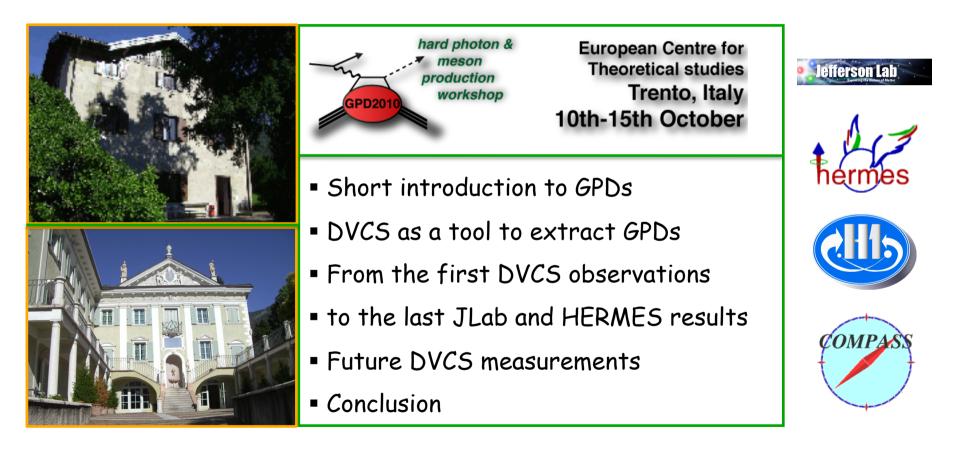
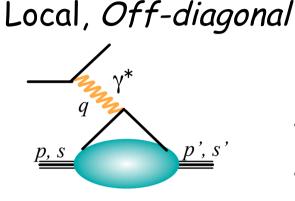
Deeply Virtual Compton Scattering experiments

Franck Sabatié
CEA Saclay

Irfu CCCC saclay



The electromagnetic probe for the nucleon structure

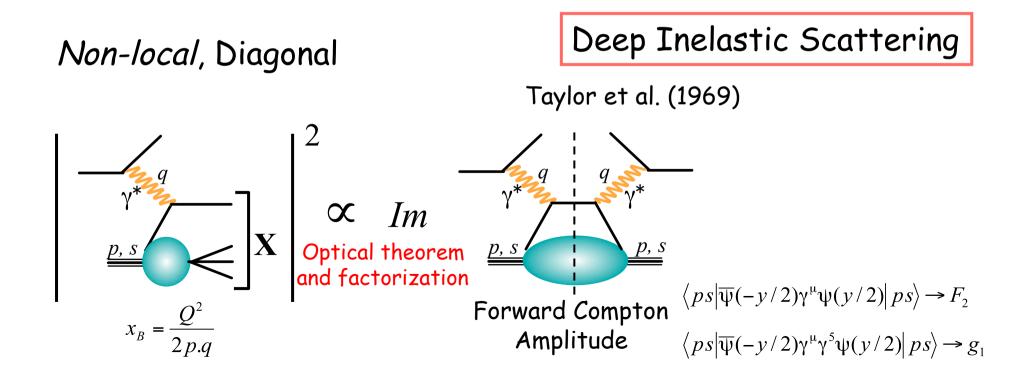


Elastic scattering

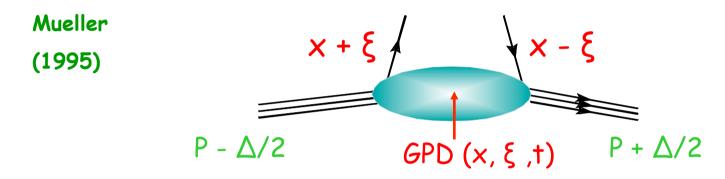
Hofstadter (1958)

 $\langle p's' | \overline{\psi}(0) \gamma^{\mu} \psi(0) | ps \rangle \rightarrow G_M, G_E$

 $\langle p's' | \overline{\psi}(0) \gamma^{\mu} \gamma^{5} \psi(0) | ps \rangle \rightarrow G_A, G_P$



A natural extension: non-local, off-diagonal matrix elements

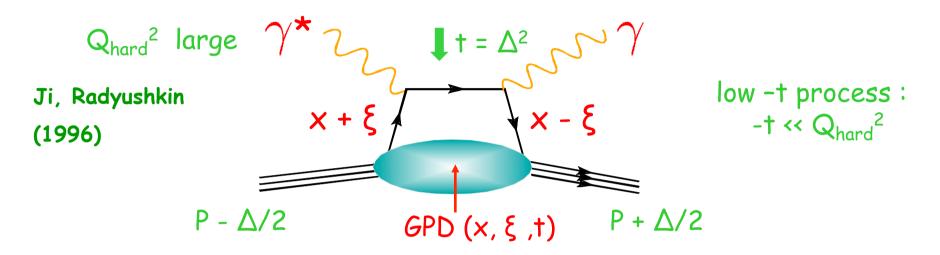


 $(x + \xi)$ and $(x - \xi)$: longitudinal momentum fractions of quarks

The structure of the nucleon can be described by 4 Generalized Parton Distributions :

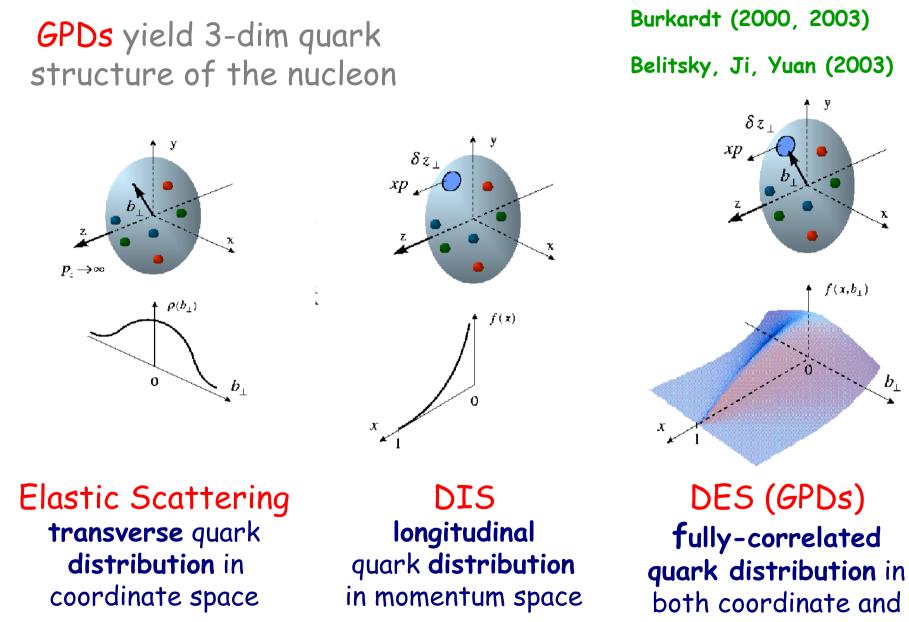
Vector : $H(x, \xi, t)$ Axial-Vector : $H(x, \xi, t)$ Tensor : $E(x, \xi, t)$ Pseudoscalar : $\tilde{E}(x, \xi, t)$

And its golden process : Deeply Virtual Compton Scattering



 $(x + \xi)$ and $(x - \xi)$: longitudinal momentum fractions of quarks

at large Q^2 : QCD factorization theorem \Rightarrow hard exclusive processes can be described by 4 Generalized Parton Distributions: Vector : $H(x, \xi, t)$ Axial-Vector : $\widetilde{H}(x, \xi, t)$ Tensor : $E(x, \xi, t)$ Pseudoscalar : $\widetilde{E}(x, \xi, t)$ Why Generalized Parton Distributions are the way to go!



momentum space

They contain what we know already through sum rules and kinematical limits: Form Factors, parton distributions

> forward limit : ordinary parton distributions

 $H^q(x, \xi = 0, t = 0) = q(x)$ unpolarized quark distributions $\tilde{H}^q(x, \xi = 0, t = 0) = \Delta q(x)$ polarized quark distributions E^q, \tilde{E}^q : do NOT appear in DIS ... new information

> first moments : nucleon electroweak form factors

$$P - \Delta/2 \qquad P + \Delta/2 \qquad \int_{-1}^{1} dx \, H^{q}(x,\xi,t) = F_{1}^{q}(t) \quad \text{Dirac}$$

$$\int_{-1}^{1} dx \, E^{q}(x,\xi,t) = F_{2}^{q}(t) \quad \text{Pauli}$$

$$\int_{-1}^{1} dx \, \tilde{H}^{q}(x,\xi,t) = G_{A}^{q}(t) \quad \text{axial}$$

$$\int_{-1}^{1} dx \, \tilde{E}^{q}(x,\xi,t) = G_{P}^{q}(t) \quad \text{pseudo-scalar}$$

They contain what we know already through sum rules and kinematical limits: Form Factors, parton distributions

Through the space-momentum correlation, they give access to the Orbital Angular Momentum (OAM) carried by partons inside the nucleon : Finally an end to the spin crisis ?

