



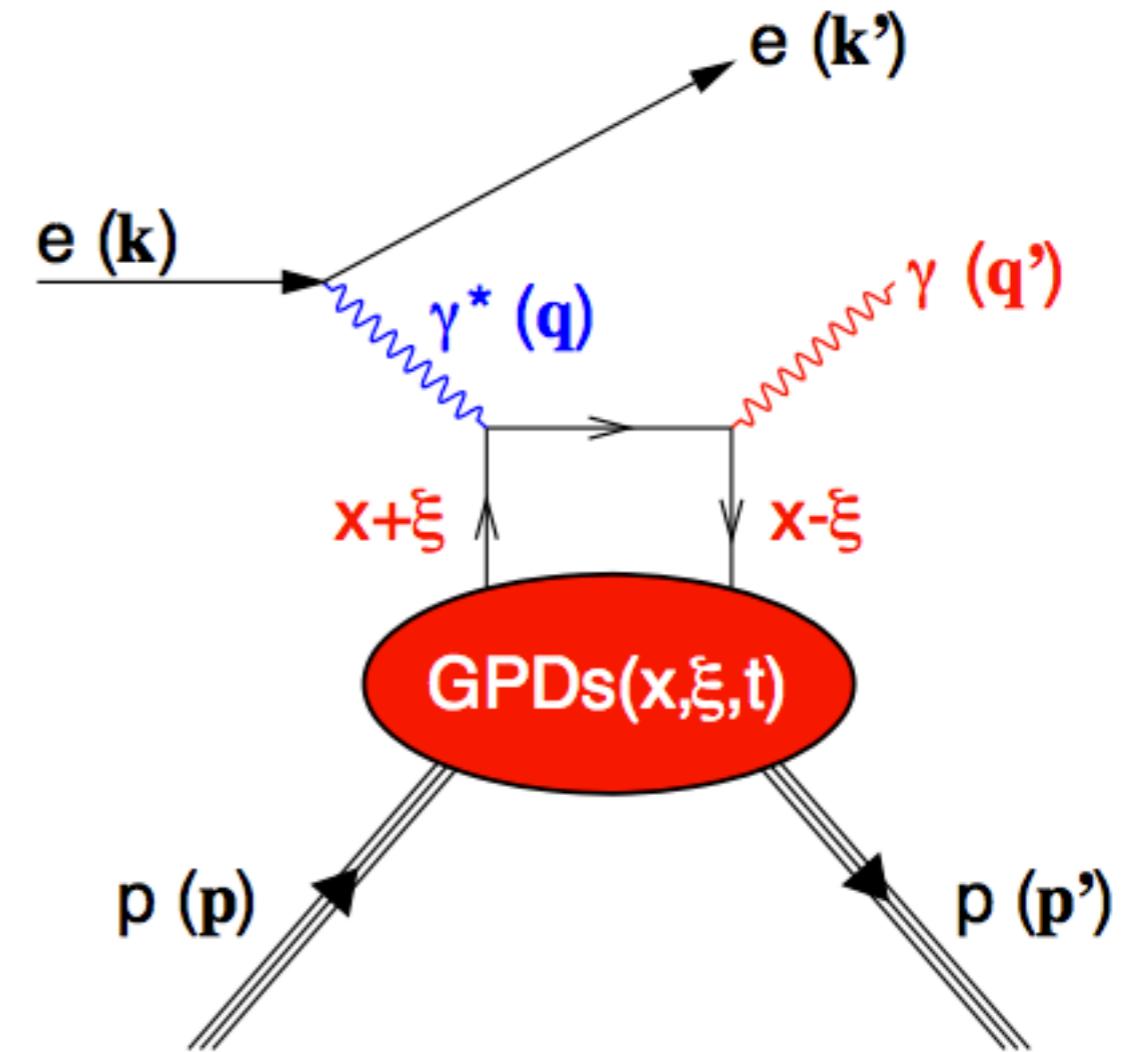
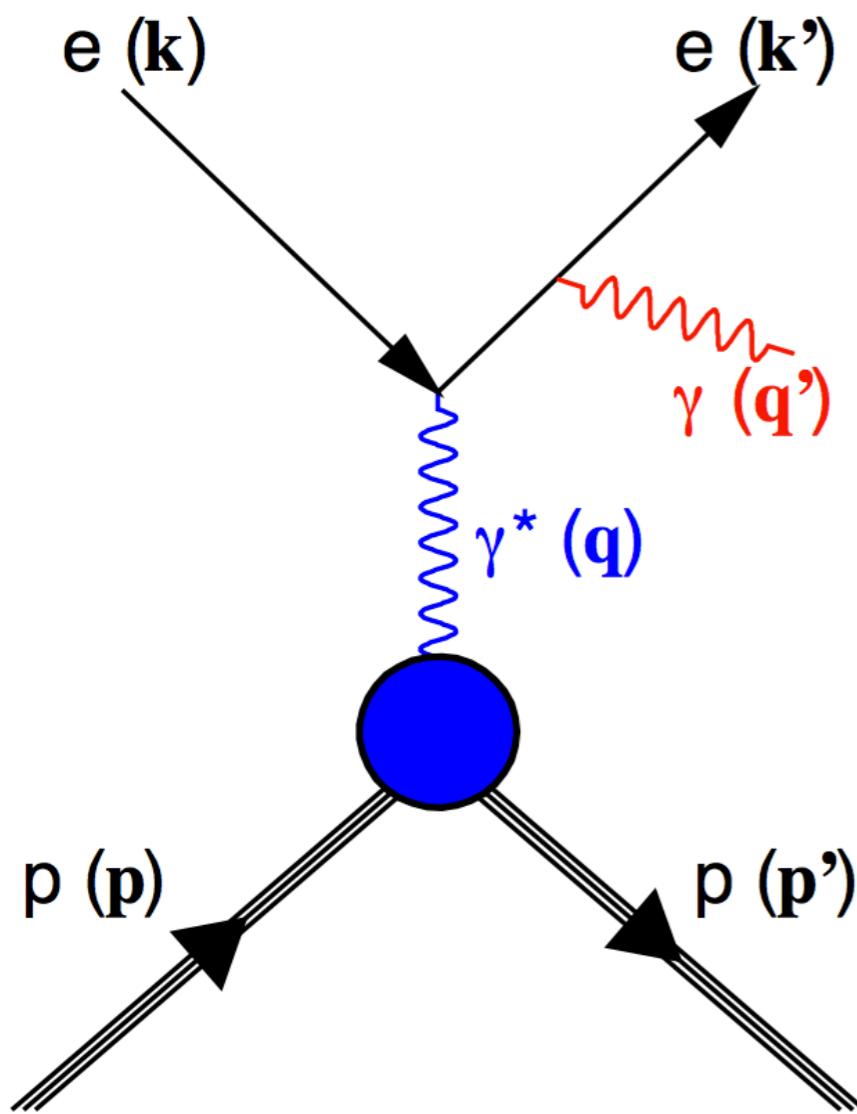
DVCS @ HERMES

M. MURRAY, UNIVERSITY OF GLASGOW
GPD 2010



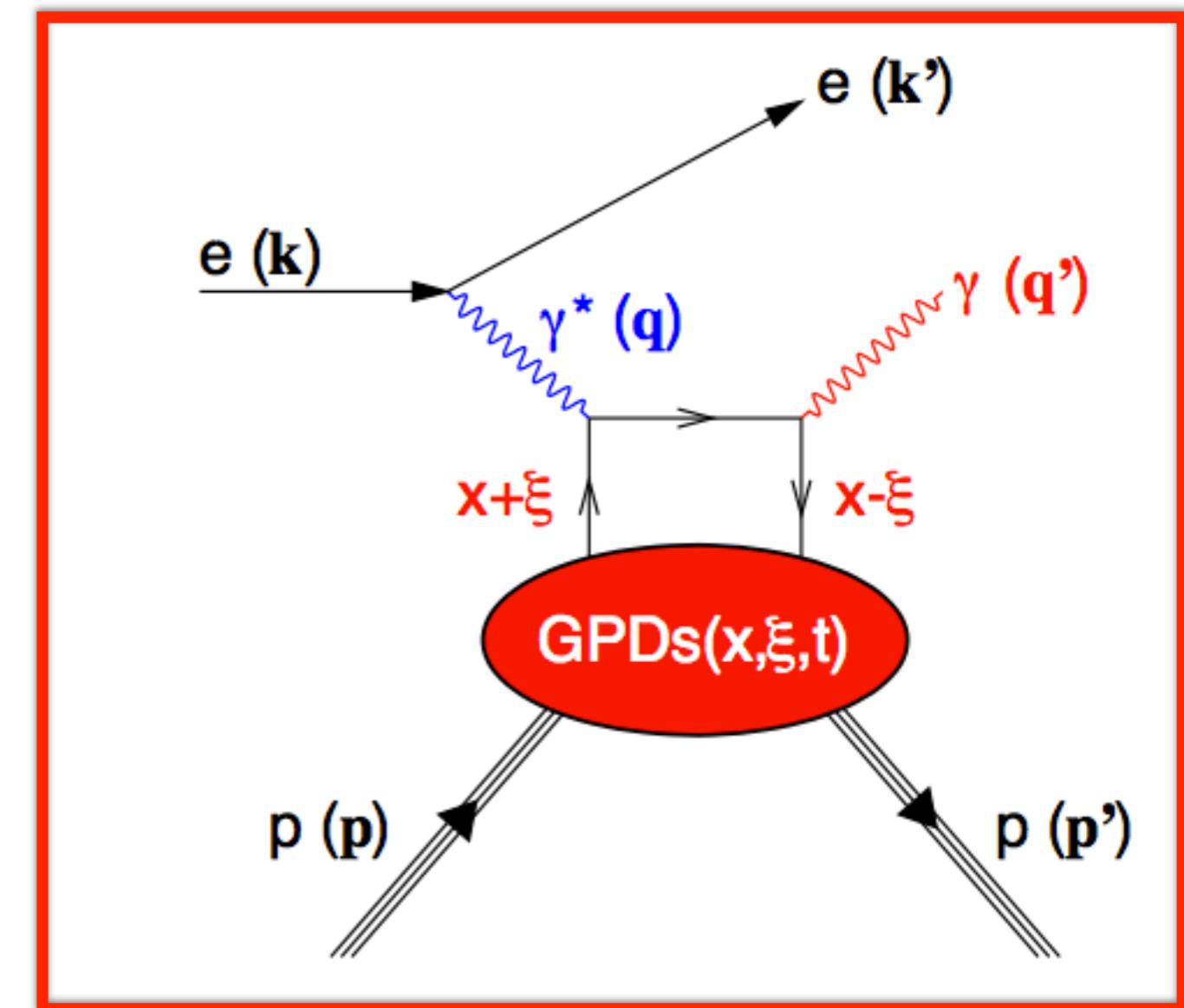
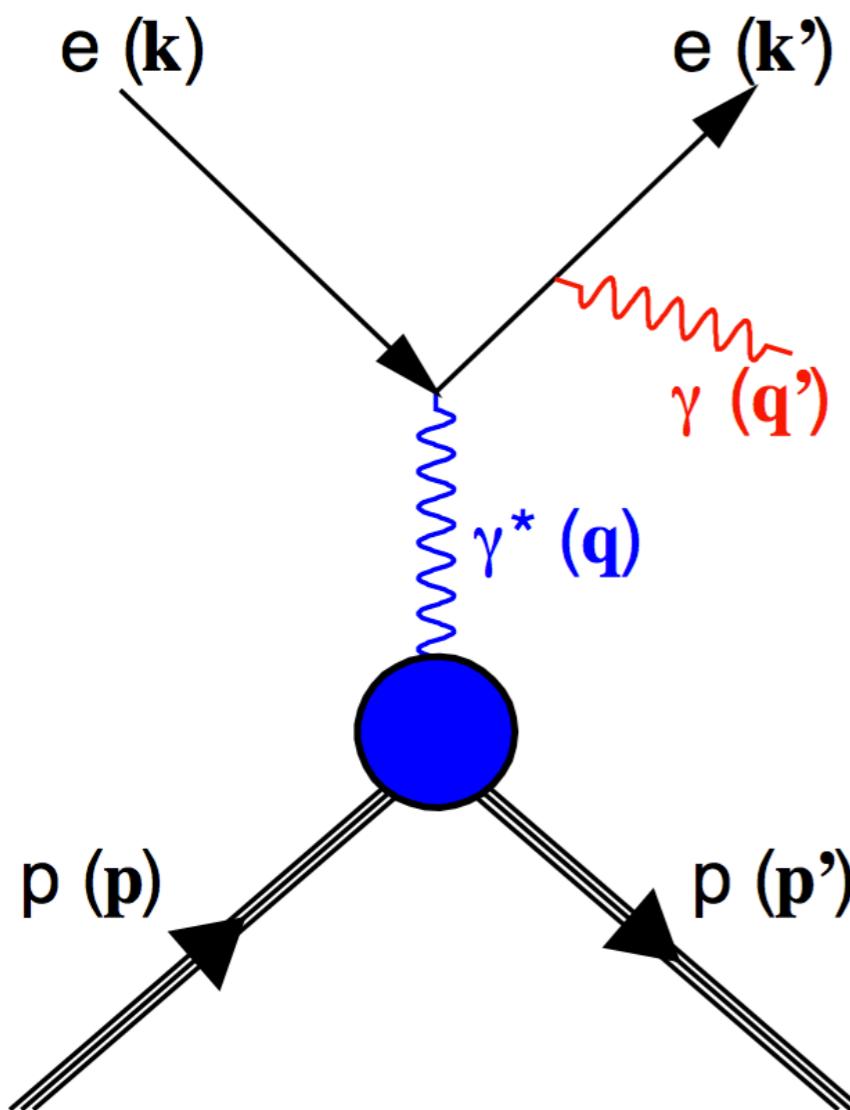
Deeply Virtual Compton Scattering

