



Fit of the $a_1(1420)$ as a Triangle Singularity

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on behalf of the COMPASS collaboration

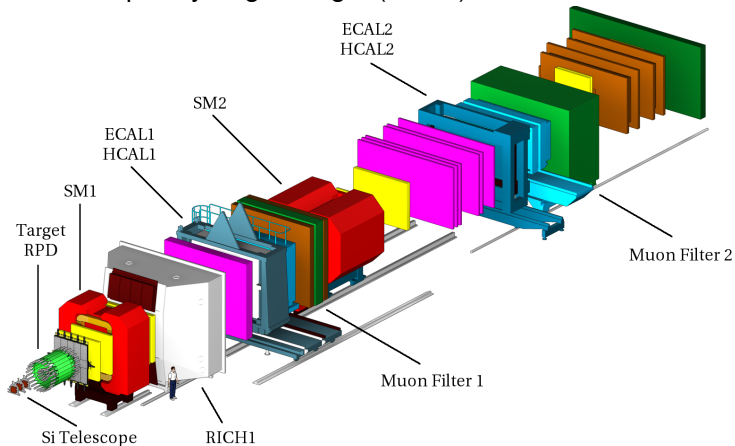
HISKP, Bonn University

September 3, 2021

DPG 2021, HK 29.7

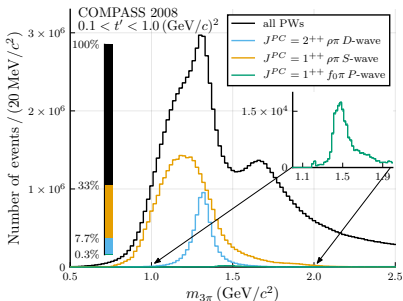
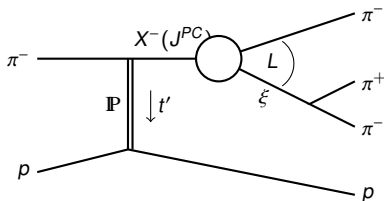
supported by BMBF

- Here: secondary hadron beam, mostly π^- ($\sim 97\%$)
- $E_{\text{beam}} = 190 \text{ GeV}$
- Fixed liquid hydrogen target (40 cm)



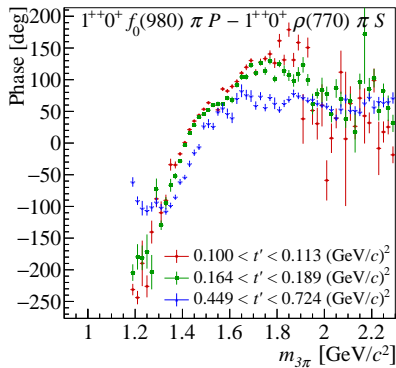
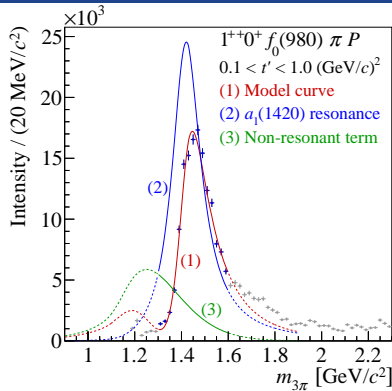
[COMPASS, NIM A779, 69-115 (2015)]

- Here: secondary hadron beam, mostly π^- ($\sim 97\%$)
- $E_{\text{beam}} = 190 \text{ GeV}$
- Fixed liquid hydrogen target (40 cm)
- $\pi^- + p \rightarrow \pi^- + \pi^- + \pi^+ + p$
- PWA with 88 waves: data binned in $m_{3\pi}, t'$



[COMPASS, PRL **127**, 082501 (2021)]

