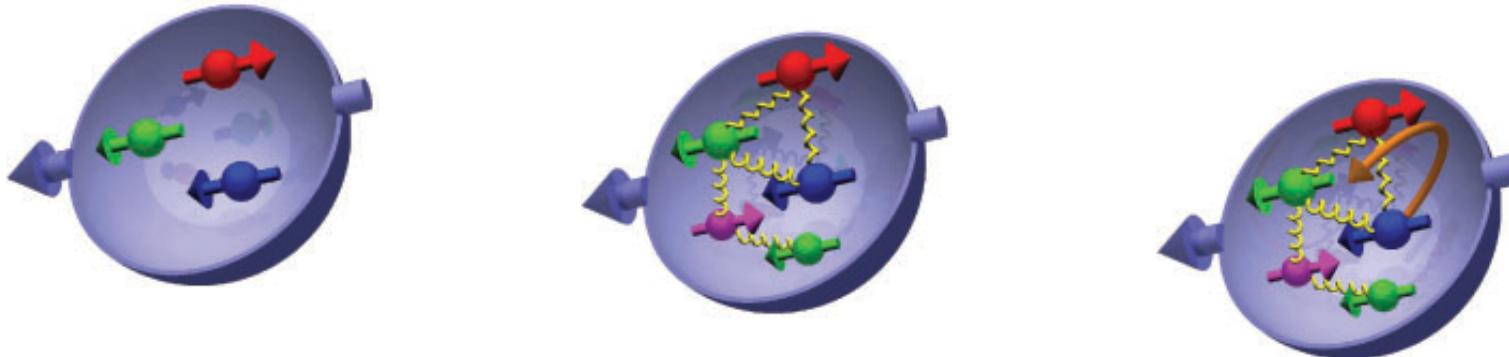


Status of polarized structure functions

C. Marchand, CEA Saclay

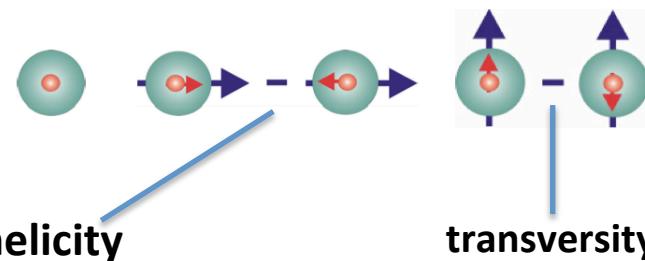


- **Polarized structure functions in longitudinal momentum space:** quark and gluon helicities from inclusive, semi-inclusive DIS and pp
- **Transverse momentum dependent (TMD) distribution functions** and relation to OAM from SIDIS, e^+e^- and pp

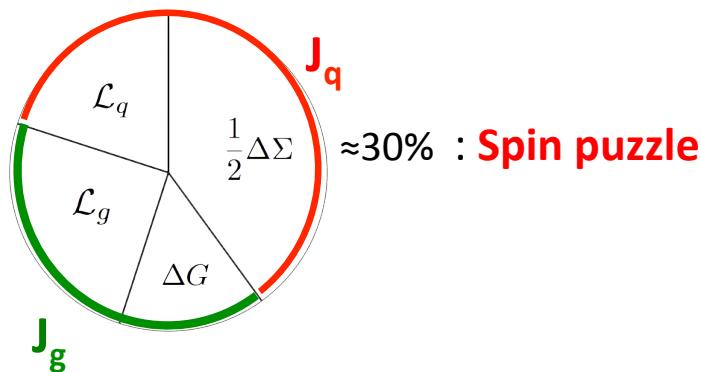
The spin of the nucleon

Three twist-2 quark DF's in collinear approximation ($\int dk_\perp$)

$$\Phi_{Coll}^{Tw-2}(x) = \frac{1}{2} \left\{ q(x) + S_L \gamma_5 \Delta q(x) + S_T \gamma_5 \gamma^1 \Delta_T q(x) \right\} n^+$$

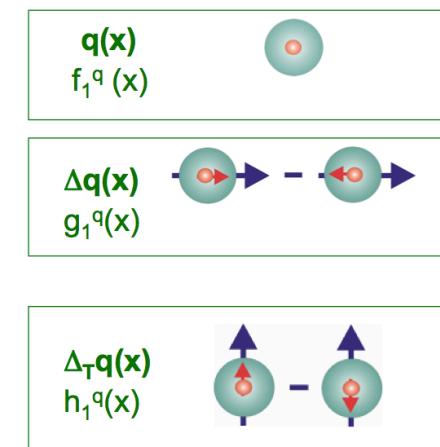
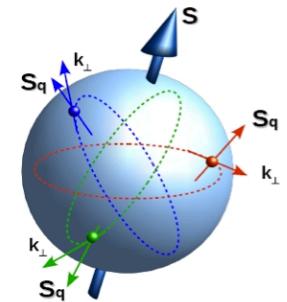


$$\frac{S_z^N}{\hbar} = \frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta G + L_z^q + L_z^g$$

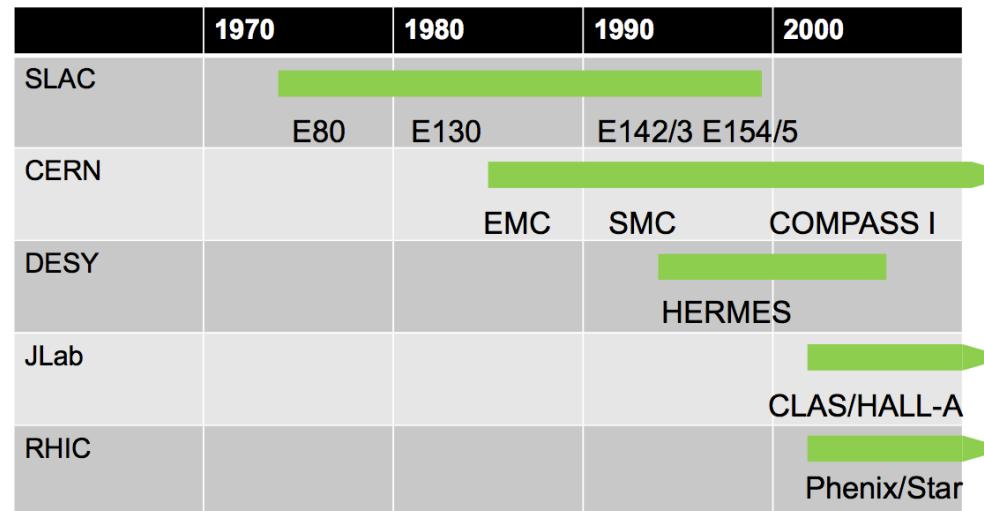
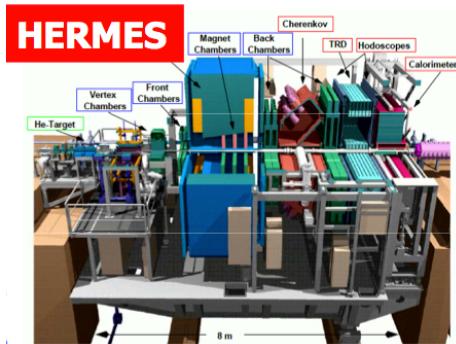


NR limit
[boost, rotat.] = 0

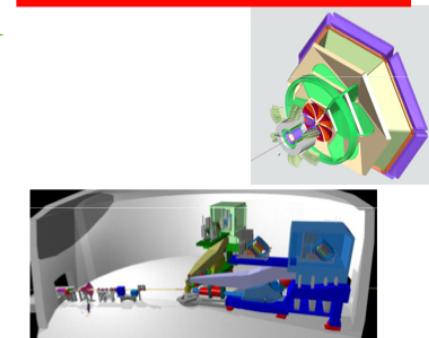
- $\Delta_T q(x, Q^2) = \Delta q(x, Q^2)$
- Expect non small $\Delta_T q$



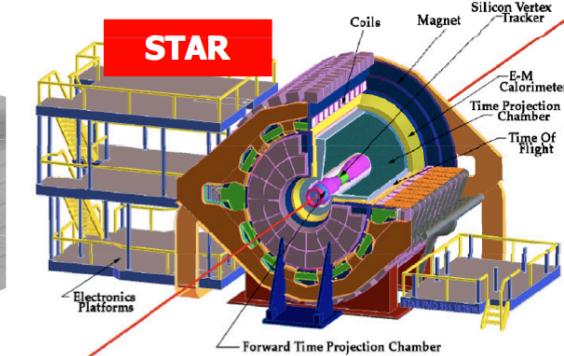
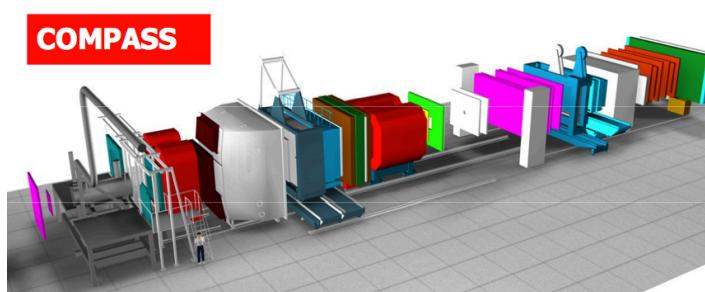
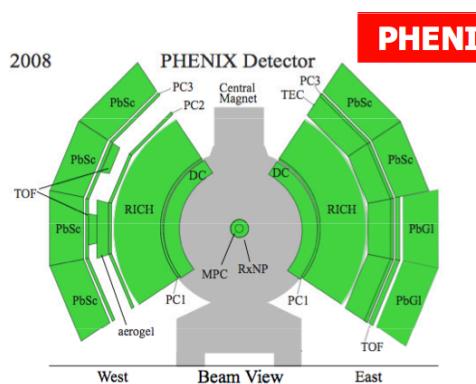
Experiments: $e(\mu)$ fixed target, pp collisions



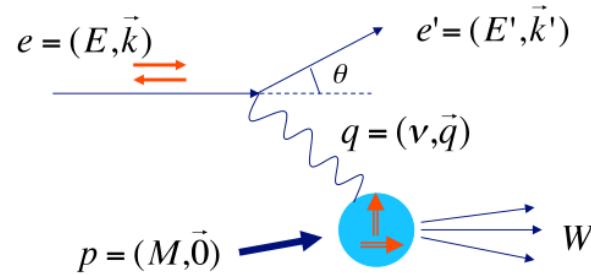
Jlab - CLAS, Hall A



◆ A worldwide effort since decades



Inclusive DIS polarized structure functions



$$A_{\parallel} = A_{LL} = \frac{1}{P_b P_T f} \frac{\vec{N} - \vec{\bar{N}}}{\vec{N} + \vec{\bar{N}}}$$

$$A_{\perp} = A_{LT} = \frac{1}{P_b P_T f} \frac{N^{\uparrow} - N^{\downarrow}}{N^{\uparrow} + N^{\downarrow}}$$

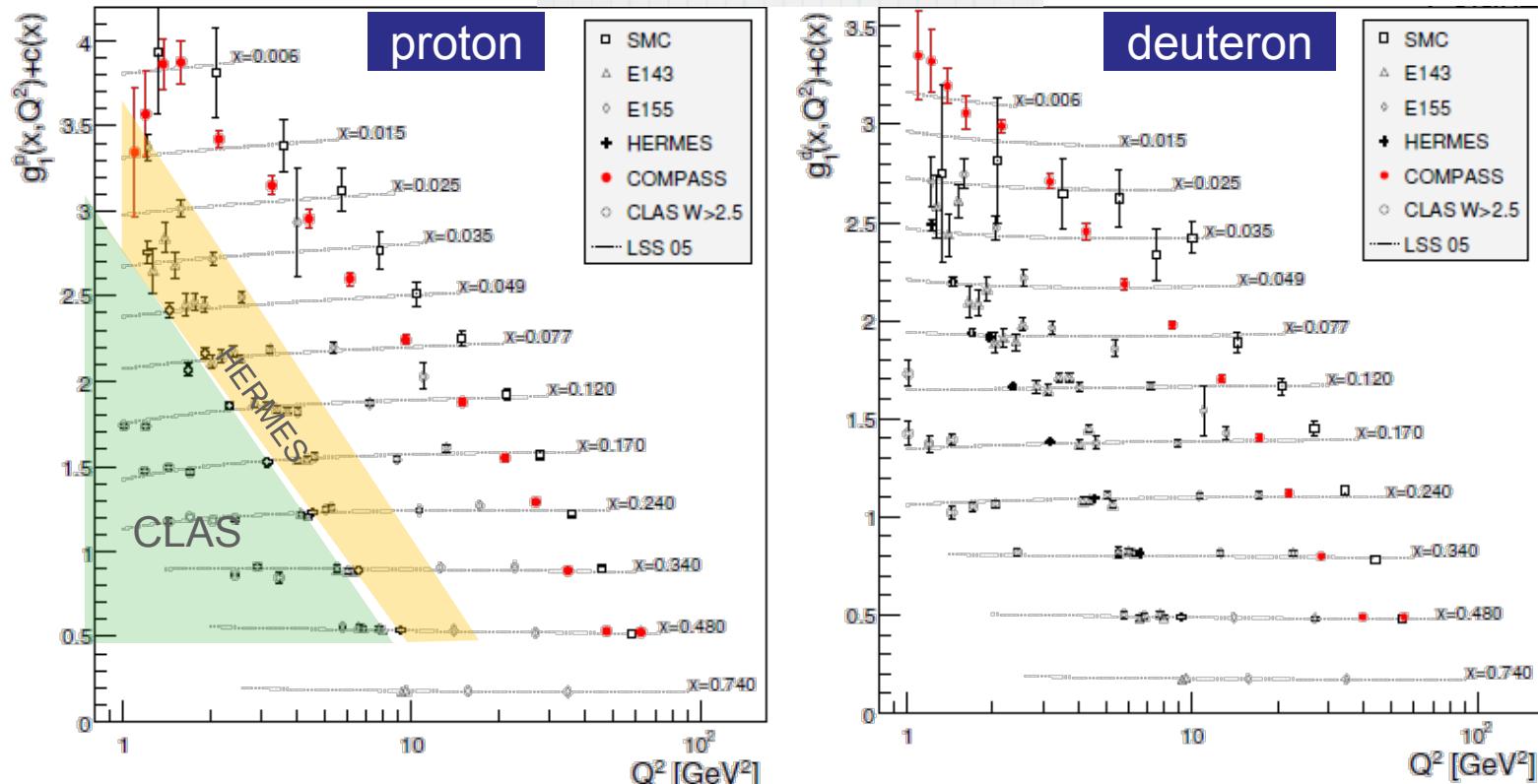
$$A_{\parallel} = D(A_1 + \eta A_2)$$

$$A_{\perp} = d(A_2 - \xi A_1)$$

$$A_1 = \frac{\sigma_{1/2}^T - \sigma_{3/2}^T}{\sigma_{1/2}^T + \sigma_{3/2}^T} = \frac{g_1 - \gamma^2 g_2}{F_1}$$

$$A_2 = \frac{2\sigma^{LT}}{\sigma_{1/2}^T + \sigma_{3/2}^T} = \gamma \frac{g_1 + g_2}{F_1}.$$

$$g_1(x, Q^2) = \frac{1}{2} \sum_q e_q^2 \Delta q(x, Q^2)$$



Inclusive DIS: extraction of quark and gluon helicities

$$g_1^p = \frac{1}{2} \left[\frac{4}{9} (\Delta u + \Delta \bar{u}) + \frac{1}{9} (\Delta d + \Delta \bar{d}) + \frac{1}{9} (\Delta s + \Delta \bar{s}) \right]$$

$$g_1^d = \frac{1}{2} \left[\frac{1}{9} (\Delta u + \Delta \bar{u}) + \frac{4}{9} (\Delta d + \Delta \bar{d}) + \frac{1}{9} (\Delta s + \Delta \bar{s}) \right]$$

Singlet: $\Delta \Sigma = [(\Delta u + \Delta \bar{u}) + (\Delta d + \Delta \bar{d}) + (\Delta s + \Delta \bar{s})]$

NS: $\Delta q_3 = [(\Delta u + \Delta \bar{u}) - (\Delta d + \Delta \bar{d})]$

NS: $\Delta q_8 = [(\Delta u + \Delta \bar{u}) + (\Delta d + \Delta \bar{d}) - 2(\Delta s + \Delta \bar{s})]$

$$\int g_1 dx \quad \Gamma_1^p = \int_0^1 g_1^p(x) dx; \quad \Gamma_1^d = \int_0^1 g_1^d(x) dx$$

► Moments $\Gamma_1^p - \Gamma_1^d = \frac{a_3}{6}(1 + \alpha^2 \text{corr})$ (Bjorken SR)

+ $a_3 = \Delta \Sigma_u - \Delta \Sigma_d = F + D = 1.267,$
 $a_8 = \Delta \Sigma_u + \Delta \Sigma_d - 2\Delta \Sigma_s = 3F - D \approx 0.58$
 from neutron and hyperon decays

$6(\Gamma_1^p - \Gamma_1^d)/(1 + \alpha^2 \text{corr}) = 1.28 \pm 0.07 \pm 0.10$

$\Delta \Sigma = a_0 = 0.33 \pm 0.03 \pm 0.05$ (evol. to $Q^2 = \infty$)

$(\Delta s + \Delta \bar{s}) = 1/3(a_0 - a_8) = -0.08 \pm 0.01 \pm 0.02$

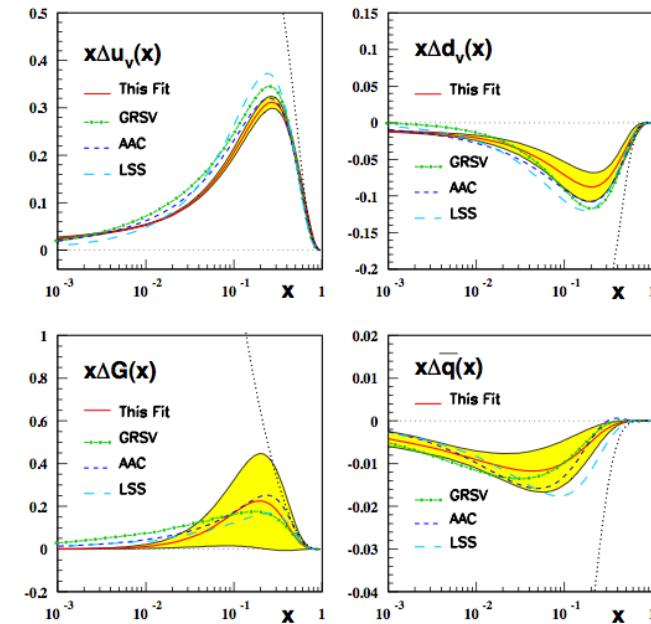
Compass

NLO: DGLAP links q and g

$$\frac{d}{d \ln Q^2} \Delta q_{NS}(x, Q^2) = \frac{\alpha_s(Q^2)}{2\pi} P_{qq}^{NS} \otimes \Delta q_{NS}$$

$$\frac{d}{d \ln Q^2} \begin{pmatrix} \Delta \Sigma \\ \Delta G \end{pmatrix} = \frac{\alpha_s(Q^2)}{2\pi} \begin{pmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{pmatrix} \otimes \begin{pmatrix} \Delta \Sigma \\ \Delta G \end{pmatrix}$$