$e p \rightarrow e p \gamma$

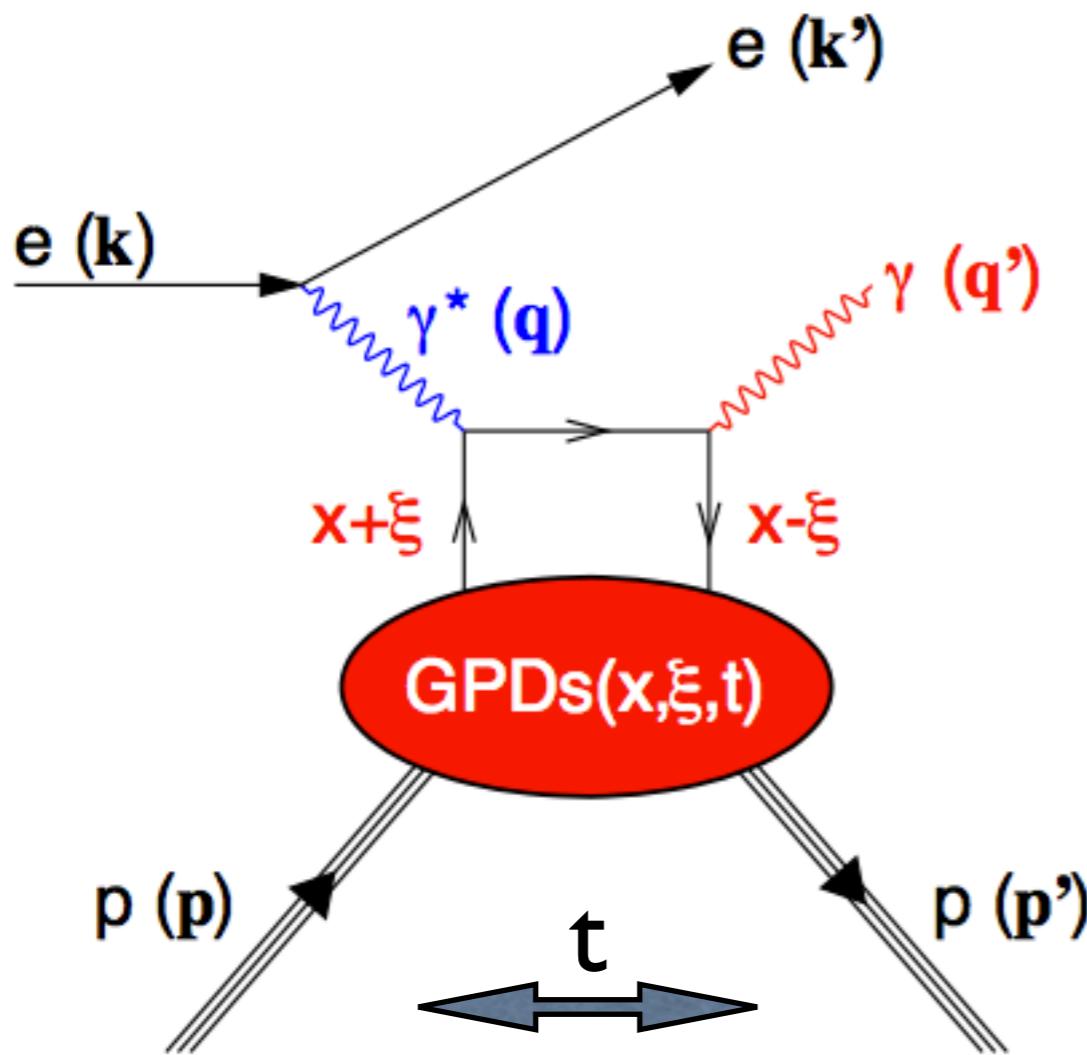


Deeply Virtual Compton Scattering

$$\frac{d\sigma}{dx_B dQ^2 d|t| d\phi} = \frac{x_B e^6 |\tau|^2}{32(2\pi)^4 Q^4 \sqrt{1 + \epsilon^2}} \rightarrow |\tau|^2 = |\tau_{\text{BH}}|^2 + |\tau_{\text{DVCS}}|^2 + \overbrace{\tau_{\text{BH}} \tau_{\text{DVCS}}^* + \tau_{\text{BH}}^* \tau_{\text{DVCS}}}^{\mathcal{I}}$$



Generalised Parton Distributions



t - Mandelstam variable
(squared momentum transfer to nucleon)

x - Fraction of nucleon's longitudinal momentum carried by active quark

ξ - half the change in the longitudinal momentum of the active quark.

GPD Physics

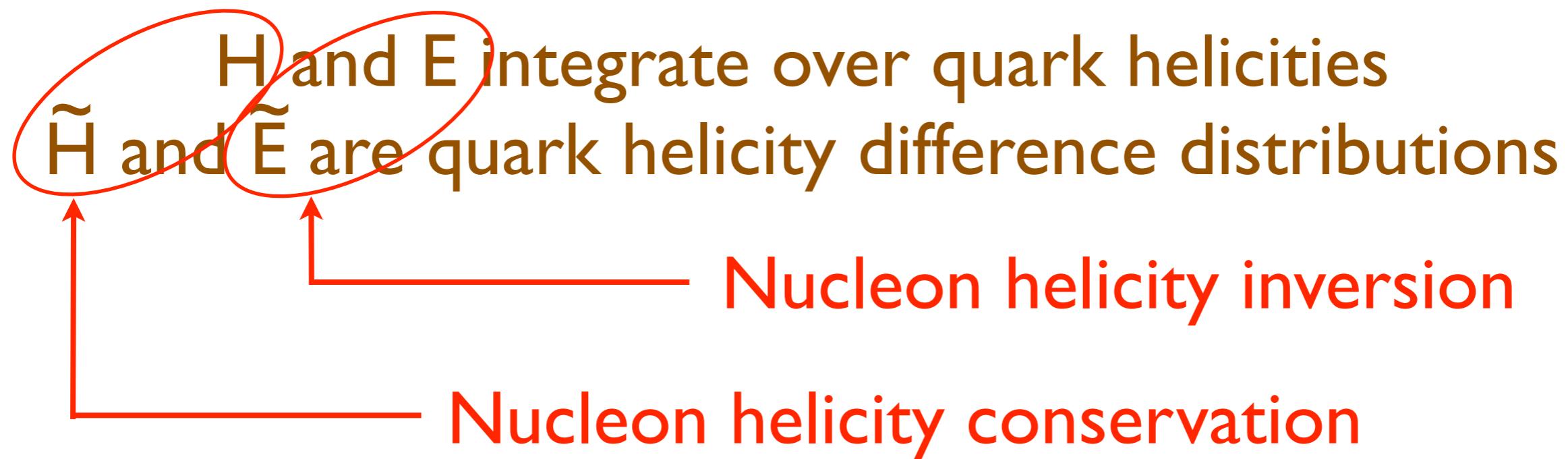
Four distributions of interest: $H, E, \tilde{H}, \tilde{E}$

H and E integrate over quark helicities
 \tilde{H} and \tilde{E} are quark helicity difference distributions

$$J_q = \frac{1}{2} \lim_{t \rightarrow 0} \int_{-1}^1 [H^q(x, \xi, t) + E^q(x, \xi, t)] x dx$$

GPD Physics

Four distributions of interest: $H, E, \tilde{H}, \tilde{E}$

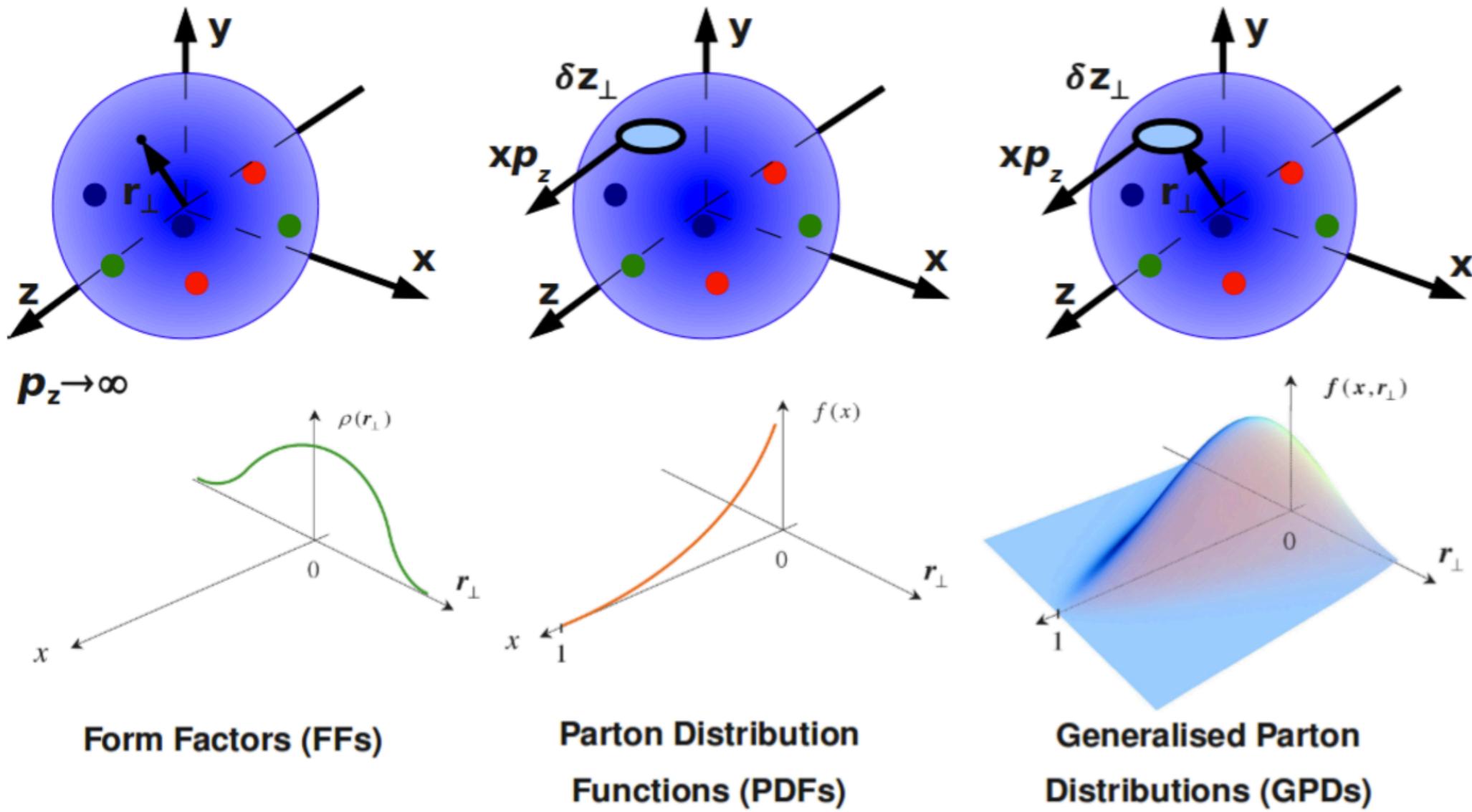


$$J_q = \frac{1}{2} \lim_{t \rightarrow 0} \int_{-1}^1 [H^q(x, \xi, t) + E^q(x, \xi, t)] x dx$$

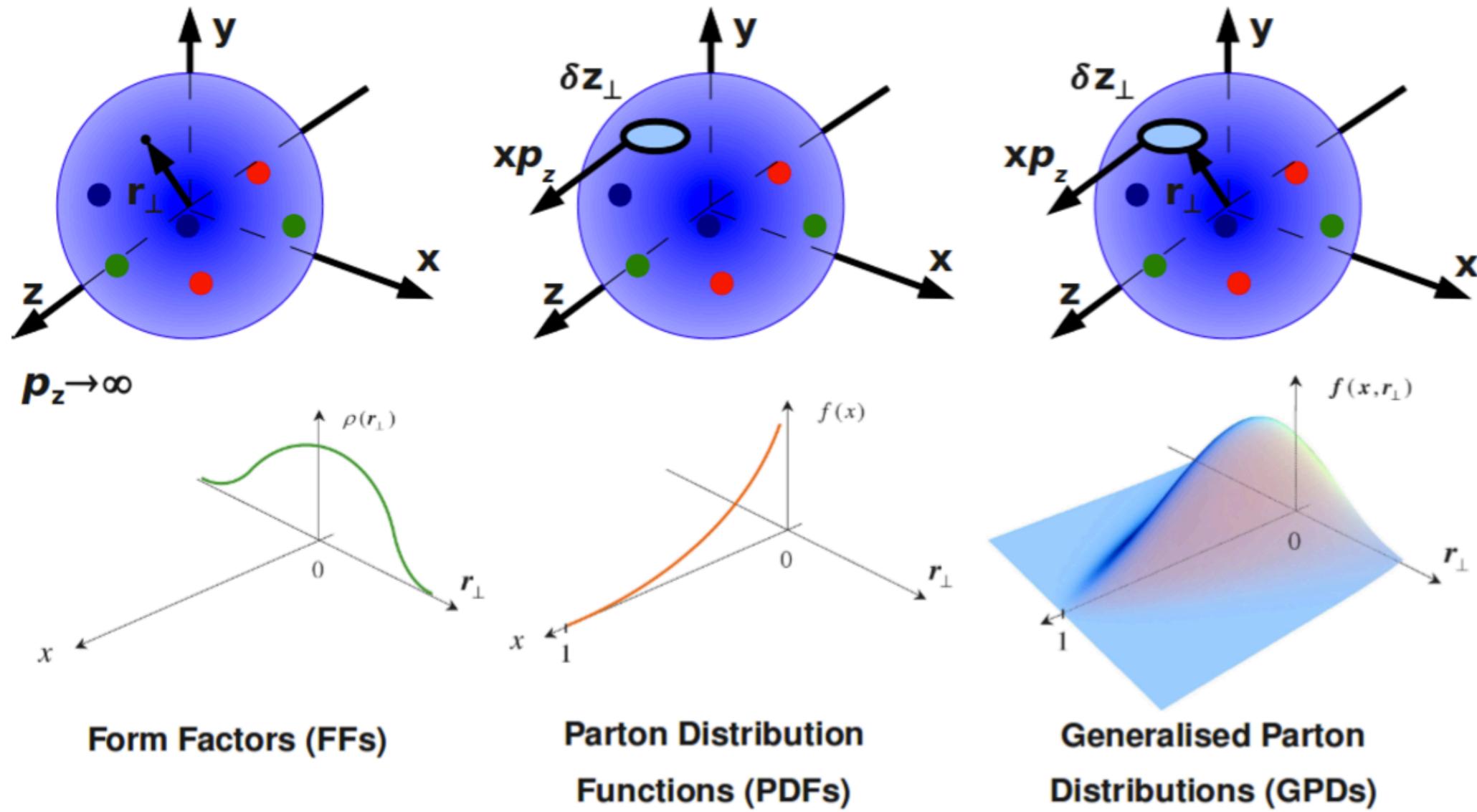
“Ji’s Relation”

Phys. Rev. Lett. 78:610, 1997

GPD Physics



GPD Physics



H - unpolarised nucleon

\tilde{H} - polarised nucleon

GPD Physics

GPDs describe only the soft part of the interaction

Accessed via cross-sections and asymmetries:
requires convolution with a hard scattering kernel

$$H \rightarrow \mathcal{H} \quad \tilde{H} \rightarrow \tilde{\mathcal{H}} \quad E \rightarrow \mathcal{E} \quad \tilde{E} \rightarrow \tilde{\mathcal{E}}$$

Results in “Compton Form Factors” accessible through DVCS, which have real and imaginary parts

GPD Physics

GPDs describe only the soft part of the interaction

Accessed via cross-sections and asymmetries:
requires convolution with a hard scattering kernel

$$\Im m \mathcal{F}(\xi, t) = F(\xi, \xi, t) \pm F(-\xi, \xi, t),$$

$$\Re e \mathcal{F}(\xi, t) = \mathcal{P}_C \int_{-1}^1 \frac{F(x, \xi, t)}{x - \xi} \pm \frac{F(x, \xi, t)}{x + \xi} dx$$

