



experiment at CERN

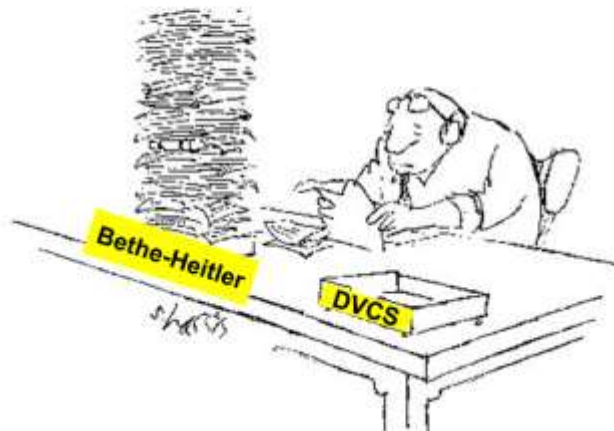
Hard Exclusive production of π^0 at COMPASS

O.Kouznetsov - JINR (Dubna) -
for the COMPASS collaboration

3 days



2008 DVCS test run: a first observation of exclusive single-photon production.



10 days



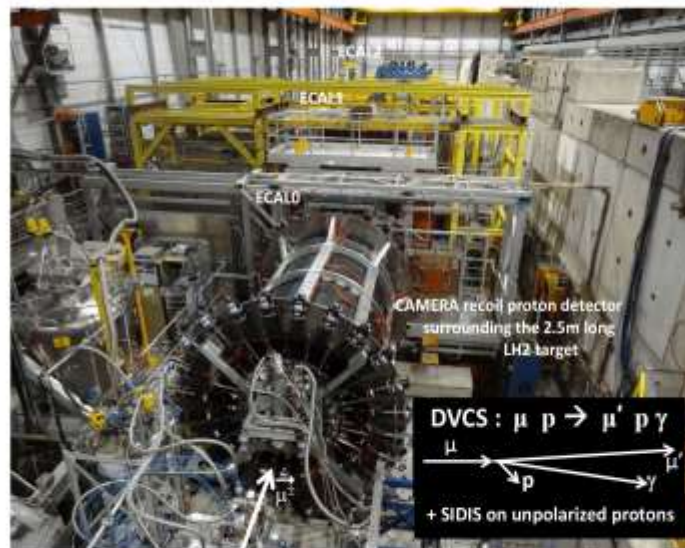
2009 DVCS test run: first estimation of pure DVCS, pure BH and DVCS-BH interference relative contributions



Exclusive π^0 signal in COMPASS

O.Kouznetsov
COMPASS Collaboration

Trento 2010

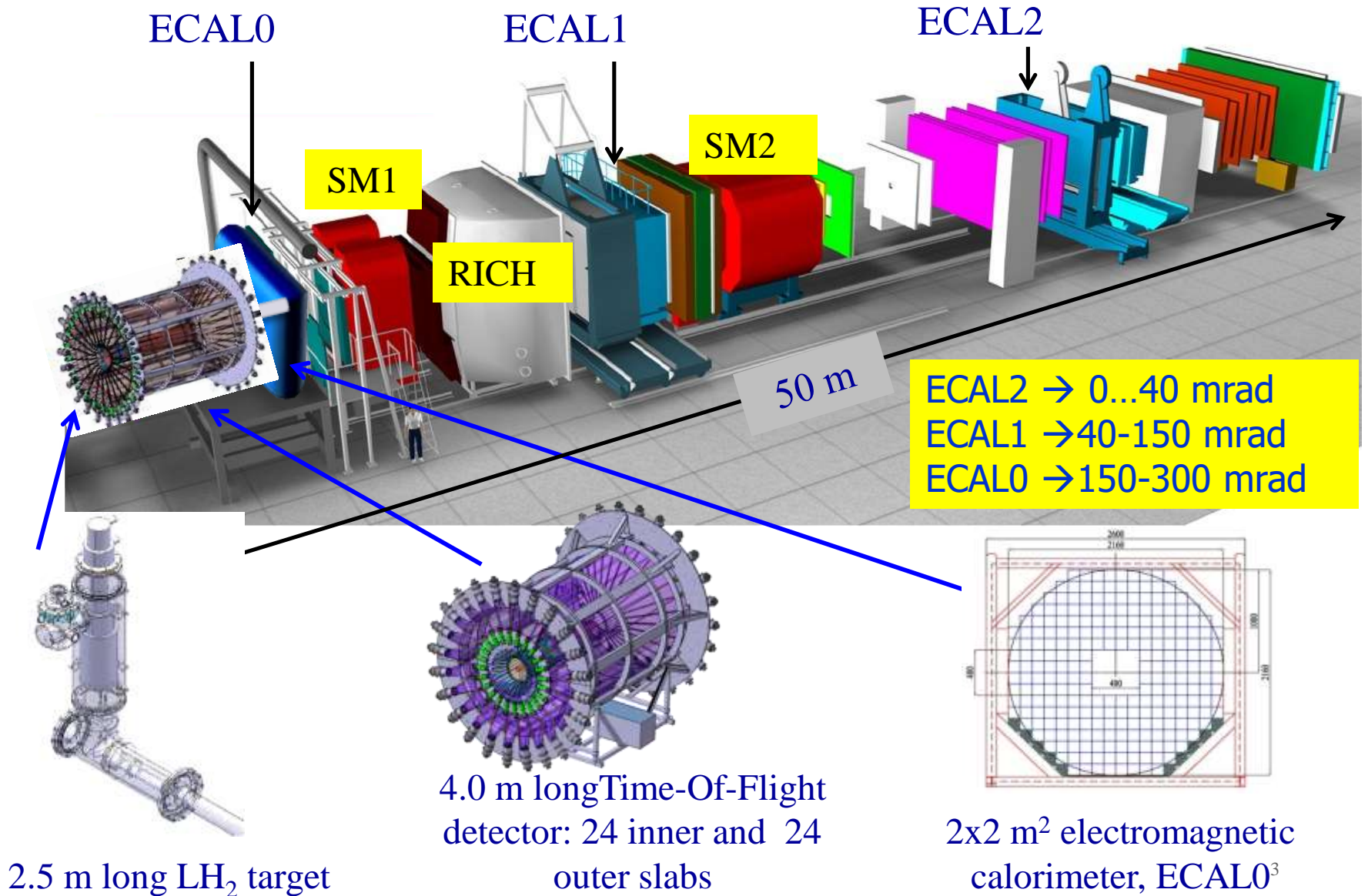


2012:
1 month pilot run

2016-17:
2 x 6 month
data taking

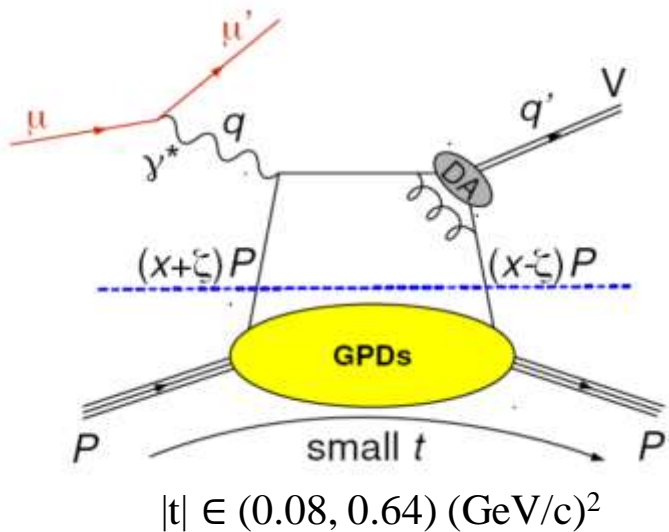
COMPASS spectrometer

COMPASS in μ runs
NIM A 577(2007) 455



Exclusive π^0 production is the main source of background for DVCS process, while it provides important information on chiral-odd GPDs. The dedicated GPD program has started with a one month pilot run in 2012, followed by two full years of data taking in 2016-2017, using 160 GeV/c positive and negative muon beams, a liquid hydrogen target and new detectors such as a recoil proton detector and a large-angle electromagnetic calorimeter ECAL0.

