

Recent COMPASS results on Transverse Spin Asymmetries in SIDIS

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On behalf of the COMPASS Collaboration



*CO*mmon
Muon and
Proton
Apparatus for
Structure and
Spectroscopy

fixed target experiment at the CERN SPS



SPS

LHC



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fixed target experiment at the CERN SPS



physics programme:

hadron spectroscopy (p , π , K)★

- light mesons, glue-balls, exotic mesons
- polarisability of pion and kaon

nucleon structure (μ)

- longitudinal spin structure
- transverse momentum and transverse spin structure



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- longitudinal spin structure
- **transverse momentum and transverse spin structure**

this talk

COMPASS spectrometer



designed to

- use high energy beams
- have large angular acceptance
- cover a broad kinematical range

COMPASS spectrometer



designed to

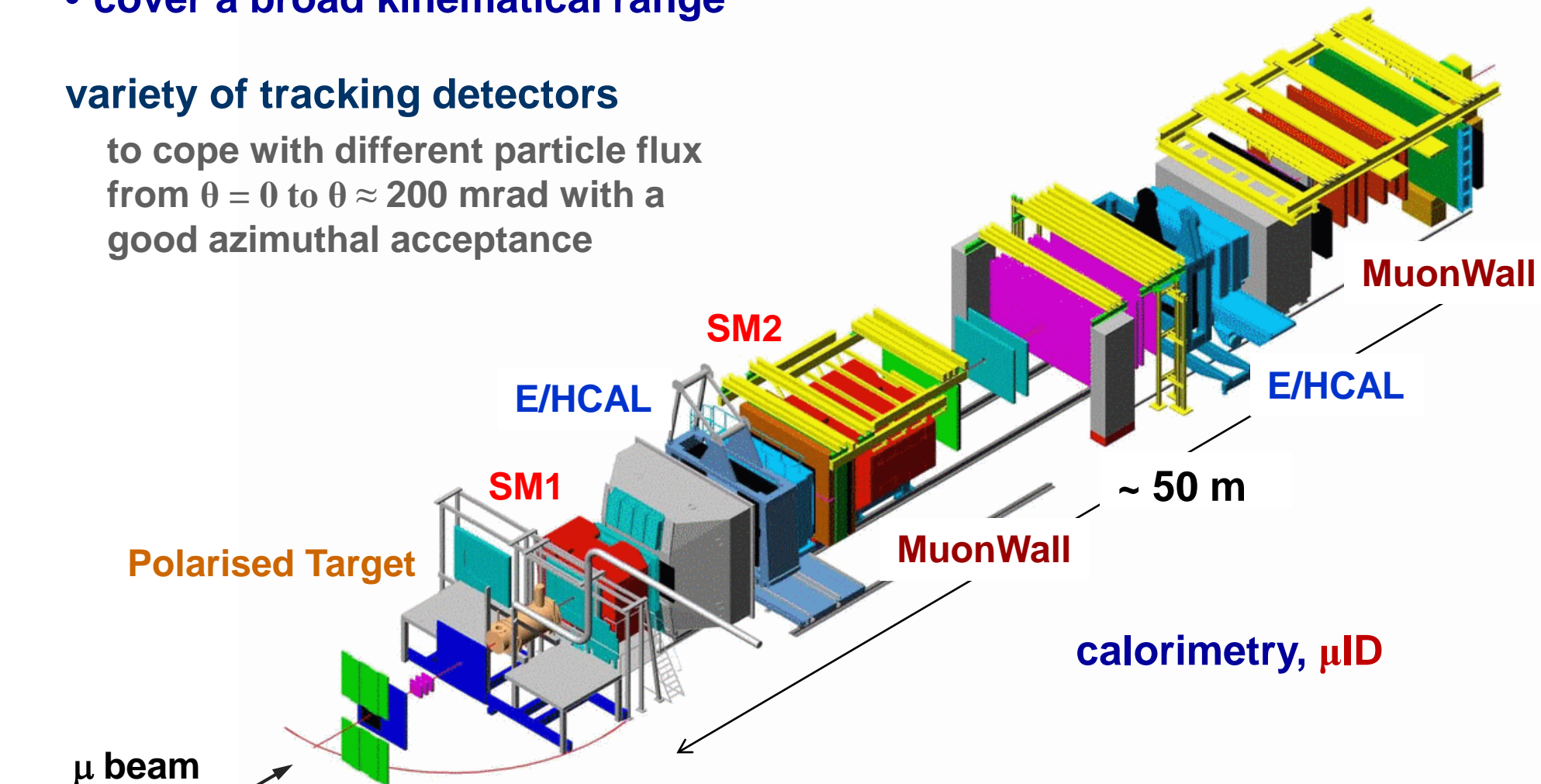
- use high energy beams
- have large angular acceptance
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two stages spectrometer

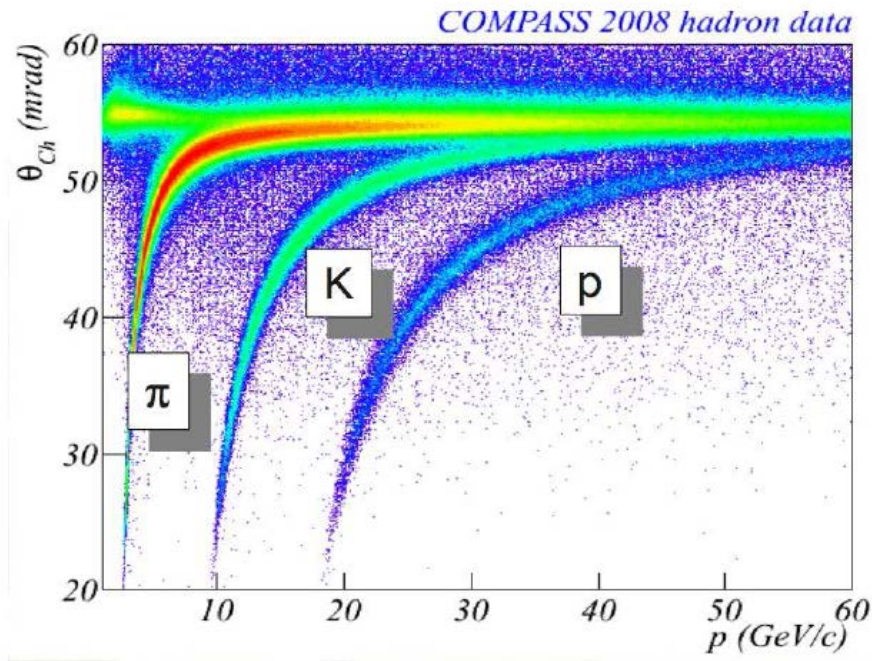
- Large Angle Spectrometer (**SM1**)
- Small Angle Spectrometer (**SM2**)

variety of tracking detectors

to cope with different particle flux from $\theta = 0$ to $\theta \approx 200$ mrad with a good azimuthal acceptance

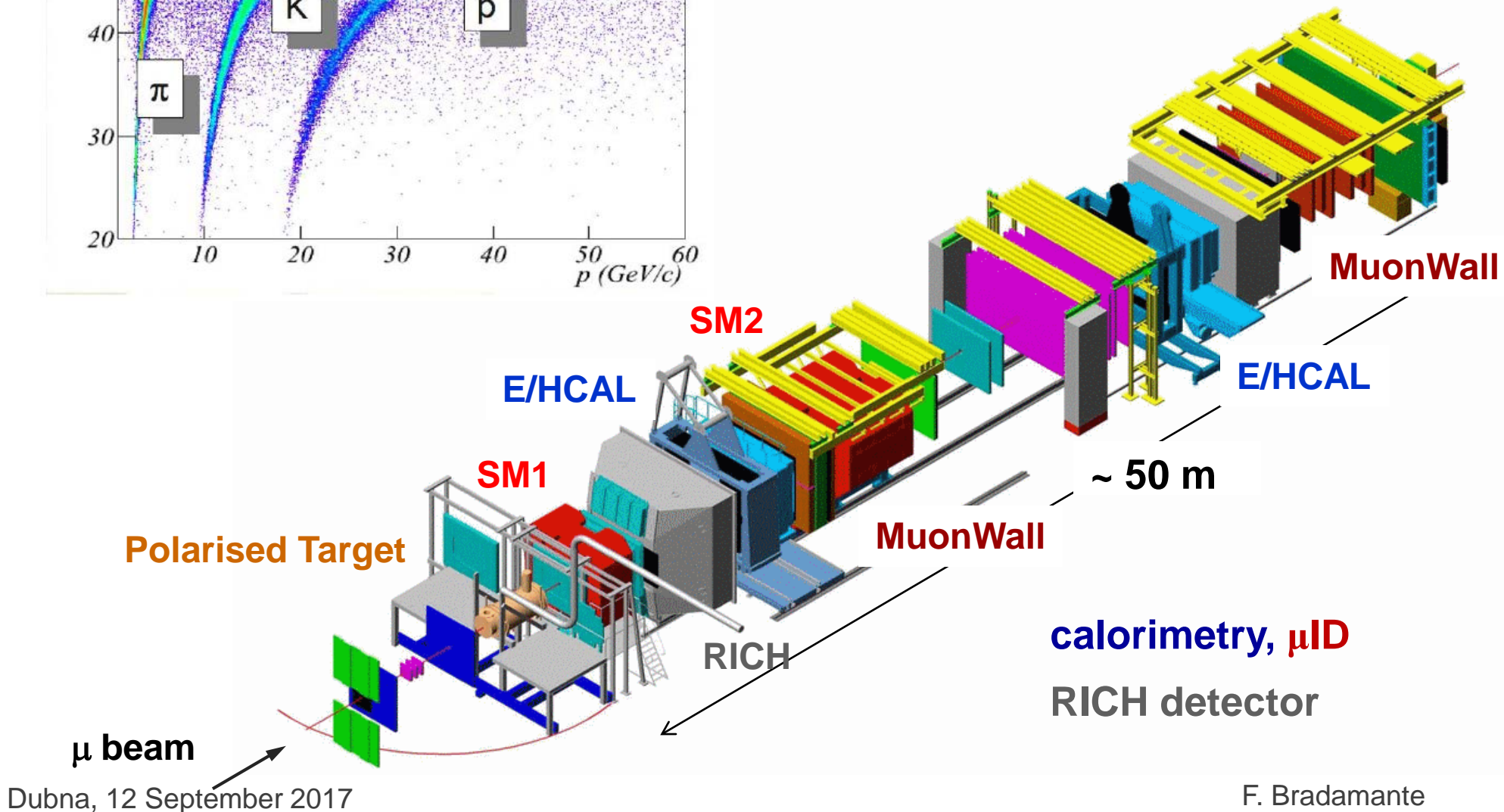


COMPASS spectrometer



two stages spectrometer

- Large Angle Spectrometer (**SM1**)
- Small Angle Spectrometer (**SM2**)

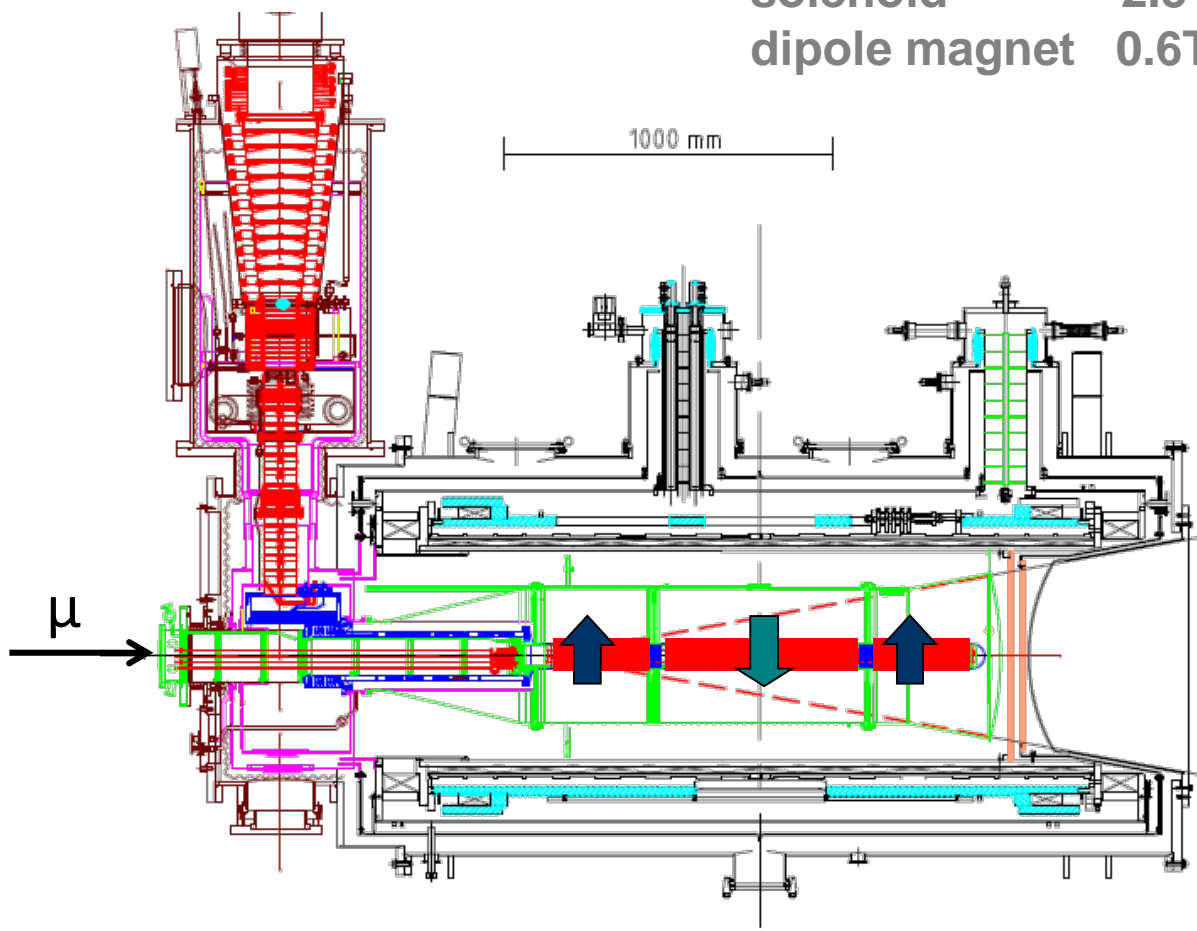




the polarized target system (>2005)

$^3\text{He} - ^4\text{He}$ dilution refrigerator ($T \sim 50\text{mK}$)

solenoid 2.5T
dipole magnet 0.6T



acceptance $> \pm 180$ mrad

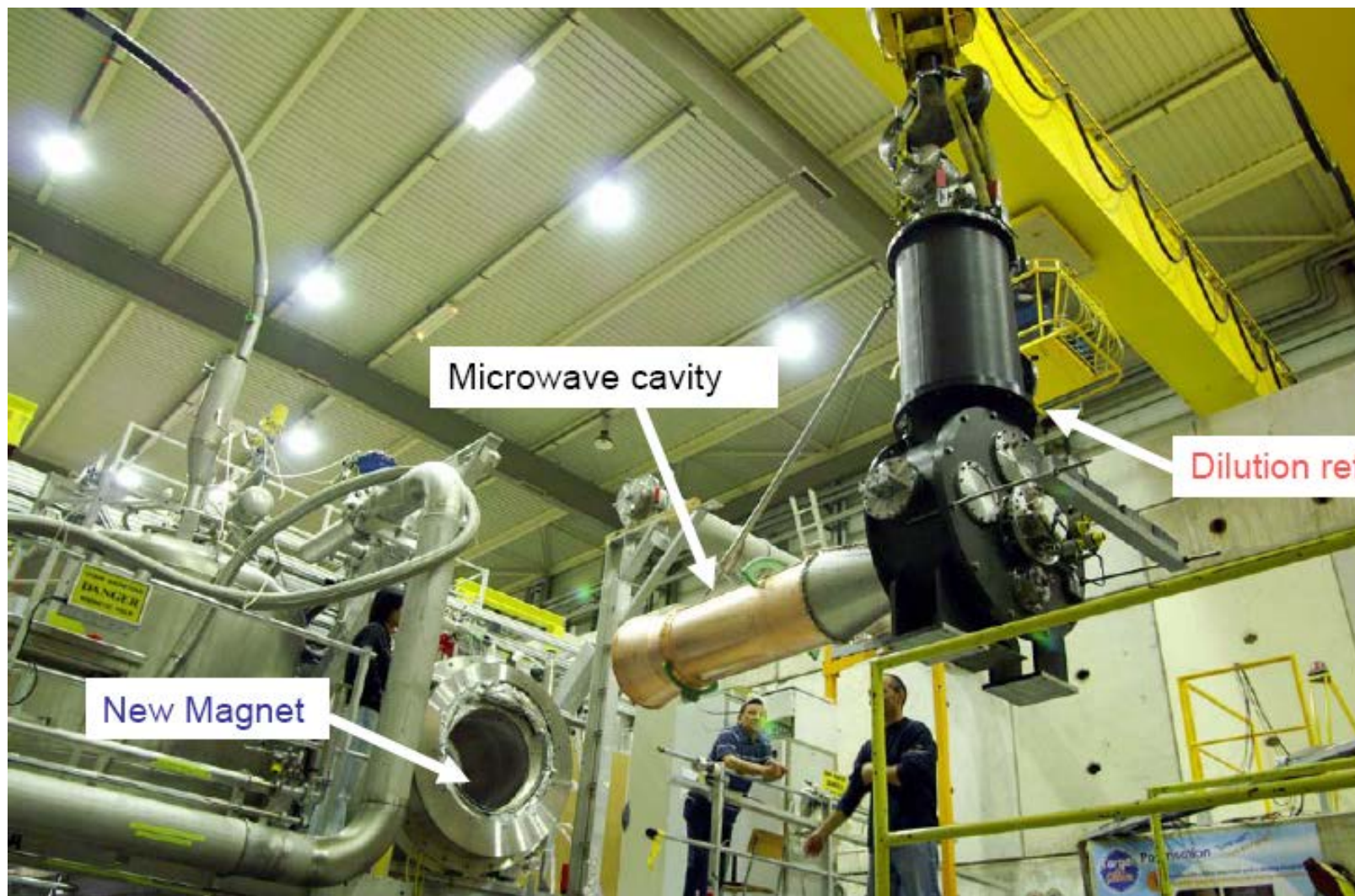
3 target cells
30, 60, and 30 cm long

opposite polarisation

	d (^6LiD)	p (NH_3)
polarization	50%	90%
dilution factor	40%	16%

no evidence for relevant nuclear effects (160 GeV)

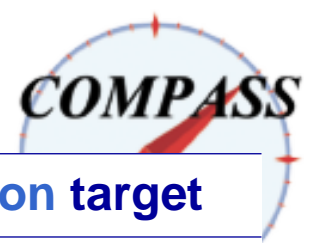
the polarized target system



COMPASS data taking



COMPASS data taking


















2002	nucleon structure with	160 GeV μ	L&T	polarised deuteron target
2003	nucleon structure with	160 GeV μ	L&T	polarised deuteron target
2004	nucleon structure with	160 GeV μ	L&T	polarised deuteron target
2005	<i>CERN accelerators shut down</i>			
2006	nucleon structure with	160 GeV μ	L	polarised deuteron target
2007	nucleon structure with	160 GeV μ	L&T	polarised proton target
2008	<i>hadron spectroscopy</i>			
2009	<i>hadron spectroscopy</i>			
2010	nucleon structure with	160 GeV μ	T	polarised proton target
2011	nucleon structure with	190 GeV μ	L	polarised proton target
2012	Primakoff & DVCS / SIDIS test			
2013	<i>CERN accelerators shut down</i>			
2014	Test beam Drell-Yan process with π beam and T polarised proton target			
2015	Drell-Yan process with π beam and T polarised proton target			
2016	DVCS / SIDIS with μ beam and unpolarised proton target			
2017	DVCS / SIDIS with μ beam and unpolarised proton target			
2018	Drell-Yan process with π beam and T polarised proton target			

MUON beam PROGRAM:

**TRANSVERSITY and
Transverse Momentum Dependent PDFs**

the structure of the nucleon

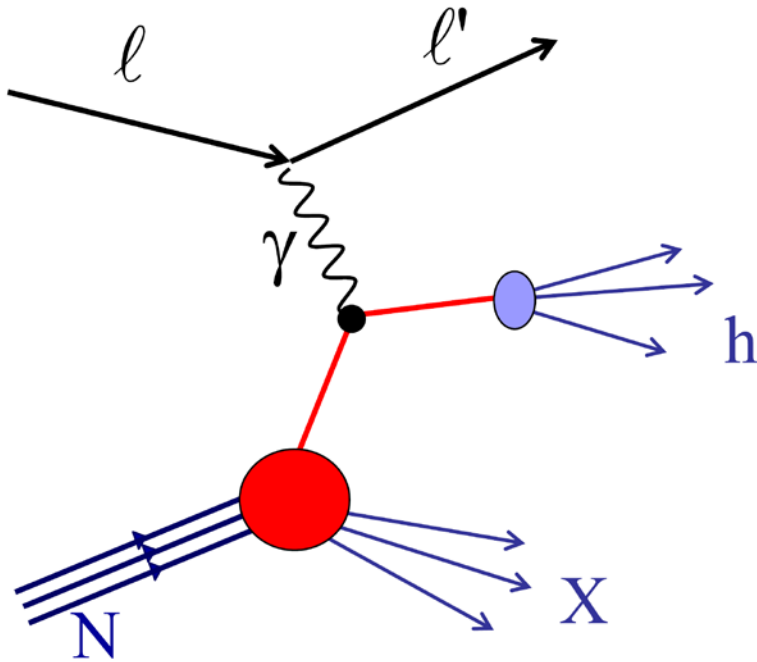
taking into account the quark intrinsic transverse momentum k_T ,
 at leading order other 6 TMD PDFs are needed for a full description
 of the nucleon structure

		nucleon polarisation			
		U	L	T	
quark polarisation	U	f_1  number density \mathbf{q}		f_{1T}^\perp  -  Sivers	$\Delta_0^T \mathbf{q}$
	L		g_1  -  helicity $\Delta \mathbf{q}$	g_{1T}  - 	
	T	h_1^\perp  -  Boer Mulders	h_{1L}^\perp  - 	h_1  -  transversity h_{1T}^\perp  - 	$\Delta_T \mathbf{q}$

