



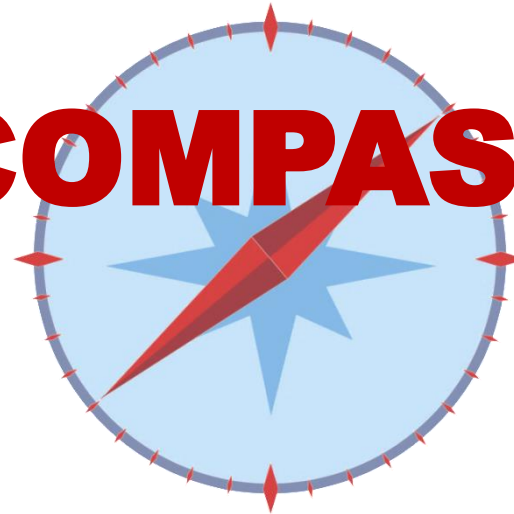
8th Workshop of the APS Topical Group on Hadronic Physics

GHP 2019

APRIL 10-12, 2019 • DENVER, CO

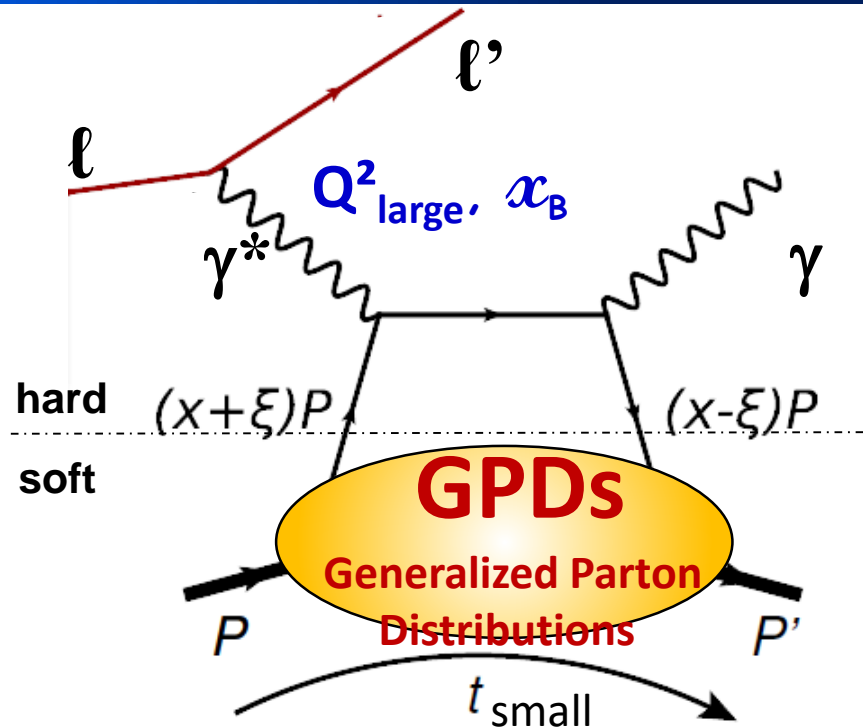
GPD at COMPASS at CERN

- 1- DVCS**
- 2- HEMP**



Nicole d'Hose – CEA – Université Paris-Saclay

Deeply virtual Compton scattering (DVCS)



D. Mueller *et al*, Fortsch. Phys. 42 (1994)

X.D. Ji, PRL 78 (1997), PRD 55 (1997)

A. V. Radyushkin, PLB 385 (1996), PRD 56 (1997)

DVCS: $\ell p \rightarrow \ell' p' \gamma$

the golden channel

because it interferes with
the Bethe-Heitler process

also meson production

$\ell p \rightarrow \ell' p' \pi, \rho, \omega$ or ϕ or $J/\psi \dots$

The GPDs depend on the following variables:

x : average long. momentum

ξ : long. mom. difference

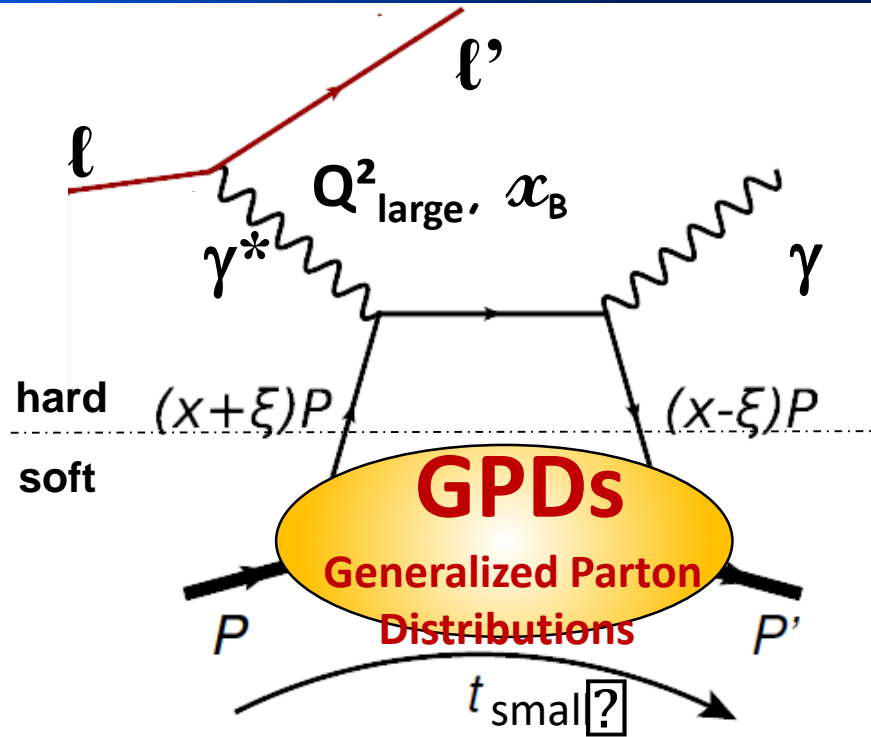
t : four-momentum transfer
related to b_{\perp} via Fourier transform

The variables measured in the experiment:

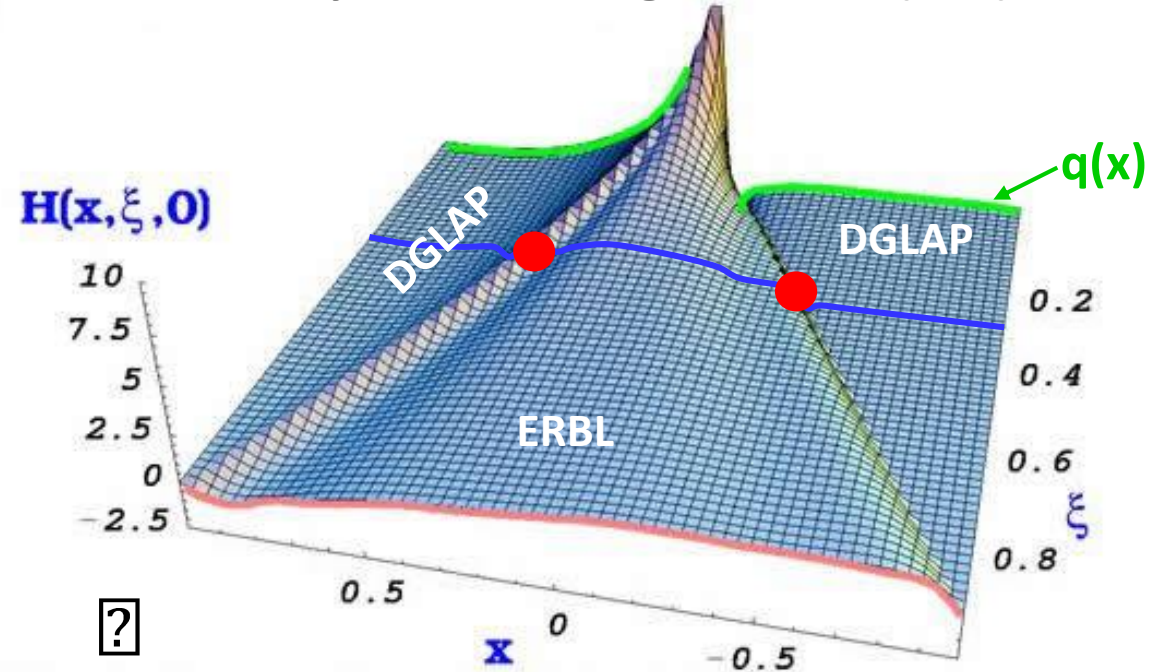
$E_{\ell}, Q^2, x_B \sim 2\xi / (1 + \xi),$

t (or $\theta_{\gamma^* \gamma}$) and ϕ ($\ell \ell'$ plane / $\gamma \gamma^*$ plane)

Deeply virtual Compton scattering (DVCS)



From Goeke, Polyakov, Vanderhaeghen, PPNP47 (2001)



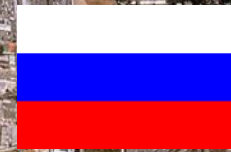
The amplitude DVCS at LT & $\xi=0$ in α_s (GPD \mathcal{H}):

$$\mathcal{H} = \int_{-1}^{+1} dx \frac{\mathcal{H}(x, \xi, t)}{x - \xi + i\epsilon} = \mathcal{P} \int_{-1}^{+1} dx \frac{\mathcal{H}(x, \xi, t)}{x - \xi} - i\pi \mathcal{H}(x \pm \xi, x, t)$$

In an experiment we measure Compton Form Factor \mathcal{H}

$$\text{Re}\mathcal{H}(\xi, t) = \int dx \frac{\text{Im}\mathcal{H}(x, t)}{x - \xi} + D(t)$$

COMPASS: Versatile facility with hadron (π^\pm , K^\pm , p ...) & lepton (polarized μ^\pm) beams of high energy ~ 200 GeV

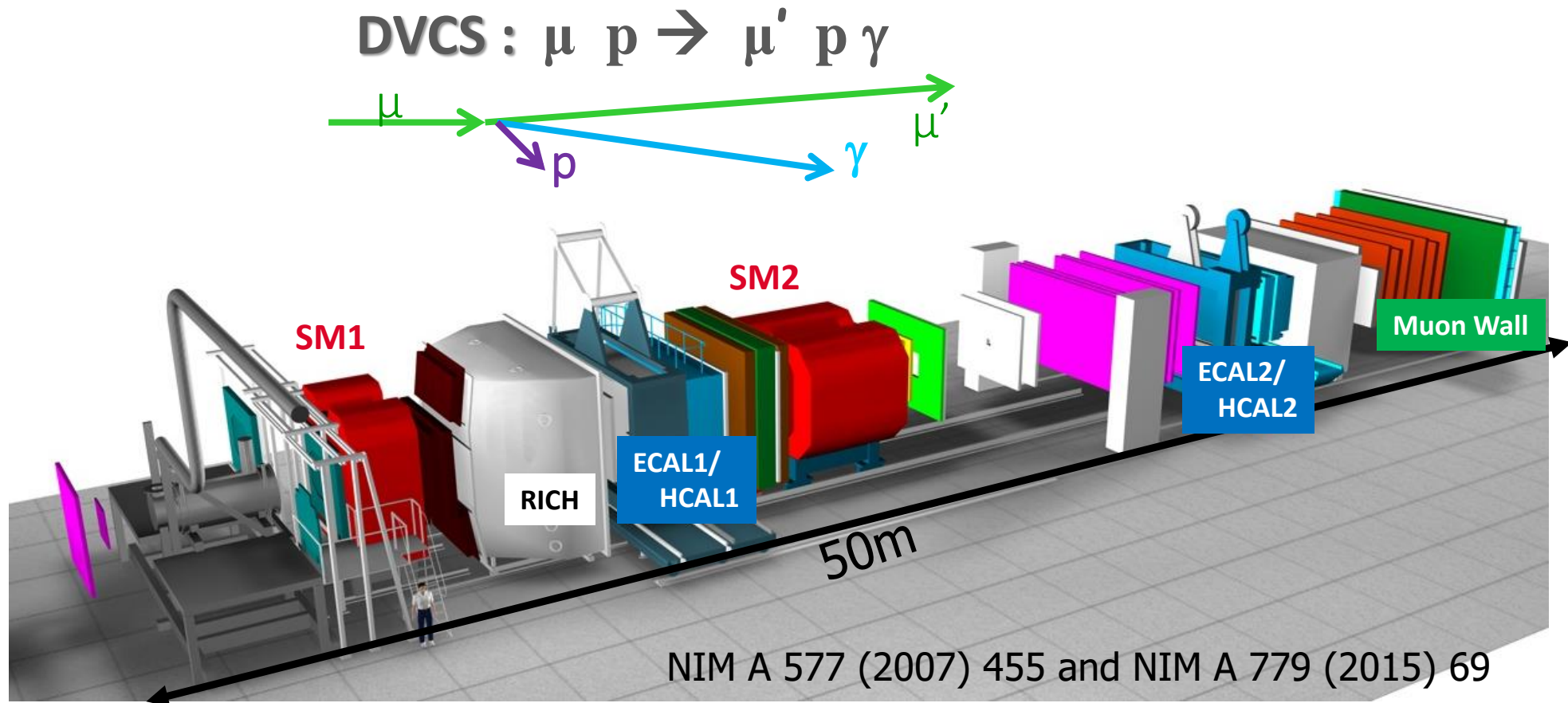


LHC

COMPASS

SPS

The DVCS experiment at COMPASS



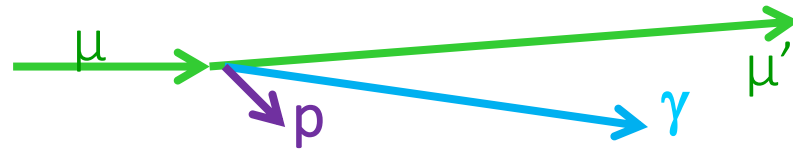
Two stage magnetic spectrometer for **large angular & momentum acceptance**

Particle identification with:

- Ring Imaging Cerenkov Counter
- Electromagnetic calorimeters (**ECAL1** and **ECAL2**)
- Hadronic calorimeters
- Hadron absorbers

The DVCS experiment at COMPASS

$$\text{DVCS : } \mu p \rightarrow \mu' p \gamma$$



New equipments:

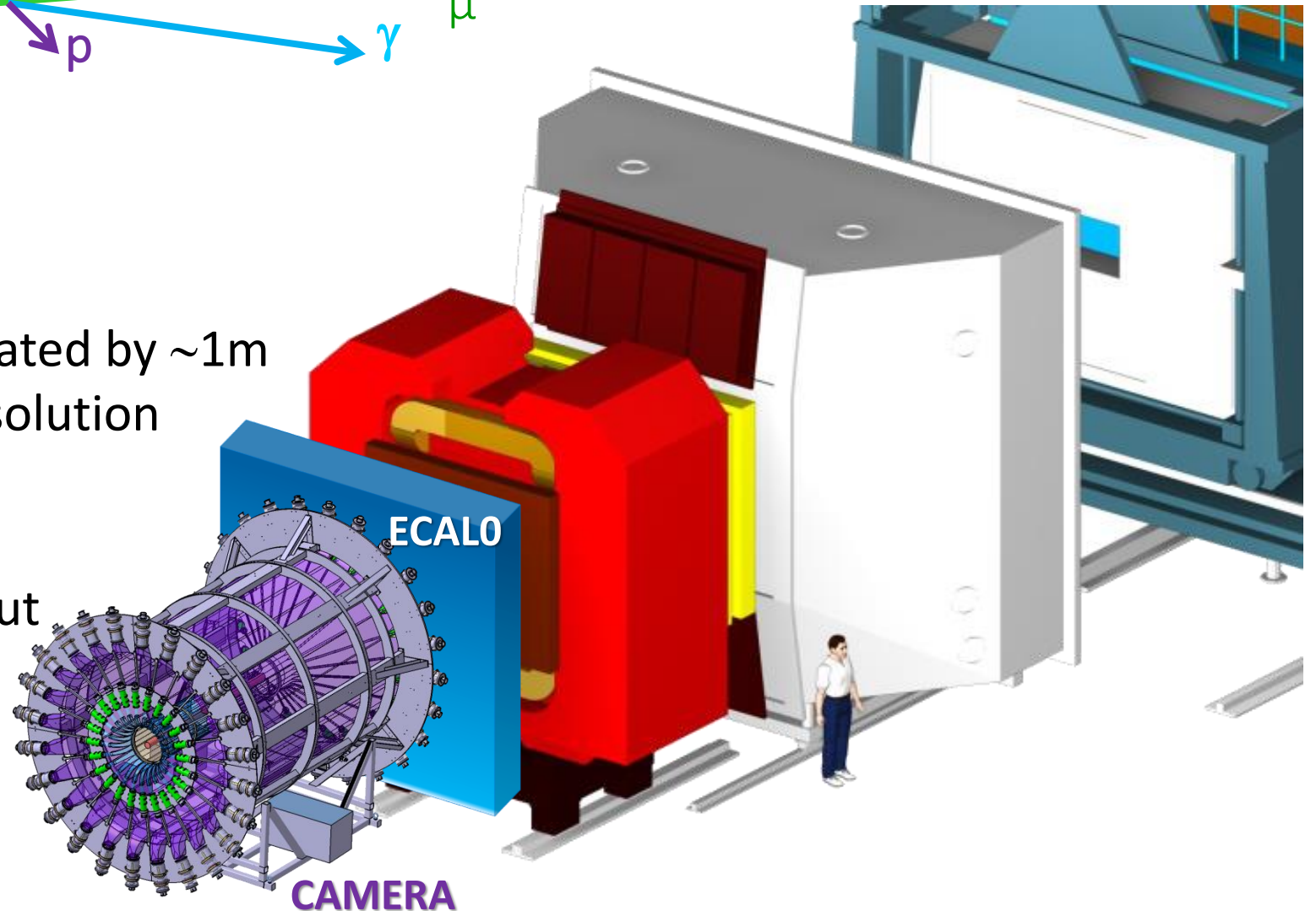
➤ **2.5m LH2 target**

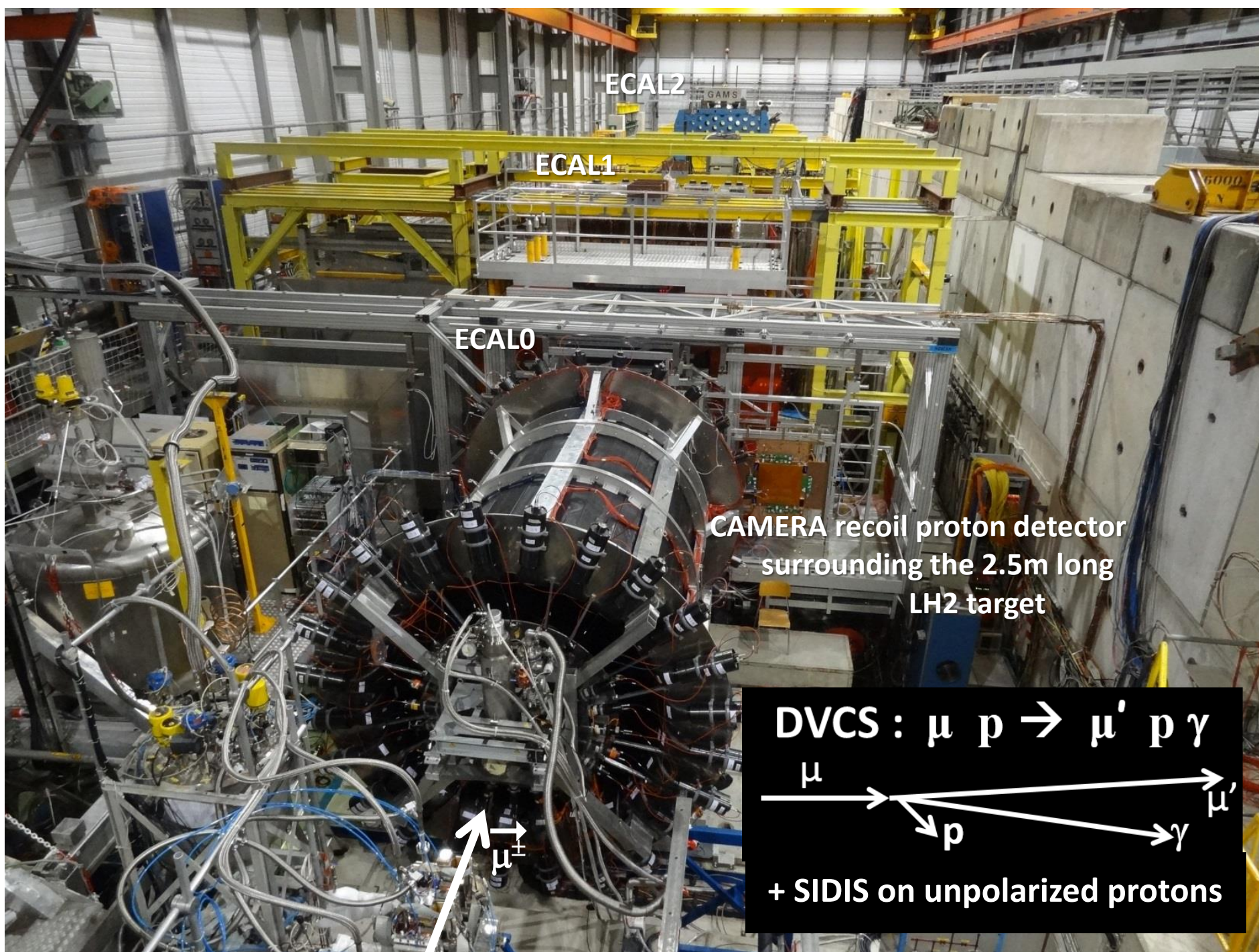
➤ **4m ToF Barrel CAMERA**

24 inner & outer scintillators separated by ~1m
1 GHz SADC readout, 330ps ToF resolution

➤ **ECALO : 2 × 2 m²**

Shashlyk modules + MAPD readout
one module is made of
9 cells (4×4 cm²)
= 194 modules or 1746 cells





2012: 1 month pilot run

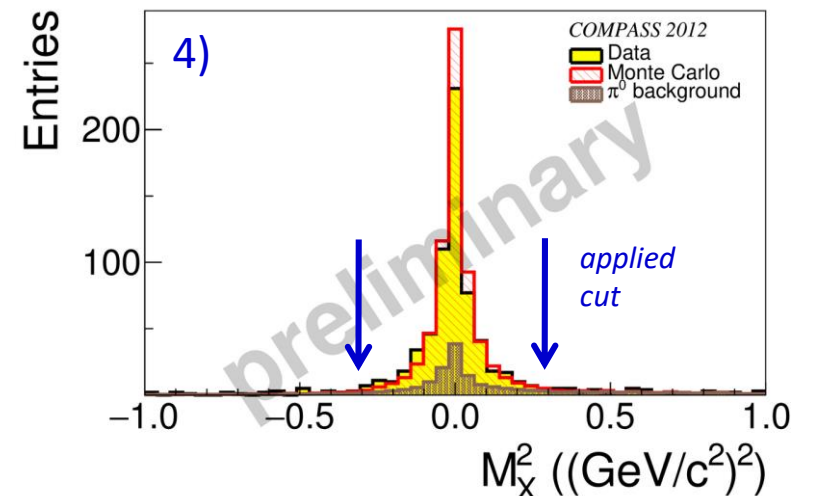
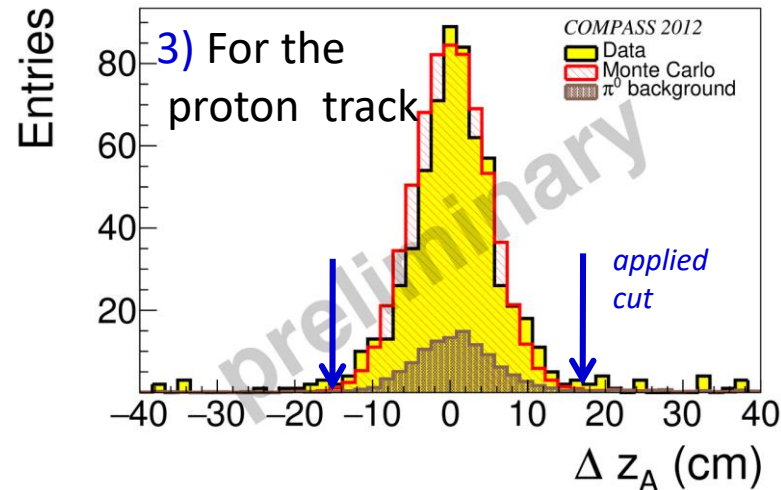
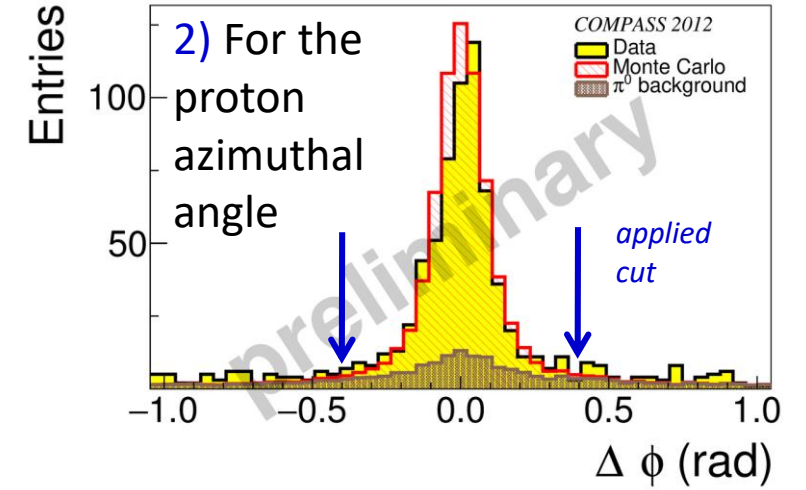
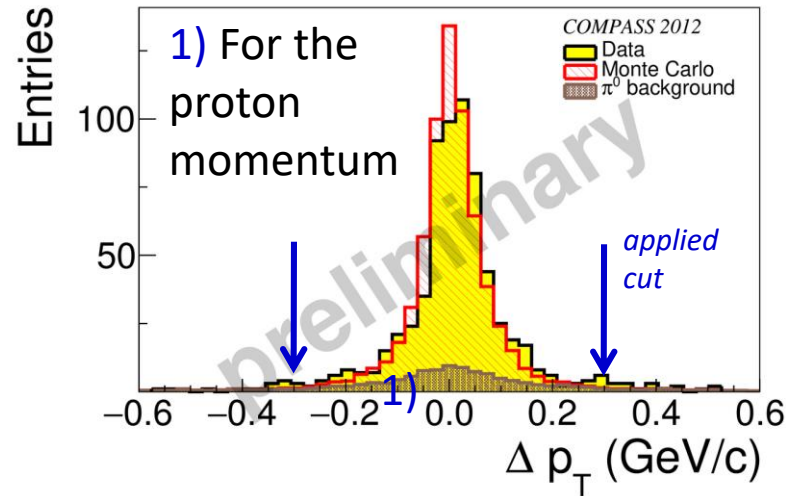
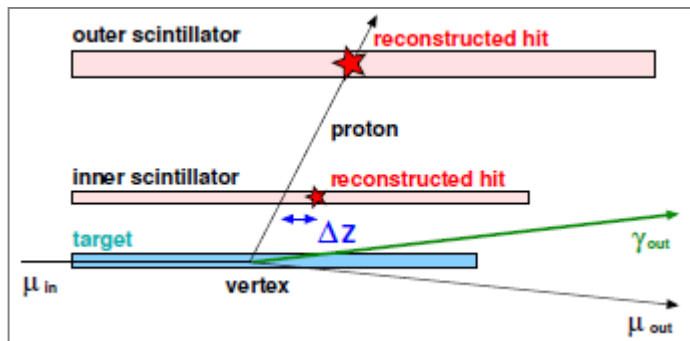
2016 -17: 2 x 6 months
data taking

$x_{Bj} > 0.03$ $10 < \nu < 32 \text{ GeV}$
with π^0 contamination

Comparison between the observables given by the spectro or by CAMERA

DVCS: $\mu p \rightarrow \mu' p \gamma$

- 1) $\Delta p_T = p_T^{\text{cam}} - p_T^{\text{spec}}$
- 2) $\Delta \varphi = \varphi^{\text{cam}} - \varphi^{\text{spec}}$
- 3) $\Delta z_A = z_A^{\text{cam}} - z_A^{\text{ZB and vertex}}$
- 4) $M_{X=0}^2 = (p_{\mu_{\text{in}}} + p_{p_{\text{in}}} - p_{\mu_{\text{out}}} - p_{p_{\text{out}}} - p_{\gamma})^2$



π^0 are one of the main background sources for excl. photon events.