Hard exclusive π^0 production on unpolarised protons and chiral-odd GPDs

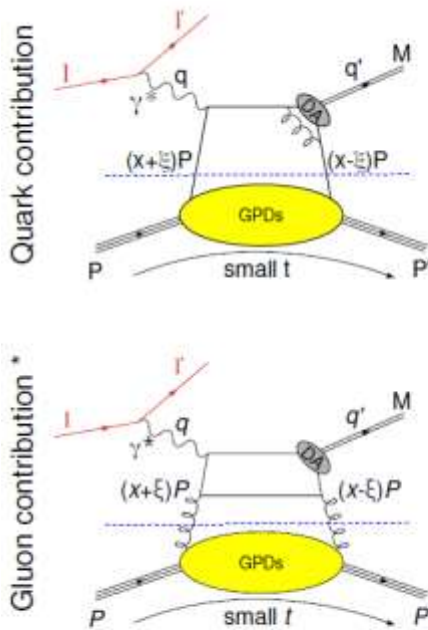


$$\mu p \rightarrow \mu' \pi^0 p'$$

μp cross-section can be reduced to $\gamma^* p$

$$\gamma^* p \rightarrow \pi^0 p'$$

GPDs and Deep Virtual Exclusive Meson Production



Chiral-even GPDs

helicity of parton conserved

$$H^{q,g}(x, \xi, t)$$

$$\tilde{H}^{q,g}(x, \xi, t)$$

$$E^{q,g}(x, \xi, t)$$

$$\tilde{E}^{q,g}(x, \xi, t)$$

Chiral-odd GPDs

helicity of parton flipped

$$H_T^q(x, \xi, t)$$

$$\tilde{H}_T^q(x, \xi, t)$$

$$E_T^q(x, \xi, t)$$

$$\tilde{E}_T^q(x, \xi, t)$$

Flavour separation

constraints for parton specific GPDs
due to different partonic content of mesons

Definition of variables:

- $q \dots \gamma^*$ four-momentum
- $x \dots$ average longitudinal momentum fraction of initial and final parton (NOT accessible)
- $\xi \dots$ difference of longitudinal-momentum fraction between initial and final parton
- $\approx x_B / (2 - x_B)$
- $t \dots$ four-momentum transfer

* Gluon contribution at same order of α_s as from quarks

In general hard exclusive π^0 is sensitive to the GPDs conserving the parton helicity (\tilde{H} , \tilde{E}) and also to the parton helicity flip or to chiral-odd GPDs (H_T and \bar{E}_T), where $\bar{E}_T = 2\tilde{H}_T + E_T$.

Hard exclusive π^0 production on unpolarised protons

Cross-section of the hard exclusive meson production, reduced to γ^*p , for the unpolarized target and polarized lepton beam

$$\frac{d^4\sigma_{\mu p}}{dQ^2 dt d\nu d\phi} = \Gamma \frac{d^2\sigma_{\gamma^* p}}{dt d\phi}$$

where $\Gamma = \Gamma(E_\mu, Q^2, \nu)$ is a transverse virtual-photon flux.

Spin independent cross-section of the hard exclusive meson production after averaging the two spin-dependent cross-sections looks following

$$\frac{d^2\sigma_{\gamma^* p}}{dt d\phi} = \frac{1}{2} \left(\frac{d^2\sigma_{\gamma^* p}^{\leftarrow}}{dt d\phi} + \frac{d^2\sigma_{\gamma^* p}^{\rightarrow}}{dt d\phi} \right) =$$

$$\frac{1}{2\pi} \left[\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \epsilon \cos(2\phi) \frac{d\sigma_{TT}}{dt} + \sqrt{\epsilon(1+\epsilon)} \cos\phi \frac{d\sigma_{LT}}{dt} \right] \Rightarrow \text{study } \phi \text{ dependence}$$

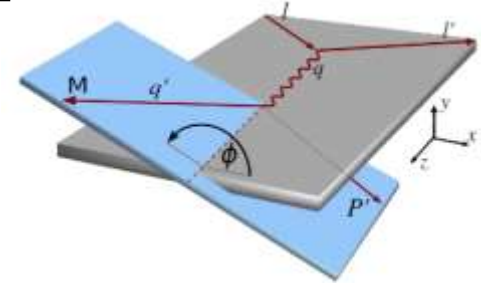
After integration in ϕ :

$$\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt}$$

\Rightarrow study t dependence

GPDs in exclusive π^0 production on unpolarised protons

$$\frac{d^2\sigma}{dt d\phi} = \frac{1}{2\pi} \left[\frac{d\sigma_T}{dt} + \varepsilon \frac{d\sigma_L}{dt} + \varepsilon \cos 2\phi \frac{d\sigma_{TT}}{dt} + \sqrt{2\varepsilon(1+\varepsilon)} \cos \phi \frac{d\sigma_{LT}}{dt} \right]$$



$$\frac{d\sigma_L}{dt} = \frac{4\pi\alpha}{k'} \frac{1}{Q^6} \left\{ (1 - \xi^2) |\langle \tilde{H} \rangle|^2 - 2\xi^2 \text{Re} [\langle \tilde{H} \rangle^* \langle \tilde{E} \rangle] - \frac{t'}{4m^2} \xi^2 |\langle \tilde{E} \rangle|^2 \right\}$$

leading twist
at JLAB only few% of $\frac{d\sigma_T}{dt}$

other contributions arise from coupling
of chiral-odd (quark helicity-flip) GPDs to twist-3 pion amplitude

$$\frac{d\sigma_T}{dt} = \frac{4\pi\alpha}{2k'} \frac{\mu_\pi^2}{Q^8} \left[(1 - \xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8m^2} |\langle \bar{E}_T \rangle|^2 \right]$$

$$\frac{\sigma_{LT}}{dt} = \frac{4\pi\alpha}{\sqrt{2}k'} \frac{\mu_\pi}{Q^7} \xi \sqrt{1 - \xi^2} \frac{\sqrt{-t'}}{2m} \text{Re} [\langle H_T \rangle^* \langle \tilde{E} \rangle]$$

$$\frac{\sigma_{TT}}{dt} = \frac{4\pi\alpha}{k'} \frac{\mu_\pi^2}{Q^8} \frac{t'}{16m^2} |\langle \bar{E}_T \rangle|^2$$

$$\text{def. } \bar{E}_T = 2\tilde{H}_T + E_T$$

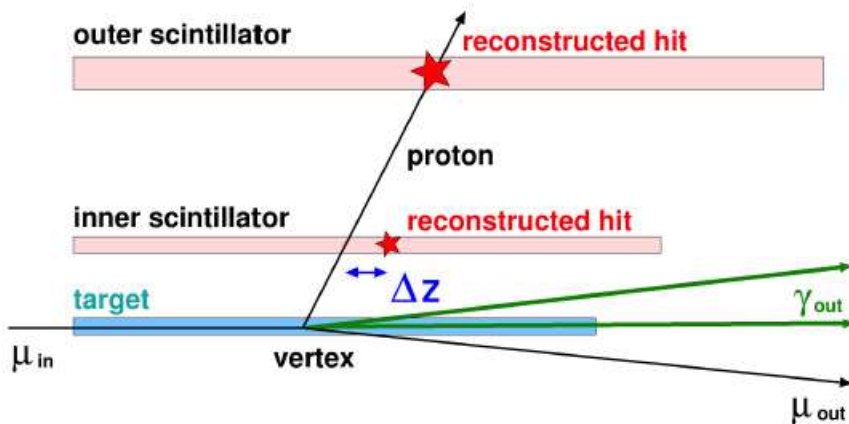
phenomenological Goloskokov&Kroll model
version 2011 based on JLab results
version 2016 update energy dependence \bar{E}_T
to fit both JLab and COMPASS results

