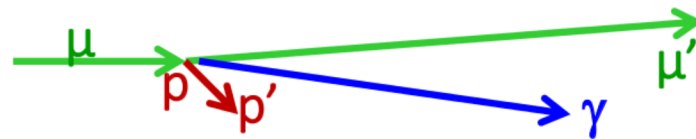


# Hard Exclusive Reactions at COMPASS at CERN

Exclusive photon (DVCS) and meson (HEMP) production at small transfer for GPD studies



$$\text{DVCS : } \mu \ p \rightarrow \mu' \ p' \ \gamma$$



$$\text{Pseudo-Scalar Meson : } \mu \ p \rightarrow \mu' \ p' \ \pi^0$$

$$\text{Vector Meson : } \mu \ p \rightarrow \mu' \ p' \ \rho \text{ or } \omega \text{ or } \phi \dots$$

*Nicole d'Hose - CEA Université Paris-Saclay for the COMPASS Collaboration*



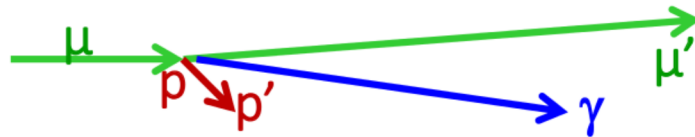
**CENTER for  
NUCLEAR FEMTOGRAPHY**

**CNF Generalized Parton Distributions and  
Global Analysis**

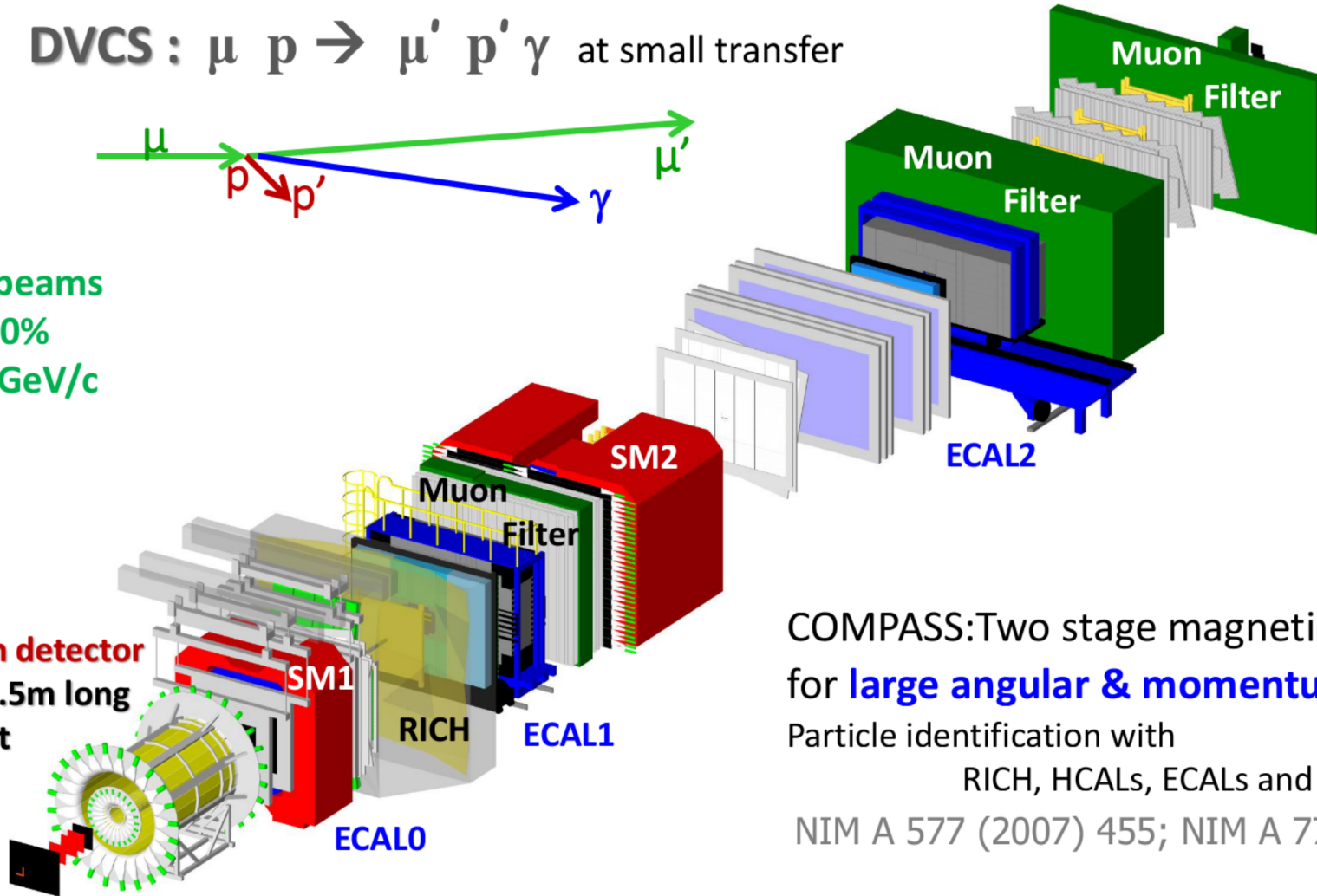
12–14 Jun 2023

# Measurement of exclusive cross sections at COMPASS

DVCS :  $\mu p \rightarrow \mu' p' \gamma$  at small transfer



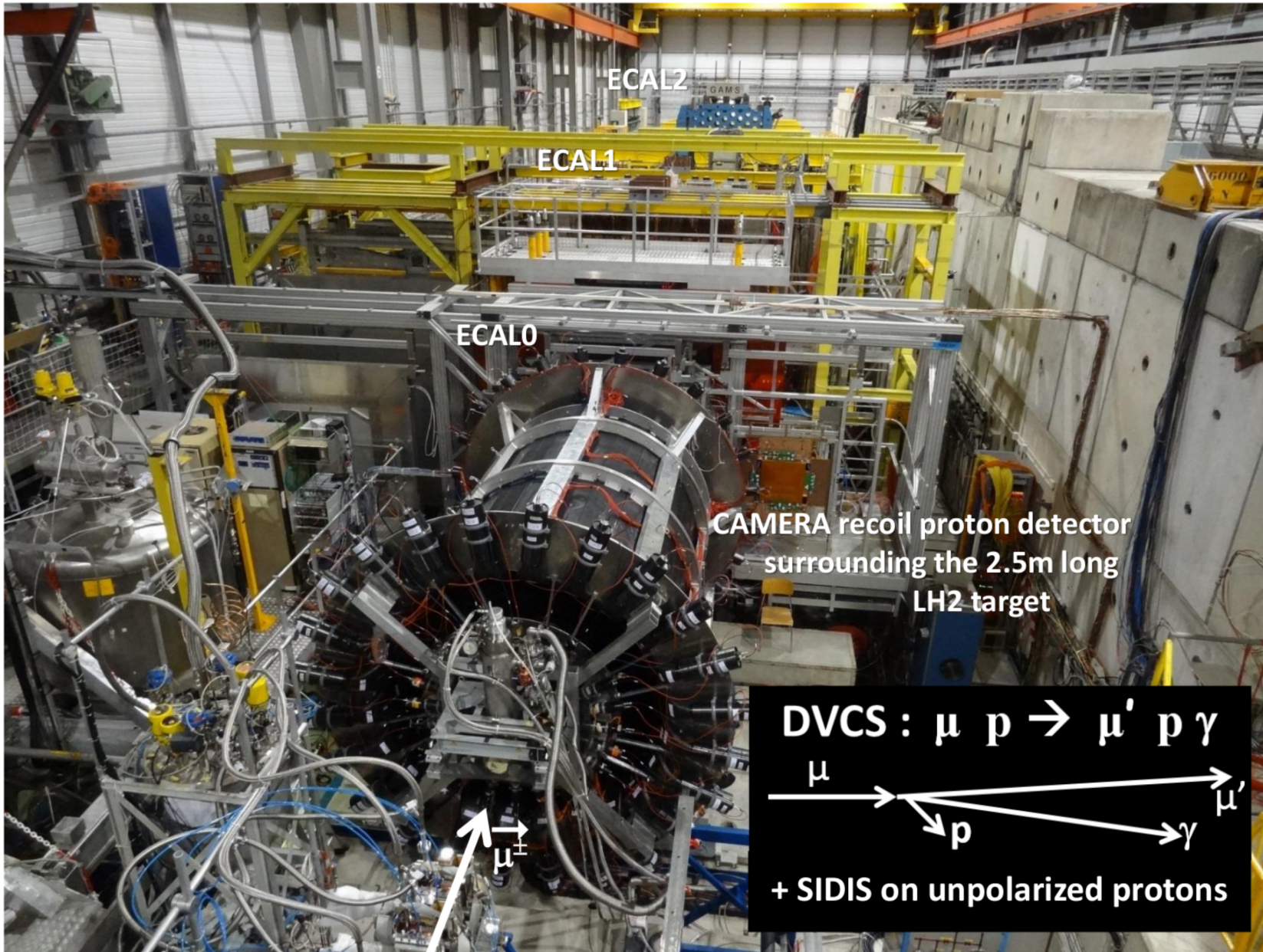
Both  $\mu^+$  and  $\mu^-$  beams  
Polarisation  $\sim \pm 80\%$   
Momentum 160 GeV/c



COMPASS: Two stage magnetic spectrometer for **large angular & momentum acceptance**

Particle identification with RICH, HCALs, ECALs and muon filters

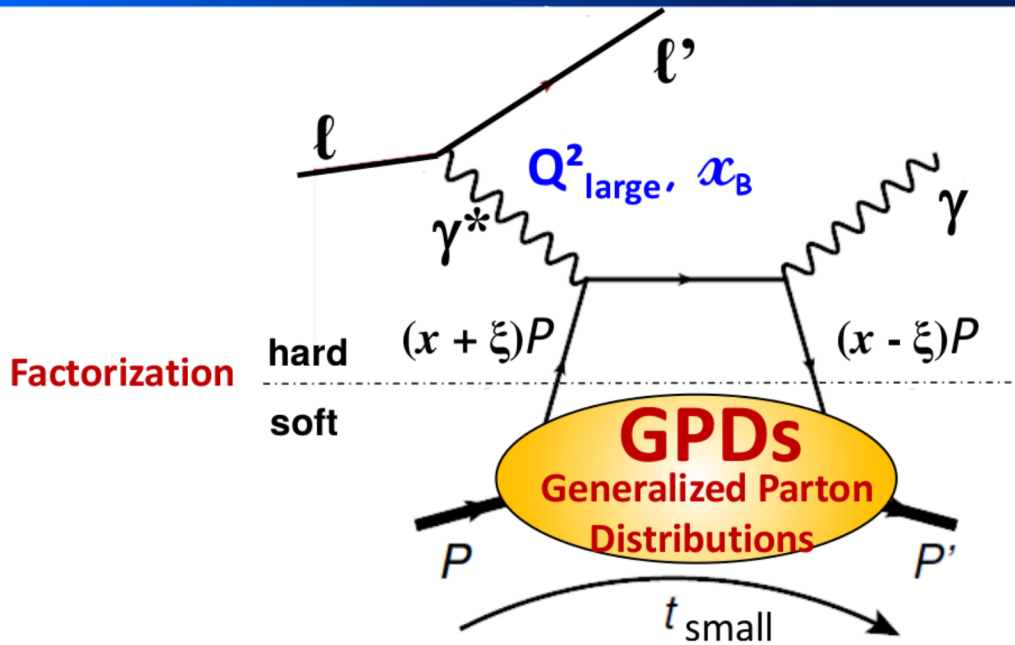
NIM A 577 (2007) 455; NIM A 779 (2015) 69