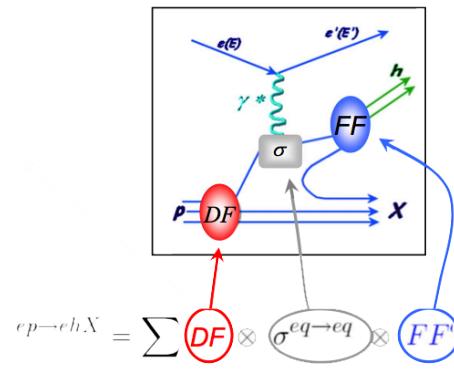
Assume SU(3) flavor symmetry: $\Delta \bar{u} = \Delta \bar{d} = \Delta \bar{s} = \Delta s$



Blümlein, Böttcher arXiv 1101.0052

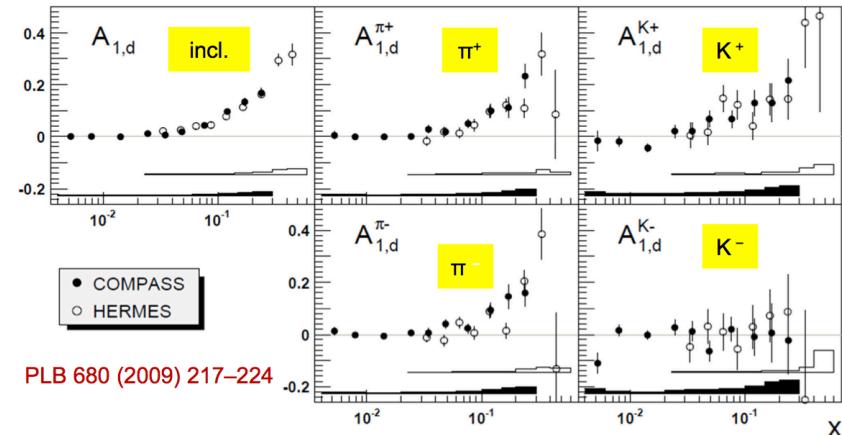
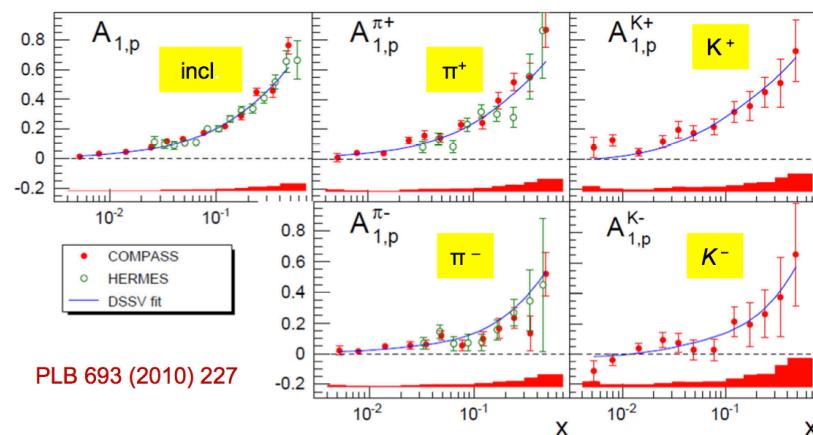
WG6: E. Nocera, A.L. Kataev

Semi-Inclusive DIS: extraction of quark helicities



$$A_1^{h(p/d)}(x, z, Q^2) \approx \frac{\sum_q e_q^2 \Delta q(x, Q^2) D_q^h(z, Q^2)}{\sum_q e_q^2 q(x, Q^2) D_q^h(z, Q^2)}$$

- Inputs needed for the extraction of $\Delta q(x, Q^2)$:
 - Unpolarised PDFs ($q(x, Q^2)$) → [MRST04](#)
 - $D_q^h(z, Q^2)$ → [DSS parameterisation](#)

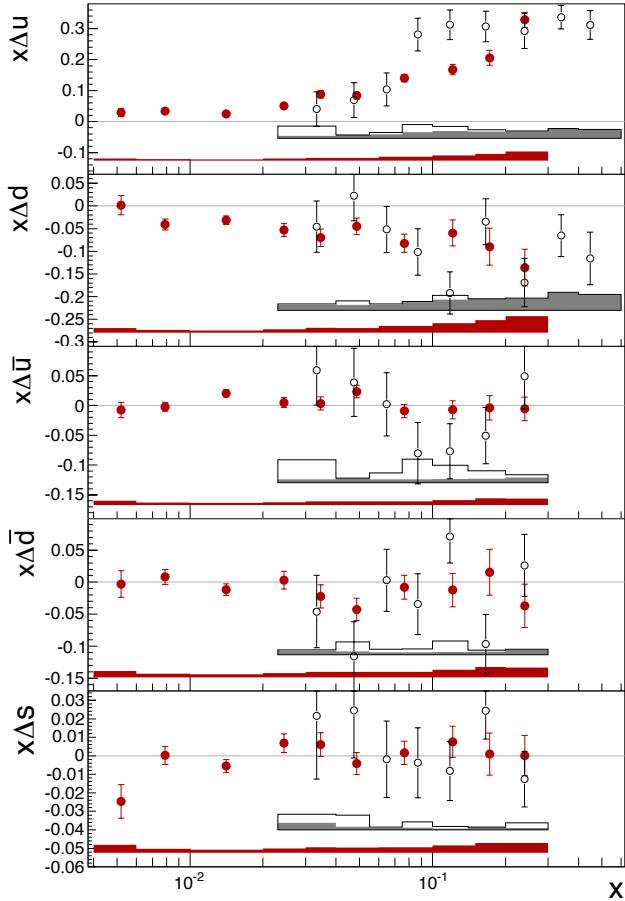


Leading Order (LO) fit of the 10 asymmetries (5d+5p)

Determine 6 flavor separated PDFs $\Delta u, \Delta d, \Delta \bar{u}, \Delta \bar{d}, \Delta s$ and $\Delta \bar{s}$

WG6: M. Stolarski

SIDIS: direct extraction of quark helicity at LO



● COMPASS

PLB693(2010)227, using DSS FF

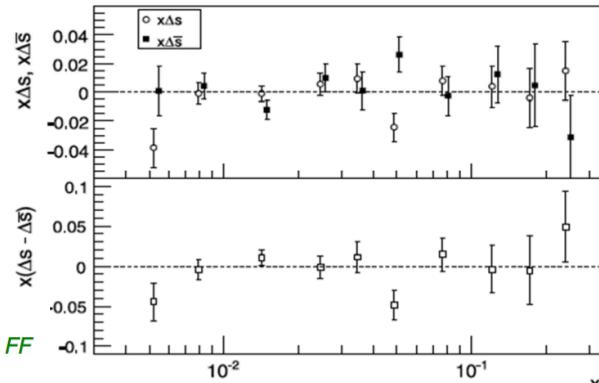
○ HERMES

PRD71(2005)012003

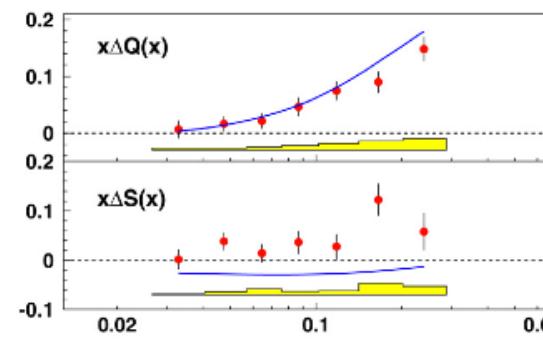
$Q^2 = 3 \text{ GeV}^2$

Check the assumption $\Delta s = \bar{\Delta s}$ (a 6 flavors fit)

COMPASS Collaboration / Physics Letters B 693 (2010) 227–235



HERMES: *PLB666(2008)446*

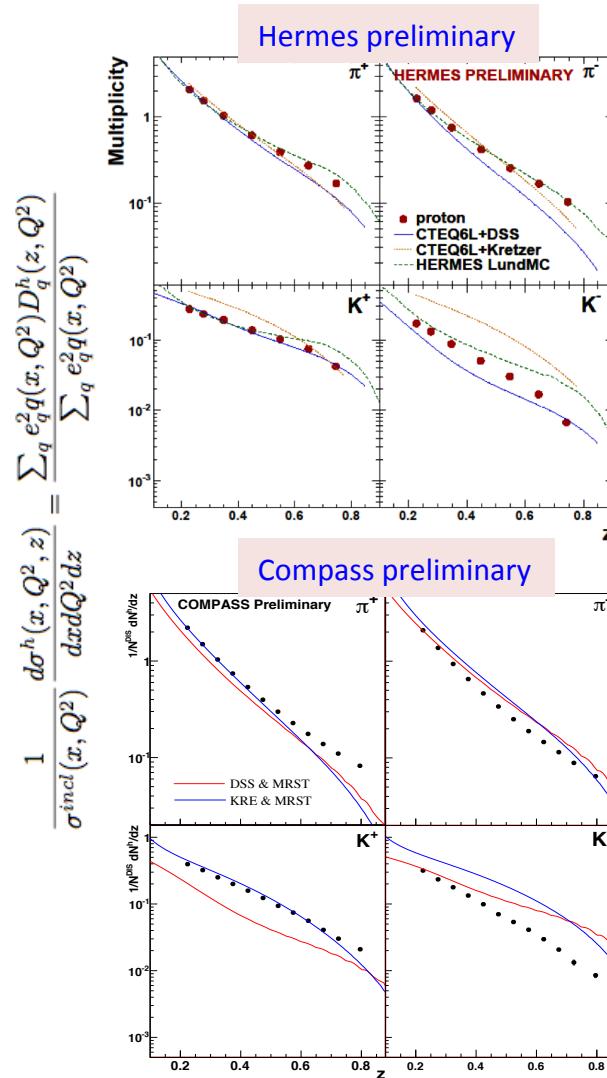
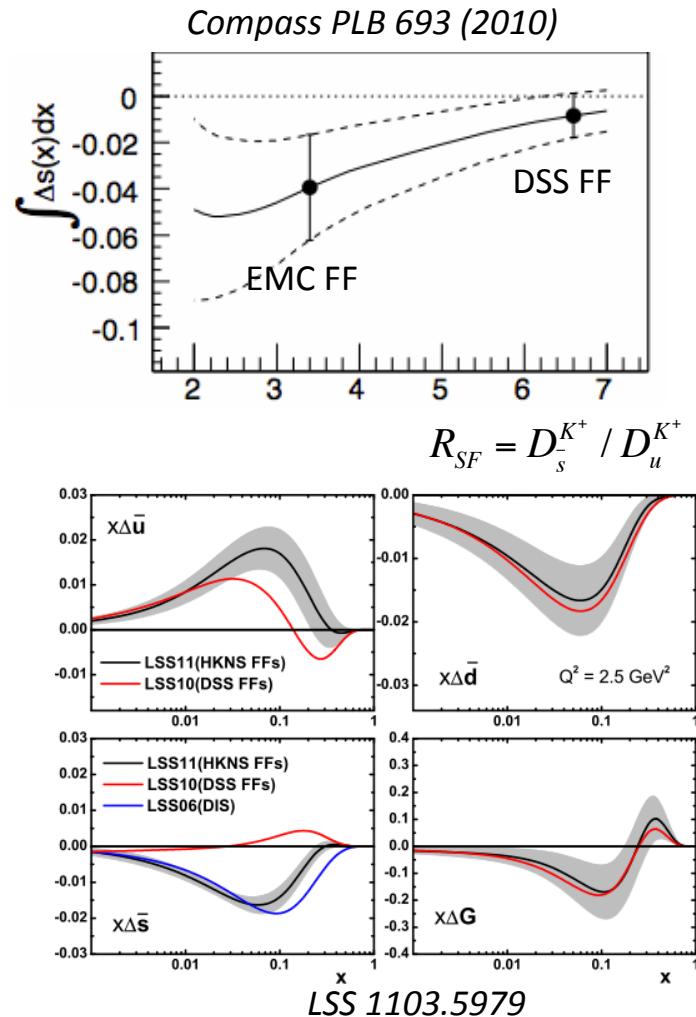


HERMES $\Delta s + \bar{\Delta s} = 0.037 \pm 0.019 \text{ (stat)} \pm 0.027 \text{ (syst)}$

COMPASS $\Delta s = -0.01 \pm 0.01 \text{ (stat)} \pm 0.01 \text{ (syst)}, 0.003 < x < 0.3$

DIS: $(\Delta s + \bar{\Delta s}) = -0.08 \pm 0.01 \pm 0.02$

Sensitivity of strange quark helicity to FF



LO interpretation

- OK with CTEQ6 pdfs + DSS FF for π^+ and K^+
- OK with CTEQ6 pdfs + Kretzer FF for π^+ and π^-
- poor agreement for K^-
- Role of unfavored FF
- Role of NLO term for negative particles

WG4: N. Makke

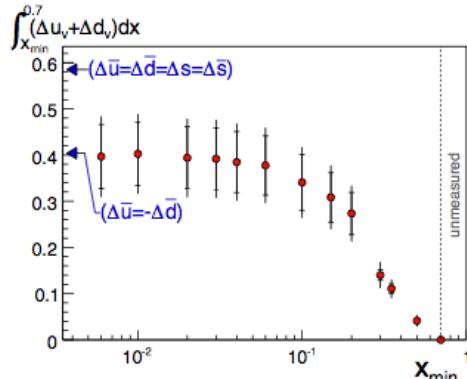
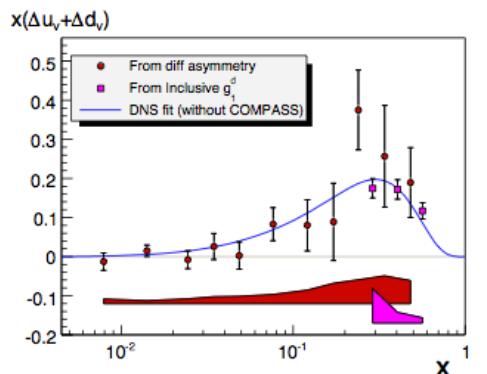
Polarized valence: difference asymmetries SIDIS

With charge conjugation symmetry in fragmentation $D_{1,q}^{h+} = D_{1,q}^{h-}$

$$A_{1d}^{h^+-h^-} = \frac{\Delta u_v + \Delta d_v}{u_v + d_v} (x)$$

K. Rith, Hera Symposium'11

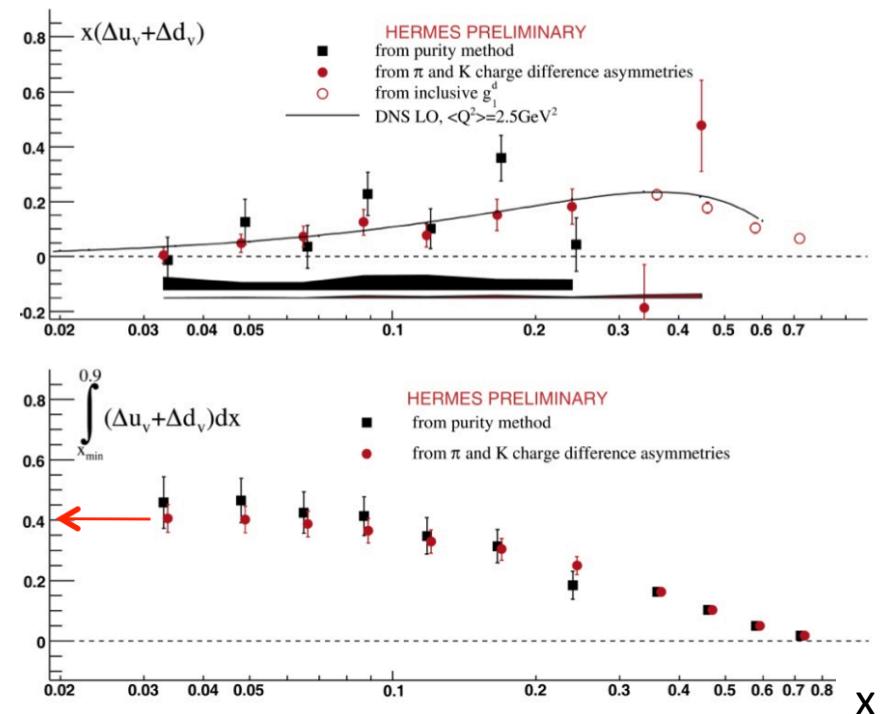
Compass d, PLB 660 (2008) 458



$$\Delta q_v = \Delta q - \Delta \bar{q}$$

$$\Gamma_v = \int_0^1 (\Delta u_v + \Delta d_v) dx = \Sigma_u + \Sigma_v - 2(\Delta \bar{u} + \Delta \bar{d})$$

$$a_8 = \Sigma_u + \Sigma_v - 2(\Delta s + \Delta \bar{s})$$

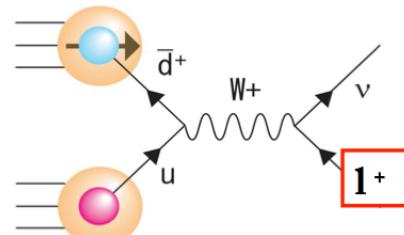


WG6: P. Kravchenko

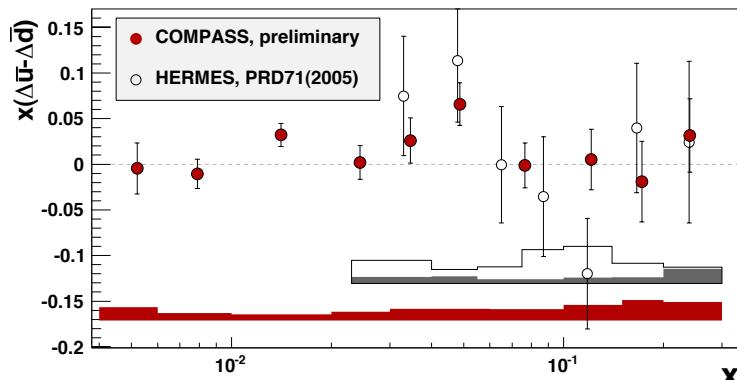
Polarized sea: parity violating W decay

$$p+p \rightarrow W^\pm \rightarrow e^\pm/\mu^\pm + \nu$$

- Parity violating W production:
 - Fixes quark helicity and flavor
- No fragmentation involved
- High Q² (set by W mass)



