

Measurement of TMD observables at COMPASS

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**COMmmon
Muon and
Proton
Apparatus for
Structure and
Spectroscopy**

Collaboration
~ 250 physicists
28 institutions
12 countries

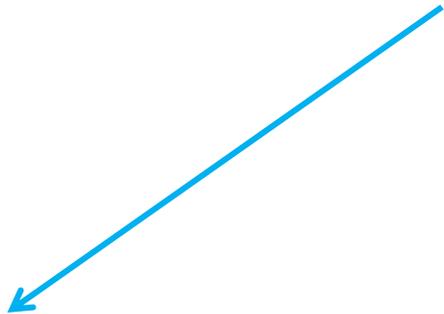


fixed target experiment at the CERN SPS



Common Muon and Proton Apparatus for Structure and Spectroscopy

wide physics program carried on using both **muon** and hadron beam



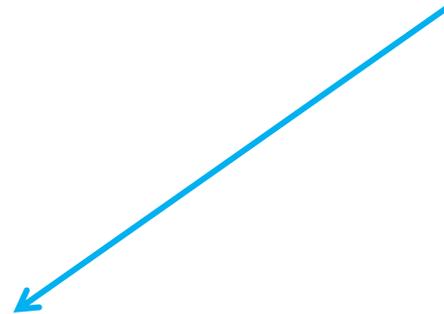
luminosity: $\sim 5 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
 beam intensity: $2 \cdot 10^8 \mu^+/\text{spill}$ (4.8s/16.2s)
 beam momentum: 160 GeV/c

longitudinally polarized muon beam	deuteron (${}^6\text{LiD}$)	2002	} L/T	hadron beam	nuclear targets	2004
	polarized target	2003				
		2004				
		2006	L		LH target	2008
	proton (NH_3)	2007	L/T	2009		
	polarized target	2010	T		2012	
		2011	L			
	H ₂ target	2012				



Common Muon and Proton Apparatus for Structure and Spectroscopy

wide physics program carried on
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longitudinally
polarized
muon beam

deuteron (${}^6\text{LiD}$)
polarized target

proton (NH_3)
polarized target

H_2 target



*in this talk
the results on the
TMDs measured from SIDIS
are presented*

SIDIS: a key process to investigate the structure of the nucleon

lepton interacts with a **single constituent** of the nucleon ($Q^2 > 1 \text{ GeV}^2/c^2$)

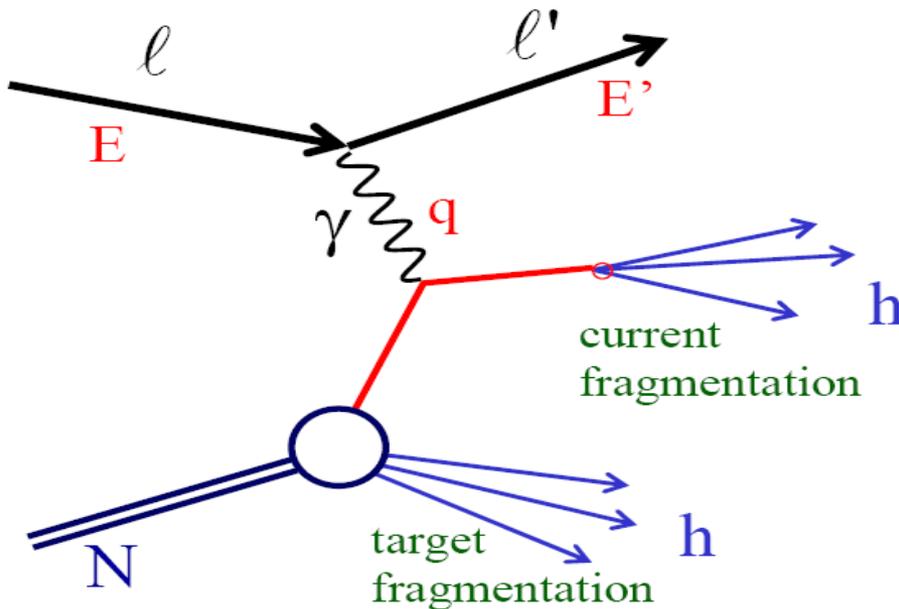
$$q = \ell - \ell'$$

$$Q^2 = -q^2 \quad W^2 = (P + q)^2$$

$$x = \frac{Q^2}{2P \cdot q} \quad \text{Bjorken scaling variable}$$

$$y = \frac{P \cdot q}{P \cdot \ell} = \frac{E - E'}{E}$$

$$z = \frac{P \cdot P_h}{P \cdot q} = \frac{E_h}{E - E'}$$



at least one hadron is detected
in the final state
(information on the **struck quark**)

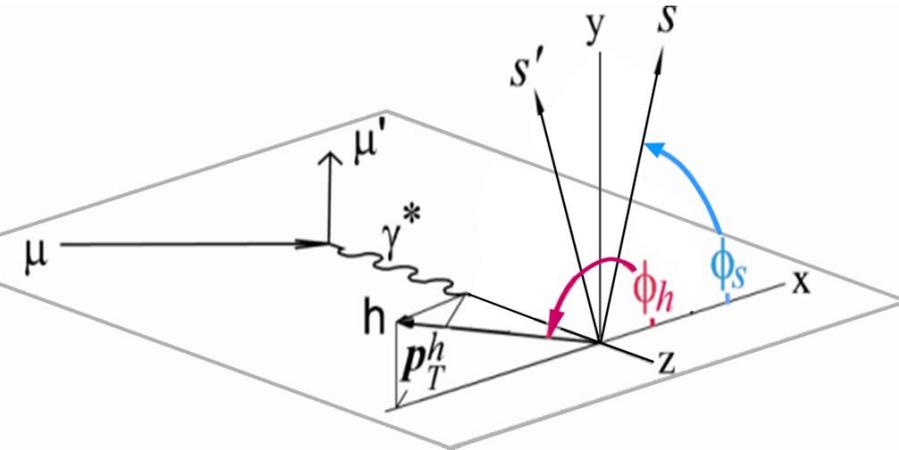
The complete **SIDIS cross section** (one photon exchange approximation)

$$\begin{aligned}
 \frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = & \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\
 & + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \\
 & + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \\
 & + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
 & + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) \\
 & + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(\phi_h - \phi_S) \\
 & \left. + |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\},
 \end{aligned}$$

14 **independent modulations** in the **azimuthal angle** of the **spin ϕ_S** and of the **hadron ϕ_h** (around the virtual photon direction)

target nucleon polarisation

beam polarisation



transverse momentum of the hadron \mathbf{p}_T^h

The complete **SIDIS cross section** (one photon exchange approximation)

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 & + \varepsilon \cos(2\phi_h) \boxed{F_{UU}^{\cos 2\phi_h}} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h \boxed{F_{LU}^{\sin\phi_h}} \\
 & + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h \boxed{F_{UL}^{\sin\phi_h}} + \varepsilon \sin(2\phi_h) \boxed{F_{UL}^{\sin 2\phi_h}} \right] + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h \boxed{F_{LL}^{\cos\phi_h}} \right] \\
 & + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(\boxed{F_{UT,T}^{\sin(\phi_h - \phi_S)}} + \varepsilon \boxed{F_{UT,L}^{\sin(\phi_h - \phi_S)}} \right) \right. \\
 & + \varepsilon \sin(\phi_h + \phi_S) \boxed{F_{UT}^{\sin(\phi_h + \phi_S)}} + \varepsilon \sin(3\phi_h - \phi_S) \boxed{F_{UT}^{\sin(3\phi_h - \phi_S)}} \\
 & + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S \boxed{F_{UT}^{\sin\phi_S}} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) \boxed{F_{UT}^{\sin(2\phi_h - \phi_S)}} \\
 & + |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) \boxed{F_{LT}^{\cos(\phi_h - \phi_S)}} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S \boxed{F_{LT}^{\cos\phi_S}} \right. \\
 & \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) \boxed{F_{LT}^{\cos(2\phi_h - \phi_S)}} \right] \left. \right\},
 \end{aligned}$$

14 **independent modulations**
in the **azimuthal angle**
of the **spin ϕ_S**
and of the **hadron ϕ_h**
(around the virtual photon direction)



each **modulation** is coupled
with a **Structure Function F**

The Structure Function F is the convolution over the

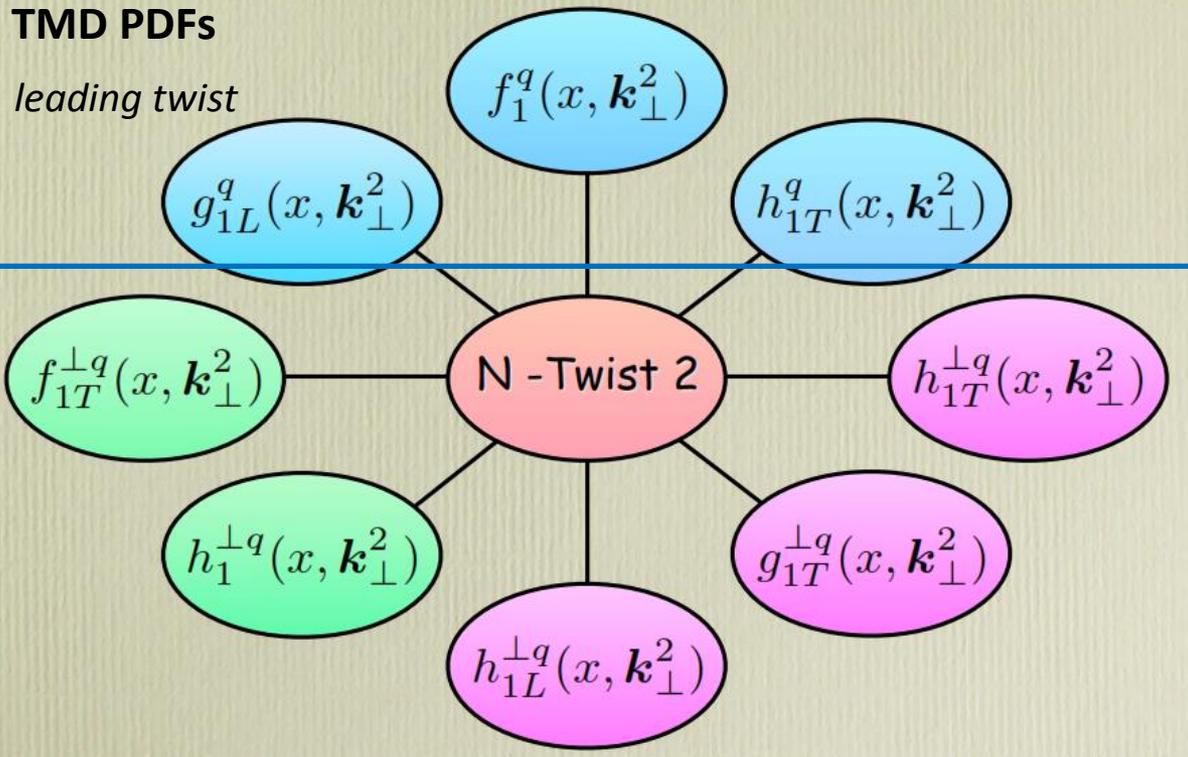
transverse momenta of a PDF and a FF

describing the fragmentation of the quark in an unpolarised hadron

$$F = \sum_q PDF_q(x) \otimes FF_q^h(z)$$

TMD PDFs

leading twist



survive after integration over k_{\perp}

correlation between spin and transverse momentum

all of the amplitudes of the azimuthal modulations of the polarised SIDIS cross section have been measured at COMPASS



the COMPASS spectrometer



- high energy beams
- large angular acceptance
- broad kinematical range

two stages spectrometer

Large Angle Spectrometer (SM1)

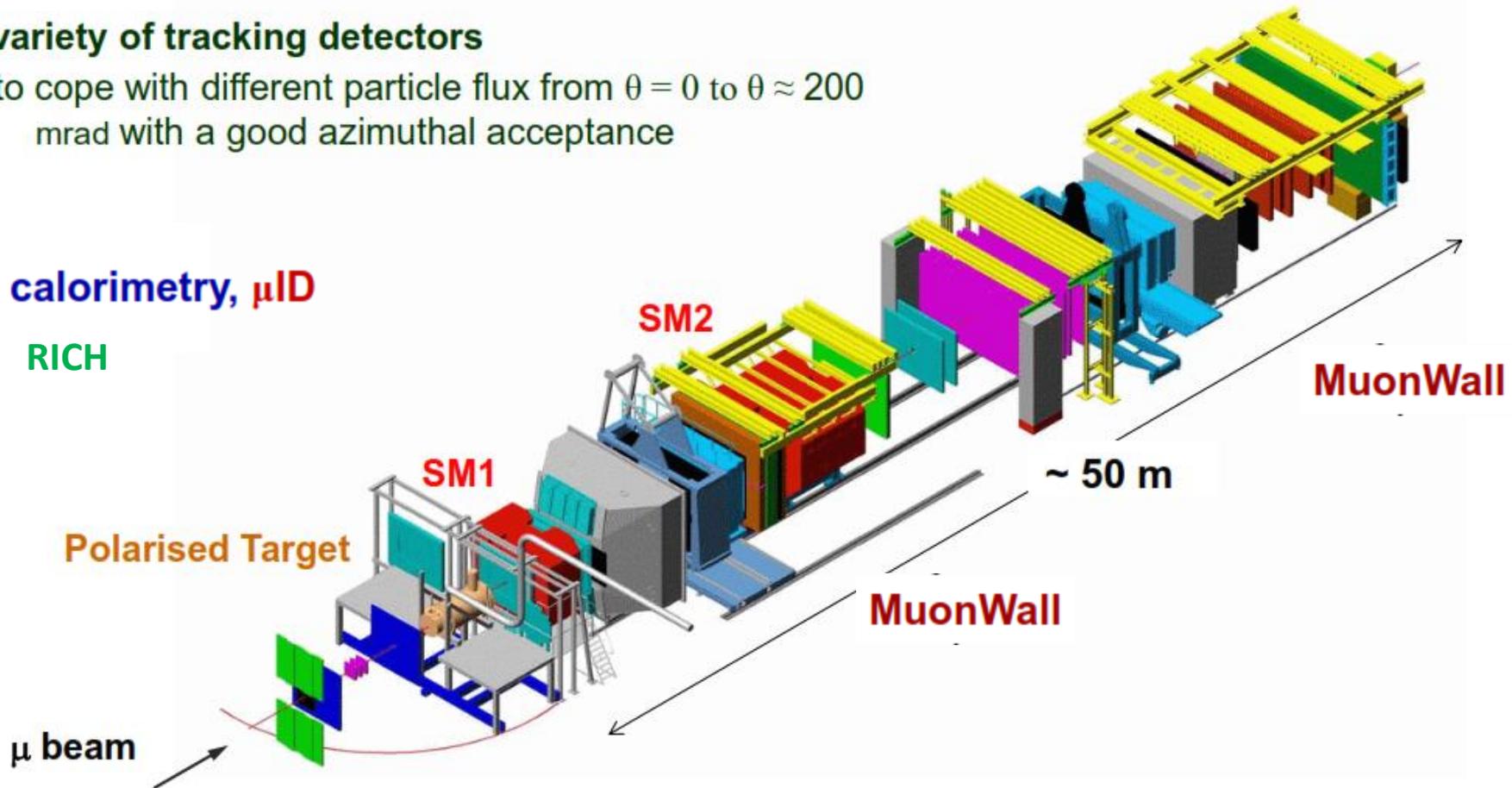
Small Angle Spectrometer (SM2)

variety of tracking detectors

to cope with different particle flux from $\theta = 0$ to $\theta \approx 200$ mrad with a good azimuthal acceptance

calorimetry, μ ID

RICH



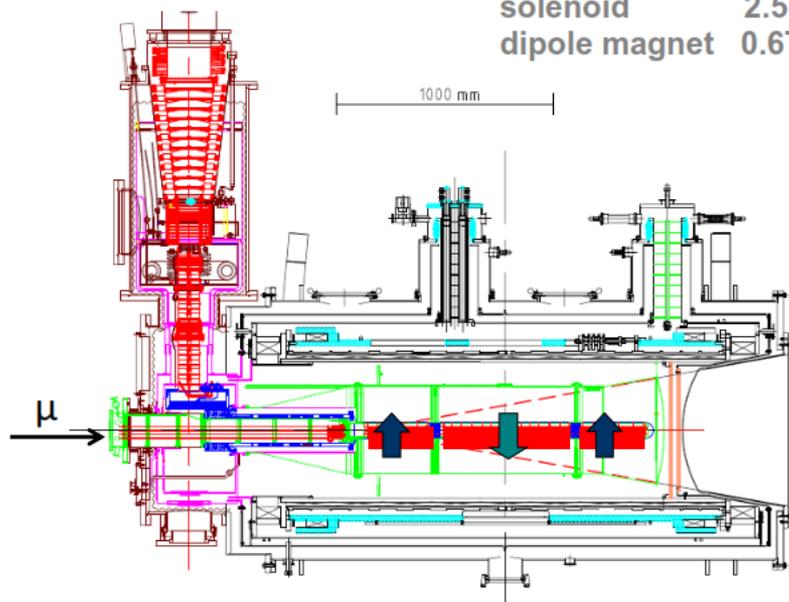
polarized target system (>2005)

solid state target operating in frozen spin mode

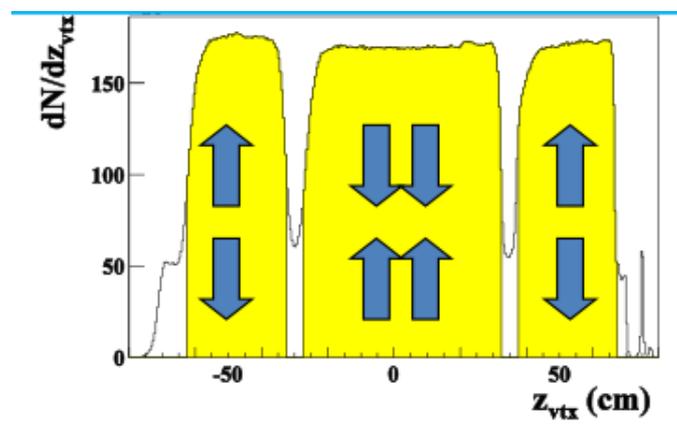
	d (⁶LiD)	p (NH₃)
polarization	50%	90%
dilution factor	40%	16%

