



#### Klaus Rith



#### COMPASS Meeting, June 2018



## (still awaiting publication)

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#### Semi-inclusive Deep-Inelastic Scattering



DF(x,Q<sup>2</sup>): Parton Distribution Function –  $f_1(x,Q^2)$ ,  $g_{1L}(x,Q^2)$ ,  $h_1(x,Q^2)$ ,...

FF( $z,Q^2$ ): Fragmentation Function -  $D_1^{q \rightarrow h}(z,Q^2)$ ,  $H_1^{\perp q \rightarrow h}(z,Q^2)$ , ...















Topic 1: ALL<sup>h</sup>

## $\cos\phi$ moment of asymmetry $A_{LL}$



Consistent with zero

No significant kinematic dependence observed

Compatible result for proton and deuteron



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## $\cos\phi$ moment of asymmetry $A_{LL}$



Consistent with zero

No significant kinematic dependence observed

Compatible result for proton and deuteron

Presented at Praha, July 2012 !!!

#### Subleading twist coso moment of asymmetry ALL



Consistent with zero

No significant kinematic dependence observed

Compatible result for proton and deuteron

Presented at Praha, July 2012 !!!

# Topic 2: $A_{U(L)T}^{h}$ ('all' TMDs)

### TMDs from transversely polarized proton target

- Data taken in 2002-2005
- 5 PhD theses in 2004-2010
- 3 publications (Sivers and Collins for pions and charged kaons) [PRL 94 (2005) 012002, PRL 1003 (2009) 152002, PLB 693 (2010) 11]; (So far ~1140 citations)
  - Since then include protons and antiprotons
    - extend analysis to  $\square$  in (x, z, P<sub>h</sub>) for all but antiprotons (important for global fits)
    - detailed study of systematics
- Simultaneous fit of the Fourier amplitudes for  $sin(\phi-\phi_s)$ ,  $sin(\phi+\phi_s)$ ,  $sin(3\phi-\phi_s)$ ,  $sin(\phi_s)$ ,  $sin(2\phi-\phi_s)$ ,  $sin(2\phi+\phi_s)$  (UT)  $cos(\phi-\phi_s)$ ,  $cos(\phi_s)$ ,  $cos(2\phi-\phi_s)$  (LT)
- >200 plots

N/q	U	L	Т
U	- f <sub>1</sub>		$\mathbf{h}_{1}^{\perp}$
L		<b>g</b> 1	h <sub>1L</sub> ⊥
т	$f_{1T}^{\perp}$	g <sup>⊥</sup> <sub>1T</sub>	$h_1$ $h_{1T}^{\perp}$

 $C[h_1^q(x) \otimes H_1^{\perp,q}(z)]$ 



Transversity DF and Collins FF non-zero, lead to large effects

Opposite in sign for charged pions

Disfavored Collins FF large and opposite in sign to favored one

N/q	U	L	Т
U	- f <sub>1</sub>		$\mathbf{h}_1^{\perp}$
L		<b>g</b> 1	h <sub>1L</sub> ⊥
т	$f_{1T}^{\perp}$	g <sup>⊥</sup> <sub>1T</sub>	$h_1$ $h_{1T}^{\perp}$

#### $C[h_1^q(x) \otimes H_1^{\perp,q}(z)]$



Positive Collins SSA for positive Kaons

Consistent with zero for negative Kaons

N/q	U	L	Т
U	- f <sub>1</sub>		$\mathbf{h_1}^{\perp}$
L		<b>g</b> 1	h <sub>1L</sub> ⊥
т	$f_{1T}^{\perp}$	g <sup>⊥</sup> <sub>1T</sub>	$h_1$ $h_{1T}^{\perp}$

#### $C[h_1^q(x) \otimes H_1^{\perp,q}(z)]$



- Positive Collins SSA for positive Kaons
- Consistent with zero for negative Kaons and (anti-)protons
  - Vanishing sea-quark transversity and baryon Collins effect ?

