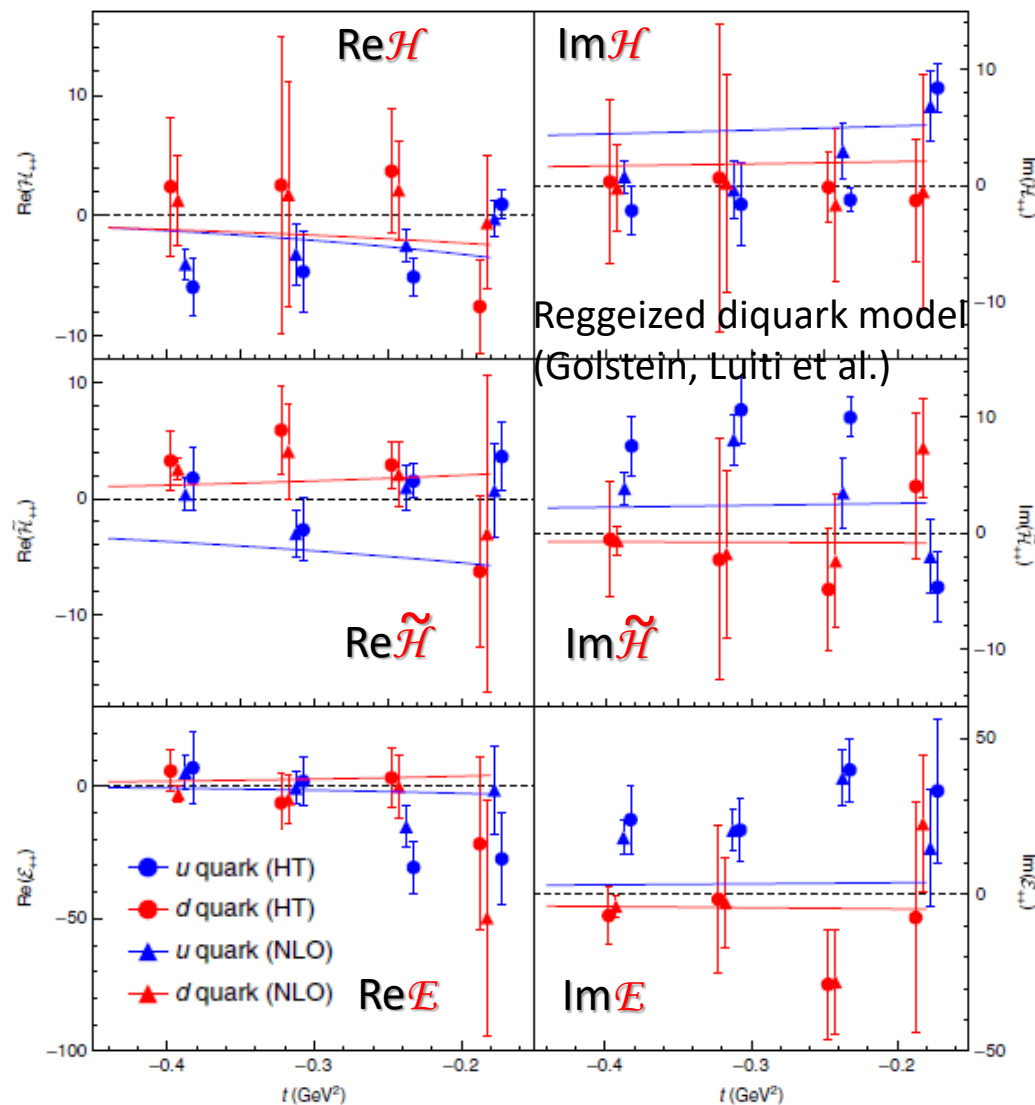
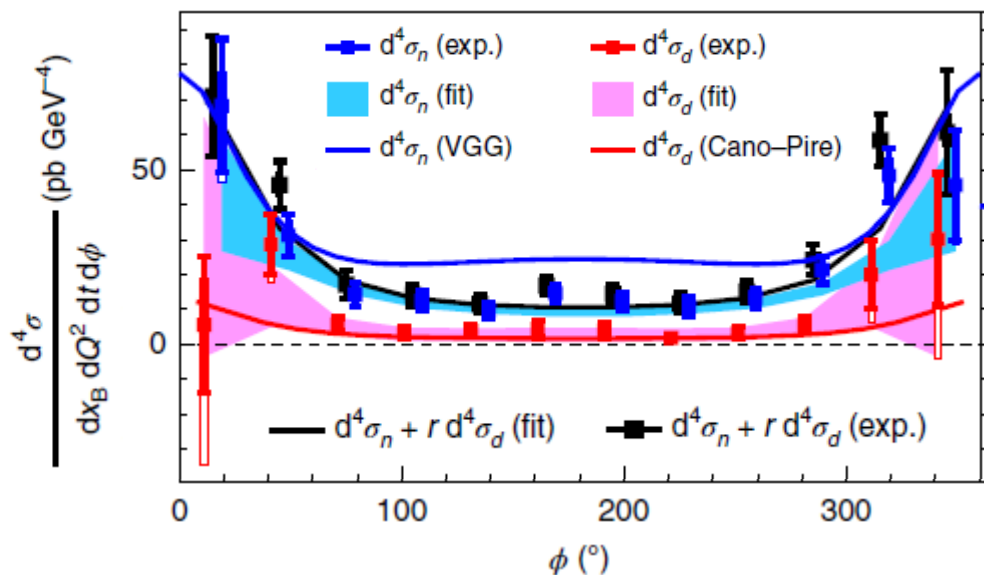
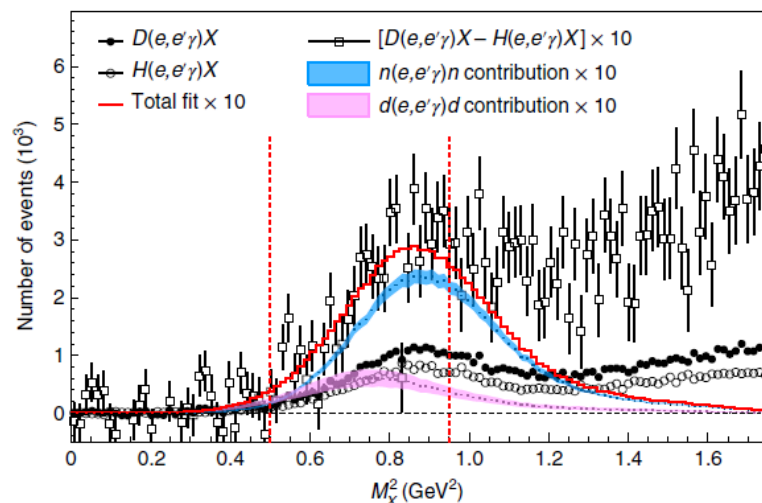
Flavor separation of CFFs when combined with p-DVCS

