

# Longitudinal Spin Structure of the Nucleon COMPASS Legacy

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# COMPASS @ *DSPIN-17*

- Longitudinal spin structure
  - Artem Ivanov: Tuesday AM  
Longitudinal target spin dependent azimuthal asymmetries in SIDIS
- Fragmentation
  - Nikolai Mitrofanov: Tuesday AM  
Multiplicities of charged hadrons, pions and kaons in DIS
- TMD
  - Franco Bradamante: TSAs in SIDIS
  - Michael Pešek: TSAs in Drell-Yan
  - Jan Matoušek:  $q_T$ -weighted TSAs in Drell-Yan
- GPD
  - Andrzej Sandacz: DVCS and exclusive  $\pi^0$
  - Bohdan Marianski TSAs in exclusive vector meson production

# Partonic structure of the nucleon

- Twist-2 PDFs, integrated over  $k_T$  ( $k_T = \text{parton intrinsic transverse momentum}$ )  
⇒ Three parton distributions:
  - Parton momentum DF  $q(x)$
  - $\Delta q(x)$  for a longitudinally polarised Nucleon via Structure Function  $g_1(x, Q^2)$
  - $\Delta_T q(x)$  for a transversally polarised Nucleon
  
- $\Delta_T q(x)$ : See talk by F. Bradamante at this conference
  
- Finite  $A_{UL}^{\sin\phi_R}$  in LSA of di-hadron production, sensitive to collinear twist-3 PDF and FF  
See talk by A. Ivanov at this conference

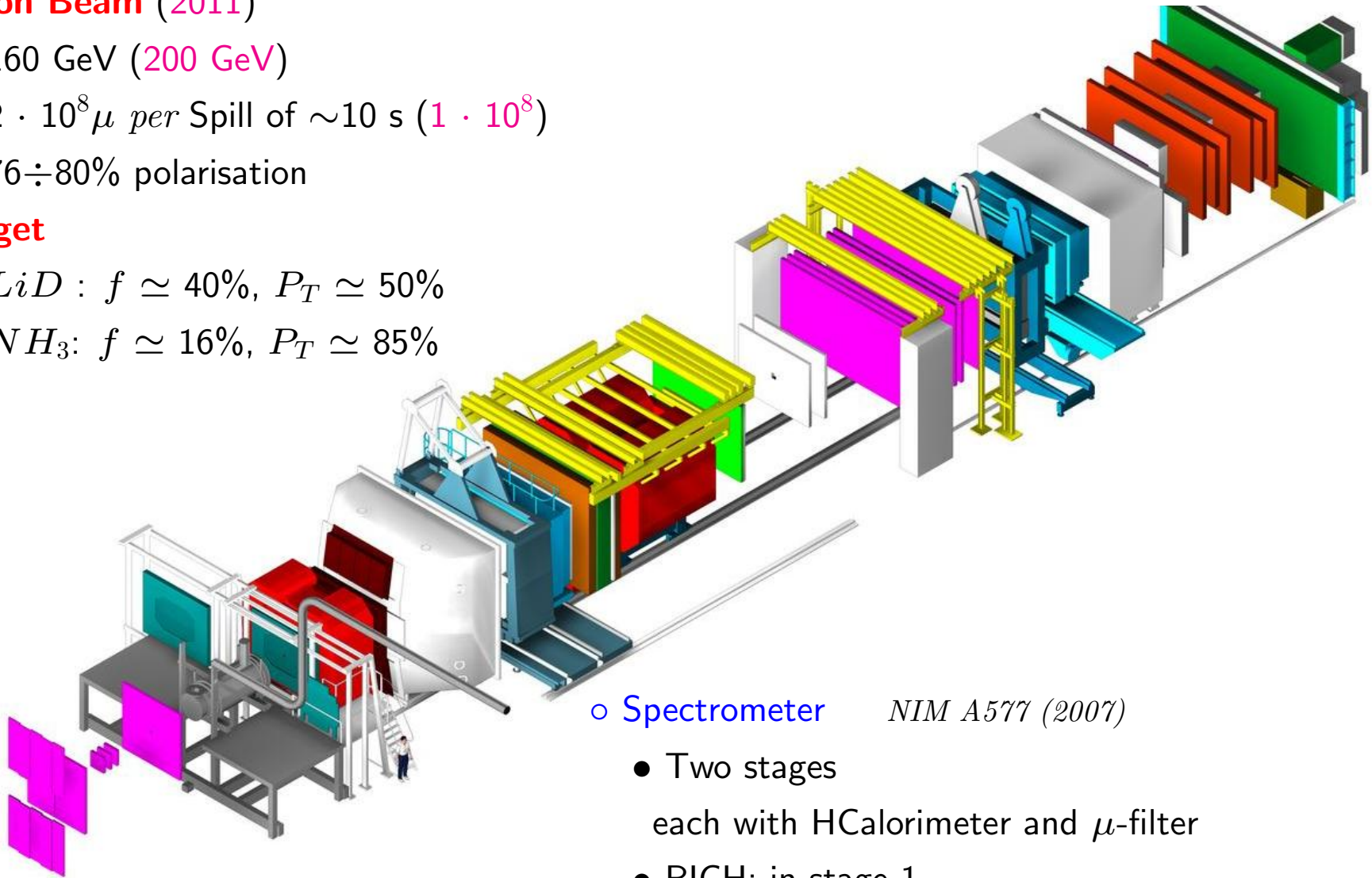
# COMPASS: Spectrometer

## ○ Muon Beam (2011)

- 160 GeV (200 GeV)
- $2 \cdot 10^8 \mu$  per Spill of  $\sim 10$  s ( $1 \cdot 10^8$ )
- 76÷80% polarisation

## ○ Target

- $LiD$  :  $f \simeq 40\%$ ,  $P_T \simeq 50\%$
- $NH_3$ :  $f \simeq 16\%$ ,  $P_T \simeq 85\%$



## ○ Spectrometer *NIM A577 (2007)*

- Two stages  
each with HCalorimeter and  $\mu$ -filter
- RICH: in stage 1
- ECalorimeters(1&2): since 2008

## Double Spin Asymmetry Measurement

- **Simultaneous** recording of the two spin states in oppositely polarised target cells. . .

. . . 2 cells:  $1/2 \uparrow 1/2 \downarrow \iff 8 \text{ hours} \implies 1/2 \downarrow 1/2 \uparrow$

. . . 3 cells:  $1/3 \uparrow 2/3 \downarrow 1/3 \uparrow \iff 24 \text{ hours} \implies 1/3 \downarrow 2/3 \uparrow 1/3 \downarrow$

- Reversal by **target-field** rotation to cancel acceptance diff

$$\frac{A^{\parallel}}{D} = \frac{1}{|P_{\mu}P_T|fD} \frac{1}{2} \left( \frac{N^{\uparrow\downarrow} - N^{\uparrow\uparrow}}{N^{\uparrow\downarrow} + N^{\uparrow\uparrow}} + \frac{N^{\uparrow\downarrow} - N^{\uparrow\uparrow}}{N^{\uparrow\downarrow} + N^{\uparrow\uparrow}} \right) \quad D = \text{Depolarisation factor}$$

$LiD: P_{\mu} \times D \times P_T \times f \simeq 80\% \times 60\% \times 50\% \times 40\% \simeq 10\%$  (typical values)