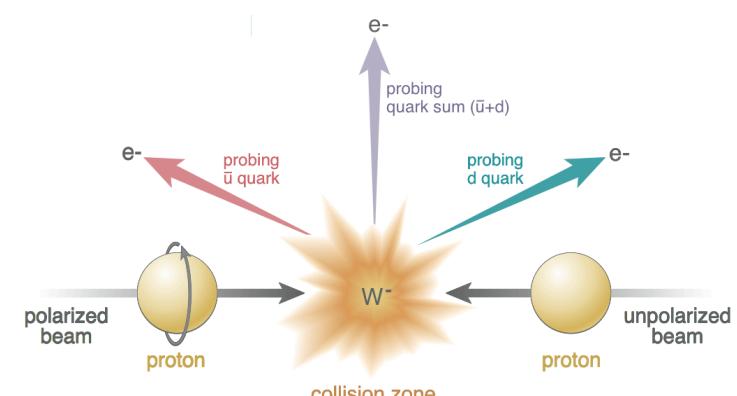
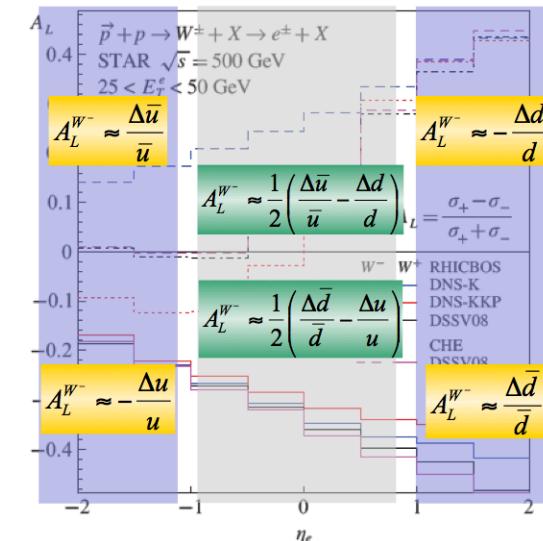
$$A_L^{W^+} = \frac{-\Delta u(x_a)\bar{d}(x_b) + \Delta \bar{d}(x_a)u(x_b)}{u(x_a)\bar{d}(x_b) + \bar{d}(x_a)u(x_b)}$$



COMPASS @ $Q^2=3(\text{GeV}/c)^2$: $\int_{0.004}^{0.3} (\Delta \bar{u} - \Delta \bar{d}) dx = 0.052 \pm 0.035(\text{stat}) \pm 0.013(\text{syst})$

HERMES @ $Q^2=2.5(\text{GeV}/c)^2$: $\int_{0.023}^{0.6} (\Delta \bar{u} - \Delta \bar{d}) dx = 0.048 \pm 0.057(\text{stat}) \pm 0.028(\text{syst})$

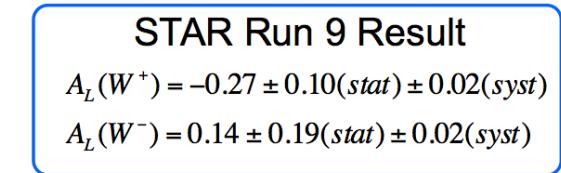
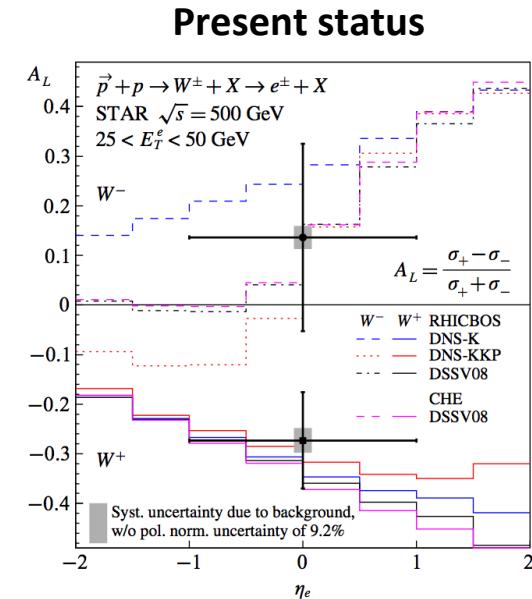
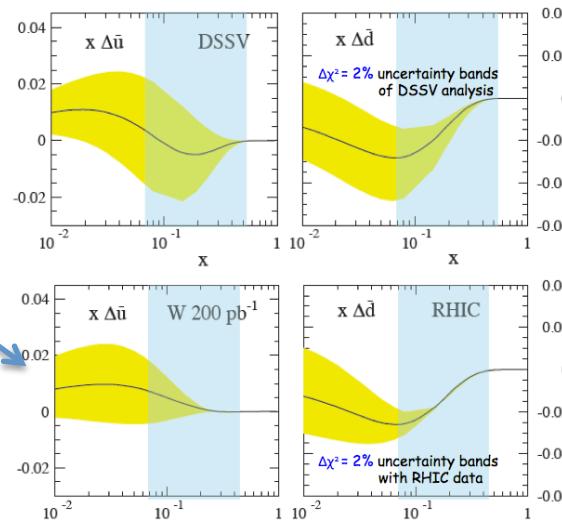
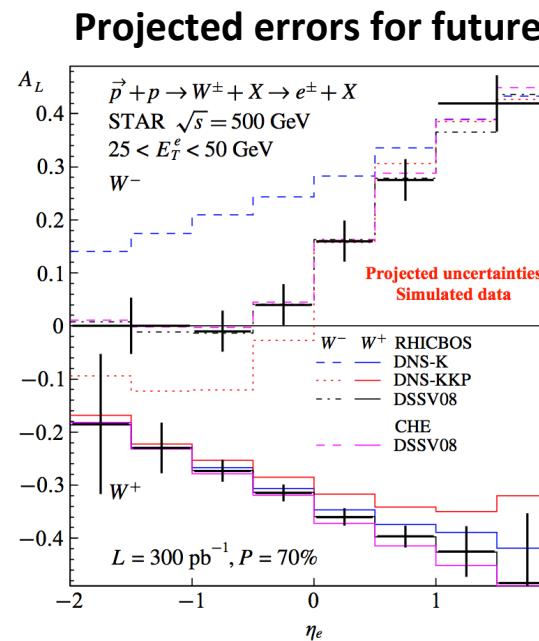
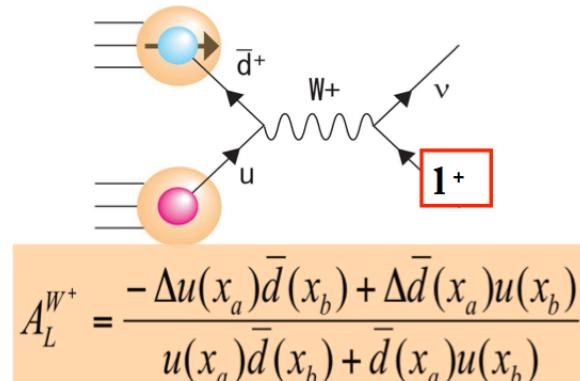
Unpolarized: $\int (\bar{u} - \bar{d}) dx = -0.118 \pm 0.012$



Polarized sea: parity violating W decay

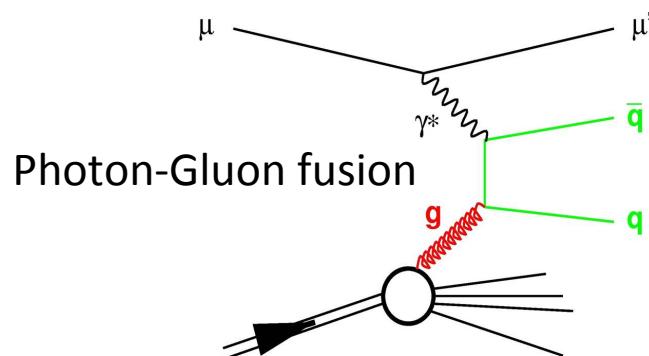
$$p+p \rightarrow W^\pm \rightarrow e^\pm/\mu^\pm + \nu$$

- Parity violating W production:
 - Fixes quark helicity and flavor
- No fragmentation involved
- High Q^2 (set by W mass)



WG6: B. Surrow, Y.J. Kim

Direct measurements of DeltaG/G: Open charm (NLO)



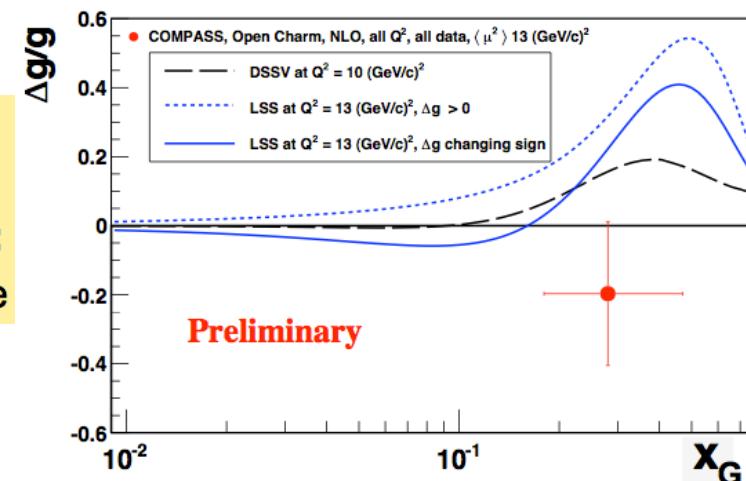
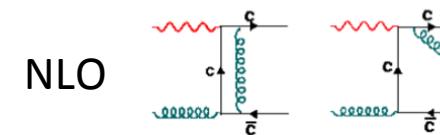
$$A_{LL}^{\mu N(LO)} = R^{PGF} a_{LL}^{PGF} \frac{\Delta g}{g}(x)$$

- First extraction of ΔG at NLO
- Constrains ΔG at larger x
- Charm result can be included in global NLO fits:
model independent asymmetries $A_{LL}(p_T, E_D)$ available

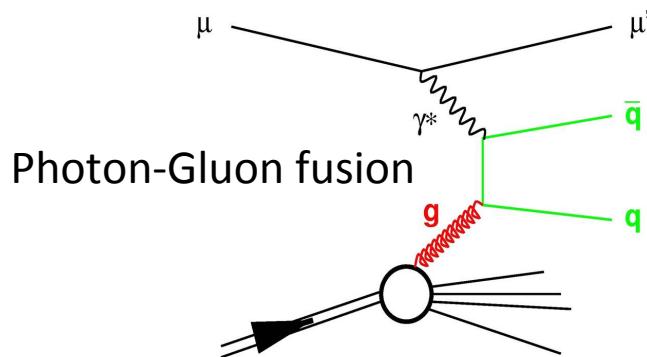
WG6: M. Stolarski

- **Open Charm production**

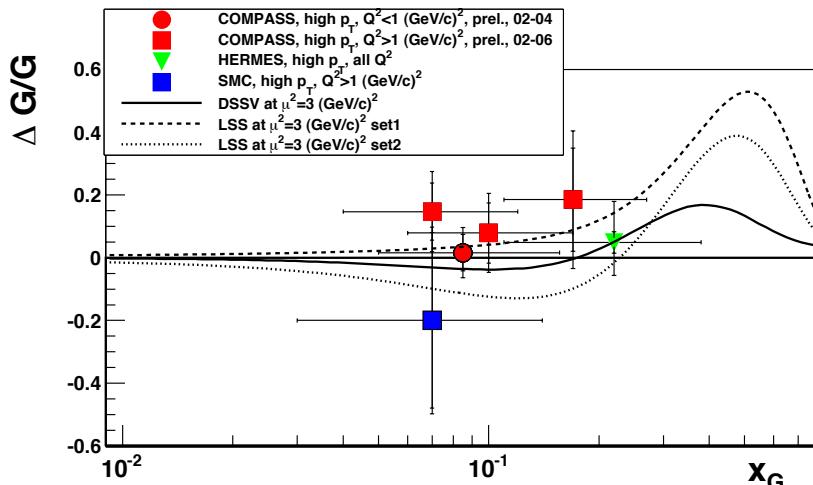
- $\gamma^* g \rightarrow c\bar{c} \Rightarrow$ reconstruct D⁰ mesons
- Hard scale: M_c^2
- No intrinsic charm in COMPASS kinematics
- No physical background
- Weakly Monte Carlo dependent
- Low statistics



Direct measurements of DeltaG/G: High pT (LO)



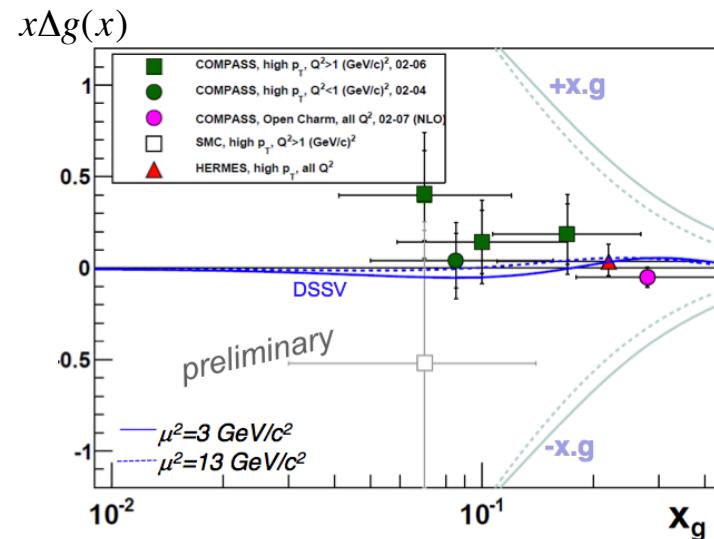
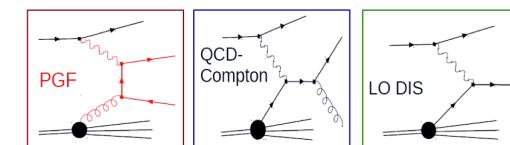
$$A_{LL}^{\mu N(LO)} = R^{PGF} a_{LL}^{PGF} \frac{\Delta g}{g}(x) + \dots$$



DSSV: D. de Florian et al., Phys. Rev. D80(2009)034030
 LSS: E. Leader, A.V. Sidorov, D.B. Stamenov, arXiv 1010.5742(2010)

- **High- p_T hadron pairs**

- $\gamma^* g \rightarrow q\bar{q} \Rightarrow$ reconstruct 2 jets or h^+h^-
- Hard scale: Q^2 or Σp_T^2 [$Q^2 > 1$ or $Q^2 < 1$ (GeV/c^2)²]
- High statistics
- Physical background
- Strongly Monte Carlo dependent

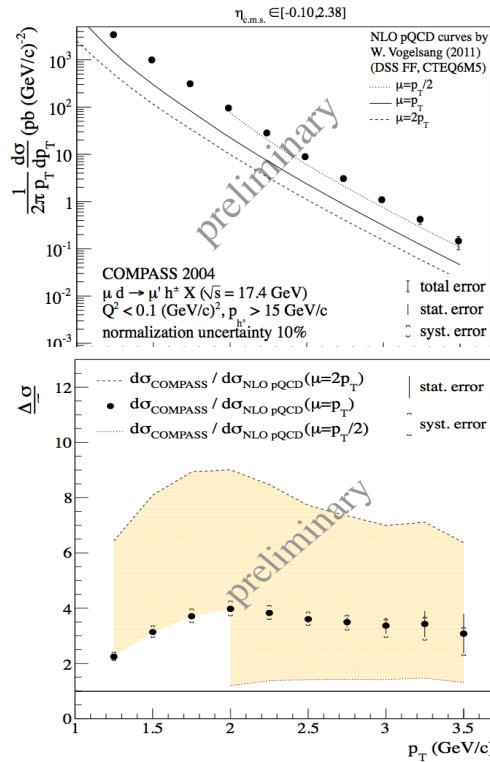


Charm at NLO, all other points at LO

DeltaG/G from hig pT: validity of NLO

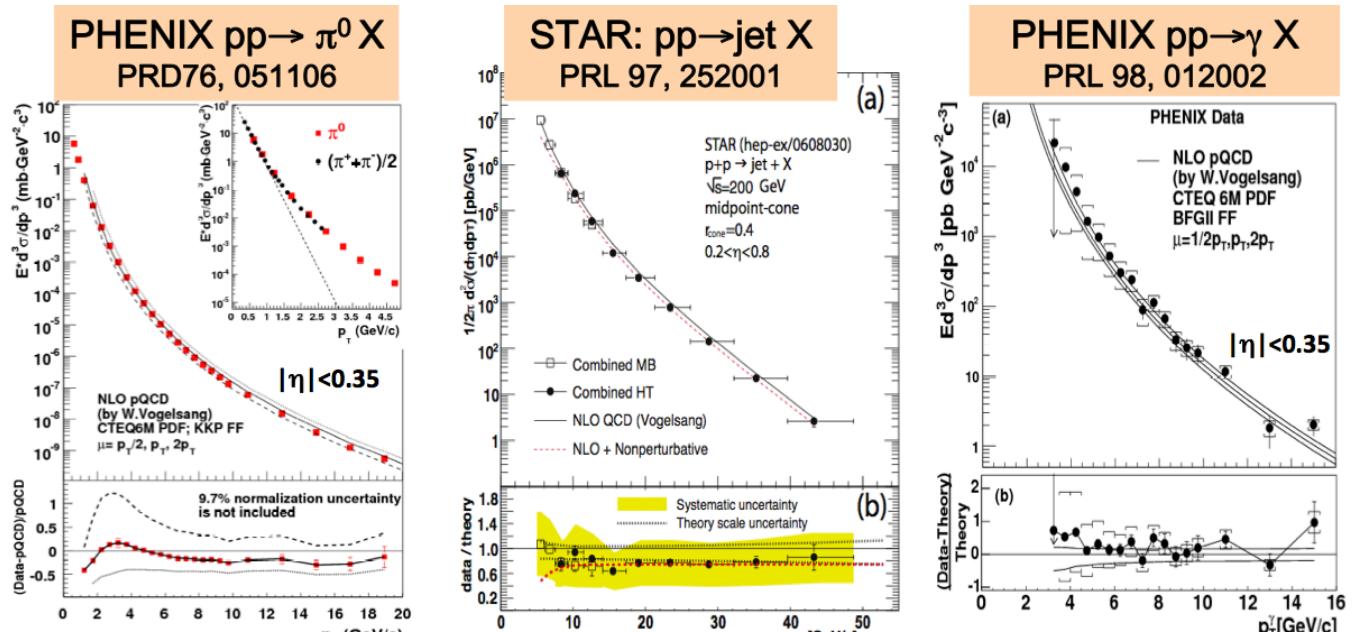
Unpolarized X-section as function of pT

COMPASS $\langle\sqrt{s}\rangle=17$ GeV



**WG6: C. Hoeppner
M. Pfeuffer**

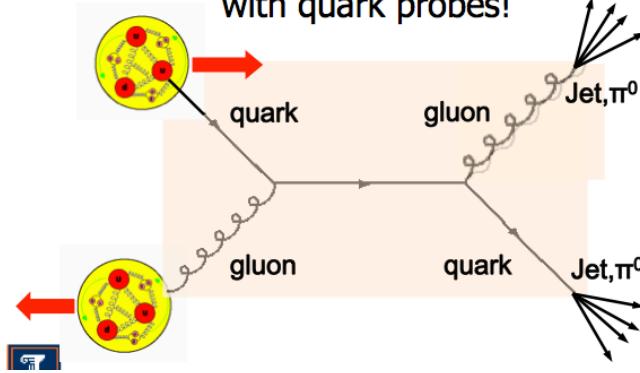
RHIC $\sqrt{s} = 200$ GeV



- CERN, HERMES data not included in NLO global fits yet
- RHIC CM energies OK for applicability of NLO to $A_{LL}(p_T)$

