

Radiative Coupling of $a_2(1320)$ and $\pi_2(1670)$

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on leave of absence from

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*for the **COMPASS** collaboration*

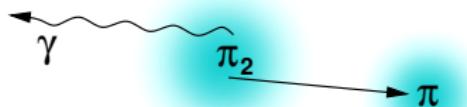
10th European Research Conference on
"Electromagnetic Interactions with Nucleons and Nuclei"
EINN2013, 1. November

Πάφος, Κύπρος





Why Radiative Couplings?



- radiative coupling “map” wave functions of the involved particles (known transition matrix element)
→ information of **shape** and **inner structure** of the meson resonances
- Related e.g. to $\pi\rho$ decay channels (vector meson dominance)



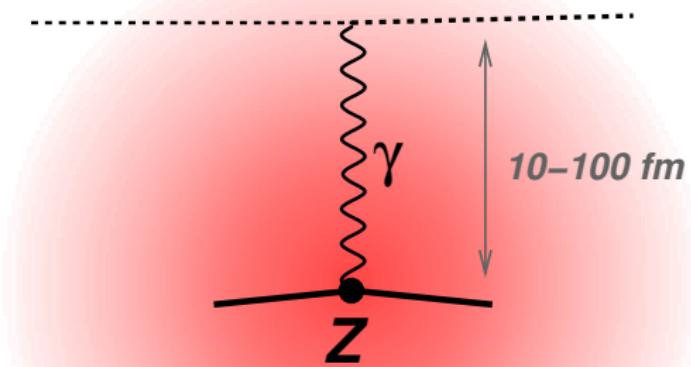
Outline



- Investigation of $\pi^- \gamma$ coupling through high- E hadron scattering
 - Primakoff reactions
 - higher-order aspects
- The apparatus: COMPASS
- Partial-wave analysis for identification of resonance contributions

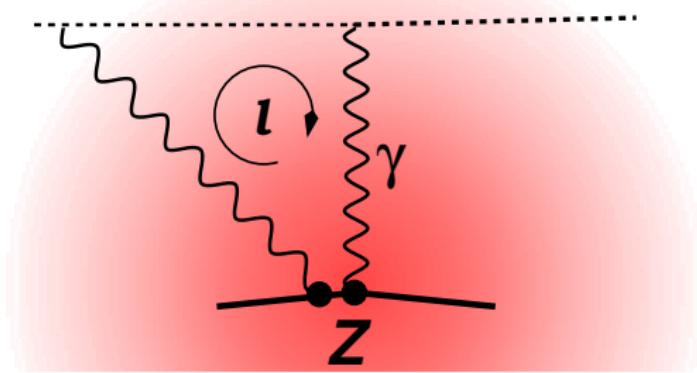


High- E Reactions: Momentum Transfer Regimes



- Small momentum transfer \Leftrightarrow large impact parameter in coordinate space
- At large distance: long-range *EM* potential dominates
- missing at this step:
EM at higher order, finite size of nucleus, strong interaction

High- E Reactions: Momentum Transfer Regimes

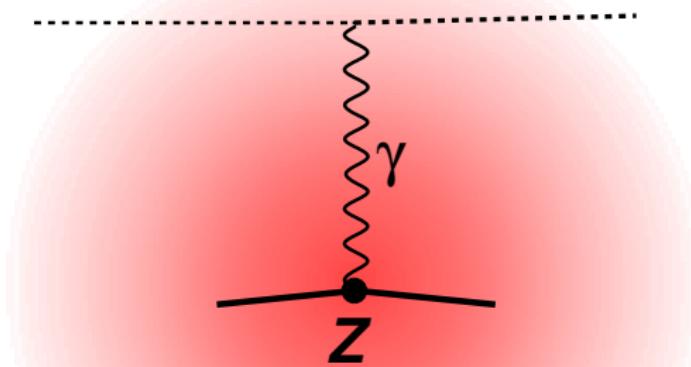


EM at higher order

- corresponds to “multiple-photon exchange”, but e.g. for lead $Z\alpha \approx 0.6$
- main effect (at low Q^2): distortion of the wave function in the potential, calculable in Glauber (eikonal) approach



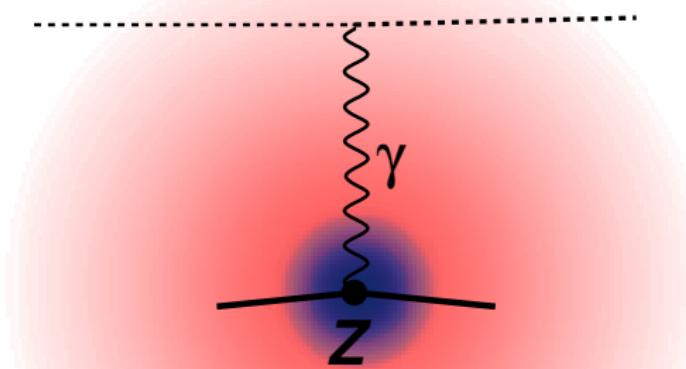
High- E Reactions: Momentum Transfer Regimes



- EM scattering becomes complex-valued, “Coulomb phase”
(\sim loops over resonance contributions)
- no observable effect for elastic scattering on static (heavy-mass) target



High- E Reactions: Momentum Transfer Regimes

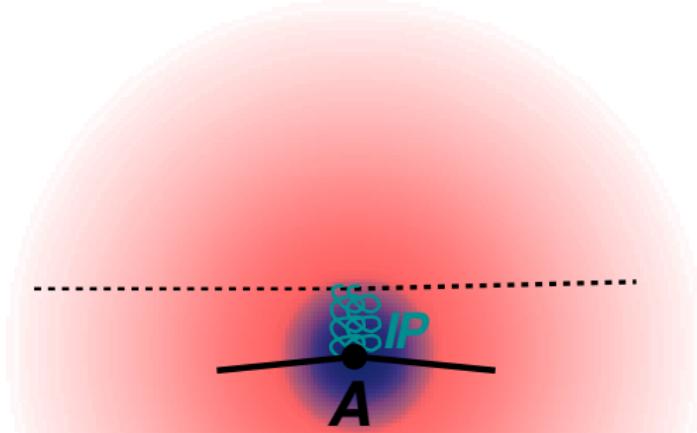


Within the Glauber model

- also the charge extension (as modified potential) can be taken into account
- drawbacks: e.g. neglect of recoil, not QFT



