TMD measurements at COMPASS

4th Workshop on the QCD Structure of the Nucleon (QCD-N'16)

Giulio Sbrizzai – Trieste INFN on behalf of the COMPASS Collaboration



COmmon Muon and Proton Apparatus for Structure and Spectroscopy

fixed target experiment at the CERN SPS

wide physics program carried on using both muon and hadron beam

luminosity: ~5 · 10³² cm⁻² s⁻¹

beam intensity: $2.10^8 \mu^+/\text{spill}$ (4.8s/16.2s)

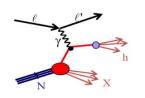
beam momentum: 160 GeV/c

nuclear 2004 2002 targets deuteron (6LID) 2003 hadron beam L/T 2008 2004 longitudinally LH target 2009 2006 polarized 2012 proton (NH₃) L/T 2007 muon beam 2010 2014 2011 2015 H₂ target T polarised DY 2012

Transversely (T) or Longitudinally (L) polarised Target

Accessing Spin and TMD PDFs and FFs

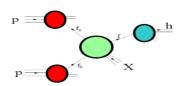
SIDIS off polarized p, d, n targets



HERMES
COMPASS
JLab

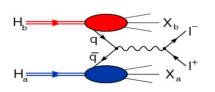
$$\sigma^{\ell p \to \ell' h X} \sim q(x) \otimes \hat{\sigma}^{lq \to lq} \otimes D_q^h(z)$$

hard polarised pp scattering



RHIC

polarised Drell-Yan

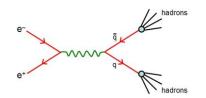


COMPASS

RHIC FNAL

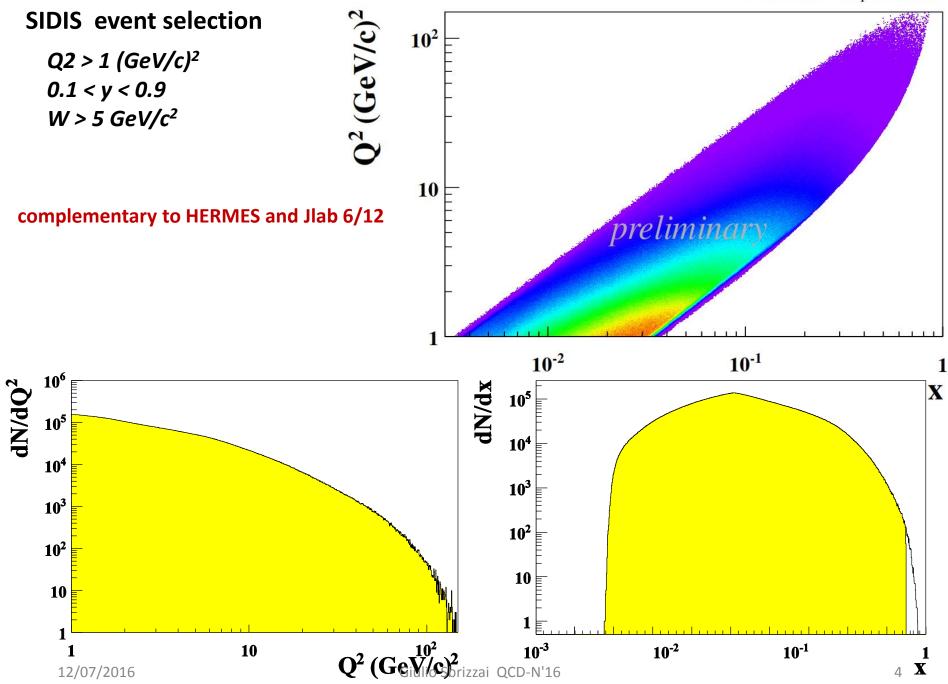
$$\sigma^{hp\to\mu\mu}{\sim}\bar{q}_h(x_1)\otimes q_p(x_2)\otimes\hat{\sigma}^{\bar{q}q\to\mu\mu}(\hat{s})$$

• $e^+e^- \rightarrow h_1h_2$

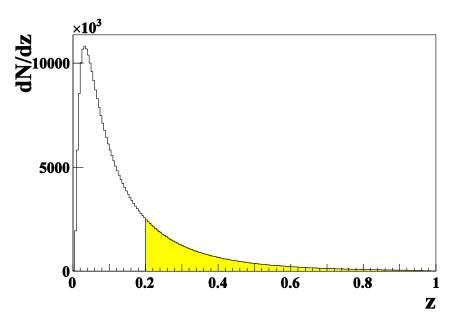


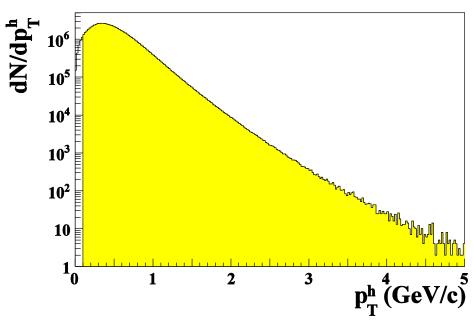
BaBar Belle Bes III

$$\sigma^{e^+e^-\to h_1h_2}\sim \hat{\sigma}^{\ell\ell\to\bar{q}q}(\hat{s})\otimes D_q^{h_1}(z_1)\otimes D_{\bar{q}}^{h_2}(z_2)$$



SIDIS event selection

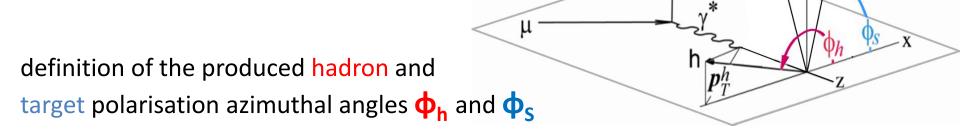




charged hadron selection

z>0.2

 $p_t^h > 0.1 \text{ GeV/c}$



polarised SIDIS azimuthal cross section

"one photon exchange approximation"