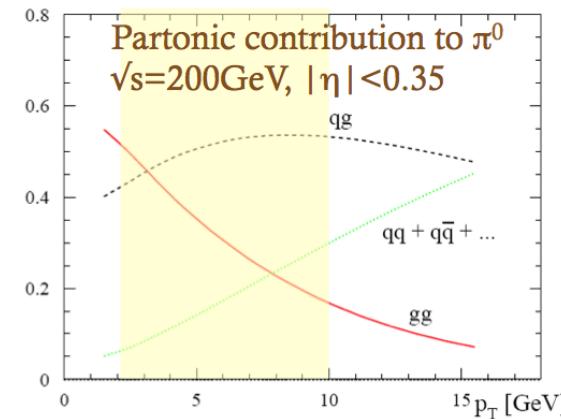
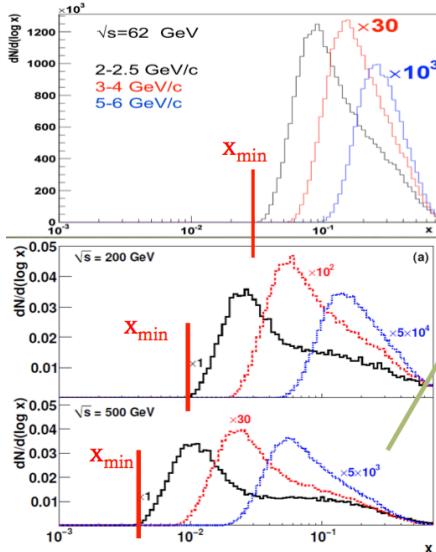
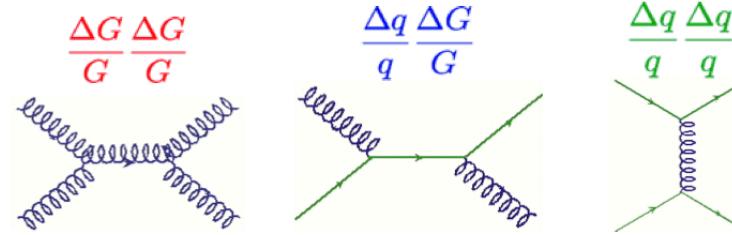
Direct measurements of DeltaG/G with pp->jets

Example: Production of neutral pions
 ~ probe gluon content with
 with quark probes!

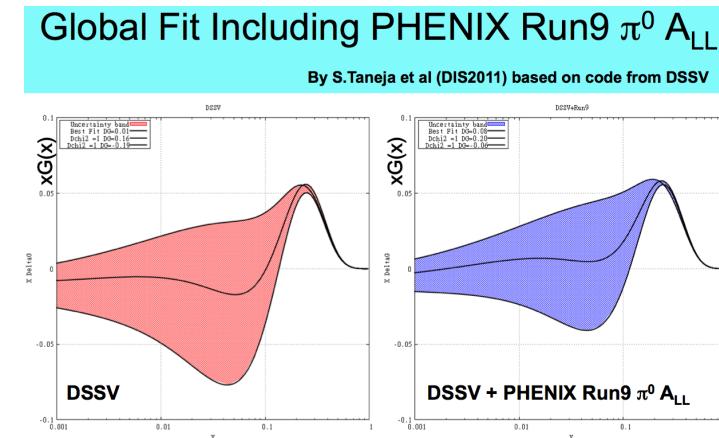
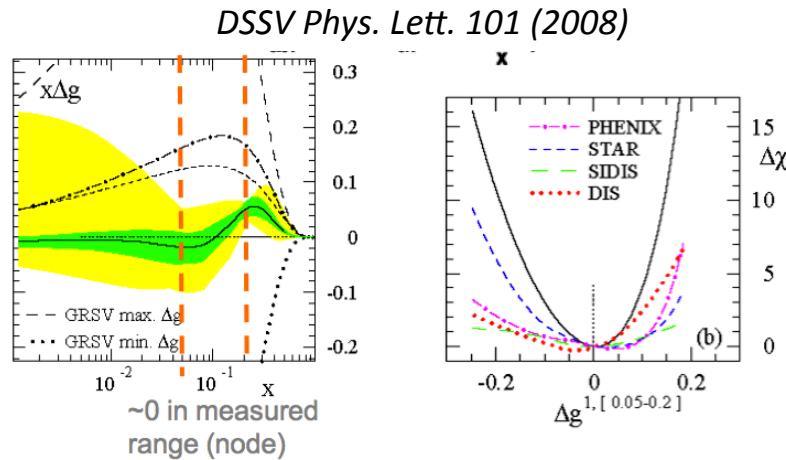
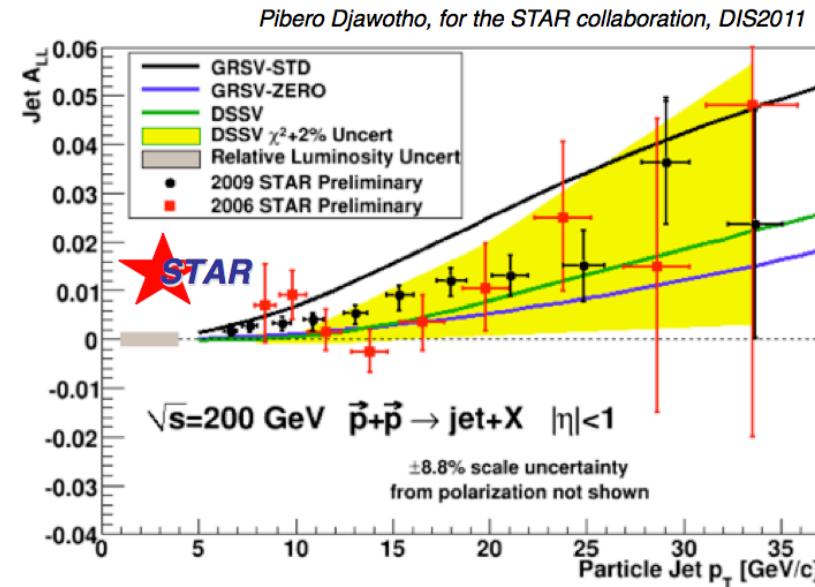
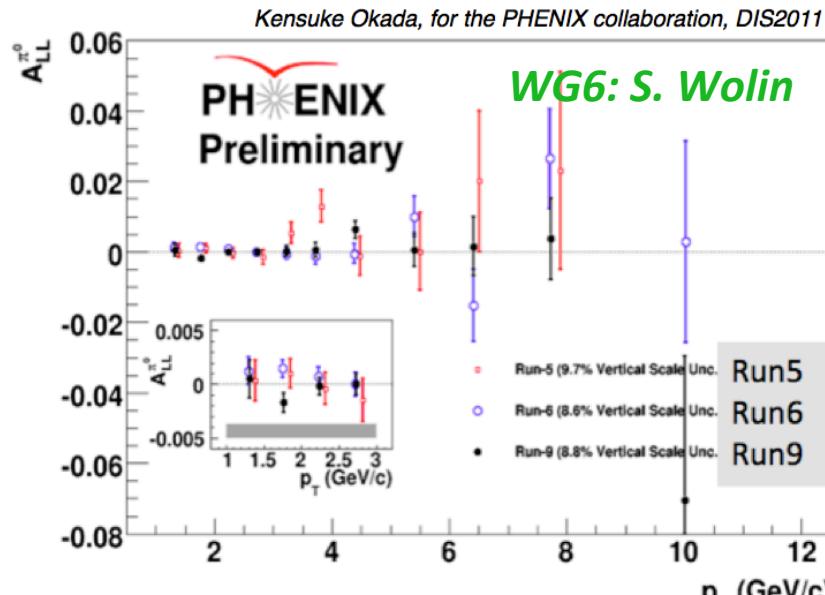


$$A_{LL} = \frac{\sigma^{++} - \sigma^{+-}}{\sigma^{++} + \sigma^{+-}} \propto \frac{\Delta f_a \Delta f_b}{f_a f_b} \hat{a}_{LL}$$

Δf : polarized parton distribution functions



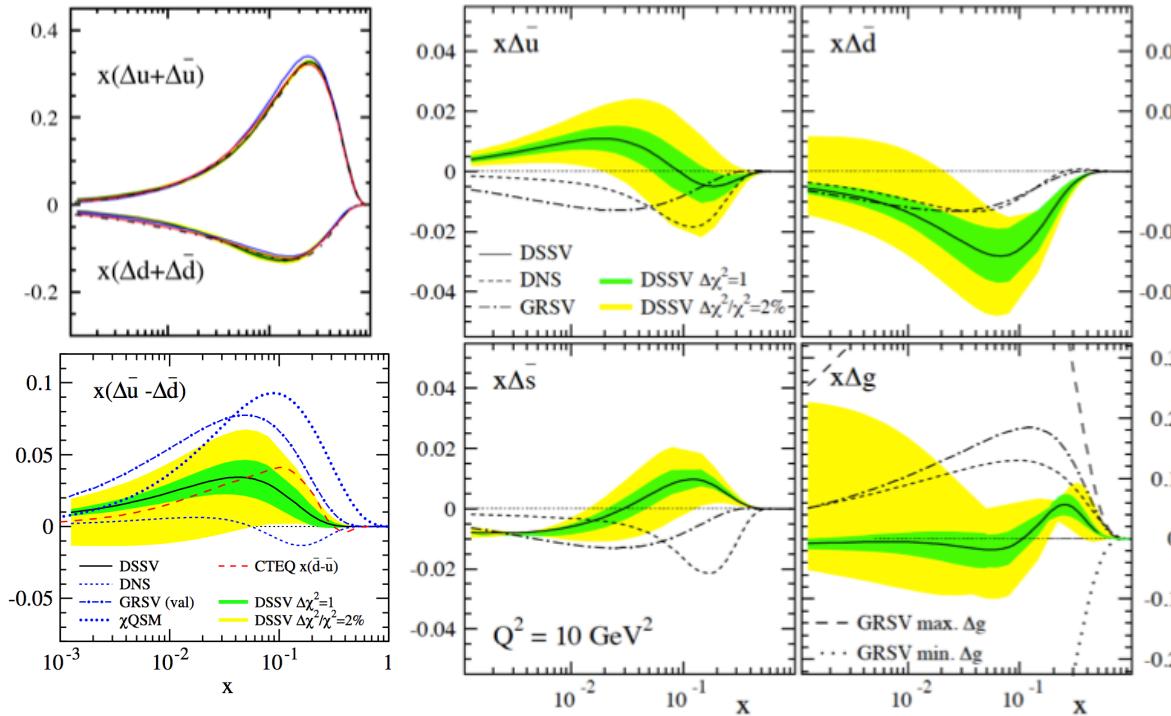
Direct measurements of DeltaG/G with pp-> jets



$$\int_{0.05}^{0.2} \Delta g(x) dx = 0.005^{+0.051}_{-0.058} (\Delta \chi^2 = 1); \int_{0.05}^{0.2} \Delta g(x) dx = 0.005^{+0.129}_{-0.164} (\Delta \chi^2 / \chi^2 = 2\%)$$

DIS+SIDIS+pp: global fit to extract q and g helicity

DSSV Phys. Lett. 101 (2008)



$$\int_{0.05}^{0.2} \Delta g(x) dx = 0.005^{+0.051}_{-0.058} (\Delta \chi^2 = 1); \int_{0.05}^{0.2} \Delta g(x) dx = 0.005^{+0.129}_{-0.164} (\Delta \chi^2 / \chi^2 = 2\%)$$

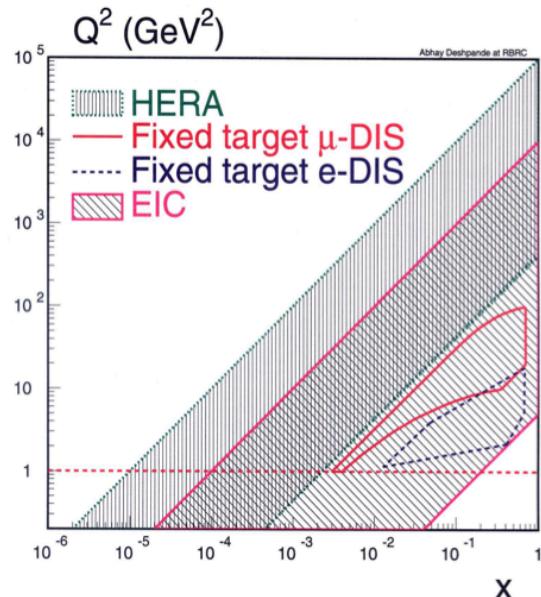
$$\int_0^1 \Delta g(x) dx = 0.013^{+0.106}_{-0.120} (\Delta \chi^2 = 1); \int_0^1 \Delta g(x) dx = 0.013^{+0.702}_{-0.314} (\Delta \chi^2 / \chi^2 = 2\%)$$

$$\frac{1}{2} = \frac{1}{2} \Delta \Sigma + \Delta G + L_z^q + L_z^g = 0.16 + \Delta G + L_z^q + L_z^g$$

Experiment	Process	N_{data}	χ^2
EMC [2]	DIS (p)	10	3.9
SMC [3]	DIS (p)	12	3.4
SMC [3]	DIS (d)	12	18.4
COMPASS [4]	DIS (d)	15	8.1
E142 [5]	DIS (n)	8	5.6
E143 [6]	DIS (p)	28	19.3
E143 [6]	DIS (d)	28	40.8
E154 [7]	DIS (n)	11	4.5
E155 [8]	DIS (p)	24	22.6
E155 [9]	DIS (d)	24	17.1
HERMES [10]	DIS (He)	9	6.3
HERMES [11]	DIS (p)	15	10.5
HERMES [11]	DIS (d)	15	16.9
HALL A [12]	DIS (n)	3	0.2
CLAS [13]	DIS (p)	10	5.9
CLAS [13]	DIS (d)	10	2.5
SMC [14]	SIDIS (p, h^+)	12	18.7
SMC [14]	SIDIS (p, h^-)	12	10.6
SMC [14]	SIDIS (d, h^+)	12	7.3
SMC [14]	SIDIS (d, h^-)	12	14.1
HERMES [15]	SIDIS (p, h^+)	9	6.4
HERMES [15]	SIDIS (p, h^-)	9	4.9
HERMES [15]	SIDIS (d, h^+)	9	11.4
HERMES [15]	SIDIS (d, h^-)	9	4.5
HERMES [10]	SIDIS (He, h^+)	9	4.7
HERMES [10]	SIDIS (He, h^-)	9	6.9
HERMES [15]	SIDIS (p, π^+)	9	9.6
HERMES [15]	SIDIS (p, π^-)	9	4.9
HERMES [15]	SIDIS (d, π^+)	9	9.4
HERMES [15]	SIDIS (d, π^-)	9	19.5
HERMES [15]	SIDIS (d, K^+)	9	6.2
HERMES [15]	SIDIS (d, K^-)	9	5.8
HERMES [15]	SIDIS (d, $K^+ + K^-$)	9	3.4
COMPASS [16]	SIDIS (d, h^+)	12	6.2
COMPASS [16]	SIDIS (d, h^-)	12	12.0
PHENIX [22]	pp (200 GeV, π^0)	10	14.2
PHENIX [23]	pp (200 GeV, π^0)	10	7.1 [13.8] ^a
PHENIX [24]	pp (62 GeV, π^0)	5	3.1 [2.8] ^a
STAR [25]	pp (200 GeV, jet)	10	8.8
STAR (prel.) [26]	pp (200 GeV, jet)	9	6.9
TOTAL:		467	392.6

