# COMPASS results on hard exclusive muoproduction



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### Main goals of the GPD program

GPD a 3-dimensional image of the partonic structure of the nucleon

$$H(x, \xi=0, t) \rightarrow H(x, r_{y,z})$$

probability interpretation (Burkardt)



this talk

t-dependence of pure DVCS cross section on unpolarised protons

Contribution to the nucleon spin puzzle
 GPD E related to the orbital angular momentum

$$2J_{q} = \int x (H^{q}(x,\xi,0) + E^{q}(x,\xi,0)) dx$$

$$1/2 = 1/2 \Delta \Sigma + \Delta G + \langle L_z^q \rangle + \langle L_z^g \rangle$$



this talk

Exclusive vector meson production on transversely polarised protons and deuterons

COMPASS experiment at CERN

Two basic ingredients of versatile COMPASS experimental setup

#### secondary beam line M2 from the SPS

- delivers: high energy polarised  $\mu^+$  or  $\mu^-$  beams
  - negative or positive hadron beams

#### two-stage forward spectrometer SM1 + SM2

≈ 300 tracking detectors planes – high redundancy
 variety of tracking detectors to cope with different particle flux
 from θ = 0 to θ ≈ 200 mrad
 + calorimetry, μID, RICH

#### A flexibility to carry out a diverse physics programs by modifying mainly the target region

- spin structure of nucleons; polarised muon-nucleon scattering
- hadron spectroscopy in diffractive and central hadron production
- Primakoff reactions and test of chiral perturbative theory
- polarised and unpolarised Drell-Yan scattering
- GPD studies; DVCS and hard exclusive meson production

# The COMPASS set-up for the GPD program (starting from 2012)

# Main new equipments



ECAL1

#### **Target TOF System**

24 inner & outer scintillators 1 GHz SADC readout goal: **310 ps** TOF resol ECALO Calorimeter Shashlyk modules + MAPD readout





ECAL2

### Mounting of Recoil Proton Detector ('CAMERA') in clean area at CERN



### Recoil particle reconstruction in CAMERA



Proton signature clearly visible after exclusivity selections







$$\frac{Beam Charge & Spin Difference}{\mathcal{D}_{CS,U} \equiv d\sigma(\mu^{+\downarrow}) - d\sigma(\mu^{-\uparrow}) = 2(e_{\mu} a^{BH} Re_{\mu} T^{DVCS} + P_{\mu} d\sigma^{DVCS}_{pol})$$

$$c_{0}^{Int} + c_{1}^{Int} \cos \phi + c_{2}^{Int} \cos 2\phi + c_{3}^{Int} \cos 3\phi$$

$$s_{1}^{DVCS} \sin \phi$$

$$r_{0,1}^{Int} \rightarrow Re(F_{1}\mathcal{H})$$

$$Re \mathcal{H}(\xi,t) = \mathcal{P}\int dx H(x,\xi,t) = \mathcal{P}\int dx H(x,x,t) + \mathcal{D}(t)$$

## Interplay of DVCS and BH at 160 GeV









### Selection of exclusive single photon events



### Estimate of $\pi^0$ background

Major source of background for exclusive photon events

Two cases:

- Visible; detected second  $\gamma$  (below DVCS threshold) => events rejected from final sample
- Invisible; one  $\gamma$  lost => estimated from MC normalised to  $\pi^0$  peak for 'visible' sample



Relative contributions from both processes to  $\pi^0$  background estimated from combined fits to the distributions of 'exclusivity variables' ( $M_x^2$ ,  $\Delta \phi$ ,  $\Delta p_T$ ) and  $E_{miss} = v - E_{\gamma} + t/(2m_p^2)$ 





for normalization of BH MC to the data beam flux measurement used

- dominant BH process at large v (small  $x_{BJ}$ ) clearly visible
- shape of  $\boldsymbol{\phi}$  distribution reproduced well by MC
- estimates of  $\pi^0$  background contributing at small v (large  $x_{BJ}$ )
- at small v (large  $x_{BJ}$ ) an excess of DVCS events above BH + background



### COMPASS acceptance for DVCS



Binning of acceptance in  $Q^2$ , v and |t|

recall: 
$$\frac{d^{3}\sigma^{\mu p}}{dQ^{2}dv\,dt} = \Gamma \frac{d\sigma^{\gamma^{*}p}}{dt}$$

with the virtual photon flux  $\Gamma = \Gamma (Q^2, v)$ 

#### DVCS cross section and t-slope



#### Comparison of t-slope B to HERA results



#### Model independent result

### GPDs and Hard Exclusive Meson Production



Factorisation proven only for  $\sigma_{\rm L}$  $\sigma_{\rm T}$  suppressed by  $1/Q^2$ 

wave function of meson (DA) additional non-perturbative term

➤at Q<sup>2</sup> ≈ few GeV<sup>2</sup> higher order pQCD terms important Chiral-even GPDs<br/>helicity of parton unchanged $H^{q,g}(x,\xi,t)$  $E^{q,g}(x,\xi,t)$  $\widetilde{H}^{q,g}(x,\xi,t)$  $\widetilde{E}^{q,g}(x,\xi,t)$ 