³He – ⁴He dilution refrigerator (T~50mK)

solenoid 2.5T
dipole magnet 0.6T



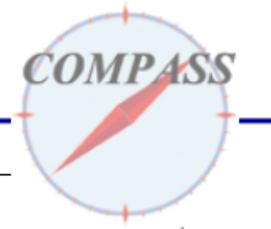
3 cells target with opposite polarizations



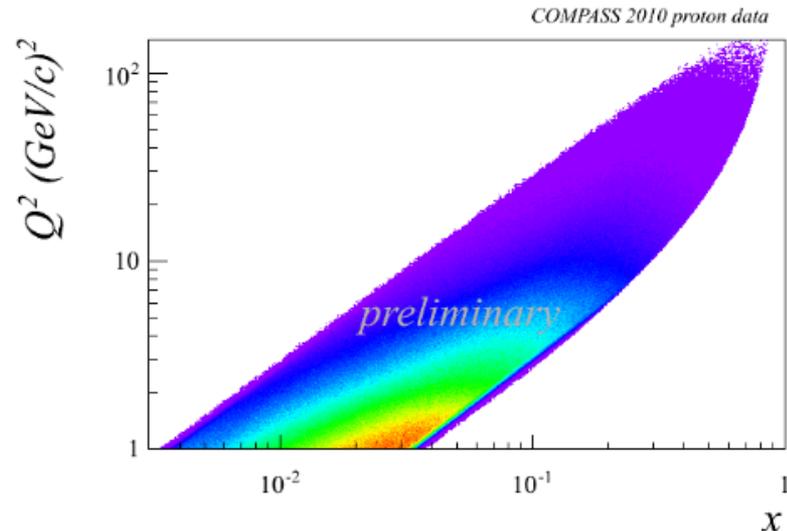
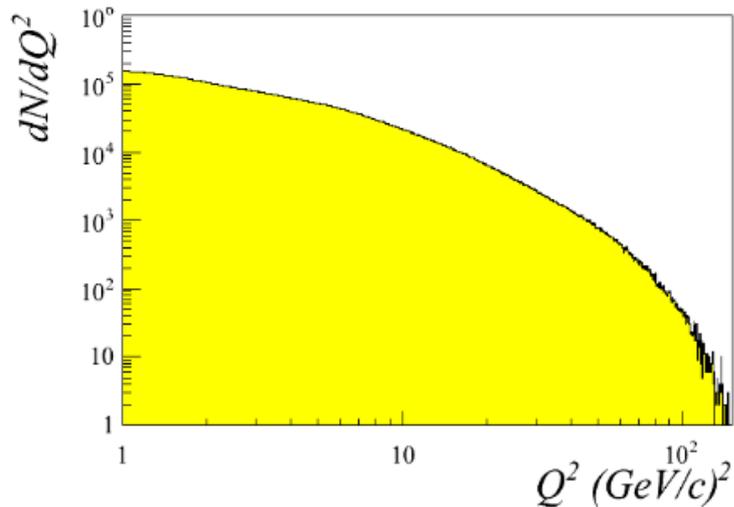
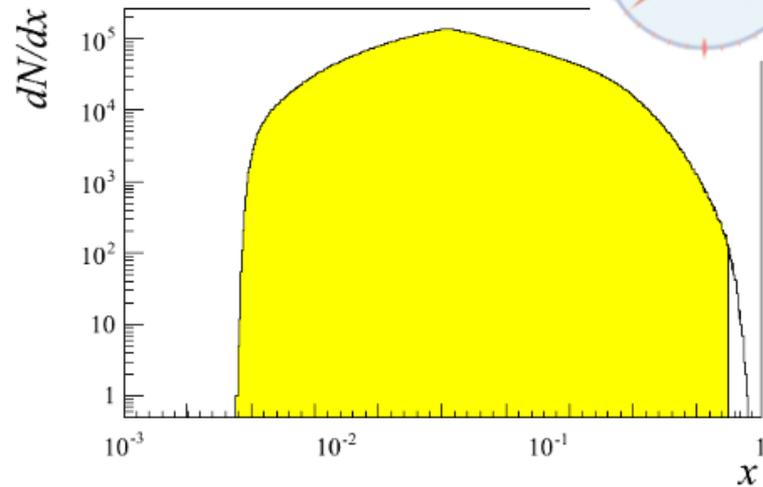
**2 configurations:
polarisation reversed each week
to minimize possible systematic errors**

SIDIS event selection

$p_\mu = 160 \text{ GeV}/c$

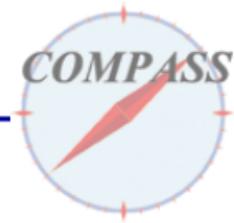


DIS cuts: $Q^2 > 1 \text{ (GeV}/c)^2$
 $0.1 < y < 0.9$
 $W > 5 \text{ GeV}/c^2$



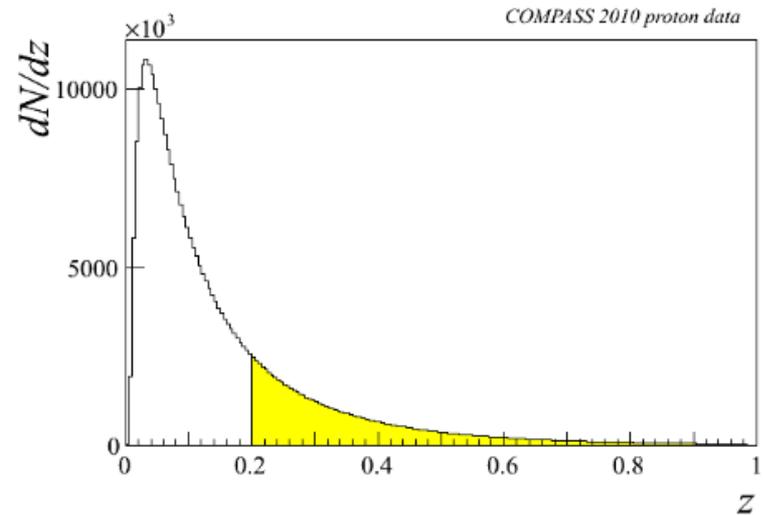
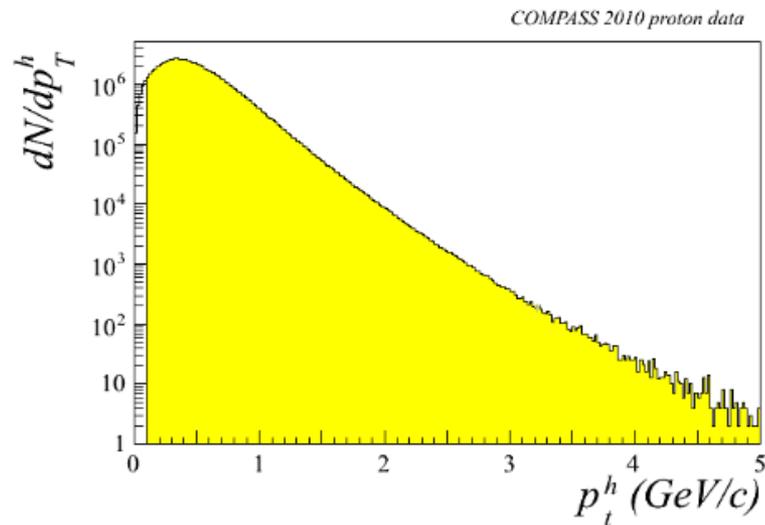
SIDIS event selection

$p_\mu = 160 \text{ GeV}/c$



DIS cuts: $Q^2 > 1 \text{ (GeV}/c)^2$
 $0.1 < y < 0.9$
 $W > 5 \text{ GeV}/c^2$

h^\pm selection: $p_t^h > 0.1 \text{ GeV}/c$
 $z > 0.2$

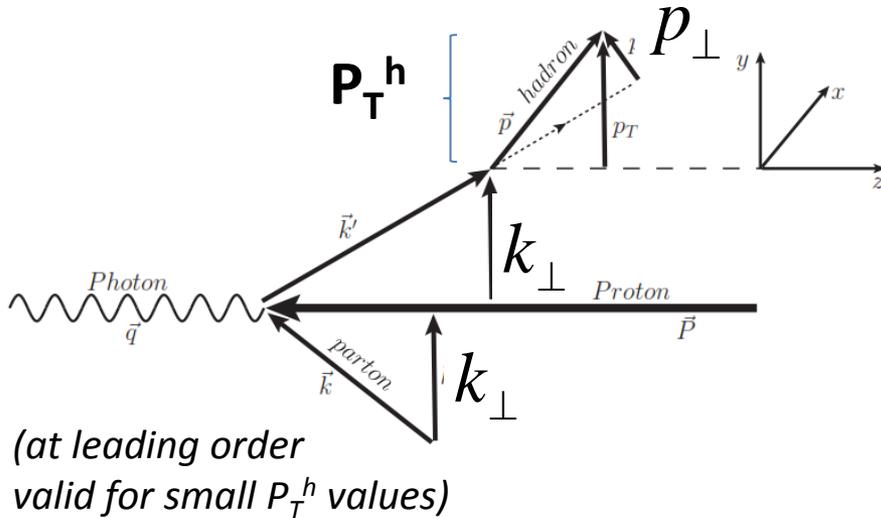


TM effects measurements

shown in this presentation

- hadron multiplicities in SIDIS off unpolarised target
- azimuthal asymmetries in SIDIS off unpolarised target
- azimuthal asymmetries in SIDIS off transversely polarised target

hadron multiplicities as function of $(P_T^h)^2$



the \mathbf{P}_T^h of the detected hadron is given by the **intrinsic TM of the quark** inside the nucleon k_\perp and the TM of the quark in the fragmentation p_\perp

differential SIDIS cross-section / differential DIS cross-section

$$\frac{d^2 n^{h^\pm}(x, Q^2, z, p_T^{h^2})}{dz dp_T^{h^2}} \approx \frac{\Delta^4 N^{h^\pm}(x, Q^2, z, p_T^{h^2}) / (\Delta x \Delta Q^2 \Delta z \Delta p_T^{h^2})}{\Delta^2 N^{DIS}(x, Q^2) / (\Delta x \Delta Q^2)}$$

hadron multiplicities as function of $(P_T^h)^2$

SIDIS events from 2004 ^6LiD data

multiplicities measured
in 4 dimensions:

23 x, Q^2 intervals

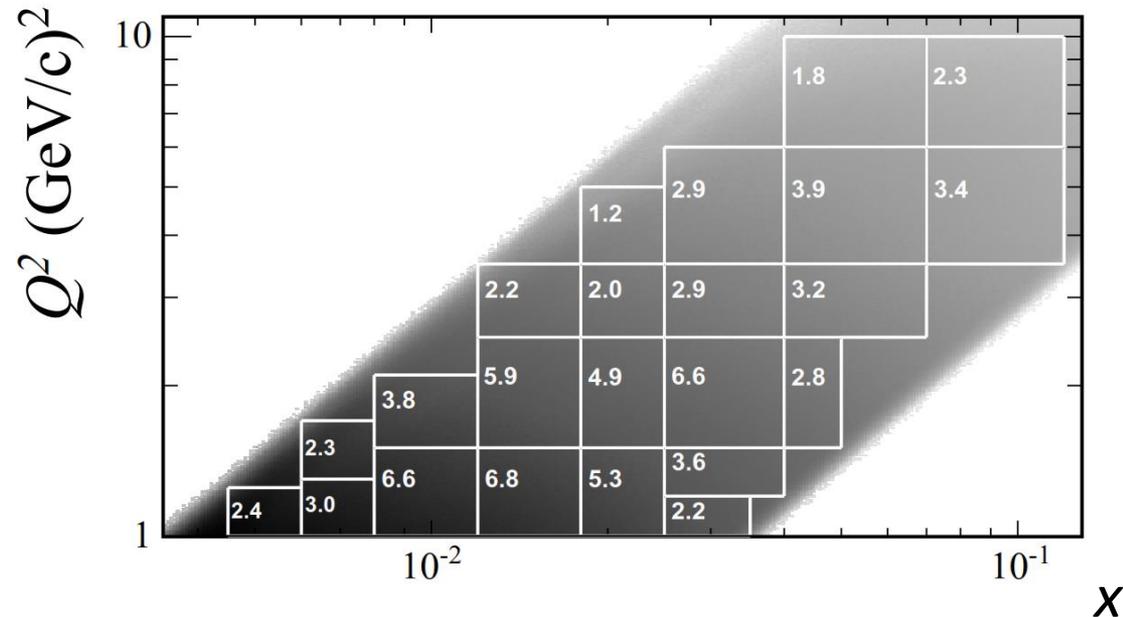
$1 < Q^2 < 10$ (GeV/c) 2

$0.004 < x < 0.12$

8x40

z and P_T^h
intervals

45.8×10^6 DIS events

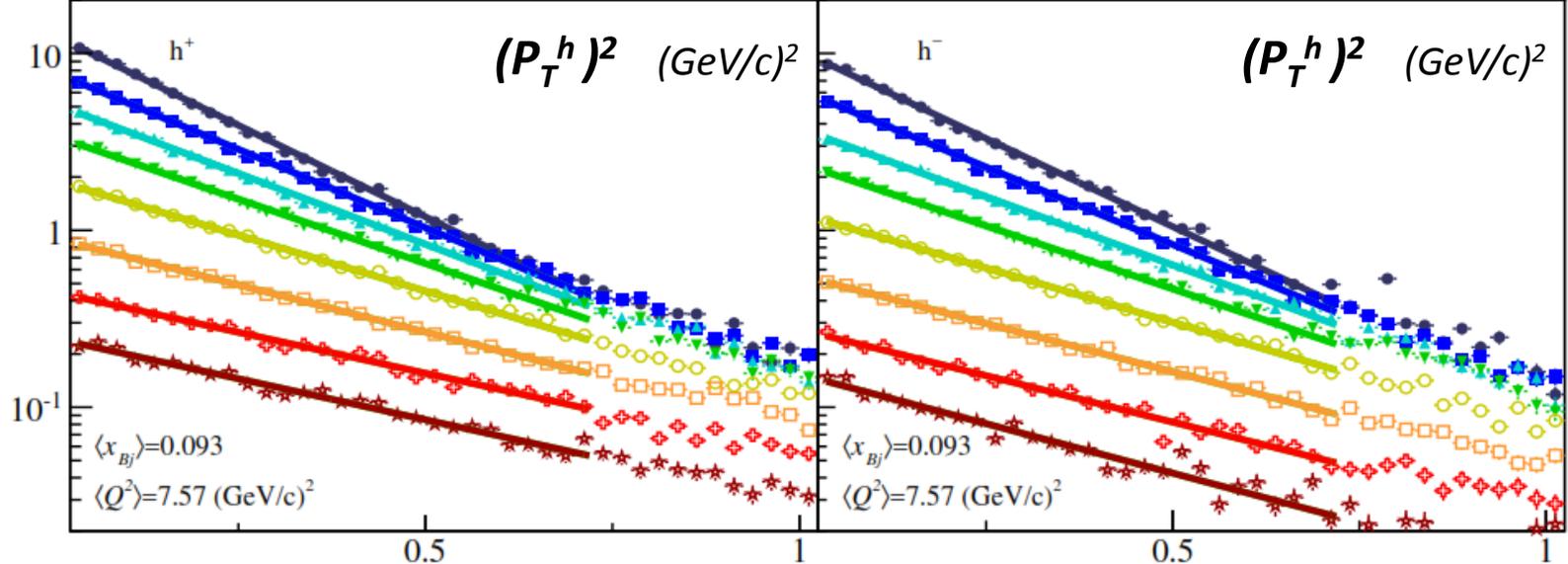
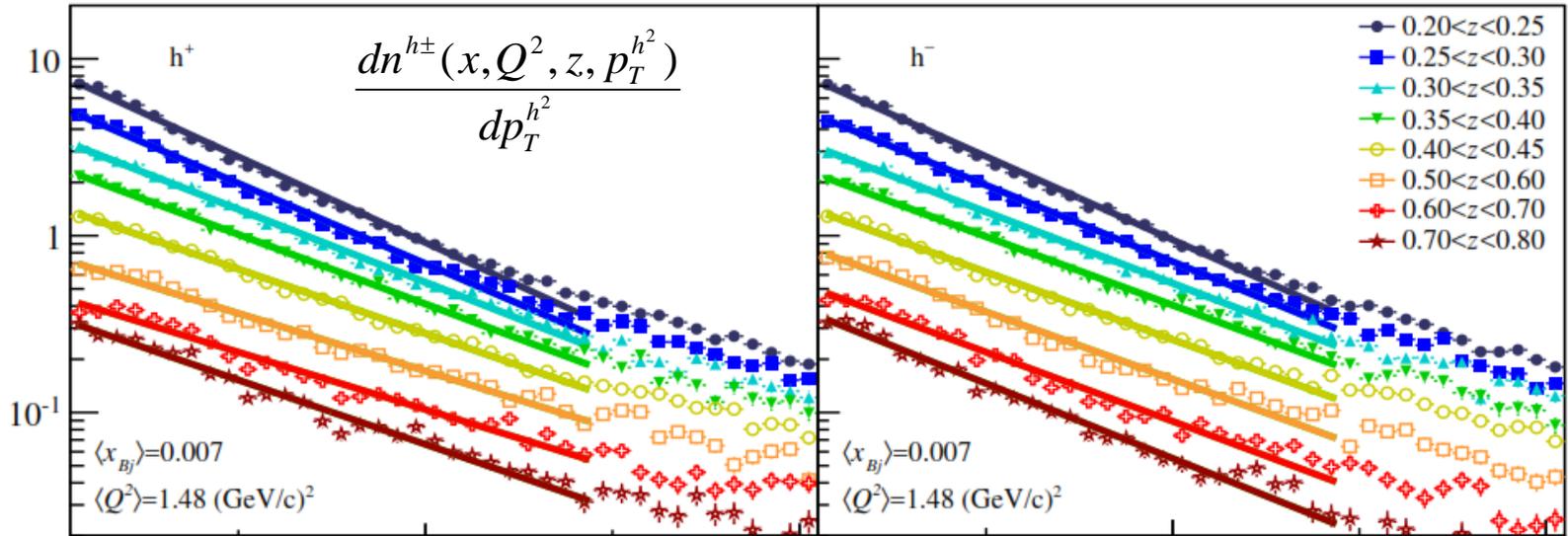


separately for h^+ and h^-

for each bin the “raw” multiplicity is **corrected for the acceptance**
calculated using **Monte Carlo simulations**

5% systematic uncertainties

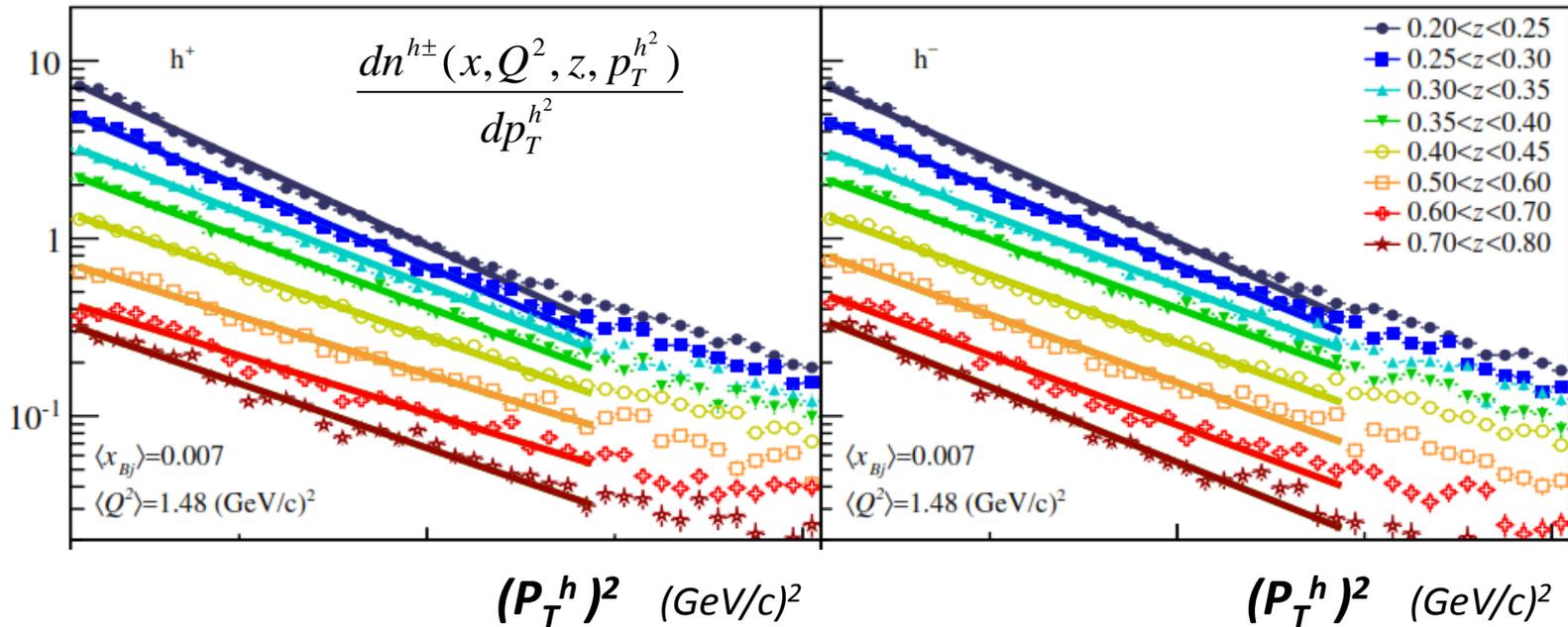
hadron multiplicities as function of $(P_T^h)^2$



$(P_T^h)^2 (GeV/c)^2$

$(P_T^h)^2 (GeV/c)^2$

hadron multiplicities as function of $(P_T^h)^2$



for each bin in \mathbf{x} , Q^2 and \mathbf{z}

fit using

$$A \cdot e^{-P_T^{h^2} / \langle P_T^{h^2} \rangle}$$

→ extract

$$\langle P_T^{h^2} \rangle = \langle P_{\perp}^2 \rangle + z^2 \langle k_{\perp}^2 \rangle$$

for $P_T^h < 0.85 \text{ GeV/c}$

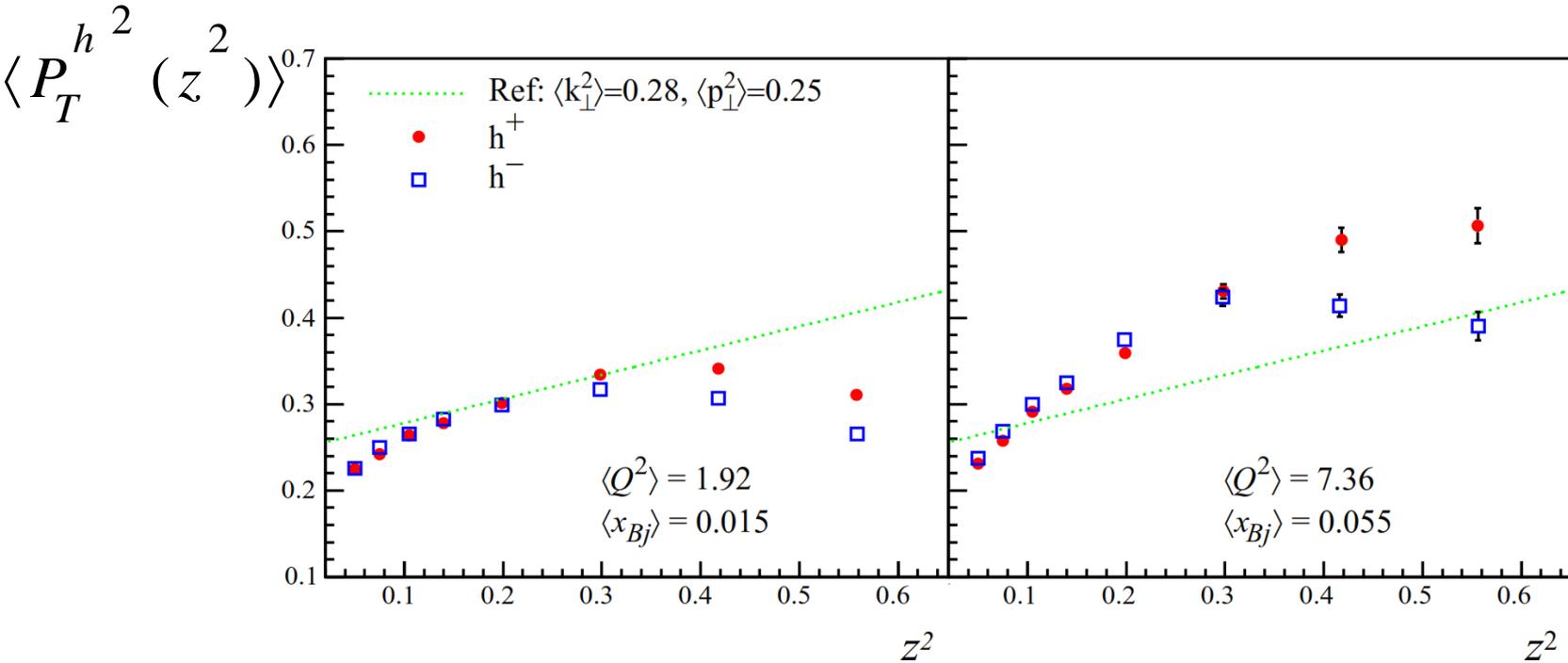
to stay away from the region

where transverse momentum

effects related to gluon radiation

become relevant

z dependence



$$\langle P_T^2(z) \rangle = \langle P_{\perp}^2(z) \rangle + z^2 \langle k_{\perp}^2 \rangle$$

this term should be calculated from DIS generator
using the best knowledge on jet fragmentation

not so easy to reproduce the results...

- hadron multiplicities in SIDIS off unpolarised target
- azimuthal asymmetries in SIDIS off unpolarised target
- azimuthal asymmetries in SIDIS off transversely polarised target

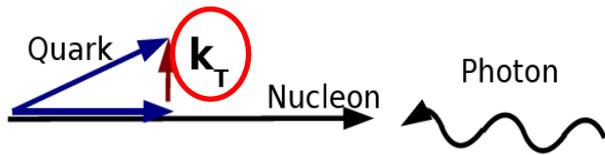
The azimuthal hadron distribution for the unpolarised target is:

$$N(\phi_h) \propto N_0 \cdot (1 + \varepsilon_1 A_{\cos \phi_h}^{UU} \cos \phi_h + \varepsilon_2 A_{\cos 2\phi_h}^{UU} \cos 2\phi_h + \lambda_l \varepsilon_3 A_{\sin \phi_h}^{LU} \sin \phi_h)$$

$$\varepsilon_1 = \frac{2(2-y)\sqrt{1-y}}{1+(1-y)^2} \quad \varepsilon_2 = \frac{2(1-y)}{1+(1-y)^2} \quad \varepsilon_3 = \frac{2y\sqrt{1-y}}{1+(1-y)^2}$$



mainly **Cahn** effect: kinematical effect proportional to the **quark transverse momentum**



$$d\sigma^{lq \rightarrow lq} \propto \hat{s}^2 + \hat{u}^2 \propto \left(1 + \varepsilon_1 \frac{k_{\perp}}{Q} \cos \varphi \right)$$

Boer-Mulders (*T-odd* !) function, one of the most famous **TMD PDF**, convoluted with the **Collins FF**



the **Boer-Mulders** function correlates the **quark transverse momentum** and the **quark spin** in an **unpolarized nucleon**

fundamental prediction pQCD
sign change between B-M TMD as measured in SIDIS and in Drell-Yan

higher twist effect proportional to beam polarization λ_l
no clear interpretation in terms of PM

The azimuthal hadron distribution for the unpolarised target is:

$$N(\phi_h) \propto N_0 \cdot (1 + \varepsilon_1 A_{\cos \phi_h}^{UU} \cos \phi_h + \varepsilon_2 A_{\cos 2\phi_h}^{UU} \cos 2\phi_h + \lambda_l \varepsilon_3 A_{\sin \phi_h}^{LU} \sin \phi_h)$$

(SIDIS events from 2004 ${}^6\text{LiD}$ data)

separately for h+ and h-
in **each bin** of x, z and P_T^h :

$Q^2 > 1 \text{ (GeV/c)}^2$
$\theta_y^{\text{lab}} < 0.06$
$0.003 < x < 0.13$
$0.2 < y < 0.9$
$W > 5 \text{ GeV/c}^2$
$0.2 < z < 0.85$
$0.1 < P_T^h < 1 \text{ GeV/c}$

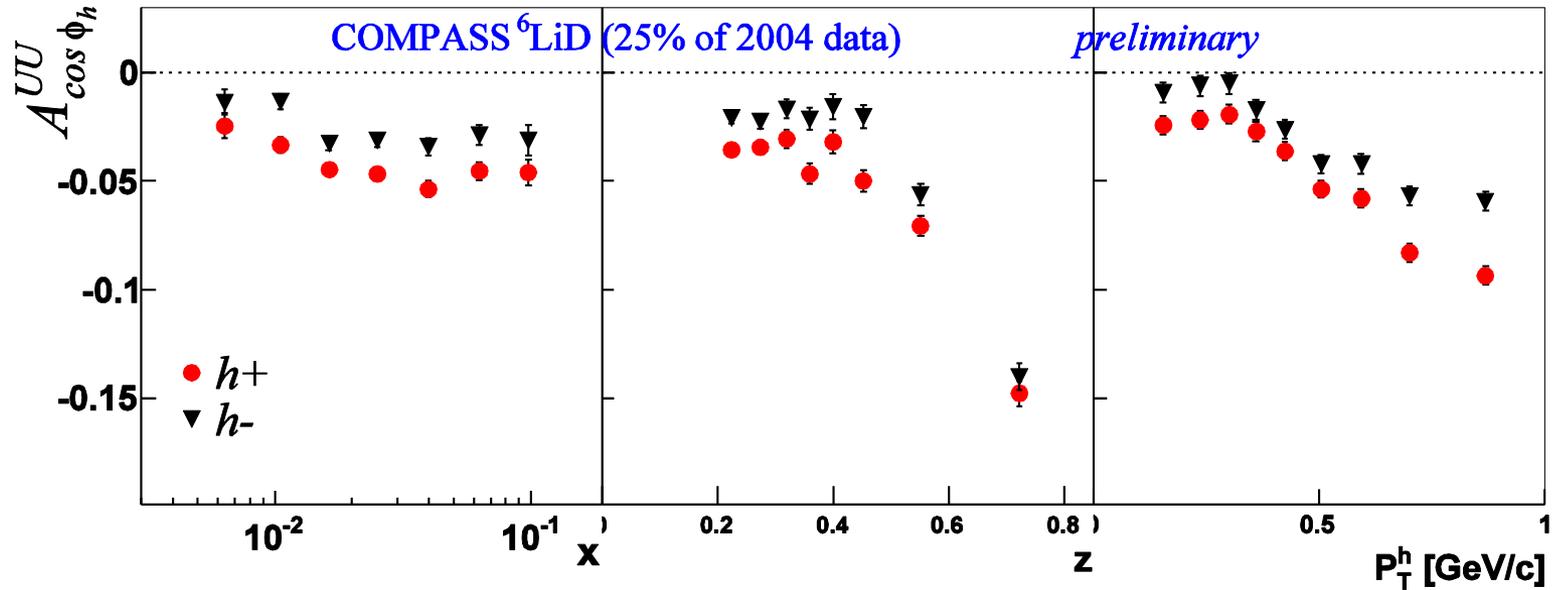
the **acceptance is calculated**, using **Monte Carlo** simulations

the “raw” **azimuthal distribution is corrected** for the apparatus **azimuthal acceptance**

$$A_{\cos \phi_h}^{UU}, \quad A_{\cos 2\phi_h}^{UU}, \quad A_{\sin \phi_h}^{LU}$$

are **extracted fitting the hadron distribution in ϕ_h**

Results extracted binning alternatively in x, z and P_T^h



$sys \approx 2 \cdot stat$

$$A_{\cos\phi_h}^{UU} = \frac{1}{Q} \text{Cahn} + \frac{1}{Q} \text{BM}$$

strong z dependence, for $z > 0.5$

to better understand the interesting and unexpected kinematical dependencies found



