

Final COMPASS results on the spin-dependent structure functions g_1^p and g_1^d in the deep-inelastic and nonperturbative region



Barbara Badelek
University of Warsaw

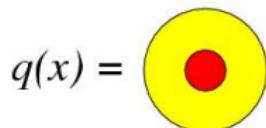
On behalf of COMPASS Collaboration

DIS2017

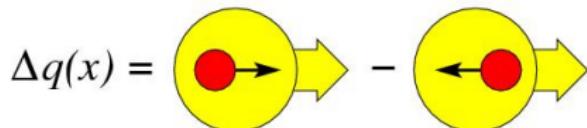
Birmingham, April 3 – 7, 2017

Partonic structure of the nucleon; distribution functions

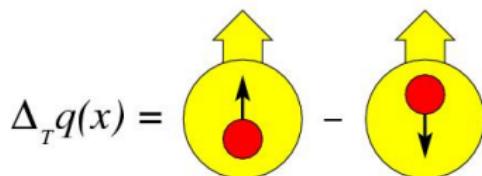
Three twist-two quark distributions in QCD and after integrating over the quark intrinsic k_t



Quark momentum DF;
well known (unpolarised DIS $\rightarrow \mathbf{F}_{1,2}(\mathbf{x}, \mathbf{Q}^2)$).



Difference in DF of quarks with spin parallel or antiparallel to the nucleon's spin in a longitudinally polarised nucleon;
less well known (polarised DIS $\rightarrow \mathbf{g}_1(\mathbf{x}, \mathbf{Q}^2)$).



Difference in DF of quarks with spin parallel or antiparallel to the nucleon's spin in a transversely polarised nucleon;
poorly known (polarised DIS $\rightarrow \mathbf{h}_1(\mathbf{x}, \mathbf{Q}^2)$).

Nonrelativistically: $\Delta_T q(x, Q^2) \equiv \Delta q(x, Q^2)$. OBS.! $\Delta_T q(x, Q^2)$ are C-odd and chiral-odd

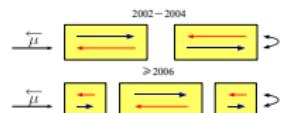
Our goal: measurement of $g_1(\mathbf{x}, \mathbf{Q}^2)$ and determination of $\Delta q(\mathbf{x}, \mathbf{Q}^2)$

Method of extraction of g_1

- Inclusive asymmetry, $A_{meas}(x, Q^2)$, γ^* -N asymmetry, $A_1(x, Q^2)$, and $g_1(x, Q^2)$:

$$A_{meas} = \frac{1}{f P_T P_B} \left(\frac{N^{\leftarrow} - N^{\rightarrow}}{N^{\leftarrow} + N^{\rightarrow}} \right) \approx D A_1 = D \frac{g_1(x, Q^2)}{F_1(x, Q^2)} \stackrel{\text{LO}}{=} D \frac{\sum_q e_q^2 \Delta q(x, Q^2)}{\sum_q e_q^2 q(x, Q^2)}$$

f, D : dilution and depolarisation factors; P_T, P_B : target and beam polarisations;
 $N^{\leftarrow, \rightarrow}$: number of $\vec{\mu}$ interactions in each target cell:
(upstream, downstream) or (outer, central)



- Then $g_1(x, Q^2)$:

$$g_1(x, Q^2) = A_1(x, Q^2) \cdot F_1(x, Q^2) = A_1(x, Q^2) \cdot \frac{F_2(x, Q^2)}{2x(1 + R(x, Q^2))}$$

- For the deuteron target:

$$(\text{per nucleon}) g_1^d = g_1^N \left(1 - \frac{3}{2}\omega_D\right) = \frac{g_1^p + g_1^n}{2} \left(1 - \frac{3}{2}\omega_D\right); \quad \omega_D = 0.05 \pm 0.01$$

Method of extraction of g_1, \dots cont'd

- Polarisation of all two/three target cells measured simultaneously...
- ... and reversed regularly by field reversal or microwave repolarisation
- Number of events in each target cell: $N_i = a_i \phi_i n_i \bar{\sigma} (1 + P_B P_T f D A_1)$
(a_i - acceptance, ϕ_i - incident flux, n_i - number of target nuclei,
 $\bar{\sigma}$ - spin independent cross section)
- Each event weighted by $w = f D P_B$
(no P_T as it changes in time \Rightarrow false asymmetries)
- Calculate A_1 from combinations of Σw ; fluxes and acceptances cancel,
if ratio of acceptances for two sets of cells does not change during field rotation
- f corrected for spin-independent radiative processes: TERAD
- A_1 corrected for spin-dependent radiative processes: POLRAD
- A_1 corrected for polarisable ^{14}N and ^7Li in respective targets

Method of extraction of g_1, \dots cont'd

- Final $g_1^{\text{p,d}}(x, Q^2)$ is extracted as:

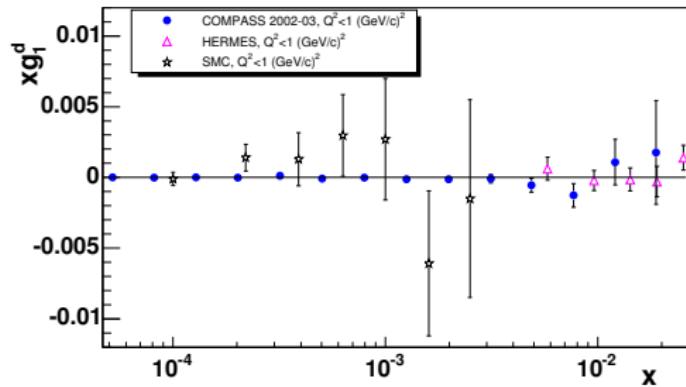
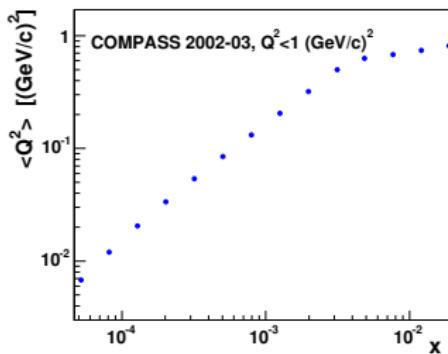
$$g_1^{\text{p,d}}(x, Q^2) = A_1^{\text{p,d}}(x, Q^2) \cdot \frac{F_2^{\text{p,d}}(x, Q^2)}{2x(1 + R(x, Q^2))}$$