#### Chiral-odd GPDs

helicity of parton changed (not probed by DVCS)

$H^q_T(x,\xi,t)$	$E_T^q(x,\xi,t)$
$\widetilde{H}^q_T(x,\xi,t)$	$\widetilde{E}_{T}^{q}(x,\xi,t)$

# Flavour separation for GPDs example:

$$E_{\rho^{0}} = \frac{1}{\sqrt{2}} \left( \frac{2}{3} E^{u(+)} + \frac{1}{3} E^{d(+)} + \frac{3}{4} E^{g} / x \right)$$
  

$$E_{\omega} = \frac{1}{\sqrt{2}} \left( \frac{2}{3} E^{u(+)} - \frac{1}{3} E^{d(+)} + \frac{1}{4} E^{g} / x \right)$$
  

$$E_{\phi} = -\frac{1}{3} E^{s(+)} + \frac{1}{4} E^{g} / x$$
  
Diehl, Vinnikov  
PLB, 2005

- contribution from gluons at the same order of  $\alpha_{\!_{\rm S}}$  as from quarks

### Spin-dependent cross section for exclusive meson leptoproduction

$$\begin{split} & \left[\frac{\alpha_{em}}{8\pi^{3}}\frac{y^{2}}{1-\epsilon}\frac{1-x_{Bj}}{x_{Bj}}\frac{1}{Q^{2}}\right]^{-1}\frac{d\sigma}{dx_{Bj}dQ^{2}dtd\phi\phi_{s}} \\ & = \frac{1}{2}(\sigma_{++}^{++}+\sigma_{+-}^{--})+\epsilon\sigma_{00}^{++}-\epsilon\cos(2\phi)\operatorname{Re}\sigma_{+-}^{++}-\sqrt{\epsilon(1+\epsilon)}\cos\phi\operatorname{Re}(\sigma_{+0}^{++}+\sigma_{+0}^{--}) \\ & -P_{\ell}\sqrt{\epsilon(1-\epsilon)}\sin\phi\operatorname{Im}(\sigma_{+0}^{++}+\sigma_{+0}^{--}) \\ & -S_{L}\left[\epsilon\sin(2\phi)\operatorname{Im}\sigma_{+-}^{++}+\sqrt{\epsilon(1+\epsilon)}\sin\phi\operatorname{Im}(\sigma_{+0}^{++}-\sigma_{+0}^{--})\right] \\ & +S_{L}P_{\ell}\left[\sqrt{1-\epsilon^{2}}\frac{1}{2}\left(\sigma_{++}^{++}-\sigma_{++}^{--}\right)-\sqrt{\epsilon(1-\epsilon)}\cos\phi\operatorname{Re}(\sigma_{+0}^{++}-\sigma_{+0}^{--})\right] \\ & -S_{T}\left[\sin(\phi-\phi_{S})\operatorname{Im}(\sigma_{+-}^{++}+\epsilon\sigma_{00}^{--})+\frac{\epsilon}{2}\sin(\phi+\phi_{S})\operatorname{Im}\sigma_{+-}^{++}+\frac{\epsilon}{2}\sin(3\phi-\phi_{S})\operatorname{Im}\sigma_{+-}^{-+} \\ & +\sqrt{\epsilon(1+\epsilon)}\sin\phi_{S}\operatorname{Im}\sigma_{+0}^{++}+\sqrt{\epsilon(1+\epsilon)}\sin(2\phi-\phi_{S})\operatorname{Im}\sigma_{+0}^{-+}\right] \\ & +S_{T}P_{\ell}\left[\sqrt{1-\epsilon^{2}}\cos(\phi-\phi_{S})\operatorname{Re}\sigma_{+0}^{++}\right] - \sqrt{\epsilon(1-\epsilon)}\cos(2\phi-\phi_{S})\operatorname{Re}\sigma_{+0}^{-+}\right]. \end{split}$$

 $\sigma_{\it mn}^{\it ij}$ : helicity-dependent photoabsorption cross sections and interference terms

$$\sigma_{mn}^{ij}(x_B,Q^2,t)\propto \sum (M_m^i)^*M_n^j$$

 $M_m^i$ : amplitude for subprocess  $\gamma^* p \to V p'$  with photon helicity *m* and target proton helicity *i* 

$$\epsilon = \frac{1 - y - \frac{1}{4}y^2\gamma^2}{1 - y + \frac{1}{2}y^2 + \frac{1}{4}\gamma^2}$$
$$\gamma = 2x_{Bi}M_P/Q$$

