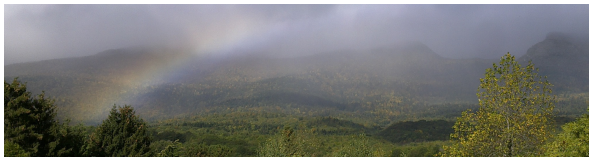


The COMPASS Measurement of the Pion Polarizability

Jan M. Friedrich

Physik-Department, TU München

COMPASS collaboration



Gordon Research Conference "Photonuclear Reactions"
August 9, 2016



Bundesministerium
für Bildung
und Forschung



Common Muon and Proton Apparatus for Structure and Spectroscopy





Common Muon and Proton Apparatus for Structure and Spectroscopy

CERN SPS: protons ~ 400 GeV (5 – 10 sec spills)

- secondary $\pi, K, (\bar{p})$: up to $2 \cdot 10^7/s$ (typ. $5 \cdot 10^6/s$)
Nov. 2004, 2008-09, 2012:
hadron spectroscopy & Primakoff reactions
- tertiary muons: $4 \cdot 10^7 / s$
2002-04, 2006-07, 2010-11: spin structure of the nucleon

LHC

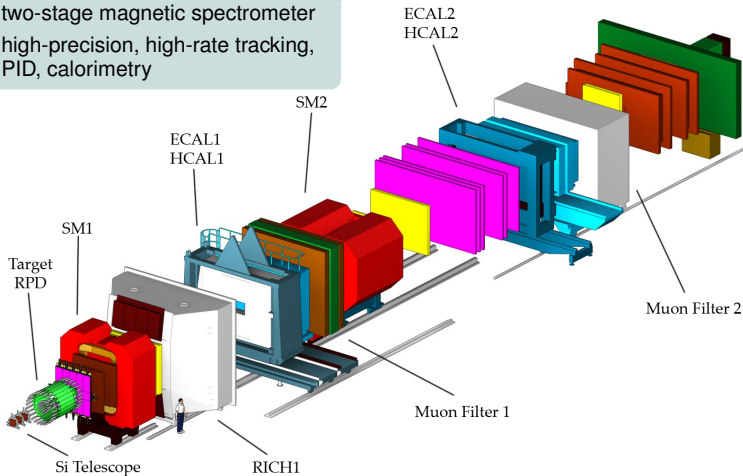
COMPASS

SPS



Fixed-target experiment

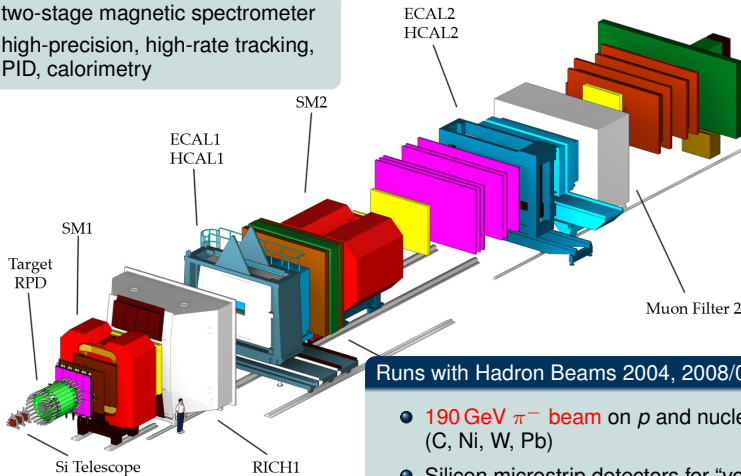
- two-stage magnetic spectrometer
- high-precision, high-rate tracking, PID, calorimetry





Fixed-target experiment

- two-stage magnetic spectrometer
- high-precision, high-rate tracking, PID, calorimetry

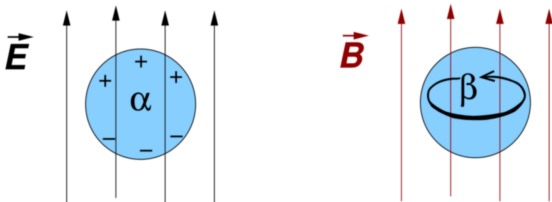


Runs with Hadron Beams 2004, 2008/09, 2012

- **190 GeV π^- beam** on p and nuclear targets (C, Ni, W, Pb)
- Silicon microstrip detectors for “vertexing”
- recoil and (digital) ECAL triggers



Electromagnetic Polarizabilities



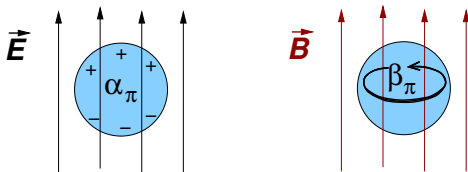
- structure-dependent response to outer e.m. fields:

$$\Delta H = -\frac{1}{2} (\alpha \cdot \vec{E}^2 + \beta \cdot \vec{B}^2)$$

- well-known for atoms and molecules
- measured on 10%-level for nucleons (also spin dependent)



Pion polarisability and ChPT



pion polarisabilities α_π, β_π in units of 10^{-4} fm^3

size of the pion $\sim 1 \text{ fm}^3$ [cf. atoms: polarisability \approx size $\approx 1 \text{ \AA}^3$]

Theory: ChPT (2-loop) prediction:

$$\begin{aligned}\alpha_\pi - \beta_\pi &= 5.7 \pm 1.0 \\ \alpha_\pi + \beta_\pi &= 0.16 \pm 0.1\end{aligned}$$

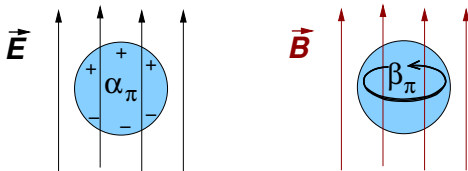
experiments for $\alpha_\pi - \beta_\pi$ lie in the range $4 \dots 14$

($\alpha_\pi + \beta_\pi = 0$ assumed)

ChPT: chiral perturbation theory: low-energy expansion of QCD



Pion polarisability and ChPT



pion polarisabilities α_π, β_π in units of 10^{-4} fm^3

size of the pion $\sim 1 \text{ fm}^3$ [cf. atoms: polarisability \approx size $\approx 1 \text{ \AA}^3$]

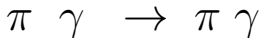
Theory: ChPT (2-loop) prediction:

$$\begin{aligned}\alpha_\pi &= 2.93 \pm 0.5 \\ \beta_\pi &= -2.77 \pm 0.5\end{aligned}$$

experiments for α_π lie in the range $2 \dots 7$

($\alpha_\pi + \beta_\pi = 0$ assumed)