Ji's sum rule :

$$J^{q} = \frac{1}{2}\Delta\Sigma + L^{q} = \lim_{t \to 0} \frac{1}{2} \int_{-1}^{1} x dx \left[\frac{H^{q}(x,\xi,t)}{1} + \frac{E^{q}(x,\xi,t)}{1} \right]$$

Related to momentum fraction carried by quarks

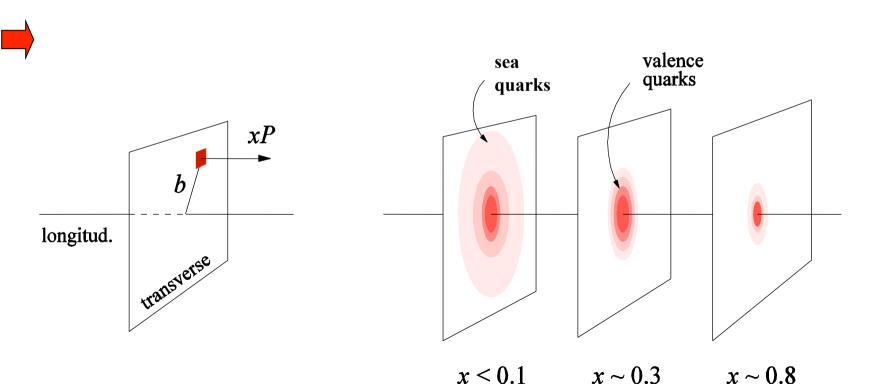
Moments of GPDs are calculable in Lattice QCD. Lowest moments have already been computed for valence quarks (Goeckeler) Best accessible using transverse polarized target, but also neutron DVCS

Properties, applications

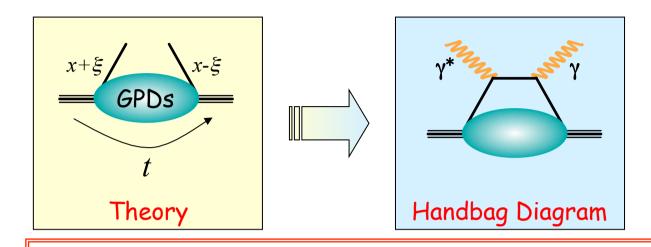
They contain what we know already through sum rules and kinematical limits: Form Factors, parton distributions

Through the space-momentum correlation they MivBurkarskt Mthiehl (2002) Ang $H^{q}(x, \mathbf{b}_{\perp}) = \int \frac{d^{2} \Delta_{\perp}}{(2\pi)^{2}} e^{i\mathbf{b}_{\perp} \cdot \Delta_{\perp}} H^{q}(x, \xi = 0, -\Delta_{\perp}^{2})$

e



GPDs from Theory to Experiment



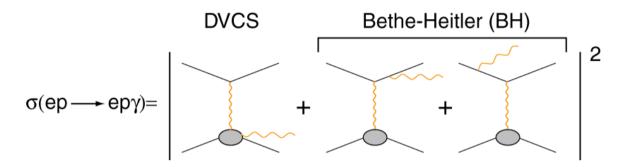
2. The GPDs enter the DVCS amplitude as an integral over x:

- GPDs appear in the real part through a PP integral over x

$$T^{DVCS} = \int_{-1}^{+1} \frac{GPD(x,\xi,t)}{x-\xi+i\varepsilon} dx + \cdots$$
$$= P \int_{-1}^{+1} \frac{GPD(x,\xi,t)}{x-\xi} dx - i\pi GPD(x=\xi,\xi,t) + \cdots$$

Experimental observables linked to GPDs

Experimentally, DVCS is undistinguishable with Bethe-Heitler



However, we know FF at low t and BH is fully calculable

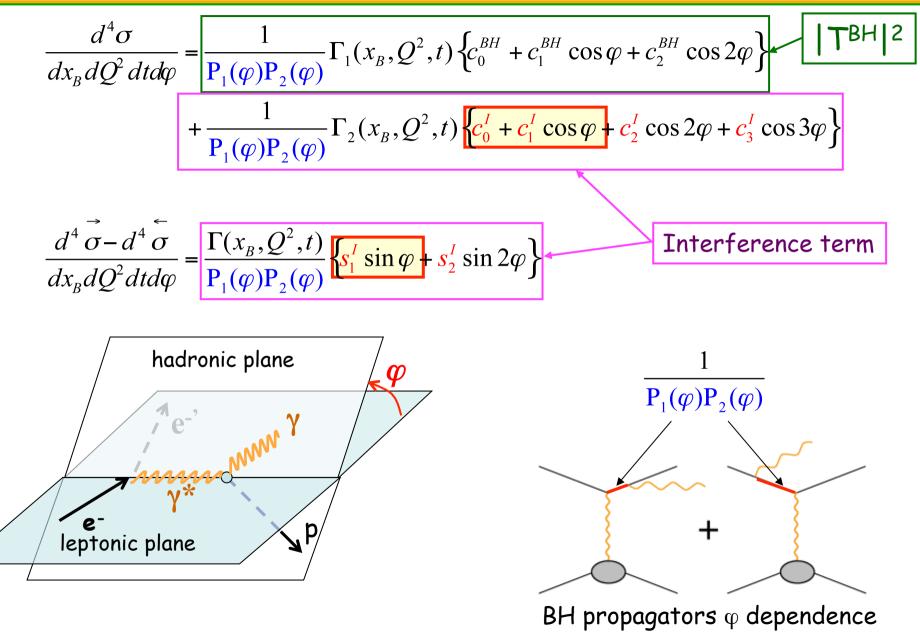
Using a polarized beam on an unpolarized target, 2 observables can be measured:

$$\frac{d^4\sigma}{dx_B dQ^2 dt d\varphi} \approx \left| T^{BH} \right|^2 + 2T^{BH} \cdot \operatorname{Re}\left(T^{DVCS} \right) + \left| T^{DVCS} \right|^2$$

$$\frac{d^{4}\vec{\sigma} - d^{4}\vec{\sigma}}{dx_{B}dQ^{2}dtd\varphi} \approx 2T^{BH} \cdot \operatorname{Im}\left(T^{DVCS}\right) + \left[\left|T^{DVCS}\right|^{2} - \left|T^{DVCS}\right|^{2}\right]$$
At low energy,
$$|T^{DVCS}|^{2} \text{ supposed small}$$

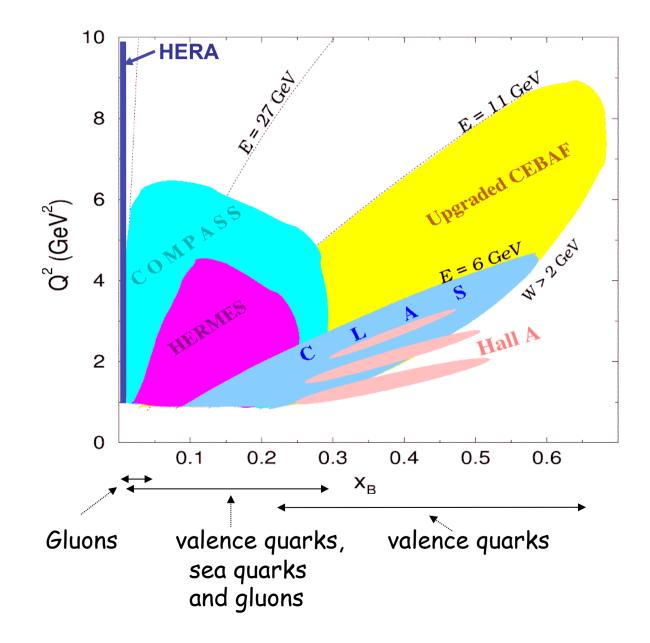
Ji, Kroll, Guichon, Diehl, Pire, ...

