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National Centre for
Nuclear Research
Warsaw, Poland

Gluon Contribution to the Sivers Effect COMPASS Results on Deuteron Target

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Spin2014

The 21st International Symposium on Spin Physics

October 20-24, 2014, Beijing, China

Beam: $2 \cdot 10^8 \mu^+$ / spill (4.8s / 16.2s)

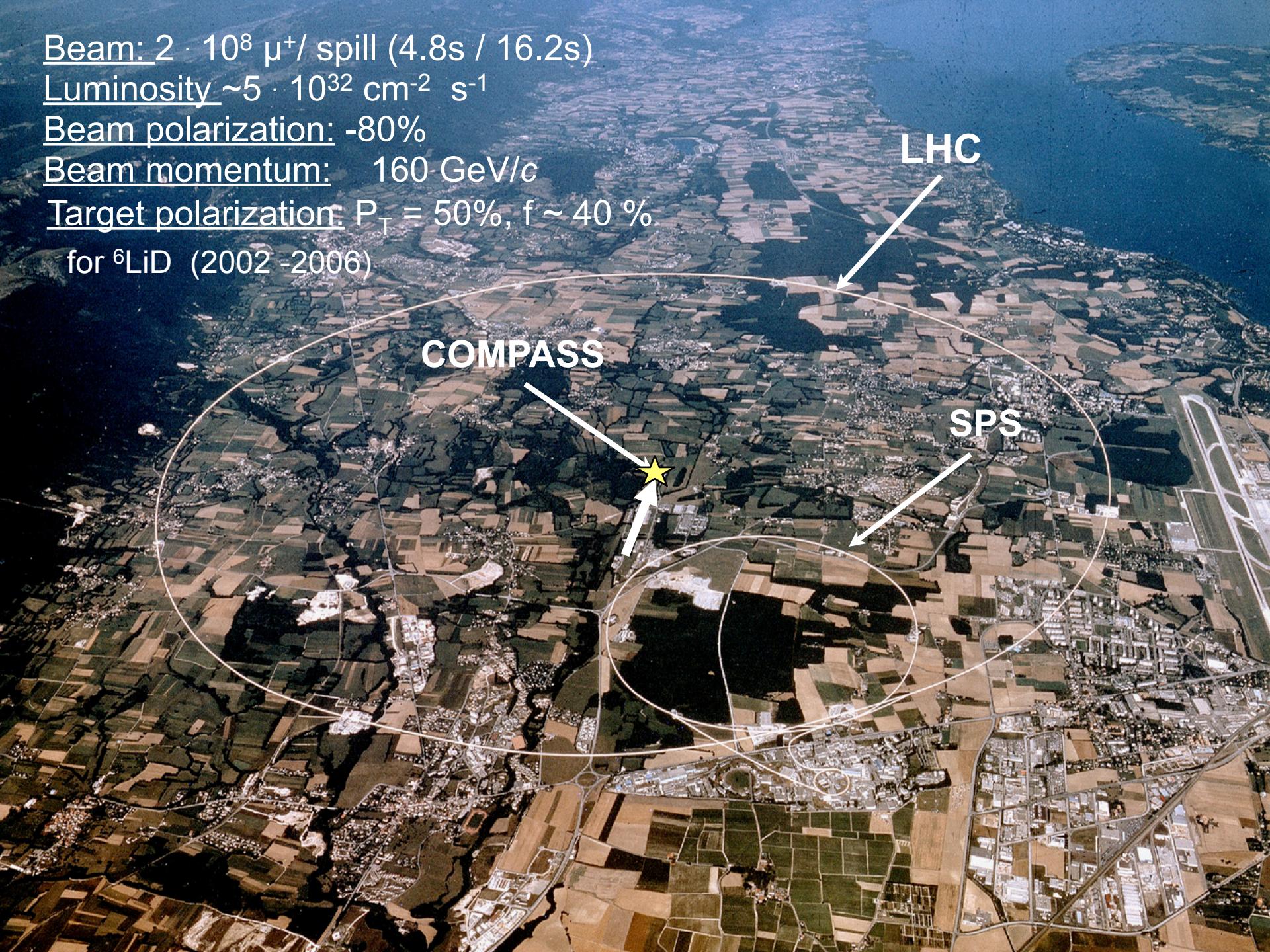
Luminosity $\sim 5 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$

