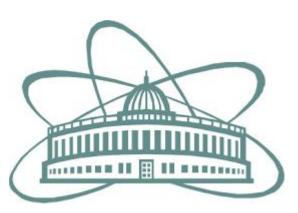
# Measuring the Primakoff reaction and $F_{K2\pi}$ with AMBER phase-II

## Andrei Maltsev, JINR, Dubna



Perceiving the EHM, AMBER@CERN, April 27 - April 30, 2021

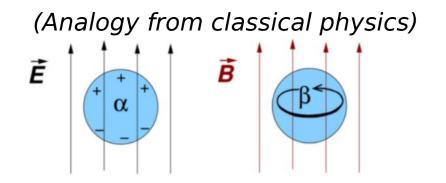


## **Motivation**

- QCD: extremely successful theory of strong interactions, but not yet possible to derive from the first principles, the fundamental properties of the bound states (masses, spectra)
- Effective QCD-based models are able to give quantitative predictions for processes at low energies (*chiral perturbation theory*, *quark confinement model*, *etc.*) → need to test applicability regions
- Simplest QCD objects: pions & kaons → obtain experimental results on their structure parameters → control the applicability region of these effective models

## **Polarizabilities**

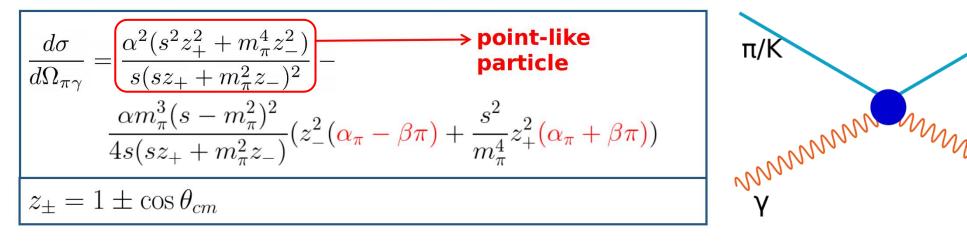
Interaction between **hadron** and **external electromagnetic field** described by parameters  $\alpha$ ,  $\beta$  (LO), encoding information about its internal structure



π/K

 $H_{em} = ... - \frac{1}{2}(\alpha \mathbf{E}^2 + \beta \mathbf{B}^2) + ...$ ,  $\alpha$ : electric polarizability,  $\beta$ : magnetic polarizability

## Additional term into compton scattering cross-section:



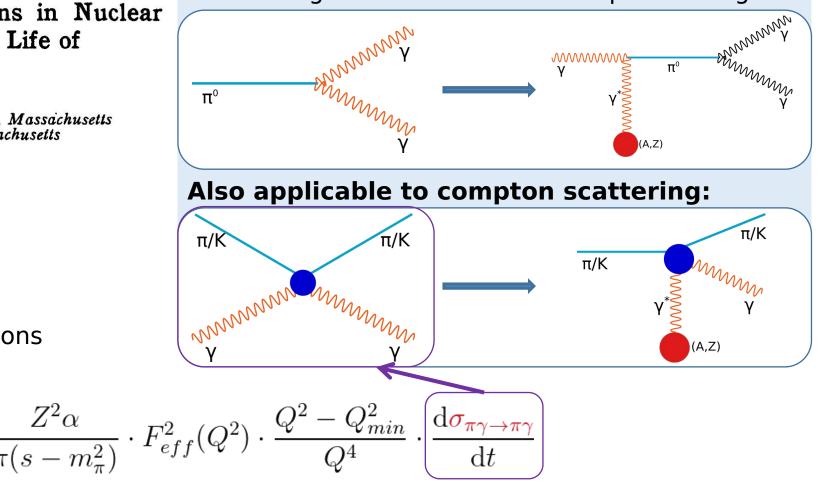
## How to access polarizabilities in experiment?

#### Photo-Production of Neutral Mesons in Nuclear Electric Fields and the Mean Life of the Neutral Meson<sup>\*</sup>

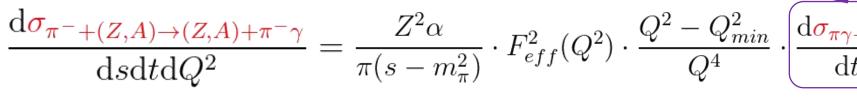
H. PRIMAKOFF† Laboratory for Nuclear Science and Engineering, Massachusetts Institute of Technology, Cambridge, Massachusetts January 2, 1951

## **Initial idea of Henry Primakoff:**

Electromagnetic field of nucleus = photon target!

