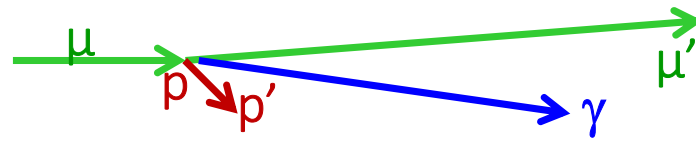


Hard Exclusive Reactions at COMPASS at CERN

Exclusive photon (DVCS) and meson (HEMP) production at small transfer for GPD studies



$$\text{DVCS : } \mu \text{ p} \rightarrow \mu' \text{ p}' \gamma$$



$$\text{Pseudo-Scalar Meson : } \mu \text{ p} \rightarrow \mu' \text{ p}' \pi^0$$

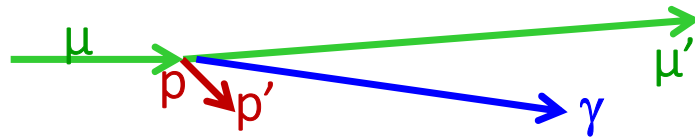
$$\text{Vector Meson : } \mu \text{ p} \rightarrow \mu' \text{ p}' \rho \text{ or } \omega \text{ or } \dots$$

Nicole d'Hose - CEA Université Paris-Saclay for the COMPASS Collaboration

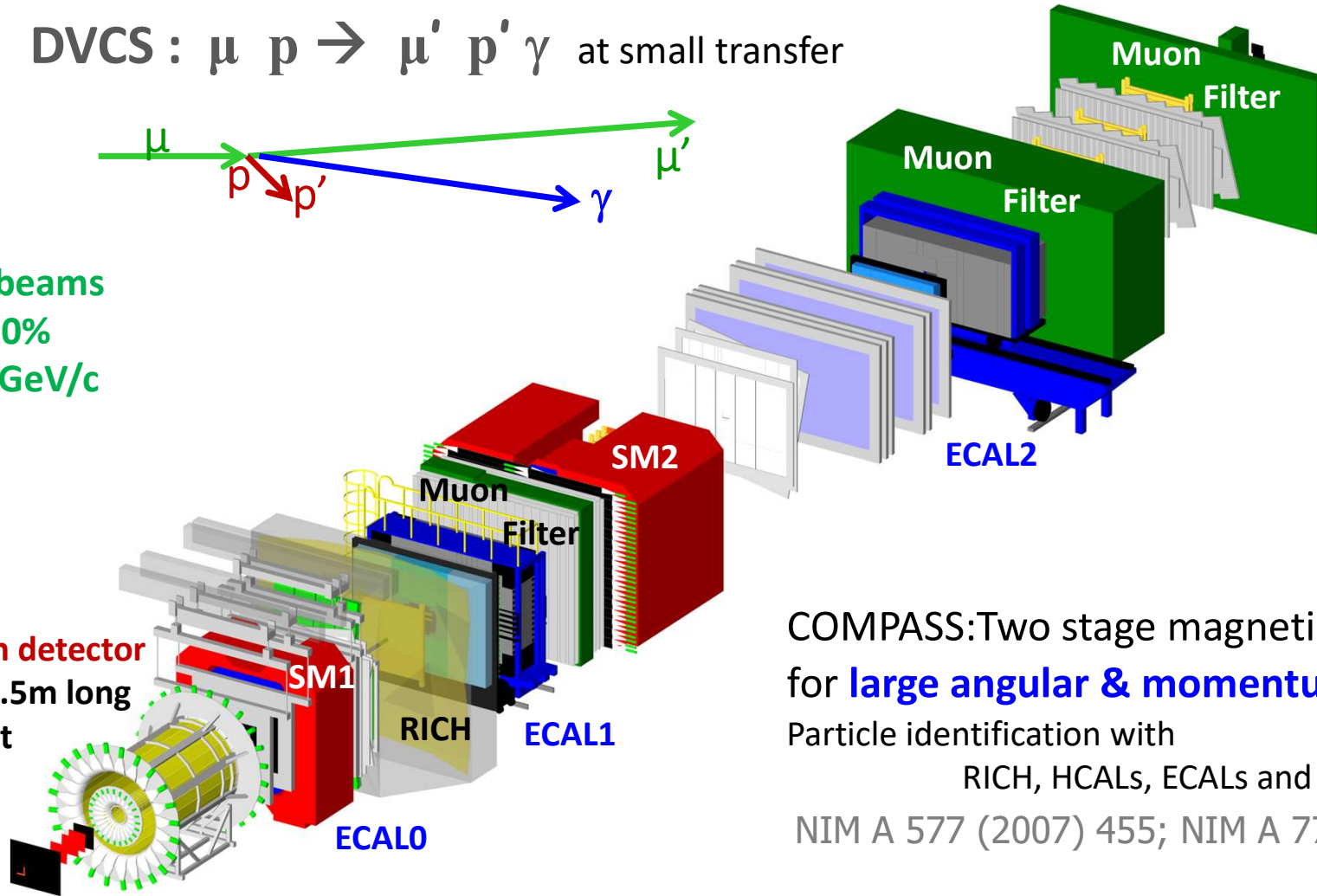


Measurement of exclusive cross sections at COMPASS

DVCS : $\mu p \rightarrow \mu' p' \gamma$ at small transfer



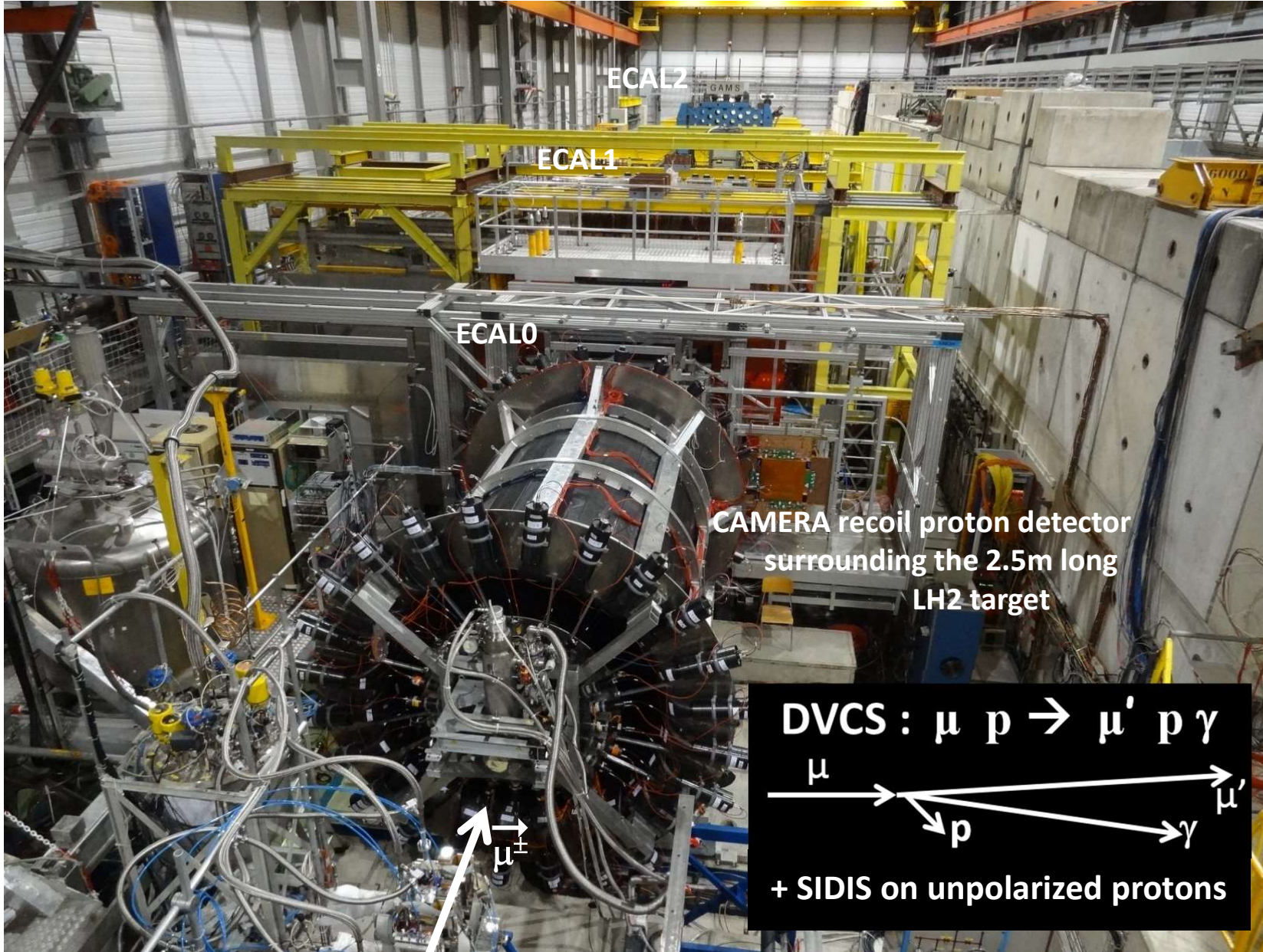
Both μ^+ and μ^- beams
Polarisation $\sim \pm 80\%$
Momentum 160 GeV/c



COMPASS: Two stage magnetic spectrometer for **large angular & momentum acceptance**

Particle identification with RICH, HCALs, ECALs and muon filters

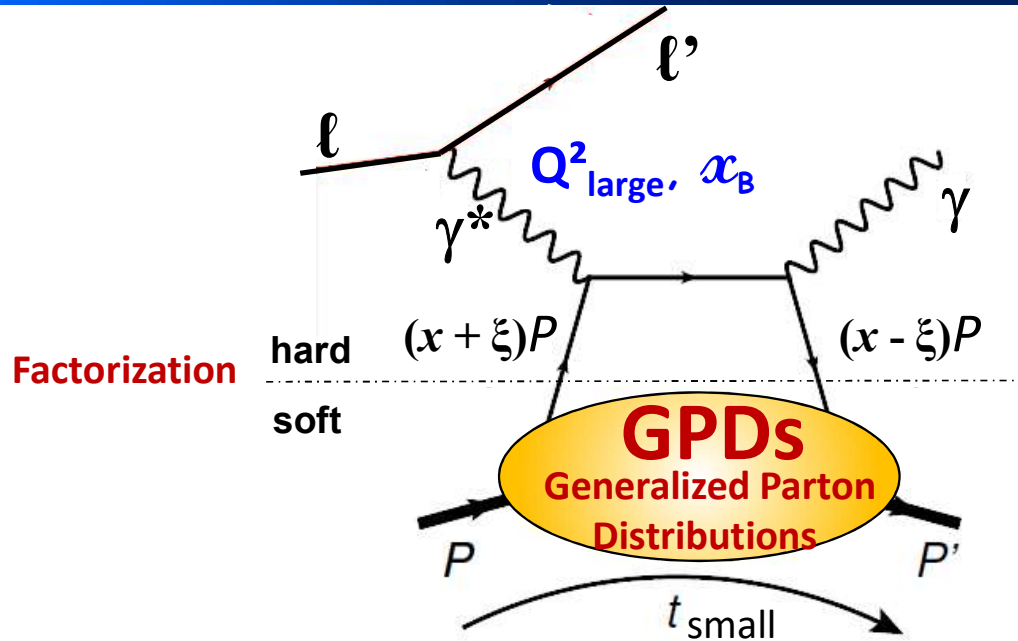
NIM A 577 (2007) 455; NIM A 779 (2015) 69



2012:
1 month pilot run

2016 -17:
2 x 6 month
data taking

Deeply virtual Compton scattering (DVCS)



D. Mueller *et al*, Fortsch. Phys. 42 (1994)

X.D. Ji, PRL 78 (1997), PRD 55 (1997)

A. V. Radyushkin, PLB 385 (1996), PRD 56 (1997)

DVCS: $\ell p \rightarrow \ell' p' \gamma$

the golden channel

because it interferes with
the Bethe-Heitler process

also meson production

$\ell p \rightarrow \ell' p' \pi, \rho, \omega$ or ϕ or J/ψ ...