Coherent deuteron

+

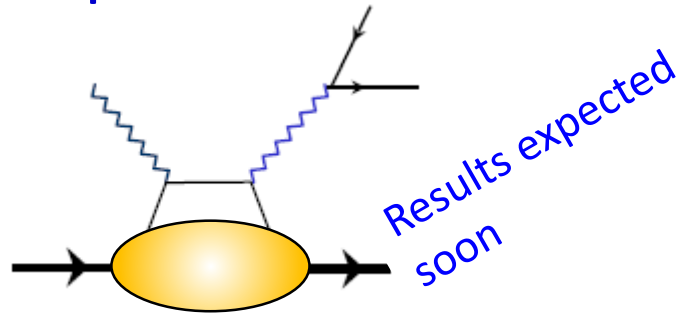
quasi-free neutron

DVCS  $\chi$ sections off LD2

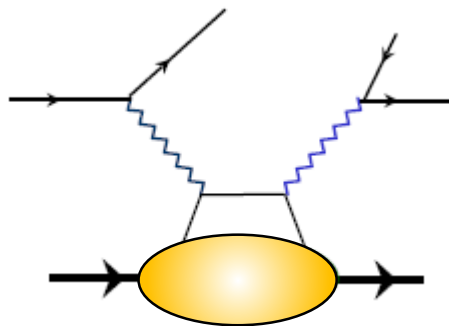


# And other paths to get GPDs

## Study of protons and neutrons



Time Like Compton Scattering



Double DVCS

Projects which start to be explored with the high luminosity of JLab12 (in Hall-C or with CLAS12 and Solid)

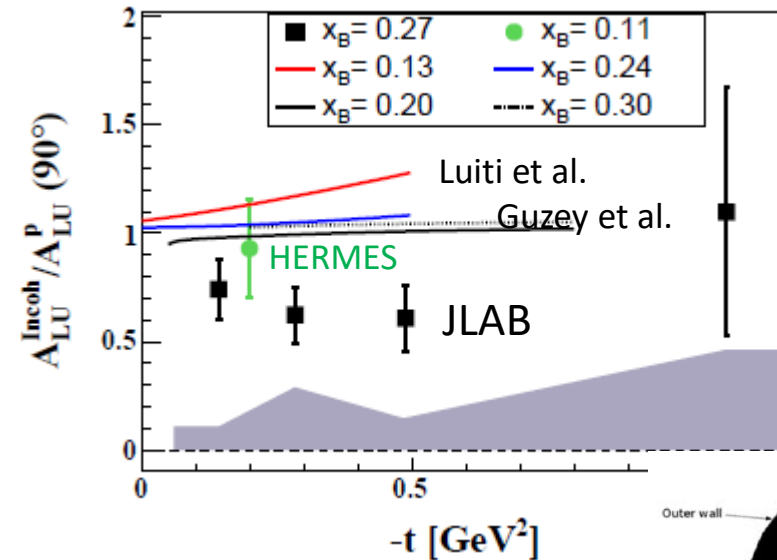
## Study of nuclei

(HERMES, JLab6, JLab12)

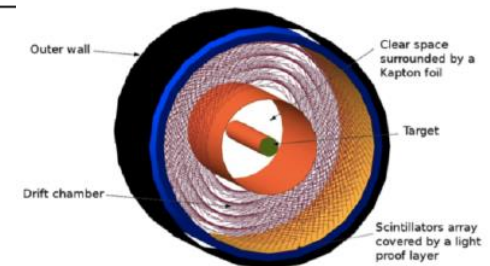
First measurement on He4: Hattawy, PRL 119 (2017)  
Spin 0 target, one chiral even GPD

Off bound protons: Hattawy, PRL 123(2019)

Ratio of the bound to the free proton at  $\phi=90^\circ$

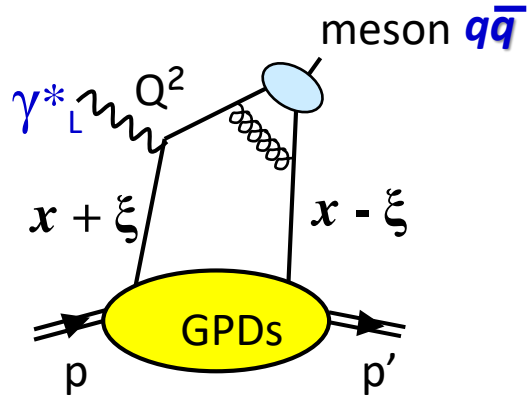


Other projects with the recoil detector ALERT

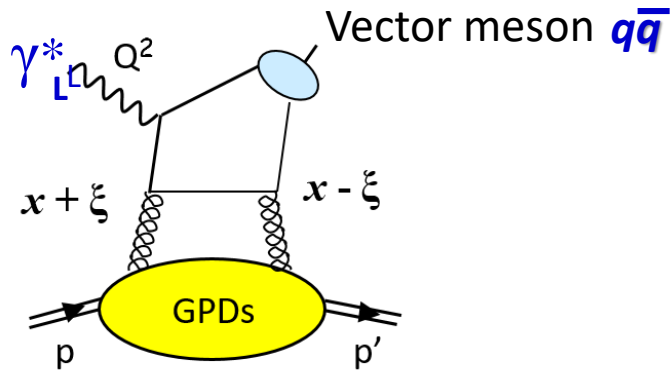


# GPDs and Hard Exclusive Meson Production

Quark contribution



Gluon contribution at the same order in  $\alpha_s$



The meson wave function  
Is an additional non-perturbative term

4 chiral-even GPDs: helicity of parton unchanged

$H^q(x, \xi, t)$	$E^q(x, \xi, t)$	<b>For Vector Meson</b>
$\tilde{H}^q(x, \xi, t)$	$\tilde{E}^q(x, \xi, t)$	<b>For Pseudo-Scalar Meson</b>