- Input $F_2^{\text{p,d}}(x, Q^2)$:** from NMC measurements at $Q^2 > 0.2 \text{ (GeV}/c)^2$ PRD 58, 112001;
at $Q^2 < 0.2 \text{ (GeV}/c)^2$ – from the model based on the GVMD concept PLB 295 (1992) 263
- Input $R(x, Q^2)$:** R1998, PLB 452 (1999) 194,
suitably extended to low Q^2 , PLB 647 (2007) 330
- Maximal systematic uncertainties on A_1 :**
 - at low Q^2 : from $D(R)$, $\sim (0.01 - 0.39)D$ (mult.)
and false asymm., $\sim 1.5 \sigma_{\text{stat.}}$ (additive)
 - at high Q^2 : from P_B, P_T , $\sim 0.05P_B, P_T$ (mult.)
and false asymm., $< 0.8 \sigma_{\text{stat.}}$ (additive)

g_1^d at low Q^2

g_1^d in the nonperturbative ($Q^2 < 1 \text{ (GeV/c)}^2$ region)

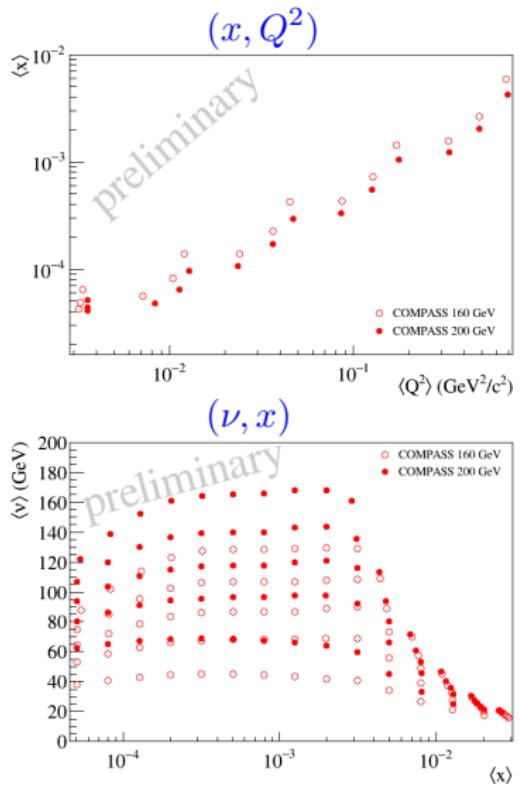
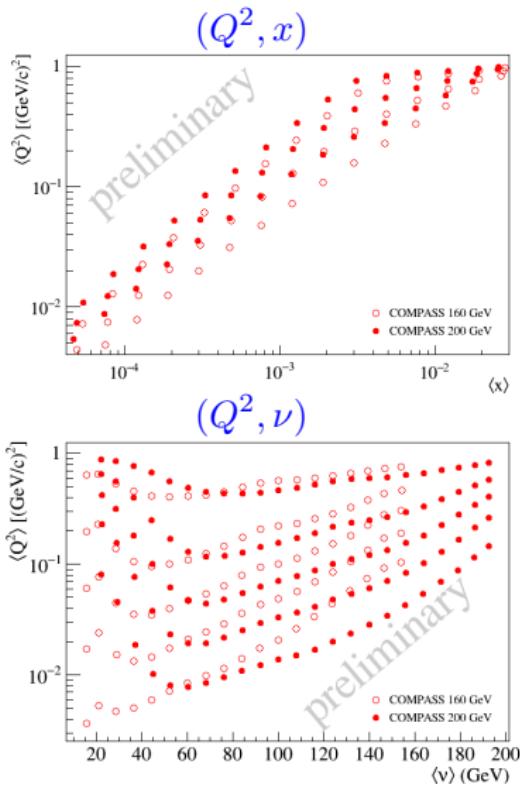
V.Yu. Alexakhin (COMPASS) *et al.* Phys. Lett. B **647** (2007) 330



- More than ten-fold improvement over the statistical precision of SMC.
- Spin effects in g_1^d at low x and Q^2 absent?

g_1^p at low Q^2

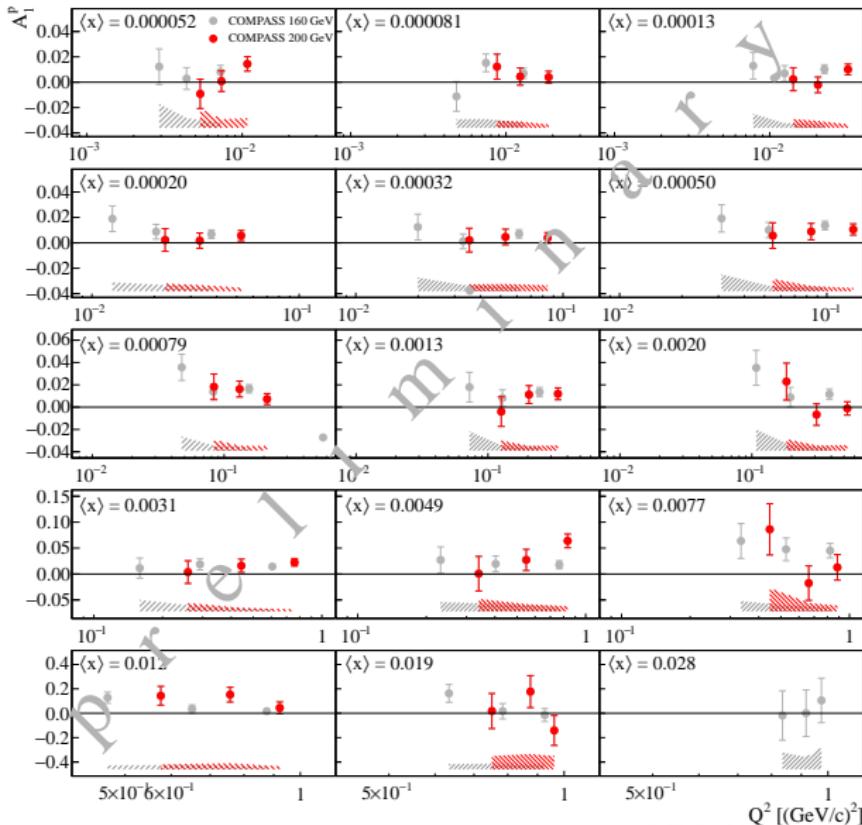
Phase space coverage of the low Q^2 , low x measurement of g_1^p