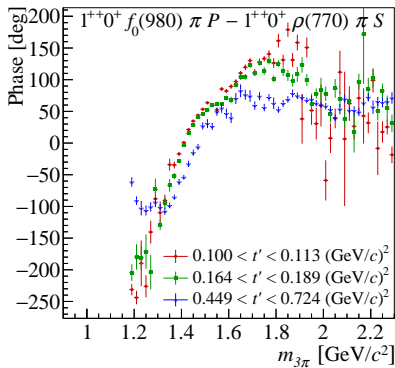
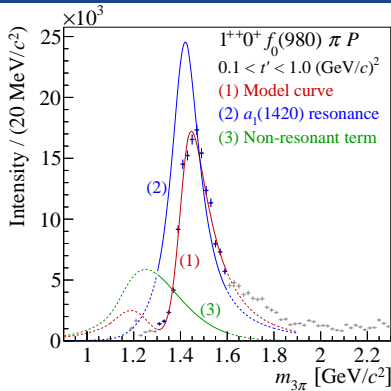
BW-fit to resonance-like signal in 1^{++} partial wave



[COMPASS, PRL **115**, 082001 (2015)]

- $a_1(1420)$ narrow peak, strong phase motion
- No ordinary radial excitation of $a_1(1260)$
 - Mass too close to ground state
 - Narrower than ground state
 - Peculiar decay mode

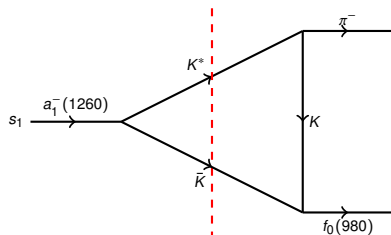
BW-fit to resonance-like signal in 1^{++} partial wave



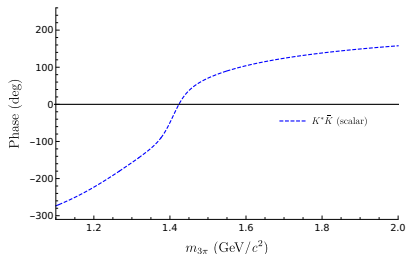
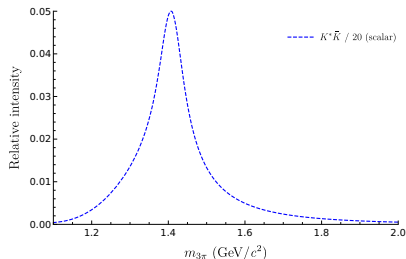
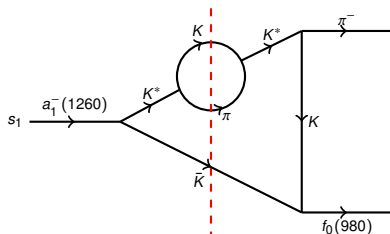
[COMPASS, PRL **115**, 082001 (2015)]

- 4-quark state [H.-X. Chen et al. (2015)], [T. Gutsche et al. (2017)]
- $K^* \bar{K}$ molecule (similar to $X(3872)$) [T. Gutsche et al. (2017)]
- Dynamic interference effect of $a_1(1260)$ with Deck-amplitude [Basdevant & Berger, PRL **114**, 192001 (2015)]
- **Triangle singularity (TS)** [Mikhasenko et al., PRD **91**, 094015 (2015)]
[Aceti et al., PRD **94**, 096015 (2016)]

- Dispersive approach

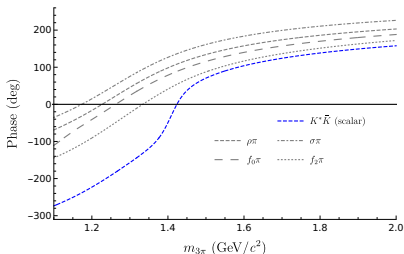
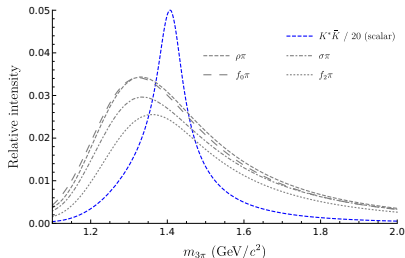
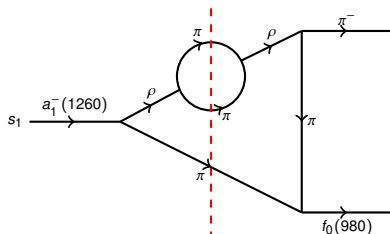