Into the harmonic structure of DVCS

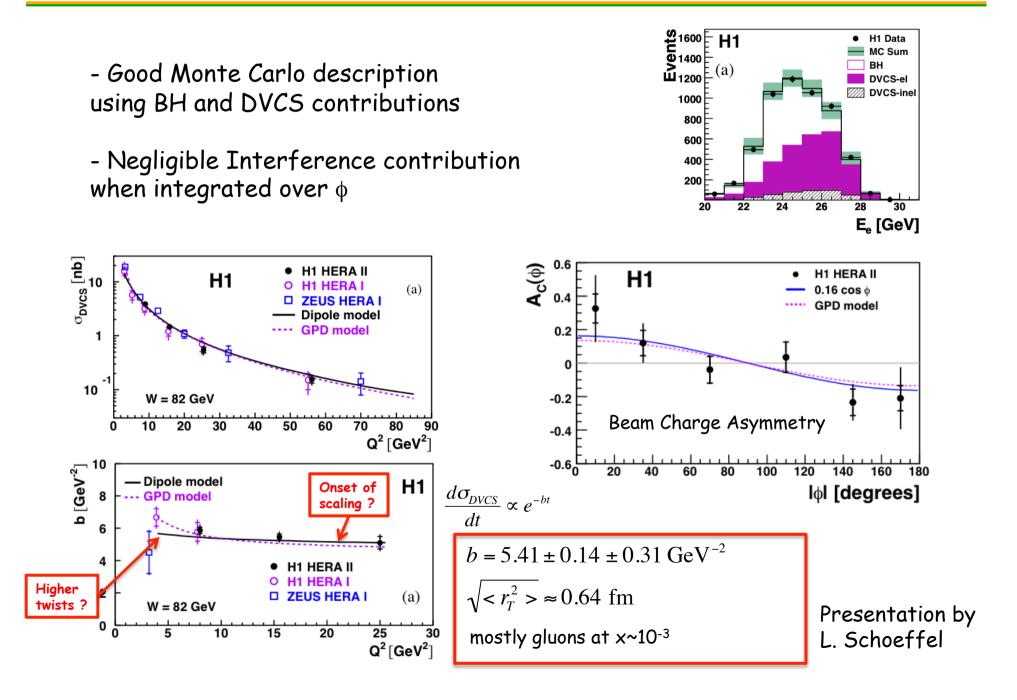


Belitsky, Mueller, Kirchner

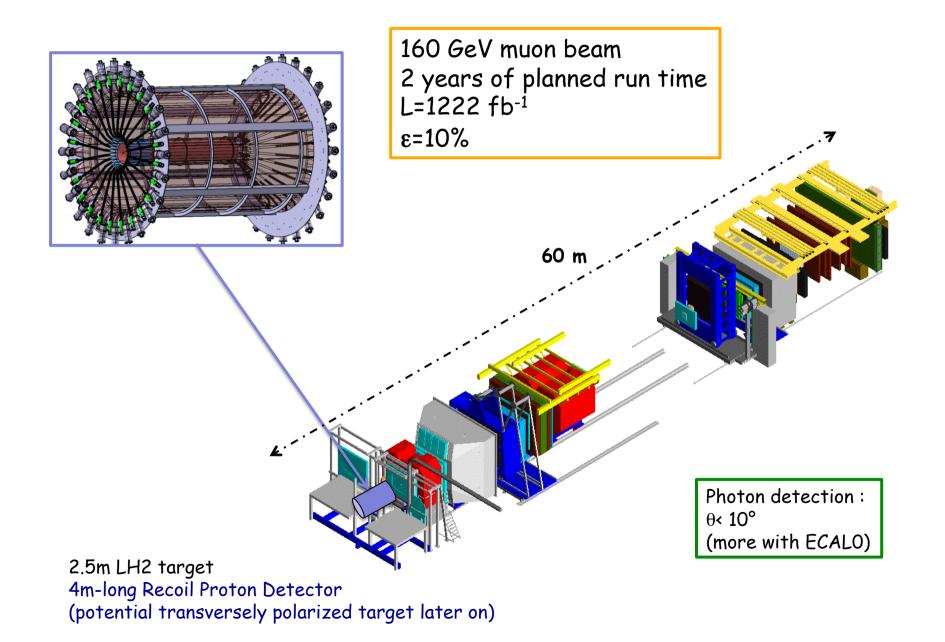
DVCS experiments and proposals worldwide



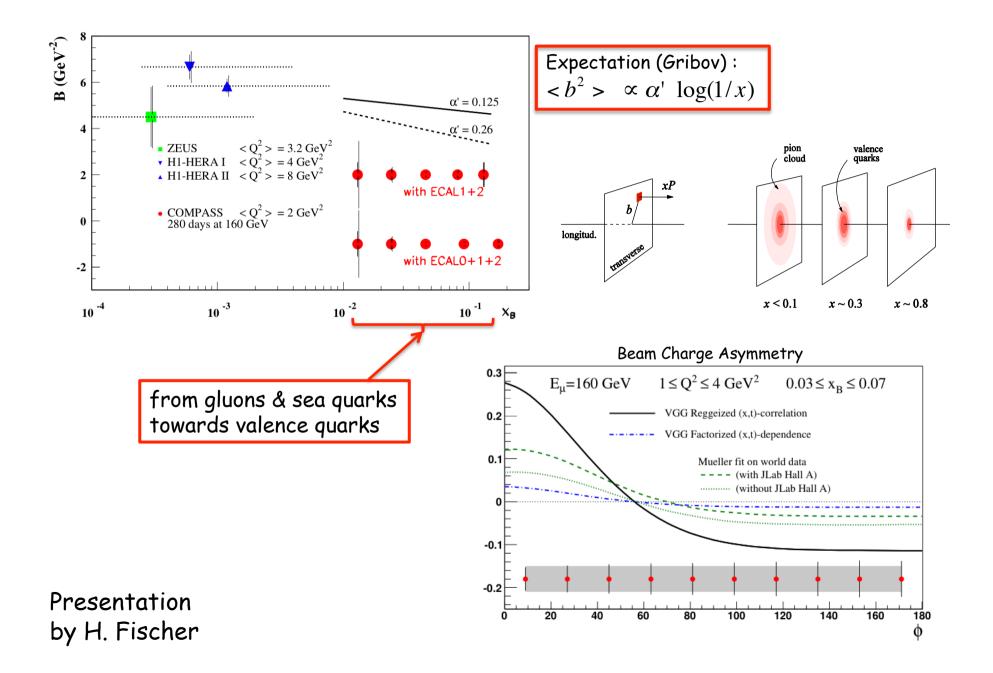
HERA data

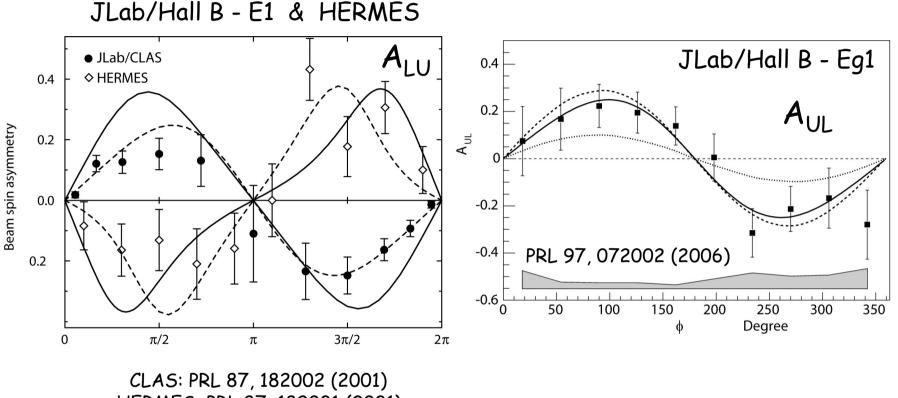


COMPASS DVCS setup



COMPASS DVCS proposal



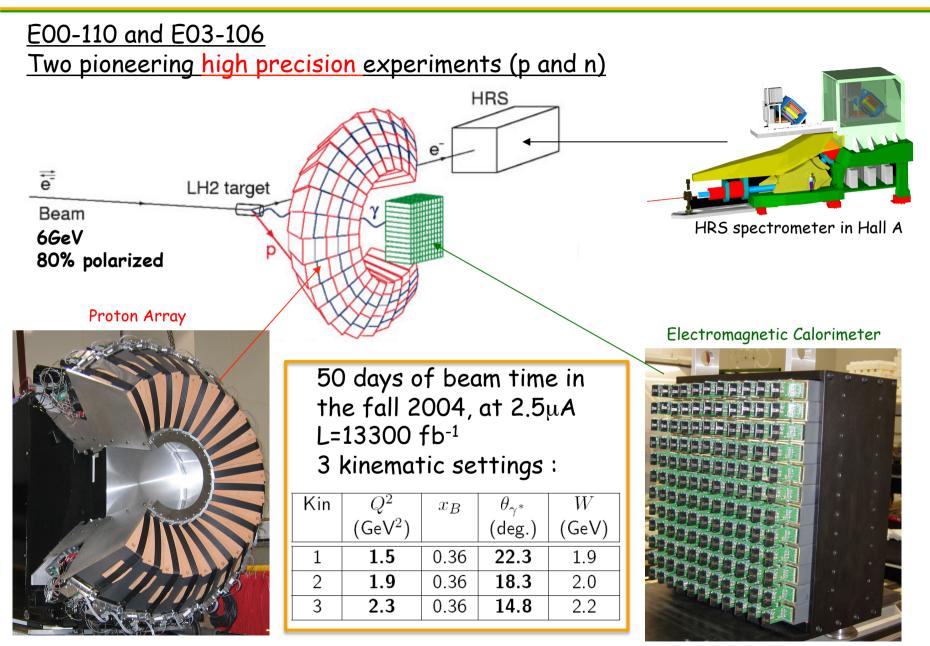


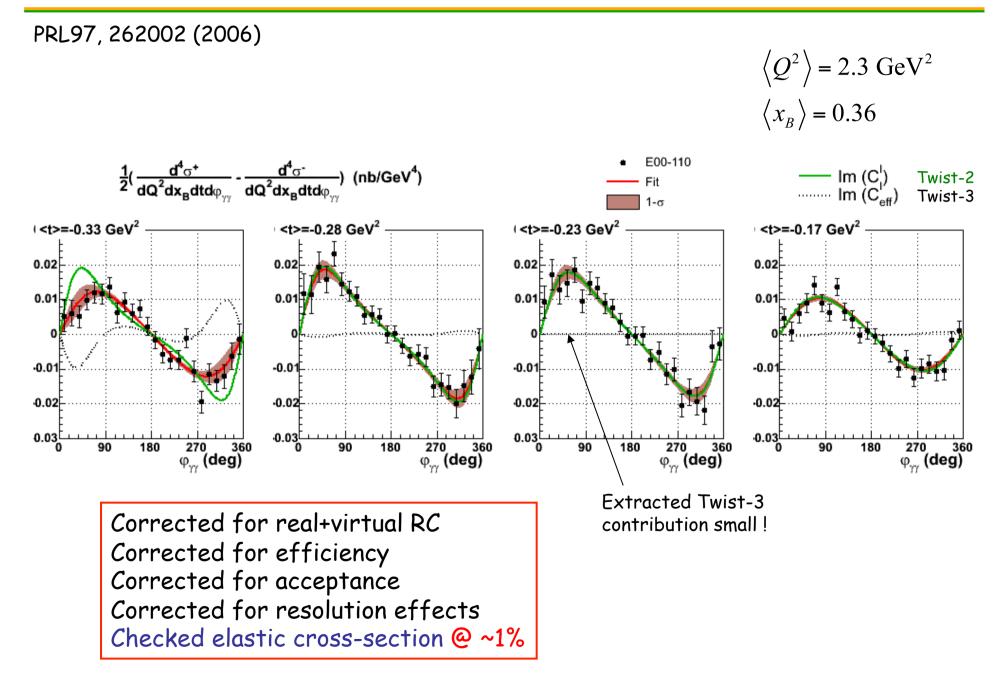
HERMES: PRL 87, 182001 (2001)

Both results show, with a limited statistics, a sin ϕ behavior (necessary condition for handbag dominance)