An impact of \bar{E}_T should be visible in $\frac{\sigma_{TT}}{dt}$

and in a dip at small t' of $\frac{d\sigma_T}{dt}$

the effect of H_T should be visible

Exclusive π^0 events selection



Exclusivity variables

→ Transverse momentum constraint:

$$\Delta p_T = p_{T,spect}^p - p_{T,recoil}^p$$

→ Azimuthal angle constraint:

$$\Delta\varphi = \varphi_{spect}^p - \varphi_{recoil}^p$$

→ Z coordinate of inner CAMERA ring:

$$\Delta z = z_{spect}^p - z_{recoil}^p$$

→ Energy-momentum conservation:

$$M_X^2 = (p_{\mu,in} + p_p - p_{\mu,out} - p_{p'} - p_{\pi^0})^2$$

Invariant mass $M_{\mu\gamma}$ cut

General selection in the phase-space

- $8.5 \text{ GeV} < v < 28 \text{ GeV}$
- $1 (\text{GeV}/c)^2 < Q^2 < 5 (\text{GeV}/c)^2$
- $0.08 (\text{GeV}/c)^2 < |t| < 0.64 (\text{GeV}/c)^2$

→

Incoming and outgoing μ from the same/unique primary vertex

→

Two clusters in the ECAL0, ECAL1 calorimeters

→

Identified proton in Time-of-flight detector CAMERA

→

Final selection

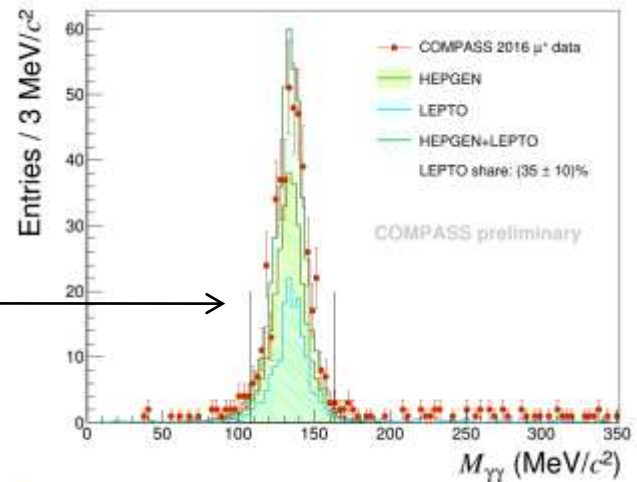
Invariant mass $M_{\gamma\gamma}$ cut

Kinematic fit

Final selection

Distribution of the $2 \gamma\gamma$ mass 2.5σ
cut $M_{\pi^0} \in (108, 163) \text{ MeV}/c^2$

Kinematic fit



Measurement of exclusive processes at COMPASS is overconstrained \rightarrow
can be used to improve precision of kinematic quantities using
kinematically constrained fit

Kinematic fit improves the resolution of the signal and lowers the
background

It works in a principle of minimisation of least square function
 $\chi^2(\vec{k}) = (\vec{k}_{fit} - \vec{k})^T \hat{C}^{-1} (\vec{k}_{fit} - \vec{k})$, where \vec{k} is a vector of measured
quantities and \hat{C} is their covariance matrix

Method used for the minimisation is Lagrange multipliers with
constraints g_i :

$$L(\vec{k}, \vec{\alpha}) = \chi^2(\vec{k}) + 2 \sum_{i=1}^N \alpha_i g_i$$

Constraints include momentum and energy conservation, common vertex
for all tracks (except proton), constraints for final proton, and mass
constraint

SIDIS background estimation

Main background for exclusive π^0 production \Rightarrow non-exclusive DIS processes

Two Monte Carlo simulations with the same π^0 selection criteria were used:

- \rightarrow LEPTO for the non-exclusive DIS background
- \rightarrow HEPGEN++ shape of distributions of exclusive π^0 production (signal contribution)

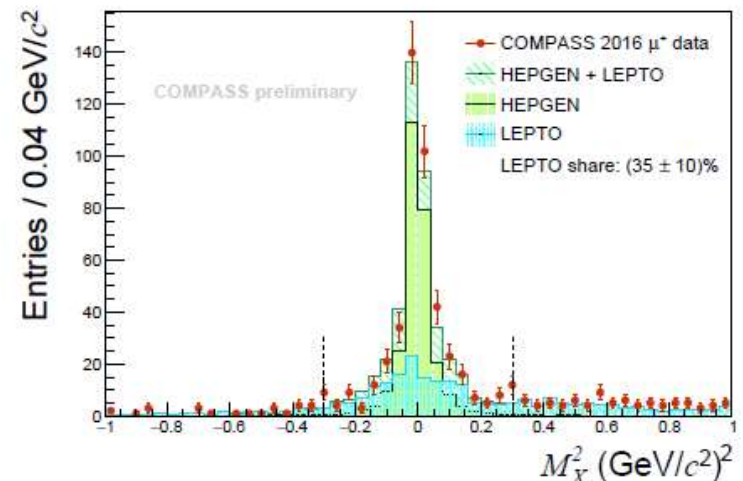
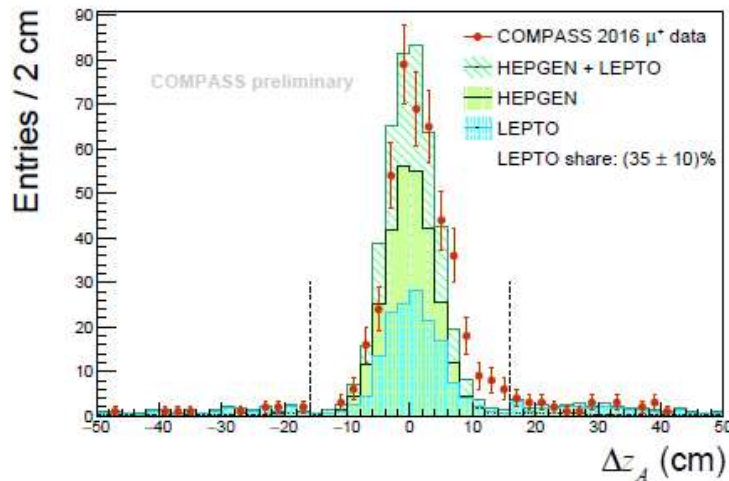
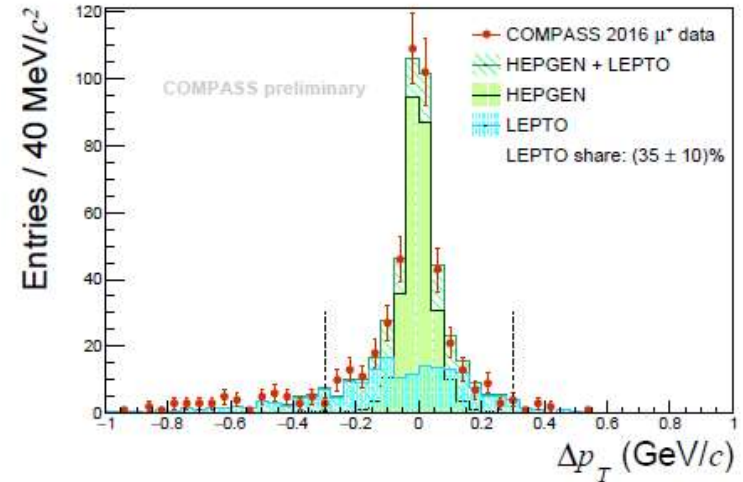
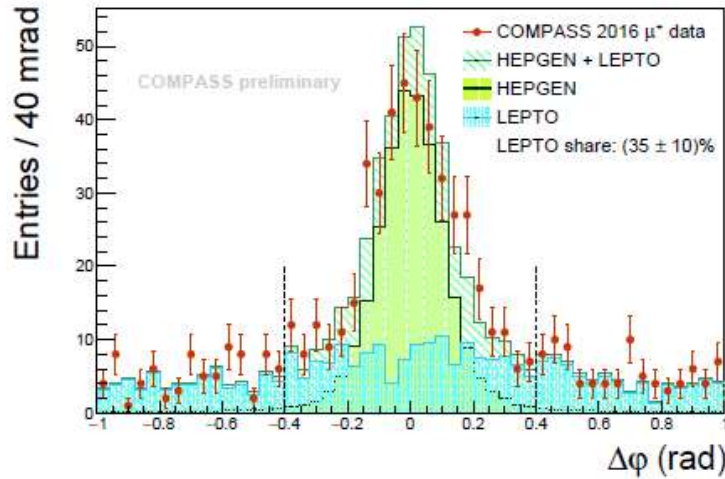
Search for best description of data fitting by mixture of both MC
Both MC samples normalized to the experimental $M_{\gamma\gamma}$ distribution
The ratio of background events r_{LEPTO} is determined by a fit on the exclusivity distributions

The fraction of non-exclusive back-ground in the 2016 data is estimated to be $(35 \pm 10)\%$
Background fit method is currently the main source of systematic uncertainty.