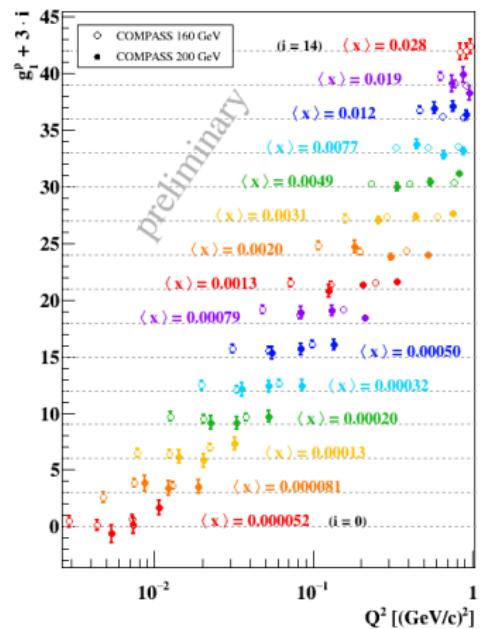
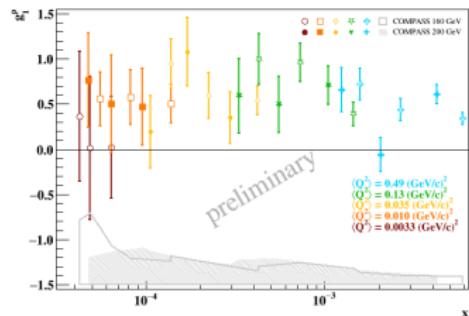
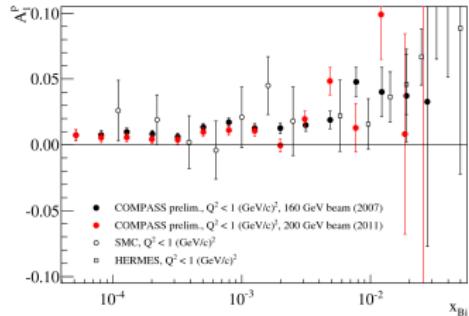
g_1^p at low Q^2 : data sample

- Reconstructed event: (incident μ) \wedge (outgoing μ) \wedge (at least 1 outtrack)
This “hadron method” does not bias inclusive asymmetry @ low x (SMC) but improves resolution of interaction point and decreases radiative bckg.
- Kinematic cuts
 - $Q^2 < 1 \text{ (GeV}/c)^2$
 - $x > 4 \times 10^{-5}$ (at smaller x uncertainties on x large)
 - $0.1 < y < 0.9$ (bad reconstruction @ low y , large rad. corr. @ large y)
 - Above cuts result in: $W \gtrsim 5 \text{ GeV}/c^2$
- Removed $\mu e \rightarrow \mu e$ events, around $x_{\mu e} = m_e/M_p = 5.45 \times 10^{-4}$
- Final sample: 447 million events @ 160 GeV + 229 million events @ 200 GeV
 $\implies \sim 150 \times$ SMC statistics

Low Q^2 : results on $A_1^p(x, Q^2)$



Low Q^2 : results on $A_1^P(x)$ and $g_1^P(x, Q^2)$



- More than ten-fold improvement over the statistical precision of SMC.
- Very clear spin effects in g_1^P at low x and Q^2

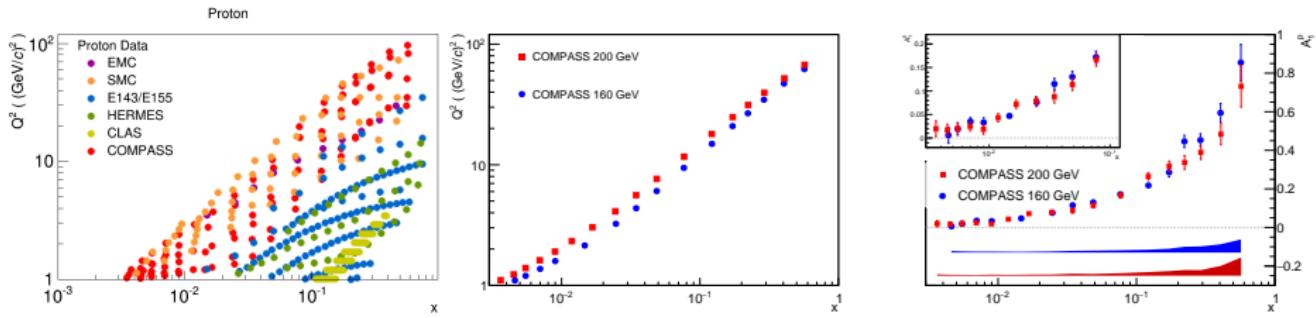
Importance of $g_1(x, Q^2)$ at low x, Q^2

- At low x : strong increase of parton density with decreasing x
 \implies parton recombination effects
- Low x effects particularly visible in g_1^{NS} $\implies \ln^2(1/x)$ terms
- BUT: in fixed-target experiments low x correlated with low Q^2 region
 \implies nonperturbative effects must be considered
 and parton mechanisms have to be suitably extended to low Q^2 .
- Attempts to describe that region phenomenologically:
 - B. Badelek, J. Kwiecinski, J. Kiryluk, Phys. Rev. D61 (2000) 014009.
 - B. Badelek, J. Kwiecinski, B. Ziaja, Eur. Phys. J. C26 (2002) 45.
 - B.I. Ermolaev, M. Greco, S.I. Troyan, Eur. Phys. J. C50 (2007) 823;
ibid. C51 (2007) 859.
 - W. Zhu, J. Ruan, Int. J. Mod. Phys. E24 (2015) 1550077.

g_1^p at $Q^2 > 1 \text{ (GeV}/c)^2$

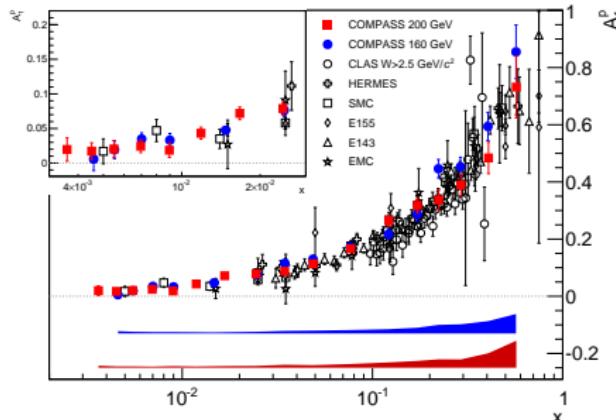
g_1^p : data samples

- Measurements at incident $E = 160 \text{ GeV}$ (from 2007): PLB 690 (2010) 466
- Measurements at incident $E = 200 \text{ GeV}$ (from 2011): PLB 753 (2016) 18
energy increased to reach lower x and higher Q^2
(and to balance the amount of data from the deuteron target)
- Final sample: 85 million events @ 160 GeV + 77 million events @ 200 GeV
- Results on A_1^p at both energies agree very well

