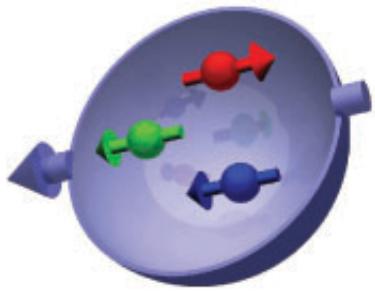




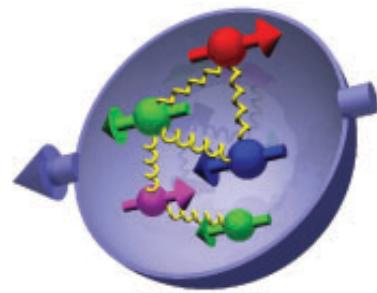
# $\Delta G/G$ from high $p_T$ hadron pairs at $Q^2 > 1$ ( $\text{GeV}/c^2$ )<sup>2</sup>

- 1 Motivation
- 2 COMPASS experiment
- 3 Direct measurements of  $\Delta G/G$
- 4 Details of analysis method
- 5 Results
- 6 Conclusion

## Contributions to the nucleon spin

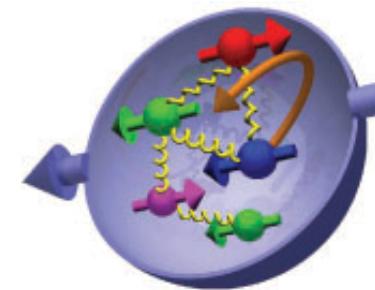


CQM (EJ):  
 $\Delta\Sigma \sim 0.6$



E&SMC,SLAC,HERMES  
CLAS,COMPASS  $\mathbf{g}_1$ :

$$\Delta\Sigma \sim 0.3$$



→ **Spin puzzle**

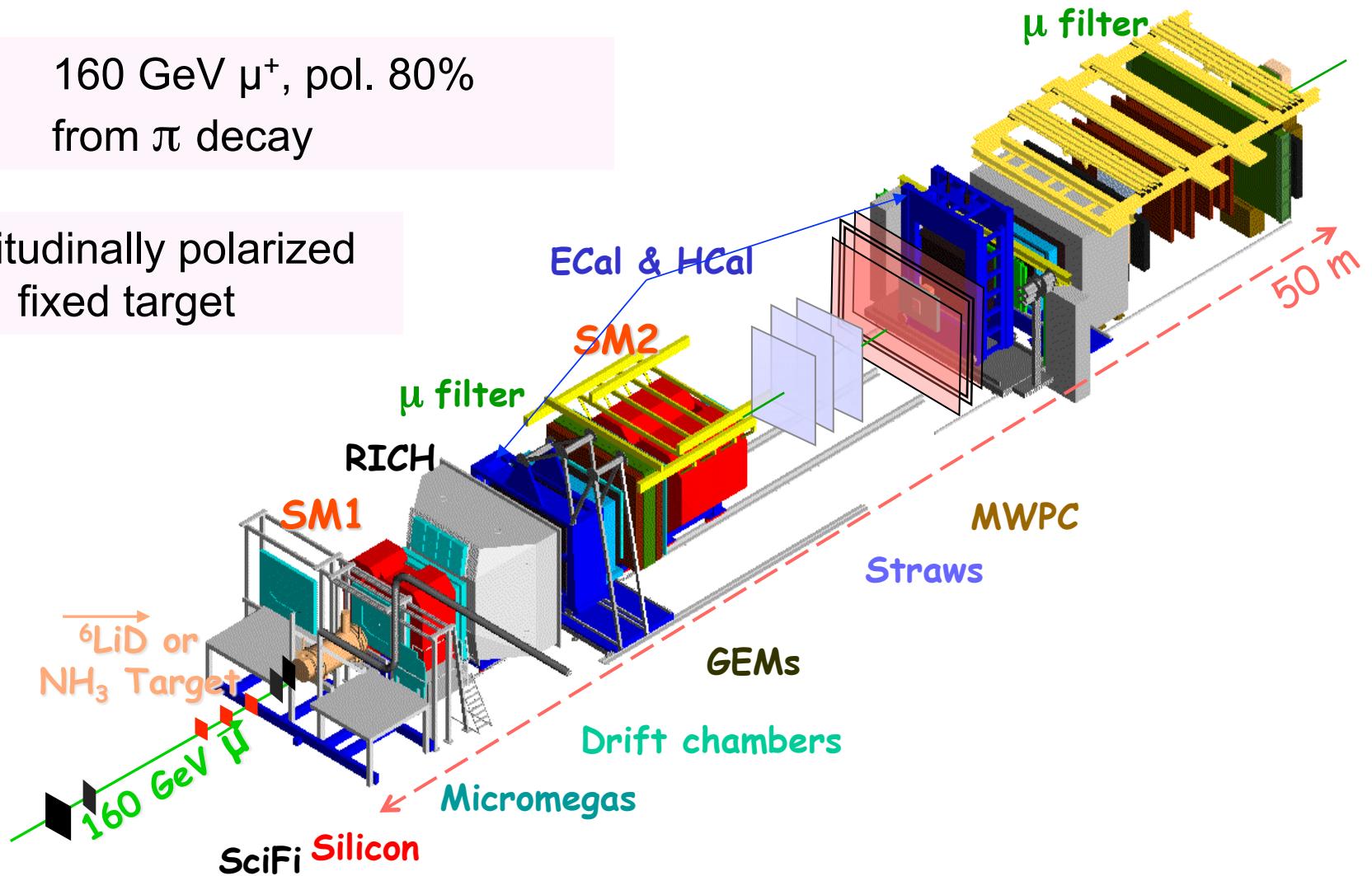
$$S_N = \frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta G + L_q + L_g$$