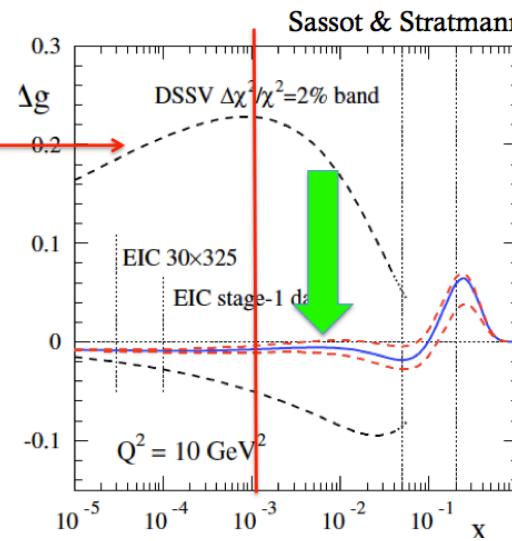
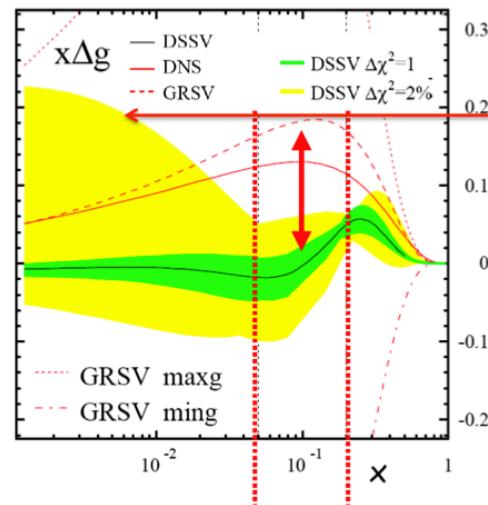
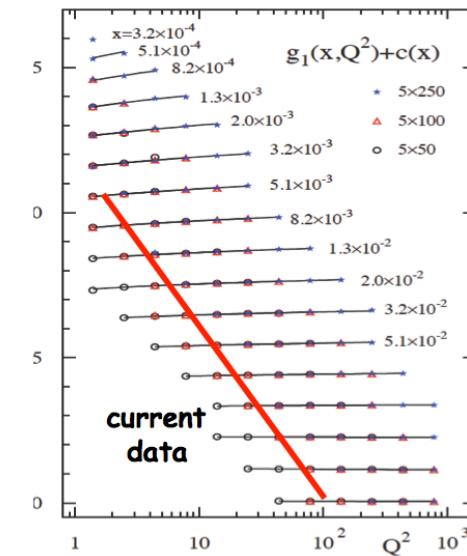
WG6: W. Vogelsang

Constrain gluon helicity at future EIC



- E_e : 5-30 GeV
- E_p : 50-325 GeV
- \sqrt{s} : 30-200 GeV
- $x_{\min} = 10^{-4}$; $Q^2_{\max} = 10^4 \text{ GeV}^2$
- Polar: ~70% (e,p,3He,D)
- Lumi: $> 10^{33} \text{ cm}^{-2}\text{s}^{-1}$

$$\frac{dg_1(x, Q^2)}{d \ln Q^2} \propto -\Delta g(x, Q^2)$$



WG6+7: M. Stratmann

g2 structure function

$$A_2 = \frac{2\sigma^{LT}}{\sigma_{1/2}^T + \sigma_{3/2}^T} = \gamma \frac{g_1 + g_2}{F_1}$$

$$g_2(x, Q^2) = g_2^{WW}(x, Q^2) + \bar{g}_2(x, Q^2)$$

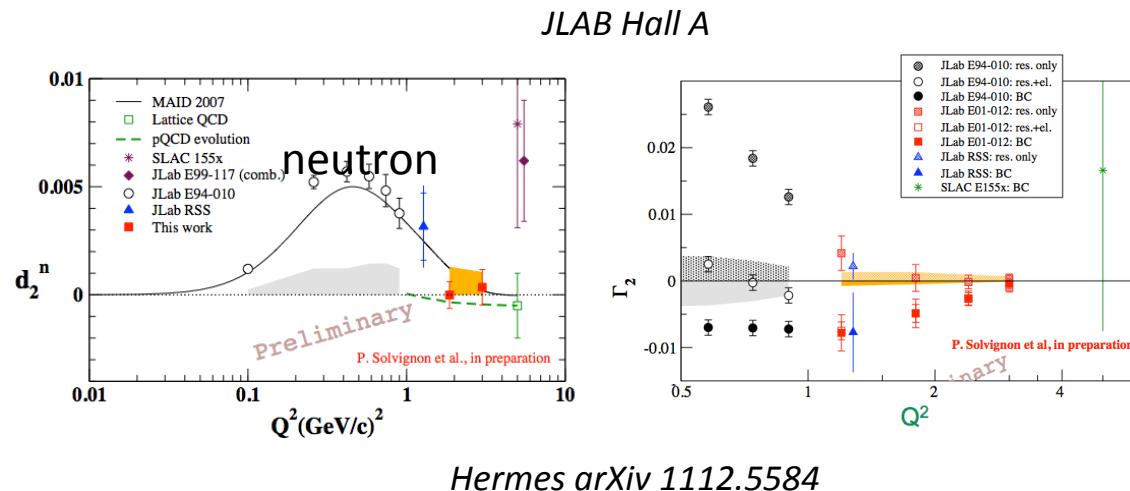
$$g_2^{WW}(x, Q^2) = -g_1(x, Q^2) + \int_x^1 g_1(y, Q^2) \frac{dy}{y}$$

Higher TW: quark-gluon correlations

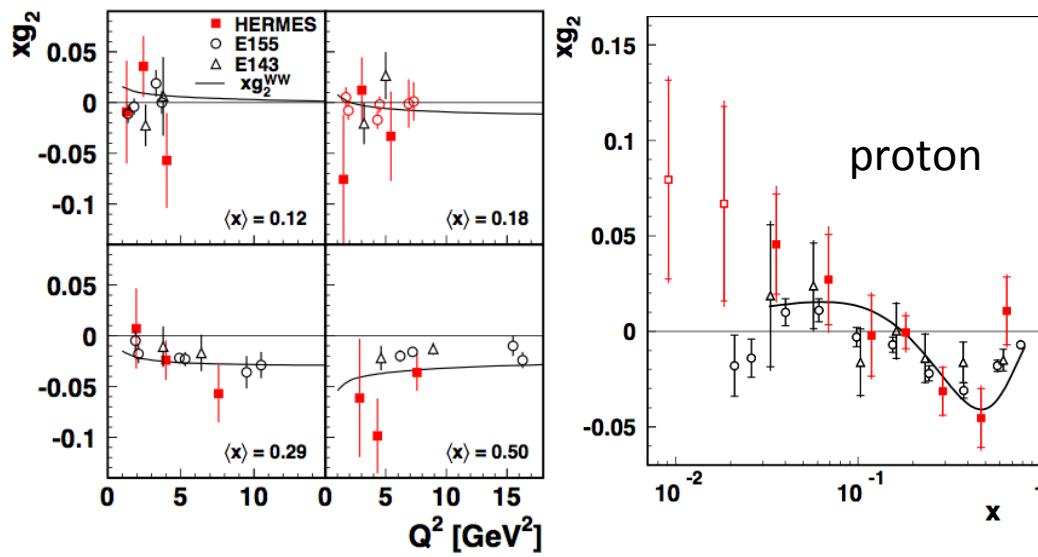
$$d_2(Q^2) = 3 \int_0^1 x^2 \bar{g}_2(x, Q^2) dx$$

Burkhardt-Cottingham Sum Rule

$$\int_0^1 g_2(x, Q^2) dx = 0$$



Hermes arXiv 1112.5584



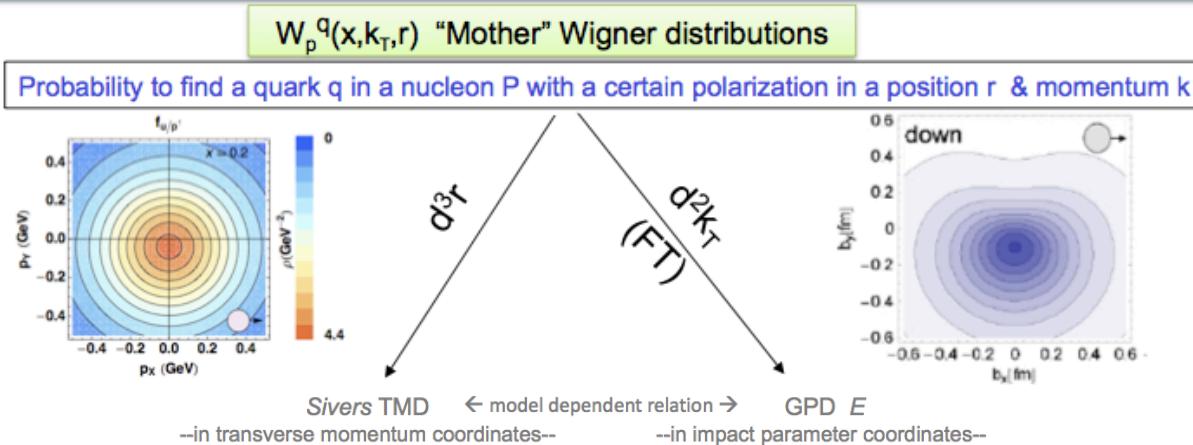
$$d_2^p(Q^2 = 5 \text{ GeV}^2) = 0.0148 \pm 0.0096(\text{stat}) \pm 0.0048(\text{syst})$$

$$\int_{0.023}^{0.9} g_2^p(x, Q^2 = 5 \text{ GeV}^2) = 0.006 \pm 0.024(\text{stat}) \pm 0.017(\text{syst})$$

Quark distributions in nuclei: 3D picture

Quantum phase-space distributions of quarks

Bacchetta et al.
PRD78(2008)



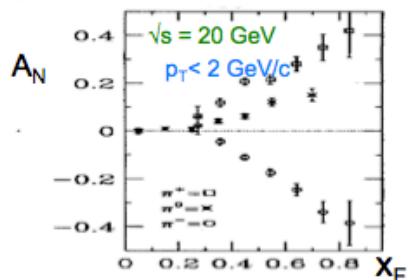
Göckeler et al.
PRL98(2007)

TMD (x, k_T)
spin-orbit
correlations

TMD PDFs: $f_p^u(x, k_T, \dots)$

Semi-inclusive measurements
Momentum transfer to quark
Direct info about momentum distribution

May explain SSA & Lam-Tung



GPDs: $H_p^u(x, \xi, t), \dots$

Exclusive Measurements
Momentum transfer to target
Direct info about spatial distribution

GPD(x, ξ, t)
2+1D imaging;
access to OAM

PDFs $f_p^u(x, \dots)$

May solve
proton spin puzzle

$$J_q = \frac{1}{2} \Delta \Sigma + L_q = \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H(x, \xi, t) + E(x, \xi, t)]$$

Transverse Momentum Dependent (TMD) DF

N/q	U	L	T
U	f_1		h_1^\perp
Twist-2 DF	L	g_1	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}^\perp	h_1
			h_{1T}^\perp

$h_1 \neq g_1$: relativistic effects and no mix with gluons in spin $\frac{1}{2}$ nucleon

Survive k_T integration

N/q	U	L	T
U	f^\perp	g^\perp	h, e
L	f_L^\perp	g_L^\perp	h_L, e_L
T	f_T, f_T^\perp	g_T, g_T^\perp	$h_T, e_T, h_T^\perp, e_T^\perp$

Sivers & BM: Naive T-odd elements:
- contain information about OAM
- sign change between SIDIS and DY

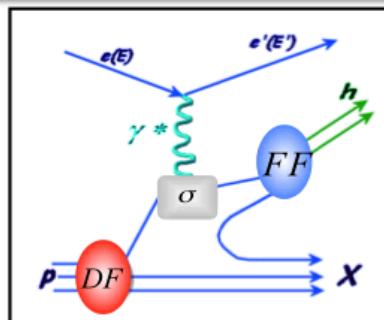
Fragmentation Functions (FF)

		quark		
		U	L	T
had.	U	D_1		H_1^\perp
		Unpol. FF		Collins FF

Chiral-Odd TMD ($\gamma_5 \gamma^1$)

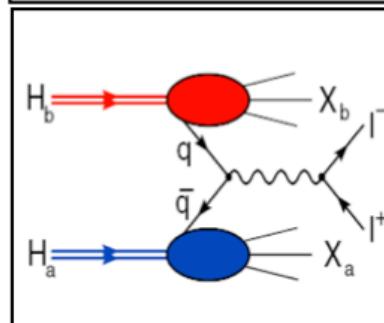
Possible ways to access TMD's

Physics reactions



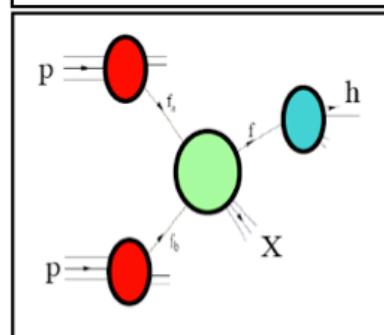
SIDIS

$$\sigma^{ep \rightarrow ehX} = \sum_q (DF) \otimes (\sigma^{eq \rightarrow eq}) \otimes (FF)$$



e⁺e⁻

$$\sigma^{e^+e^- \rightarrow hhX} = \sum_q (\sigma^{qq \rightarrow ee}) \otimes (FF) \otimes (FF)$$



DY

$$\sigma^{pp \rightarrow eeX} = \sum_q (DF) \otimes (DF) \otimes (\sigma^{qq \rightarrow ee})$$



Hadron reactions: challenging for theory (ISI + FSI)

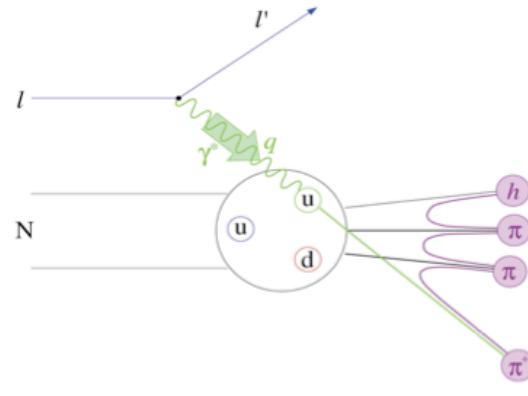
pp

$$\sigma^{pp \rightarrow hX} = \sum_q (DF) \otimes (DF) \otimes (\sigma^{qq \rightarrow qq}) \otimes (FF)$$

WG6: Y. Makdisi



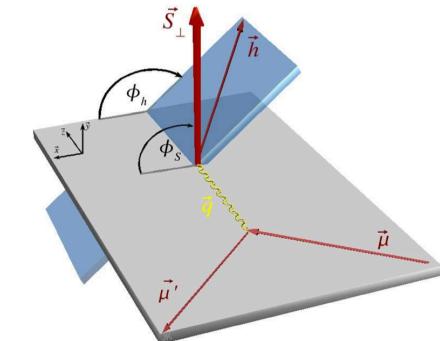
TMD formalism in SIDIS



WG6: P. Mulders, J.W. Qui

Courtesy: Jiang, Wehai'11

$$\frac{d\sigma}{dxdy d\phi_S dz d\phi_h dP_{h\perp}^2} = \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)}.$$

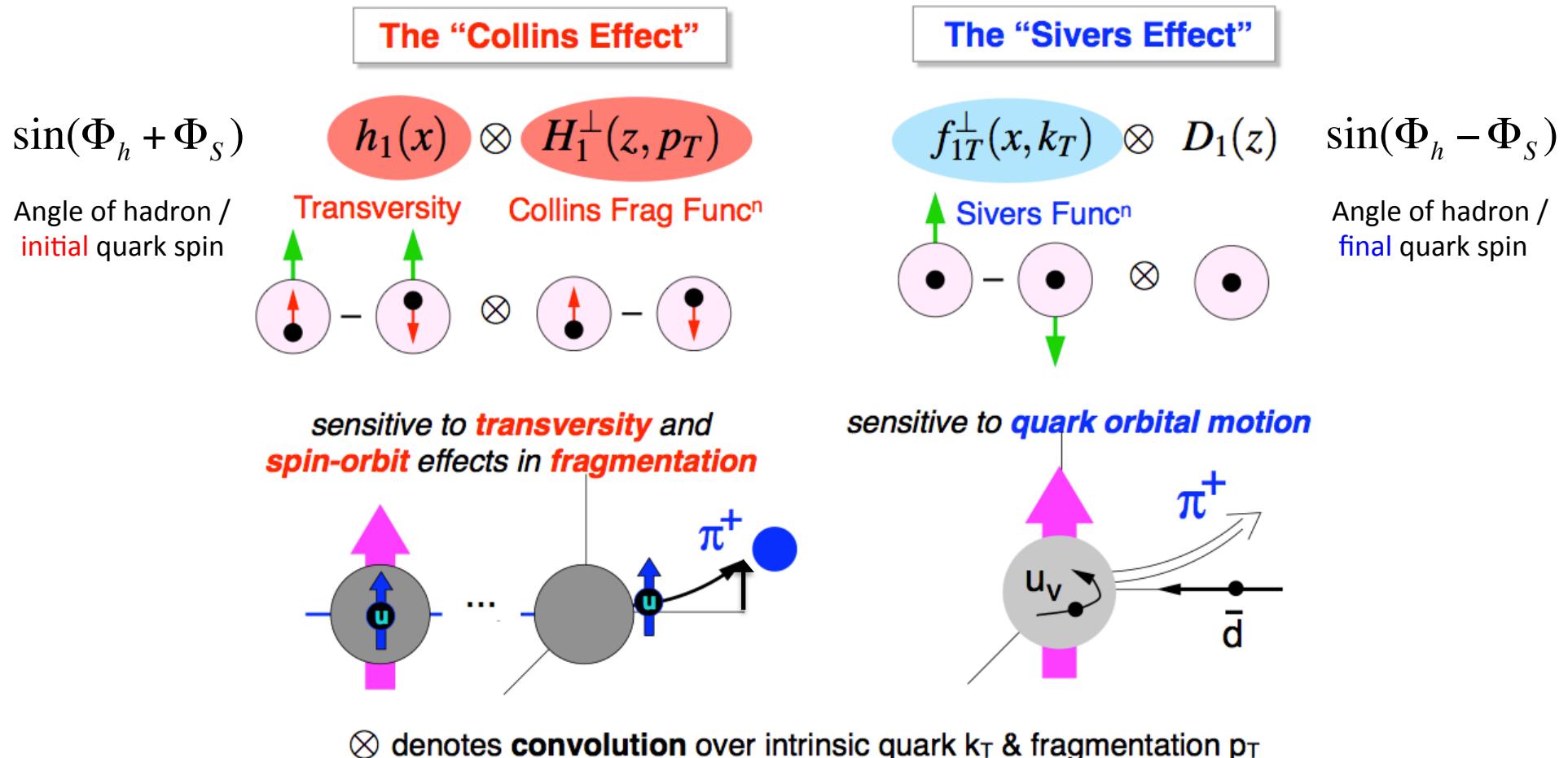


	$f_1 =$	$\{F_{UU,T} + \dots$	
Boer-Mulders	$h_1^\perp =$	$+ \varepsilon \cos(2\phi_h) \cdot F_{UU}^{\cos(2\phi_h)} + \dots$	Unpolarized
Worm-Gear	$h_{1L}^\perp =$	$+ S_L [\varepsilon \sin(2\phi_h) \cdot F_{UL}^{\sin(2\phi_h)} + \dots]$	
Transv/Collins	$h_{1T}^\perp =$	$+ S_T [\varepsilon \sin(\phi_h + \phi_S) \cdot F_{UT}^{\sin(\phi_h + \phi_S)} + \dots]$	Polarized Target
Sivers	$f_{1T}^\perp =$	$+ \sin(\phi_h - \phi_S) \cdot (F_{UL}^{\sin(\phi_h - \phi_S)} + \dots)$	
Pretzelosity	$h_{1T}^\perp =$	$+ \varepsilon \sin(3\phi_h - \phi_S) \cdot F_{UT}^{\sin(3\phi_h - \phi_S)} + \dots]$	
Helicity	$g_{1L} =$	$+ S_L \lambda_e [\sqrt{1 - \varepsilon^2} \cdot F_{LL} + \dots]$	DSA
Worm-Gear	$g_{1T} =$	$+ S_T \lambda_e [\sqrt{1 - \varepsilon^2} \cos(\phi_h - \phi_S) \cdot F_{LT}^{\cos(\phi_h - \phi_S)} + \dots]$	Polarized Beam and Target