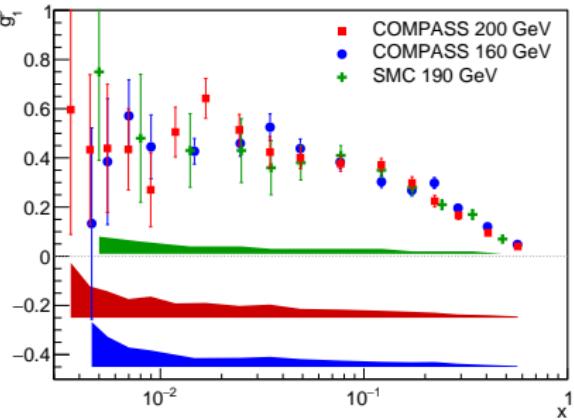


Results on $A_1^P(x)$ and $g_1^P(x)$

- $A_1^P(x)$ and $g_1^P(x)$ shown at the measured values of Q^2
- Bands of systematic uncertainties for each energy separately

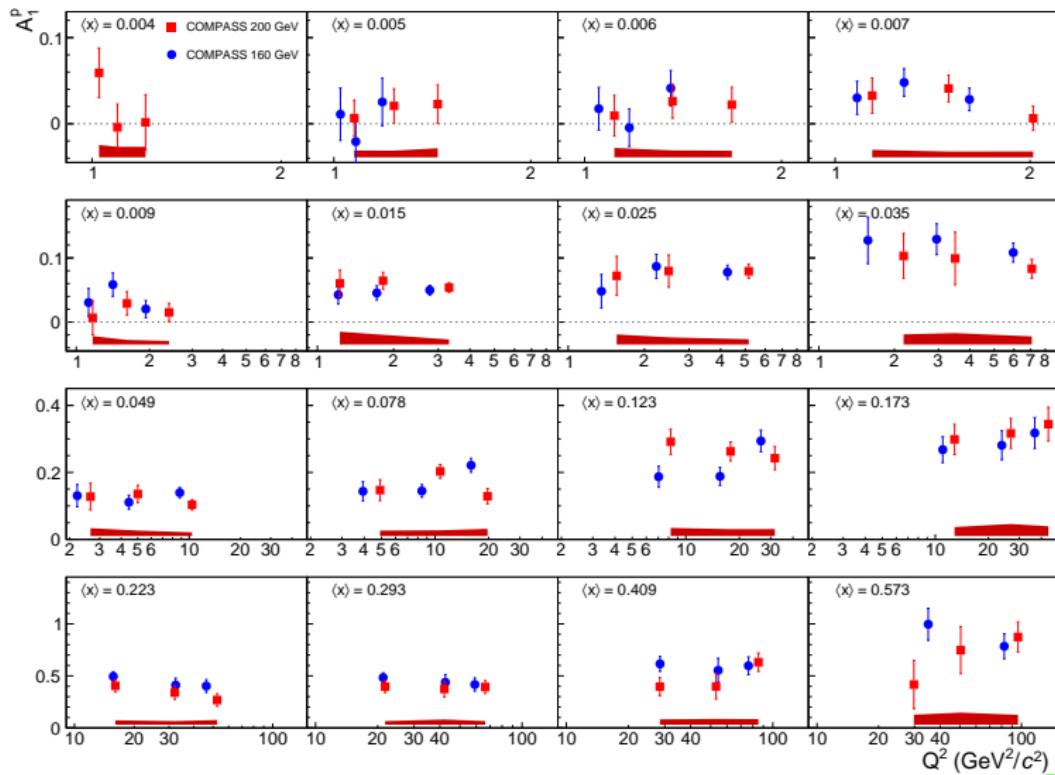


PLB 753 (2016) 18



- Good agreement of $A_1^P(x)$ and $g_1^P(x)$ with world data
- $g_1^P(x)$ clearly positive at lowest measured values of x

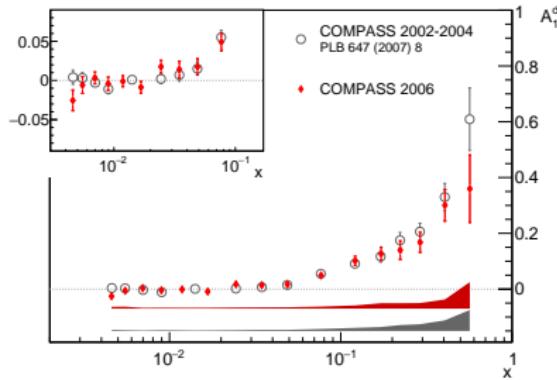
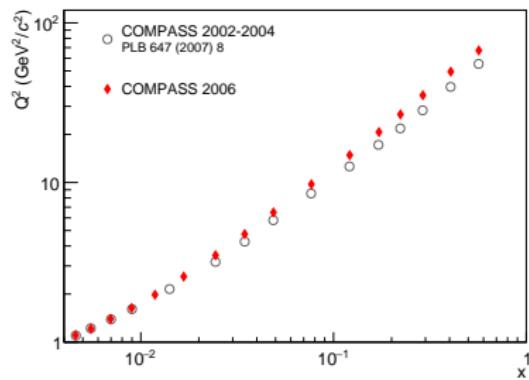
Results on $A_1^p(x, Q^2)$ at 160 and 200 GeV



g_1^d at $Q^2 > 1 \text{ (GeV}/c)^2$

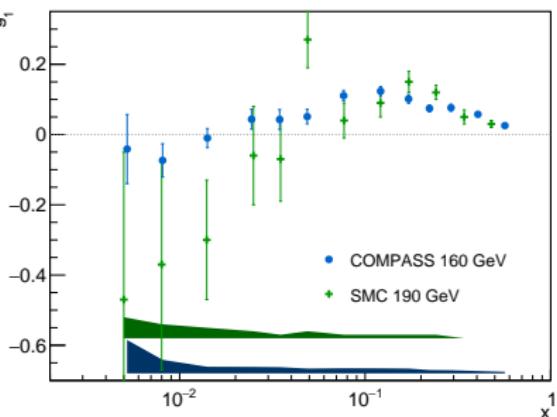
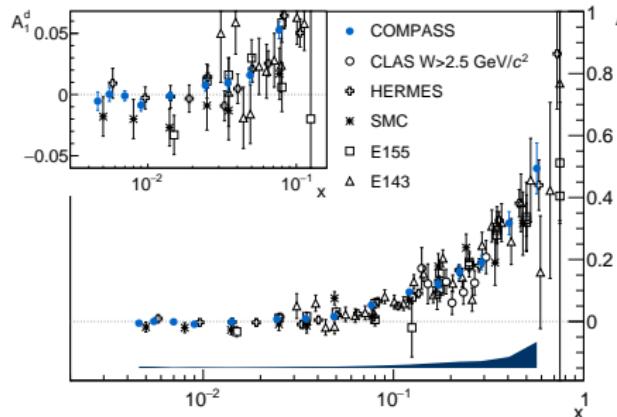
g_1^d : data samples