N/q	U	L	Т
U	- f <sub>1</sub>		$\mathbf{h}_1^{\perp}$
L		<b>g</b> 1	h <sub>1L</sub> ⊥
Т	$f_{1T}^{\perp}$	g <sup>⊥</sup> <sub>1T</sub>	$h_1$ $h_{1T}^{\perp}$

#### $C[h_1^q(x) \otimes H_1^{\perp,q}(z)]$

Z. Kang et al., PRD 93 (2016) 014009

M. Anselmino et al., PRD 87 (2013) 094019 (M. Radici et al., JHEP 1505 (2015) 123)



Pion SIDIS data from COMPASS (PLB 673 (2009) 127, PLB 744 (2015) 250) HERMES (PRL 1003 (2009) 152002) JLAB - Hall A (PRL 107 (2011) 072003) e+e- data from BELLE (PRD 78 (2008) 032011) BABAR (PRD 90 (2014) 052003



- $\pi^-$  amplitudes increasing with x at large  $P_{h\perp}$ , increasing with z
- Other hadrons: no such clear kinematic dependencies
- No 3D for antiprotons

N/q	U	L	Т
U	f <sub>1</sub>		$h_1^{\perp}$
L		<b>g</b> 1	$\mathbf{h}_{1L}^{\perp}$
Т	$f_{1T}^{\perp}$	g <sup>⊥</sup> <sub>1T</sub>	$h_1  h_{1T}^{\perp}$

#### Sivers amplitudes

#### $C[f_{1T}^{\perp,q}(x) \otimes D_1^{q}(z)]$



- Positive Sivers amplitude for positive pions
  and kaons
- Experimental evidence for orbital angular momentum L<sub>q</sub> of quarks



Positive Sivers amplitude for (anti-)protons, Similar magnitude as for  $\pi^+$ 



Largest at large x and z, region of purest "u-quark probe"

N/q	U	L	Т
U	$ \mathbf{f_1}  $		$\mathbf{h}_1^{\perp}$
L		<b>g</b> 1	h <sub>1L</sub> ⊥
Т	$f_{1T}^{\perp}$	g <sup>⊥</sup> <sub>1T</sub>	$h_1  h_{1T}^{\perp}$

#### Sivers amplitudes



A. Airapetian et al., PRL 103 (2009) 152002







N/q	U	L	Т
U	- f <sub>1</sub>		$\mathbf{h_1}^{\perp}$
L		<b>g</b> 1	h <sub>1L</sub> ⊥
т	$f_{1T}^{\perp}$	g <sup>⊥</sup> <sub>1T</sub>	$h_1 $ $h_{1T}^{\perp}$

Pretzelosity DF  $h_{1T}^{\perp}$ 



Sideways transversely polarised quarks in transversely polarised nucleon

$$F_{UT}^{\sin(3\phi_h - \phi_S)} = \mathcal{C}\left[\frac{2\left(\hat{\boldsymbol{h}} \cdot \boldsymbol{p}_T\right)\left(\boldsymbol{p}_T \cdot \boldsymbol{k}_T\right) + \boldsymbol{p}_T^2\left(\hat{\boldsymbol{h}} \cdot \boldsymbol{k}_T\right) - 4\left(\hat{\boldsymbol{h}} \cdot \boldsymbol{p}_T\right)^2\left(\hat{\boldsymbol{h}} \cdot \boldsymbol{k}_T\right)}{2M^2 M_h} H_1^{\perp}\right]$$

- leading-twist
- Related to parton orbital motion: requires interference between wave functions with OAM difference by 2 units
- Expected to scale with  $(p_T)^2 k_T$
- Suppressed w.r.t Collins and Sivers (these scale with  $k_T$ ,  $p_T$ )
  - Cahn, Boer-Mulders ( $\langle cos \phi \rangle$  scales with  $k_T$ ,  $p_T$ )
  - Boer-Mulders ( $\langle cos2\phi \rangle$  scales with  $k_Tp_T$ )

N/q	U	L	т
U	- f <sub>1</sub>		$\mathbf{h}_{1}^{\perp}$
L		<b>g</b> 1	h₁∟⊥
т	$f_{1T}^{\perp}$	g <sup>⊥</sup> <sub>1T</sub>	$h_1  h_{1T}^{\perp}$

#### Pretzelosity DF $h_{1T}^{\perp}$

 $C[h_1^{\perp, q}(x) \otimes H_1^{\perp, q}(z)]$ 







- Compatible with zero within uncertainties
  - h<sub>1T</sub> $^{\perp}$  might be non-zero, look at higher p<sub>T</sub>