S_L, S_T : Target Polarization; λ_e : Beam Polarization

Collins and Sivers TMD's

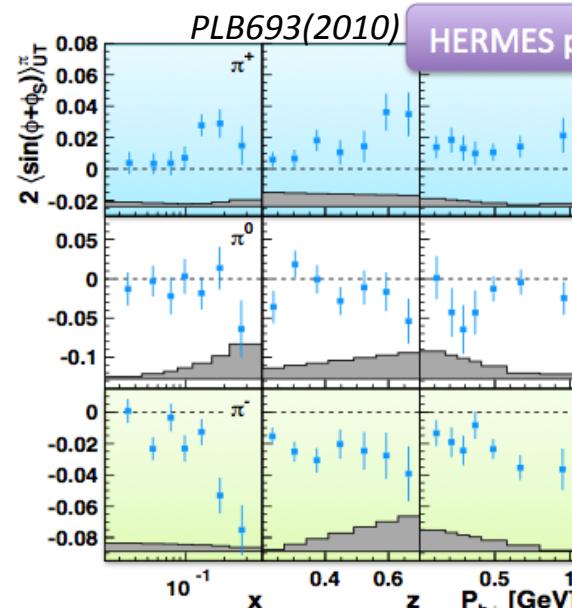
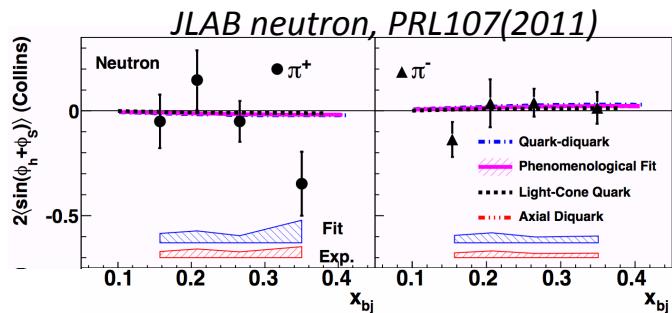
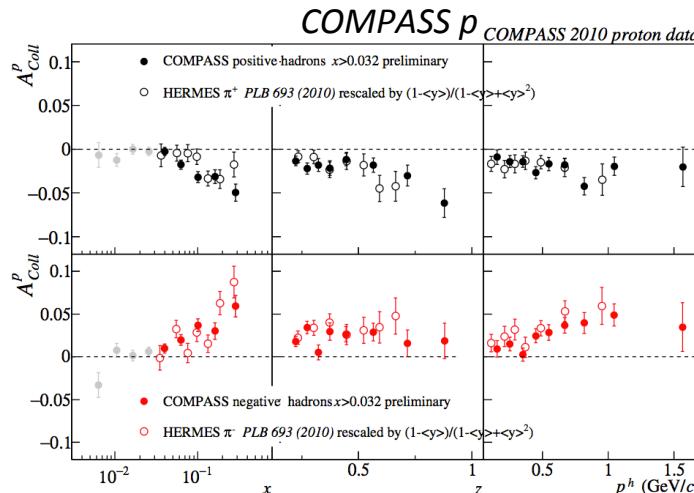
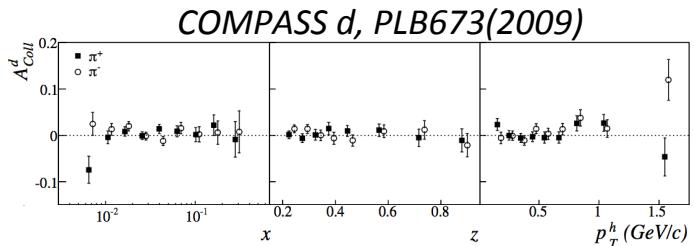
Courtesy: N. Makins, PANIC'11



- correlation between parton transverse polarization in a transversely polarized nucleon and transverse momentum of the produced hadron

- correlation between parton transverse momentum and nucleon transverse polarization
- requires orbital angular momentum

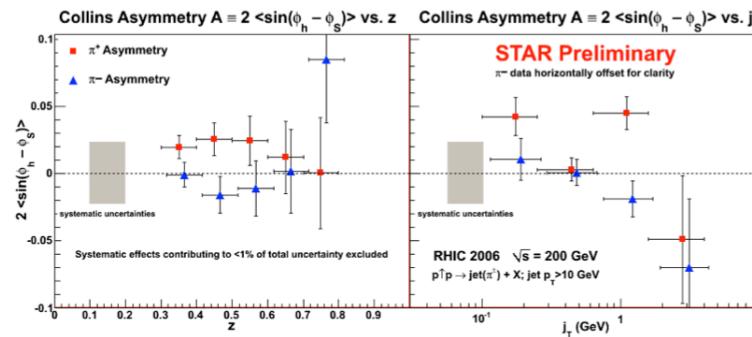
TMD: Collins SIDIS & pp SSA for π



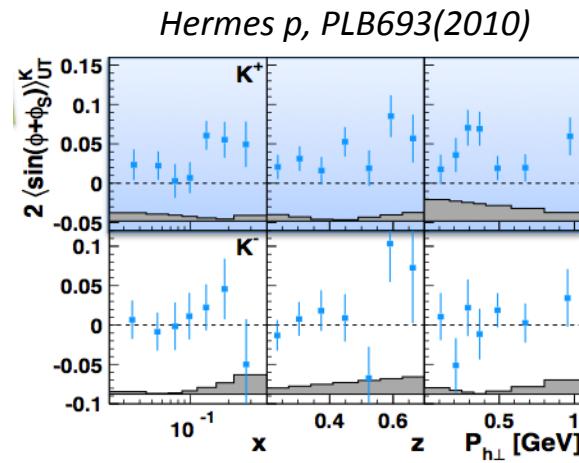
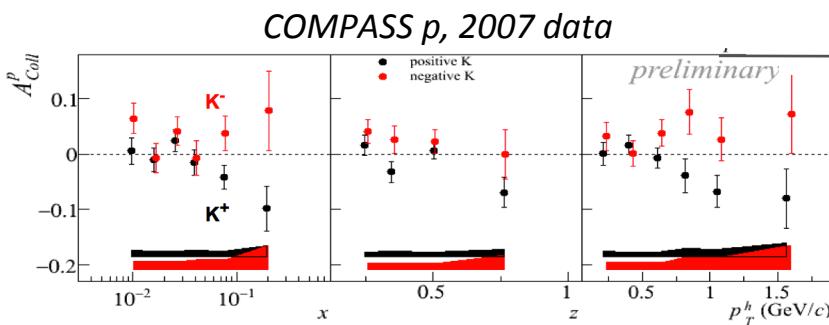
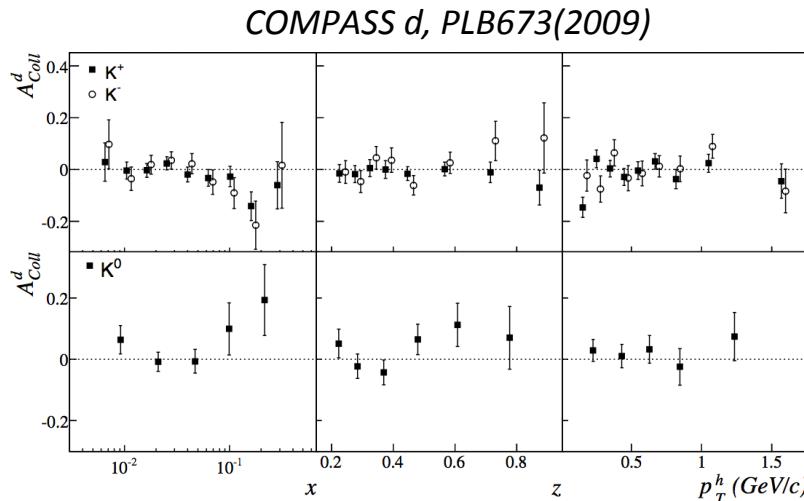
- positive amplitude for π^+
- compatible with zero amplitude for π^0
- large negative amplitude for π^-
- increase in magnitude with x
- transversity mainly receives contribution from valence quarks
- increase with z
- in qualitative agreement with BELLE results
- positive for π^+ and negative for π^-
- role of disfavored Collins FF:
 $H_1^{\perp, \text{disfav}} \approx -H_1^{\perp, \text{fav}}$
 $u \Rightarrow \pi^+; \quad d \Rightarrow \pi^- (\text{fav})$
 $u \Rightarrow \pi^-; \quad d \Rightarrow \pi^+ (\text{disfav})$
 $h_1^u > 0$
 $h_1^d < 0$

WG6: C. Adolf

STAR pp, mid rapidity



TMD: Collins SIDIS SSA for K



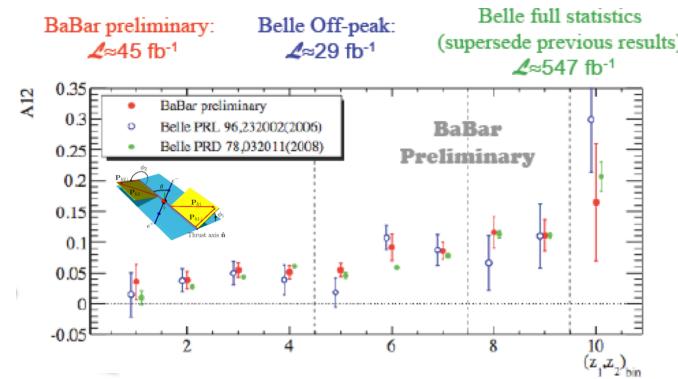
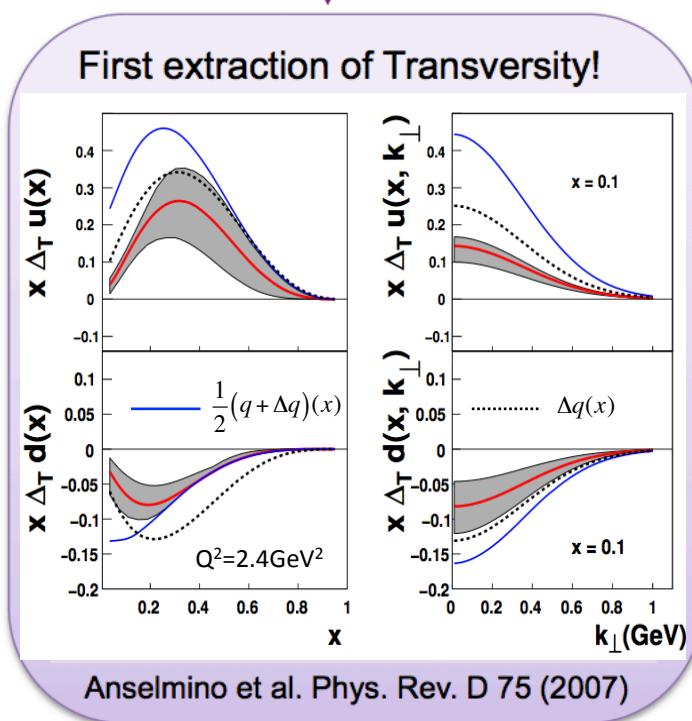
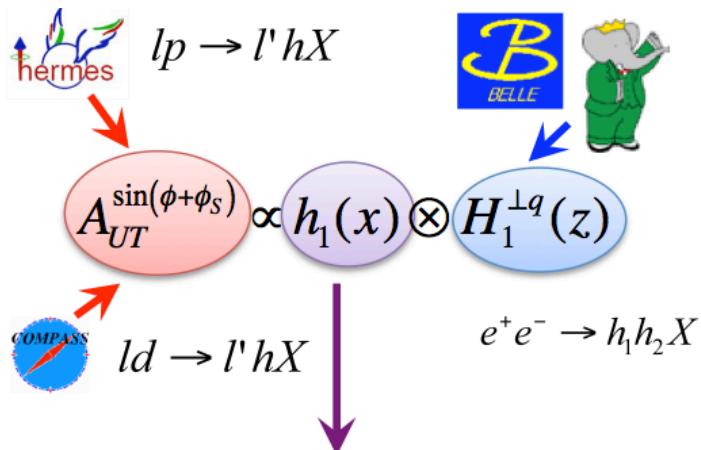
K^+

- ➡ K^+ amplitudes are similar to π^+ as expected from the u-quark dominance
- ➡ K^+ are larger than π^+

K^-

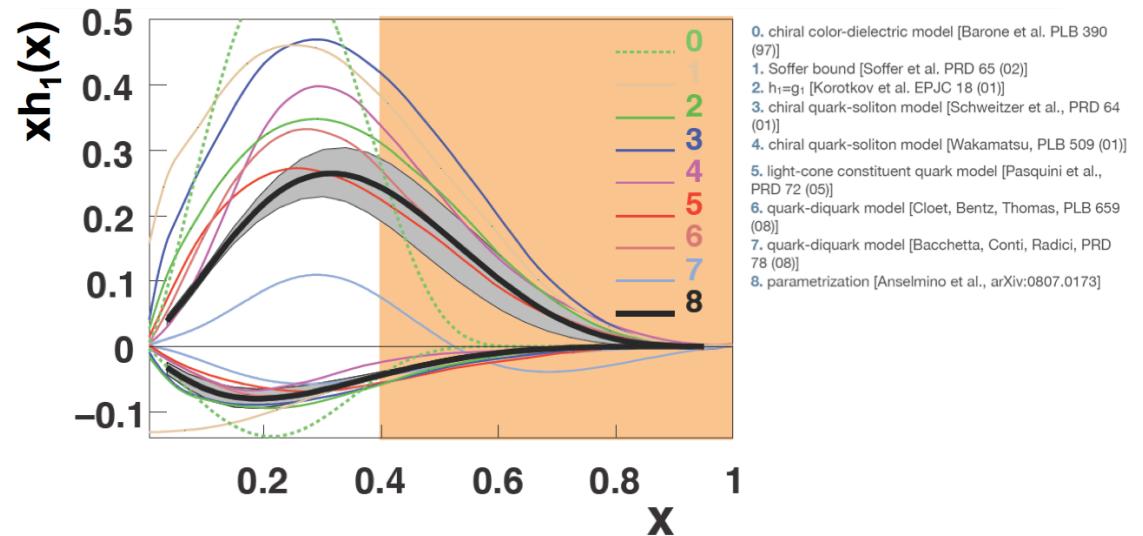
- ➡ consistent with zero amplitudes
- ➡ $K^- (\bar{u}s)$ is all sea object

TMD: Global fits to Collins 1h

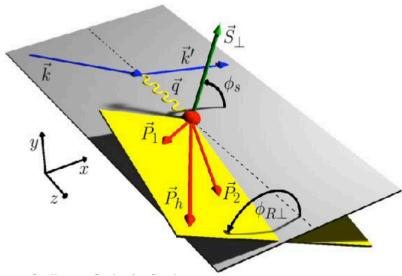


WG6: M. Leitgab

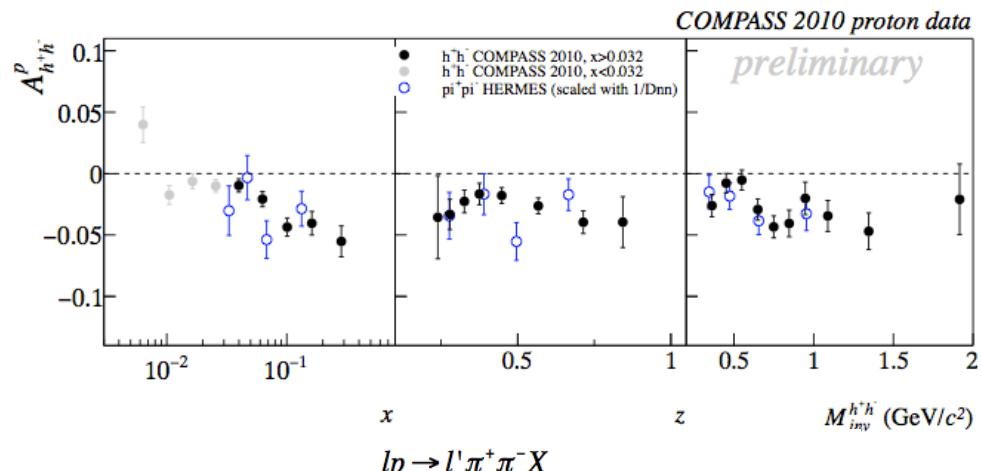
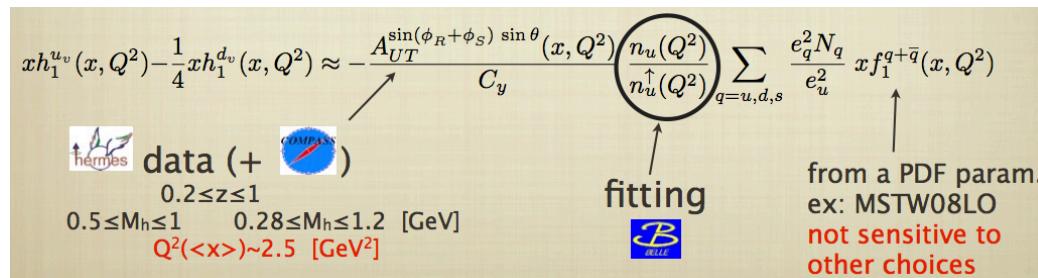
- Fit to COMPASS d + HERMES p (no COMPASS p in fit yet)
- $h_1^{u,d}$ slightly smaller than $g_1^{u,d}$ ($\Delta u, \Delta d$)



