

Deeply Virtual Compton Scattering at COMPASS

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on behalf of the COMPASS collaboration

September 10, 2022

1. Generalized Parton Distribution functions
2. Measurements at COMPASS and experimental setup
3. DVCS cross-section and steps for its extraction
4. Analysis of the cross section and its t -dependence

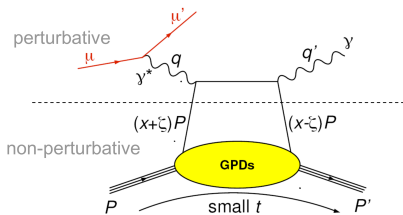
Generalized Parton Distribution functions (GPDs)

- ▶ Provide comprehensive description of 3-D partonic structure of the nucleon.
- ▶ Can be viewed as correlation functions between different partonic states.
- ▶ Allows to study decomposition of the total nucleon spin.

Generalized Parton Distribution functions (GPDs)

Deeply Virtual Compton Scattering

$$\gamma^* + N \rightarrow \gamma + N'$$



$q = (p_\mu - p_{\mu'})$: 4-mom. of virt. photon

$Q^2 = -q^2$: virt. photon virtuality

$t = (p_P - p_{P'})^2$: mom. transfer to nucleon squared

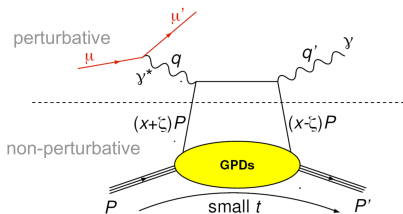
x : avg. long. mom. fraction

ξ : half of long. mom. fraction transfer

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No nucleon spin flip

$$H^f(x, \xi, t)$$

$$\tilde{H}^f(x, \xi, t)$$

With nucleon spin flip

$$E^f(x, \xi, t)$$

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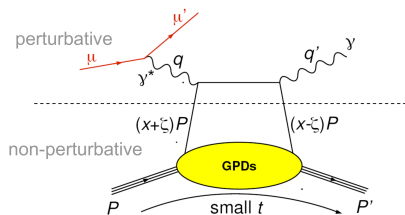
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Generalized Parton Distribution functions (GPDs)

Deeply Virtual Compton Scattering

$$\gamma^* + N \rightarrow \gamma + N'$$



GPDs are not experimentally accessible, but related to Compton Form Factors (CFFs)

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$$\tilde{H}^f(x, \xi, t)$$

With nucleon spin flip

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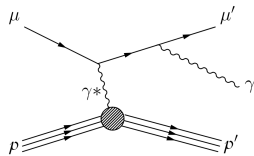
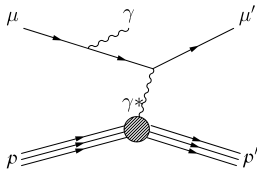
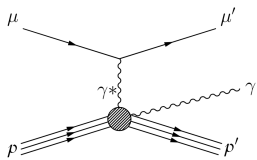
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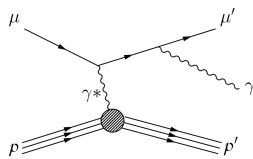
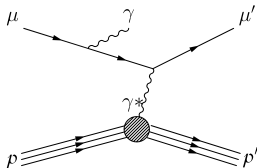
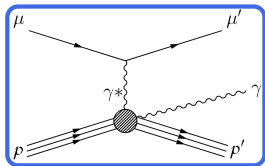
CFFs are observables in cross section measurements

$$\mathcal{H}(\xi, t) = \int_{-1}^1 \frac{H(x, \xi, t)}{x - \xi - i\epsilon} dx$$

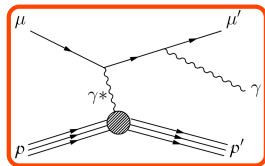
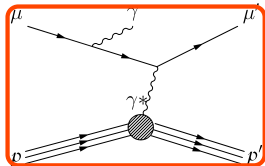
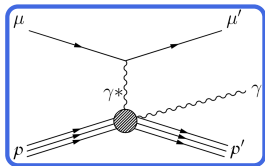
$\mu p \rightarrow \mu' p' \gamma$ processes