The GPDs depend on the following variables:

x : average } quark longitudinal
 ξ : transferred } momentum fraction

t : proton momentum transfer squared
related to b_{\perp} via Fourier transform

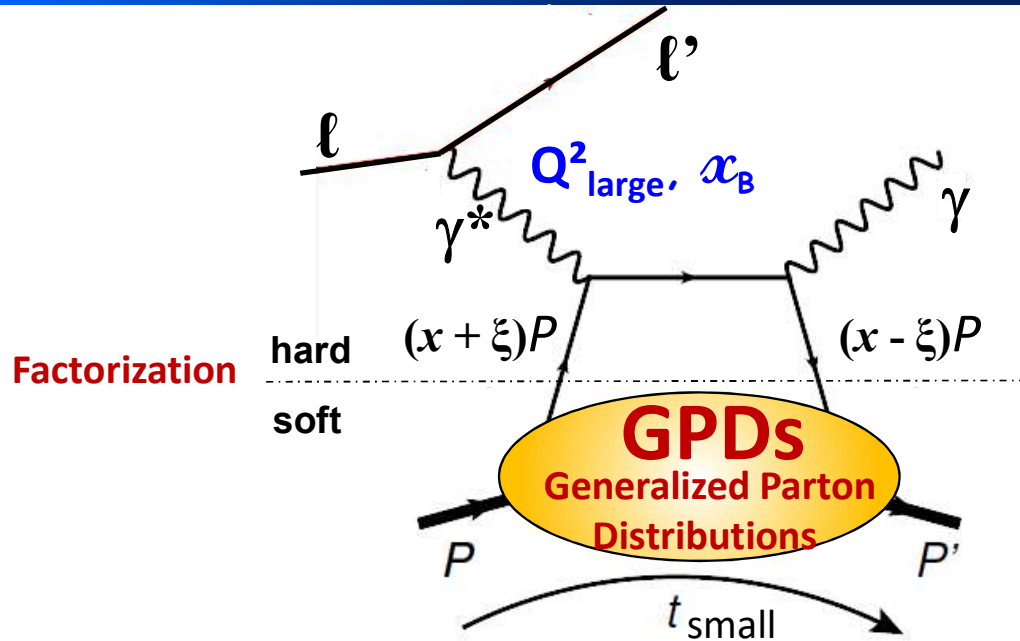
Q^2 : virtuality of the virtual photon

The variables measured in the experiment:

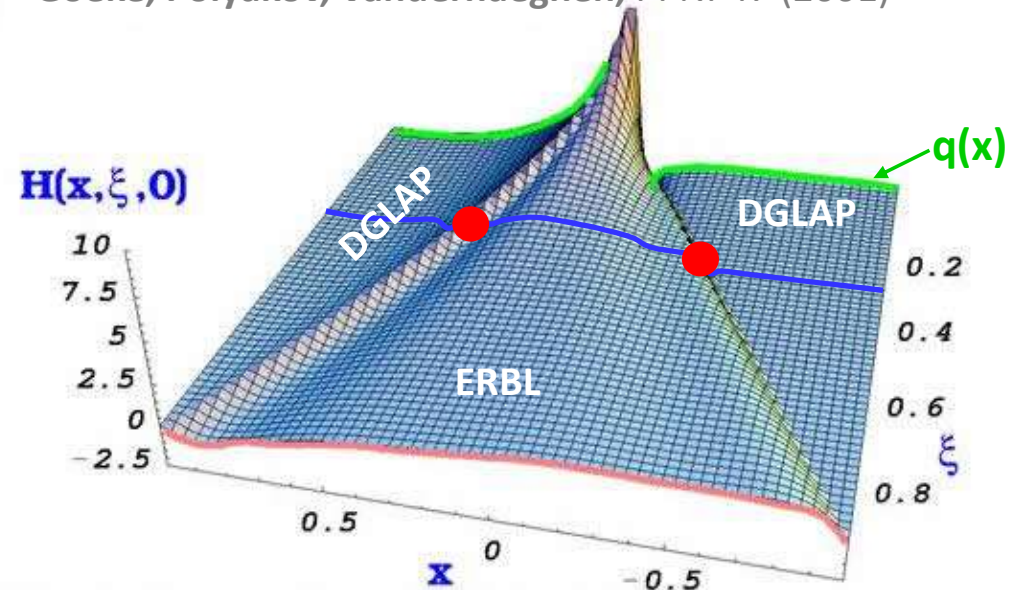
$E_{\ell}, Q^2, x_B \sim 2\xi / (1 + \xi),$

t (or $\theta_{\gamma^* \gamma}$) and ϕ ($\ell \ell'$ plane / $\gamma \gamma^*$ plane)

Deeply virtual Compton scattering (DVCS)



Goeke, Polyakov, Vanderhaeghen, PPNP47 (2001)



The amplitude DVCS at LT & LO in α_s (GPD \mathbf{H}):

Real part

Imaginary part

$$\mathcal{H} = \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi + i\varepsilon} = \mathcal{P} \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi} - i\pi H(x = \pm \xi, \xi, t)$$

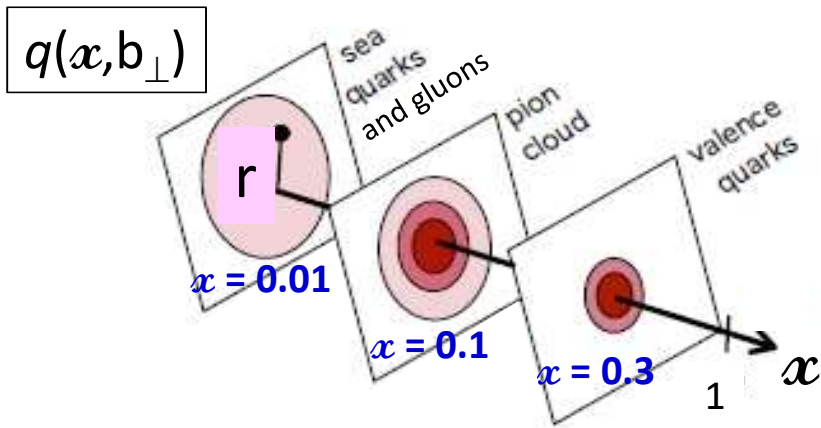
In an experiment we measure

Compton Form Factor \mathcal{H}

Deeply virtual Compton scattering (DVCS)

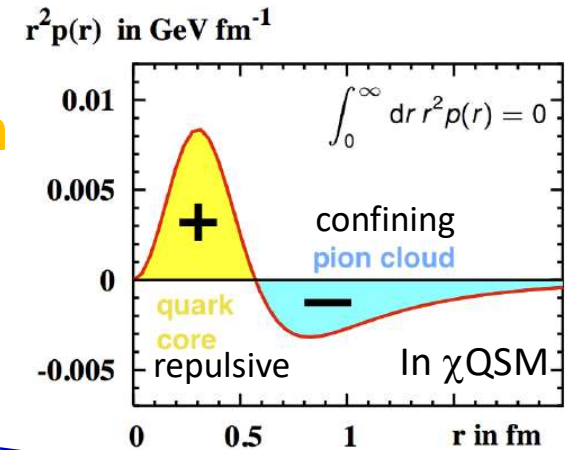
M. Burkardt, PRD66(2002)

Mapping in the transverse plane



M. Polyakov, P. Schweitzer, Int.J.Mod.Phys. A33 (2018)

Pressure Distribution



FT of $H(x, \xi=0, t)$

The amplitude DVCS at LT & LO in α_s (GPD \mathbf{H}):

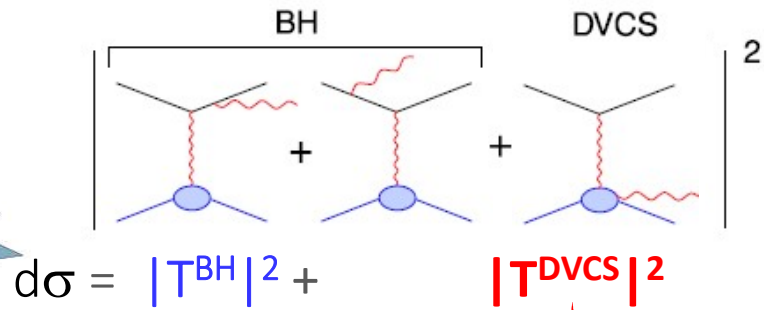
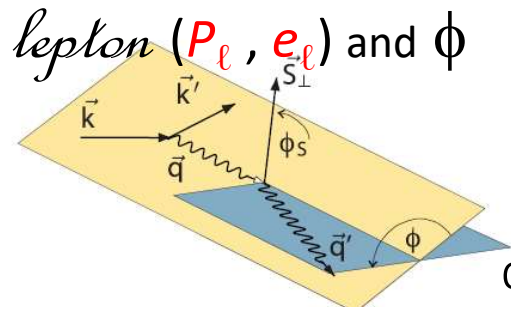
$$\mathcal{H} = \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi + i\varepsilon} = \mathcal{P} \int_{-1}^{+1} dx \frac{\mathbf{H}(x, \xi, t)}{x - \xi} - i\pi \mathbf{H}(x = \pm \xi, \xi, t)$$

In an experiment we measure Compton Form Factor \mathcal{H}

$$\text{Re}\mathcal{H}(\xi, t) = \pi^{-1} \int_0^1 dx \frac{2x \text{Im}\mathcal{H}(x, t)}{x^2 - \xi^2} + \Delta(t)$$

$d_1(t)$
D-term

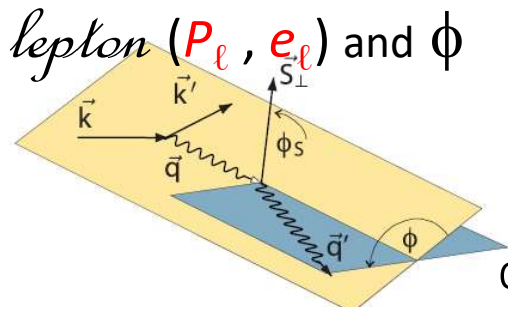
Deeply virtual Compton scattering (DVCS)



$$d\sigma = |T^{BH}|^2 + |T^{DVCS}|^2 + \text{Interference Term}$$

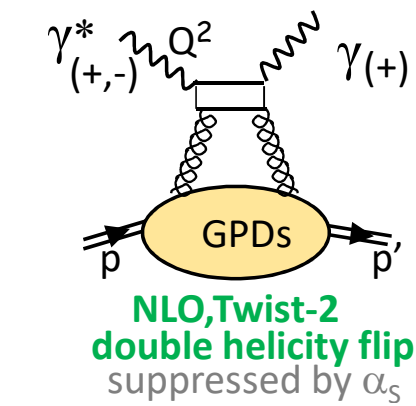
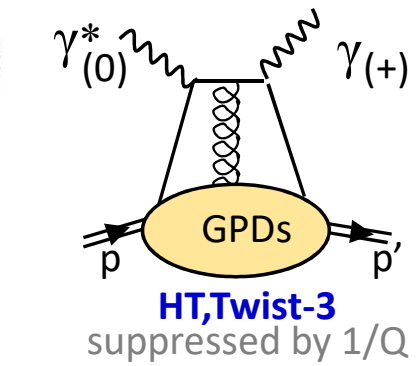
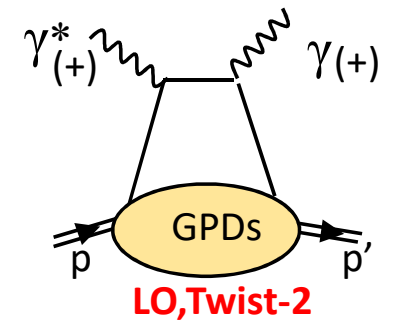
$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$