- Data at $E = 160$ GeV (from 2002–2004) published PLB 647 (2007) 8
- Data at $E = 160$ GeV (from 2006)
combined results in hep-ex/1612.00620, accepted PLB
- Final sample: 135 million events
- Results on A_1^d in both samples agree very well



Results on $A_1^d(x)$ and $g_1^d(x)$

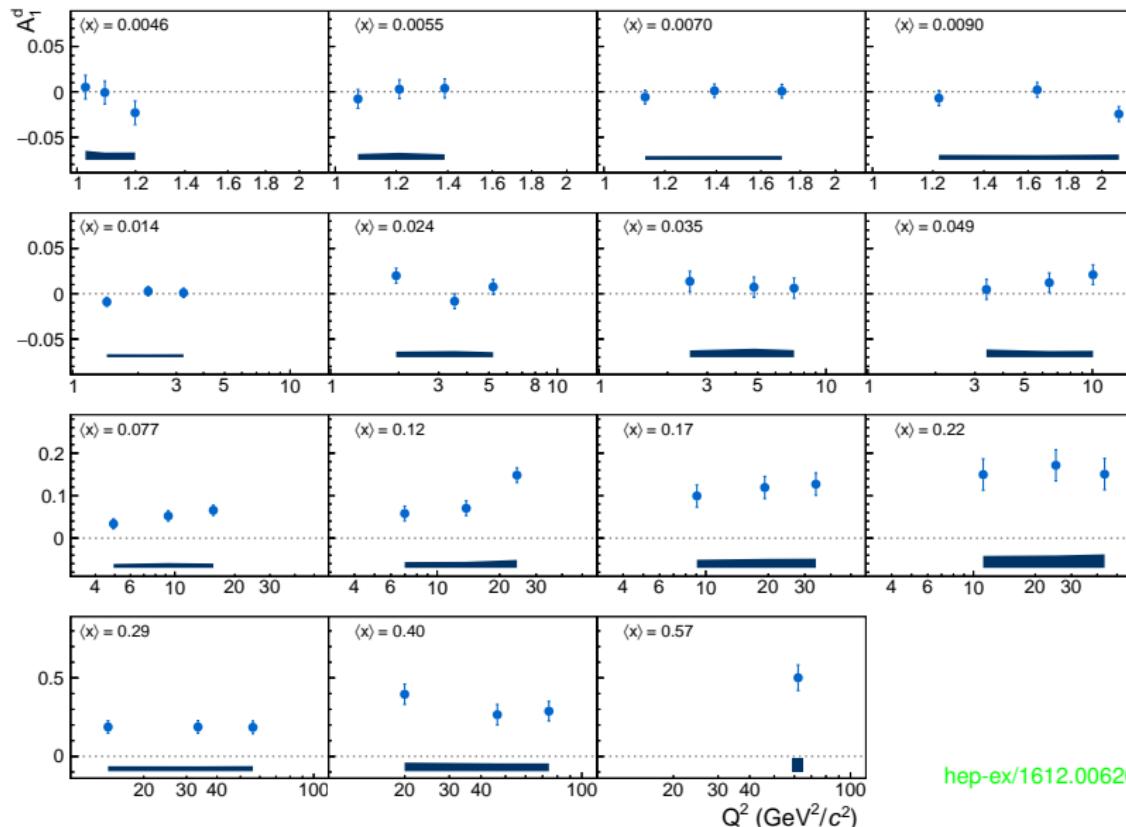
- Combined results from 2002–2004 and 2006
- $A_1^d(x)$ and $g_1^d(x)$ shown at the measured values of Q^2



hep-ex/1612.00620, accepted PLB

- Good agreement of $A_1^d(x)$ and $g_1^d(x)$ with world data
- $g_1^d(x)$ compatible with zero at lowest measured values of x , contrary to hints from SMC

Results on $A_1^d(x, Q^2)$



hep-ex/1612.00620, accepted PLB

NLO QCD fit to p,d, ${}^3\text{He}$ world data

PLB 753 (2016) 18

NLO QCD fit: conditions

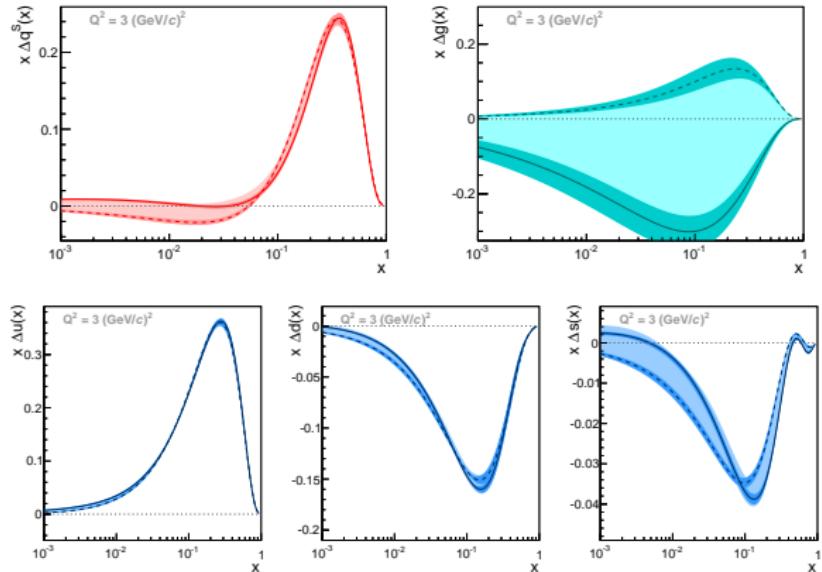
- A fit of the g_1^p , g_1^d , $g_1^{^3\text{He}}$ inclusive data
- $\overline{\text{MS}}$ scheme
- $W^2 < 10 (\text{GeV}/c^2)^2$ excluded
- Number of data points total/COMPASS: 495/138
- Fitted: Δg , $\Delta q^s = \Delta(u + \bar{u}) + \Delta(d + \bar{d}) + \Delta(s + \bar{s})$, gluons, singlet
 $\Delta q_3 = \Delta(u + \bar{u}) - \Delta(d + \bar{d})$ nonsinglet
 $\Delta q_8 = \Delta(u + \bar{u}) + 2\Delta(d + \bar{d}) - \Delta(s + \bar{s})$ nonsinglet
- Parameterisation (at $Q_0^2 = 1 (\text{GeV}/c)^2$):

$$\Delta f_k(x) = \eta_k \frac{x^{\alpha_k} (1-x)^{\beta_k} (1+\gamma_k x)}{\int_0^1 x^{\alpha_k} (1-x)^{\beta_k} (1+\gamma_k x) dx}, \quad (k = s, 3, 8, g)$$