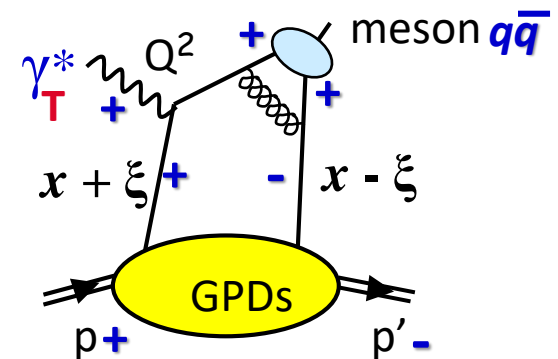
+ 4 chiral-odd or transversity GPDs: helicity of parton changed  
(not possible in DVCS)

$H_T^q(x, \xi, t)$	$E_T^q(x, \xi, t)$	$\bar{E}_T^q = 2 \tilde{H}_T^q + E_T^q$ (as Boer-Mulders)
$\tilde{H}_T^q(x, \xi, t)$	$\tilde{E}_T^q(x, \xi, t)$	

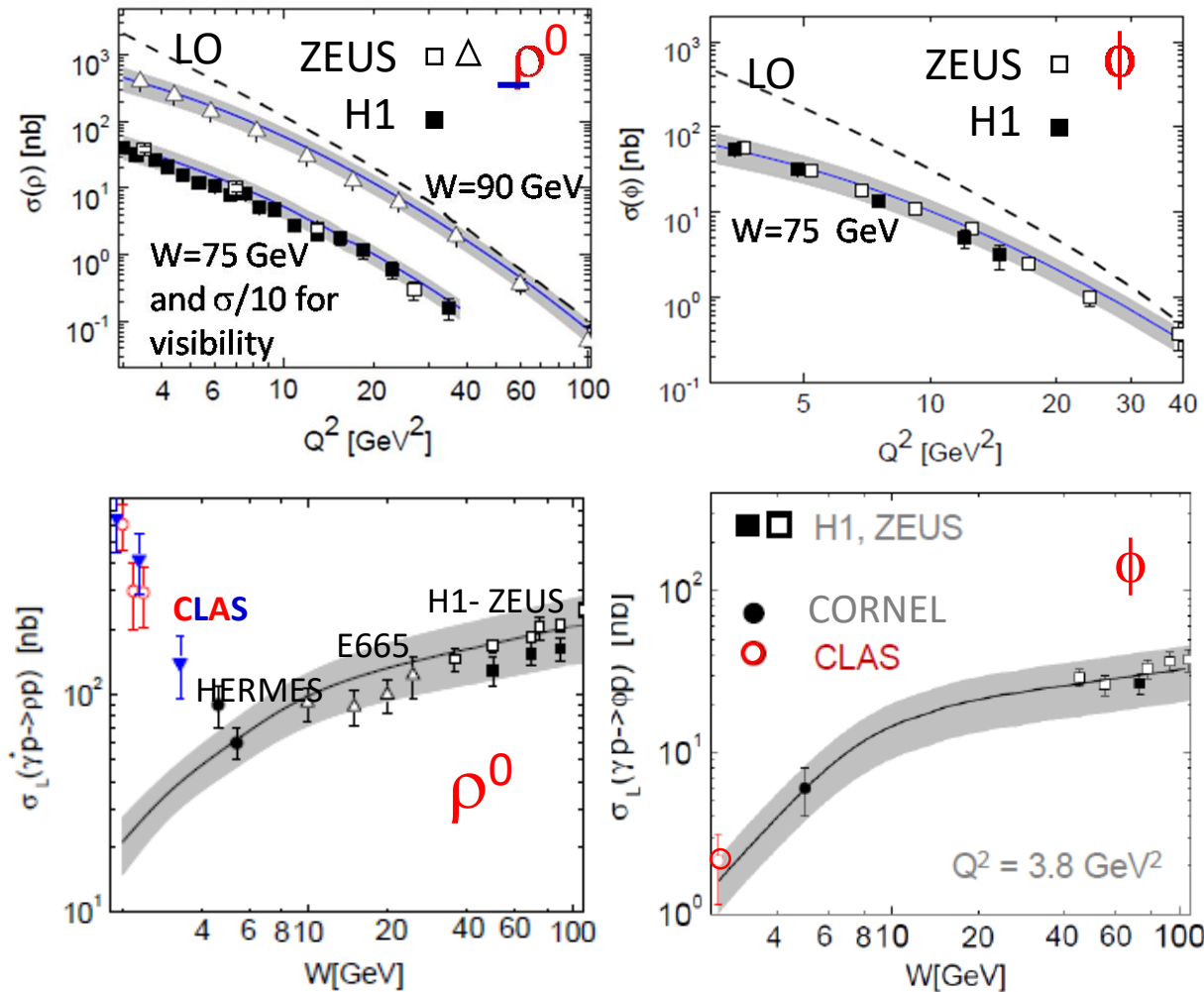
Factorisation proven only for  $\sigma_L$

$\sigma_T$  is asymptotically suppressed by  $1/Q^2$  but large contribution observed  
model of  $\sigma_T$  with transversity GPDs - divergencies regularized by  $k_T$  of  $q$   
and  $\bar{q}$  and Sudakov suppression factor

$\mathcal{M}_{0-, ++}$  sensitive to  $H_T^q$   
and to a twist-3 meson wave function