ChPT: chiral perturbation theory: low-energy expansion of QCD



- Two kinematic variables, in CM: total energy \sqrt{s} , scattering angle θ_{cm}

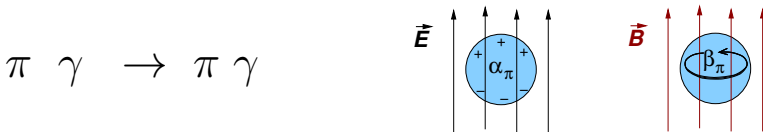
$$\frac{d\sigma_{\pi\gamma}}{d\Omega_{cm}} = \frac{\alpha^2 (s^2 z_+^2 + m_\pi^4 z_-^2)}{s (sz_+ + m_\pi^2 z_-)^2} - \frac{\alpha m_\pi^3 (s - m_\pi^2)^2}{4s^2 (sz_+ + m_\pi^2 z_-)} \cdot \mathcal{P}$$

$$\mathcal{P} = z_-^2 (\alpha_\pi - \beta_\pi) + \frac{s^2}{m_\pi^4} z_+^2 (\alpha_\pi + \beta_\pi) - \frac{(s - m_\pi^2)^2}{24s} z_-^3 (\alpha_2 - \beta_2)$$

$$z_\pm = 1 \pm \cos \theta_{cm}$$



Pion Compton Scattering



- Two kinematic variables, in CM: total energy \sqrt{s} , scattering angle θ_{cm}

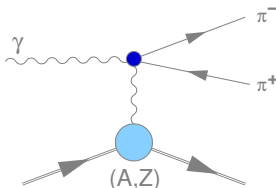
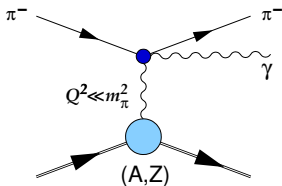
$$\frac{d\sigma_{\pi\gamma}}{d\Omega_{cm}} = \frac{\alpha^2 (s^2 z_+^2 + m_\pi^4 z_-^2)}{s (s z_+ + m_\pi^2 z_-)^2} - \frac{\alpha m_\pi^3 (s - m_\pi^2)^2}{4s^2 (s z_+ + m_\pi^2 z_-)} \cdot \mathcal{P}$$

$$\mathcal{P} = z_-^2 (\alpha_\pi - \beta_\pi) + \frac{s^2}{m_\pi^4} z_+^2 (\alpha_\pi + \beta_\pi) - \frac{(s - m_\pi^2)^2}{24s} z_-^3 (\alpha_2 - \beta_2)$$

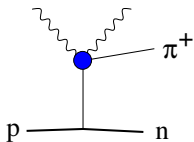
$$z_\pm = 1 \pm \cos \theta_{cm}$$



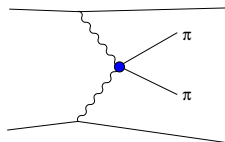
Pion Compton scattering: embedding the process



Primakoff processes



Radiative pion photoproduction

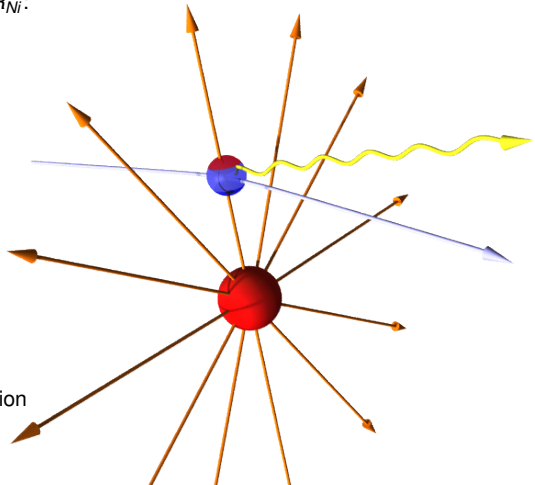


Photon-Photon fusion



Polarizability effect in Primakoff technique

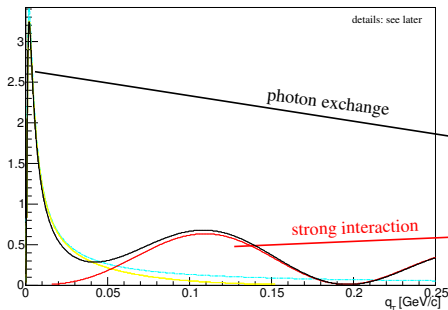
- Charged pions traverse the nuclear **electric field**
 - typical field strength at $d = 5R_N$:
 $E \approx 300 \text{ kV/fm}$
- Bremsstrahlung process:
 - particles scatter off **equivalent photons**
 - tiny momentum transfer
 $Q^2 \approx 10^{-5} \text{ GeV}^2/c^2$
 - pion/muon (quasi)-real Compton scattering
- Polarisability contribution
 - Compton cross-section typically diminished
 - corresponding charge separation
 $\approx 10^{-5} \text{ fm} \cdot e$





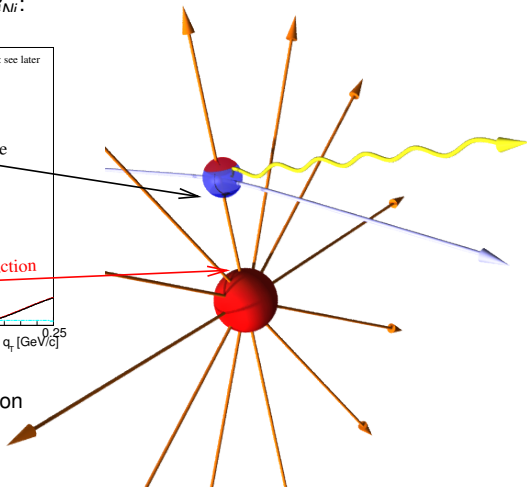
Polarisability effect in Primakoff technique

- Charged pions traverse the nuclear **electric field**
 - typical field strength at $d = 5R_{MI}$:



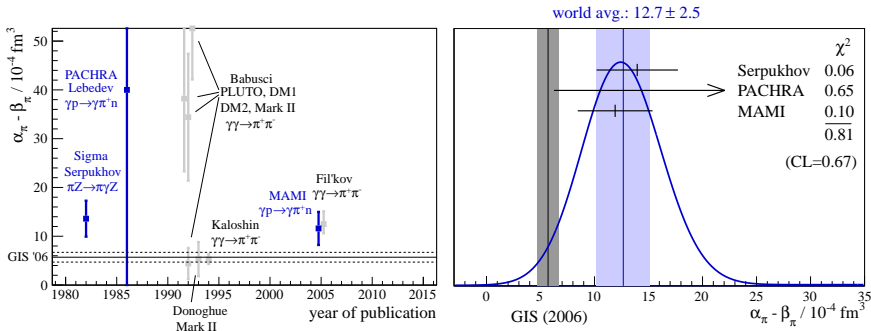
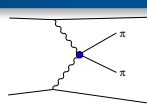
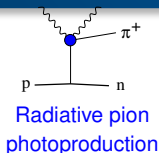
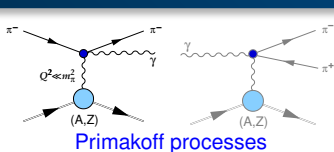
typically diminished

- corresponding charge separation $\approx 10^{-5} \text{ fm} \cdot e$





Pion polarisability: world data before COMPASS



GIS'06: ChPT prediction, Gasser, Ivanov, Sainio, NPB745 (2006), plots: T. Nagel, PhD
 Fil'kov analysis objected by Pasquini, Drechsel, Scherer PRC81, 029802 (2010)



PRL 114, 062002 (2015)

PHYSICAL REVIEW LETTERS

week ending
13 FEBRUARY 2015

Measurement of the Charged-Pion Polarizability

C. Adolph,⁸ R. Akhunzyanov,⁷ M. G. Alexeev,²⁷ G. D. Alexeev,⁷ A. Amoroso,^{27,29} V. Andrieux,²² V. Anosov,⁷
... [213 authors]
(COMPASS Collaboration)

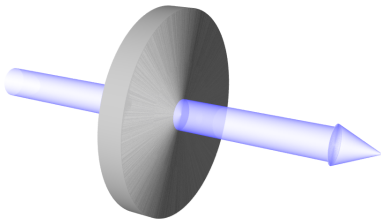
(Received 2 June 2014; revised manuscript received 24 December 2014; published 10 February 2015)

The COMPASS collaboration at CERN has investigated pion Compton scattering, $\pi^- \gamma \rightarrow \pi^- \gamma$, at center-of-mass energy below 3.5 pion masses. The process is embedded in the reaction $\pi^- \text{Ni} \rightarrow \pi^- \gamma \text{Ni}$, which is initiated by 190 GeV pions impinging on a nickel target. The exchange of quasireal photons is selected by isolating the sharp Coulomb peak observed at smallest momentum transfers, $Q^2 < 0.0015 \text{ (GeV}/c^2)^2$. From a sample of 63 000 events, the pion electric polarizability is determined to be $\alpha_\pi = (2.0 \pm 0.6_{\text{stat}} \pm 0.7_{\text{sys}}) \times 10^{-4} \text{ fm}^3$ under the assumption $\alpha_\pi = -\beta_\pi$, which relates the electric and magnetic dipole polarizabilities. It is the most precise measurement of this fundamental low-energy parameter of strong



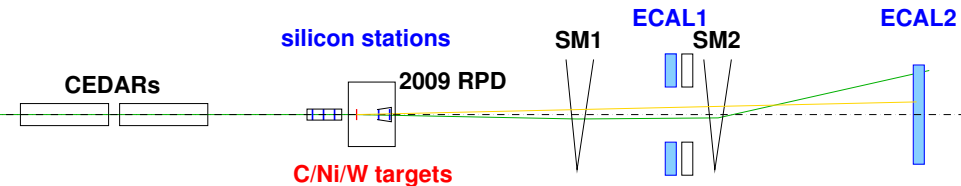


- high-energetic pion beam on 4mm nickel disk
- observe scattered **pions** in coincidence with produced **hard photons**
- study of the cross-section shape





Principle of the measurement



spatial resolution of tracks $\sim 10\mu\text{m}$
 angular resolution of photons $\sim 30\mu\text{rad}$



ECAL2: 3000 cells of different types



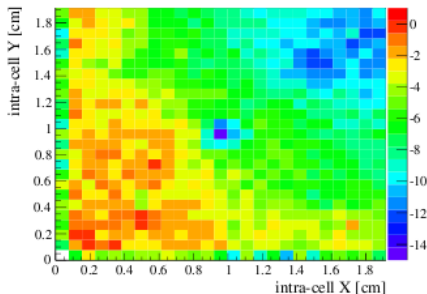


Figure 3.5: Profile of energy deviations shown for 1/4 of a shashlik block and for muon data photons within the range $133 \text{ GeV} < E_\gamma < 152 \text{ GeV}$.

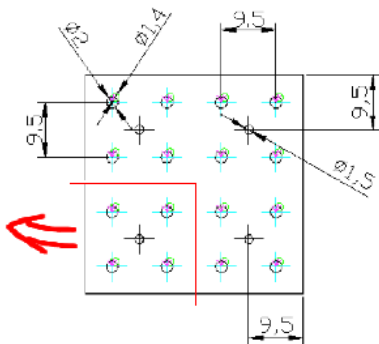


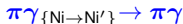
Figure 3.6: Technical drawing of a full shashlik cell to be compared with the figure to the left.

from: Th. Nagel, PhD thesis TUM 2012



Extraction of the pion polarizability

- Identify exclusive reactions



at smallest momentum transfer $< 0.001 \text{ GeV}^2/c^2$

- Assuming $\alpha_\pi + \beta_\pi = 0$, from the cross-section

$$R = \frac{\sigma(x_\gamma)}{\sigma_{\alpha_\pi=0}(x_\gamma)} = \frac{N_{meas}(x_\gamma)}{N_{sim}(x_\gamma)} = 1 - \frac{3}{2} \cdot \frac{m_\pi^3}{\alpha} \cdot \frac{x_\gamma^2}{1-x_\gamma} \alpha_\pi$$

is derived, depending on $x_\gamma = E_{\gamma(lab)}/E_{Beam}$.

Measuring R the polarizability α_π can be concluded.

- Control systematics by



and





Extraction of the pion polarizability

- Identify **exclusive reactions**



at smallest momentum transfer $< 0.001 \text{ GeV}^2/c^2$

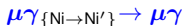
- Assuming $\alpha_\pi + \beta_\pi = 0$, from the cross-section

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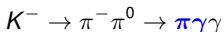
is derived, depending on $x_\gamma = E_{\gamma(\text{lab})}/E_{\text{Beam}}$.

Measuring R the polarizability α_π can be concluded.