In the A_{LU} result, DD models (VGG) tend to over-estimate the data

DVCS experiments at 6 GeV in the Jefferson Lab Hall A





PRL97, 262002 (2006) <-t>=0.26 GeV², <x_B>=0.36 $C^{I}(F) = F_{1}\mathbf{H} + \frac{x_{B}}{2 - x_{B}}(F_{1} + F_{2})\tilde{\mathbf{H}} - \frac{t}{4M^{2}}F_{2}\mathbf{E}$ 3.5 3 2.5 2 1.5 Im C^I ♦ Im C^I_{eff} Supposedly mostly H (~80%) 0.5

1.6

0

1.4

No Q² dependence using BMK separation: strong indication for scaling behavior and handbag dominance

2

1.8

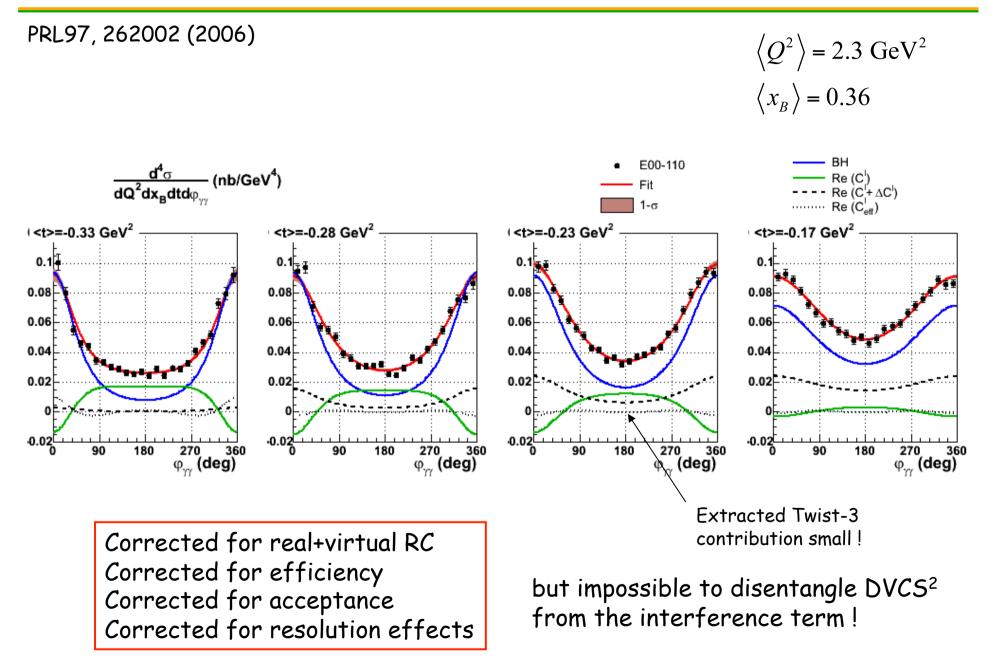
2.2

2.4 Q² (GeV²) Twist-2

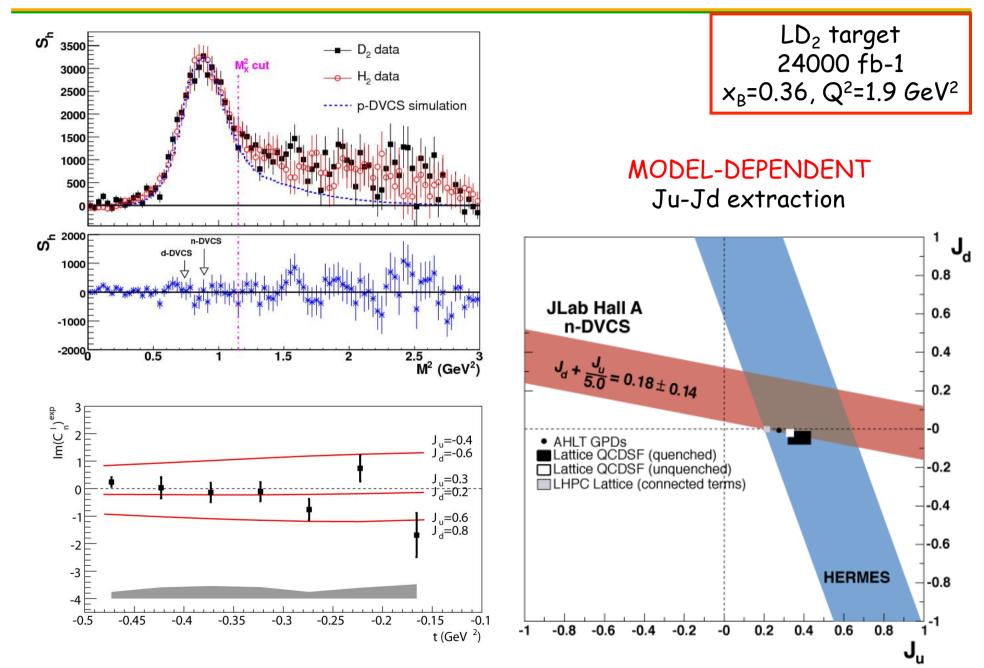
Twist-3

Twist 4+ contributions are smaller than 10%

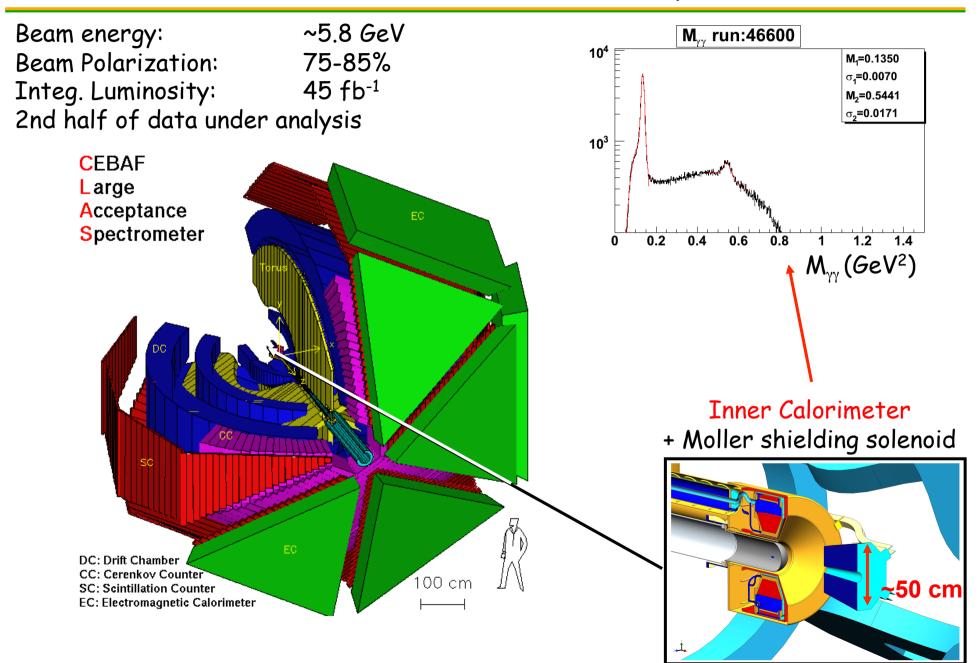
Total cross-section



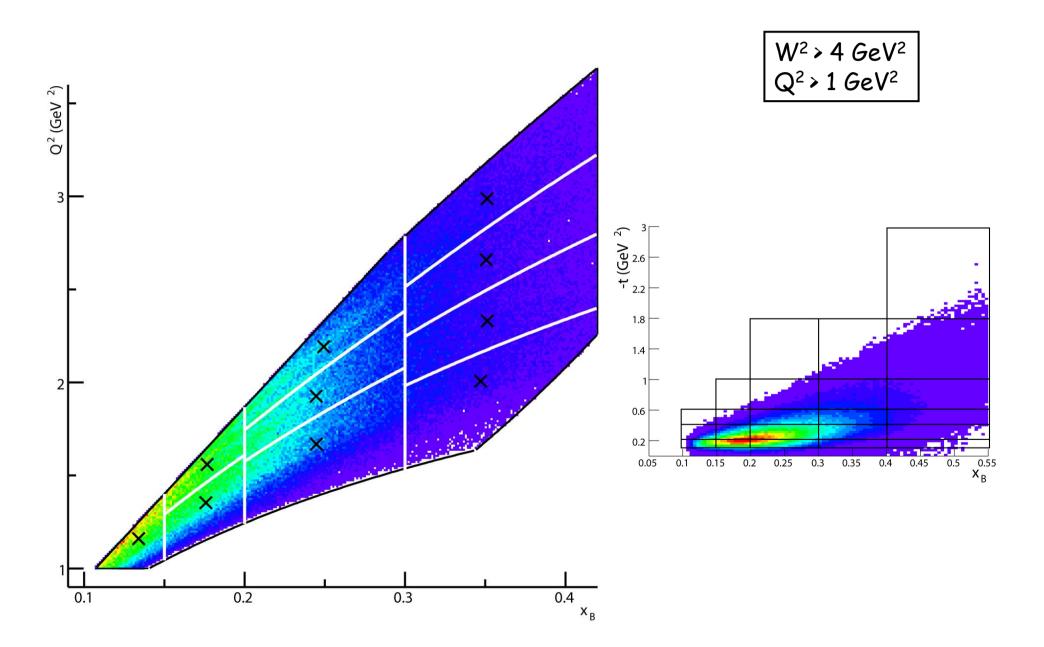
DVCS on the neutron in JLab/Hall A: E03-106