Two possible case:

- **Visible** (both γ detected \rightarrow subtracted)

the DVCS photon after all exclusivity cuts is combined with all detected photons below the DVCS threshold: 4,5,10 GeV in ECALO, 1, 2

- **Invisible** (one γ lost \rightarrow estimated by MC)

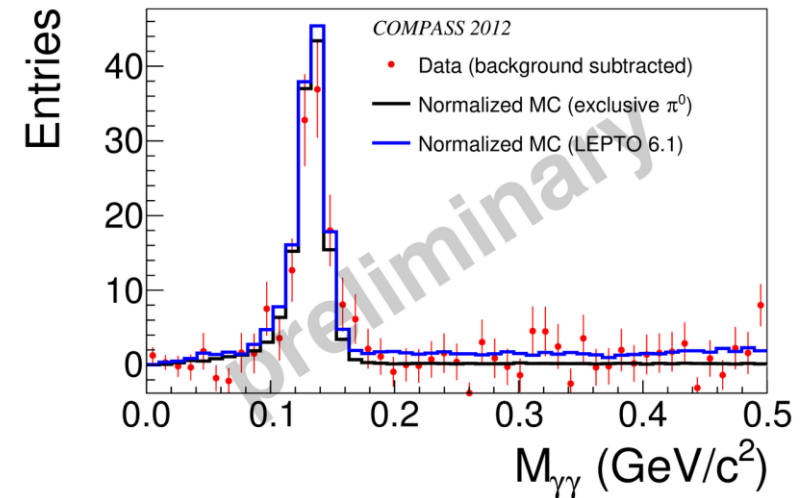
➤ **Semi-inclusive LEPTO 6.1**

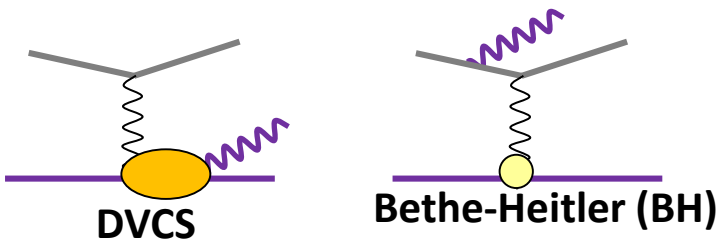
➤ **Exclusive HEPGEN π^0**
(Goloskokov-Kroll model)

Comparing the two components to the data allows the determination of their relative normalisation.

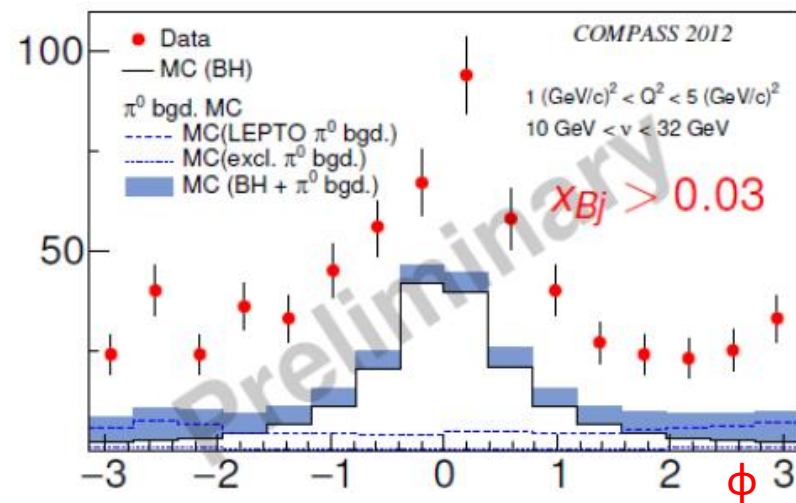
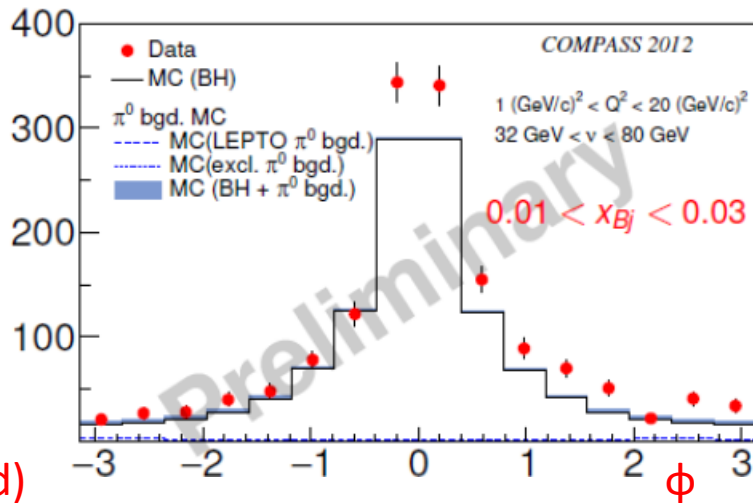
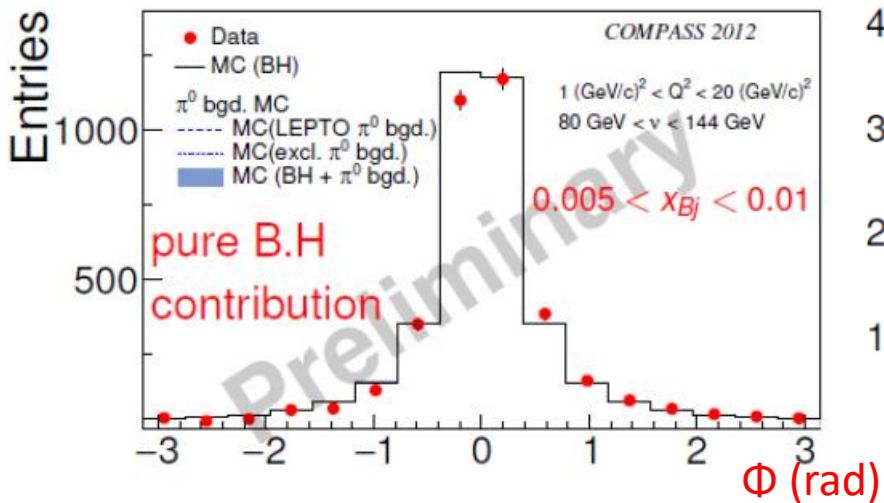
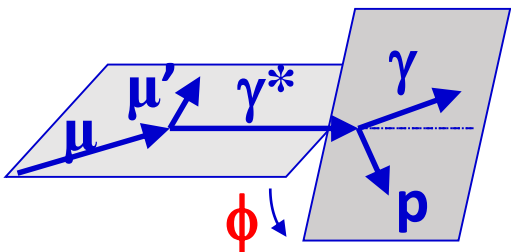
The sum of the 2 components is normalized to the visible π^0 contamination in the $M_{\gamma\gamma}$ peak

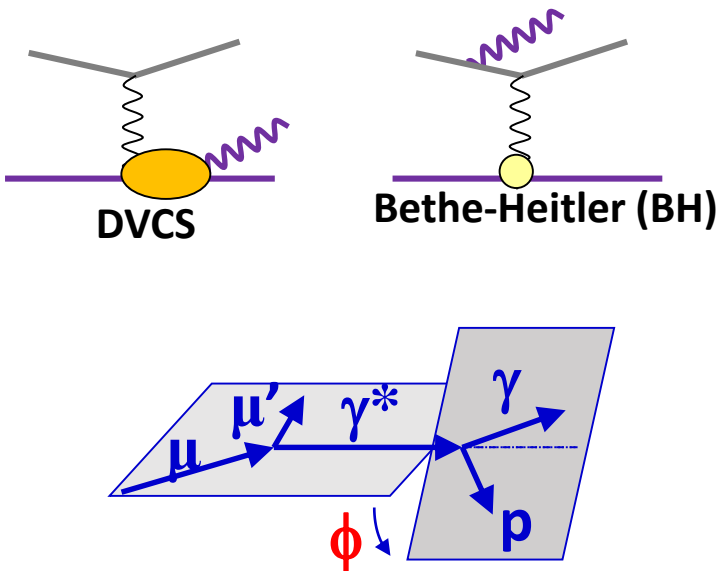
Visible leaking π^0 in the data



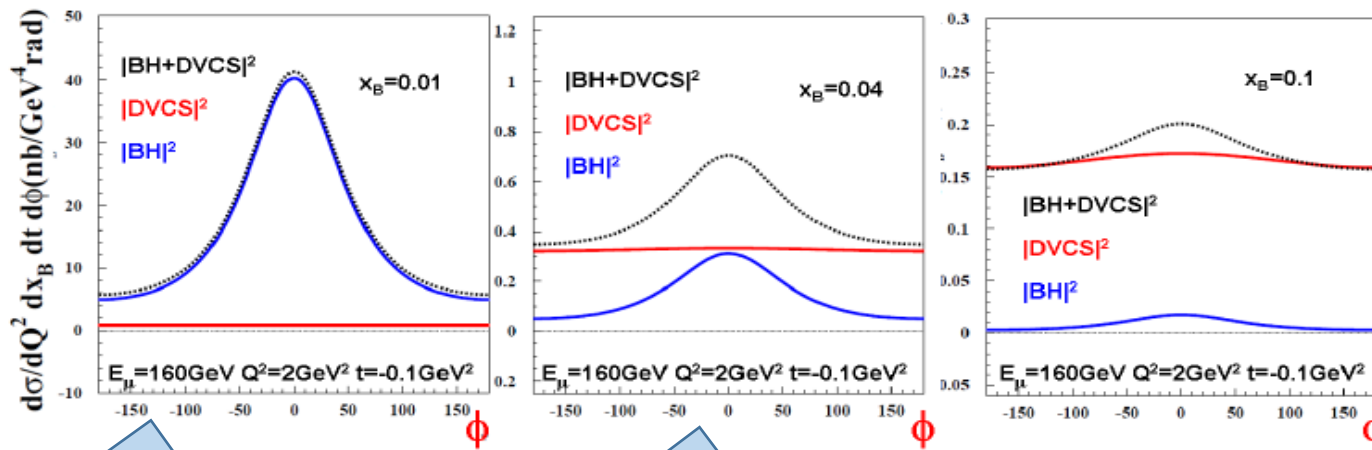


$$d\sigma \propto |T^{BH}|^2 + \text{Interference Term} + |T^{DVCS}|^2$$





$$d\sigma \propto |T^{BH}|^2 + \text{Interference Term} + |T^{DVCS}|^2$$



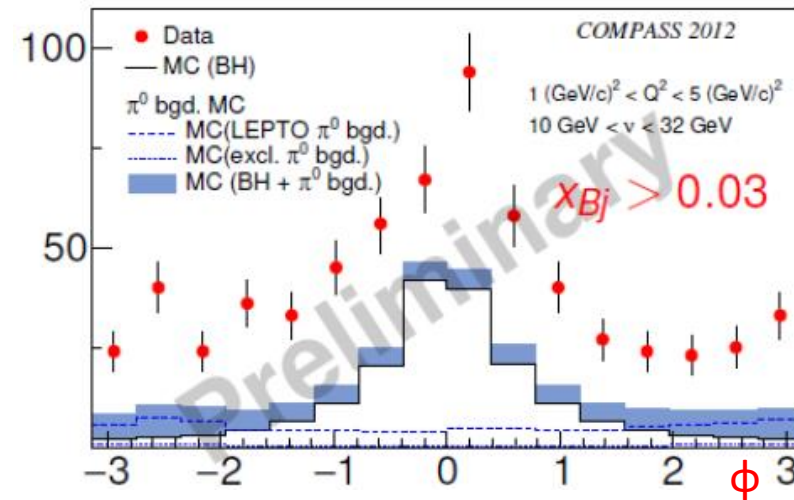
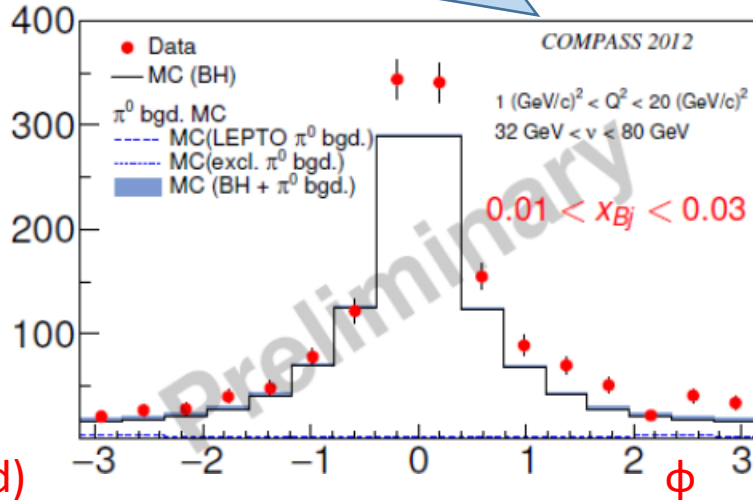
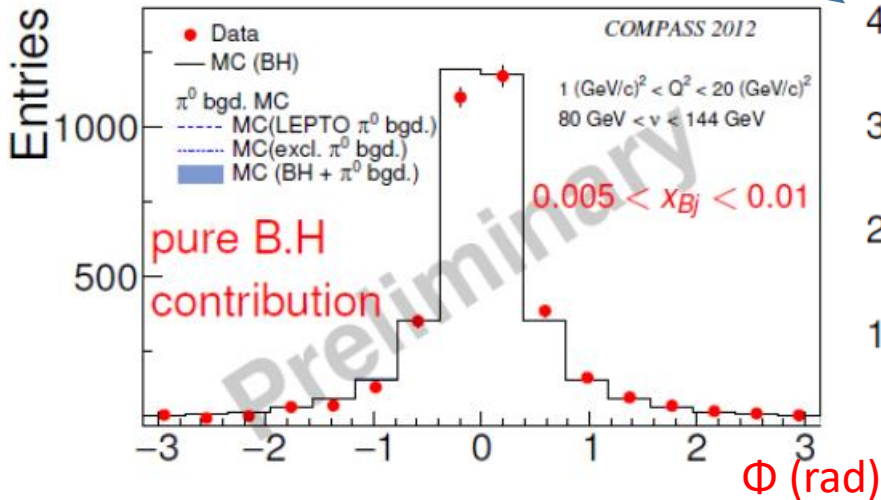
BH dominates
Reference yield