- Control systematics by



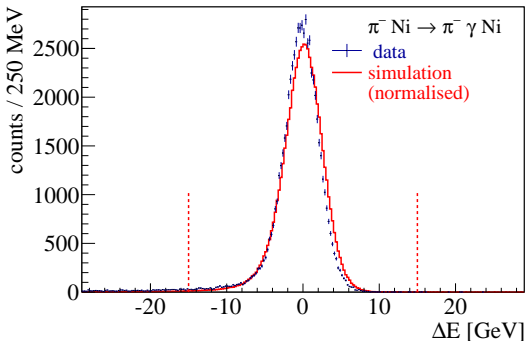
and





Identifying the $\pi\gamma \rightarrow \pi\gamma$ reaction

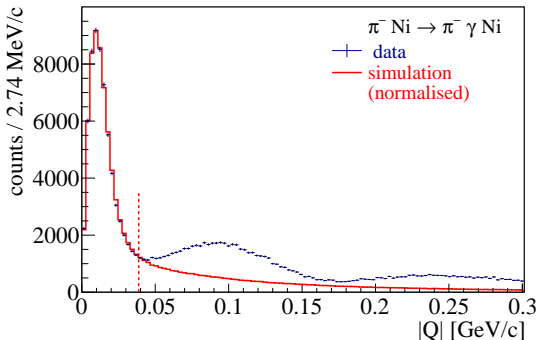
Phys. Rev. Lett. 114, 062002 (2015)



- Energy balance $\Delta E = E_\pi + E_\gamma - E_{\text{Beam}}$
- Exclusivity peak $\sigma \approx 2.6 \text{ GeV}$ (1.4%)
- ~ 63.000 exclusive events ($x_\gamma > 0.4$) (Serpukhov ~ 7000 for $x_\gamma > 0.5$)



Phys. Rev. Lett. 114, 062002 (2015)

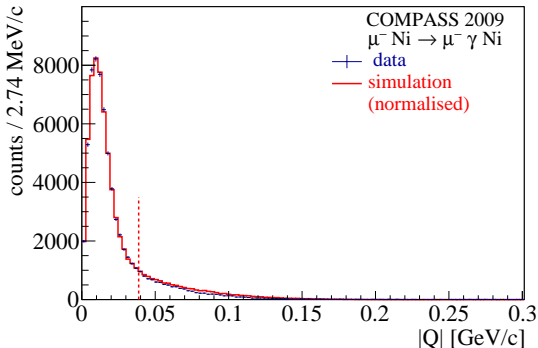


- $\Delta Q_T \approx 12 \text{ MeV}/c$ (190 GeV/c beam \rightarrow requires few- μrad angular resolution)
- first diffractive minimum on Ni nucleus at $Q \approx 190 \text{ MeV}/c$
- data a little more narrow than simulation \rightarrow negative interference?



Primakoff peak: muon data

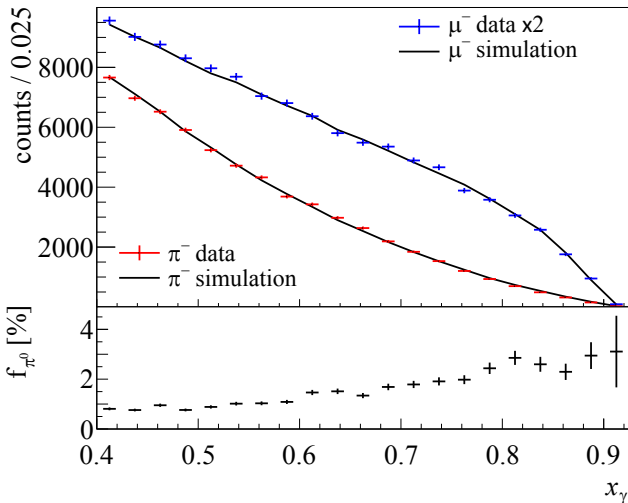
Phys. Rev. Lett. 114, 062002 (2015)



- **muon control measurement:** pure electromagnetic interaction
- e.m. nuclear effects well understood



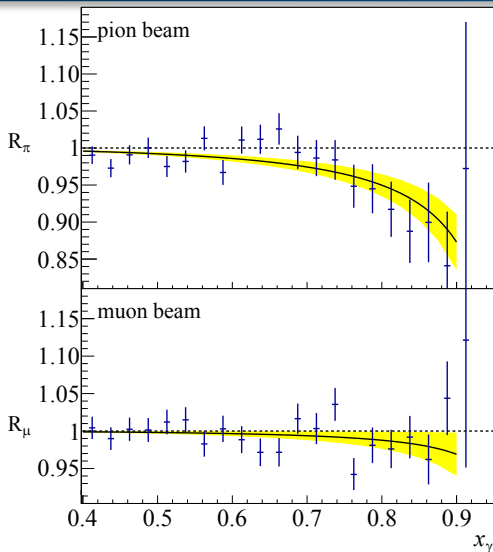
Photon energy spectra for muon and pion beam



Phys. Rev. Lett. 114, 062002 (2015)



Pion polarisability: COMPASS result



$$\alpha_\pi = (2.0 \pm 0.6_{\text{stat}}) \times 10^{-4} \text{ fm}^3$$

(assuming $\alpha_\pi = -\beta_\pi$)

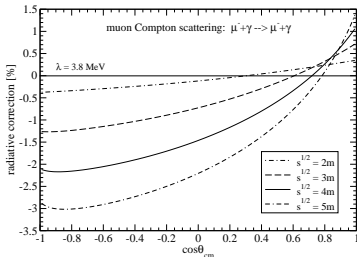
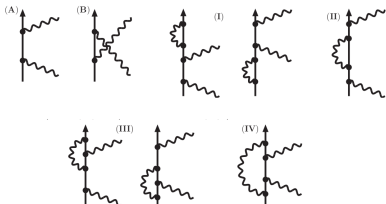
“false polarisability” from muon data:

$$(0.5 \pm 0.5_{\text{stat}}) \times 10^{-4} \text{ fm}^3$$

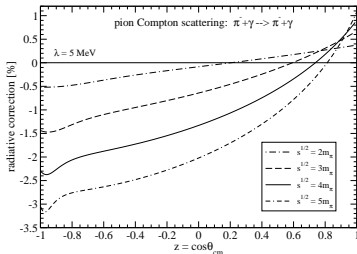
Phys. Rev. Lett. 114, 062002 (2015)



Radiative corrections (Compton scattering part)



Nucl.Phys. A837 (2010)



Eur.Phys.J. A39 (2009) 71



source of systematic uncertainty	estimated magnitude CL = 68 % [10 ⁻⁴ fm ³]
determination of tracking-detector efficiencies	0.5
treatment of radiative corrections	0.3
subtraction of π^0 background	0.2
strong interaction background	0.2
pion-electron elastic scattering	0.2
contribution of muons in the beam	0.05
quadratic sum	0.7



