# Spin-orbit correlations in Monte Carlo simulations 

Focus on recent work on the polarized quark fragmentation process

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## nucleon structure at leading twist

The collinear nucleon structure at leading twist is described by unpolarised $f_{1}$, helicity $g_{1}$, transversity $h_{1}$

|  | nucleon |  |  |
| :---: | :---: | :---: | :---: |
| quark | U | L | T |
| U | $\begin{gathered} f_{1}(x \quad) \\ \text { (unpolarized) } \end{gathered}$ |  |  |
| L |  | $\begin{aligned} & g_{1}(x) \\ & \text { (helicity) } \end{aligned}$ |  |
| T |  |  | $\begin{gathered} h_{1}(x \quad) \\ \text { (transversity) } \end{gathered}$ |



Anselmino et al., PRD 92, 114023 (2015)

transversity is the less
known

## nucleon structure at leading twist

The collinear nucleon structure at leading twist is described by unpolarised $f_{1}$, helicity $g_{1}$, transversity $h_{1}$

|  |  | nucleon |  |  |
| :---: | :---: | :---: | :---: | :---: |
| quark | $\mathbf{U}$ | $\mathbf{L}$ | $\mathbf{T}$ |  |
| $\mathbf{y y y}$ | $\mathbf{U}$ | $f_{1}\left(x, k_{T}^{2}\right)$ <br> (unpolarized) |  |  |
| $\mathbf{L}$ |  | $f_{1 T}\left(x, k_{T}^{2}\right)$ <br> (Sivers) |  |  |
|  |  | $g_{1}\left(x, k_{T}^{2}\right)$ <br> (helicity) | $g_{1 T}\left(x, k_{T}^{2}\right)$ <br> (worm-gear) |  |
| $\mathbf{T}$ | $h_{1}^{\perp}\left(x, k_{T}^{2}\right)$ <br> (Boer-Mulders) | $h_{1 L}^{\perp}\left(x, k_{T}^{2}\right)$ <br> (worm-gear) | $h_{1}\left(x, k_{T}^{2}\right)$ <br> (transversity) <br> $h_{1 T}^{\perp}\left(x, k_{T}^{2}\right)$ <br> (pretzelosity) |  |

+5 other TMD PDFs when the intrinsic $k_{T}$ is taken into account $\rightarrow$ correlations between the nucleon spin and the spin and $k_{T}$ of partons
~ basically unknown, but Sivers

In SIDIS some TMDs are coupled to the unpolarized FF $D_{1}$

$$
\begin{array}{ll}
A_{U U}^{\cos \phi_{h}} & \sim f_{1} \otimes D_{1} \\
A_{U T}^{\sin \left(\phi_{h}-\phi_{S}\right)} & \sim f_{1 T} \otimes D_{1} \\
A_{L T}^{\cos \left(\phi_{h}-\phi_{S}\right)} & \sim g_{1 T} \otimes D_{1}
\end{array}
$$

others coupled to the Collins FF $H_{1}^{\perp}$

| $A_{U T}^{\sin \left(\phi_{h}+\phi_{S}-\pi\right)}$ | $\sim h_{1} \otimes H_{1}^{\perp}$ | Collins asymmetry |
| :--- | :--- | :--- |
| $A_{U U}^{\cos 2 \phi_{h}}$ | $\sim h_{1}^{\perp} \otimes H_{1}^{\perp}$ | $\sim$ Boer-Mulders |
| $A_{U L}^{\sin \left(2 \phi_{h}\right)}$ | $\sim h_{1 L}^{\perp} \otimes H_{1}^{\perp}$ | Kotzinian-Mulders |
| $A_{U T}^{\sin \left(3 \phi_{h}-\phi_{S}\right)}$ | $\sim h_{1 T}^{\perp} \otimes H_{1}^{\perp}$ | pretzelosity |

## MC implementation of spin-orbit effects

Monte Carlo event generators (MCEGs) are important tools
event generation in the full phase space, access to correlations, multi-dimensional studies, phenomenology

A systematic implementation of all leading order effects related to the intrinsic transverse momentum of quarks and to their spin degree of freedom in complete MCEGs is still missing