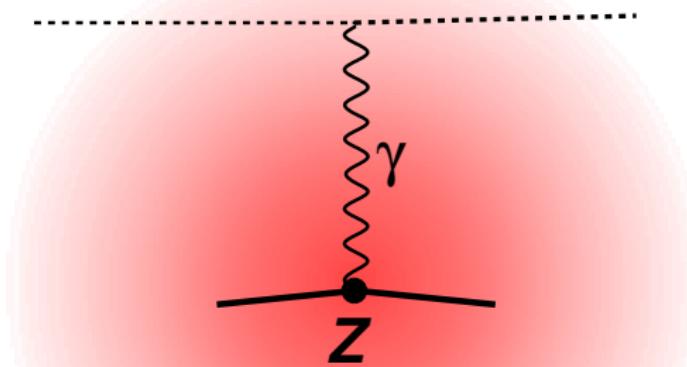
High- E Reactions: Momentum Transfer Regimes



- at the same level, the (absorptive) nuclear potential can be taken into account → diffractive regime
- exchange “particle” (Pomeron) \Leftrightarrow Fourier transform of potential



High- E Reactions: Momentum Transfer Regimes



- with a high-energy pion beam on nuclei, $\pi\gamma$ reactions become accessible
- identification as a strong increase of cross-section at smallest Q^2 ,
Primakoff effect



COmmon Muon and Proton Apparatus for Structure and Spectroscopy

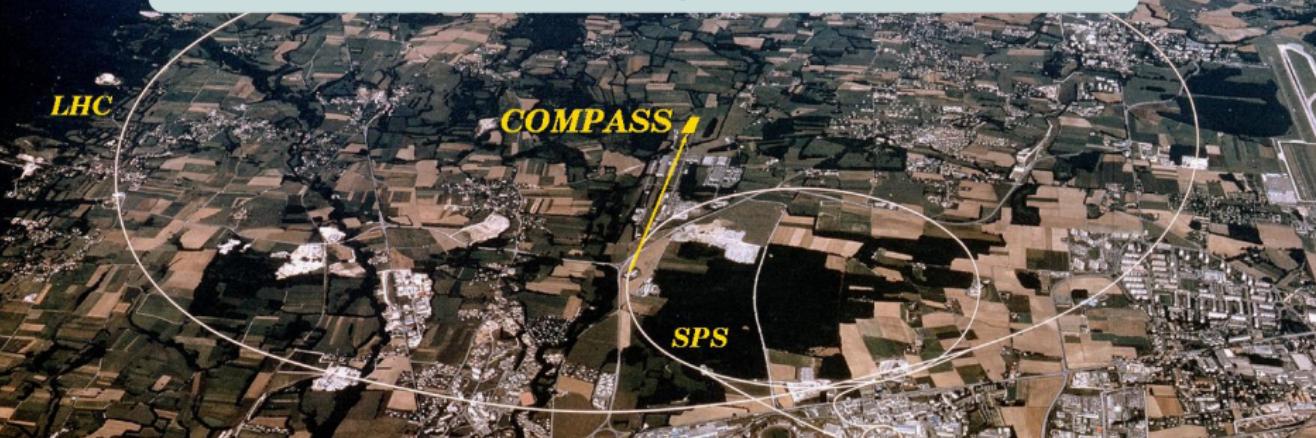




COmmon Muon and Proton Apparatus for Structure and Spectroscopy

CERN SPS: protons ~ 400 GeV (5 – 10 sec spills)

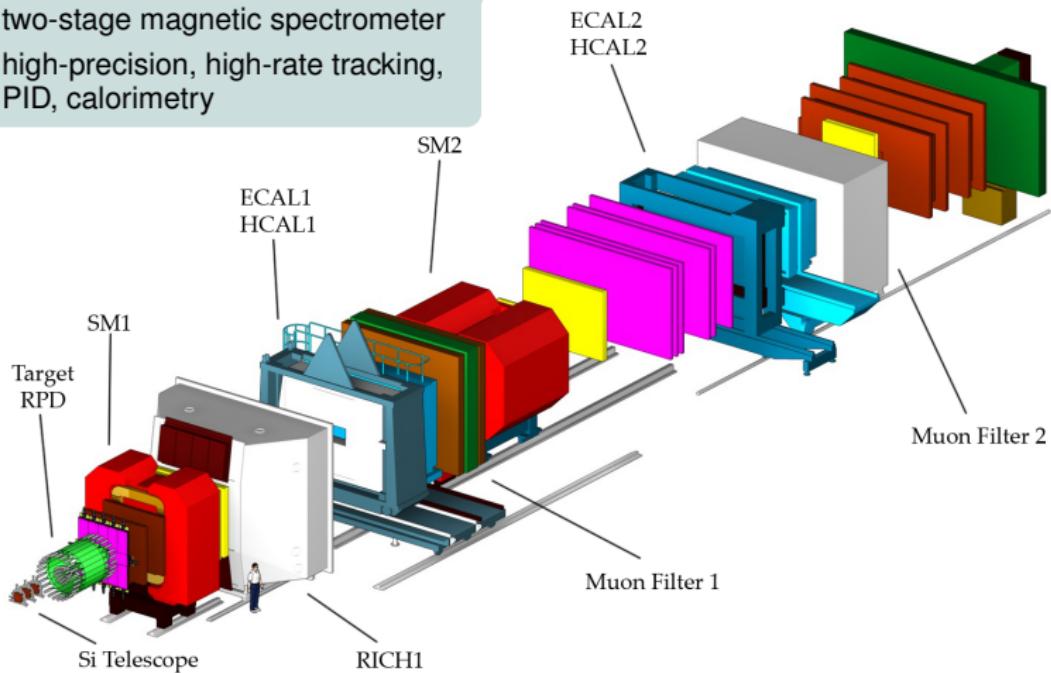
- secondary $\pi, K, (\bar{p})$: up to $2 \cdot 10^7$ / s
Nov. 2004, 2008-09, 2012:
hadron spec. & Primakoff reactions
- tertiary muons: $4 \cdot 10^7$ / s
2002-04, 2006-07, 2010-11: spin structure of the nucleon





Fixed-target experiment

- two-stage magnetic spectrometer
- high-precision, high-rate tracking, PID, calorimetry