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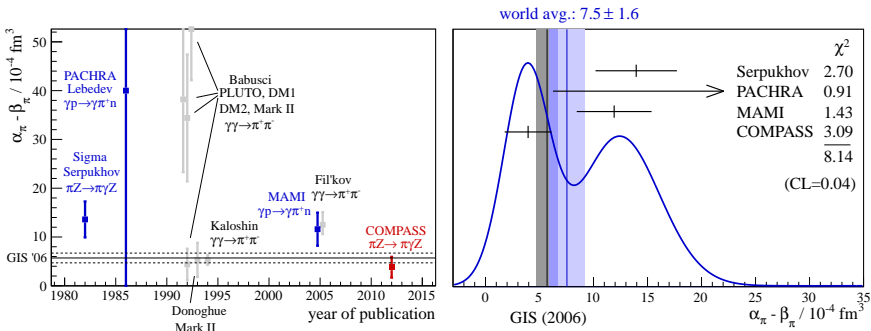
COMPASS result for the pion polarizability:

$$\alpha_\pi = (2.0 \pm 0.6_{\text{stat}} \pm 0.7_{\text{syst}}) \times 10^{-4} \text{ fm}^3$$

with $\alpha_\pi = -\beta_\pi$ assumed



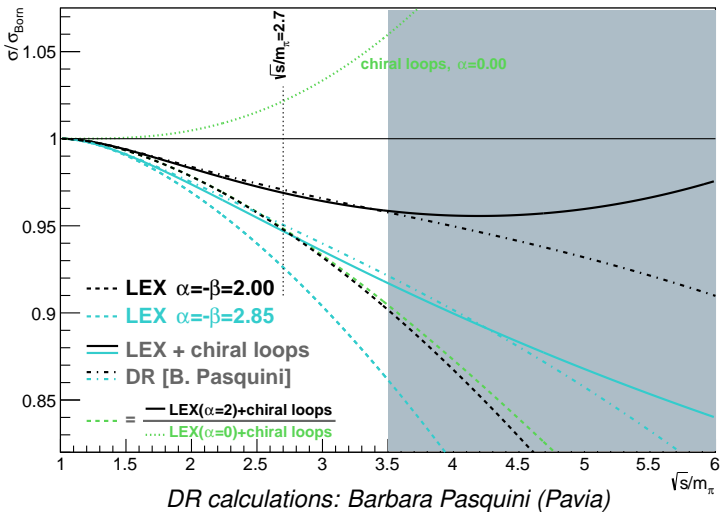
Pion polarisability: world data including COMPASS



- The new COMPASS result is in significant tension with the earlier measurements of the pion polarizability
- The expectation from ChPT is confirmed within the uncertainties



Polarisability and Loop Contributions $z=-1.0$



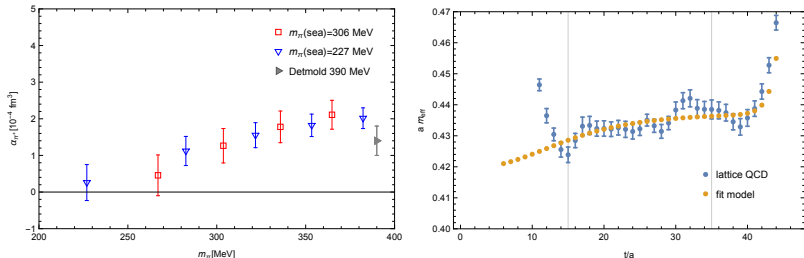
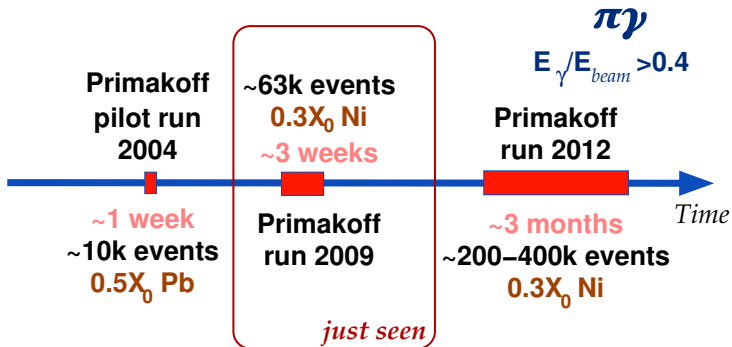


FIGURE 3. Left: electric polarizability for the charged pions as a function of the valence quark mass. The data for $m_\pi = 390$ MeV is taken from [5]. Right: effective mass for a charged pion correlator together with the scalar particle correlator determined from the fit. The fitting range is indicated by the vertical bars.

Alexandru *et al.*, Pion electric polarizability from lattice QCD, arXiv:1501.06516



Pion polarizability at COMPASS: further efforts



$$\mathcal{P} = z_-^2 (\alpha_{\pi} - \beta_{\pi}) + \frac{s^2}{m_{\pi}^4} z_+^2 (\alpha_{\pi} + \beta_{\pi}) - \frac{(s - m_{\pi}^2)^2}{24s} z_-^3 (\alpha_2 - \beta_2)$$

$$z_{\pm} = 1 \pm \cos \theta_{cm}$$

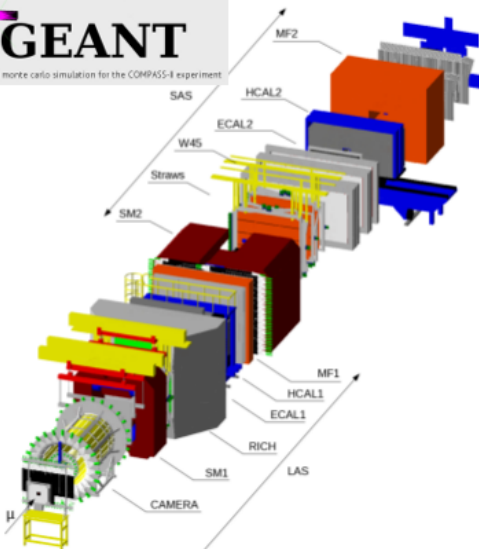


Transition to GEANT4

<https://ea58-project-tgeant.web.cern.ch>

TGEANT

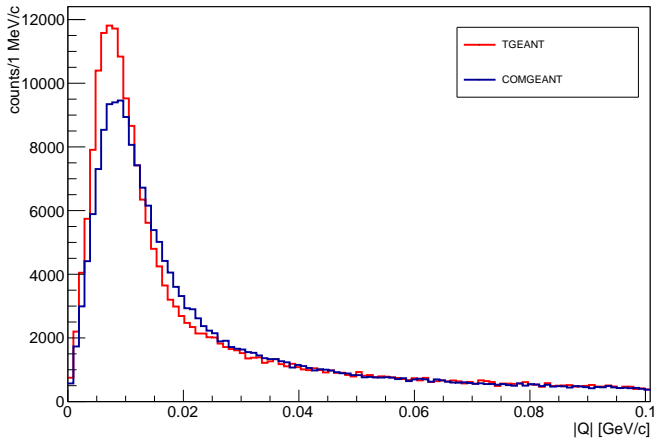
A Geant4 based monte carlo simulation for the COMPASS-E experiment