DVCS ampl. via interference

DVCS dominates - Study of $d\sigma^{DVCS}/dt$

Only for H1, ZEUS, COMPASS

Jlab, HERMES,
H1, COMPASS



when BH is not dominant

At COMPASS using polarized positive and negative muon beams:

$$S_{CS,U} \equiv d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-} = 2[d\sigma^{BH} + d\sigma_{unpol}^{DVCS} + \text{Im } I]$$

$$= 2[d\sigma^{BH} + c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi + s_1^I \sin \phi + s_2^I \sin 2\phi]$$

calculable
can be subtracted

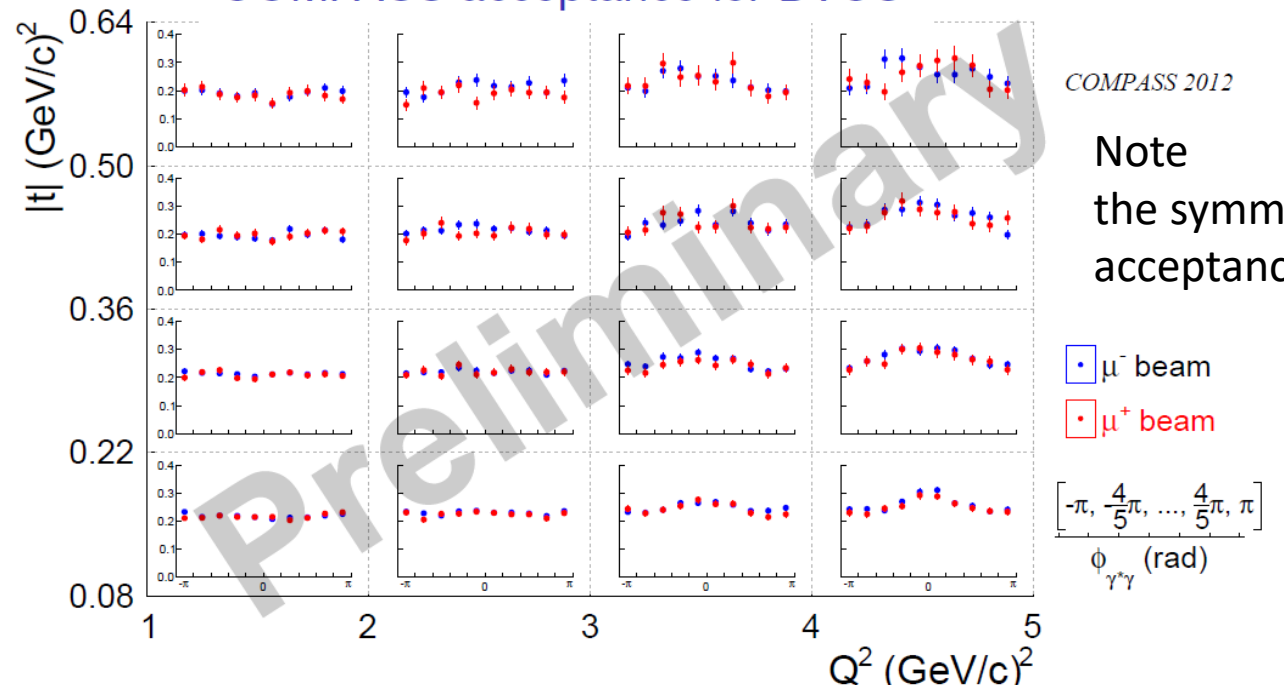
All the other terms are cancelled in the integration over ϕ

$$\frac{d^3\sigma_T^{\mu p}}{dQ^2 d\nu dt} = \int_{-\pi}^{\pi} d\phi (d\sigma - d\sigma^{BH}) \propto c_0^{DVCS}$$

$$\frac{d\sigma^{\gamma^* p}}{dt} = \frac{1}{\Gamma(Q^2, \nu, E_\mu)} \frac{d^3\sigma_T^{\mu p}}{dQ^2 d\nu dt}$$

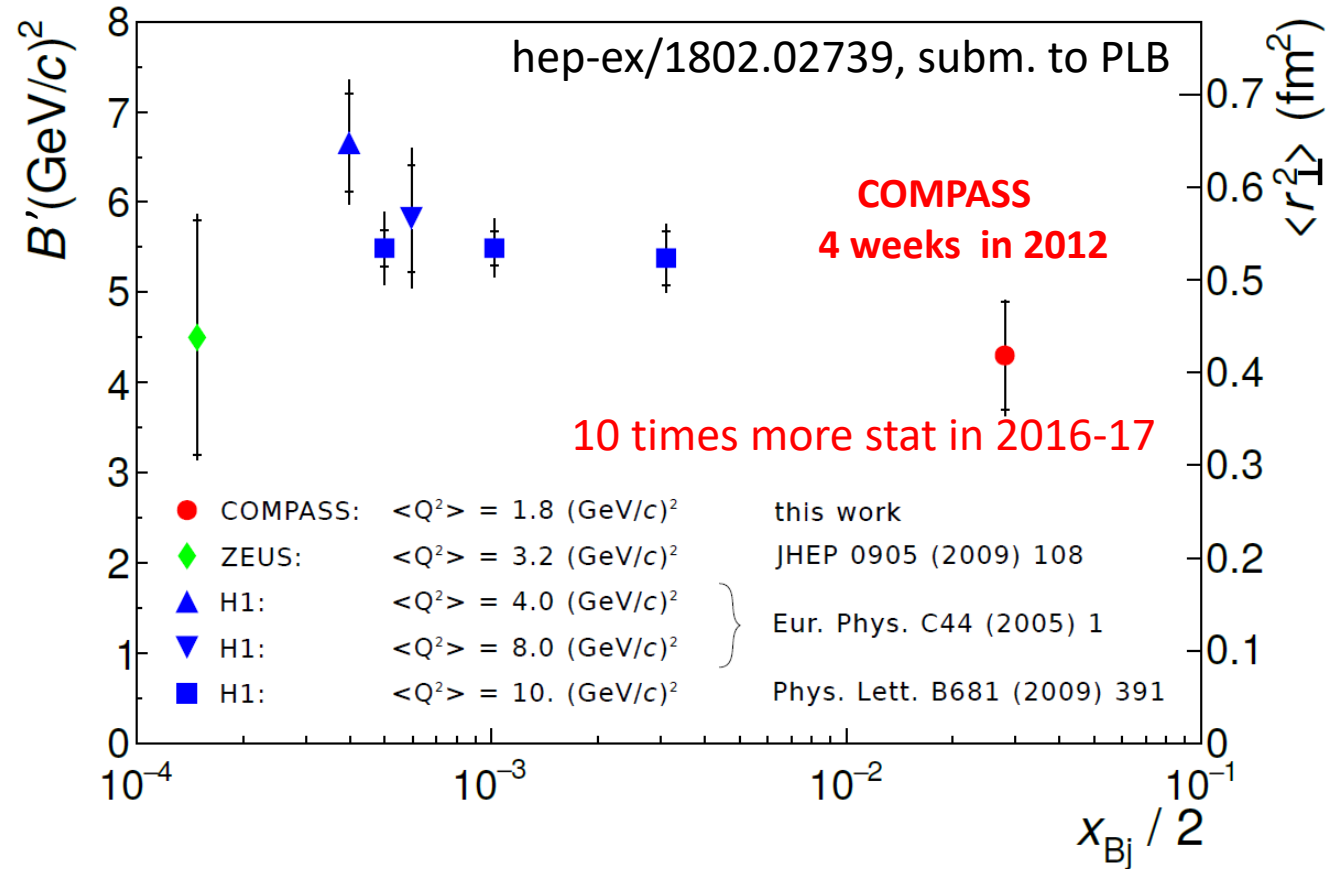
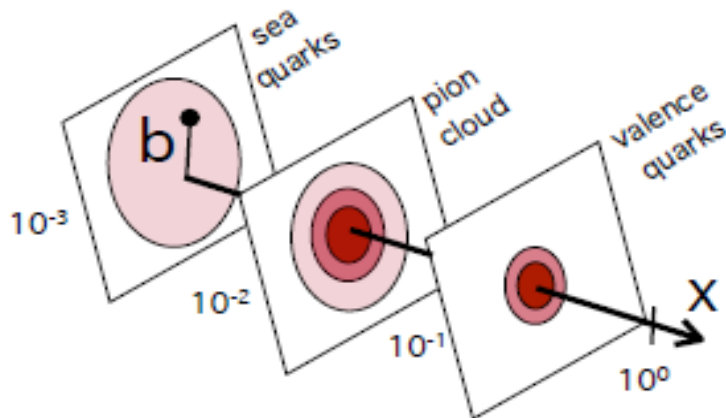
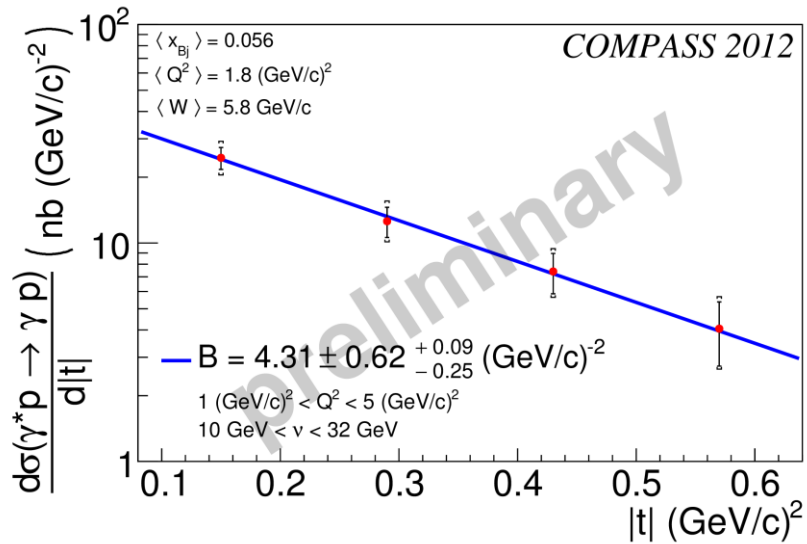
Flux for transverse virtual photons

COMPASS acceptance for DVCS



COMPASS 2012 Transverse extension of partons in the sea quark range

$$d\sigma^{DVCS}/dt = e^{-B'|t|} = c_0^{DVCS}$$



$$B = (4.3 \pm 0.6_{\text{stat}} \pm 0.1_{\text{sys}}) \text{ (GeV/c)}^{-2}$$

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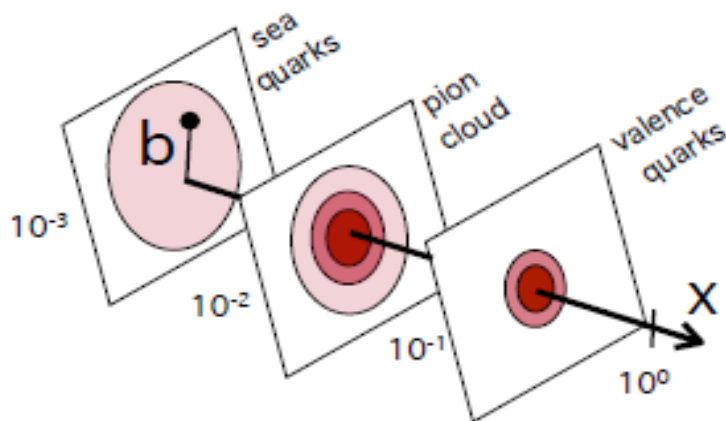
At COMPASS:

$\langle x_{Bj} \rangle = 0.056$; $\langle Q^2 \rangle = 1.8 \text{ GeV}^2$;
 t varies from 0.08 to 0.64 GeV^2

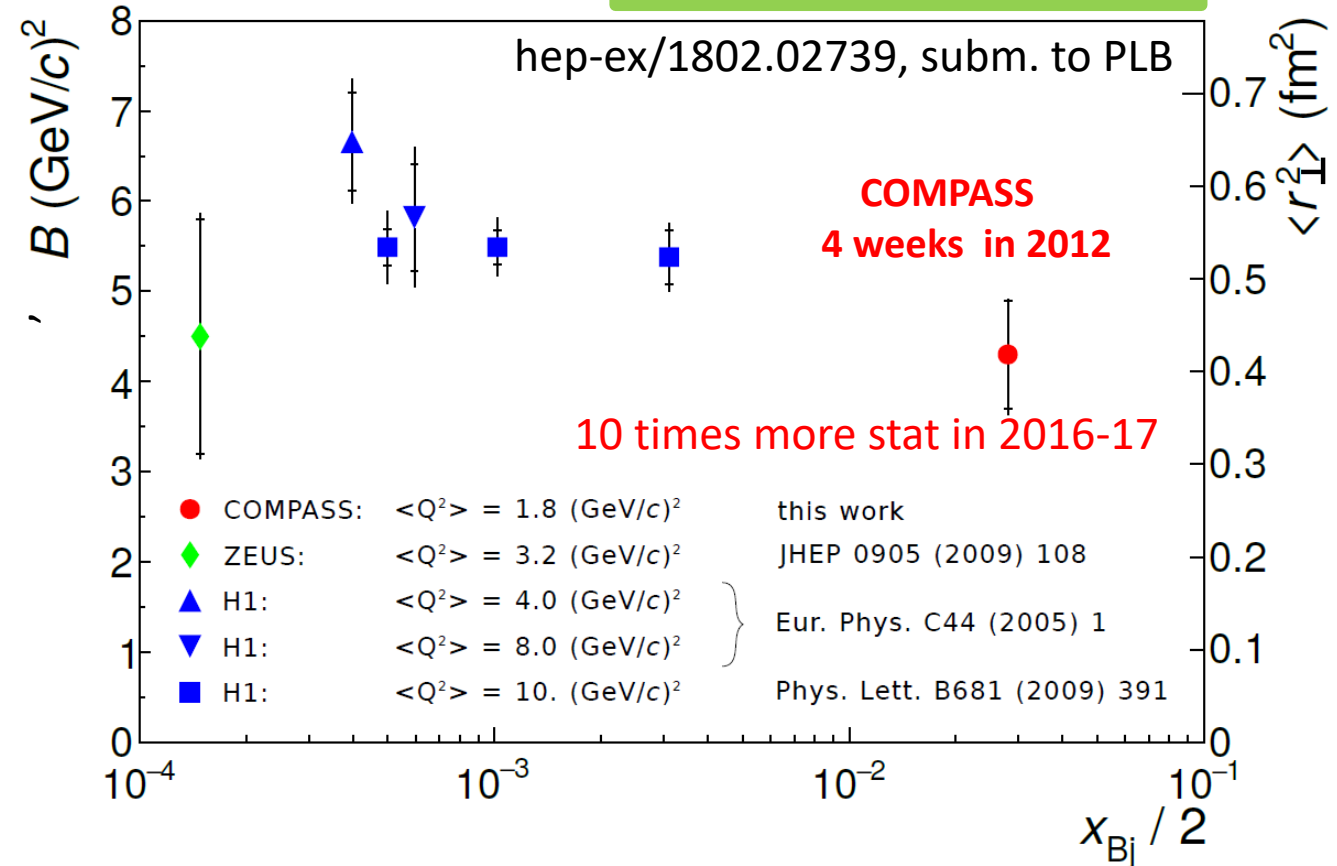
At small x_{Bj} and small t :

$$c_0^{DVCS} \propto 4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) - \frac{t}{M^2} \mathcal{E}\mathcal{E}^*$$

Dominance of $Im\mathcal{H}$
 (with respect of $Re\mathcal{H}$ and other CFP)



$$\langle r_{\perp}^2(x_B) \rangle \approx 2B'(x_B)$$



$$B = (4.3 \pm 0.6_{\text{stat}} \pm 0.1_{\text{sys}}) (\text{GeV}/c)^{-2}$$

$$\sqrt{\langle r_{\perp}^2 \rangle} = (0.58 \pm 0.04_{\text{stat}} \pm 0.01_{\text{sys}} \pm 0.04_{\text{model}}) \text{ fm}$$