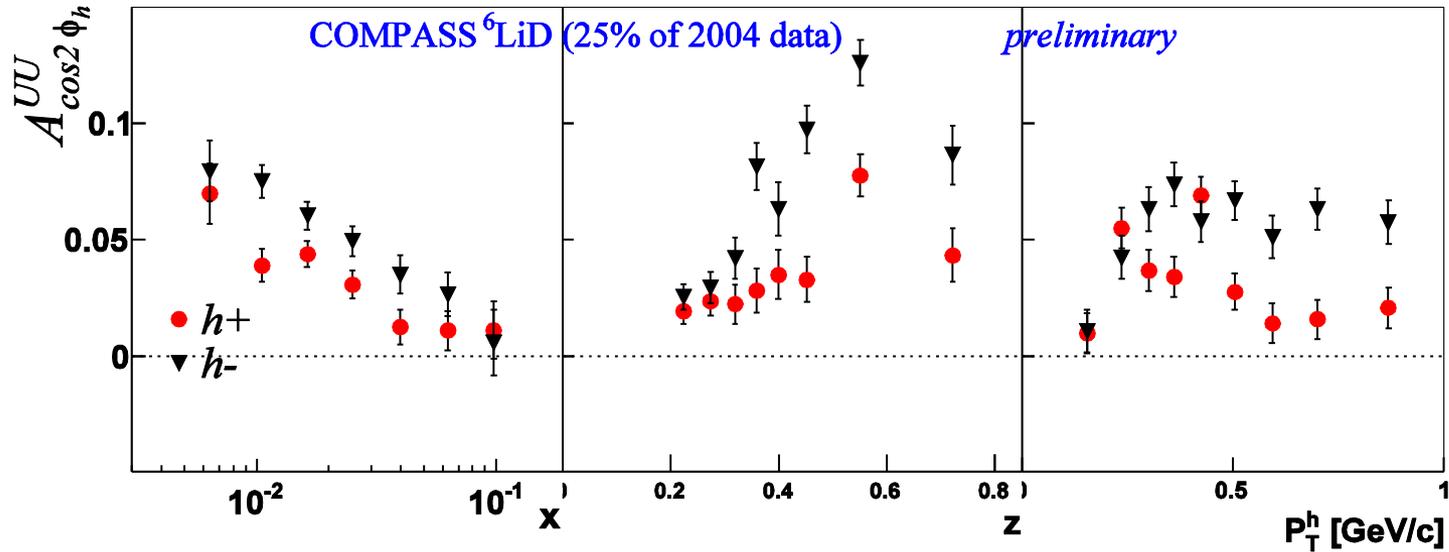
a multi dimensional analysis has been done

binning simultaneously in x , z and P_T^h

x	P_T^h	z
0.003 - 0.012	0.1 - 0.3	0.2 - 0.25
0.012 - 0.02	0.3 - 0.5	0.25 - 0.32
0.02 - 0.038	0.5 - 0.64	0.32 - 0.40
0.038 - 0.13	0.64 - 1.0	0.40 - 0.55
		0.55 - 0.70
		0.70 - 0.85

results shown at SPIN2012 in Dubna

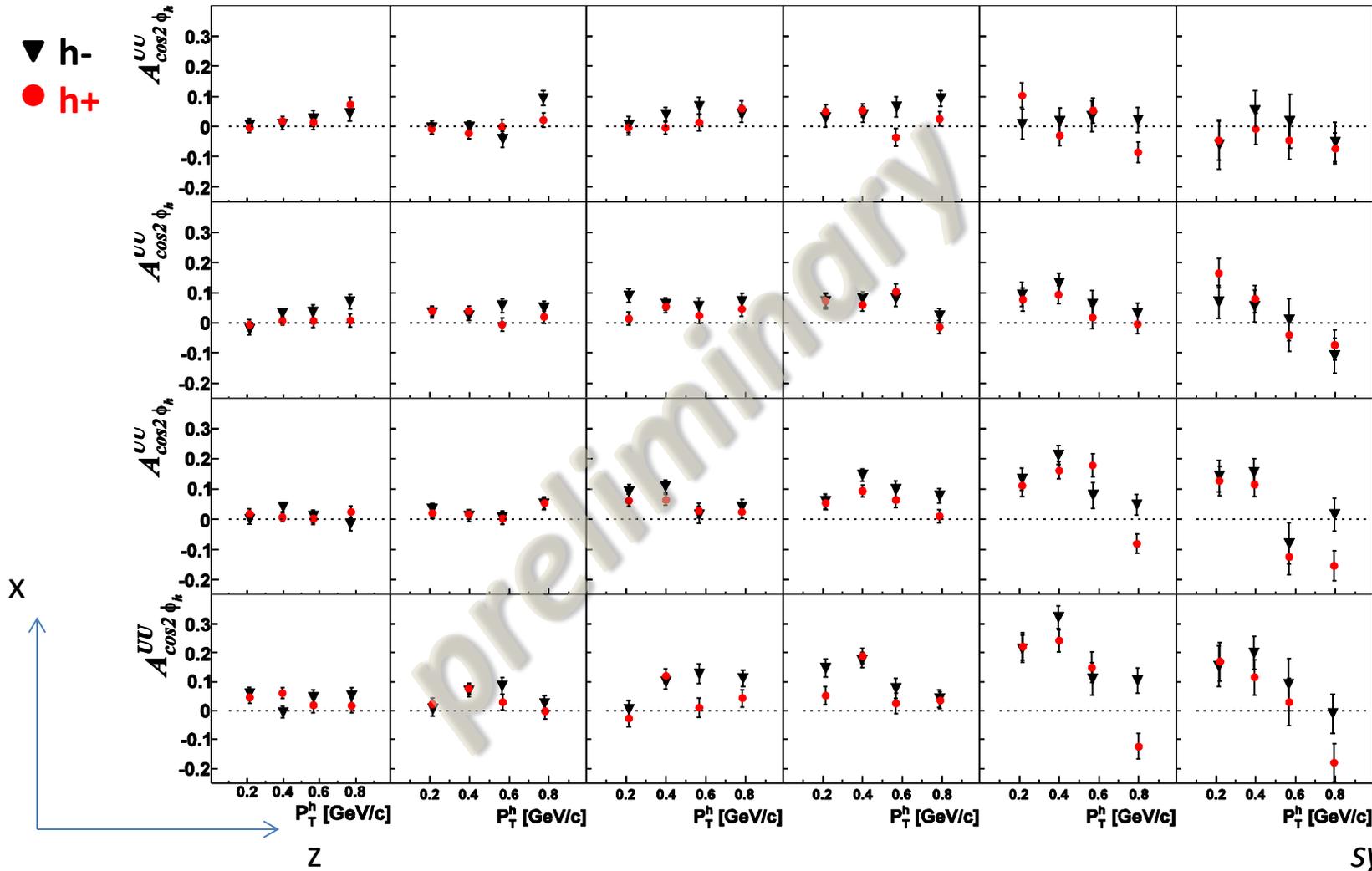
Results extracted binning alternatively in x , z and P_T^h



$sys \approx 2 \cdot stat$

$$A_{\cos 2 \phi_h}^{UU} = BM + \frac{1}{Q^2} Cahn$$

P_T^h dependence difficult to be reproduced (PRD81, Barone, Melis, Prokudin)



the P_T^h trend difficult to be reproduced by calculations is there for large z and low x

increasing interest for unpolarised SIDIS: theoretical work in progress...

- hadron multiplicities in SIDIS off unpolarised target
- azimuthal asymmetries in SIDIS off unpolarised target
- azimuthal asymmetries in SIDIS off transversely polarised target

Azimuthal hadrons distribution on transversely polarised target

direction of the polarization

kinematical factors $D_{NN} = \frac{1-y}{1-y-\frac{y^2}{2}}$

$$N_h^\pm(\phi_h, \phi_S) = N_h^0(\phi_h) \left[1 \pm f P_T D_{NN} A_{Coll} \sin(\phi_h + \phi_S + \pi) \pm f P_T A_{Siv} \sin(\phi_h - \phi_S) \pm \dots \right]$$

target polarisation P_T
dilution factor f

Collins angle

Sivers angle

amplitudes of the azimuthal modulations **extracted**

fitting the azimuthal distribution in ϕ_h and ϕ_S

in the different bins of \mathbf{x} , \mathbf{z} and \mathbf{p}_T^h

Azimuthal hadrons distribution on transversely polarised target

there are also **other 6 modulations related to different TMDs**
 they all have been measured at COMPASS

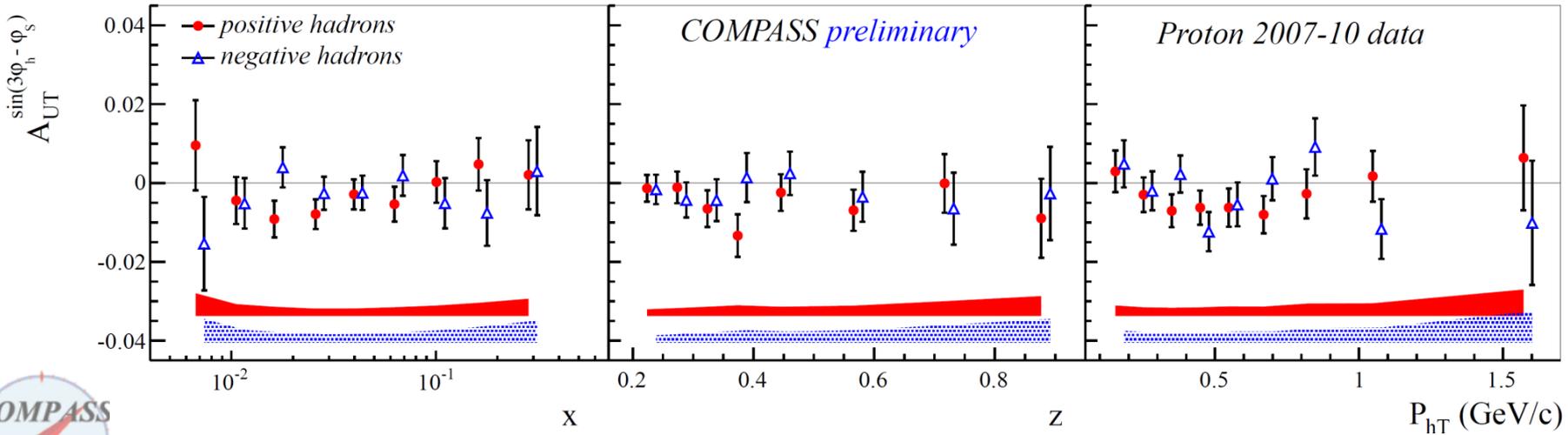
$$\begin{aligned}
 & + |\mathbf{S}_\perp| \left[\boxed{\sin(\phi_h - \phi_S)} \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
 & + \varepsilon \boxed{\sin(\phi_h + \phi_S)} F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \boxed{\sin(3\phi_h - \phi_S)} F_{UT}^{\sin(3\phi_h - \phi_S)} \\
 & \left. + \sqrt{2\varepsilon(1+\varepsilon)} \boxed{\sin\phi_S} F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \boxed{\sin(2\phi_h - \phi_S)} F_{UT}^{\sin(2\phi_h - \phi_S)} \right] \\
 & + |\mathbf{S}_\perp| \lambda_e \left[\sqrt{1-\varepsilon^2} \boxed{\cos(\phi_h - \phi_S)} F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \boxed{\cos\phi_S} F_{LT}^{\cos\phi_S} \right. \\
 & \left. + \sqrt{2\varepsilon(1-\varepsilon)} \boxed{\cos(2\phi_h - \phi_S)} F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \Bigg\},
 \end{aligned}$$

sivers
collins
pretzelosity
worm-gear

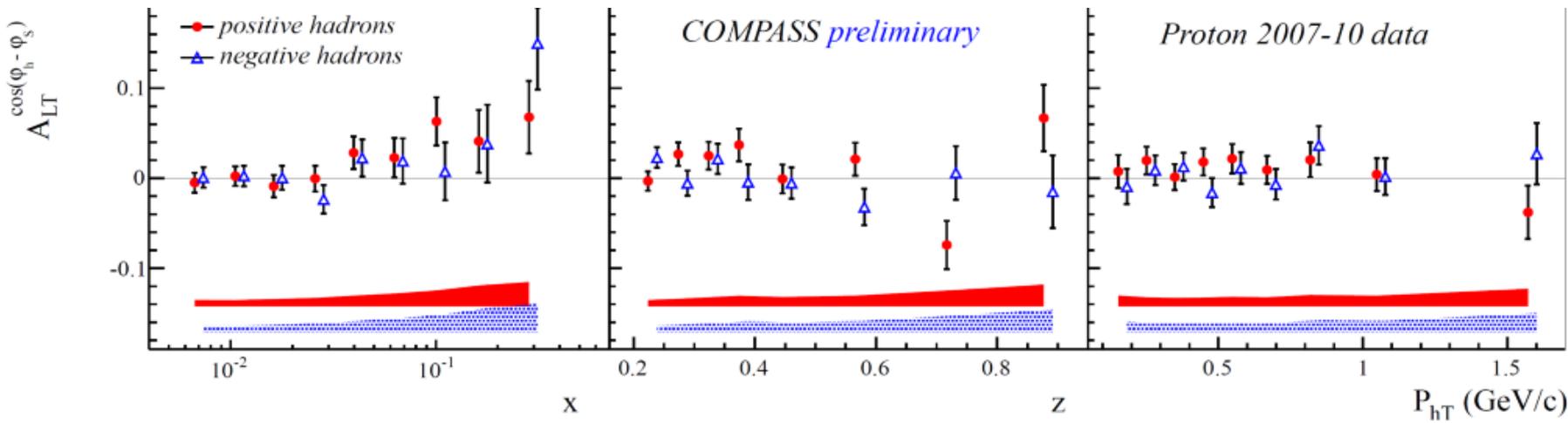
higher twist effects

sensitive to the D-wave component
non spherical shape of the nucleon

$$A_{UT}^{\sin(3\phi_h - \phi_s)} \propto h_{1T}^{\perp q} \otimes H_{1q}^{\perp h}, \text{ pretzelosity } h_{1T}^{\perp q} : \text{---} \odot \rightarrow \text{---} \odot \rightarrow$$



$$A_{LT}^{\cos(\phi_h - \phi_s)} \propto g_{1T}^q \otimes D_{1q}^h, \text{ worm-gear } g_{1T}^q : \text{---} \odot \rightarrow \text{---} \odot \leftarrow$$



Azimuthal hadrons distribution on transversely polarised target

there are also other 6 modulations related to different TMDs
they all have been measured at COMPASS

$$\begin{aligned}
 & + |\mathbf{S}_\perp| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
 & + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\
 & \left. + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right] \\
 & + |\mathbf{S}_\perp| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right. \\
 & \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \Bigg\},
 \end{aligned}$$

both on deuteron and on proton
in general the **asymmetries are small and compatible with zero**

The Sivers asymmetry

$$A_{Siv} \approx \frac{\sum_q e_q^2 \cdot f_{1T}^{\perp q}(k_{\perp}^2, x) \otimes D_{1q}^h(p_{\perp}^2, z)}{\sum_q e_q^2 \cdot f_1^q(k_{\perp}^2, x) \otimes D_{1q}^h(p_{\perp}^2, z)}$$

convolution on the intrinsic
transverse momentum of the quark



sivers PDF
(time odd)



correlation between the
nucleon transverse polarisation
and the **quark transverse momentum** k_{\perp}

*fundamental prediction pQCD (like B-M TMD)
sign change between Sivers TMD measured in SIDIS and in Drell-Yan
(future measurements from COMPASS)*