$NH_3: \dots P_T \times f \simeq \dots 85\% \times 16\% \simeq 6\%$

- Reversal *via* re-polarisation once per year to cancel **target-field**/acceptance correlation
- $f$  corrected for spin-independent radiative processes: TERAD
- $A_1 \simeq A^{\parallel}/D$  corrected for spin-dependent radiative processes: POLRAD

## Method of extraction

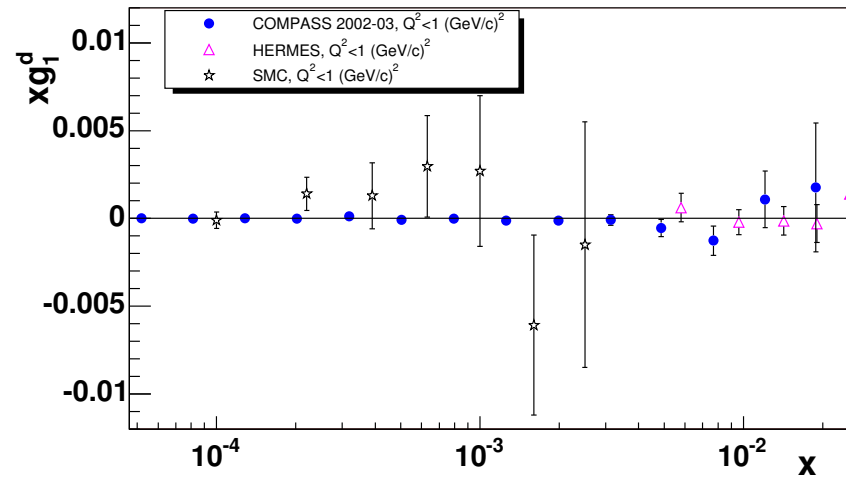
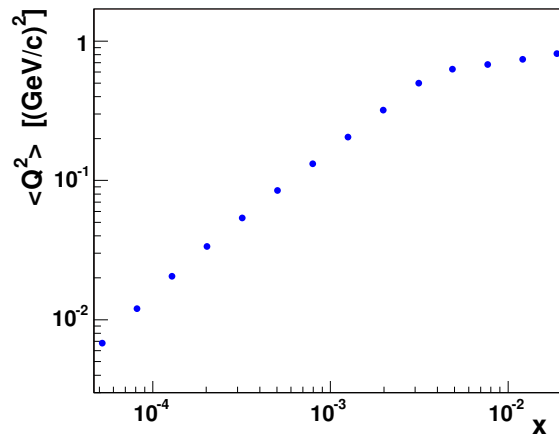
- $g_1^{\text{p,d}}(x, Q^2)$  is extracted as:

$$g_1^{\text{p,d}}(x, Q^2) = A_1^{\text{p,d}}(x, Q^2) \frac{F_2^{\text{p,d}}(x, Q^2)}{2x(1 + R(x, Q^2))}$$

- Input  $F_2^{\text{p,d}}(x, Q^2)$ : from NMC parametrisation at  $Q^2 > 0.2 \text{ GeV}^2$  (*PRD 58 (1998) 112001*)  
at  $Q^2 < 0.2 \text{ GeV}^2$ : from the model based on the GVMD concept (*PLB 295 (1992) 263*)
- $R(x, Q^2)$ : R1998 from SLAC (*PLB 452 (1999) 194*)  
suitably extended to low  $Q^2$ , (*PLB 647 (2007) 330*)
- Maximum systematic uncertainties on  $A_1$ 
  - At low  $Q^2$ : from  $D(R)$ ,  $\sim(0.01 - 0.39)D$  (mult.)  
and false asym.,  $1.5\sigma_{\text{stat.}}$  (additive)
  - At high  $Q^2$ : from  $P_B, P_T$ ,  $\sim 0.05P_B, P_T$  (mult.)  
and false asym.,  $< 0.8\sigma_{\text{stat.}}$  (additive)

## $g_1^d$ at low $Q^2$

- COMPASS Coll., Phys. Lett. B **647** (2007)



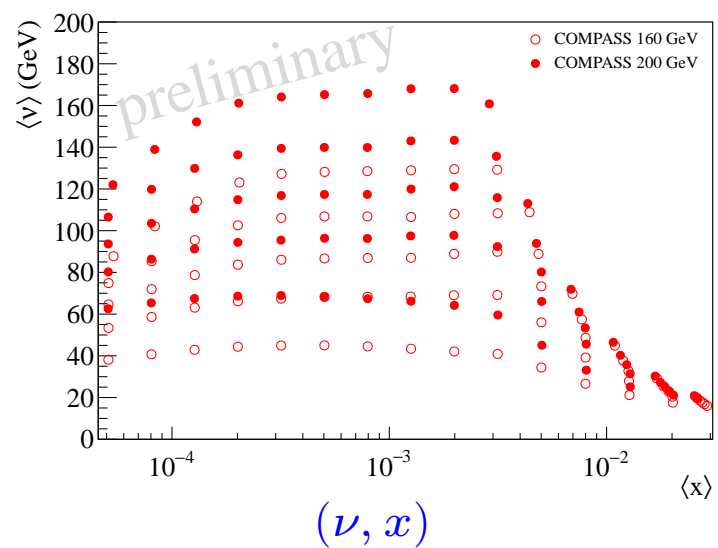
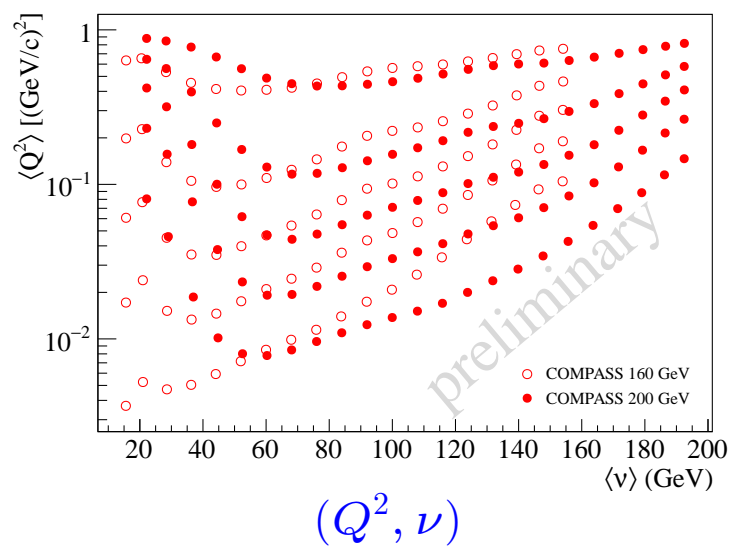
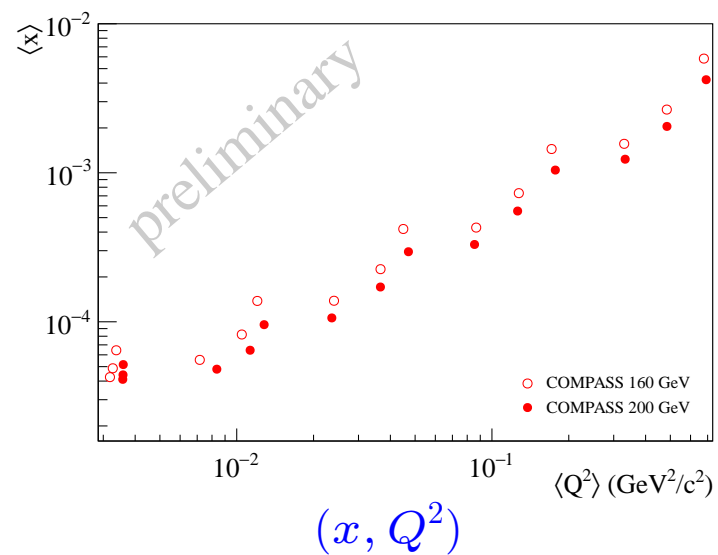
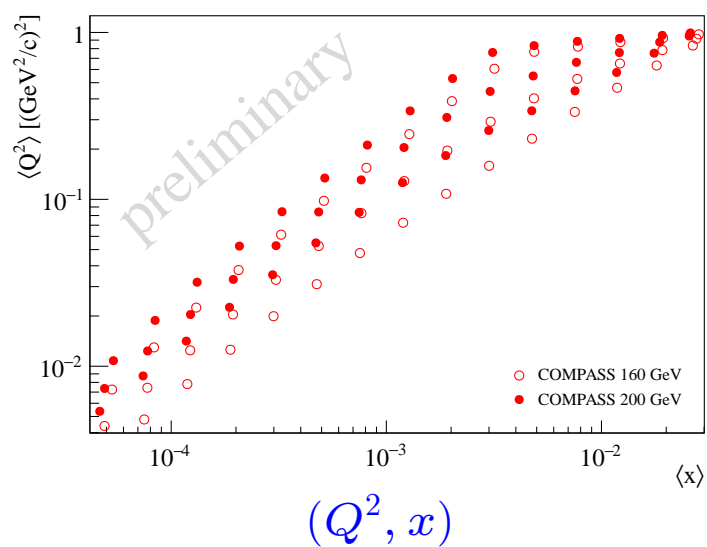
- More than ten-fold improvement over SMC
- Consistent with 0 for  $5 \times 10^{-5} < x \lesssim 3 \times 10^{-2}$
- Tend to indicate that  $\Delta G$  is small. Would then rule out large  $\Delta G$  solution to the spin crisis,