- Dispersive approach
- Include finite width of K^*



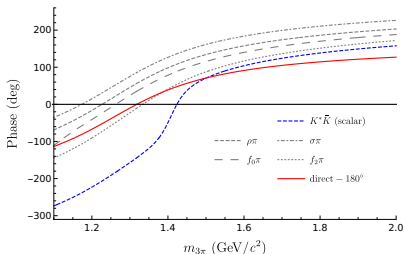
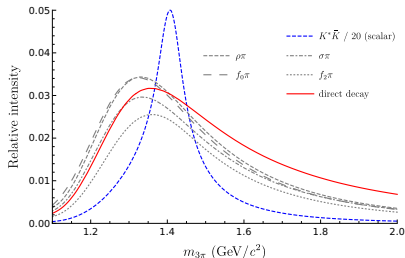
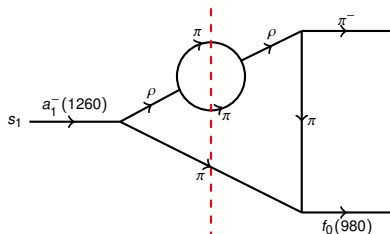
[COMPASS, PRL **127**, 082501 (2021)]

- Dispersive approach
- Include finite width of K^*
- Negligible contribution from other triangles



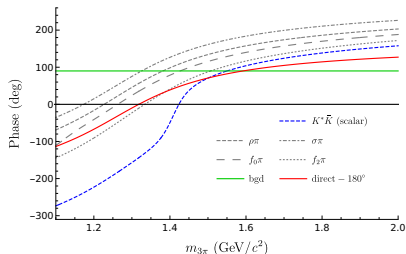
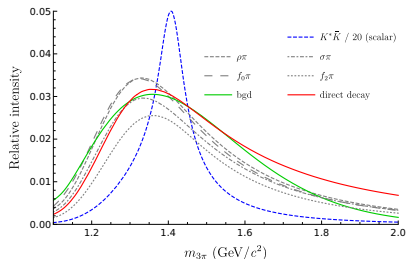
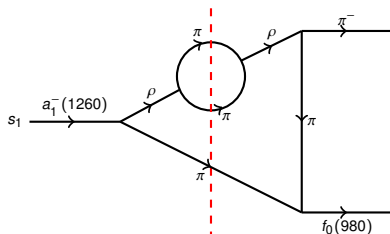
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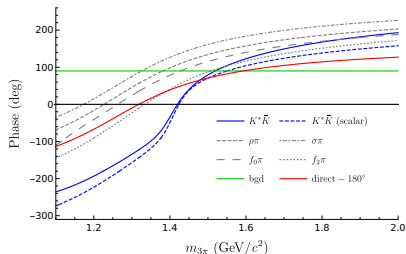
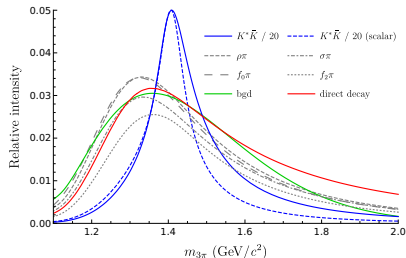
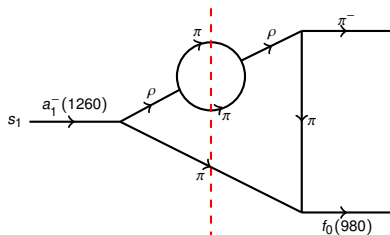
[COMPASS, PRL **127**, 082501 (2021)]