For the moment the contributions of other background sources are considered to be negligible. For example, the production of single ω mesons, where the ω decays into a π^0 and a photon that remains undetected, was found in Monte Carlo studies to contribute at the level of 1% (2012 analysis).

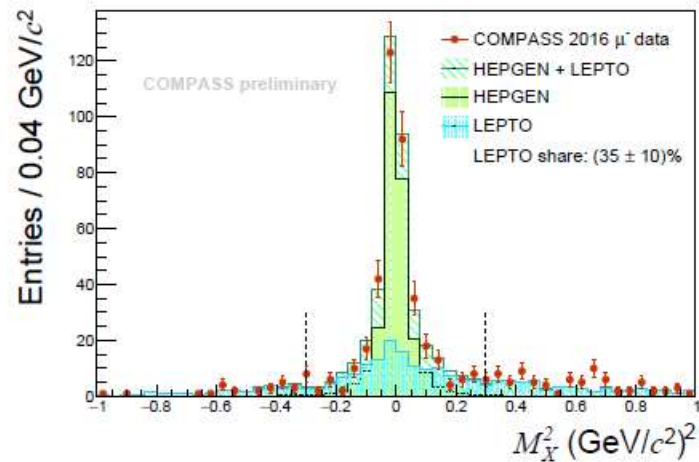
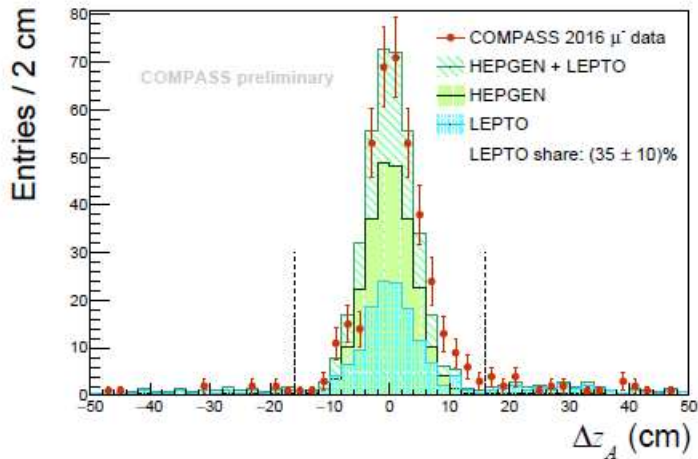
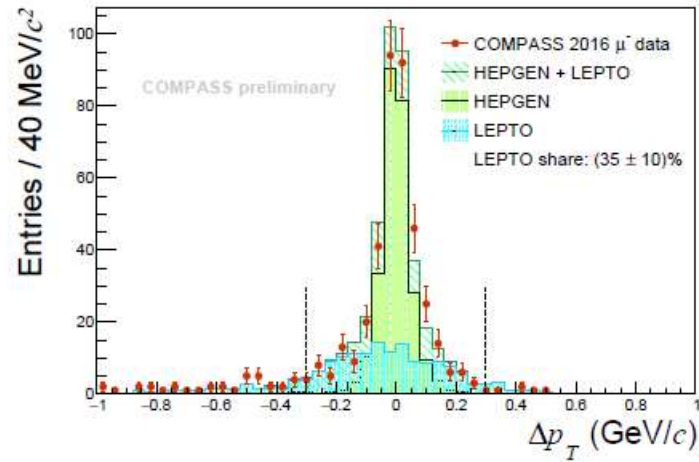
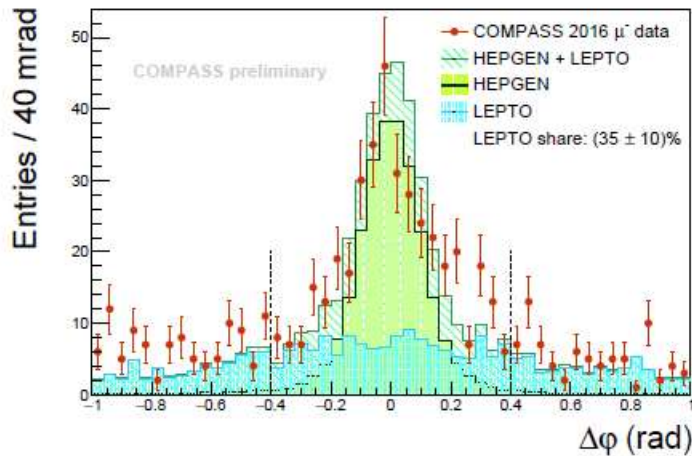
Distributions of the exclusivity variables for μ^+

Cuts: $|\Delta\phi| < 0.4$ rad $|\Delta p_T| < 0.3$ GeV/c $|\Delta Z_A| < 16$ cm $|M_X^2| < 0.3$ (GeV/c²)²



Distributions of the exclusivity variables for μ^-

Cuts: $|\Delta\phi| < 0.4$ rad $|\Delta p_T| < 0.3$ GeV/c $|\Delta Z_A| < 16$ cm $|M_X^2| < 0.3$ (GeV/c²)²



Exclusive π^0 production: COMPASS acceptance

The cross section is determined presently in the same phase space as for the 2012 data analysis:

- $8.5 \text{ GeV} < \nu < 28 \text{ GeV}$
- $1 \text{ (GeV/c)}^2 < Q^2 < 5 \text{ (GeV/c)}^2$
- $0.08 \text{ (GeV/c)}^2 < |t| < 0.64 \text{ (GeV/c)}^2$

For the acceptance determination, the HEPGen- π^0 MC simulation is used.