$$\text{via: } \Sigma = \Sigma_{naive\,QM} - N_f \alpha_S(Q^2) / 2\pi \Delta G(Q^2)$$

(see *e.g.* A. Thomas, Int. J. Mod. Phys. E **18** (2009))

# $g_1^p$ at low $Q^2$

- Kinematic coverage:

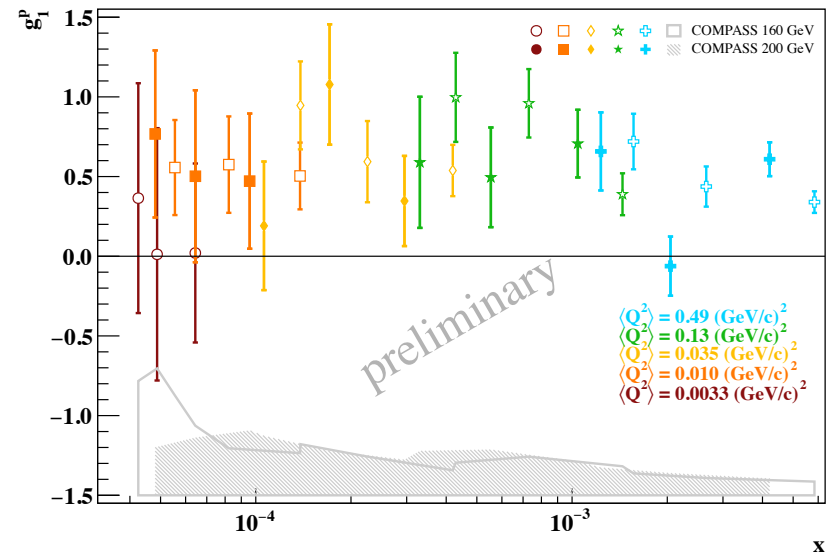
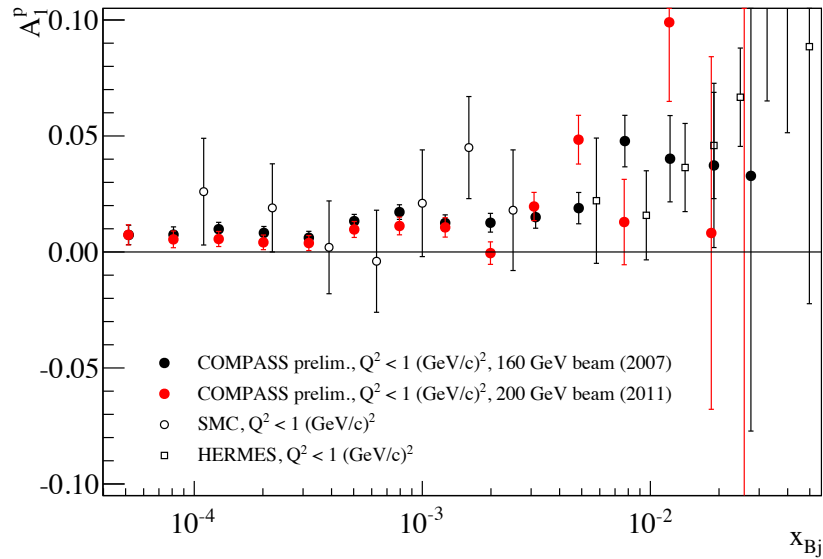