Bacchetta et al. JHEP 0702:093,2007

$$\begin{split} \frac{d\sigma}{dx\,dy\,d\psi\,dz\,d\phi_h\,dP_{h\perp}^2} &= \\ \frac{\alpha^2}{xy\,Q^2}\,\frac{y^2}{2\,(1-\varepsilon)}\,\left(1+\frac{\gamma^2}{2x}\right) \left\{F_{UU,T}+\varepsilon F_{UU,L}+\sqrt{2\,\varepsilon(1+\varepsilon)}\cos\phi_h\,F_{UU}^{\cos\phi_h}\right. \\ &+\varepsilon \cos(2\phi_h)F_{UU}^{\cos2\phi_h}+\lambda_e\,\sqrt{2\,\varepsilon(1-\varepsilon)}\sin\phi_h\,F_{LU}^{\sin\phi_h} \\ &+S_{\parallel}\left[\sqrt{2\,\varepsilon(1+\varepsilon)}\sin\phi_h\,F_{UL}^{\sin\phi_h}+\varepsilon\sin(2\phi_h)F_{UL}^{\sin2\phi_h}\right] +S_{\parallel}\lambda_e\left[\sqrt{1-\varepsilon^2}\,F_{LL}+\sqrt{2\,\varepsilon(1-\varepsilon)}\cos\phi_h\,F_{LL}^{\cos\phi_h}\right. \\ &+|S_{\perp}|\left[\sin(\phi_h-\phi_S)\left(F_{UT,T}^{\sin(\phi_h-\phi_S)}+\varepsilon\,F_{UT,L}^{\sin(\phi_h-\phi_S)}\right)\right. \\ &+\left.\left.\left.\left(1+\varepsilon\right)\left[\sin(\phi_h-\phi_S)F_{UT}^{\sin(\phi_h-\phi_S)}\right]\right.\right\} \\ &+\sqrt{2\,\varepsilon(1+\varepsilon)}\left[\sin\phi_S\,F_{UT}^{\sin\phi_S}+\sqrt{2\,\varepsilon(1+\varepsilon)}\left[\sin(2\phi_h-\phi_S)F_{UT}^{\sin(3\phi_h-\phi_S)}\right]\right. \\ &+\left.\left.\left.\left(1+\varepsilon\right)\left[\sin(\phi_h-\phi_S)F_{UT}^{\sin\phi_S}\right]\right.\right\} \\ &+\left.\left.\left.\left(1+\varepsilon\right)\left[\cos(\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \right\} \\ &+\left.\left.\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \\ &+\left.\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \\ &+\left.\left.\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \right\} \\ &+\left.\left.\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \\ &+\left.\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \\ &+\left.\left.\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \\ &+\left.\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \\ &+\left.\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \\ &+\left.\left.\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \\ &+\left.\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \\ &+\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \\ &+\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \\ &+\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \\ &+\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.\right\} \\ &+\left.\left(1+\varepsilon\right)\left[\cos(2\phi_h-\phi_S)F_{LT}^{\cos\phi_h-\phi_S}\right]\right.$$

14 independent modulations

beam polarisation

polarised SIDIS azimuthal cross section

"one photon exchange approximation"

Bacchetta et al. JHEP 0702:093,2007

$$\frac{d\sigma}{dx\,dy\,d\psi\,dz\,d\phi_h\,dP_{h\perp}^2} = \\ \frac{\alpha^2}{xy\,Q^2} \frac{y^2}{2\left(1-\varepsilon\right)} \left(1+\frac{\gamma^2}{2x}\right) \left\{F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)}\cos\phi_h F_{UU}^{\sin\phi_h}\right\} \\ + \varepsilon \cos(2\phi_h F_{UU}^{\cos2\phi_h}) + \lambda_{\varepsilon} \sqrt{2\varepsilon(1-\varepsilon)}\sin\phi_h F_{UU}^{\sin\phi_h}\right\} \\ + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)}\sin\phi_h F_{UU}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin2\phi_h}\right] + S_{\parallel}\lambda_{\varepsilon} \left[\sqrt{1-\varepsilon^2}\,F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)}\cos\phi_h F_{LL}^{\cos\phi_h}\right] \\ + |S_{\perp}| \left[\sin(\phi_h - \phi_S)\left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)}\right) + \varepsilon \sin(3\phi_h - \phi_S)F_{UT}^{\sin(3\phi_h - \phi_S)}\right] \\ + |S_{\perp}| \left[\sin(\phi_h + \phi_S)\left(F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S)F_{UT}^{\sin(3\phi_h - \phi_S)}\right) + \sqrt{2\varepsilon(1+\varepsilon)}\sin\phi_h F_{UT}^{\sin(\phi_h - \phi_S)}\right] \\ + \sqrt{2\varepsilon(1+\varepsilon)}\sin\phi_h F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)}\sin(2\phi_h - \phi_S)F_{UT}^{\sin(2\phi_h - \phi_S)}\right] \\ + |S_{\perp}|\lambda_{\varepsilon} \left[\sqrt{1-\varepsilon^2}\cos(\phi_h - \phi_S)F_{UT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)}\cos\phi_s F_{LT}^{\cos\phi_S}\right] \\ + |S_{\perp}|\lambda_{\varepsilon} \left[\sqrt{1-\varepsilon^2}\cos(2\phi_h - \phi_S)F_{LT}^{\cos(2\phi_h - \phi_S)}\right] \\ + |S_{\perp}|\lambda_{\varepsilon} \left[\sqrt{1-\varepsilon^2}\cos($$

OUTLINE

- → Hadron Transverse Momentum (P_T^h) dependent Multiplicities
- → Unpolarised Azimuthal Asymmetries (Boer-Mulders, Cahn effect)

- → Transverse Spin dependent Asymmetries
 - → Collins, Sivers
 - → P_T^h weighted Sivers asymmetries **NEW**