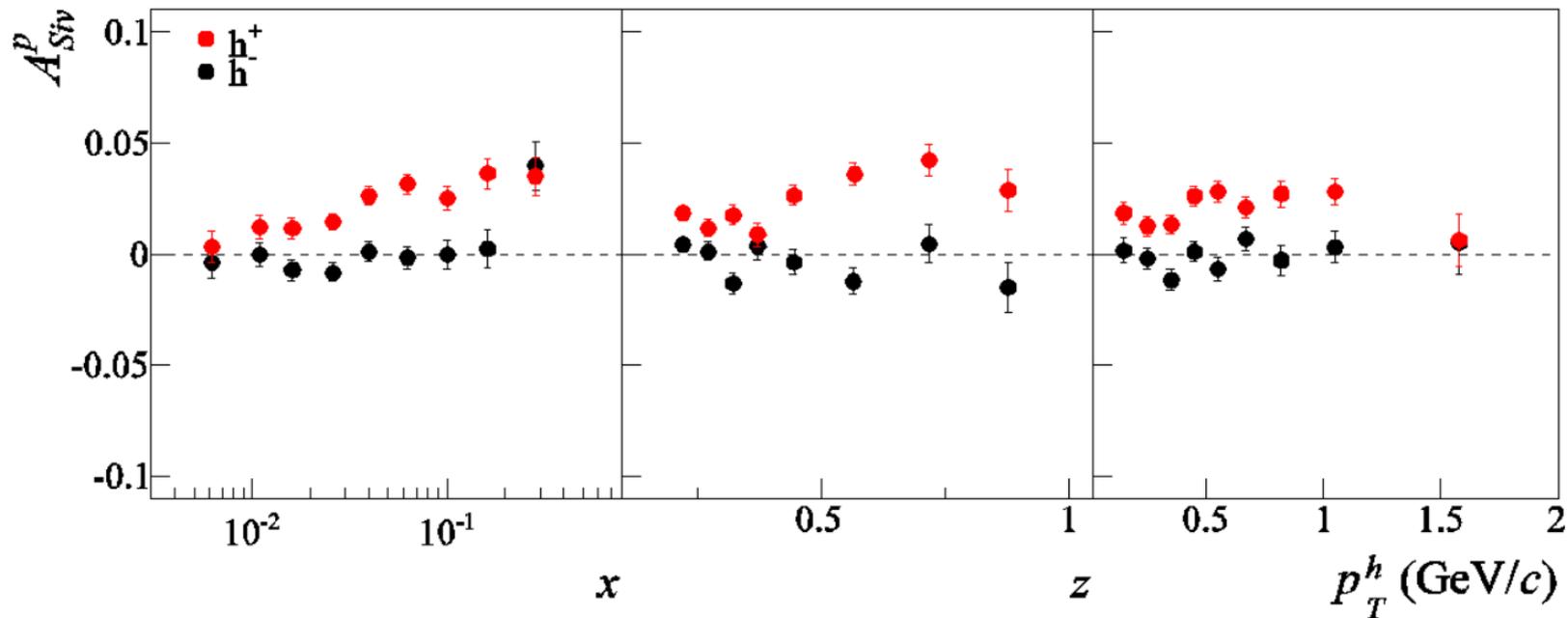
related with quark orbital angular momentum

The Sivers asymmetry *charged hadrons*

combined **2007** - PLB 692 (2010) 240 - and **2010** - PLB 717 (2012) 383 -

measurements on proton

very good agreement between two independent data set



h^+ : clear signal down to low x , in the previously unmeasured region

results on deuteron (2002-2004 data) compatible with zero

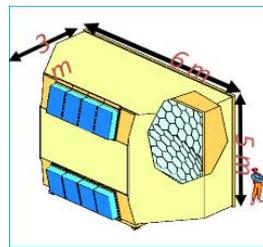
The Sivers asymmetry *pions and kaons*



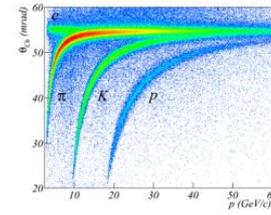
$\sim h^+, h^-$

$\sim 85\%$ of hadrons
are pions

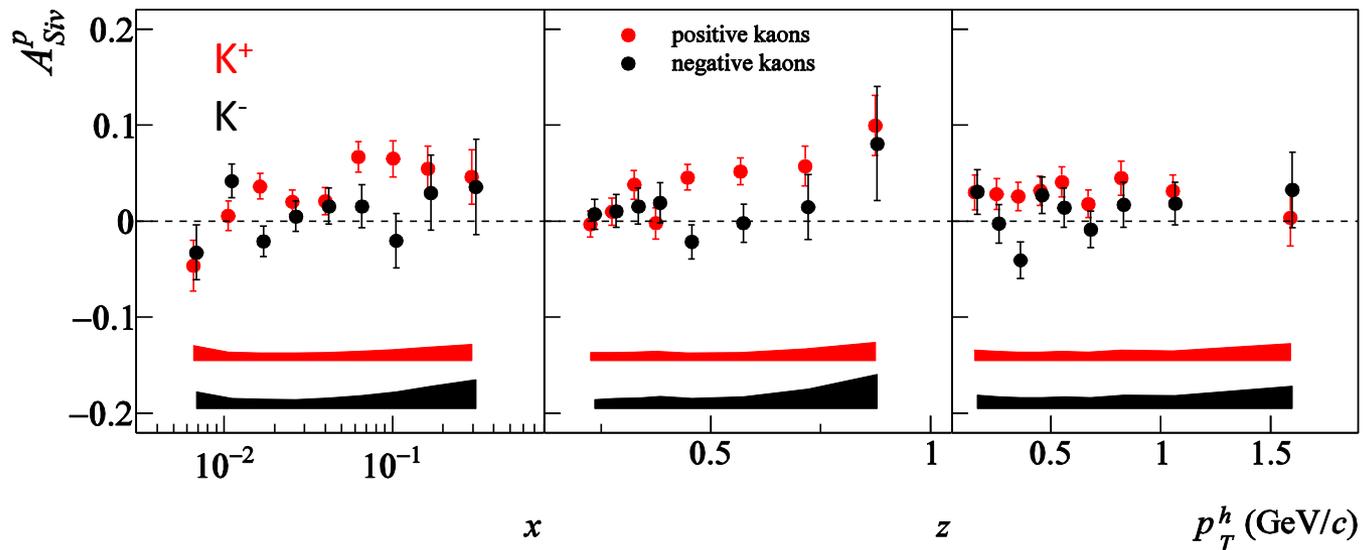
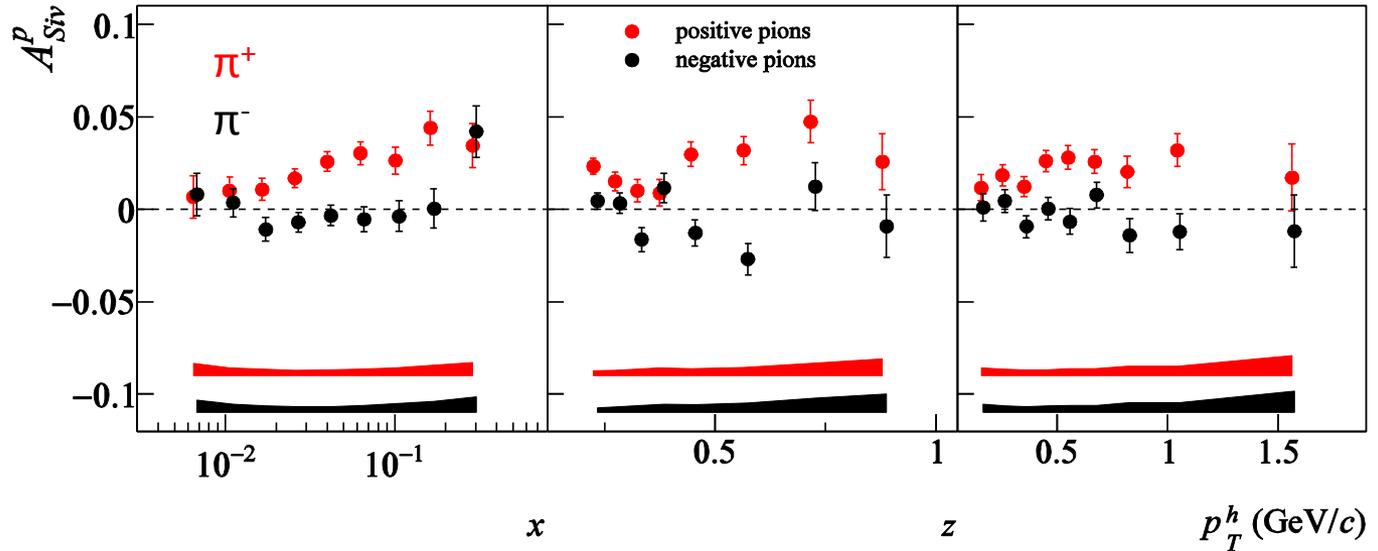
larger than for π^+

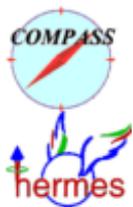


Hadron
identification:
 π^\pm, K^\pm
based on RICH-1
response
(likelihood algorithm)

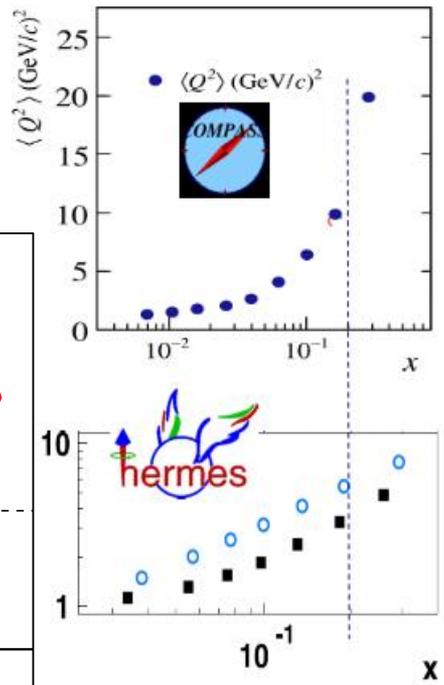
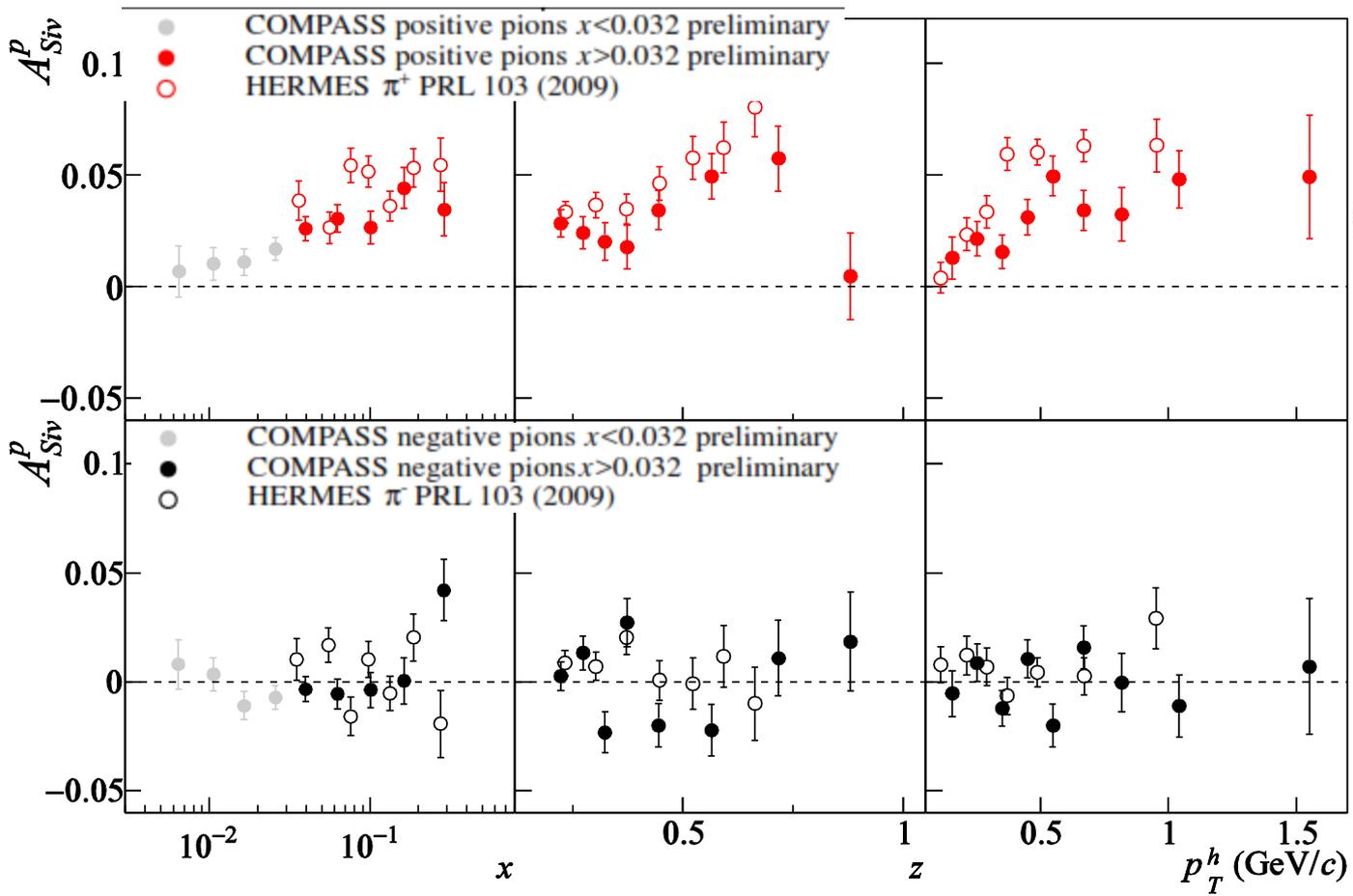


Cherenkov
thresholds
 $\pi \sim 3 \text{ GeV}/c$
 $K \sim 9 \text{ GeV}/c$
 $p \sim 18 \text{ GeV}/c$





comparison with HERMES
interesting, different Q^2



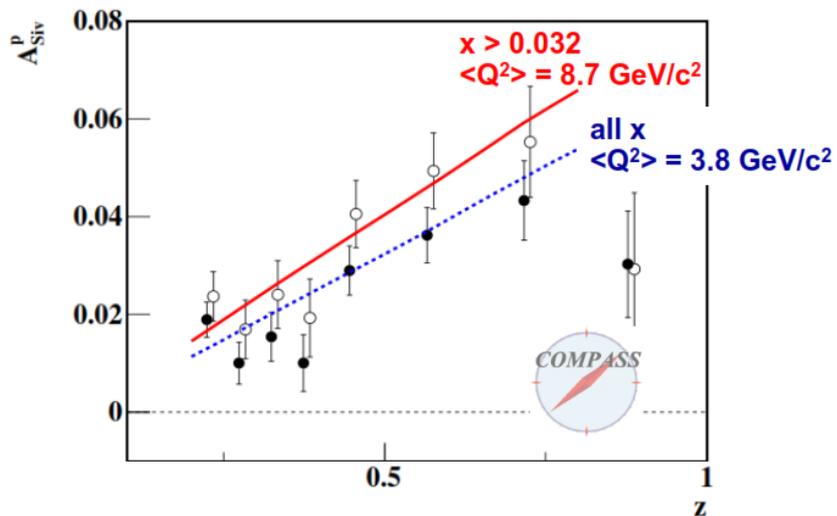
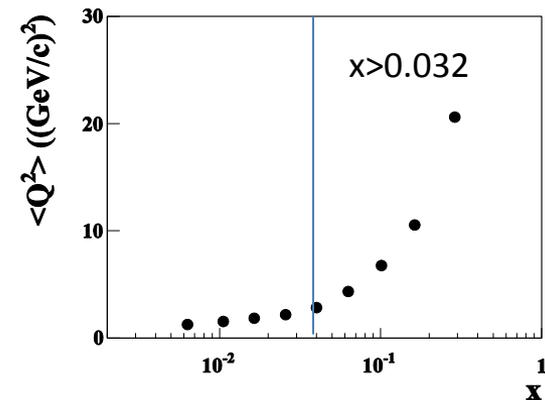
→ Q^2 TMD evolution ?