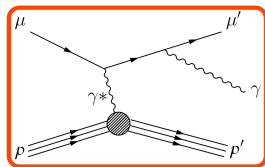
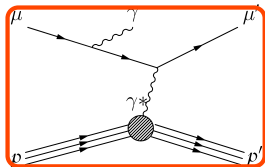
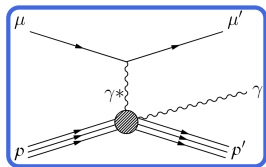
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The observable products of these reactions are identical, therefore they can not be separated on per-event basis.

$$\sigma_{\gamma} \propto |A_{\text{DVCS}}|^2 + |A_{\text{BH}}|^2 + \text{Interference Term}$$

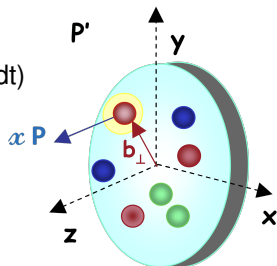
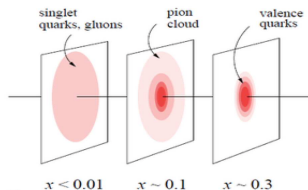
Interference Term $\equiv A_{\text{DVCS}} A_{\text{BH}}$ - allows to study DVCS on the amplitude level.

Most attractive goals

3D tomography via GPD H

$$H(x, \xi = 0, t) \rightarrow H(x, b_{\perp}) \sim \rho(x, b_{\perp})$$

probability interpretation (Burkardt)

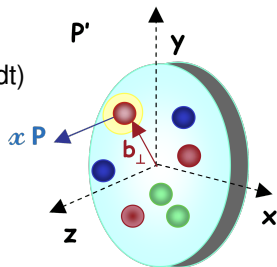
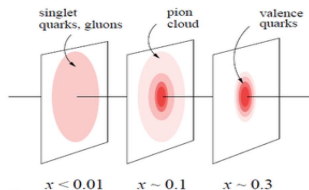


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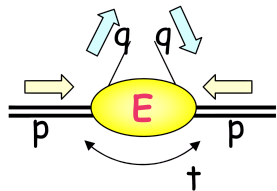


Contribution to the nucleon spin puzzle

$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta G + \langle L_z^q \rangle + \langle L_z^g \rangle$$

by constraining GPD H and E

$$J^q = \frac{1}{2} \lim_{t \rightarrow 0} \int_{-1}^{+1} x [H^q(x, \xi, t) + E^q(x, \xi, t)] dx$$



GPD E related to the orbital angular momentum

Measurement at COMPASS

$$\text{Diff. cross section } \frac{d\sigma^4}{dQ^2 d\nu d|t| d\phi}$$

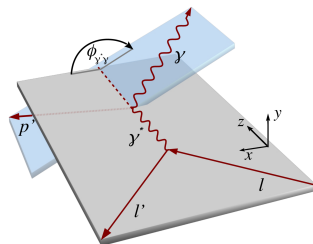
Kinematic dependence:

$Q^2 = -q^2$: virtual photon virtuality

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ϕ : angle between scattering and production planes



$$\mu p \rightarrow \mu' p' \gamma$$

Identify exclusive photon events:

μ : beam muon

μ' : scattered muon

P' : recoil proton

γ : real photon

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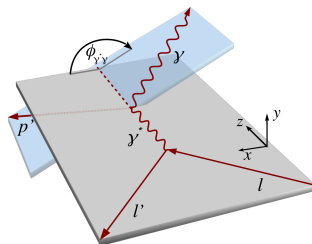
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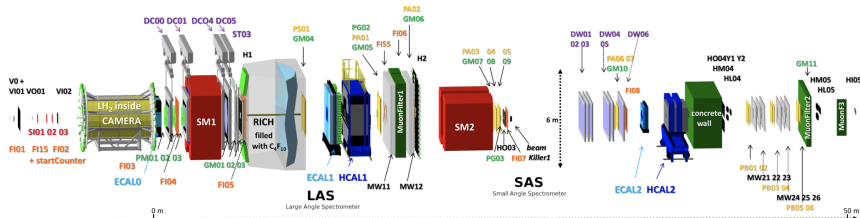
► 2012 pilot run for 4 weeks