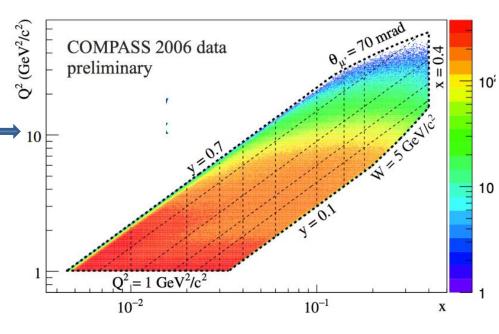
Polarised SIA

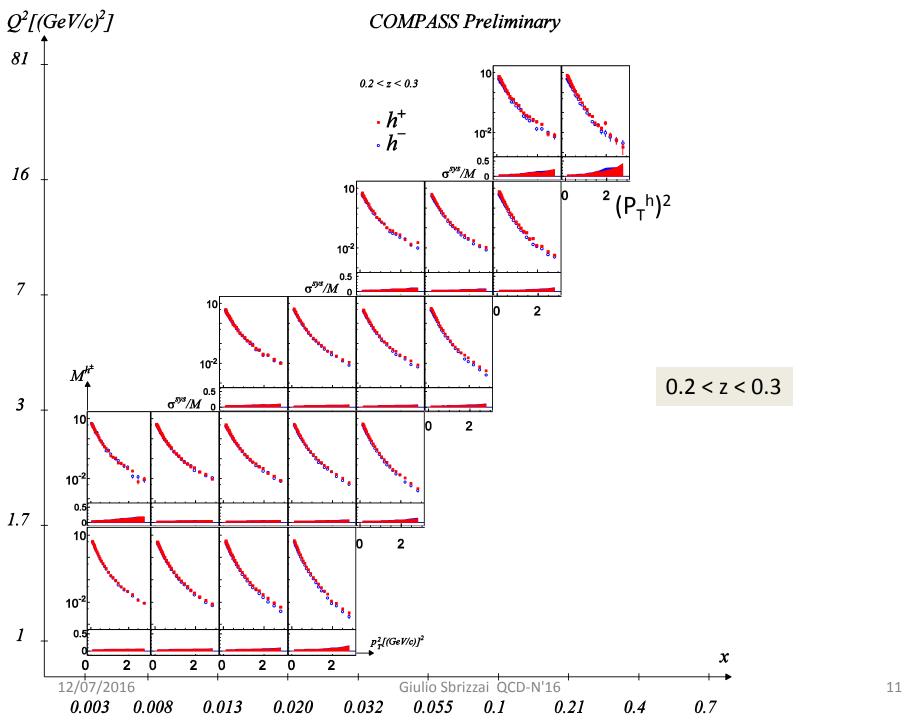
Unpolarised SIDIS

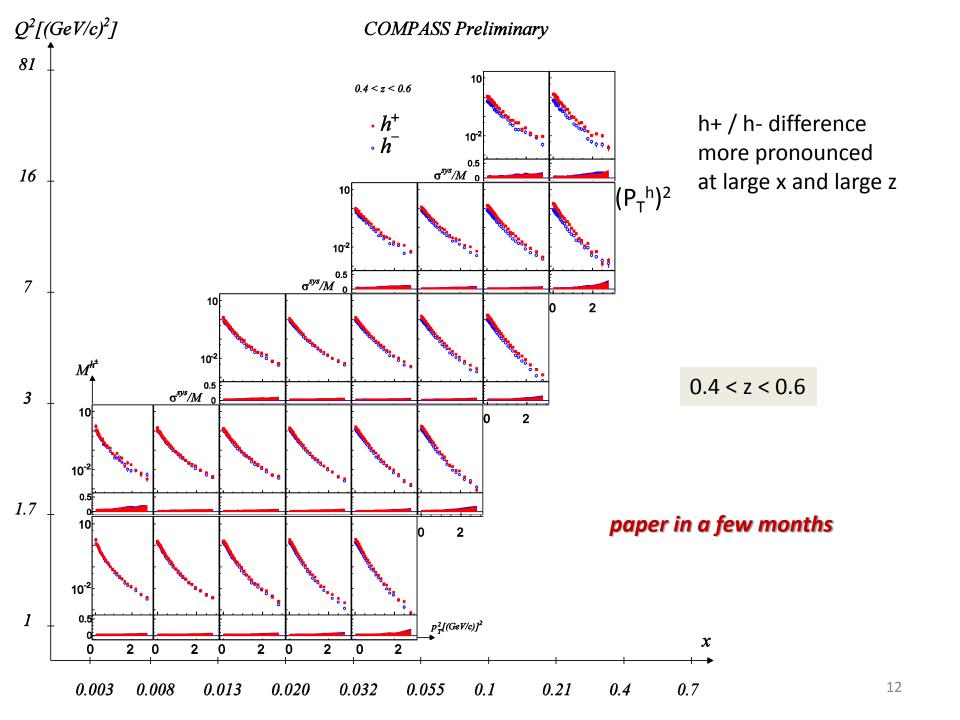
Relevance of unpolarised SIDIS for the TMDs

- The cross-section dependence from P_{hT} results from:
 - intrinsic k_{\perp} of the quarks
 - p_{\perp} generated in the quark fragmentation
 - A Gaussian ansatz for k_{\perp} and p_{\perp} leads to
 - $\langle P_{hT}^2 \rangle = z^2 \langle k_{\perp}^2 \rangle + \langle p_{\perp}^2 \rangle$

Results have been produced from 2004/2006 deuteron data multi dimensional analysis performed



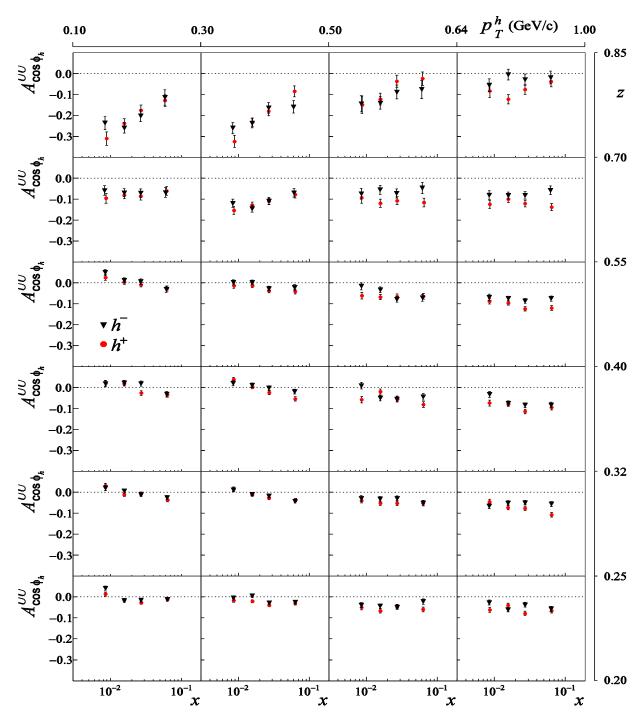




Unpolarised Azimuthal Asymmetries measured from 2004 deuteron data

$$A_{\cos\phi_h}^{UU} pprox rac{1}{Q} Cahn + rac{1}{Q} BM$$
 convoluted with the Collins FF