**2012:**  
1 month pilot run

**2016 -17:**  
2 x 6 month  
data taking

# Deeply virtual Compton scattering (DVCS)



D. Mueller *et al*, Fortsch. Phys. 42 (1994)  
 X.D. Ji, PRL 78 (1997), PRD 55 (1997)  
 A. V. Radyushkin, PLB 385 (1996), PRD 56 (1997)

DVCS:  $\ell p \rightarrow \ell' p' \gamma$   
 the golden channel  
 because it interferes with  
 the Bethe-Heitler process  
  
 also meson production  
 $\ell p \rightarrow \ell' p' \pi, \rho, \omega$  or  $\phi$  or  $J/\psi \dots$

The GPDs depend on the following variables:

$x$ : average } quark longitudinal  
 $\xi$ : transferred } momentum fraction

$t$ : proton momentum transfer squared  
 related to  $b_{\perp}$  via Fourier transform

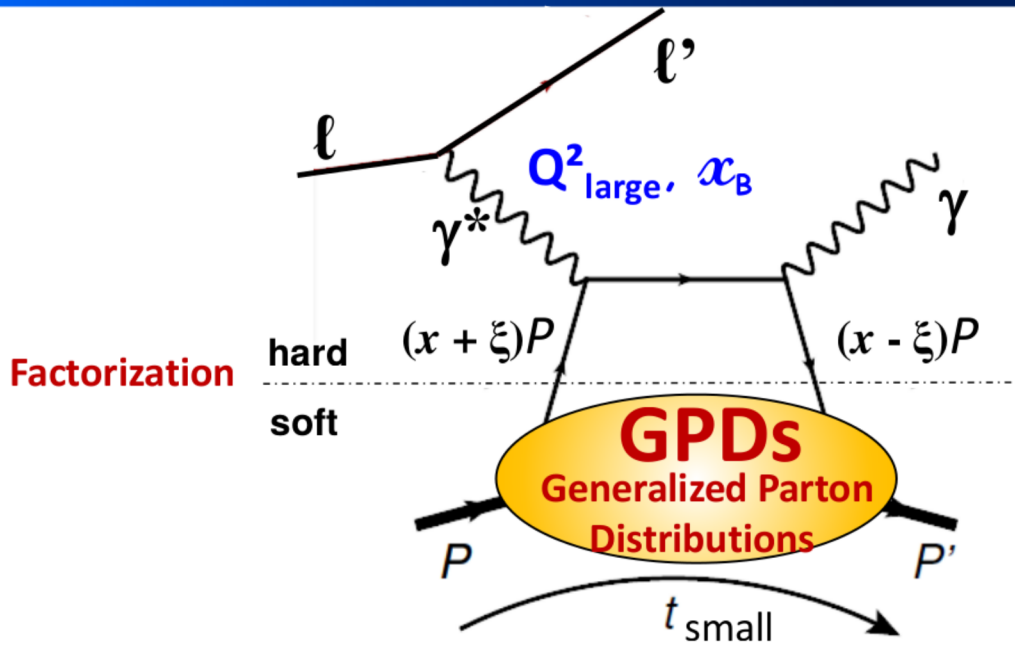
$Q^2$ : virtuality of the virtual photon

The variables measured in the experiment:

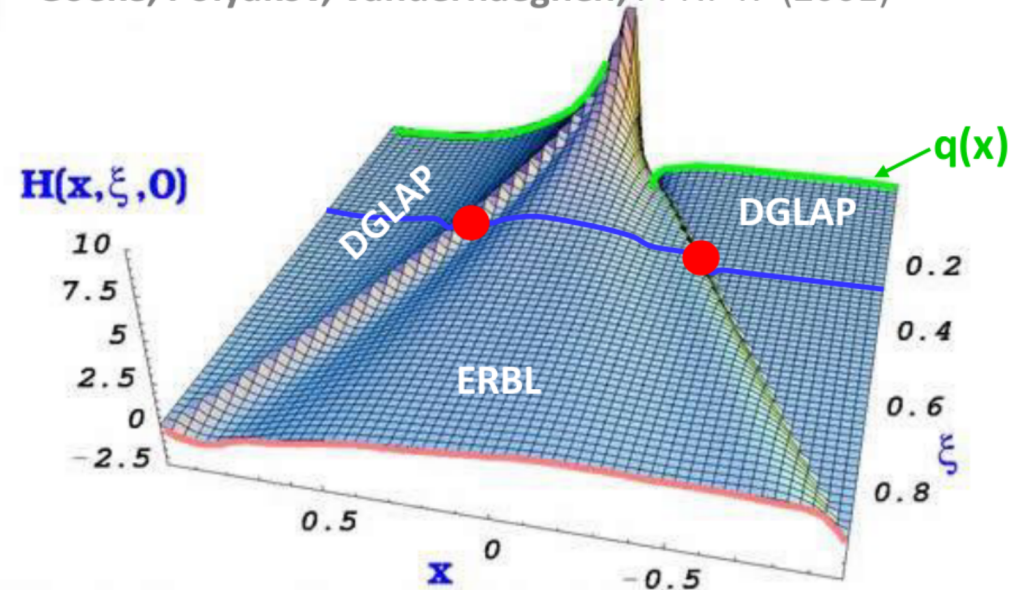
$E_{\ell}, Q^2, x_B \sim 2\xi / (1 + \xi),$

$t$  (or  $\theta_{\gamma^* \gamma}$ ) and  $\phi$  ( $\ell \ell'$  plane /  $\gamma \gamma^*$  plane)

# Deeply virtual Compton scattering (DVCS)



Goeke, Polyakov, Vanderhaeghen, PPNP47 (2001)



The amplitude DVCS at LT & LO in  $\alpha_s$  (GPD  $\mathbf{H}$ ):

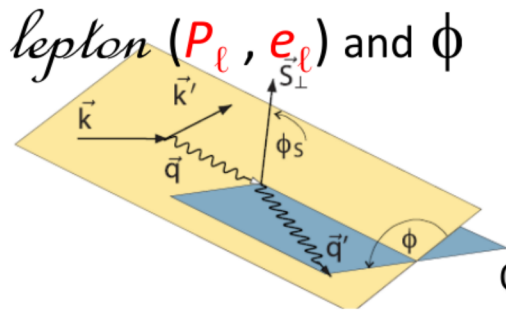
Real part

Imaginary part

$$\mathcal{H} = \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi + i\varepsilon} = \mathcal{P} \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi} - i\pi H(x = \pm \xi, \xi, t)$$

In an experiment we measure Compton Form Factor  $\mathcal{H}$

# Deeply virtual Compton scattering (DVCS)



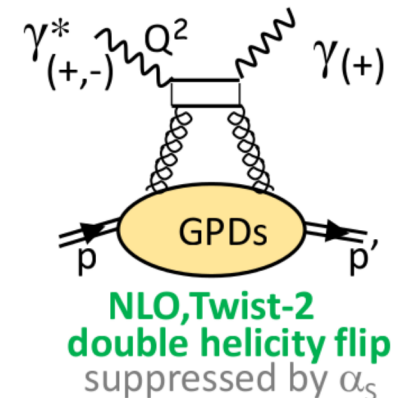
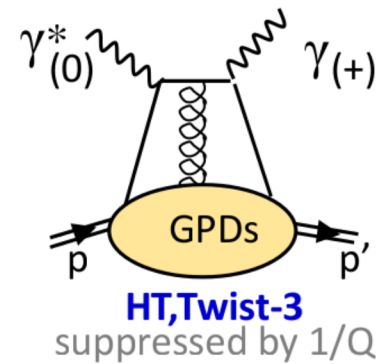
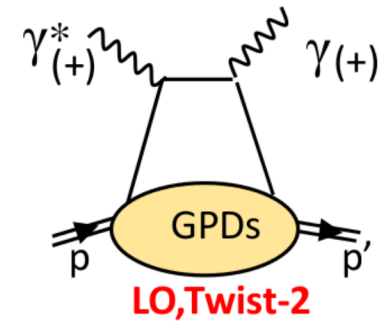
$$d\sigma = \left| \begin{array}{c} \text{BH} \\ \text{DVCS} \end{array} \right|^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underset{\text{Well known}}{d\sigma^{BH}} + \left( d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) - (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$

With unpolarized target:

Belitsky, Müller, Kirner, NPB629 (2002)

$$\begin{aligned} d\sigma^{BH} &\propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi \\ d\sigma_{unpol}^{DVCS} &\propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi \\ d\sigma_{pol}^{DVCS} &\propto s_1^{DVCS} \sin \phi \\ \text{Re } I &\propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi \\ \text{Im } I &\propto s_1^I \sin \phi + s_2^I \sin 2\phi \end{aligned}$$



# Deeply virtual Compton scattering (DVCS)

With both  $\mu^{\leftarrow}$  and  $\mu^{\rightarrow}$  beams we can build:

① beam charge-spin sum

$$\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$$

$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$$

$$d\sigma_{unpol}^{DVCS} \propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$$

② difference

$$\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$$

$$d\sigma_{pol}^{DVCS} \propto s_1^{DVCS} \sin \phi$$

$$\text{Re } I \propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$$

$$\text{Im } I \propto s_1^I \sin \phi + s_2^I \sin 2\phi$$

$$\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-} \rightarrow s_1^I \propto \text{Im } \mathcal{F}$$

$$\text{and } c_0^{DVCS} \propto (\text{Im } \mathcal{H})^2$$

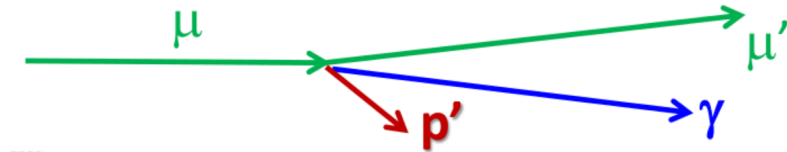
$$\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-} \rightarrow c_1^I \propto \text{Re } \mathcal{F}$$

$$\mathcal{F} = F_1 \mathcal{H} + \xi (F_1 + F_2) \tilde{\mathcal{H}} - t/4m^2 F_2 \mathcal{E}$$

for proton  
 $\rightarrow$   $F_1 \mathcal{H}$   
 at small  $x_B$   
 COMPASS domain