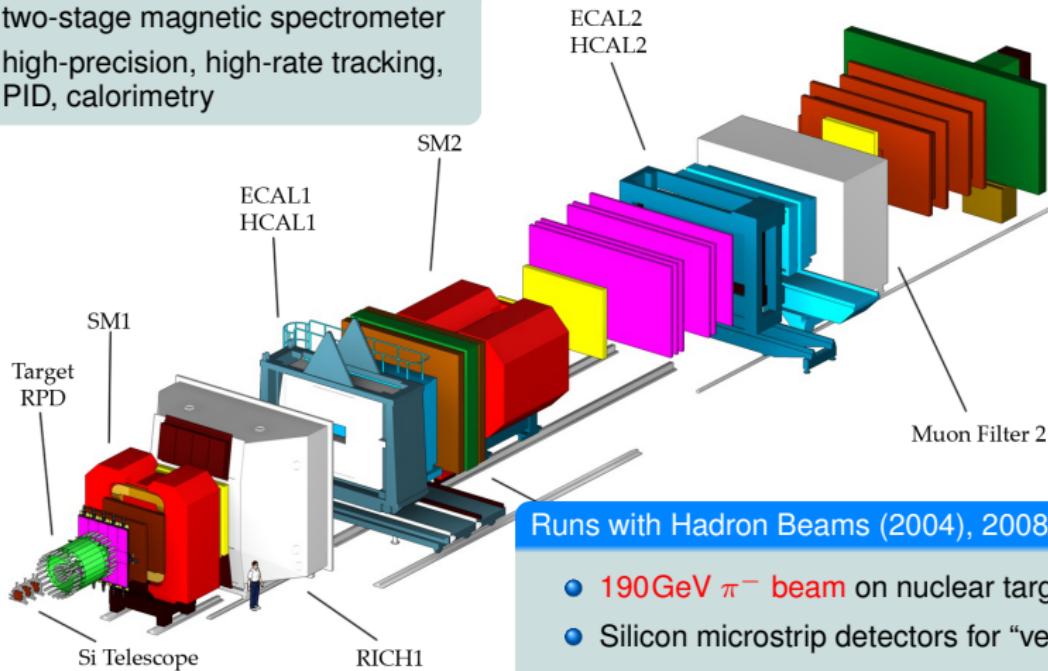


COMPASS

Experimental Setup

Fixed-target experiment

- two-stage magnetic spectrometer
- high-precision, high-rate tracking, PID, calorimetry

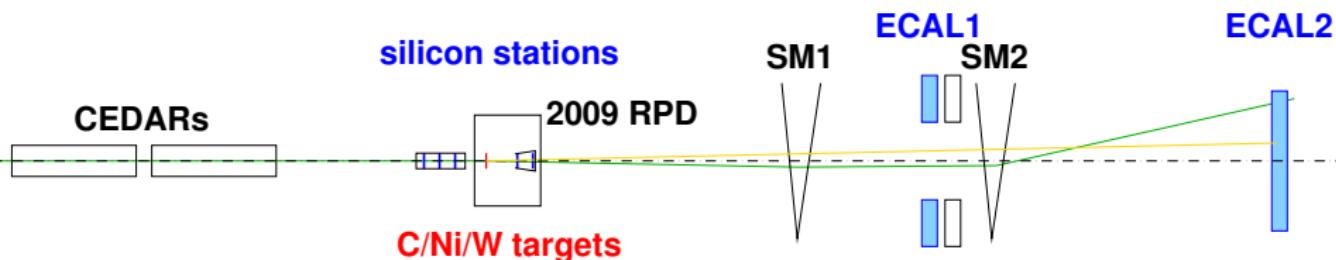


Runs with Hadron Beams (2004), 2008/09, 2012

- **190 GeV π^- beam** on nuclear targets (Ni, W)
- Silicon microstrip detectors for “vertexing”
- (digital) ECAL trigger



Principle of the measurement



- Trigger on **multiplicity** of charged forward tracks or on energy deposit in **EM calorimeter**
- Kaon component of the beam identified in Cerenkov detectors
- Free decay of beam K^- into $\pi^- \pi^0$, $\pi^- \pi^0 \pi^0$ and $\pi^- \pi^- \pi^+$ allows for cross-section normalization



Primakoff reactions accessible at COMPASS

Access to $\pi + \gamma$ reactions via the **Primakoff effect**:

At smallest momentum transfers to the nucleus, high-energetic particles scatter predominantly off the **electromagnetic field quanta** ($\sim Z^2$)

$$\pi^- + \gamma \rightarrow \left\{ \begin{array}{l} \pi^- + \gamma \\ \pi^- + \pi^0 / \eta \\ \pi^- + \pi^0 + \pi^0 \\ \pi^- + \pi^- + \pi^+ \\ \pi^- + \pi^- + \pi^+ + \pi^- + \pi^+ \\ \pi^- + \dots \end{array} \right. \quad \Leftarrow$$

analogously: Kaon-induced reactions $K^- + \gamma \rightarrow \dots$

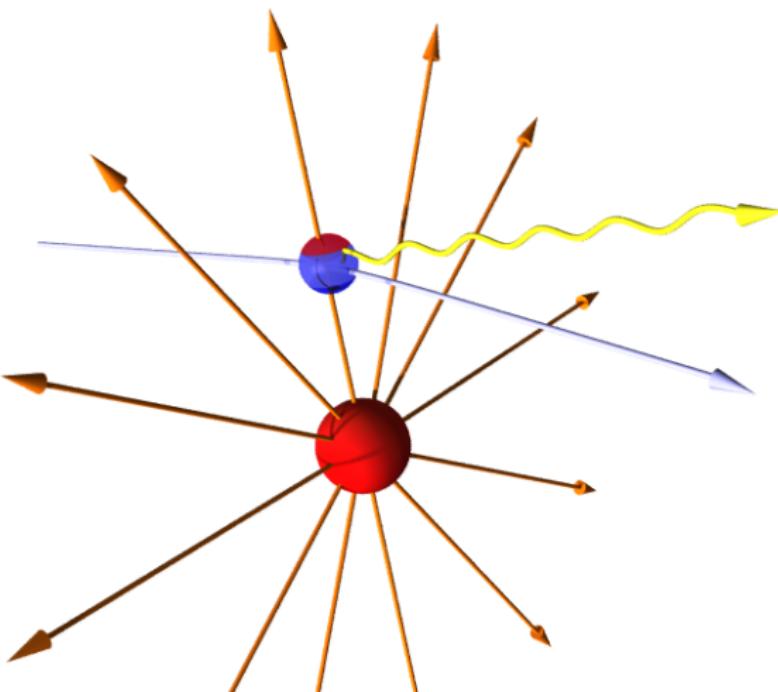


Primakoff technique: pion polarisability

- Charged pion traversing the nuclear **electric field**
 - typical field strength at $r = 5R_{Ni}$: $E \sim 300\text{ kV/fm}$

