

Freed-isobar Partial-Wave Analysis of the Spin-Exotic $J^{PC} = 1^{-+}$ Wave at COMPASS

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Plan of the talk

- Conventional PWA for 3π final states
- PWA with model-independent (free) parametrization of $\pi^-\pi^+$ isobars
 - Formulation of the technics
 - Continuous ambiguities in 3-body amplitudes - Zero Modes
 - Approaches for resolving the ambiguities and extracting physics information from the experimental data
- Results for free-isobarred analysis of $J^{PC} = 1^{-+}$ amplitude in COMPASS $\pi^-\pi^+\pi^-$ system

This Meeting: $\pi^-\pi^+\pi^-$ COMPASS results

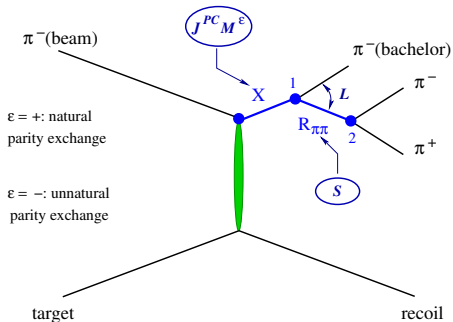
HK1 Boris Grube, SYCC Bernhard Ketzer, HK30 Fabian Krinner

"Resonance production and $\pi\pi$ S-wave in $\pi^-p \rightarrow \pi^-\pi^-\pi^+p_{recoil}$ at 190 GeV/c"
Phys. Rev. D95(2017) no.3, 032004, COMPASS collaboration

arXiv:1710.09849 [hep-ph]

"Ambiguities in model-independent partial-wave analysis" (F. Krinner et al.)

The reaction



- Reggeon exchange, naturality $\eta = P_R(-1)^J$
- Gottfried-Jackson frame: SCM of X: $Z_{GJ} \parallel \vec{p}_{beam}^*$, $Y_{GJ} = [\vec{p}_{recoil}^* \times \vec{p}_{beam}^*]$
- Reflectivity basis for system of mesons:
 $|JM\epsilon\rangle = |JM\rangle - \epsilon P(-1)^{J-M} |J-M\rangle$
- At high beam energies: reflectivity ϵ equal to naturality η
- unpolarised target: $\epsilon = \pm 1$ states do not interfere

PWA with fixed shapes of isobars vs. freed isobares

PWA with established isobars:

Partial waves are labelled as $J^{PC} M^{\epsilon} \xi \pi L$

The mass-independent PWA events density:

$$\mathcal{I}(\tau) = \sum_{\epsilon} \sum_r |\sum_i T_{ir}^{\epsilon} \psi_i^{\epsilon}(\tau)|^2$$

The decay amplitude $\psi_i^{\epsilon}(\tau)$ contains angular part and $\pi^{-}\pi^{+}$ isobar Breit-Wigner function and is bose-symmetrized by (1) \leftrightarrow (3) of $\pi_{(1)}^{-}\pi_{(2)}^{+}\pi_{(3)}^{-}$ system:

$$\psi_i^{\epsilon}(\tau) = A(\Omega_{12}, \Omega_1^*) BW(m_{12}) + A(\Omega_{32}, \Omega_3^*) BW(m_{32})$$

PWA with free isobars:

The fixed amplitude of $\pi^{-}\pi^{+}$ isobar is replaced by sum of step-like functions with complex coefficients: $BW(m)^a \rightarrow \sum_{\beta} \omega_{\beta}^a \Pi_{\beta}(m)$

The full free-isobarred amplitude for $J^{PC} M^{\epsilon}$ sector:

$$\sum_a \sum_{\beta} \omega_{\beta}^a (A_{12}^a \Pi_{\beta}(m_{12}) + A_{32}^a \Pi_{\beta}(m_{32}))$$

Switch off bose-symmetrization \rightarrow components of freed amplitude will be orthogonal for each different indices a and β

Bose-symmetrization breaks orthogonality in 2-fold way:

- Step-like functions with different β - intersect on the Dalitz plot
- Different angular functions are cross non-orthogonal

Analysis with freed isobars - Zero Modes

We have found sets of linear dependences between "free"-components inside same $J^{PC} M^\epsilon$, they are always real-valued functions z_β^a so that $\omega_\beta^a \rightarrow \omega_\beta^a + C_z z_\beta^a$ will not change the whole amplitude.