# Compass setup 2002-2007



Beam: 160 GeV  $\mu^+$ , pol. 80%  
from  $\pi$  decay

Longitudinally polarized  
fixed target



# Asymmetry measurement



$$A^{\mu N} = \frac{1}{f P_\mu P_T} \frac{N_u^{\leftarrow} - N_d^{\leftarrow}}{N_u^{\leftarrow} + N_d^{\leftarrow}}$$

$$A^{\mu N} \approx D A_1^{\gamma N}$$

**Target:**  ${}^6\text{LiD}$  (02-06) -  $\text{NH}_3$  (2007)

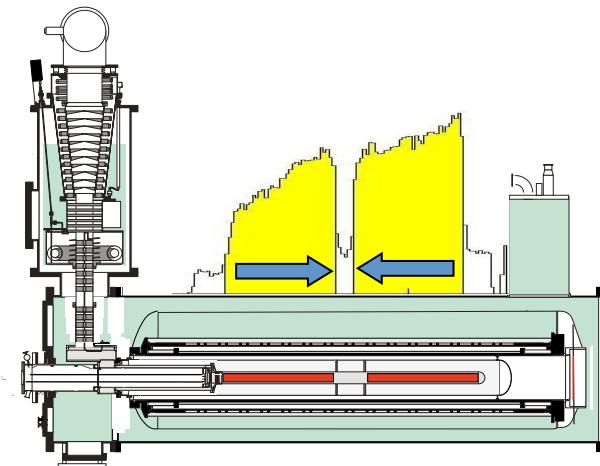
$P_T \sim 50\% / 90\%$

$f$  (dilution factor)  $\sim 40\% / 16\%$

$T \sim 50 \text{ mK}$  ( ${}^3\text{He} / {}^4\text{He}$ )

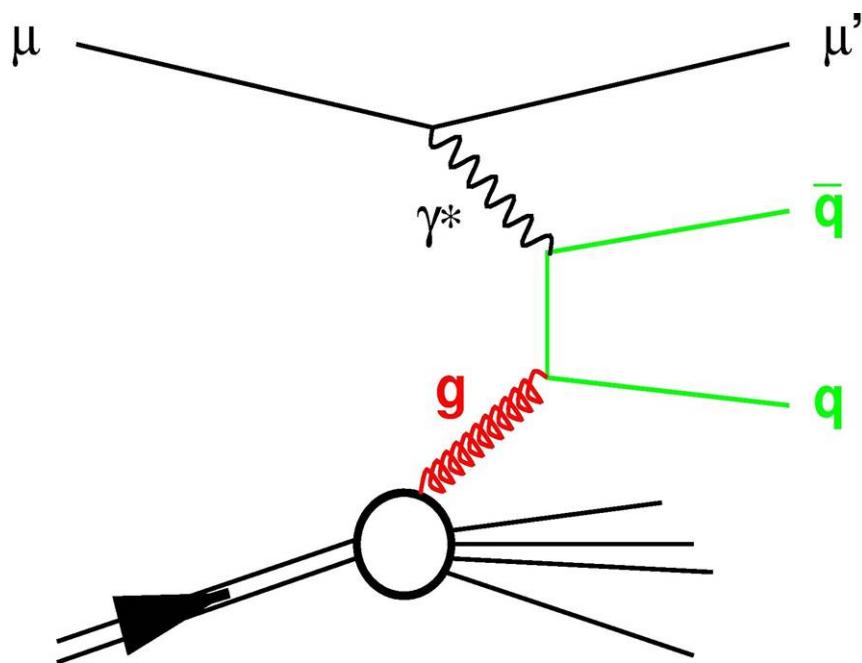
**Weighting each event with  $\omega = (f P_\mu D)$ :**

$$A_1^{\gamma N} = \frac{1}{P_T} \left( \frac{\sum_u \omega - \sum_d \omega}{\sum_u \omega + \sum_d \omega} \right)$$





## Photon-Gluon Fusion process



$$A_{LL}^{\mu N} = R^{PGF} a_{LL}^{PGF} \frac{\Delta G}{G}$$

**There are two methods to tag this process:**

- **Open Charm production**

- $\gamma^*g \rightarrow c\bar{c} \Rightarrow$  reconstruct  $D^0$  mesons
- Hard scale:  $M_c^2$
- No intrinsic charm in COMPASS kinematics
- No physical background
- Weakly Monte Carlo dependent
- Low statistics

- **High- $p_T$  hadron pairs**

- $\gamma^*g \rightarrow q\bar{q} \Rightarrow$  reconstruct 2 jets or  $h^+h^-$
- Hard scale:  $Q^2$  or  $\Sigma p_T^2$  [ $Q^2 > 1$  or  $Q^2 < 1$   $(\text{GeV}/c)^2$ ]
- High statistics
- Physical background
- Strongly Monte Carlo dependent