Sivers asymmetries measured for **all x** and for **x>0.032**

(different Q^2 values)



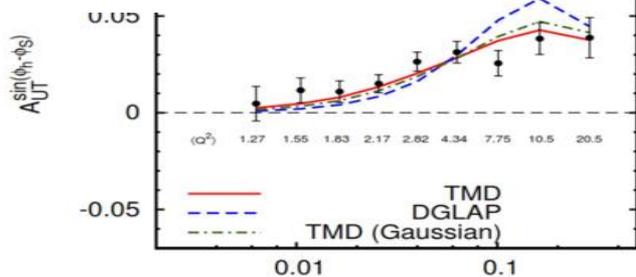
compared with the calculations based on the **TMD evolution**



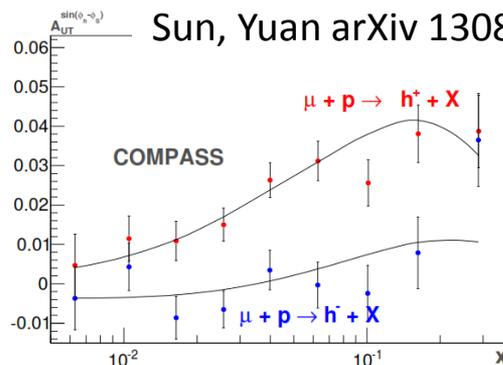
COMPASS PROTON

solid lines:
calculations from
Aybat, Prokudin, Rogers
PRL 108 (2012) 242003

Anselmino, Boglione, Melis
PRD 86 (2012) 014028



Sun, Yuan arXiv 1308.5003



*a lot of theoretical
work ongoing ...*

The Collins asymmetry



transversity PDF

(chiral odd
cannot be measured in DIS)



gives the difference between the number of quarks which have parallel and antiparallel transverse polarisation

collinear PDF: it survives after integration on the transverse momenta

$$A_{Coll} \approx \frac{\sum_q e_q^2 \cdot h_1^q(k_\perp^2, x) \otimes H_{1q}^{\perp h}(p_\perp^2, z)}{\sum_q e_q^2 \cdot f_1^q(k_\perp^2, x) \otimes D_{1q}^h(p_\perp^2, z)}$$

convolution on the intrinsic
transverse momentum of the quark

Collins FF (chiral odd)

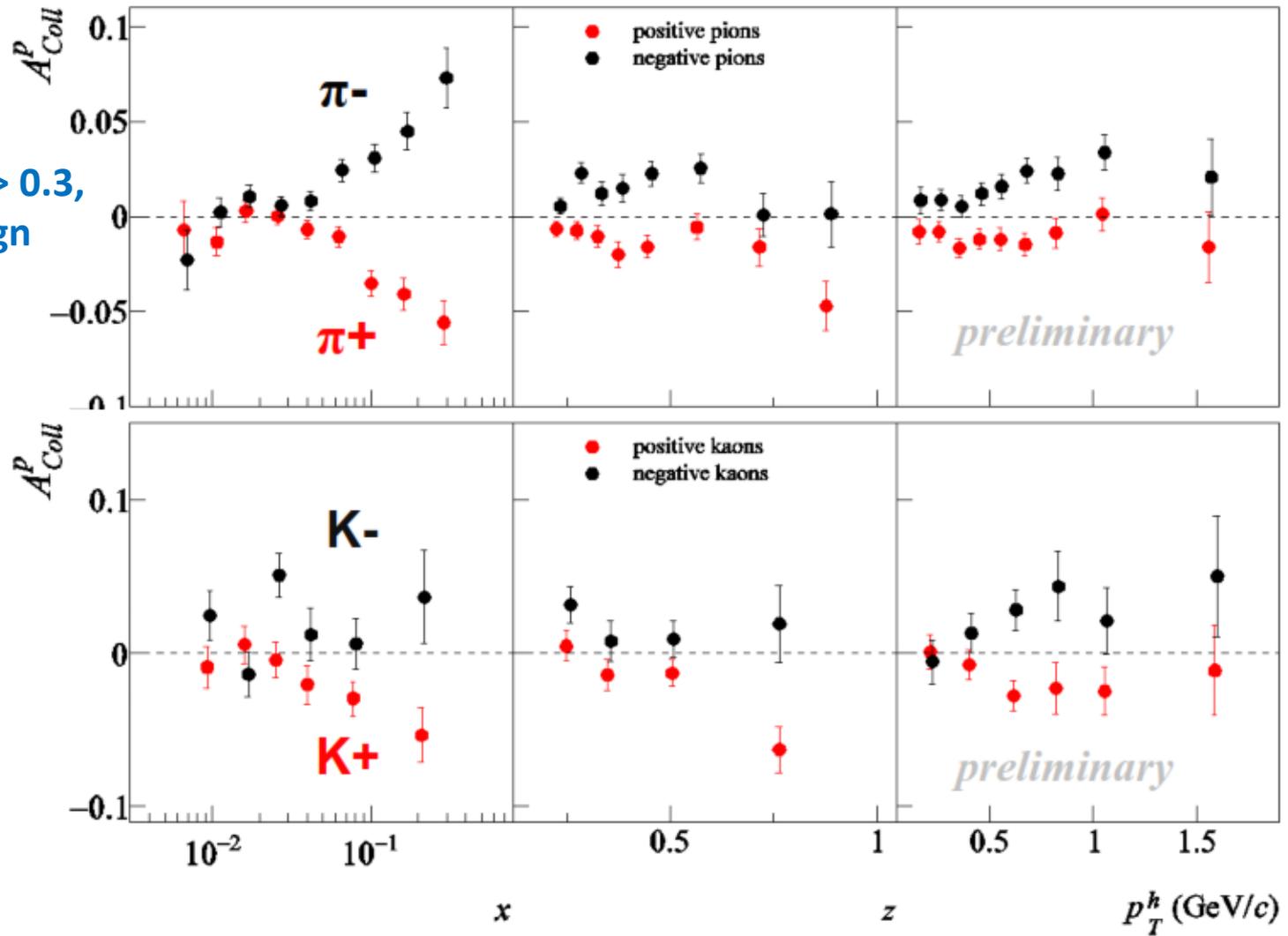
describes the hadronisation of a transversely polarised quark

The Collins asymmetry *charged pions and kaons*

combined 2007 - SPIN 2010 - and 2010 - SPIN 2012 -

measurements on transversely polarised proton

clear signal at $x > 0.3$,
with opposite sign
for h^+ and h^-

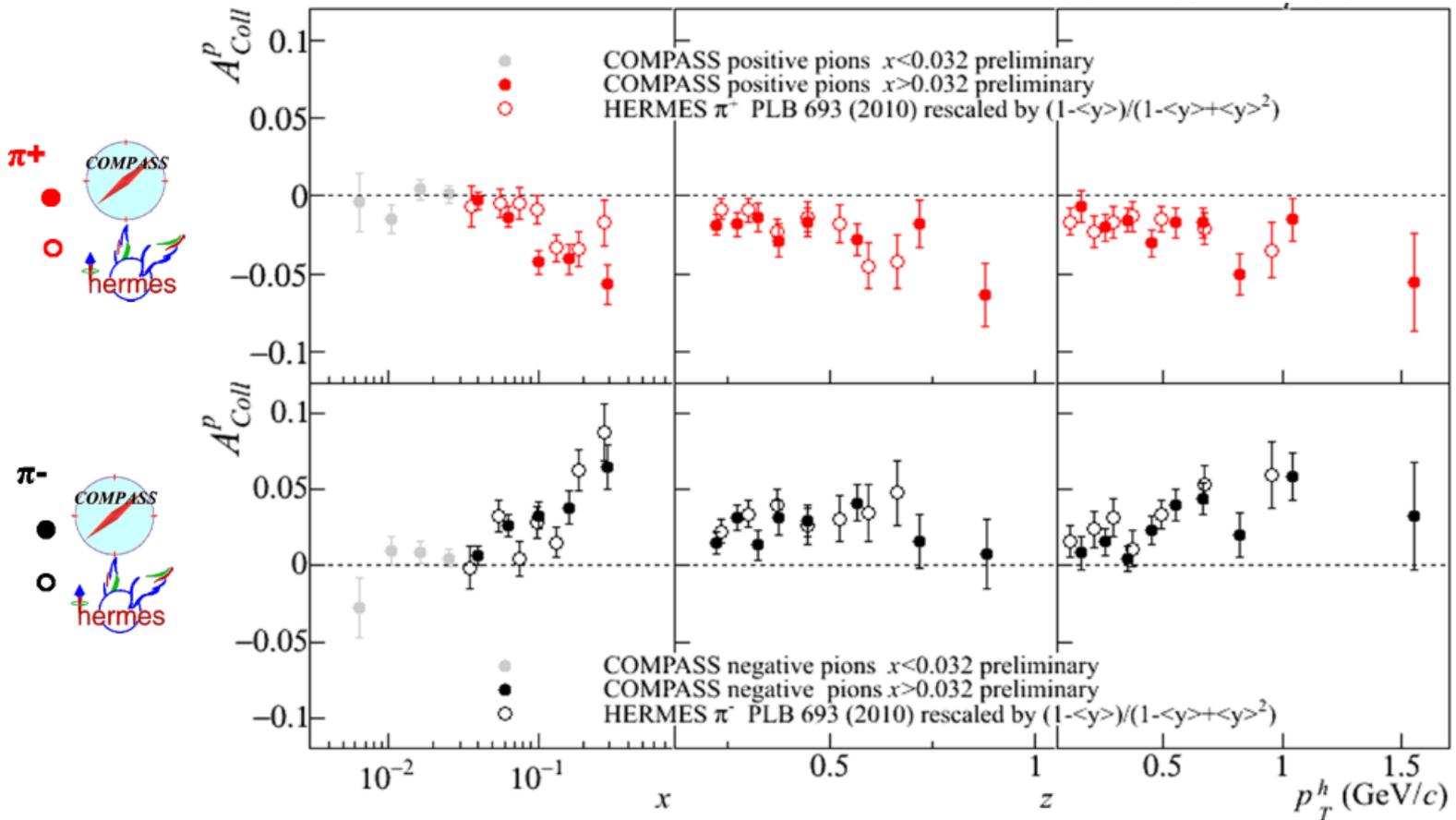


compatible
with π^+ , π^-

results on deuteron (2002-2004 data) compatible with zero

The Collins asymmetry

comparison with HERMES results



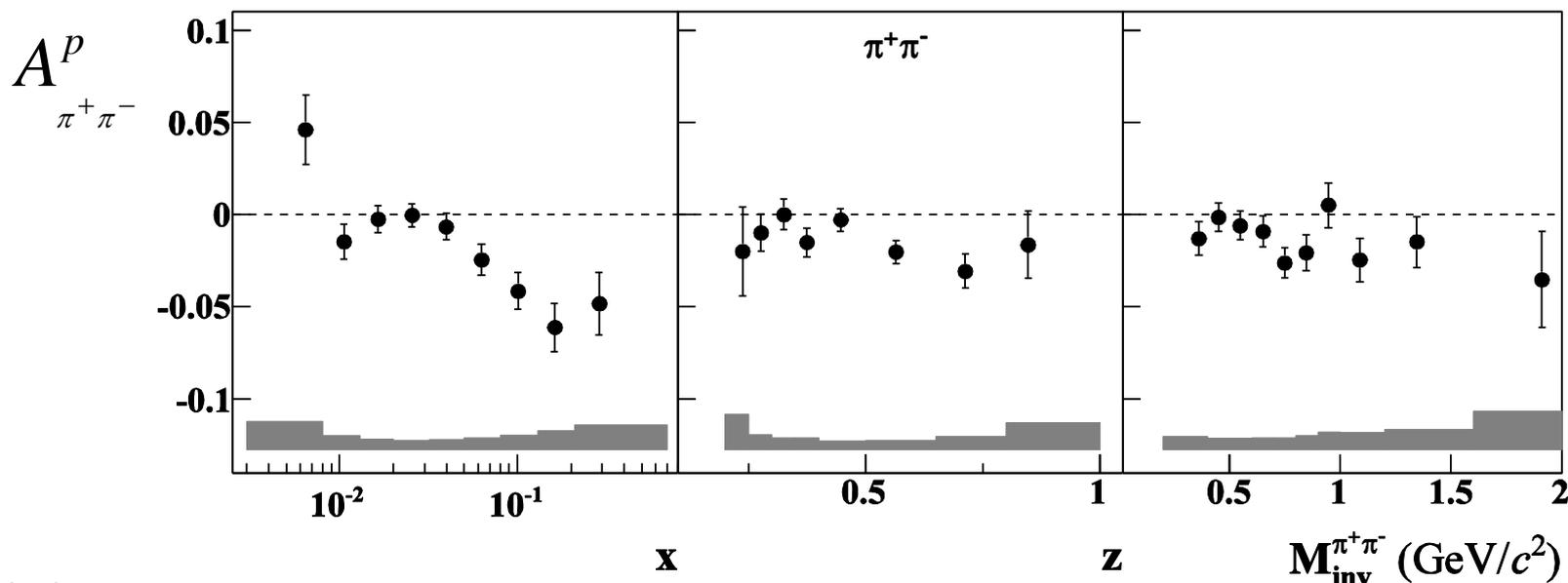
same strength: **no strong dependence on Q^2**

HERMES p + COMPASS $p&d$ + BELLE \rightarrow extraction of transversity for u and d quarks

The **transversity PDF** can be measured coupled with the 2-hadron interference fragmentation function

it gives rise to another azimuthal asymmetry in the 2h cross section, on an angle ϕ_{RS} defined as $\phi_R + \phi_S - \pi$

$$A_{h^+h^-}^p \approx \frac{\sum_q e_q^2 \cdot h_1^q(x) \cdot H_{1q}^{\angle}(z, M_{hh}^2)}{\sum_q e_q^2 \cdot f_1^q(x) \cdot D_{1q}^h(z, M_{hh}^2)}$$



Conclusions

The exploration of the **3-dimensional structure of the nucleon** is one of the major issues in high energy hadron physics.

embedded in the *Transverse Momentum Dependent Distributions Functions* (TMD PDFs)

→ SIDIS polarised azimuthal cross section is a very useful tool to measure them

COMPASS has given an important contribution with the measurement at high energy of all the structure functions which are expected to appear in the SIDIS cross section

- ***both transversity and Sivers PDFs are clearly different from zero***
- the results give an ***important input to TMD Q^2 evolution***
- ***interesting kinematic dependences*** have been found in ***unpolarised SIDIS***

outlook

more results are coming

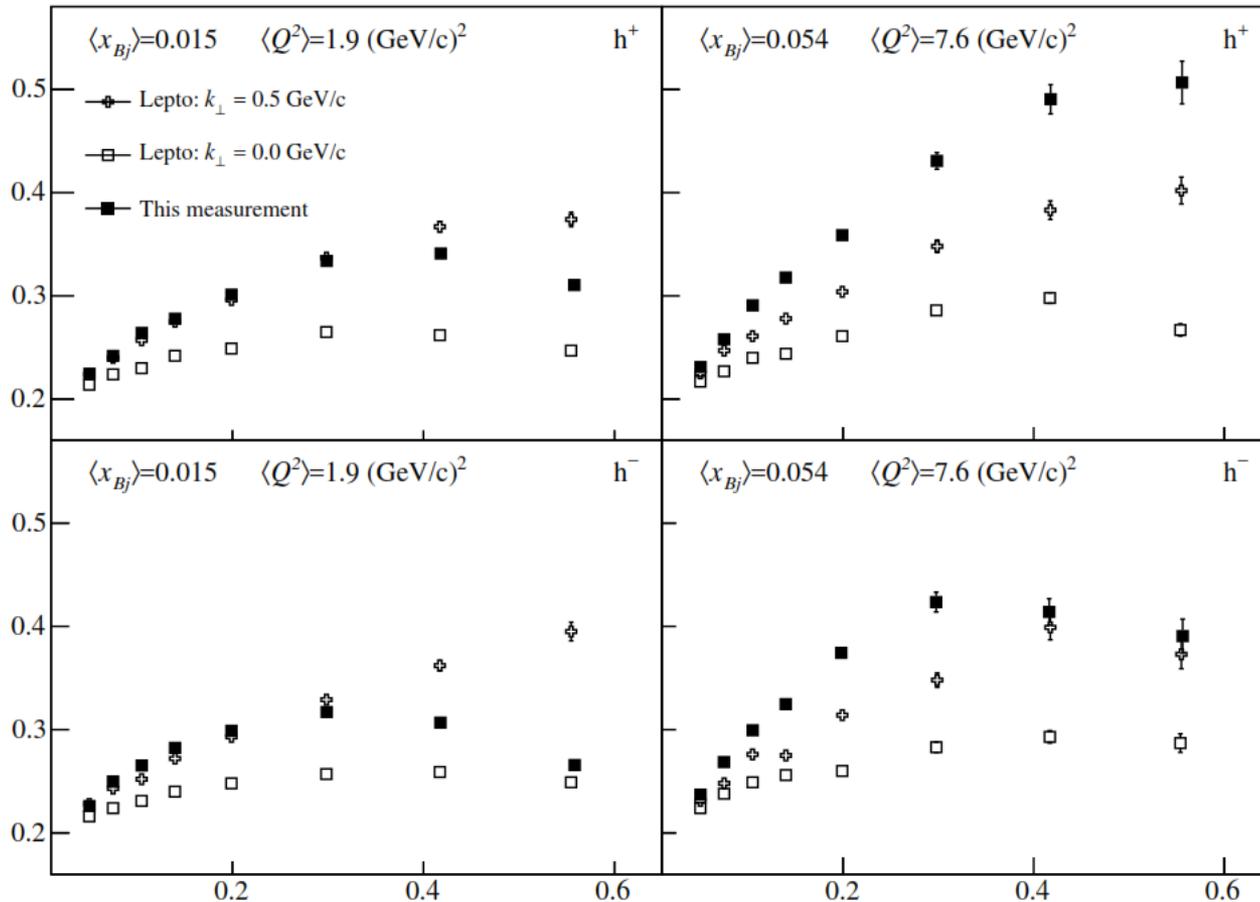
- ***hadron multiplicities*** in SIDIS off ***unpolarised*** target from ***2006 deuteron*** data
- ***azimuthal asymmetries*** in SIDIS off ***unpolarised*** target from ***2006 deuteron*** data
- ***multi dimensional analysis*** for the azimuthal asymmetries in SIDIS off ***transversely polarised*** target from the ***2010 proton data***

also (COMPASS-II, 2014-2017)