The current PWA includes 12 different $J^{PC} M^\epsilon \xi \pi L$ free-isobarred amplitudes:

$J^{PC} M^\epsilon$	ξ	L	ZM		
$0^{-+}0^+$	$(\pi^+\pi^-)_S \pi S$;	$(\pi^+\pi^-)_P \pi P$	+		
$1^{++}0^+$	$(\pi^+\pi^-)_S \pi P$;	$(\pi^+\pi^-)_P \pi S$	+		
$1^{++}1^+$		$(\pi^+\pi^-)_P \pi S$			
$1^{-+}1^+$		$(\pi^+\pi^-)_P \pi P$	+		
$2^{-+}0^+$	$(\pi^+\pi^-)_S \pi D$;	$(\pi^+\pi^-)_P \pi P$;	$(\pi^+\pi^-)_P \pi F$;	$(\pi^+\pi^-)_D \pi S$	+
$2^{-+}1^+$		$(\pi^+\pi^-)_P \pi P$			
$2^{++}1^+$		$(\pi^+\pi^-)_P \pi D$			

+72 waves with fixed isobares

Strategy for resolving the ambiguity:

- find the PWA solution for 2π amplitude $\hat{\omega}_\beta^a$ (which is ambiguous, so has huge correlations)
- perform model-dependent fit to 2π amplitude:

$$C_a B W_\beta^a(M_0, \Gamma_0) + \sum_z C_z z_\beta^a \text{ to } \hat{\omega}_\beta^a$$

- Subtract zero mode(s) $\sum_z C_z z_\beta^a$ from $\hat{\omega}_\beta^a$

Explanation of Zero Mode in 1^{-+} amplitude

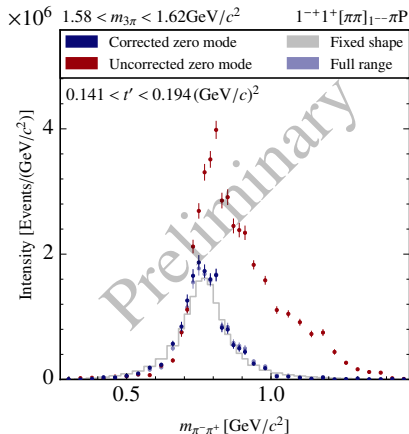
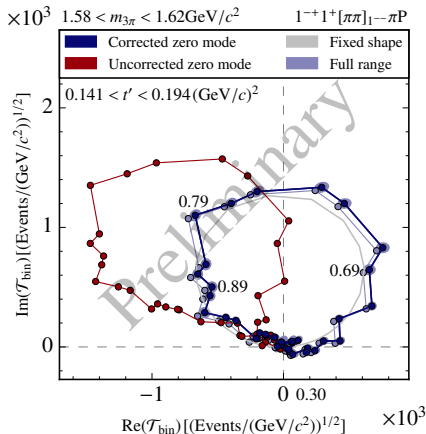
Non-relativistic (Zemach) tensor formalism used to derive the decay amplitude:

$$\psi_{1^{-+}}(\tau)_k = [\vec{p}_1 \times \vec{p}_2]_k (BW(m_{12}) + z(m_{12})) + [\vec{p}_3 \times \vec{p}_2]_k (BW(m_{32}) + z(m_{32}))$$

$$\text{as } [\vec{p}_3 \times \vec{p}_2] = -[\vec{p}_1 \times \vec{p}_2]$$

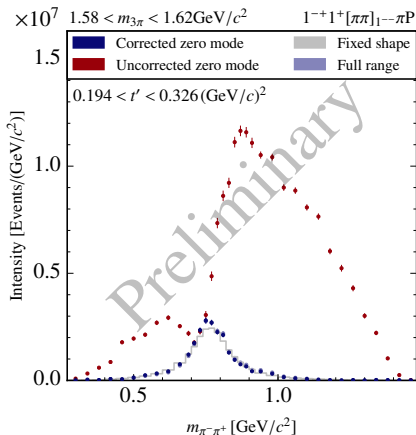
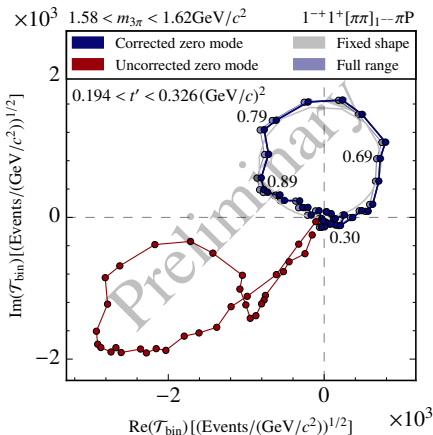
vanishing reads: $z(m_{12}) - z(m_{32}) = 0 \rightarrow z(m) = \text{const.}$