→ analysis finished and published

► 2016/17 long runs (2 × 6 months) dedicated to DVCS

→ analysis ongoing, preliminary results

COMPASS experiment setup



Common Muon and Proton Apparatus for Structure and Spectroscopy

2.5m long Liquid Hydrogen target

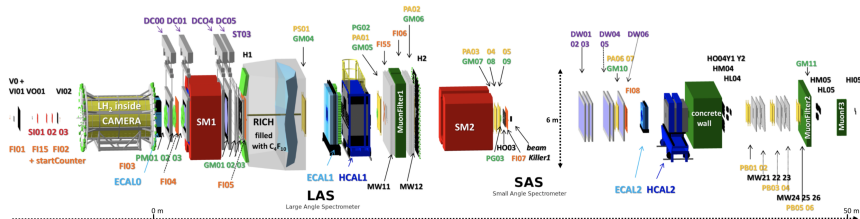
Beam energy is 140 – 180 GeV

Beam polarisations: $\mu^{+\downarrow}$ and $\mu^{-\uparrow}$

$$\pi^+ \rightarrow \mu^+ + \nu_\mu \quad P_{\mu^+} \approx -80\%$$

$$\pi^- \rightarrow \mu^- + \bar{\nu}_\mu \quad P_{\mu^-} \approx +80\%$$

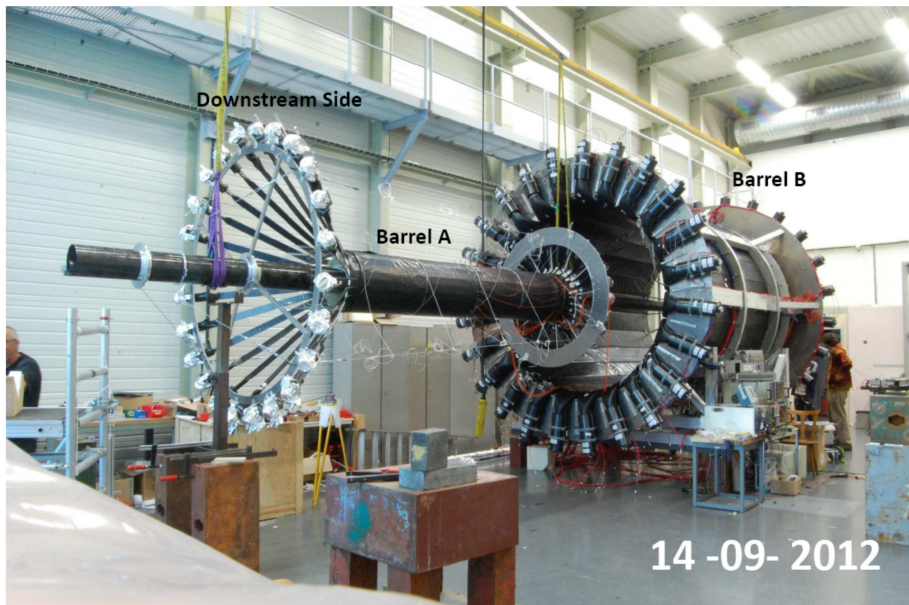
COMPASS experiment setup



Two stage forward spectrometer **SM1** + **SM2**

- ▶ Beam flux determined using true Random trigger, $\approx 1\%$ precision
- ▶ **ECAL0**, **ECAL1** and **ECAL2** for photon detection
- ▶ Muon trigger system ~ 300 tracking detector planes
- ▶ μ ID for muon identification

CAMERA



Selection of events with a single γ topology

Vertex candidates:

μ Beam muon

μ' Scattered muon

Real photon candidate γ :

Single photon with the energy above DVCS threshold $E_\gamma > 4, 5, 10$ GeV in ECAL0, 1, 2

Recoil proton candidate P' :

$$|t|_{\max}^{\text{exp}} = 0.64(\text{GeV}/c)^2$$

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Exclusivity selections:

- $|\Delta p_T| < 0.3 \text{ GeV}/c$
- $|\Delta\phi| < 0.4 \text{ rad}$
- $|\Delta z_A| < 16 \text{ cm}$
- $|M_{\text{Undet}}^2| < 0.3(\text{GeV}/c^2)^2$
- $0.08(\text{GeV}/c)^2 < |t_{\text{fit}}| < 0.64(\text{GeV}/c)^2$

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Perform kinematic fit:

- constrain on kinematic variables $\chi^2 < 10$

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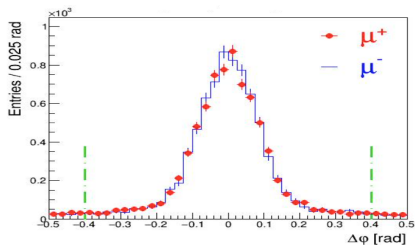
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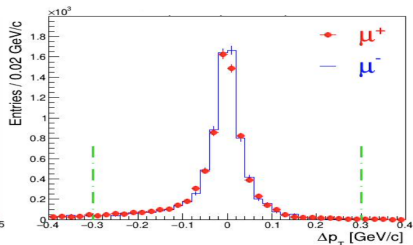
Only events with single valid combination

Vertex candidate \times Real photon candidate \times Recoil proton candidate

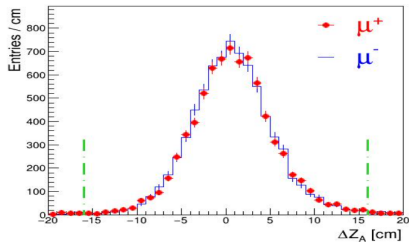
Exclusive selections (COMPASS 2016 preliminary results)



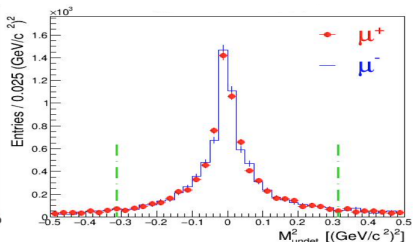
$$\Delta\phi = \phi^{\text{cam}} - \phi^{\text{spect}}$$



$$\Delta p_T = |p_T^{\text{cam}}| - |p_T^{\text{spect}}|$$



$$\Delta Z_A = Z_A^{\text{cam}} - Z_A^{\text{extrapolated}}$$

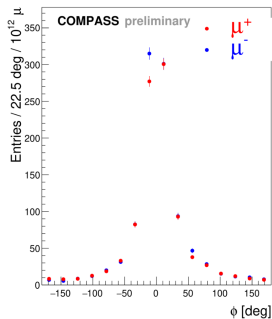


$$M_{\text{undet}}^2 = (p_\mu + p_P - p_{\mu'} - p_{P'} - p_\gamma)^2$$

ϕ distribution of single γ events at different ν

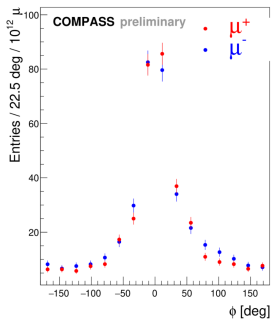
Bethe-Heitler dominant

$80 < \nu \text{ [GeV]} < 144$



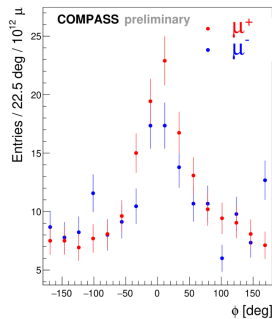
Interference

$32 < \nu \text{ [GeV]} < 80$



DVCS dominant

$10 < \nu \text{ [GeV]} < 32$



Number of reconstructed single-photon events as a function of angle ϕ .