Different effects considered separately

- TMD PDFs
$\boldsymbol{k}_{T}$ dependence in $f_{1} \rightarrow$ «primordial kT» in PYTHIA? (but PDFs in PYTHIA do not depend on $\boldsymbol{k}_{T}$ )
Cahn and Sivers effects implemented in LEPTO Kotzinian '05


## - TMD FFs

unpolarized quarks, $\mathrm{D}_{1} \rightarrow$ Lund String Model, Cluster Model .. (is relation with theory OK?)
polarized quarks $\rightarrow(T)$ Collins effect $H_{1}^{\perp},(\mathrm{L})$ jet handedness
recursive string+3P0 model $\rightarrow$ this talk
stand alone MC
inclusion in PYTHIA 8.2
extended NJL-jet model stand alone MC

AK et al., PRD 100 (2019) 1, 014003, PRD 97 (2018) 7, 074010
AK and L. Lönnblad, PoS DIS2O19 (2019) 179

Matevosyan et al., PRD 95 (2017) 1, 014021

## the classical string $+{ }^{3} P_{0}$ model

The string+3P0 model with pseudoscalar (PS) meson production

$\pi^{+}$and $\pi^{-}$are emitted on opposite sides

- qualitative agreement with data

- predicts also a di-hadron asymmetry

COMPASS PLB 717 (2012) 376

A quantum mechanical formulation of the string+ ${ }^{3} P_{0}$ model is used for simulations

## elementary splitting: emission of a PS

string decay $=$ recursive repetition of the elementary splitting $q \rightarrow h+q^{\prime}$
X. Artru, Z. Belghobsi DSPIN-2011, 2013

$$
\begin{array}{cl}
\text { hadron }=\{\mathrm{h}, & \text { hadron type with transverse mass } \varepsilon_{h}=\left(m_{h}^{2}+\boldsymbol{p}_{T}^{2}\right)^{1 / 2} \\
Z=p^{+} / k^{+}, & \text {longitudinal momentum fraction } \\
\left.\boldsymbol{p}_{T}=\boldsymbol{k}_{T}-\boldsymbol{k}_{T}^{\prime}\right\} & \text { transverse momentum w.r.t string axis }
\end{array}
$$

quark $^{\prime}=\left\{q^{\prime}, k^{\prime}, \rho\left(\boldsymbol{S}_{q^{\prime}}\right)\right\}$


| quark $=\{q$, | flavour |
| ---: | :--- | ---: |
| $k$, | 4-momentum |
| $\left.\rho\left(\boldsymbol{S}_{q}\right)\right\}$ | spin density matrix | ( $2 \times 2$, Pauli spinors)

Splitting in flavour $\otimes$ momentum $\otimes$ spin space
$\rightarrow$ Transition Amplitude

$$
T_{q^{\prime}, h, q}\left(Z, \boldsymbol{p}_{T} \mid \boldsymbol{k}_{T}\right)
$$

Splitting Probability
$\rightarrow$ Splitting Function

$$
F_{q^{\prime} h q}\left(Z, \boldsymbol{p}_{T} \mid \boldsymbol{k}_{T}, \boldsymbol{S}_{q}\right)=\operatorname{tr} T_{q^{\prime}, h, q} \rho\left(\boldsymbol{S}_{q}\right) T_{q^{\prime}, h, q}^{\dagger}
$$

Spin transfer to $q^{\prime}$
$\rightarrow$ spin density matrix of $q^{\prime} \quad \rho\left(\boldsymbol{S}_{q^{\prime}}\right) \propto T_{q^{\prime}, h, q} \rho\left(\boldsymbol{S}_{q}\right) T_{q^{\prime}, h, q}^{\dagger}$
the very basic ingredients for the MC simulations