multiD
performed
both on
cosф
and cos2ф

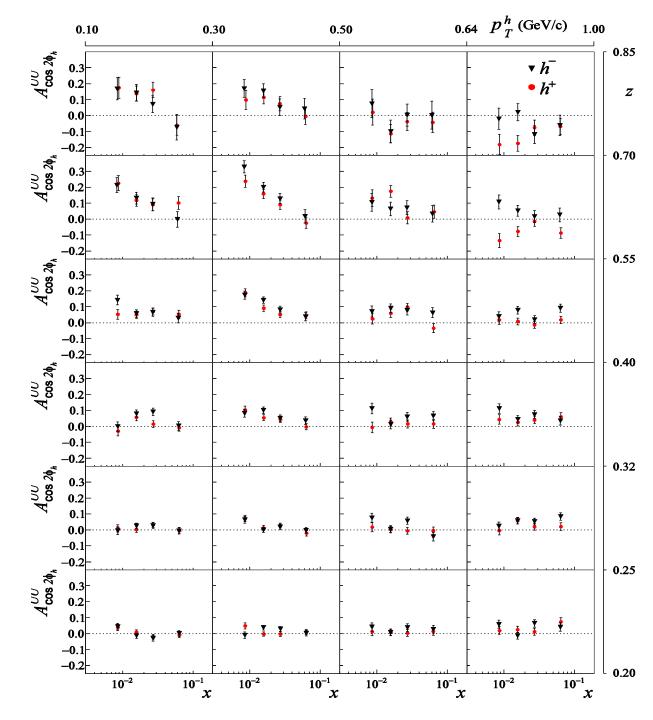


Unpolarised Azimuthal Asymmetries measured from 2004 deuteron data

$$A_{\cos 2\phi_h}^{UU} \approx \frac{1}{Q^2} Cahn + BM$$

convoluted with the Collins FF

multiD
performed
both on
cosф
and cos2ф



Polarised SIDIS

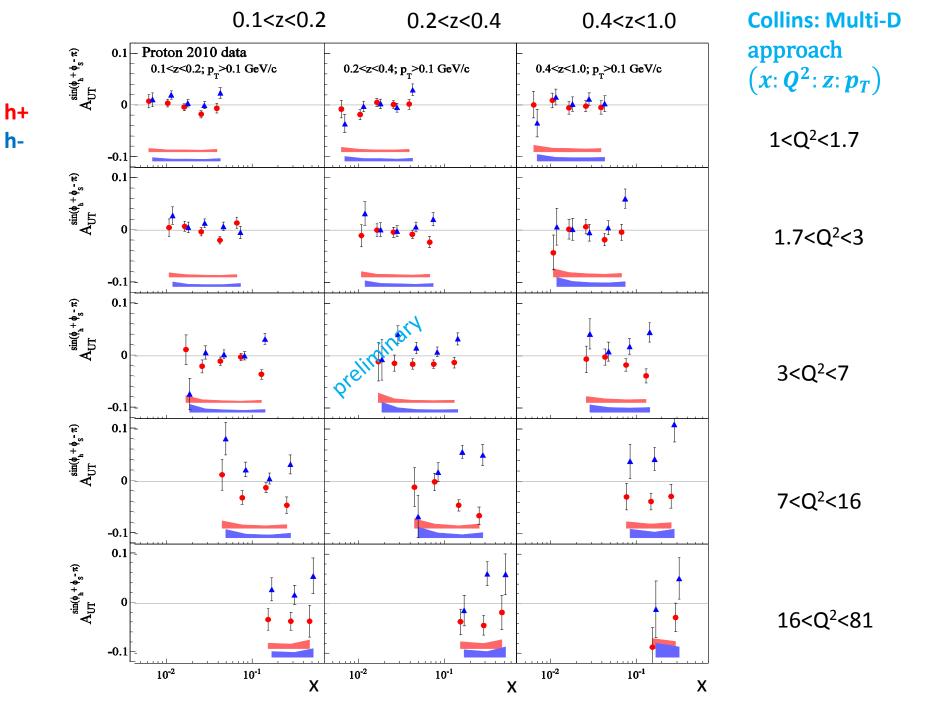
The Collins asymmetry



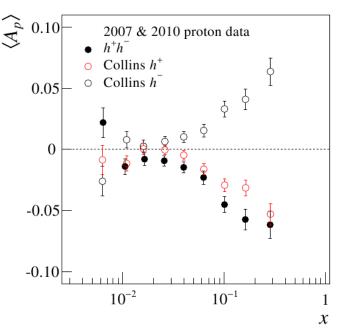
- results on deuteron (2002-2004 data) compatible with zero

 HERMES $p + COMPASS p\&d + BELLE \rightarrow extraction of transversity for u and d quarks$
- combined 2007 PLB 692 (2010) 240 and 2010 PLB 717 (2012) 376 PLB 744 (2015) 250
 published measurements on transversely polarised proton
 very good agreement between two independent data set

- multi dimensional analysis
- more recent:
 - comparison with di-hadron asymmetries → interplay



Interesting studies comparing with the Di-hadron Transverse Spin Asymmetries



$$lN \to l'h^+h^-X$$
 $A_{UT}^{\sin\phi_{RS}} \approx \frac{\sum_q e_q^2 \cdot h_1^q(x) \cdot H_{1q}^2(z, M_{hh}^2)}{\sum_q e_q^2 \cdot f_1^q(x) \cdot D_{1q}^h(z, M_{hh}^2)}$

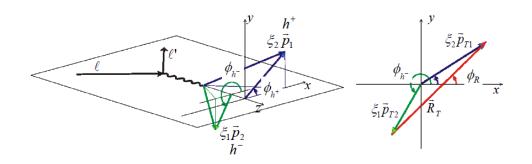
- Collins asymmetry for h+ and for h-: "mirror symmetry"
- dihadron asymmetry vs Collins asymmetry:
 only somewhat larger

analysis of the single hadron and dihadron asymmetries performed on a common data sample (2010 transversely polarised proton)

standard COMPASS SIDIS sample but wth h⁺ h⁻ at least detected (each hadron with z >0.1)

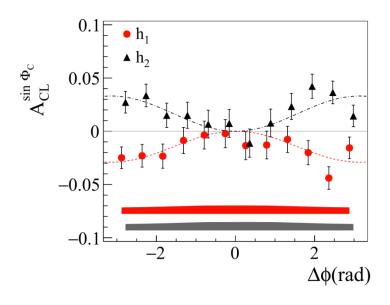
Study of the azimuthal correlations of the two hadrons

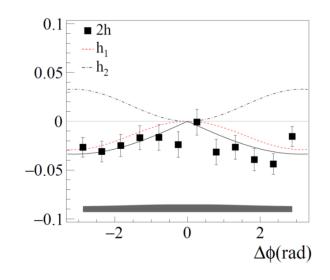
$$\Delta \phi = \phi_{h^+} - \phi_{h^-}$$



Interesting studies comparing with the Di-hadron Transverse Spin Asymmetries

$$\frac{d\sigma^{h^+h^-}}{d\phi_{h^+}d\phi_{h^-}d\phi_{S}} = \sigma_{U}^{h^+h^-} + S_T \cdot \left[\sigma_{1C}^{h^+h^-} \sin(\phi_{h^+} + \phi_{S} - \pi) + \sigma_{2C}^{h^+h^-} \sin(\phi_{h^-} + \phi_{S} - \pi)\right]$$





PLB 753 (2016) 406

$$A_{2CL}^{\sin(\phi_{h^{-}}+\phi_{S}-\pi)} = -\frac{\sigma_{1C}^{h^{+}h^{-}}}{\sigma_{U}^{h^{+}h^{-}}} \cdot (1 - \cos \Delta \phi)$$

$$A_{1CL}^{\sin(\phi_{h^{+}} + \phi_{S} - \pi)} = \frac{\sigma_{1C}^{h^{+}h^{-}}}{\sigma_{U}^{h^{+}h^{-}}} \cdot (1 - \cos \Delta \phi)$$

$$A_{2h,CL}^{\sin(\phi_{2h}+\phi_{S}-\pi)} = \frac{\sigma_{1C}^{h^{+}h^{-}}}{\sigma_{U}^{h^{+}h^{-}}} \cdot \sqrt{2 \cdot (1 - \cos \Delta \phi)}$$

"a common origin"

The Sivers asymmetry



• results on deuteron (2002-2004 data) compatible with zero

combined 2007 - PLB 692 (2010) 240 - and 2010 - PLB 717 (2012) 383 -

measurements on proton

very good agreement between two independent data set

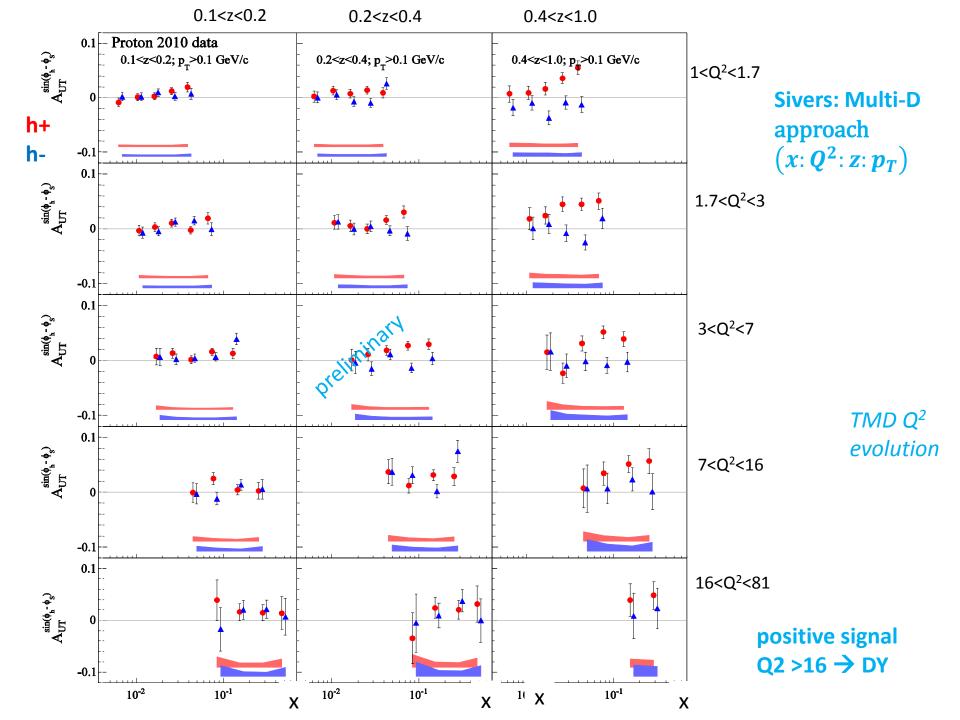
multi dimensional analysis

Recent Measurements

- Gluon Sivers
- Sivers from J/ψ
- Weighted Sivers asymmetries

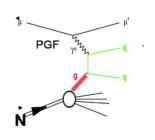


NEW!

