

A model for F_L at low Q^2 - revisited

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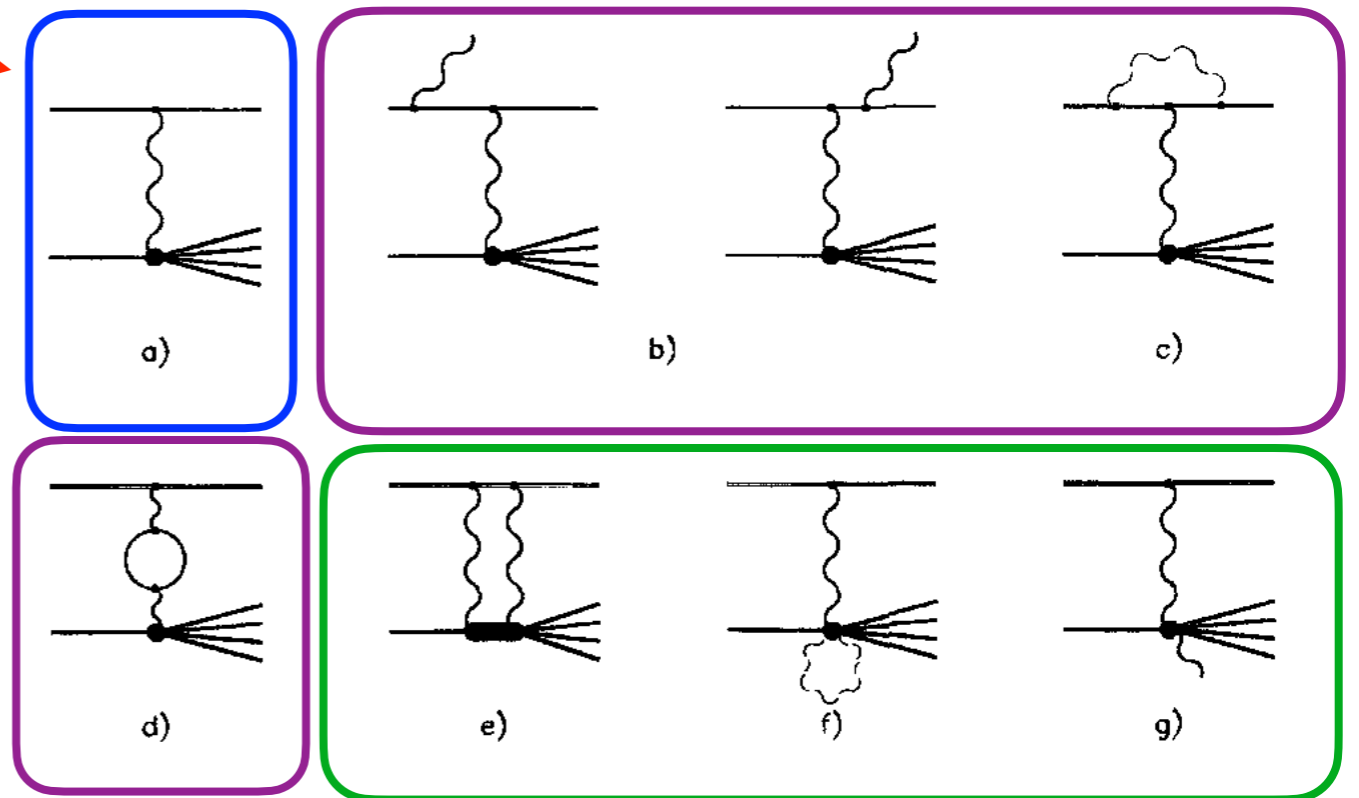
DIS2021, April 14, 2021

Outline

- Motivation for model of structure functions in the low Q^2 region: DIS radiative corrections (essential for EIC)
- Model for F_L at low Q^2 : k_T factorization + higher twist
- Results: x and Q^2 behavior
- Comparison with JLAB, SLAC and HERA data
- Outlook

Radiative corrections

Extraction of the one - photon cross section from the measurement requires the correction for the **radiative corrections**



RC factor:

$$\eta(x, y) = \frac{\sigma_{1\gamma}}{\sigma_{\text{meas}}}$$

Examples:

Mo-Tsai scheme **b)-d)**

Dubna scheme **b)-d)**

and **e)-g)** hadron current corrections calculated using Quark Parton Model

Comparison of schemes

Badelek, Bardin, Kurek, Scholz Z. Phys. C 66 , 591 (1995)

Radiative corrections

Need following input for

$$x_{\text{meas}} < x < 1$$

$$0 < Q^2 < Q^2_{\text{max}}$$

Structure function

$$F_2(x, Q^2)$$

And structure function

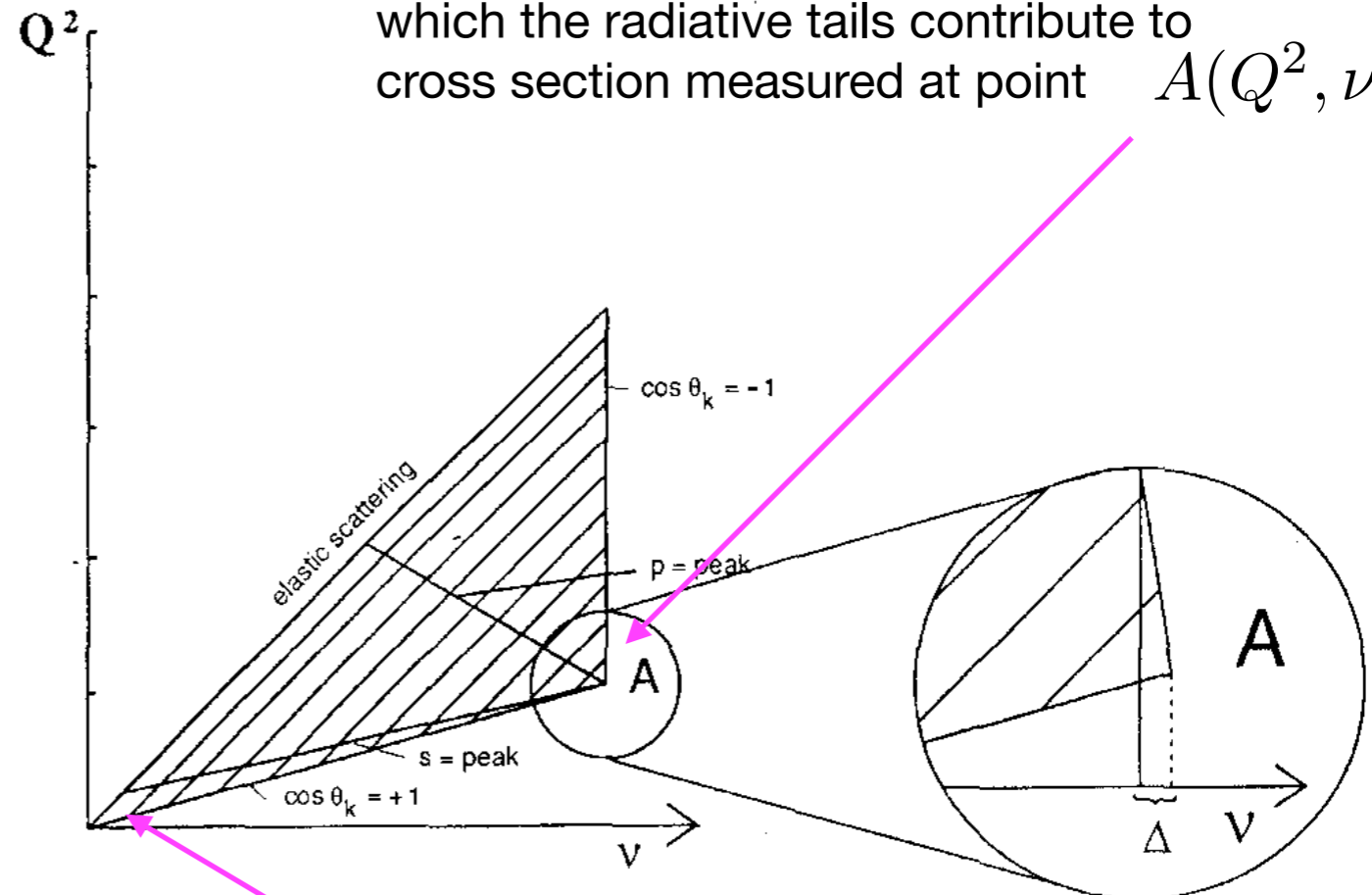
$$F_L(x, Q^2)$$

Or the ratio $R(x, Q^2) = F_L(x, Q^2) / F_T(x, Q^2)$

Plot from

Badelek, Bardin, Kurek, Scholz Z. Phys. C 66, 591 (1995)