N/q	U	L	т
U	- f <sub>1</sub>		$\mathbf{h}_{1}^{\perp}$
L		<b>g</b> 1	h₁∟⊥
т	$f_{1T}^{\perp}$	g <sup>⊥</sup> <sub>1T</sub>	$h_1  h_{1T}^{\perp}$

#### **Pretzelosity DF** $h_{1T}^{\perp}$









•  $h_{1T}^{\perp}$  might be non-zero, look at higher  $p_T$ C. Lefky, A. Prokudin, PRD 91 (2015) 034010



Data from COMPASS, HERMES, JLAB



Worm-gear DF  $g_{1T}^{\perp,q}$ 

longitudinally polarised quarks in transversely polarised nucleon

 $C[g_{1T}^{\perp, q}(\mathbf{x}) \otimes D_1^{q}(\mathbf{z})]$ 







Related to parton orbital motion: requires interference between wave functions with OAM difference by 1 unit

$$g_{1T}^{\perp,q} = -h_{1L}^{\perp,q} \text{ (supported by many models)}$$

$$g_{1T}^{\perp,q} \approx \times \int_{x}^{1} \frac{dy}{y} g_{1}^{q}(y) \text{ (Wandzura-Wilczek type approximation)}$$

wanazura-wiiczek type approximation)

Slightly non-zero

for pions, kaons, protons





# Topic 4: A<sub>LU</sub><sup>h</sup>



- Agreement between H and D data
- Positive asymmetries for pions



No clear kinematic dependencies in 3D



Consistent behavior for charged pions/hadrons at HERMES/COMPASS for isoscalar targets



CLAS probes higher x region; more sensitive to x  $e^q \otimes H_1^{\perp,q}$ ?



## Topic 5: **A** Polarization

$$\vec{l} + N \rightarrow \vec{l} + \vec{\lambda} + X$$
  $\Lambda \rightarrow p + \pi^{-}$ 

A polarization A from proton angular distribution in A rest frame

$$\propto (1 + \alpha \overrightarrow{P^{\wedge}} \cdot \overrightarrow{p}) = (1 + \alpha \sum_{i} P_{i}^{\wedge} \cos \theta_{i}) \quad (i = X, Y, Z)$$



- P<sub>Z</sub><sup>A</sup> = P<sub>B</sub>  $\sqrt{1-\epsilon^2}$  D<sub>LZ</sub>(x,z) P<sub>B</sub> = |P<sub>B</sub>| $\lambda_e$  longitudinal beam polarization
  P<sub>X</sub><sup>A</sup> = P<sub>B</sub>  $2\sqrt{2\epsilon(1-\epsilon)}$  D<sub>LX</sub>(x,z)
  P<sub>y</sub><sup>A</sup> =  $2\sqrt{2\epsilon(1+\epsilon)}$  D<sub>LY</sub>(x,z) beam-helicity independent transverse  $\Lambda$  pol.
- D<sub>LX</sub>(x,z), D<sub>LY</sub>(x,z), D<sub>LZ</sub>(x,z): "spin-transfer coefficients"
- Asymmetry for helicity balanced data set: Py^ drops out

Relations between spin-transfer coefficients and TMDs: (P. Mulders and R. Tangerman, Nucl. Phys. B461 (1996) 197)

$$D_{LZ}(x,z) = \frac{\sum_{q} e_{q}^{2} \times f_{1}^{q}(x) G_{1}^{q \to \Lambda}(z)}{\sum_{q} e_{q}^{2} \times f_{1}^{q}(x) D_{1}^{q \to \Lambda}(z)}$$

Production of longitudinally polarized  $\Lambda s$ from originally unpolarized quarks and longitudinally polarized virtual photons

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Production of longitudinally polarized  $\Lambda s$  from originally unpolarized quarks and longitudinally polarized virtual photons

$$\mathsf{D}_{\mathsf{LX}}(\mathsf{x},\mathsf{z}) = \frac{-\frac{\mathsf{M}}{\mathsf{Q}} \{ \sum_{\mathsf{q}} e_{\mathsf{q}}^{2} \mathsf{x}^{2} e^{\mathsf{q}}(\mathsf{x}) \, \mathsf{H}_{1}^{\perp,\mathsf{q}} \, \wedge \, (\mathsf{z}) + \frac{\mathsf{M}^{\Lambda}}{\mathsf{M}} \sum_{\mathsf{q}} e_{\mathsf{q}}^{2} \mathsf{x} f_{1}^{\mathsf{q}}(\mathsf{x}) \, \tilde{\mathsf{G}}^{\perp,\mathsf{q}} \, \wedge \, (\mathsf{z})/\mathsf{z} \}}{\sum_{\mathsf{q}} e_{\mathsf{q}}^{2} \mathsf{x} f_{1}^{\mathsf{q}}(\mathsf{x}) \, \mathsf{D}_{1}^{\mathsf{q}} \, \wedge \, (\mathsf{z})}$$