Beam polarization: -80%

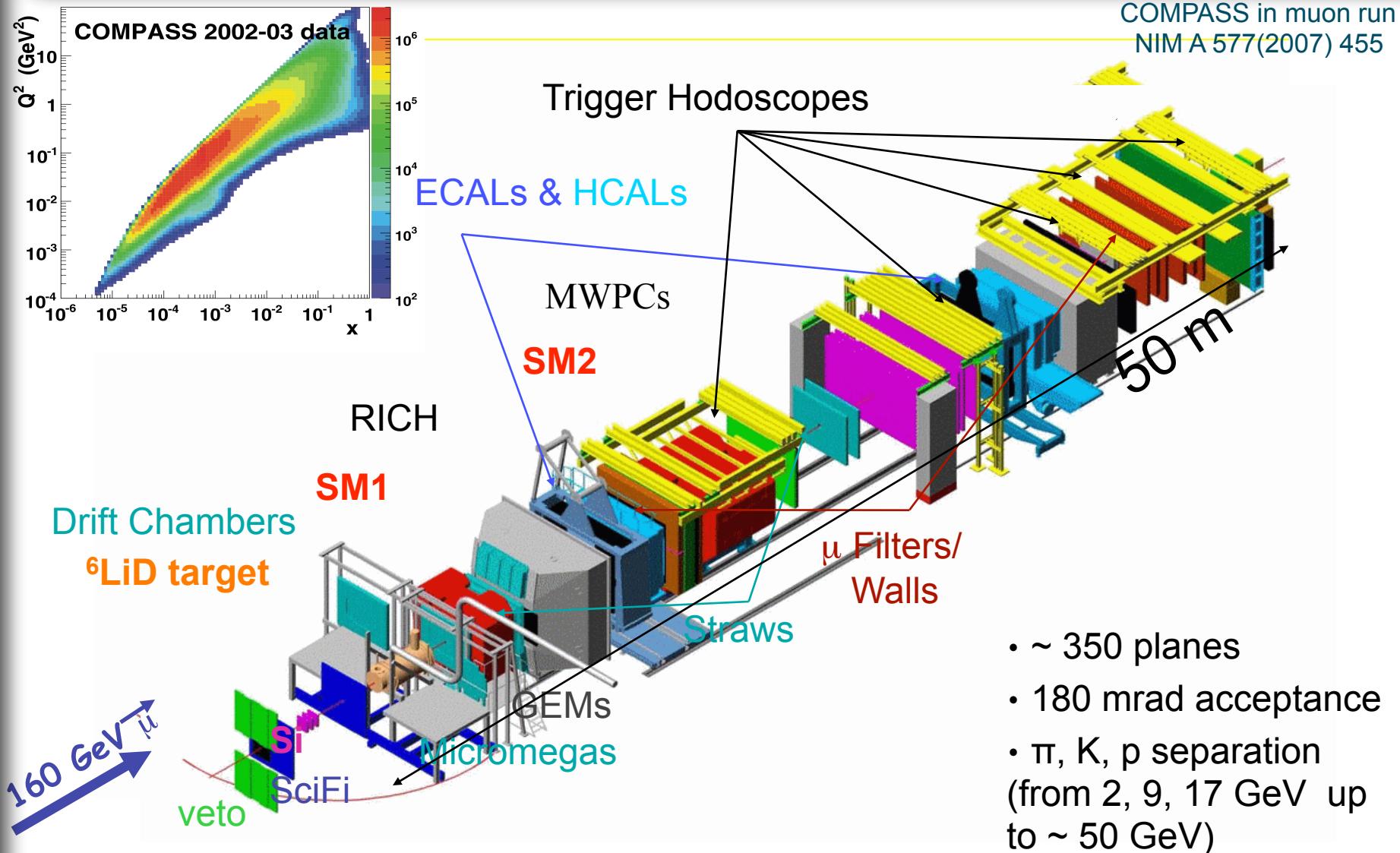
Beam momentum: 160 GeV/c

Target polarization: $P_T = 50\%$, $f \sim 40\%$

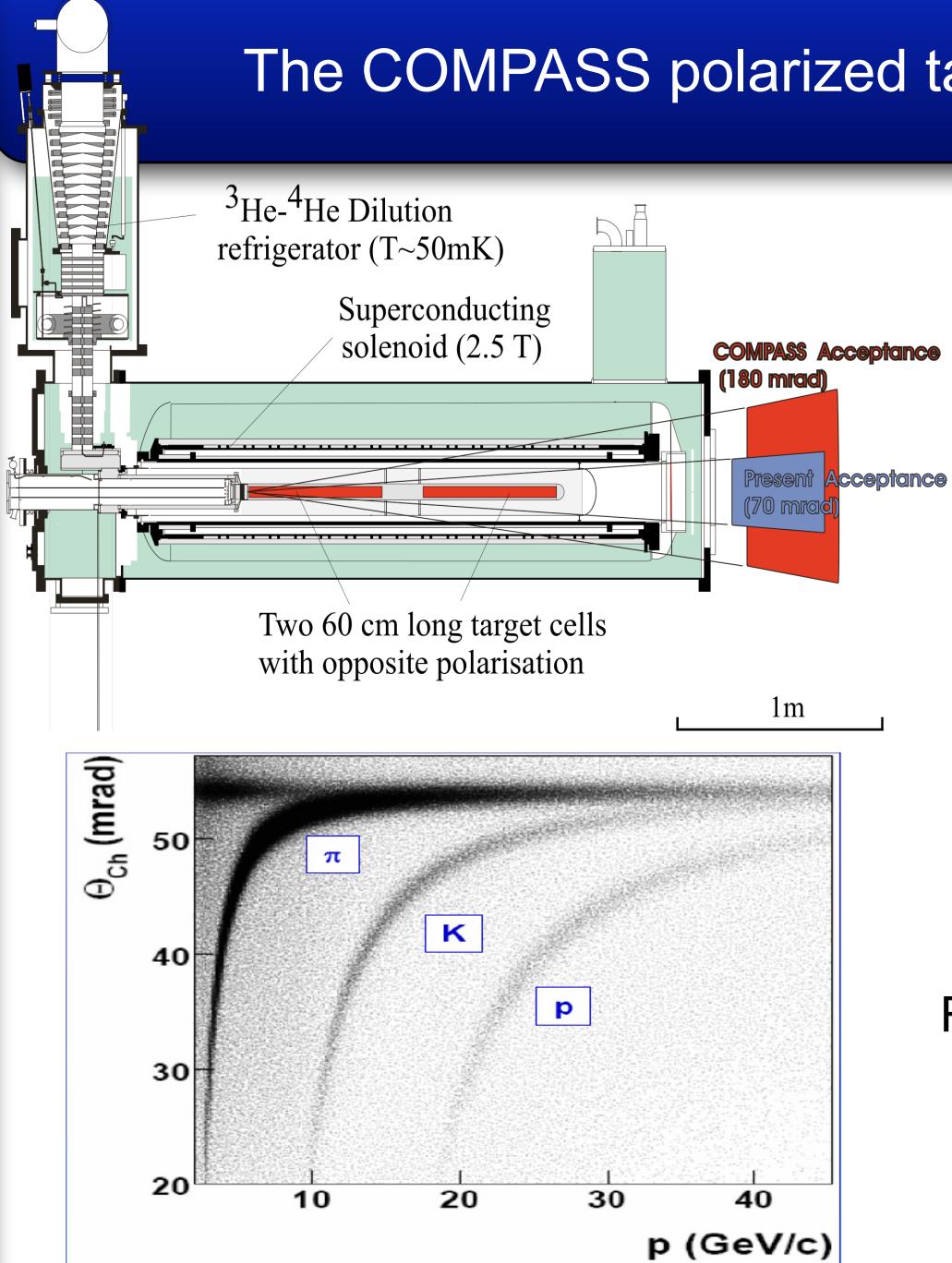
for ${}^6\text{LiD}$ (2002 -2006)



The COMPASS spectrometer

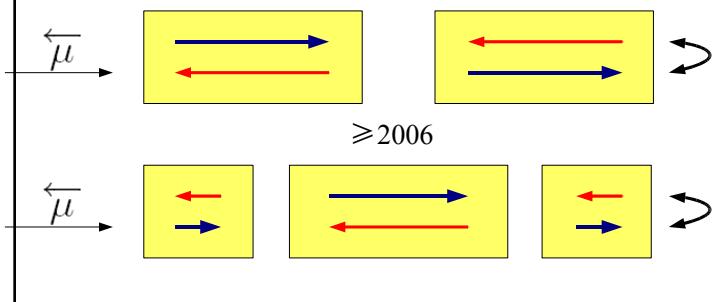


The COMPASS polarized target and PID



Target material: ${}^6\text{LiD}$
 Polarisation: >50%
 Dilution factor: ~0.4
 Dynamic Nuclear Polarization

2006 - new solenoid
 with acceptance 180 mrad
 3 target cells
 (reduce false asymmetries)
 2002 – 2004



RICH 2006 upgrade : better PID

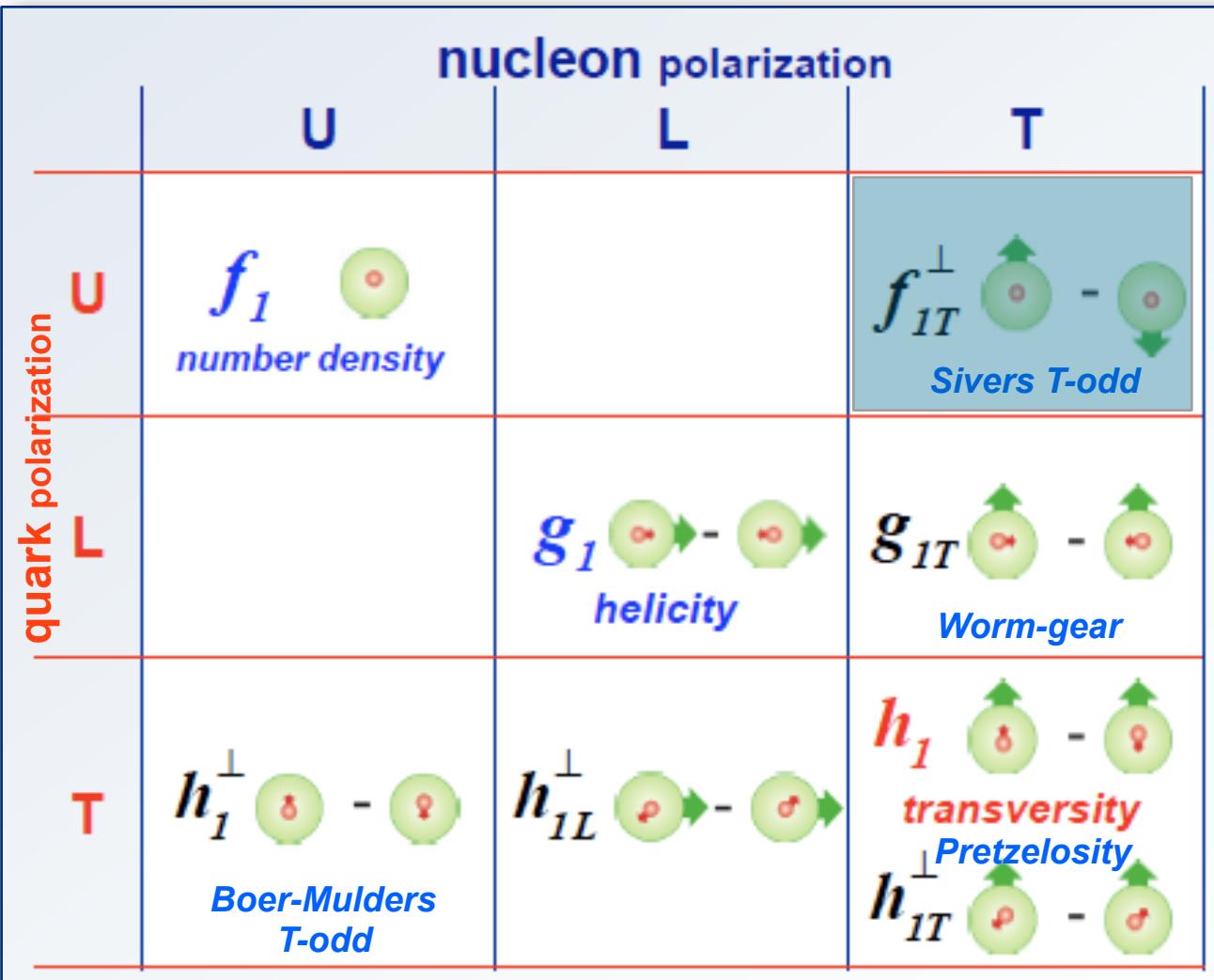
MAPMTs in central region

APV electronics in periphery

- Introduction
- Gluon „Sivers efect” measurement @ COMPASS
- Artificial Neural Network approach
- Data and Monte Carlo comparison
- Validation of the method - MC data
- Data selection
- Preliminary results on deuteron target
- Systematic studies
- Summary