SIDIS gives access to all of them

Semi-Inclusive Deep Inelastic Scattering

hard interaction of a lepton with a nucleon via virtual photon exchange

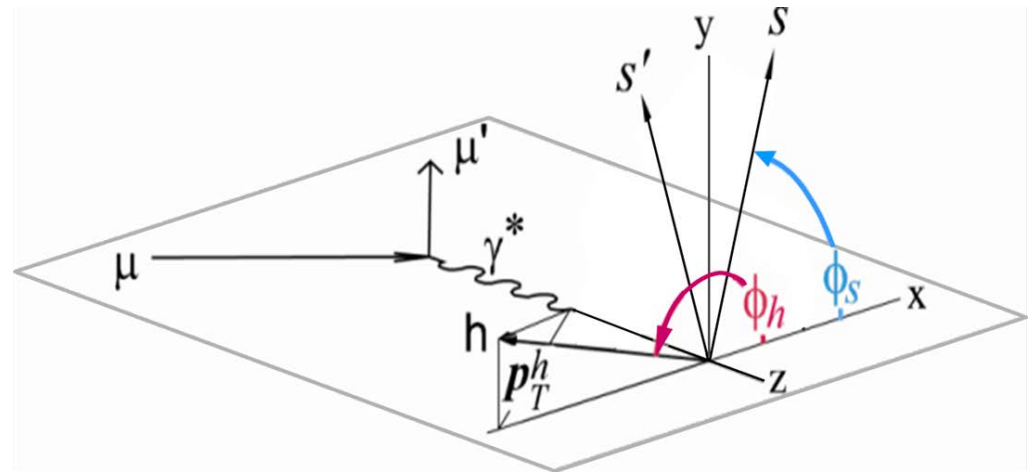


$$x = \frac{Q^2}{2P \cdot q} \quad y = \frac{P \cdot q}{P \cdot \ell} =_{LAB} \frac{E - E'}{E}$$

$$Q^2 = -q^2 \quad W^2 = (P + q)^2$$

$$z = \frac{P \cdot P_h}{P \cdot q} =_{LAB} \frac{E_h}{E - E'}$$

$$\sigma^{lN \rightarrow lhX} \propto \sum_q f(x) \otimes \sigma^{lq \rightarrow lq} \otimes D_q^h(z)$$



Semi-Inclusive Deep Inelastic Scattering

$$\begin{aligned}
 \frac{d\sigma}{dx dy dz d\psi d\phi_h dP_{h\perp}^2} = & \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\
 & + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \\
 & + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \\
 & + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
 & + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\
 & + \left. \left. \left. \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right) \right] \right. \\
 & + |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right. \\
 & \left. \left. \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right) \right] \right\},
 \end{aligned}$$

unpol target

→ pol target

↑ pol target

18 structure functions

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 \end{aligned}$$

14 independent azimuthal modulations

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14 independent azimuthal modulations

amplitudes of the modulations
→ TMD PDFs

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 \end{aligned}$$

14 independent azimuthal modulations

amplitudes of the modulations
→ TMD PDFs

SIDIS

- allows to disentangle the effects related to the different TMD PDFs and to access all of them
- by identifying the final state hadrons and using different targets allows for flavour separation
→ *very powerful tool*

all the amplitudes (AA) have been measured in COMPASS

TMDs in unpolarised SIDIS

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ \begin{aligned} & k_T F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \\ & + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \\ & h_i^\perp H_i^\perp \\ & + \dots \end{aligned} \right\}$$

unpolarised SIDIS

Relevance for TMDs:

- the cross-section **dependence on p_{Th}** comes from:
 - intrinsic k_T of the quarks
 - p_{\perp} generated in the quark fragmentation

$$\langle p_{Th}^2 \rangle = \langle p_{\perp}^2 \rangle + z^2 \langle k_T^2 \rangle$$

- the **azimuthal modulations** in the unpolarized cross-sections comes from:
 - intrinsic k_T of the quarks
 - Boer-Mulders PDF

combined analysis should allow to disentangle the different effects

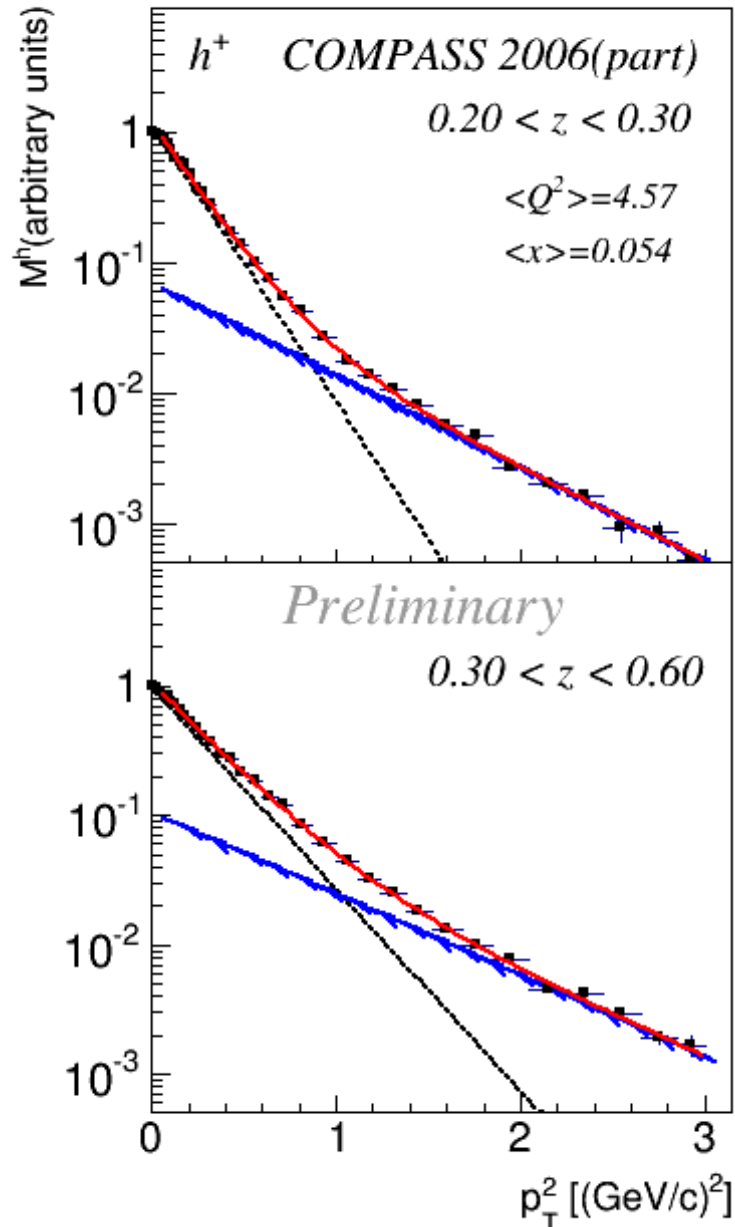
COMPASS

- has produced results on ${}^6\text{LiD}$ ($\sim d$) from 2004/6 data
- will measure SIDIS on LH_2 in parallel with DVCS

unpolarised SIDIS – p_{Th} distributions



deuteron



Fit distributions with

- 1 exponential for $p_{Th}^2 \in [0.05, 0.68]$
- 2 exponentials for $p_{Th}^2 \in [0.05, 3]$

↓
needed to describe the shape
of p_{Th}^2 the COMPASS data

Transversity 2014

unpolarised SIDIS – p_{Th} distributions



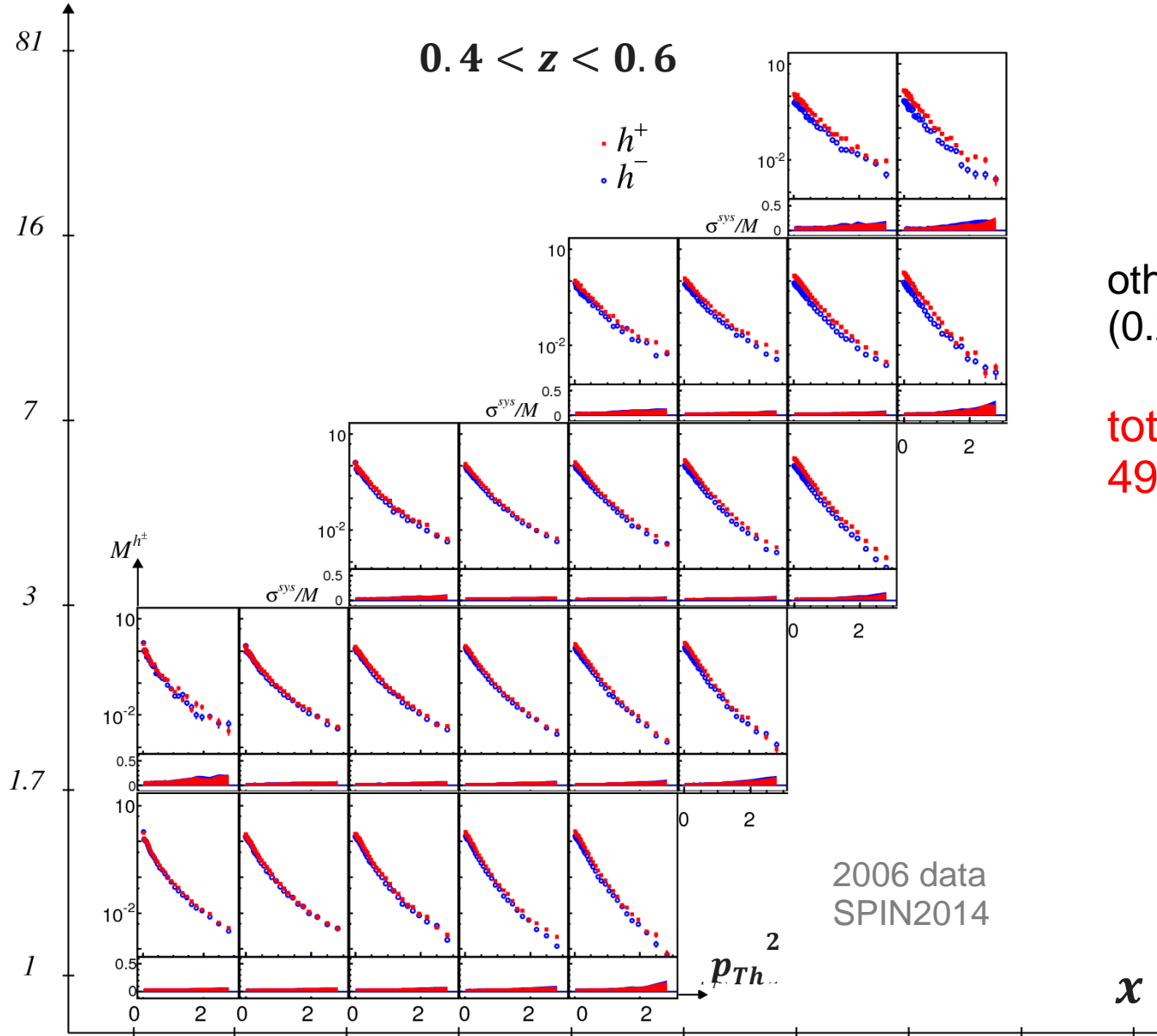
Q^2 [GeV/c]²

COMPASS Preliminary

deuteron

$0.4 < z < 0.6$

h^+
 h^-



other 3 z bins
(0.2-0.3, 0.3-0.4, 0.6-0.8)

total:
4918 data points

paper ready,
to be sent to PRD

2006 data
SPIN2014

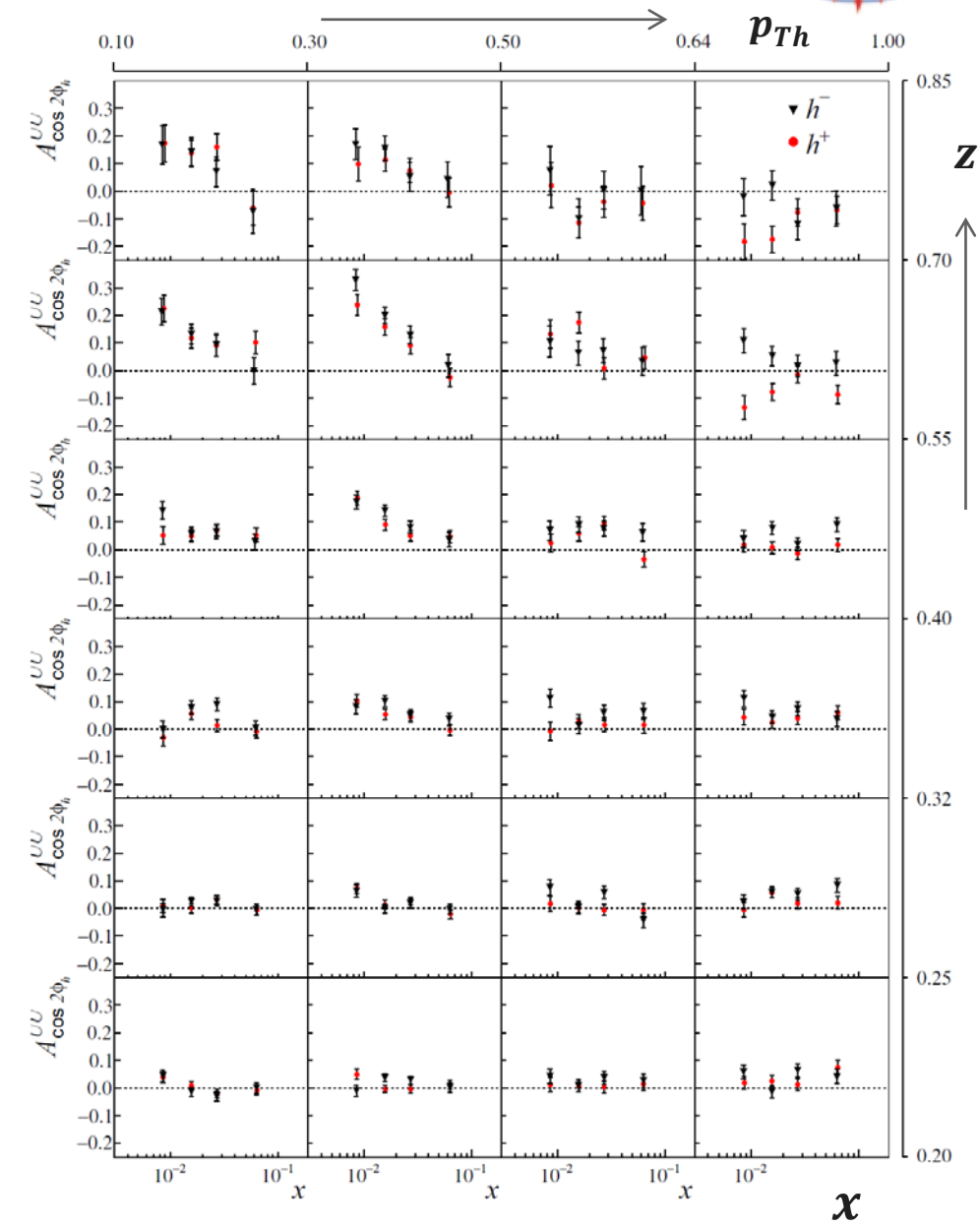
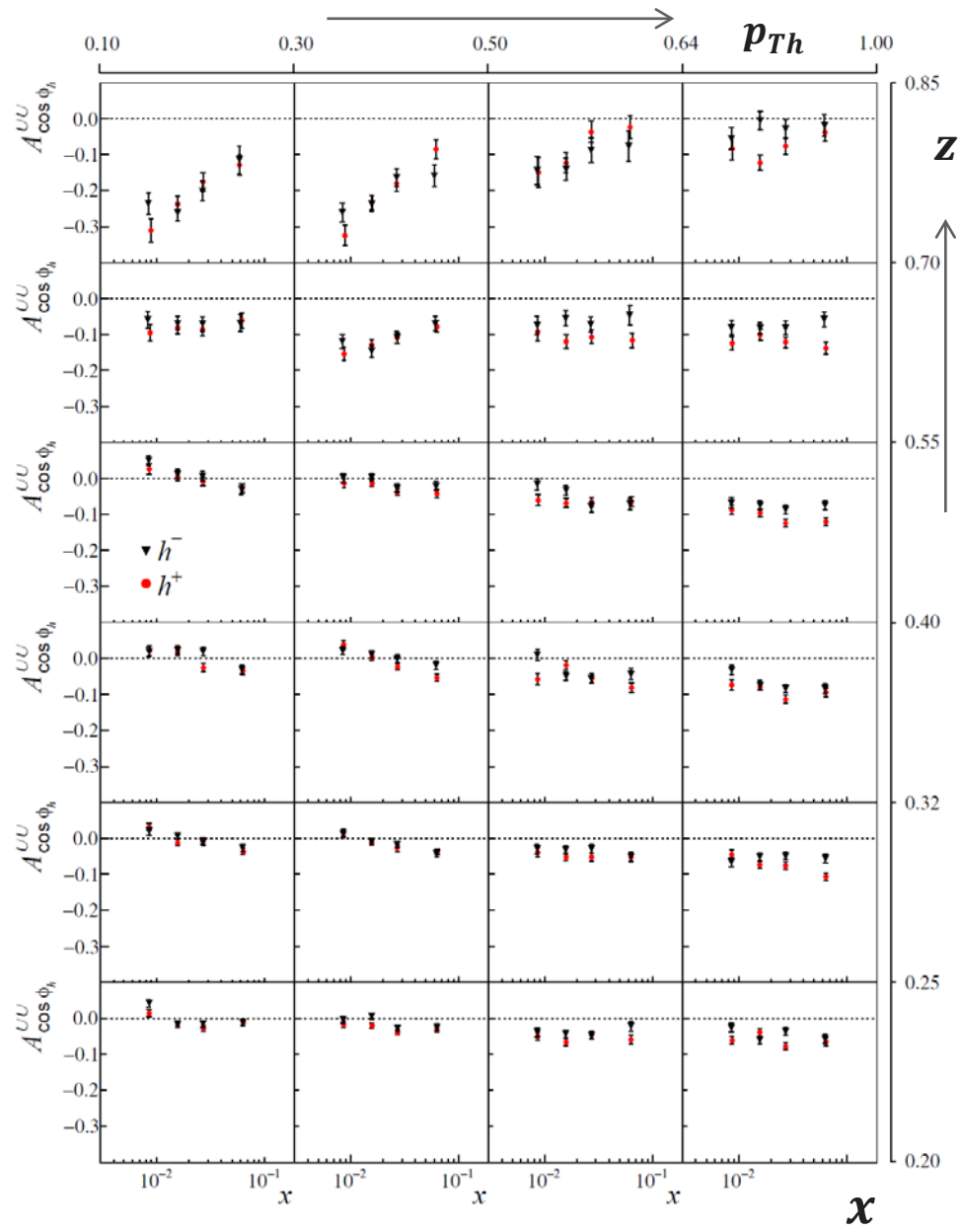
unpolarised SIDIS – azimuthal asymmetries