- Bremsstrahlung emission
 - particle scatters off **equivalent photons** (Weizsäcker-Williams)
 - pion (or muon) Compton scattering

- Polarisability contribution
 - Compton cross-section typically diminished
 - Theory prediction: $\alpha_\pi^{\text{ChPT}} = 2.9 \pm 0.5 \cdot 10^{-4} \text{ fm}^3$
 - expected charge separation $\sim 10^{-5} \text{ fm} \cdot e$





Pion polarisability at COMPASS

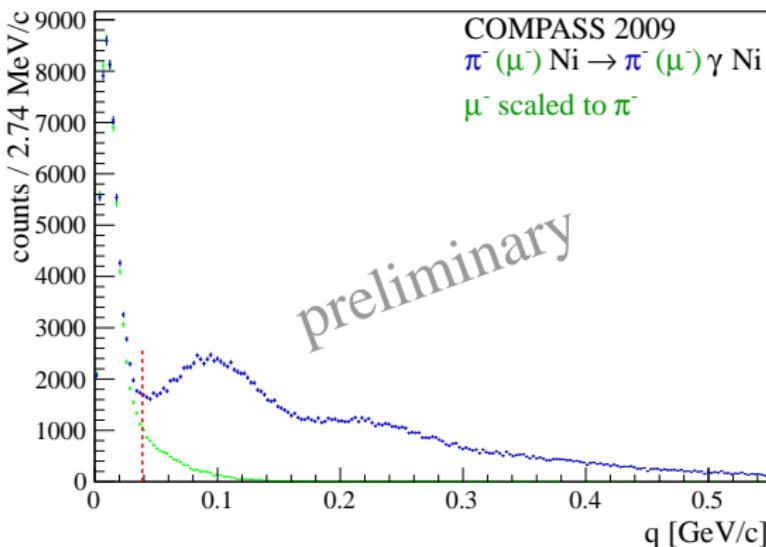
source of systematic uncertainty	estimated magnitude CL = 68 % [10 ⁻⁴ fm ³]
tracking	0.6
radiative corrections	0.3
background subtraction in Q	0.4
pion electron scattering	0.2
quadratic sum	0.8

COMPASS preliminary:

$$\alpha_\pi = 1.9 \pm 0.7_{\text{stat}} \pm 0.8_{\text{syst}} \times 10^{-4} \text{ fm}^3$$



Pion Compton Scattering: Primakoff peak

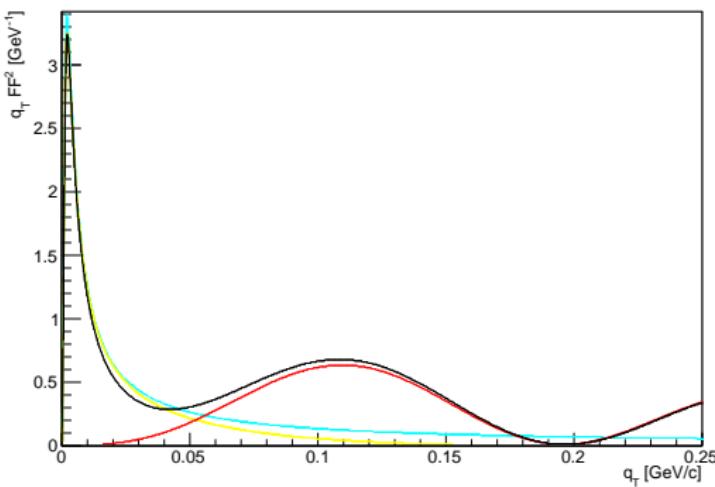


- $\Delta Q_T \approx 12 \text{ MeV}/c$ (190 GeV/c beam \rightarrow requires few- μrad angular resolution)
- first diffractive minimum on Ni nucleus at $Q \approx 190 \text{ MeV}/c$



Coulomb-nuclear interference

Photon density squared form factor



- Calculation following a 2009 paper of Göran Fäldt (Uppsala)
- Eikonal approximation: pions cross Coulomb and strong-interaction potentials



Primakoff reactions accessible at COMPASS

Access to $\pi + \gamma$ reactions via the **Primakoff effect**:

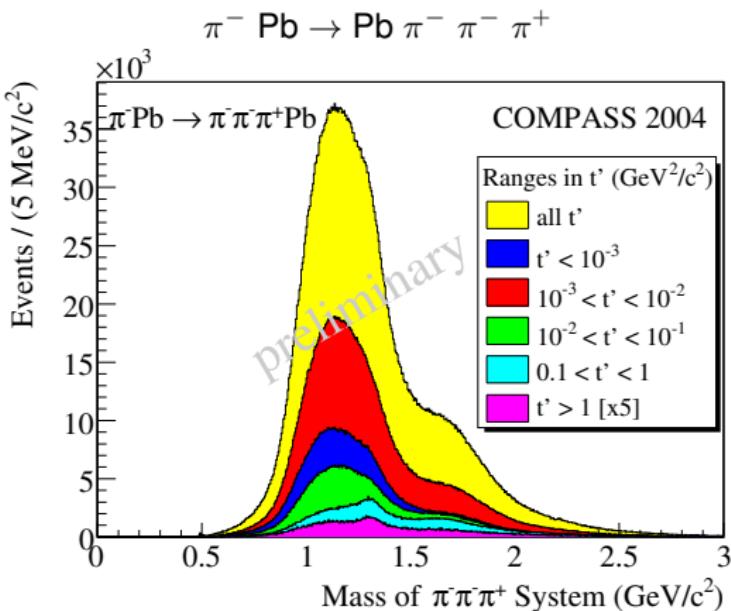
At smallest momentum transfers to the nucleus, high-energetic particles scatter predominantly off the **electromagnetic field quanta** ($\sim Z^2$)

$$\pi^- + \gamma \rightarrow \left\{ \begin{array}{l} \pi^- + \gamma \\ \pi^- + \pi^0 / \eta \\ \pi^- + \pi^0 + \pi^0 \\ \pi^- + \pi^- + \pi^+ \\ \pi^- + \pi^- + \pi^+ + \pi^- + \pi^+ \\ \pi^- + \dots \end{array} \right.$$