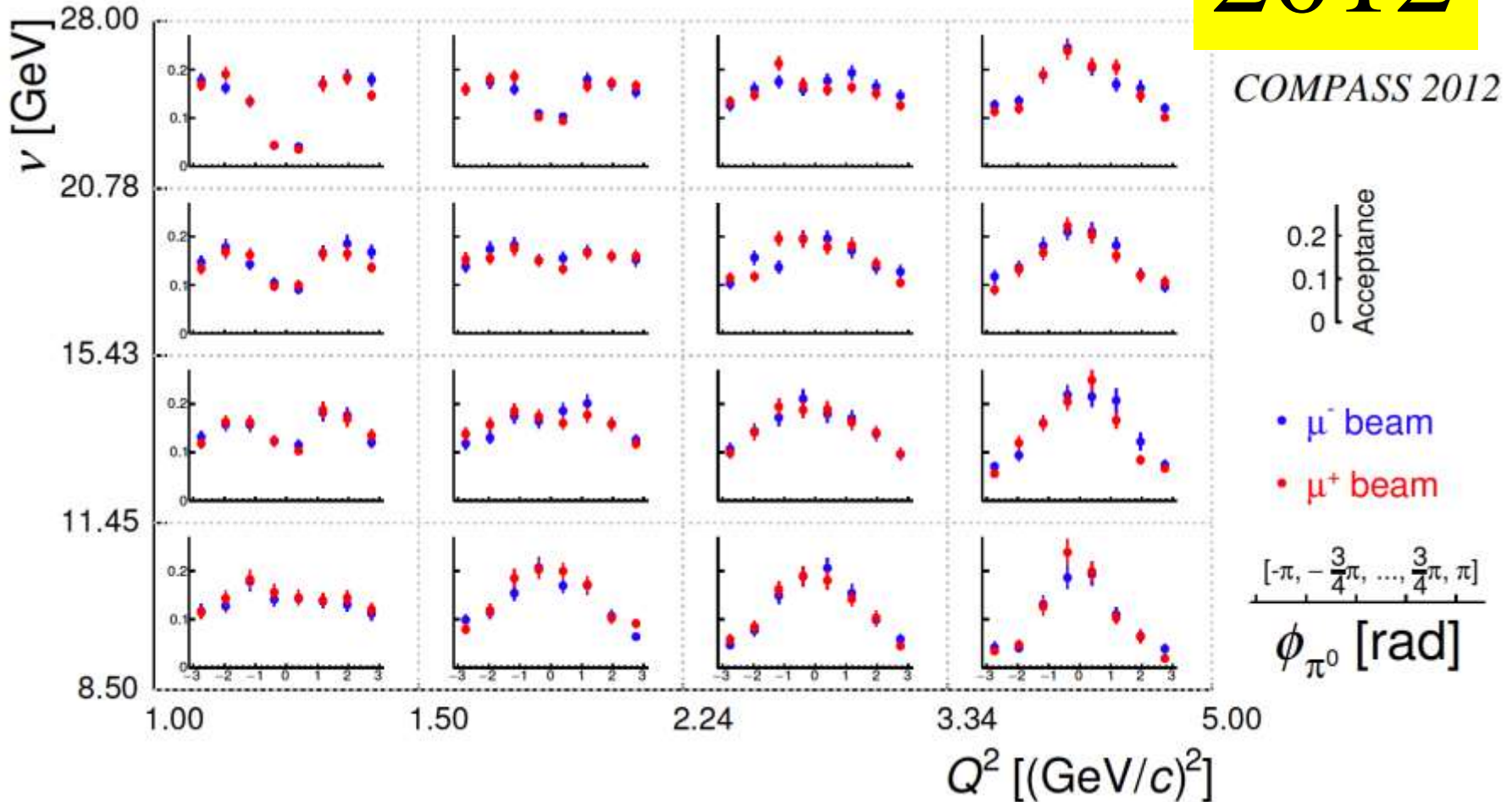
4D acceptance in bins of ϕ_{π^0} , ν , $|t|$, Q^2

- 5 bins in $|t|$ with binning $[0.08, 0.15]$, $[0.15, 0.22]$, $[0.22, 0.36]$, $[0.36, 0.5]$, $[0.5, 0.64] \text{ (GeV/c)}^2$,
- 8 bins in ϕ equally spaced from $-\pi$ to $+\pi$
- 4 bins in Q^2 with binning $[1, 1.5]$, $[1.5, 2.24]$, $[2.24, 3.34]$, $[3.34, 5] \text{ (GeV/c)}^2$,
- 4 bins in ν with binning $[8.5, 11.45]$, $[11.45, 15.43]$, $[15.43, 20.78]$, $[20.78, 28] \text{ (GeV)}$.

Exclusive π^0 production: COMPASS acceptance $|t|$ averaged

4D acceptance in bins of ϕ_{π^0} , ν , $|t|$, Q^2
 figure shows 3D projection, as a function of ϕ_{π^0}

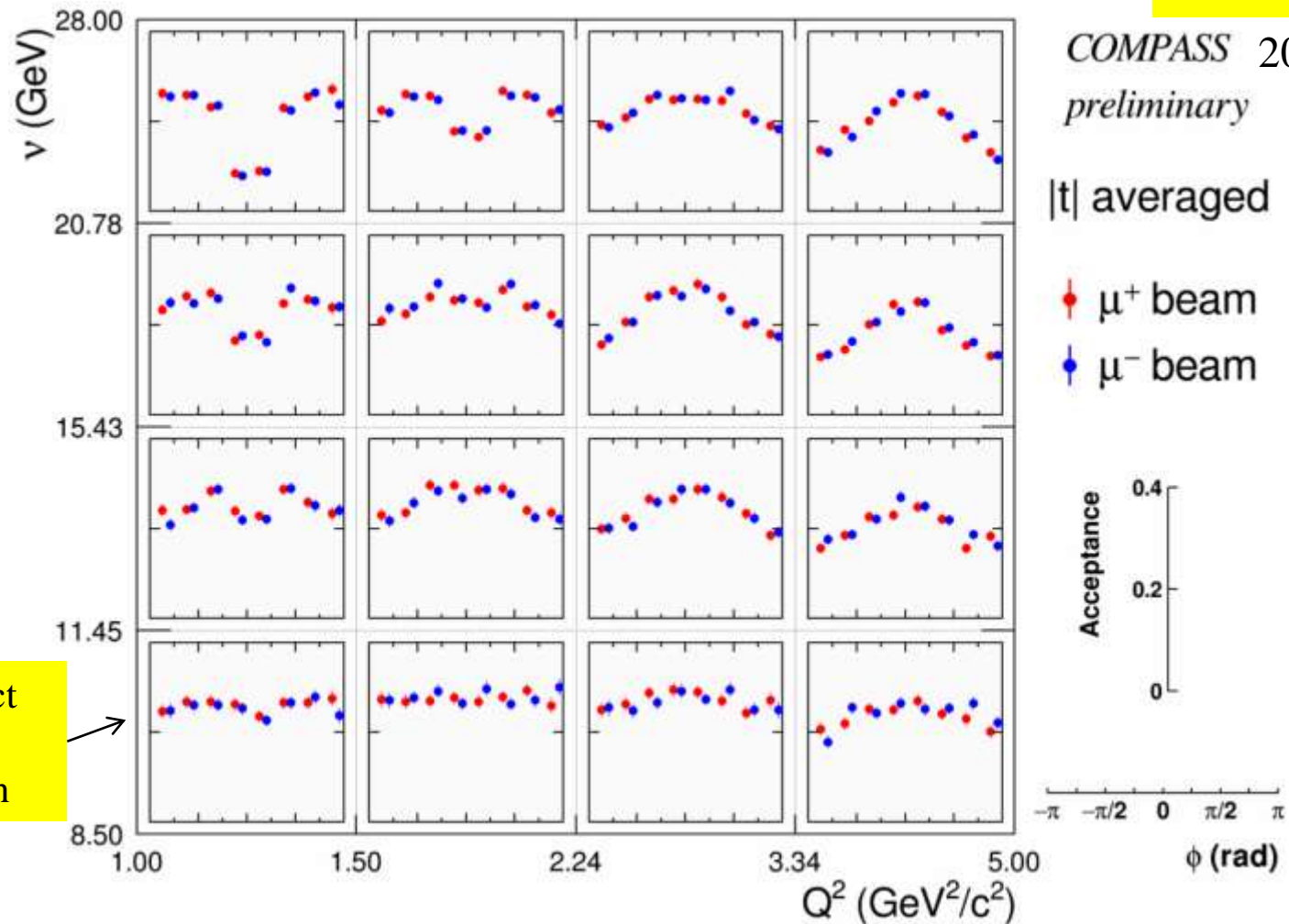
2012



Exclusive π^0 production: COMPASS acceptance $|t|$ averaged

4D acceptance in bins of ϕ_{π^0} , ν , $|t|$, Q^2
figure shows 3D projection, as a function of ϕ_{π^0}