- Dispersive approach
- Include finite width of K^*
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[COMPASS, PRL **127**, 082501 (2021)]

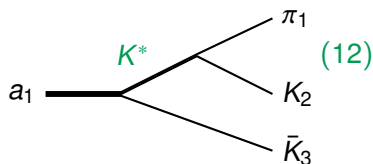
- Dispersive approach
- Include finite width of K^*
- Negligible contribution from other triangles
- Inclusion of spin distorts shape



[COMPASS, PRL **127**, 082501 (2021)]

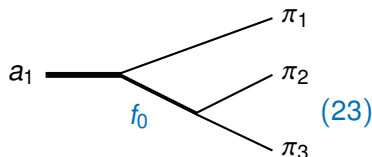
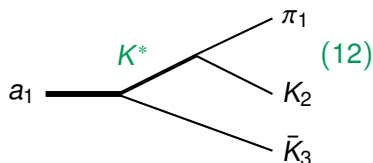
Include spin via partial-wave projection:

1. Look at the partial wave for $a_1(1260) \rightarrow K\bar{K}\pi$ with isobar K^*



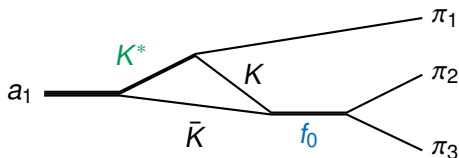
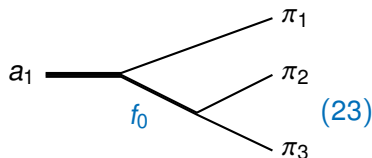
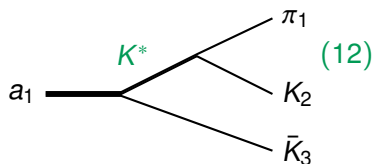
Include spin via partial-wave projection:

1. Look at the partial wave for $a_1(1260) \rightarrow K\bar{K}\pi$ with isobar K^*
2. Project it onto the 3π final state with isobar $f_0(980)$

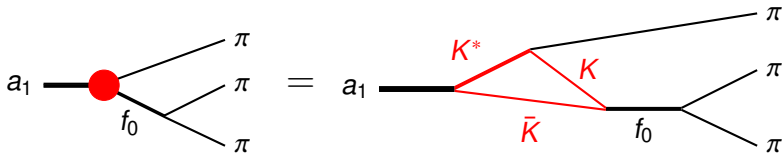


Include spin via partial-wave projection:

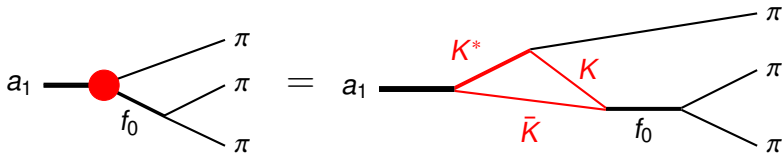
1. Look at the partial wave for $a_1(1260) \rightarrow K\bar{K}\pi$ with isobar K^*
2. Project it onto the 3π final state with isobar $f_0(980)$
3. Obtain the first order approximation of the Khuri-Treiman approach



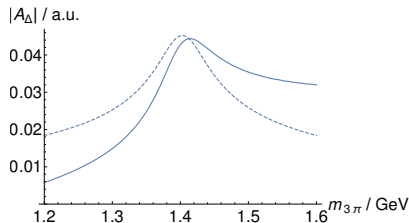
$$F(s_{23}) = C_{a_1} t_{f_0}(s_{23}) \frac{1}{2\pi} \int_{4m_K^2}^{\infty} d\tilde{s}_{23} \frac{\rho(\tilde{s}_{23}) \int dZ_{f_0}(\tilde{s}_{23}) t_{K^*}(s_{12}) Z_{K^*}^*(\Omega_{3,12})}{\tilde{s}_{23} - s_{23}}$$