η_k = first moment of $\Delta f_k(x)$ at Q_0^2

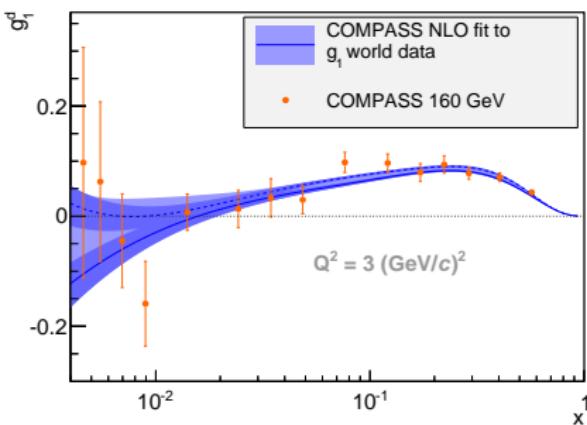
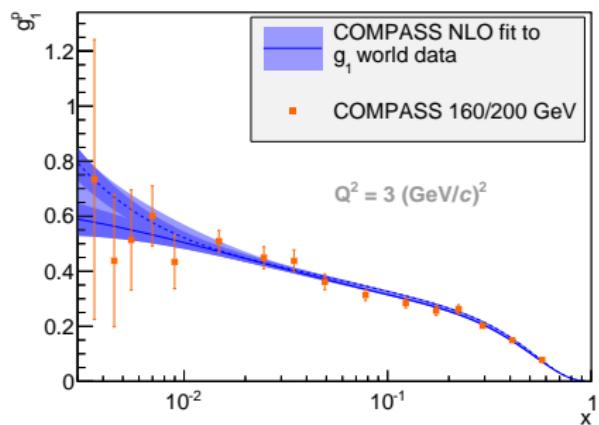
- Number of free parameters in the fitted pdfs: 11
- Only statistical errors taken, normalisations of each data set varied
- Required at every iteration:
 $|\Delta q(x) + \Delta \bar{q}(x)| \leq q(x) + \bar{q}(x), \quad |\Delta g(x)| \leq g(x)$ at $Q^2 = 1 (\text{GeV}/c)^2$
- Studied: parameterisation $\Delta f_k(x)$ dependence and Q_0^2 dependence of the fit.

NLO QCD fit: results



- Statistical uncertainties (dark bands) \ll systematic (light bands)
- Gluon polarisation poorly constraint \implies “direct” methods
- Quark spin contribution to the nucleon spin: $0.26 < \Delta\Sigma < 0.36$ (due to poor Δg)

NLO QCD fit: results...cont'd



- g_1^P clearly positive at low x and raising with decreasing x
- g_1^d consistent with zero at low x ?

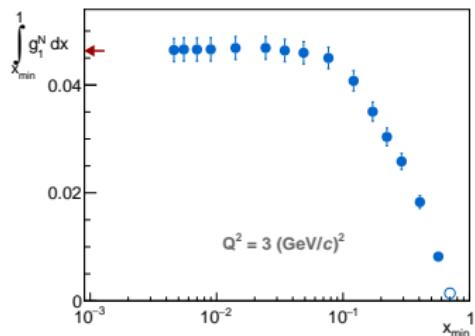
First moments of g_1 and singlet axial charge a_0

hep-ex/1612.00620, accepted PLB

- First moments $\Gamma_1^p, \Gamma_1^d, \Gamma_1^N$
where $\Gamma_1^i = \int_0^1 g_1^i(x, Q^2) dx$

- In particular:

$$\begin{aligned}\Gamma_1^N(Q^2) &= \frac{1}{36} [4a_0 C_S(Q^2) + a_8 C_{NS}(Q^2)] \\ &= \int_0^1 \frac{g_1^d(x, Q^2)}{1 - 1.5\omega_D} dx\end{aligned}$$



- In the $\overline{\text{MS}}$: $a_0 = \Delta\Sigma = (\Delta u + \Delta \bar{u}) + (\Delta d + \Delta \bar{d}) + (\Delta s + \Delta \bar{s})$
- Γ_1^N approaches asymptotic value already at $Q^2 = 3 \text{ (GeV}/c)^2$
- From COMPASS data alone:
 $\Gamma_1^N(Q^2 = 3 \text{ (GeV}/c)^2) = 0.046 \pm 0.002_{\text{stat.}} \pm 0.004_{\text{syst.}} \pm 0.005_{\text{evol.}}$
- From COMPASS data alone (and a_8 from PRD 82 (2010) 114018):
 $a_0(Q^2 = 3 \text{ (GeV}/c)^2) = 0.32 \pm 0.02_{\text{stat.}} \pm 0.04_{\text{syst.}} \pm 0.05_{\text{evol.}}$
(consistent with value from the COMPASS NLO QCD fit of world data).

Non-singlet structure function, $g_1^{\text{NS}}(x, Q^2)$

- Non-singlet structure function:

$$\begin{aligned} g_1^{\text{NS}} &= g_1^{\text{P}}(x, Q^2) - g_1^{\text{N}}(x, Q^2) \\ &= 2 \left[g_1^{\text{P}}(x, Q^2) - g_1^{\text{N}}(x, Q^2) \right] \end{aligned}$$

- Its moment connected to the Bjorken sum rule:

$$\Gamma_1^{\text{NS}}(Q^2) = \int_0^1 g_1^{\text{NS}}(x, Q^2) dx = \frac{1}{6} \left| \frac{g_A}{g_V} \right| C_1^{\text{NS}}(Q^2)$$

- g_1^{NS} calculated, NLO QCD fitted (only Δq_3), evolved to $Q^2 = 3 (\text{GeV}/c)^2$ and fit-extrapolated $x \rightarrow 0, 1$:

$$\Gamma_1^{\text{NS}} = 0.192 \pm 0.007_{\text{stat.}} \pm 0.015_{\text{syst.}}$$