# COMPASS 2016 data Selection of exclusive single photon production

Comparison between the observables given by the spectro or by CAMERA



**DVCS:  $\mu p \rightarrow \mu' p \gamma$**

1)  $\Delta\varphi = \varphi^{\text{cam}} - \varphi^{\text{spec}}$

2)  $\Delta p_T = p_T^{\text{cam}} - p_T^{\text{spec}}$

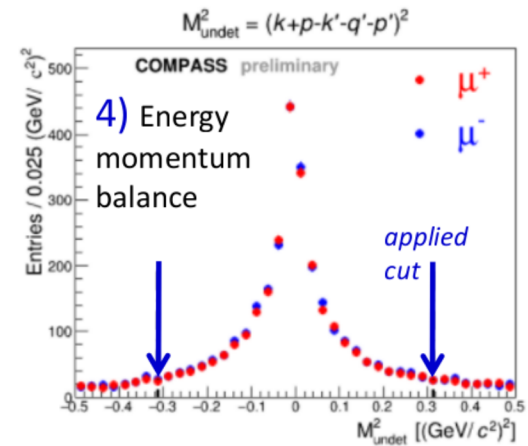
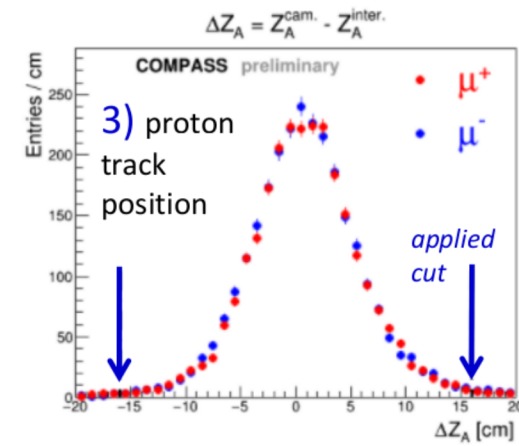
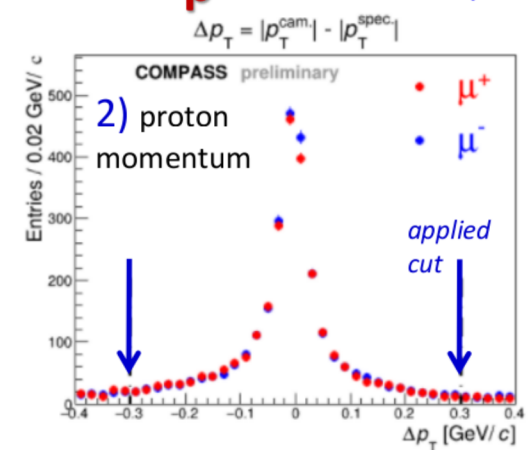
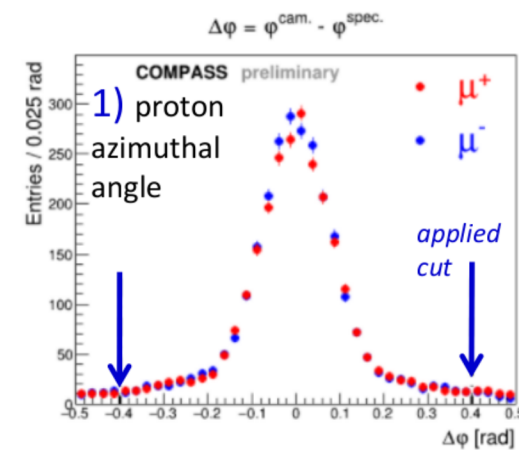
3)  $\Delta z_A = z_A^{\text{cam}} - z_A^{\text{ZB and vertex}}$

4)  $M_{X=0}^2 = (p_{\mu_{\text{in}}} + p_{p_{\text{in}}} - p_{\mu_{\text{out}}} - p_{p_{\text{out}}} - p_{\gamma})^2$

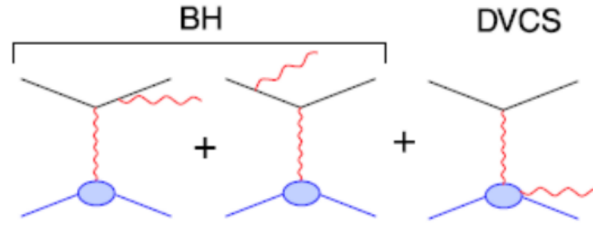
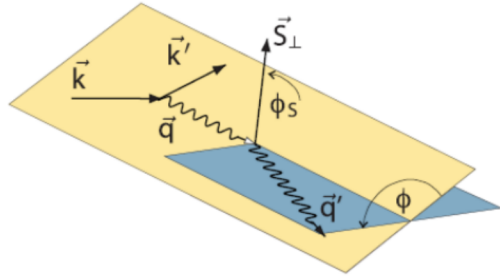
Good agreement between  $\vec{\mu}^+$  and  $\vec{\mu}^-$  yields  
 Important achievement for:

①  $\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$  **Easier, done first**

②  $\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$  **Challenging, but promising**



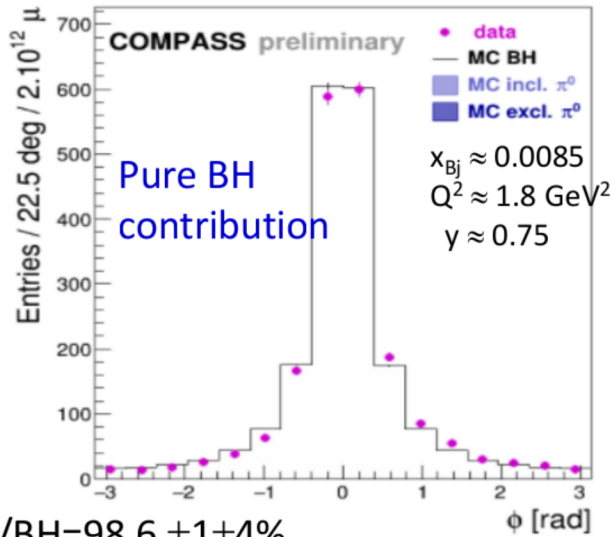




$$\Sigma = d\sigma(\mu^+) + d\sigma(\mu^-)$$

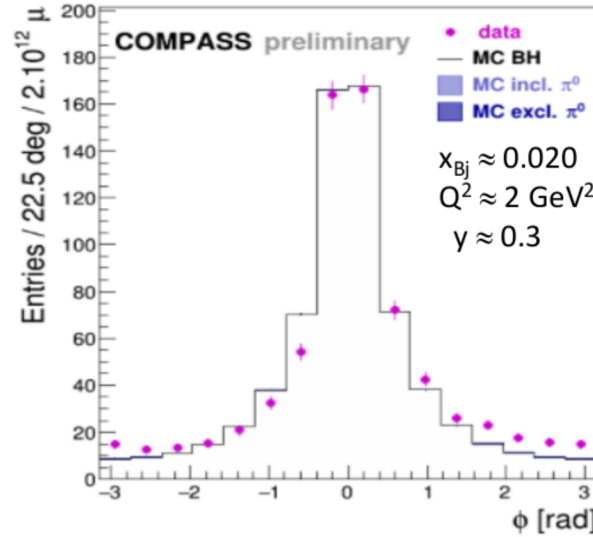
$$d\sigma \propto |T^{BH}|^2 + \text{Interference Term} + |T^{DVCS}|^2$$

80 <  $v$  [GeV] < 144

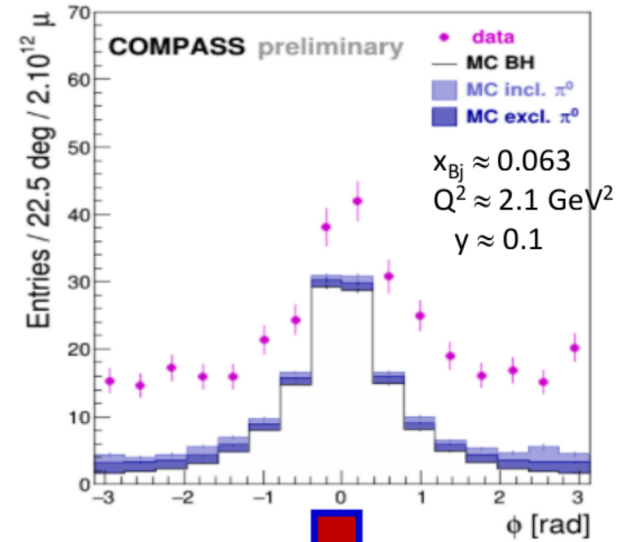


Data/BH =  $98.6 \pm 1 \pm 4\%$

32 <  $v$  [GeV] < 80



10 <  $v$  [GeV] < 32



**DVCS** above the **BH** contrib.

MC: BH contribution evaluated for the integrated luminosity  
 $\pi^0$  background contribution from SIDIS (LEPTO) + exclusive production (HEPGEN)

At COMPASS using polarized positive and negative muon beams:

$$\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-} = 2[d\sigma^{BH} + d\sigma_{unpol}^{DVCS} + \text{Im } I]$$

$$= 2[d\sigma^{BH} + c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi + s_1^I \sin \phi + s_2^I \sin 2\phi]$$