TMD: Collins asymmetry 2 h SIDIS

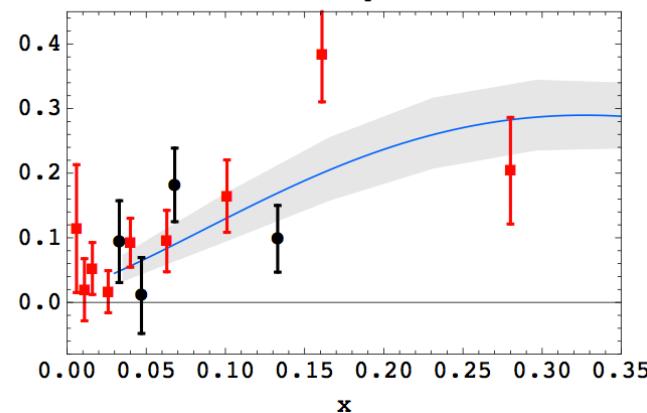


$$A_{UT}^{\sin(\phi_R + \phi_S) \sin \theta} \propto \frac{\sum_q e_q^2 h_1(x, Q^2) H_1^q(z, M_h^2, Q^2)}{\sum_q e_q^2 f_1(x, Q^2) D_1^q(z, M_h^2, Q^2)}$$



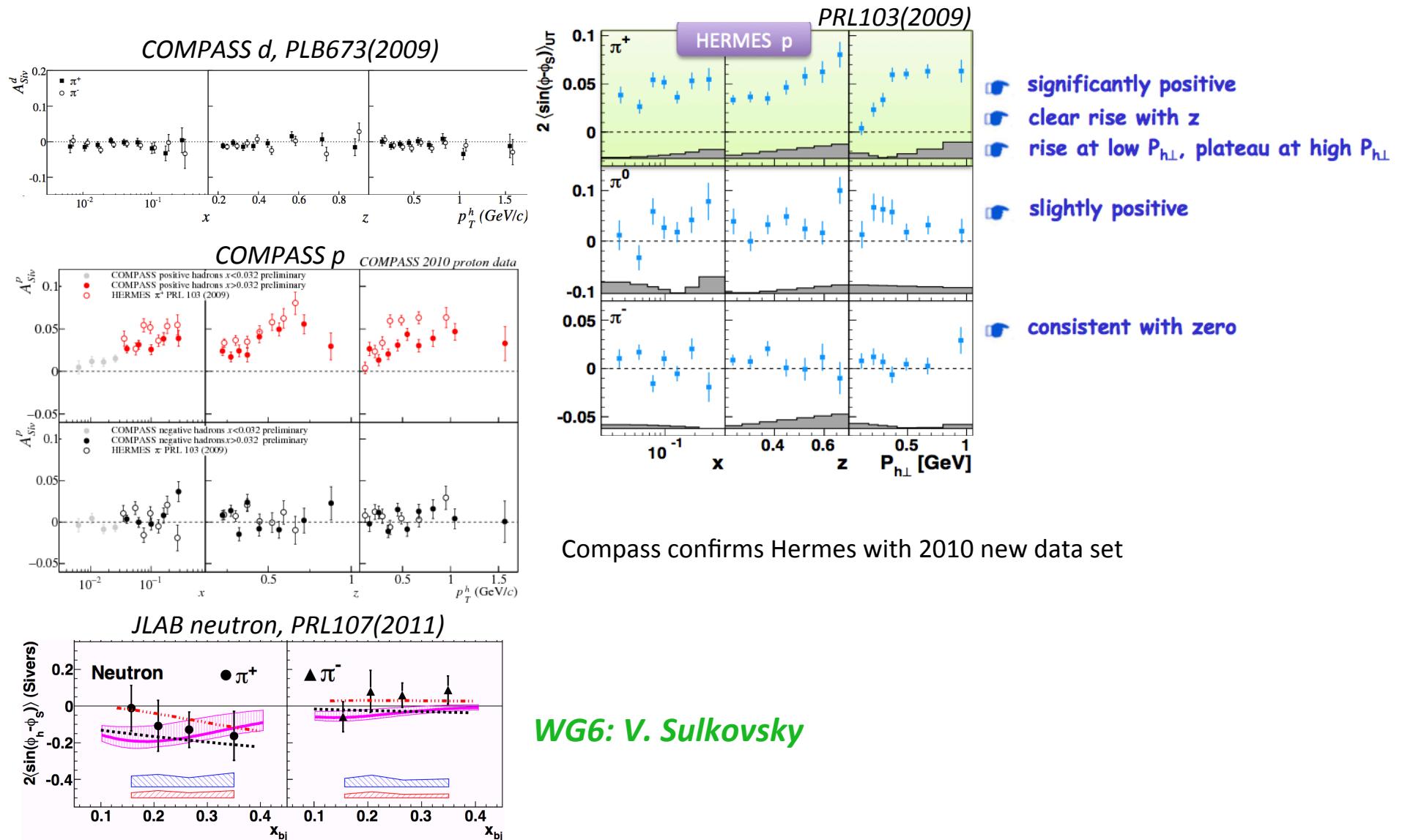
$$A_{UT}^{\sin(\phi_R + \phi_S) \sin \theta} \propto \sin \theta h_1(x) \otimes H_1^{q\bar{q}}(z)$$

$$x h_1^{uv}(x) - \frac{1}{4} h_1^{dv}(x)$$

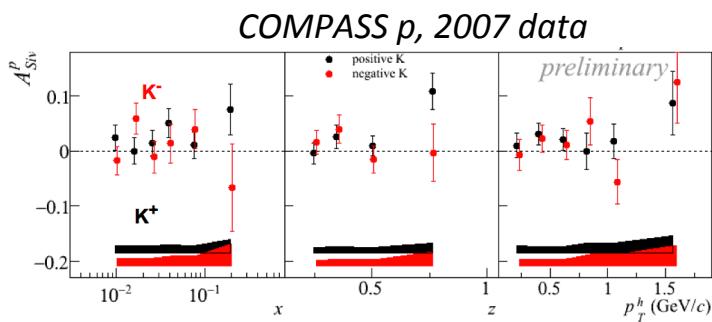
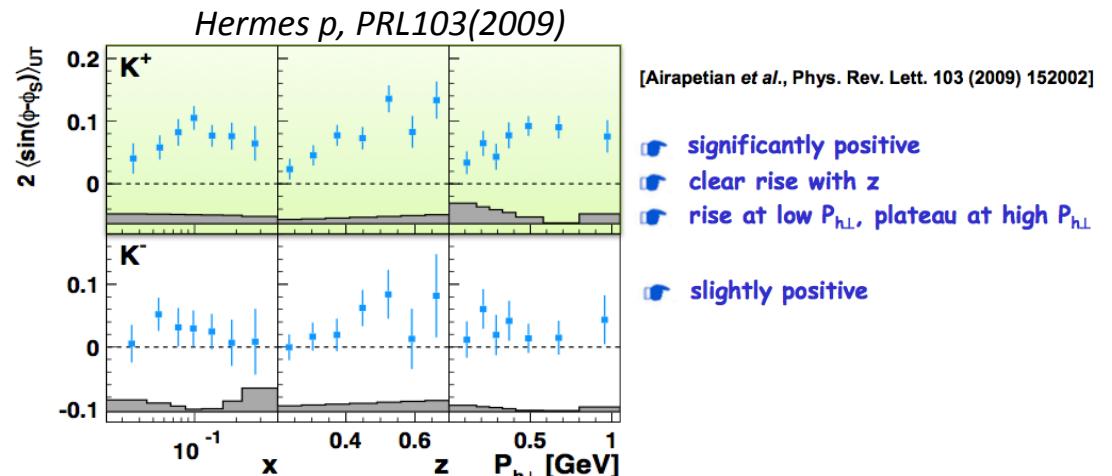
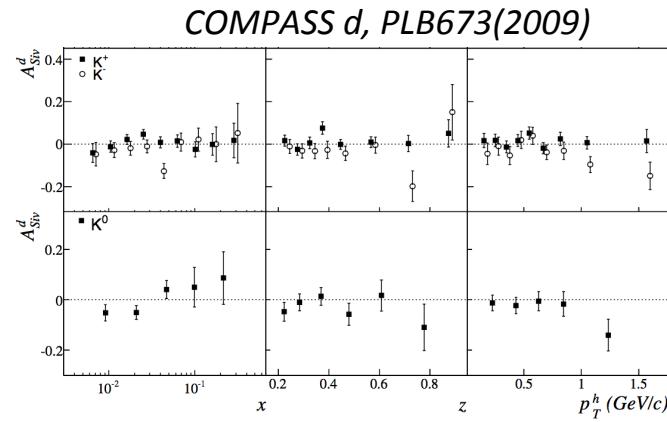


- Independent access to h_1^q
- Combining p+d $\rightarrow h_1^u, h_1^d$
(not shown here)

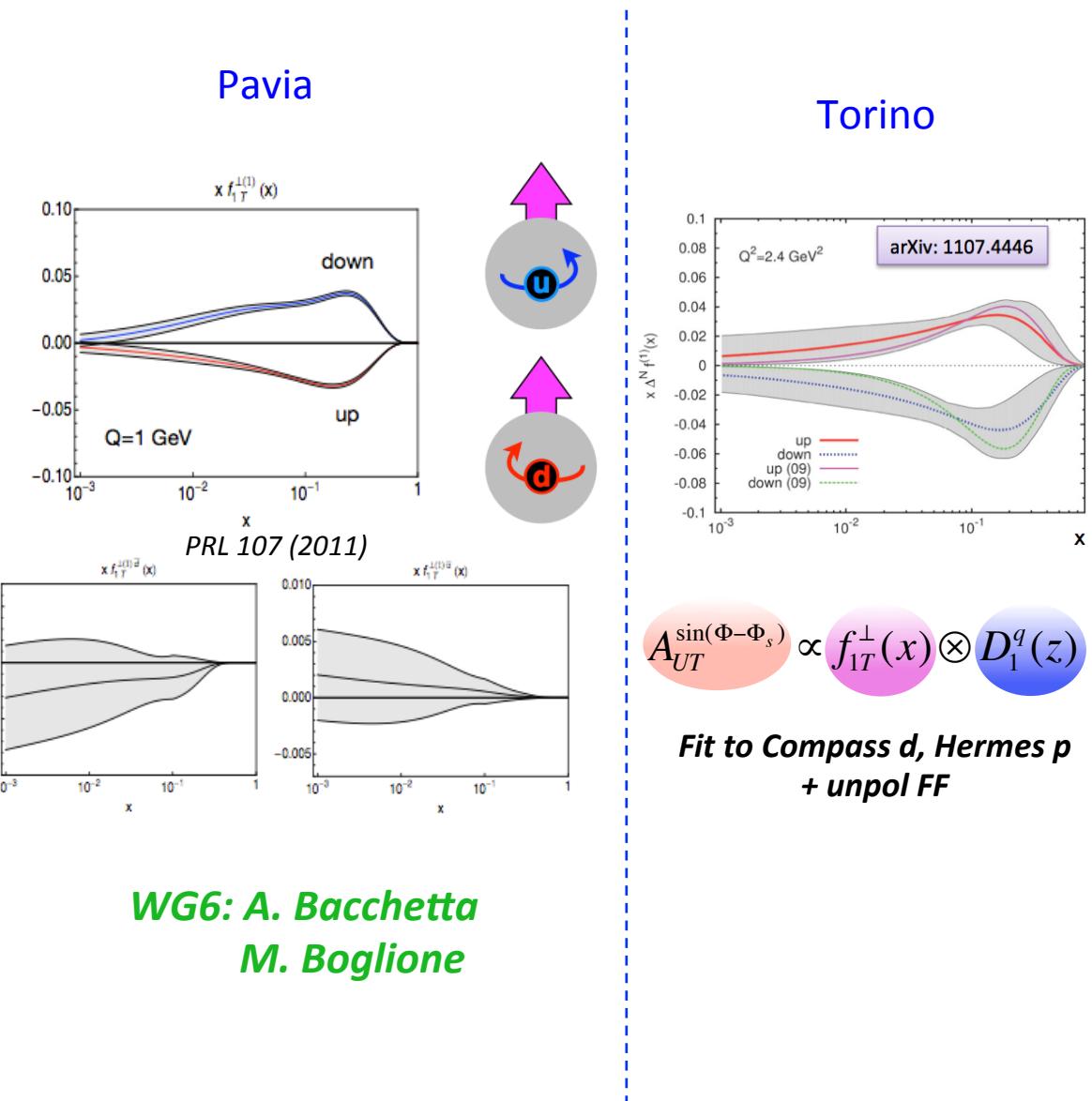
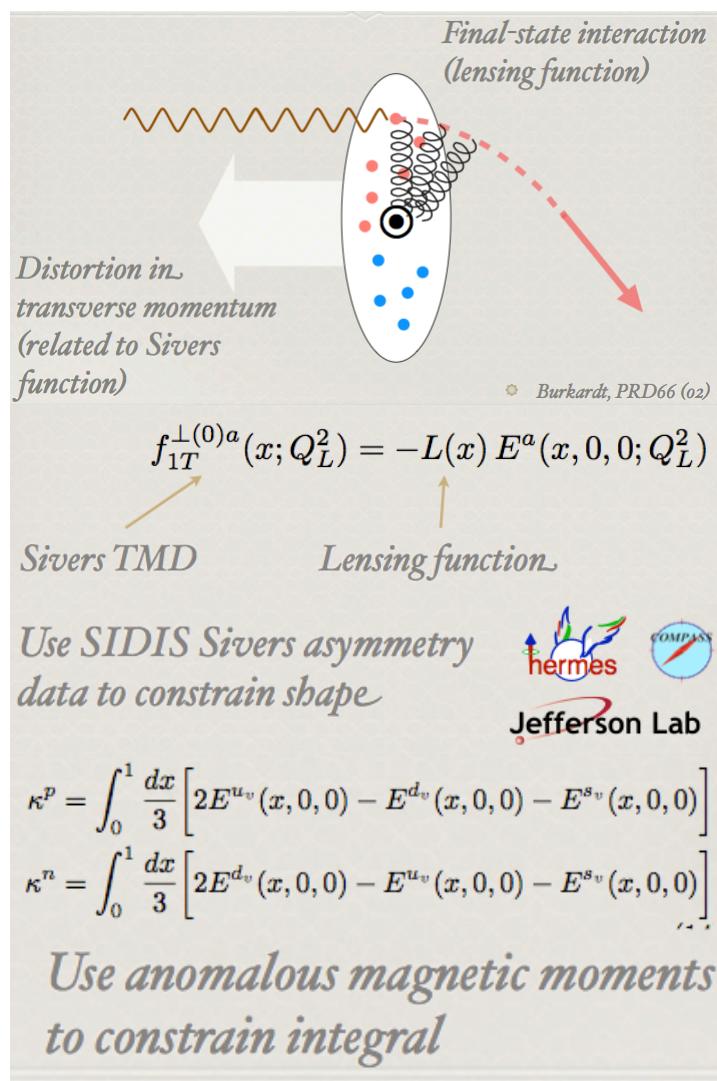
TMD: Sivers SIDIS SSA for π



TMD: Sivers SIDIS SSA for K



Sivers asymmetry SIDIS and OAM



Courtesy Bacchetta, PINAN'11

Sivers asymmetry SIDIS and quark OAM

Courtesy Bacchetta, PINAN'11

$$J^u = 0.233 \pm 0.002^{+0.008}_{-0.012},$$

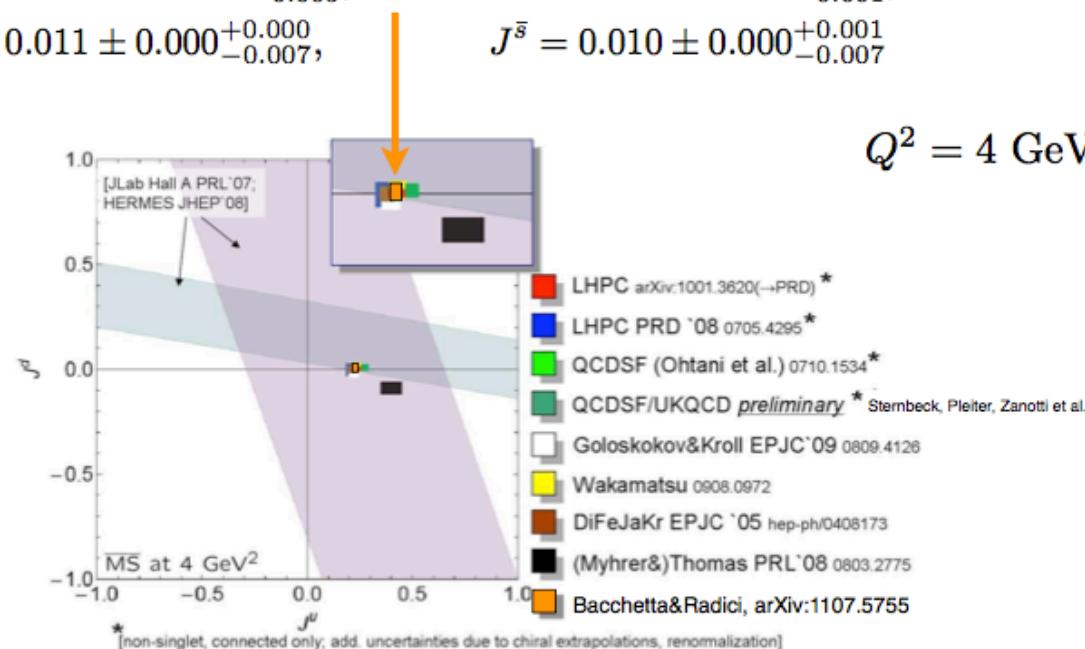
$$J^d = -0.003 \pm 0.002^{+0.020}_{-0.005},$$

$$J^s = 0.011 \pm 0.000^{+0.000}_{-0.007},$$

$$J^{\bar{u}} = -0.019 \pm 0.003^{+0.002}_{-0.001},$$

$$J^{\bar{d}} = 0.026 \pm 0.003^{+0.002}_{-0.001},$$

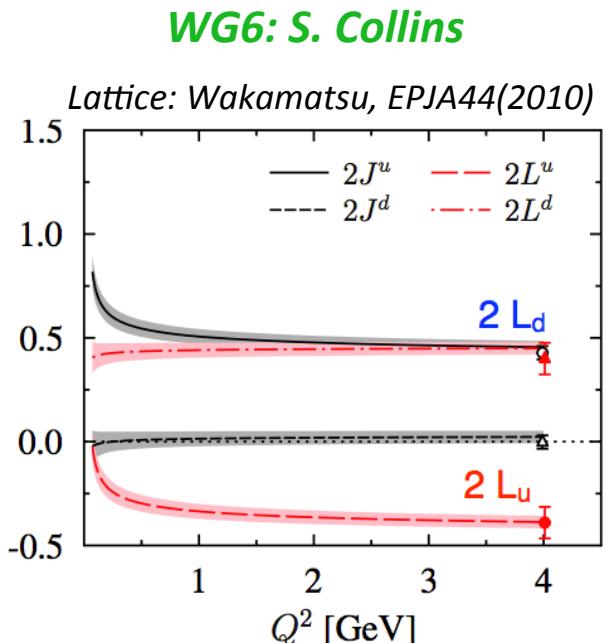
$$J^{\bar{s}} = 0.010 \pm 0.000^{+0.001}_{-0.007}$$