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$$

quark ${ }^{\prime}=\left\{q^{\prime}, k^{\prime}, \rho\left(\boldsymbol{S}_{q^{\prime}}\right)\right\}$


| quark $=\{q$, | flavour |
| ---: | :--- |
| $k$, | 4-momentum |

- Expression for the Transition Amplitude
$T_{q^{\prime}, h, q}=C_{q^{\prime}, h, q} \times\left(\frac{1-Z}{\varepsilon_{h}^{2}}\right)^{a / 2} e^{\left.-\frac{b_{L}}{2} \frac{\varepsilon_{h}^{2}}{Z} \times e^{-\frac{b_{T}}{2} \boldsymbol{k}_{T}^{\prime 2}} \times \check{g}\left(\varepsilon_{h}^{2}\right) \times\left[\mu+\sigma_{z} \sigma \cdot \boldsymbol{k}_{T}^{\prime}\right] \times \sigma_{z} \times \hat{u}^{-1 / 2}\left(\boldsymbol{k}_{T}\right),{ }^{2}\right)}$
$T_{q^{\prime}, h, q}=$ Lund String Fragmentation Model $\times{ }^{3} P_{0}$ operator $\times$ PS coupling $\times .$.
- Few free parameters
$a, b_{L}, \mathrm{~b}_{\mathrm{T}} \rightarrow$ string fragmentation dynamics (Lund Model, e.g. PYTHIA, LEPTO) $\mu$ complex mass from ${ }^{3} P_{0}$ mechanism $\rightarrow$ responible for spin effects ( $\operatorname{Im}(\mu) \rightarrow$ transverse)
- Input function $\check{g} \rightarrow$ governs spin-independent $\mathbf{k}_{\mathrm{T}}-\mathbf{k}_{\mathrm{T}}^{\prime}$ correlations correlations $\rightarrow$ Model M18 PRD 97 (2018) 7, 074010 NO correlations $\rightarrow$ Model M19 (much simpler) PRD100 (2019) no.1, 014003


## Stand alone simulations (M19)

M18 and M19 have been implemented in stand alone MC programs give very similar results, in spite for M19 being much simpler

Next slide, simulated Collins analysing power $a^{u \uparrow \rightarrow h+X}$ with initial conditions

## $u$ quarks fully transversely polarized along $\hat{y}$

Energy calculated from a $\left\{x_{B}, Q^{2}\right\}$ sample of SIDIS events no primordial KT

Values of the free parameters

$$
\begin{aligned}
& a=0.9 \\
& b_{L}=0.5\left(\mathrm{GeV} / c^{2}\right)^{-2} \\
& b_{T}=5.17(\mathrm{GeV} / c)^{-2} \\
& \mu=(0.42+\mathrm{i} 0.76) \mathrm{GeV} / \mathrm{c}^{2}
\end{aligned}
$$

see AK, X. Artru, Z. Belghobsi, F. Bradamante, A. Martin PRD 100 (2019) 1, 014003

## Stand alone simulations: comparison with SIDIS data (M19)



Di-hadron asymmetry for $h^{+} h^{-}$pairs MC $\rightarrow$ same scale factor $\lambda$
same mechanism as for Collins


Collins asymmetry on protons $A_{\text {Coll }}^{p} \simeq \frac{h_{1}^{u}}{f_{1}^{u}} a^{u \uparrow \rightarrow \pi+X}$
MC $\rightarrow$ Collins analysing power $a^{u \uparrow \rightarrow \pi+X}$ scaled by $\lambda \sim\left\langle h_{1}^{u} / f_{1}^{u}\right\rangle=0.055 \pm 0.010$ only u quarks fully transversely polarized $\lambda$ is estimated by comparison with $A_{\text {Coll }}^{p}$ for $\pi^{-}$

## Improving the model

The model gives already a good description of the main properties of data few parameters, same mechanism for Collins and dihadron asymmetries, jet handedness (not shown here)

We have improved it further following two directions ..
a) Exploit the true predictive power of the model via a more complete simulation of the event
$\rightarrow$ interface M19 with PYTHIA 8.2 for SIDIS
$\rightarrow$ introduction of transversity PDF
in collaboration with L. Lönnblad
the first step towards a systematic implementation of spin effects in PYTHIA!
b) Improve the description of the polarized fragmentation process $\rightarrow$ extend M19 by introducing vector mesons $\rightarrow$ the NEW model M20