deuteron

$\cos \phi_h$

$\cos 2\phi_h$



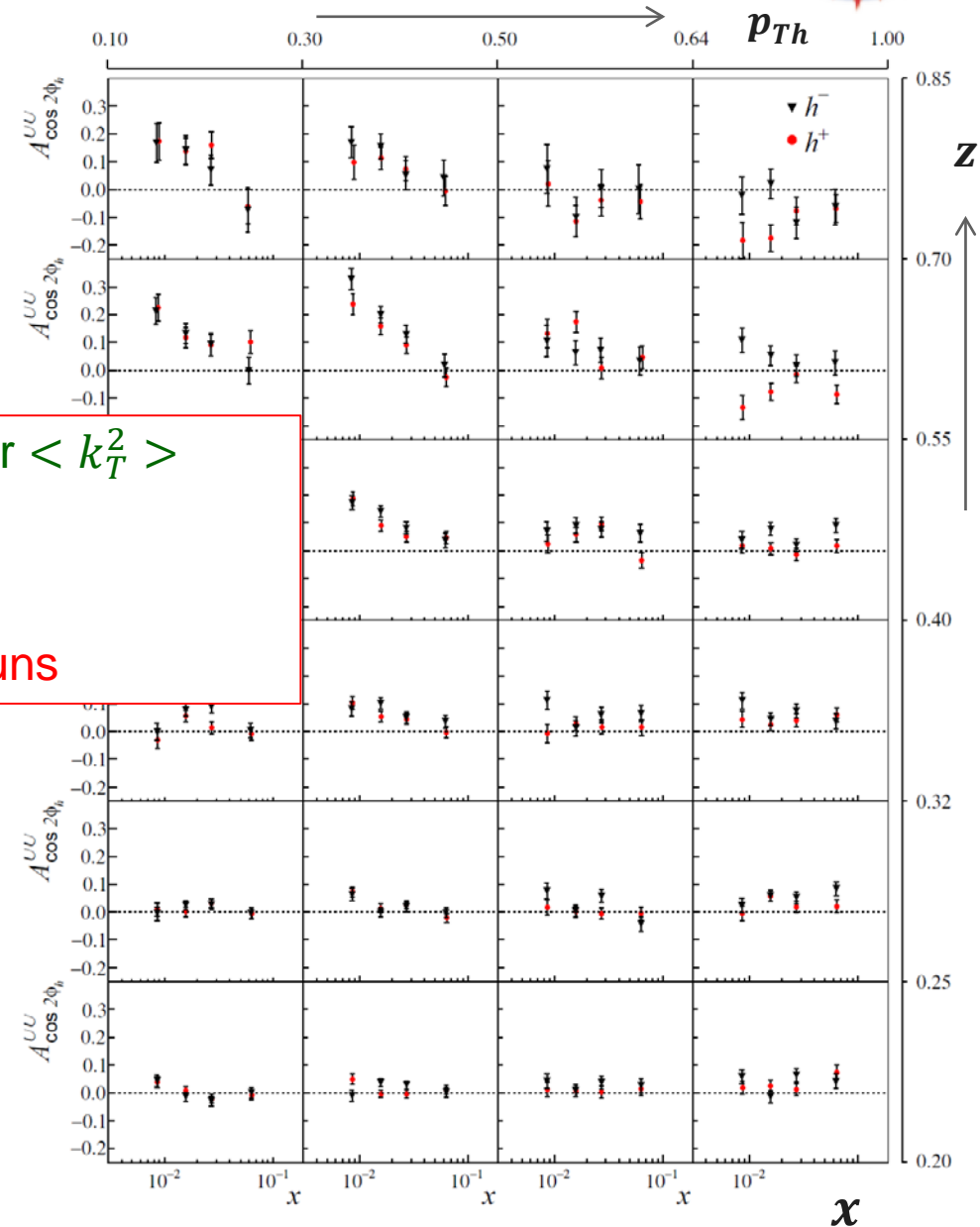
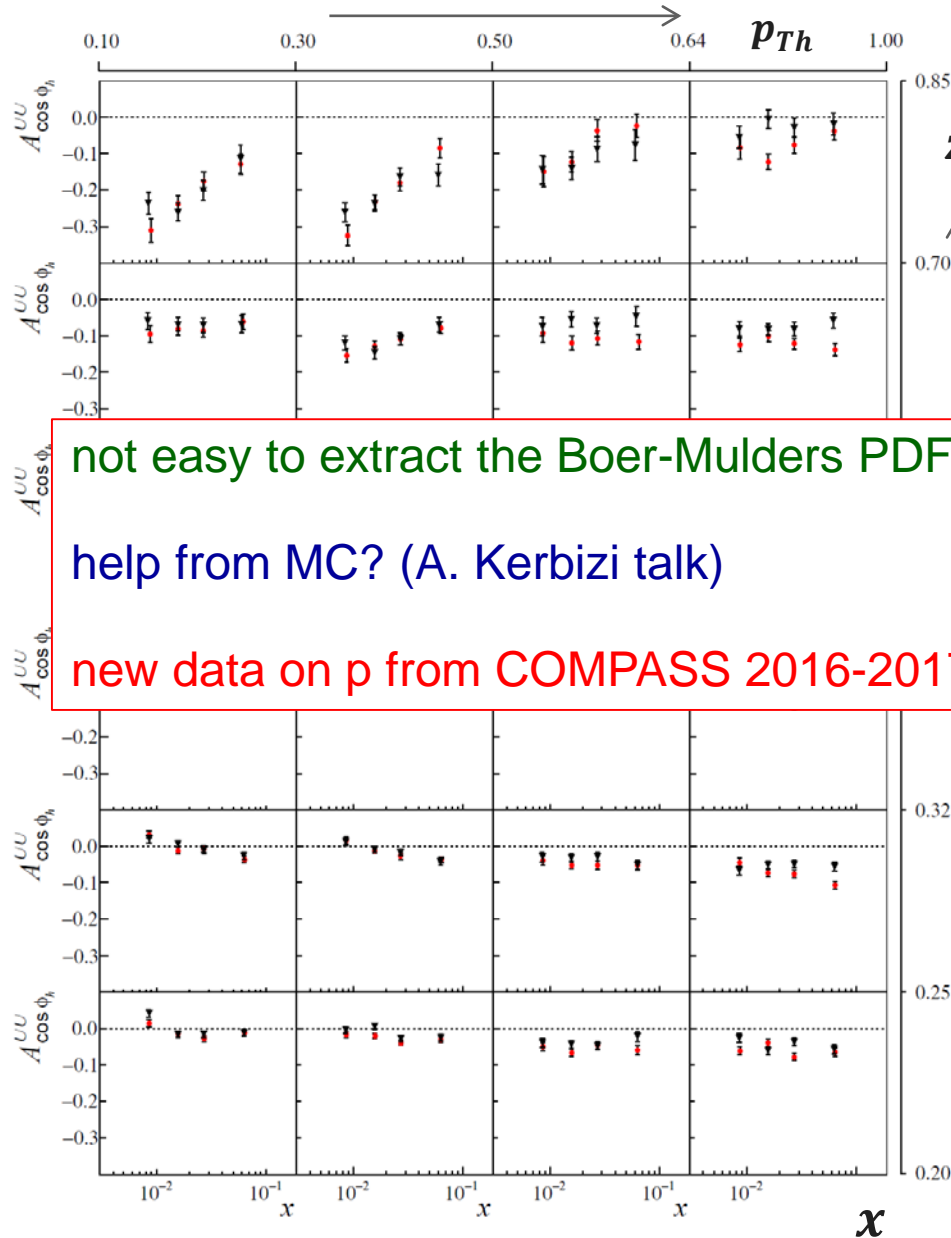
unpolarised SIDIS – azimuthal asymmetries



deuteron

$\cos \phi_h$

$\cos 2\phi_h$



not easy to extract the Boer-Mulders PDF nor $\langle k_T^2 \rangle$

help from MC? (A. Kerbizi talk)

new data on p from COMPASS 2016-2017 runs

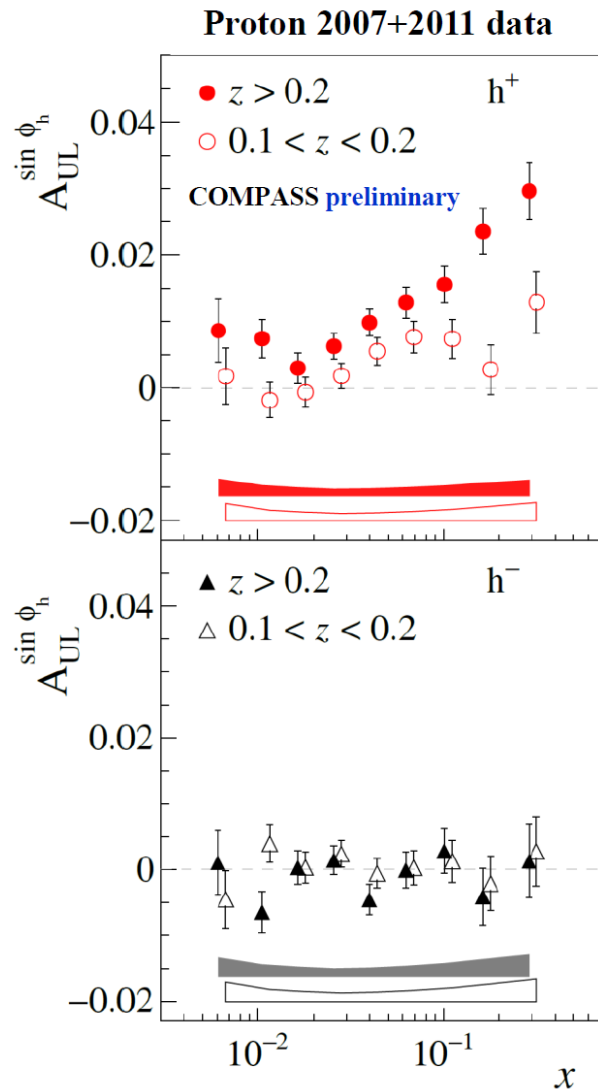
TMDs in SIDIS off longitudinally polarised N

$$\begin{aligned}
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 & \dots \\
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 \end{aligned}$$

SIDIS off longitudinally polarised p



$$A_{UL}^{\sin \phi_h} = F_{UL}^{\sin \phi_h} / F_{UU}$$

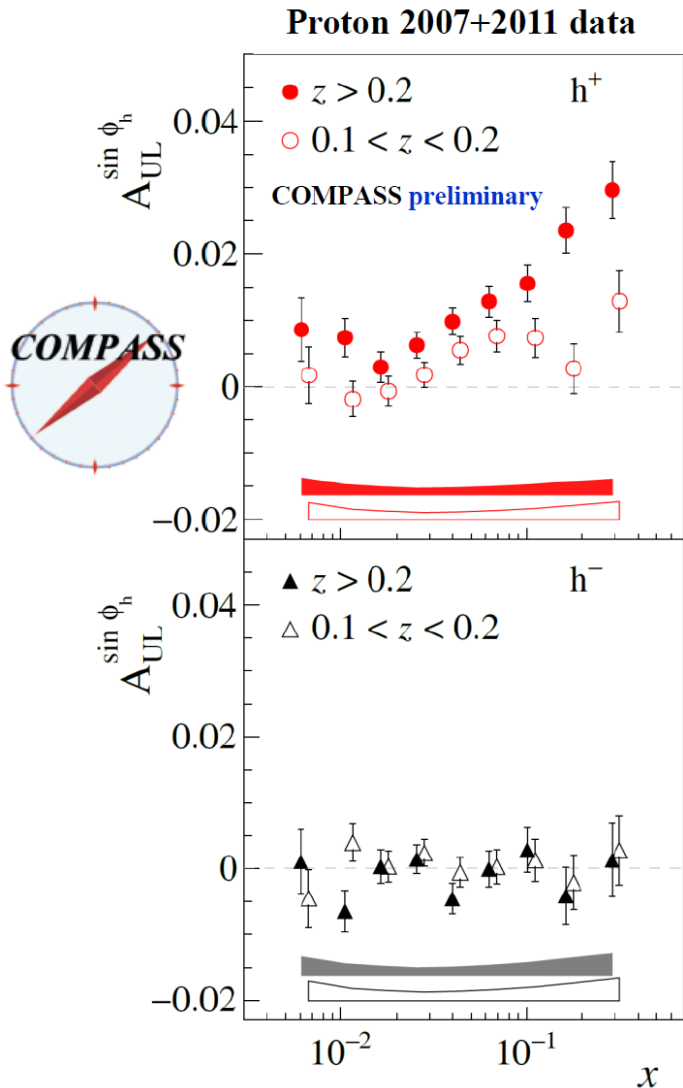


SIDIS off longitudinally polarised p

$$A_{UL}^{\sin \phi_h} = F_{UL}^{\sin \phi_h} / F_{UU}$$

$$F_{UL}^{\sin \phi_h} = \frac{2M}{Q} \mathcal{C} \left\{ -\frac{\hat{h} \cdot \mathbf{p}_T}{M_h} \left(xh_L^q H_{1q}^{\perp h} + \frac{M_h}{M} g_{1L}^q \frac{\tilde{G}_q^{\perp h}}{z} \right) + \frac{\hat{h} \cdot \mathbf{k}_T}{M} \left(xf_L^{\perp q} D_{1q}^h - \frac{M_h}{M} h_{1L}^{\perp q} \frac{\tilde{H}_q^h}{z} \right) \right\}$$

Q-suppressed,
different “twist” contributions

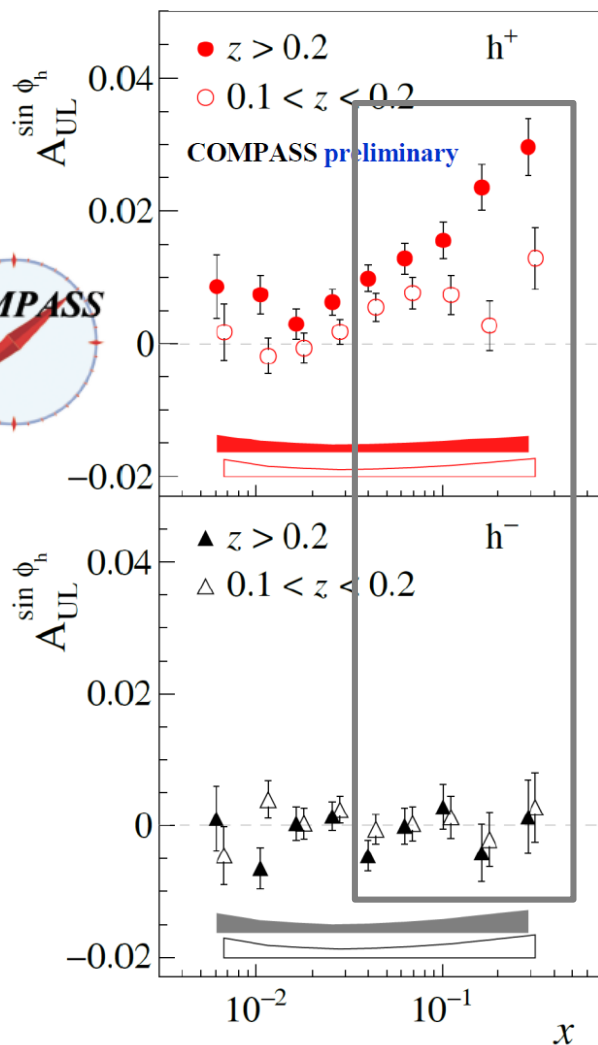


SIDIS off longitudinally polarised p

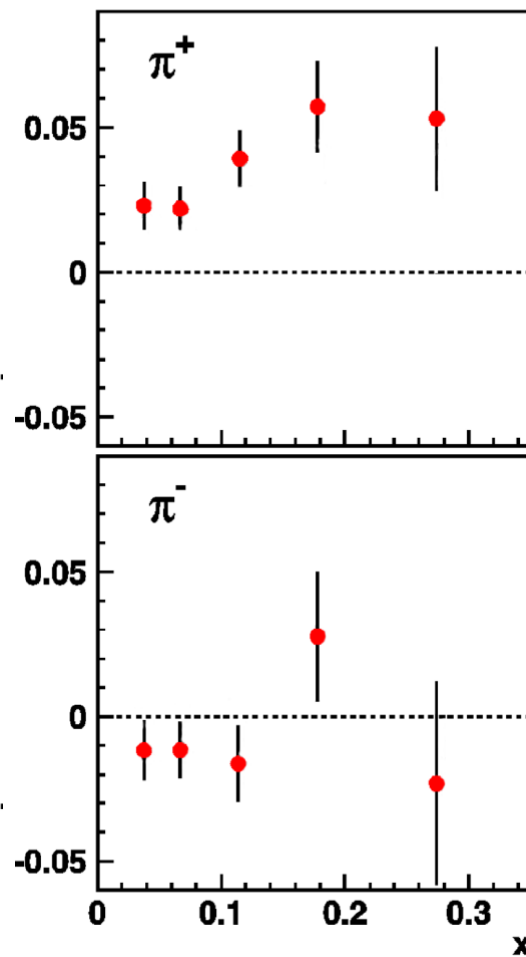
$$A_{UL}^{\sin \phi_h} = F_{UL}^{\sin \phi_h} / F_{UU}$$

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Proton 2007+2011 data



HERMES PLB 622 (2005) 14



Q-suppressed,
different "twist" contributions

$$\sqrt{2 \epsilon (1 + \epsilon)}$$

SIDIS off transversely polarised N

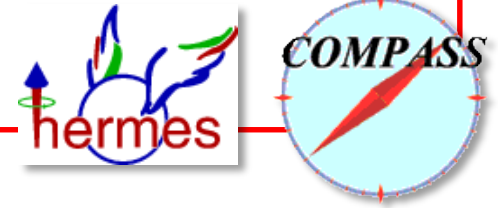
$$\begin{aligned}
 \frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = & \dots \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ \right. \\
 & + |\mathbf{S}_\perp| \left[\begin{aligned} & \overset{f_{IT}^\perp D_I}{\sin(\phi_h - \phi_S)} \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \\ & + \varepsilon \overset{h_{IT}^\perp H_I^\perp}{\sin(\phi_h + \phi_S)} F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \overset{h_{IT}^\perp H_I^\perp}{\sin(3\phi_h - \phi_S)} F_{UT}^{\sin(3\phi_h - \phi_S)} \\ & + \sqrt{2\varepsilon(1+\varepsilon)} \overset{h_I H_I^\perp}{\sin \phi_S} F_{UT}^{\sin \phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \overset{h_I H_I^\perp}{\sin(2\phi_h - \phi_S)} F_{UT}^{\sin(2\phi_h - \phi_S)} \end{aligned} \right] \\
 & + |\mathbf{S}_\perp| \lambda_e \left[\begin{aligned} & \overset{g_{IT} D_I}{\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S)} F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{\varepsilon(1-\varepsilon)} \overset{g_{IT} D_I}{\cos \phi_S} F_{LT}^{\cos \phi_S} \\ & + \sqrt{2\varepsilon(1-\varepsilon)} \overset{g_{IT} D_I}{\cos(2\phi_h - \phi_S)} F_{LT}^{\cos(2\phi_h - \phi_S)} \end{aligned} \right] \left. \right\},
 \end{aligned}$$

Semi-Inclusive Deep Inelastic Scattering

MAJOR RESULT:

in the past 10 years 2 of these new PDF's have been measured and shown to be different from zero

by COMPASS and HERMES



the transversity PDF

amplitude of the sine modulation in $\phi_h + \phi_s - \pi$
Collins asymmetry $\sim h_1^\perp \otimes H_1^\perp$

the Sivers PDF

amplitude of the sine modulation in $\phi_h - \phi_s$
Sivers asymmetry $\sim f_{1T}^\perp \otimes D_1$

A STEP TOWARDS THE 3-D STRUCTURE OF THE NUCLEON

Collins asymmetry

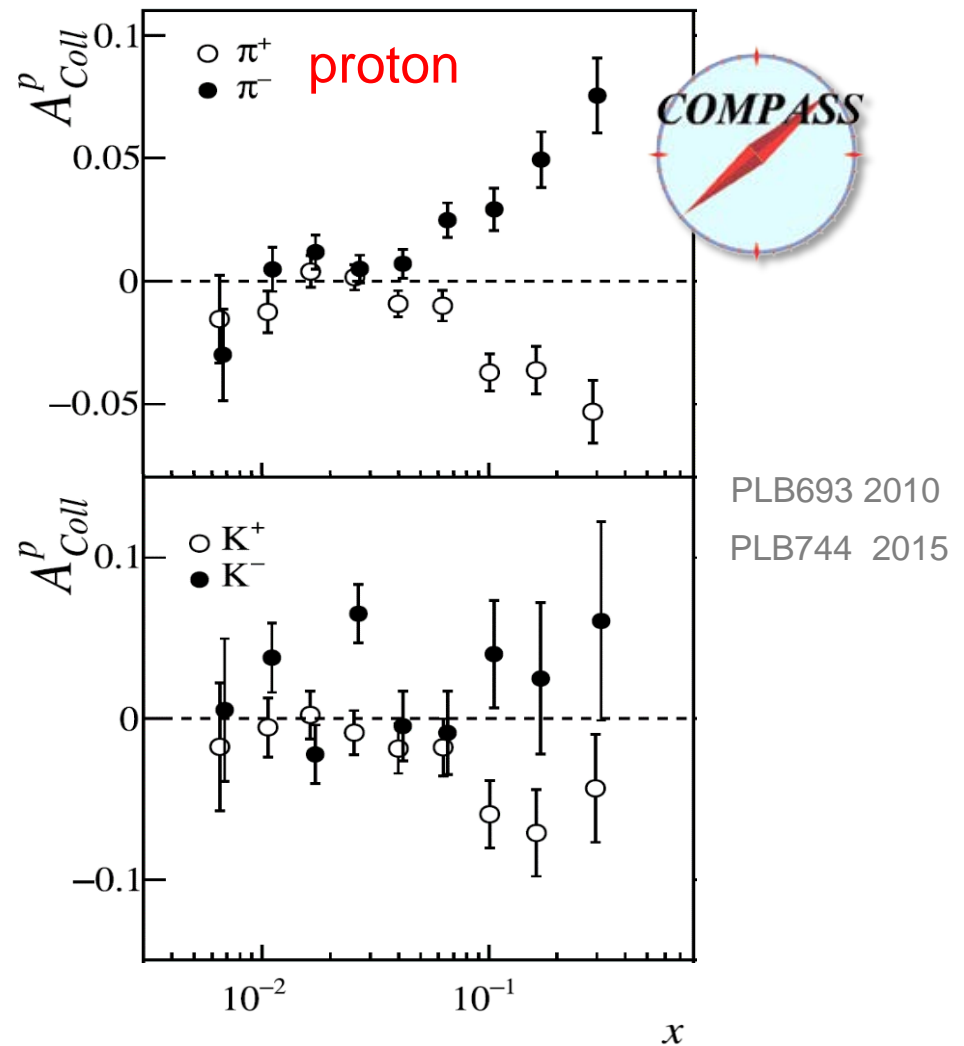
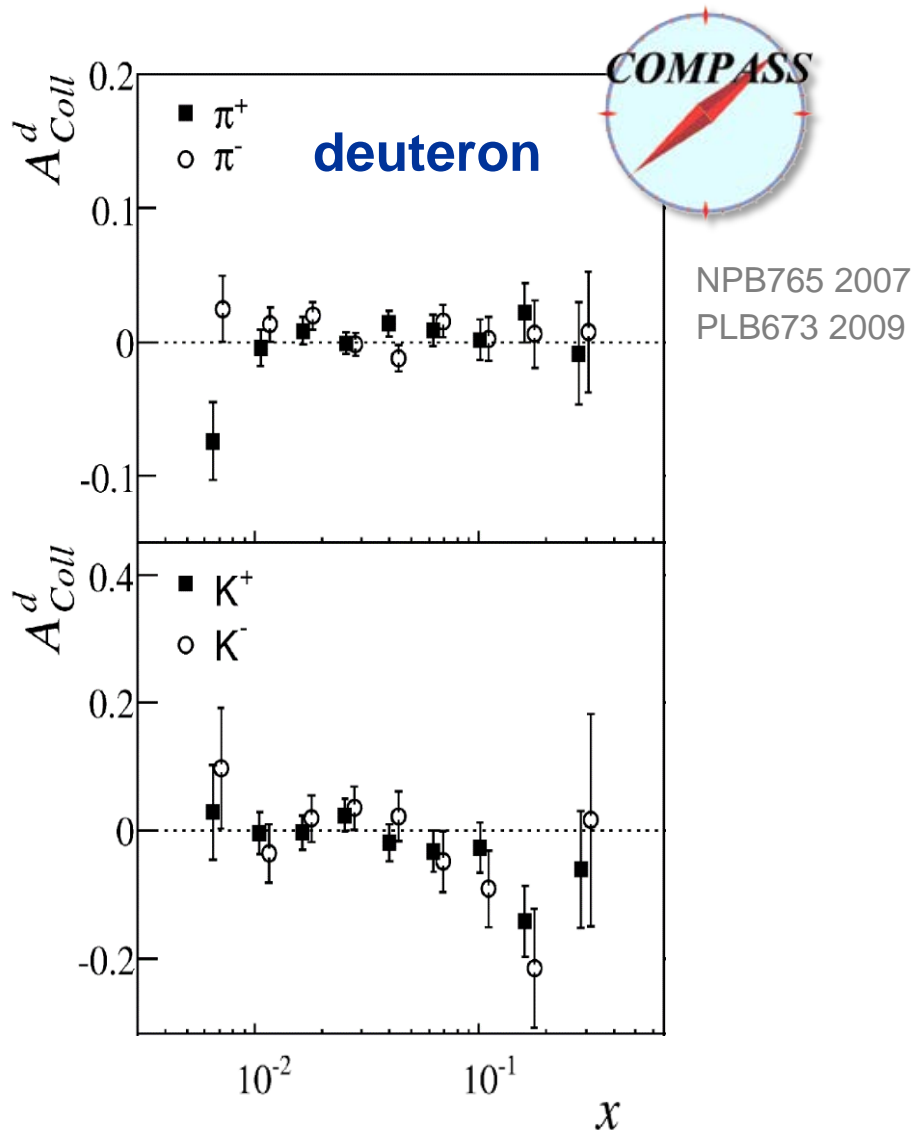
Collins asymmetry