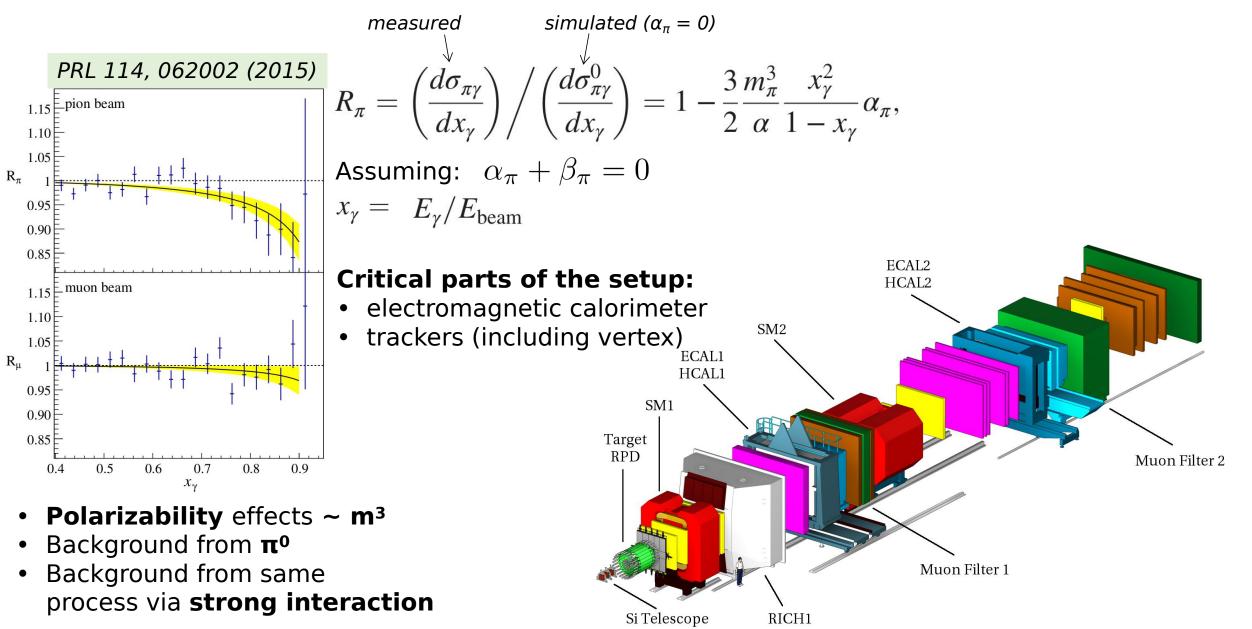


Scattering off equivalent photons Assumption:  $Q^2 \ll m_{\pi/K}^2$ 

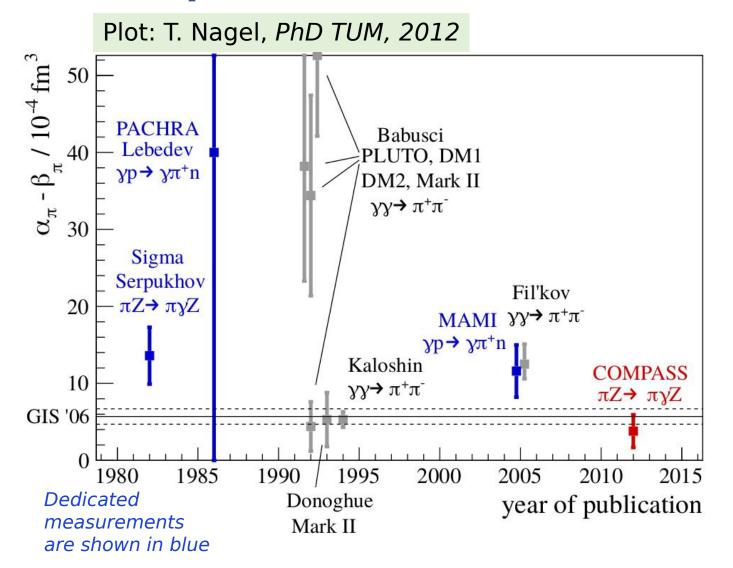


 $Q_{\min} = (s - m_{\pi}^2)/2E_{\text{beam}}$  Q is the virtual photon 4-momentum, F(Q<sup>2</sup>) - nucleus form factor