left. $80\text{GeV} < \nu < 144\text{GeV}$

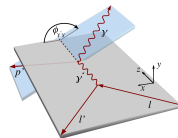
BH dominates, DVCS - negligible

middle. $32\text{GeV} < \nu < 80\text{GeV}$

BH and DVCS are comparable

right. $10\text{GeV} < \nu < 32\text{GeV}$

DVCS dominates BH



The binned DVCS cross section

DVCS cross section in bins of t , ϕ , Q^2 , ν :

$$\left\langle \frac{d\sigma_{\text{DVCS}}}{d|t|d\phi dQ^2 d\nu} \right\rangle_{t_i \phi_j Q_k^2 \nu_l}^{\pm} = \frac{1}{\mathcal{L}^{\pm} \Delta t_i \Delta \phi_j \Delta Q_k^2 \Delta \nu_l} \left[\left(a_{ijkl}^{\pm} \right)^{-1} \left(\text{data} - \text{BH}_{\text{MC}} - \pi_{\text{MC}}^0 \right) \right]$$

a_{ijkl}^{\pm} Acceptance

BH_{MC} Exclusive single photon MC sample

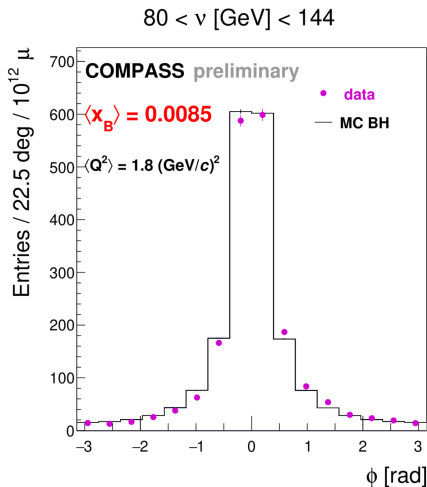
π_{MC}^0 π^0 MC sample (background estimation)

The Bethe-Heitler contribution

Bethe-Heitler process is well known, pure QED

→ evaluated using **Monte-Carlo sample** for BH

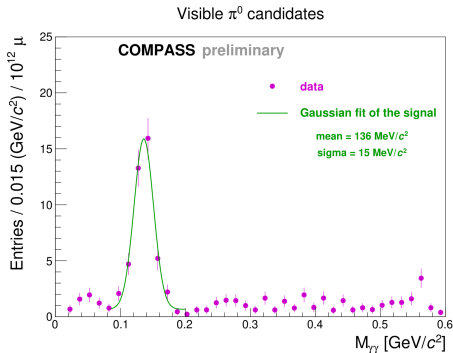
- ▶ Kinematic range where **BH is dominant**
→ **use data/MC luminosities for absolute normalization**
- ▶ **BH subtracted** from the data in the DVCS region (small ν)



The π^0 background contamination

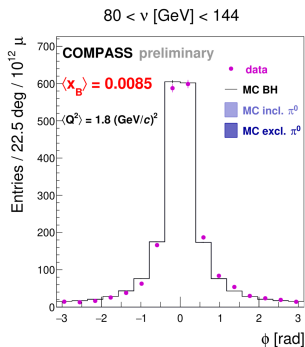
- ▶ Photons from π^0 decay
- ▶ One photon identified as exclusive photon event
→ above DVCS energy threshold in ECALS

- ▶ **Visible** (both γ are detected)
- *subtracted*
Combine γ_{he} and γ_{1e} (below DVCS energy threshold)
- ▶ **Invisible** (second γ lost)
- *estimated by MC*
 - ▶ **Inclusive:** LEPTO
 - ▶ **Exclusive:** HEPGEN π^0

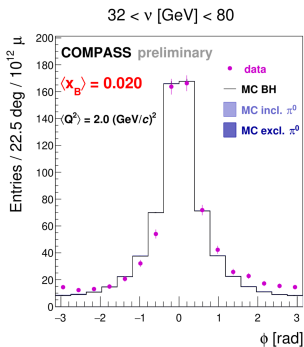


ϕ distribution of exclusive photon events

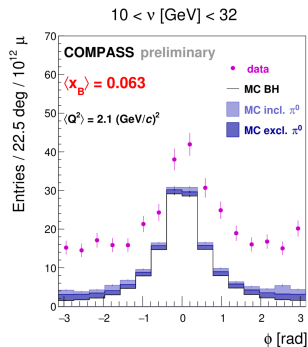
$$1 < Q^2 < 10(\text{GeV}/c)^2$$



▶ 64% of events in data



▶ 24% of events in data



- ▶ 12% of events in data
- ▶ 37% BH contribution
- ▶ 10% inv. π^0 contribution

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$$\pi_{\text{MC}}^0 = (1 - R) \times \pi_{\text{HEPGEN}}^0 + R \times \pi_{\text{LEPTO}}^0$$