# High $p_T$ pairs asymmetries $Q^2 > 1$ (GeV/C) $^2$



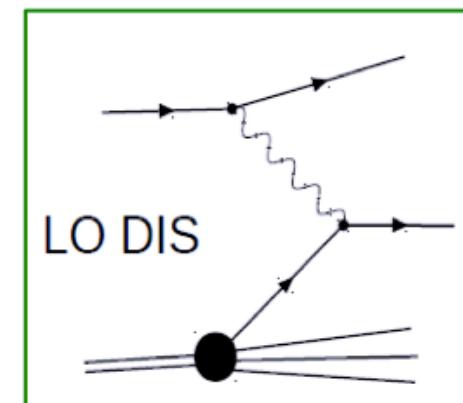
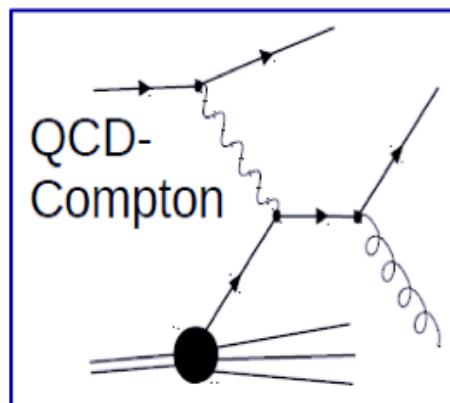
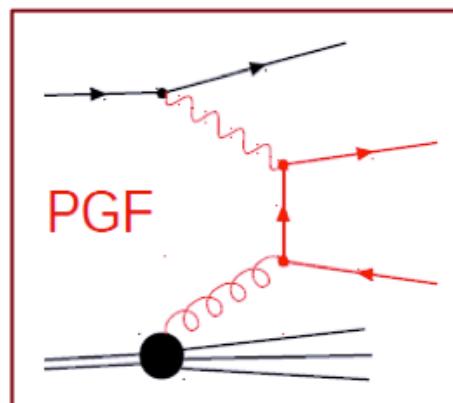
- Two samples are considered:

Inclusive asymmetry

$$A_1^d(x) = \frac{\Delta G}{G}(x_g) \left( a_{LL}^{PGF, inc} \frac{\sigma^{PGF, inc}}{\sigma^{Tot, inc}} \right) + A_1^{LO}(x_C) \left( a_{LL}^{C, inc} \frac{\sigma^{C, inc}}{\sigma^{Tot, inc}} \right) + A_1^{LO}(x_{Bj}) \left( D \frac{\sigma^{LO, inc}}{\sigma^{Tot, inc}} \right)$$

$$A_{LL}^{2h}(x) = \left( \frac{A_{exp}}{f P_\mu P_T} \right) = \frac{\Delta G}{G}(x_g) \left( a_{LL}^{PGF} \frac{\sigma^{PGF}}{\sigma^{Tot}} \right) + A_1^{LO}(x_C) \left( a_{LL}^C \frac{\sigma^C}{\sigma^{Tot}} \right) + A_1^{LO}(x_{Bj}) \left( D \frac{\sigma^{LO}}{\sigma^{Tot}} \right)$$

high- $p_T$  hadron pairs ( $p_{T1} / p_{T2} > 0.7 / 0.4$  GeV/c)  $\Rightarrow$  enhancement of the PGF contribution



## Extraction of $\Delta G/G$ from high $p_T$ pairs



- The gluon polarisation is determined from two asymmetry samples: the **two high- $p_T$  hadrons** and the **inclusive** data samples. **The final formula is:**

$$\frac{\Delta g}{g}(x_g) = \frac{1}{\beta} [A_{LL}^{2h}(x) + A_{corr}]$$

$$A_{corr} = - \left( A_1(x_{Bj}) D \frac{R_{LO}}{R_{inc}} - A_1(x_C) \beta_1 + A_1(x_C') \beta_2 \right)$$

$$\beta = a_{LL}^{PGF} R_{PGF} - a_{LL}^{PGF, inc} R_{PGF}^{incl} \frac{R_{LO}}{R_{inc}} - a_{LL}^{PGF, incl} \frac{R_C R_{PGF}^{inc}}{R_{LO}^{inc}} \frac{a_{LL}^C}{D}$$

- $\beta_1$  and  $\beta_2$  are factors depending on  $a_{LL}^i$  and  $R_i$
- Each event is weighted with  $\omega = f D P_\mu \beta$  → **statistical improvement!**
- The following parameters are obtained from Monte Carlo, and then they are parameterised event-by-event by a Neural Network (to allow for their use in data):

$$R_{PGF}, R_C, R_{LO}, R_{PGF}^{inc}, R_C^{inc}, R_{LO}^{inc}, a_{LL}^{PGF}, a_{LL}^C, a_{LL}^{LO}, a_{LL}^{PGF, inc}, a_{LL}^{C, inc} \text{ and } a_{LL}^{LO, inc}$$

## Event selection



- Interaction vertex which contains an incoming and an outgoing **muon**
- For DIS variables:  $Q^2 > 1 \text{ (GeV/c)}^2$  and  $0.1 < y < 0.9$

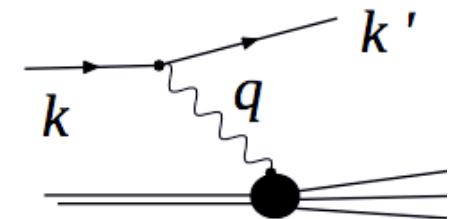
⇒ This selection constitutes the **inclusive sample**

- Events with at least **2** outgoing **hadrons** are selected
- The **hadrons** of the **high  $p_T$**  pair are required to have:

$$p_{T1} > 0.7 \text{ GeV/c}, \quad p_{T2} > 0.4 \text{ GeV/c}$$

$$z_1 + z_2 < 0.95$$

⇒ All this selection produces the **high  $p_T$**  sample



$$Q^2 = -q^2$$

$$q = k - k'$$

$$\nu = E - E'$$

$$y = \frac{\nu}{E}$$

$$x = \frac{Q^2}{2M\nu}$$

Years	2002	2003	2004	2006	all years
Statistics	450 K	1.3 M	2.8 M	2.7 M	<b>7.3 M</b>