#### Azimuthal asymmetries of cross section for exclusive meson leptoproduction



 $\sigma_{0}$  - 'unpolarised cross section'

$$\sigma_0 = \frac{1}{2} \left( \sigma_{++}^{++} + \sigma_{++}^{--} \right) + \epsilon \sigma_{00}^{++} = \sigma_L + \epsilon \sigma_T$$

### COMPASS polarised target



#### Exclusive $\omega$ production on $p^{\uparrow}$ at COMPASS

(Selections similar for  $\rho^0$  sample)

 $\begin{array}{c|c} \mu \ N \rightarrow \mu \ \omega \ N \end{array} \quad \text{Trans} \\ \hline & & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ &$ 

Transversely polarised proton target (NH<sub>3</sub>), 2010 data

note: there was no Recoil Detector for these data

only two hadron tracks of opposite charge associated to the primary vertex only two ECAL clusters time-correlated with beam and not associated to a charged particle



#### Extraction of asymmetries and subtraction of non-exclusive background

- $\rho^{\rm 0}$  analysis
  - 1D (deuteron) and 2D (proton) binned maximum likelihood estimator with subtraction of background in ( $\phi$ ,  $\phi$ <sub>s</sub>) bins
- $\omega$  analysis
  - Unbinned maximum likelihood estimator with simultaneous fit of signal and background asymmetries

# Background rejection:

For each target cell and polarization state



shape of semi-inclusive background from MC (LEPTO with COMPASS tuning + simulation of spectrometer response + reconstruction as for real data)

MC weighted using ratio between real data and MC for wrong charge combination sample  $(h^+h^+\gamma\gamma + h^-h^-\gamma\gamma)$ 

$$w(E_{miss}) = \frac{N_{RD}^{h+h+\gamma\gamma}(E_{miss}) + N_{RD}^{h-h-\gamma\gamma}(E_{miss})}{N_{MC}^{h+h+\gamma\gamma}(E_{miss}) + N_{MC}^{h-h-\gamma\gamma}(E_{miss})}$$

Normalization of MC to the real data using two component fit Gaussian function (signal) + shape from MC (bkg)



•  $A_{UT}^{sin(\phi-\phi_s)}$  for transversely polarised protons and deuterons small, compatible with 0

- for the proton agreement with HERMES results COMPASS results with statistical errors improved by factor 3 and extended kinematic range
- for the deuteron the first measurement
- reasonable agreement with predictions of the GPD model of Goloskokov Kroll

[EPJ C59 (2009) 809]

small values expected due to approximate cancellation of contributions from  $E^u$  and  $E^d$ ,  $E^u \approx -E^d$ (cf. upper-right plot)

### Complete set of transverse target spin asymmetries for exlusive $\rho^0$ production on $p^{\uparrow}$



#### Single spin asymmetries

- Improved method of extraction (2D)
- Simultaneous extraction of
   5 single spin asymmetries and
   3 double spin asymmetries
   for transversely polarised protons

→ PLB 731 (2014) 19

- predictions of GPD model of Goloskokov-Kroll
- reasonable agreement with GK model (also for not-shown double spin asym.)



indication of H<sub>T</sub>, 'transversity' GPD, contribution

larger effects for some asymmetries expected for exclusive oproduction

### Azimuthal asymmetries for exlusive () production on p<sup>1</sup>



#### Single spin asymmetries

- new result, to be published
- unbinned maximum likelihood method
- extraction of 8 transverse spin asymmetries



when 'global' comparison to the data no clear preference for any version

## Comparison to HERMES asymmetries for $\omega$ production on p^{\uparrow}



✓ Note: contribution of pion pole decreases with W

-> each experiment to be compared to corresp. predictions

COMPASS uncertainties smaller by a factor > 2

✓ within large errors combined HERMES data compatible with all 3 scenarios

(authors conclusion 'data favor a positive  $\pi\omega$  transition form factor')

Future measurements at JLab12
 EPJ A48 (2012) 187
 expected to resolve the issue of πω transition form factor

## **COMPASS-II** time lines

#### Part of the COMPASS-II proposal approved and scheduled by CERN

- > 2012: pion and kaon polarisabilities (Primakoff) + comissioning and pilot run for DVCS
- > 2013-2014: long SPS/LHC shutdown
- > 2014-2015: Drell-Yan measurements with transversely polarised protons (NH<sub>3</sub> target)
- > 2016-2017: stage 1 of GPD program and in parallel SIDIS (LH target)
- > 2018: Drell-Yan measurements with transversely polarised protons (NH<sub>3</sub> target)

#### Measurements to be pursued at COMPASS-II > 2020 (subject to a new proposal)

- ✓ stage 2 of GPD program with transversely polarised NH<sub>3</sub> target and RPD
- ✓ SIDIS (high statistics) from transversely polarised deuteron and proton targets
- ✓ Drell-Yan on transversely polarised deuterons, unpolarised protons and nuclear targets
- ✓ hadron spectroscopy program with high-intensity separated kaon and antiproton beams



Backup





Azimuthal dependence of exclusive photon xsec.