$$\sim h_1 \otimes H_1^\perp$$

2004: first evidence for non-zero Collins asymmetry on p from HERMES



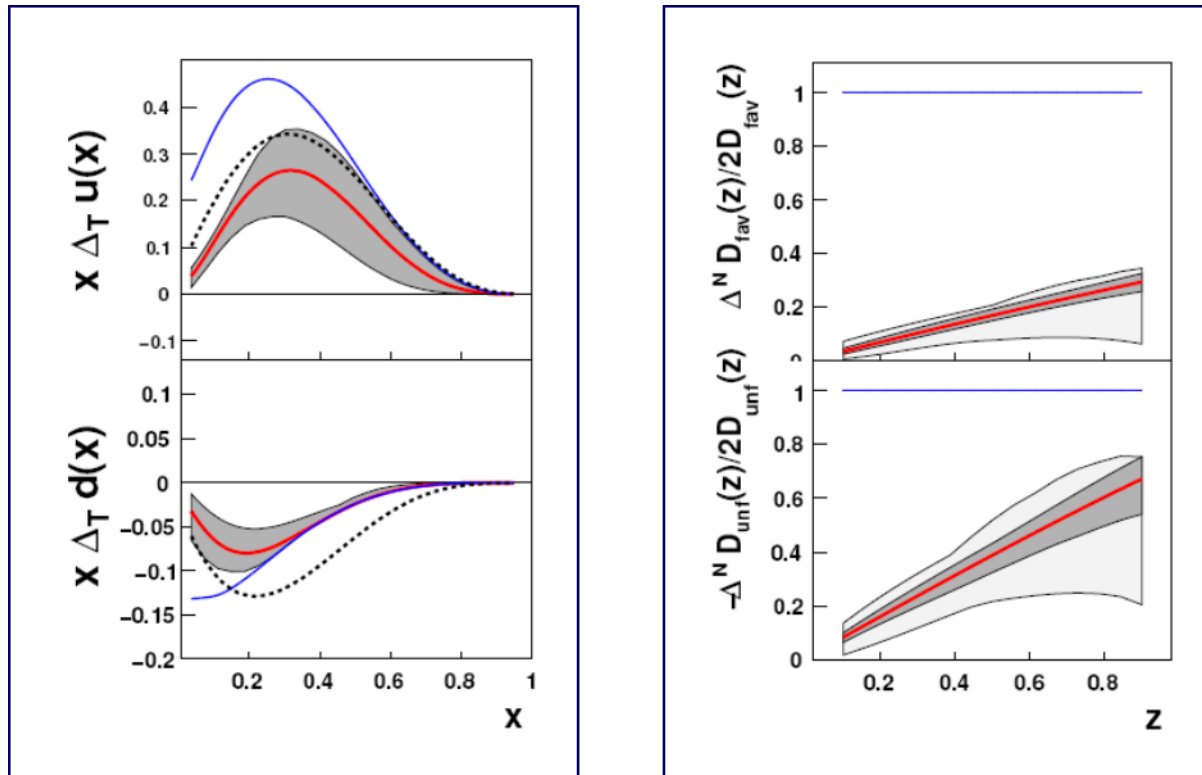
final COMPASS results



Transversity from SIDIS

M. Anselmino et al., Nucl. Phys. Proc. Suppl. 2009

fit to HERMES p, COMPASS d, Belle e+e- data



and many others ...

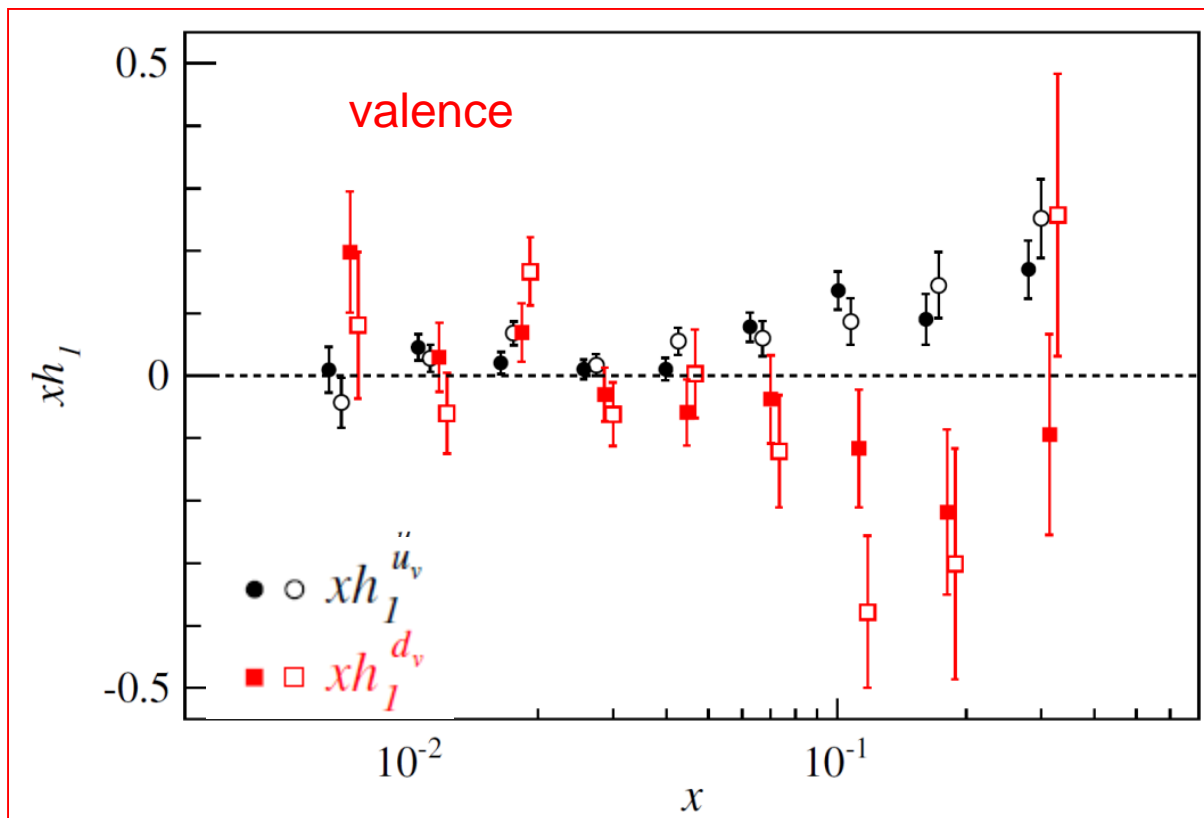
Transversity from SIDIS

Collins and di-hadron asymmetries

point by point extraction

one can use directly the COMPASS p and d asymmetries, and the Belle data to evaluate the analysing power (with some “reasonable” assumptions)

advantage: no MC nor parametrisation is needed



open points: dihadron

closed points: Collins

large uncertainties
on the d distribution
due to the poor
deuteron/neutron
data

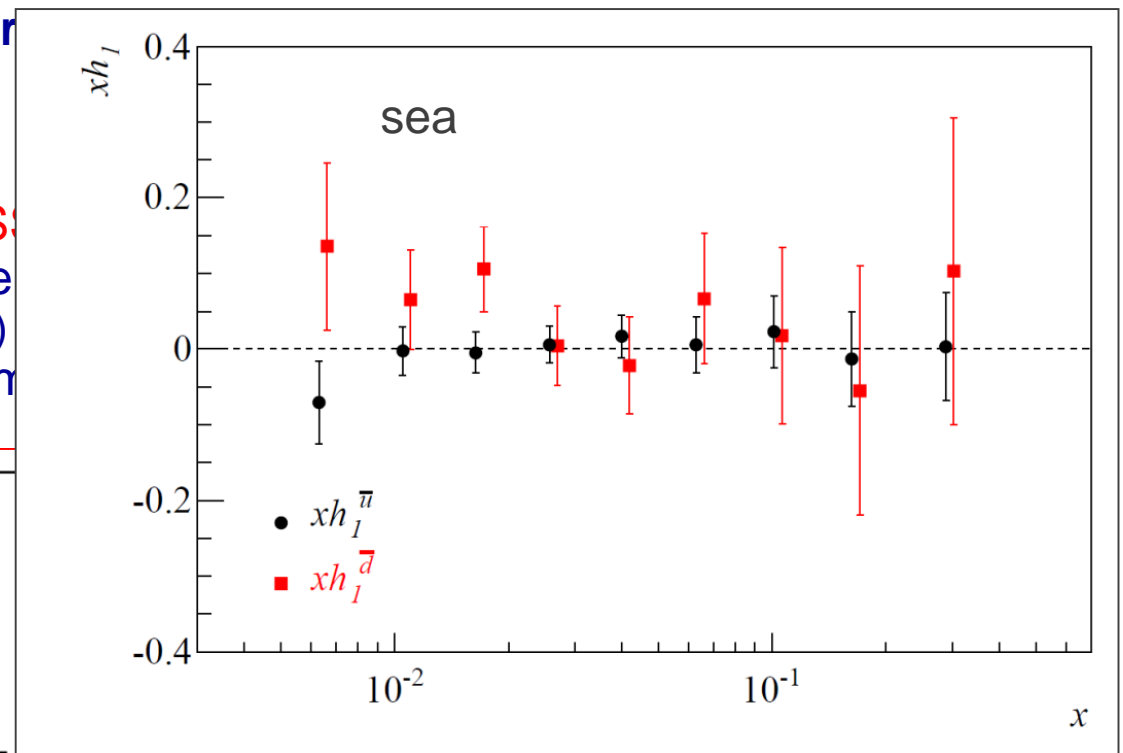
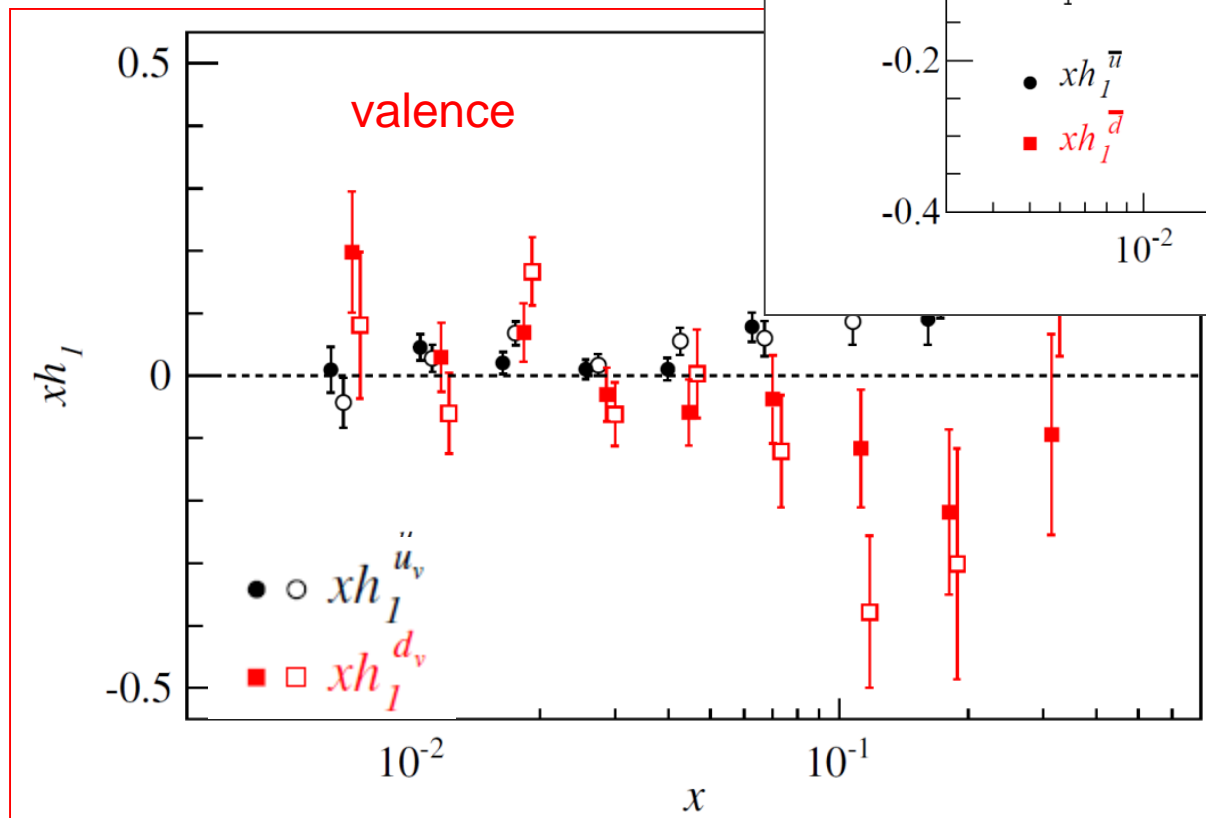
A. Martin F. B. V. Barone
PRD91 2015

Transversity from SIDIS

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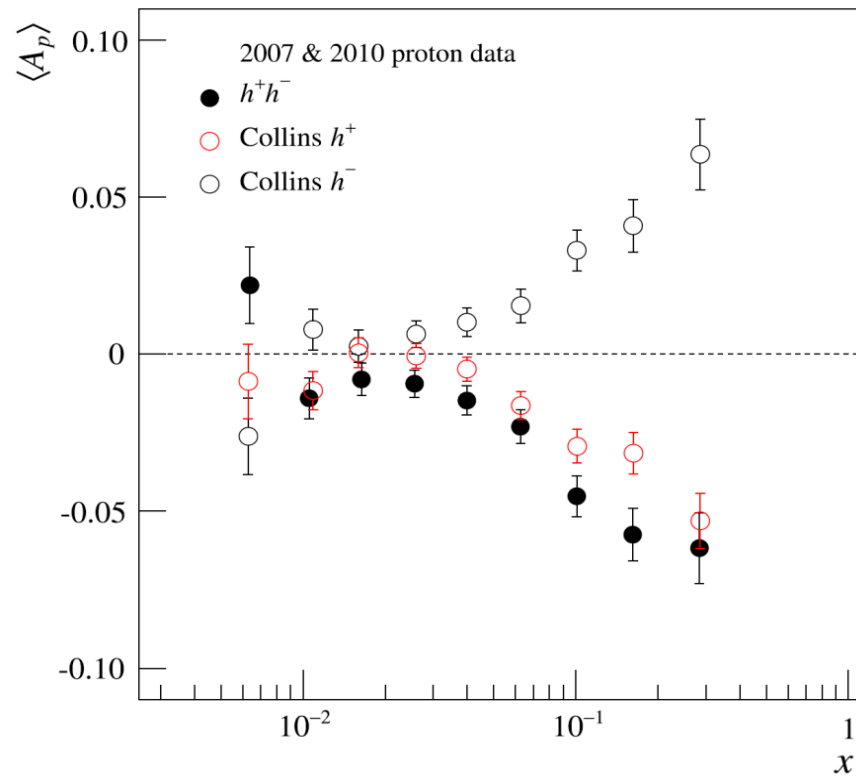
large uncertainties
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due to the poor
deuteron/neutron
data

A. Martin F. B. V. Barone
PRD91 2015

Transversity from SIDIS

not shown here

- di-hadron asymmetries [PLB 713 (2012) 10, PLB 736 (2014) 124]
- interplay among transversity induced asymmetries [PLB 753 (2016) 406]



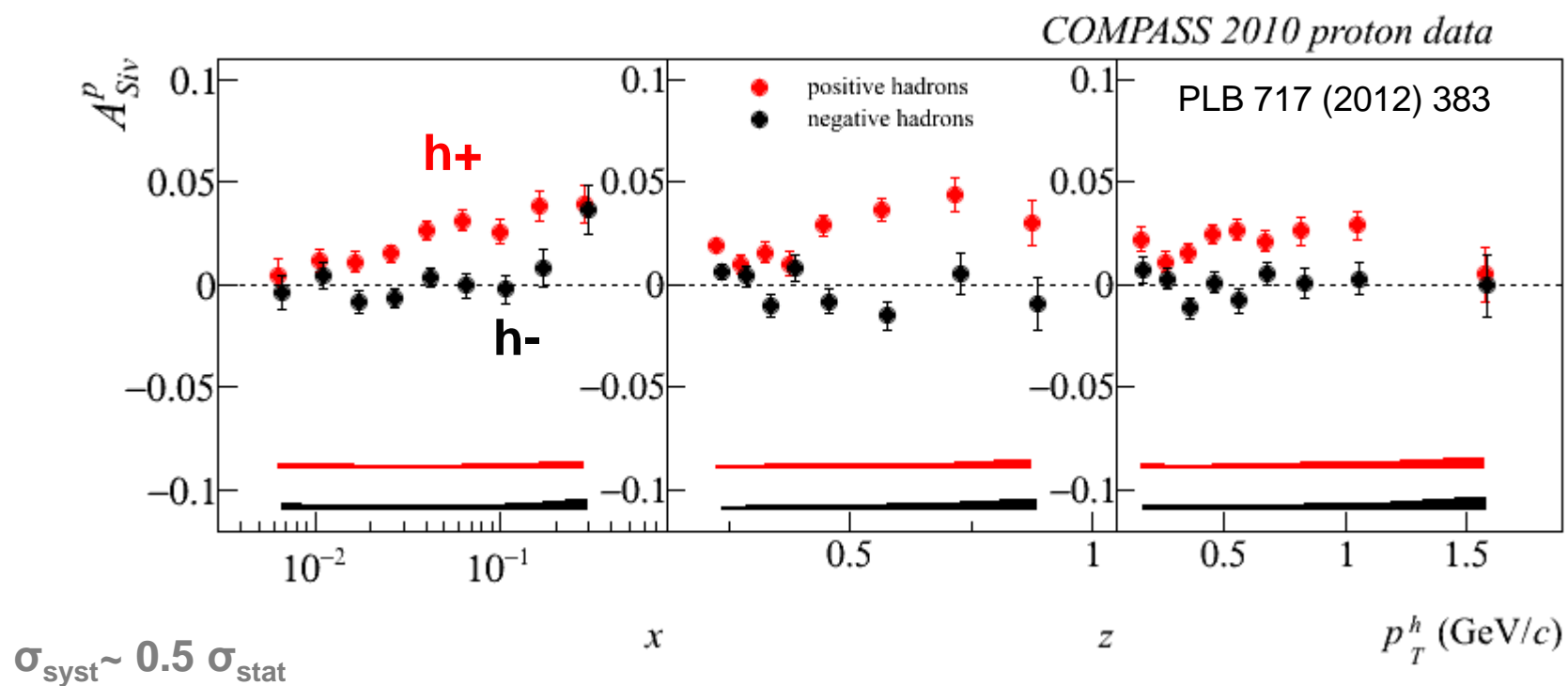
Sivers asymmetries



Sivers asymmetries on **proton**

charged hadrons

2010 data



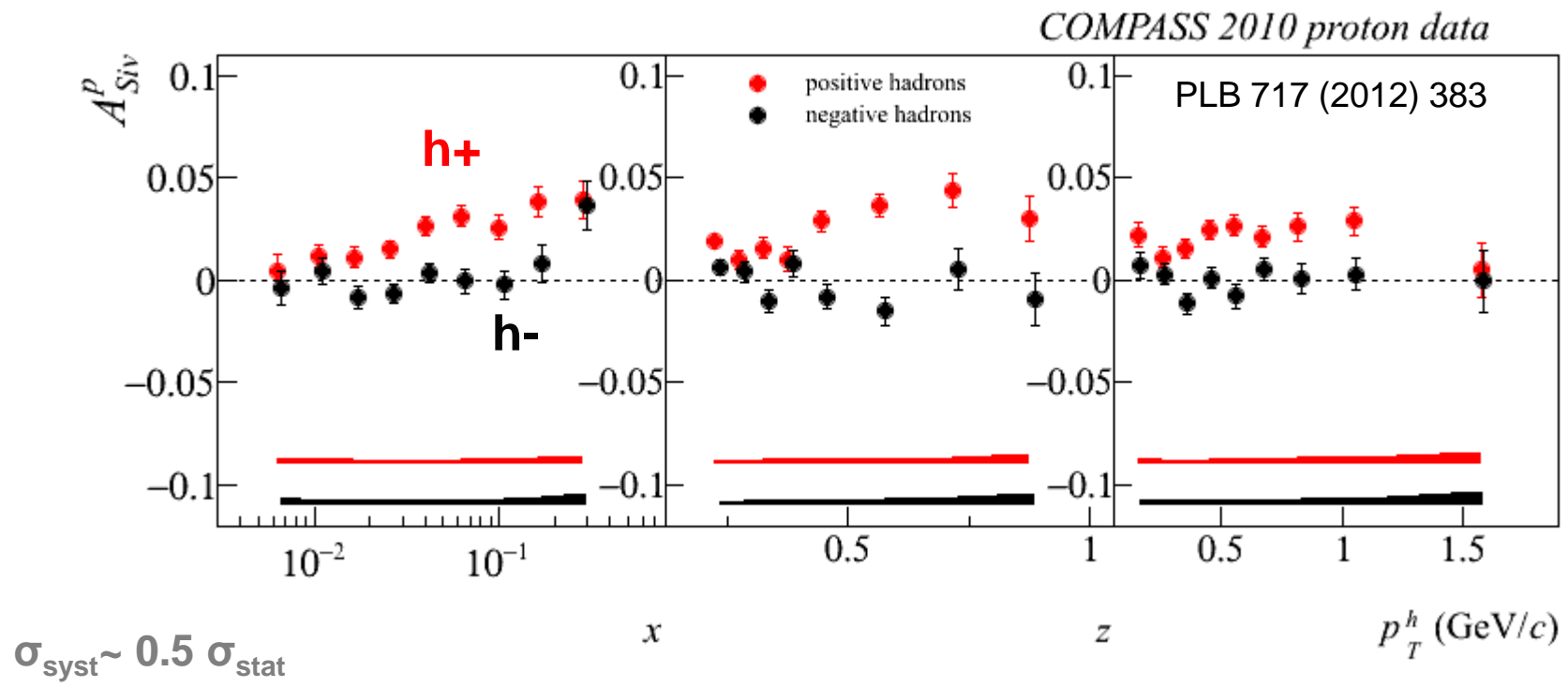
clear evidence for a positive signal for h^+ , which extends to small x