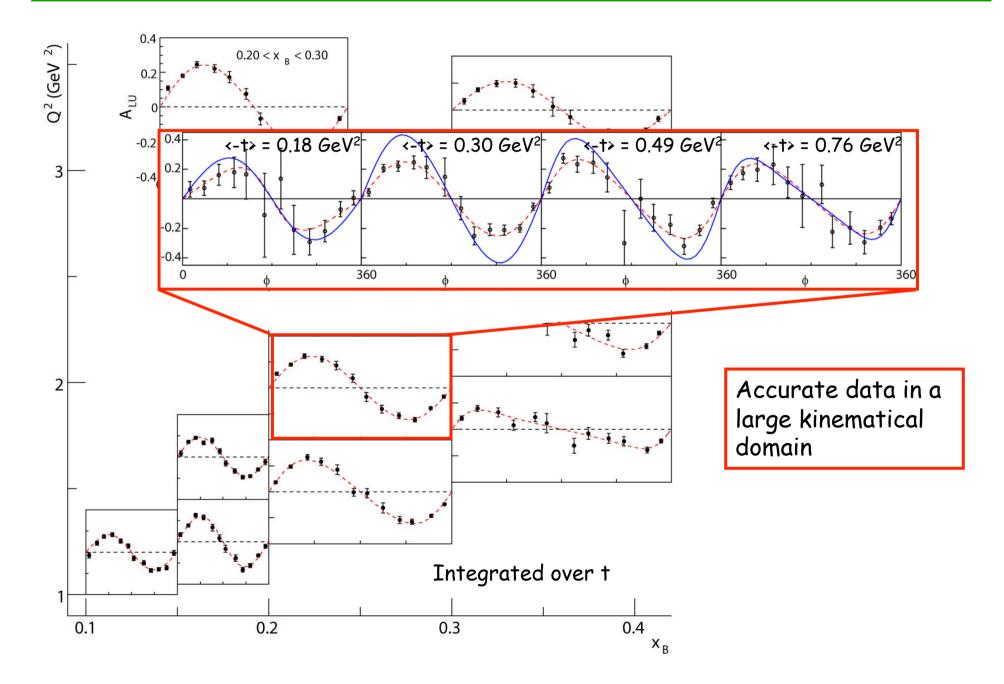
E1-DVCS with CLAS : a dedicated DVCS experiment in Hall B



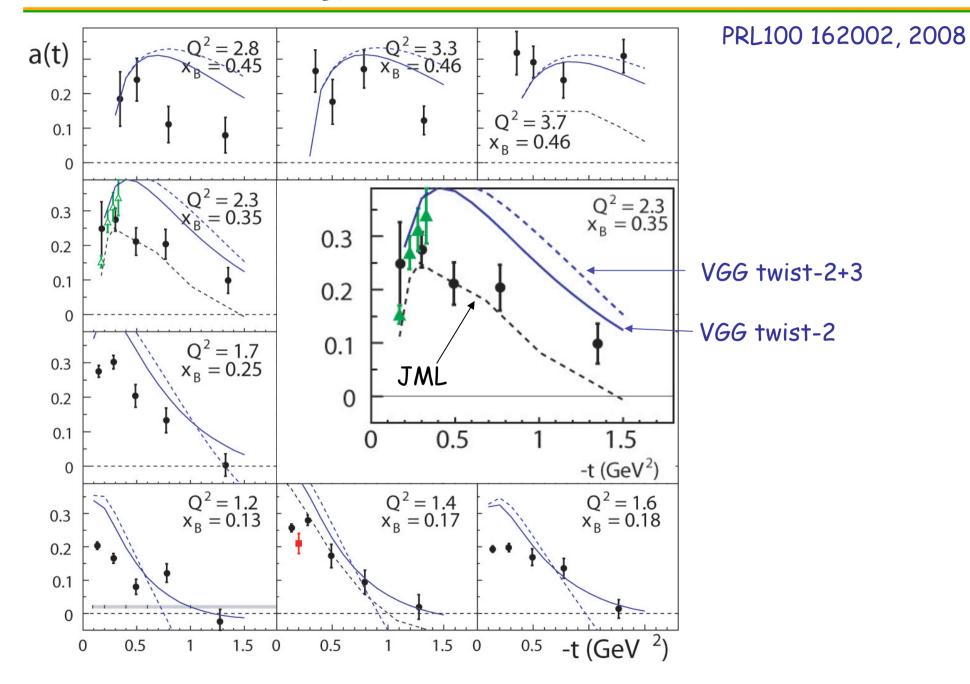
E1-DVCS kinematical coverage and binning



E1-DVCS : Asymmetry as a function of x_B and Q^2

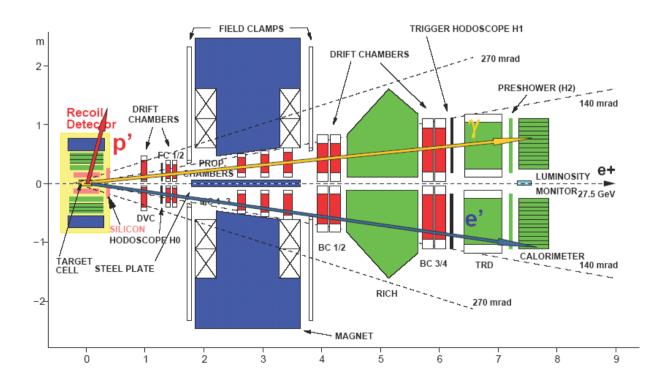


E1-DVCS : $A_{LU}(90^{\circ})$ as a function of -t + models



HERMES

HERMES has a very broad coverage of observables with:
Beam Charge (e⁺ and e⁻ beams)
Beam spin direction (both beam helicities)
Target spin direction (parallel, transverse, unpolarized)
Different targets (H, D, He, N, Ne, Kr, Xe)
Recoil and spectator proton detection (results soon ?)



Interesting aspect : reversing both beam charge and spin

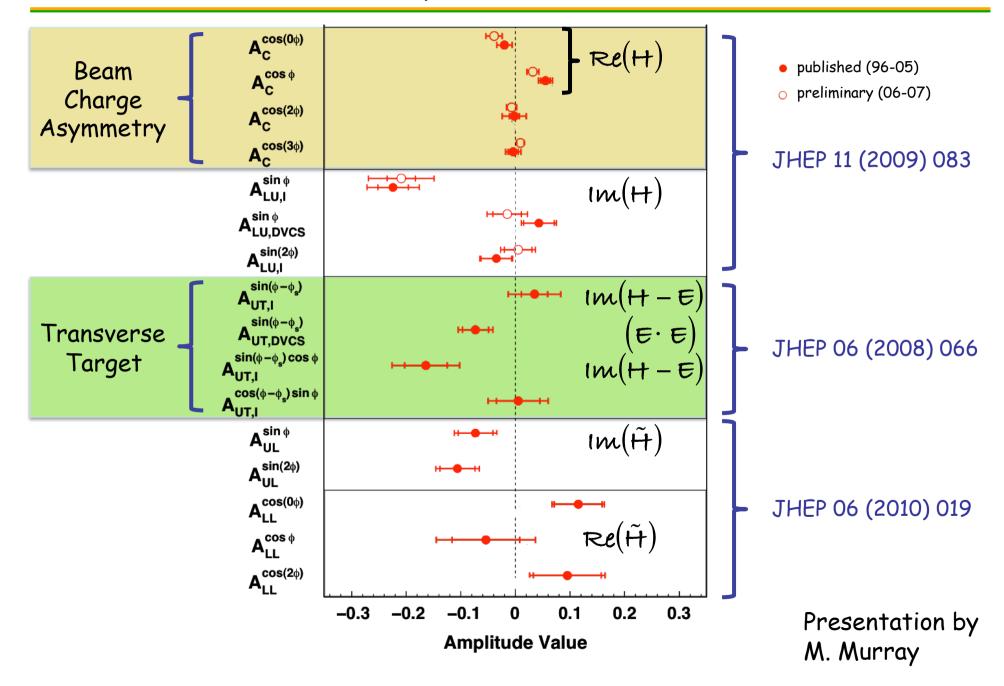
Asymmetry of
interference term
$$\mathcal{A}_{\mathrm{LU}}^{\mathrm{I}}(\phi) \equiv \frac{(d\sigma^{+\rightarrow} - d\sigma^{+\leftarrow}) \bigcirc (d\sigma^{-\rightarrow} - d\sigma^{-\leftarrow})}{(d\sigma^{+\rightarrow} + d\sigma^{+\leftarrow}) + (d\sigma^{-\rightarrow} + d\sigma^{-\leftarrow})}$$

$$= \frac{-\frac{K_{\mathrm{I}}}{\mathcal{P}_{1}(\phi)\mathcal{P}_{2}(\phi)} \left[\sum_{n=1}^{2} s_{n}^{\mathrm{I}} \sin(n\phi)\right]}{\frac{K_{\mathrm{BH}}}{\mathcal{P}_{1}(\phi)\mathcal{P}_{2}(\phi)} \sum_{n=0}^{2} c_{n}^{\mathrm{BH}} \cos(n\phi) + \frac{1}{Q^{2}} \sum_{n=0}^{2} c_{n}^{\mathrm{DVCS}} \cos(n\phi)}$$