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$$\mathsf{D}_{\mathsf{LZ}}(\mathsf{x},\mathsf{z}) = \frac{\sum_{q} e_{q}^{2} \mathsf{x} \mathsf{f}_{1}^{q}(\mathsf{x}) \, \mathcal{G}_{1}^{q \to \Lambda}(\mathsf{z})}{\sum_{q} e_{q}^{2} \mathsf{x} \mathsf{f}_{1}^{q}(\mathsf{x}) \, \mathsf{D}_{1}^{q \to \Lambda}(\mathsf{z})}$$

Production of longitudinally polarized  $\Lambda s$  from originally unpolarized quarks and longitudinally polarized virtual photons

$$D_{LX}(x,z) = \frac{-\frac{M}{Q} \{\sum_{q} e_{q}^{2} \times 2e^{q}(x) H_{1}^{\perp,q} \rightarrow \Lambda(z) + \frac{M}{M} \sum_{q} e_{q}^{2} \times f_{1}^{q}(x) \widetilde{G}^{\perp,q} \rightarrow \Lambda(z)/z \}}{\sum_{q} e_{q}^{2} \times f_{1}^{q}(x) D_{1}^{q} \rightarrow \Lambda(z)}$$
Same combinations of PDFs and FF as in Twist-3  $\langle sin(\phi) \rangle_{LU}$   
 $\propto C[h_{1}^{\perp,q} \otimes \widetilde{E}^{q}, x e^{q} \otimes H_{1}^{\perp,q}; x g^{\perp,q} \otimes D_{1}^{q}, f_{1}^{q} \otimes \widetilde{G}^{\perp,q}]$ 





## Topic 6: Kaons from nuclei

#### Fragmentation in nuclear matter



Courtesy of J. Rubin)

Useful for understanding the fundamental aspects of hadronisation

Input for calculations of nuclear PDFs and FF

- typical hadronisation length  $\propto$  (1-z) v is of order of nucleus size (1-10 fm)
- time development of hadronisation can be studied with nuclei of increasing size
- struck quark or qq-pair propagate through ,cold' nuclear matter
- interaction signature: reduction of the numer of hadrons per DIS event and per nucleon

$$R_A^h(\nu, Q^2, z, p_t^2) = \frac{\left(\frac{N^h(\nu, Q^2, z, p_t^2)}{N^e(\nu, Q^2)}\right)_A}{\left(\frac{N^h(\nu, Q^2, z, p_t^2)}{N^e(\nu, Q^2)}\right)_D}$$

Multi-dimensional study

#### Fragmentation in nuclear matter



- p: weak v-dependence
- p: R<sub>A</sub> exceeding unity at high v and low z (apart from hadronisation other production mechanisms contribute)

#### Charged Kaon production in nuclear matter

#### N.B. Chang et al., PRC 89 (2014) 034911



FIG. 1: (color online) Parton fragmentation functions to  $K^-$  in vacuum at  $Q^2 \approx 10 \text{ GeV}^2$  from the HKN parametrization [27].



FIG. 2: (color online) Ratios of mFF's to the vacuum ones for  $K^-$  from different partons with initial energy  $\nu=10~{\rm GeV}$  and  $Q^2=10~{\rm GeV}^2$  in SIDIS off Pb.



FIG. 3: (color online) The nuclear modification factor for (a)  $K^+$  and (b)  $K^-$  for different initial quark energy  $\nu$  in SIDIS at  $x_B = 0.1$ .





Prediction: enhanced medium modified FF for K<sup>-</sup> and enhanced K-/K+ production ratio at large × and z

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Prediction: enhanced medium modified FF for K<sup>-</sup> and enhanced K-/K+ production ratio at large x and z

Analysis completed. Results to be released in a few weeks