Range of kinematical variables from which the radiative tails contribute to cross section measured at point $A(Q^2, \nu)$



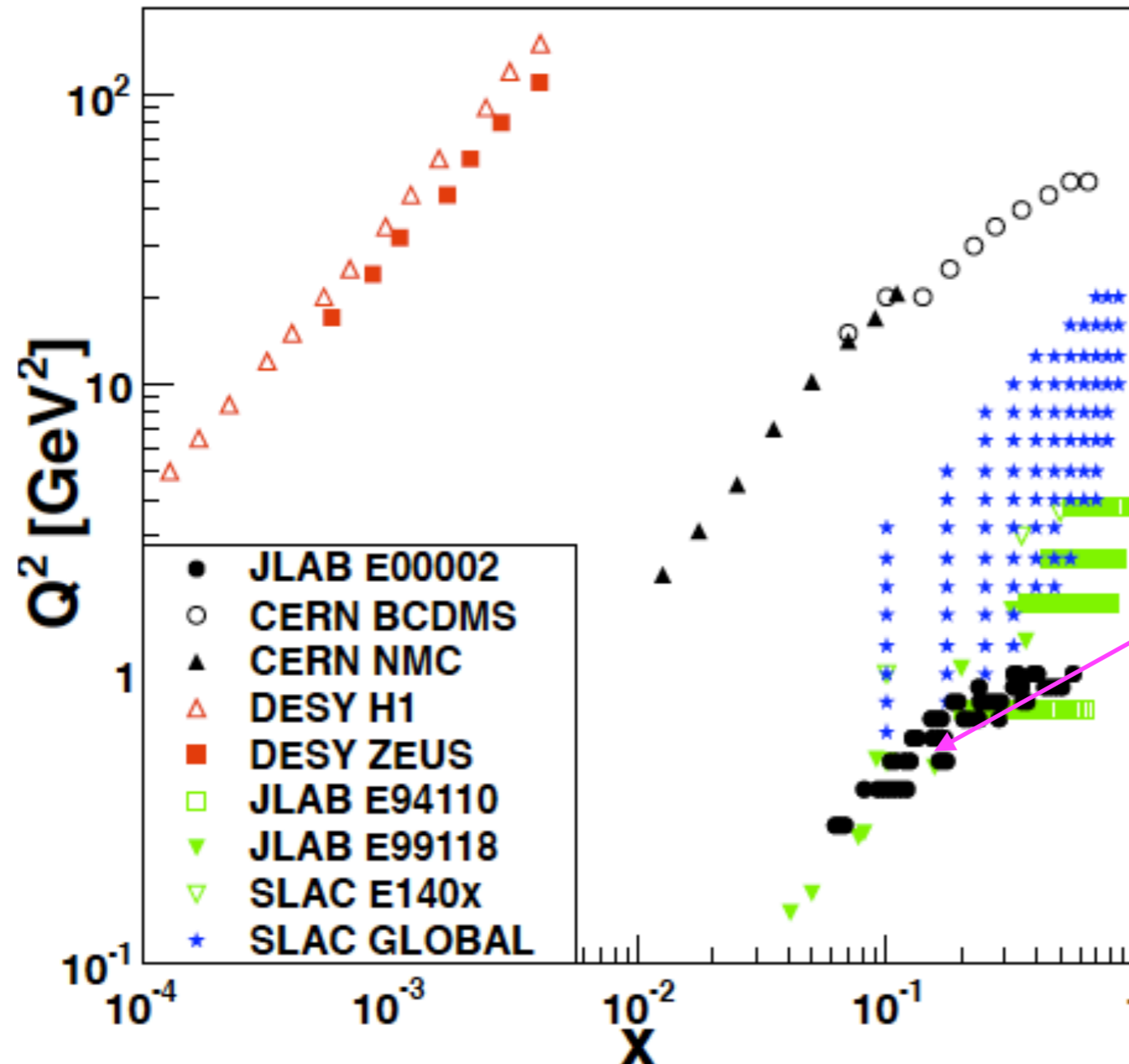
Needed down to

$$Q^2 \rightarrow 0$$

Essential for DIS

Phase space (x, Q^2) for F_L

JLab E00002 Phys. Rev. C 97, 045204 (2018)



Unlike for F_2 there are not many data for the longitudinal structure function

Lowest Q^2 data

Model for low Q^2

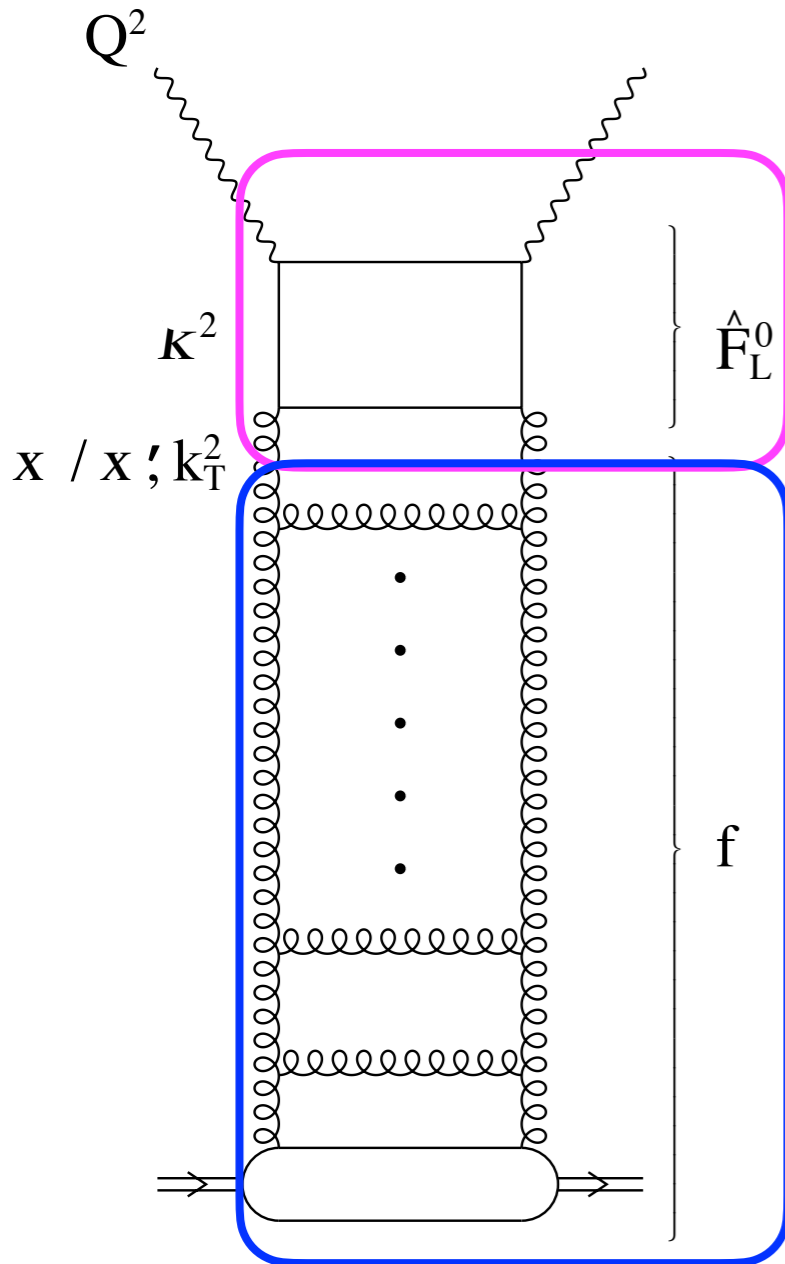
- Construct physically motivated model rather than pure parametrization
- Use of k_T factorization formula with off shell gluon
- Need unintegrated gluon distribution. Construct from the standard DGLAP collinear gluon distribution
- Extrapolation to low values of Q^2
- Introduce the cutoff parameter on the quark transverse momenta
- Higher twist contribution from the low quark transverse momenta. Match the normalization to the ‘soft’ contribution from the F_2 structure function