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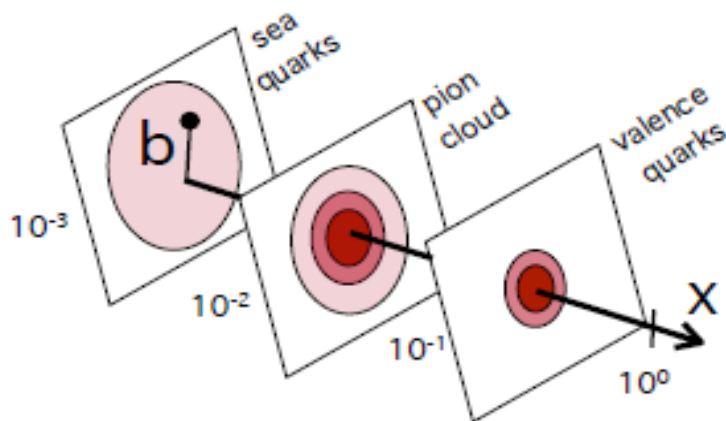
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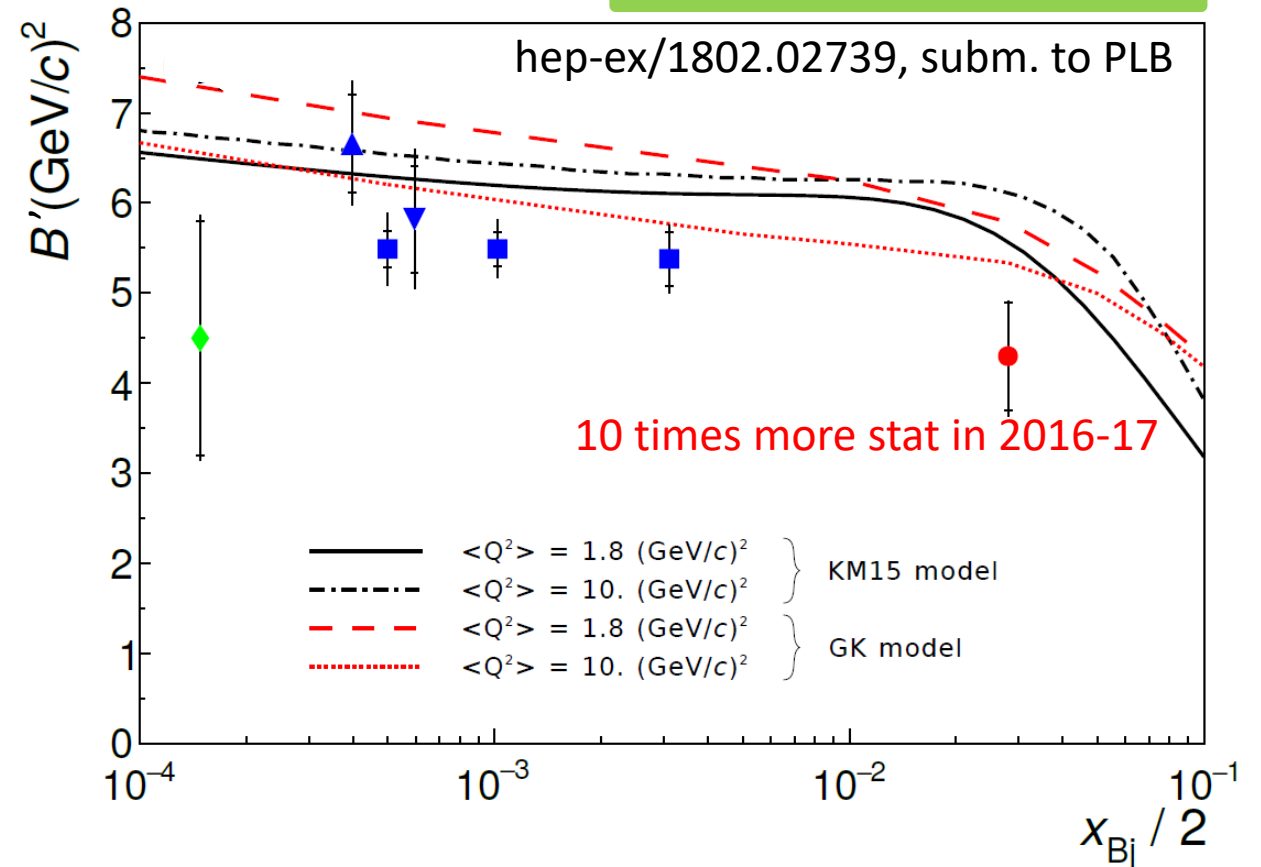
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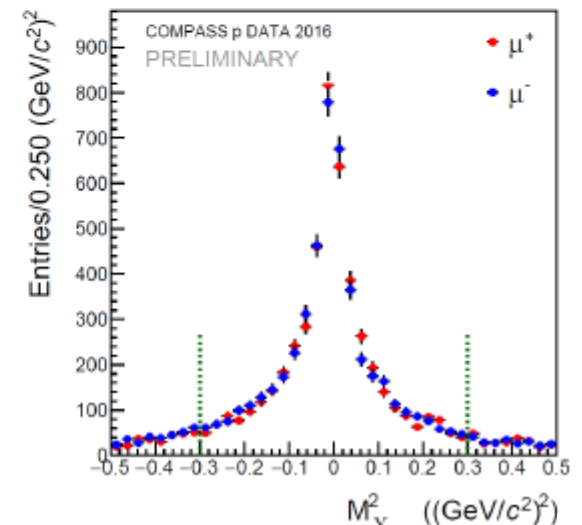
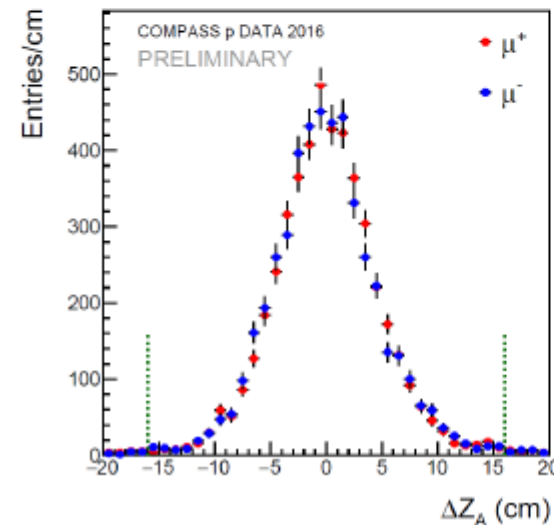
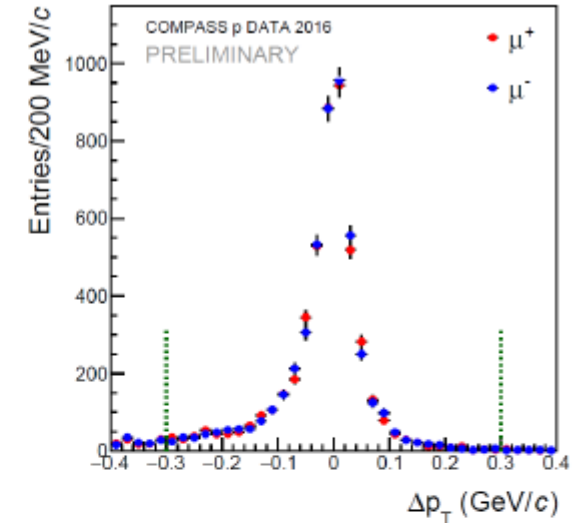
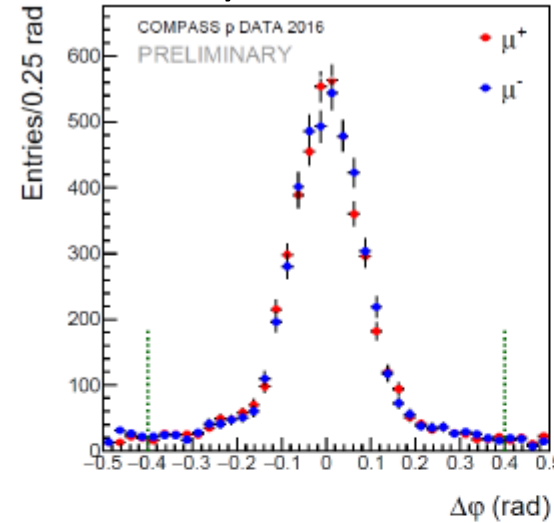
Comparison between the observables given by the spectro or by CAMERA

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- 1) $\Delta p_T = p_T^{\text{cam}} - p_T^{\text{spec}}$
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- 4) $M_{X=0}^2 = (p_{\mu_{\text{in}}} + p_{p_{\text{in}}} - p_{\mu_{\text{out}}} - p_{p_{\text{out}}} - p_{\gamma})^2$

$$d\sigma^{\leftarrow+} + d\sigma^{\rightarrow-} \quad t\text{-slope}$$

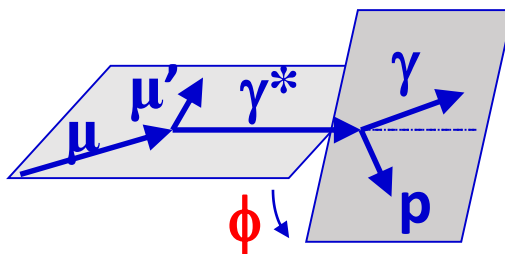
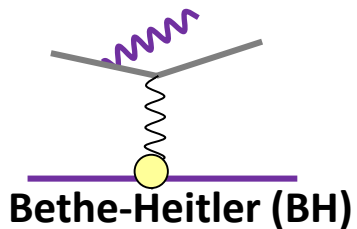
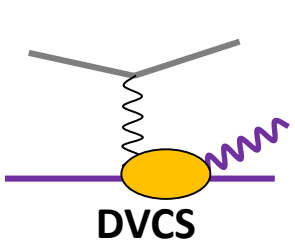
$$d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-} \quad d\text{-term}$$



COMPASS 2016-17

First insight

Distributions in ϕ

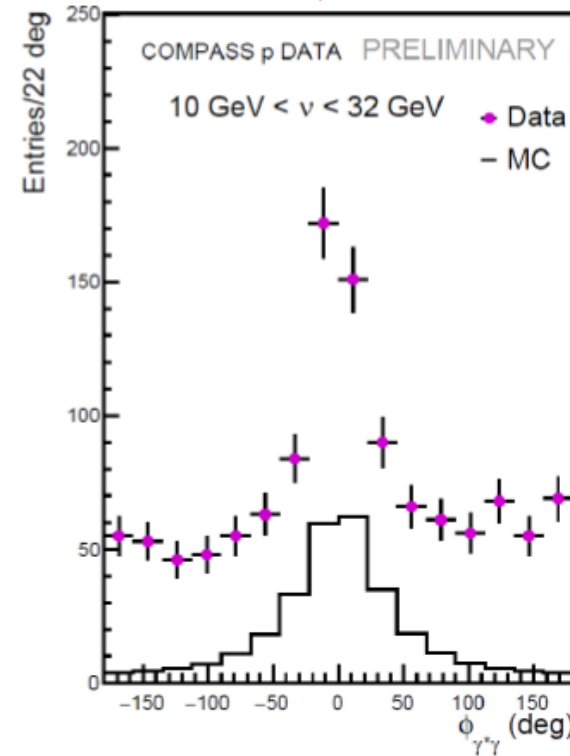
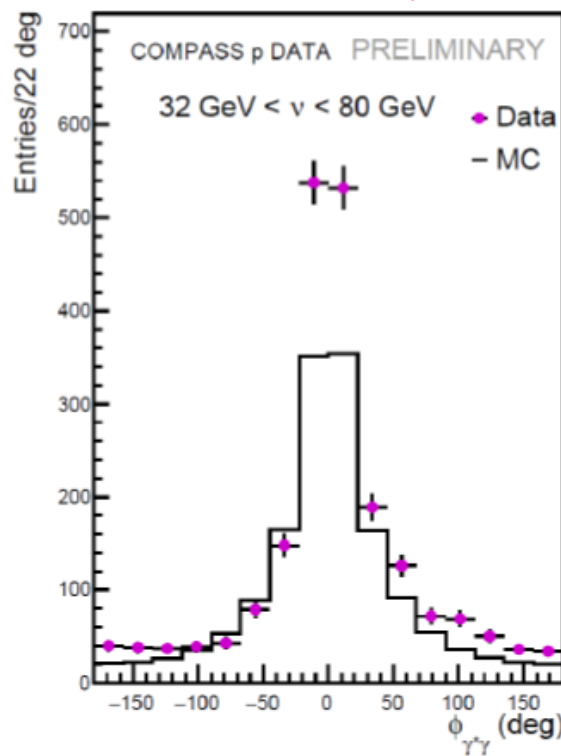
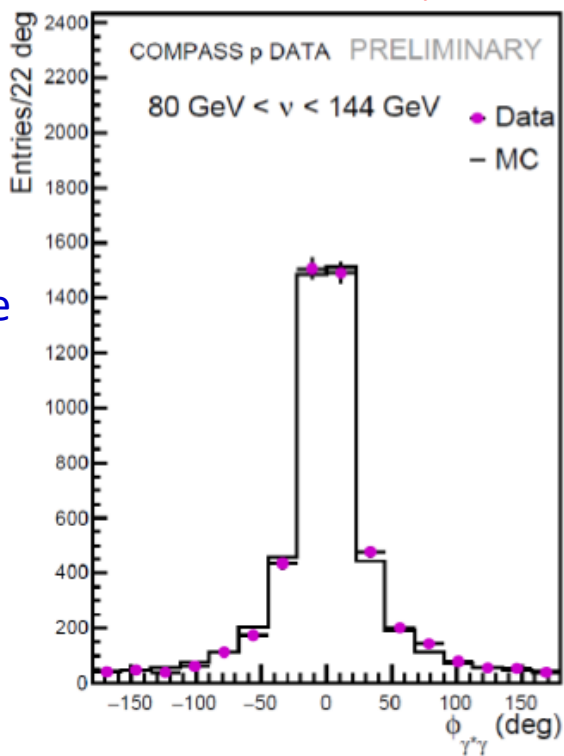


Only 13% of 2016-17 data

$$0.005 < x_{Bj} < 0.01$$

$$0.01 < x_{Bj} < 0.03$$

$$x_{Bj} > 0.03$$



No "invisible" π^0 still to be removed

a significant DVCS contribution will allow to study $d\sigma^{DVCS}/dt = e^{-B'|t|} = c_0^{DVCS}$

BH expected to contribute only
BH MC is normalized to this bin

BLUE WATERS

This research is part of the Blue Waters sustained-petascale computing project, which is supported by the National Science Foundation (awards OCI-0725070 and ACI-1238993) and the state of Illinois. Blue Waters is a joint effort of the University of Illinois at Urbana-Champaign and its National Center for Supercomputing Applications. This work is also part of the "Mapping Proton Quark Structure using Petabytes of COMPASS Data" PRAC allocation supported by the National Science Foundation (award number OCI 1713684).