GPD Physics

GPDs describe only the soft part of the interaction

Accessed via cross-sections and asymmetries:
requires convolution with a hard scattering kernel

$$\Im m \mathcal{F}(\xi, t) = F(\xi, \xi, t) \pm F(-\xi, \xi, t),$$

$$\Re e \mathcal{F}(\xi, t) = \mathcal{P}_C \int_{-1}^1 \frac{F(x, \xi, t)}{x - \xi} \pm \frac{F(x, \xi, t)}{x + \xi} dx$$

Limited x access

DVCS @ HERMES

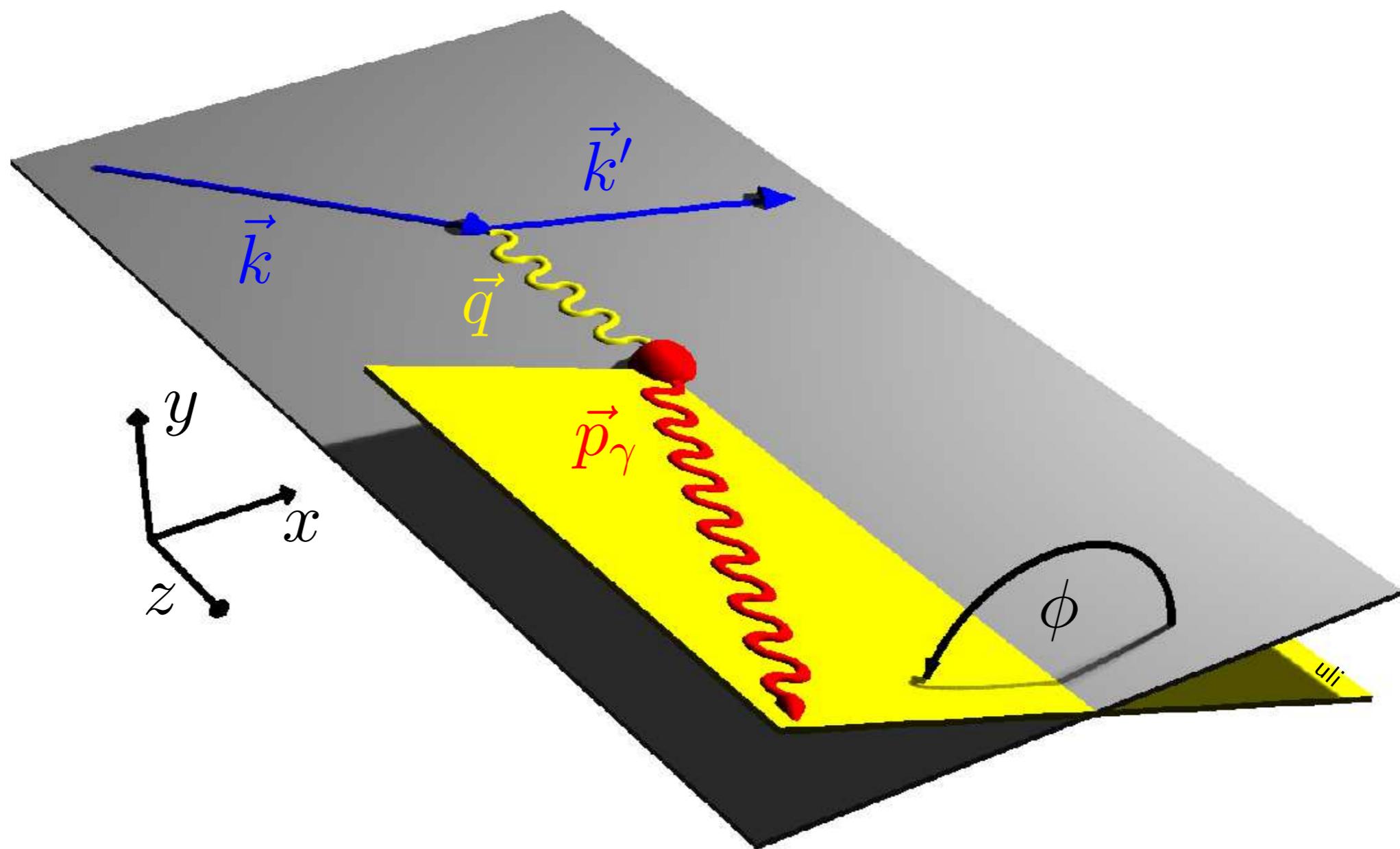
$$\mathcal{A}_C(\phi) \equiv \frac{d\sigma^+(\phi) - d\sigma^-(\phi)}{d\sigma^+(\phi) + d\sigma^-(\phi)} \propto \text{Re}(\mathcal{H})$$

$$\mathcal{A}_{LU}(\phi) \equiv \frac{[\sigma^{\rightarrow\leftarrow}(\phi) + \sigma^{\rightarrow\Rightarrow}(\phi)] - [\sigma^{\leftarrow\leftarrow}(\phi) + \sigma^{\leftarrow\Rightarrow}(\phi)]}{[\sigma^{\rightarrow\leftarrow}(\phi) + \sigma^{\rightarrow\Rightarrow}(\phi)] + [\sigma^{\leftarrow\leftarrow}(\phi) + \sigma^{\leftarrow\Rightarrow}(\phi)]} \propto \text{Im}(\mathcal{H})$$

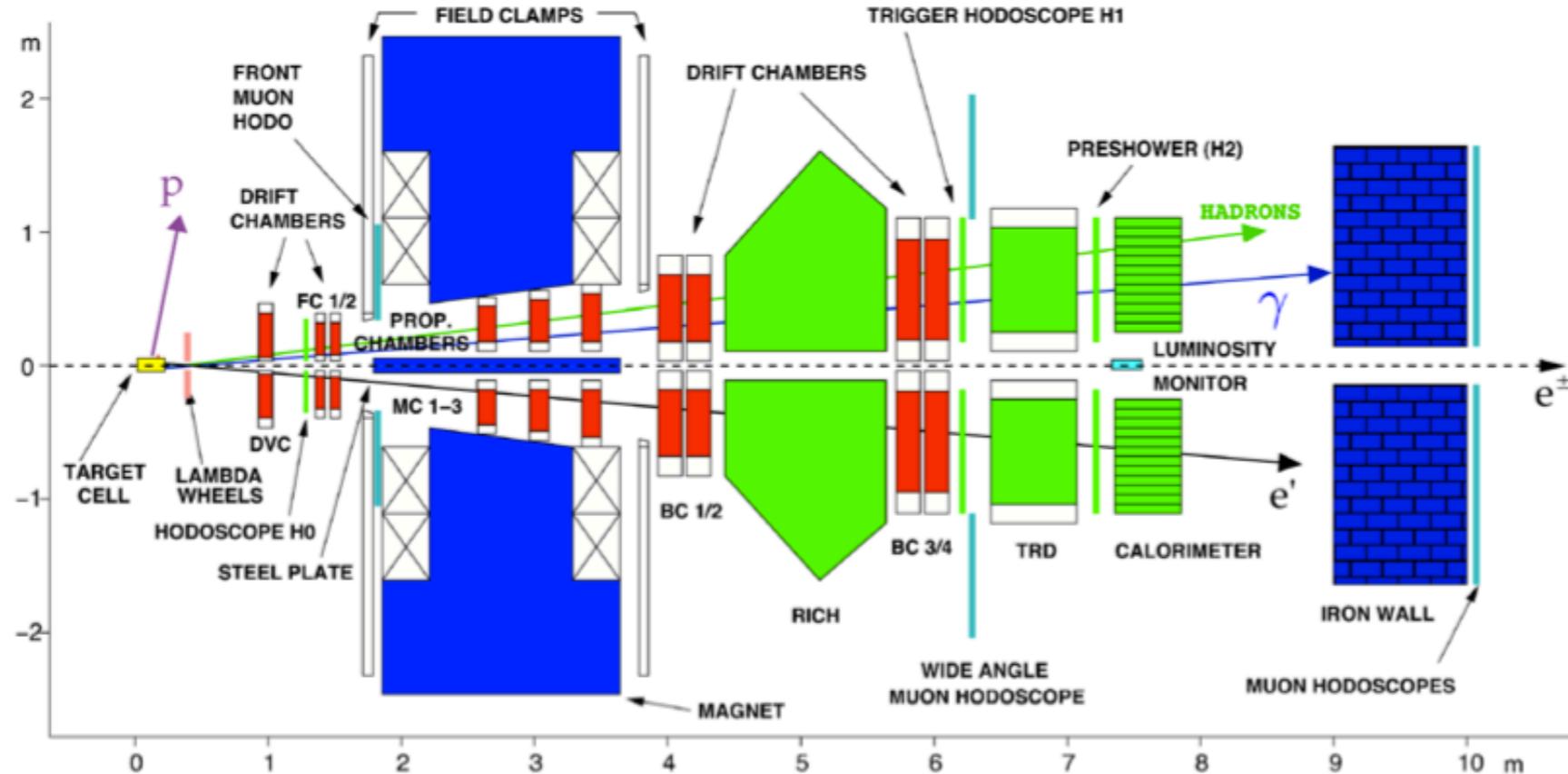
$$\mathcal{A}_{UL}(\phi) \equiv \frac{[\sigma^{\leftarrow\Rightarrow}(\phi) + \sigma^{\rightarrow\Rightarrow}(\phi)] - [\sigma^{\leftarrow\leftarrow}(\phi) + \sigma^{\rightarrow\leftarrow}(\phi)]}{[\sigma^{\leftarrow\Rightarrow}(\phi) + \sigma^{\rightarrow\Rightarrow}(\phi)] + [\sigma^{\leftarrow\leftarrow}(\phi) + \sigma^{\rightarrow\leftarrow}(\phi)]} \propto \text{Im}(\tilde{\mathcal{H}})$$

$$\mathcal{A}_{LL}(\phi) \equiv \frac{[\sigma^{\rightarrow\Rightarrow}(\phi) + \sigma^{\leftarrow\leftarrow}(\phi)] - [\sigma^{\leftarrow\Rightarrow}(\phi) + \sigma^{\rightarrow\leftarrow}(\phi)]}{[\sigma^{\rightarrow\Rightarrow}(\phi) + \sigma^{\leftarrow\leftarrow}(\phi)] + [\sigma^{\leftarrow\Rightarrow}(\phi) + \sigma^{\rightarrow\leftarrow}(\phi)]} \propto \text{Re}(\tilde{\mathcal{H}})$$