## $g_1^p$ at low $Q^2$ : data sample

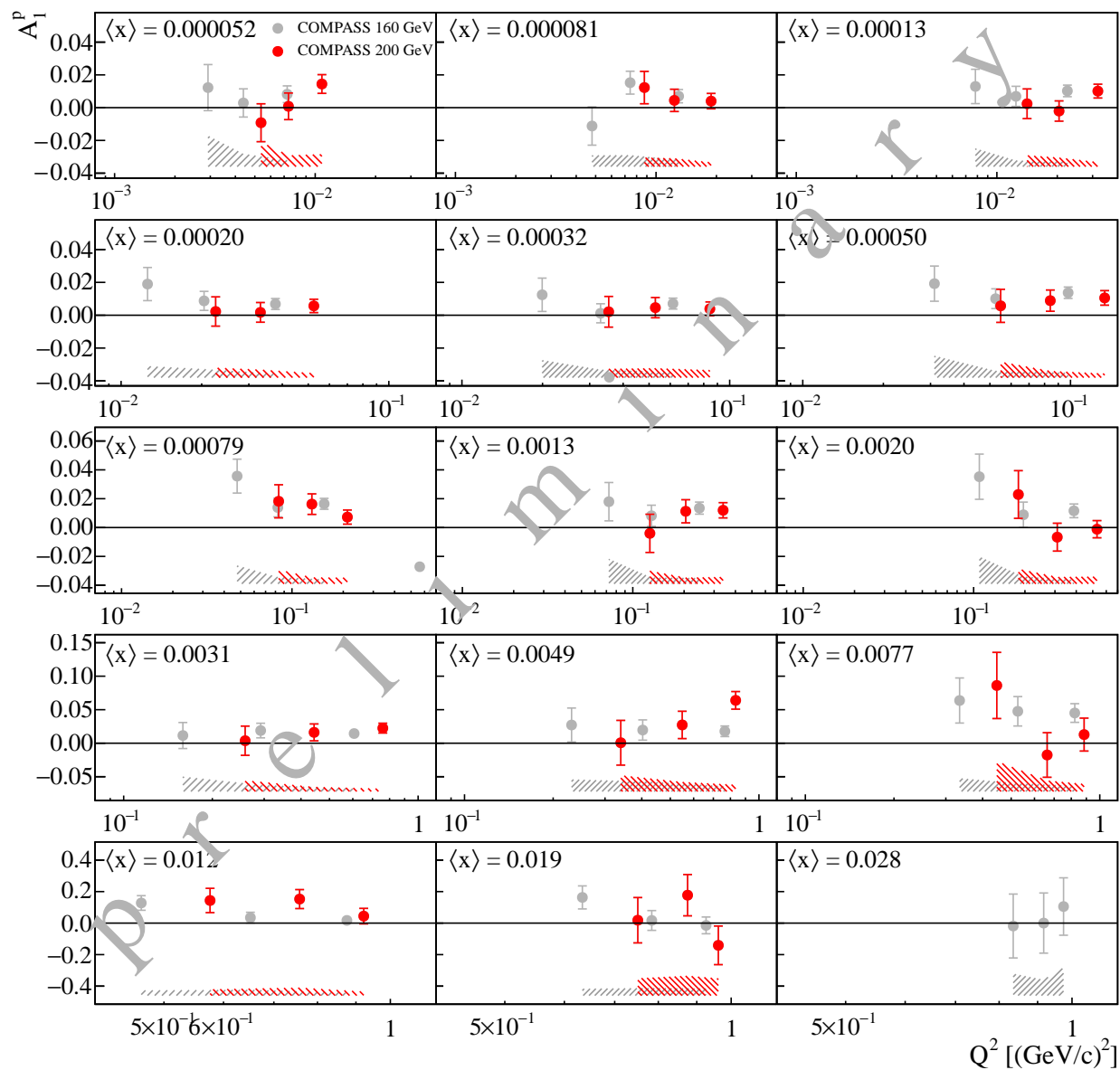
- **Reconstructed event**: Small scattering angle
  - ⇒ No way to be strictly inclusive and stay within polarised target cells
  - ⇒ Require one extra outgoing charged particle
  - Doesn't bias asymmetry at low  $x$  (see SMC Coll., Phys. Rev. D **58** (1998) 112001)
- Kinematic cuts
  - $x > 4 \times 10^{-5}$  (at smaller  $x$  uncertainties become too large)
  - $0.1 < y < 0.9$  (bad reconstruction @ low  $y$ , large rad. corr. at high  $y$ )
  - Above cuts result in:  $W \gtrsim 5$  GeV
- Removed  $\mu e \rightarrow \mu e$  events, around  $x = m_e/M_p = 5.45 \times 10^{-4}$
- Final sample: 447 Mevnts @ 160 GeV + 229 Mevents @ 200 GeV
  - ~ 150 × SMC statistics

## Low $Q^2$ : results on $A_1^P(x, Q^2)$



- More than ten-fold improvement over the statistical precision of SMC.
- Clear spin effects at low  $x$  and  $Q^2$

# Low $Q^2$ : results on $A_1^P(x, Q^2)$



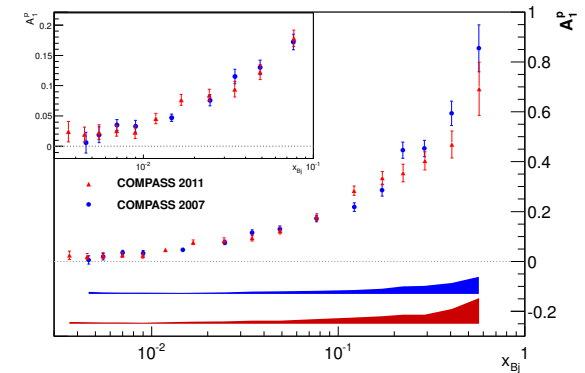
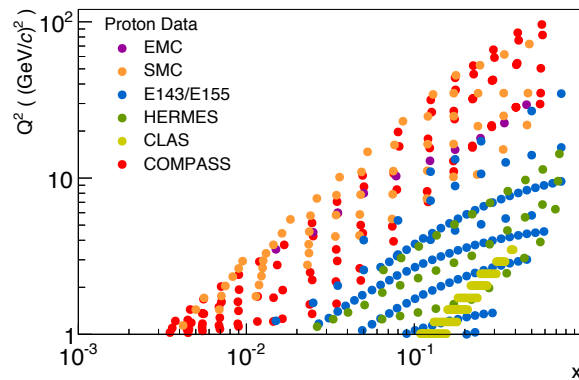
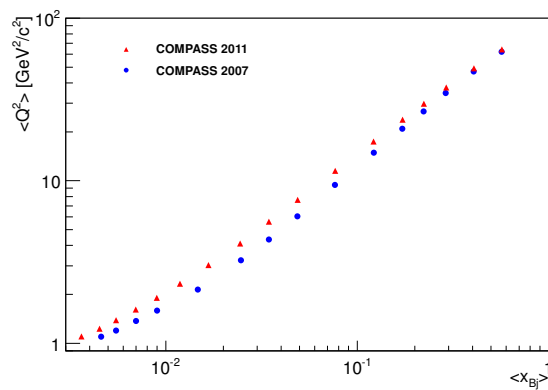
## Importance of $g_1^P$ at low $x$ , $Q^2$ : results on $A_1^P(x, Q^2)$

- At low  $x$ : strong increase of parton density with decreasing  $x$   
⇒ parton recombination effects
- BUT: in fixed-target experiments low  $x$  correlated with low  $Q^2$  region  
⇒ non-perturbative effects must be considered
- Attempts to describe that region phenomenologically:
  - B. Badelek, J. Koryluk, J. Kwiecinski, Phys. Rev. D**61** (2000) 014009.
  - B. Badelek, J. Kwiecinski, B. Ziaja, Eur. Phys. J. C**26** (2002) 45.
  - B.I. Ermolaev, M. Greco, S.I. Troyan, Eur. Phys. J. C**50** (2007) 823;  
*ibid.* C**51** (2007) 859.
  - W. Zhu, J. Ruan, Int. J. Mod. Phys. E **24** (2015) 1550077.
- $g_1$  at low  $Q^2$  needed for Radiative Corrections

## $g_1^p$ at high $Q^2$ : data samples

- Measurements at incident  $E = 200$  GeV: Phys. Lett. B **753** (2016) 18
- Complementing earlier data at  $E = 160$  GeV: (PLB 690 (2010) 466)