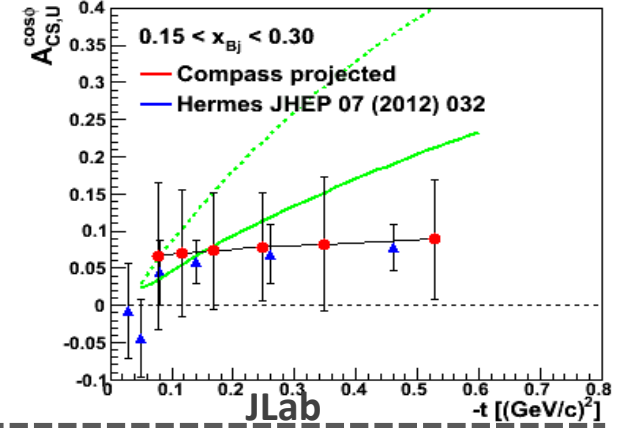
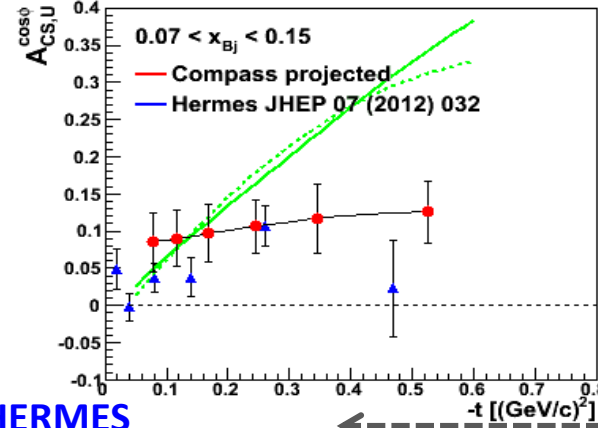
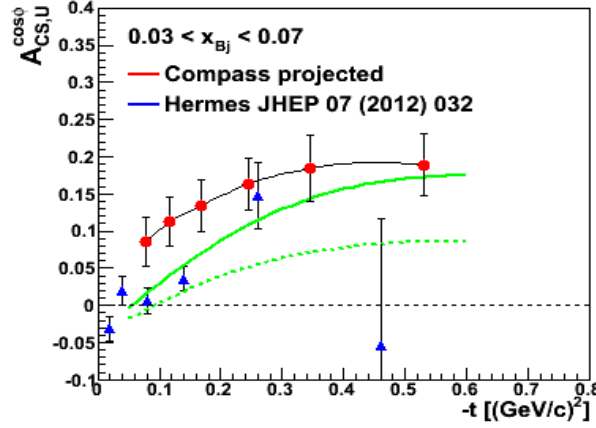
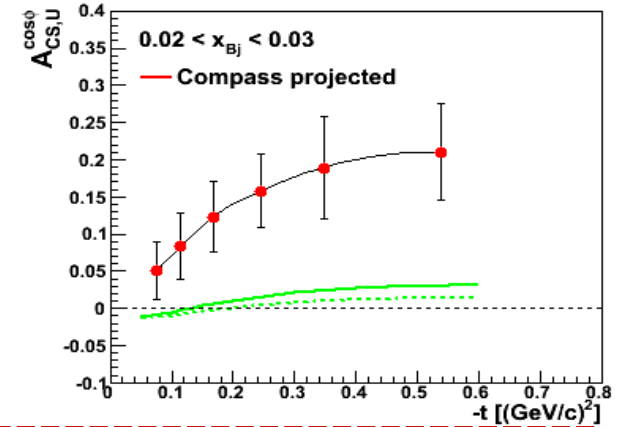
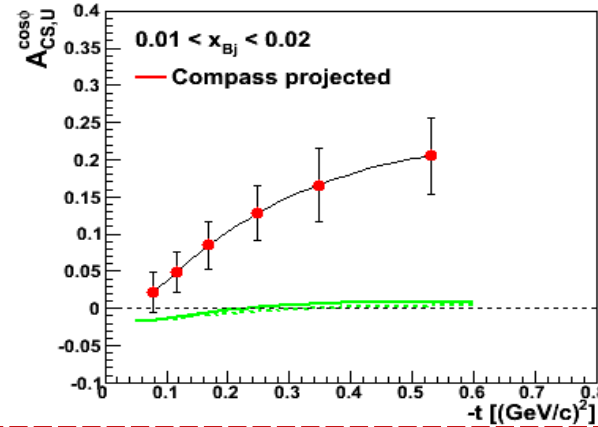
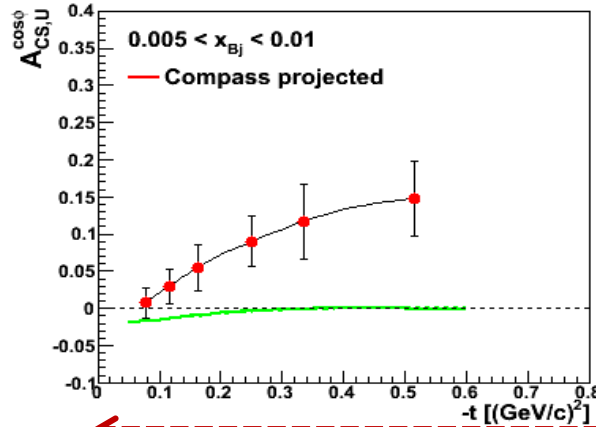
Beam Charge and Spin Diff. @ COMPASS

$$\mathcal{D}_{CS,U} \equiv d\sigma^{\leftarrow+} - d\sigma^{\rightarrow-} = 2[d\sigma_{pol}^{DVCS} + \text{Re } I] \xrightarrow{L.T.} c_0^I + c_1^I \cos \phi$$

$\text{Re } \mathcal{H} > 0$ at H1
 < 0 at HERMES
 Value of x_{Bj} for the node?

$$c_1^I = \text{Re } F_1 \mathcal{H}$$

Predictions with
VGG
KM10



The knowledge of
 $\text{Re } F_1 \mathcal{H}$ and $\text{Im } F_1 \mathcal{H}$
 is essential to play with
 the dispersion relation
 to extract the D-term

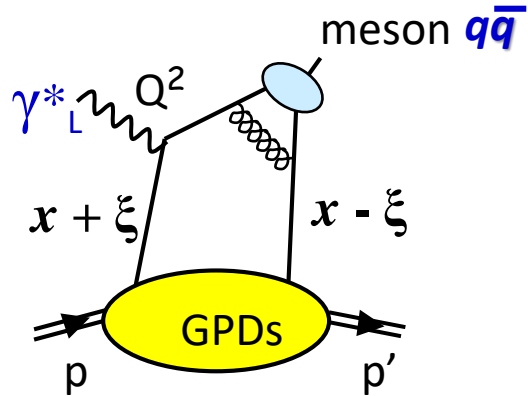
COMPASS 2 years of data $E\mu = 160 \text{ GeV}$ $1 < Q^2 < 8 \text{ GeV}^2$

HERMES

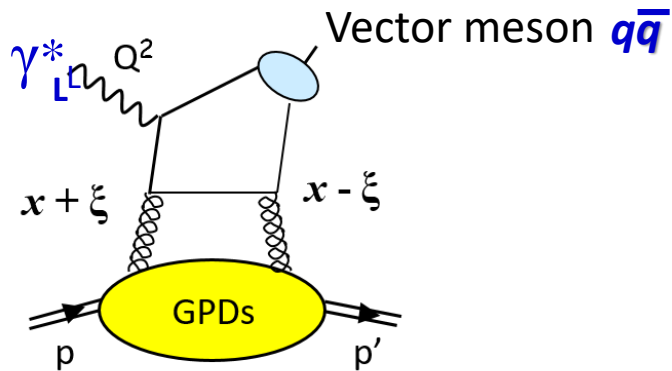
JLab

GPDs and Hard Exclusive Meson Production

Quark contribution



Gluon contribution at the same order in α_s



The meson wave function
Is an additional non-perturbative term

4 chiral-even GPDs: helicity of parton unchanged

$H^q(x, \xi, t)$	$E^q(x, \xi, t)$	For Vector Meson
$\tilde{H}^q(x, \xi, t)$	$\tilde{E}^q(x, \xi, t)$	For Pseudo-Scalar Meson

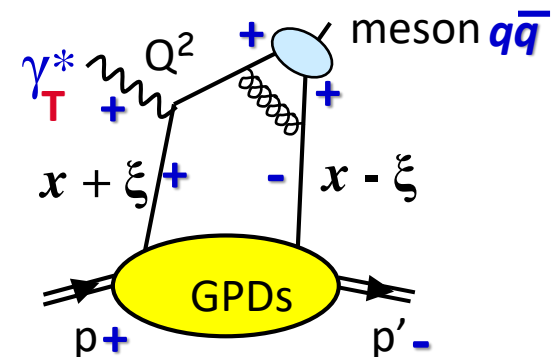
+ 4 chiral-odd or transversity GPDs: helicity of parton changed
(not possible in DVCS)

$H_T^q(x, \xi, t)$	$E_T^q(x, \xi, t)$	$\bar{E}_T^q = 2 \tilde{H}_T^q + E_T^q$
$\tilde{H}_T^q(x, \xi, t)$	$\tilde{E}_T^q(x, \xi, t)$	

Factorisation proven only for σ_L

σ_T is asymptotically suppressed by $1/Q^2$ but large contribution observed
model of σ_T with transversity GPDs - divergencies regularized by k_T of q
and \bar{q} and Sudakov suppression factor

$\mathcal{M}_{0-, ++}$ sensitive to H_T^q
and to a twist-3 meson wave function



$e p \rightarrow e \pi^0 p$

$$\frac{d^2\sigma}{dt d\phi_\pi} = \frac{1}{2\pi} \left[\left(\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} \right) + \epsilon \cos 2\phi_\pi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_\pi \frac{d\sigma_{LT}}{dt} \right]$$

$$\frac{d\sigma_L}{dt} = \frac{4\pi\alpha}{k'} \frac{1}{Q^6} \left\{ (1 - \xi^2) |\langle \tilde{H} \rangle|^2 - 2\xi^2 \text{Re} [\langle \tilde{H} \rangle^* \langle \tilde{E} \rangle] - \frac{t'}{4m^2} \xi^2 |\langle \tilde{E} \rangle|^2 \right\}$$

Leading twist should be dominant
but \approx only a few % of $\frac{d\sigma_T}{dt}$

The other contributions arise from coupling between chiral-odd (quark helicity flip) GPDs to the twist-3 pion amplitude

$$\frac{d\sigma_T}{dt} = \frac{4\pi\alpha}{2k'} \frac{\mu_\pi^2}{Q^8} \left[(1 - \xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8m^2} |\langle \bar{E}_T \rangle|^2 \right]$$

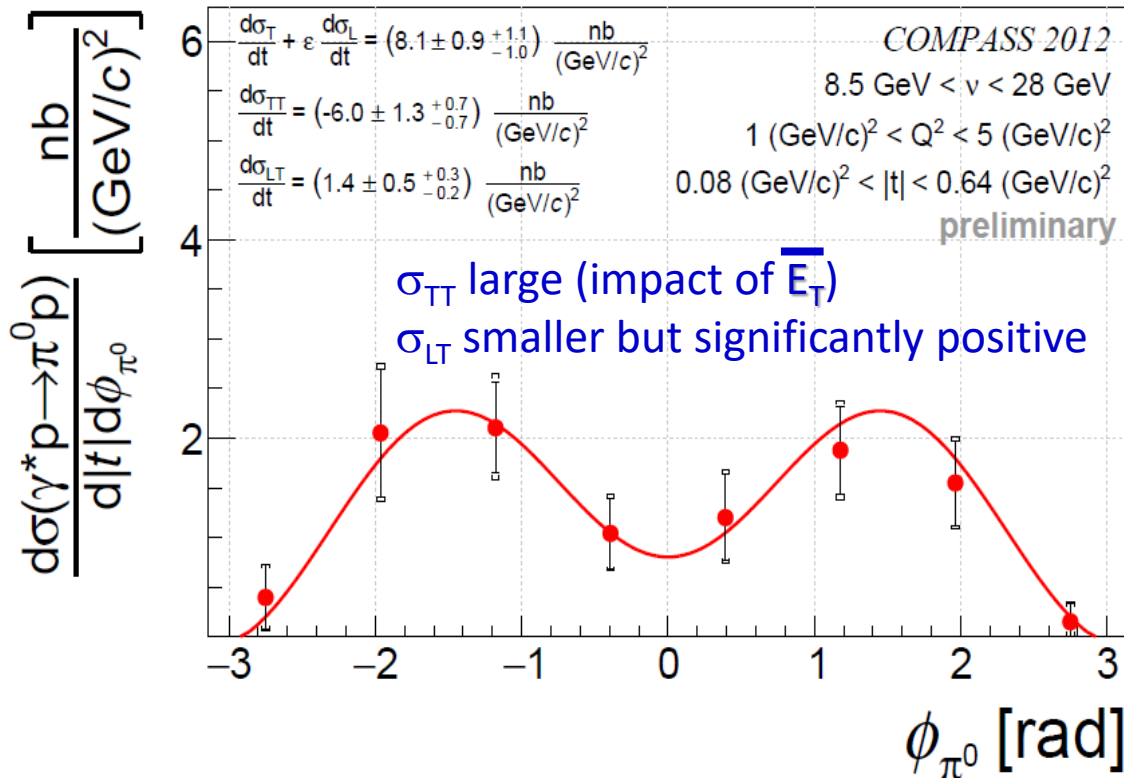
$$\frac{d\sigma_{LT}}{dt} = \frac{4\pi\alpha}{\sqrt{2}k'} \frac{\mu_\pi}{Q^7} \xi \sqrt{1 - \xi^2} \frac{\sqrt{-t'}}{2m} \text{Re} [\langle H_T \rangle^* \langle \tilde{E} \rangle]$$

$$\frac{d\sigma_{TT}}{dt} = \frac{4\pi\alpha}{k'} \frac{\mu_\pi^2}{Q^8} \frac{t'}{16m^2} |\langle \bar{E}_T \rangle|^2$$