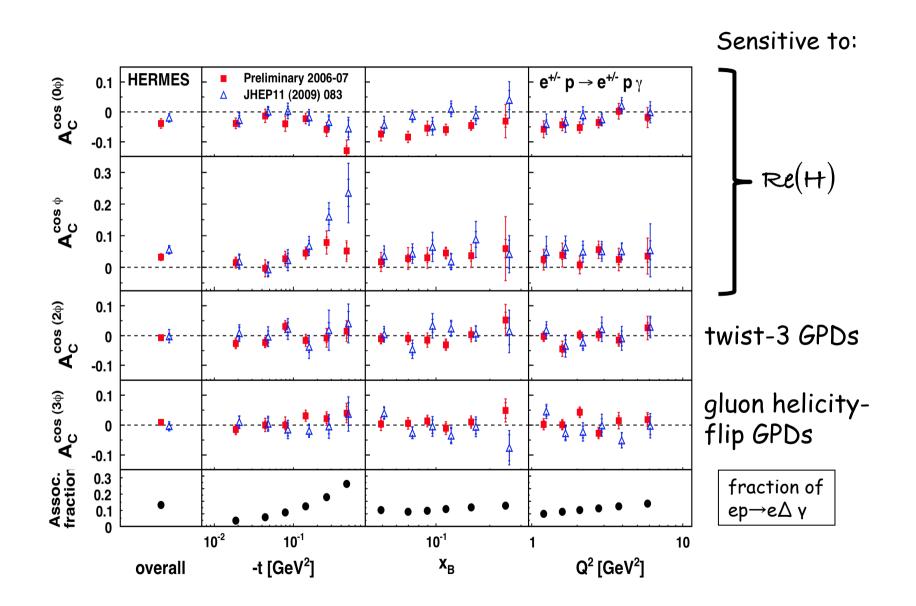
Cancels terms proportional to BH

Asymmetry of DVCS $\mathcal{A}_{\mathrm{LU}}^{\mathrm{DVCS}}(\phi) \equiv \frac{(d\sigma^{+\rightarrow} - d\sigma^{+\leftarrow}) \bigoplus (d\sigma^{-\rightarrow} - d\sigma^{-\leftarrow})}{(d\sigma^{+\rightarrow} + d\sigma^{+\leftarrow}) + (d\sigma^{-\rightarrow} + d\sigma^{-\leftarrow})}$

HERMES proton measurements



HERMES Beam Charge Asymmetries



HERMES Transverse Target Spin Asymmetries

 $A_{UT}^{sin(\varphi-\varphi_S)}$ 8.1% scale uncertainty 0.2 $A_{UT,I}$ $J_u = 0.6 - 0.4 - 0.2 \dots$ Im(H - E)ſ -0.2 ${\bf A}_{{\rm UT}}^{\sin(\varphi-\varphi_{\rm S})\cos\phi}$ $(\mathbf{E} \cdot \mathbf{E})$ -0.4 $A_{UT}^{cos(\varphi-\varphi_S)sin\varphi}$ 0.2 $Im(\tilde{H} - \tilde{E})$ 0 ΡIJ -0.2 .4 0.6 0 -t (GeV²) $\frac{6}{0} \frac{8}{8} \frac{10}{10}$ Q² (GeV²) 0.2 0.1 0.2 0.4 0.3 overall X_B

Sensitive to:

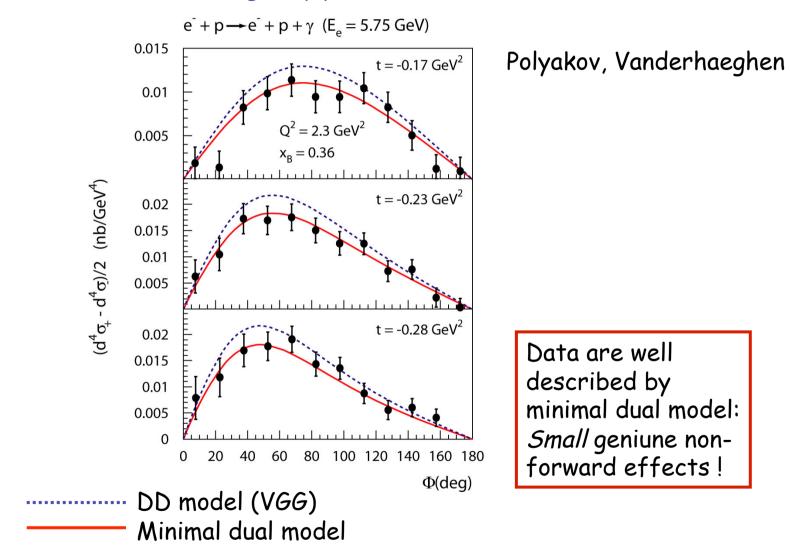
<u>A few interesting attempts to understand the data in terms of GPDs:</u>

- (minimal) Dual Model
- Local fits of Compton Form Factors
- Global fits using Dual Model
- « Dispersion relations » global fits
- ... (not exhaustive)

More information? don't miss the following talks: Moutarde, Kumericky, Semenov, Goldstein, Liuti

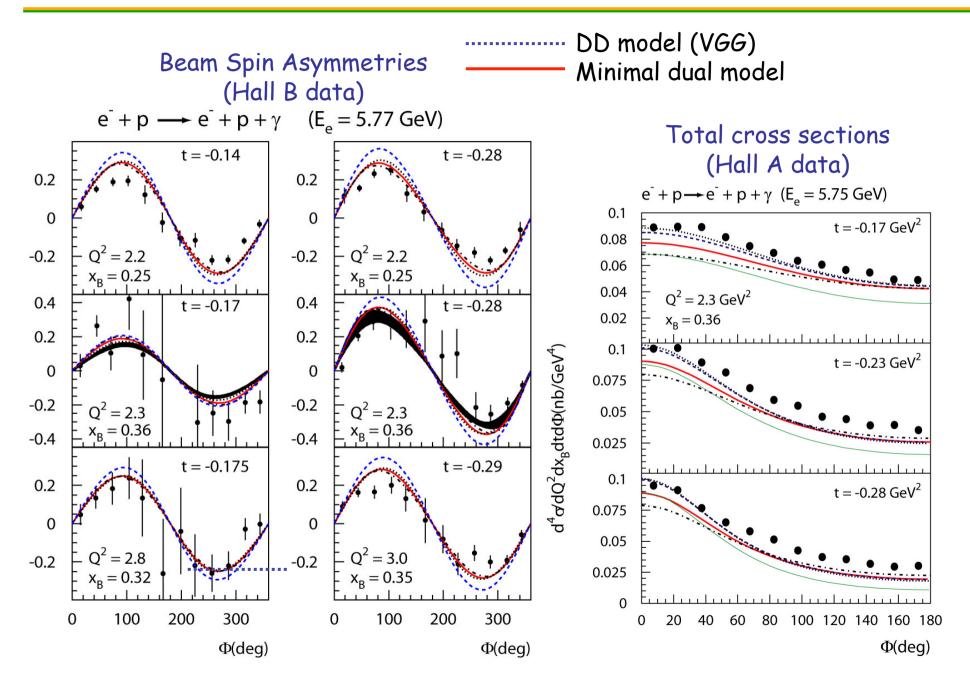
Minimal dual model (i.e. forward model) of DVCS data

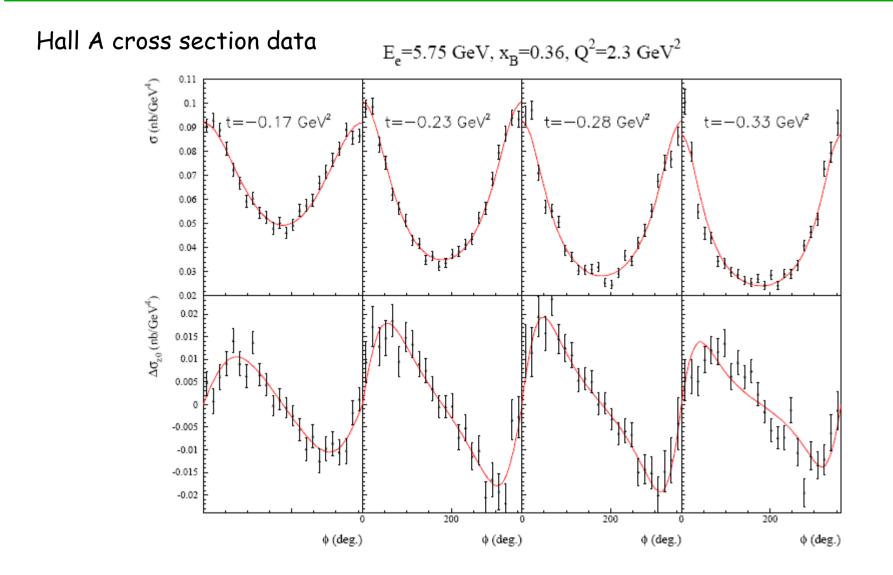
Difference of cross sections (imaginary part of interference term)



<u>Data</u>: Hall A polarized cross section data at Q²=2.3 GeV²

Minimal dual model (2)