Deeply virtual Compton scattering (DVCS)



$$d\sigma = \left[\underbrace{\text{BH}}_{|\mathbb{T}^{BH}|^2} + \underbrace{\text{DVCS}}_{|\mathbb{T}^{DVCS}|^2} \right]^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$



With unpolarized target:

Belitsky, Müller, Kirner, NPB629 (2002)

$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$$

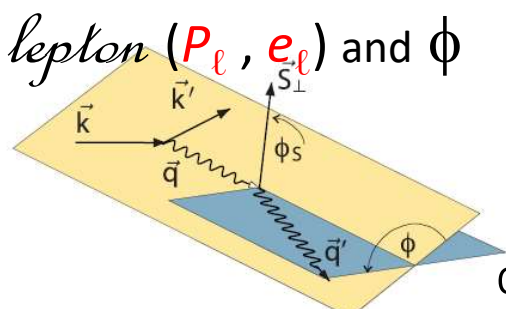
$$d\sigma_{unpol}^{DVCS} \propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$$

$$d\sigma_{pol}^{DVCS} \propto s_1^{DVCS} \sin \phi$$

$$\text{Re } I \propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$$

$$\text{Im } I \propto s_1^I \sin \phi + s_2^I \sin 2\phi$$

Deeply virtual Compton scattering (DVCS)



$$d\sigma = \left| \overbrace{\text{BH}} + \overbrace{\text{DVCS}} \right|^2 + \text{Interference Term}$$

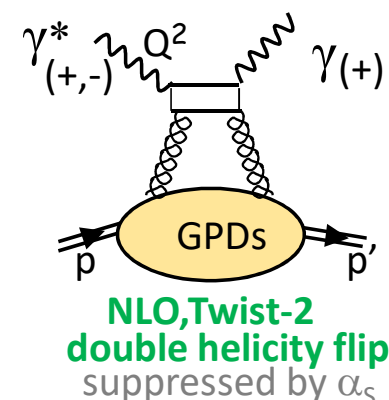
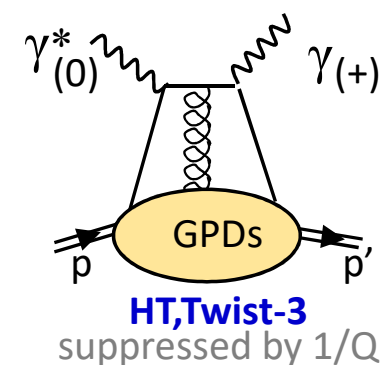
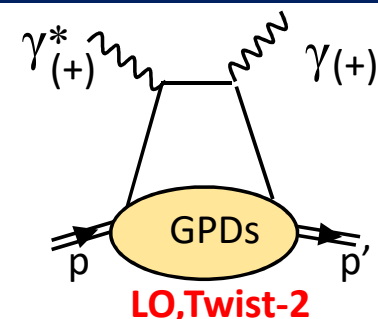
$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \underbrace{\left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right)}_{|T^{DVCS}|^2} + \underbrace{\left(e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I \right)}_{\text{Interference Term}}$$

With both μ^{\uparrow} and μ^{\rightarrow} beams we can build:

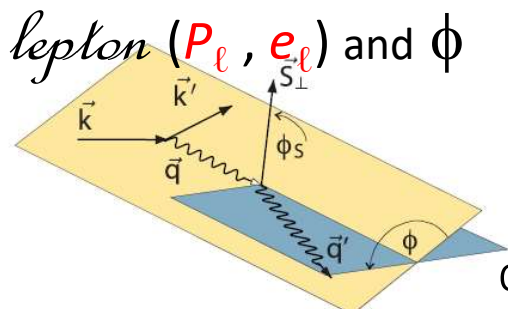
① beam charge-spin sum

$$\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$$

$$\begin{aligned} d\sigma^{BH} &\propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi \\ d\sigma_{unpol}^{DVCS} &\propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi \\ d\sigma_{pol}^{DVCS} &\propto s_1^{DVCS} \sin \phi \\ \text{Re } I &\propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi \\ \text{Im } I &\propto s_1^I \sin \phi + s_2^I \sin 2\phi \end{aligned}$$

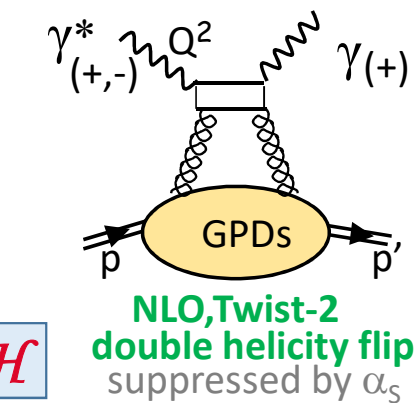
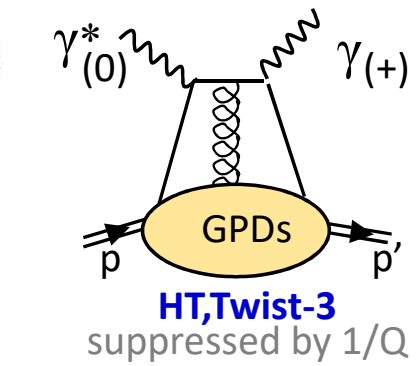
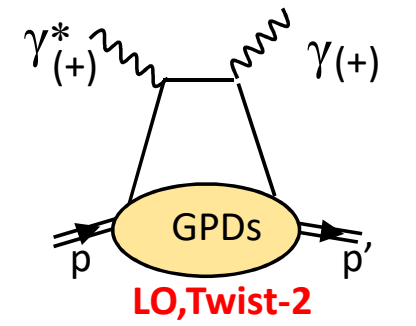


Deeply virtual Compton scattering (DVCS)



$$d\sigma = \left| \mathcal{T}^{BH} \right|^2 + \left| \mathcal{T}^{DVCS} \right|^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(\ell p \rightarrow \ell p \gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{BH}}_{\text{Well known}} + \left(d\sigma_{unpol}^{DVCS} + P_\ell d\sigma_{pol}^{DVCS} \right) + (e_\ell \text{Re } I + e_\ell P_\ell \text{Im } I)$$



With both μ^+ and μ^- beams we can build:

① beam charge-spin sum

$$\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$$

② difference

$$\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$$

$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$$

$$d\sigma_{unpol}^{DVCS} \propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$$

$$d\sigma_{pol}^{DVCS} \propto s_1^{DVCS} \sin \phi$$

$$\text{Re } I \propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$$

$$\text{Im } I \propto s_1^I \sin \phi + s_2^I \sin 2\phi$$

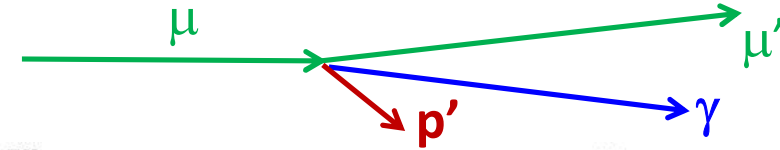
$$s_1^I \propto \text{Im } \mathcal{F}$$

$$c_1^I \propto \text{Re } \mathcal{F}$$

$$\mathcal{F} = F_1 \mathcal{H} + \xi (F_1 + F_2) \tilde{\mathcal{H}} - t/4m^2 F_2 \mathcal{E} \xrightarrow{\text{for proton}} F_1 \mathcal{H}$$

COMPASS 2016 data Selection of exclusive single photon production

Comparison between the observables given by the spectro or by CAMERA



DVCS : $\mu p \rightarrow \mu' p \gamma$

1) $\Delta\varphi = \varphi^{\text{cam}} - \varphi^{\text{spec}}$

2) $\Delta p_T = p_T^{\text{cam}} - p_T^{\text{spec}}$

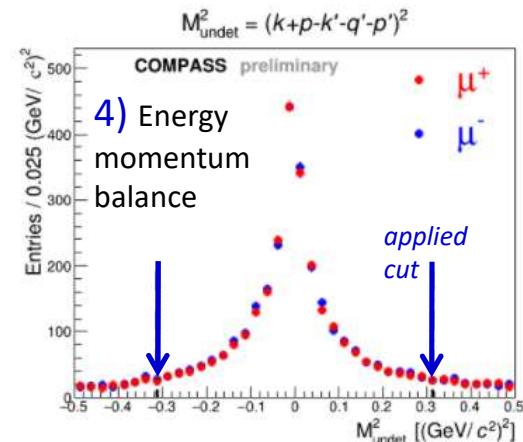
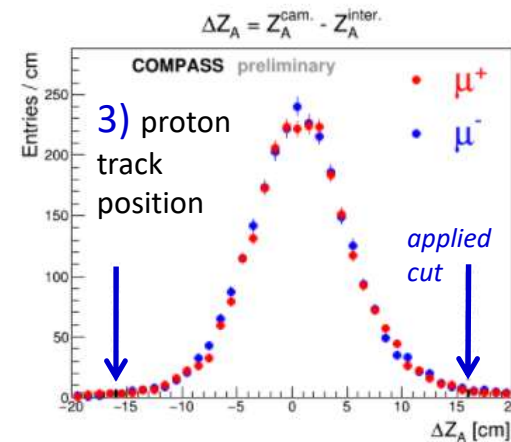
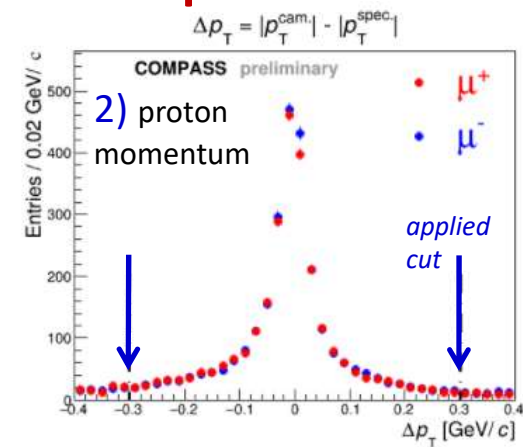
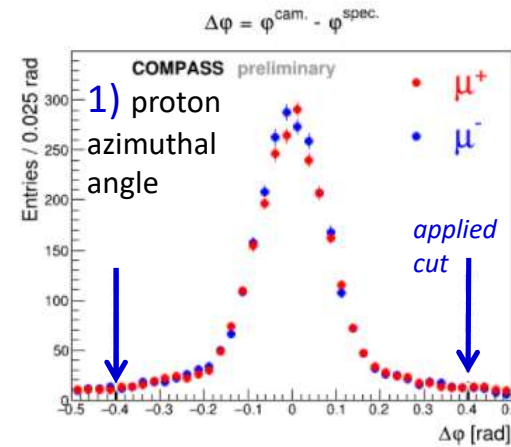
3) $\Delta z_A = z_A^{\text{cam}} - z_A^{\text{ZB and vertex}}$