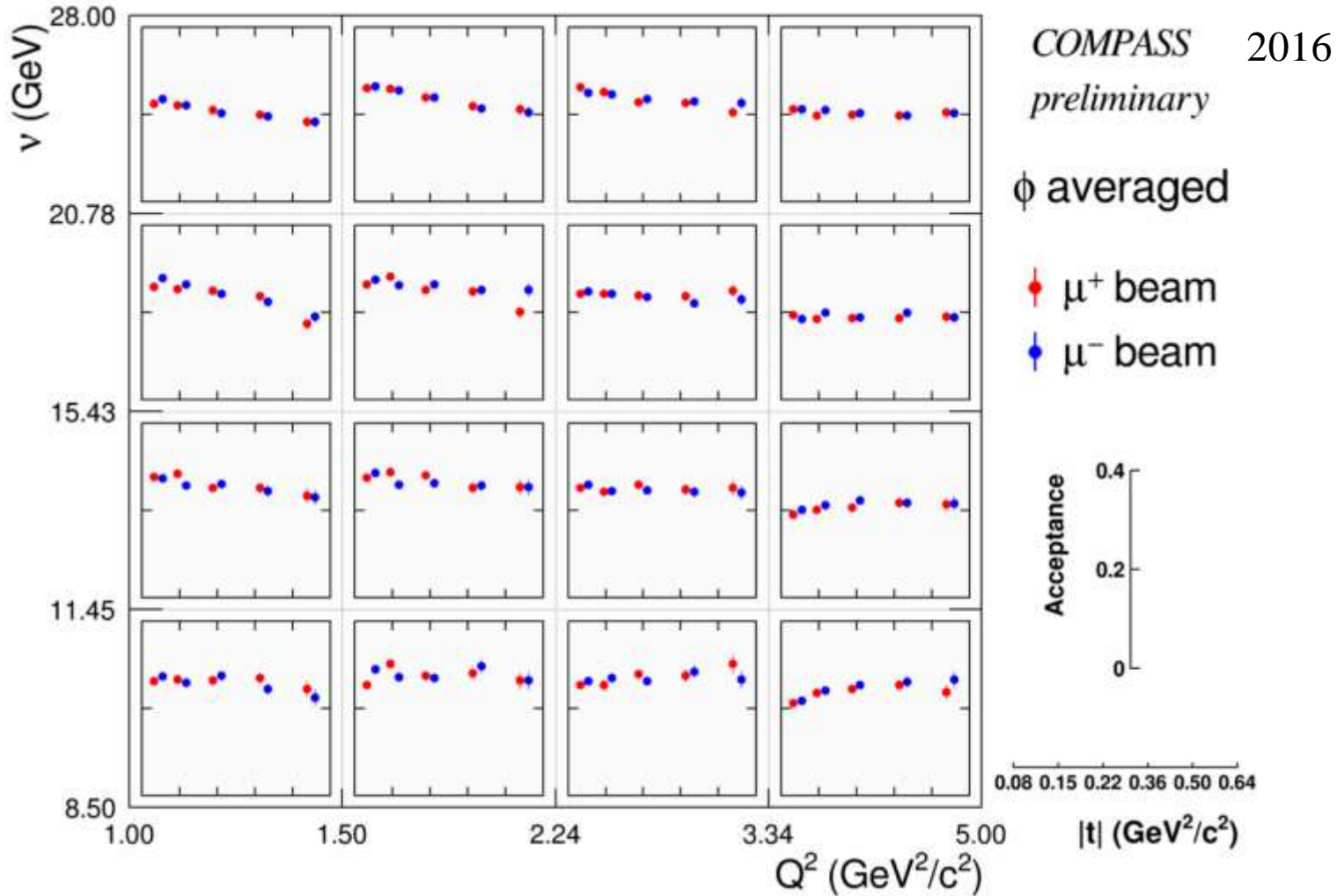
2016



Exclusive π^0 production: COMPASS acceptance ϕ averaged

4D acceptance in bins of ϕ_{π^0} , ν , $|t|$, Q^2
figure shows 3D projection, as a function of $|t|$

2016



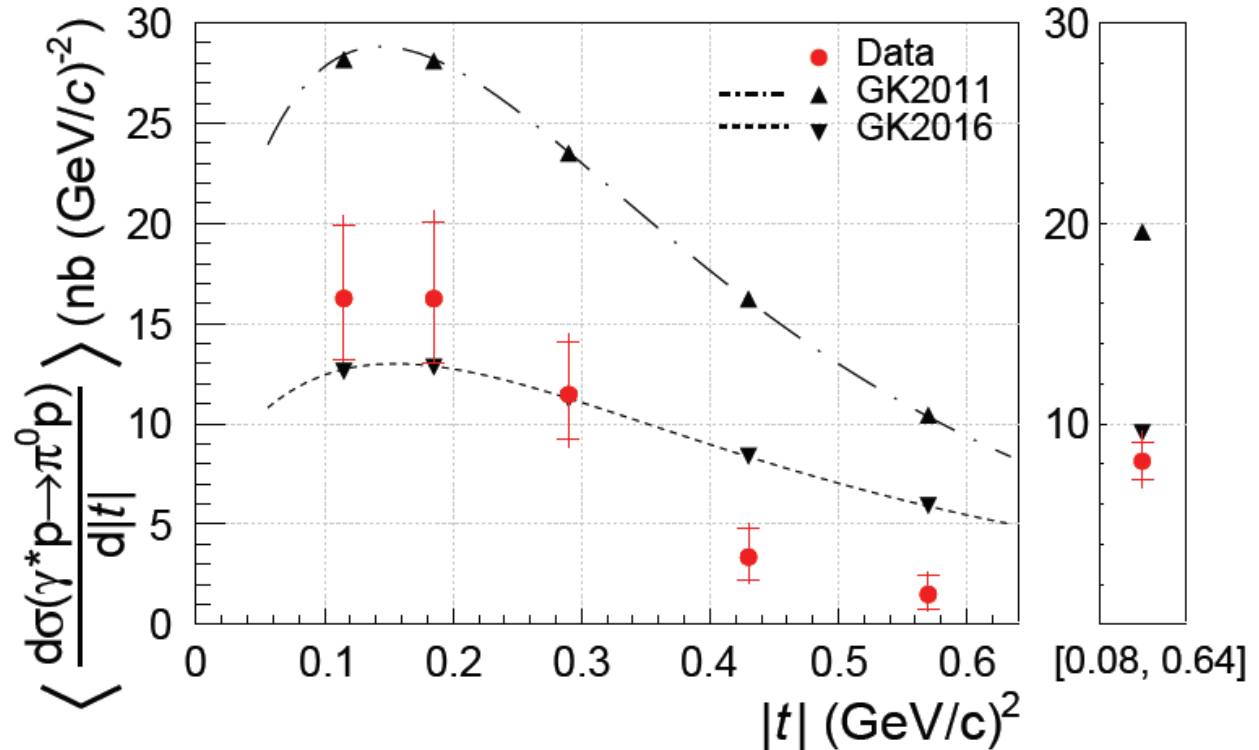
Exclusive π^0 production cross sections as a function of $|t|$