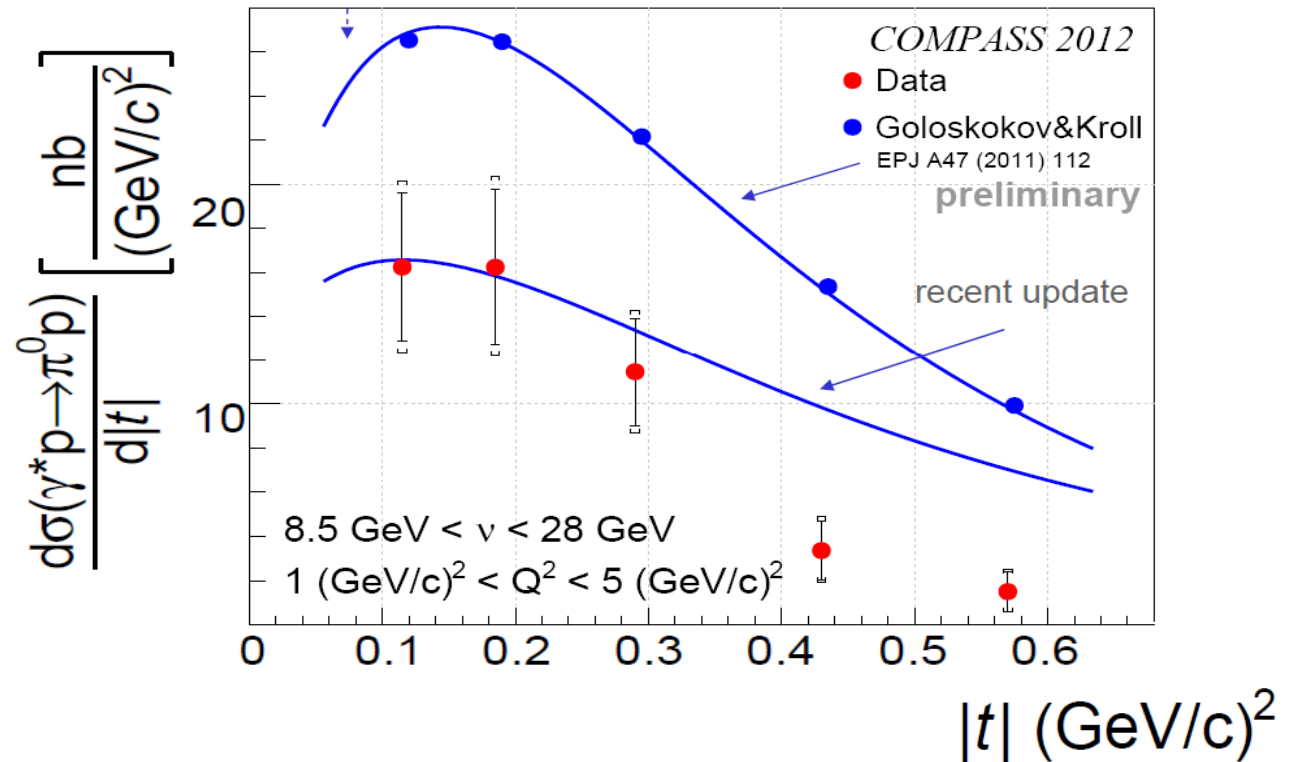
A large impact of \bar{E}_T should be clearly visible in σ_{TT} and in the dip at small $|t|$ of σ_T



$$\frac{d^2\sigma}{dt d\phi_\pi} = \frac{1}{2\pi} \left[\left(\epsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} \right) + \epsilon \cos 2\phi_\pi \frac{d\sigma_{TT}}{dt} + \sqrt{2\epsilon(1+\epsilon)} \cos \phi_\pi \frac{d\sigma_{LT}}{dt} \right]$$



A dip at small t would indicate a large impact of \overline{E}_T



SCHC ($\lambda_\gamma = \lambda_V$)

SCHC implies:

- $r_{1-1}^1 + \text{Im} r_{1-1}^2 = 0$

$= -0.010 \pm 0.032 \pm 0.047$ OK

- $\text{Re} r_{10}^5 + \text{Im} r_{10}^6 = 0$

$= 0.014 \pm 0.011 \pm 0.013$ OK

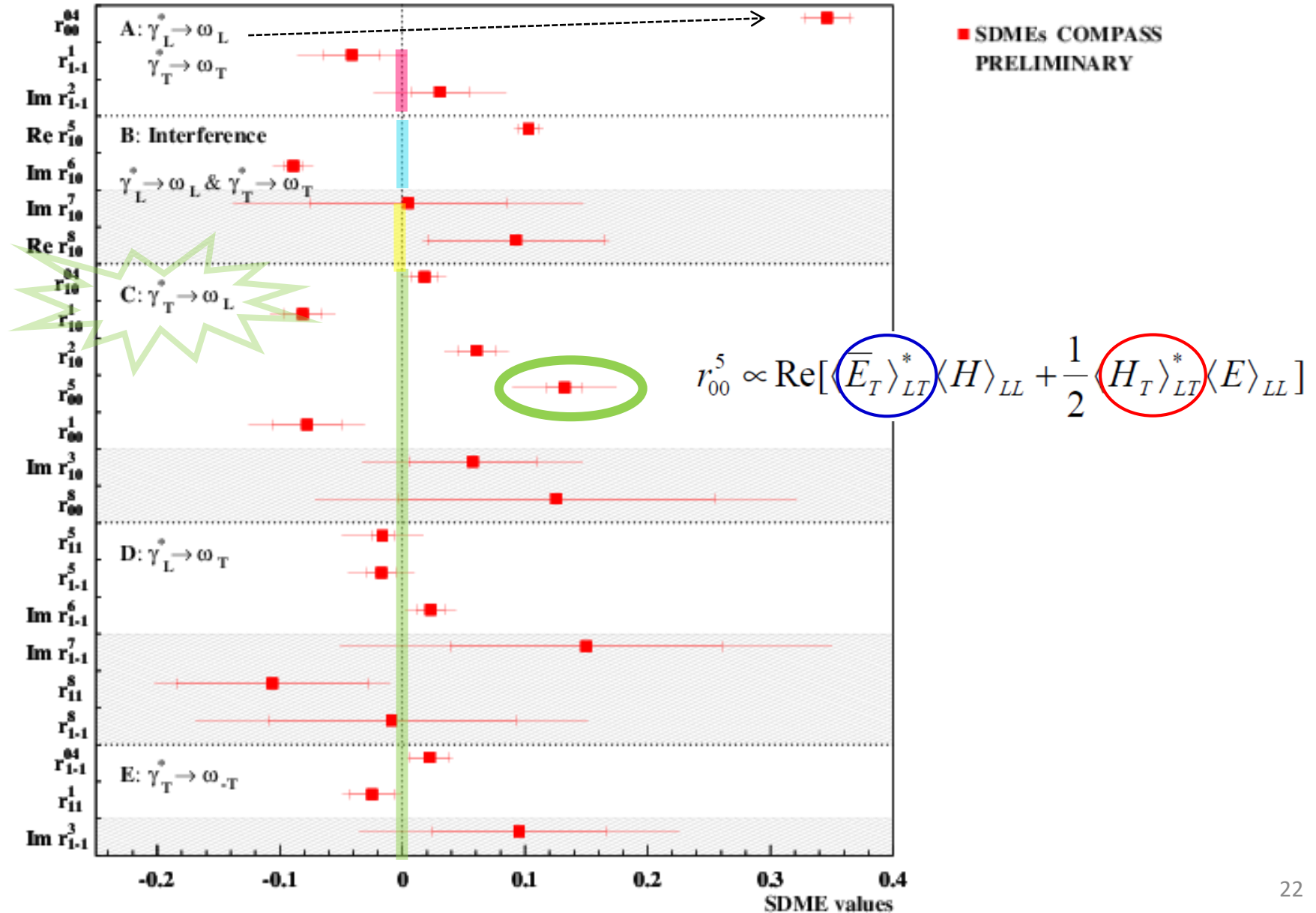
- $\text{Im} r_{10}^7 - \text{Re} r_{10}^8 = 0$

$= -0.088 \pm 0.110 \pm 0.196$ OK

- all elements of classes C, D, E should be 0

for $\gamma_L^* \rightarrow \omega_T$ and $\gamma_T^* \rightarrow \omega_T$ OK within errors

not obeyed for transitions $\gamma_T^* \rightarrow \omega_L$



Conclusions

From 2016-17 data

sum and difference of DVCS x-sections with polarized μ^+ and μ^-

- transverse extension of partons as a function of x_{Bj}
- $\text{Im}\mathcal{H}(\xi,t)$ and $\text{Re}\mathcal{H}(\xi,t)$ for D-term and pressure distribution

HEMP $\pi^0, \rho, \omega, \phi, J/\psi$ → universality of GPDs - transverse GPDs - flavor decomposition



COMPASS++/AMBER starting in 2022

Letter of Intent Draft 1.0: <https://arXiv.org/abs/1808.00848>

New collaborators are welcome: <https://nqf-m2.web.cern.ch>



A new QCD facility
at the M2 beam line of the CERN SPS

Letter of Intent - Draft 1.0: <https://arXiv.org/abs/1808.00848>

COMPASS++/AMBER starting in 2022

Program	Physics Goals	Beam Energy [GeV]	Beam Intensity [s^{-1}]	Trigger Rate [kHz]	Beam Type	Target	Earliest start time, duration	Hardware Additions
μp elastic scattering	Precision proton-radius measurement	100	$4 \cdot 10^6$	100	μ^\pm	high-pressure H2	2022 1 year	active TPC, SciFi trigger, silicon veto,
Hard exclusive reactions	GPD E	160	$2 \cdot 10^7$	10	μ^\pm	NH_3^\dagger	2022 2 years	recoil silicon, modified PT magnet
Input for Dark Matter Search	\bar{p} production cross section	20-280	$5 \cdot 10^5$	25	p	LH2, LHe	2022 1 month	LHe target
\bar{p} -induced Spectroscopy	Heavy quark exotics	12, 20	$5 \cdot 10^7$	25	\bar{p}	LH2	2022 2 years	target spectr.: tracking, calorimetry
Drell-Yan	Pion PDFs	190	$7 \cdot 10^7$	25	π^\pm	C/W	2022 1-2 years	
Drell-Yan (RF)	Kaon PDFs & Nucleon TMDs	~ 100	10^8	25-50	K^\pm, \bar{p}	NH_3^\dagger , C/W	2026 2-3 years	"active absorber", vertex det.
Primakoff (RF)	Kaon polarisability & pion life time	~ 100	$5 \cdot 10^6$	> 10	K^-	Ni	non-exclusive 2026 1 year	
Prompt Photons (RF)	Meson gluon PDFs	≥ 100	$5 \cdot 10^6$	10-100	K^\pm π^\pm	LH2, Ni	non-exclusive 2026 1-2 years	hodoscope
K -induced Spectroscopy (RF)	High-precision strange-meson spectrum	50-100	$5 \cdot 10^6$	25	K^-	LH2	2026 1 year	recoil TOF, forward PID
Vector mesons (RF)	Spin Density Matrix Elements	50-100	$5 \cdot 10^6$	10-100	K^\pm, π^\pm	from H to Pb	2026 1 year	

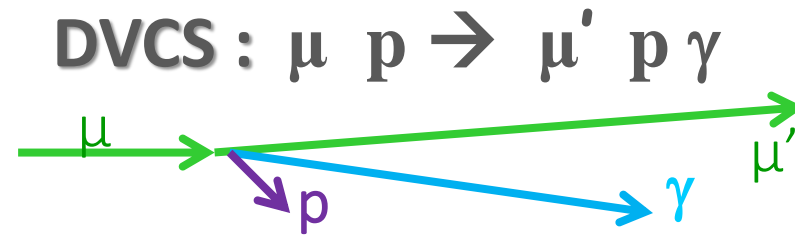
Beam line unique with polarised μ^+ and μ^- and high intensity pion beam

Possible RF separated beam for high intensity antiproton and K beams

Versatile apparatus (Upgrade ++)

Proton Radius
Meson PDF – gluon PDF
Proton spin structure
3D imaging (TMDs and GPDs)
Hadron spectroscopy
Anti-matter cross section

The DVCS experiment at COMPASS



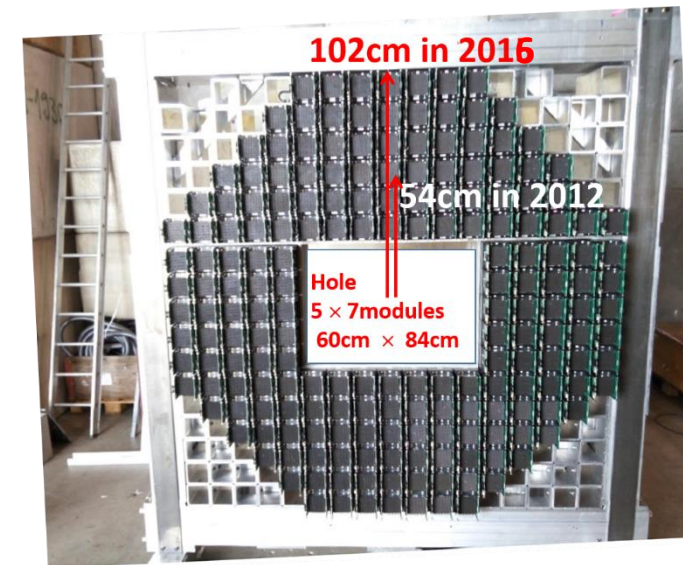
New equipments:

- 2.5m LH2 target
- 4m ToF Barrel CAMERA
- ECALO



CAMERA
L=4m
Ø=2m

24 inner & outer scintillators separated by about 1m
1 GHz SADC readout, 330ps ToF resolution



ECALO: 2 × 2 m²

Shashlyk modules + MAPD readout
one module is made of 9 cells (4×4 cm²)
= 194 modules or 1746 cells

COMPASS 2012 Transverse extension of partons in the sea quark range

$$d\sigma^{DVCS}/dt = e^{-B'} |t| = c_0^{DVCS}$$

$$c_{0,\text{unp}}^{DVCS} = 2(2 - 2y + y^2) C_{\text{unp}}^{DVCS}(\mathcal{F}, \mathcal{F}^*)$$

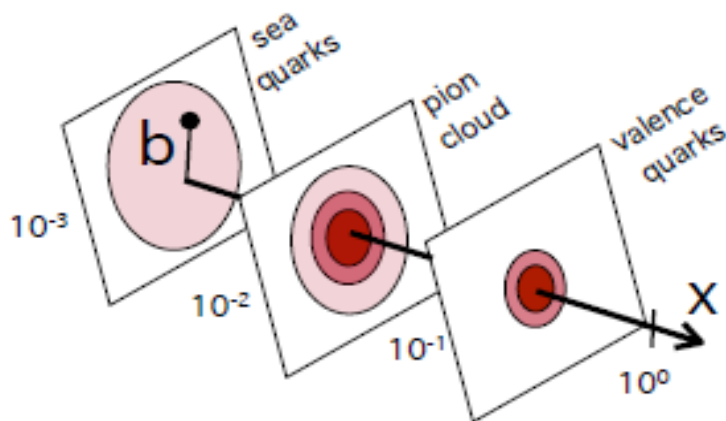
$$C_{\text{unp}}^{DVCS}(\mathcal{F}, \mathcal{F}^*) = \frac{1}{(2 - x_B)^2} \left\{ 4(1 - x_B) (\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) - x_B^2 (\mathcal{H}\mathcal{E}^* + \mathcal{E}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{E}}^* + \tilde{\mathcal{E}}\tilde{\mathcal{H}}^*) - \left(x_B^2 + (2 - x_B)^2 \frac{\Delta^2}{4M^2} \right) \mathcal{E}\mathcal{E}^* - x_B^2 \frac{\Delta^2}{4M^2} \tilde{\mathcal{E}}\tilde{\mathcal{E}}^* \right\}$$