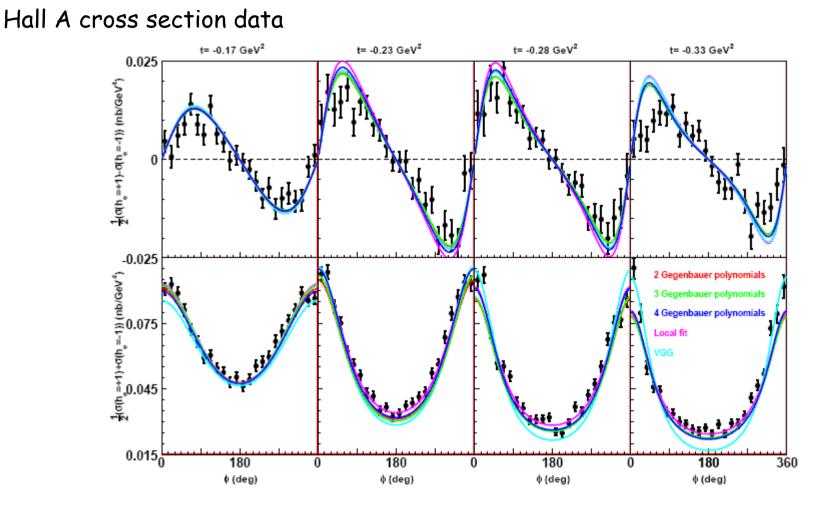




Includes H, \widetilde{H} and E

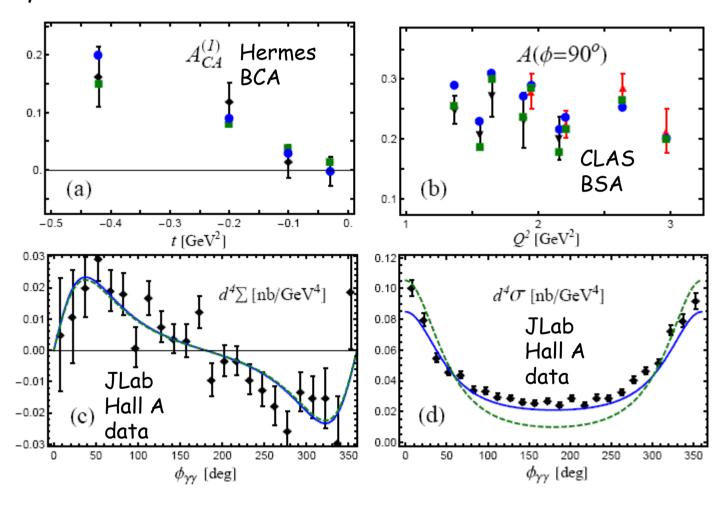
M. Guidal

Compton Form Factor (local) or dual model global fits (H only)



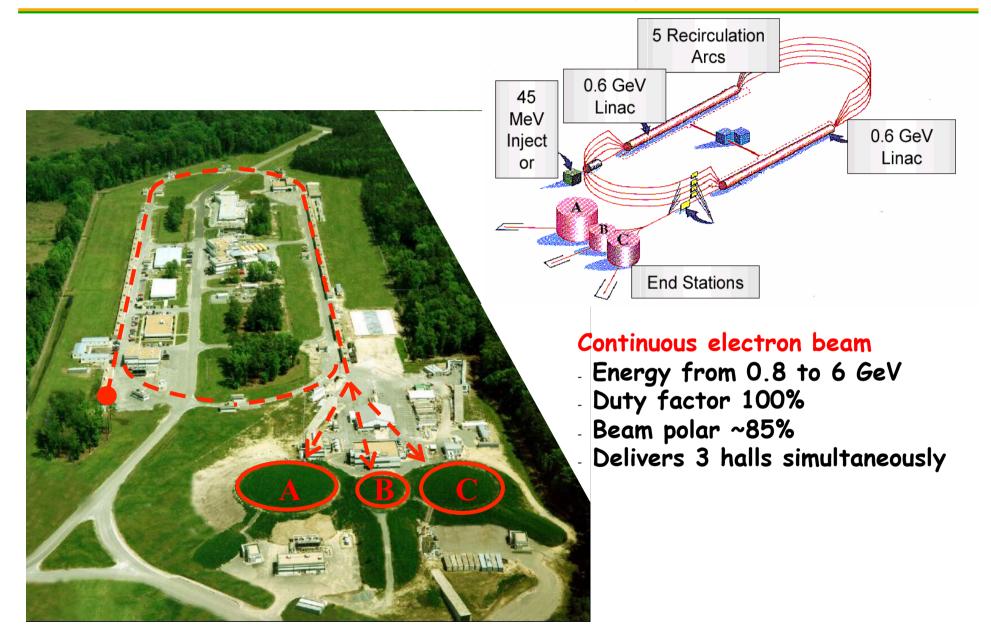
Presentation by H. Moutarde

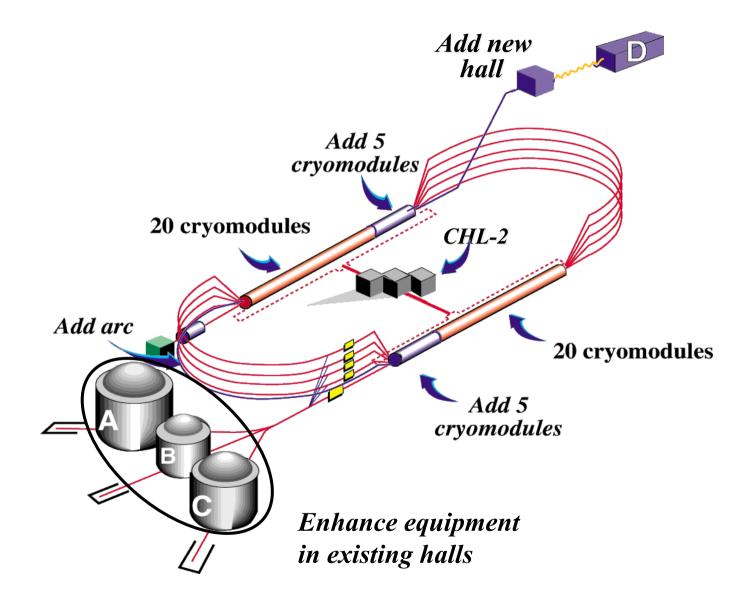
Presentation by K. Kumericky



• Fit of H1/ZEUS+CLAS+Hermes with only H • + Hall A and H + \tilde{H}

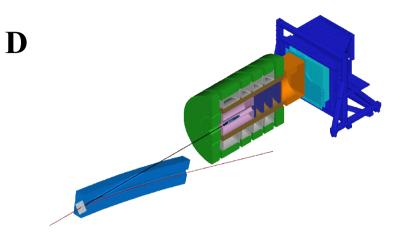
Jefferson Lab today





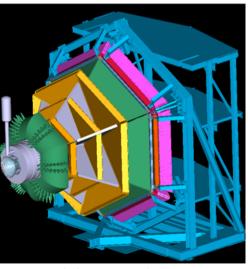
New equipments for 11-12 GeV beam

C

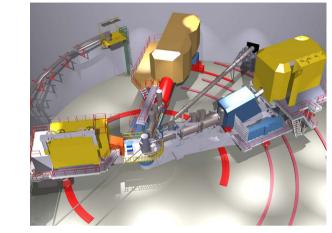


9 GeV tagged polarized photons and a 4π hermetic detector

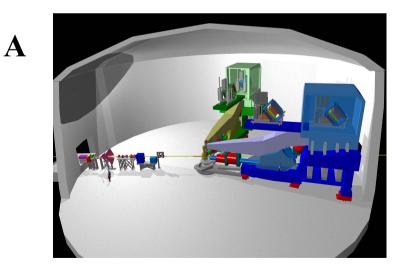
B



CLAS12 with new detectors and higher luminosity (10³⁵/cm²-s)

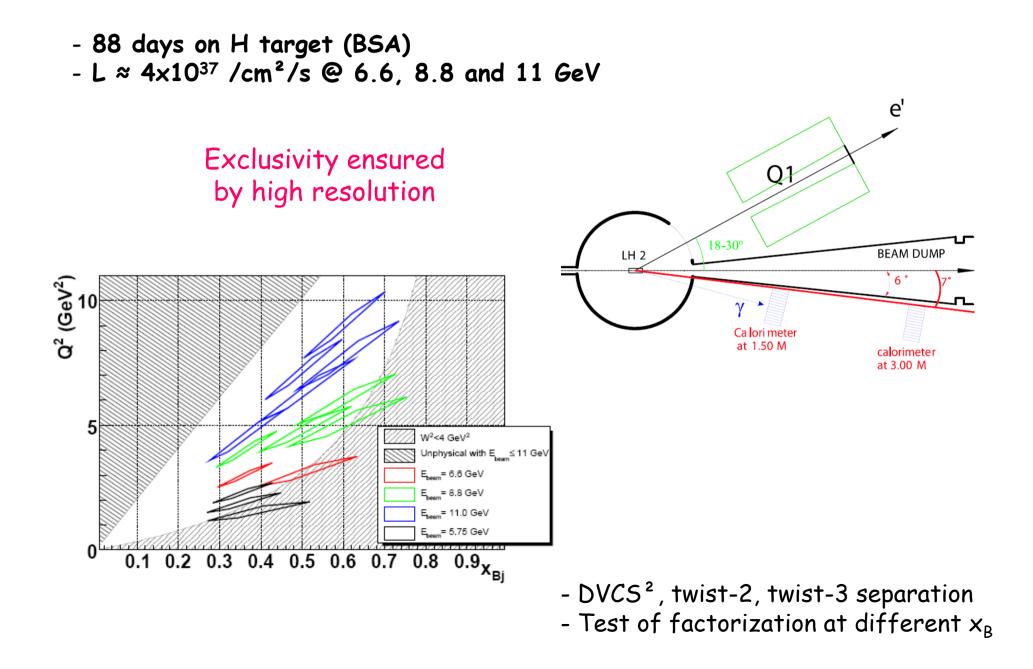


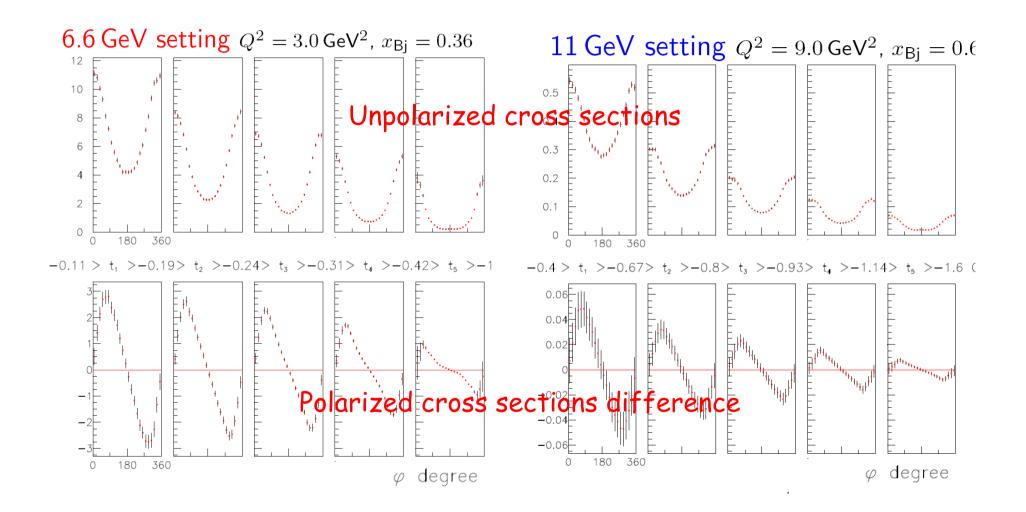
Super High Momentum Spectrometer (SHMS) at high luminosity and forward angles