analogously: Kaon-induced reactions $K^- + \gamma \rightarrow \dots$



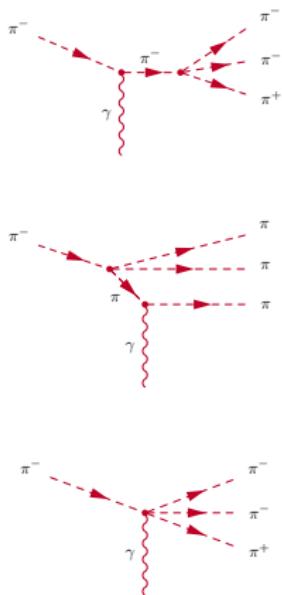
Three charged-pions final state



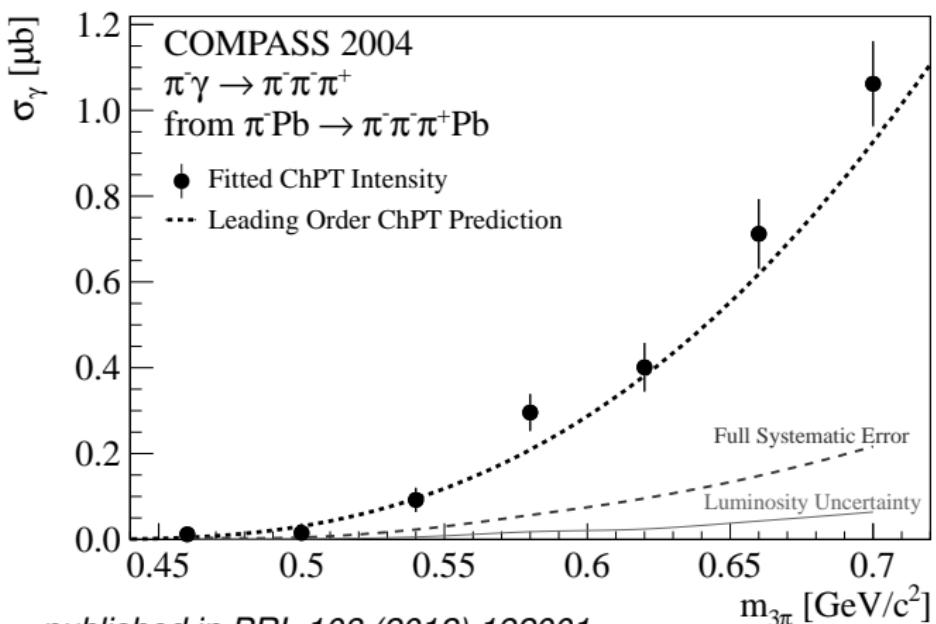
- "Low t' ": $10^{-3} (\text{GeV}/c)^2 < t' < 10^{-2} (\text{GeV}/c)^2$ ~ 2000000 events
- "Primakoff region": $t' < 10^{-3} (\text{GeV}/c)^2$ ~ 1000000 events



First Measurement of $\pi\gamma \rightarrow 3\pi$ Absolute Cross-Section



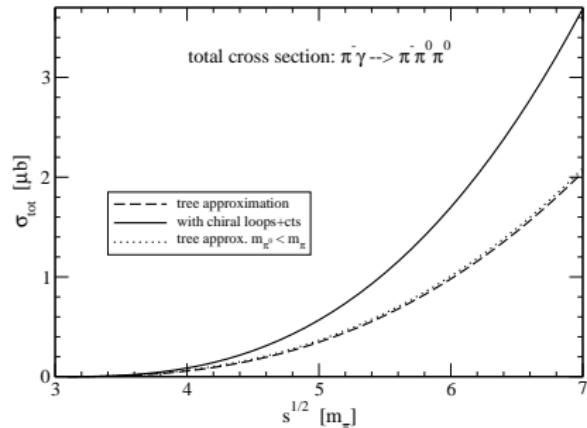
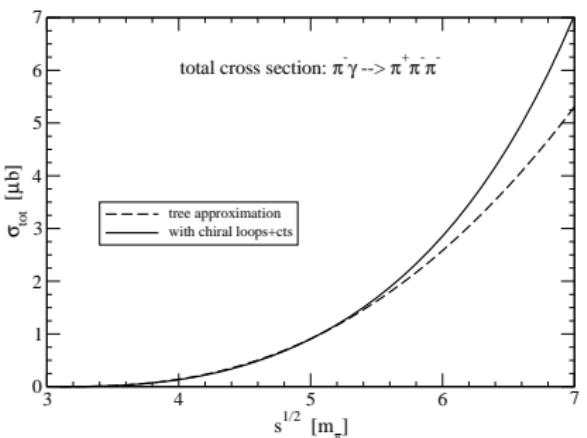
Measured absolute cross-section of $\pi^- \gamma \rightarrow \pi^- \pi^- \pi^+$



published in *PRL* 108 (2012) 192001

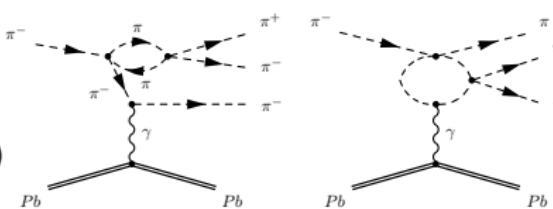


Higher-order effects

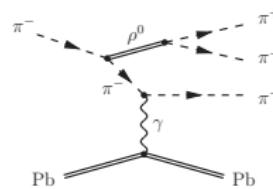


Chiral loops, e.g.

(N. Kaiser,
NPA848 (2010) 198)

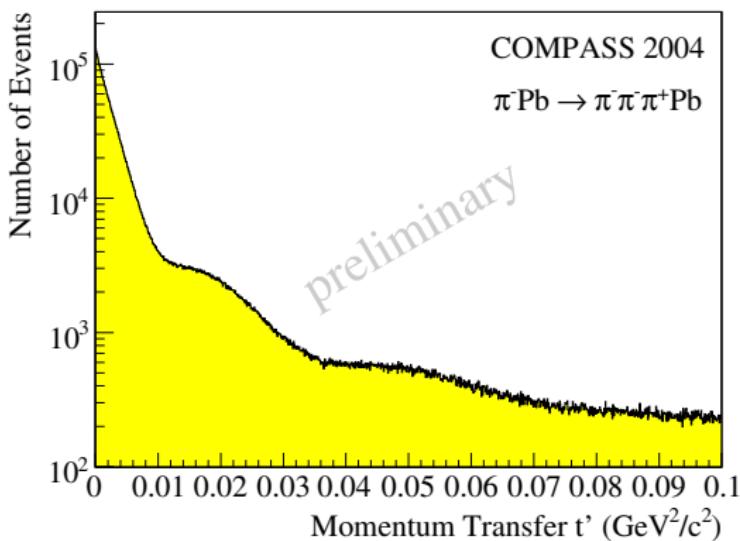


not (yet)
included:





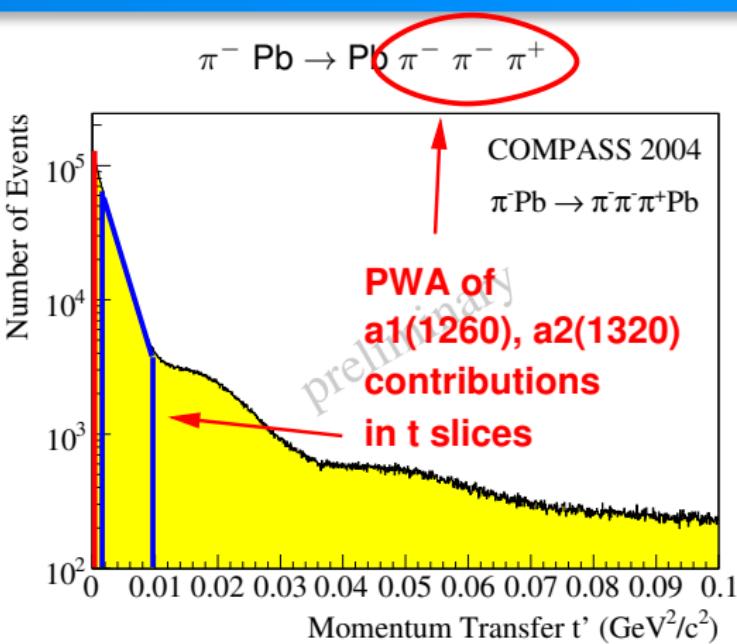
2004 Primakoff results on the resonances



- "Low t' ": $10^{-3} (\text{GeV}/c)^2 < t' < 10^{-2} (\text{GeV}/c)^2$ ~ 2000000 events
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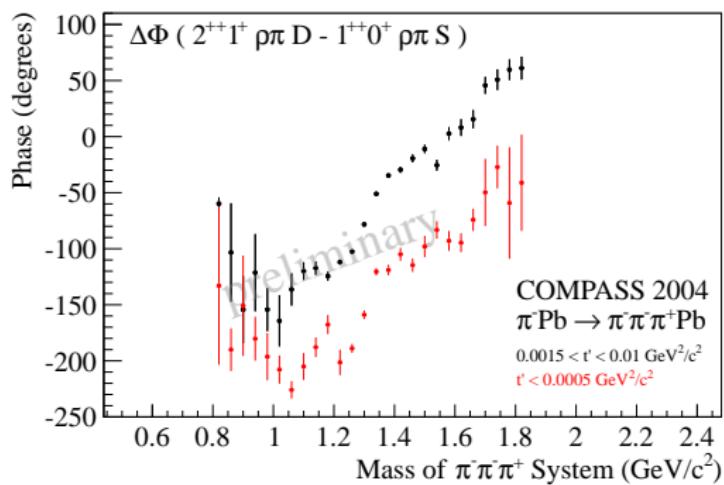
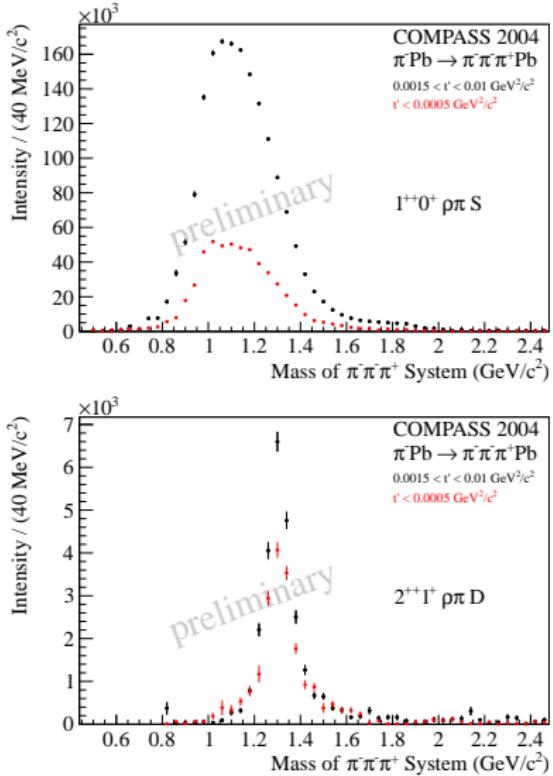
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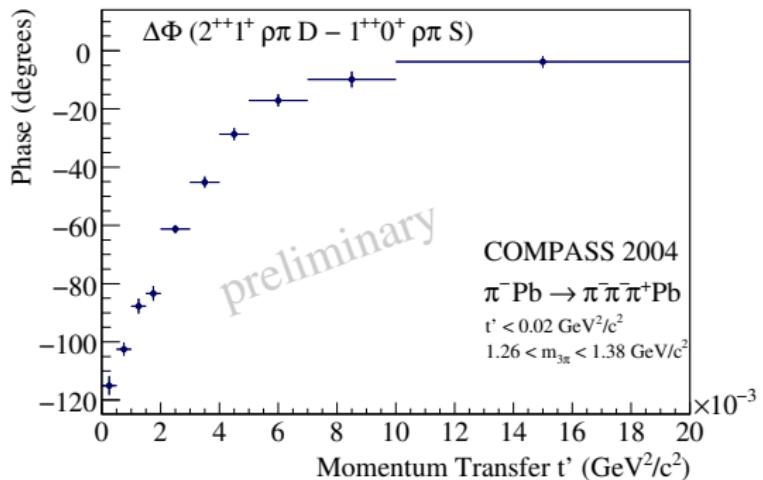
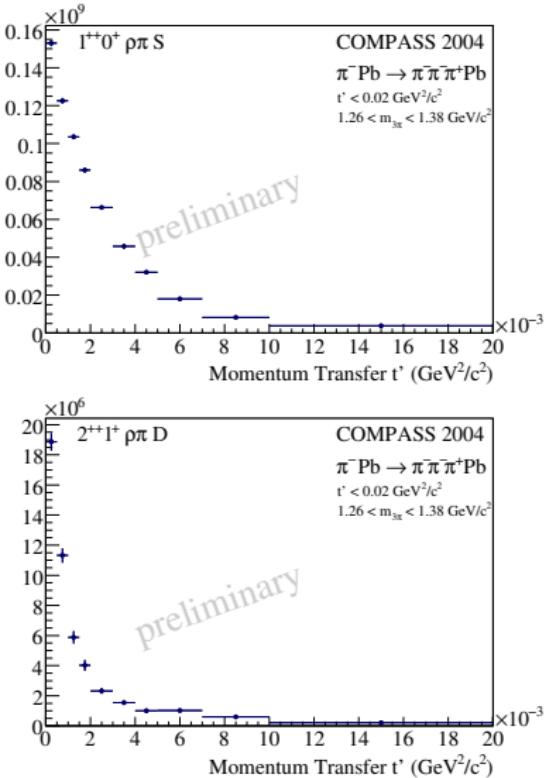