Beyond collinear approximation - k_T dependence

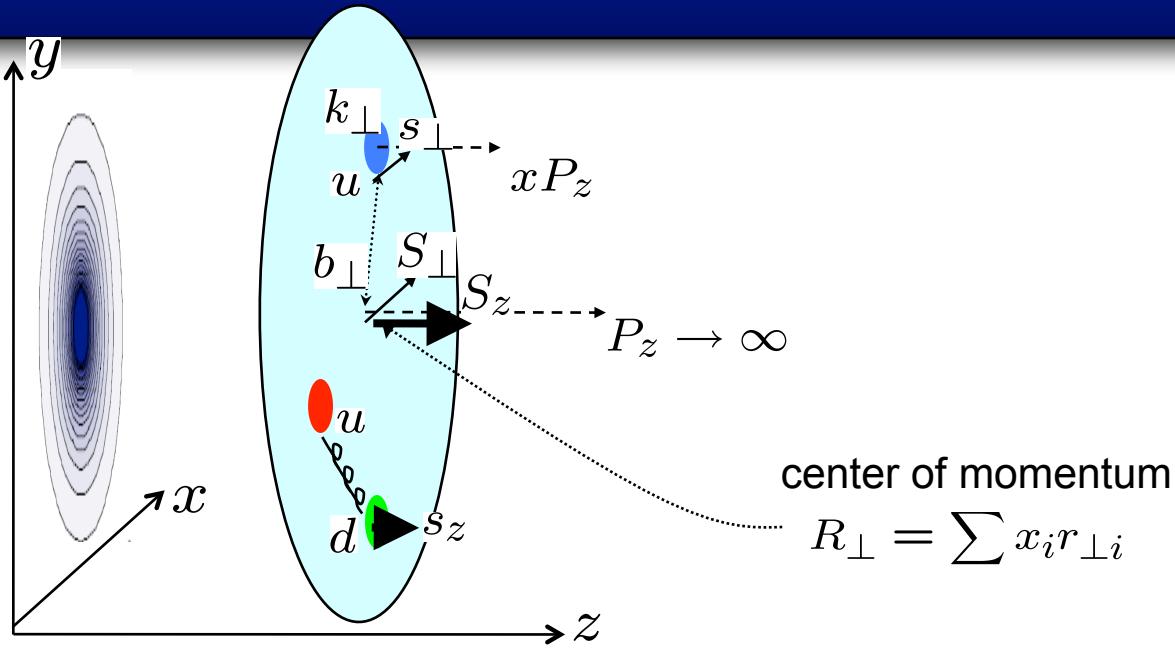
Introduction



LO, twist-2 - 8 independent functions to parameterize structure

Spatial deformation & GPD

Burhardt,2000



Introduction

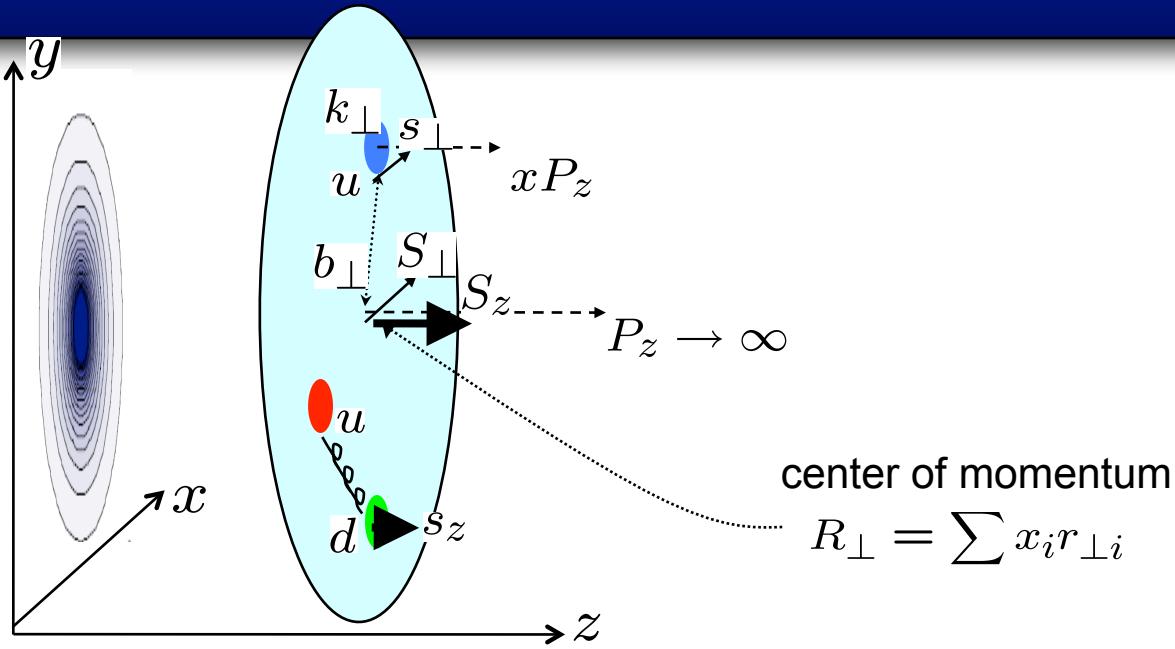
For a transversely polarized nucleon (e.g. polarized in the $+\hat{x}$ -direction) the IPD $q_{\hat{x}}(x, \vec{b}_{\perp})$ is no longer symmetric due to the non-zero value of the spin-flip GPD E . This deformation is described by the gradient of the Fourier transform of E :

$$q_{\hat{x}}(x, \vec{b}_{\perp}) = \mathcal{H}(x, \vec{b}_{\perp}) - \frac{1}{2M} \frac{\partial}{\partial b_y} \mathcal{E}(x, \vec{b}_{\perp}).$$

non-zero spin-flip GPD E - existence of non-zero orbital momentum

Spatial deformation & GPD

Burhardt,2000



Introduction

For a transversely polarized nucleon (e.g. polarized in the $+\hat{x}$ -direction) the IPD $q_{\hat{x}}(x, \vec{b}_\perp)$ is no longer symmetric due to the non-zero value of the spin-flip GPD E . This deformation is described by the gradient of the Fourier transform of E :

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non-zero spin-flip GPD E - existence of non-zero orbital momentum