Events



- The purpose of the **MC tuning** is to **correct** the shapes of the **hadron variables** (momenta) and **fragmentation** (multiplicity).
- In **LEPTO** this can be **achieved** by changing **JETSET** parameters:

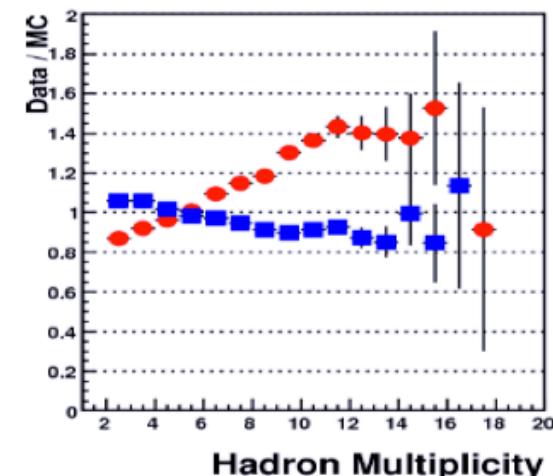
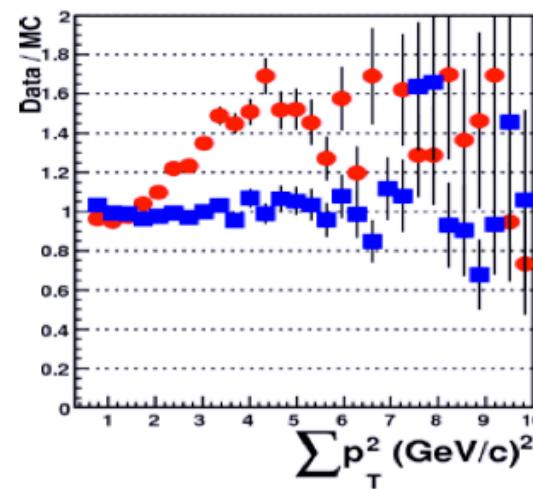
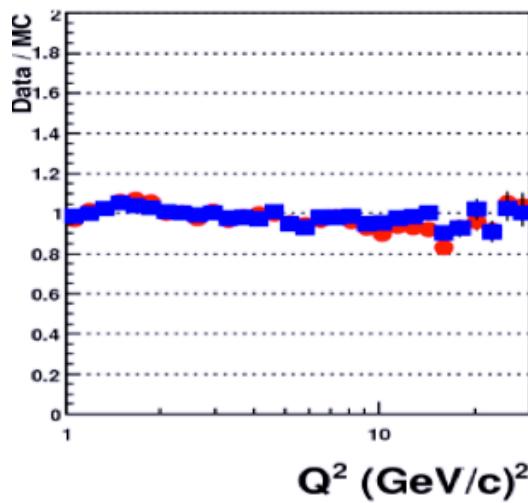
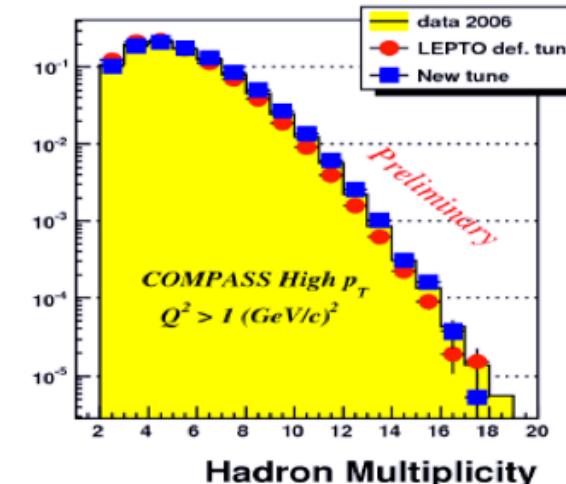
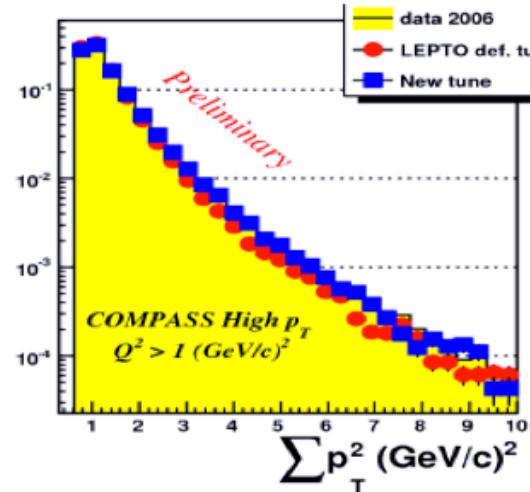
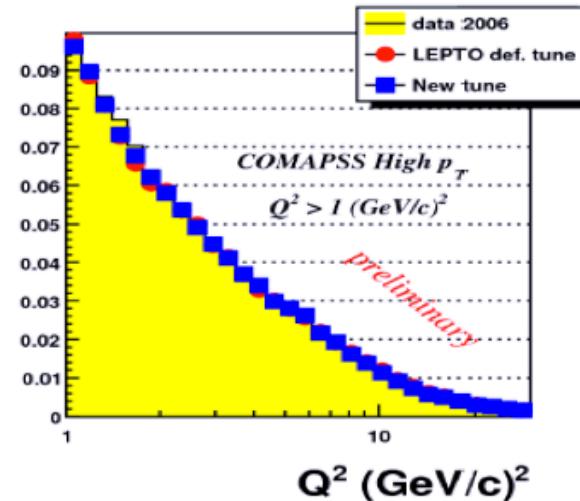
PARJ(21)	PARJ(23)	PARJ(24)	PARJ(41)	PARJ(42)
Transverse momentum of the hadron fragmentation			Fragmentation function	

- These **parameters** can be **divided** into **two sets** regarding the **component** of the **trajectory** of the particles: **Transverse** and **longitudinal** variable components.
- The **sets** can be **tuned** independently.  
⇒ The tuning improves substantially the Data-MC agreement.