Revisiting original model:

Badelek, Kwiecinski, Stasto Z. Phys. C 74 , 297 (1997)

k_T factorization

- k_T factorization, appropriate for small x is convolution of:
 - photon-gluon impact factor — off-shell matrix element
 - unintegrated gluon density



$$F_L(x, Q^2) = \int_x^1 \frac{dx'}{x'} \int \frac{dk_T^2}{k_T^2} \hat{F}_L^0(x', Q^2, k_T^2) f\left(\frac{x}{x'}, k_T^2\right)$$

$$\hat{F}_L^0(x', Q^2, k_T^2) = \frac{Q^4}{\pi^2 k_T^2} \sum_q e_q^2 \int_0^1 d\beta \int d^2 \kappa'_T x' \delta \left(x' - \left(1 + \frac{\kappa_T'^2 + m_q^2}{\beta(1-\beta)Q^2} + \frac{k_T^2}{Q^2} \right)^{-1} \right) \times$$

$$\times \alpha_s \beta^2 (1-\beta)^2 \left(\frac{1}{D_{1q}} - \frac{1}{D_{2q}} \right)^2$$

Exact kinematics

Denominators

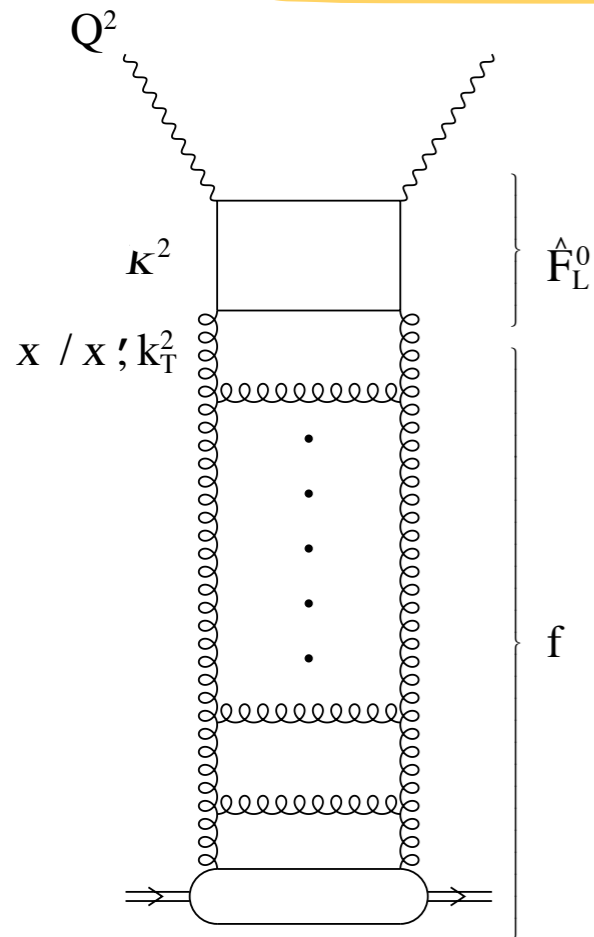
$$D_{1q} = \kappa_T^2 + \beta(1-\beta)Q^2 + m_q^2$$

$$D_{2q} = (\kappa_T - \mathbf{k}_T)^2 + \beta(1-\beta)Q^2 + m_q^2$$

Quark transverse momentum

$$\kappa_T' = \kappa_T - (1-\beta)\mathbf{k}_T$$

Model for higher twist



Large quark transverse momenta: $\kappa^2 > \kappa_0^2$ ← Cutoff

- Use standard kT factorization formula

Small quark transverse momenta: $\kappa^2 < \kappa_0^2$

- Make on-shell approximation $k_T^2 \rightarrow 0$
- Make ansatz $\alpha_s x g(x, Q^2) \rightarrow A$

$$F_L^{HT} = 2A \sum_q e_q^2 \frac{Q^4}{\pi} \int_0^1 d\beta \beta^2 (1-\beta)^2 \int_0^{\kappa_{0T}^{\prime 2}} d\kappa_T^{\prime 2} \frac{\kappa_T^{\prime 2}}{D_q^4}$$

$$D_q = \kappa_T^{\prime 2} + \beta(1-\beta)Q^2 + m_q^2.$$

Higher twist contribution: $\sim 1/Q^2$ when $Q^2 \rightarrow \infty$
 $\sim Q^4$ when $Q^2 \rightarrow 0$

Model for HT cont'd

Constant A is not free parameter.

Estimate it from the F_2 assuming the non-perturbative contribution also comes from the low quark transverse momenta.

$$F_T(x, Q^2) = 2 \sum_q e_q^2 \frac{Q^2}{4\pi} \alpha_s \int_0^1 d\beta \int d\kappa_T'^2 \frac{x}{x'} g\left(\frac{x}{x'}, Q^2\right) \times$$

$$\times \left[\frac{\beta^2 + (1 - \beta)^2}{2} \left(\frac{1}{D_q^2} - \frac{2\kappa_T'^2}{D_q^3} + \frac{2\kappa_T^2 \kappa_T'^2}{D_q^4} \right) + \frac{m_q^2 \kappa_T'^2}{D_q^4} \right]$$

Integrate over low quark transverse momentum $\kappa^2 < \kappa_0^2$

Again assume $\alpha_s x g(x, Q^2) \rightarrow A$

$$F_2^{Bg} = A \times \frac{\sum_q e_q^2}{\pi} \int_0^\infty dt \int_0^{\kappa_{0T}^2} d\kappa_T'^2 \left[\frac{1}{2} \left(\frac{1}{D_q^2} - \frac{2\kappa_T'^2}{D_q^3} + \frac{2\kappa_T^2 \kappa_T'^2}{D_q^4} \right) + \frac{m_q^2 \kappa_T'^2}{D_q^4} \right]$$

Set the background (soft) contribution (from separate analysis)

$$F_2^{\text{bg}} \rightarrow 0.4$$

Input for model

Different models for unintegrated gluon density:

- Small x linear BFKL
- Small x with nonlinear corrections, Balitsky-Kovchegov
- DGLAP motivated: Kimber-Martin-Ryskin formalism

Unintegrated gluon from integrated gluon PDF:

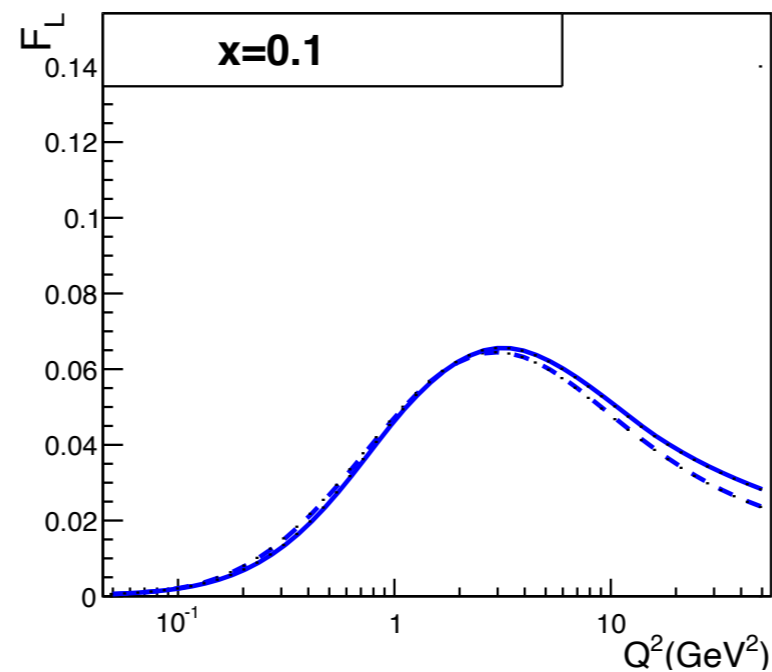
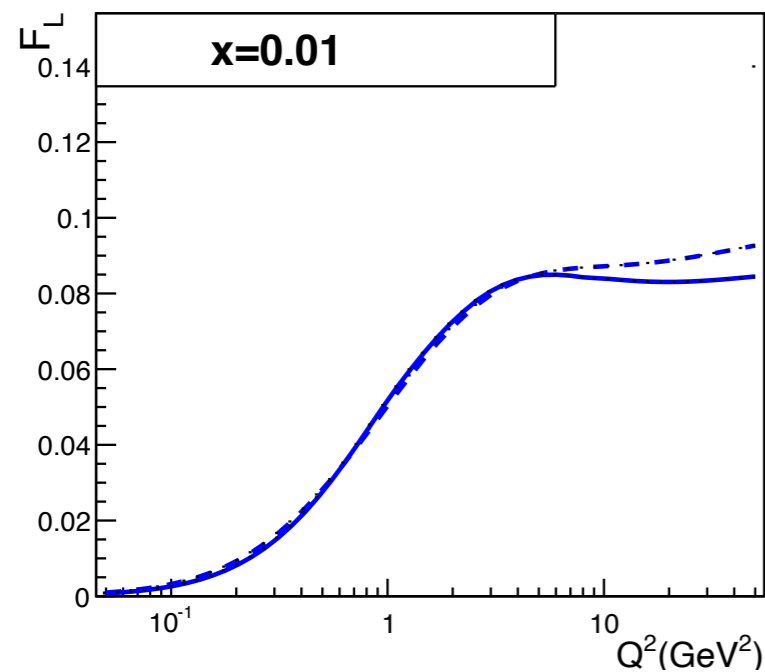
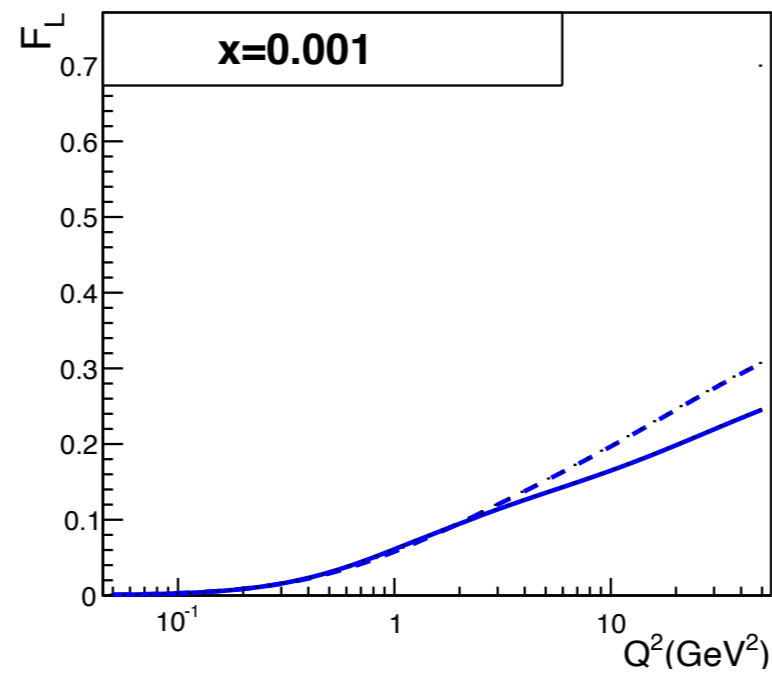
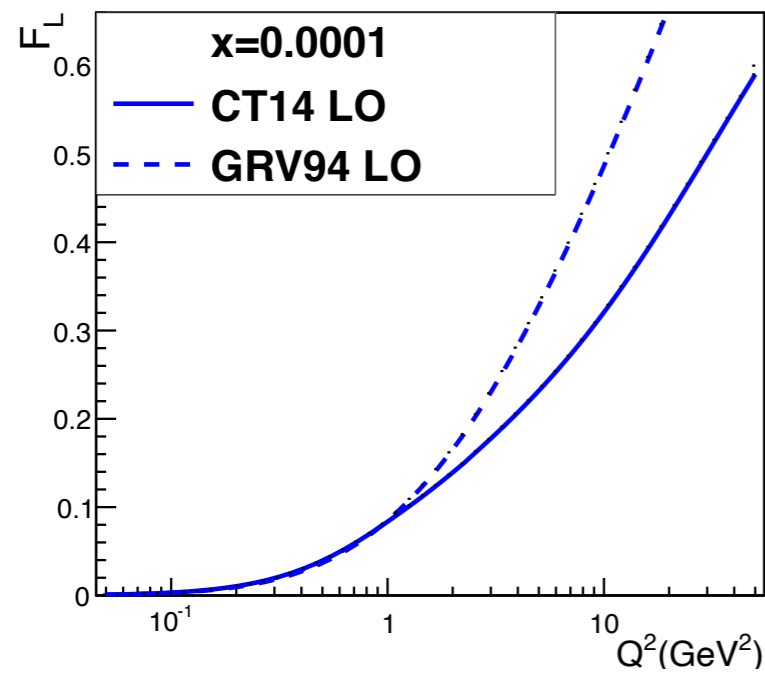
$$f(y, k_T^2) = y \frac{\partial g^{\text{AP}}(y, Q^2)}{\partial \ln Q^2} \Big|_{Q^2 = k_T^2}$$

Use standard PDFs, LO, NLO, not much sensitivity in the low Q^2

Cutoff variation $\kappa_0^2 = 0.8 - 1.5 \text{ GeV}^2$ Results not very sensitive