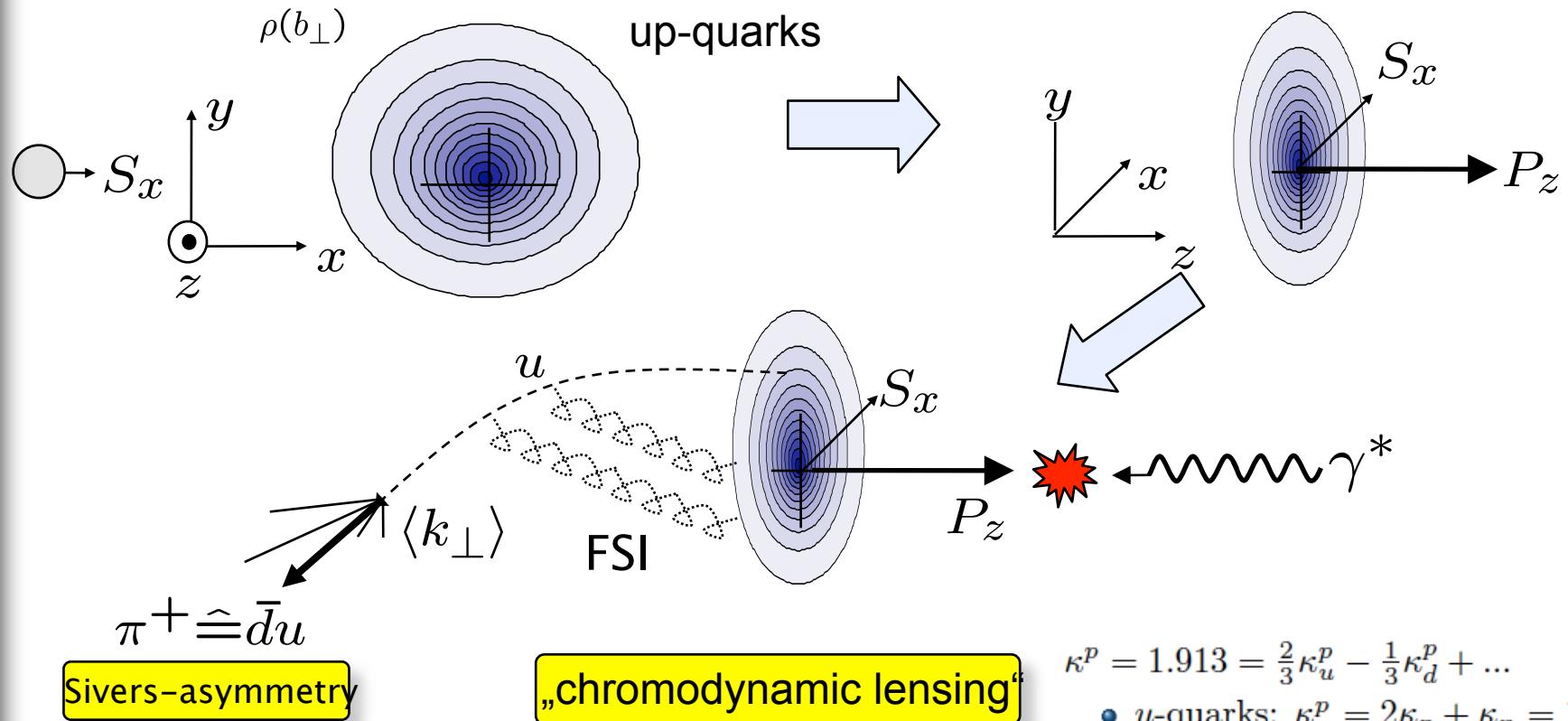
Sivers function and spatial deformation



M.Burhardt 2002/2003

Dynamical origin of quark transverse momentum

Introduction



Deformation details are model-dependent but the size and directions is determined by anomalous magnetic moments of proton and neutron.

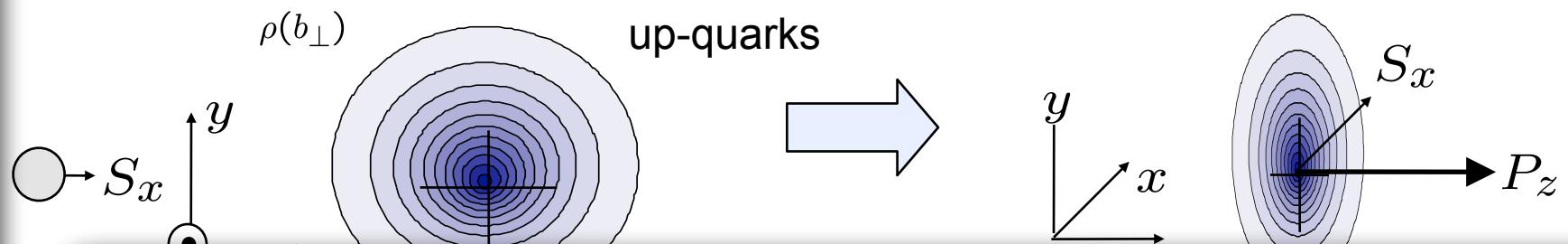
$$\kappa^p = 1.913 = \frac{2}{3}\kappa_u^p - \frac{1}{3}\kappa_d^p + \dots$$

- u -quarks: $\kappa_u^p = 2\kappa_p + \kappa_n = 1.673$
- ↪ shift in $+\hat{y}$ direction
- d -quarks: $\kappa_d^p = 2\kappa_n + \kappa_p = -2.033$
- ↪ shift in $-\hat{y}$ direction
- $\langle b_y^q \rangle = \mathcal{O}(\pm 0.2 \text{ fm})$!!!!

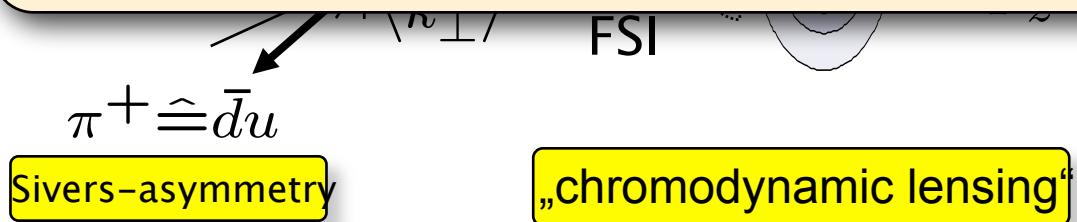
Sivers function and spatial deformation

M.Burhardt 2002/2003

Dynamical origin of quark transverse momentum



Q: Is gluon's spacial distribution in transverse plane also deformed?



Deformation details are model-dependent but the size and directions is determined by anomalous magnetic moments of proton and neutron.

- $\kappa^p = 1.913 = \frac{2}{3}\kappa_u^p - \frac{1}{3}\kappa_d^p + \dots$
- u -quarks: $\kappa_u^p = 2\kappa_p + \kappa_n = 1.673$
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- $\langle b_y^q \rangle = \mathcal{O}(\pm 0.2 \text{ fm})$!!!!

The idea: selection of high- p_T hadron pairs to increase photon-gluon fusion (PGF) similar to Δg determination with longitudinally polarised target.

difficulties comparing to Δg

- asymmetry in azimuthal angle
- gluon „simulated” from pair of hadrons from PGF
- Sivers effect due to final interactions (no analyzing powers)

Single-spin asymmetry is measured:

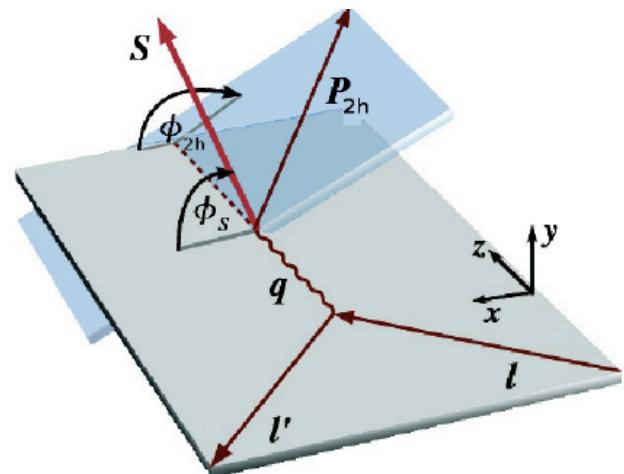
$$A_T^h \sim \frac{d^6\sigma^\uparrow - d^6\sigma^\downarrow}{d^6\sigma^\uparrow + d^6\sigma^\downarrow}.$$

8 asymmetries; concentrated on Sivers

The statistically weighted method was used, similar to open-charm and all- p_T methods used in Δg determination

J.Pretz, J-M Le Goff NIM A 602 (2009) 594

COMPASS open-charm: Phys. Rev. D (2013) 052018



$$\mathbf{P}_{2h} = \mathbf{p}_1 + \mathbf{p}_2$$

Sivers angle:

$$\phi = \phi_{2h} - \phi_s$$

The weighted method

The number of events:

$$n_c(\vec{x}) = \alpha_c(\vec{x})(1 + \beta_c(\vec{x})A_{UT}^{\sin(\phi_{2h}-\phi_s)}(\vec{x}))$$

where:

$$\vec{x} = (x_{Bj}, y, t, \phi, \dots)$$

$$\alpha_c(\vec{x}) = a_c \Phi n_c \sigma,$$

$$\beta_c(\vec{x}) = P_T f \sin(\phi_{2h} - \phi_s)$$

$$c = u, d, u', d'$$

Every event is weighted by the weight ω

$$\begin{aligned} p_c := \int \omega(\vec{x}) n_c(\vec{x}) d\vec{x} &= \int \omega(\vec{x}) \alpha_c(\vec{x}) d\vec{x} \\ &+ \int \omega(\vec{x}) \alpha_c(\vec{x}) \beta_c(\vec{x}) A_{UT}^{\sin(\phi_{2h}-\phi_s)}(\vec{x}) d\vec{x} \approx \sum_{i=1}^{N_c} \omega_i, \\ \sum_{i=1}^{N_c} \omega_i &= \tilde{\alpha}_c (1 + \{\beta_c\}_{\omega} \boxed{A_{UT}^{\sin(\phi_{2h}-\phi_s)}}_{\omega \beta_c}), \end{aligned}$$