► BH_{MC} : BH MC sample

- π_{HEPGEN}^0 : exclusive π^0 MC sample
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▶ R : relative contrib. of LEPTO ($\approx 40\%$)

▶ a_{ijkl}^{\pm} : acceptance

Acceptance studies limit to region with mostly flat acceptance

avg. acc. $\approx 40\%$, good agreement between μ^+ and μ^-

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- ▶ π_{HEPGEN}^0 : exclusive π^0 MC sample
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Binning and kinematic range:

- ▶ 4 bins in $|t|$ between 0.08 and 0.64 (GeV/c)² (equistatics)
- ▶ 4 bins ν between 10 and 32 GeV (equidistant)
- ▶ 4 bins Q^2 between 1 and 5 (GeV/c)² (equidistant)
- ▶ 8 bins ϕ between $-\pi$ and $+\pi$ (equidistant)

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Calculate the t -dependence of the cross section

From μp to $\gamma^* p$:
$$\frac{d\sigma^{\mu p}}{dt d\phi dQ^2 d\nu} = \Gamma(Q^2, \nu) \times \frac{d\sigma^{\gamma^* p}}{dt d\phi dQ^2 d\nu}$$

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by **weighting each event** in data and MC by the inverse **flux of the transverse polarized photons**

$$\Gamma(Q^2, \nu) = \frac{\alpha_{EM} (1 - x_{Bj})}{2\pi Q^2 y E} \left[y^2 \left(1 - 2 \frac{2m_\mu^2}{Q^2} \right) + \frac{2}{1 + (Q^2/\nu^2)} \left(1 - y - \frac{Q^2}{4E^2} \right) \right]$$

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t-dependence for μ^+ and μ^-

$$\left\langle \frac{d\sigma_{\text{DVCS}}}{d|t|} \right\rangle_{t_i}^{\pm} = (\Delta Q^2 \Delta \nu)^{-1} \sum_{k,l} \left\langle \frac{d\sigma_{\text{DVCS}}}{d|t| dQ^2 d\nu} \right\rangle_{t_i Q_k^2 \nu_l}^{\pm} \Delta Q_l^2 \Delta \nu_l$$

Calculate the t -dependence of the cross section

$$\text{From } \mu p \text{ to } \gamma^* p: \quad \frac{d\sigma^{\mu p}}{dt d\phi dQ^2 d\nu} = \Gamma(Q^2, \nu) \times \frac{d\sigma^{\gamma^* p}}{dt d\phi dQ^2 d\nu}$$

by **weighting each event** in data and MC by the inverse **flux of the transverse polarized photons**

$$\Gamma(Q^2, \nu) = \frac{\alpha_{EM} (1 - x_{Bj})}{2\pi Q^2 y E} \left[y^2 \left(1 - 2 \frac{2m_\mu^2}{Q^2} \right) + \frac{2}{1 + (Q^2/\nu^2)} \left(1 - y - \frac{Q^2}{4E^2} \right) \right]$$

t-dependence for μ^+ and μ^-

$$\left\langle \frac{d\sigma_{\text{DVCS}}}{d|t|} \right\rangle_{t_i}^{\pm} = (\Delta Q^2 \Delta \nu)^{-1} \Sigma_{k,l} \left\langle \frac{d\sigma_{\text{DVCS}}}{d|t| dQ^2 d\nu} \right\rangle_{t_i Q_k^2 \nu_l}^{\pm} \Delta Q_l^2 \Delta \nu_l$$