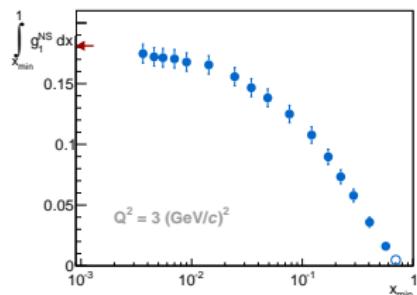
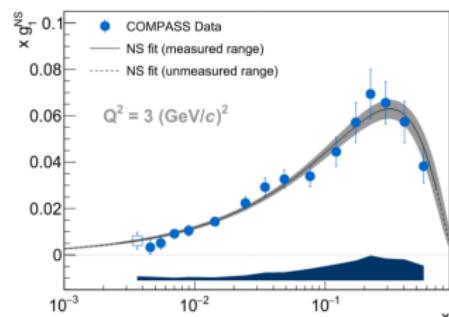
$$\left| \frac{g_A}{g_V} \right| = 1.29 \pm 0.05_{\text{stat.}} \pm 0.10_{\text{syst.}}$$

- Neutron β decay gives: $|g_A/g_V| = 1.2701 \pm 0.002$

PDG, PRD86 (2012) 010001

- This validates the Bjorken sum rule with an accuracy of 9%

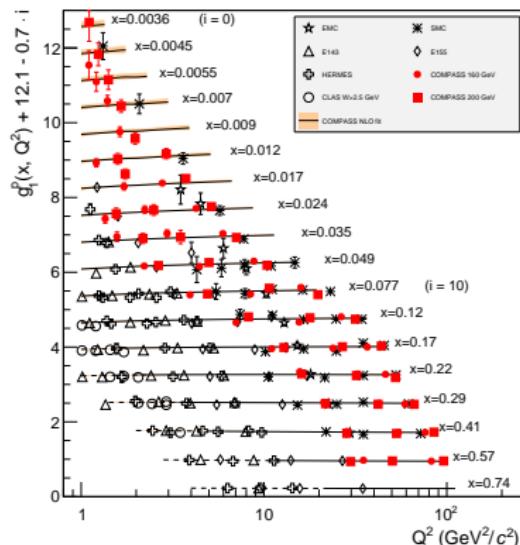
PLB 753 (2016) 18



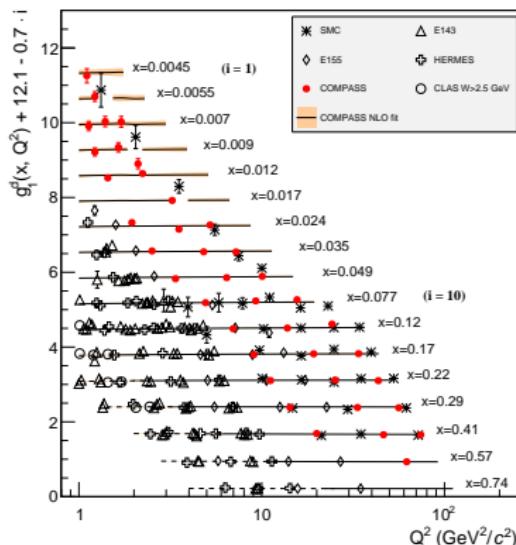
g_1^p and g_1^d , $Q^2 > 1 \text{ (GeV}/c^2)^2$, COMPASS full statistics

COMPASS NLO QCD fit to the world data at $W^2 > 10 \text{ (GeV}/c^2)^2$
 dashed line: extrapolation to $W^2 < 10 \text{ (GeV}/c^2)^2$

proton



deuteron



Phys.Lett.B753(2016)18

hep-ex/1612.00620, accepted PLB

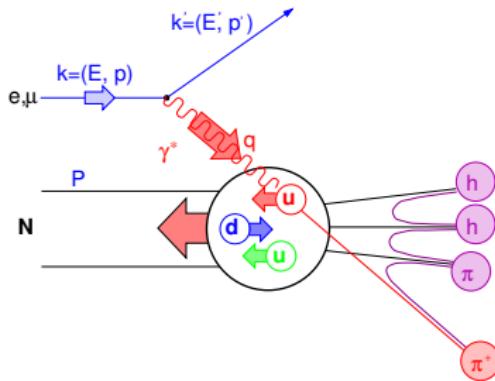
COMPASS measurements at high Q^2 important for the QCD analysis! but little sensitive to Δg

To remember

- COMPASS legacy on $g_1^p(x, Q^2)$ and $g_1^d(x, Q^2)$ presented for DIS ($Q^2 > 1 \text{ (GeV}/c)^2$) and nonperturbative ($Q^2 < 1 \text{ (GeV}/c)^2$) regions
- The $g_1^p(x)$ at low x and low Q^2 is clearly positive ($g_1^d(x)$ is consistent with zero) \Rightarrow first observation of the spin effect at such low x
- The $g_1^p(x)$ at low x is positive (and $g_1^d(x) \approx 0$) also at $Q^2 > 1 \text{ (GeV}/c)^2$
- NLO QCD fit of g_1 world data gave well constrained quark distributions; gluons poorly determined. Quark helicity contribution to nucleon spin:
 $0.26 < \Delta\Sigma < 0.36$
- From the COMPASS data alone:
 - first moments determined and Bjorken sum rule verified to 9% accuracy
 - flavour-singlet axial charge a_0 extracted:
 $a_0(Q^2 = 3 \text{ (GeV}/c)^2) = 0.32 \pm 0.02_{\text{stat.}} \pm 0.04_{\text{syst.}} \pm 0.05_{\text{evol.}}$
(in $\overline{\text{MS}}$ identified with total contribution of quark helicities to the nucleon spin)

SPARES

Nucleon spin structure in DIS: $\mu + N \rightarrow \mu' + X$

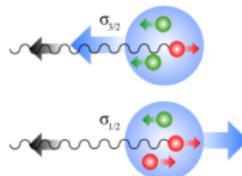


- $\frac{d^2\sigma}{d\Omega dE'} = \frac{\alpha^2}{2Mq^4} \frac{E'}{E} L_{\mu\nu} W^{\mu\nu}$
- Symmetric part of $W^{\mu\nu}$ – unpolarised DIS, antisymmetric – polarised DIS
- Nominally $F_{1,2}$, $q(x, Q^2) \rightarrow g_{1,2}$, $\Delta q(x, Q^2)$ where $q = q^+ + q^-$, $\Delta q = q^+ - q^-$, but...
- ...anomalous gluon contribution to $g_1(x, Q^2)$
- ... $g_2(x, Q^2)$ has no interpretation in terms of partons.