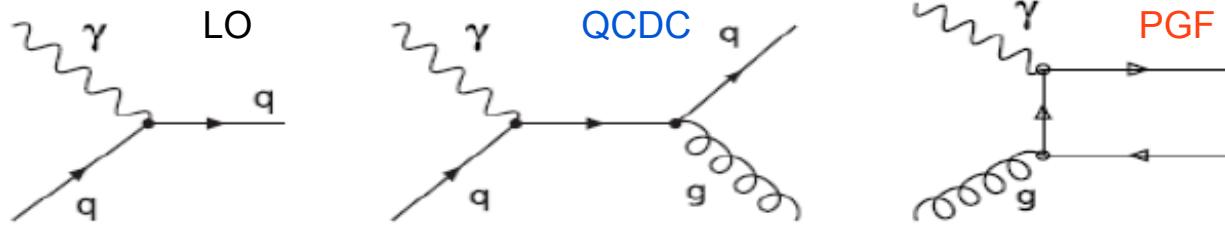
$$\omega(\vec{x}) = \frac{\beta(\vec{x})}{P_T} = f \sin(\phi_{2h} - \phi_s)$$

where

$$\begin{aligned} \tilde{\alpha}_c &= \int \omega(\vec{x}) \alpha_c(\vec{x}) d\vec{x}, \\ \{A_{UT}^{\sin(\phi_{2h}-\phi_s)}\}_{\omega \beta_c} &= \frac{\int A_{UT}^{\sin(\phi_{2h}-\phi_s)}(\vec{x}) \omega(\vec{x}) \beta_c \alpha_c(\vec{x}) d\vec{x}}{\int \omega(\vec{x}) \beta_c \alpha_c(\vec{x}) d\vec{x}}, \\ \{\beta\}_{\omega} &= \frac{\int \beta(\vec{x}) \omega(\vec{x}) \alpha_c(\vec{x}) d\vec{x}}{\int \omega(\vec{x}) \alpha_c(\vec{x}) d\vec{x}} \approx \frac{\sum_i \beta_i \omega_i}{\sum_i \omega_i}. \end{aligned}$$

The weighted method

Physical model: three basic processes @LO



$$A_{UT}^{\sin(\phi_{2h} - \phi_s)} = R_{PGF} A_{PGF}^{\sin(\phi_{2h} - \phi_s)}(< x_G >) + R_{LP} A_{LP}^{\sin(\phi_{2h} - \phi_s)}(< x_{Bj} >) \\ + R_{QCD} A_{QCD}^{\sin(\phi_{2h} - \phi_s)}(< x_C >)$$

$$\begin{aligned} \omega_{PGF} &\equiv \omega^G = R_{PGF} f \sin(\phi_{2h} - \phi_s) = \frac{\beta^G}{P_T}, \\ \omega_{LP} &\equiv \omega^L = R_{LP} f \sin(\phi_{2h} - \phi_s) = \frac{\beta^L}{P_T}, \\ \omega_{QCD} &\equiv \omega^C = R_{QCD} f \sin(\phi_{2h} - \phi_s) = \frac{\beta^C}{P_T}. \end{aligned}$$

The weighted method

Physical model: three basic processes @LO

leads to 12 eqs.:

$$\begin{aligned}
 p_c^j &= \sum_{i=1}^{N_c} \omega_i^j = \tilde{\alpha}_c^j (1 + \{\beta_c^G\}_{\omega^j} A_{PGF}^{\sin(\phi_{2h}-\phi_s)} (\langle x_G \rangle) \\
 &\quad + \{\beta_c^L\}_{\omega^j} A_{LP}^{\sin(\phi_{2h}-\phi_s)} (\langle x_B \rangle) + \{\beta_c^C\}_{\omega^j} A_{QCDC}^{\sin(\phi_{2h}-\phi_s)} (\langle x_C \rangle)) \\
 &= \tilde{\alpha}_c^j (1 + A_{PGF} \{\beta^G\}_{\omega^j} + A_{LP} \{\beta^L\}_{\omega^j} + A_{QCDC} \{\beta^C\}_{\omega^j})
 \end{aligned}$$

with 15 unknowns: (3 asymmetries + 12 acceptances) but thanks to it is reduced to 12. *Here j stands for LO, QCDC and PGF, respectively

$$\frac{\tilde{\alpha}_u^j \tilde{\alpha}_{d'}^j}{\tilde{\alpha}_d^j \tilde{\alpha}_{u'}^j} = 1,$$

To determine asymmetries the minimization procedure has been used:

$$\chi^2 = (\vec{N_{exp}} - \vec{N_{obs}})^T \text{Cov}^{-1} (\vec{N_{exp}} - \vec{N_{obs}})$$

$$\sim \sum_{N_c} \omega_x \omega_y.$$

$$\vec{N_{obs}} = \left(\sum_{i=0}^{N_u} \omega_i^G, \dots, \sum_{i=0}^{N_d} \omega_i^C \right),$$

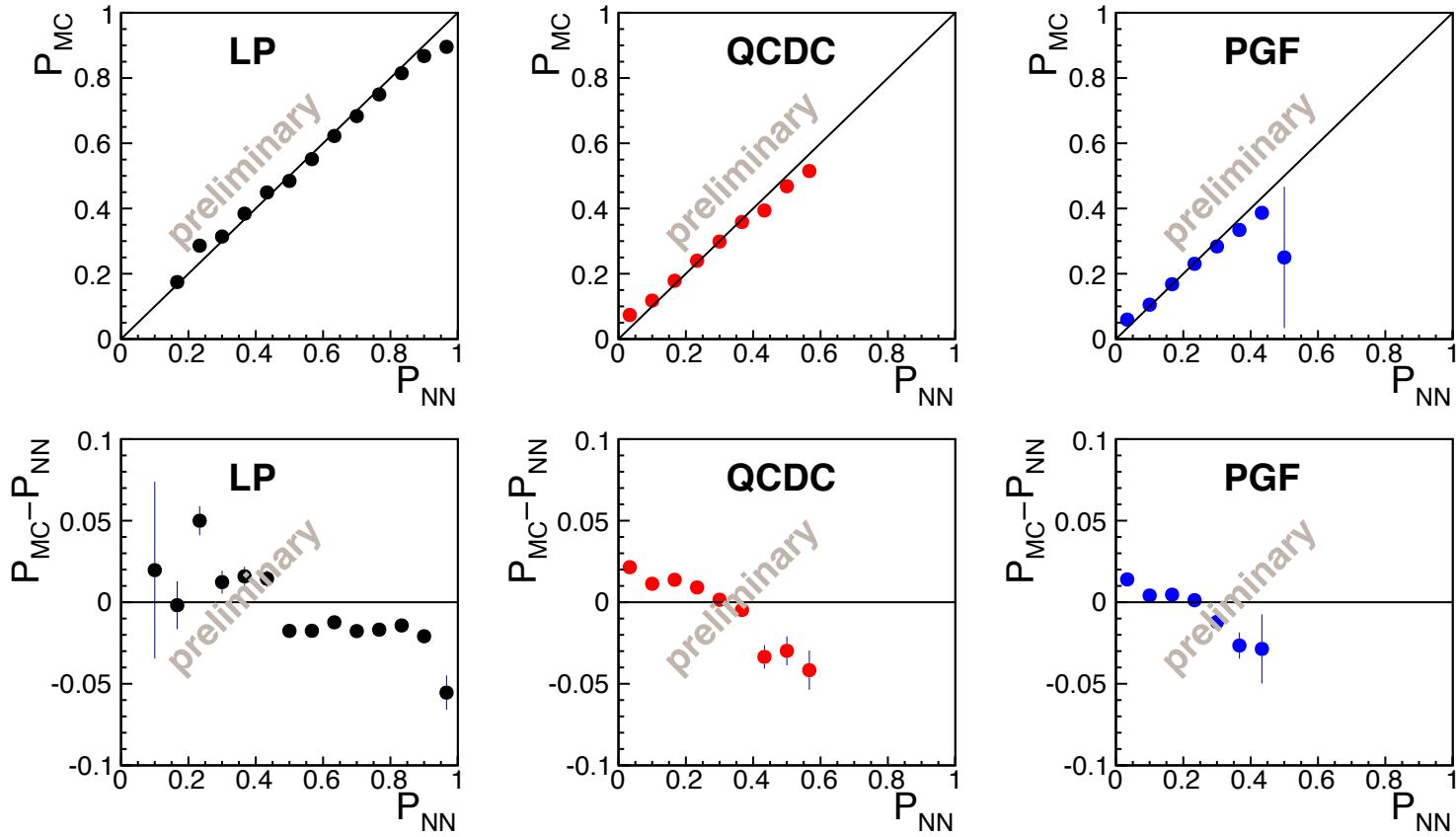
$$\vec{N_{exp}} = (N_{exp,G}^u, \dots, N_{exp,C}^{d'}), \quad N_{exp,i}^c = \tilde{\alpha}_c^j (1 + A_{PGF} \{\beta^G\}_{\omega^j} + A_{LP} \{\beta^L\}_{\omega^j} + A_{QCDC} \{\beta^C\}_{\omega^j})$$