DVCS @ HERMES



DVCS @ HERMES



Forward spectrometer ⇒
measure asymmetries directly

$$\langle Q^2 \rangle \approx 2.4 \text{ GeV}^2$$

- $1 \text{ GeV}^2 < Q^2 \equiv -q^2 < 10 \text{ GeV}^2$

$$\langle x_B \rangle \approx 0.1$$

- $0.03 < x_B < 0.3$

$$\langle -t \rangle \approx 0.1 \text{ GeV}^2$$

- $0 \text{ GeV}^2 < -t \equiv -(p-p')^2 < 0.7 \text{ GeV}^2$

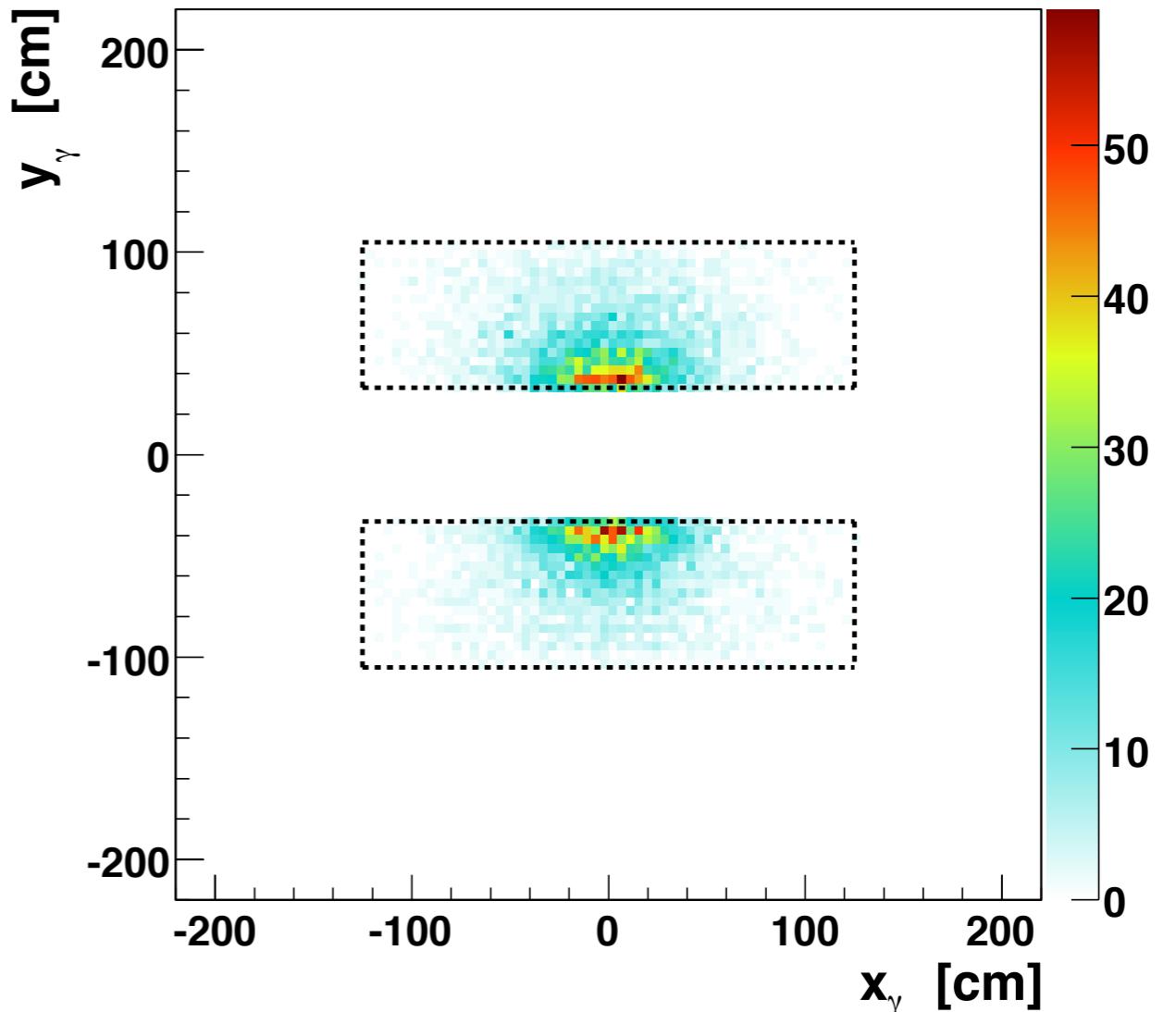
DVCS @ HERMES

Forward
spectrometer \Rightarrow
measure
asymmetries
directly

$$\langle Q^2 \rangle \approx 2.4 \text{ GeV}^2$$

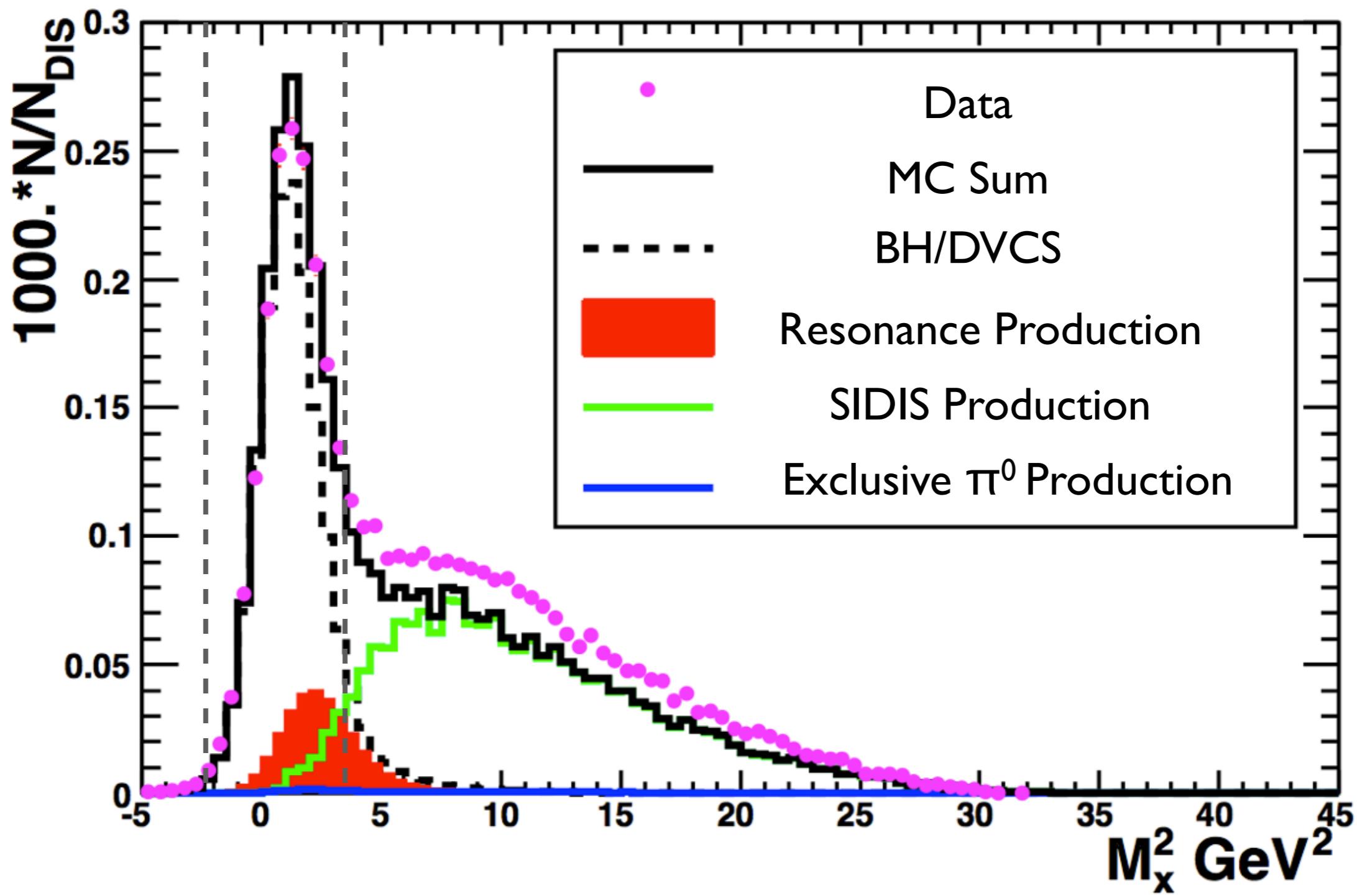
$$\langle x_B \rangle \approx 0.1$$

$$\langle -t \rangle \approx 0.1 \text{ GeV}^2$$

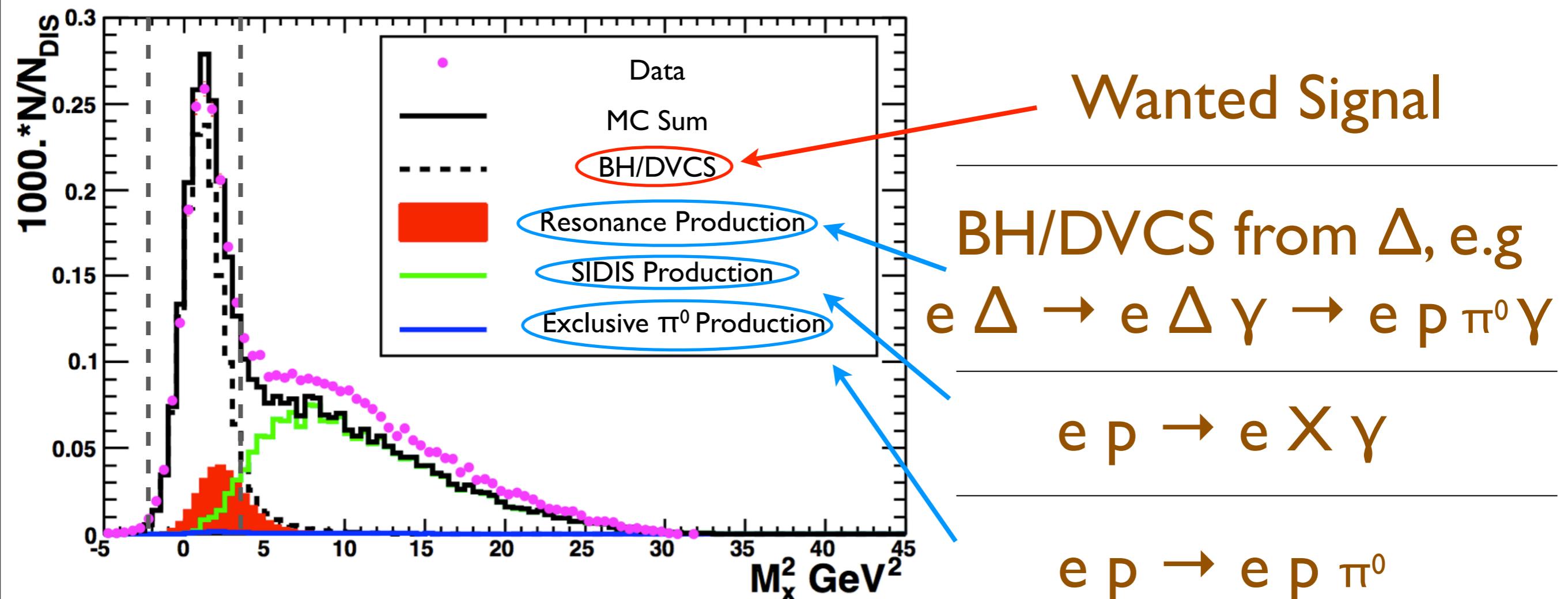


- $1 \text{ GeV}^2 < Q^2 \equiv -q^2 < 10 \text{ GeV}^2$
- $0.03 < x_B < 0.3$
- $0 \text{ GeV}^2 < -t \equiv -(p-p')^2 < 0.7 \text{ GeV}^2$