# Polarizability measurement @ COMPASS



# Pion polarizabilities: world data



assuming  $\alpha_{\pi} + \beta_{\pi} = 0$ : **COMPASS:**  $\alpha_{\pi} = (2.0 \pm 0.6_{stat} \pm 0.7_{syst}) \times 10^{-4} \text{ fm}^3$ 

## **Theory predictions:**

**χPT (two-loop, pions):**   $α_{π}-β_{π} = (5.7 \pm 1.0) \times 10^{-4} \text{ fm}^{3}$   $α_{π}+β_{π} = 0.16 \times 10^{-4} \text{ fm}^{3}$ 

Most other low-energy models (chiral quark model, dispersion relations):  $8 \times 10^{-4} \text{ fm}^3 < \alpha_{\pi} - \beta_{\pi} < 12 \times 10^{-4} \text{ fm}^3$ 

 $\chi$ PT (chiral perturation theory): low-energy expansion of QCD in particle momenta and quark masses 6

# Kaon polarizabilities: world data

 $|\alpha_{\kappa}| < 200 \times 10^{-4} \, \text{fm}^3$  (90% confidence) (from kaonic atoms spectrum)

G. Backenstoss et. al, *Phys.Lett.43B*, 5 (1973)

**Theory predictions:** 

 $\chi$ PT (one-loop):  $\alpha_{K}$ - $\beta_{K}$  = 1.16×10<sup>-4</sup> fm<sup>3</sup>

Quark confinement model:  $\alpha_{K}$ - $\beta_{K} = 3.6 \times 10^{-4} \text{ fm}^{3}$  $\alpha_{K}$ + $\beta_{K} = 2.3 \times 10^{-4} \text{ fm}^{3}$ 

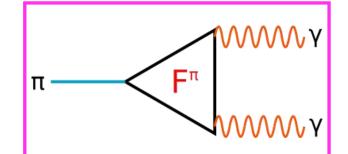
**Pion data:** nice precision, COMPASS measurement in nice agreement with  $\chi$ PT **Kaon data:** only vague estimates

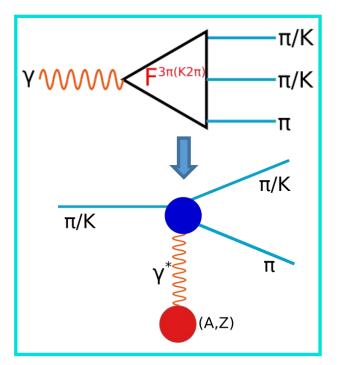
χPT (chiral perturation theory): low-energy expansion of QCD in particle momenta and quark masses Will it hold true also for strange quark sector?

### There is a large demand for high-precision kaon data on polarizabilities

# Chiral anomaly

Describes  $\pi^0 \rightarrow \gamma \gamma$  ( $F^\pi$ ) decay width and  $\gamma \pi \rightarrow \pi \pi$  ( $F^{3\pi}$ ),  $\gamma K \rightarrow \pi K$  ( $F^{K2\pi}$ ),  $\gamma \pi \rightarrow \pi \eta$ ,  $\gamma K \rightarrow K \eta$  vertices  $\frac{F^{3\pi}(0,0,0)}{F^{\pi}(0,0)} = \frac{1}{ef_{\pi}^{2}}$ f\_{\pi}: pion decay constant





Method of equivalent photons:

$$\frac{d\sigma_{\pi^- N \to N \pi^- \pi^0}}{ds dt dq^2} = \frac{Z^2 \alpha}{\pi} \cdot \frac{q^2 - q_{min}^2}{q^4} \cdot \frac{1}{s - m_\pi^2} \cdot |F(q^2)|^2 \cdot \frac{d\sigma_{\gamma \pi \to \pi \pi}}{dt}$$
$$\frac{d\sigma_{\gamma \pi \to \pi \pi}}{dt} = \frac{F_{3\pi}(s, t, u)^2}{128\pi} \cdot \frac{s - 4m_\pi^2}{4} \sin^2 \theta$$

Problems with accessing  $F^{3\pi}$ ,  $F^{\kappa 2\pi}$ :

- need to bridge the gap between  $s=t=q^2=0$  and physical region
- $\rightarrow \chi PT$ , dispersive framework
- large background from  $\gamma \pi \rightarrow \pi \pi \pi$  into  $\gamma \pi \rightarrow \pi \pi$  (Pomeron exchange) at high energies, and from  $\rho/K^*(892)$  production