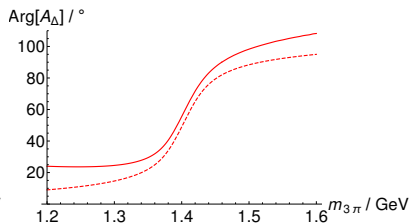
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— vector - - - scalar



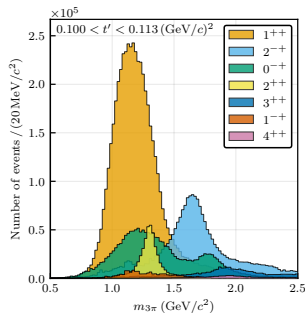
— vector - - - scalar



Minimal fit model \leadsto choose 3 of the 88 waves of the PWA

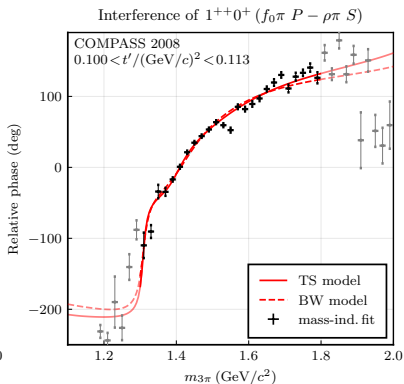
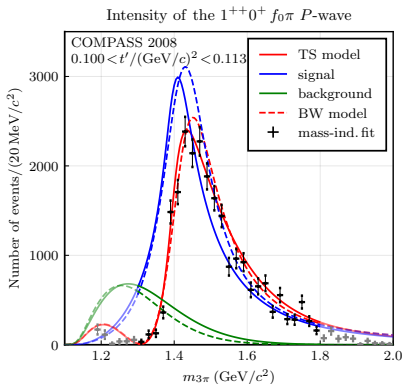
Notation: $J^{PC} M^{\xi} \pi L$

- $1^{++} 0^{+} \rho\pi$ *S*-wave:
Contains source $a_1(1260)$, but
huge non-res. background
- $1^{++} 0^{+} f_0(980)\pi$ *P*-wave:
Signal of interest $a_1(1420)$
- $2^{++} 1^{+} \rho\pi$ *D*-wave:
Clean $a_2(1320)$ with almost no
non-res. background



[B. Ketzner, B. Grube, D. Ryabchikov,
PPNP **113**, 0146-6410 (2020)]

Note: Fit all t' -slices with common resonance parameters.
Show only fit of first slice.

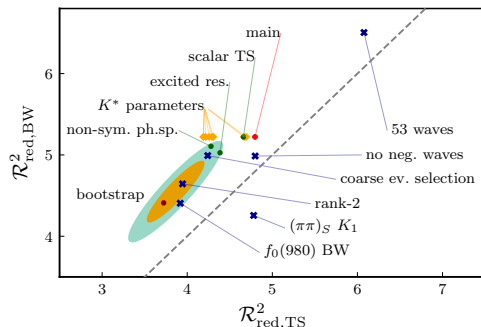


[COMPASS, PRL **127**, 082501 (2021)]

- Comparison between TS model (solid) and BW model (dashed)
- Similar fit quality

$$\text{Compare } \mathcal{R}_{\text{red}}^2 = \sum_i \left(\frac{y_i - f(x_i)}{\sigma_i} \right)^2,$$

but sum only over $f_0\pi$ P -intensity and its rel. phase to $\rho\pi$ S



$$--- \mathcal{R}_{\text{red,TS}}^2 = \mathcal{R}_{\text{red,BW}}^2$$

● main fit

× syst. studies of PWA

● syst. studies of model

◆ changing K^* resonance parameters

1σ and 2σ ellipses for bootstrap of data points

[COMPASS, PRL **127**, 082501 (2021)]

(Almost) all studies show a better fit quality for the TS model.