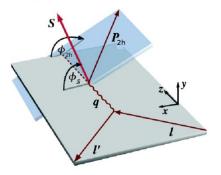


Gluon Sivers from high-P_T^h two hadrons pairs





 $\ell + N \rightarrow \ell' + 2h + X$

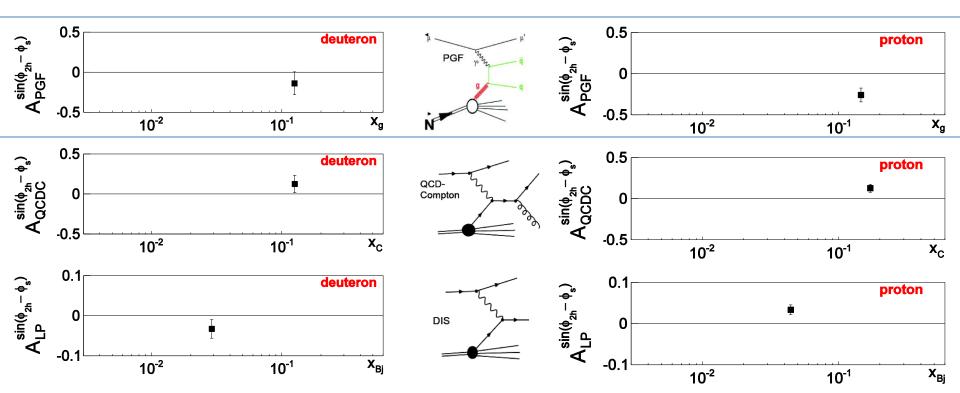


$$\vec{P}_{2h} = \vec{p}_1 + \vec{p}_2$$

$$\phi = \phi_{2h} - \phi_{S}$$

NN MC to tag the different contributing processes

 φ_{2h} correlated to φ_{gluon}



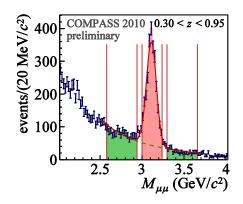
Gluon Sivers from J/ψ

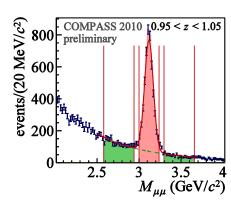
$$\mu^{+} + N \rightarrow \mu^{+} + J/\Psi + X \rightarrow 2\mu^{+} + \mu^{-} + X$$

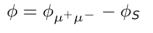
 $P_{J/\Psi}$ ϕ_s q l'

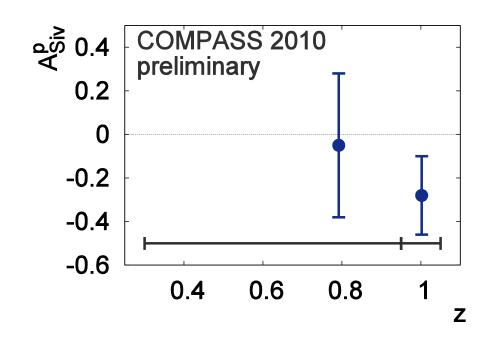
 $\mathbf{P}_{J/\Psi} = \mathbf{p}_{\mu^+} + \mathbf{p}_{\mu^-} \ \phi_{\mu^+\mu^-} = \phi_{J/\Psi} = \phi_{\mathbf{g}}$

from 2010 transversely polarised proton data in two z bins 0.30<z<0.95 and 0.95<z<1.05









New measurement: P_T^h weighted Sivers Asymmetry

A. Kotzinian and P. J. Mulders, PLB 406 (1997) 373

D. Boer and P. J. Mulders, PRD 57 (1998) 5780

J. C. Collins et al. PRD 73 (2006) 014021

$$A_{Siv}^{w} = \frac{1}{M} \frac{\sigma_{S}^{w}}{\sigma_{U}} \qquad \sigma_{S}^{w} = \int \sigma_{S}(P_{T}^{h}) \cdot \frac{P_{T}^{h}}{z} dP_{T}^{h}$$

$$\sigma(\varphi_{Siv}) = \sigma_{U} + \int \sigma_{S}(P_{T}^{h}) \cdot \frac{P_{T}^{h}}{z} dP_{T}^{h} \cdot \sin(\varphi_{Siv})$$

assuming polarisation = +1 and dilution factor = 1

standard Sivers Asymmetry:
$$A_{Siv} = \frac{\sigma_S}{\sigma_U}$$
 $\sigma(\varphi_{Siv}) = \sigma_U + \sigma_S \sin(\varphi_{Siv})$

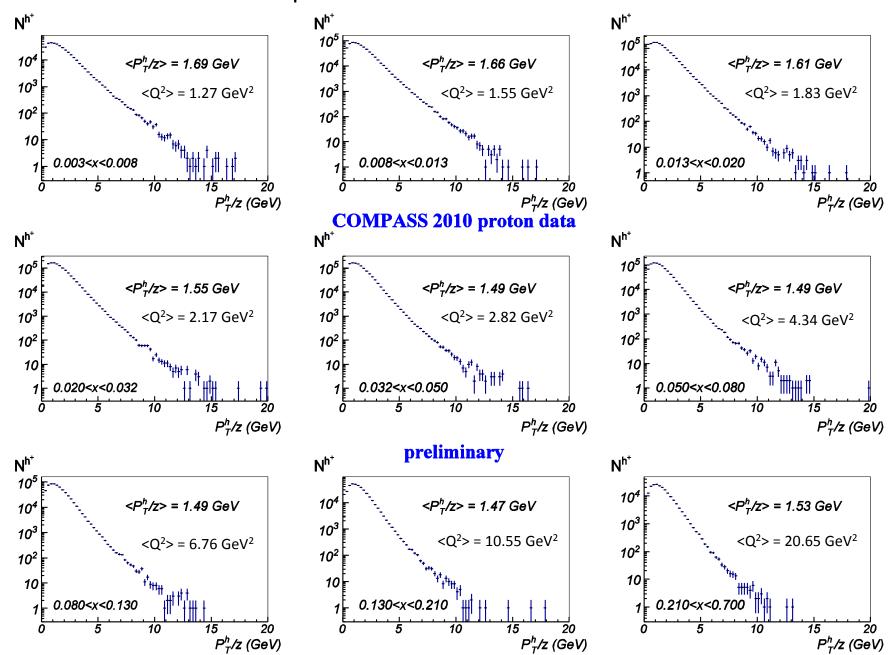
convolution over TM — product of Sivers (first moment) and FF

$$A_{Siv}^{w} = 2 \cdot \frac{\sum_{quarks} e_{q}^{2} \cdot f_{1T,q}^{\perp (1)}(x) \cdot D_{1,q}^{h}(z)}{\sum_{quarks} e_{q}^{2} \cdot f_{1,q}(x) \cdot D_{1,q}^{h}(z)}$$