# **Chiral anomaly**

## Theoretical prediction from chiral anomaly:

$$F_{K2\pi} = F_{3\pi} = \frac{e}{4\pi^2 F_\pi^3} = (9.78 \pm 0.05) \text{ GeV}^{-3}$$

Additional interest in  $\gamma\pi \rightarrow \pi\pi$  and  $\gamma K \rightarrow K\pi$  vertex studies:

- access to ρ/K\*(892) radiative widths
- help link lattice QCD calculations to physical parameters

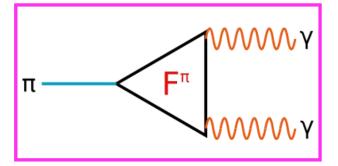
M. Niehus, M. Hoferichter, B. Kubis, PoS CD2018, 076 (2019)

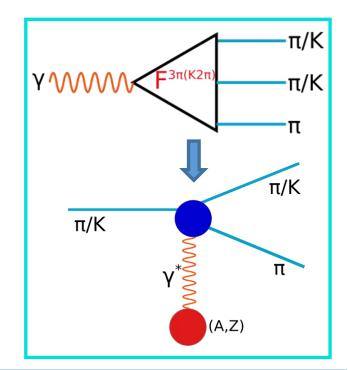
## **Experimental results:**

SIGMA (Serpukhov, 1980-s): π-Z→Zπ-π<sup>0</sup> F<sub>3π</sub> = (10.7 ± 1.2) GeV<sup>-3</sup>

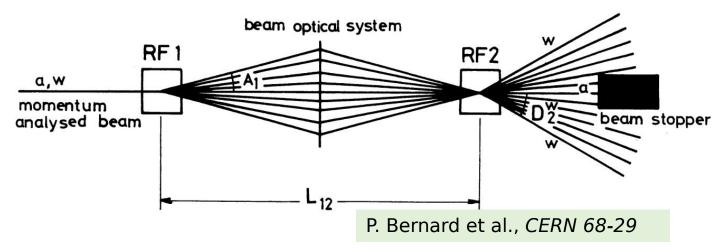
Y. M. Antipov et al., *Phys.Rev.D36, 21(1987)* L. Ametller et al., *Phys.Rev.D64, 094009(2001)*   $\frac{CERN SPS: \pi - e^{-} \rightarrow \pi - e^{-} \pi^{0}}{F_{3\pi}}$  = (9.6 ± 1.1) GeV<sup>-3</sup>

S. R. Amendolia et al., *Phys.Lett.B155, 457(1985)*  **F<sup>κ2π</sup>: no measurement yet** (planned at IHEP, Protvino)





## Polarisabilities with RF separated kaon-enriched beam at AMBER

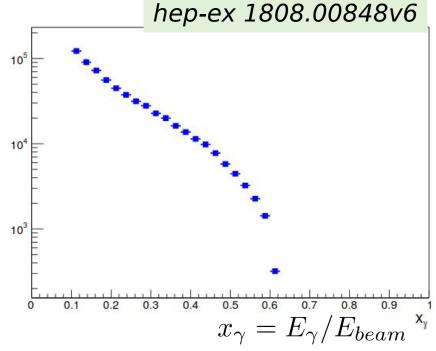


**Kaon enriched beam:** momentum  $p_K \sim 80$  GeV, intensity ~ 5×10<sup>6</sup> s<sup>-1</sup> (now: **kaons** @ COMPASS: ~ 10<sup>5</sup> s<sup>-1</sup>)

### **Critical parts of the setup:**

- electromagnetic calorimeter
- trackers (including vertex)
- performance similar to COMPASS