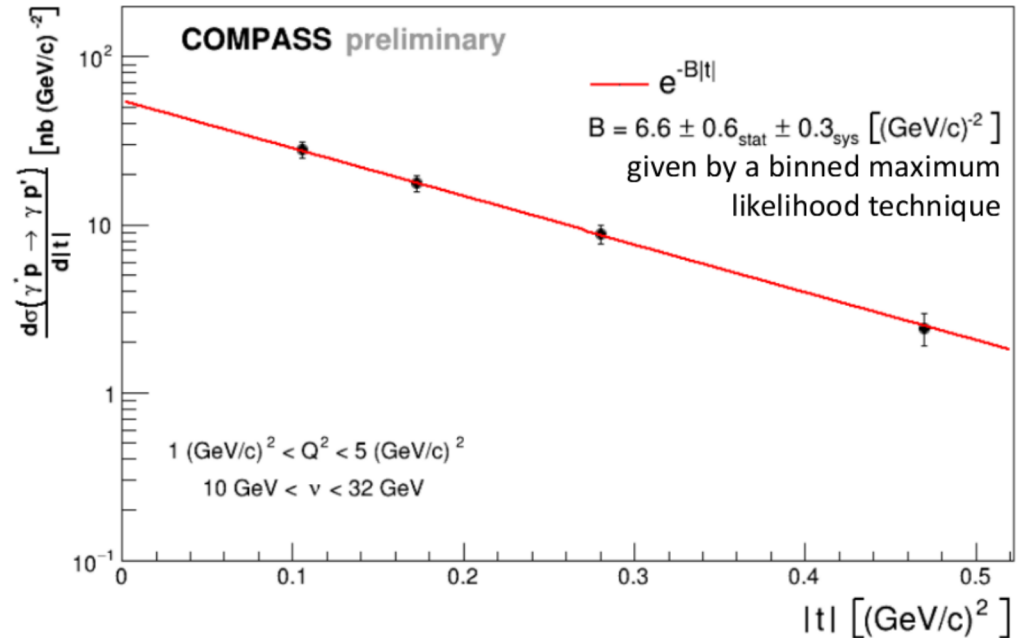
calculable  
can be subtracted

All the other terms are cancelled in the integration over  $\phi$

$$\frac{d^3\sigma_T^{\mu p}}{dQ^2 d\nu dt} = \int_{-\pi}^{\pi} d\phi (d\sigma - d\sigma^{BH}) \propto c_0^{DVCS}$$

$$\frac{d\sigma^{\gamma^* p}}{dt} = \frac{1}{\Gamma(Q^2, \nu, E_\mu)} \frac{d^3\sigma_T^{\mu p}}{dQ^2 d\nu dt}$$

Flux for transverse  
virtual photons



# COMPASS 12-16 Transverse extension of partons in the sea quark range

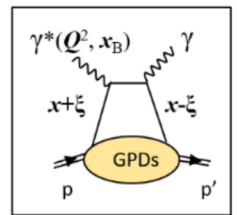
$$d\sigma^{DVCS}/dt = e^{-B|t|} = c_0^{DVCS} \propto (Im\mathcal{H})^2$$

$$c_0^{DVCS} \propto 4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) + \frac{t}{M^2}\mathcal{E}\mathcal{E}^*$$

In the COMPASS kinematics,  $x_B \approx 0.06$ , dominance of  $Im\mathcal{H}$   
 97% (GK model) 94% (KM model)

$$Im\mathcal{H} = H(x=\xi, \xi, t)$$

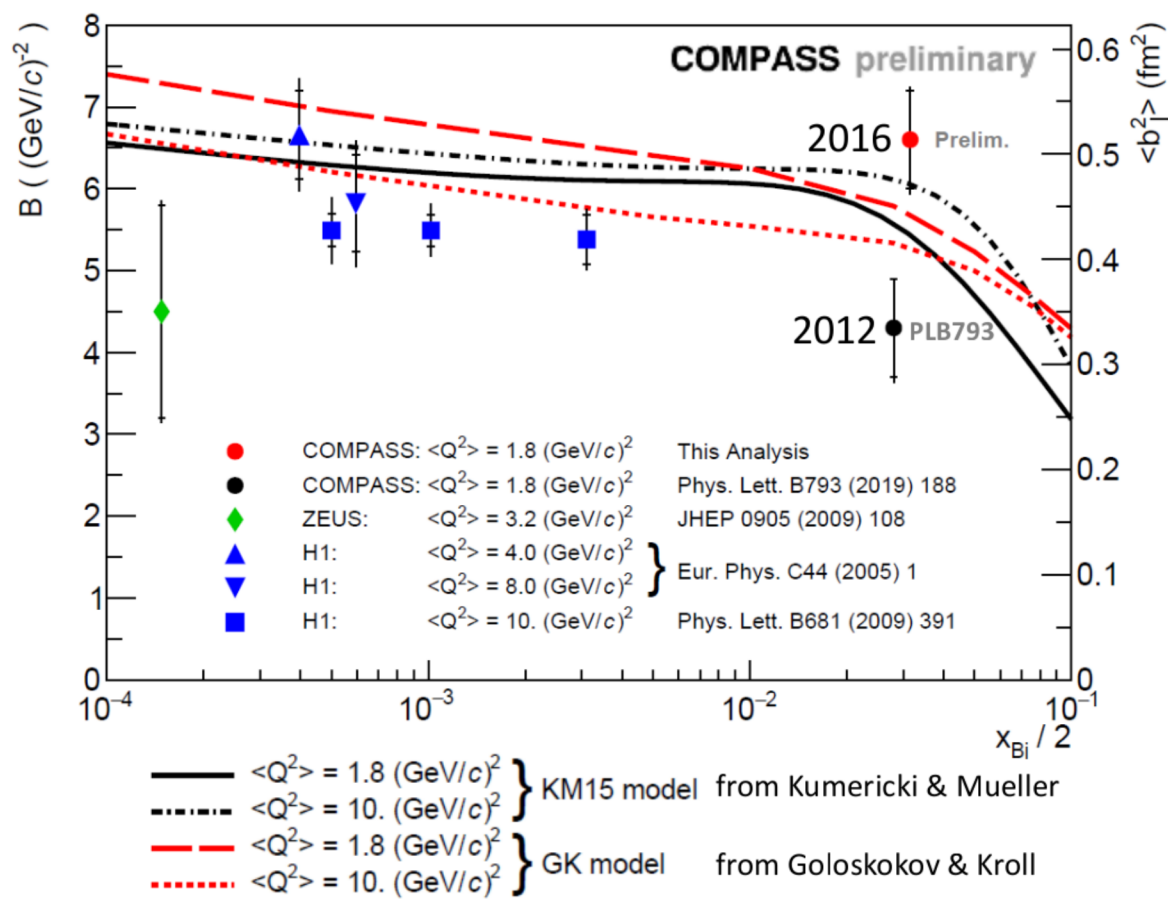
$$x = \xi \approx x_B/2 \text{ close to } 0$$



$$q(x, b_\perp) = \int \frac{d^2\Delta_\perp}{(2\pi)^2} e^{-ib_\perp \cdot \Delta_\perp} H^q(x, 0, -\Delta_\perp^2)$$

$$\langle b_\perp^2 \rangle_x^f = \frac{\int d^2b_\perp b_\perp^2 q_f(x, b_\perp)}{\int d^2b_\perp q_f(x, b_\perp)} = -4 \frac{\partial}{\partial t} \log H^f(x, \xi=0, t) \Big|_{t=0}$$

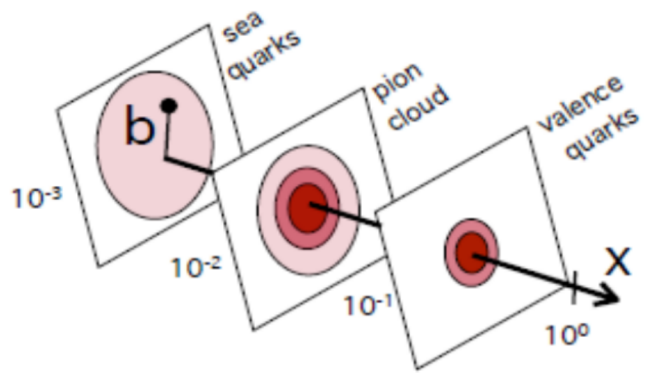
$$\langle b_\perp^2(x) \rangle \approx 2B(\xi)$$



# COMPASS 12-16 Transverse extension of partons in the sea quark range

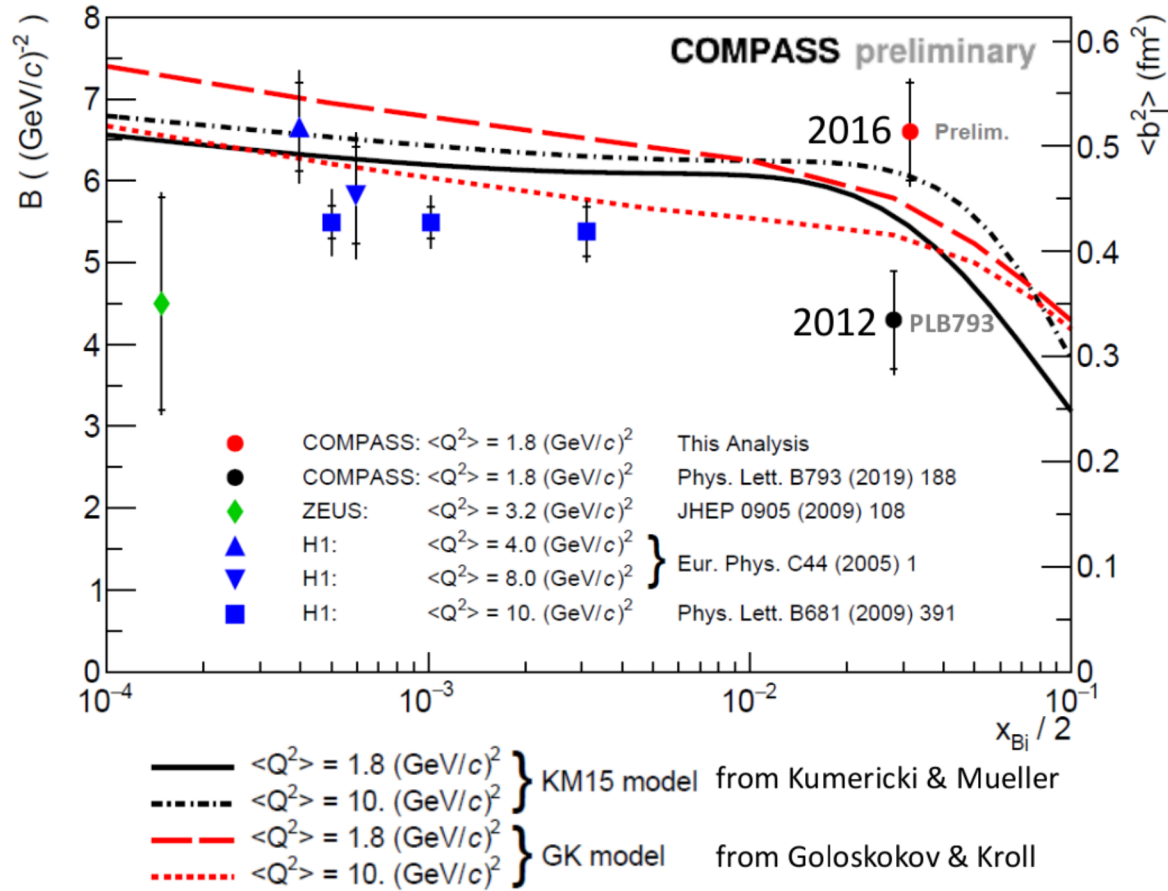
$$d\sigma^{DVCS}/dt = e^{-B|t|} = c_0^{DVCS} \propto (Im\mathcal{H})^2$$

$$\langle b_{\perp}^2(x) \rangle \approx 2B(\xi)$$



- 3 $\sigma$  difference between 2012 and 2016 data
- more advanced analysis with 2016 data
  - $\pi^0$  contamination with different thresholds
  - binning with 3 variables (t, Q<sup>2</sup>, v) or 4 variables (t,  $\phi$ , Q<sup>2</sup>, v)