# Data vs Monte Carlo comparison



Monte Carlo (PS on): LEPTO generator with PDFs from MSTW2008LO

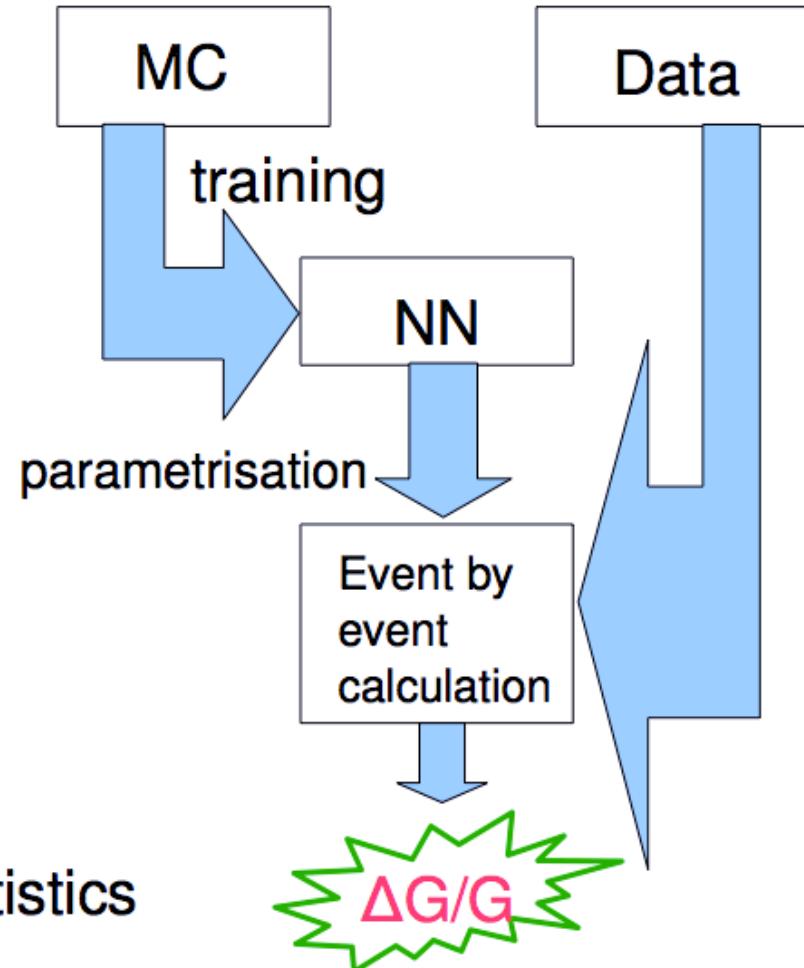


The impact of this tuning is included in the systematic error



A Neural Network is used to assign to each event a probability to be originated from each of the three processes (LO, PGF or Compton).

- A **MC** sample is used to train the Neural Network (NN).
- A parametrisation is constructed for all variables involved in the weight.
- A **Data** sample is weighted on an event-by-event basis.