## from Belitsky, Kirchner, Müller :

e,

polarized beam off unpolarized target

$$d\sigma_{(\mu\rho \to \mu\rho\gamma)} = d\sigma^{BH} + d\sigma^{DVCS}_{unpol} + P_{\mu} d\sigma^{DVCS}_{pol} + e_{\mu} a^{BH} Re A^{DVCS} + e_{\mu} P_{\mu} a^{BH} Im A^{DVCS}$$

$$d\sigma^{BH} = \frac{\Gamma(x_{\beta}, Q^{2}, t)}{P_{1}(\varphi)P_{2}(\varphi)} (c_{0}^{BH} + C_{1}^{BH} \cos \varphi + c_{2}^{BH} \cos 2\varphi) \leftarrow \text{Known expression}$$

$$d\sigma^{DVCS}_{unpol} = \frac{e^{6}}{\gamma^{2}Q^{2}} (c_{0}^{DVCS} + C_{1}^{DVCS} \cos \varphi + c_{2}^{DVCS} \cos 2\varphi)$$

$$P_{\mu} \times d\sigma^{DVCS}_{pol} = \frac{e^{6}}{\gamma^{2}Q^{2}} (s_{0}^{DVCS} \sin \varphi)$$

$$e_{\mu} \times a^{BH} \Re e A^{DVCS} = \frac{e^{6}}{x\gamma^{3}tP_{1}(\varphi)P_{2}(\varphi)} (c_{0}^{Int} + c_{1}^{Int} \cos \varphi + c_{2}^{Int} \cos 2\varphi + c_{3}^{Int} \cos 3\varphi)$$

$$e_{\mu}P_{\mu} \times a^{BH} \Im m A^{DVCS} = \frac{e^{6}}{x\gamma^{3}tP_{1}(\varphi)P_{2}(\varphi)} (s_{1}^{Int} \sin \varphi + s_{2}^{Int} \sin 2\varphi)$$
Twist-2 >> Twist-3 Twist-2 gluon

 $\gamma^*$ 

A

### COMPASS acceptance for DVCS (1)



Symmetric acceptance in  $\phi$  leads to cancellation of the interference terms when integrated over  $\phi$ 

### Beam Charge&Spin Difference of cross sections

$$\mathcal{D}_{CS,U} \equiv d\sigma(\mu^{+\downarrow}) - d\sigma(\mu^{-\uparrow}) = c_0^{Int} + c_1^{Int} \cos \phi + c_2^{Int} \cos 2\phi + c_3^{Int} \cos 3\phi + s_1^{DVCS} \sin \phi$$

$$c_{0,1}^{Int} \rightarrow \mathcal{R}e(F_1\mathcal{H})$$

![](_page_32_Figure_2.jpeg)

t-slope measurement for exclusive  $\rho^0$  production

![](_page_33_Figure_1.jpeg)

![](_page_33_Figure_2.jpeg)

 $(=Q^2 \text{ for DVCS})$ 

### Exclusive $\rho^{0}$ production on $p^{\uparrow}$ and $d^{\uparrow}$ at COMPASS

![](_page_34_Picture_1.jpeg)

Transversely polarised proton target (NH<sub>3</sub>), 2007, 2010 Transversely polarised deuteron target (<sup>6</sup>LiD), 2003-2004

note: there was no RPD for these data

only two hadron tracks of opposite charge associated to the primary vertex

![](_page_34_Figure_5.jpeg)

### Role of pion exchange

![](_page_35_Figure_1.jpeg)

- Effect known since early photoproduction experiments
- At COMPASS kinematics:
  - small for ρ<sup>0</sup> production
  - sizable for  $\omega$  production
- Unnatural parity exchange process
   → impact on helicity-dependent observables
- Crucial for description of SDMEs for excl. ω production
   → Goloskokov and Kroll, Eur. Phys. J. A50 (2014) 9, 146
- Sign of  $\pi\omega$  form factor not resolved from SDMEs data  $\rightarrow$  azimuthal asymmetries more sensitive

![](_page_35_Figure_9.jpeg)