- ***complementary information on TMDs*** will come from (polarised) Drell-Yan
→Sivers, Boer-Mulders
- ***precise measurements from SIDIS*** on a long liquid H₂ target will be performed in parallel to the DVCS measurement
→ Cahn effect, Boer-Mulders

backup

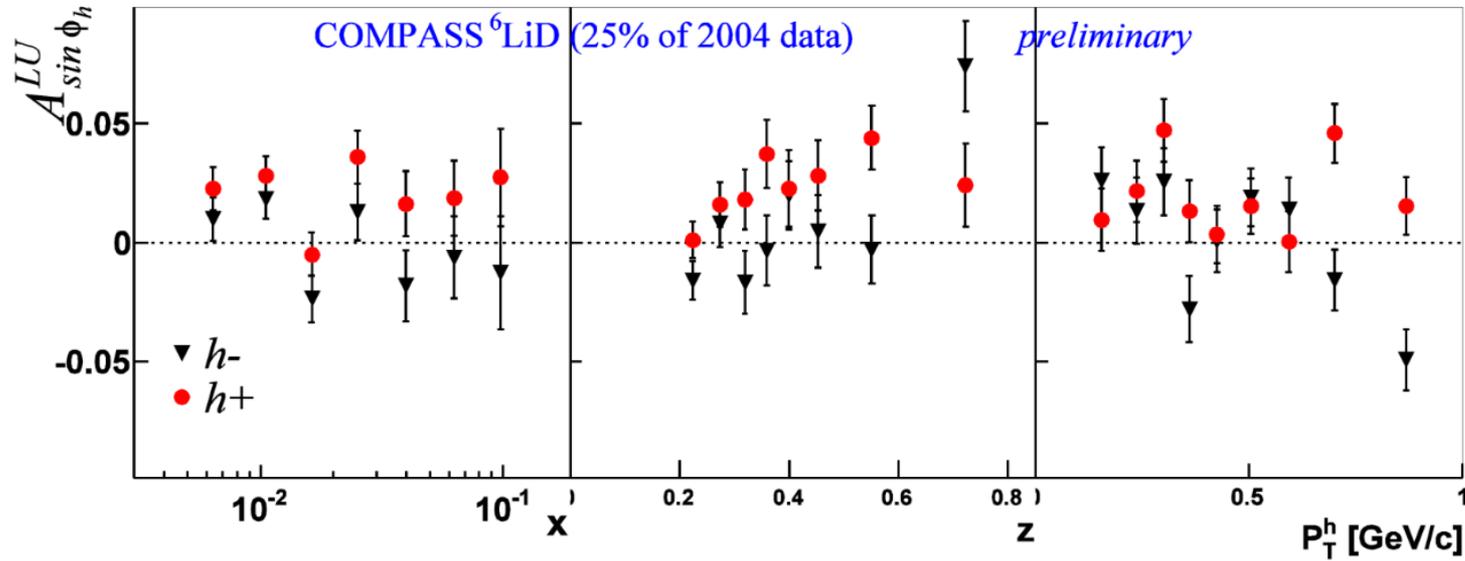
$$\langle P_T^{h^2} \rangle$$



the more relevant features (like the z dependence for h^+ and for h^-)
are not reproduced by the Monte Carlo



more work is needed to extract $\langle k_{\perp}^2 \rangle$



$A_{\sin \phi_h}^{LU}$ higher twist effect

$sys \approx 2 \cdot stat$

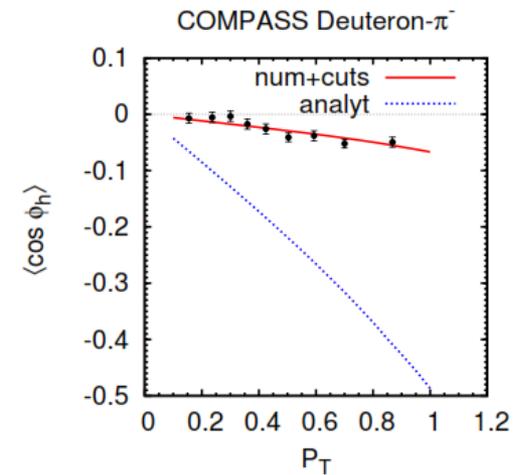
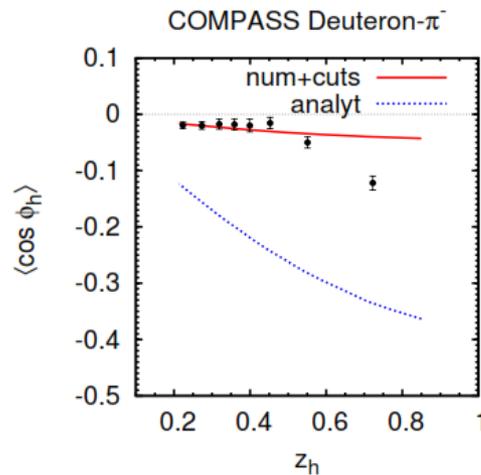
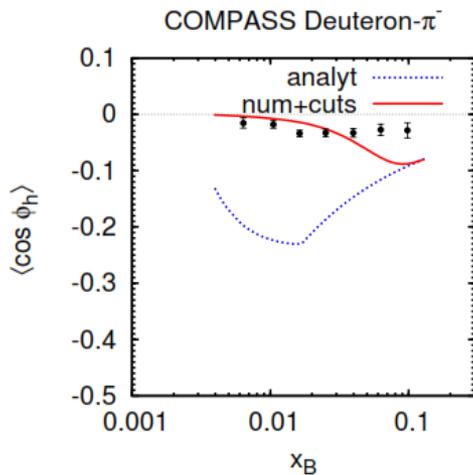
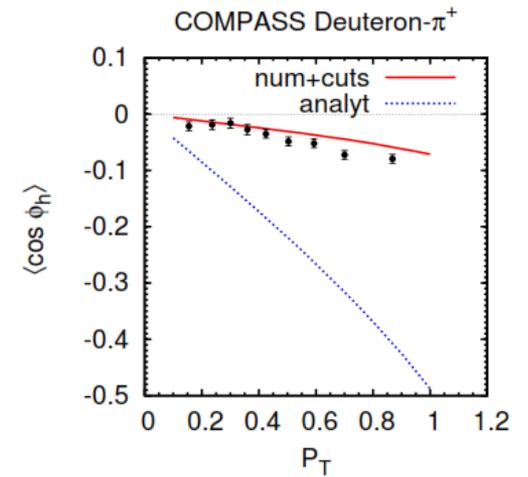
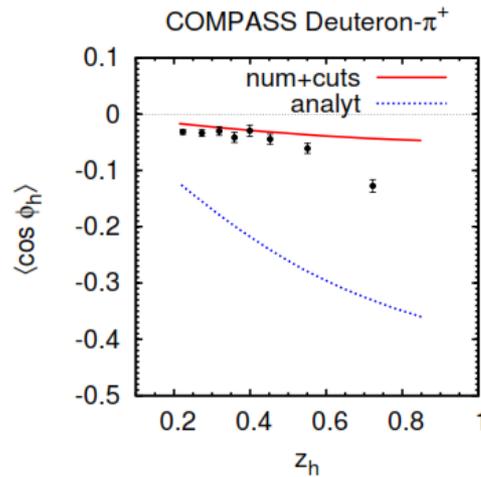
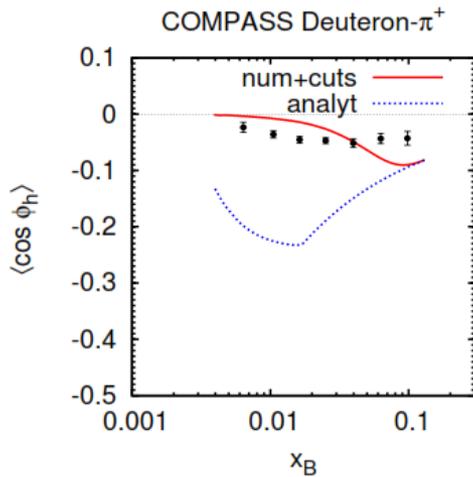
slightly positive signal for h^+
 h^- compatible with zero

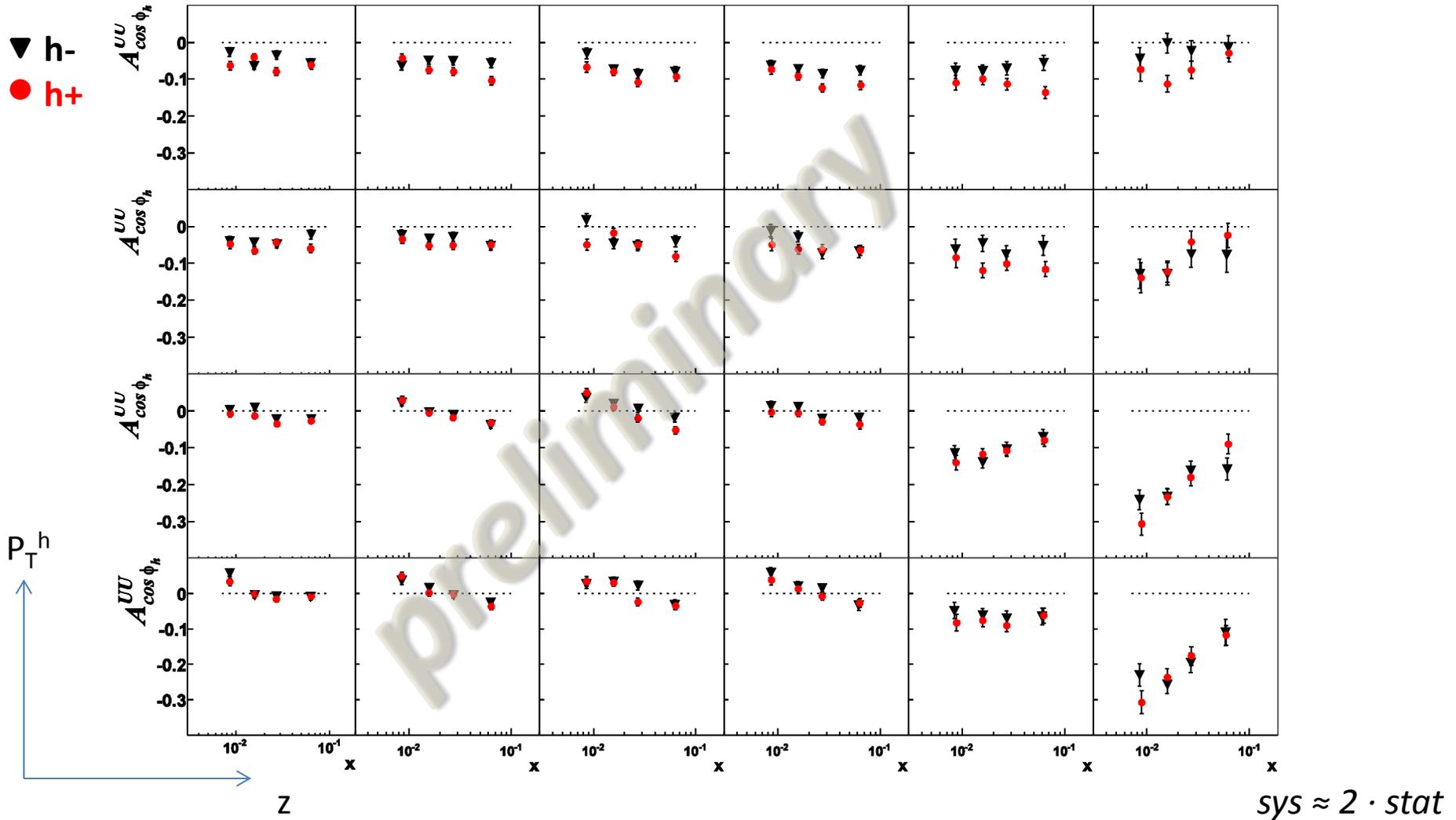
important input to phenomenological predictions



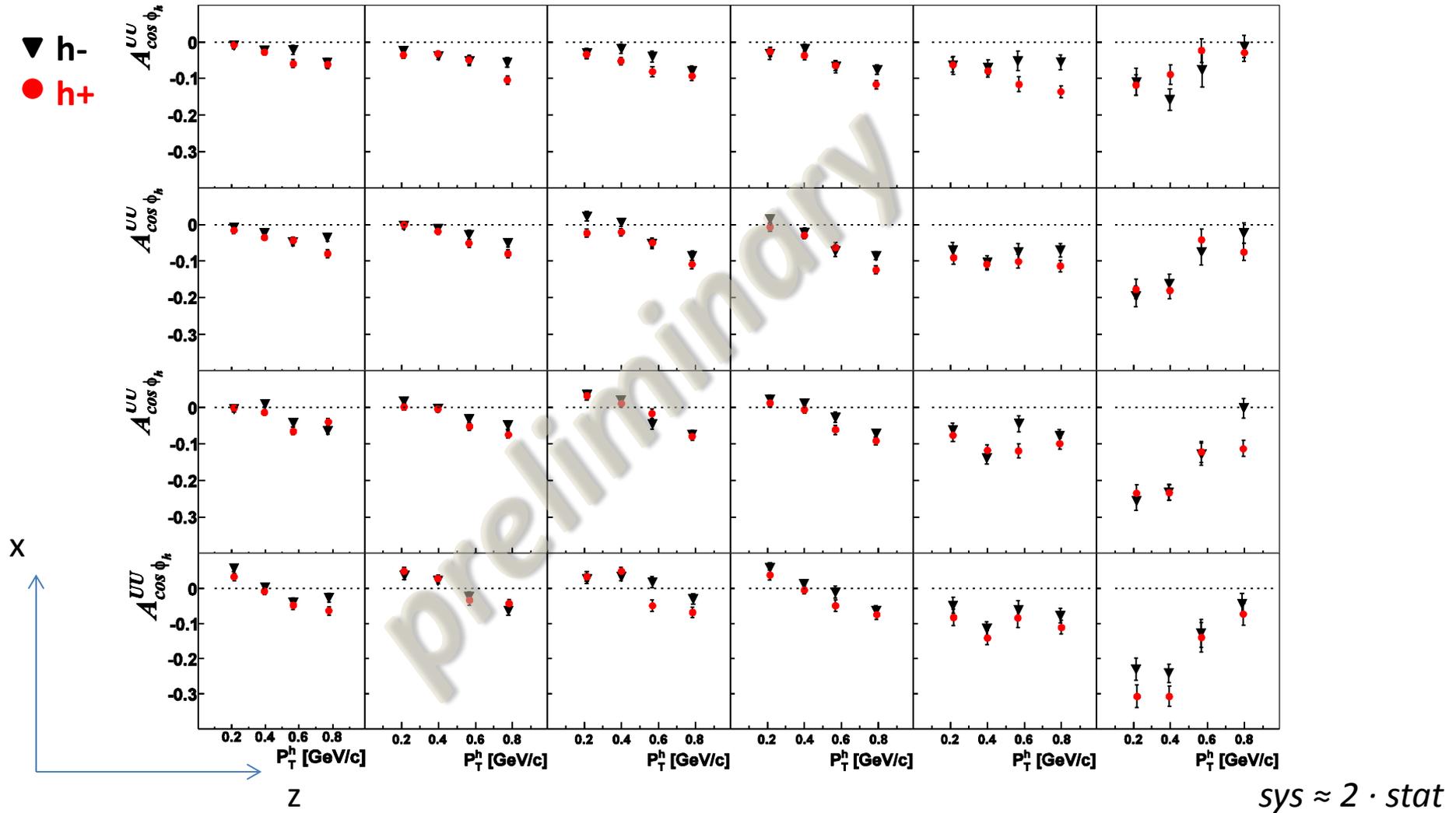
kinematical cuts had to be introduced for the quark intrinsic transverse momentum

$$f_{q/p}(x, k_{\perp}) = f_{q/p}(x) \frac{1}{1 - e^{-(k_{\perp}^{\max})^2 / \langle k_{\perp}^2 \rangle}} \frac{e^{-k_{\perp}^2 / \langle k_{\perp}^2 \rangle}}{\pi \langle k_{\perp}^2 \rangle}$$