2012 statistics = Ref  
 2016 analysed statistics = 2.3 × Ref  
 2016+2017 expected statistics = 10 × Ref



## Possible next steps for DVCS

- ✓ DVCS and the sum  $\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$ 
  - $c_0 \sim (\text{Im}\mathcal{H})^2$  final conclusion using all the data sets 2012, 2016, 2017
  - $s_1 \sim \text{Im}\mathcal{H}$   
constrain on  $\text{Im}\mathcal{H}$  and Transverse extension of partons
- ✓ DVCS and the difference  $\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$ 
  - $c_1$  and constrain on  $\text{Re}\mathcal{H}$  (>0 as H1 or <0 as HERMES)  
for D-term and pressure distribution

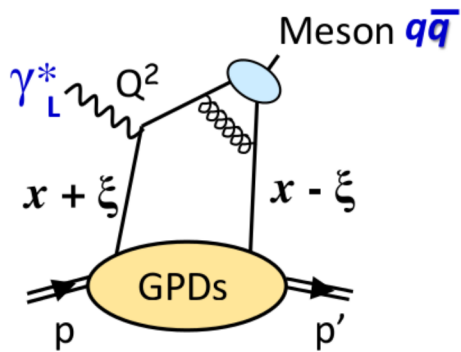
# GPDs and Hard Exclusive Meson Production

Factorisation proven only for  $\sigma_L$

The meson wave function

Is an additional non-perturbative term

Quark contribution



**For Pseudo-Scalar Meson, as  $\pi^0$**

chiral-even GPDs: helicity of parton unchanged

$$\tilde{H}^q(x, \xi, t) \quad \tilde{E}^q(x, \xi, t)$$

+ chiral-odd or transversity GPDs: helicity of parton changed

$$H_T^q(x, \xi, t) \quad (\text{as the transversity TMD})$$

$$\bar{E}_T^q = 2 \tilde{H}_T^q + E_T^q \quad (\text{as the Boer-Mulders TMD})$$

$\sigma_T$  is asymptotically suppressed by  $1/Q^2$  but large contribution observed

GK model:  $k_T$  of  $q$  and  $\bar{q}$  and Sudakov suppression factor are considered

Chiral-odd GPDs with a twist-3 meson wave function

# COMPASS 2012 - 16 Exclusive $\pi^0$ production on unpolarized proton



$F\pi^0 = 2/3 F^u + 1/3 F^d$

$$\frac{d^2\sigma}{dt d\phi_\pi} = \frac{1}{2\pi} \left[ \left( \epsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} \right) + \epsilon \cos 2\phi_\pi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_\pi \frac{d\sigma_{LT}}{dt} \right]$$

$$\frac{d\sigma_L}{dt} \propto |\langle \tilde{H} \rangle|^2 - \frac{t'}{4m^2} |\langle \tilde{E} \rangle|^2$$

$$\frac{d\sigma_T}{dt} \propto |\langle H_T \rangle|^2 - \frac{t'}{8m^2} |\langle \bar{E}_T \rangle|^2$$

$$\frac{\sigma_{TT}}{dt} \propto \frac{t'}{16m^2} |\langle \bar{E}_T \rangle|^2$$

$$\frac{\sigma_{LT}}{dt} \propto \frac{\sqrt{-t'}}{2m} \text{Re} [\langle H_T \rangle^* \langle \tilde{E} \rangle]$$

$$\left\langle \frac{d\sigma_T}{d|t|} + \epsilon \frac{d\sigma_L}{d|t|} \right\rangle = (8.2 \pm 0.9_{\text{stat}} \pm 1.2_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{TT}}{d|t|} \right\rangle = (-6.1 \pm 1.3_{\text{stat}} \pm 0.7_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{LT}}{d|t|} \right\rangle = (1.5 \pm 0.5_{\text{stat}} \pm 0.3_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

COMPASS

$Q^2 = 2.0 \text{ GeV}^2$

$x_B = 0.093$

$|t| \sim 0.26 \text{ GeV}^2$

$\epsilon$  close to 1

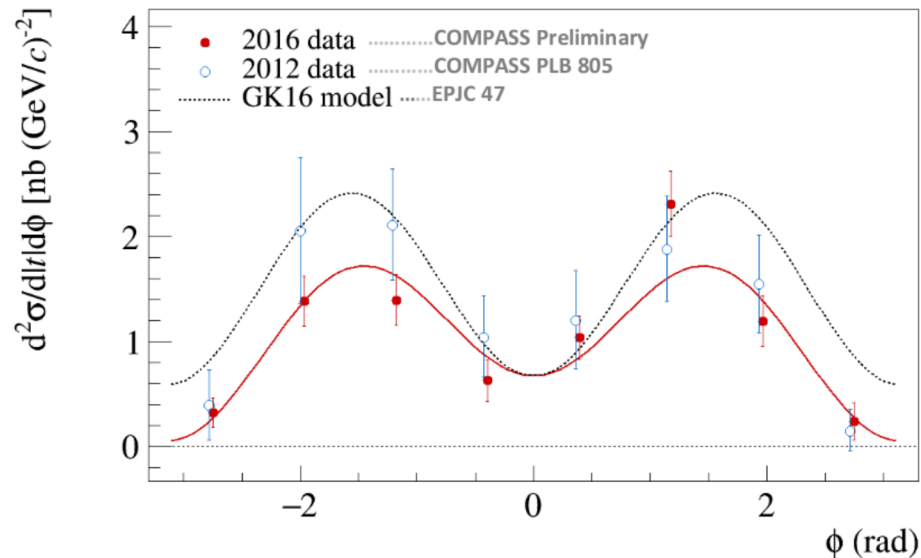
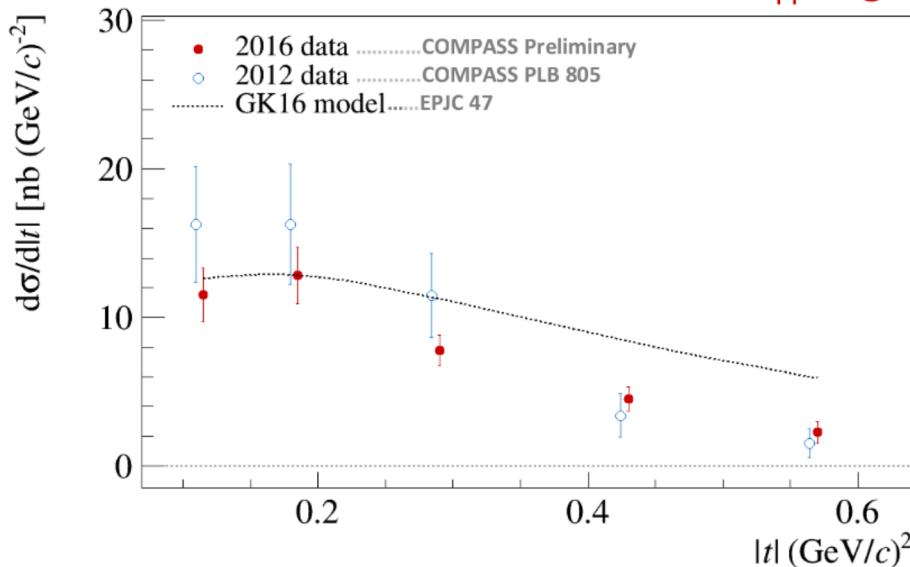
$8.5 < \nu < 26 \text{ GeV}$

$1 < Q^2 < 5 \text{ GeV}^2$

PLB 805 (2020)

$\sigma_{TT}$  large - impact of  $\bar{E}_T$

$\sigma_{LT}$  small but significantly positive as at CLAS



Data: **COMPASS**, PLB 805 (2020) Models: **GK** Kroll Goloskokov EPJC47 (2011) Also **GGL**: Golstein Gonzalez Liuti PRD91 (2015) 15

## Next steps for pi0

Analysis of the 2016 data set should be completed by the end of the month

Extended kinematical domain at small and large  $\nu$  to provide  $x_B$  evolution

$$8.5 < \nu < 26 \text{ GeV}$$

$$6.4 \checkmark$$

$$\checkmark 40 \text{ GeV}$$

The 2017 data set will still increase the statistics



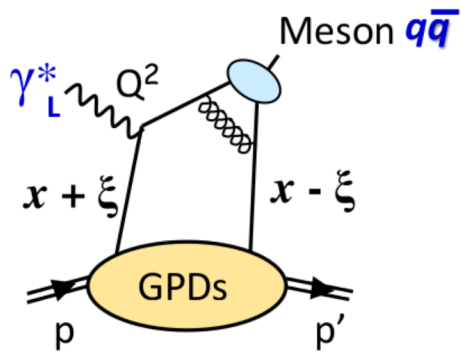
# GPDs and Hard Exclusive Meson Production

**Factorisation proven only for  $\sigma_L$**

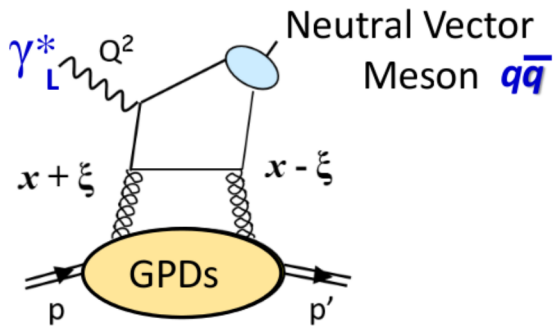
The meson wave function

Is an additional non-perturbative term

**Quark contribution**



**Gluon contribution at the same order in  $\alpha_s$**



**For Vector Meson, as  $\rho, \omega, \phi...$**

chiral-even GPDs: helicity of parton unchanged

$$H^q(x, \xi, t) \quad \mathbf{E}^q(x, \xi, t)$$

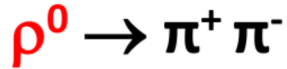
+ chiral-odd or transversity GPDs: helicity of parton changed

$$H_T^q(x, \xi, t) \quad (\text{as the transversity TMD})$$

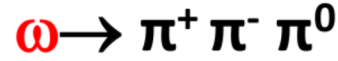
$$\overline{E}_T^q = 2 \tilde{H}_T^q + E_T^q \quad (\text{as the Boer-Mulders TMD})$$