- COMGEANT: original Monte Carlo Simulation for COMPASS based on GEANT3
- Since 2012 efforts to upgrade to C++ based on GEANT4



Transition to GEANT4





Access to $\pi + \gamma$ reactions via the **Primakoff effect**

$$\pi^- + \gamma \rightarrow \left\{ \begin{array}{l} \pi^- + \gamma \\ \pi^- + \pi^0 / \eta \\ \pi^- + \pi^0 + \pi^0 \\ \pi^- + \pi^- + \pi^+ \\ \pi^- + \pi^- + \pi^+ + \pi^- + \pi^+ \\ \pi^- + \dots \end{array} \right.$$

analogously: Kaon-induced reactions $K^- + \gamma \rightarrow \dots$



- explorative study with 2009 data on “double tungsten” target
- ongoing study with 2012 data on “(side-)segmented tungsten” target
- best future opportunity: measure π^0 decays

$$\pi^0 \rightarrow \gamma\gamma \quad \text{and} \quad \pi^0 \rightarrow \gamma e^+ e^-$$

from the (exclusive) $\pi^- \pi^0$ final state

on a double tungsten target with *varying distance*

- inspired by and extending the 1980’s “CERN direct measurement of the π^0 lifetime”
- complementary to the PRIMEX measurement



- Measurement of the **pion polarisability** at COMPASS
 - Via the Primakoff reaction, COMPASS has determined

$$\alpha_{\pi} = (2.0 \pm 0.6_{\text{stat}} \pm 0.7_{\text{syst}}) \times 10^{-4} \text{ fm}^3$$

- most direct access to the $\pi\gamma \rightarrow \pi\gamma$ process
- Most precise experimental determination
- Systematic control: $\mu\gamma \rightarrow \mu\gamma$, $K^- \rightarrow \pi^- \pi^0$
- Related topics at COMPASS: radiative widths and chiral dynamics in $\pi^- \gamma \rightarrow \pi^- \pi^0$ and $\pi\gamma \rightarrow \pi\pi\pi$ reactions
- High-statistics run 2012
 - separate determination of α_{π} and β_{π}
 - s -dependent quadrupole polarisabilities
 - First measurement of the kaon polarisability

LIGHT UNFLAVORED MESONS
($S = C = B = 0$)

For $l = 1$ ($\pi, \eta, h, \eta', \rho, \omega$): $a\vec{a}, (a\vec{b}-d\vec{d})/\sqrt{2}, \dots$
for $l = 0$ ($\eta, \eta', h, \eta', \omega, \phi, f, f'$): $c_1(a\vec{b} + d\vec{d}) + c_2(s\vec{s})$

π^{\pm} $f_0^G(f_0^P) = 1 \cdot (0^-)$

π ELECTRIC POLARIZABILITY α_{π}

See HOLSTEIN 14 for a general review on hadron polarizability.

VALUE (10^{-4} fm^3)	EVTS	DOCUMENT ID	TECH	COMMENT
$2.0 \pm 0.6 \pm 0.7$	63k	¹ ADORNIN	15A SPEC	$\pi^+ \gamma \rightarrow \pi^+ \gamma$ Compton scatt.

¹Value is derived assuming $\alpha_{\pi} = -\beta_{\pi}$.

π^{\pm} REFERENCES

We have omitted some papers that have been superseded by later exper-



Thank you for your attention!

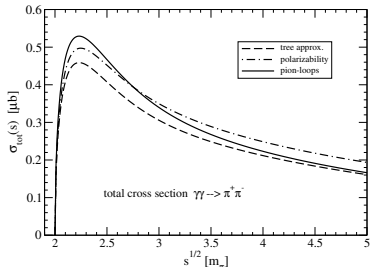




- Planned measurements at ALICE and JLab

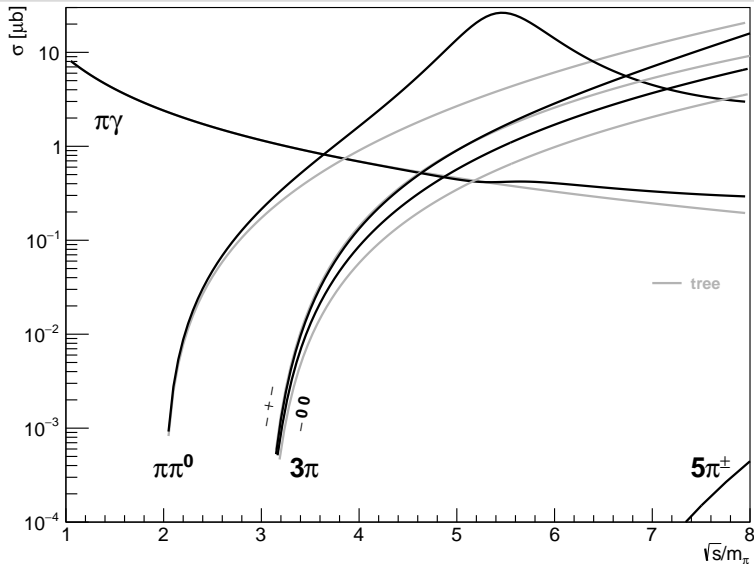
$$\sigma_{tot}(s) = \frac{2\pi\alpha^2}{\hat{s}^3 m_\pi^2} \left\{ [4 + \hat{s} + \hat{s} |C(\hat{s})|^2] \sqrt{\hat{s}(\hat{s} - 4)} + 8[2 - \hat{s} + \hat{s} \operatorname{Re} C(\hat{s})] \ln \frac{\sqrt{\hat{s}} + \sqrt{\hat{s} - 4}}{2} \right\},$$

$$C(\hat{s}) = -\beta_\pi \frac{m_\pi^3}{2\alpha} \hat{s} - \frac{m_\pi^2}{(4\pi f_\pi)^2} \left\{ \frac{\hat{s}}{2} + 2 \left[\ln \frac{\sqrt{\hat{s}} + \sqrt{\hat{s} - 4}}{2} - \frac{i\pi}{2} \right]^2 \right\}$$