PLB 805 (2020) 135454

$$\frac{d\sigma}{dt} = \frac{d\sigma_T}{dt} + \varepsilon \frac{d\sigma_L}{dt}$$

2012 data

Dip would indicate contribution of \bar{E}_T

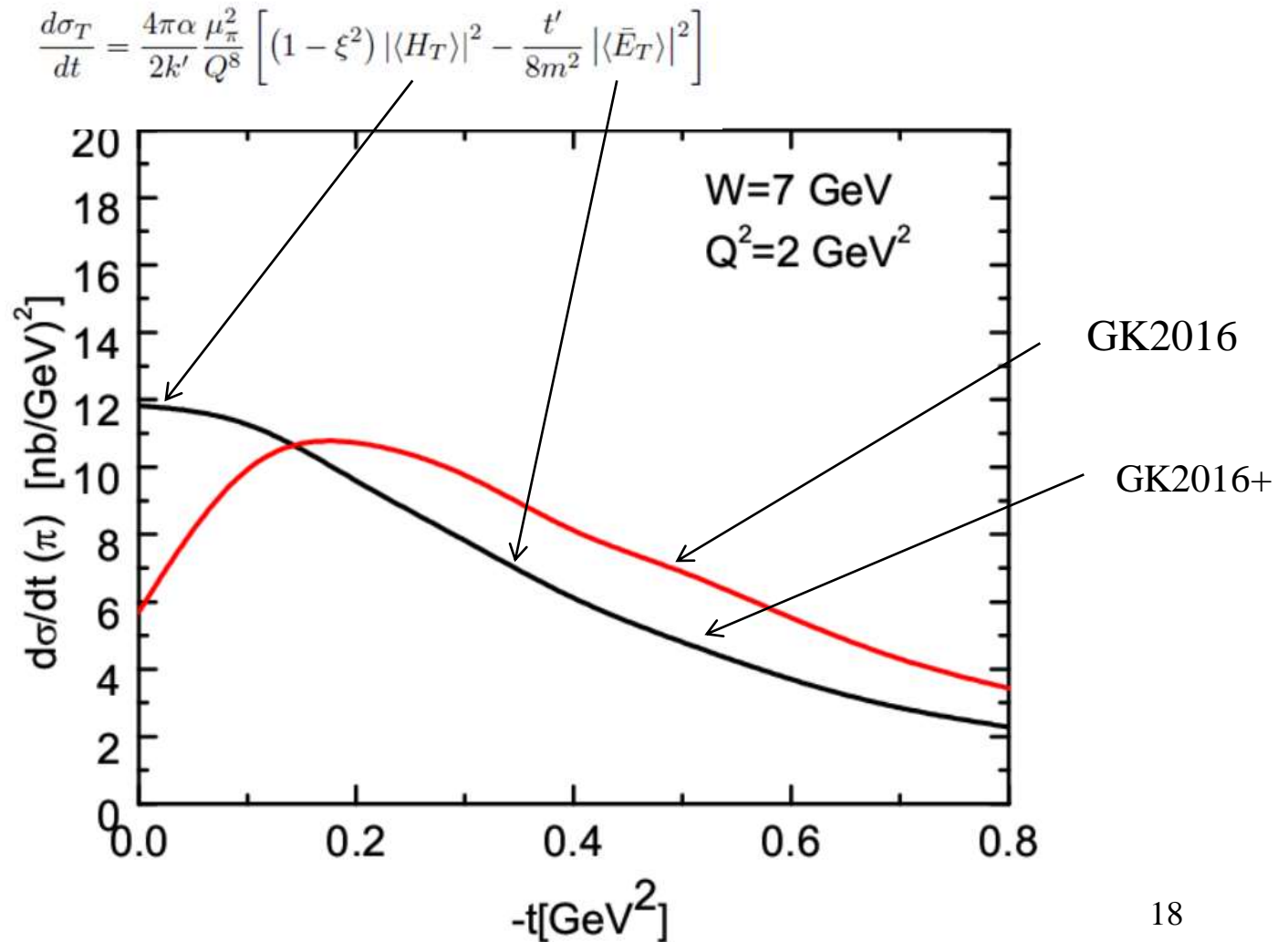


$$\langle Q^2 \rangle = 2.0 \text{ (GeV/c)}^2, \langle \nu \rangle = 12.8 \text{ GeV}, \langle x_{Bj} \rangle = 0.093 \text{ and } \langle -t \rangle = 0.256 \text{ (GeV/c)}^2$$