Sivers asymmetries on **proton**

charged hadrons

2010 data



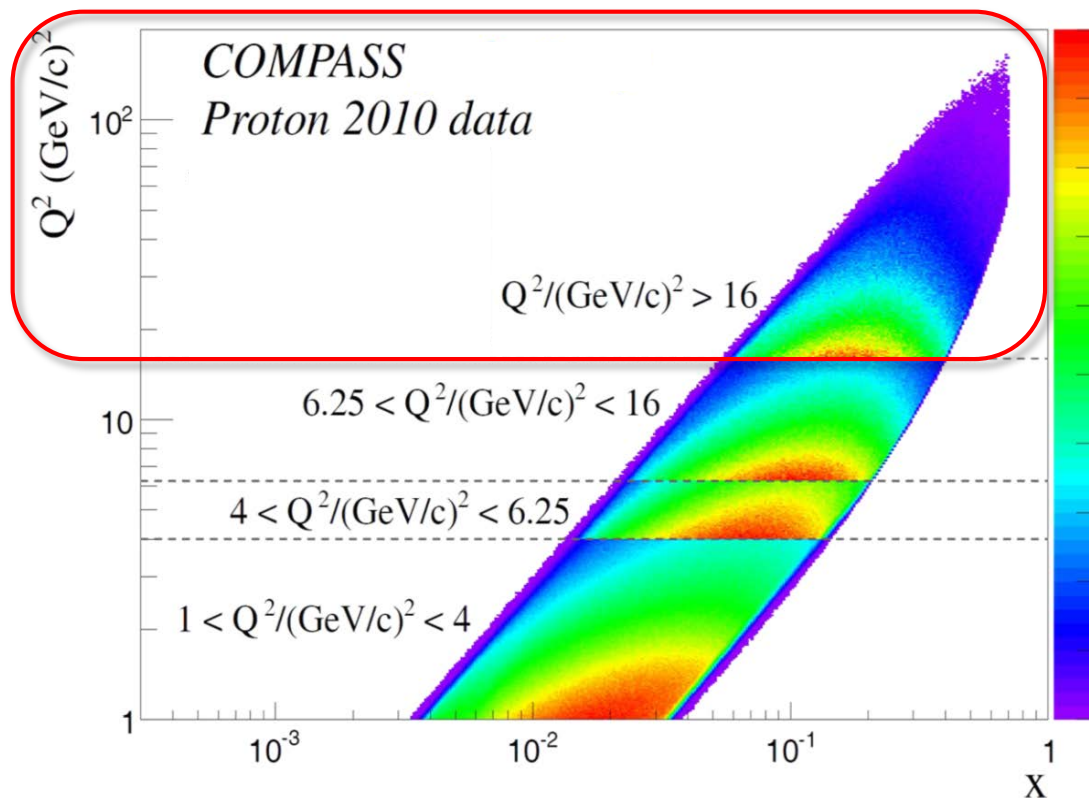
clear evidence for a positive signal for h^+ , which extends to small x

d data compatible with zero but with large statistical uncertainties



Sivers asymmetries on **proton**

COMPASS has measured the SIDIS TSA in the four Q^2 ranges of the Drell-Yan measurement



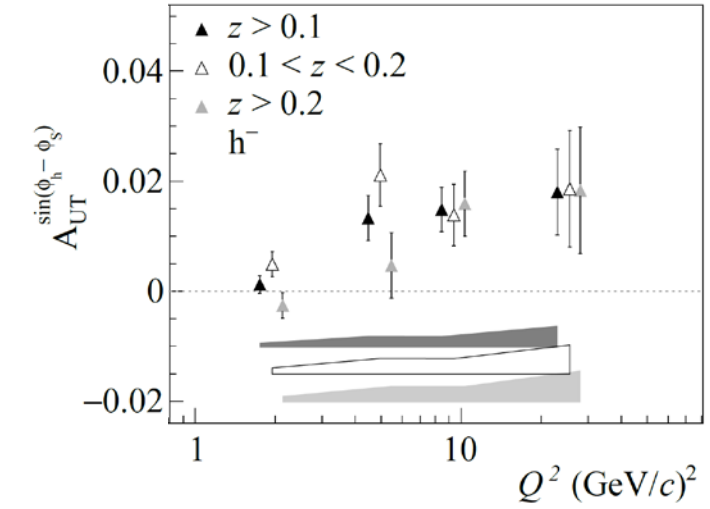
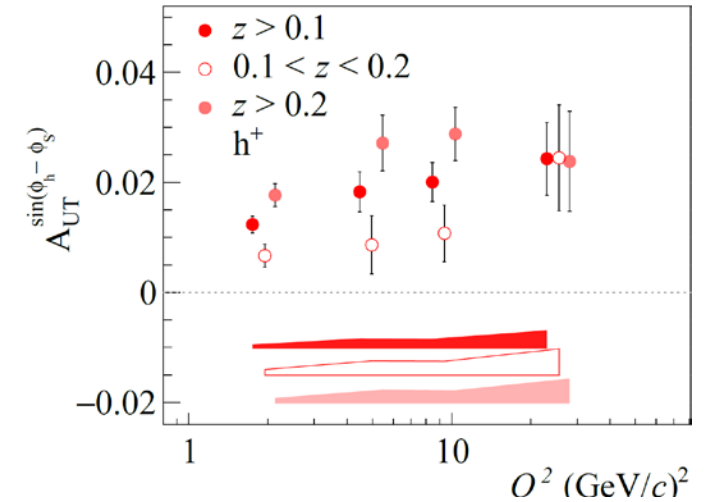
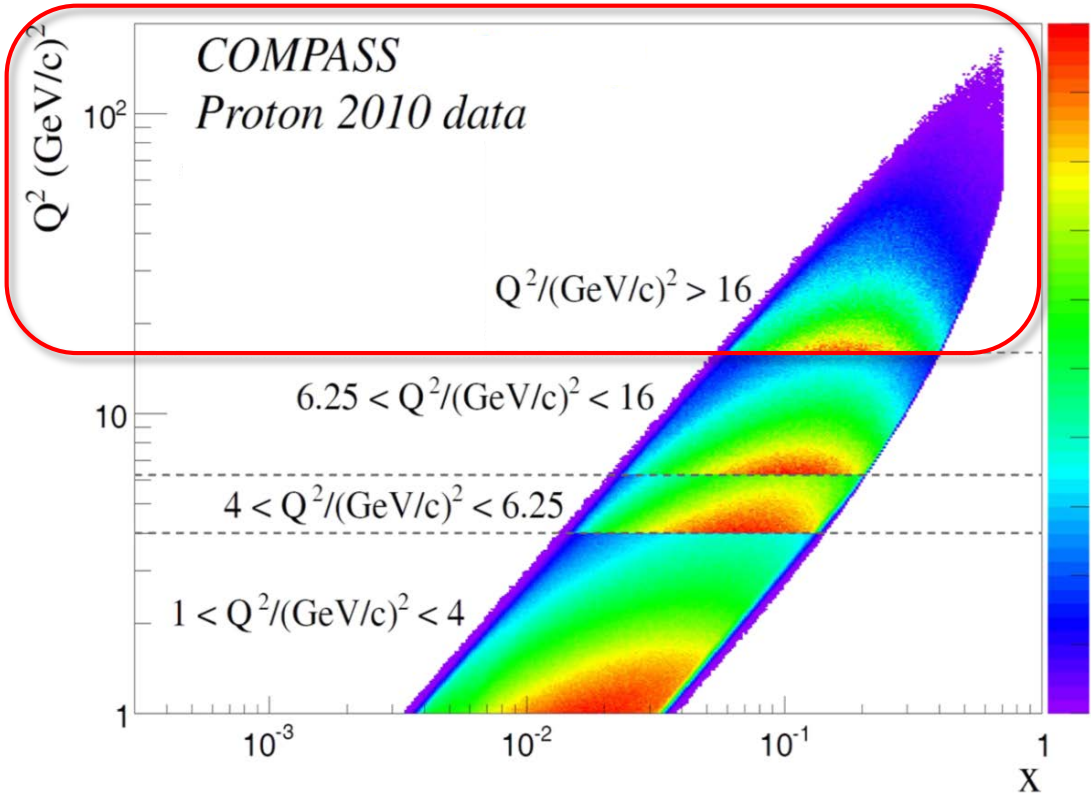
“golden” region for DY: $Q^2 > 16 \text{ GeV}^2$



Sivers asymmetries on **proton**

COMPASS has measured the SIDIS TSA in the four Q^2 ranges of the Drell-Yan measurement

PLB 770 (2017) 138



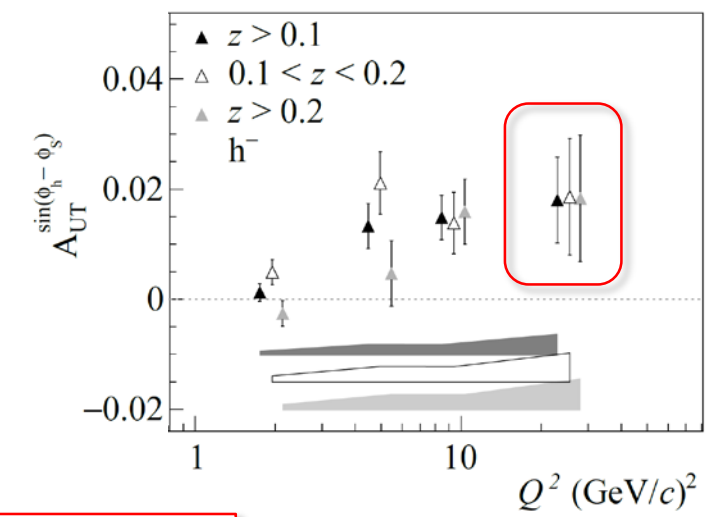
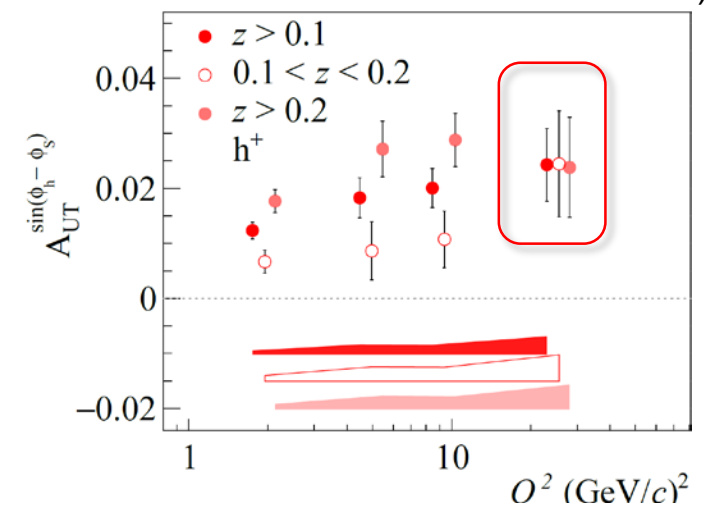
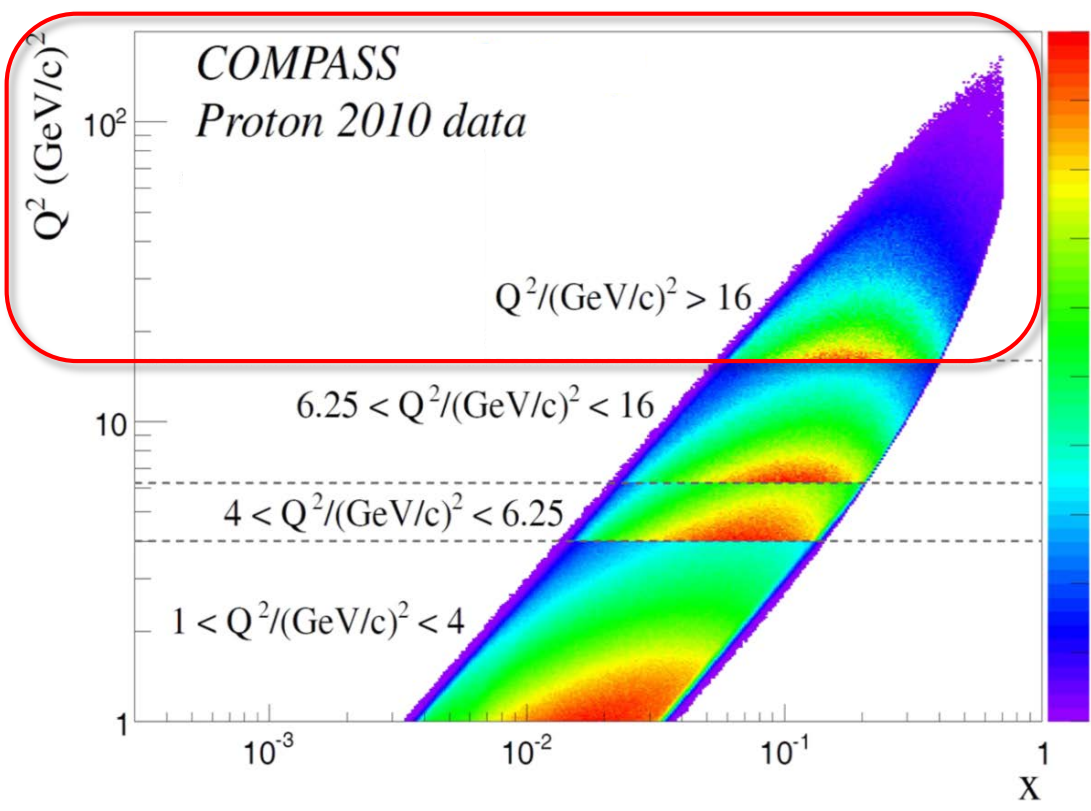
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Sivers asymmetries on **proton**

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PLB 770 (2017) 138



“golden” region for DY: $Q^2 > 16 \text{ GeV}^2$

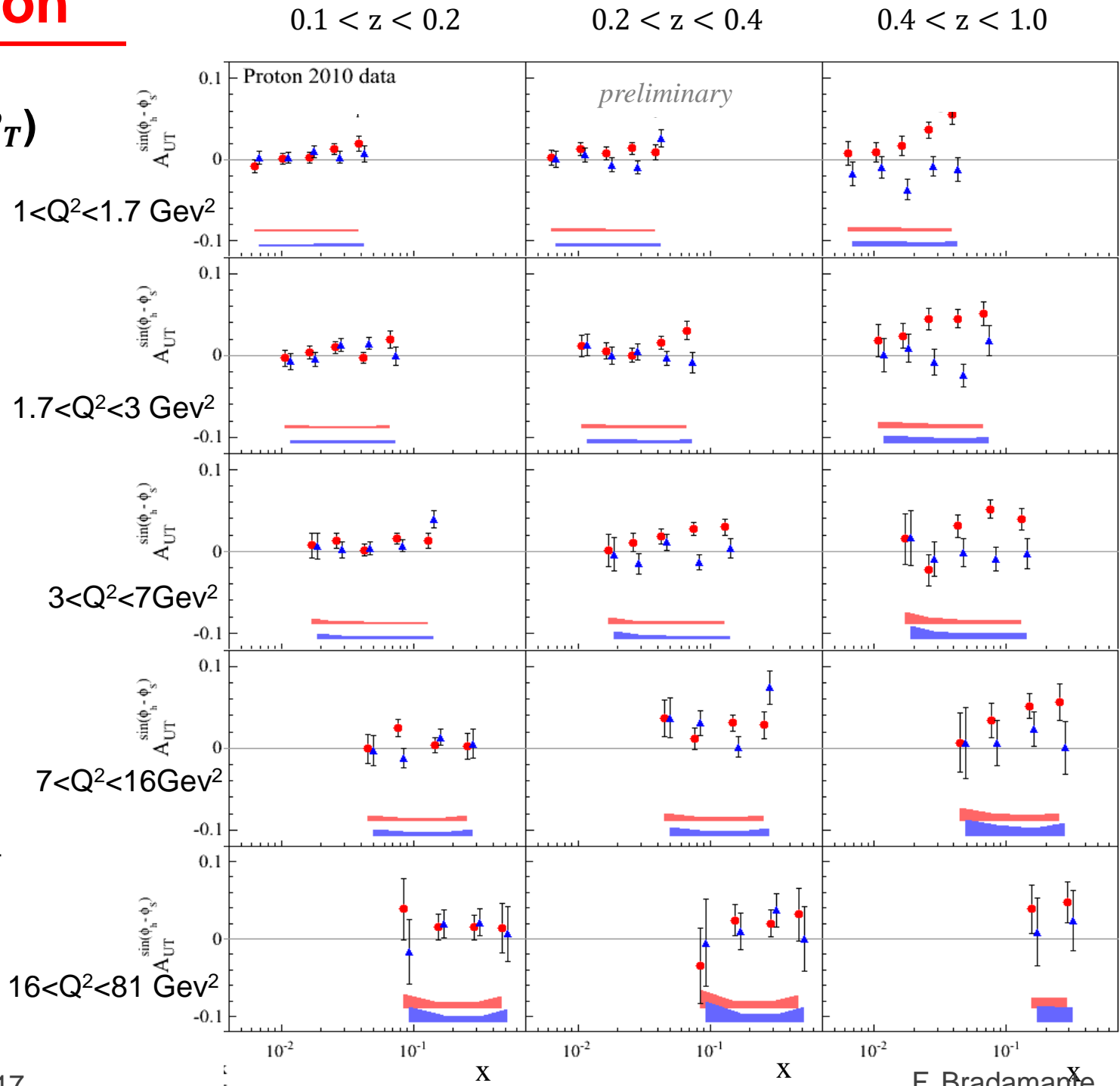
clearly positive
test of change of sign feasible

TSA on **proton**

multiD ($x, Q^2; z, P_T$)
analysis

an example:
Sivers
asymmetry

$P_T > 0.1 \text{ GeV}/c$



Sivers asymmetries



not shown here: gluon Sivers

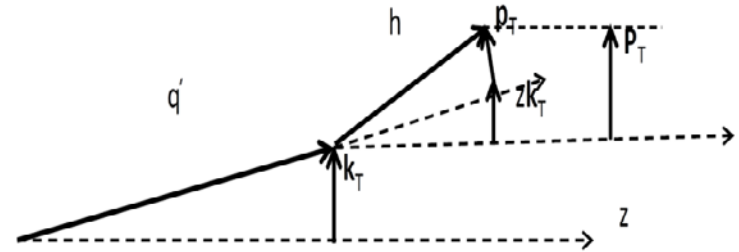
- J/Ψ asymmetry [J. Matousek, DSPIN-15]
- high p_T hadron pair asymmetry [K. Kurek, DSPIN-15, PLB 772 (2017) 854]

new:
weighted
Sivers asymmetries
in SIDIS

weighted Siverts asymmetries - why

“standard” Siverts asymmetry

$$A_{Siv}(x, z) = \frac{\sum_q e_q^2 x f_{1T}^{\perp q}(x) \otimes D_{1q}(z)}{\sum_q e_q^2 x f_1^q(x) D_{1q}(z)}$$



to evaluate the convolution, models are needed
e.g. Gaussian model

with several assumptions, one gets

$$A_{Siv,G}(x, z) \cong \frac{\pi M}{2 \langle P_T \rangle} \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) z D_{1q}(z)}{\sum_q e_q^2 x f_1^q(x) D_{1q}^h(z)}$$

$$f_{1T}^{\perp(1)q}(x) = \int d^2 \vec{k}_T \frac{k_T^2}{2M^2} f_{1T}^{\perp q}(x, k_T^2)$$

weighted Sivers asymmetries

by **weighting** the spin dependent part of the cross-section with P_T
one can solve the convolution

there are two slightly different possibilities:

$$w = P_T/zM \quad A_{Siv}^w(x, z) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) D_{1q}(z)}{\sum_q e_q^2 x f_1^q(x) D_{1q}(z)}$$

easier interpretation

$$w' = P_T/M \quad A_{Siv}^{w'}(x, z) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) z D_{1q}(z)}{\sum_q e_q^2 x f_1^q(x) D_{1q}(z)}$$

easier comparison with the “gaussian ansatz”

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easier comparison with the “gaussian ansatz”