Results for COMPASS $\pi^- \pi^+ \pi^-$ data



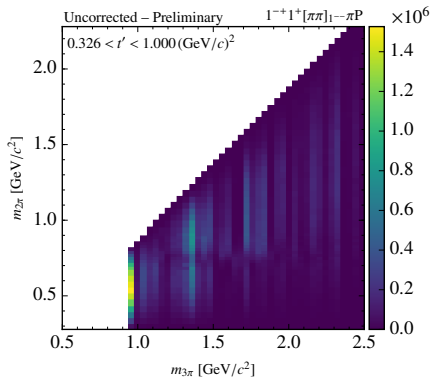
Argand diagram for $(\pi\pi)_P$ -amplitude. Intensity of $(\pi\pi)_P$ -amplitude.
 Red - ambiguity unresolved. Blue - ambiguity resolved.

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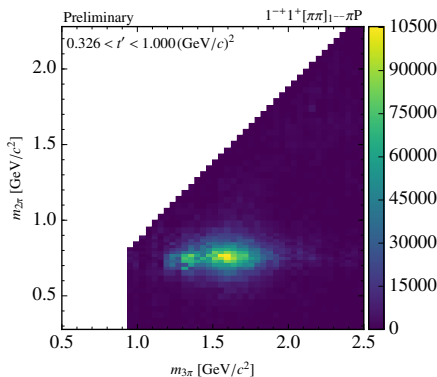
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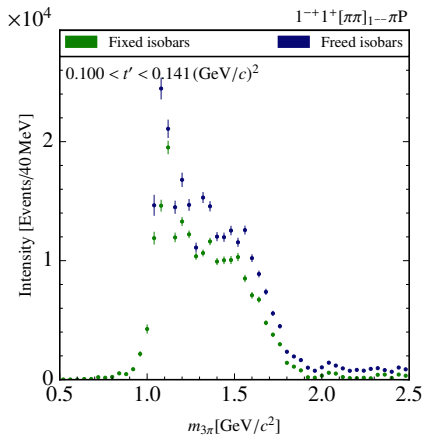
Unresolved ambiguity

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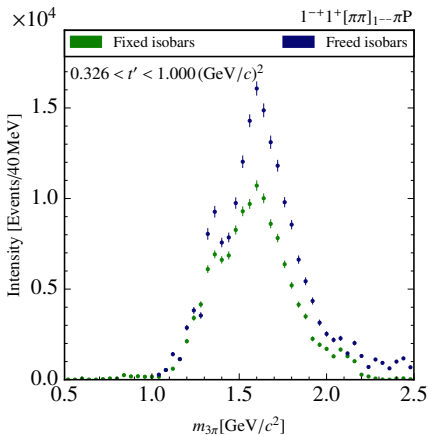


Resolved ambiguity

Results for COMPASS $\pi^-\pi^+\pi^-$ data



The lowest t' interval



The highest t' -interval - peak of spin-exotic $\pi_1(1600)$ (?)

CONCLUSIONS:

- Analysis with freed isobars was first time performed for $J^{PC} M^{\epsilon} = 1^{-+}1^{+}(\pi\pi)_{P}\pi P$
- The continuous ambiguities were resolved by applying Breit-Wigner model for $(\pi\pi)_{P}$ freed amplitude
- The data shows clean signal of $\rho(770)$ as $(\pi\pi)_{P}$ isobar in 3π amplitude with $1^{-+}1^{+}$ quantum numbers for various $m(3\pi)$ and t' bins

OUTLOOK:

- Extract physical quantities: intensities and coupling phases of the model components (so called fixed-shapes) and compare with PWA results with fixed isobars
- Study other $J^{PC} M^{\epsilon}$ freed amplitudes