→ Integrating over ϕ dependence **removes interference and ϕ -dependent DVCS contribution**

$$\mathcal{S}_{CS,U} = d\sigma^{+\downarrow} + d\sigma^{-\uparrow} = 2 \left[d\sigma^{\text{BH}} + c_0^{\text{DVCS}} + c_1^{\text{DVCS}} \cos\phi + c_2^{\text{DVCS}} \cos 2\phi + s_1^{\text{I}} \sin\phi + s_2^{\text{I}} \sin 2\phi \right]$$

Calculate the t -dependence of the cross section

From μp to $\gamma^* p$:
$$\frac{d\sigma^{\mu p}}{dt dt \phi dQ^2 d\nu} = \Gamma(Q^2, \nu) \times \frac{d\sigma^{\gamma^* p}}{dt dt \phi dQ^2 d\nu}$$

by **weighting each event** in data and MC by the inverse **flux of the transverse polarized photons**

$$\Gamma(Q^2, \nu) = \frac{\alpha_{EM} (1 - x_{Bj})}{2\pi Q^2 y E} \left[y^2 \left(1 - 2 \frac{2m_\mu^2}{Q^2} \right) + \frac{2}{1 + (Q^2/\nu^2)} \left(1 - y - \frac{Q^2}{4E^2} \right) \right]$$

t-dependence for μ^+ and μ^-

$$\left\langle \frac{d\sigma_{\text{DVCS}}}{d|t|} \right\rangle_{t_i}^{\pm} = (\Delta Q^2 \Delta \nu)^{-1} \Sigma_{k,l} \left\langle \frac{d\sigma_{\text{DVCS}}}{d|t| dQ^2 d\nu} \right\rangle_{t_i Q_k^2 \nu_l}^{\pm} \Delta Q_l^2 \Delta \nu_l$$

→ Integrating over ϕ dependence **removes interference and ϕ -dependent DVCS contribution**

$$\begin{aligned} S_{CS,U} &= d\sigma^{+\downarrow} + d\sigma^{-\uparrow} = \\ &2 \left[d\sigma^{\text{BH}} + c_0^{\text{DVCS}} + c_1^{\text{DVCS}} \cos\phi + c_2^{\text{DVCS}} \cos 2\phi + s_1^{\text{I}} \sin\phi + s_2^{\text{I}} \sin 2\phi \right] \end{aligned}$$

t-dependence of the cross section:

$$\left\langle \frac{d\sigma_{\text{DVCS}}}{d|t|} \right\rangle_{t_i} = \frac{1}{2} \left(\left\langle \frac{d\sigma_{\text{DVCS}}}{d|t|} \right\rangle_{t_i}^+ + \left\langle \frac{d\sigma_{\text{DVCS}}}{d|t|} \right\rangle_{t_i}^- \right)$$

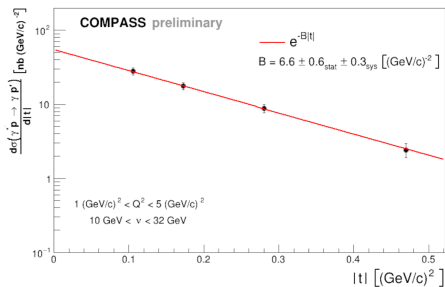
Analyse the cross section t -slope

$$d\sigma^{\text{DVCS}}/dt \sim e^{-B|t|} \propto c_0^{\text{DVCS}} = (\text{Im}\mathcal{H})^2$$

Perform binned maximum Likelihood-fit.

$$B = (6.6 \pm 0.6_{\text{stat}} \pm 0.3_{\text{sys}}) (\text{GeV}/c)^{-2}$$

Dominant source of systematics:
MC normalisation to visible π^0 in data.