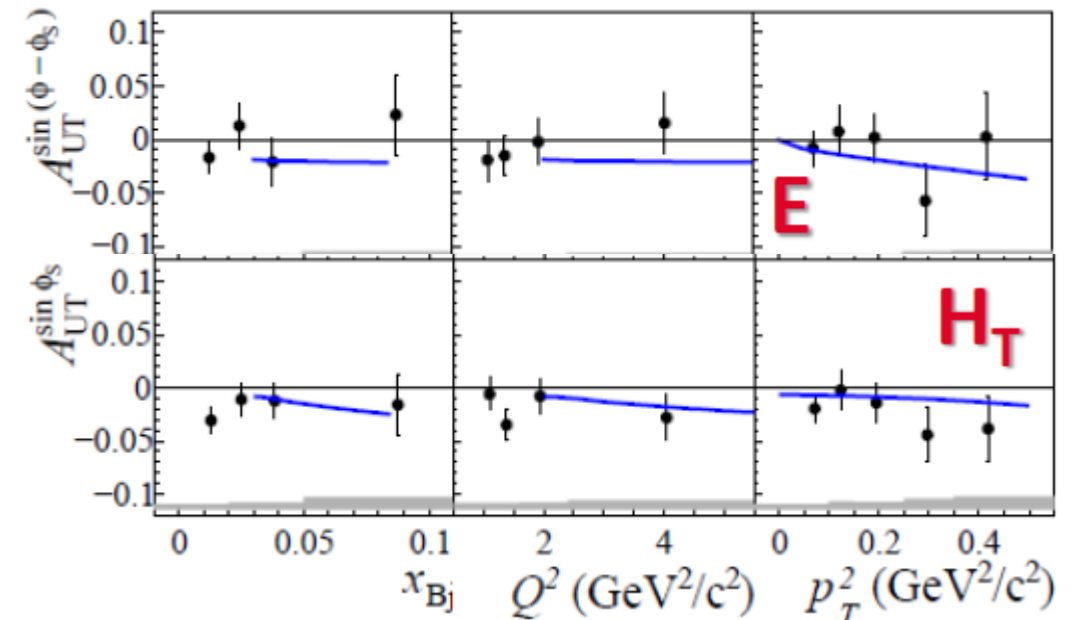


# GPDs and Hard Exclusive Vector Meson Production



$\rho^0$  ( $\rightarrow \pi^+\pi^-$ ) production at COMPASS with Transversely Polarized Target

COMPASS, NPB 865 (2012) 1-20, PLB731 (2014) 19

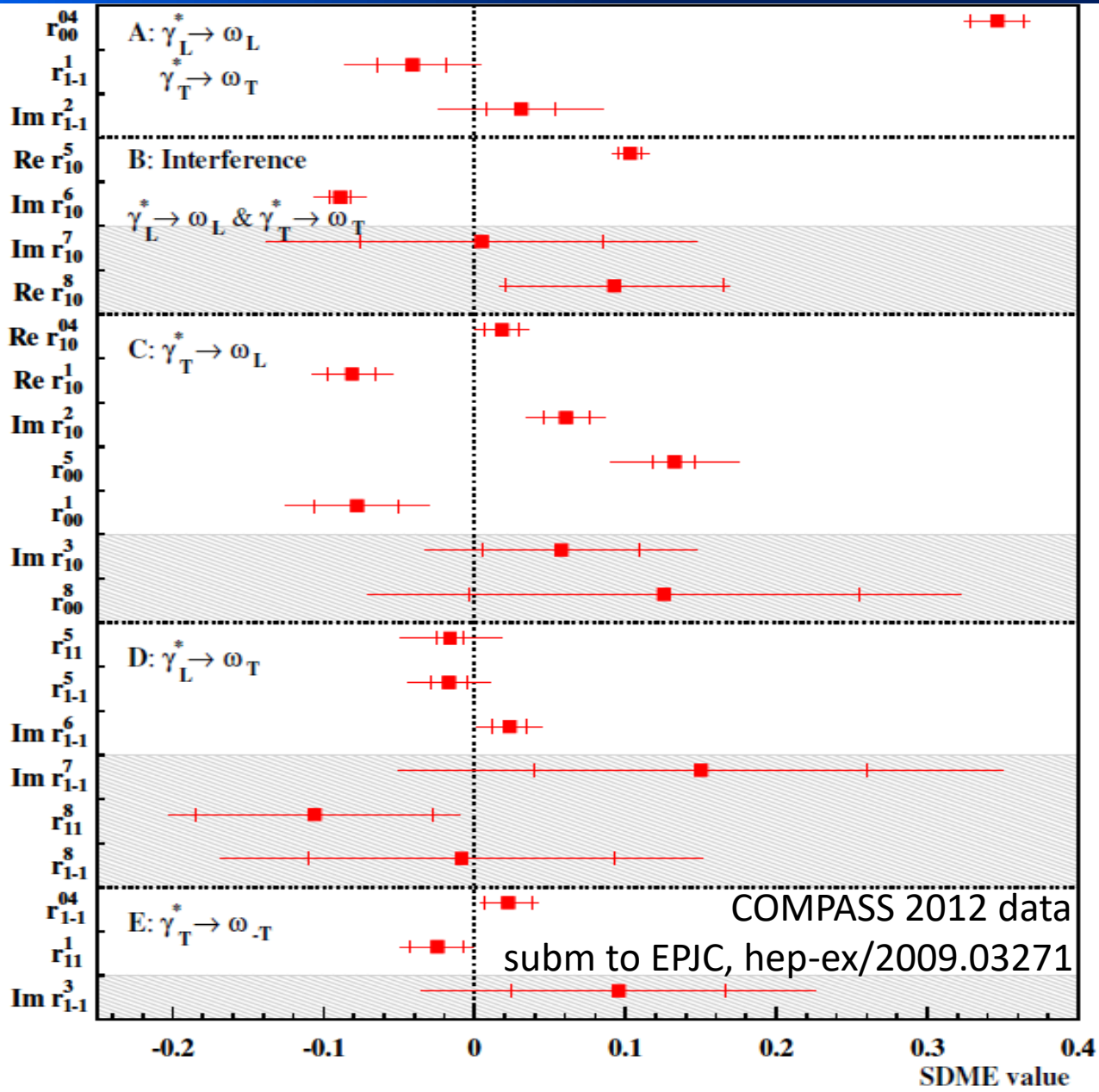


**GK** Goloskokov, Kroll, EPJC42,50,53,59,65,74 GPD model constrained by HEMP at small  $x_B$  (or large  $W$ )

dominant (longitudinal)  $\gamma_L^* p \rightarrow M p$  and transv. polar.  $\gamma_T^* p \rightarrow M p$

quark and gluon contributions (GPDs  $H$ ,  $E$ ,  $H_T$ ) and beyond leading twist