High Resolution Spectrometer (HRS) Pair, and specialized large installation experiments

Extension of the Hall A experiments

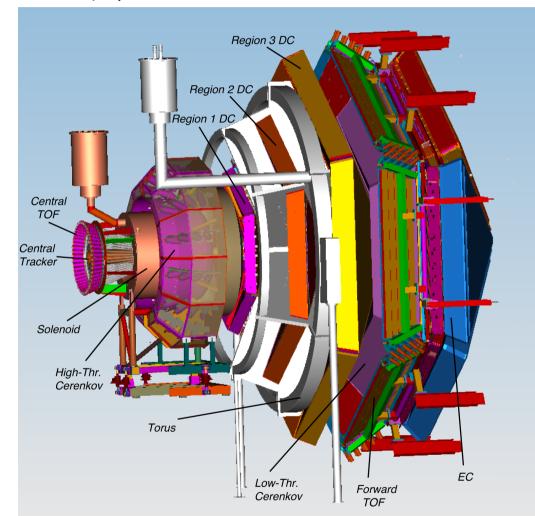




 \rightarrow Access the real and imaginary parts of T^{DVCS} separately

Experimental Setup and proposed experiments at 11 GeV

Use of base CLAS12 equipment,

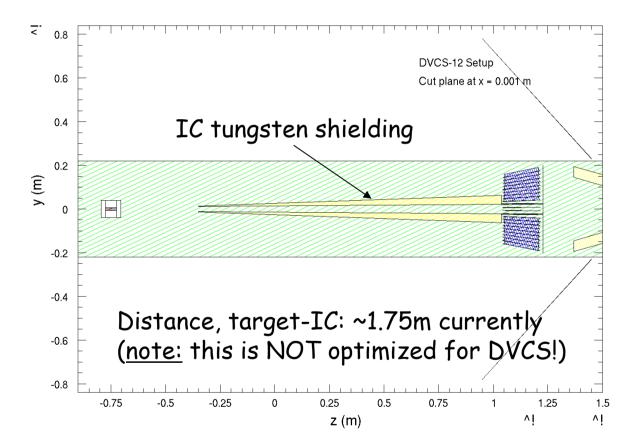


Detection of the full (e,p, γ) final state Perform 2 experiments for the extraction of the BSA and the TSA

Experimental Setup and proposed experiments at 11 GeV

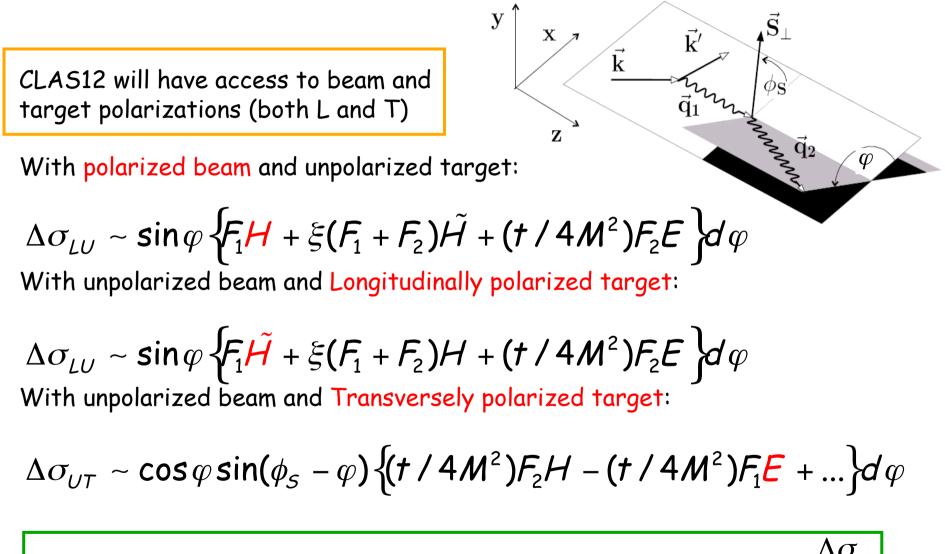
Use of base CLAS12 equipment, including Inner Calorimeter (IC)

2006/06/15 14.37



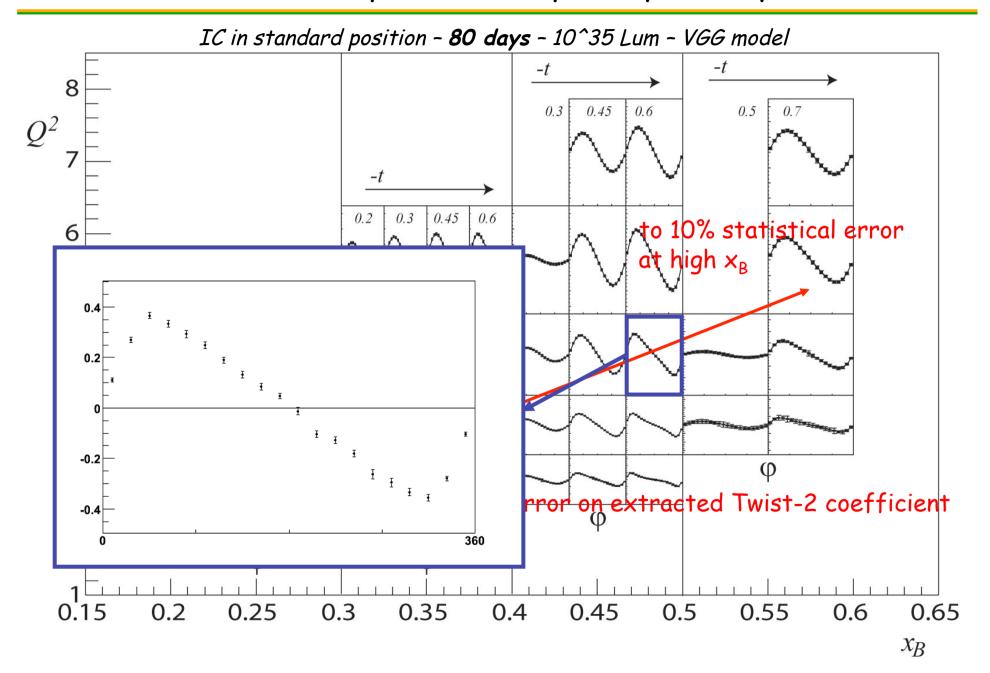
Detection of the full (e,p, γ) final state Perform 2 experiments for the extraction of the BSA and the TSA

Measuring DVCS using all polarization settings !

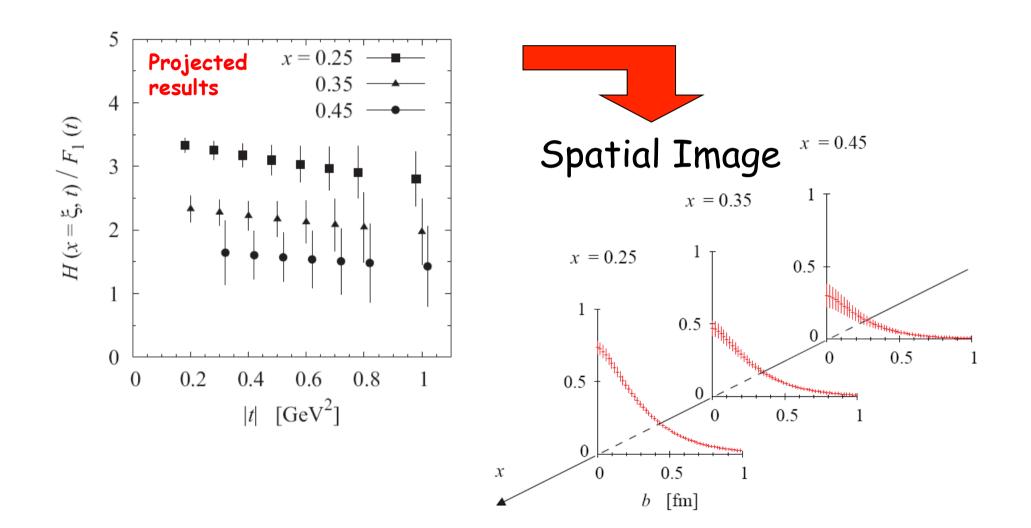


Depending on the experiment, measurement of $\Delta\sigma$ and σ or $A = \frac{\Delta\sigma}{2\sigma}$

One example : Beam Spin Asymmetry



Extraction of the GPD $H(x=\xi)$ from projected CLAS12 data



> A large set of data from HERA, HERMES, JLab Hall A/B

More data to come soon from HERMES (recoil detector), CLAS (2nd part of eldvcs experiment and Longitudinal + Transverse Target data), Hall A (full separation of cross section at 6 GeV)

> Even more data later on : COMPASS, CLAS12

> A few groups are working hard on ways to parametrize and extract GPDs from these data and show promising results and progress

The 3D pictures of the nucleon are not so far ahead, stay tuned !