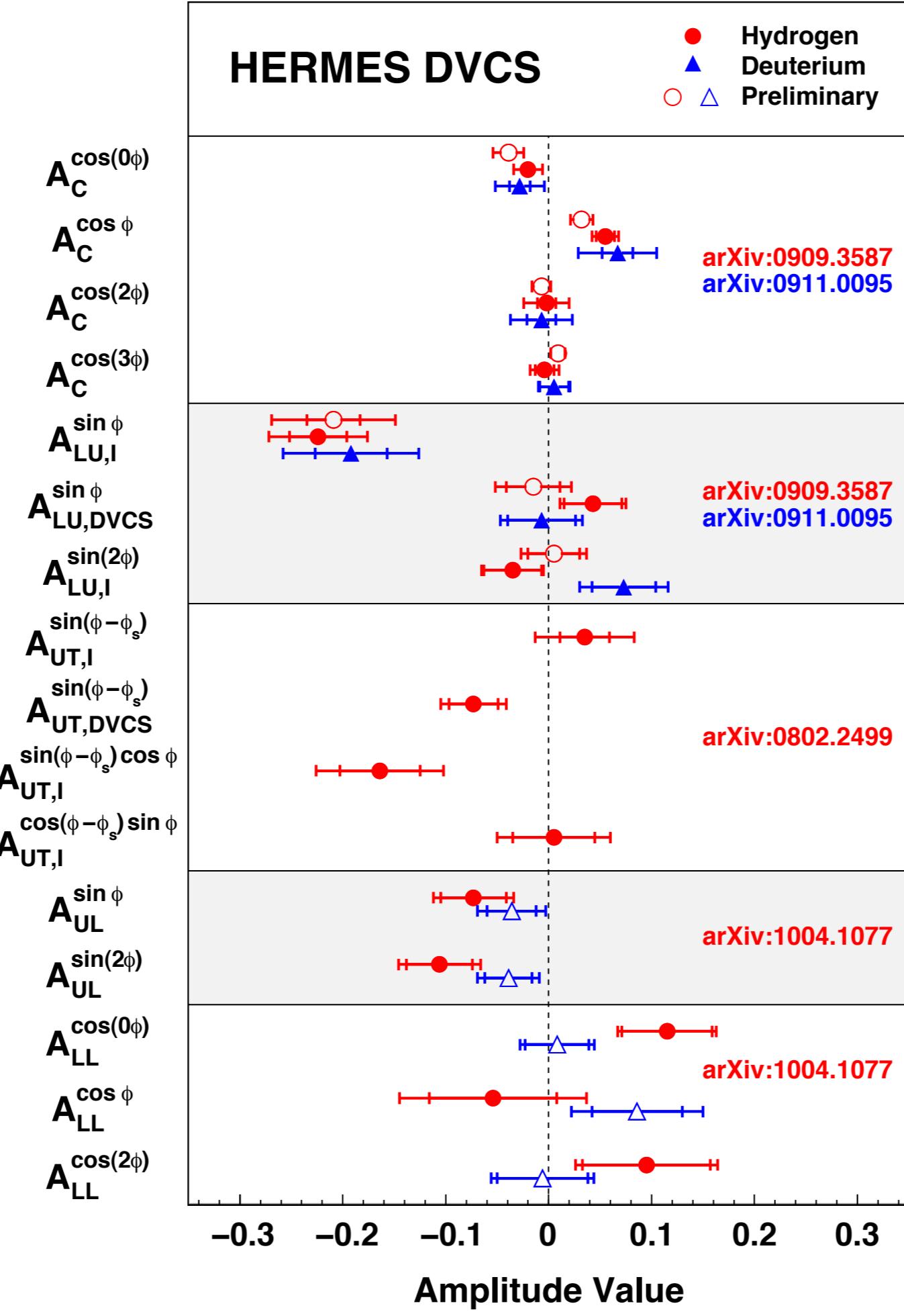
DVCS @ HERMES



DVCS @ HERMES

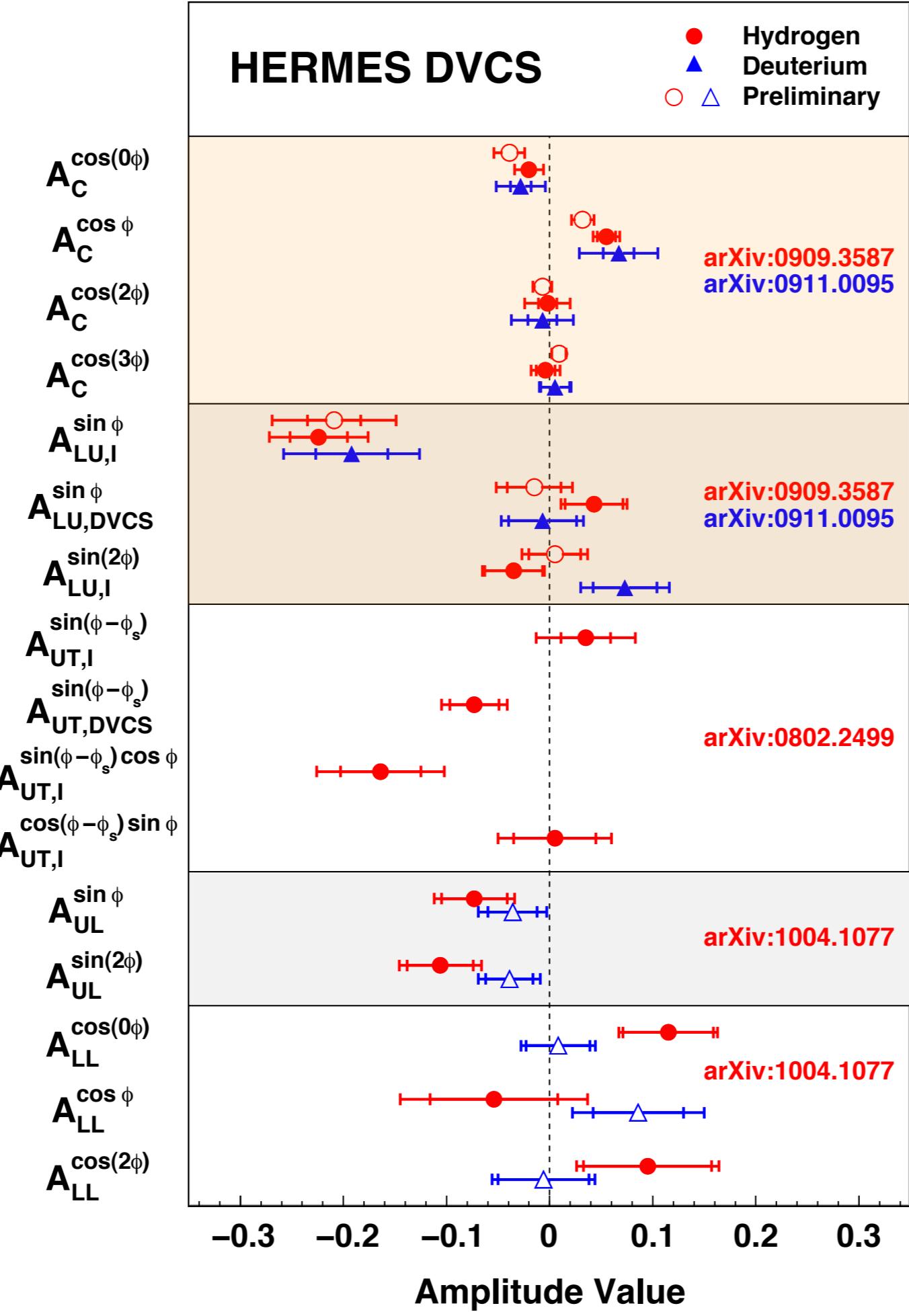


D
V
C
S
@



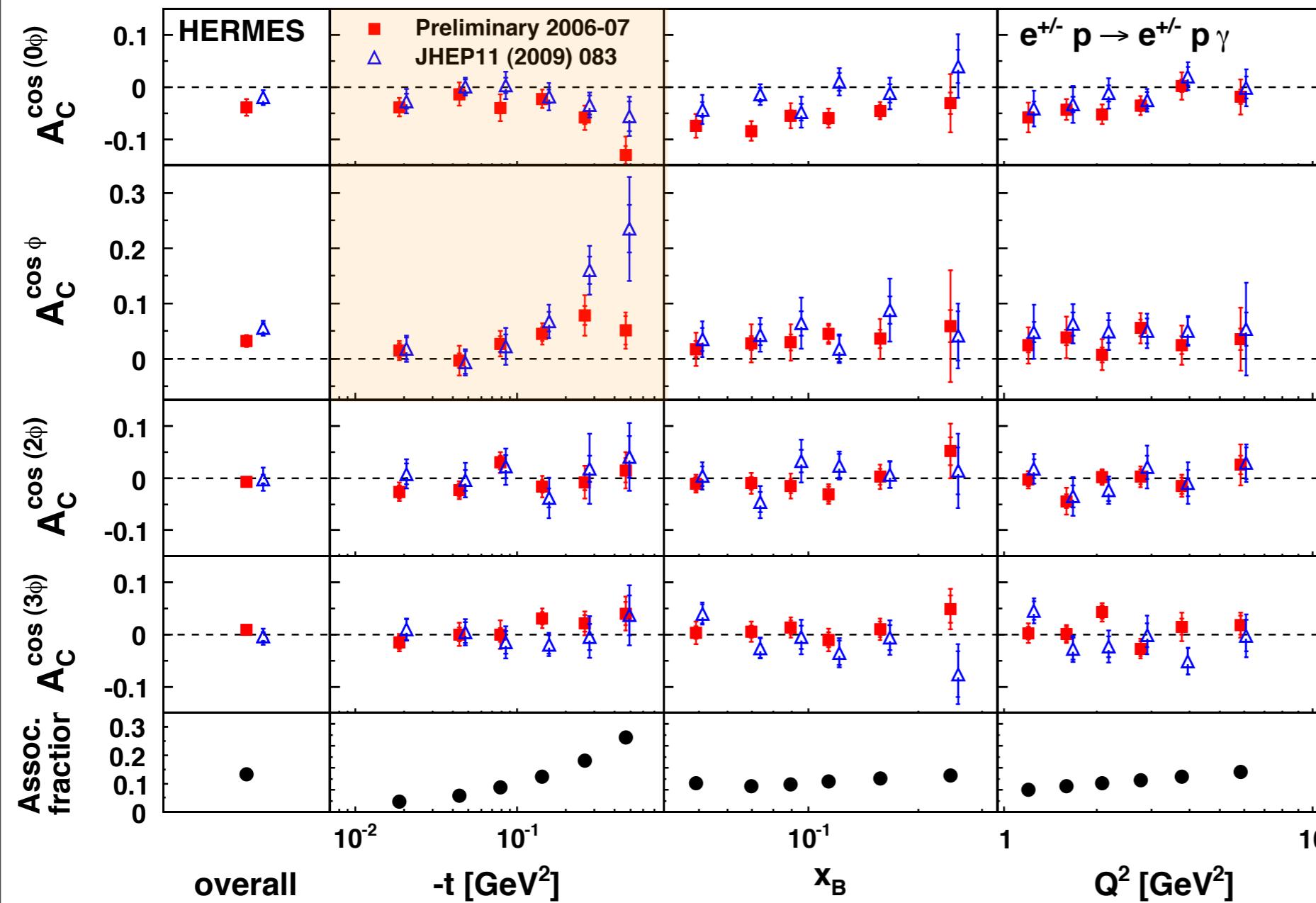
H
E
R
M
E
S

D
V
C
S
@



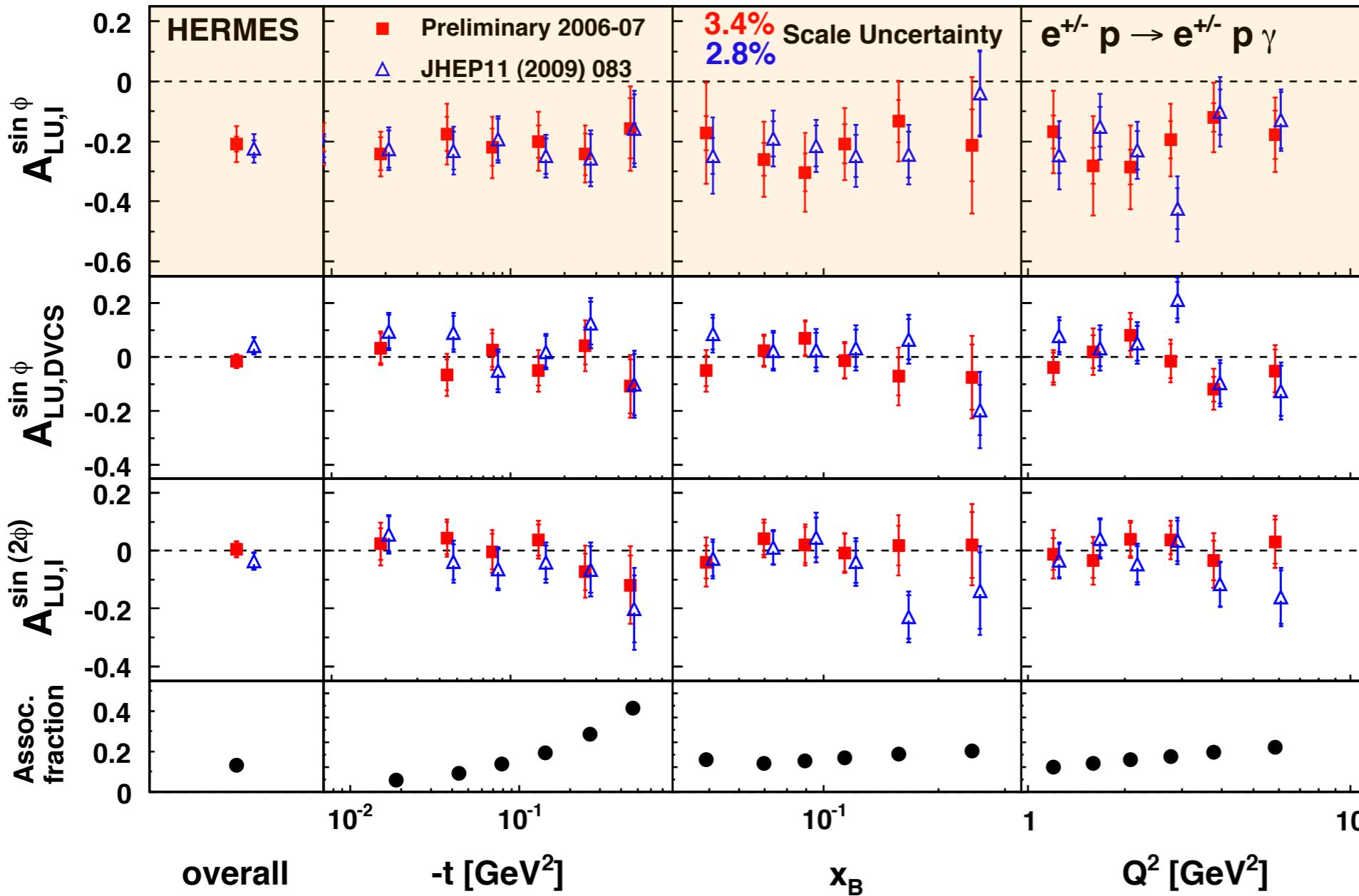
H
E
R
M
E
S

Beam Asymmetries



Beam Charge
Asymmetries
access $\text{Re}(\mathcal{H})$

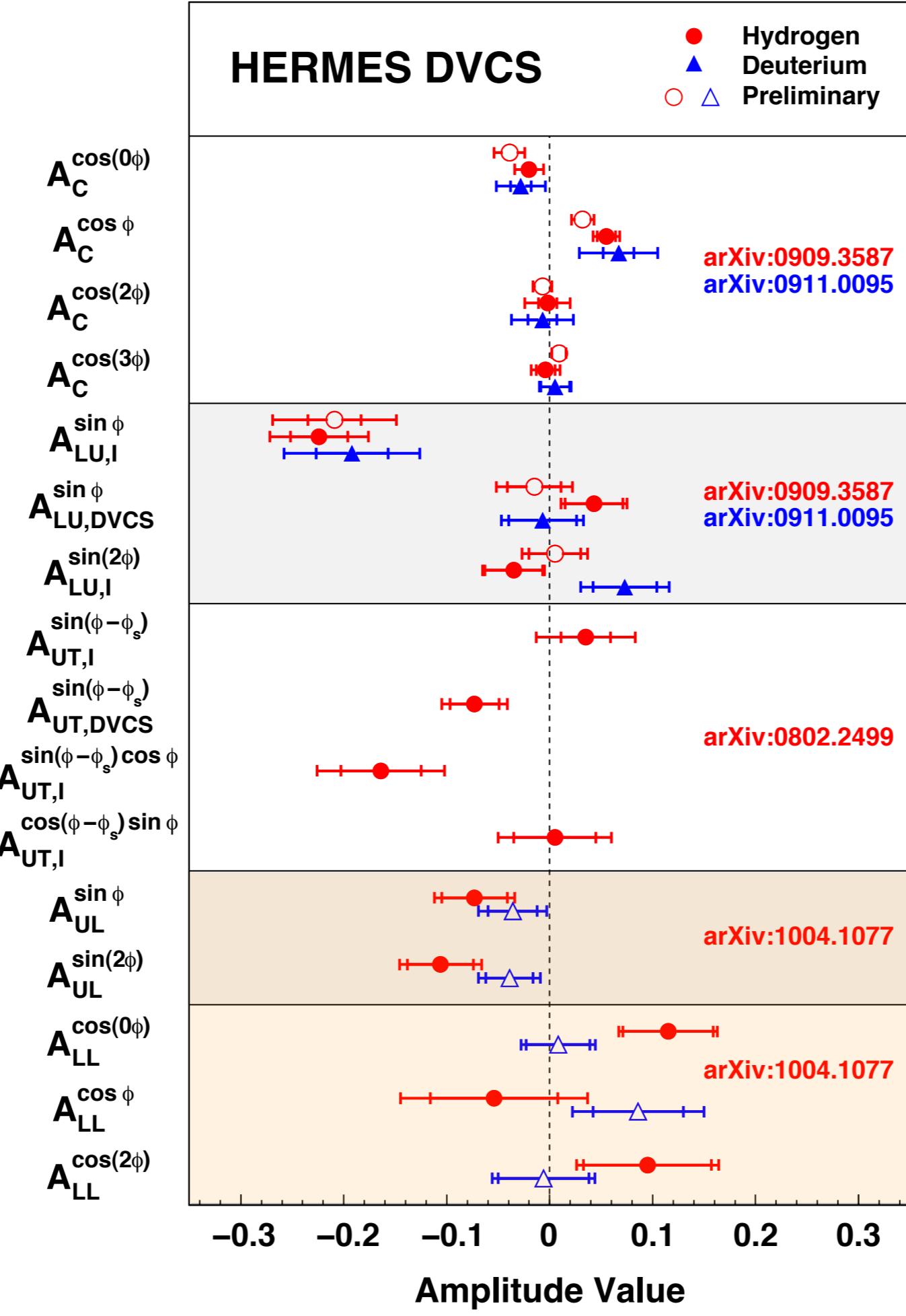
Beam Asymmetries



Beam Helicity
Asymmetries
access $\text{Im}(\mathcal{H})$

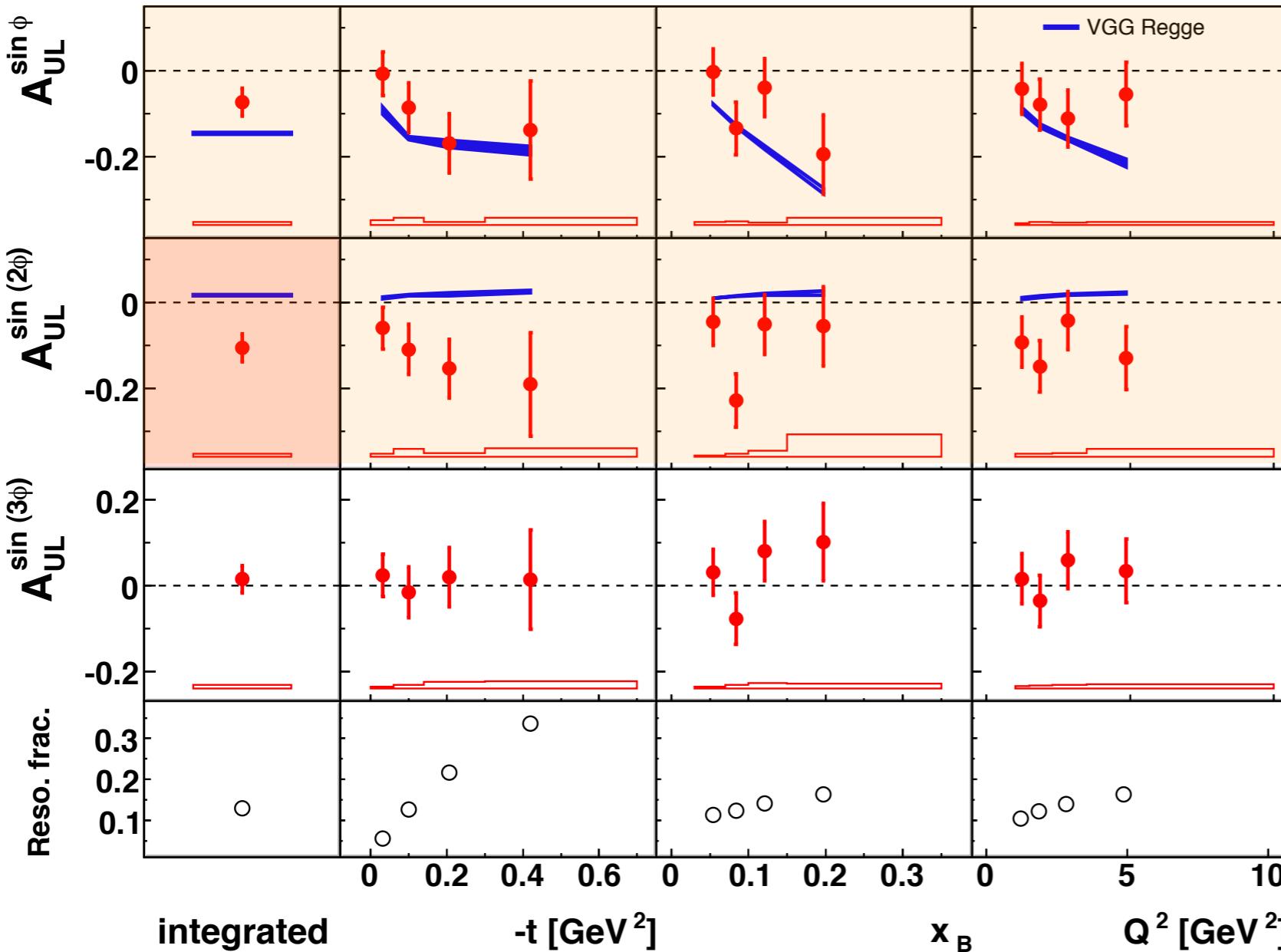
Larger values
for the BHA
than BCA -
node in the
small- x region
near to the
HERMES
kinematics?