# HEMP with Transversely Polarized Target without RPD

Gparity:  $G(\pi)=-1$ ;  $G(\rho)=+1$ ;  $G(\omega)=-1$



$$E_{\rho^0} = \frac{1}{\sqrt{2}} \left( \frac{2}{3} E^u \oplus \frac{1}{3} E^d + \frac{3}{4} \frac{E_g}{x} \right)$$

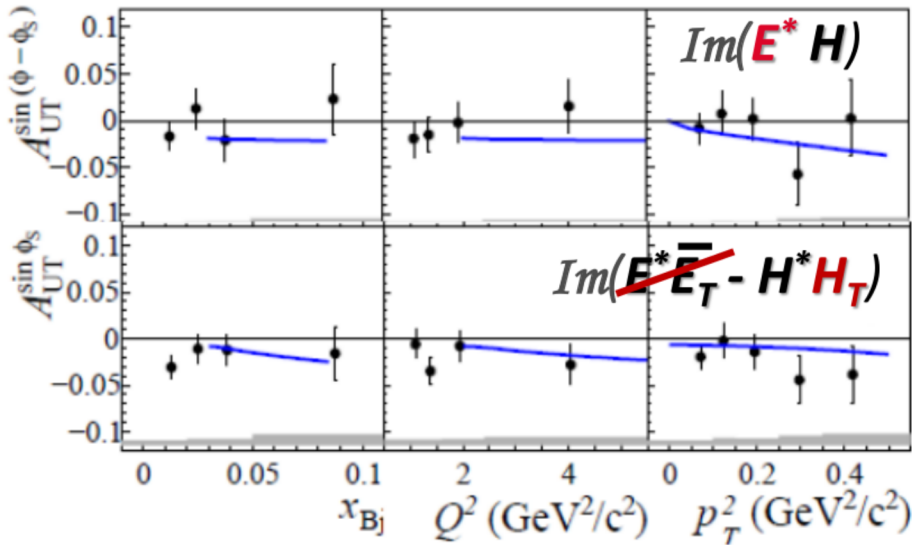


$$E_{\omega} = \frac{1}{\sqrt{2}} \left( \frac{2}{3} E^u \ominus \frac{1}{3} E^d + \frac{1}{4} \frac{E_g}{x} \right)$$

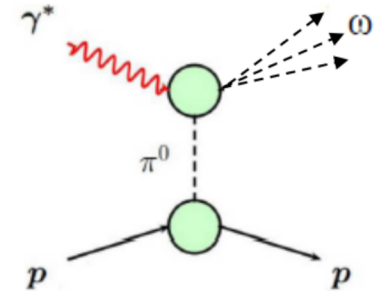
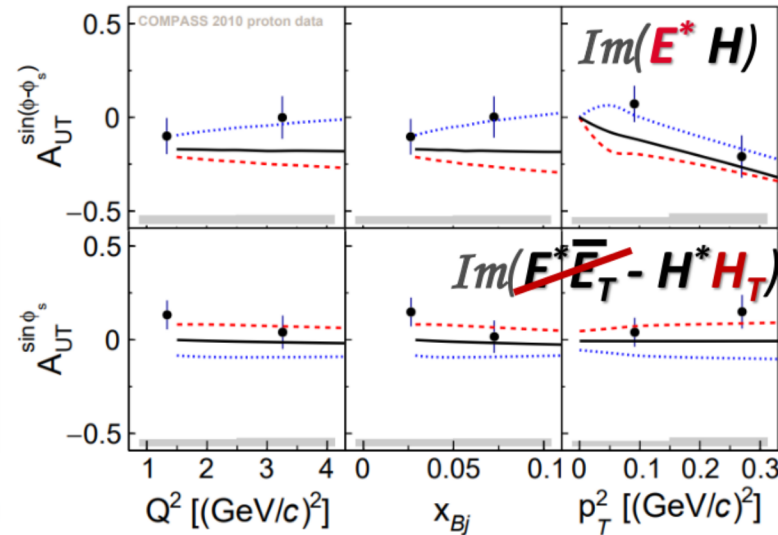
$E^u$  and  $E^d$  of opposite sign  
 $\omega$  is more promising  
(see the larger scale)  
but there is the inherent  
pion pole contribution

$\Gamma(\omega \rightarrow \pi^0 \gamma) = 9 \times \Gamma(\rho^0 \rightarrow \pi^0 \gamma)$   
Same for  $\pi\omega$  FF but sign unknown

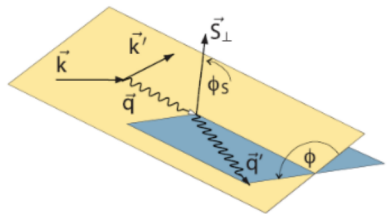
COMPASS, NPB 865 (2012) 1-20, PLB731 (2014) 19



COMPASS, NPB 915 (2017)



- ▶ positive  $\pi\omega$  form factor
- ▶ no pion pole
- ▶ negative  $\pi\omega$  form factor



GK EPJC42,50,53,59,65,74

# exclusive VM production with Unpolarised Target and SDME

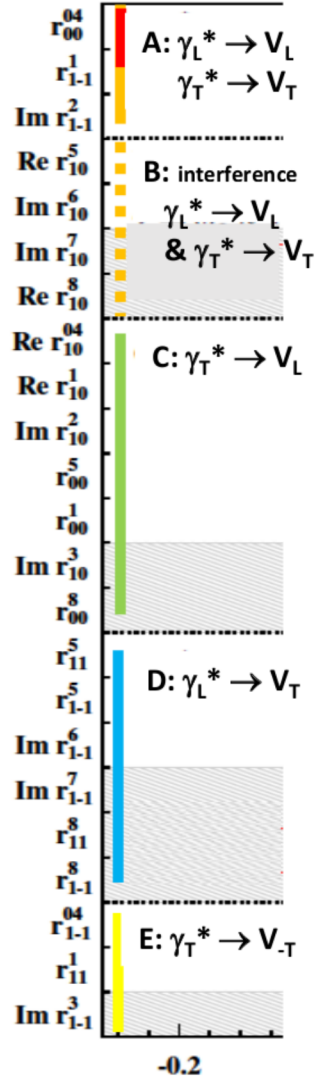
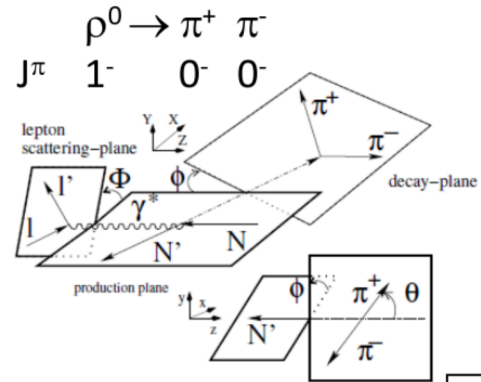
experimental angular distributions:

$$\mathcal{W}^{U+L}(\Phi, \phi, \cos \Theta) = \mathcal{W}^U(\Phi, \phi, \cos \Theta) + P_b \mathcal{W}^L(\Phi, \phi, \cos \Theta)$$