Definitions of DIS variables...

$$\begin{aligned} Q^2 &= -q^2 && \gamma^* \text{ virtuality} \\ x &= Q^2/(2Pq) && \text{Bjorken variable} \\ y &= Pq/(Pk) && \text{relative } \gamma^* \text{ energy} \\ W &= P + q && \gamma^*-N \text{ cms energy} \end{aligned}$$

...and of the γ^*-N asymmetry (e.g. for γ^*-p):



$$A_1(x, Q^2) = \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}}$$

COmmon Muon and Proton Apparatus for Structure and Spectroscopy



NA58, at the CERN SPS
~ 250 physicists
~ 30 institutes



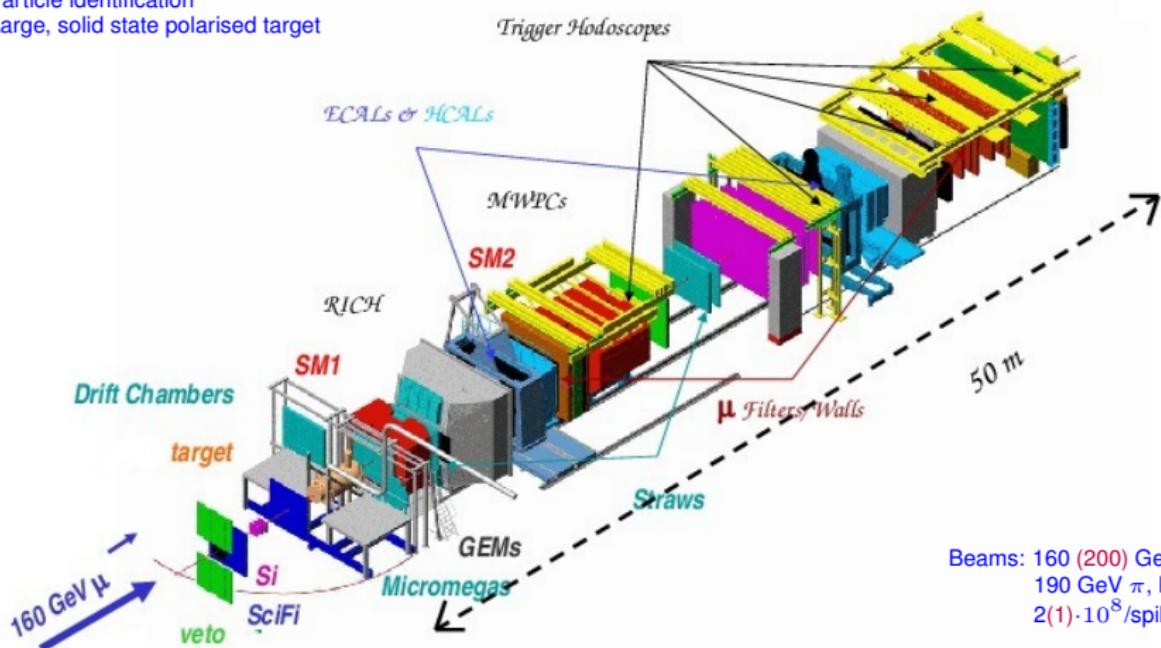
Muon programme	Hadron programme
<p>Spin dependent structure functions g_1</p> <p>Gluon polarisation in the nucleon</p> <p>Quark polarisation distributions</p> <p>Transversity</p> <p>Vector meson production</p> <p>Λ polarisation</p> <p>DVCS/GPD</p>	<p>Primakoff effect, π and K polarisabilities</p> <p>Exotic states, glueballs</p> <p>(Double) charmed barions</p> <p>Multiquark states</p> <p>Drell-Yan process on a polarised target</p>

COMPASS Spectrometer

(muon run)

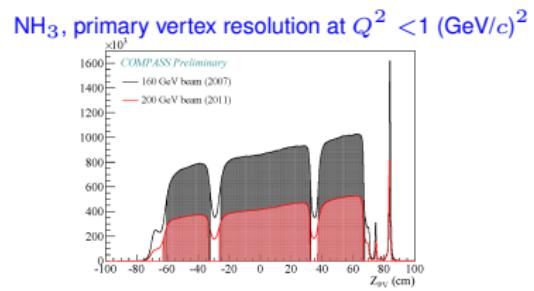
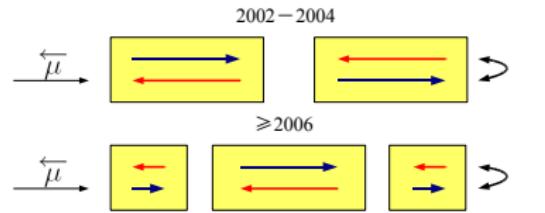
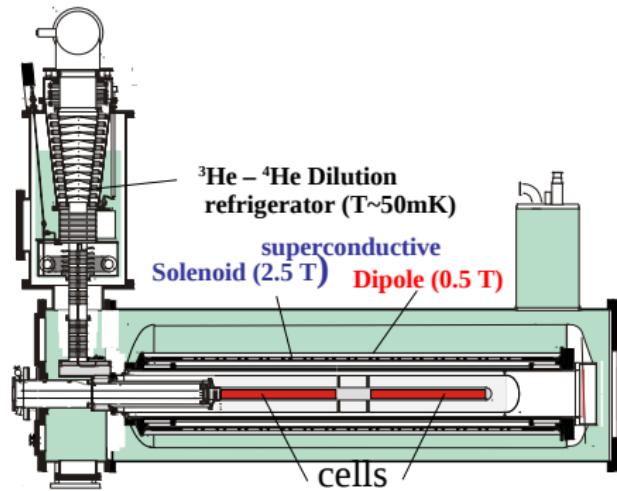
Nucl. Instr. Meth. A577 (2007) 455

Two stages, ~ 350 planes
Calorimetry
Particle identification
Large, solid state polarised target



Beams: 160 (200) GeV μ^\pm
190 GeV π , K, p
 $2(1) \cdot 10^8$ /spill

COMPASS polarised targets: NH₃ and ⁶LiD



- * Two (three) target cells, oppositely polarised
- * Polarisation reversed every 8 h (less frequent after 2005) by field rotation
- * Material: solid ⁶LiD (NH₃)
- * Polarisation: ~ 50% (~90%), by the Dynamical Nuclear Polarisation
- * Dilution: f~0.4 (~0.15)
- * Polar acceptance: ~70 mrad (~180 mrad after 2005)