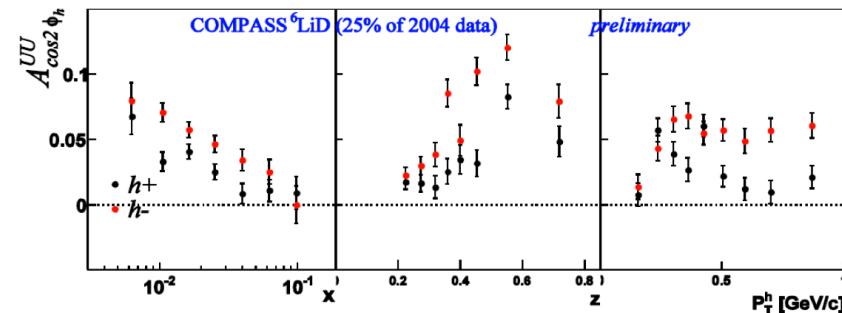
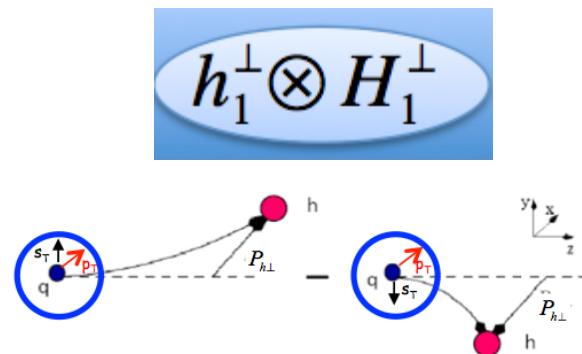
INT writeup, Boer et al., arXiv:1108.1713 (fig. 4.3)

$$J_q = \frac{1}{2} \Delta \Sigma + L_q = \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H(x, \zeta, t) + E(x, \zeta, t)]$$

WG6: F. Yuan

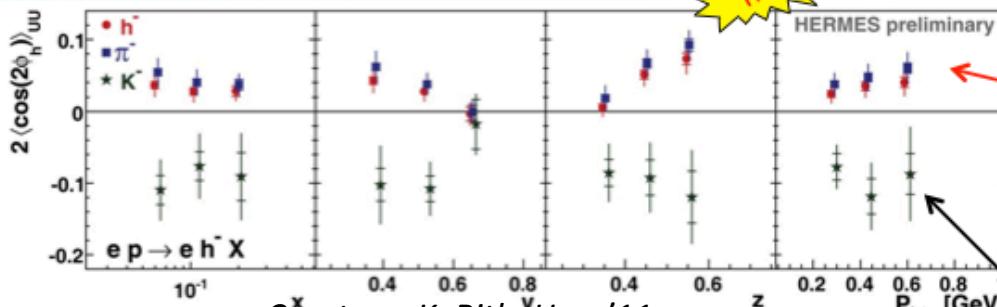


TMD: Boer-Mulders in SIDIS

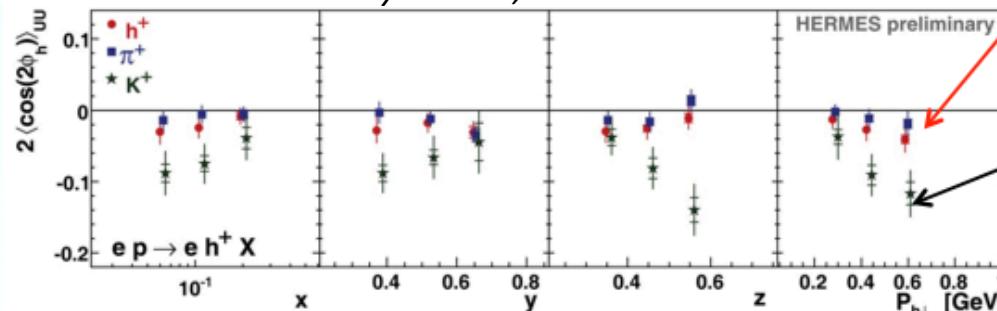


transversely polarised quarks
with p_T in unpolarised nucleon

WG6: F. Giordano



Courtesy: K. Rith, Hera'11



h_1^\perp is chiral-odd and
naive T-odd (like f_{1T^\perp})
requires FSI/ISI

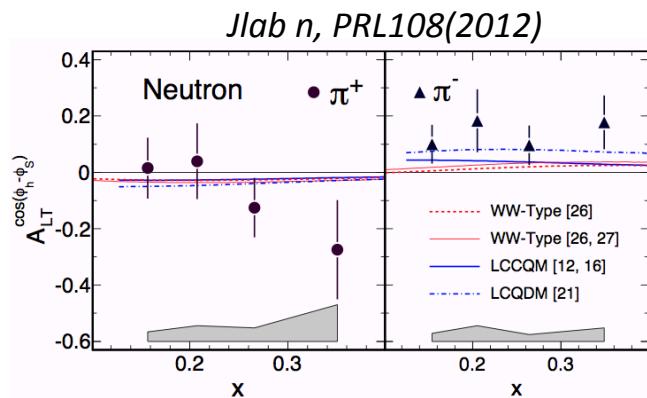
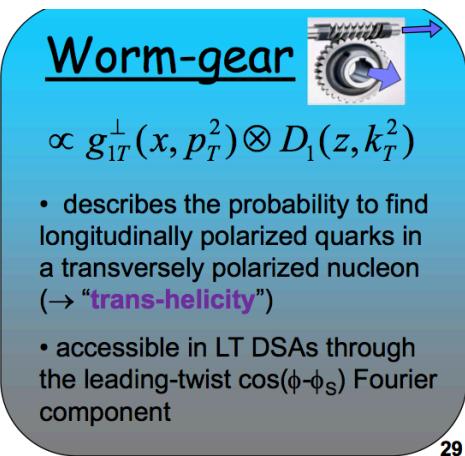
Opposite sign for
 π^+ and π^- , larger
magnitude for π^-

$h_1^{\perp,u}$ and $h_1^{\perp,d}$
have same sign

Large signal with
same sign for K^\pm

sea fragmentation
important

TMD: Worm-Gear in SIDIS



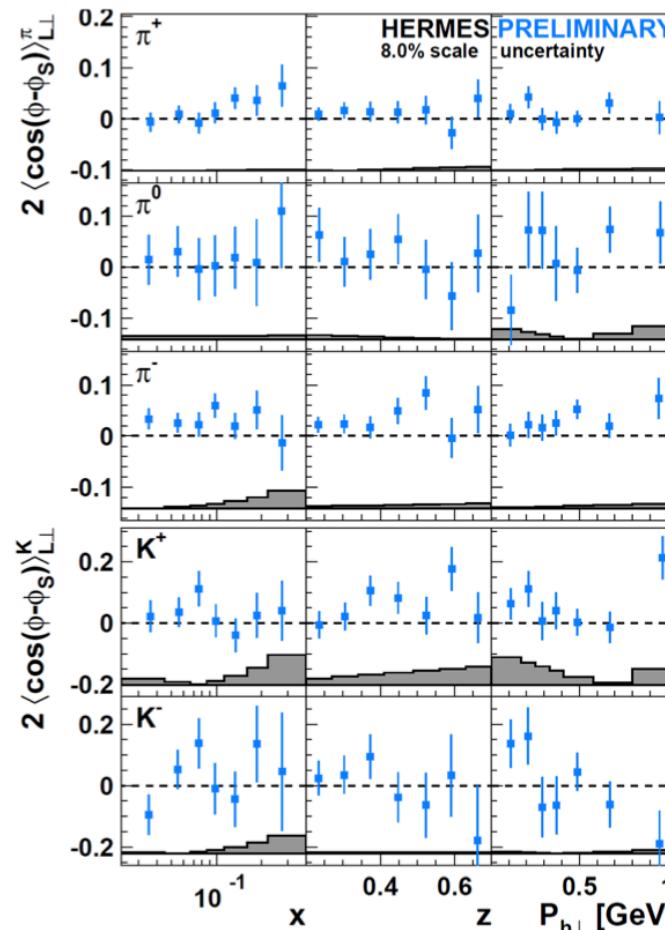
The only TMD that is both chiral-even and naïve-T-even

requires interference between wave funct. components that differ by 1 unit of OAM

➤ Many models support simple relations among g_{1T}^{\perp} and other TMDs:

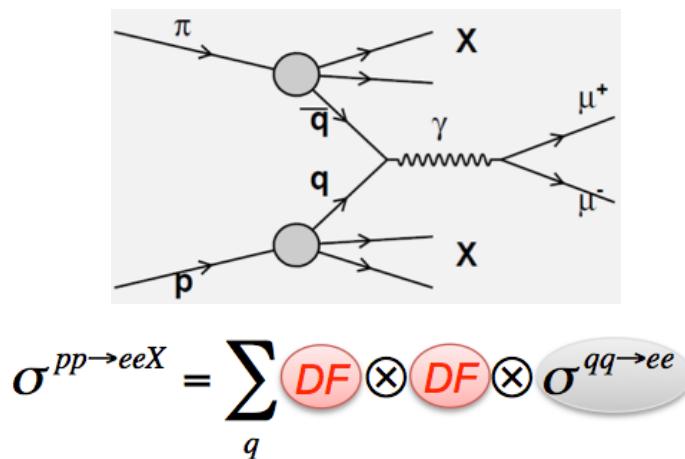
- $g_{1T}^q = -h_{1L}^{\perp q}$ (also supported by Lattice QCD and first data)

- $g_{1T}^{q(1)}(x) \stackrel{WW-type}{\approx} x \int_x^1 \frac{dy}{y} g_1^q(y)$ (Wandzura-Wilczek appr.)



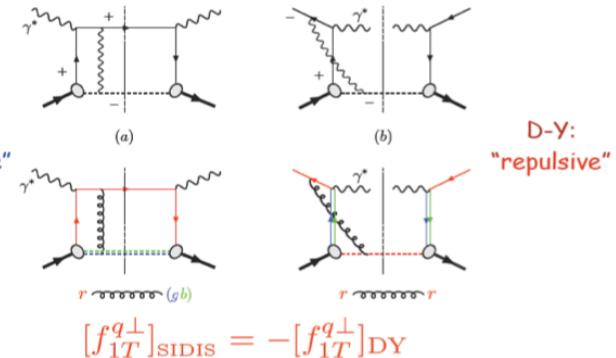
consistent positive signal for π^+ on p,n

TMD: Sivers (&BM) asymmetry DY (future)

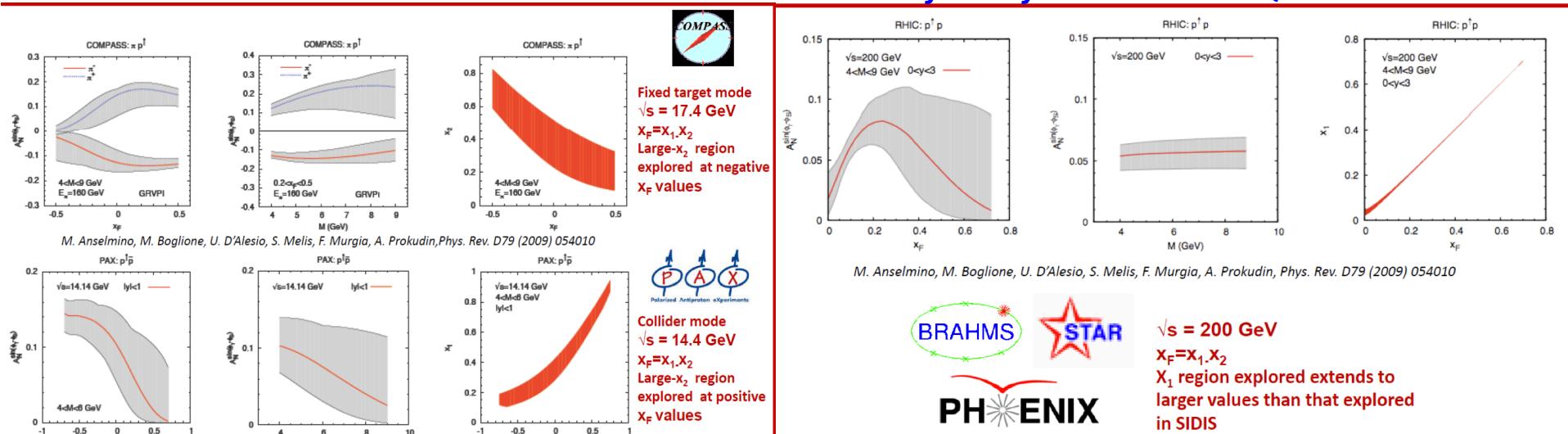


Crucial role of gauge-links in TMDs
Brodsky, Hwang, Schmidt;
Collins; Belitsky, Ji, Yuan;
Boer, Mulders, Pijlman

process-dependence of Sivers functions



Test of TMD factorization at QCD level



M. Anselmino et al., Phys. Rev. D79:054010, 2009

WG7: G. Mallot, K. Boyle, M. Lamont

Summary & Outlook

❑ Helicity distributions for quark and gluons:

- Valence $\Delta u_v, \Delta d_v$ quite well known, but $\Sigma(\Delta q)$ only 33% of nucleon spin
- Separated $\Delta q, \Delta \bar{q}$ obtained, but still issues with sea polarization
- $\Delta G/G(x)$ from direct measurements, but still limited in x range

❑ Transverse momentum dependent DF for quarks:

- Transversity (Collins): clean opposite signals for π^+, π^- , $\Delta_T q \leq \Delta q$
- T-odd Sivers: confirmed signal $\neq 0$ for π^+ , $= 0$ for π^- on p, relation to OAM
- Many other SSA and SSA measured, more to come

❑ Expected improvements for future:

- Gluon sector: RHIC, EIC
- Quark sector (helicity, TMD-DY, GPD): COMPASS, JLAB12, RHIC, EIC,

WG7: *G. Mallot, H. Wollny (COMPASS)*

K. Griffioen, K. Allada (JLAB12)

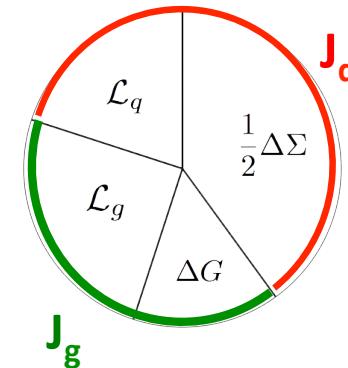
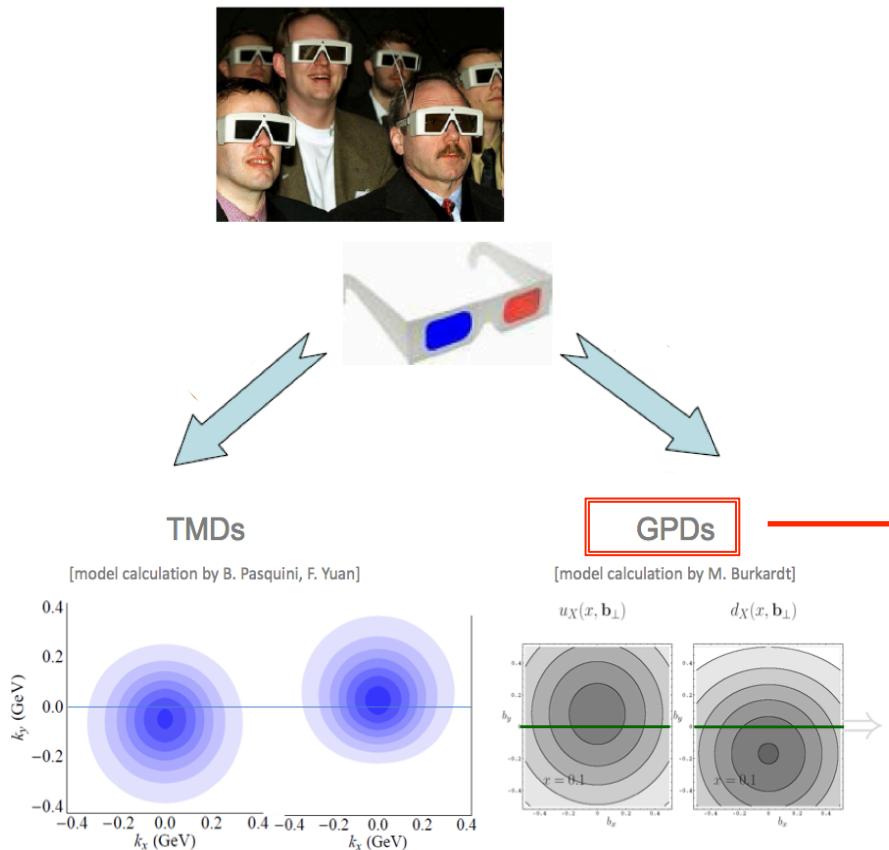
K. Boyle, M. Lamont (RHIC)

A. Bazilewski (EIC)

WG6+7: *F. Yuan, M. Stratmann (EIC)*

Stay tuned for further « surprises »

3D imaging of nucleon and OAM



Theory:

WG6: F. Yuan, S. Liuti, J. Wagner
WG2+6: M. Diehl
WG6+7: D. Mueller

Experimental:

WG6+7: S. Fazio
WG2+6: H. Moutarde, M. Murray
D. Sokhan