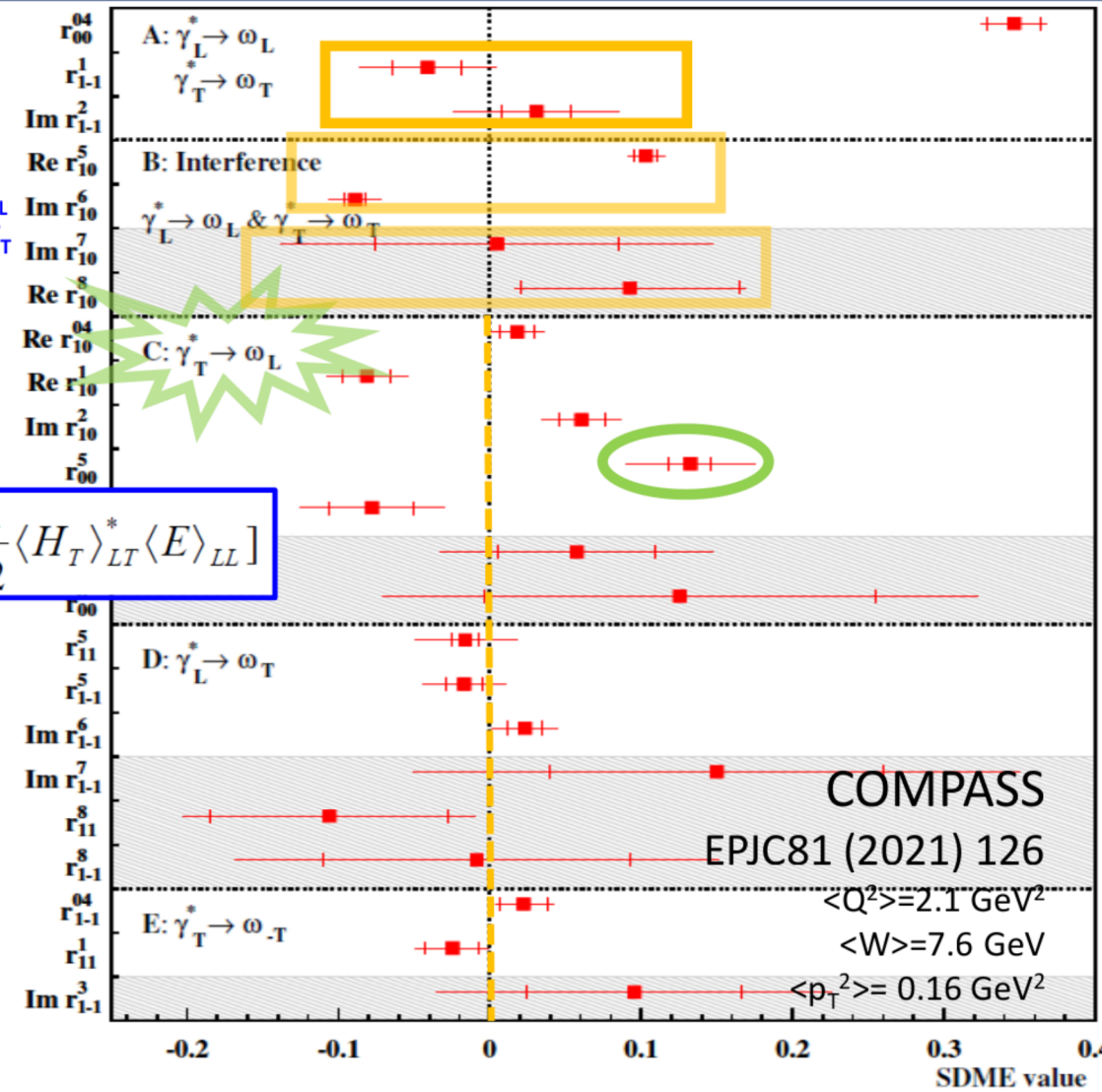
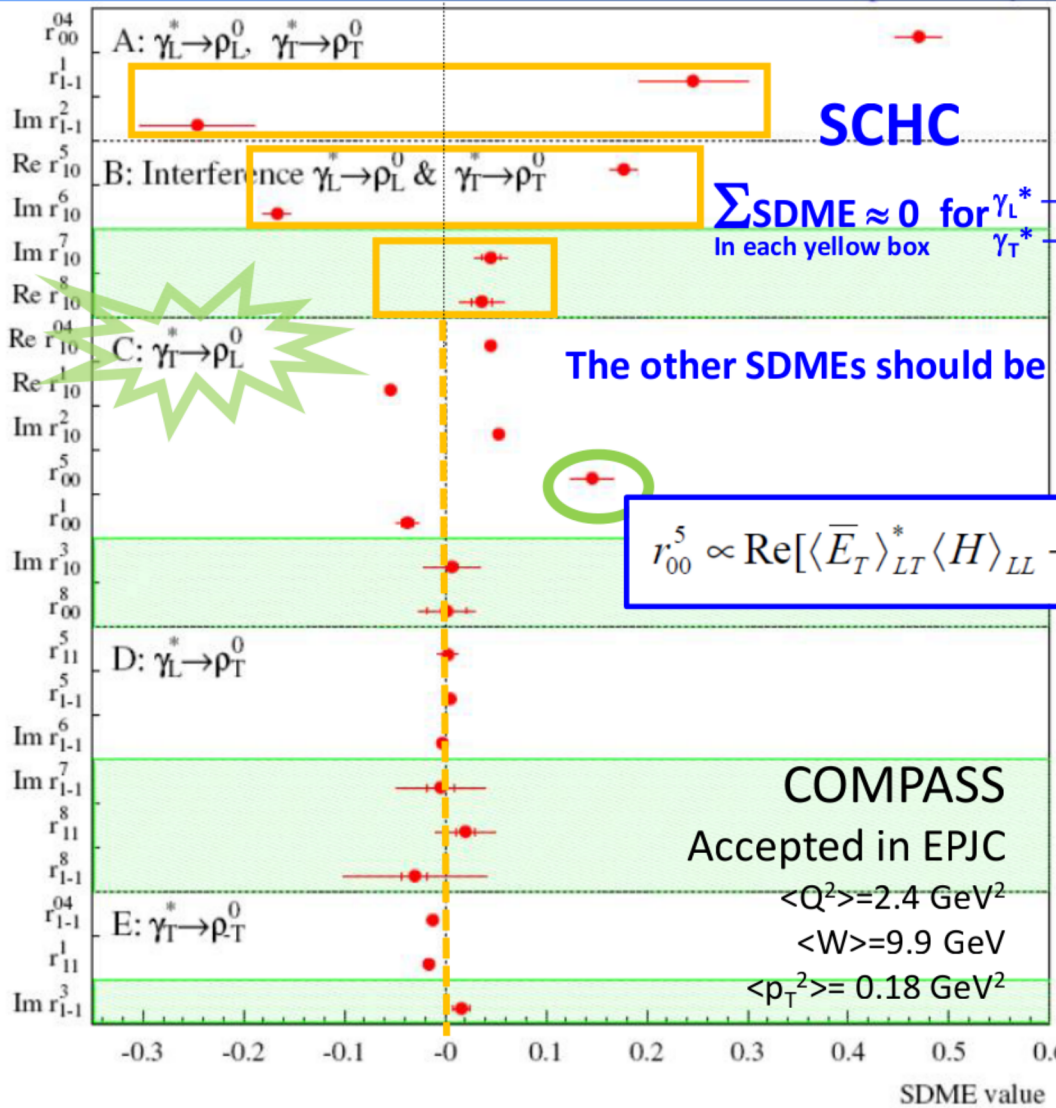
15 'unpolarized' and 8 'polarized' SDMEs

$$\begin{aligned} \mathcal{W}^U(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[ \frac{1}{2}(1 - r_{00}^{04}) + \frac{1}{2}(3r_{00}^{04} - 1) \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^{04}\} \sin 2\Theta \cos \phi - r_{1-1}^{04} \sin^2 \Theta \cos 2\phi \right. \\ & - \epsilon \cos 2\Phi \left( r_{11}^1 \sin^2 \Theta + r_{00}^1 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^1\} \sin 2\Theta \cos \phi - r_{1-1}^1 \sin^2 \Theta \cos 2\phi \right) \\ & \left. - \epsilon \sin 2\Phi \left( \sqrt{2}\text{Im}\{r_{10}^2\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^2\} \sin^2 \Theta \sin 2\phi \right) \right. \\ & + \sqrt{2\epsilon(1+\epsilon)} \cos \Phi \left( r_{11}^5 \sin^2 \Theta + r_{00}^5 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^5\} \sin 2\Theta \cos \phi - r_{1-1}^5 \sin^2 \Theta \cos 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \sin \Phi \left( \sqrt{2}\text{Im}\{r_{10}^6\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^6\} \sin^2 \Theta \sin 2\phi \right) \right], \\ \mathcal{W}^L(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[ \sqrt{1-\epsilon^2} \left( \sqrt{2}\text{Im}\{r_{10}^3\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^3\} \sin^2 \Theta \sin 2\phi \right) \right. \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos \Phi \left( \sqrt{2}\text{Im}\{r_{10}^7\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^7\} \sin^2 \Theta \sin 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \sin \Phi \left( r_{11}^8 \sin^2 \Theta + r_{00}^8 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^8\} \sin 2\Theta \cos \phi - r_{1-1}^8 \sin^2 \Theta \cos 2\phi \right) \right] \end{aligned}$$

$\epsilon$  close to 1,  
small  $\mathcal{W}^L$   
no L/T separation

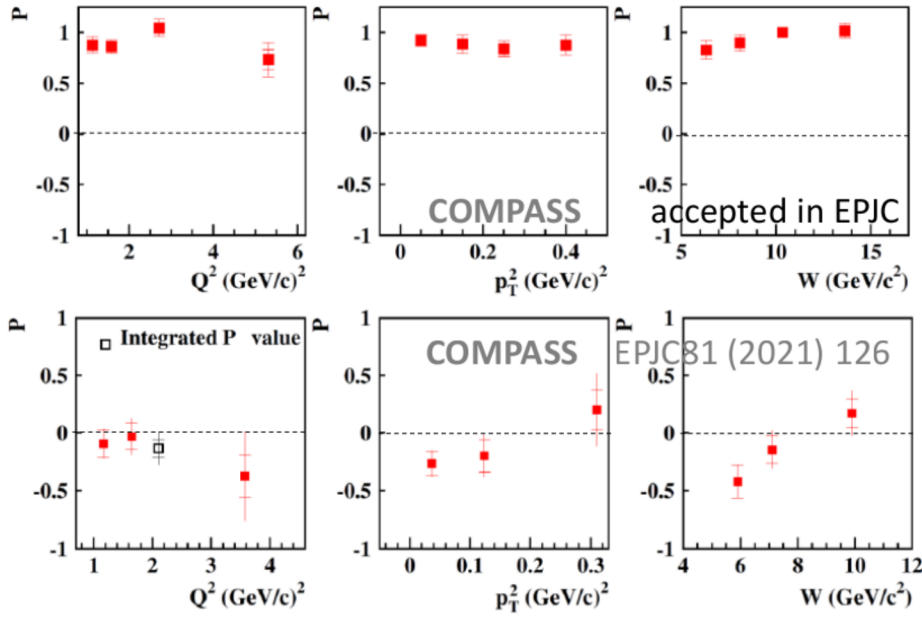


# COMPASS 2012 Exclusive $\rho^0$ and $\omega$ production on unpolarized proton



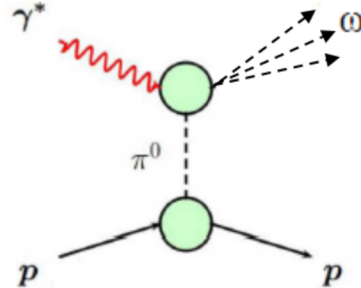
Natural (N) to Unnatural (U)  
Parity Exchange for  $\gamma_T^* \rightarrow V_T$

$$P = \frac{2r_{1-1}^1}{1 - r_{00}^{04} - 2r_{1-1}^{04}} \approx \frac{d\sigma_T^N(\gamma_T^* \rightarrow V_T) - d\sigma_T^U(\gamma_T^* \rightarrow V_T)}{d\sigma_T^N(\gamma_T^* \rightarrow V_T) + d\sigma_T^U(\gamma_T^* \rightarrow V_T)}$$

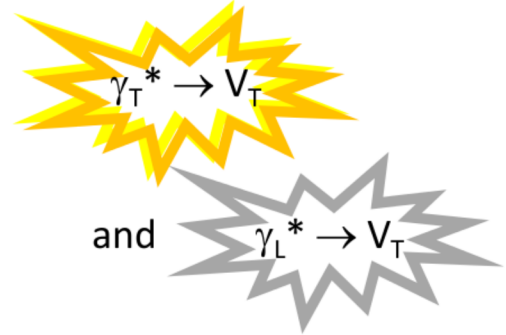


The pion pole exchange (UPE) is large for  $\omega$  compared to  $\rho^0$

$$\Gamma(\omega \rightarrow \pi^0 \gamma) = 9 \times \Gamma(\rho^0 \rightarrow \pi^0 \gamma) \text{ as for } \pi^0 \text{ Vector Meson FF}$$



It plays an important role in  $\omega$  production for:



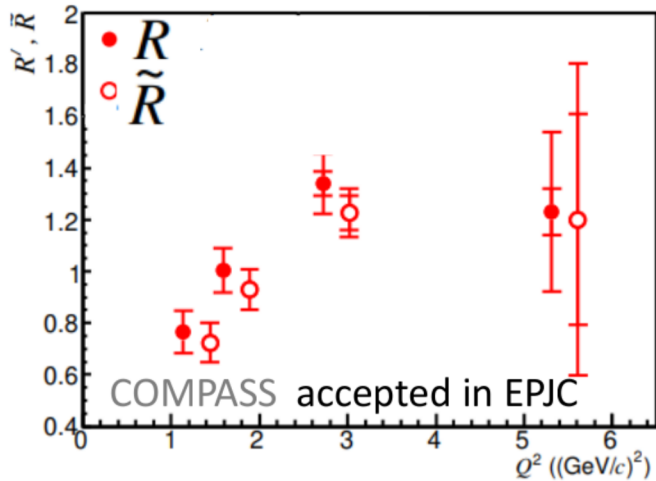
$\rho^0$ :  $P \sim 1 \rightarrow$  NPE dominance  $P \sim 1$   
NPE with GPDs  $H, E$