## PYTHIA+3PO

M19 is interfaced with PYTHIA 8.2 for the simulation of polarized SIDIS
$\rightarrow$ spin effects introduced for the first time in the hadronization part of a complete MCEG
$\rightarrow$ parameterizations for $u^{v}$ and $d^{v}$ transversity PDFs implemented
$\rightarrow$ PYTHIA+3P0 allows to simulate the Collins and dihadron asymmetries
Simulation of SIDIS off protons @ COMPASS kinematics
ISR/FSR switched OFF, no intrinsic $\boldsymbol{k}_{\perp}$
complex mass retuned to $\mu=(0.78+i 0.38) \mathrm{GeV} / \mathrm{c}^{2}$

Collins asymmetry PYTHIA+3P0


## Nice description of data!

 trend vs $Z_{h}$ is modified in PYTHIA also good description of di-hadron asymmetriesCollins asymmetry M19


PYTHIA+3PO will be available soon!!
$A K$ and L. Lönnblad, in preparation

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the first step towards a systematic implementation of spin effects in PYTH $/$ A!
b) Improve the description of the polarized fragmentation process
$\rightarrow$ extend M19 by introducing vector mesons $\rightarrow$ the NEW model M20
spin effects with PS and VM production are treated systematically, for the first time

## elementary splitting: emission of a VM

AK, PhD thesis, 2004.00524 paper in preparation
quark' $\left\{q^{\prime}, k^{\prime}, \rho\left(\boldsymbol{S}_{q^{\prime}}\right)\right\}$
$\rho, K^{*}, \omega, \phi$
inv. mass spin matrix

quark $\left\{q, k, \rho\left(\boldsymbol{S}_{q}\right)\right\}$

## NEW Transition Amplitude

Splitting Probability
$\rightarrow$ NEW Splitting Function
Spin transfer to $h$
$\rightarrow$ spin density matrix of $h$
Decay (e.g. VM $\rightarrow h_{1} h_{2}$ )
$\rightarrow$ angular distribution
Bring decay information back to $q^{\prime}$
$\rightarrow$ decay matrix
$\widetilde{D}_{s_{h}^{\prime} s_{h}}=\mathcal{M}_{s_{h}}^{\dagger}, \mathcal{M}_{s_{h}}$
Transfer spin information to $q^{\prime}$
$\rightarrow$ spin density matrix of $q^{\prime}$
$T_{q^{\prime}, h, q}^{s_{h}}\left(M, Z, \boldsymbol{p}_{T}, s_{h} \mid \boldsymbol{k}_{T}\right)$
$d N / d \Omega \propto \mathcal{M}_{s_{h}} \hat{\rho}_{s_{h} s_{h}}, \mathcal{M}_{s_{h^{\prime}}}^{\dagger}$

$F_{q^{\prime} h q}\left(M, Z, \boldsymbol{p}_{T} \mid \boldsymbol{k}_{T}, \boldsymbol{S}_{q}\right)=\operatorname{tr} T_{q^{\prime}, h, q}^{S_{h}} \rho\left(\boldsymbol{S}_{q}\right) T_{q^{\prime}, h, q}^{S_{h} \dagger}$
$\hat{\rho}_{s_{h} s_{h^{\prime}}}(h) \propto \operatorname{tr} T_{q^{\prime}, h, q}^{s_{h}} \rho\left(\boldsymbol{S}_{q}\right) \mathrm{T}_{q^{\prime}, h, q}^{s_{h^{\prime}} \dagger}$
$\rho\left(\boldsymbol{S}_{q^{\prime}}\right)=\widetilde{D}_{s_{h} s_{h}^{\prime}} T_{q^{\prime}, h, q}^{s_{h}^{\prime}} \rho\left(\boldsymbol{S}_{q}\right) \mathrm{T}_{q^{\prime}, h, q}^{s_{h} \dagger}$