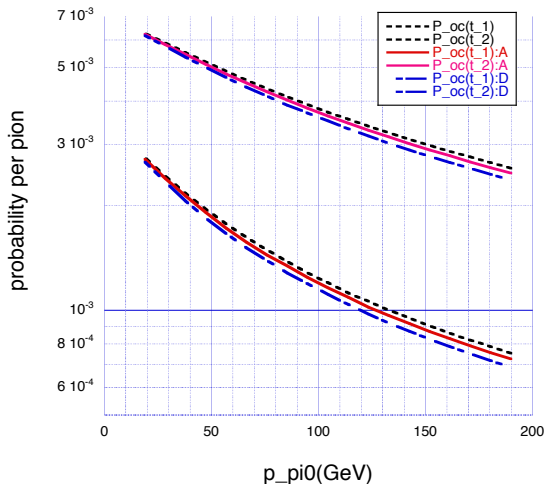
courtesy Norbert Kaiser (TUM)

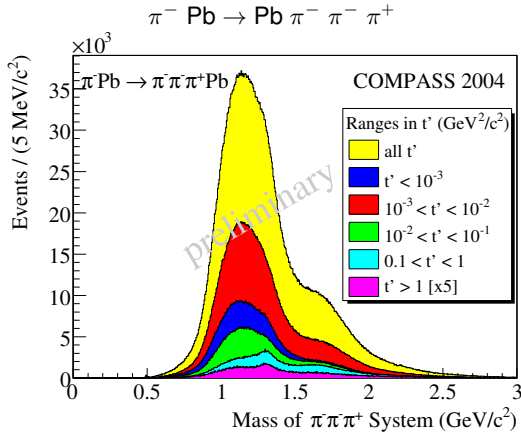
limited sensitivity to the polarisability contribution



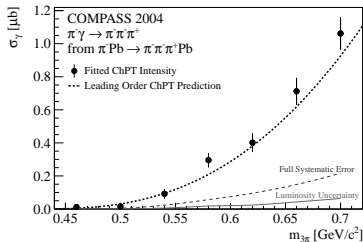
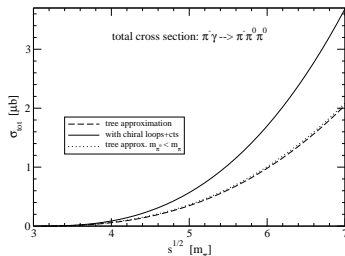
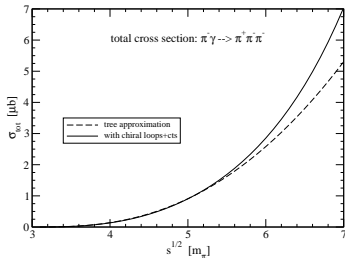


Compass projected pair probability $t_1=25u, t_2=50u$ W targets

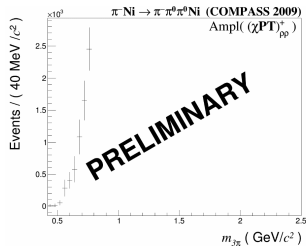




- "Low t' ": $10^{-3} (\text{GeV}/c)^2 < t' < 10^{-2} (\text{GeV}/c)^2 \sim 2\,000\,000$ events
- "Primakoff region": $t' < 10^{-3} (\text{GeV}/c)^2 \sim 1\,000\,000$ events



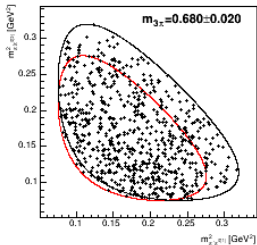
published in PRL 108 (2012) 192001



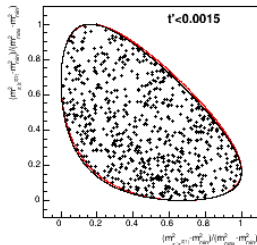
normalization: analysis ongoing



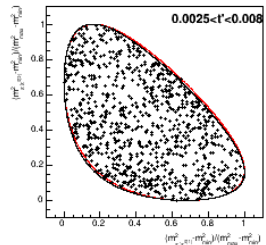
Dalitz plot for $\pi\pi^0\pi^0$



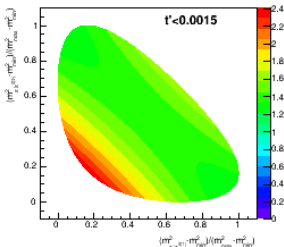
Reduced Dalitz plot for $\pi\pi^0\pi^0$



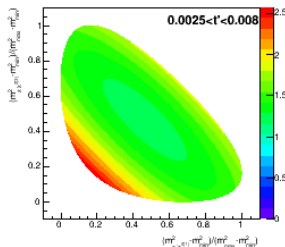
Reduced Dalitz plot for $\pi\pi^0\pi^0$



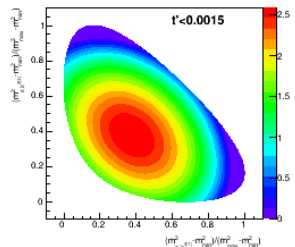
Fitted Reduced Dalitz plot, Primakoff

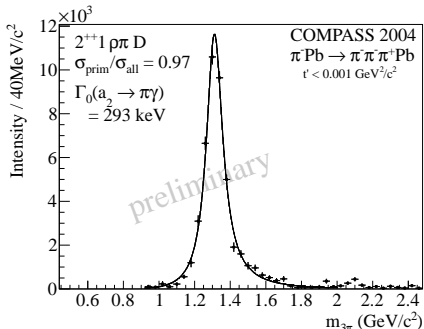


Fitted Reduced Dalitz plot, diffractive sidebin

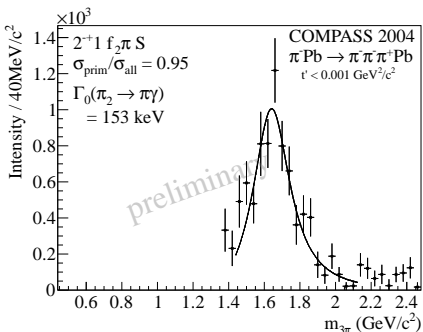


Fitted Reduced Dalitz plot, Primakoff-diffractive





$\Gamma_0(a_2(1320) \rightarrow \pi\gamma)$ **M2**



$\Gamma_0(\pi_2(1670) \rightarrow \pi\gamma)$ **E2**

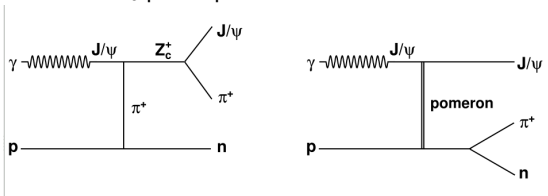
\Leftrightarrow meson w.f.'s: $\Gamma_{i \rightarrow f} \propto |\langle \Psi_f | e^{-i\vec{q} \cdot \vec{r}} \hat{\epsilon} \cdot \vec{p} | \Psi_i \rangle|^2$, VMD