At COMPASS:

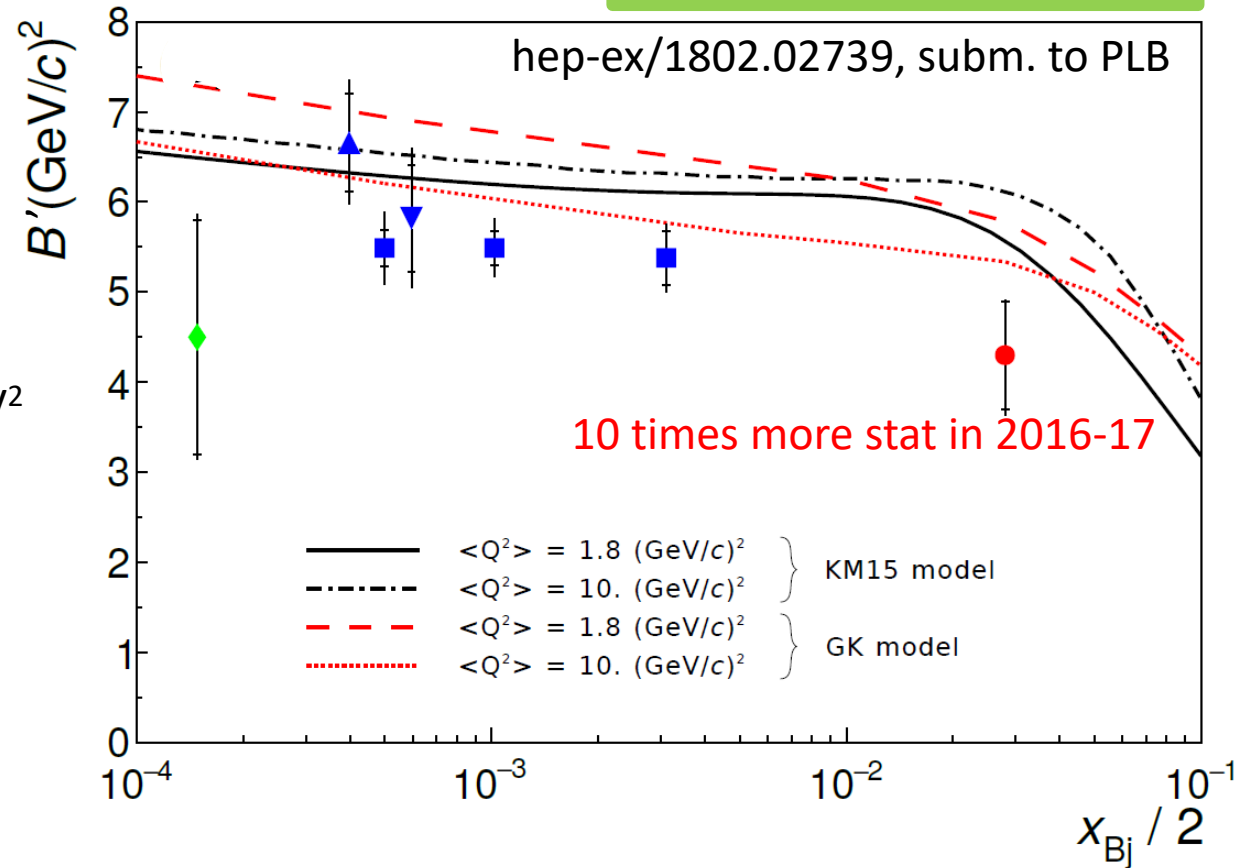
$\langle x_{Bj} \rangle = 0.056$; $\langle Q^2 \rangle = 1.8 \text{ GeV}^2$; t varies from 0.08 to 0.64 GeV^2

Due to the small value of x_{Bj} and t it remains only:

$$c_0^{DVCS} \propto 4(\mathcal{H}\mathcal{H}^* + \tilde{\mathcal{H}}\tilde{\mathcal{H}}^*) - \frac{t}{M^2} \mathcal{E}\mathcal{E}^*$$

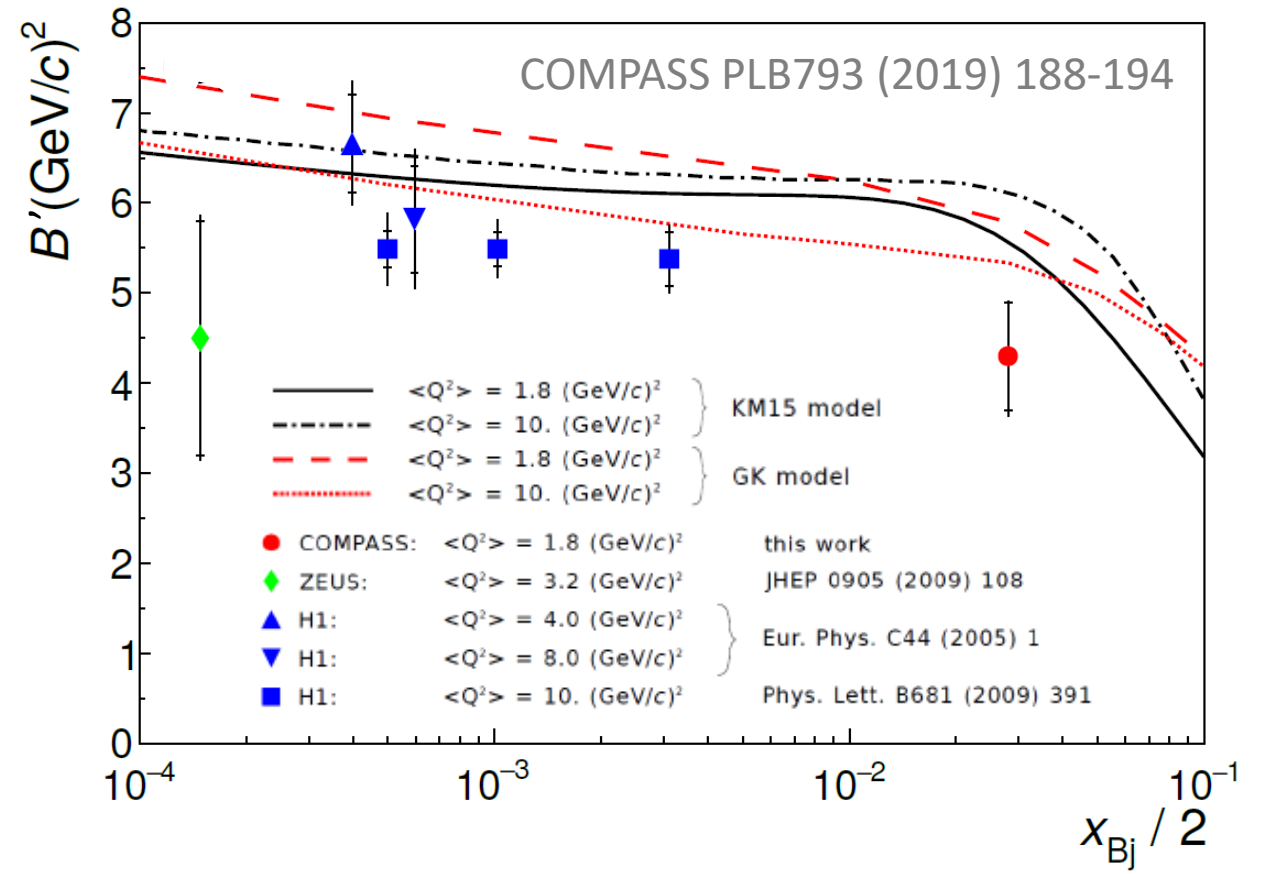


$$\langle r_{\perp}^2(x_B) \rangle \approx 2B'(x_B)$$



$$B = (4.31 \pm 0.62_{\text{stat}} \pm 0.09_{\text{sys}}) (\text{GeV}/c)^{-2}$$

$$\sqrt{\langle r_{\perp}^2 \rangle} = (0.58 \pm 0.04_{\text{stat}} \pm 0.01_{\text{sys}}) \text{ fm}$$



Dominance of $Im\mathcal{H}$ (with respect of $Re\mathcal{H}$ and other CTF) at small x_B

Figure from Kumericki, Mueller

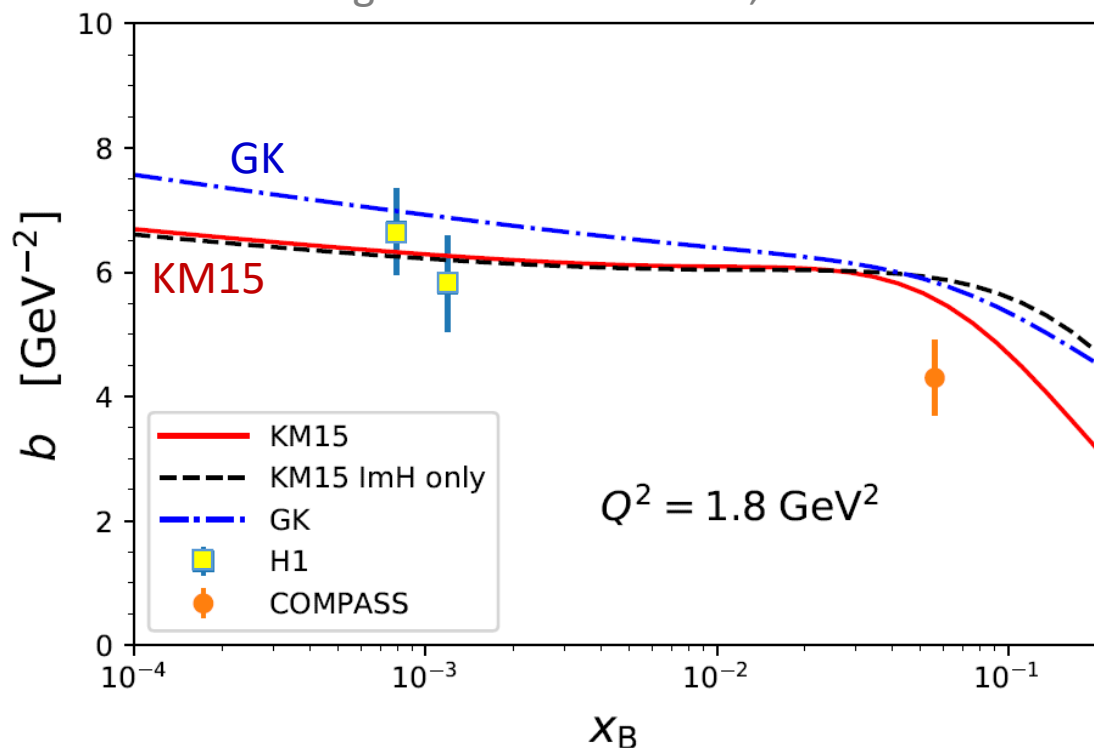
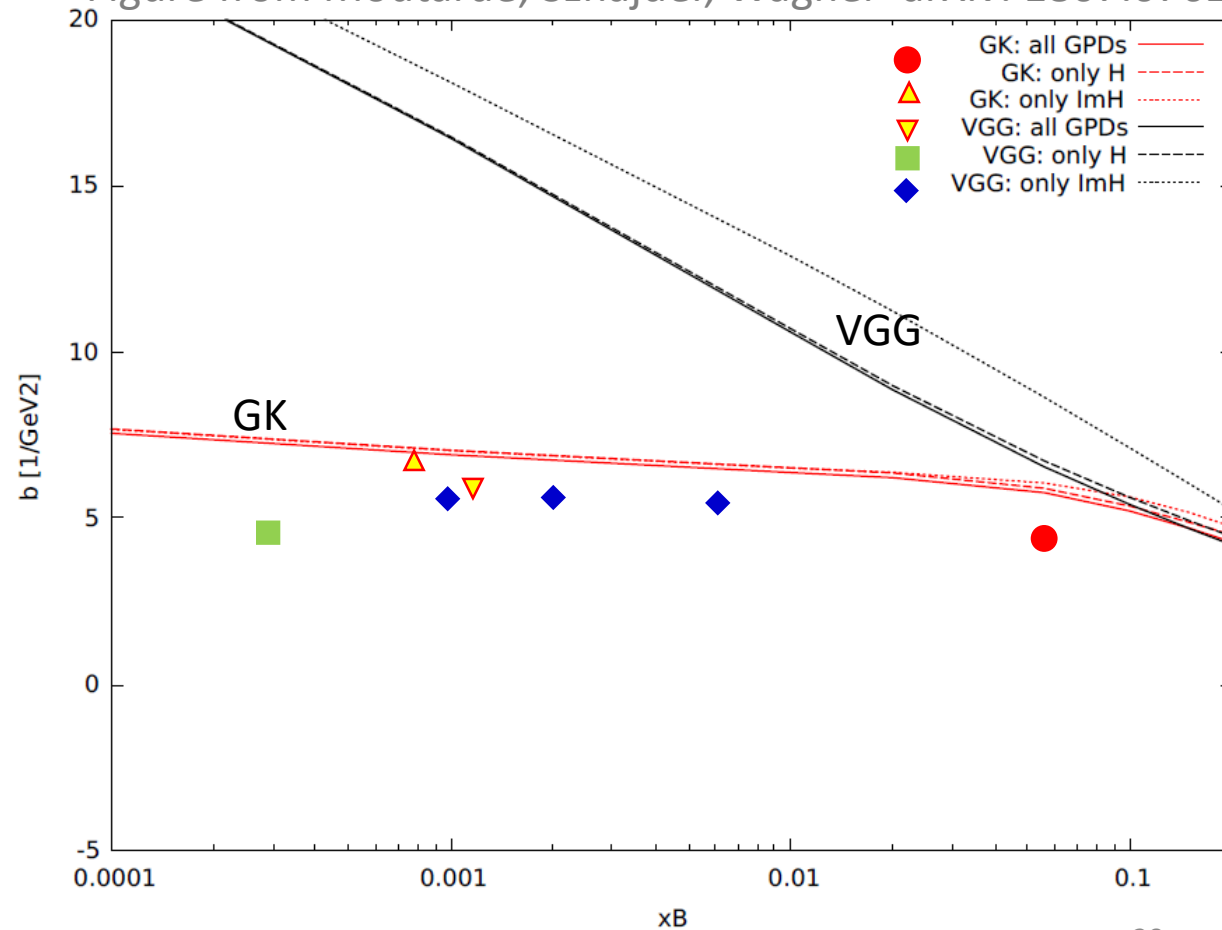


Figure from Moutarde, Sznajder, Wagner arXiv: 1807.07620



The past and future DVCS experiments