Masses for quarks (u,d,s,c): 0.35, 0.35, 0.5, 1.2-1.5 GeV

$F_L(Q^2)$ bins in x



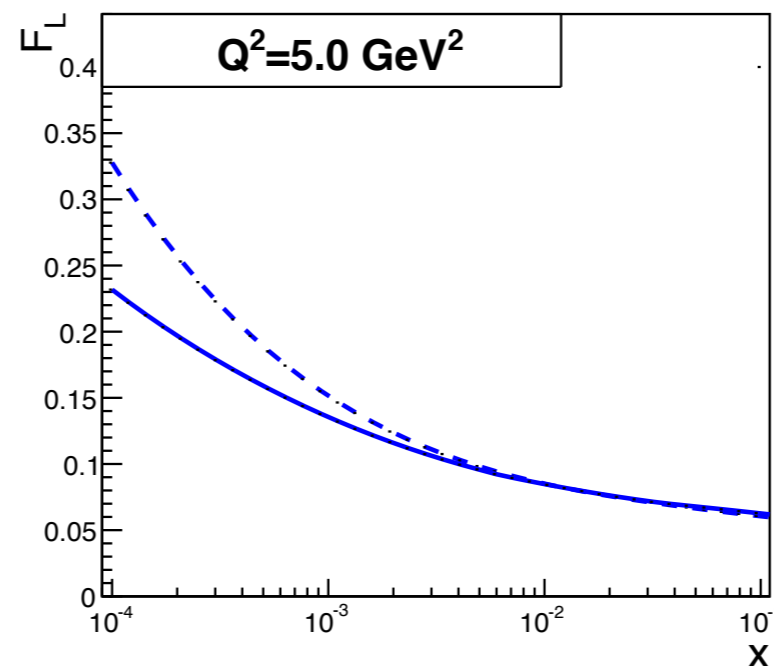
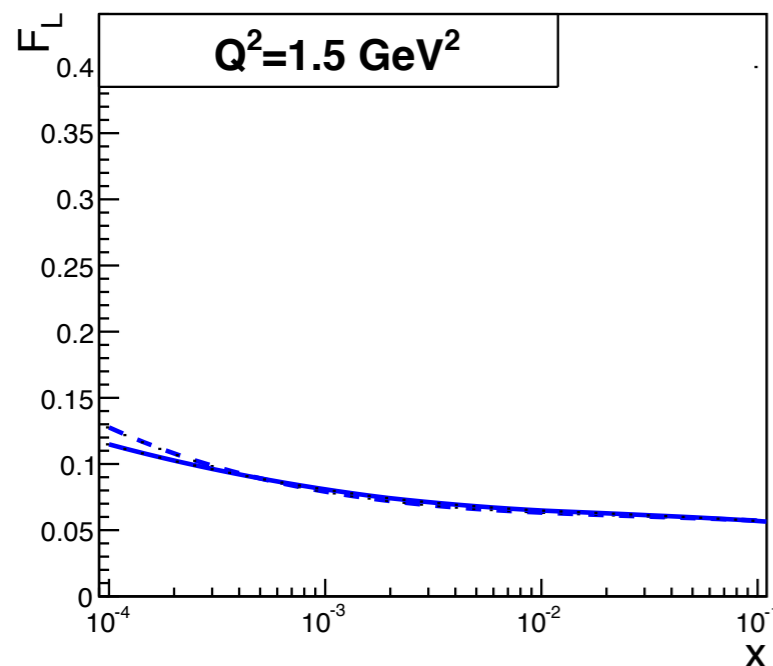
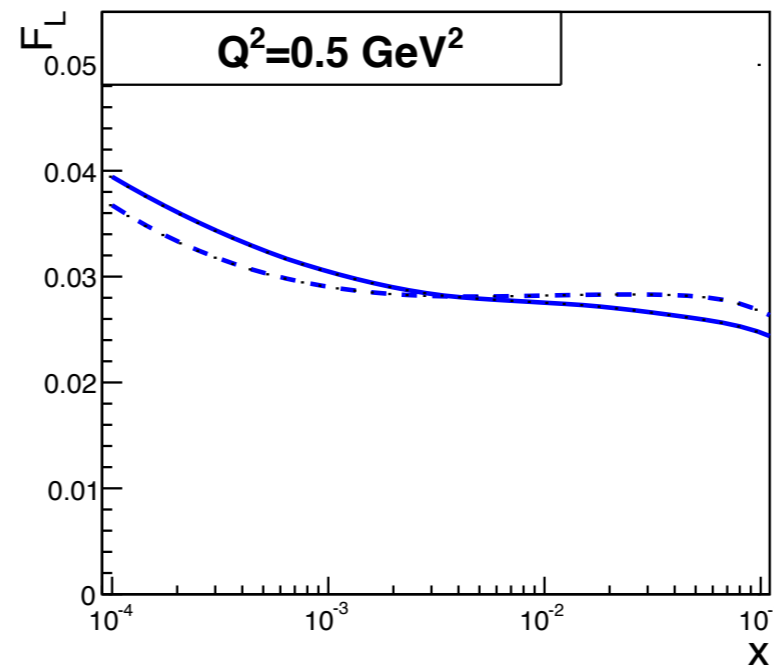
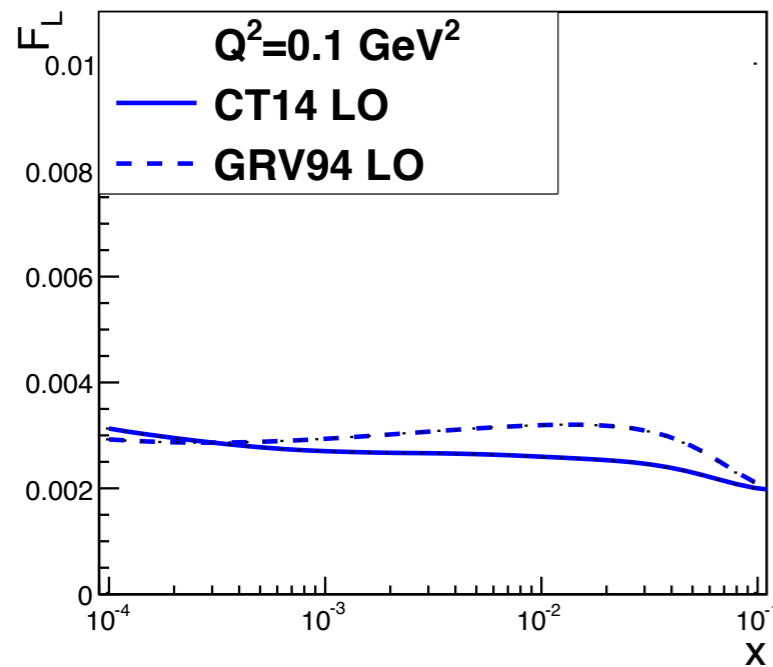
Comparison of old (GRV94) vs updated (CT14) calculation

Largest difference in moderate to large Q^2 and low x

No difference in low Q^2 region

Observe different vertical scale on different panels

$F_L(x)$ bins in Q^2



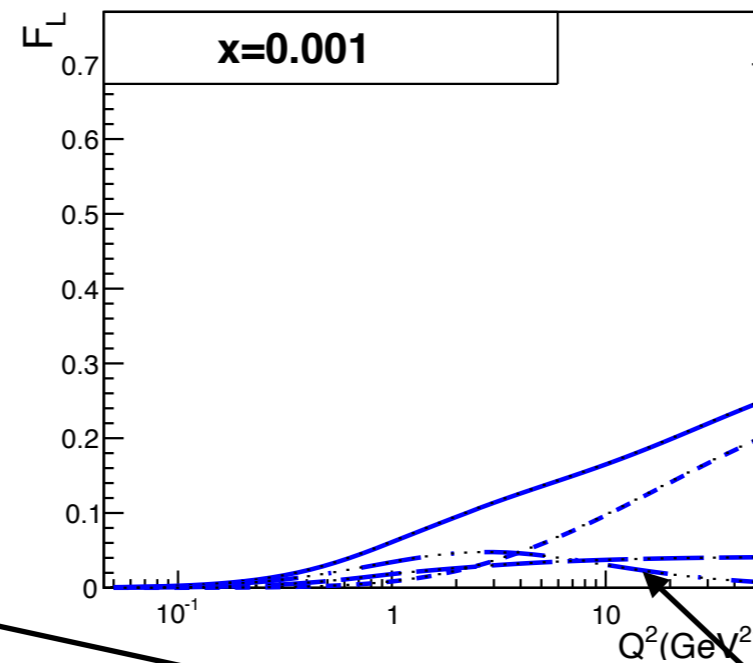
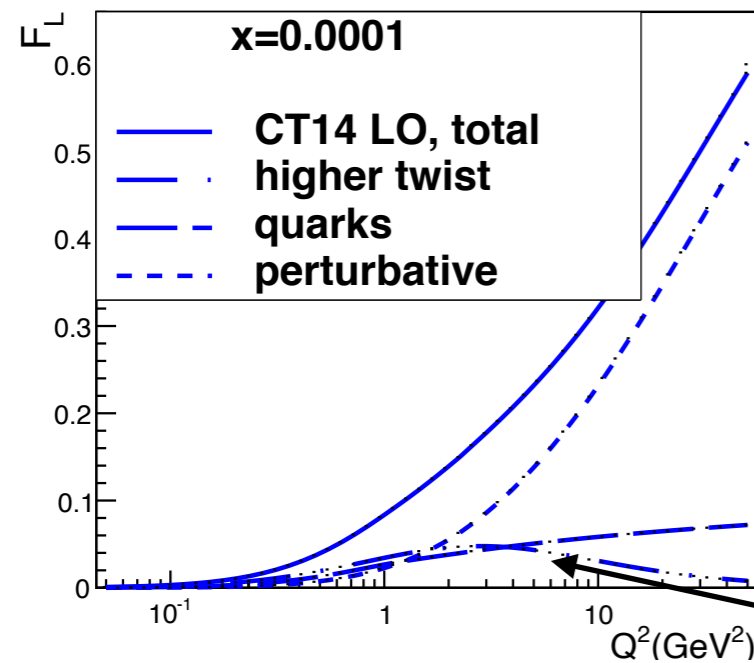
Comparison of old (GRV94) vs updated (CT14) calculation