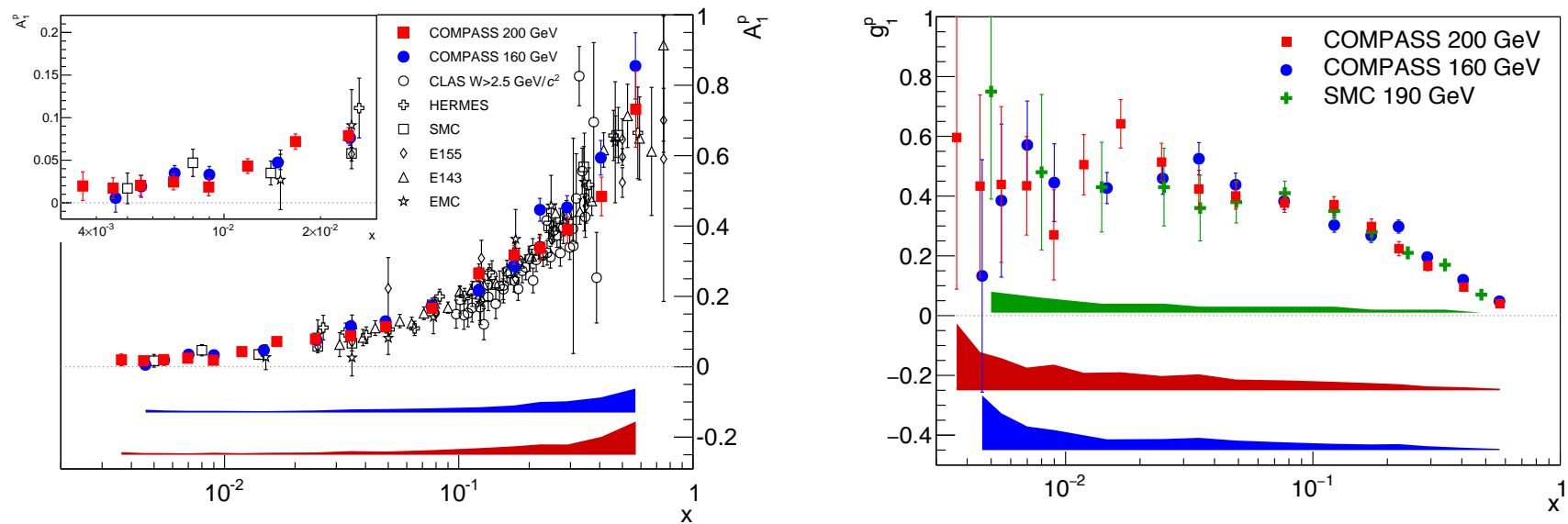
Reaching lower  $x$  and higher  $Q^2$



- Results on  $A_1^p$  at both energies agree well

# Results on $A_1^P$ and $g_1^P$

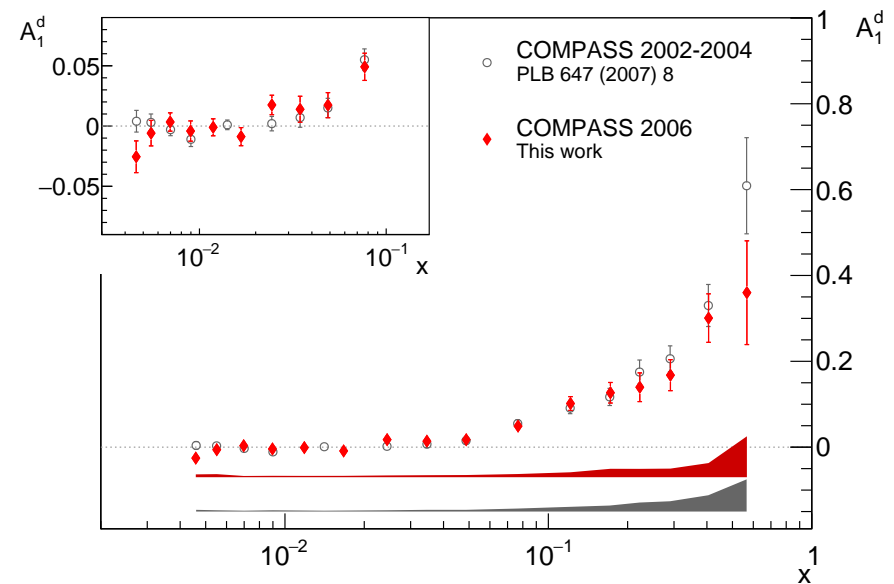
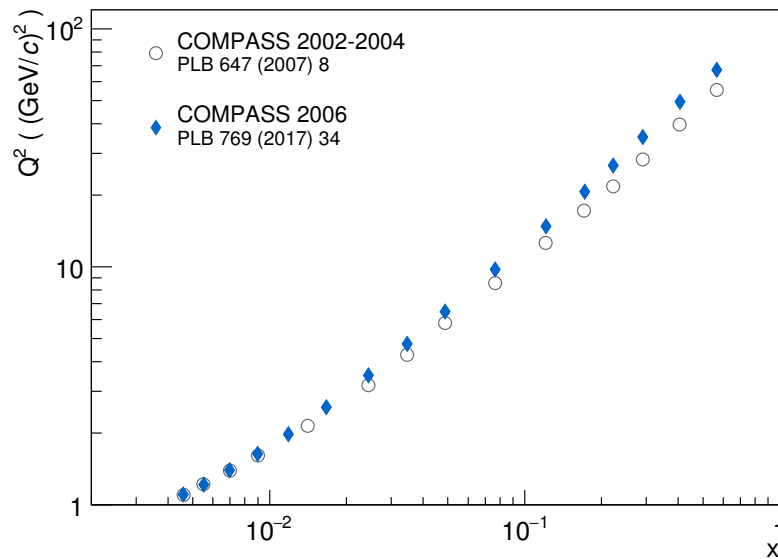
- $A_1^P$  and  $g_1^P$  shown at the measured values of  $Q^2$



- Good agreement with world data
- $g_1^P$  clearly positive at lowest measured values of  $x$

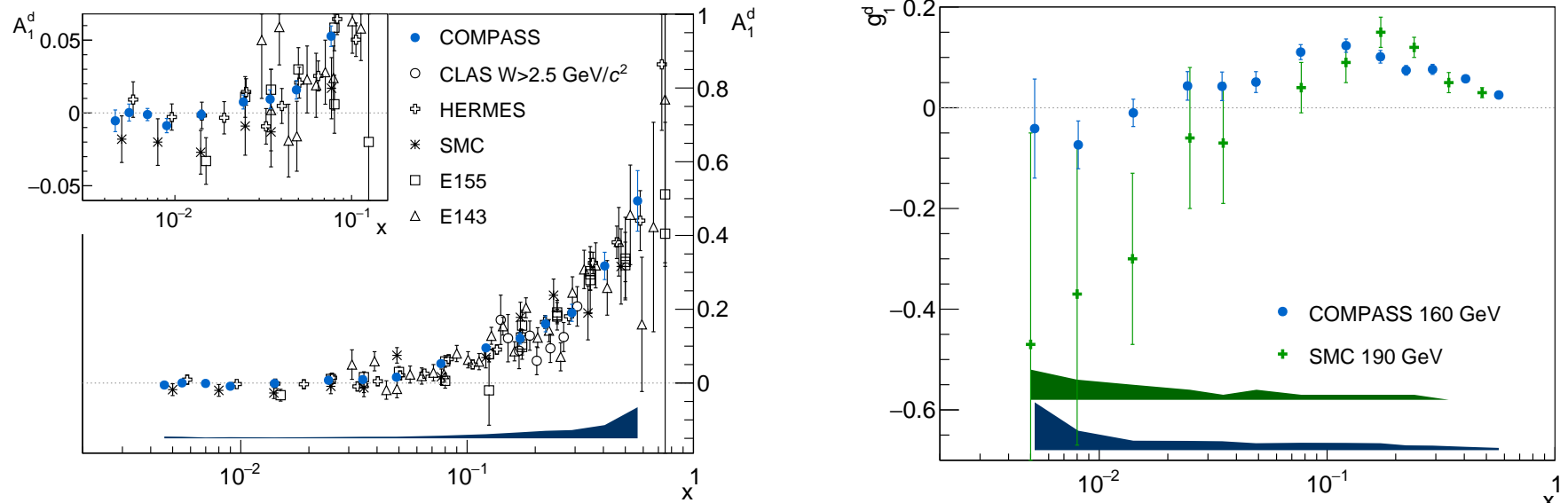
## $g_1^d$ at high $Q^2$ : data samples

- Combined data, all @ 160 GeV: Phys. Lett. B **769** (2017) 34
- Results on  $A_1^d$  agree very well