GK2016 + an updated model with another energy dependence of \bar{E}_T

Courtesy of S.Goloskokov & P.Kroll

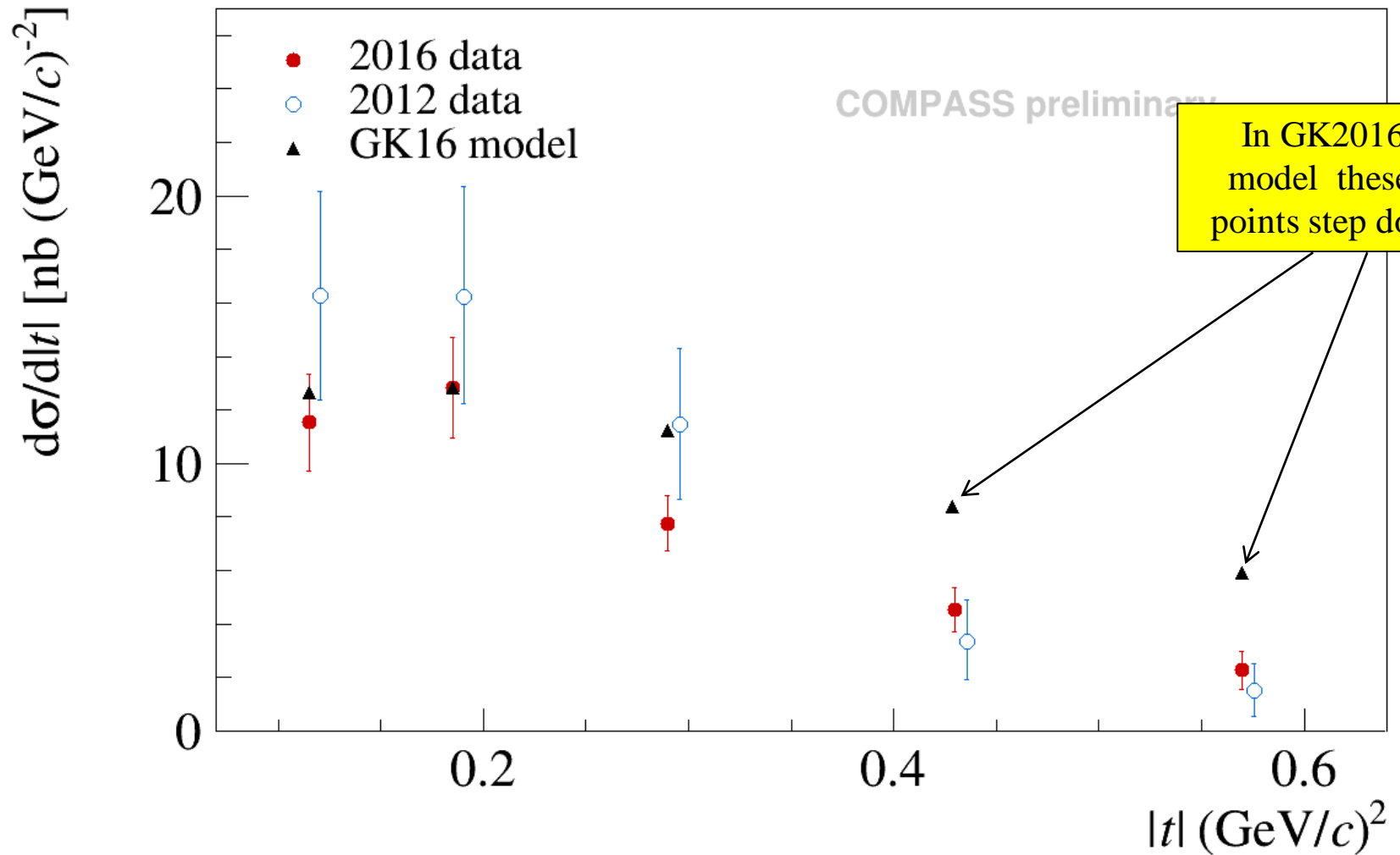
goloskqv@theor.jinr.ru



Exclusive π^0 production cross sections as a function of $|t|$

2012 data PLB 805 (2020) 135454

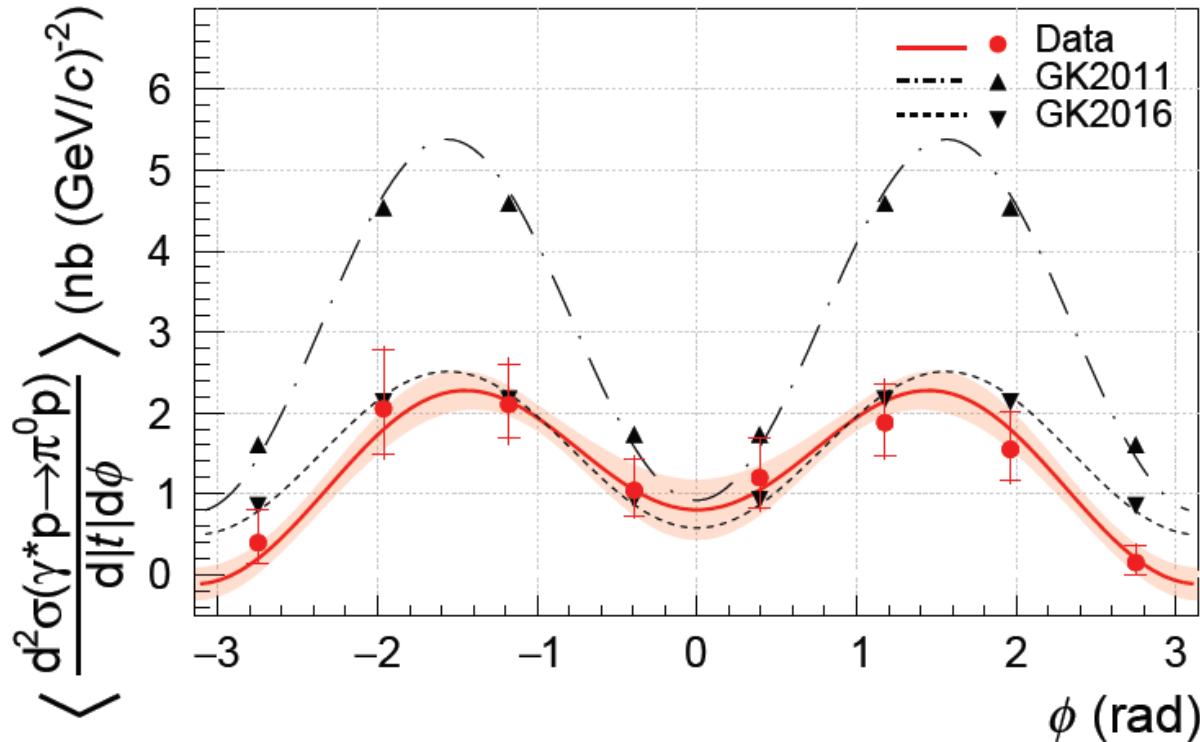
GK16 model EPJ A47 (2011) 112 + private communications



Exclusive π^0 production cross sections as a function of ϕ

PLB 805 (2020) 135454

$$\frac{d^2\sigma}{dt d\phi} = \frac{1}{2\pi} \left[\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \epsilon \cos 2\phi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi \frac{d\sigma_{LT}}{dt} \right]$$



2012 data

$$\left\langle \frac{d\sigma_T}{d|t|} + \epsilon \frac{d\sigma_L}{d|t|} \right\rangle = (8.1 \pm 0.9_{\text{stat}} + 1.1_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{TT}}{d|t|} \right\rangle = (-6.0 \pm 1.3_{\text{stat}} + 0.7_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

$$\left\langle \frac{d\sigma_{LT}}{d|t|} \right\rangle = (1.4 \pm 0.5_{\text{stat}} + 0.3_{\text{sys}}) \frac{\text{nb}}{(\text{GeV}/c)^2}$$

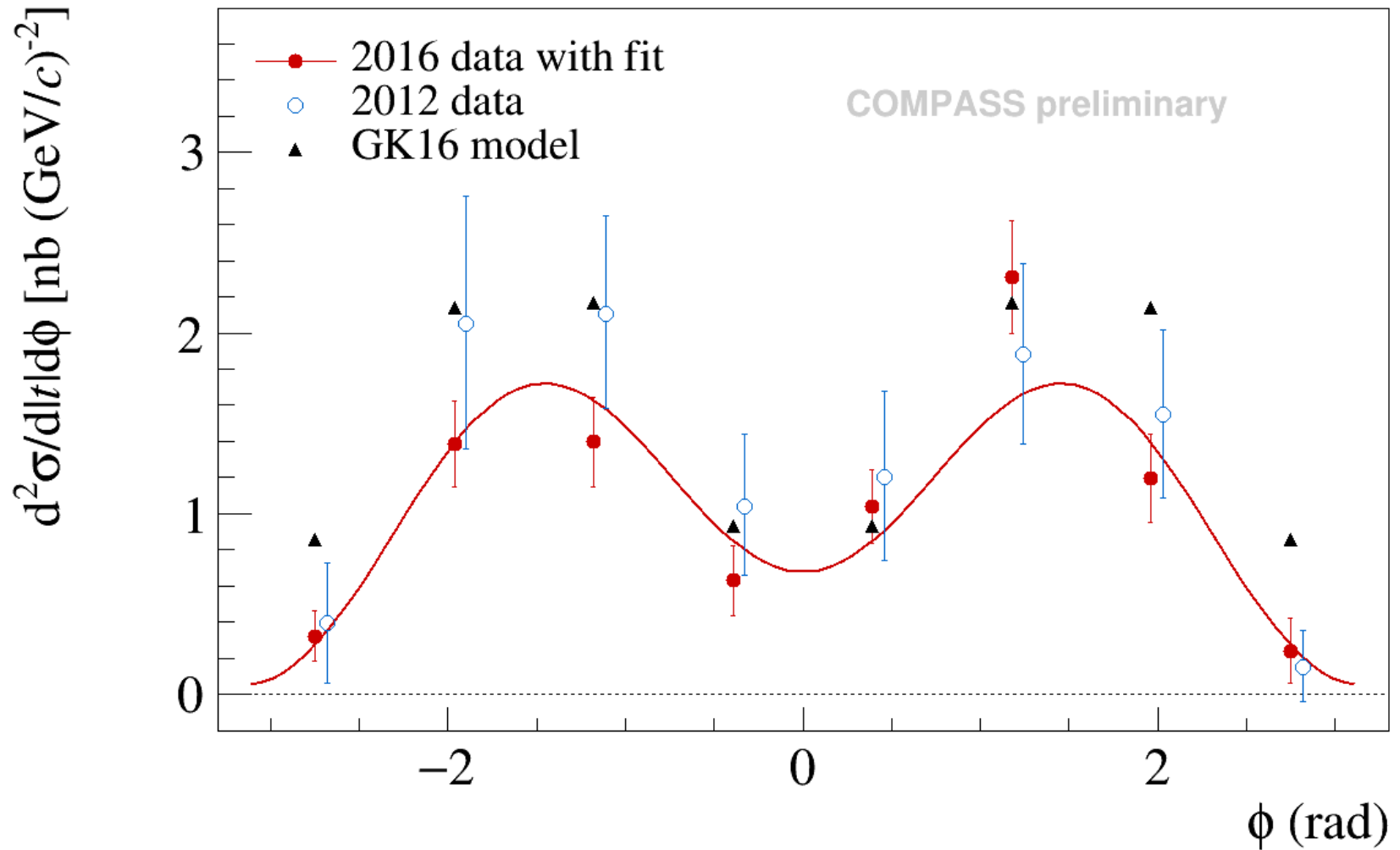
Large impact of \bar{E}_T visible in $\frac{d\sigma_{TT}}{dt} \sim \bar{E}_T$

positive result for $\frac{d\sigma_{LT}}{dt}$

Exclusive π^0 production cross sections as a function of ϕ

2012 data PLB 805 (2020) 135454

GK16 model EPJ A47 (2011) 112 + private communications



Summary

- The differential virtual-photon proton cross sections are extracted from the 2016 data as a function of the squared four-momentum transfer t , and of the azimuthal angle ϕ between the scattering plane and the π^0 production plane.
- The average differential cross sections from the 2016 data are compared to the published results of the 2012 data and no significant difference is observed.
- A slightly different t -shape is seen in 2016 data (as in 2012 one) with respect to GK2016 model prediction, which however can be reduced in the GK2016 + model with another energy dependence of \bar{E}_T
- From the results we observe a large contribution of σ_{TT} and a small positive contribution of σ_{LT} . This supports the expectation of the exclusive π^0 cross section to be dominated by transverse polarized virtual photon, which indicates a significant effect of the chiral-odd GPD \bar{E}_T
- Total statistics 2017+2016 data is estimated to be 10 times higher than the 2012 one
- Input to constrain phenomenological models such as those provided by Goloskokov&Kroll, Goldstein&Luiti, etc