# GPDs and Hard Exclusive $\omega$ Vector Meson Production

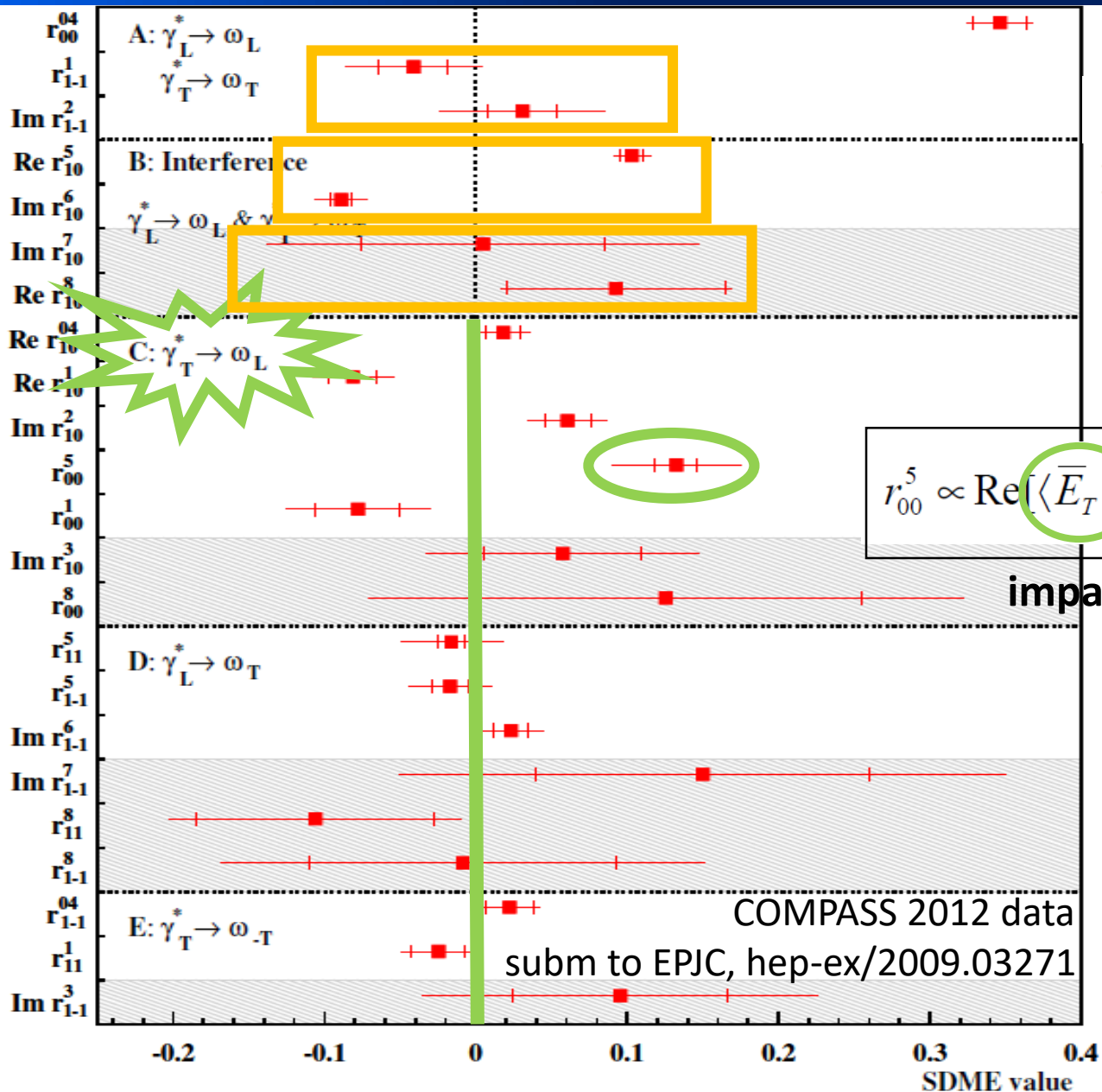


**COMPASS**

**23 SDMEs in 5 classes A, B, C, D, E  
 depending on helicity transitions**

SDMEs dependent on beam polarisation  
 shown within shaded areas

# GPDs and Hard Exclusive $\omega$ Vector Meson Production



If sCHC ( $\lambda_\gamma = \lambda_V$ )

$$r_{1-1}^1 + \text{Im}\{r_{1-1}^2\} = 0 \quad = -0.010 \pm 0.032 \pm 0.047$$

$$\text{Re}\{r_{10}^5\} + \text{Im}\{r_{10}^6\} = 0 \quad = 0.014 \pm 0.011 \pm 0.013$$

$$\text{Im}\{r_{10}^7\} - \text{Re}\{r_{10}^8\} = 0 \quad = -0.088 \pm 0.110 \pm 0.196$$

measurements:

All the other SDME in classes C, D, E should be 0  
not observed for class C

$$r_{00}^5 \propto \text{Re}[\langle \bar{E}_T \rangle_{IT}^* \langle H \rangle_{LL} + \frac{1}{2} \langle H_T \rangle_{LT}^* \langle E \rangle_{LL}]$$

impact of  $\bar{E}_T$  and  $H_T$

Already clearly observed  
in exclusive  $\rho^0, \omega$  production  
with transvers. polar. target  
COMPASS 2004-7-10 data