## Results on $A_1^d$ and $g_1^d$

- $A_1^d$  and  $g_1^d$  shown at the measured values

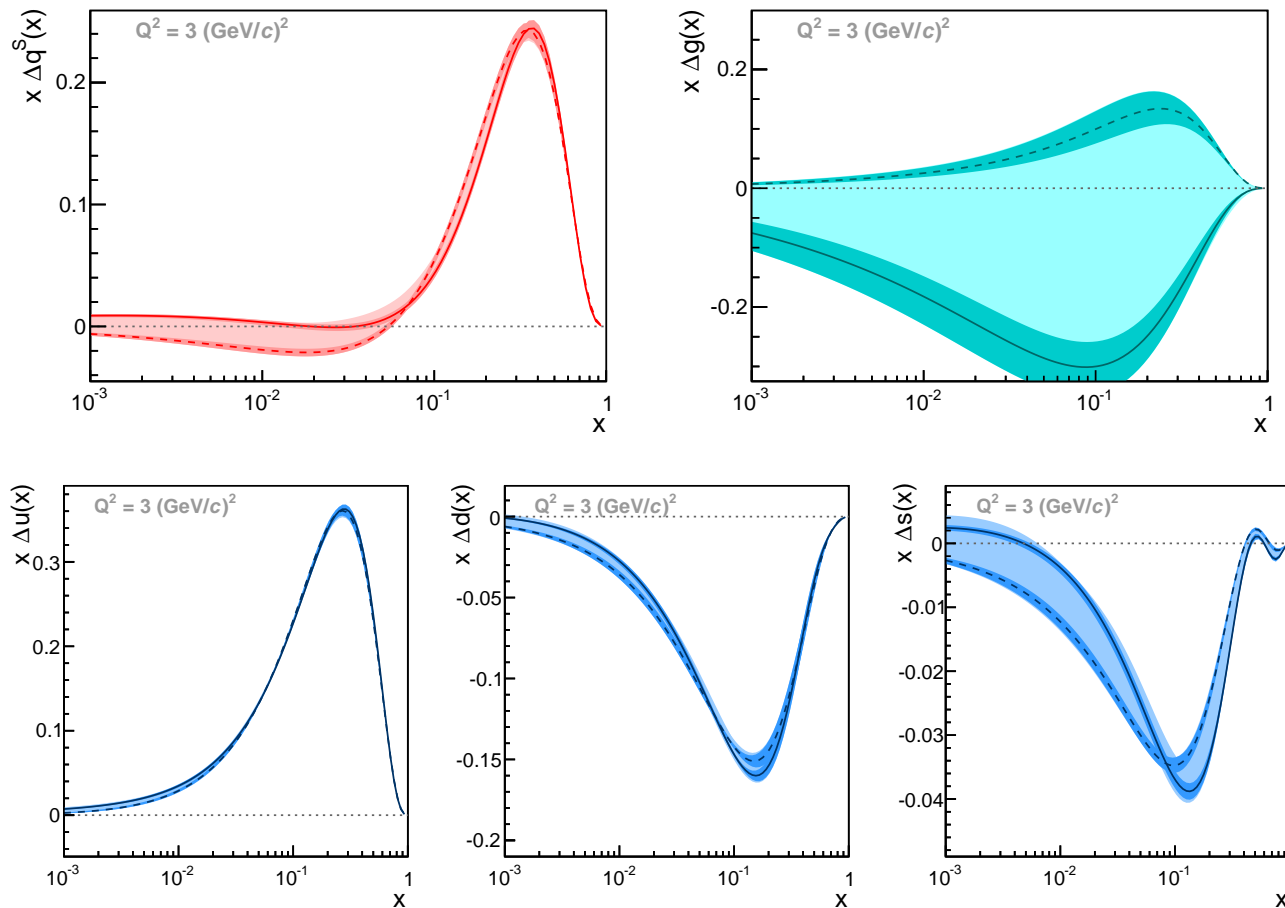


- Good agreement of  $A_1^d$  and  $g_1^d$  with world data
- $g_1^d$  compatible with zero at lowest measured values of  $x$ ,  
contrary to hints from SMC





## NLO QCD fit: results



- Statistical uncertainties (dark bands)  $\ll$  systematic (light bands)
- Gluon polarisation poorly constrained  $\Rightarrow$  Need “direct” methods
- Quark spin contribution to the nucleon spin:  $0.26 < \Delta\Sigma < 0.36$  (due to poor  $\Delta G$ )

## First moments of $g_1$ and singlet axial charge $a_0$

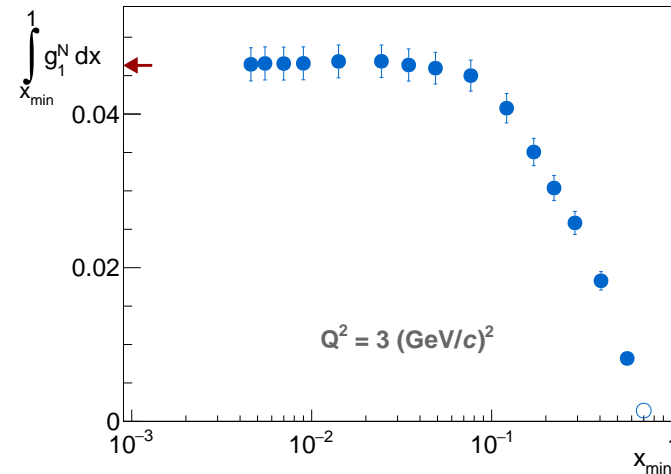
- First moments  $\Gamma_1^p, \Gamma_1^d, \Gamma_1^N$

$$\text{where } \Gamma_1^i(Q^2) = \int_0^1 g_1^i(x, Q^2) dx$$

- In particular:

$$\Gamma_1^N(Q^2) = 1/36 [4a_0 C_S(Q^2) + a_8 C_{NS}(Q^2)]$$

$$= \int_0^1 \frac{g_1^d(x, Q^2)}{1 - 1.5\omega_D} dx$$



- In the  $\overline{\text{MS}}$ :  $a_0 = \Delta\Sigma = \Delta(u + \bar{u}) + \Delta(d + \bar{d}) + \Delta(s + \bar{s})$

- From COMPASS data alone:

$$\Gamma_1^N(Q^2 = 3\text{GeV}^2) = 0.046 \pm 0.002_{\text{stat.}} \pm 0.004_{\text{syst.}} \pm 0.005_{\text{evol.}}$$

- From COMPASS data alone (and still  $a_8$  from Phys. Rev. D **82** (2010):

$$a_0(Q^2 = 3\text{GeV}^2) = 0.32 \pm 0.02_{\text{stat.}} \pm 0.04_{\text{syst.}} \pm 0.05_{\text{evol.}}$$

(consistent with value from the COMPASS NLO QCD fit of world data).