Largest difference in moderate to large Q^2 and low x . Update gluon has slower x dependence

No difference in low Q^2 region

Observe different vertical scale on different panels

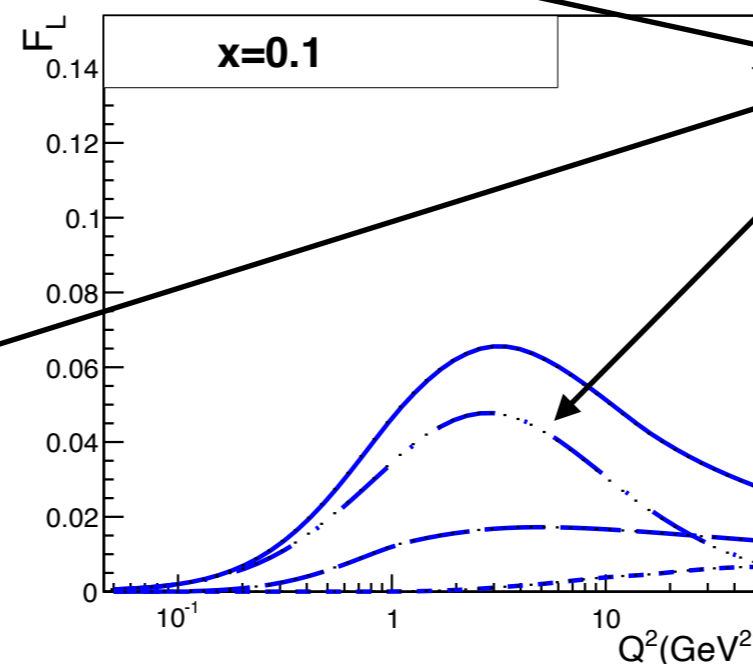
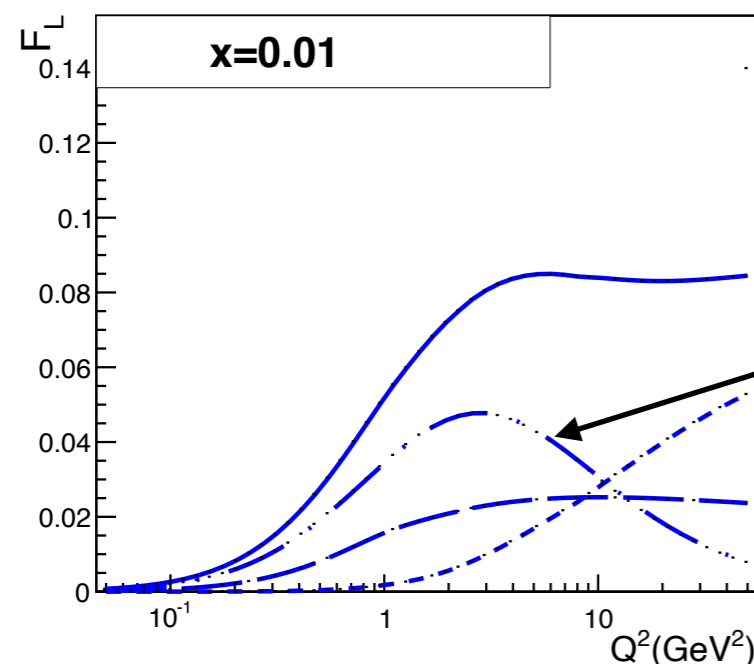
$F_L(Q^2)$ bins in x



Components of the model

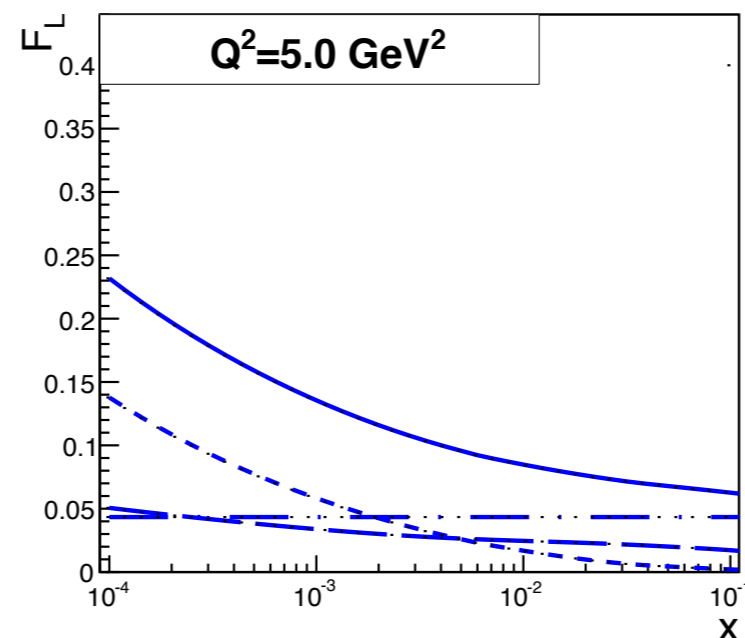
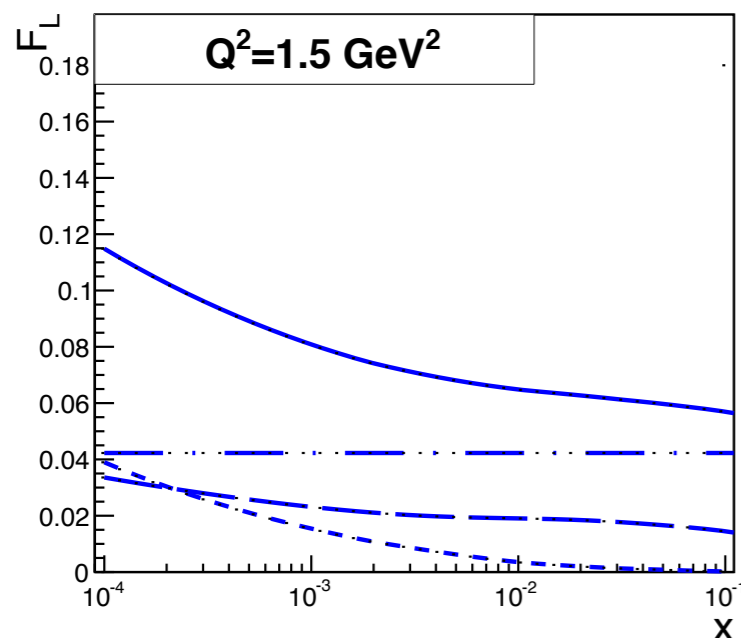
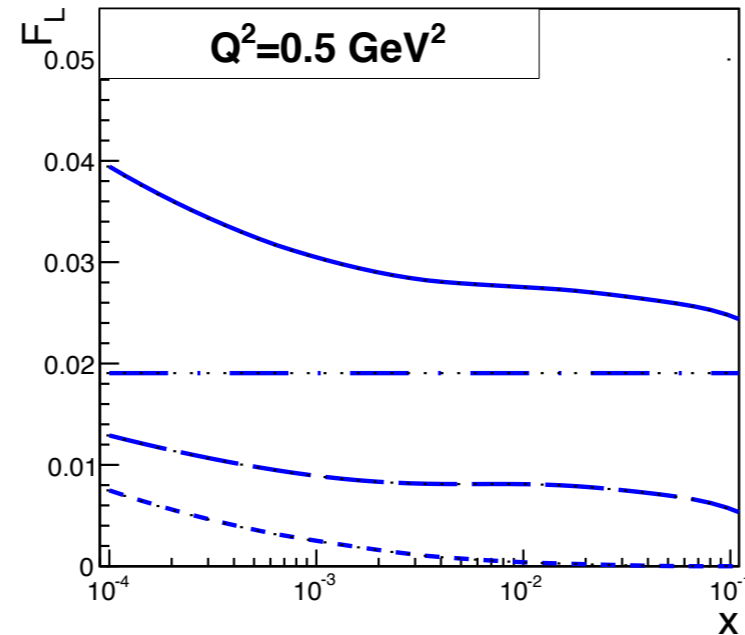
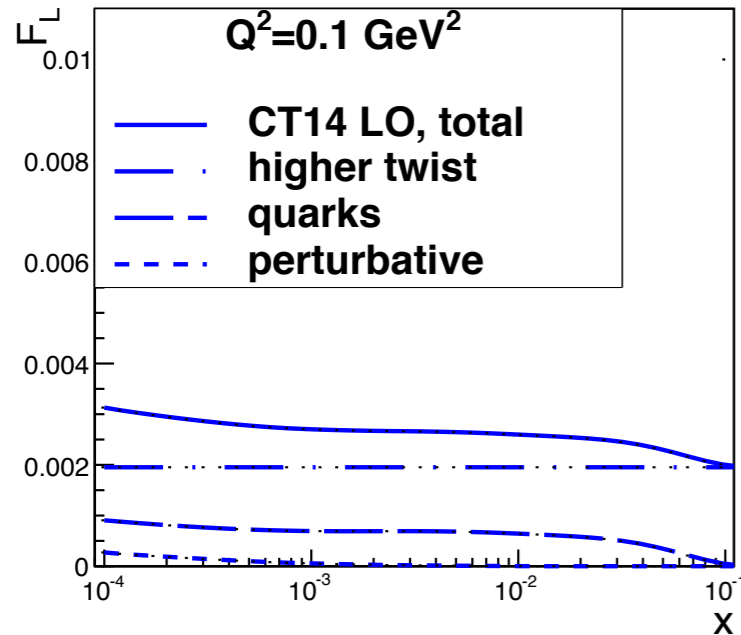
Perturbative component dominates at large Q^2 and low x