Artificial Neural Network approach - validation



To find R's (fractions or probabilities of three processes) the ANN approach has been used (as in longitudinal high- p_T analysis for gluon polarisation measurement, see: Phys. Lett. B 718 (2013) 922)

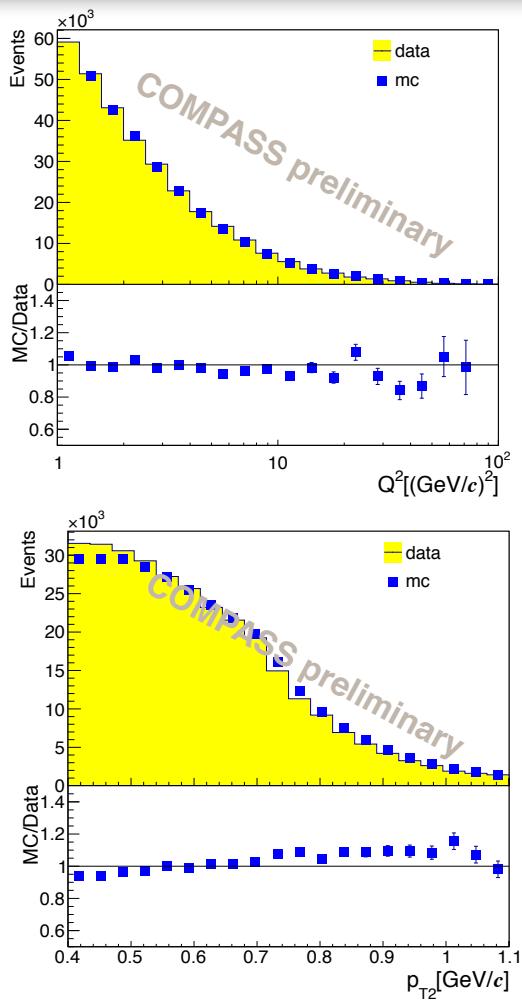
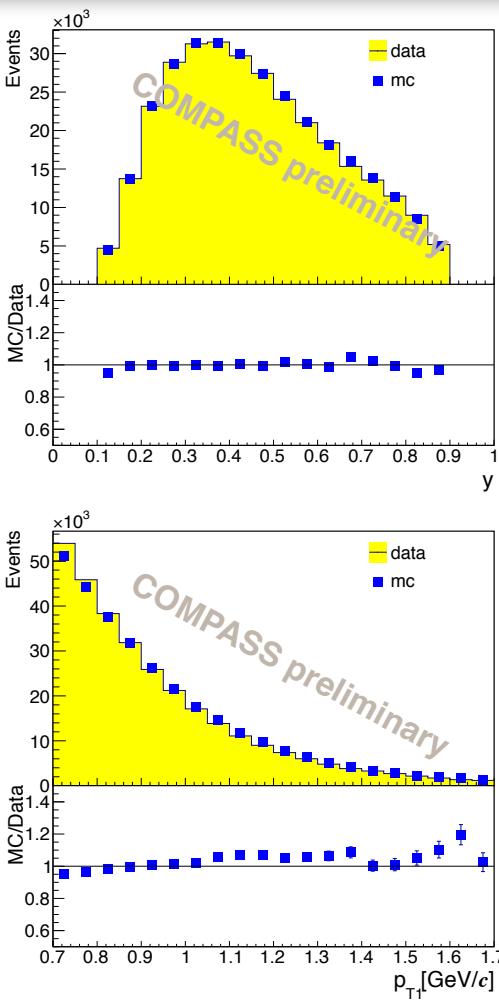
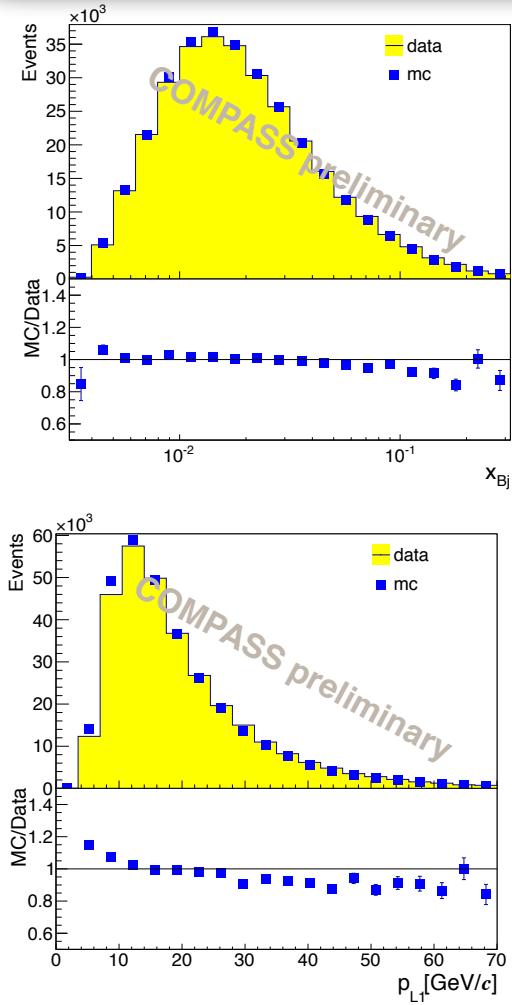
Artificial Neural Network approach



training vector: p_T, p_L, x_{Bjk}, Q^2

Data-MC comparison

Data and Monte Carlo comparison



good agreement (COMPASS tuning high- p_T)

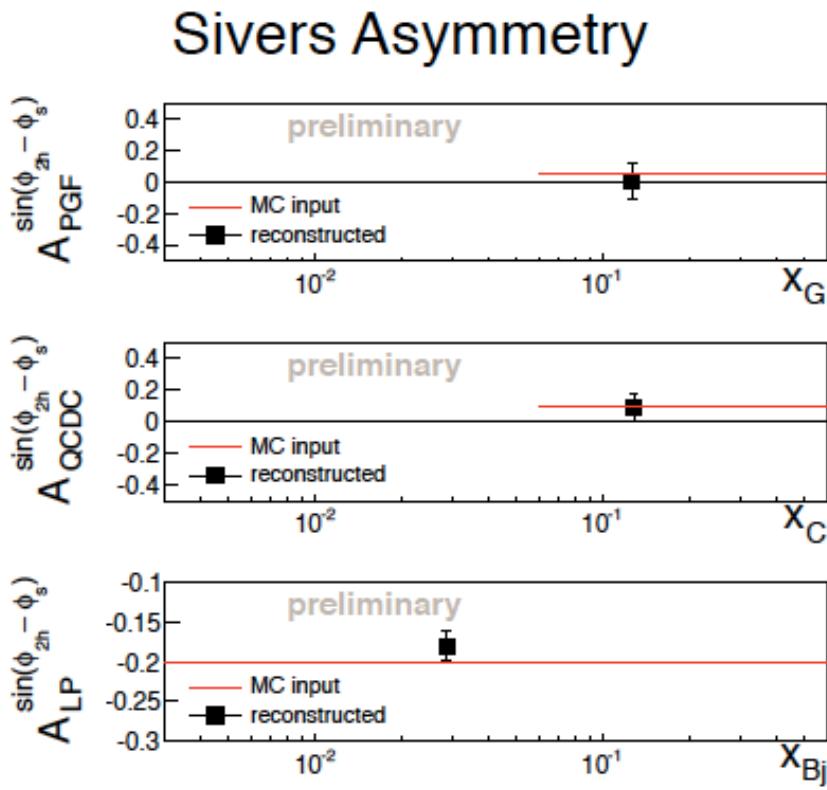
Validation of the method

- MC (LEPTO + COMGEANT, high- p_T tuning) events have no azimuthal asymmetries therefore we weight every event by $1+A\sin(\phi_{2h} - \phi_s)$.

ϕ_{2h} - azimuthal angle of the vector sum of the 2 leading hadron momenta.