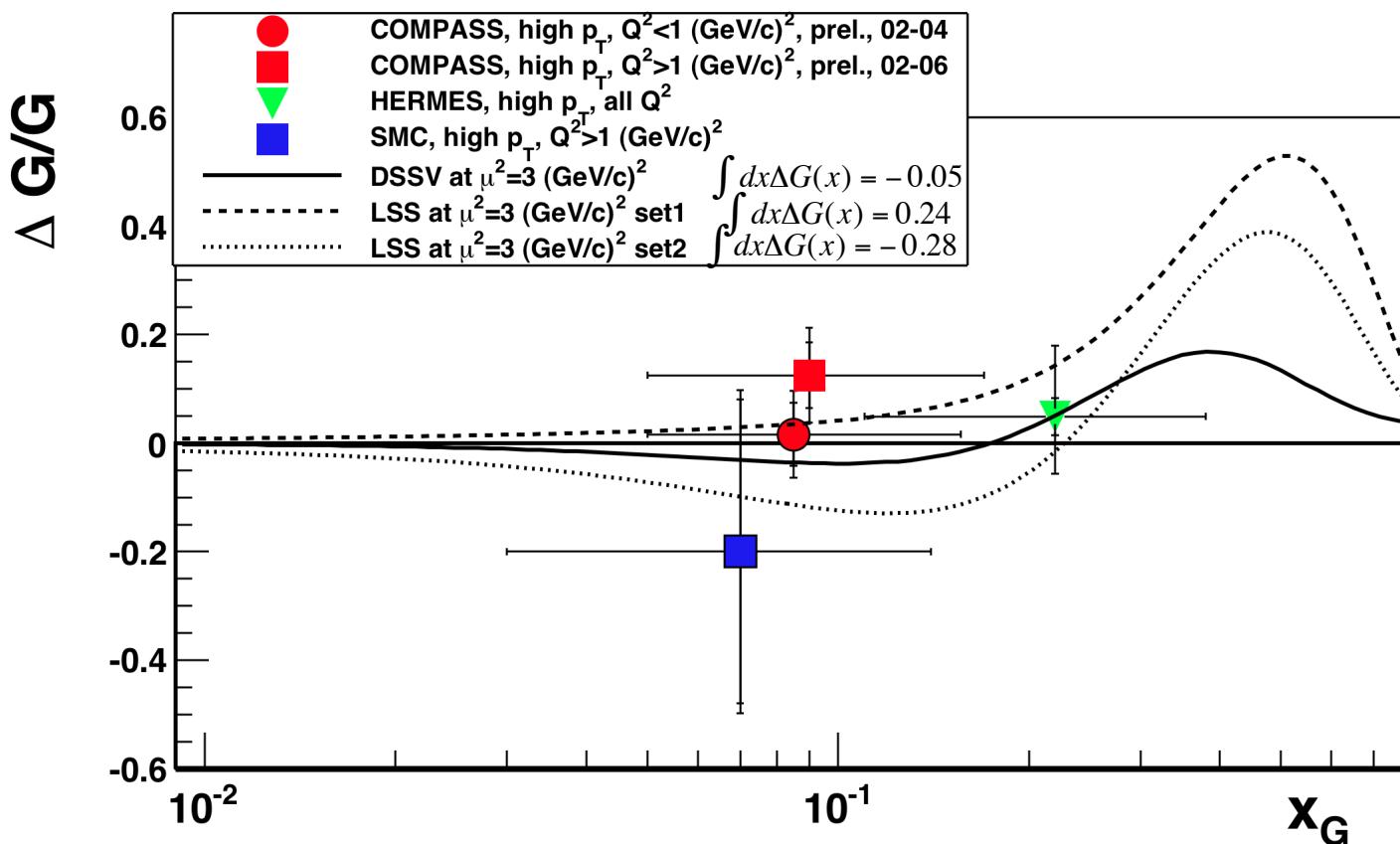


Optimal usage of the data sample statistics

# LO $\Delta G/G$ results high $p_T Q^2 > 1$ ( $\text{GeV}/c$ ) $^2$



$$\frac{\Delta G}{G} = 0.125 \pm 0.060 \pm 0.063 \quad x_g = 0.09^{+0.08}_{-0.04} \quad \langle \mu^2 \rangle = 3.4 (\text{GeV}/c)^2$$



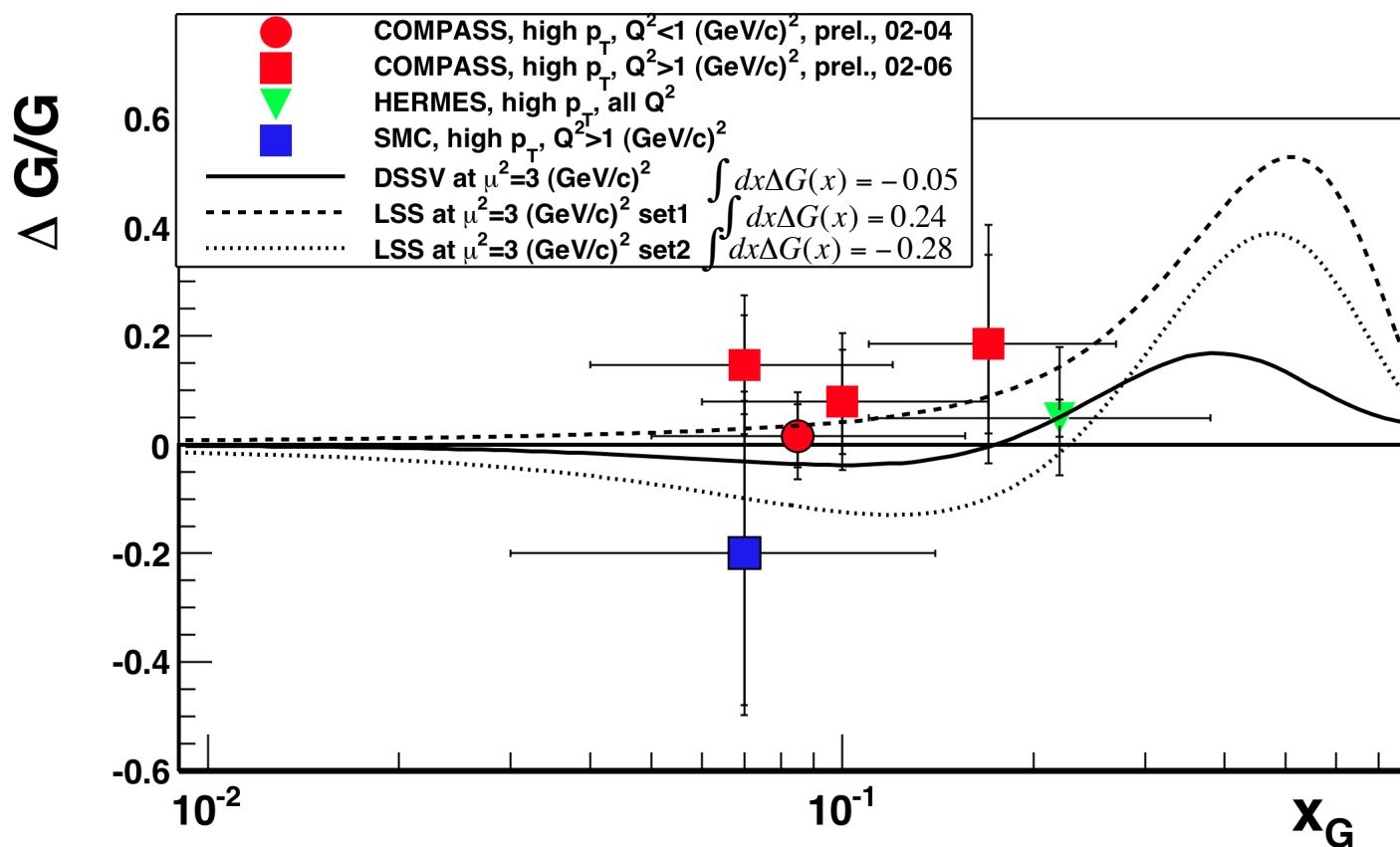
DSSV: D. de Florian et al., Phys. Rev. D80(2009)034030