D
V
C
S
@



H
E
R
M
E
S

Target Asymmetries



Long. Pol. target
asymmetries
access $\text{Im}(\tilde{\mathcal{H}})$

<http://arxiv.org/abs/1004.0177>

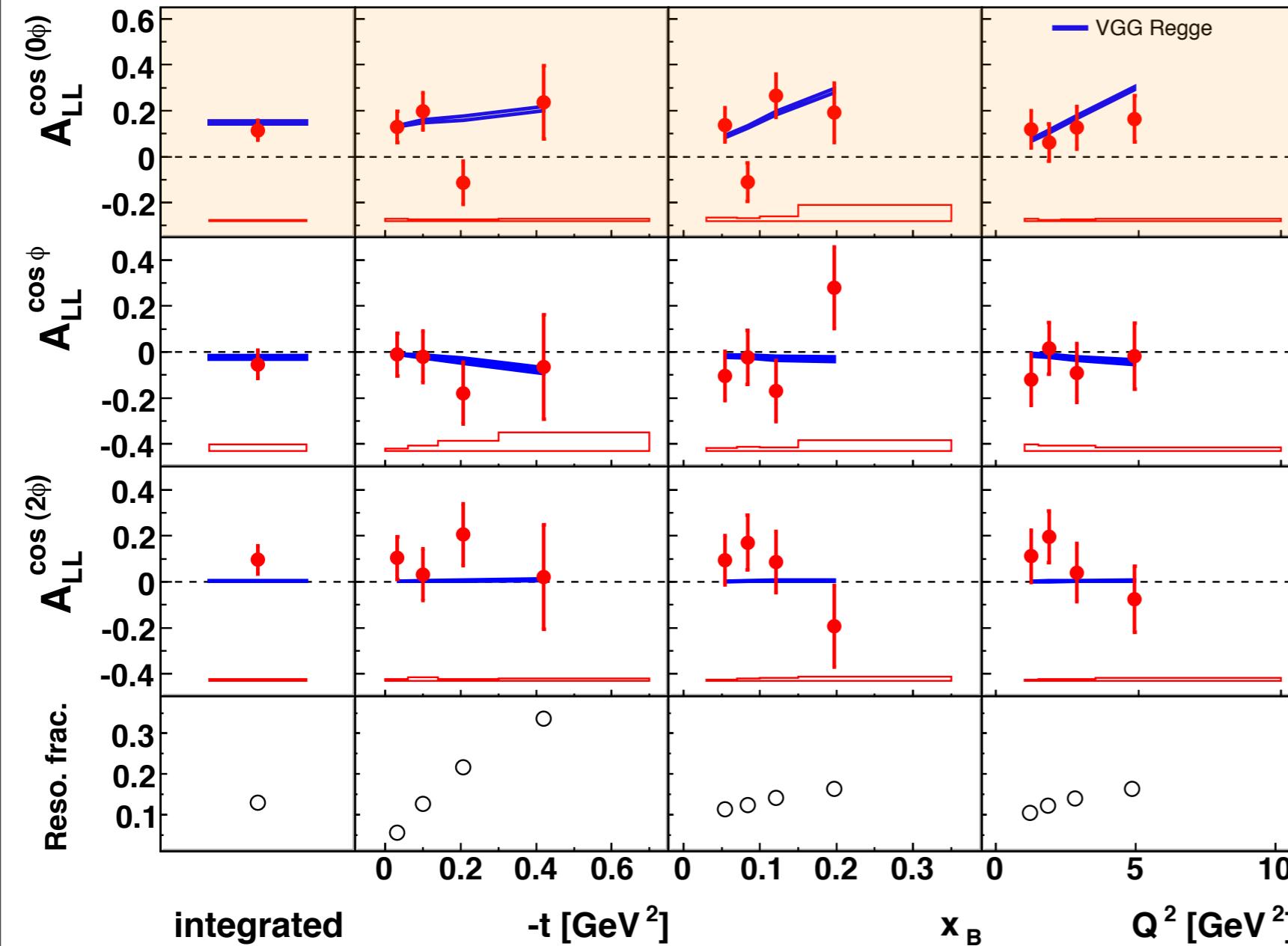
A. Airapetian et al, JHEP 06 (2010) 019

VGG Model:

<http://arxiv.org/abs/hep-ph/9905372>

Phys.Rev. D60 (1999) 094017

Double Spin Asymmetries



Long. Pol. target /
Long. Pol. Beam
access $\text{Re}(\tilde{\mathcal{H}})$

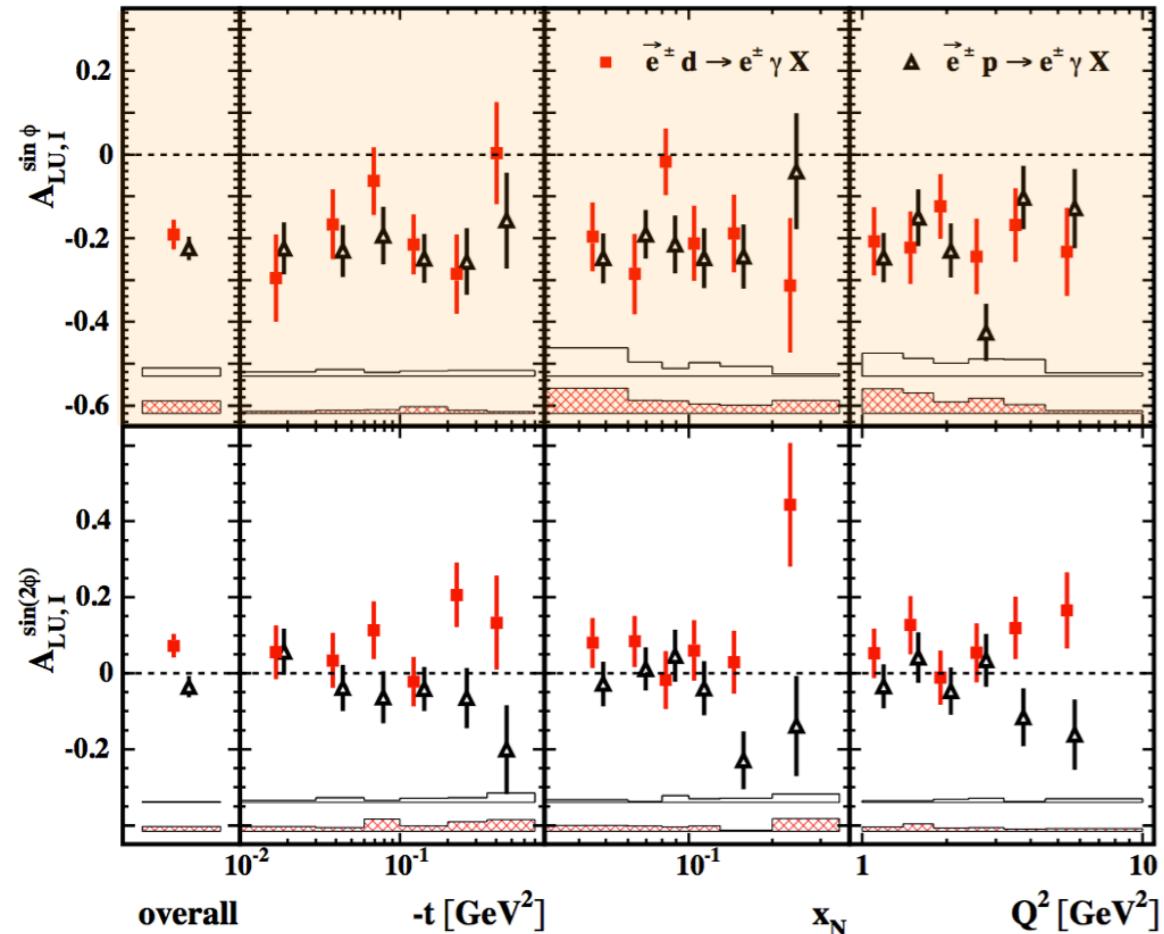
Caveat! Relatively
large BH
contribution to
these asymmetries!

<http://arxiv.org/abs/1004.0177>

A. Airapetian *et al*, JHEP 06 (2010) 019

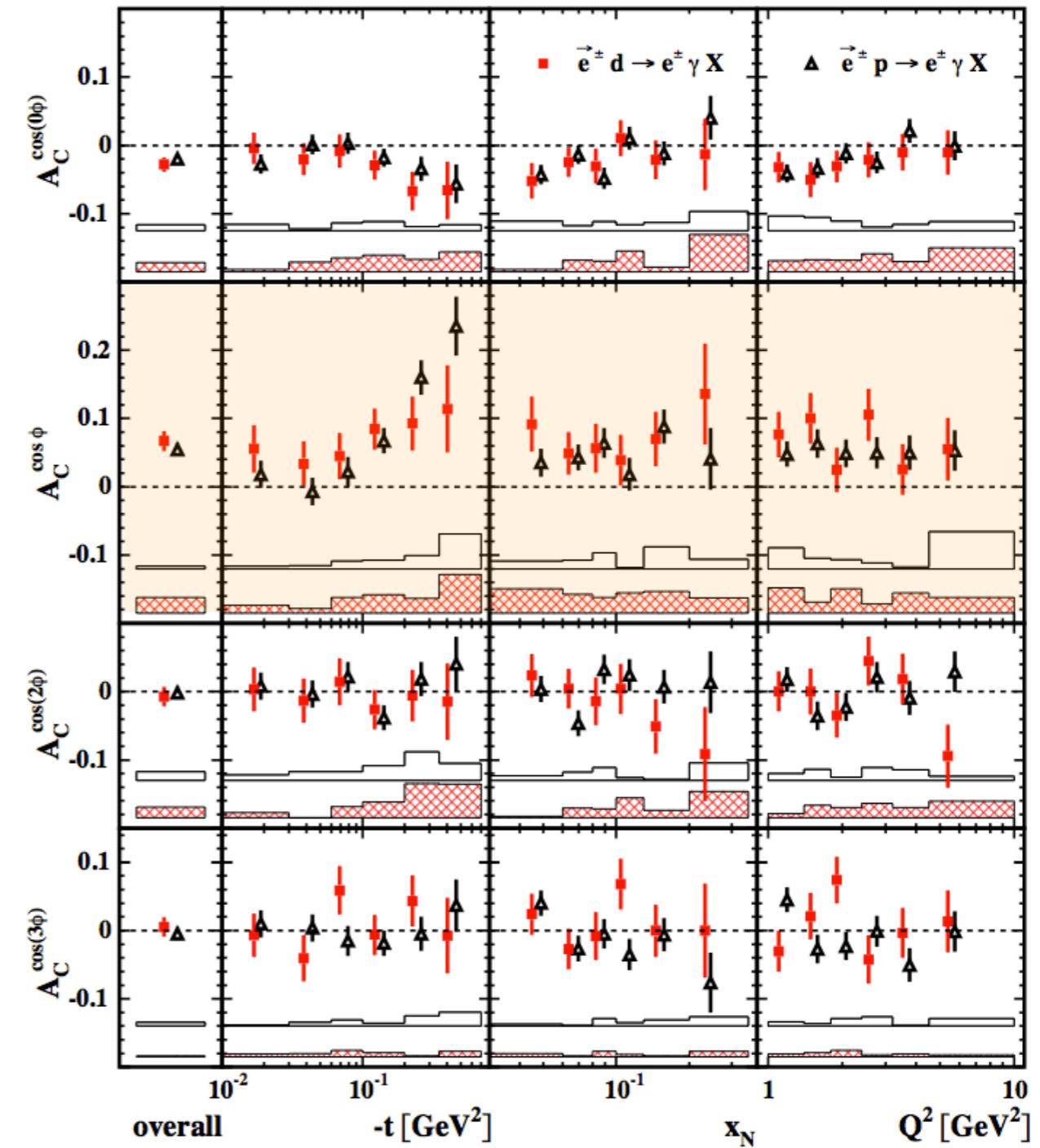
Beam Asymmetries

A. Airapetian *et al*, Nucl. Phys. B 829 (2010) 1-27

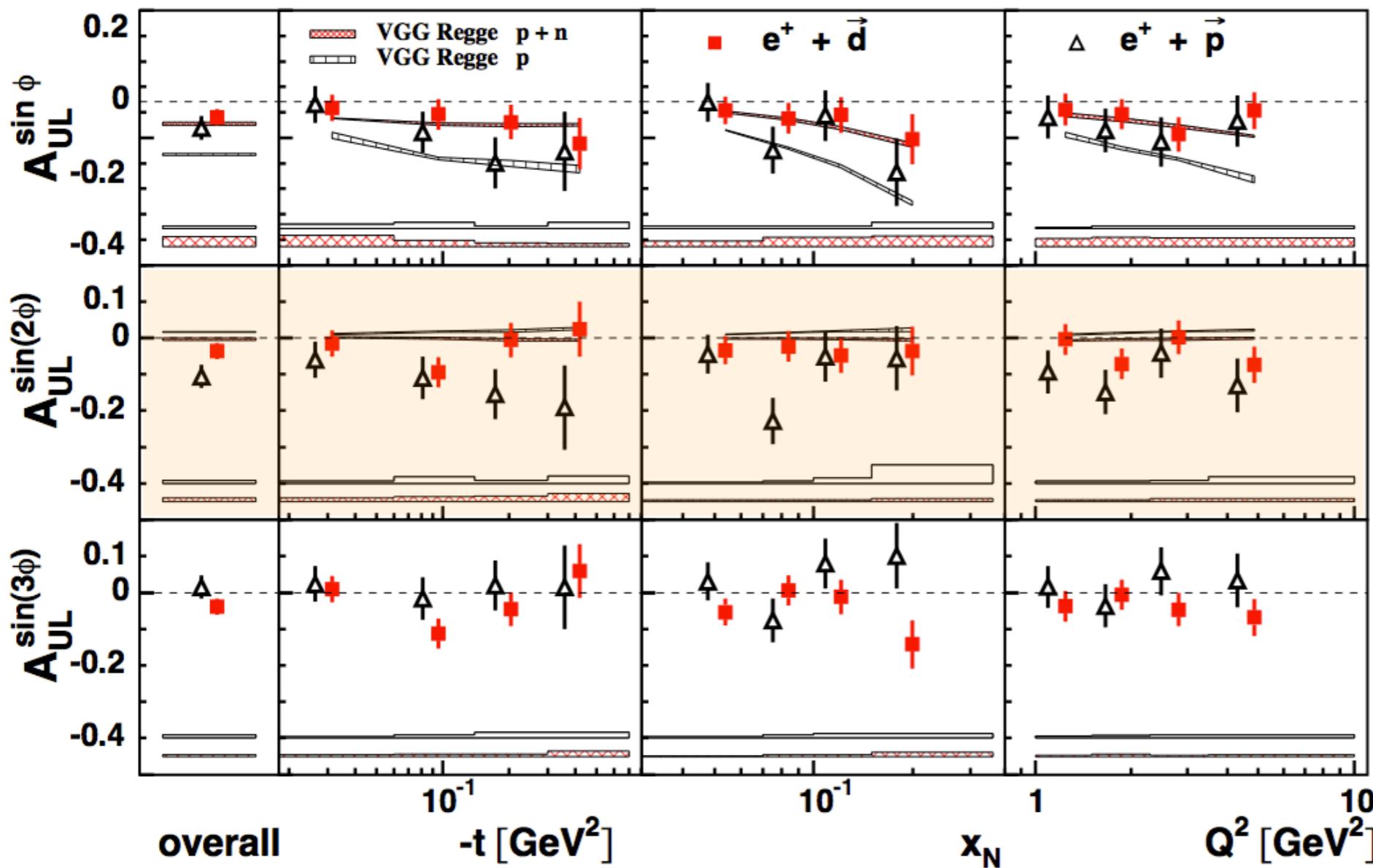


<http://www.arxiv.org/abs/0911.0095>

Deuterium is governed by different GPDs - but the asymmetry data is not so different!