Conclusion:

- Reproduce features with scalar approximation
- ⇒ Good starting point for first investigation
- $a_1(1420)$ fully explainable with rescattering
- Similar fit quality as with Breit-Wigner
- No free parameters needed to fix the position and shape!
- Triangle singularity expected to be present
- Systematic studies also prefer the TS model
- Occam's razor: No need for a new genuine resonance
- ⇒ Have to include rescattering into future analyses!

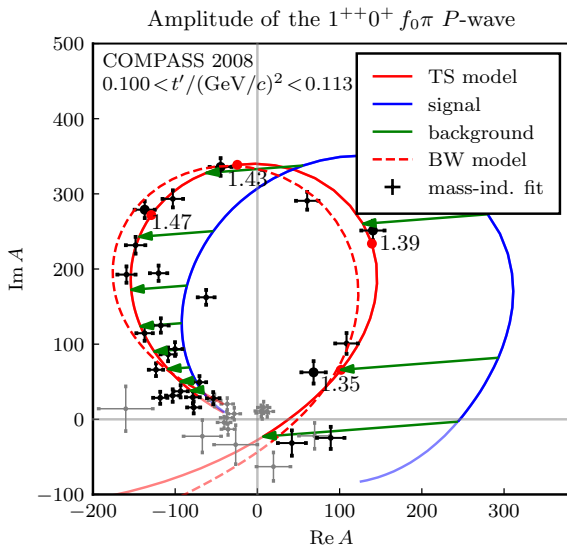
Outlook:

- Investigate $K\bar{K}\pi$ spectrum
- Look into $\tau \rightarrow 3\pi$ e.g. at BELLE (II), no Deck-like background

Thank you for your
attention!