LSS: E. Leader, A.V. Sidorov, D.B. Stamenov, arXiv 1010.5742(2010)

# LO $\Delta G/G$ results high $p_T$ $Q^2 > 1$ ( $\text{GeV}/c$ ) $^2$



1<sup>st</sup> point :  $0.147 \pm 0.091_{\text{stat}} \pm 0.088_{\text{sys}}$  @  $x_g = 0.07^{+0.05}_{-0.03}$   
 2<sup>nd</sup> point :  $0.079 \pm 0.096_{\text{stat}} \pm 0.082_{\text{sys}}$  @  $x_g = 0.10^{+0.07}_{-0.04}$   
 3<sup>rd</sup> point :  $0.185 \pm 0.165_{\text{stat}} \pm 0.143_{\text{sys}}$  @  $x_g = 0.17^{+0.10}_{-0.06}$



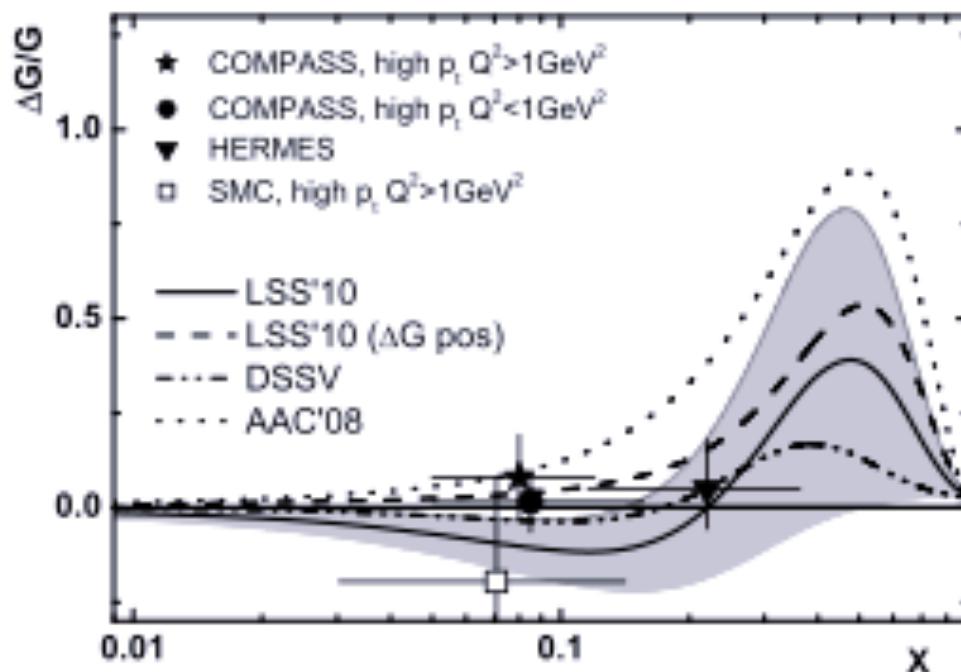
# Comparison to NLO fits to DIS & SIDIS



LSS: E. Leader, A.V. Sidorov, D.B. Stamenov, arXiv 1010.5742(2010)

TABLE IV: First moments of polarized PDFs at  $Q^2 = 4 \text{ GeV}^2$ . The corresponding DSSV values are also presented.

Fit	$\Delta\bar{s}$	$\Delta G$	$\Delta\Sigma$
LSS10 (pos $x\Delta G$ )	$-0.063 \pm 0.004$	$0.316 \pm 0.190$	$0.207 \pm 0.034$
LSS10 (node $x\Delta G$ )	$-0.055 \pm 0.006$	$-0.339 \pm 0.458$	$0.254 \pm 0.042$
DSSV (node $x\Delta G$ )	$-0.056$	$-0.096$	$0.245$



$$\begin{aligned} J_z &= \frac{1}{2} = \frac{1}{2}\Delta\Sigma(Q^2) + \Delta G(Q^2) + L_z(Q^2) \\ &= -0.21 \pm 0.46 + L_z(Q^2) \quad (\text{node } \Delta G) \\ &= 0.42 \pm 0.19 + L_z(Q^2) \quad (\text{pos } \Delta G). \end{aligned}$$

## Conclusion



- New extraction of  $\Delta G/G$  at LO from high  $p_T$  pairs at  $Q^2 > 1$  (GeV/c) $^2$  and  $x_G$  [0.04-0.27]
- Compatible with 0 and global QCD fits at NLO:
  - large  $\Delta G$  (2-3) excluded (  $a_0 = \Delta \Sigma - n_f \frac{\alpha_s}{2\pi} \Delta G$  )
  - still large error bars, and no NLO treatment
    - not constraining the spin puzzle
- Still to come:
  - $\Delta G/G$  from single high  $p_T$  hadrons at  $Q^2 < 1$  (GeV/c) $^2$  with NLO treatment at COMPASS energies.