NPB865 [2012] 1-20, PLB731 (2014) 19  
NPB915 (2017) 454-475



# GPDs and Hard Exclusive $\pi^0$ Production

$e p \rightarrow e \pi^0 p$

$$\frac{d^2\sigma}{dt d\phi_\pi} = \frac{1}{2\pi} \left[ \left( \epsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} \right) + \epsilon \cos 2\phi_\pi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_\pi \frac{d\sigma_{LT}}{dt} \right]$$

$$\left| \langle \tilde{H} \rangle \right|^2 - \frac{t'}{4m^2} \left| \langle \tilde{E} \rangle \right|^2$$

$$\left| \langle H_T \rangle \right|^2 - \frac{t'}{8m^2} \left| \langle \bar{E}_T \rangle \right|^2$$

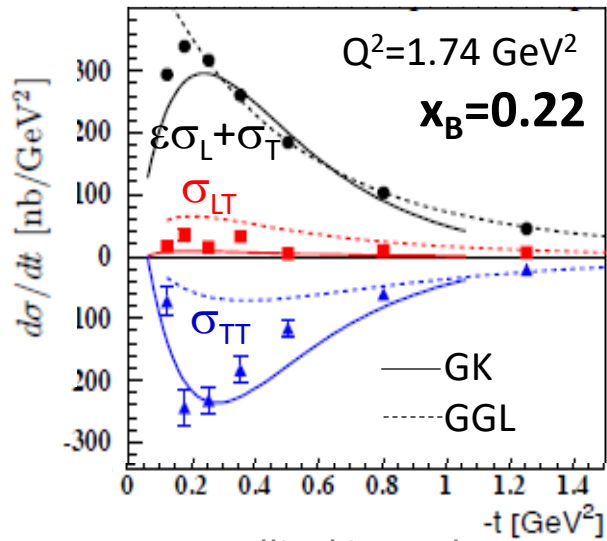
$$\frac{t'}{16m^2} \left| \langle \bar{E}_T \rangle \right|^2$$

$$\frac{\sqrt{-t'}}{2m} \text{Re} \left[ \langle H_T \rangle^* \langle \tilde{E} \rangle \right]$$

Jlab 6 GeV CLAS

$\sigma_{TT}$  large (impact of  $\bar{E}_T$ )

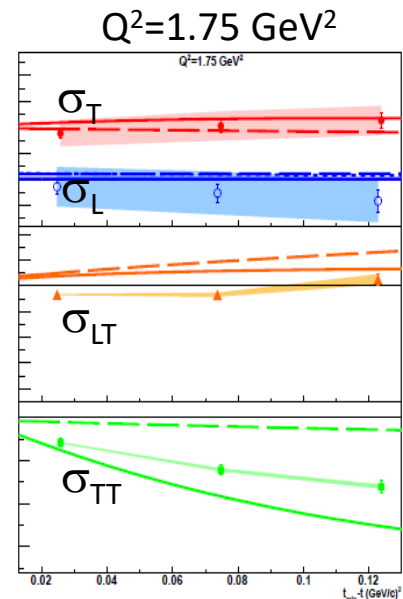
$\sigma_{LT}$  smaller



Bedlinskiy et al, PRL109 (2012), PRC90 (2014)

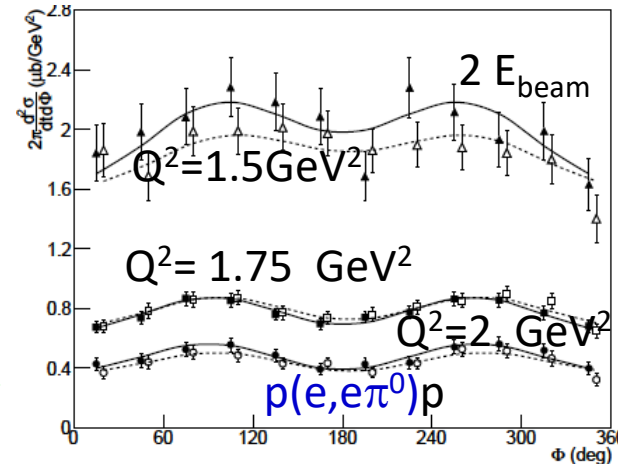
Jlab 6 GeV Hall-A LH2 target  $\rightarrow$  proton

Different beam energies  $\rightarrow$  L/T separation



Defurne et al, PRL117 (2016)

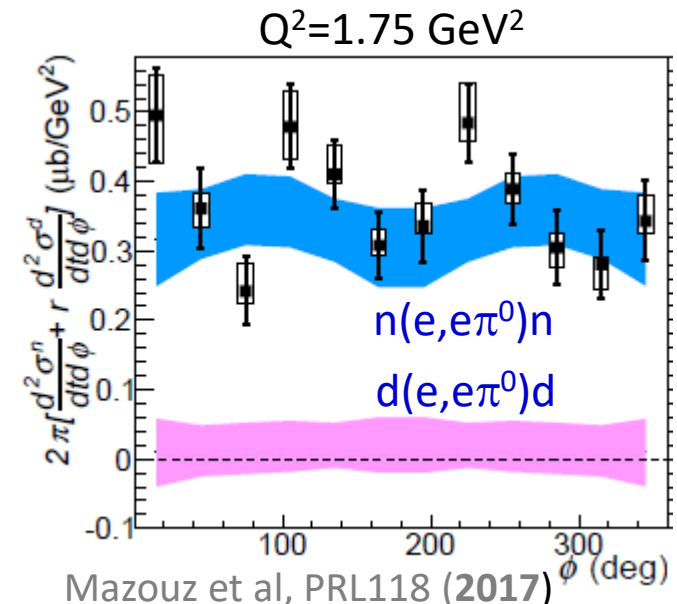
$x_B=0.36$   $t \sim 0.18$  GeV<sup>2</sup>