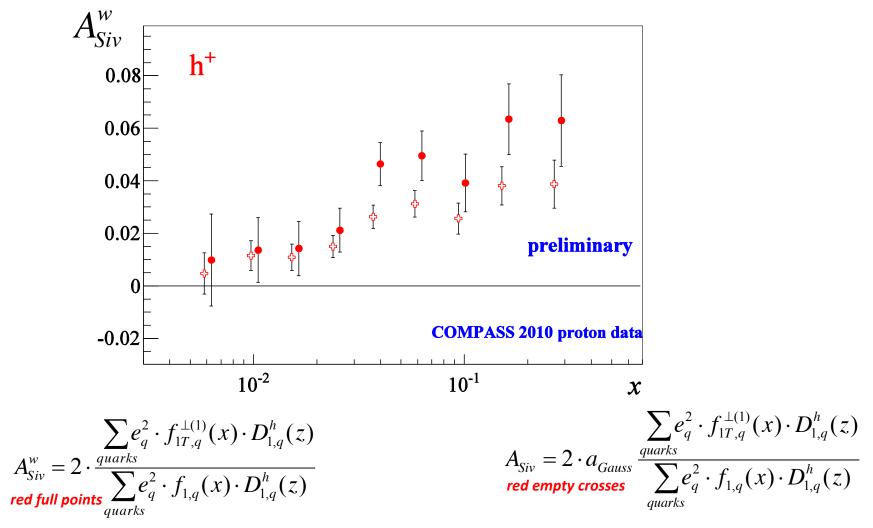
in a *model independent way* (no assumption on the shape of PDFs and FFs)

$$f_{1T}^{\perp (1)}(x) = \int d^2k_T \frac{k_T^2}{2M^2} f_{1T}^{\perp}(x, k_T^2)$$

The P_T^h/z distributions for each bin of x



Final results • compared with the standard Sivers Asymmetries 中 (PLB 717 (2012) 383)



Gassian assumption and a_{Gauss} needed to extract standard Sivers Function

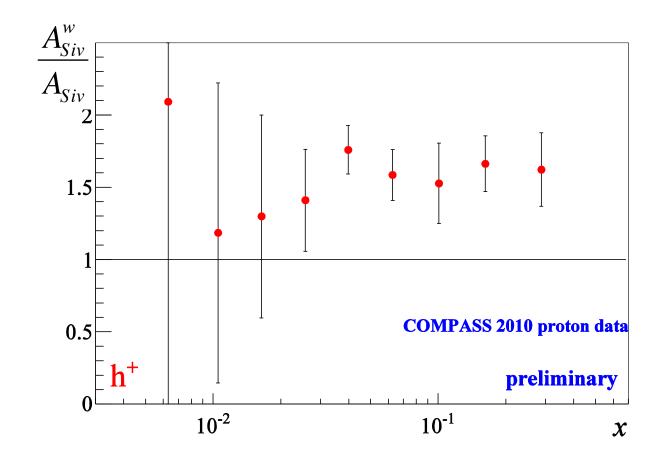
from the ratio

$$\frac{A_{Siv}^{w}}{A_{Siv}}$$

one can measure

$$a_{Gauss} \propto \frac{1}{\sqrt{k_{T,Sivers}^2 + p_{\perp}^2 / z^2}}$$

J. C. Collins et al. PRD 73 (2006) 014021

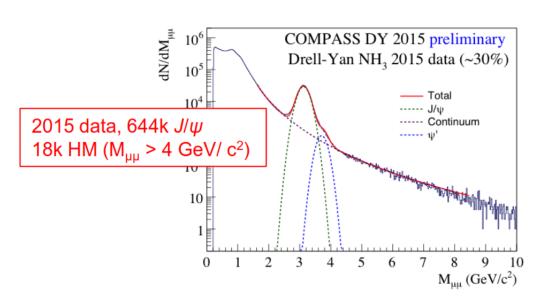


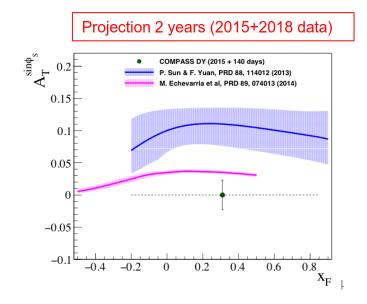
Special: plots from Drell-Yan 2015 polarised proton data

Pion beam on transversely polarized nucleon

COMPASS assets

- SIDIS and DY experiments: large acceptance, same kinematic region
- Unique hadron beam (π, K, p) with valence antiquarks
- Polarized target





Conclusions and Outlook

Many important results produced by COMPASS to investigates TMDs in SIDIS higher statistics data on transversely polarised d data still needed

DY data promising, more results soon to come

New data in the near future 2016-2017 unpolarised SIDIS on p, in parallel with DVCS

COMPASS III

backup

the COMPASS spectrometer

high energy beams

12/07/2016

- large angular acceptance
- broad kinematical range

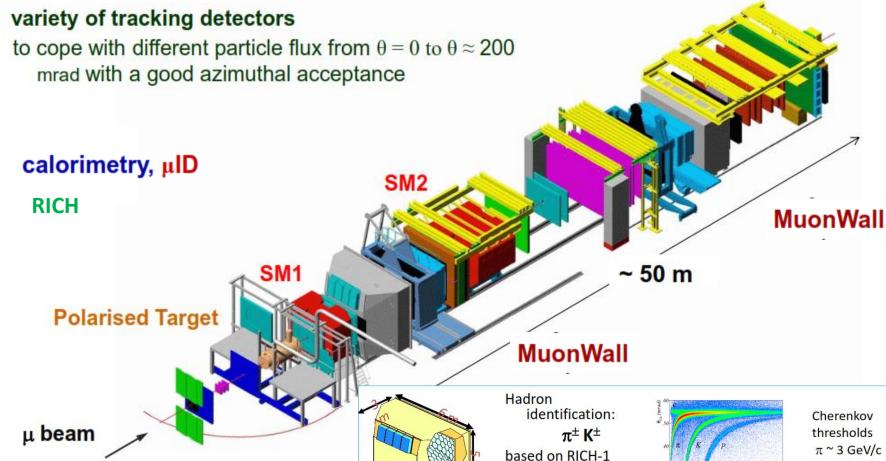
two stages spectrometer

Large Angle Spectrometer (SM1)

Small Angle Spectrometer (SM2)

response

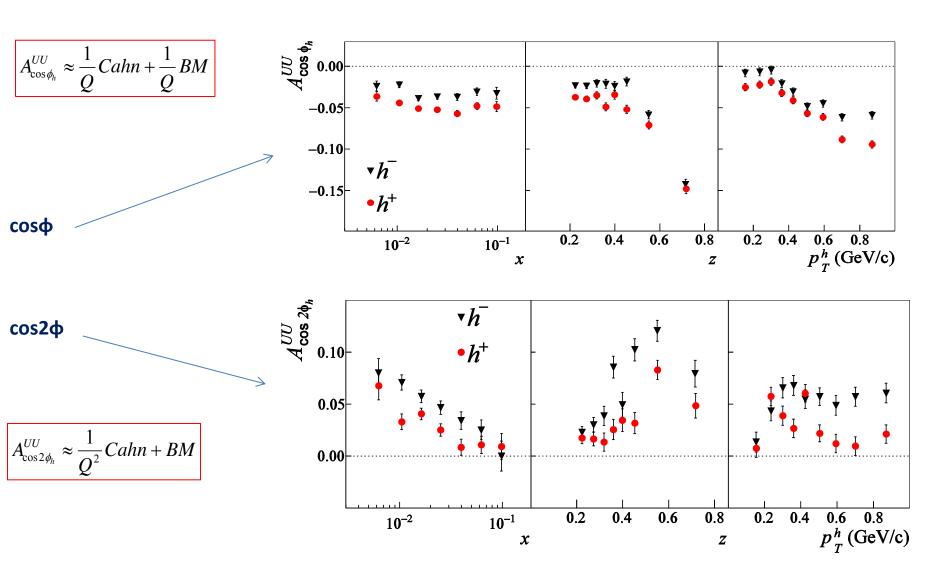
(likelihood algorithm)