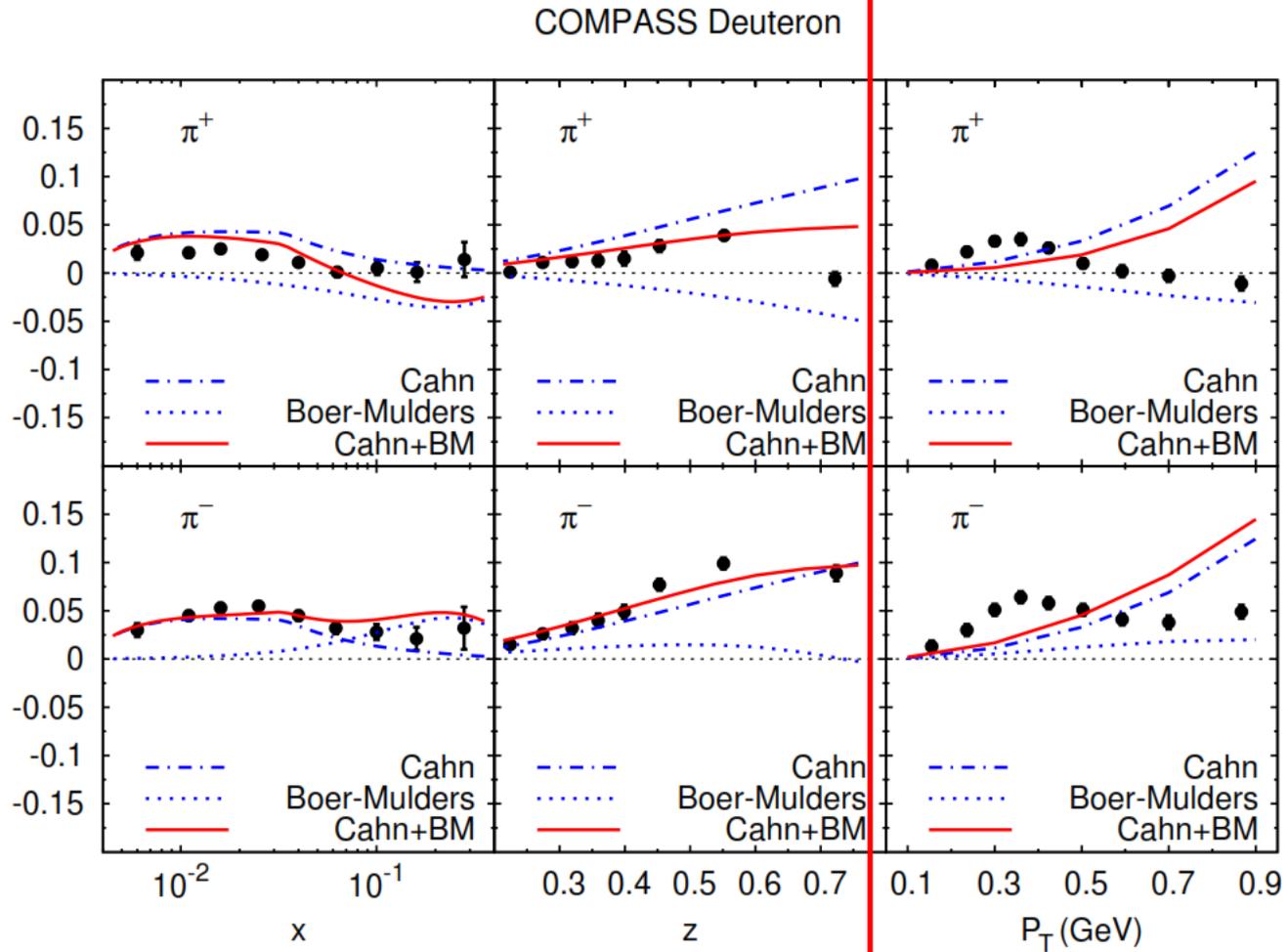


largest difference between positive and negative hadrons at large P_T^h
 x trend changes going from small to large z values

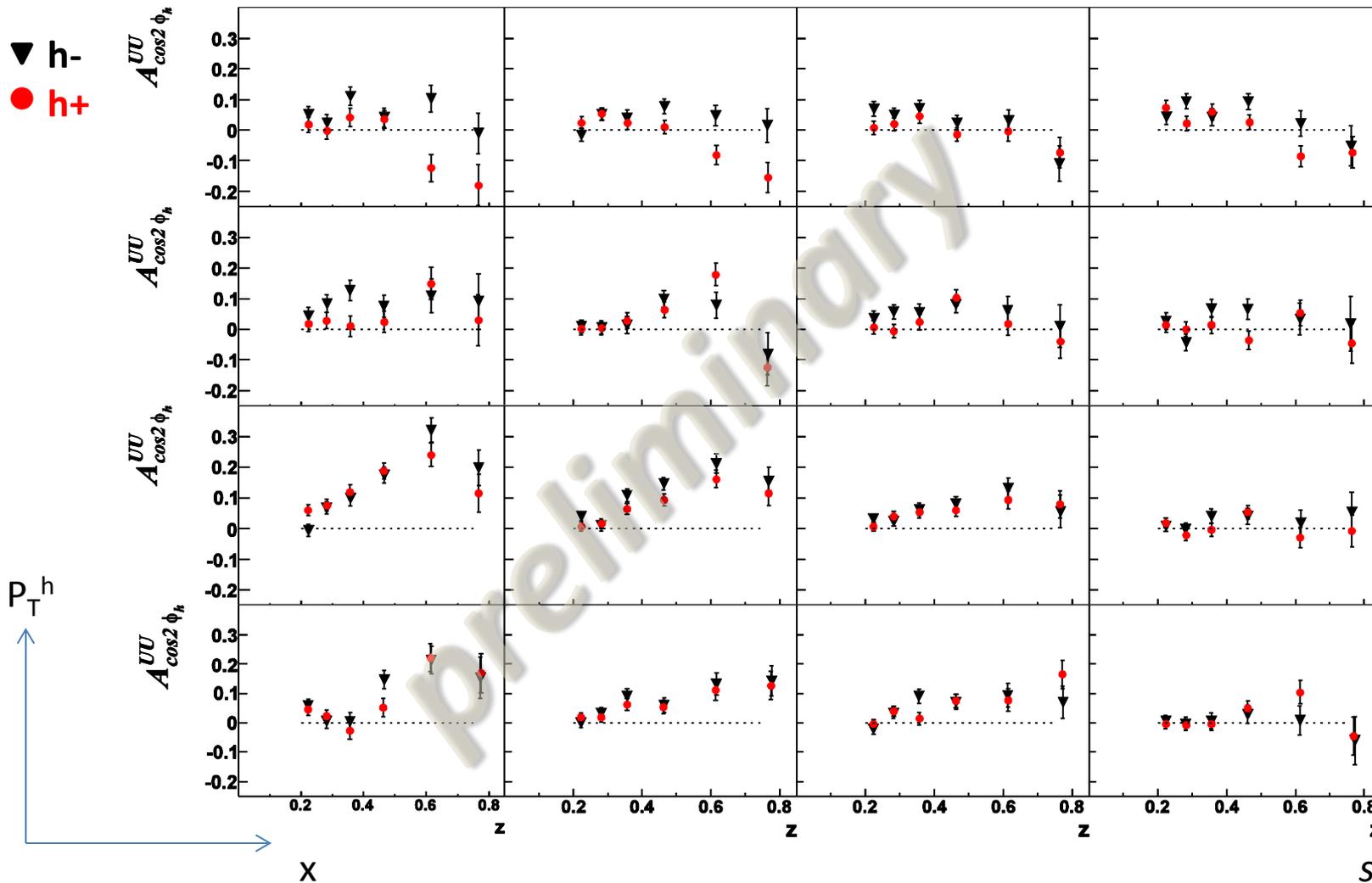


P_T^h trend changes going from small to large z values
 and it is roughly the same for all x intervals

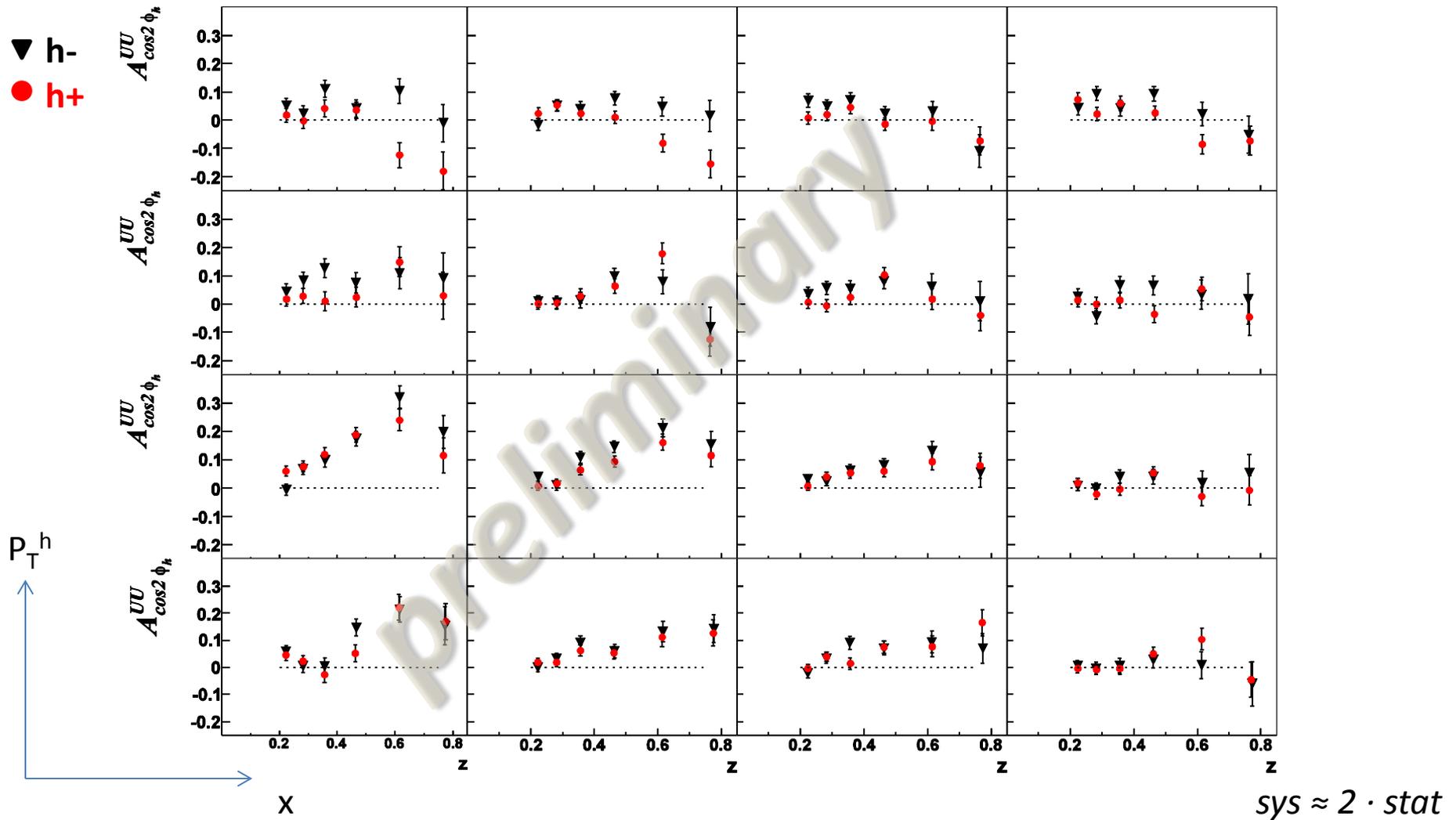
$$A_{\cos 2\phi_h}^{UU}$$



P_T^h dependence difficult to reproduce (PRD81, Barone, Melis, Prokudin)



strongest effect at low x and low P_T^h



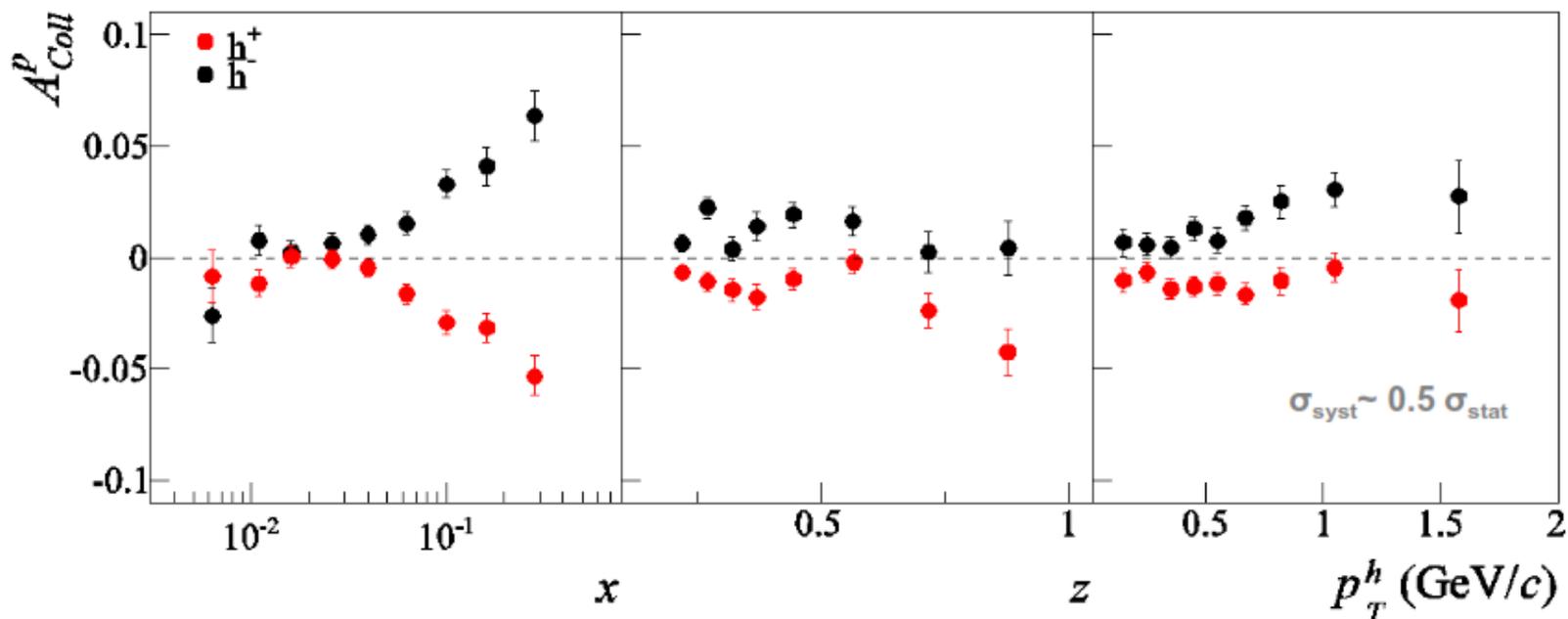
strongest effect at low x and low P_T^h

The Collins asymmetry *charged hadrons*

combined **2007** - PLB 692 (2010) 240 - and **2010** - PLB 717 (2012) 376 -

published measurements on transversely polarised proton

very good agreement between two independent data set



- **precise measurements**
- **clear signal at $x > 0.3$, with opposite sign for h^+ and h^-**
in agreement with the HERMES results

results on deuteron (2002-2004 data) compatible with zero

HERMES p + COMPASS $p&d$ + BELLE \rightarrow extraction of transversity for u and d quarks