## First moments of $g_1^{\text{NS}}$ and Bjorken sum rule

- Non-singlet structure function:

$$g_1^{\text{NS}} = g_1^{\text{p}}(x, Q^2) - g_1^{\text{n}}(x, Q^2)$$

- Bjorken sum rule:

$$\Gamma_1^{\text{NS}}(Q^2) = \frac{1}{6} \left| \frac{g_A}{g_V} \right| C_1^{\text{NS}}(Q^2)$$

- $g_1^{\text{NS}}$  calculated, NLO QCD fitted (only  $\Delta q_3$ ), evolved to  $Q^2 = 3\text{GeV}^2$

and fit-extrapolated  $x \rightarrow 0, 1$ :

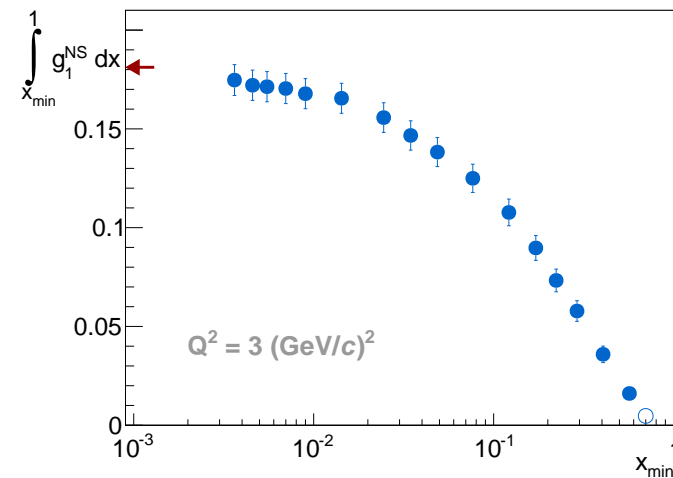
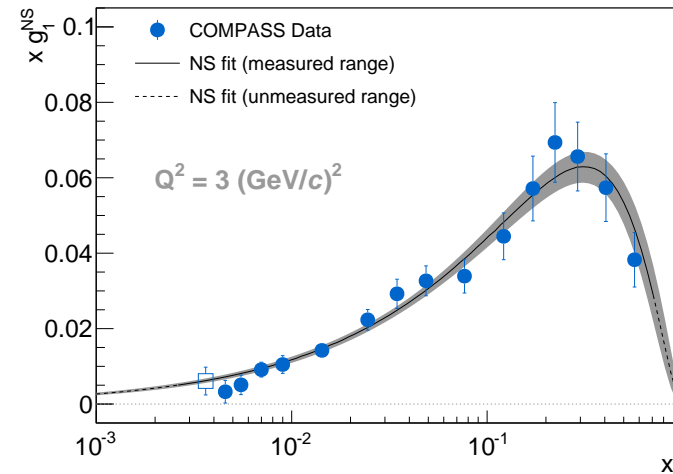
$$\Gamma_1^{\text{NS}}(Q^2) = 0.192 \pm 0.007_{\text{stat.}} \pm 0.015_{\text{syst.}}$$

$$\left| \frac{g_A}{g_V} \right| = 1.29 \pm 0.05_{\text{stat.}} \pm 0.10_{\text{syst.}}$$

- Neutron  $\beta$  decay gives:  $|g_A/g_V| = 1.2701 \pm 0.002$

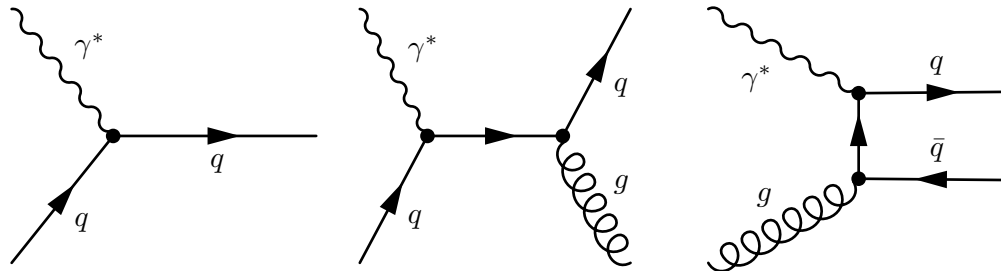
PDG, PRD86 (2012) 010001

- This validates the Bjorken sum rule with an accuracy of 9%



## $\Delta g$ : “direct” measurement

- Re-analysis of polarised deuteron data: COMPASS coll., Eur. Phys. J. C **77** (2017) no.4, 209

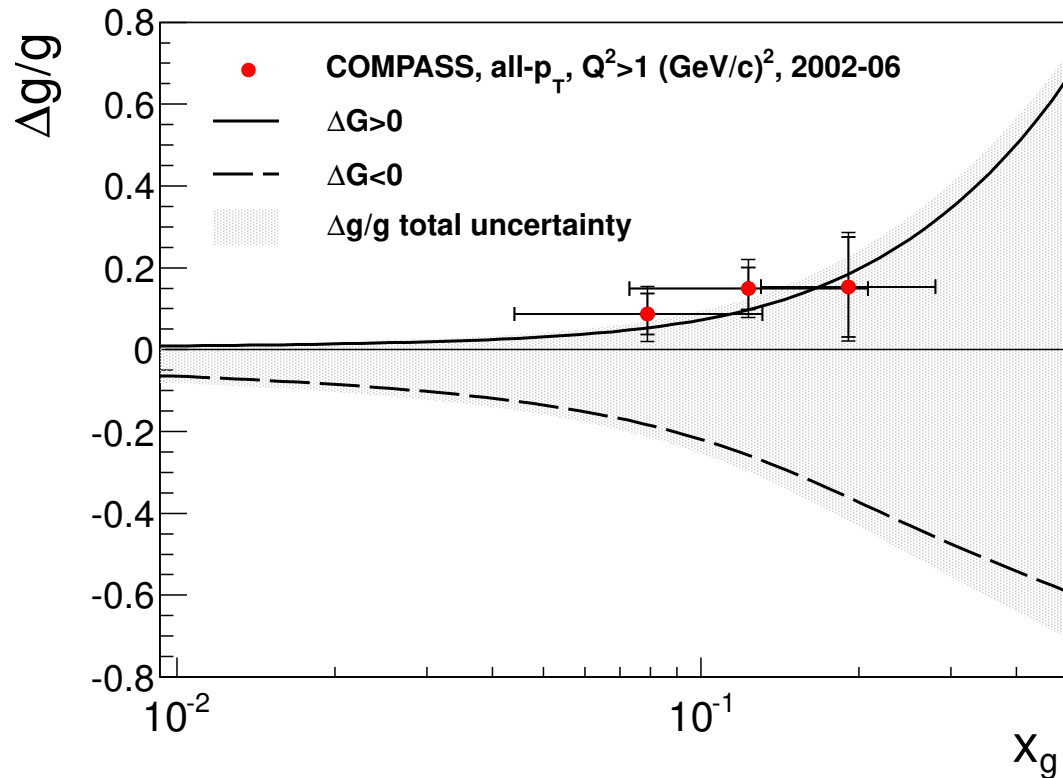


- $\Delta g$  accessed *via* PGF
- Double spin asymmetries of single inclusive hadron in DIS  
over a wide range in  $p_T$ ,  $0.05 < p_T < 2.5$  GeV/c
- LO order interpretation:
  - Hadron level  $\rightarrow$  parton level *via* LEPTO Monte Carlo
  - 1.6 gain in statistical and systematics precision
  - $\Delta g/g = 0.113 \pm 0.038_{\text{stat.}} \pm 0.035_{\text{syst.}}$ ,  $\langle Q^2 \rangle = 3\text{GeV}^2$ ,  $\langle x_g \rangle = 0.10$

## $\Delta g$ : “direct” measurement (*cont’d*)

- LO “direct” measurement, here in 3  $x_g$  bins, compared to NLO QCD fit

(*Note: 1st and 2nd bin are correlated.*)

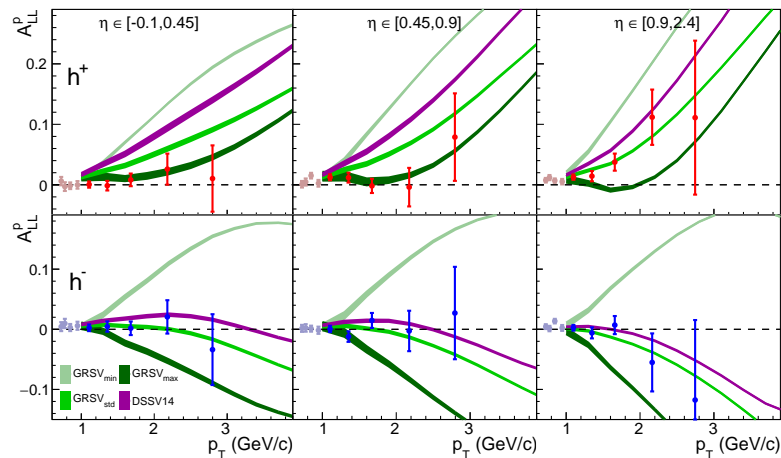


⇒ Illustrates the potential of “direct” asymmetry data.