 $\pi \sim 3 \text{ GeV/c}$ K~9 GeV/c p ~ 18 GeV/c

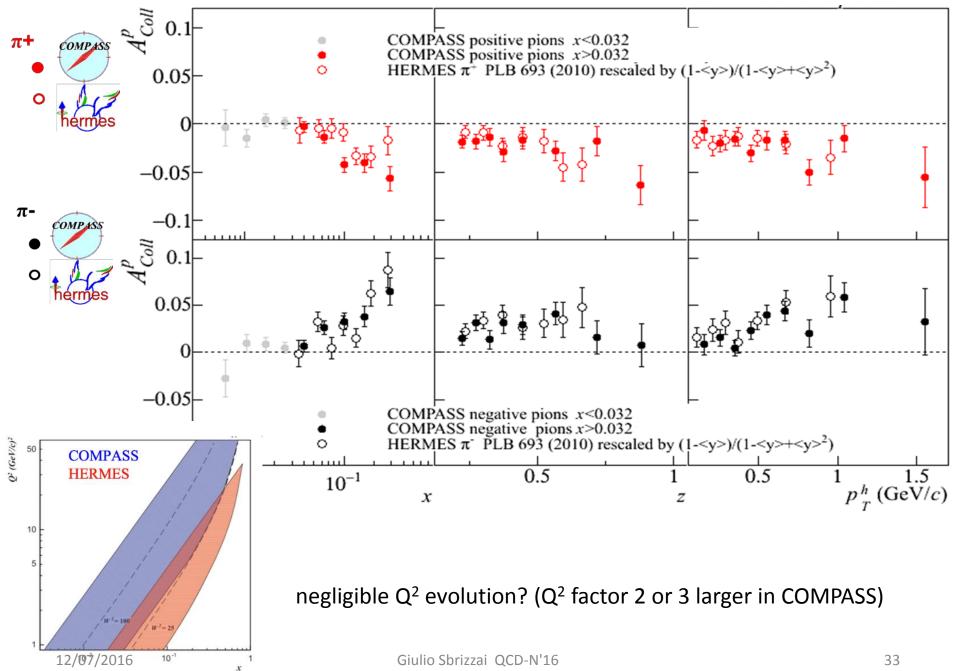
COMPA

Unpolarised Azimuthal Asymmetries measured from 2004 deuteron data



multi dimensional analysis performed to further investigate the interesting dependencies found

COMPASS and HERMES results

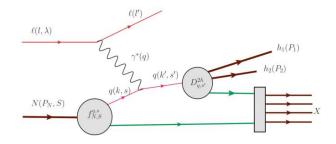


Interesting studies comparing with the Di-hadron Transverse Spin Asymmetries

$$A_{\!UT}^{\sin\phi_{\!RS}}$$

$$lN \rightarrow l'h^+h^-X$$

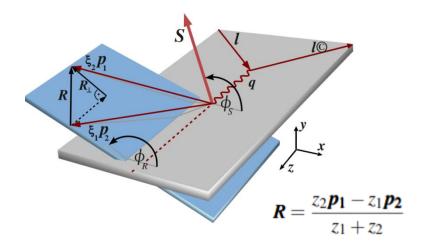
$$N_{h^+h^-}(\phi_{RS}) = N_h^0 \left[1 \pm f P_T D_{NN} A_{UT}^{\sin \phi_{RS}} \sin(\phi_{RS}) \right]$$



on oppositely charged hadrons pairs

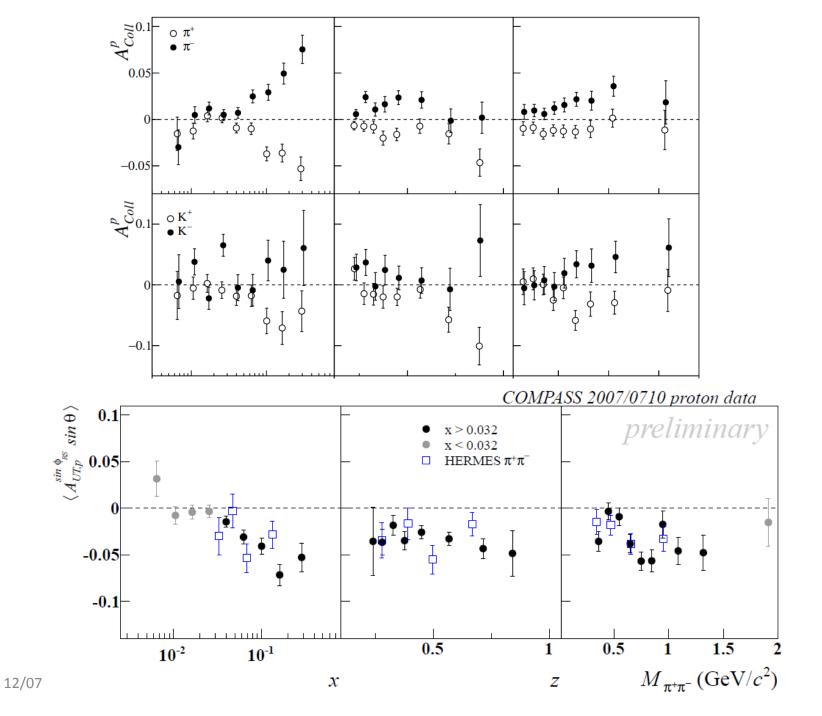
$$A_{UT}^{\sin\phi_{RS}} \approx \frac{\sum_{q} e_{q}^{2} \cdot h_{1}^{q}(x) \cdot H_{1q}^{2}(z, M_{hh}^{2})}{\sum_{q} e_{q}^{2} \cdot f_{1}^{q}(x) \cdot D_{1q}^{h}(z, M_{hh}^{2})}$$

collinear!