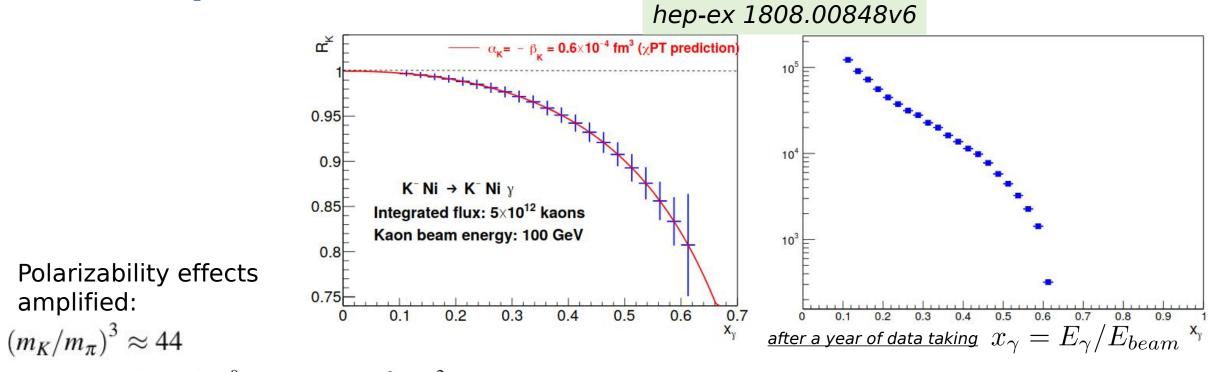
# AMBER setup will provide all the needed elements for high-precision kaon polarizability measurements



Estimated number of  $K^-Z \rightarrow ZK^-\gamma$ ("polarizability") events after **1 year** of data taking

(Assuming trigger rate of **100 kHz**)

# Kaon polarizabilities at AMBER



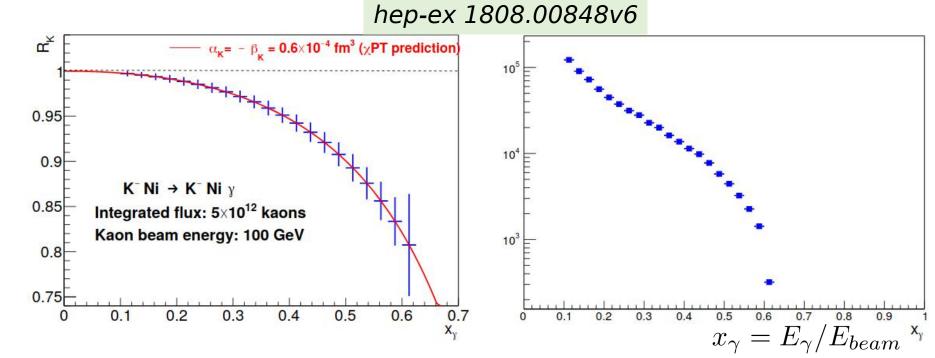
$$R_{\pi} = \left(\frac{d\sigma_{\pi\gamma}}{dx_{\gamma}}\right) \left/ \left(\frac{d\sigma_{\pi\gamma}^{0}}{dx_{\gamma}}\right) = 1 - \frac{3}{2} \frac{m_{\pi}^{3}}{\alpha} \frac{x_{\gamma}^{2}}{1 - x_{\gamma}} \alpha_{\pi},$$

- Expected statistical accuracy on  $\alpha_{\kappa} \beta_{\kappa}$ :  $\sigma_{stat} = 0.03 \times 10^{-4} \text{ fm}^3 (\alpha_{\kappa} + \beta_{\kappa} = 0)$ :
- No competitors so far

## **Theory predictions:**

 $\chi$ PT (one-loop):  $\alpha_{K}$ - $\beta_{K} = 1.16 \times 10^{-4}$  fm<sup>3</sup> QCM:  $\alpha_{K}$ - $\beta_{K} = 3.6 \times 10^{-4}$  fm<sup>3</sup>

# Kaon polarizabilities at AMBER



## More possible measurements (also for pions):

- separate measurements of  $\alpha_K$  and  $\beta_K$
- quadrupole polarizabilities

# AMBER setup will allow to measure charged kaon polarizabilities with and unprecedented precision

# Chiral anomaly in $\gamma K \rightarrow \pi K$

Idea: access  $\gamma K \rightarrow K\pi$  via  $K^{\pm} + (Z,A) \rightarrow (Z,A) + K^{\pm} + \pi^{0}$ similarly to:  $\mathbf{\gamma \pi} \rightarrow \mathbf{\pi \pi}$  via  $\pi^{\pm} + (Z,A) \rightarrow (Z,A) + \pi^{\pm} + \pi^{0}$ 

Anomalous contributions into  $\gamma \pi \rightarrow \pi \pi$  and  $\gamma K \rightarrow K \pi$  are equal:

$$F_{K2\pi} = F_{3\pi} = \frac{e}{4\pi^2 F_{\pi}^3} = (9.78 \pm 0.05) \text{ GeV}^{-3}$$