COMPASS has measured both, in bins of x and z

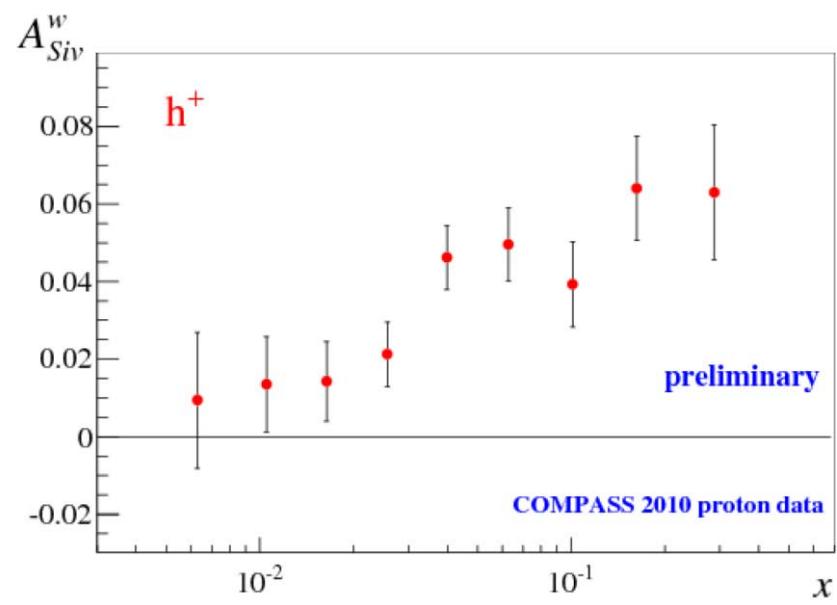


P_T/zM weighted Sivers asymmetry

$$A_{Siv}^w(x) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) \int D_{1q}(z) dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz}$$

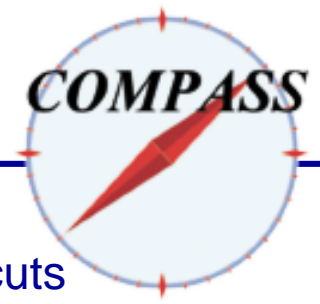
$$w = P_T/zM$$

standard cuts
 $z > 0.2$
arXiv:1702.00621



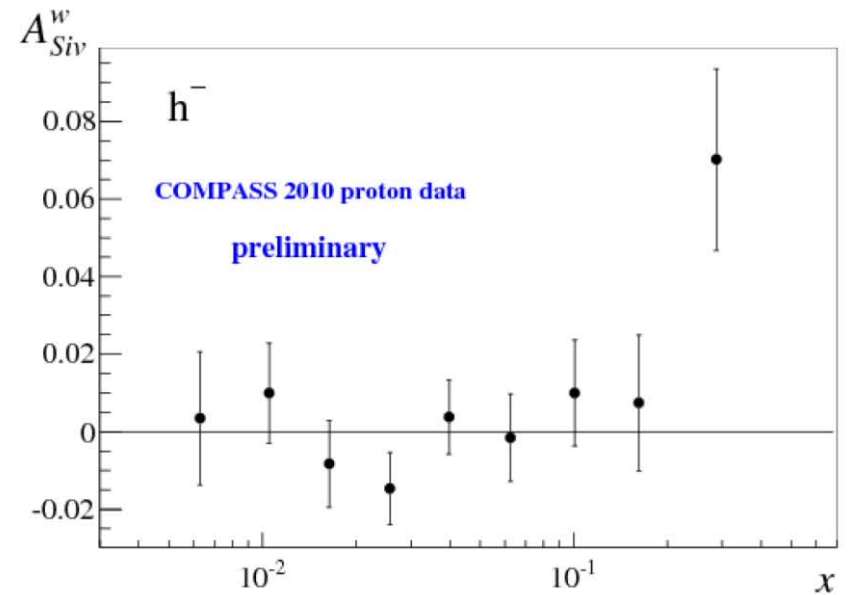
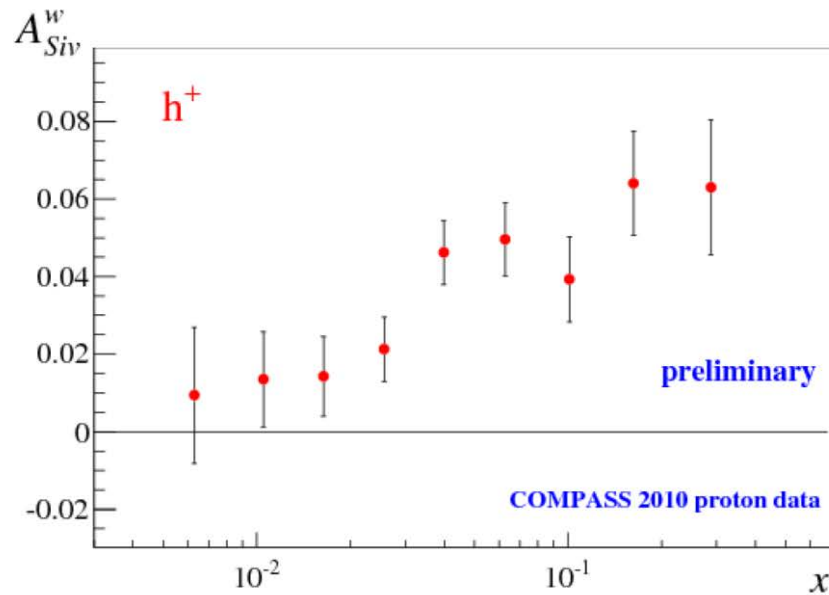
$$\sim 2 \frac{f_{1T}^{\perp(1)u}(x)}{f_1^u(x)}$$

P_T/zM weighted Sivers asymmetry



$$A_{Siv}^w(x) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) \int D_{1q}(z) dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz} \quad w = P_T/zM$$

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 $z > 0.2$
 arXiv:1702.00621



$$\sim 2 \frac{f_{1T}^{\perp(1)u}(x)}{f_1^u(x)}$$

both $f_{1T}^{\perp(1)u}$ and $f_{1T}^{\perp(1)d}$ contribute

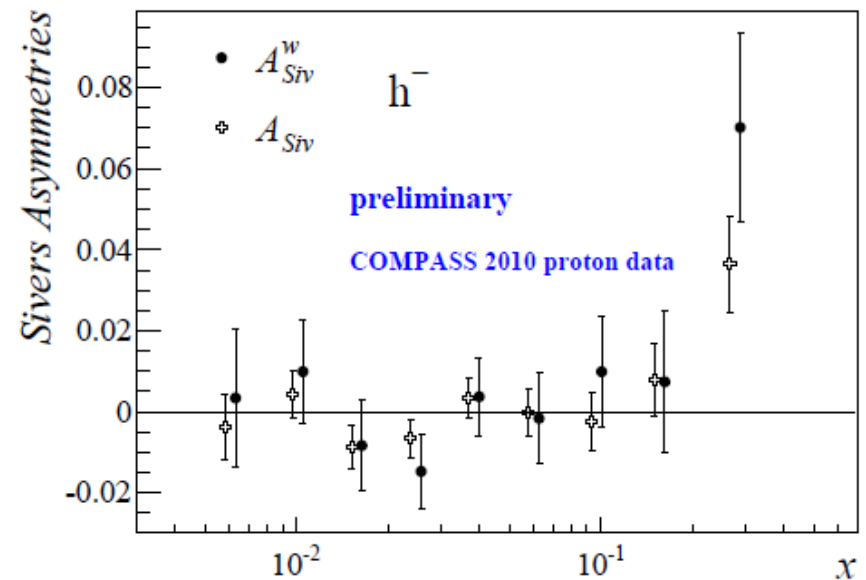
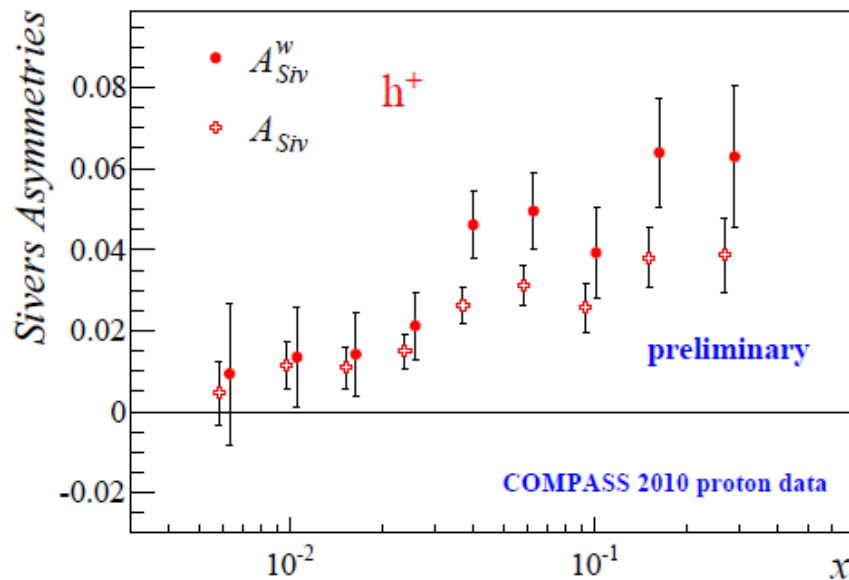


P_T/zM weighted Sivers asymmetry

$$A_{Siv}^w(x) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) \int D_{1q}(z) dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz} \quad w = P_T/zM$$

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comparison with the published “standard” asymmetries



open points: standard asymmetries, PLB 717 (2012) 383



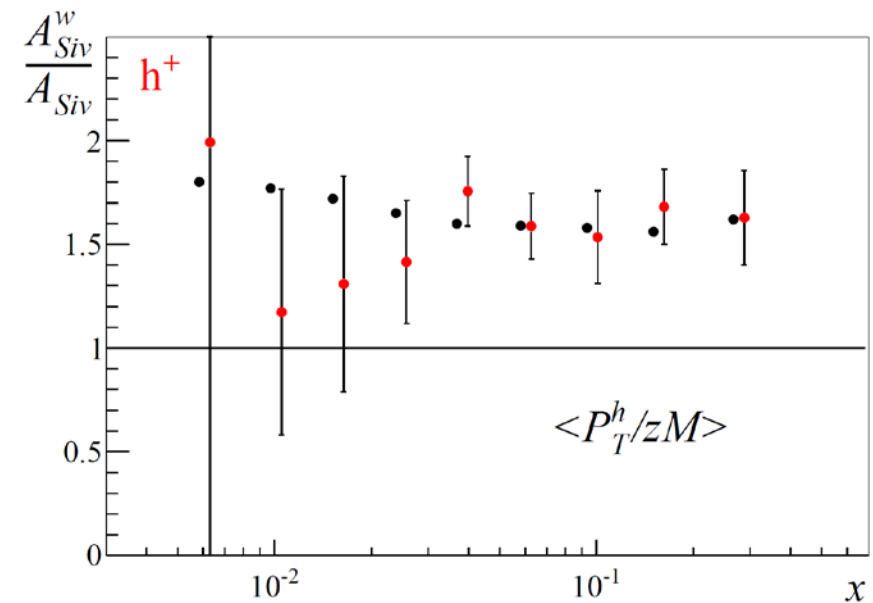
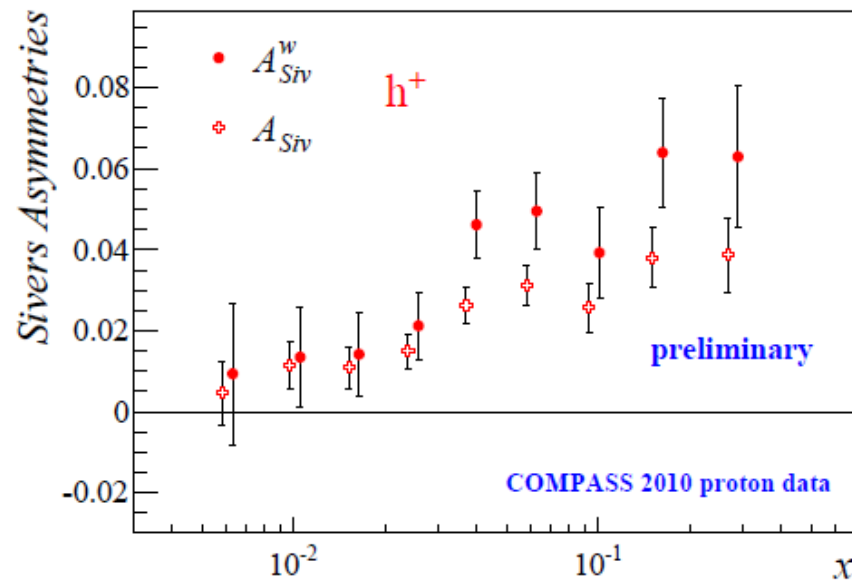
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P_T/zM weighted Sivers asymmetry

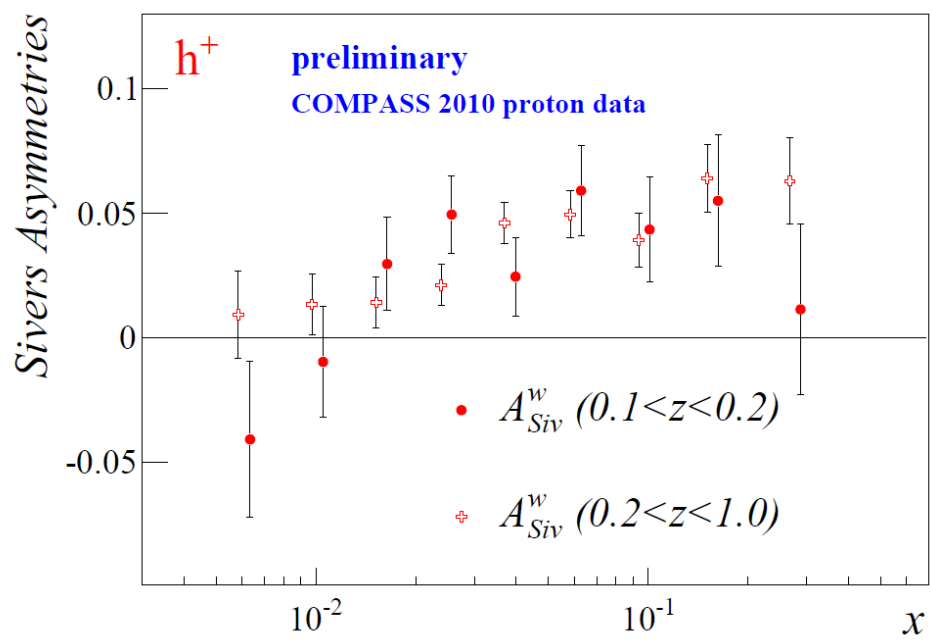
$$A_{Siv}^w(x) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) \int D_{1q}(z) dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz}$$

$$w = P_T/zM$$

standard cuts

arXiv:1702.00621

different z ranges



$$\sim 2 \frac{f_{1T}^{\perp(1)u}(x)}{f_1^u(x)}$$

ok



P_T/zM weighted Siverts asymmetry

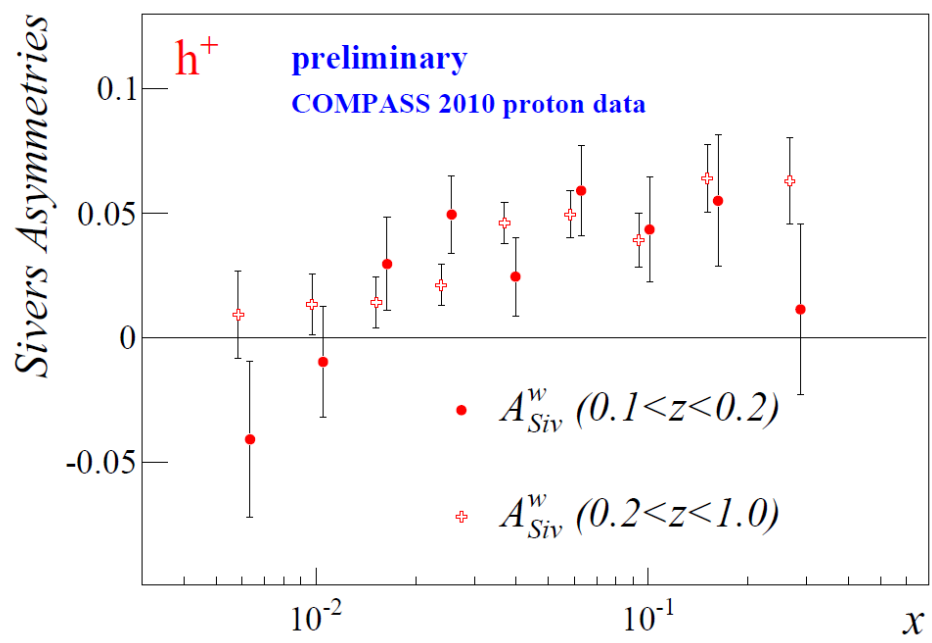
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$w = P_T/zM$

standard cuts

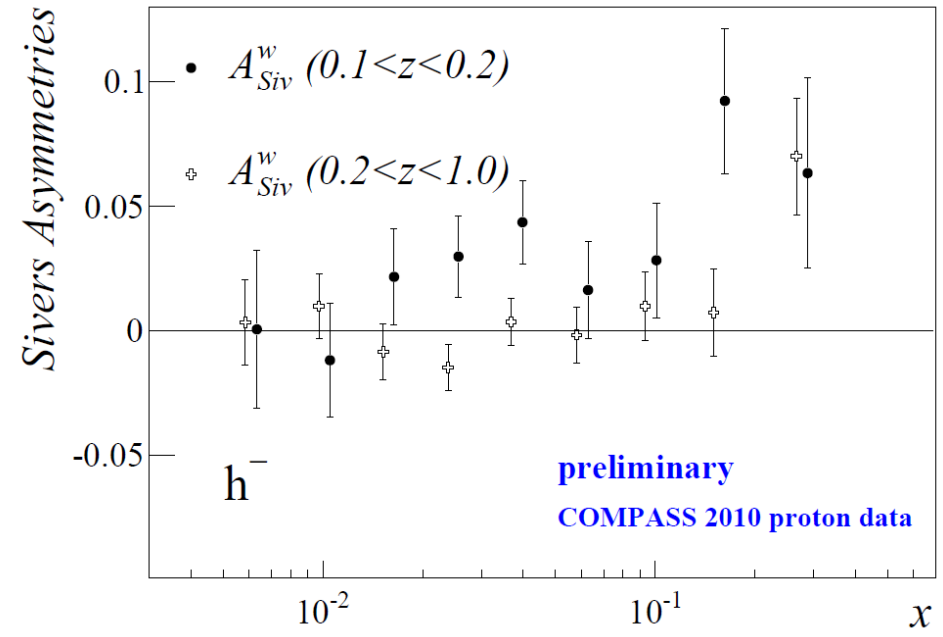
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different z ranges



$$\sim 2 \frac{f_{1T}^{\perp(1)u}(x)}{f_1^u(x)}$$

ok



both $f_{1T}^{\perp(1)u}$ and $f_{1T}^{\perp(1)d}$ contribute
the relative contribution depends on z

$0.1 < z < 0.2$ $h^+ \sim h^-$



P_T/zM weighted Sivers asymmetry

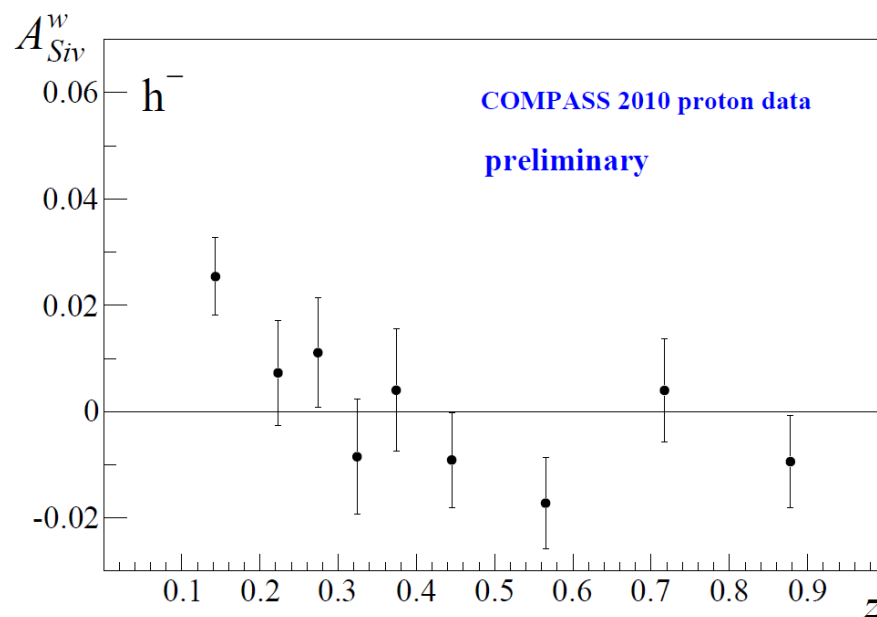
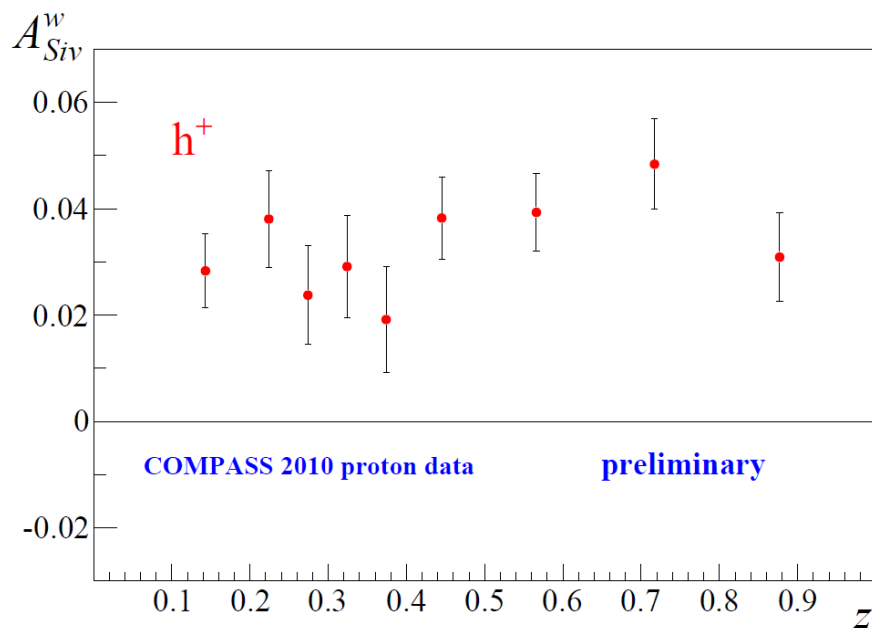
$$A_{Siv}^w(x) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) \int D_{1q}(z) dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz}$$

$$w = P_T/zM$$

standard cuts

arXiv:1702.00621

z dependence



$$\sim 2 \frac{\int C(x) f_{1T}^{\perp(1)q}(x) dx}{\int C(x) f_1^q(x) dx}$$

ok

both $f_{1T}^{\perp(1)u}$ and $f_{1T}^{\perp(1)d}$ contribute
the relative contribution depends on z