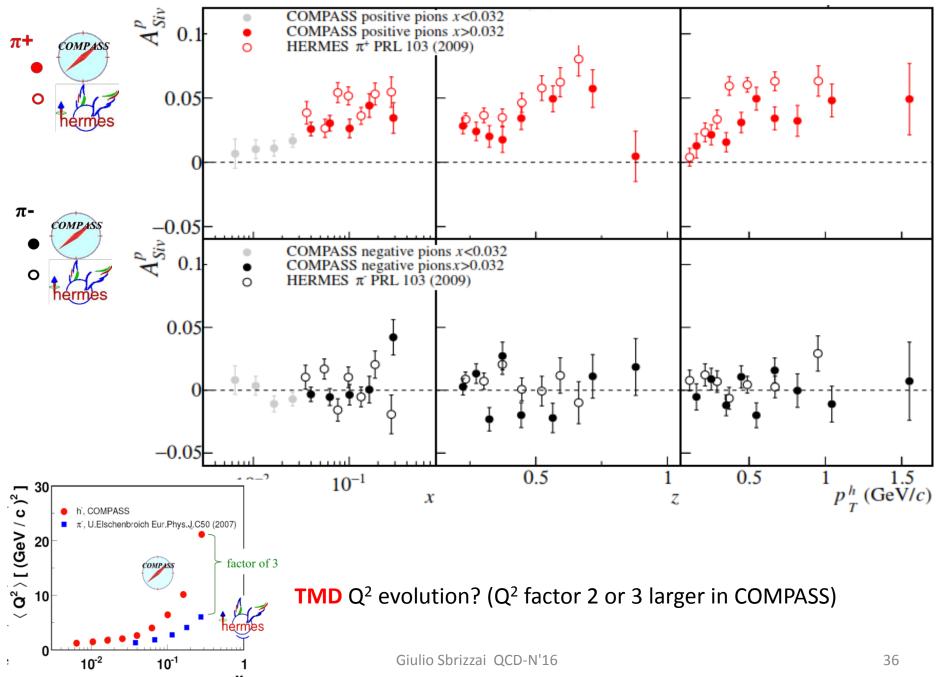


the azimuthal distribution of the hadrons pairs shows a modulation in the azimuthal angle:

$$\phi_{RS} = \phi_{R} + \phi_{S} - \pi$$

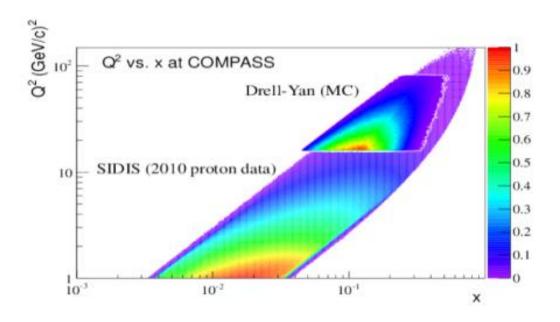


COMPASS and HERMES results



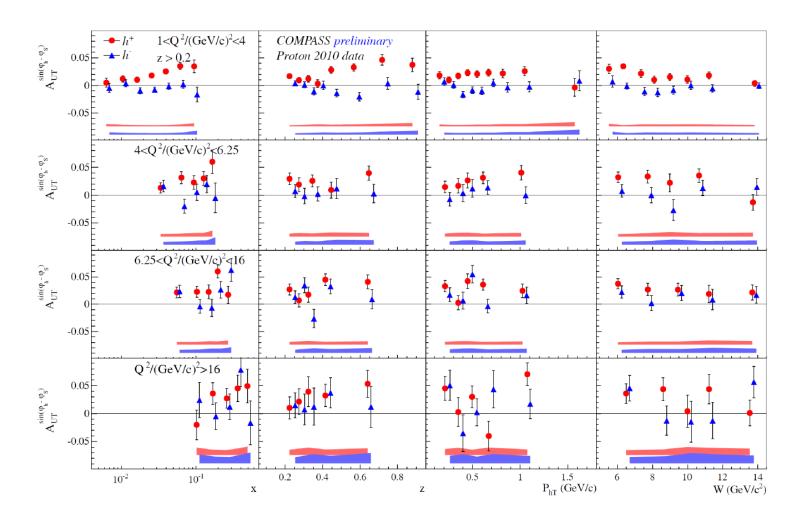
$$f_{1T}^{\perp}(SIDIS) = -f_{1T}^{\perp}(DY)$$

COMPASS is taking Drell-Yan data with transversely polarised target (full year dedicated)



superposition DY – SIDIS kinematical region at COMPASS

Sivers in DY range



one period (week) number of hadrons in sub-period 1

in each φ_{Siv} bin

number of hadrons in sub-period 2

 $N_1^{}$ cell **a+d** first sub-period $N_2^{}$ cell **b+c** first sub-period

 N_1 cell **a+d** second sub-period N_2 cell **b+c** second sub-period

data taking period (target transverse polarisation reversal)

weighted counts are

defined as:

$$N_1^w = \sum_{k=1}^{N_1} \frac{P_{T,k}^h}{z_k \cdot M_{Pr}}$$

and the same for N_2^w, N_1^w, N_2^w

Since only the spin dependent part of the cross section is weighted we used different method from the standard ones (DR, UML)

First method implemented:

$$\Delta^{w} = N_{1}^{w} - N_{1}^{'w} + N_{2}^{'w} - N_{2}^{w}$$

$$\Sigma = N_{1} + N_{1}' + N_{2}' + N_{2}$$

$$R'(\Phi_{Siv}) = \frac{\Delta^{w}}{\Sigma}$$

$$R'(\Phi_{Siv}) \simeq \bar{S}_T \cdot \epsilon^w \sin \Phi_{Siv}$$

assuming azimuthal acceptance to be the same for the two sub-periods

$$\epsilon^w = \frac{\sigma_{0S,I}^w}{\sigma_{U,I}} = 2A_{Siv}^{(1)}$$

calculated in 16 bins of ϕ_{Siv} and fitted using a $p_0+p_1\sin(\phi_{Siv})$ function

The method chosen to extract the final results is

$$R(\Phi_{Siv}) = \frac{\Delta^w}{\sqrt{\Sigma^w \Sigma}}$$

$$\Delta^{w} = N_{1}^{w} N_{2}^{'w} - N_{1}^{'w} N_{2}^{w}$$

$$\Sigma^{w} = N_{1}^{w} N_{2}^{'w} + N_{1}^{'w} N_{2}^{w}$$

$$\Sigma = N_{1} N_{2}^{\prime} + N_{1}^{\prime} N_{2}$$

quantity calculated in each bin of ϕ_{Siv}

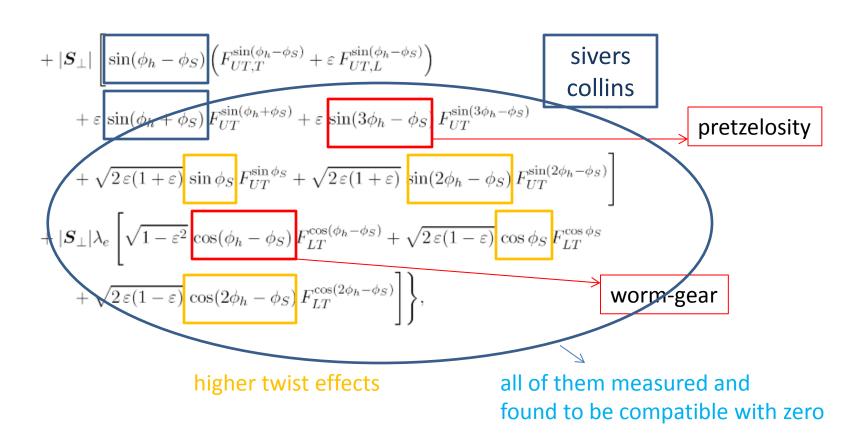
Azimuthal acceptance cancels if the Reasonable Assumption holds

$$a_1 a_2' = a_1' a_2$$

Other transverse spin dependent asymmetries

just a reminder

there are also other 6 modulations related to different TMDs they all have been measured at COMPASS





sensitive to the D-wave component non spherical shape of the nucleon