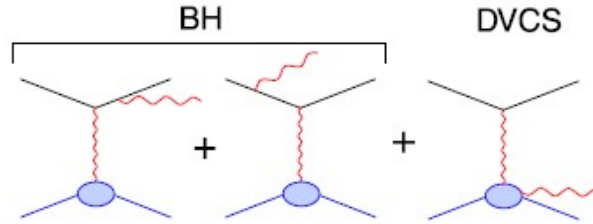
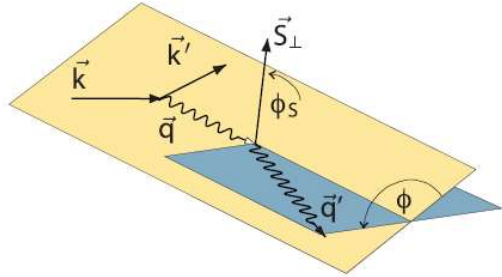
4) $M_{X=0}^2 = (p_{\mu_{\text{in}}} + p_{p_{\text{in}}} - p_{\mu_{\text{out}}} - p_{p_{\text{out}}} - p_{\gamma})^2$

Good agreement between $\mu^{\leftarrow+}$ and $\mu^{\rightarrow-}$ yields
 Important achievement for:

① $\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$ **Easier, done first**

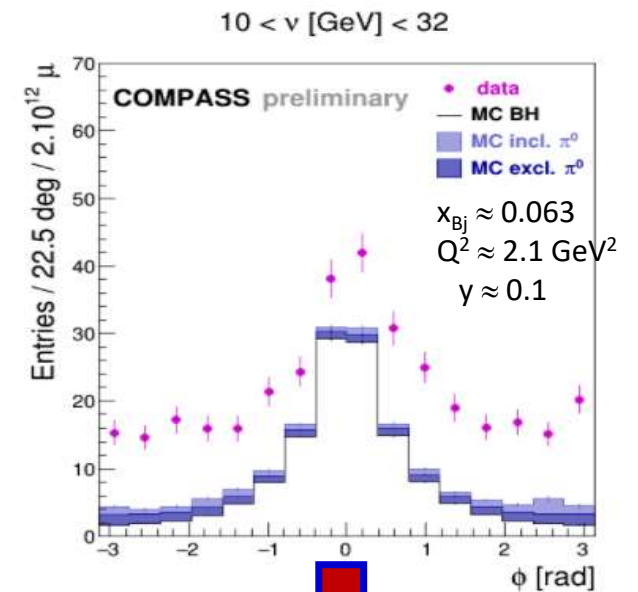
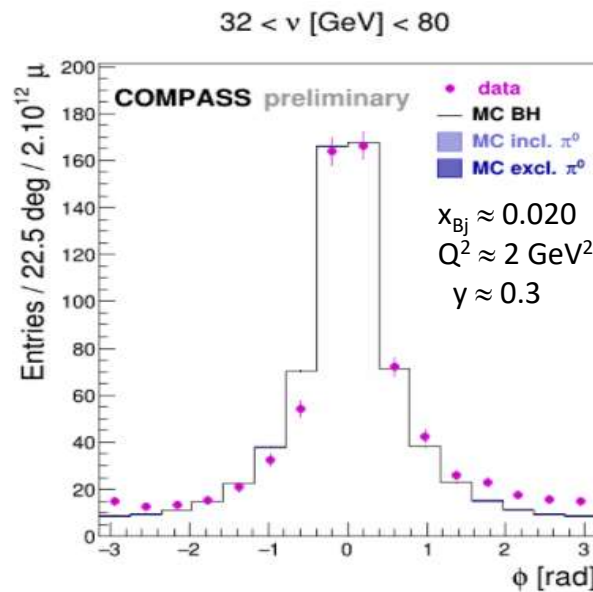
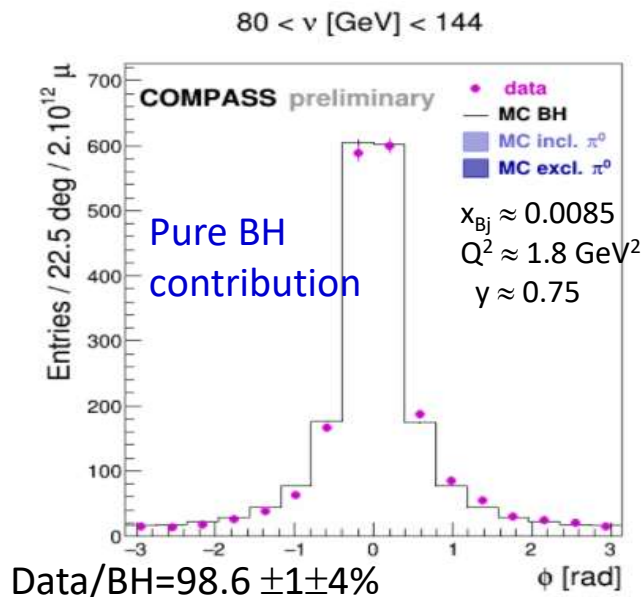
② $\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$ **Challenging, but promising**





$$\Sigma = d\sigma(\mu^+) + d\sigma(\mu^-)$$

$$d\sigma \propto |T^{BH}|^2 + \text{Interference Term} + |T^{DVCS}|^2$$



DVCS above the **BH** contrib.

MC: BH contribution evaluated for the integrated luminosity
 π^0 background contribution from SIDIS (LEPTO) + exclusive production (HEPGEN)

At COMPASS using polarized positive and negative muon beams:

$$S_{CS,U} \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-} = 2[d\sigma^{BH} + d\sigma_{unpol}^{DVCS} + \text{Im } I]$$

$$= 2[d\sigma^{BH} + c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi + s_1^I \sin \phi + s_2^I \sin 2\phi]$$

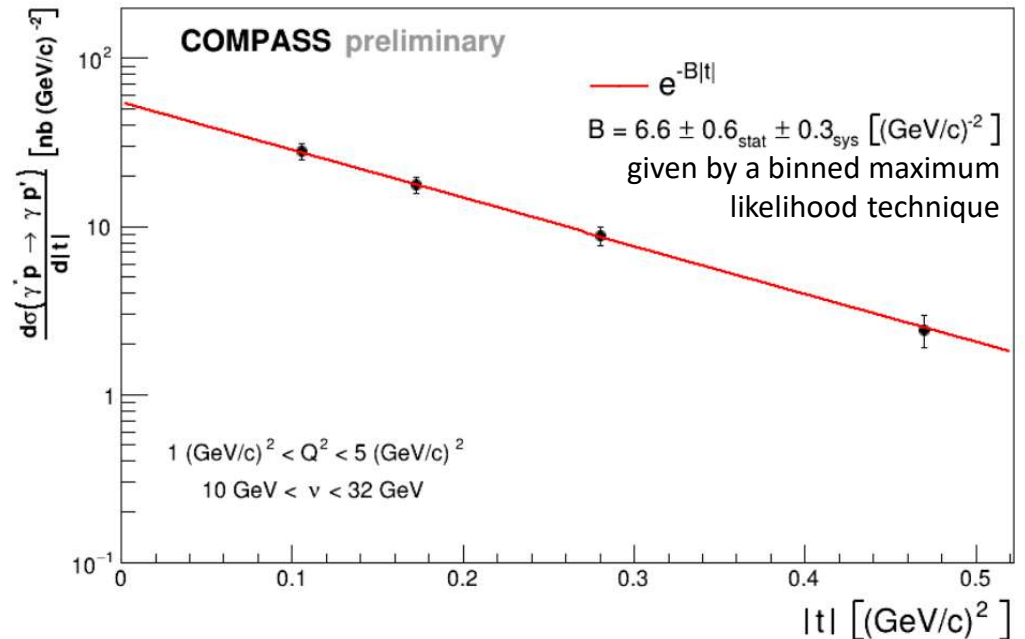
calculable
can be subtracted

All the other terms are cancelled in the integration over ϕ

$$\frac{d^3\sigma_T^{\mu p}}{dQ^2 d\nu dt} = \int_{-\pi}^{\pi} d\phi (d\sigma - d\sigma^{BH}) \propto c_0^{DVCS}$$

$$\frac{d\sigma^{\gamma^* p}}{dt} = \frac{1}{\Gamma(Q^2, \nu, E_\mu)} \frac{d^3\sigma_T^{\mu p}}{dQ^2 d\nu dt}$$

Flux for transverse virtual photons



COMPASS 12-16 Transverse extension of partons in the sea quark range

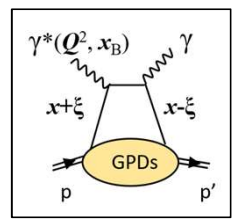
$$d\sigma^{DVCS}/dt = e^{-B|t|} = c_0^{DVCS} \propto (Im\mathcal{H})^2$$

$$c_0^{DVCS} \propto 4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) + \frac{t}{M^2}\mathcal{E}\mathcal{E}^*$$

In the COMPASS kinematics, $x_B \approx 0.06$, dominance of $Im\mathcal{H}$
 97% (GK model) 94% (KM model)

$$Im\mathcal{H} = H(x=\xi, \xi, t)$$

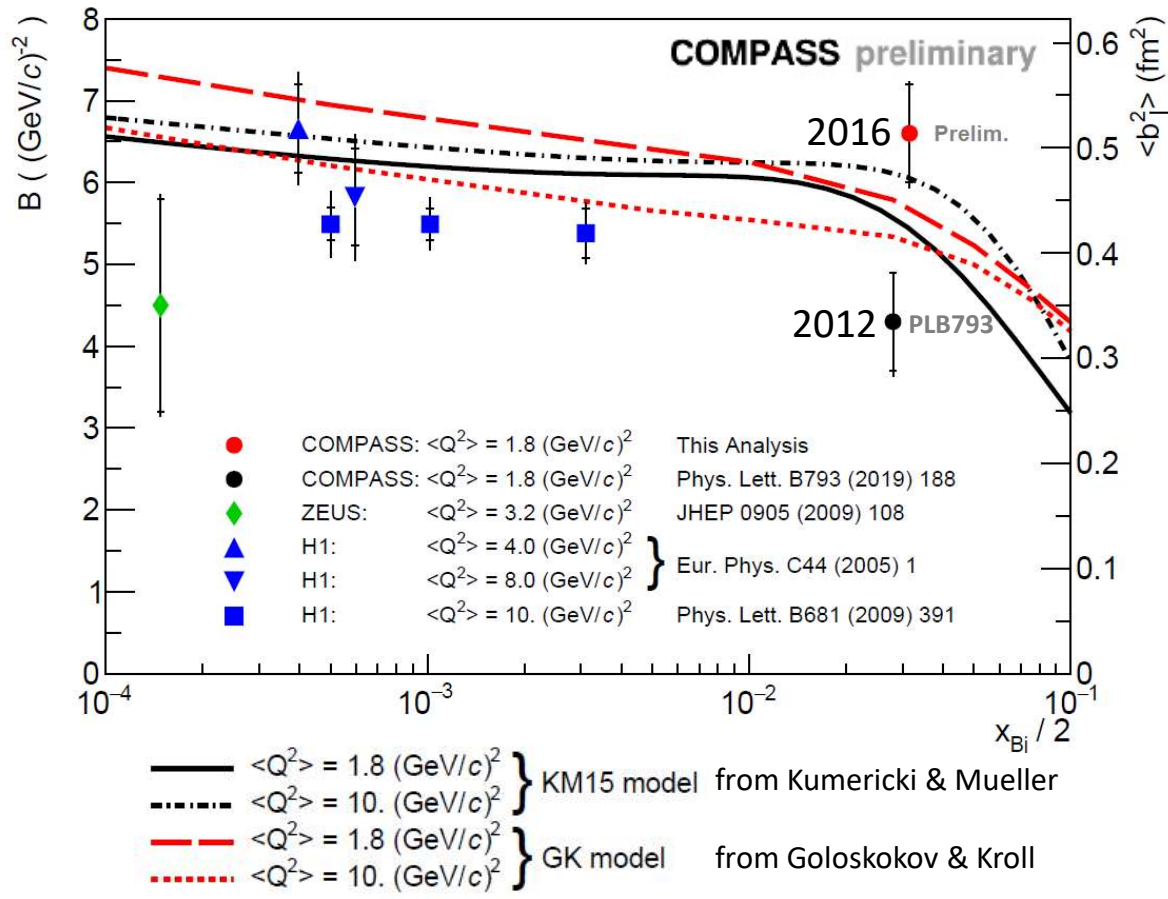
$$x = \xi \approx x_B/2 \text{ close to } 0$$



$$q(x, \mathbf{b}_\perp) = \int \frac{d^2\Delta_\perp}{(2\pi)^2} e^{-i\mathbf{b}_\perp \cdot \Delta_\perp} H^q(x, 0, -\Delta_\perp^2)$$

$$\langle b_\perp^2 \rangle_x^f = \frac{\int d^2b_\perp b_\perp^2 q_f(x, b_\perp)}{\int d^2b_\perp q_f(x, b_\perp)} = -4 \frac{\partial}{\partial t} \log H^f(x, \xi=0, t) \Big|_{t=0}$$