## elementary splitting: emission of a VM

AK, PhD thesis, 2004.00524 paper in preparation
quark' $\left\{q^{\prime}, k^{\prime}, \rho\left(\boldsymbol{S}_{q^{\prime}}\right)\right\}$


quark $\left\{q, k, \rho\left(\boldsymbol{S}_{q}\right)\right\}$

## Complicated recipe!

respects entanglement $q^{\prime} \leftrightarrow$ momenta of decay hadrons (Collins '88, Knowles '88)
Form of the NEW splitting amplitude
$T_{q^{\prime}, h, q}^{S_{h}}=$ relativistic BW $\times$ Lund String Fragmentation Model $\times{ }^{3} P_{0}$ operator $\times$ VM coupling $\times$..
VM coupling $\rightarrow$ complex free parameters
$G_{L} \rightarrow$ coupling of q to VM with linear L polarisation along the string axis
$G_{T} \rightarrow$ coupling of q to VM with linear T polarisation w.r.t the string axis
Only two new parameters in the model

$$
\begin{array}{ll}
\left|G_{L}\right| /\left|G_{T}\right| & \rightarrow \text { global Collins effect of the VM (depends on VM polarisation!) } \\
\theta_{L T}=\arg \left(\mathrm{G}_{\mathrm{L}} / \mathrm{G}_{\mathrm{T}}\right) & \rightarrow \text { oblique polarisation (LT) }
\end{array}
$$

## Stand alone MC implementation of M20

For each event define initial quark $q_{A} \equiv q_{1}$, i.e. flavour ( $u, d, s$ ), momentum, density matrix $\rho\left(q_{A}\right)$

1. Generate a $q_{2} \bar{q}_{2}$ pair and form the hadron $h_{1}\left(q_{A} \bar{q}_{2}\right)$, VM with prob. $\frac{f_{V M}}{f_{V M}+f_{P S}}$
2. Generate $M_{h_{1}}$ (if VM), $\boldsymbol{k}_{2 T}, Z_{1}$, using $F_{q_{2} h_{1} q_{A}} \rightarrow \operatorname{construct} p_{1}$
3. If $h_{1}=\mathrm{PS}$ go to 4 .

If $h_{1}=\mathrm{VM} \rightarrow$ calculate $\hat{\rho}(h)$
a) generate decay hadrons in VM rest frame and boost to string frame
b) construct the decay matrix $\breve{D}$
4. Calculate the spin density matrix of $q_{2}$

- Iterate points 1-4 until the exit condition (enough renamining mass to produce at least one baryonic resonance)



## Stand alone simulations with M20

Values of the free parameters used in simulations
all mesons

$$
\begin{aligned}
& a=0.9 \\
& b_{L}=0.5\left(\mathrm{GeV} / c^{2}\right)^{-2} \\
& b_{T}=8.43(\mathrm{GeV} / c)^{-2} \\
& \mu=(0.42+\mathrm{i} 0.76) \mathrm{GeV} / \mathrm{c}^{2}
\end{aligned}
$$

as in model M19

VM production

$$
f_{V M} / f_{P S}=0.62(0.725) \text { for } \mathrm{u}, \mathrm{~d}(\mathrm{~s}) \quad \text { as in PYTHIA } 8
$$

$$
\begin{aligned}
& \left|G_{L}\right| /\left|G_{T}\right|=1 \\
& \theta_{L T}=0
\end{aligned} \longrightarrow \begin{aligned}
& \text { following the NR quark model } \begin{array}{c}
\text { Czyzewski '96 } \\
\text { sensitivity to parameters values also explored }
\end{array}
\end{aligned}
$$

Initial conditions
$u$ quarks fully transversely polarized along $\hat{y}$
Energy calculated from a $\left\{x_{B}, Q^{2}\right\}$ sample of SIDIS events
no primordial KT

## Collins analysing power as function of rank


classical picture




- classical picture reproduced $\rho^{+}$have opposite effect w.r.t $\pi^{+}$
- quark spin information decays along the chain faster decay in M20


## Collins analysing power for $\rho$ and decay $\pi$


$\rho^{+}$analysing power
opposite to $\pi^{+}$
$\sim 3$ times smaller than $\pi^{+}$
as expected from the M20 prediction