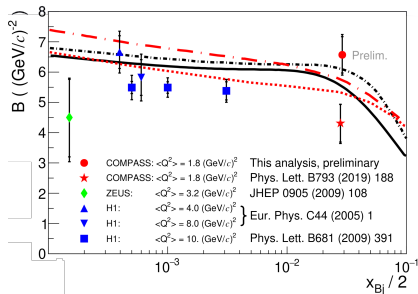
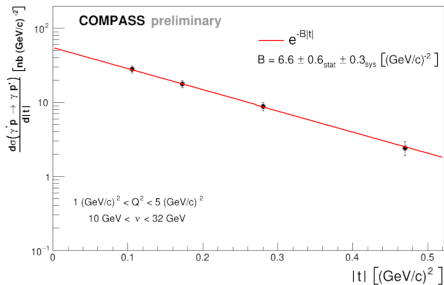
Analyse the cross section t -slope

$$d\sigma^{\text{DVCS}}/dt \sim e^{-B|t|} \propto c_0^{\text{DVCS}} = (\text{Im}\mathcal{H})^2$$

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Dominant source of systematics:
MC normalisation to visible π^0 in data.



2012 results PLB 793 (2019) 188

$$B = (4.3 \pm 0.6_{\text{stat}} \pm 0.1_{\text{sys}}) (\text{GeV}/c)^{-2}$$

- ▶ **Analyse full statistics** of 2016 and 2017 (3 times more data than 2016)
- ▶ More detailed studies of **systematic uncertainties**
- ▶ Cross section study in few x_{Bj} regions \rightarrow tomography
- ▶ **Study the azimuthal dependence** of the cross section
 \rightarrow Determine c_0^{DVCS} , c_1^{DVCS} , c_2^{DVCS} , s_1^{I} and s_2^{I}
- ▶ **Cross section difference** $\mathcal{D}_{CS,U} = d\sigma^{+\downarrow} - d\sigma^{-\uparrow}$
 \rightarrow Access to $Re\mathcal{H}$

Thank you for your attention!

Backup

Cross section

$$d\sigma \propto |A_{BH}|^2 + |A_{DVCS}|^2 + \text{Interference Term}$$

$$\frac{d^4\sigma(lp \rightarrow lp\gamma)}{dx_{Bj}dQ^2d|t|d\phi} = d\sigma^{BH} + \left(d\sigma_{unpol}^{DVCS} + P_l d\sigma_{pol}^{DVCS} \right) + \left(e_l \text{Re}I + e_l P_l \text{Im}I \right)$$

$$\mathcal{D}_{CS,U} = d\sigma^{+\downarrow} - d\sigma^{-\uparrow} = 2 \left[e_\mu a^{BH} \text{Re}A^{DVCS} + |P_\mu| d\sigma_{pol}^{DVCS} + \text{Re}I \right]$$

$$= 2 \left[s_1^{DVCS} \sin \phi + c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi \right]$$

$$\mathcal{S}_{CS,U} = d\sigma^{+\downarrow} + d\sigma^{-\uparrow} = 2 \left[d\sigma^{BH} + d\sigma_{unpol}^{DVCS} - |P_\mu| \text{Im}I \right]$$

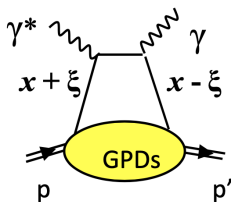
$$= 2 \left[d\sigma^{BH} + c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi + s_1^I \sin \phi + s_2^I \sin 2\phi \right]$$

Cross section

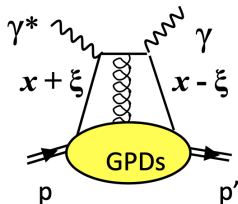
$$S_{CS,U} = d\sigma^{+\downarrow} + d\sigma^{-\uparrow}$$

$$D_{CS,U} = d\sigma^{+\downarrow} - d\sigma^{-\uparrow}$$

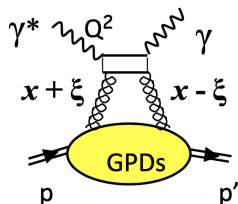
$S_{CS,U}$	$d\sigma^{BH}$	\propto	$c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$
$S_{CS,U}$	$d\sigma_{unpol}^{DVCS}$	\propto	$c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$
$D_{CS,U}$	$d\sigma_{pol}^{DVCS}$	\propto	$s_1^{DVCS} \sin \phi$
$D_{CS,U}$	ReI	\propto	$c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$
$S_{CS,U}$	ImI	\propto	$s_1^I \sin \phi + s_2^I \sin 2\phi$



LO, Twist-2



LO, Twist-3



NLO, Twist-2