A - assumed asymmetry for LO, QCDC and PGF

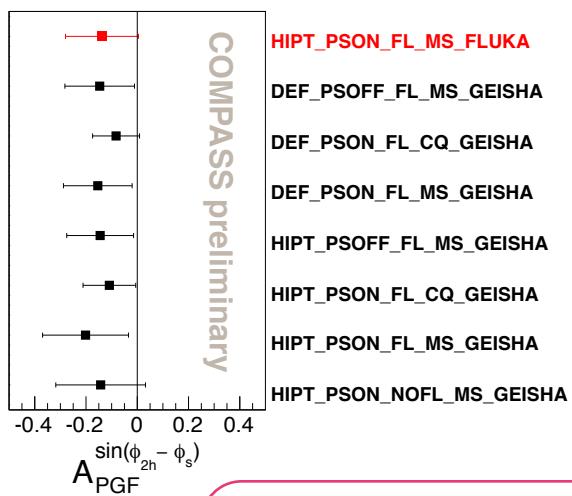
- For each MC event we get R_{LP} , R_{QCDC} , R_{PGF} and x_C , x_G from NN



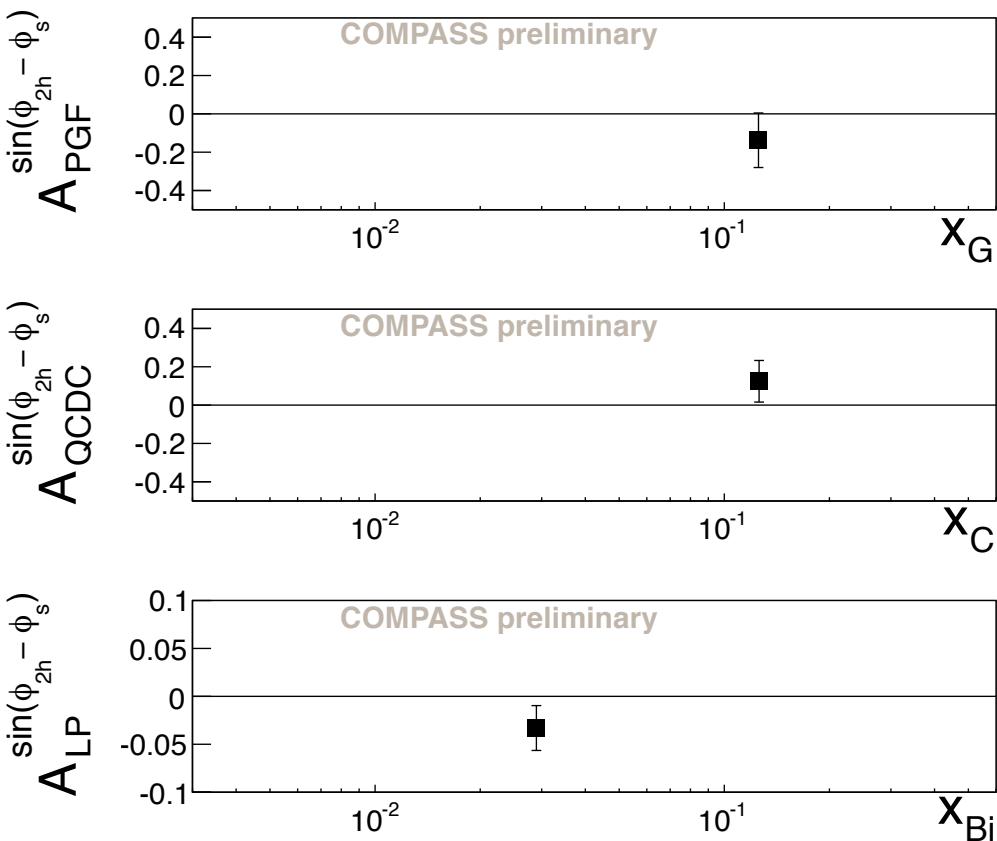
Data selection & preliminary results

COMPASS 2003-2004 data taken on deuteron polarised target

- Inclusive cuts:
 - $Q^2 > 1 \text{ (GeV}/c)^2$
 - $0.003 < x_{\text{Bj}} < 0.7$
 - $0.1 < y < 0.9$
- hadronic cuts
 - $p_{T1} > 0.7 \text{ GeV}/c$
 - $p_{T2} > 0.4 \text{ GeV}/c$
 - $z_1 > 0.1$
 - $z_2 > 0.1$



Sivers Asymmetry



$$A_{\text{PGF}}^{\sin(\phi_{2h} - \phi_s)} = 0.14 \pm 0.15(\text{stat.}) \pm 0.06(\text{syst.})$$

$$\langle x_G \rangle = 0.126.$$

- For the first time preliminary result for Sivers asymmetry for gluons has been obtained from COMPASS deuteron target data:

$$A_{\text{PGF}}^{\sin(\phi_{2h} - \phi_S)} = 0.14 \pm 0.15(\text{stat.}) \pm 0.06(\text{syst.})$$

The result is compatible with zero.

S.J.Brodsky & S. Gardner, Phys.Lett. B643 (2006) 22-28

- Analysis on proton data (more interesting, much larger statistics) is ongoing

Thank you for your attention



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Spin2014

The 21st International Symposium on Spin Physics, October 20-24, 2014, Beijing, China

Backup slides



Spares

SIDIS cross section decomposition LO

$$\begin{aligned}
 \frac{d\sigma}{dy d\phi_S d\phi_h dz_h dP_{hT}^2} = & \frac{e^4}{32\pi^2 x Q^2} \frac{y}{1-\varepsilon} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} \right. \\
 & + \sqrt{2\varepsilon(1-\varepsilon)} (\cos \phi_h F_{UU}^{\cos \phi_h} + h_l \sin \phi_h F_{LU}^{\sin \phi_h}) \\
 & + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_h F_{UL}^{\sin \phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \\
 & + S_{\parallel} h_l \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_h F_{LL}^{\cos \phi_h} \right] \\
 & + |S_{\perp}| \left[\sin(\phi_h - \phi_S) (F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)}) \right. \\
 & + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} \\
 & \left. + \sqrt{2\varepsilon(1+\varepsilon)} (\sin \phi_S F_{UT}^{\sin \phi_S} + \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)}) \right] \\
 & + |S_{\perp}| h_l \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} \right. \\
 & \left. + \sqrt{2\varepsilon(1-\varepsilon)} (\cos \phi_S F_{LT}^{\cos \phi_S} + \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)}) \right] \left. \right\},
 \end{aligned}$$

SIDIS cross section decomposition LO

$$\frac{d\sigma}{dy d\phi_S d\phi_h dz_h dP_{hT}^2} = \frac{e^4}{32\pi^2 x Q^2} \frac{y}{1-\varepsilon} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} \right. \\
 + \sqrt{2\varepsilon(1-\varepsilon)} (\cos \phi_h F_{UU}^{\cos \phi_h} + h_l \sin \phi_h F_{LU}^{\sin \phi_h}) \quad \text{unpolarised target} \\
 + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_h F_{UL}^{\sin \phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \\
 + S_{\parallel} h_l \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_h F_{LL}^{\cos \phi_h} \right] \\
 + |S_{\perp}| \left[\sin(\phi_h - \phi_S) (F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)}) \right. \\
 + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} \\
 + \sqrt{2\varepsilon(1+\varepsilon)} (\sin \phi_S F_{UT}^{\sin \phi_S} + \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)}) \left. \right] \\
 + |S_{\perp}| h_l \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} \right. \\
 \left. + \sqrt{2\varepsilon(1-\varepsilon)} (\cos \phi_S F_{LT}^{\cos \phi_S} + \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)}) \right] \}, \quad \text{Cahn & Boer-Mulders}$$

SIDIS cross section decomposition LO

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 + \sqrt{2\varepsilon(1-\varepsilon)} (\cos\phi_h F_{UU}^{\cos\phi_h} + h_l \sin\phi_h F_{LU}^{\sin\phi_h}) \quad \text{unpolarised target} \\
 + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \quad \text{longitudinally polarised target} \\
 + S_{\parallel} h_l \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \\
 + |S_{\perp}| \left[\sin(\phi_h - \phi_S) (F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)}) \right. \\
 + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} \\
 + \sqrt{2\varepsilon(1+\varepsilon)} (\sin\phi_S F_{UT}^{\sin\phi_S} + \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)}) \left. \right] \\
 + |S_{\perp}| h_l \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} \right. \\
 \left. + \sqrt{2\varepsilon(1-\varepsilon)} (\cos\phi_S F_{LT}^{\cos\phi_S} + \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)}) \right] \right\},$$

Cahn & Boer-Mulders

longitudinaly polarised target

Spares

SIDIS cross section decomposition LO

Spares

$$\frac{d\sigma}{dy d\phi_S d\phi_h dz_h dP_{hT}^2} = \frac{e^4}{32\pi^2 x Q^2} \frac{y}{1-\varepsilon} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} \right. \\
 + \sqrt{2\varepsilon(1-\varepsilon)} (\cos \phi_h F_{UU}^{\cos \phi_h} + h_l \sin \phi_h F_{LU}^{\sin \phi_h}) \quad \text{unpolarised target} \\
 + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_h F_{UL}^{\sin \phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \quad \text{longitudinally polarised target} \\
 + S_{\parallel} h_l \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_h F_{LL}^{\cos \phi_h} \right] \\
 + |S_{\perp}| \left[\sin(\phi_h - \phi_S) (F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)}) \right. \\
 + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} \\
 + \sqrt{2\varepsilon(1+\varepsilon)} (\sin \phi_S F_{UT}^{\sin \phi_S} + \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)}) \left. \right] \quad \text{transversely polarised target} \\
 + |S_{\perp}| h_l \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} \right. \\
 \left. + \sqrt{2\varepsilon(1-\varepsilon)} (\cos \phi_S F_{LT}^{\cos \phi_S} + \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)}) \right] \}, \\$$