Higher twist significant at low to moderate Q^2 . Dominates at large x



Observe different vertical scale on different panels

$F_L(x)$ bins in Q^2



Components of the model

Perturbative component dominates at large Q^2 and low x

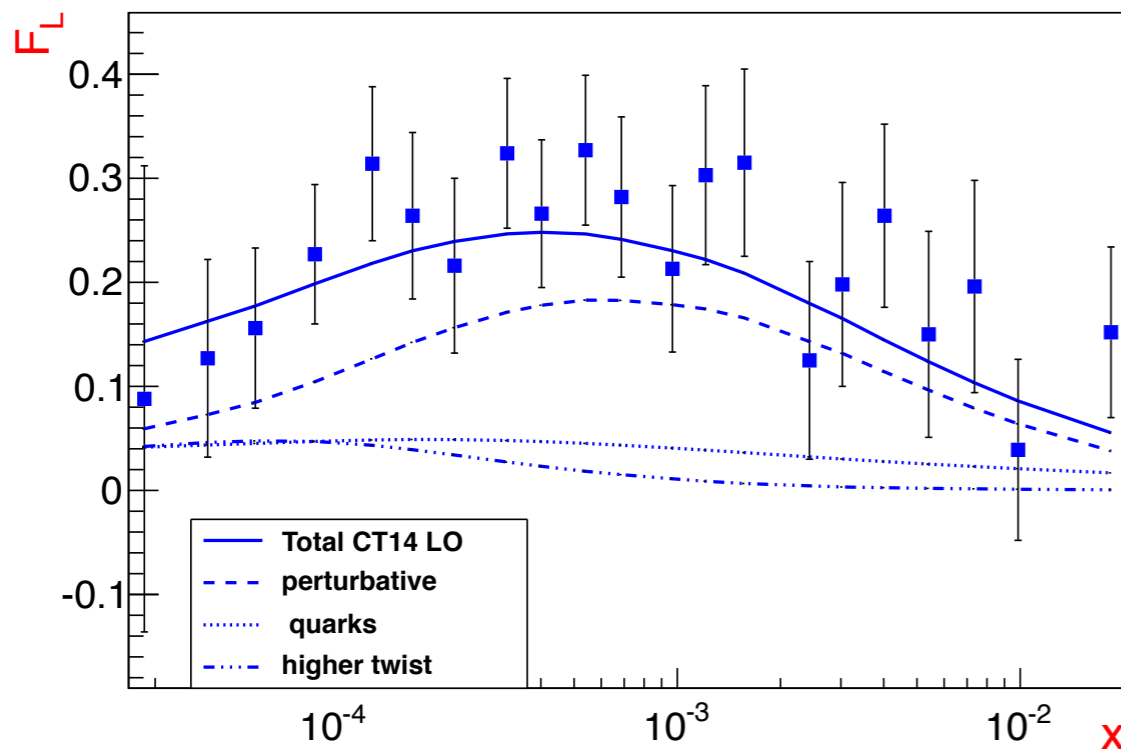
Higher twist flat as a function of x . Soft Pomeron-like behavior with intercept equal to unity

Observe different vertical scale on different panels

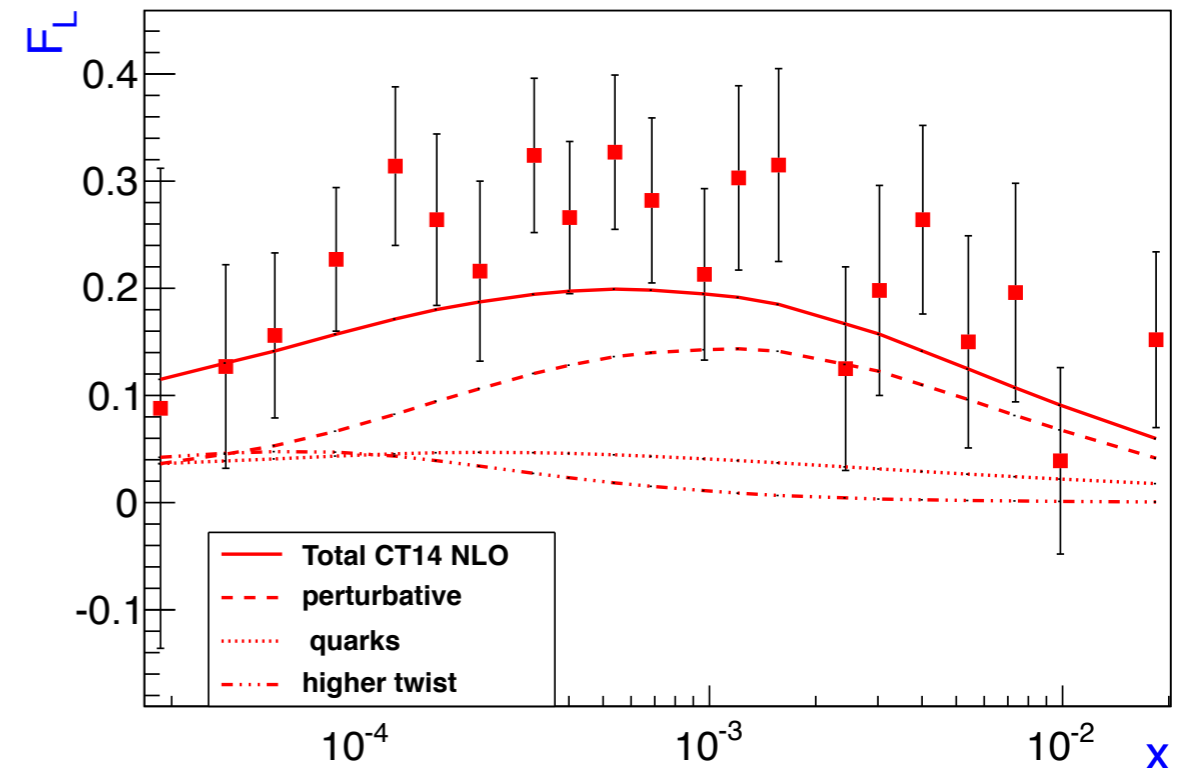
Comparison with H1 data

Q^2 1.5 2.0 2.5 3.0 3.5 4.0 4.5 5.0 6.0 8.0 12 15 20 25 35 45 60 90 120 150 200 250 346 636

$F_L(x, Q^2)$, H1 measurements (points @ different Q^2)