PWA: a_1 , a_2 and $\Delta\Phi$ in separated t' regions





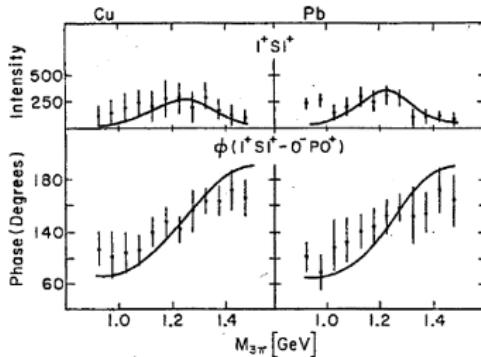
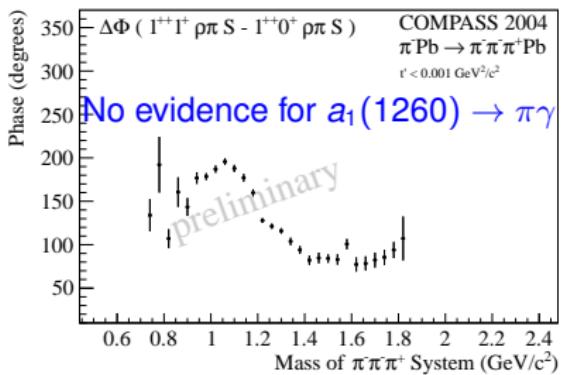
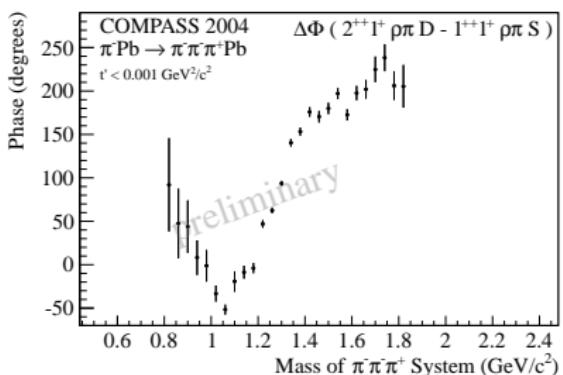
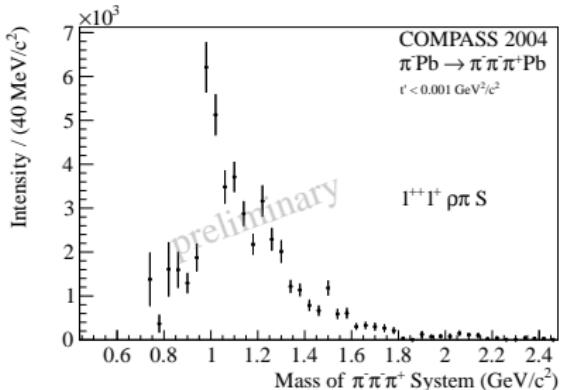
Phase $a_2 - a_1$ in detail: t' dependence



- transition of $\pi\gamma$ to $\pi IP \rightarrow a_2$ production
- work in progress
- interference can be used to map details of resonances and production mechanisms



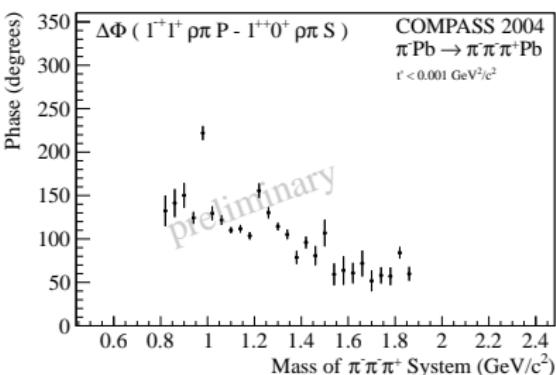
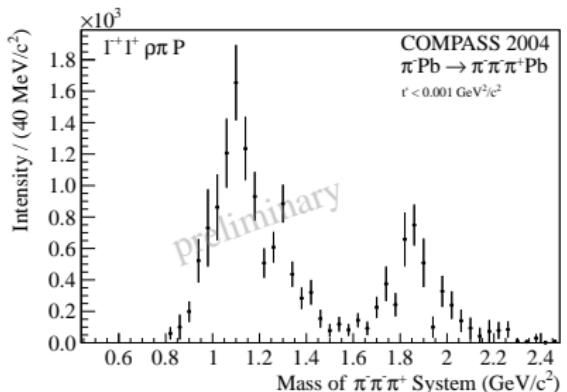
Primakoff production of $a_1(1260)$ vs. E272 result