$\omega$ :  $P \sim 0 \rightarrow$  NPE  $\sim$  UPE  
UPE dominance at small  $W$  and  $p_T^2$   
UPE with GPDs  $\tilde{H}, \tilde{E}$  and the dominant pion pole

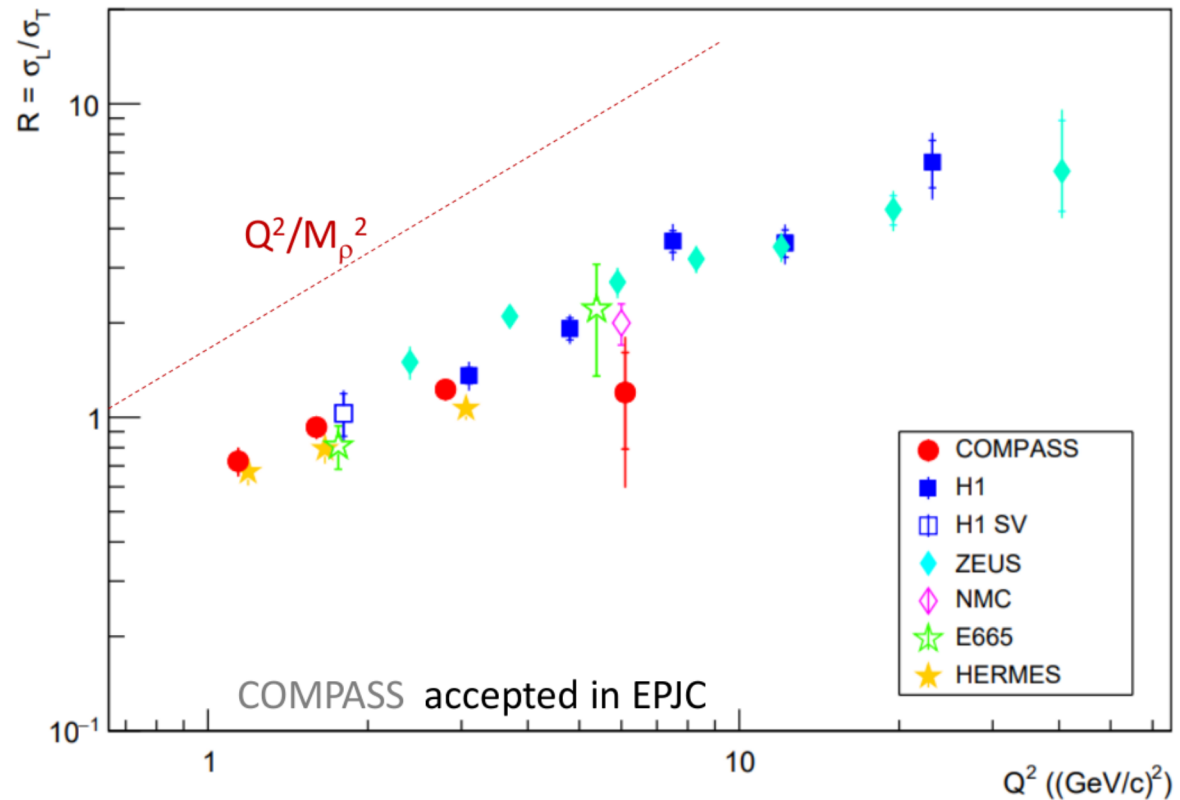
$$R = \frac{\sigma_L(\gamma_L^* \rightarrow V)}{\sigma_T(\gamma_T^* \rightarrow V)}$$

$$R = \frac{1}{\epsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}} \quad \text{only if SCHC}$$

In COMPASS domain evaluation of  $R$  and  $\tilde{R}$  considering violation of SCHC (and only NPE)



for all the experiments with  $Q^2 > 1 \text{ GeV}^2$



Deviation from the pQCD LO prediction in  $Q^2/M_\rho^2$ : QCD evolution and  $q_T$  Transverse size effects of the meson smaller for  $\sigma_L$  than for  $\sigma_T$

## Next steps for Vector Mesons

Analysis of the exclusive  $\phi$  production is currently in progress

(with cross section and SDMEs)

✓ **DVCS** and the sum  $\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$   
 →  $c_0$  and  $s_1$  and constrain on  $\text{Im}\mathcal{H}$  and Transverse extension of partons

✓ **DVCS** and the difference  $\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$   
 →  $c_1$  and constrain on  $\text{Re}\mathcal{H}$  (>0 as H1 or <0 as HERMES)  
 for D-term and pressure distribution

Importance of  $e^+$  beam  
 For Jlab 20+ GeV

✓ On-going analysis (Cross section, SDME) for HEMP of  $\pi^0$ ,  $\rho^0$ ,  $\omega$ ,  $\phi$ ,  $J/\psi$

- ✓ Transversity GPDs
- ✓ Gluon GPDs
- ✓ Flavor decomposition

Importance of large luminosity  
 For DVCS, TCS, DDVCS,  $J/\psi$  ...





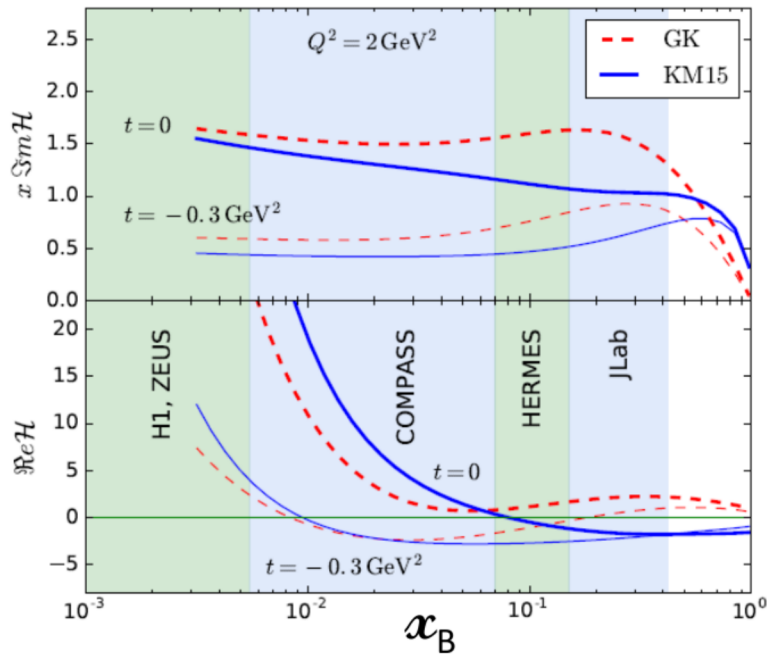


# ImH and ReH using global fits of the world data

## Global Fit KM15

Compared to GK Model GK

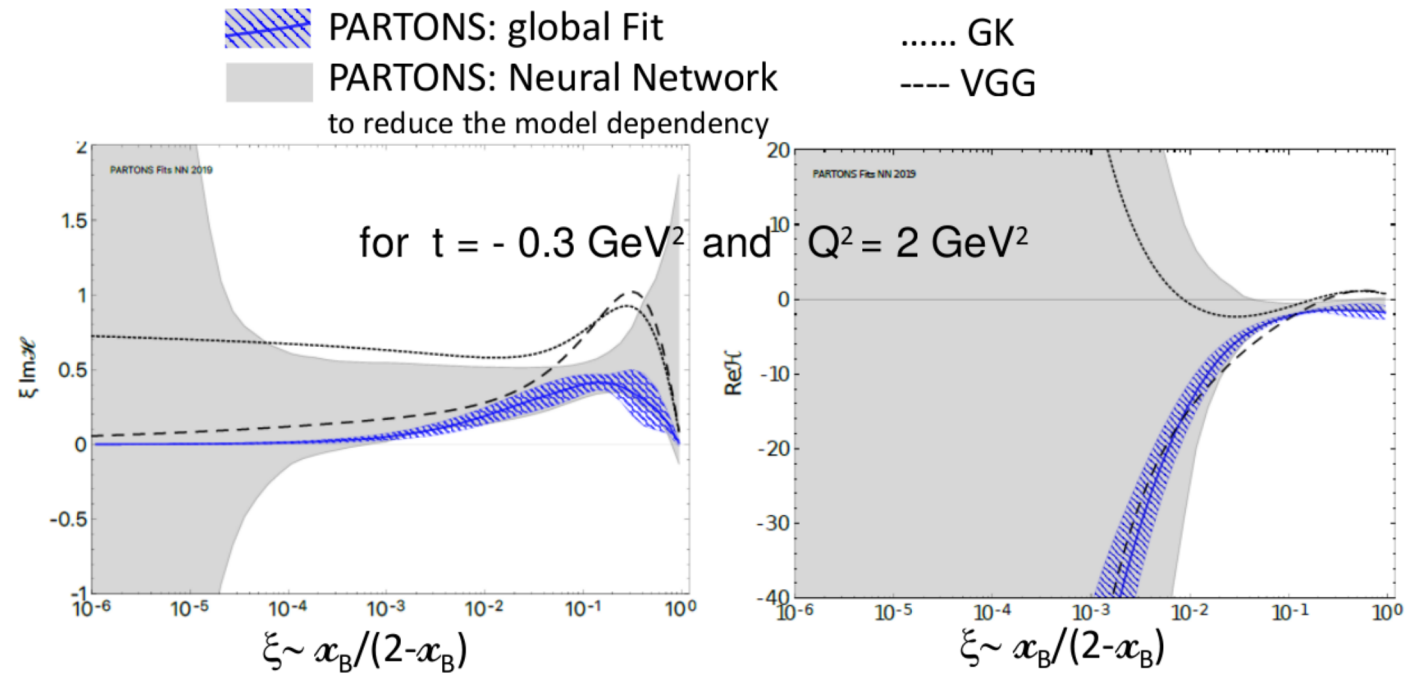
Kumericki, Mueller, NPB (2010) 841, private com.



## Global Fits using PARTONS framework

Compared to GK and VGG Models

Moutarde, Sznajder, Wagner, Eur. Phys. J. C 79 (2019) 7, 614



**ReH** is still poorly known (importance of DVCS with  $\mu^\pm$  at COMPASS,  $e^\pm$  at JLab or TCS at JLab and EIC)