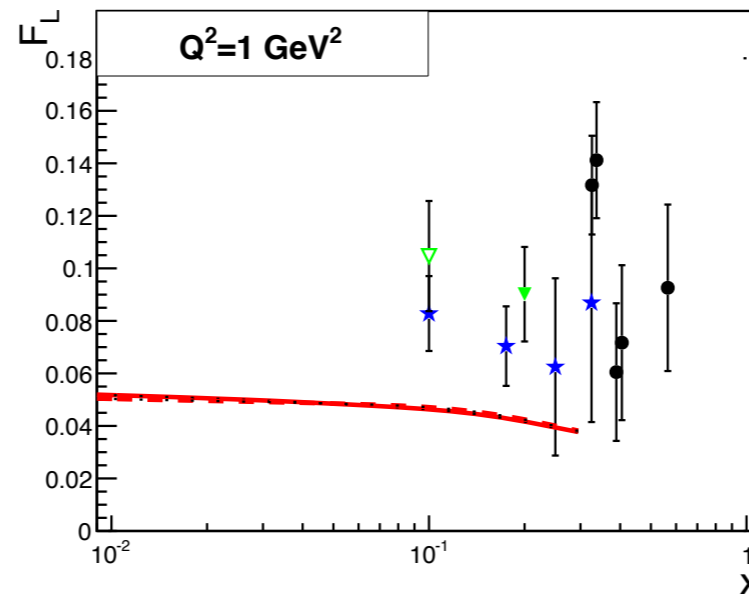
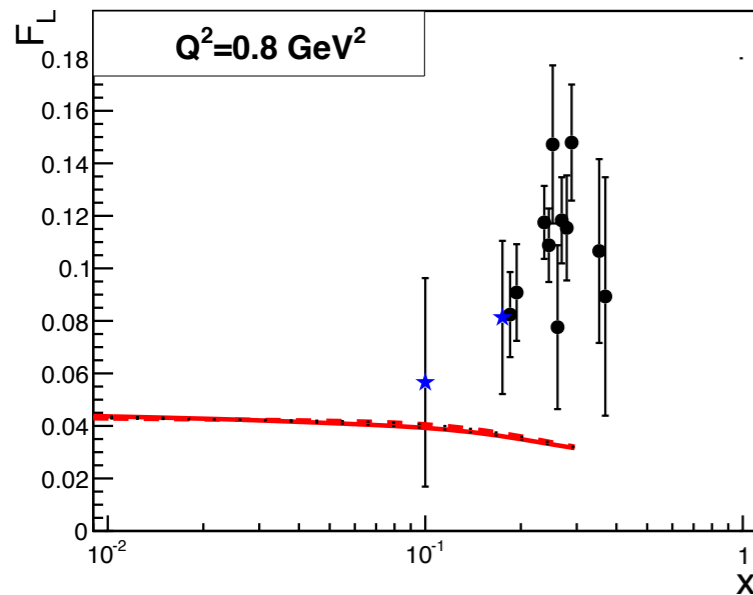
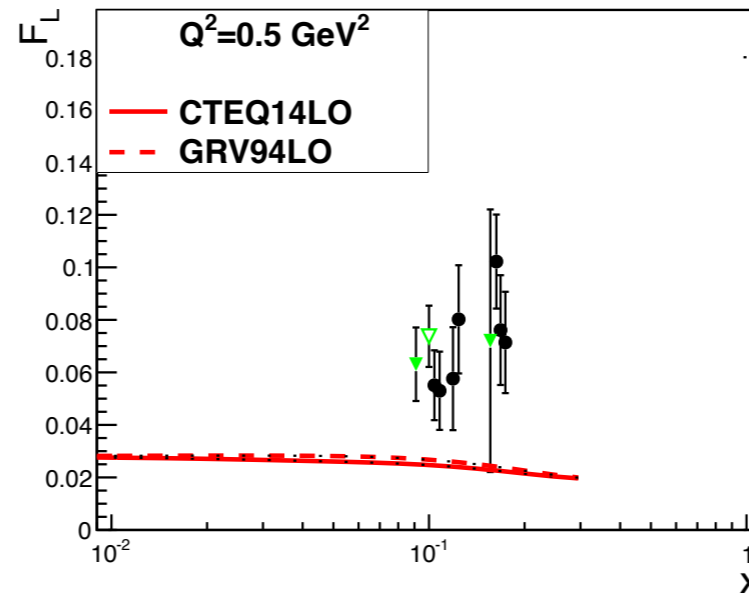
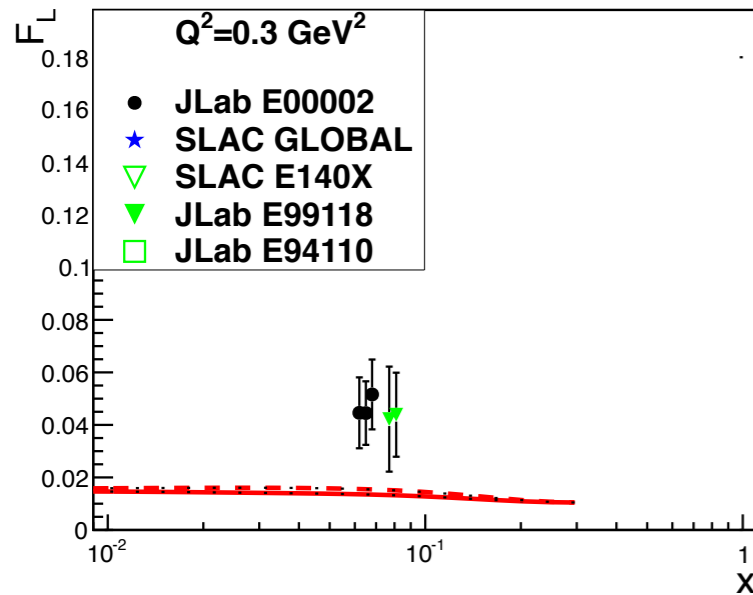


$F_L(x, Q^2)$, H1 measurements (points @ different Q^2)



- Good description of H1 data
- Mostly small x , and not low Q^2
- Most sensitive to perturbative component, variation of PDF

Comparison with JLAB and SLAC



- Comparison with low Q^2 data: SCLAC, JLAB
- Model underestimates the data
- However, model constructed for low x , beyond the region of applicability since data mostly at large x (additional contributions)
- Black points: new JLab data
- Caveat : SLAC and JLab(old) data are rebinned with models assumptions

Summary

- Knowledge of F_2 and F_L in the low Q^2 region necessary for radiative corrections
- Will be crucial for EIC precision measurement
- Revisited model for F_L at low Q^2 : k_T factorization + higher twist
- In the kinematic region of the model (low x and low Q^2), dependence on the PDFs is negligible
- Comparison with HERA: good matching to data, mostly sensitive to perturbative part, can be matched correctly by varying PDFs
- Comparison with JLab: model is underestimating the data, which are mostly at large x , beyond the region of applicability of the model
- Outlook: compare with other models (low x specific), other sources of HT, target mass corrections...