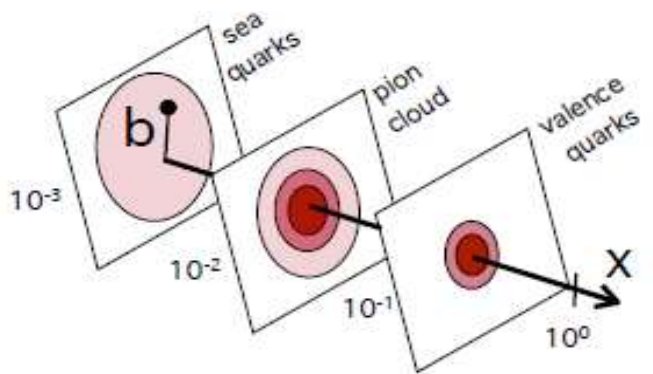
$$\langle b_\perp^2(x) \rangle \approx 2B(\xi)$$



COMPASS 12-16 Transverse extension of partons in the sea quark range

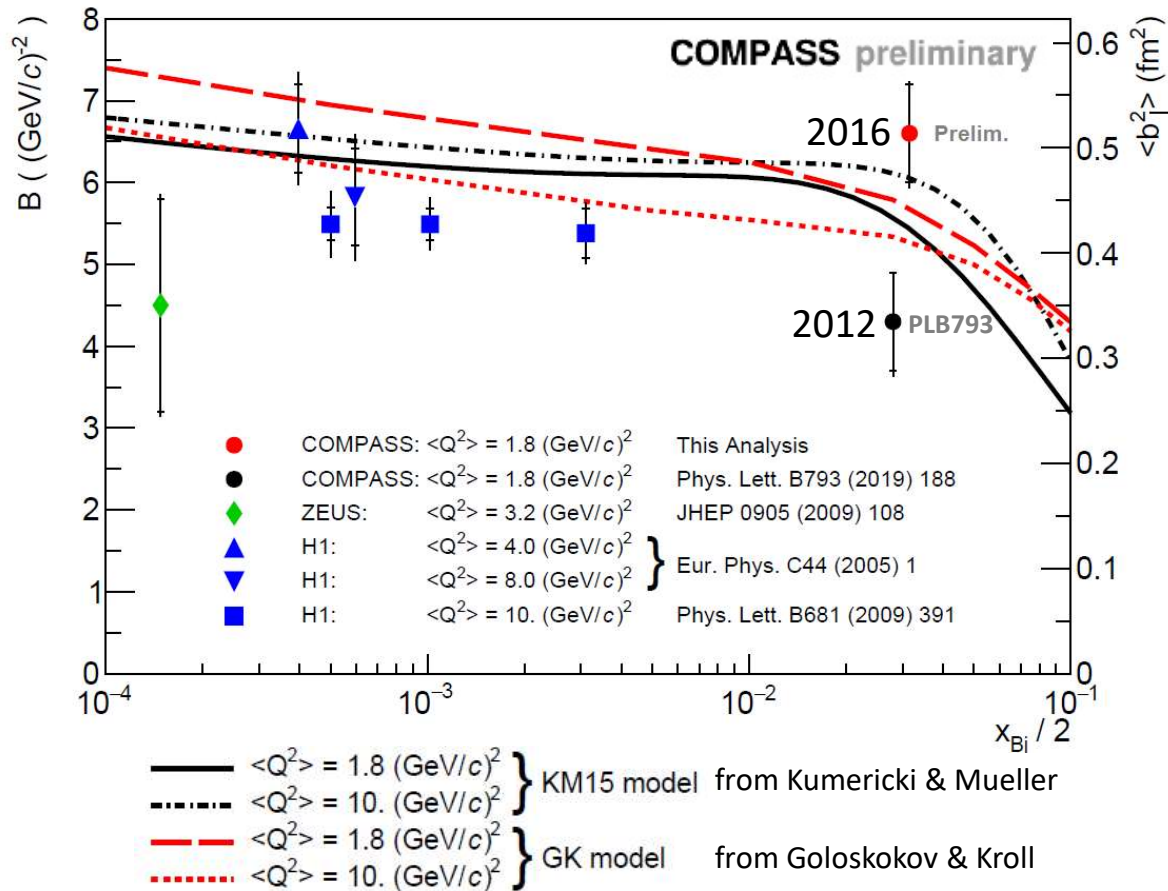
$$d\sigma^{DVCS}/dt = e^{-B|t|} = c_0^{DVCS} \propto (Im\mathcal{H})^2$$

$$\langle b_{\perp}^2(x) \rangle \approx 2B(\xi)$$



- 3 σ difference between 2012 and 2016 data
- more advanced analysis with 2016 data
 - π^0 contamination with different thresholds
 - binning with 3 variables (t,Q²,v) or 4 variables (t, ϕ ,Q²,v)

2012 statistics = Ref
 2016 analysed statistics = 2.3 \times Ref
 2016+2017 expected statistics = 10 \times Ref

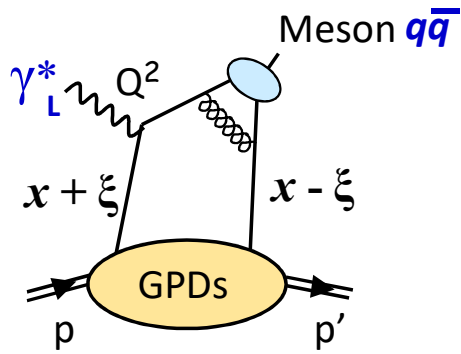


GPDs and Hard Exclusive Meson Production

Factorisation proven only for σ_L

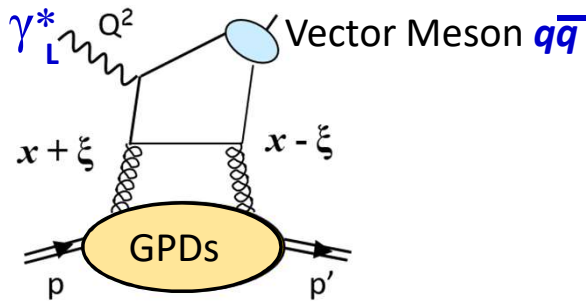
4 chiral-even GPDs: helicity of parton unchanged

Quark contribution



$\mathbf{H}^q(x, \xi, t)$	$\mathbf{E}^q(x, \xi, t)$ (as Sivers with OAM)	For Vector Meson
$\tilde{\mathbf{H}}^q(x, \xi, t)$	$\tilde{\mathbf{E}}^q(x, \xi, t)$	For Pseudo-Scalar Meson

Gluon contribution at the same order in α_s



+ Non-perturbative Meson wave function

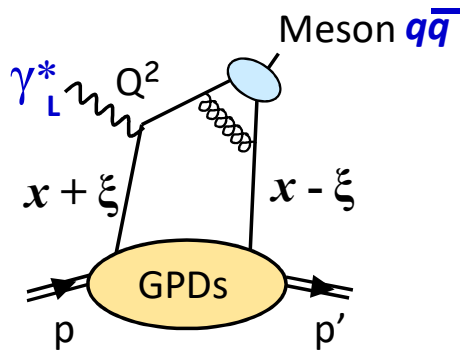
$$\rho^0 = [|u\bar{u}\rangle - |d\bar{d}\rangle] / \sqrt{2}$$

$$\omega = [|u\bar{u}\rangle + |d\bar{d}\rangle] / \sqrt{2}$$

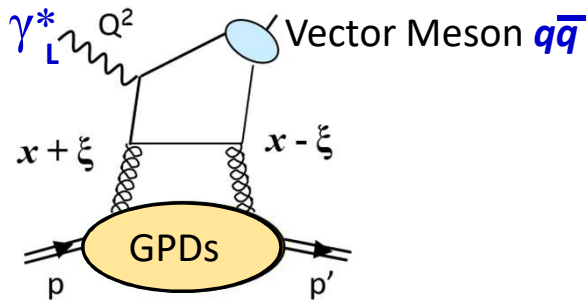
GPDs and Hard Exclusive Meson Production

Factorisation proven only for σ_L

Quark contribution



Gluon contribution at the same order in α_s



+ Non-perturbative Meson wave function

$$\rho^0 = [|u\bar{u}\rangle - |d\bar{d}\rangle] / \sqrt{2}$$

$$\omega = [|u\bar{u}\rangle + |d\bar{d}\rangle] / \sqrt{2}$$

4 chiral-even GPDs: helicity of parton unchanged

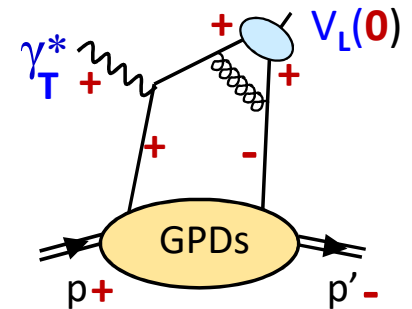
$H^q(x, \xi, t)$	$E^q(x, \xi, t)$ (as Sivers with OAM)	For Vector Meson
$\tilde{H}^q(x, \xi, t)$	$\tilde{E}^q(x, \xi, t)$	For Pseudo-Scalar Meson

+ 4 chiral-odd or transversity GPDs: helicity of parton changed (not possible in DVCS)

$H_T^q(x, \xi, t)$ (as transversity)	$E_T^q(x, \xi, t)$	$\bar{E}_T^q = 2 \tilde{H}_T^q + E_T^q$ (as Boer-Mulders)
$\tilde{H}_T^q(x, \xi, t)$	$\tilde{E}_T^q(x, \xi, t)$	

σ_T is asymptotically suppressed by $1/Q^2$ but large contribution observed
GK model: k_T of q and \bar{q} and Sudakov suppression factor are considered

$\mathcal{M}_{0-, ++}^{\gamma_T^* \rightarrow V_L}$ sensitive to H_T^q
and to a twist-3 meson wave function