LD2 target  $\rightarrow$  neutron+deuteron

$D(e,e\pi^0)X - p(e,e\pi^0)p = n(e,e\pi^0)n + d(e,e\pi^0)d$

$\rightarrow$  Flavor decomposition on u and d quarks



Mazouz et al, PRL118 (2017)

# GPDs and Hard Exclusive $\pi^0$ Production

$e p \rightarrow e \pi^0 p$

$$\frac{d^2\sigma}{dt d\phi_\pi} = \frac{1}{2\pi} \left[ \left( \epsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} \right) + \epsilon \cos 2\phi_\pi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_\pi \frac{d\sigma_{LT}}{dt} \right]$$

$$\left| \langle \tilde{H} \rangle \right|^2 - \frac{t'}{4m^2} \left| \langle \tilde{E} \rangle \right|^2$$

$$\left| \langle H_T \rangle \right|^2 - \frac{t'}{8m^2} \left| \langle \bar{E}_T \rangle \right|^2$$

$$\frac{t'}{16m^2} \left| \langle \bar{E}_T \rangle \right|^2$$

$$\frac{\sqrt{-t'}}{2m} \text{Re} \left[ \langle H_T \rangle^* \langle \tilde{E} \rangle \right]$$

$$\left\langle \frac{d\sigma_T}{d|t|} + \epsilon \frac{d\sigma_L}{d|t|} \right\rangle = (8.2 \pm 0.9_{\text{stat}} \pm 1.2_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

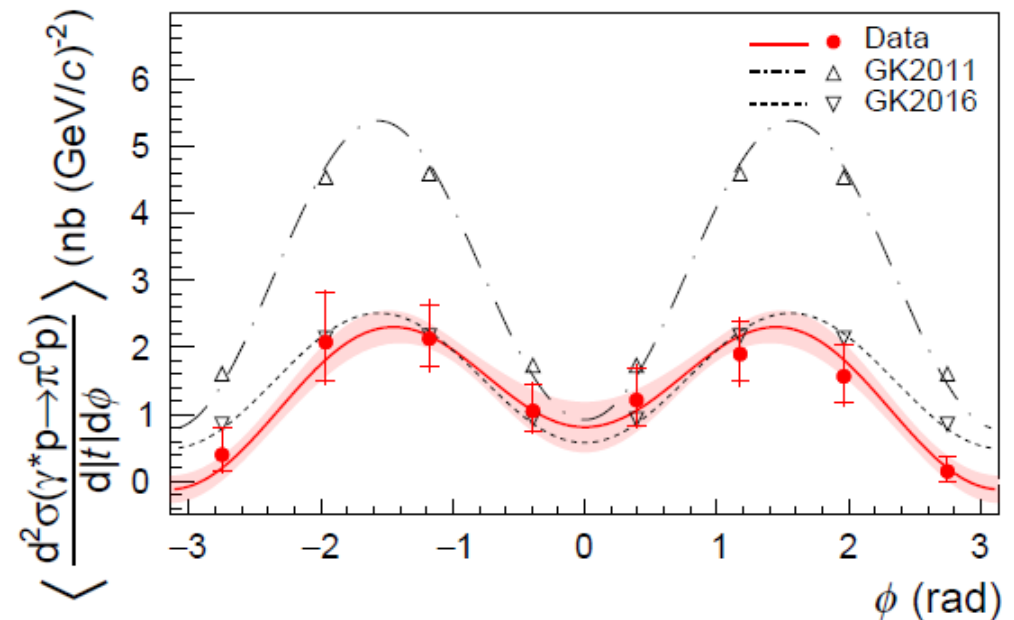
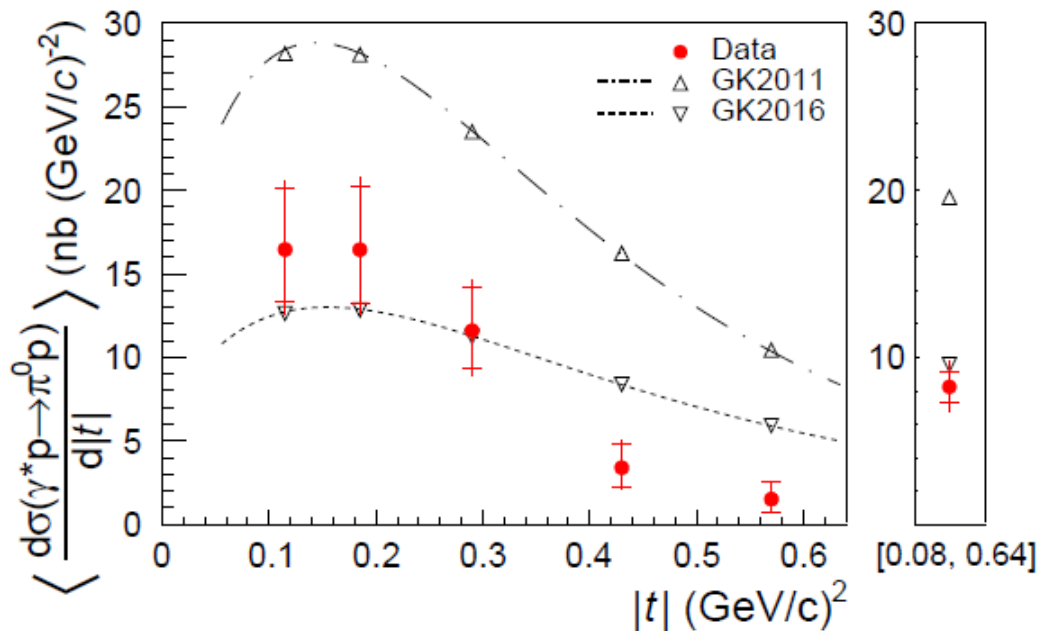
$$\left\langle \frac{d\sigma_{TT}}{d|t|} \right\rangle = (-6.1 \pm 1.3_{\text{stat}} \pm 0.7_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{LT}}{d|t|} \right\rangle = (1.5 \pm 0.5_{\text{stat}} \pm 0.3_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

COMPASS  
PLB 805 (2020)  
 $Q^2=2.0 \text{ GeV}^2$   
 $x_B=0.093$   
 $t \sim 0.26 \text{ GeV}^2$