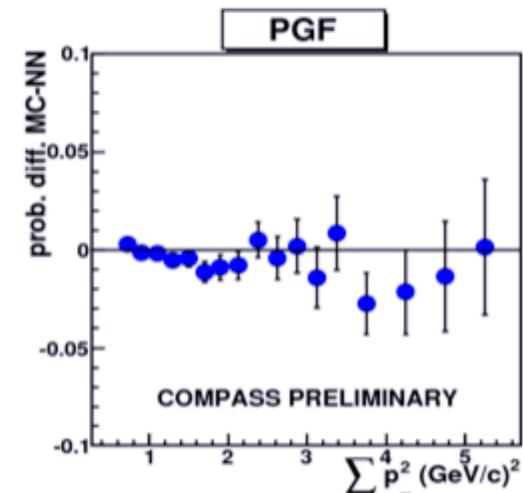
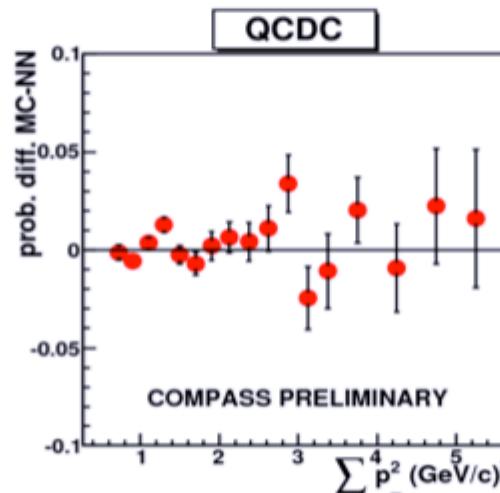
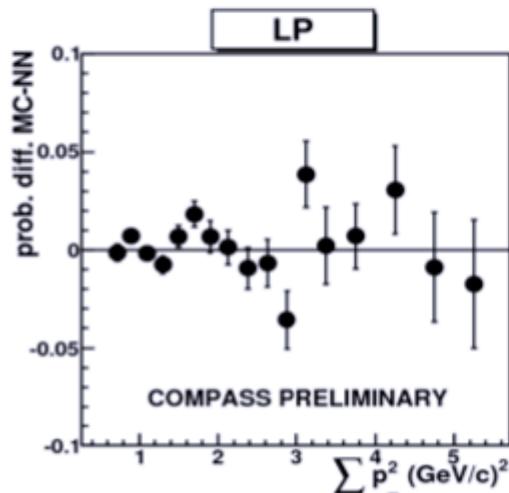
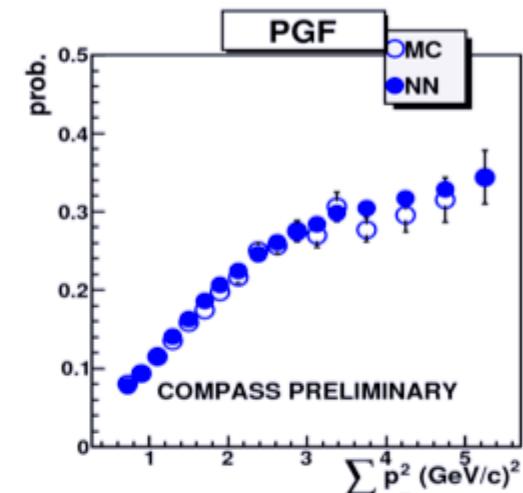
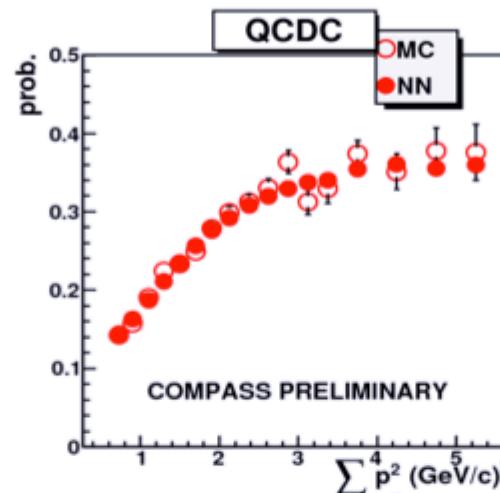
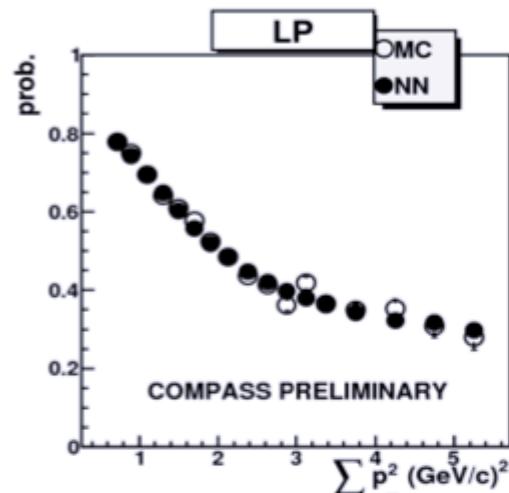
# Spares



# Stability checks of Neural Network



We parametrise the  $R^i$  fractions as probabilities.



# DSSV NLO fits to DIS & SIDIS