Target Asymmetries



No good idea
how to model
long. pol.
deuterium
GPDs. Currently
use a proton/
neutron hybrid

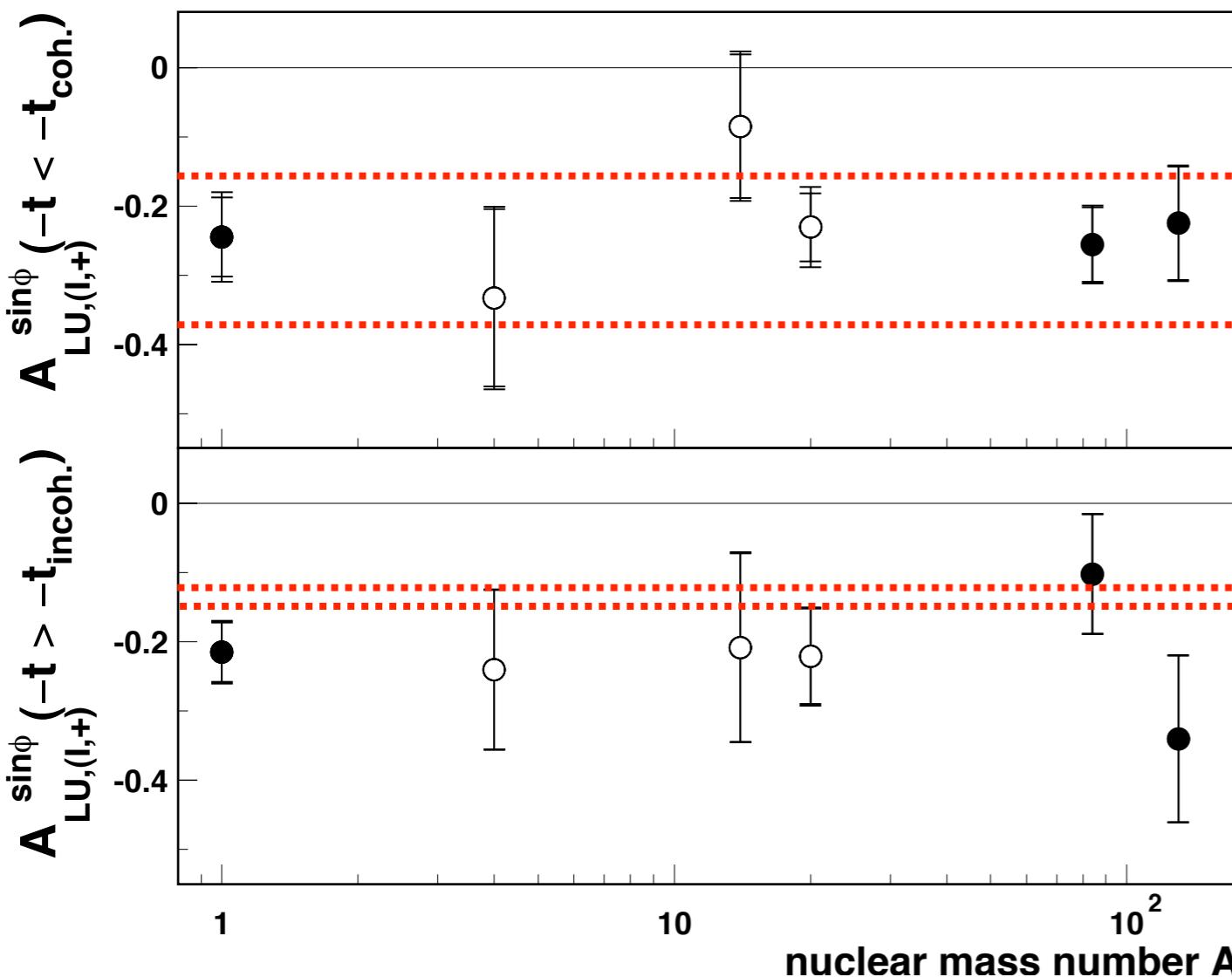
<http://www.arxiv.org/abs/1008.3996>

Will appear in Nucl. Phys. B

Nuclear Mass Dependence

<http://arxiv.org/abs/0911.0091>

A. Airapetian et al. *Phys. Rev. C* 81, 035202 (2010)



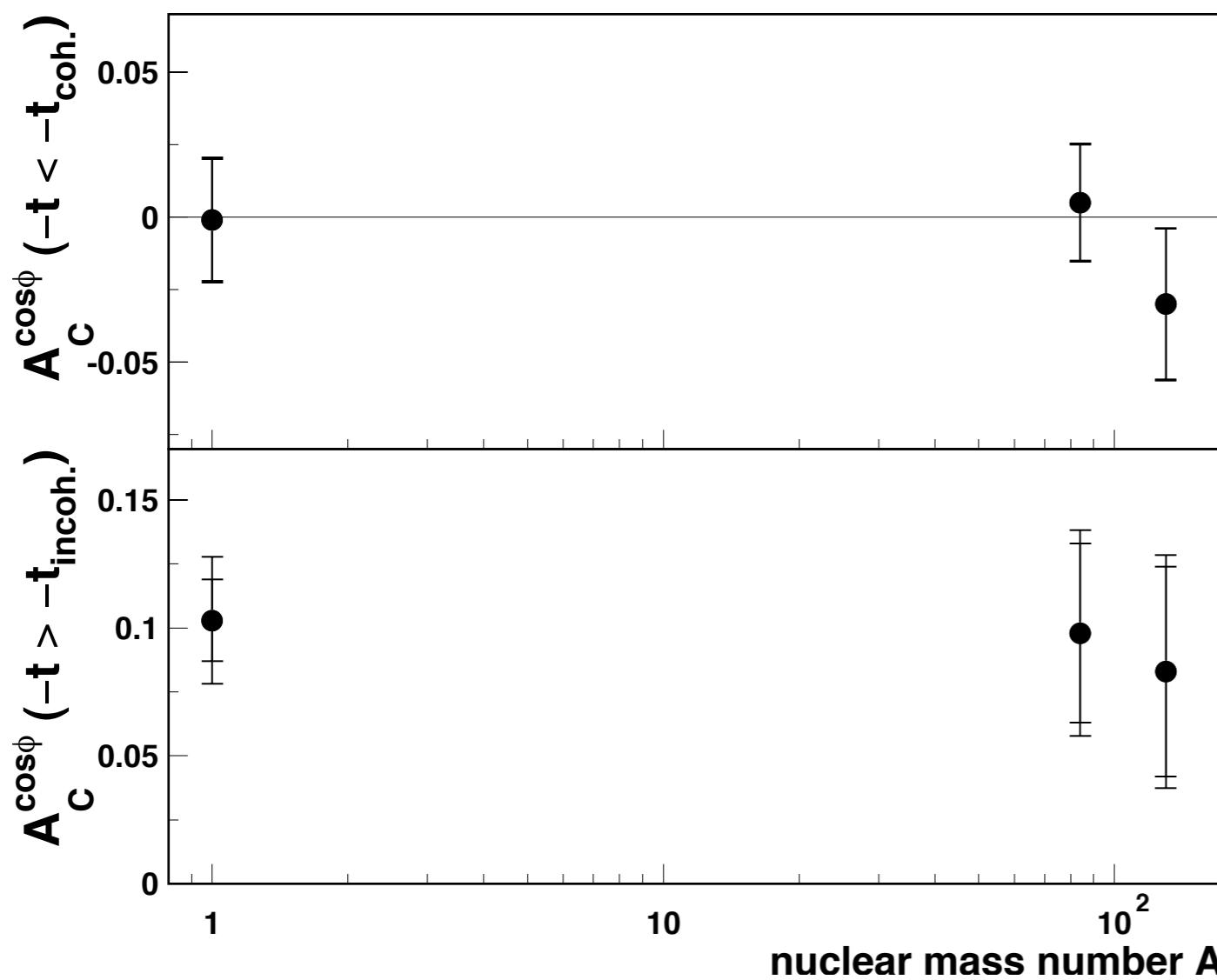
- Interference Term
- Interference + $|DVCS|^2$

The data shows
no significant difference
between coherent and
incoherent DVCS
processes

Nuclear Mass Dependence

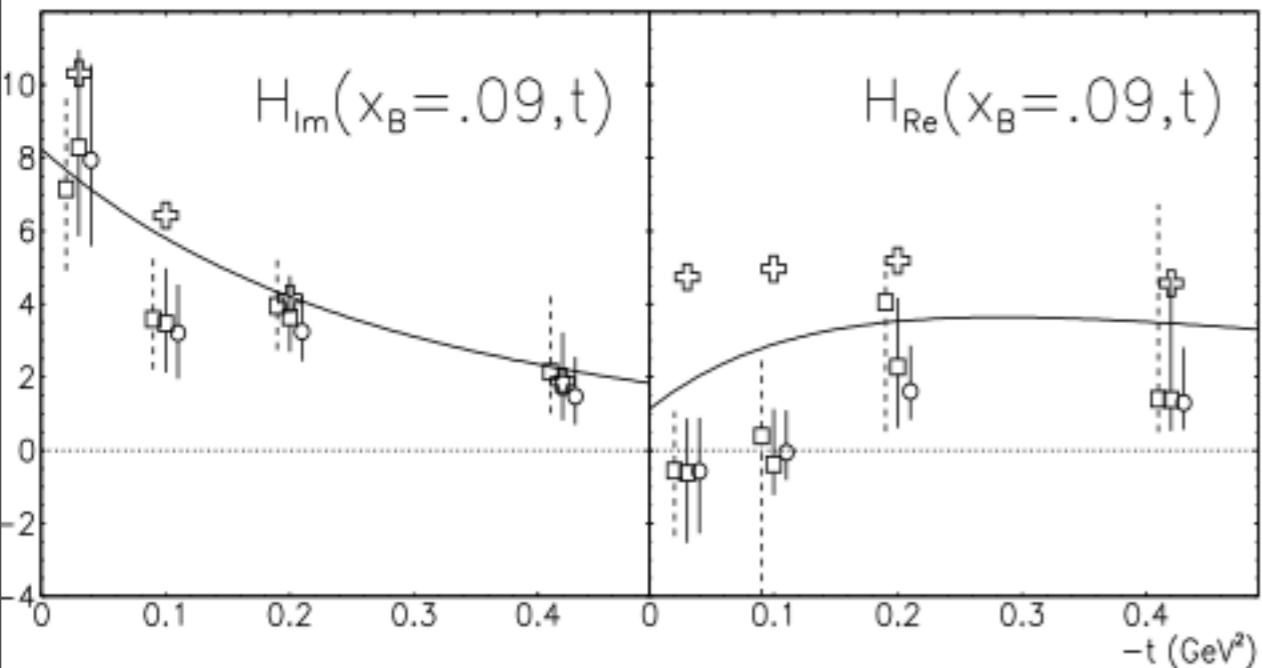
<http://arxiv.org/abs/0911.0091>

A. Airapetian *et al.* Phys. Rev. C 81, 035202 (2010)



Beam Charge Asymmetries are essentially 0 for coherent scattering
Consistent with the Proton asymmetry for heavier targets