M. Zielinski et al, Phys. Rev. Lett. vol. 52, 14, 1195



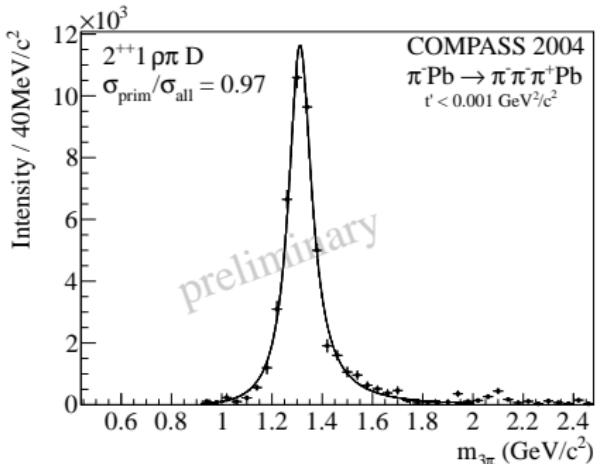
Spin-exotic 1⁻⁺



No evidence for $\pi_1(1600)$ Primakoff production



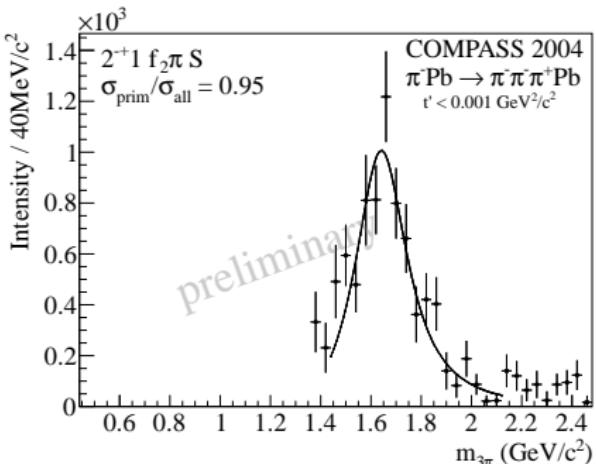
Radiative Coupling of $a_2(1320)$ and $\pi_2(1670)$



$\Gamma_0(a_2(1320) \rightarrow \pi\gamma)$ M2

↔ meson w.f.'s: $\Gamma_{i \rightarrow f} \propto | \langle \Psi_f | e^{-i\vec{q} \cdot \vec{r}} \hat{\epsilon} \cdot \vec{p} | \Psi_i \rangle |^2$, VMD

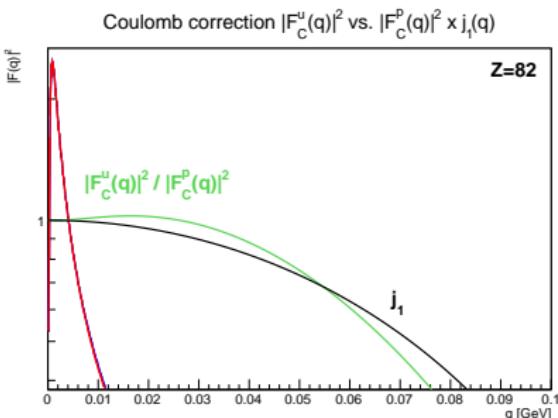
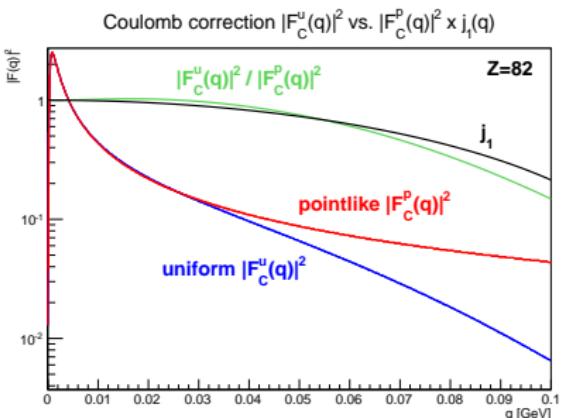
- normalization via beam kaon decays
- large Coulomb correction



$\Gamma_0(\pi_2(1670) \rightarrow \pi\gamma)$ E2



Coulomb and extended-charge correction



Calculations: G. Faeldt *et al.*,
 PhysRevC.79.014607 (2009), 82.037603 (2010), 87.029903 (2013)

Effect on $a_2(1320)$ radiative width:	form factor alone	+6.5%
	Coulomb correction alone	+26.0%
	combined	+22.2%



Radiative widths - experiment and theory

	$a_2(1320)$	$\pi_2(1670)$
COMPASS <i>preliminary</i>	$(358 \pm 6 \pm 42)$ keV	$(181 \pm 11 \pm 27)$ keV $\cdot \frac{0.56}{\text{BR}_{f_2 \pi}}$
COMPASS $F_{\text{eff}}^2 = j_1^2$ <i>prel.</i>	(312 ± 6) keV	(151 ± 9) keV $\cdot \frac{0.56}{\text{BR}_{f_2 \pi}}$
SELEX ¹ S. Cihangir <i>et al.</i> ² E. N. May <i>et al.</i> ³	$(284 \pm 25 \pm 25)$ keV (295 ± 60) keV (460 ± 110) keV	
VMD model ⁴ Relativ. Quark model ⁵ Cov. Osc. Quark model ⁶ Cov. Osc. Quark model ⁷	(375 ± 50) keV 324 keV 235 keV 237 keV	335 keV / 521 keV

¹V. V. Molchanov *et al.*, Phys. Lett. B 521, (2001) 171-180.

²S. Cihangir *et al.*, Phys. Lett. 117 B, (1982) 119-122.

³E. N. May *et al.*, Phys. Rev. D 16, (1977) 1983-1985.

⁴J. L. Rosner, Phys. Rev. D 23, (1981) 1127.

⁵I. G. Aznauryan and K. A. Oganesyan, Sov. J. Nucl. Phys. 47, (1988) 1097.

⁶S. Ishida, K. Yamada, and M. Oda, Phys. Rev. D 40, (1989) 1497-1512.

⁷T. Maeda, K. Yamada, M. Oda, and S. Ishida,
to be published in Prog. Theor. Exp. Phys, arXiv:1013.7507 [hep-ph].



Summary and Outlook



- Measurement of the pion polarisability at COMPASS
 - Most precise experimental determination
 - Systematic control: $\mu\gamma \rightarrow \mu\gamma$, $K^- \rightarrow \pi^-\pi^0$
- Chiral dynamics in $\pi\gamma \rightarrow \pi\pi\pi$ reactions
 - Charged-channel $\pi\gamma \rightarrow \pi^-\pi^-\pi^+$ tree-level ChPT prediction confirmed,
 - Neutral-channel $\pi\gamma \rightarrow \pi^-\pi^0\pi^0$ analysis ongoing
- Radiative coupling of meson resonances
 - Radiative width for $a_2(1320)$
roughly in agreement with earlier experimental determinations
 - First measurement of the radiative width for $\pi_2(1320)$,
order of magnitude as expected from theoretical models
 - no significant radiative coupling for $a_1(1260)$ or $\pi_1(1600)$
- Future
 - analysis of 2009 data for $\pi^-\pi^0\pi^0$
 - investigation of Coulomb-nuclear interference
 - COMPASS 2012 run with high statistics for neutral channels:
separate determination of α_π and β_π , quadrupole polarisabilities,
measurement of α_K