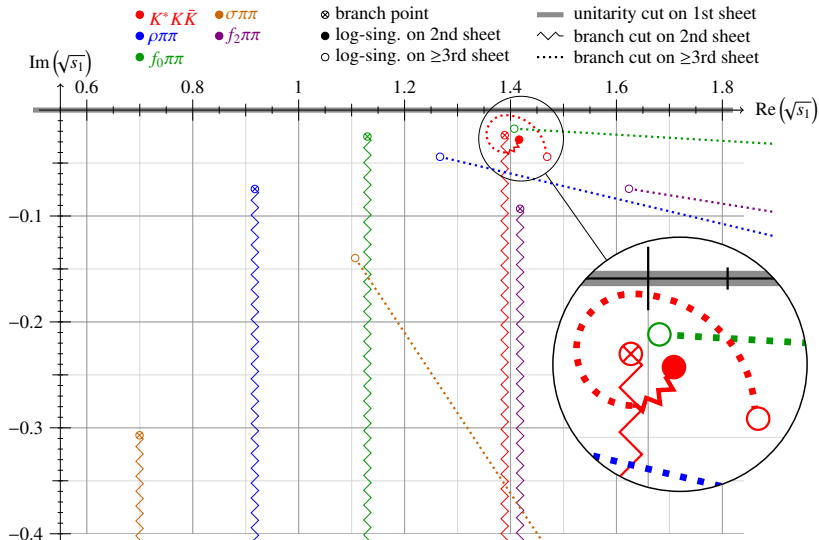
Back-up

Argand Diagram

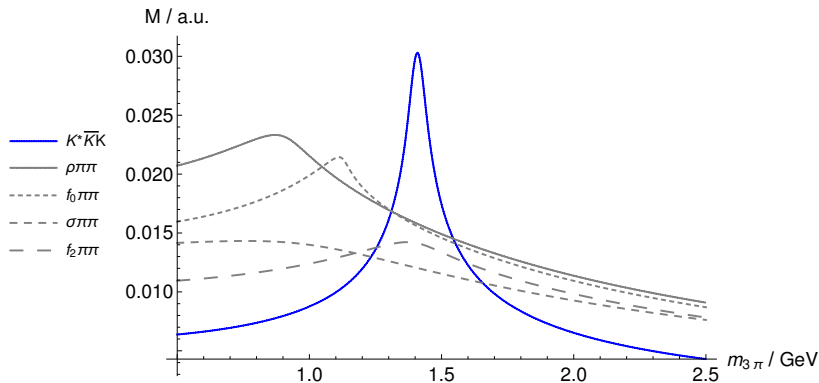


[COMPASS, PRL **127**, 082501 (2021)]

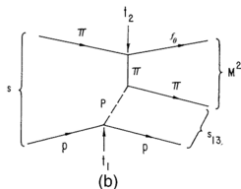
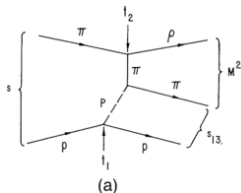
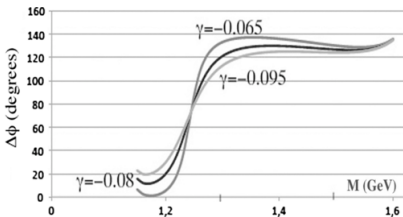
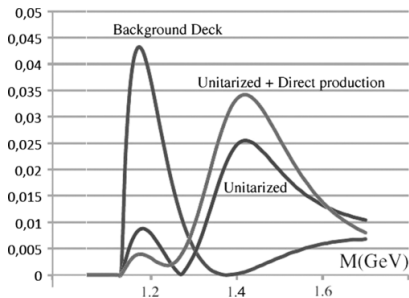
Pole Positions



Other Amplitudes



Basdevant-Berger Model

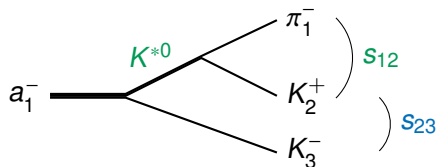


[Basdevant & Berger, PRL **114**, 192001 (2015)]

$$A(\tau) = \sum_{w=(JMLS)} \left[F_w(s_{12}) Z_w^*(\Omega_{3,12}) + F_w(s_{23}) Z_w^*(\Omega_{1,23}) \right]$$

Simple model: $F_w(s_{12}) = C_{a_1} \cdot t_{K^*}(s_{12})$

- $A(\tau)$: full amplitude of kinematic variables τ
- $F_w(s_{ij})$: isobar amplitude of decay with isobar in (ij) -channel
- $Z_w(\Omega_{k,ij})$: angular dependence of amplitude in (ij) -channel



$$A(\tau) = \sum_{w=(JMLS)} \left[F_w(s_{12}) Z_w^*(\Omega_{3,12}) + F_w(s_{23}) Z_w^*(\Omega_{1,23}) \right]$$

Projection to channel (23):

$$\begin{aligned} A_w(s_{23}) &= \int d\Omega_{1,23} Z_w(\Omega_{1,23}) A(\tau) \\ &= F_w(s_{23}) + \hat{F}_w(s_{23}) \end{aligned}$$

with $\hat{F}_w(s_{23}) := \int dZ_w(s_{23}) \sum_{w'} F_{w'}(s_{12}) Z_{w'}^*(\Omega_{3,12})$

unitarity for PW amplitude A_w :

$$\Rightarrow F_w(s_{23}) = t_\xi(s_{23}) \left[C_w + \frac{1}{2\pi} \int_{s_{\text{th}}}^{\infty} \frac{\rho(\tilde{s}_{23}) \hat{F}_w(\tilde{s}_{23})}{\tilde{s}_{23} - s_{23}} d\tilde{s}_{23} \right]$$