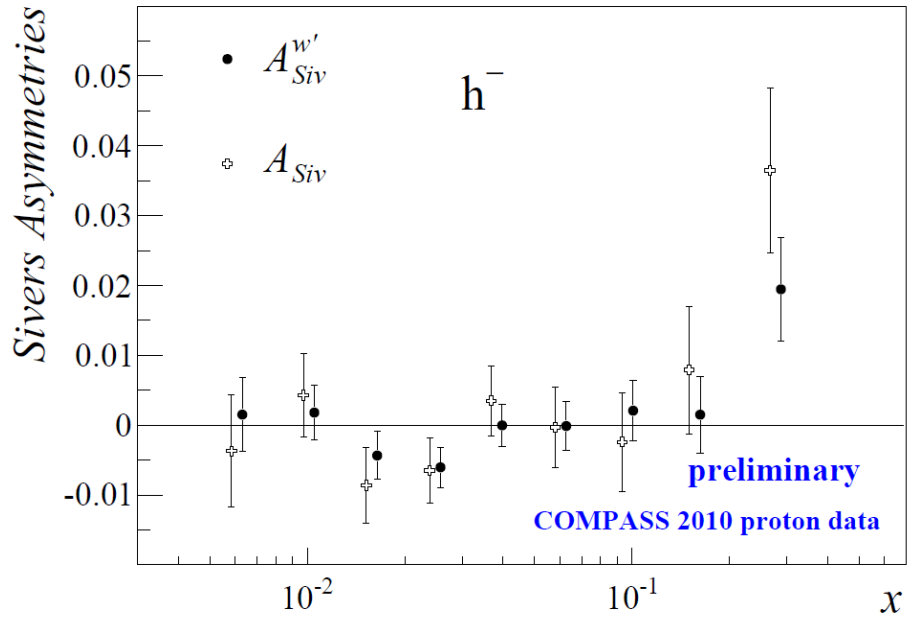
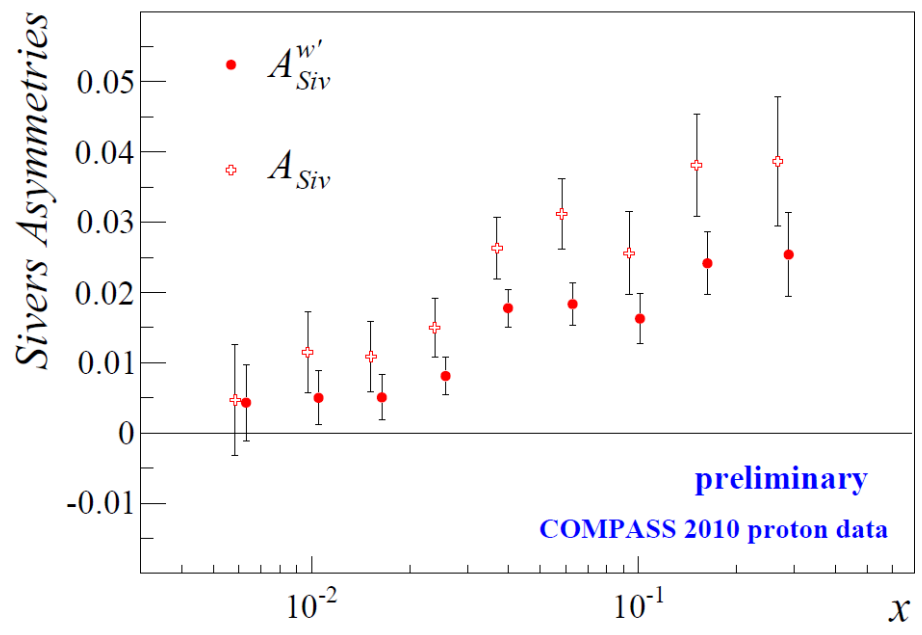


P_T/M weighted Siverts asymmetry

$$A_{Siv}^{w'}(x) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) \int z D_{1q}(z) dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz} \quad w' = P_T/M$$

standard cuts
 $z > 0.2$

comparison with the published "standard" asymmetries'



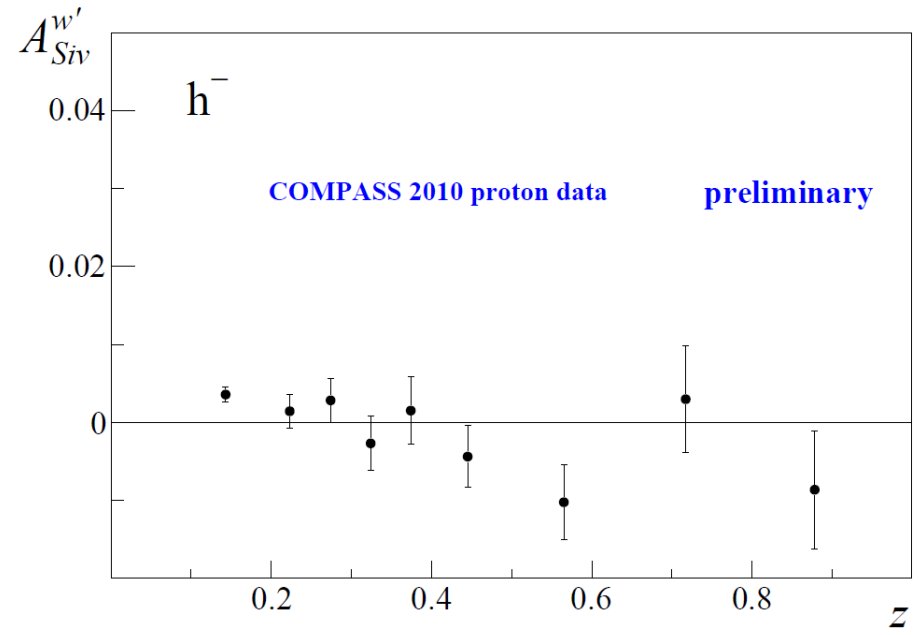
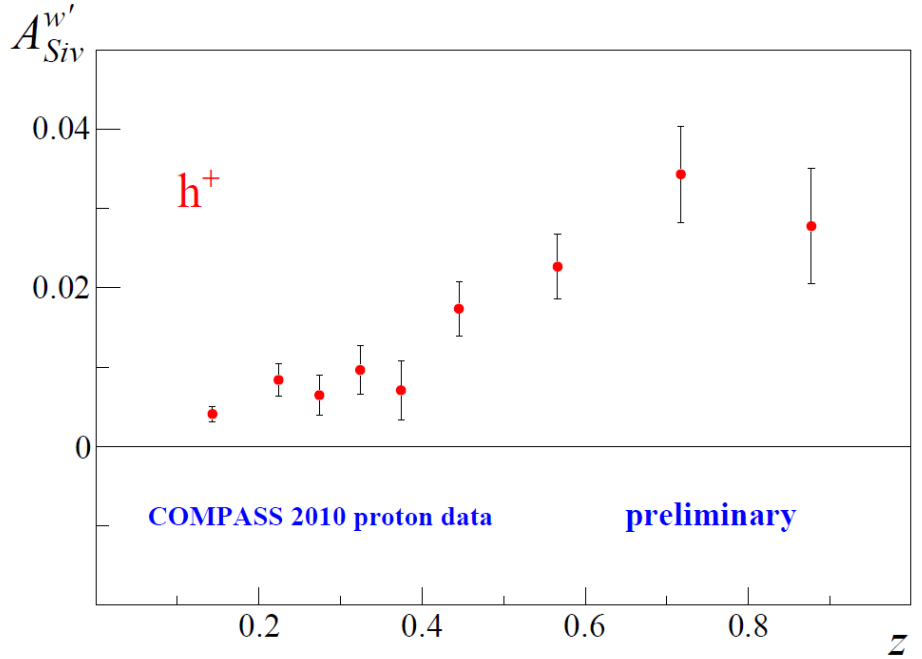


P_T/M weighted Sivers asymmetry

$$A_{Siv}^{w'}(z) = 2 \frac{\sum_q e_q^2 \int C(x) f_{1T}^{\perp(1)q}(x) dx}{\sum_q e_q^2 \int C(x) f_1^q(x) dx} z D_{1q}(z)$$

standard cuts
 $w' = P_T/M$

z dependence



$$\sim 2z \frac{\int C(x) f_{1T}^{\perp(1)q}(x) dx}{\int C(x) f_1^q(x) dx}$$

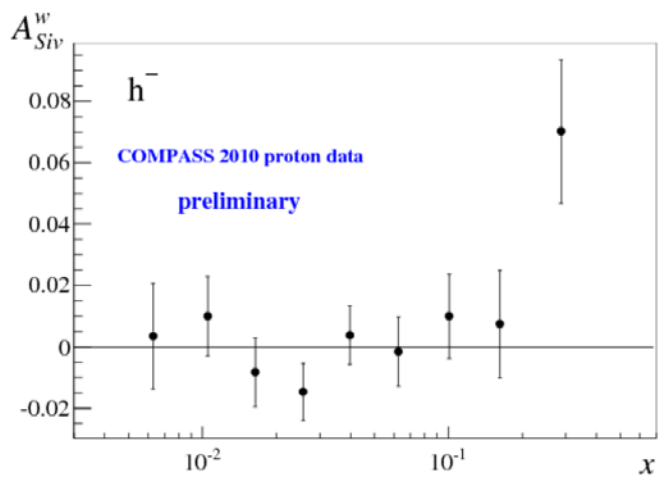
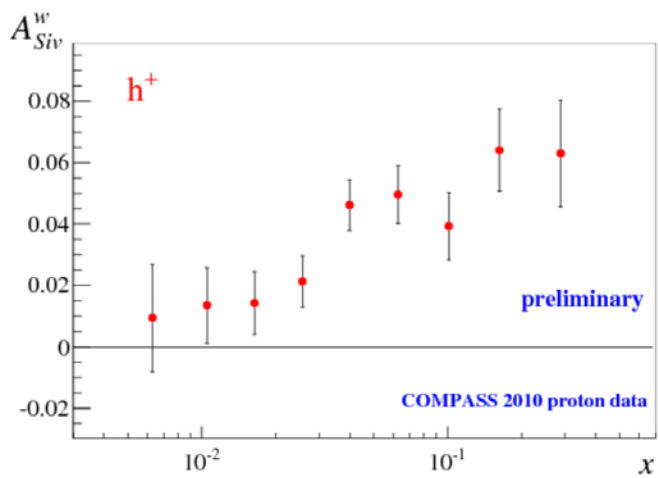
ok



1st moment of the Sivers functions

from P_T/zM weighted Sivers asymmetry

$$A_{Siv}^w(x) = 2 \frac{\sum_q e_q^2 x f_{1T}^{\perp(1)q}(x) \int D_{1q}(z) dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz}$$



used to extract the u- and d-quark Sivers functions

(J. Matousek talk)

assumptions:

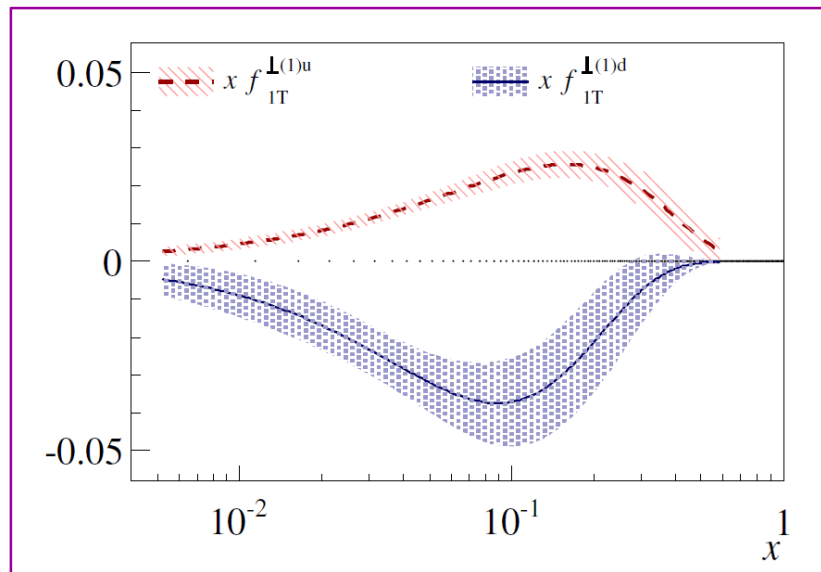
- the Sivers functions of all the sea quarks are zero

$$A_{Siv}^w(x) = 2 \frac{x f_{1T}^{\perp(1)u_v} \int D_{1u} dz + x f_{1T}^{\perp(1)d_v} \int D_{1d} dz}{\sum_q e_q^2 x f_1^q(x) \int D_{1q}(z) dz}$$

- $f_{1T}^{\perp(1)q} = a_q x^{b_q} (1-x)^{c_q}$ with $q = u_v, d_v$

1st moment of the Siverts functions

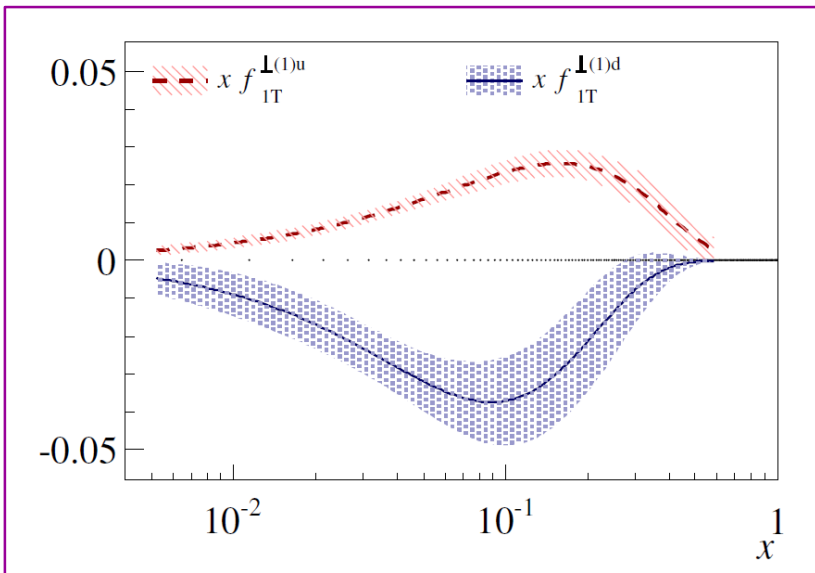
from P_T/zM weighted Siverts asymmetry



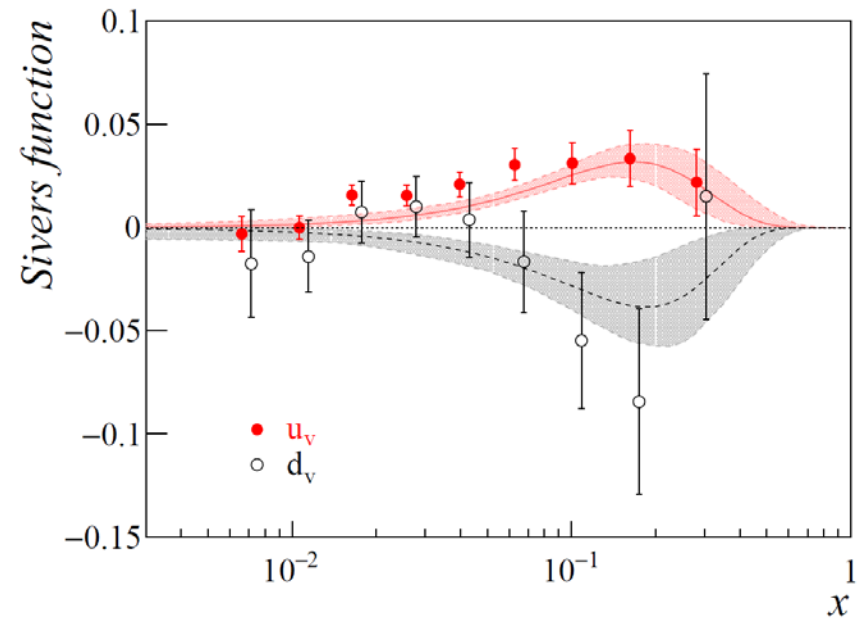
bands: 1 σ statistical errors only

1st moment of the Sivers functions

from P_T/zM weighted Sivers asymmetry



bands: 1σ statistical errors only



point by point extraction
COMPASS standard
Sivers asymmetries
p and d, pion and K
M B B 2017

curves: fit of COMPASS
and HERMES data,
Anselmino et al 2012



P_T weighted Sivers asymmetry

to summarise:

- **the results for the asymmetries for p look very interesting**
in particular
at first order no deviation from naïve expectation based on factorisation
- a first simple attempt to extract the **1st moment of the Siver functions** gives quite reasonable results
- **new d data needed also in this case to complete the exploratory COMPASS program**

**measurement of the
transversity transmitted
 Λ polarisation**

NEW!

Λ polarisation

the Λ polarisation can be measured from the angular distribution of the proton produced in the decay $\Lambda \rightarrow p \pi^-$

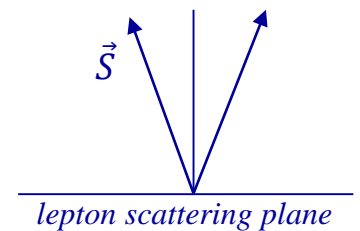
in the Λ c.m.s., the proton angular distribution is $\frac{dN}{d \cos\theta} \propto 1 + \alpha P_\Lambda \cos\theta$

where $\alpha = 0.642 \pm 0.013$ and

θ is the angle between the proton direction and the direction of the Λ polarisation

in SIDIS off transversely polarised nucleons, the Λ polarisation measured using the “reflected” direction of the nucleon spin can be written as

$$P_\Lambda = \frac{\sum_q e_q^2 h_1^q H_1^{\Lambda/q}}{\sum_q e_q^2 f_1^q D_1^{\Lambda/q}}$$



“transversity transmitted Λ polarisation”

the “transverse polarisation” is measured using as axis the normal to the lepton scattering plane

Λ polarisation

$$P_{\Lambda} = \frac{\sum_q e_q^2 h_1^q H_1^{\Lambda/q}}{\sum_q e_q^2 f_1^q D_1^{\Lambda/q}}$$

today transversity is somewhat known, while $H_1^{\Lambda/q}$ is completely unknown

with different assumptions,

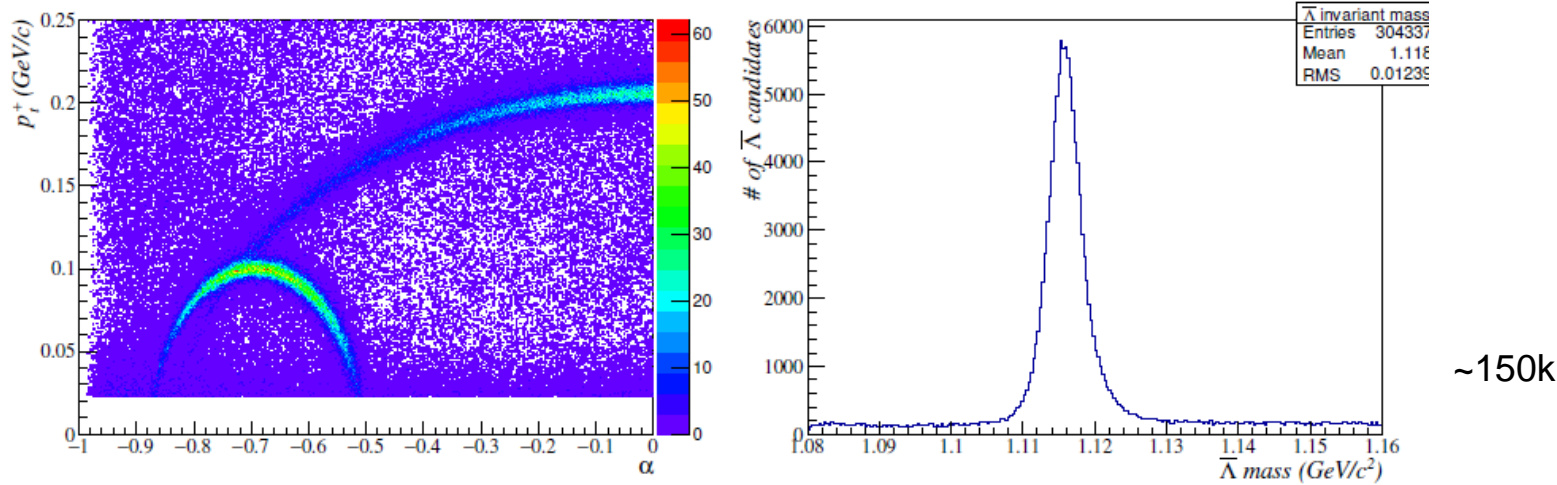
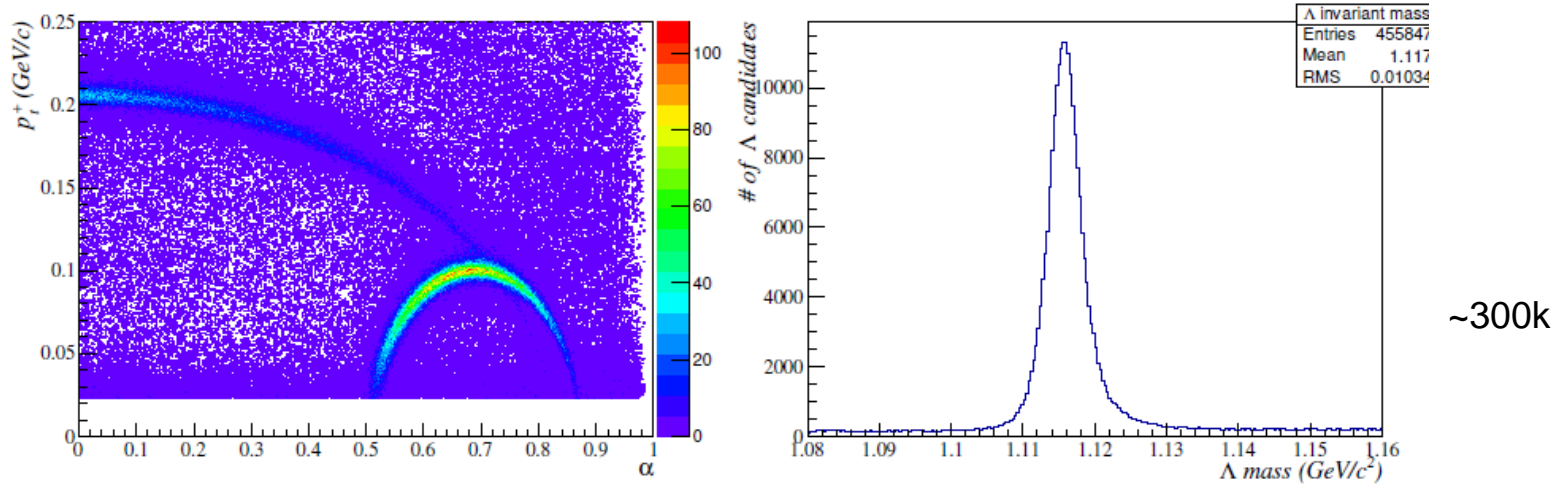
this measurement can give information either on h_1^s or on $H_1^{\Lambda/q} / D_1^{\Lambda/q}$