COMPASS 2012 - 16 Exclusive π^0 production on unpolarized proton

$$\mu p \rightarrow \mu \pi^0 p$$

$$F_{\pi^0} = 2/3 F^u + 1/3 F^d$$

$$\frac{d^2\sigma}{dt d\phi_\pi} = \frac{1}{2\pi} \left[\left(\epsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} \right) + \epsilon \cos 2\phi_\pi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_\pi \frac{d\sigma_{LT}}{dt} \right]$$

$$\frac{d\sigma_L}{dt} \propto \left| \langle \tilde{H} \rangle \right|^2 - \frac{t'}{4m^2} \left| \langle \tilde{E} \rangle \right|^2$$

$$\frac{d\sigma_T}{dt} \propto \left| \langle H_T \rangle \right|^2 - \frac{t'}{8m^2} \left| \langle \bar{E}_T \rangle \right|^2$$

$$\frac{\sigma_{TT}}{dt} \propto \frac{t'}{16m^2} \left| \langle \bar{E}_T \rangle \right|^2$$

$$\frac{\sigma_{LT}}{dt} \propto \frac{\sqrt{-t'}}{2m} \text{Re} \left[\langle H_T \rangle^* \langle \tilde{E} \rangle \right]$$

$$\left\langle \frac{d\sigma_T}{d|t|} + \epsilon \frac{d\sigma_L}{d|t|} \right\rangle = (8.2 \pm 0.9_{\text{stat}} \pm 1.2_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{TT}}{d|t|} \right\rangle = (-6.1 \pm 1.3_{\text{stat}} \pm 0.7_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{LT}}{d|t|} \right\rangle = (1.5 \pm 0.5_{\text{stat}} \pm 0.3_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

COMPASS

$Q^2 = 2.0 \text{ GeV}^2$

$x_B = 0.093$

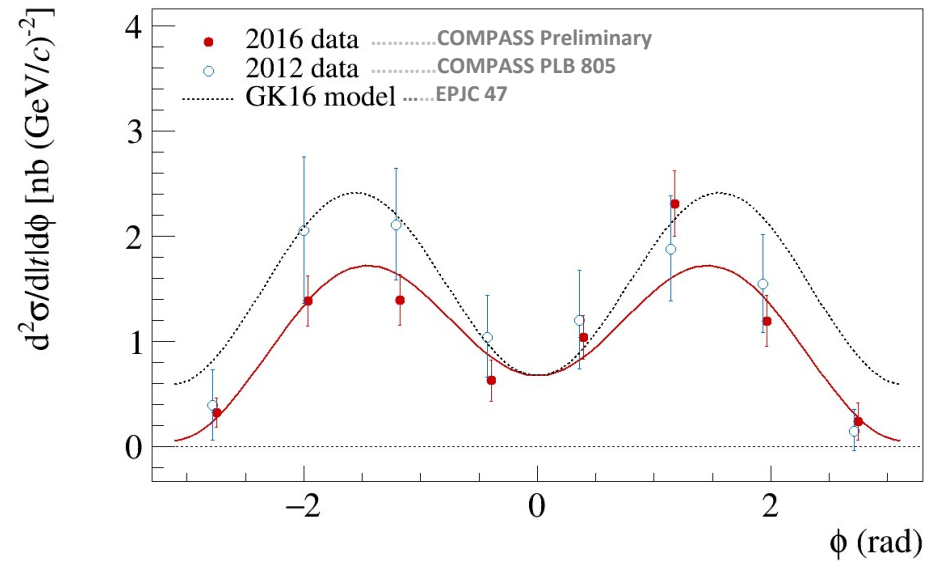
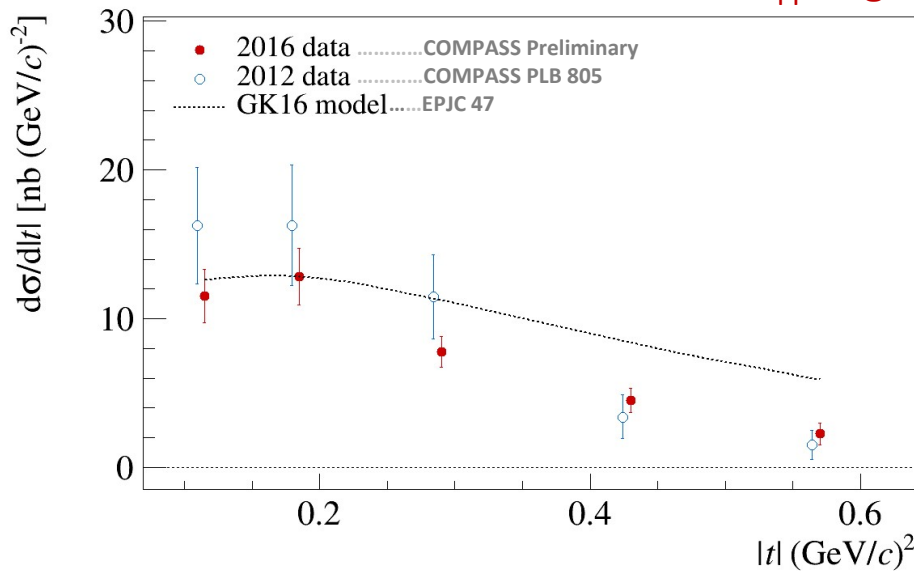
$|t| \sim 0.26 \text{ GeV}^2$

ϵ close to 1

PLB 805 (2020)

σ_{TT} large - impact of \bar{E}_T

σ_{LT} small but significantly positive as at CLAS



Data: **COMPASS**, PLB 805 (2020)

Models: **GK** Kroll Goloskokov EPJC47 (2011) Also **GGL**: Golstein Gonzalez Liuti PRD91 (2015)

HEMP with Transversely Polarized Target without RPD

Gparity: $G(\pi)=-1$; $G(\rho)=+1$; $G(\omega)=-1$

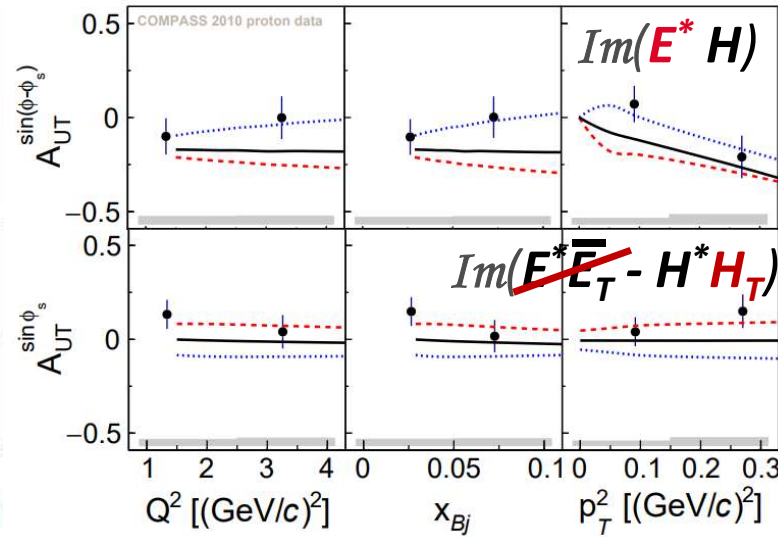
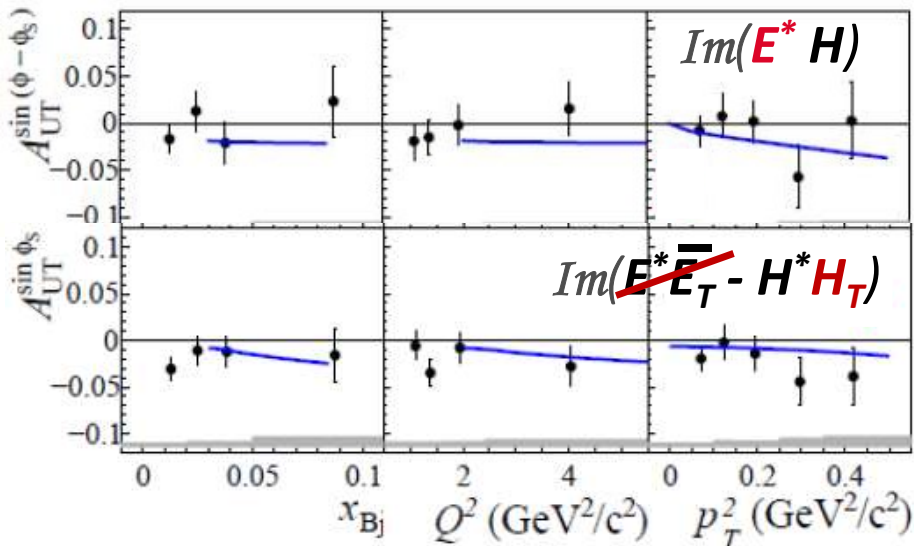
$$E_{\rho^0} = \frac{1}{\sqrt{2}} \left(\frac{2}{3} E^u \oplus \frac{1}{3} E^d + \frac{3}{4} \frac{E^g}{x} \right)$$

$$E_{\omega} = \frac{1}{\sqrt{2}} \left(\frac{2}{3} E^u \ominus \frac{1}{3} E^d + \frac{1}{4} \frac{E^g}{x} \right)$$

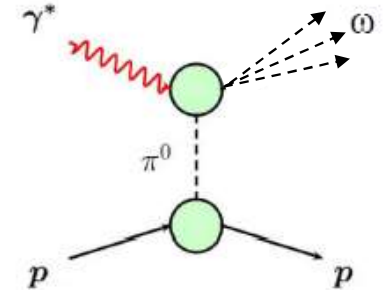
E^u and E^d of opposite sign
 ω is more promising
(see the larger scale)
but there is the inherent
pion pole contribution

COMPASS, NPB 865 (2012) 1-20, PLB731 (2014) 19

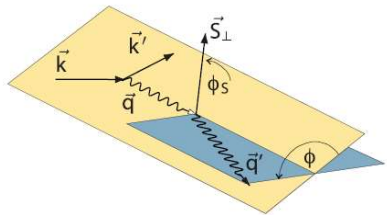
COMPASS, NPB 915 (2017)



$\Gamma(\omega \rightarrow \pi^0 \gamma) = 9 \times \Gamma(\rho^0 \rightarrow \pi^0 \gamma)$
Same for $\pi\omega$ FF but sign unknown



- ▶ positive $\pi\omega$ form factor
- ▶ no pion pole
- ▶ negative $\pi\omega$ form factor



GK Goloskokov, Kroll, EPJC42,50,53,59,65,74 GPD model constrained by HEMP at small x_B
longitudinal $\gamma_L^* p \rightarrow M p$ and transv. polar. $\gamma_T^* p \rightarrow M p$
quark and gluon contributions (GPDs H, E, H_T, E_T) and beyond leading twist

exclusive VM production with Unpolarised Target and SDME

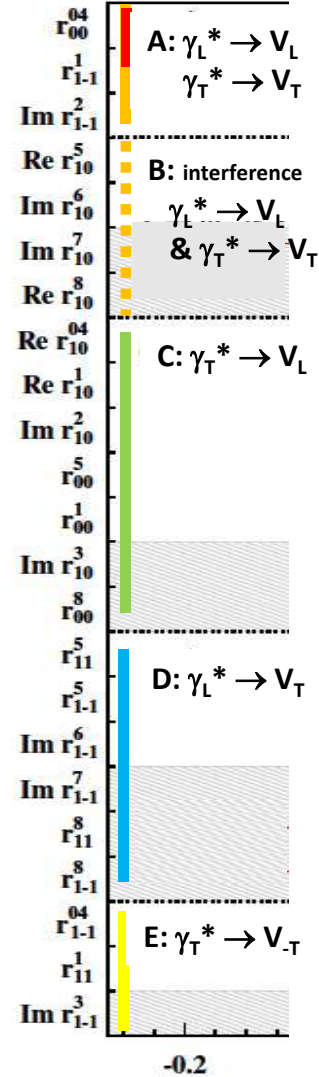
experimental angular distributions:

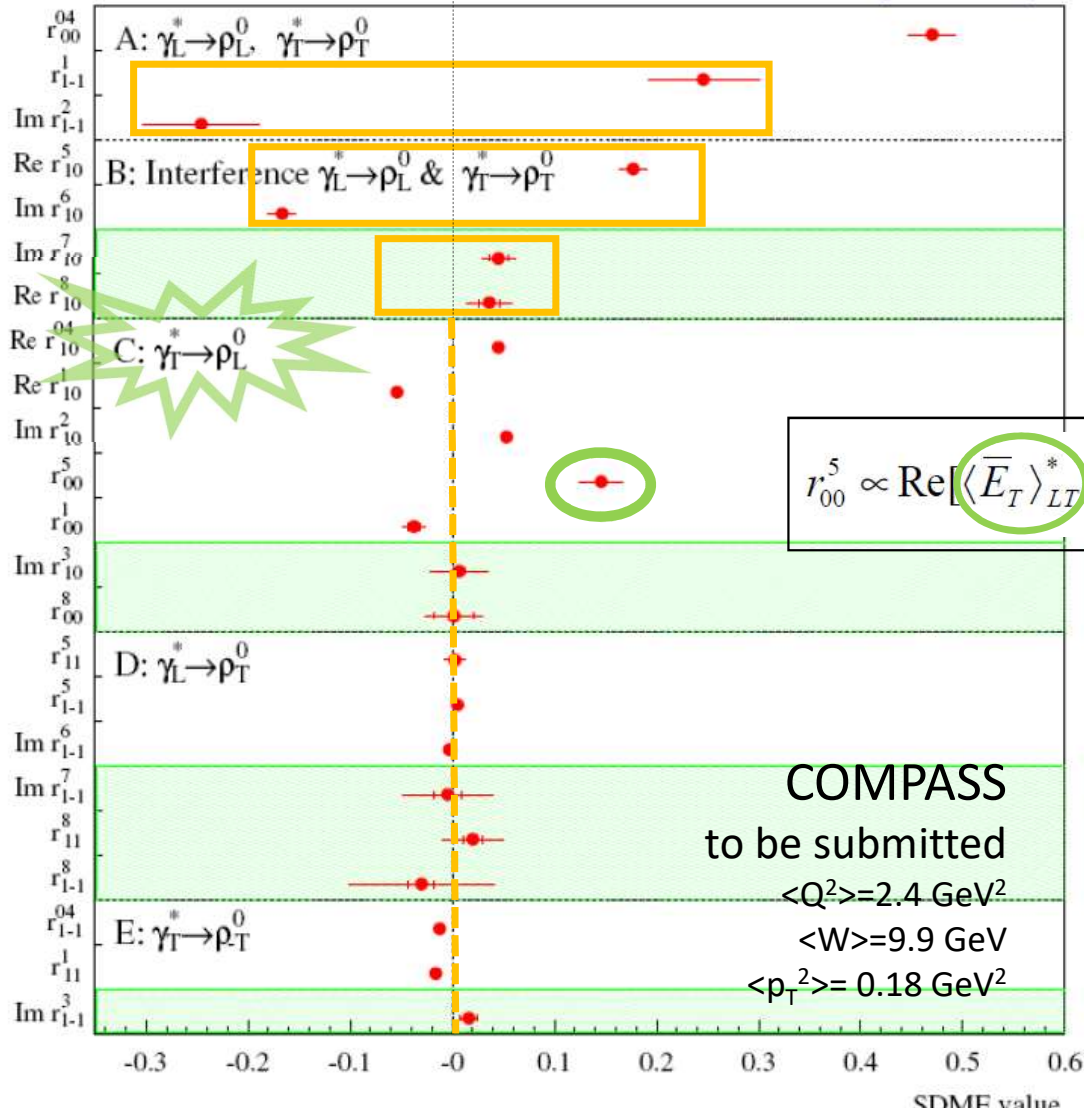
$$\mathcal{W}^{U+L}(\Phi, \phi, \cos \Theta) = \mathcal{W}^U(\Phi, \phi, \cos \Theta) + P_b \mathcal{W}^L(\Phi, \phi, \cos \Theta)$$

15 'unpolarized' and 8 'polarized' SDMEs

$$\begin{aligned} \mathcal{W}^U(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[\frac{1}{2}(1 - r_{00}^{04}) + \frac{1}{2}(3r_{00}^{04} - 1) \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^{04}\} \sin 2\Theta \cos \phi - r_{1-1}^{04} \sin^2 \Theta \cos 2\phi \right. \\ & - \epsilon \cos 2\Phi \left(r_{11}^1 \sin^2 \Theta + r_{00}^1 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^1\} \sin 2\Theta \cos \phi - r_{1-1}^1 \sin^2 \Theta \cos 2\phi \right) \\ & - \epsilon \sin 2\Phi \left(\sqrt{2}\text{Im}\{r_{10}^2\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^2\} \sin^2 \Theta \sin 2\phi \right) \\ & + \sqrt{2\epsilon(1+\epsilon)} \cos \Phi \left(r_{11}^5 \sin^2 \Theta + r_{00}^5 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^5\} \sin 2\Theta \cos \phi - r_{1-1}^5 \sin^2 \Theta \cos 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \sin \Phi \left(\sqrt{2}\text{Im}\{r_{10}^6\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^6\} \sin^2 \Theta \sin 2\phi \right) \right], \\ \mathcal{W}^L(\Phi, \phi, \cos \Theta) = & \frac{3}{8\pi^2} \left[\sqrt{1-\epsilon^2} \left(\sqrt{2}\text{Im}\{r_{10}^3\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^3\} \sin^2 \Theta \sin 2\phi \right) \right. \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos \Phi \left(\sqrt{2}\text{Im}\{r_{10}^7\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^7\} \sin^2 \Theta \sin 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \sin \Phi \left(r_{11}^8 \sin^2 \Theta + r_{00}^8 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^8\} \sin 2\Theta \cos \phi - r_{1-1}^8 \sin^2 \Theta \cos 2\phi \right) \right] \end{aligned}$$

ϵ close to 1,
small \mathcal{W}^L
no L/T separation





If SCHC ($\lambda_\gamma = \lambda_\nu$)

measurements:

$$r_{1-1}^1 + \text{Im}\{r_{1-1}^2\} = 0 \quad = 0.000 \pm 0.005 \pm 0.003,$$

$$\text{Re}\{r_{10}^5\} + \text{Im}\{r_{10}^6\} = 0 \quad = 0.011 \pm 0.002 \pm 0.002,$$

$$\text{Im}\{r_{10}^7\} - \text{Re}\{r_{10}^8\} = 0 \quad = 0.009 \pm 0.014 \pm 0.028.$$

All the other SDME in classes C, D, E should be 0
 not observed for class C

$$r_{00}^5 \propto \text{Re} \left[\langle \bar{E}_T \rangle_{LT}^* \langle H \rangle_{LL} + \frac{1}{2} \langle H_T \rangle_{LT}^* \langle E \rangle_{LL} \right]$$

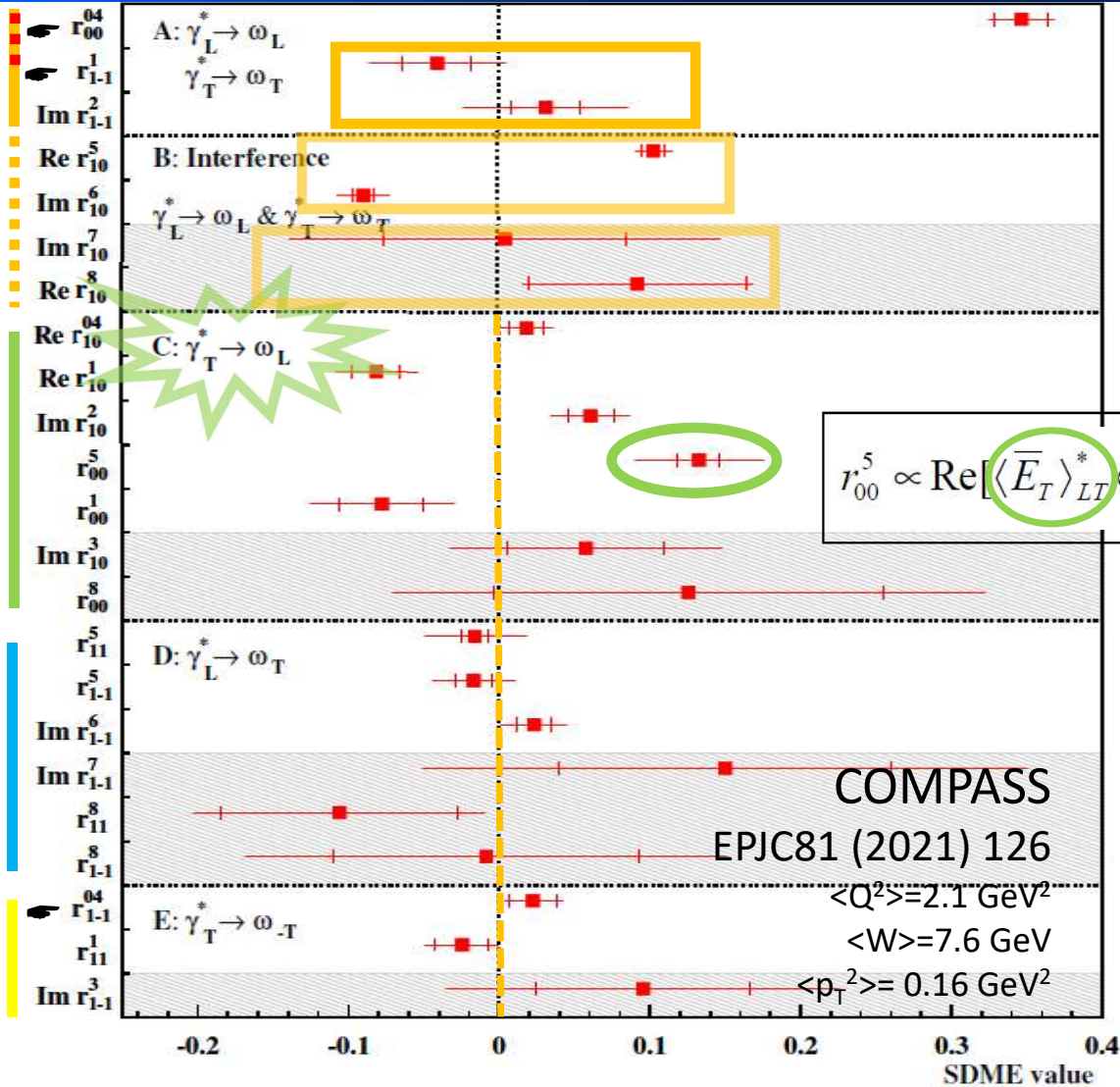
From Goloskokov and Kroll, EPJC74 (2014) 2725

$$F_{\rho^0} = 2/3 F^u + 1/3 F^d$$

→ The first term dominates and r_{00}^5 probes \bar{E}_T

- (H^u, H^d) of same sign
- $(\bar{E}_T^u, \bar{E}_T^d)$ of same sign
- (H_T^u, H_T^d) of opposite sign
- (E^u, E^d) of opposite sign