Problem: \hat{F} depends on F as well! \leadsto solve iteratively

$$F_w(\mathbf{s}, s_{23}) = t_{\xi}(s_{23}) \left[C_w(\mathbf{s}) + \frac{1}{2\pi} \int_{s_{\text{th}}}^{\infty} \frac{\rho(\tilde{s}_{23}) \hat{F}_w(\mathbf{s}, \tilde{s}_{23})}{\tilde{s}_{23} - s_{23}} d\tilde{s}_{23} \right]$$

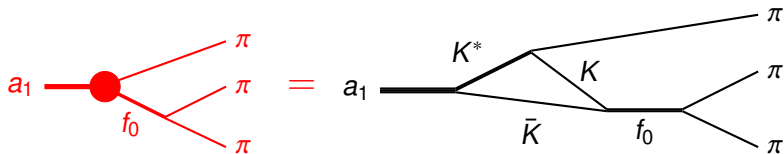
KT:

- calculate effects of rescattering on the 2-body subsystem invariant-mass dependence s_{ij}
- Iterative framework to include rescattering to any order

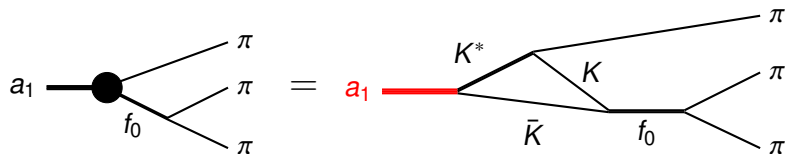
Our method:

- calculate effects of rescattering on the 3-body invariant-mass dependence $s = m_{3\pi}^2$
- Stop after the first iteration

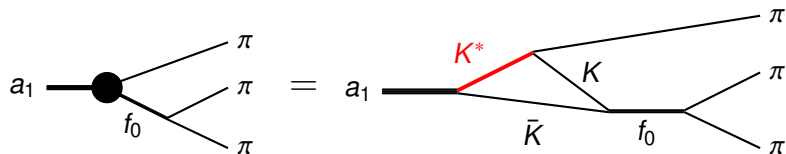
$$F(\mathbf{s}_{23}) = t_{f_0}(\mathbf{s}_{23}) \frac{1}{2\pi} \int_{4m_K^2}^{\infty} d\tilde{\mathbf{s}}_{23} \frac{\rho(\tilde{\mathbf{s}}_{23}) \int dZ_{f_0}(\tilde{\mathbf{s}}_{23}) C_{a_1} t_{K^*}(\mathbf{s}_{12}) Z_{K^*}^*(\Omega_{3,12})}{\tilde{\mathbf{s}}_{23} - \mathbf{s}_{23}}$$



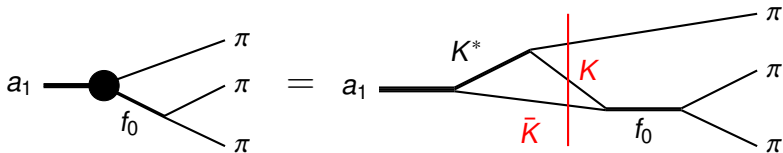
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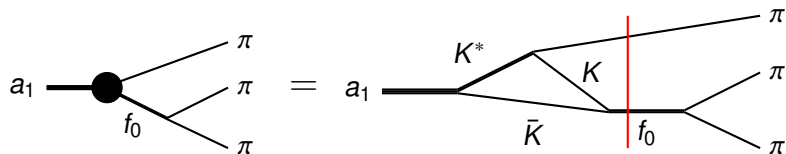
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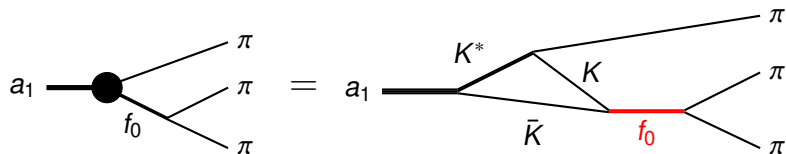
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