**COMPASS has measured the Λ and $\bar{\Lambda}$ polarisation
from the complete proton data set (2007 and 2010)**



$\Lambda/\bar{\Lambda}$ polarisation

$$Q^2 > 1 \text{ (GeV/c)}^2$$



$\Lambda/\bar{\Lambda}$ polarisation



the polarisation has been measured in x , z and P_T bins
from the complete p data set (2007 and 2010)

for all Λ and $\bar{\Lambda}$ candidates and

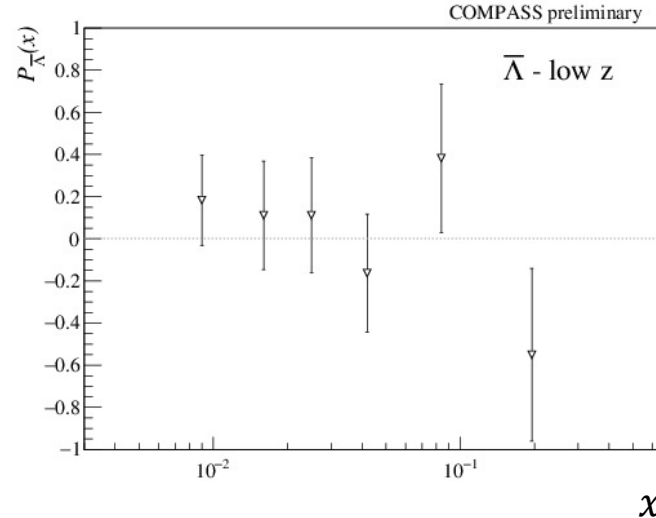
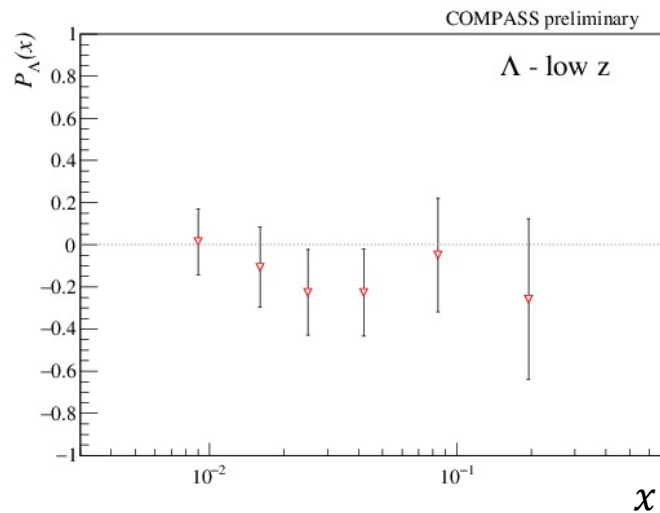
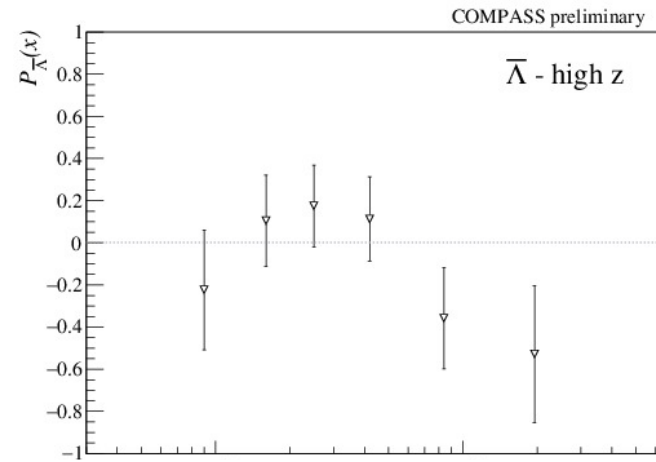
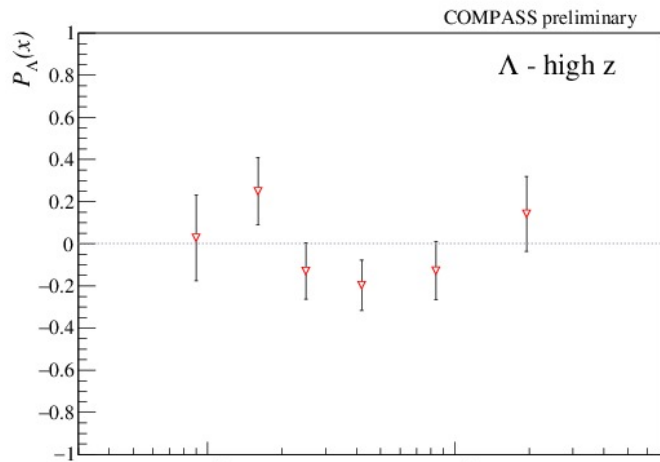
for

- $z > 0.2$ *and* $x_F > 0$ high z region
- $z > 0.2$ *or* $x_F > 0$ low z region

- $x > 0.032$ high x region
- $x < 0.032$ low x region

- $P_T > 0.5$ GeV/c high P_T region
- $P_T < 0.5$ GeV/c low P_T region

$\Lambda/\bar{\Lambda}$ polarisation



statistically limited
still the only existing measurement
interpretation work ongoing in COMPASS

COMPASS

Common Muon and Proton Apparatus for Structure and Spectroscopy

Long-Term plans



1. **COMPASS QCD facility**
2. **Beyond 2020 Workshop (March 2016)**
3. **Long term plans**
 - RF separated beam
 - Spectroscopy
 - Drell-Yan
 - Exclusive measurements with muon and hadron beams
4. **Shorter term plans**
 - SIDIS
 - Drell-Yan
 - Astrophysics
5. **Summary**

→ Lol

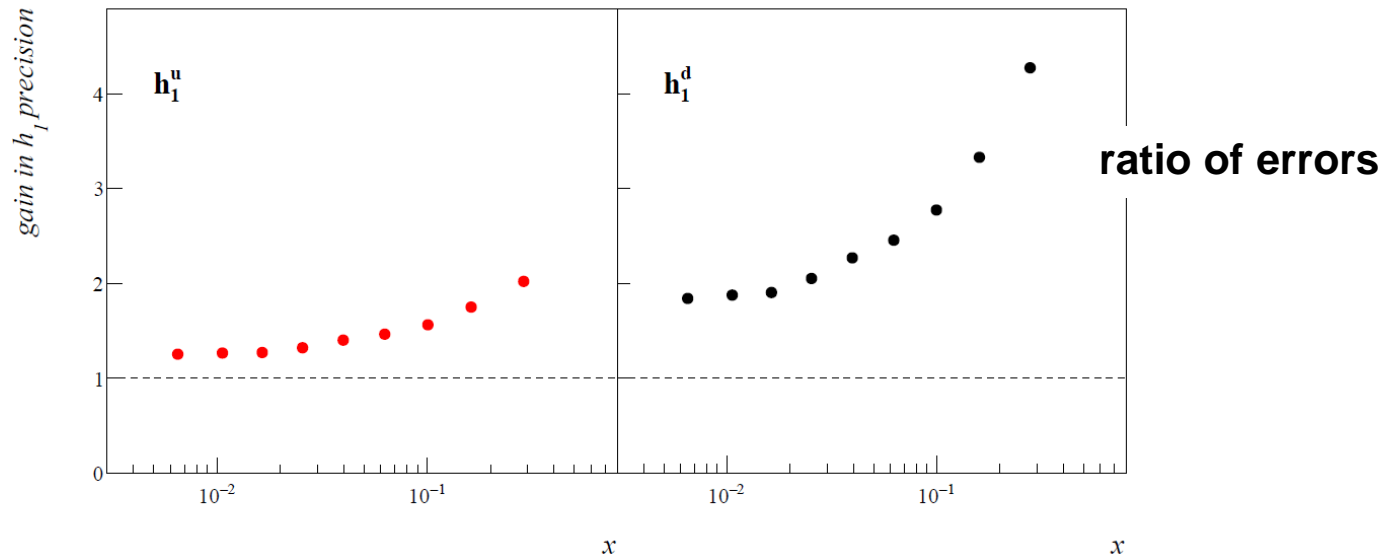
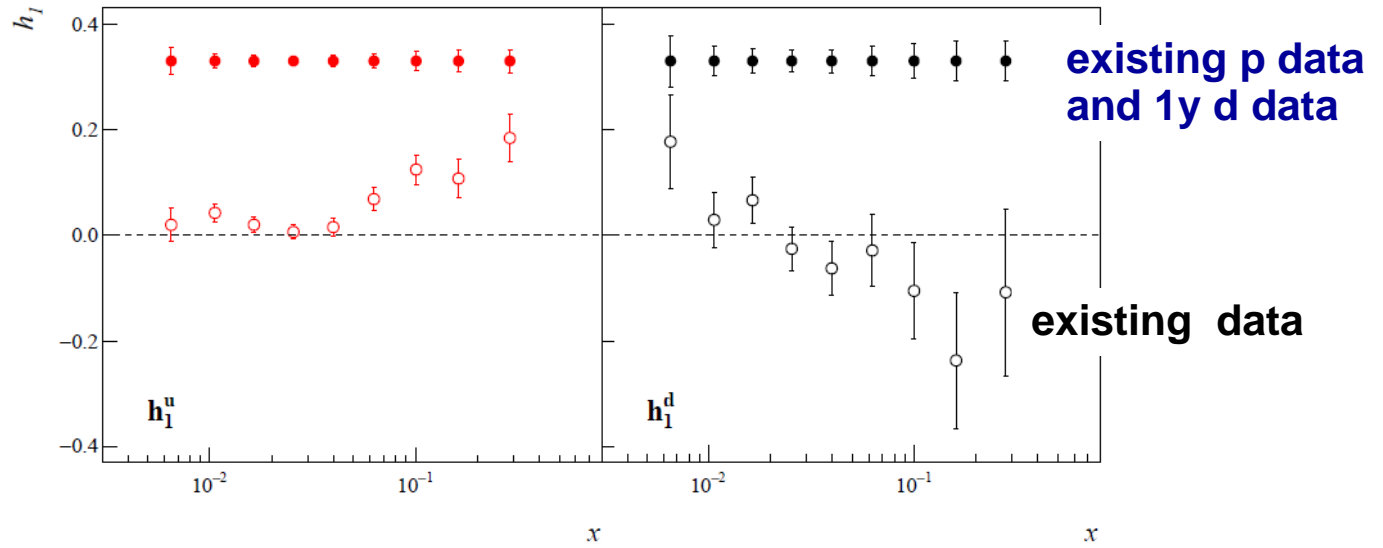


Oleg Denisov INFN(Torino)/CERN for the COMPASS Coll.

latest developments: a proposal for 1. a one year measurement of SIDIS on transversely polarised deuteron



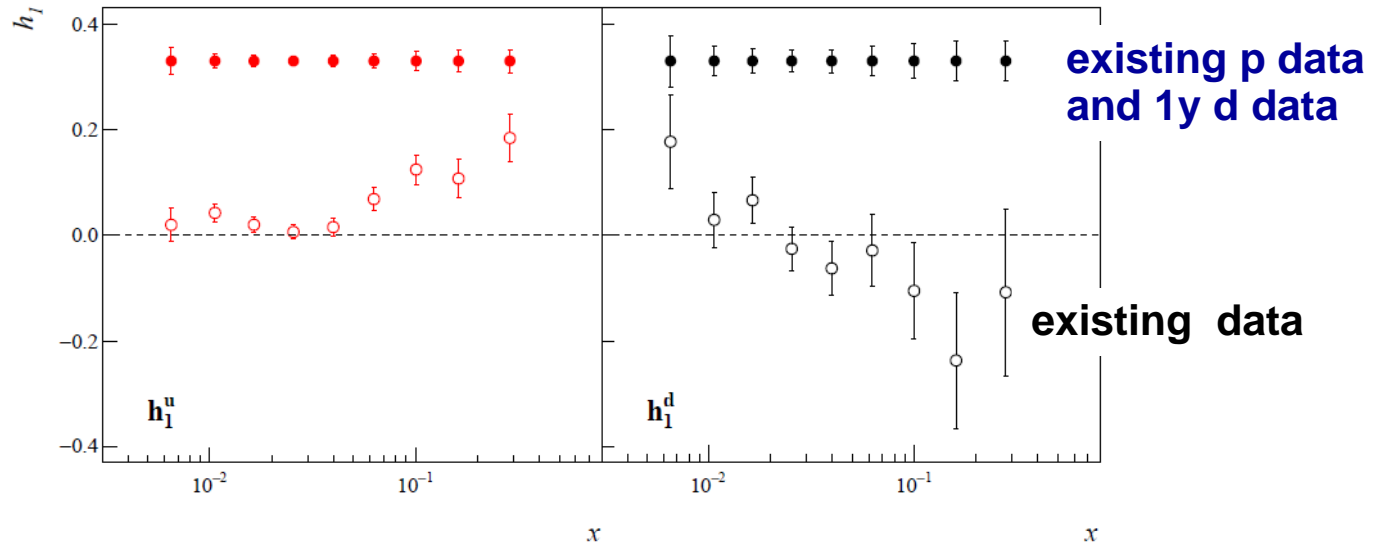
u and d quark
Transversity PDFs



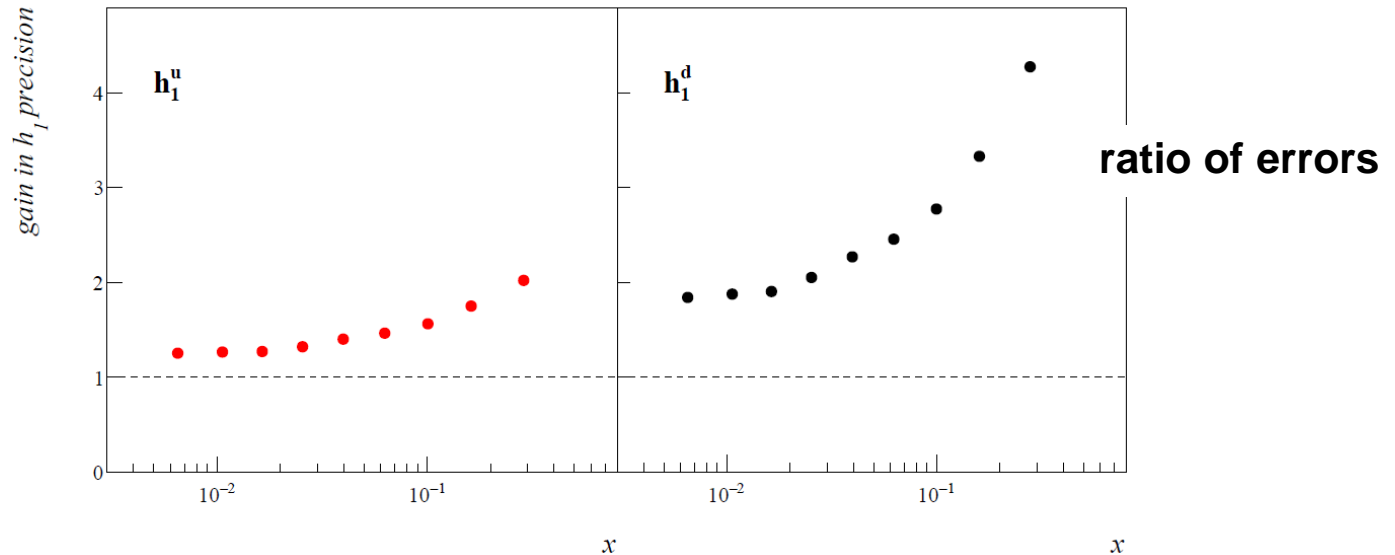
latest developments: a proposal for 1. a one year measurement of SIDIS on transversely polarised deuteron



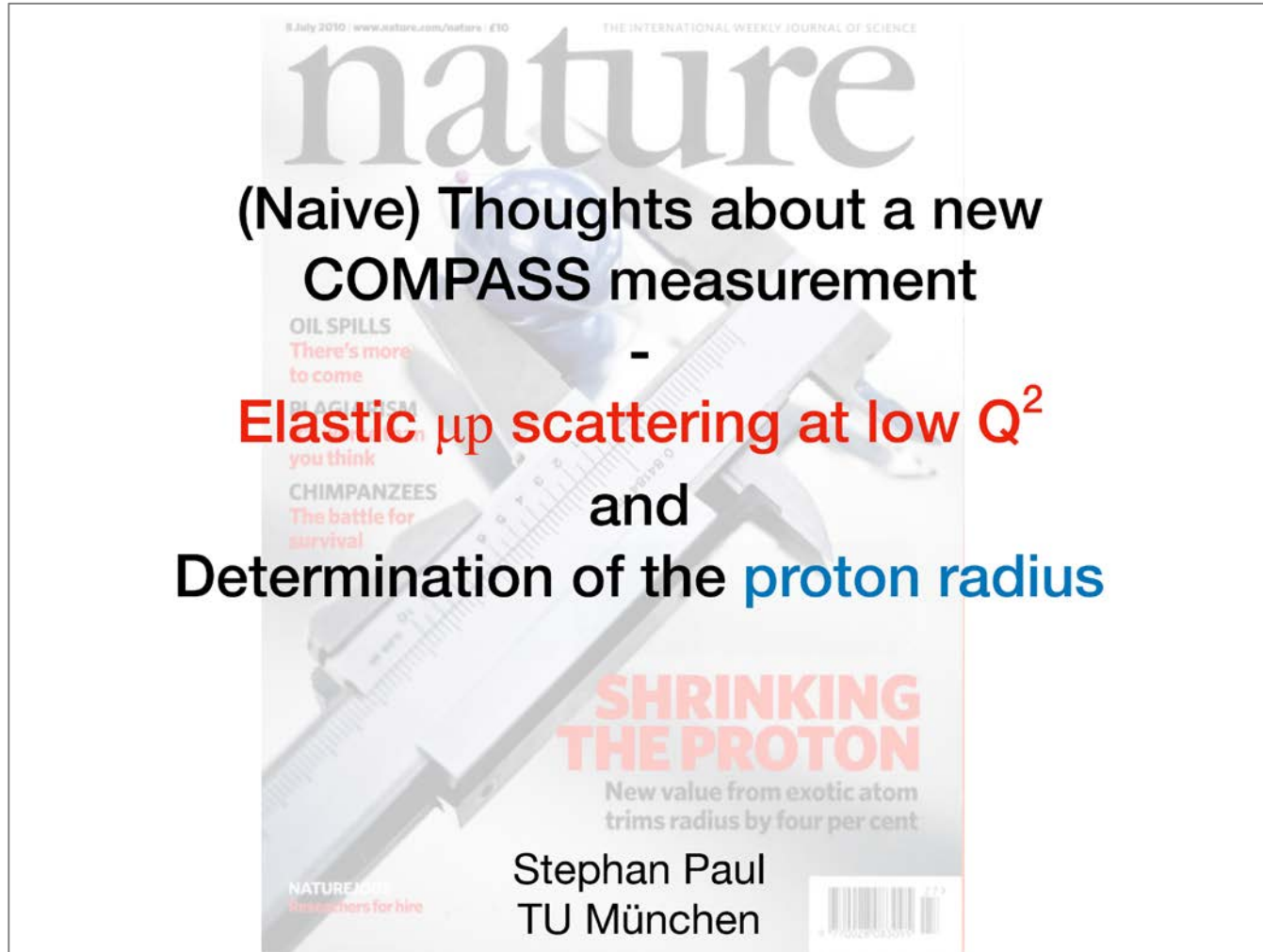
u and d quark
Transversity PDFs



needed to complete
the COMPASS
transverse
spin program



latest developments: a proposal for
2. a measurement of μp elastic scattering



in parallel with an Lol for the longer term projects