(*But cannot be included in NLO fit.*)

## $\Delta g$ : “direct” measurement @ low $Q^2$

- “Direct” measurement at low  $Q^2$ : Phys. Lett. B **753** (2016) 573  
Double Spin asymmetry of single inclusive hadron @  $Q^2 < 1 \text{ GeV}^2$
- Extra resolved subprocesses contribute:  $gg, qg. \dots$



$p$  data compared with NLO calculation  
from Jäger, Stratmann, Vogelsang,  
Eur. Phys. J. C **44** (2005)

- New calculation with resummation of large logarithms:  
Uebler, Schäfer, Vogelsang, arXiv:1708.08284 [hep-ph]  
⇒ Reasonable agreement with pPDF = DSSV2014 and FF = DSS14( $\pi$ )+DSS17( $K$ )

## Conclusions

- COMPASS legacy on  $g_1^p(x, Q^2)$  and  $g_1^d(x, Q^2)$  presented for DIS and nonperturbative regions
  - $g_1^p$  at low  $x$  and low  $Q^2$  is clearly positive ( $g_1^d$  is consistent with zero)  
First observation of the spin effect at such low  $x$
  - $g_1^d \approx 0$  also at  $Q^2 > 1 \text{ GeV}^2$
- From the COMPASS data alone:  
First moments determined and Bjorken sum rule verified to 9 %
- NLO QCD fit of  $g_1$  world data gave well constrained quark distributions.  
Quark helicity contribution to nucleon spin:  $0.26 < \Delta\Sigma < 0.36$
- “Direct”  $\Delta g$  measurement
  - In DIS, LO extraction:  $\Delta g/g = 0.113 \pm 0.038_{\text{stat.}} \pm 0.035_{\text{syst.}}$ ,  $\langle Q^2 \rangle = 3 \text{ GeV}^2$ ,  $\langle x_g \rangle = 0.10$
  - At low  $Q^2$ , rich asymmetry data, consistent with DSSV2014



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# Spares

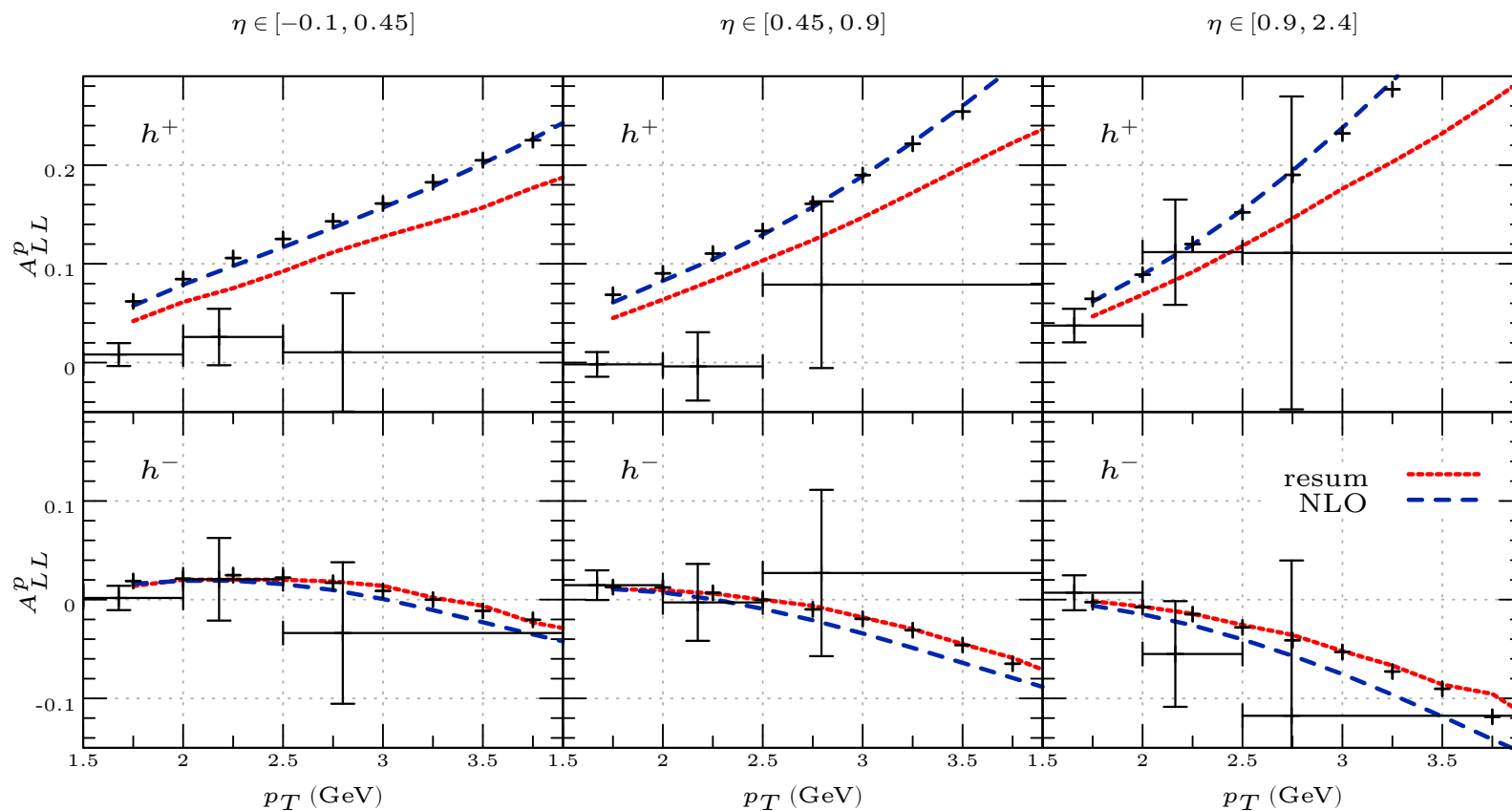
# Uebler, Schäfer, Vogelsang

- arXiv:1708.08284 [hep-ph]

NLO+resummation with pPDF = DSSV2014 and FF = DSS14( $\pi$ )+DSS17( $K$ )

Still tension for  $h^+$  off  $p$ , despite resummation of resolved processes.

Improvement for  $h^-$  off  $p$  thanks to DSS14



# Uebler, Schäfer, Vogelsang (*cont''d*)

- arXiv:1708.08284 [hep-ph]

NLO+resummation with FF = DSS14( $\pi$ )+DSS17( $K$ )