## BUT:

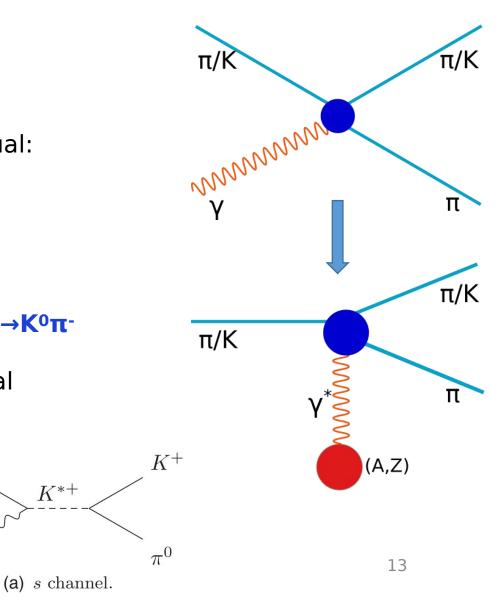
with kaons, two processes are possible:  $\mathbf{K}^{-}\mathbf{\gamma} \rightarrow \mathbf{K}^{-}\pi^{0}$  and  $\mathbf{K}^{-}\mathbf{\gamma} \rightarrow \mathbf{K}^{0}\pi^{-}$ 

 $K^+$ 

 $K^{*+}$ 

- **only**  $K^{-}\gamma \rightarrow K^{-}\pi^{0}$  is influenced by the chiral anomaly
- anomaly contributions to  $\mathbf{K}^{-}\gamma \rightarrow \mathbf{K}^{-}\pi^{0}$  and  $\pi^{-}\gamma \rightarrow \pi^{-}\pi^{0}$  are equal ۲
- physical region is further away from s=t=u=0 for kaons ٠

### How to incorporate K<sup>\*</sup> production into analysis?



# Chiral anomaly in $\gamma K \rightarrow \pi K$

M. I. Vysotsky and E. V. Zhemchugov Phys. Rev. D93, 094029(2016)

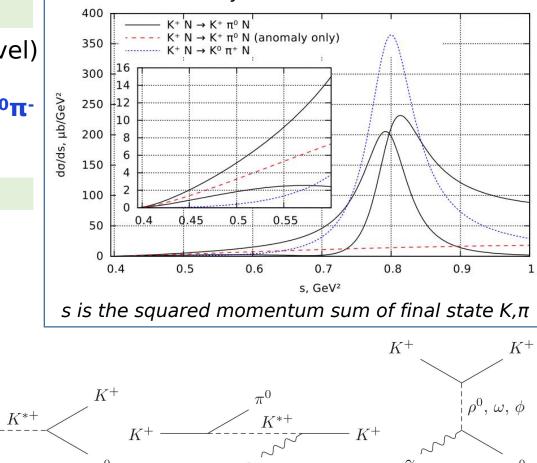
- Anomaly contribution + vector meson exchange (tree level)
- Anomaly dominates at low energies
- Extract anomaly from difference in K-γ→K-π<sup>0</sup> and K-γ→K<sup>0</sup>πcross sections in the low-energy region

M. Dax, D. Stamen, B. Kubis Eur. Phys. J. C (2021)81:221

- Dispersive framework for all charge channels
- Fix subtraction constants using data on
  K-γ→K-π<sup>0</sup> and K-γ→K<sup>0</sup>π<sup>-</sup> → extract chiral anomaly
- Allows to utilize data up to s ~ (1.2 GeV)<sup>2</sup>
- Allows to obtain radiative coupling of K\*(892)

#### Phys.Rev.D93, 094029(2016)

*Two solid lines: different interference phase between anomaly and resonance terms* 



(b) *u* channel.

## AMBER experiment could provide an opportunity to conduct precision measurements of $\gamma K \rightarrow \pi K$ cross sections

 $K^+$ 

(a) s channel.

14

(c) t channel.

# Summary

- Pion and kaon **polarizabilities**, as well as coupling constants of  $\gamma \pi \rightarrow \pi \pi$  and  $\gamma K \rightarrow K \pi$  processes are of interest as a way to test the predictions of **low-energy phenomenological models** with the goal of controlling their regions of applicability.
- There are very few data on processes involving kaons at the moment, mainly due to absense of high-intensity kaon beams.
- At **AMBER experiment**, the new RF separated kaon-enriched beam will allow to measure **kaon polarizabilities with an unprecedented precision**, as well as study the  $\gamma K \rightarrow \pi K$  vertex to extract precise information on the chiral anomaly and K\*(892) radiative width.