Cahn & Boer-Mulders
 Sivers
 Collins
 Pretzelosity
 Worm Gear

SIDIS cross section decomposition LO

$$\frac{d\sigma}{dy d\phi_S d\phi_h dz_h dP_{hT}^2} = \frac{e^4}{32\pi^2 x Q^2} \frac{y}{1-\varepsilon} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} \right.$$

Cahn & Boer-Mulders

$$+ \sqrt{2\varepsilon(1-\varepsilon)} (\cos\phi_h F_{UU}^{\cos\phi_h} + h_l \sin\phi_h F_{LU}^{\sin\phi_h}) \quad \text{unpolarised target}$$

$$\left. + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \right\}$$

Spares

$$F_{UU,T} \sim \sum_q e_q^2 \cdot f_1^q \otimes D_q^h,$$

$$F_{LT}^{\cos(\phi_h - \phi_S)} \sim \sum_q e_q^2 \cdot g_{1T}^q \otimes D_q^h,$$

$$F_{LL} \sim \sum_q e_q^2 \cdot g_{1L}^q \otimes D_q^h,$$

$$F_{UT,T}^{\sin(\phi_h - \phi_S)} \sim \sum_q e_q^2 \cdot f_{1T}^{\perp q} \otimes D_q^h,$$

$$F_{UU}^{\cos 2\phi_h} \sim \sum_q e_q^2 \cdot h_1^{\perp q} \otimes H_1^{\perp q},$$

$$F_{UT}^{\sin(\phi_h + \phi_S)} \sim \sum_q e_q^2 \cdot h_1^q \otimes H_1^{\perp q},$$

$$F_{UL}^{\sin 2\phi_h} \sim \sum_q e_q^2 \cdot h_{1L}^{\perp q} \otimes H_1^{\perp q},$$

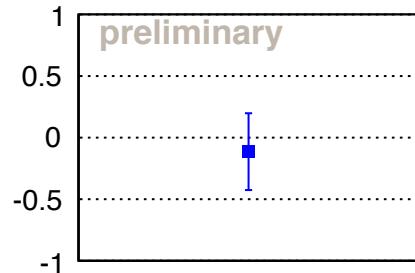
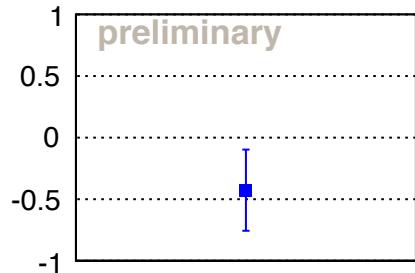
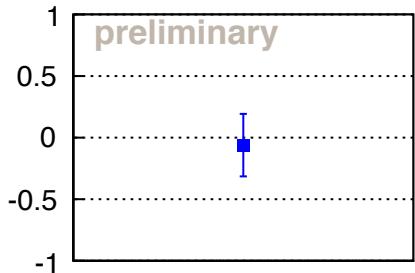
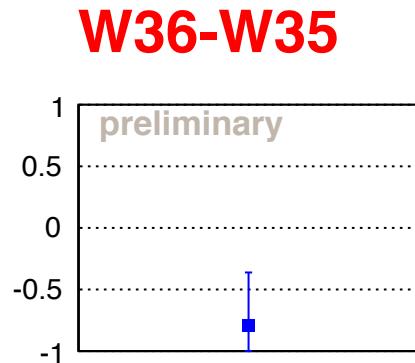
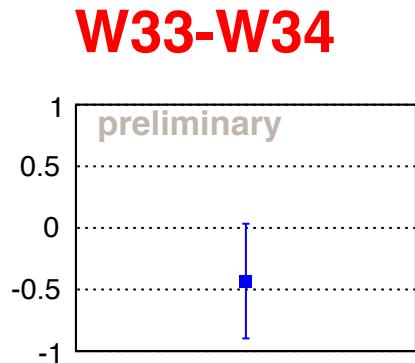
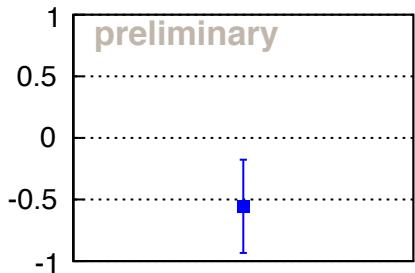
$$F_{UT}^{\sin(3\phi_h - \phi_S)} \sim \sum_q e_q^2 \cdot h_{1T}^{\perp q} \otimes H_1^{\perp q}.$$

get

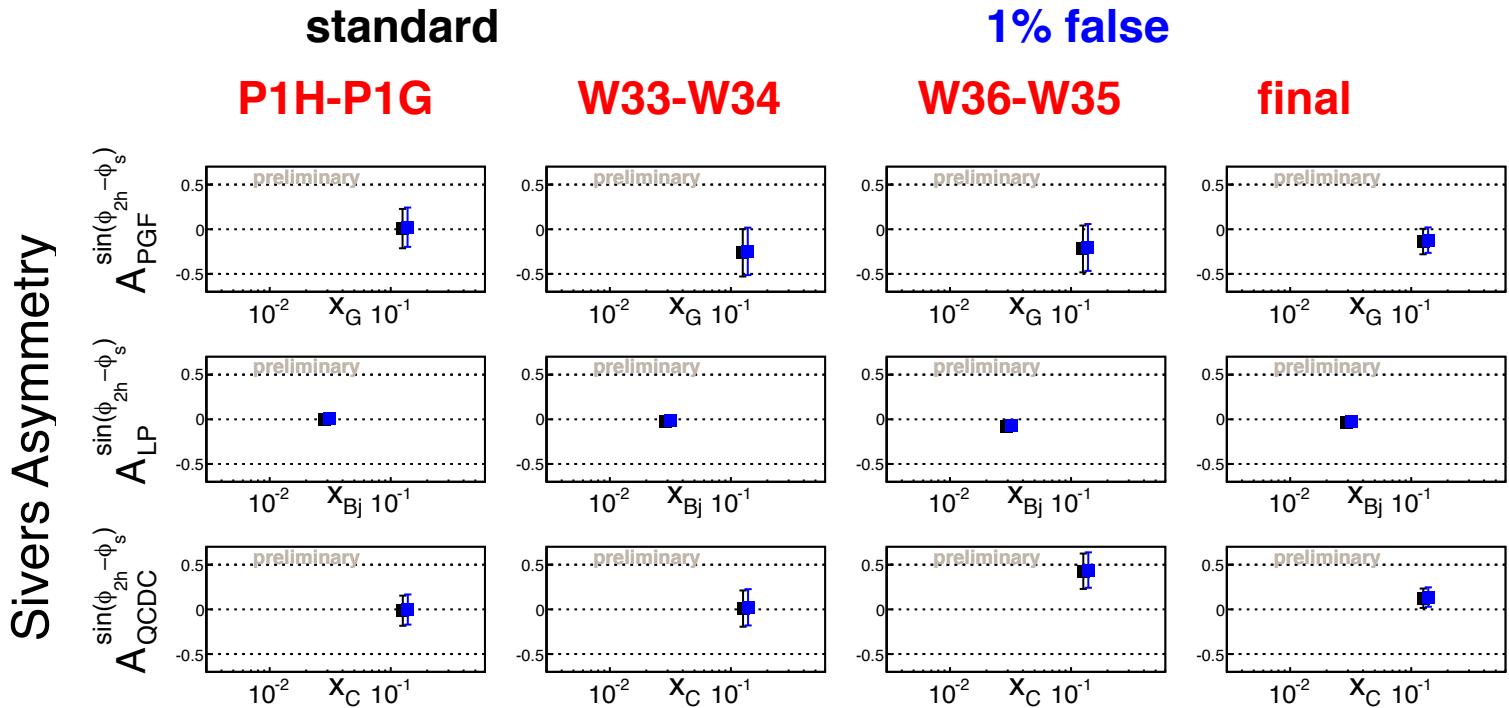
Falses



False PGF Sivers Asymmetry
Spares
up



1% false test for asymmetry



Correlations