COMPASS
 to be submitted
 $\langle Q^2 \rangle = 2.4 \text{ GeV}^2$
 $\langle W \rangle = 9.9 \text{ GeV}$
 $\langle p_T^2 \rangle = 0.18 \text{ GeV}^2$



If SCHC ($\lambda_\gamma = \lambda_\nu$)

$r_{1-1}^1 + \text{Im}\{r_{1-1}^2\} = 0$	measurements:	$= -0.010 \pm 0.032 \pm 0.047$
$\text{Re}\{r_{10}^5\} + \text{Im}\{r_{10}^6\} = 0$		$= 0.014 \pm 0.011 \pm 0.013$
$\text{Im}\{r_{10}^7\} - \text{Re}\{r_{10}^8\} = 0$		$= -0.088 \pm 0.110 \pm 0.196$

All the other SDME in classes C, D, E should be 0
not observed for class C

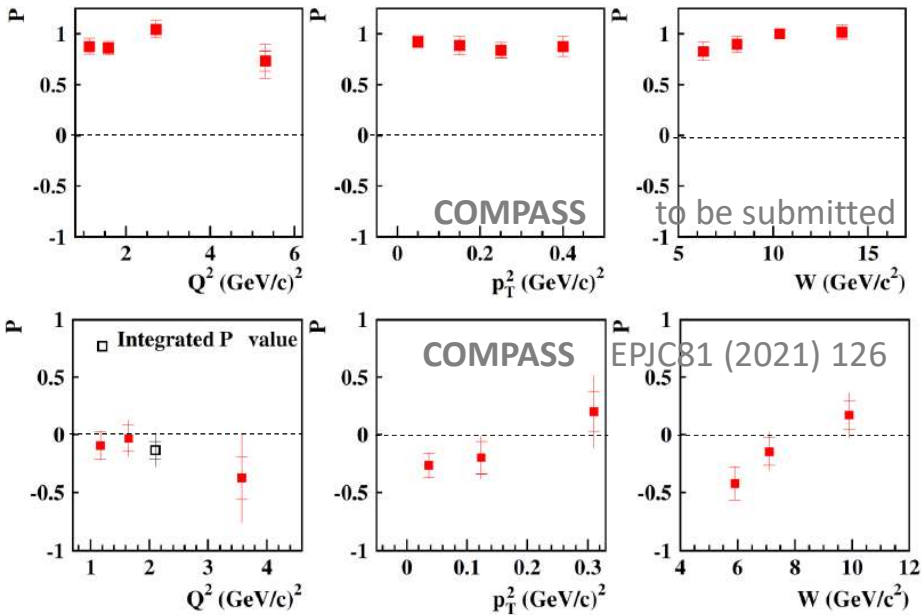
$r_{00}^5 \propto \text{Re} \left[\langle \bar{E}_T \rangle_{LT}^* \langle H \rangle_{LL} + \frac{1}{2} \langle H_T \rangle_{LT}^* \langle E \rangle_{LL} \right]$ From Goloskokov and Kroll, EPJC74 (2014) 2725

$F_\omega = 2/3 F^u - 1/3 F^d$
→ Both 2 terms are important

- (H^u, H^d) of same sign
- $(\bar{E}_T^u, \bar{E}_T^d)$ of same sign
- (H_T^u, H_T^d) of opposite sign
- (E^u, E^d) of opposite sign

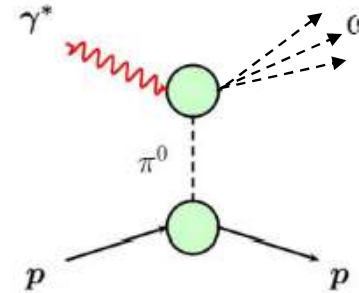
Natural (N) to Unnatural (U)
Parity Exchange for $\gamma_T^* \rightarrow V_T$

$$P = \frac{2r_{1-1}^1}{1 - r_{00}^{04} - 2r_{1-1}^{04}} \approx \frac{d\sigma_T^N(\gamma_T^* \rightarrow V_T) - d\sigma_T^U(\gamma_T^* \rightarrow V_T)}{d\sigma_T^N(\gamma_T^* \rightarrow V_T) + d\sigma_T^U(\gamma_T^* \rightarrow V_T)}$$

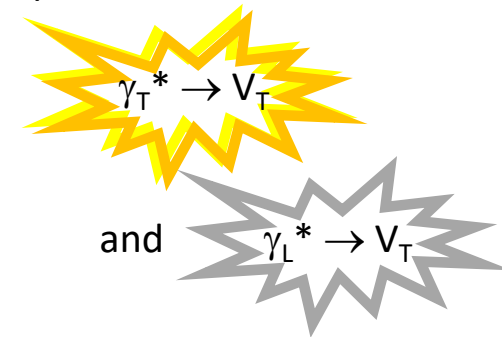


The pion pole exchange (UPE) is large for ω compared to ρ^0

$$\Gamma(\omega \rightarrow \pi^0 \gamma) = 9 \times \Gamma(\rho^0 \rightarrow \pi^0 \gamma) \text{ as for } \pi^0 \text{ Vector Meson FF}$$



It plays an important role in ω production for:



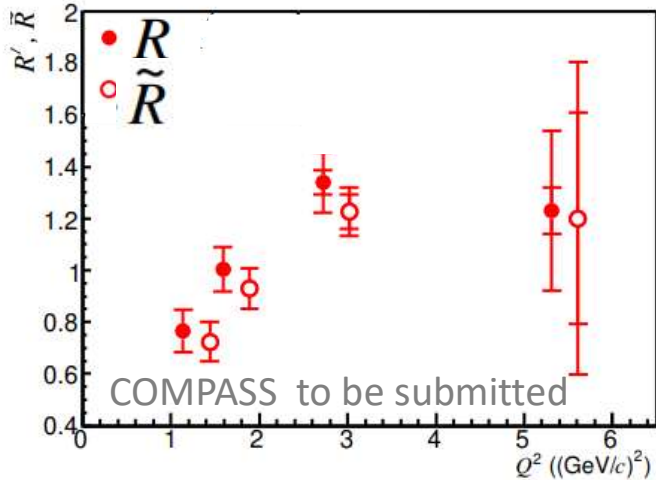
ρ^0 : $P \sim 1 \rightarrow$ NPE dominance $P \sim 1$
NPE with GPDs H, E

ω : $P \sim 0 \rightarrow$ NPE \sim UPE
UPE dominance at small W and p_T^2
UPE with GPDs \tilde{H}, \tilde{E} and the dominant pion pole

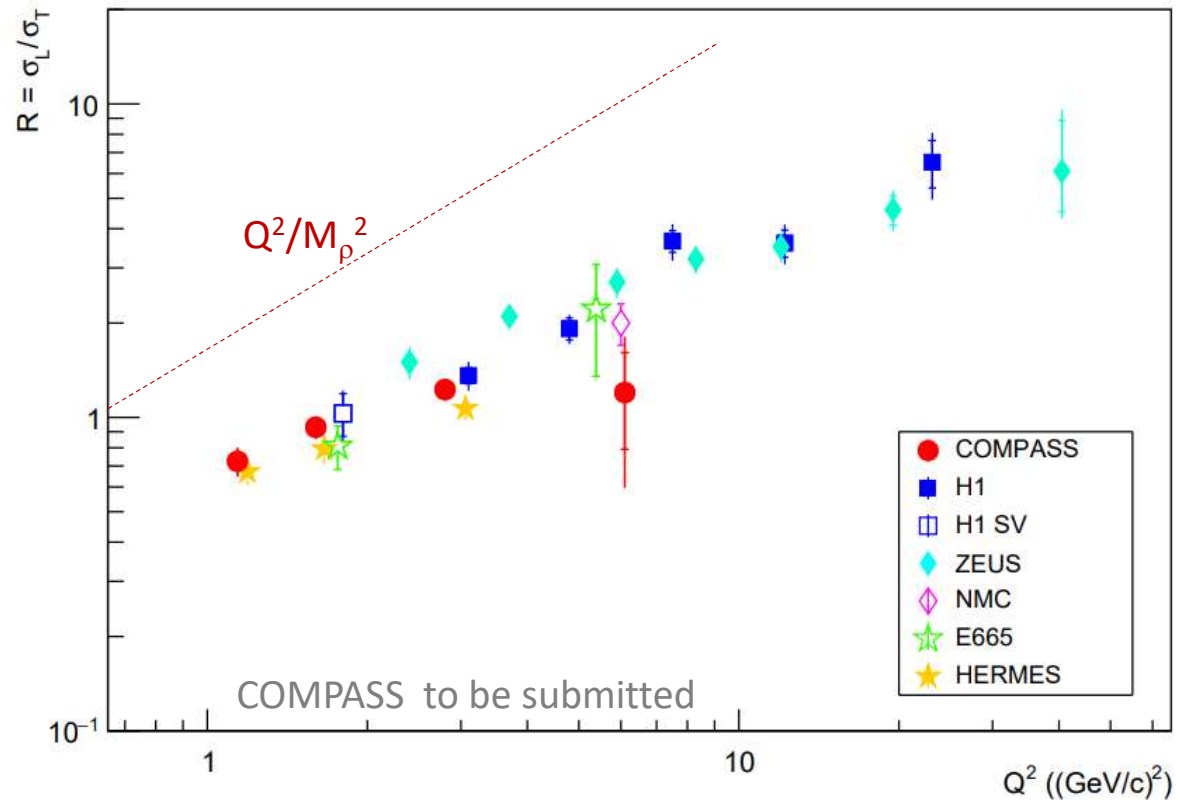
$$R = \frac{\sigma_L(\gamma_L^* \rightarrow V)}{\sigma_T(\gamma_T^* \rightarrow V)}$$

$$R = \frac{1}{\epsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}} \quad \text{only if SCHC}$$

In COMPASS domain evaluation of R and \tilde{R} considering violation of SCHC (and only NPE)



for all the experiments with $Q^2 > 1 \text{ GeV}^2$



Deviation from the pQCD LO prediction in Q^2/M_ρ^2 : QCD evolution and q_T Transverse size effects of the meson smaller for σ_L than for σ_T

- ✓ DVCS and the sum $\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$
→ c_0 and s_1 and constrain on $\text{Im}\mathcal{H}$ and Transverse extension of partons
- ✓ DVCS and the difference $\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$
→ c_1 and constrain on $\text{Re}\mathcal{H}$ (>0 as H1 or <0 as HERMES)
for D-term and pressure distribution
- ✓ On-going analysis (Cross section, SDME) for HEMP of $\pi^0, \rho^0, \omega, \phi, J/\psi$
 - ✓ Transversity GPDs
 - ✓ Gluon GPDs
 - ✓ Flavor decomposition



- ✓ DVCS and the sum $\Sigma \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-}$
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- ✓ DVCS and the difference $\Delta \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-}$
 - c_1 and constrain on $\text{Re}\mathcal{H}$ (>0 as H1 or <0 as HERMES)
- ✓ On-going analysis (Cross section, SDME) for HEMP of $\pi^0, \rho^0, \omega, \phi, J/\psi$
 - ✓ Transversity GPDs
 - ✓ Gluon GPDs
 - ✓ Flavor decomposition

For Jlab 20+ GeV

Importance of e^+ beam

for D-term and pressure distribution

Importance of large luminosity
For DVCS, TCS, DDVCS, J/ψ ...