AK et al., PRD 100 (2019) 1, 014003,
model M19 $\rightarrow$ only PS

$$
\left.\frac{a^{u \uparrow \rightarrow \rho+X}}{a^{u \uparrow \rightarrow \pi+X}}\right|_{r a n k=1}=-\frac{\left|G_{L}\right|^{2}}{2\left|G_{T}\right|^{2}+\left|G_{L}\right|^{2}}
$$

$\rho^{-}, \rho^{0} \sim$ to $\rho^{+}$for $\left|G_{L}\right| /\left|G_{T}\right|=1$
decay $\pi^{+} \rightarrow$ larger analysing power at large $z_{h}$ and large $p_{T}$ w.r.t $\rho^{+}$
large $z_{h} \rightarrow \pi^{+}$emitted along $\hat{z}$ from longitudinally polarized $\rho^{+}$
large $p_{T} \rightarrow \pi^{+}$emitted along $\vec{p}_{T}^{\rho}$ from rank $1 \rho^{+}$ emitted along $\hat{z} \times \vec{p}_{T}^{\rho}$ from rank $2 \rho^{+}$
small $p_{T} \rightarrow \pi^{+}$emitted along $-\vec{p}_{T}^{\rho}$ from rank $1 \rho^{+}$

## Effect of VM decays on transverse spin asymmetries

Collins analysing power


Large effect on Collins analysing power w.r.t M19

- different trends
- average analysing power diluted by 50\%

Di-hadron analysing power

di-hadron analysing power calculated using the relative transverse momentum

$$
\begin{gathered}
\boldsymbol{R}_{T}=z_{2} \boldsymbol{p}_{1 T} / z-z_{1} \boldsymbol{p}_{2 T} / z \\
a^{u \uparrow \rightarrow h^{+} h^{-}+X}=2\left\langle\sin \left(\phi_{R}-\phi_{S_{A}}\right)\right\rangle
\end{gathered}
$$

50\% dilution w.r.t M19

- effect at $\rho^{0}$ peak due to
$\rho^{0} \rightarrow \pi^{+} \pi^{-}$symmetric w.r.t $\boldsymbol{R}_{T} \leftrightarrow-\boldsymbol{R}_{T}$


## Comparison with SIDIS data



MC scaled by a factor $\lambda$ depending on $\left|G_{L}\right| /\left|G_{T}\right|$ and $\theta_{L T}$

Large variations for $\pi^{+}$due to different values of $\left|G_{L}\right| /\left|G_{T}\right|$ and $\theta_{L T}$ somewhat smaller for $\pi^{-}$
$\rightarrow$ both parameters are important
hint for $\frac{\left|G_{L}\right|}{\left|G_{T}\right|}>1, \theta_{L T}<0$ ?
more precise data would help to fix the free parameters

## Sensitivity to free parameters: Collins effect for $\rho$ mesons




COMPASS results: NEW! Measurements feasible
they could be used to fix the parameters

Hint for $\frac{\left|G_{L}\right|}{\left|G_{T}\right|}>1$ (in particular from $p_{T}$ )

Strong dependence on $\left|G_{L}\right| /\left|G_{T}\right|$ both size and shapes change

## Conclusions

The string+3P0 model with PS meson emission (M18, M19) implemented in a stand alone MC
$\rightarrow$ describes the main features of Collins and di-hadron asymmetries!

M19 has been interfaced to PYTHIA 8
$\rightarrow$ parameterisations for the transversity PDF implemented
$\rightarrow$ more complete description of TSA
the code will available very soon (AK and L. Lönnblad, in preparation)
For the first time implementation of the string+3PO model with PS and VM production (M20) in a stand alone MC
(paper in preparation)
$\rightarrow$ detailed study of Collins effect for VM
$\rightarrow$ only 4 free parameters for spin effects, to be fixed from comparison with data $\mu$ OK, hints for $\left|G_{L}\right| /\left|G_{L}\right|>1$ and $\theta_{L T}<0$
more precise data would help ( COMPASS 2021-2022 d run, JLab12 .. )
ongoing work, promising results ...