GPD Discovery



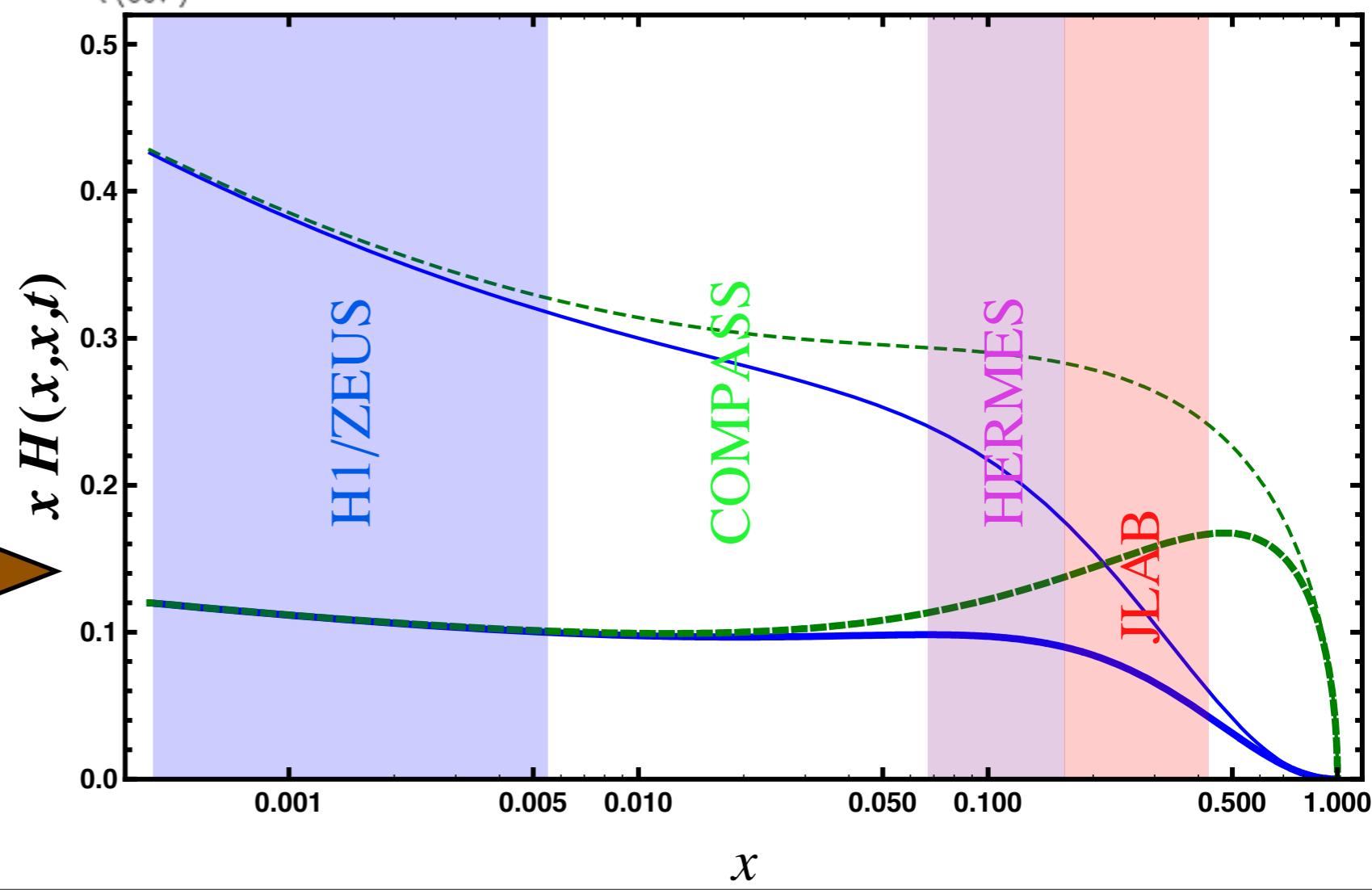
<http://arxiv.org/abs/1005.4922>

M. Guidal

New CFF Fit Result
incorporating AUL moments



Postulate GPDs from
first principle models



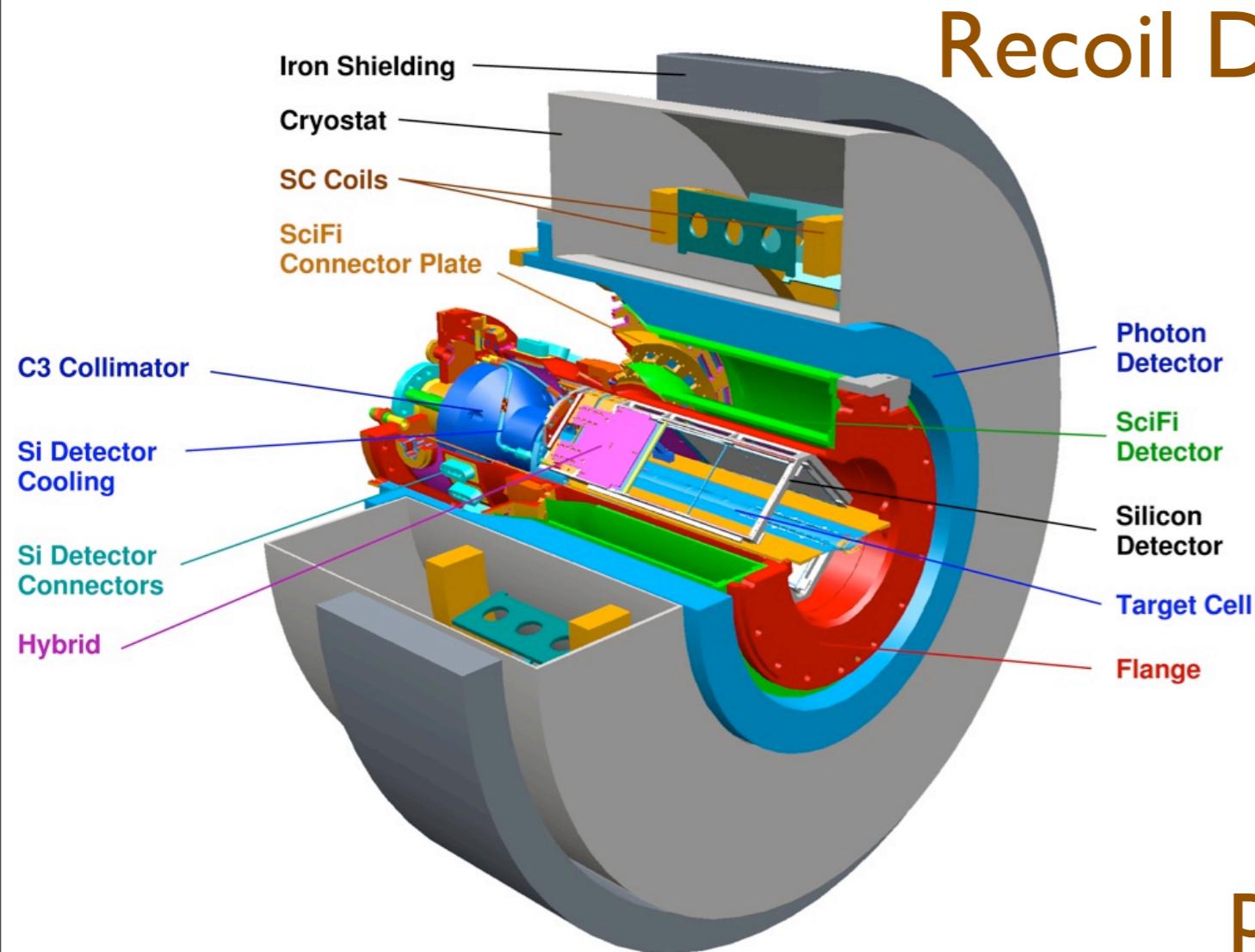
<http://arxiv.org/abs/0904.0458>

Kumerički and Müller

Nucl. Phys. B841 1 (2010)

x

Future Data



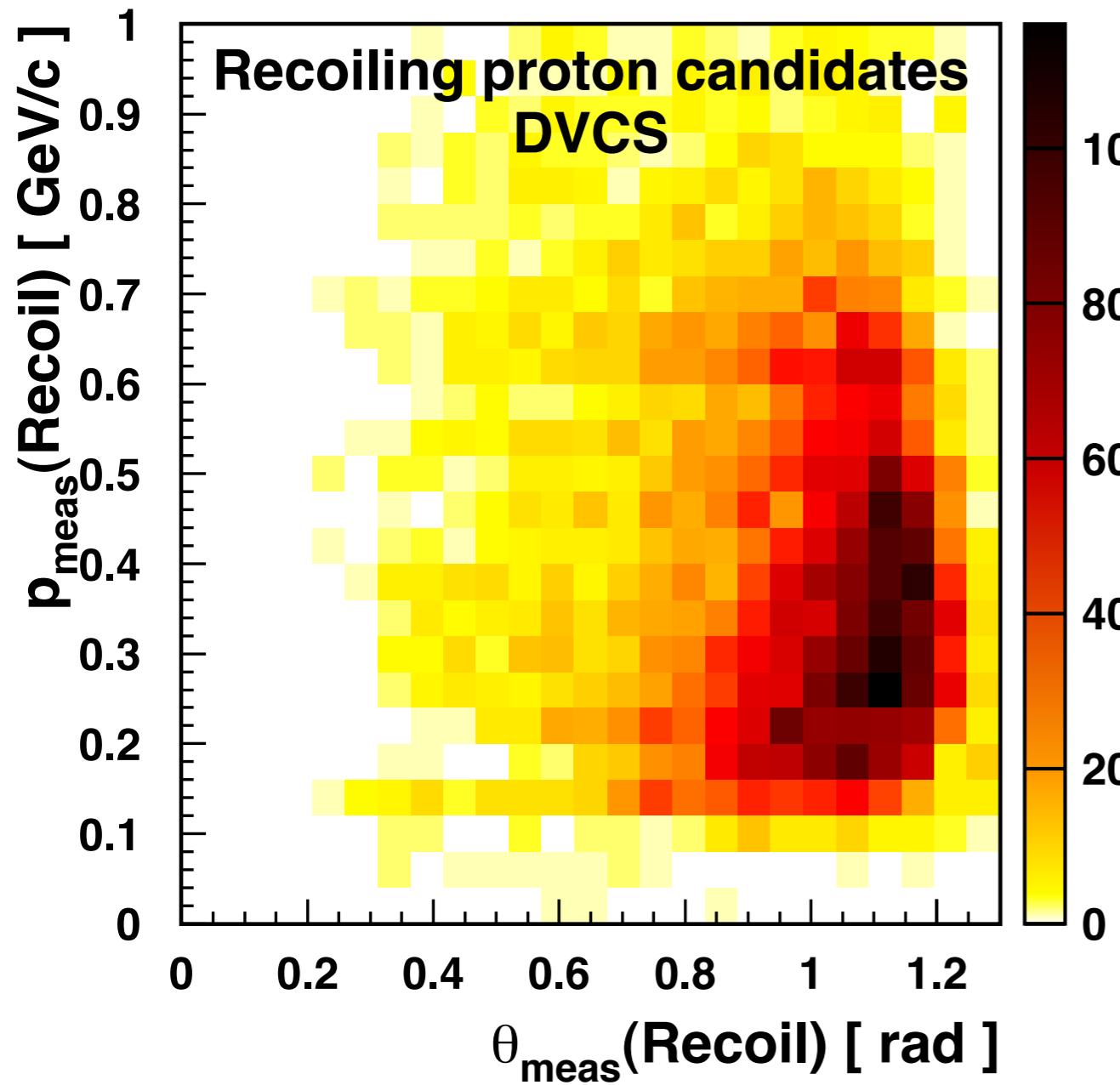
Recoil Detector - True Exclusivity!

Almost pure BH/DVCS
signal - background <1%

Potential for measuring
other exclusive physics
processes - Δ -DVCS, HEMP
physics, DVCS on neutron...

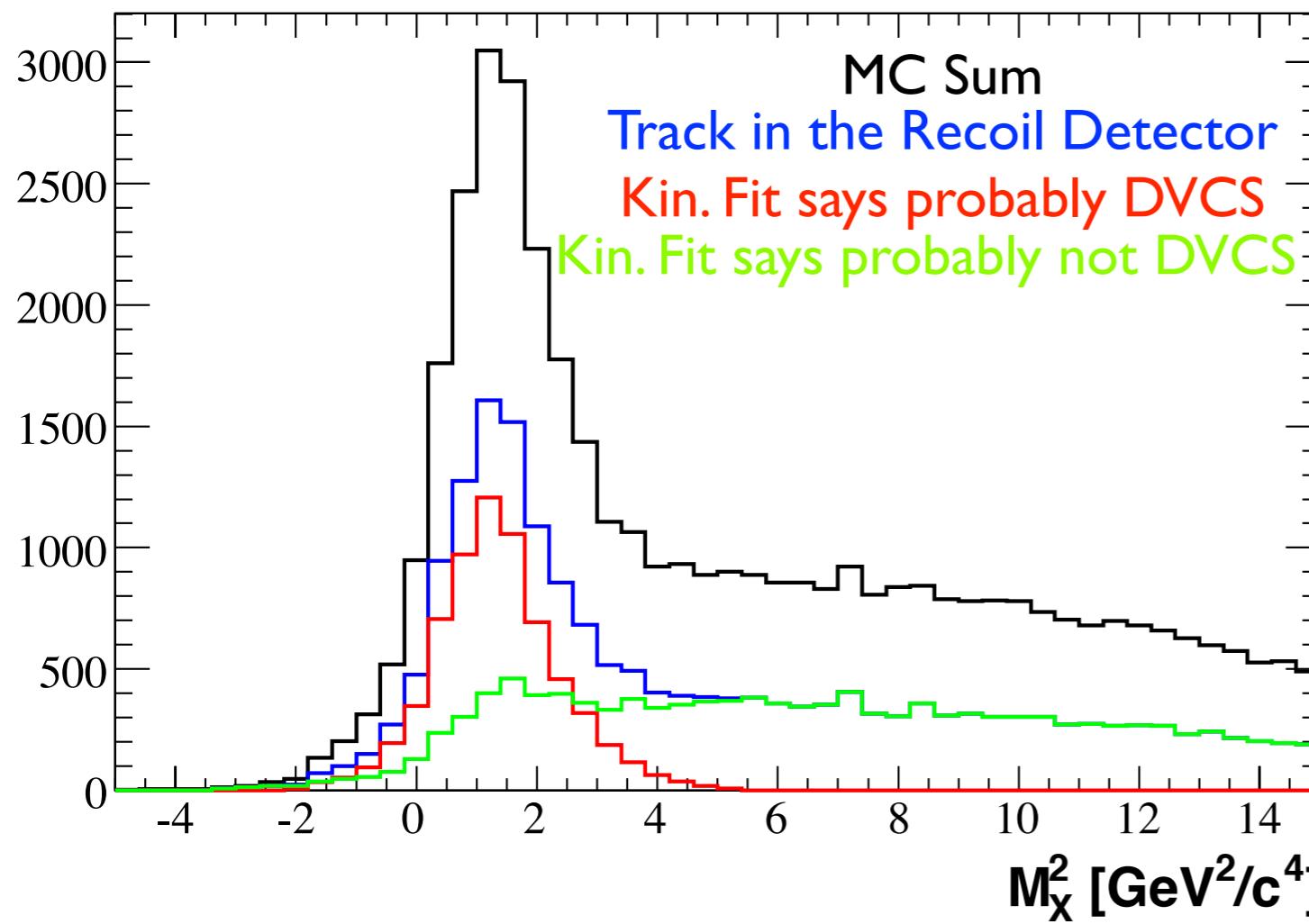
Future DVCS Data

Hermes 2007 data



Calibrations finalised
Kinematic Fitting routines implemented
Event selection method finalised - asymmetry extraction techniques to be finalised and systematic studies still to be completed.

Future DVCS Data



Significant improvement
in the purity of signal

No hard-and-fast idea on
the feasibility of a
 Δ -DVCS measurement

Expect first publicly-
available results in

Q2 2011

Conclusions

- DVCS can be used to access information on **Generalised Parton Distributions**
- That information can tell us **unique** things about **nucleon structure**
- HERMES has the most **diverse** DVCS measurements of any experiment.

Conclusions

- There is still **no clearly correct idea** about how the **nuclear medium** modifies GPD-dependent behaviour. More data?
- Already, **GPDs can be constrained** - but there is much left to do!
- Experimental extractions of GPDs and related **physical quantities** can work in harmony with **Lattice QCD**.