$\sigma_{TT}$  large (impact of  $\bar{E}_T$ )

$\sigma_{LT}$  smaller but significantly positive as at CLAS

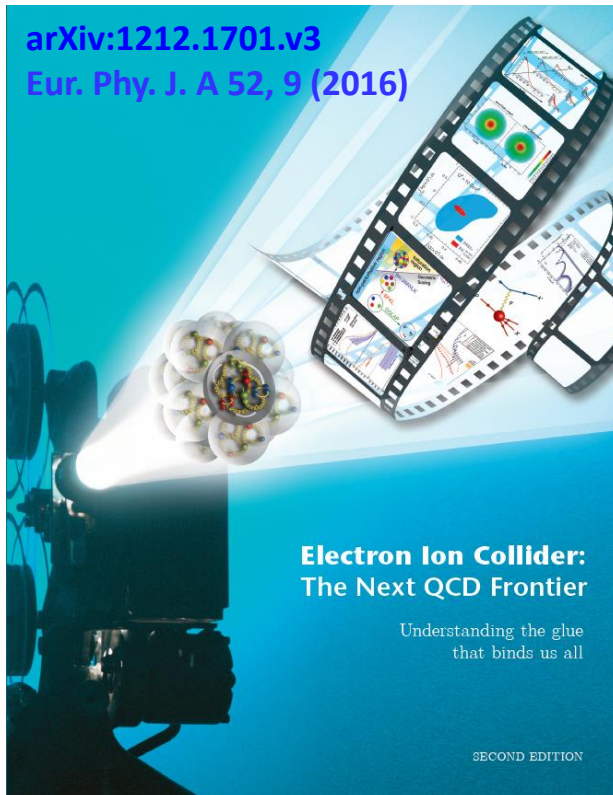


# Future: Key measurements for imaging partons with EIC

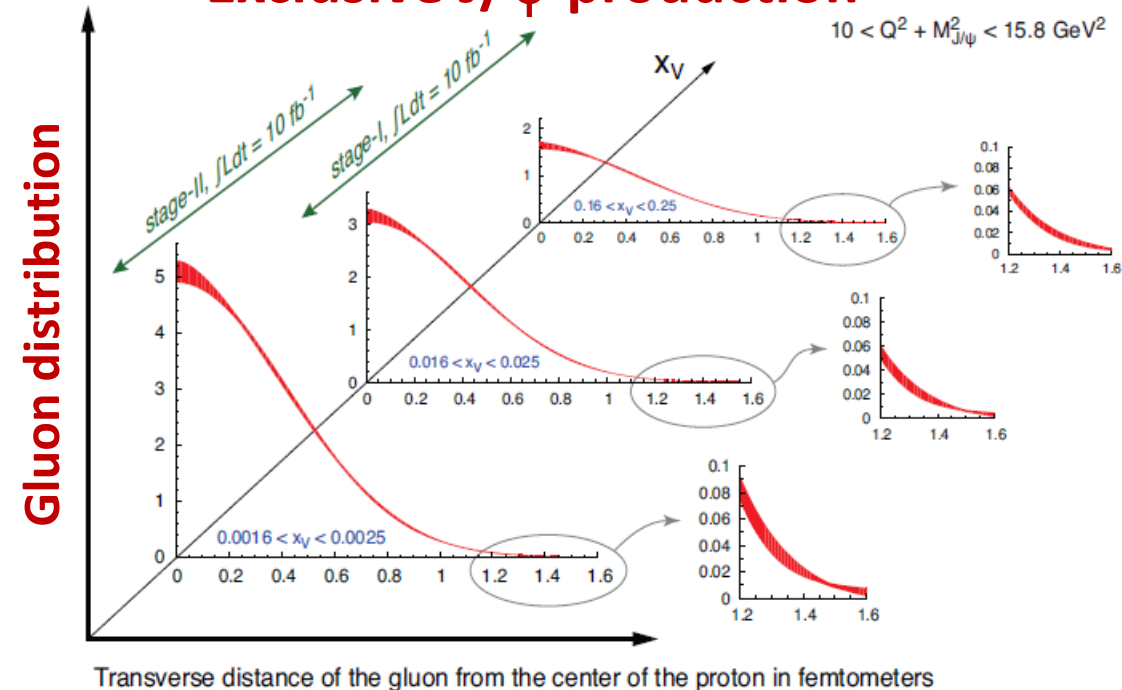
Stage 2  
Ee=20 GeV Ep=250 GeV

Stage 1  
Ee=5 GeV Ep=100 GeV

Deliverables	Observables	What we learn	Requirements
GPDs of sea quarks and gluons	DVCS and $J/\Psi, \rho^0, \phi$ production cross section and polarization asymmetries	transverse spatial distrib. of sea quarks and gluons; total angular momentum and spin-orbit correlations	$\int dt L \sim 10$ to $100 \text{ fb}^{-1}$ ; Roman Pots; polarized $e^-$ and $p$ beams; wide range of $x_B$ and $Q^2$ ;
GPDs of valence and sea quarks	electroproduction of $\pi^+, K$ and $\rho^+, K^*$	dependence on quark flavor and polarization	range of beam energies; $e^+$ beam valuable for DVCS



## Exclusive $J/\psi$ production



# Conclusions

Not an exhaustive compilation of all results and projections, also JPARC in the game...

Jlab 12 GeV with the high luminosity electron beam is at the beginning of a very exciting time with a high precision era for valence quarks at large  $x_B$

COMPASS, with high energy muon beams at CERN and RHIC with UPC will provide first results of sea quarks and gluons at small  $x_B$

They are the foundations for the preparation of new experiments at EIC

*For example preparation of the EIC Yellow Report  
for the detector requirement to study DVCS and  $\pi^0$  using:*

✓ the PARTONS framework

with KM20 CFF tables provided by Kumericki

and GK16 model from Goloskokov and Kroll based on the COMPASS results

✓ an update of the MC event generator MILOU developed for H1 and ZEUS