- normalization via beam kaon decays
- large Coulomb correction

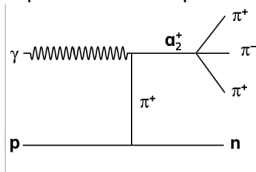
published in EPJ A50 (2014) 79



- does (charged-)pion exchange play a role at COMPASS collision energies?
- search for Z_c photo-production:



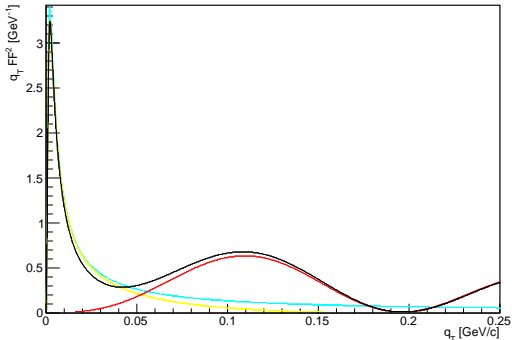
- a_2 photo-production via pion exchange:



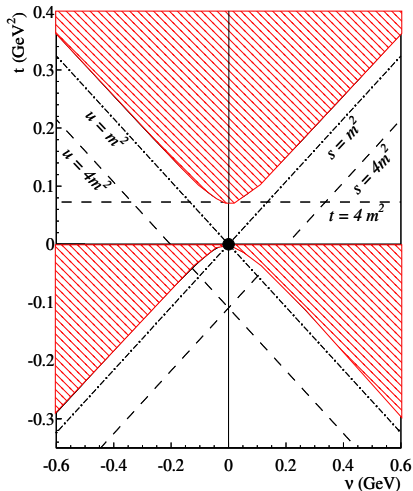
- constrain contribution from pionic trajectory



Photon density squared form factor



- calculation following G. Fäldt (Phys. Rev. C79, 014607)
- eikonal approximation: pions traverse Coulomb and strong-interaction potentials

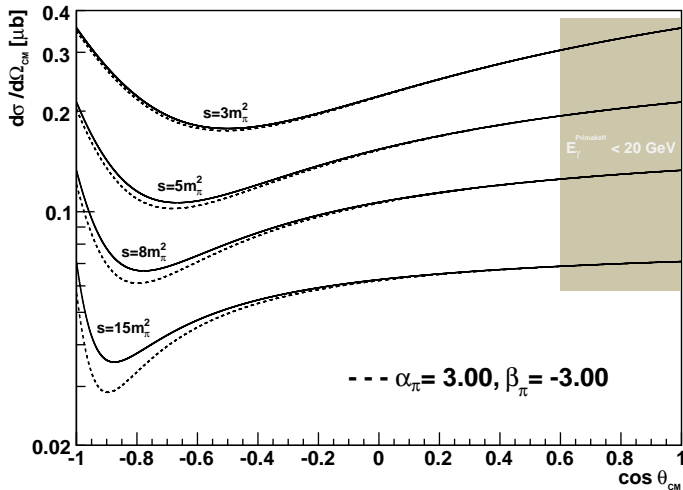


- ▶ **red hatched:**
physical regions
 $\gamma + \gamma \rightarrow \pi + \pi$
 $\gamma + \pi \rightarrow \gamma + \pi$
- ▶ two-pion thresholds
at $s = 4m_\pi^2$, $u = 4m_\pi^2$,
 $t = 4m_\pi^2$
- ▶ DR integration paths
 $t = 0$ (forward),
 $\theta = 180^\circ$ (backward)
 $u = m_\pi^2$, $s = m_\pi^2$, ...

from: D. Drechsel, talk at IWHSS 2011 Paris

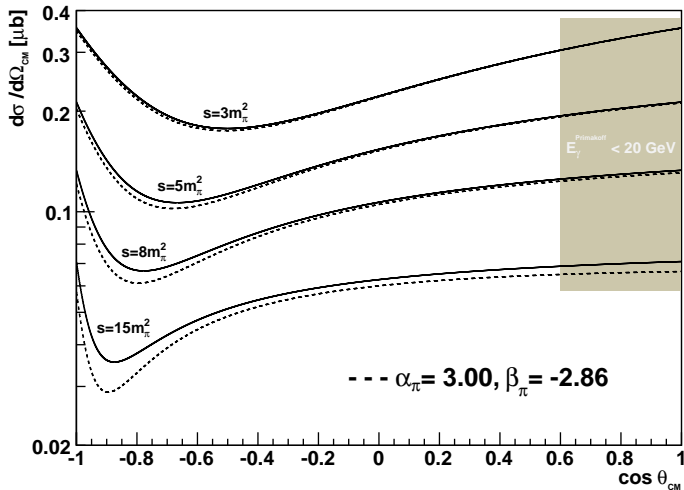


loop effects not shown



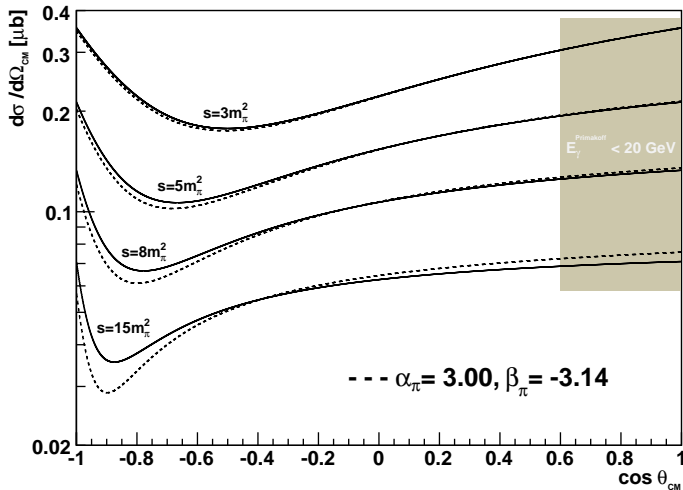


loop effects not shown



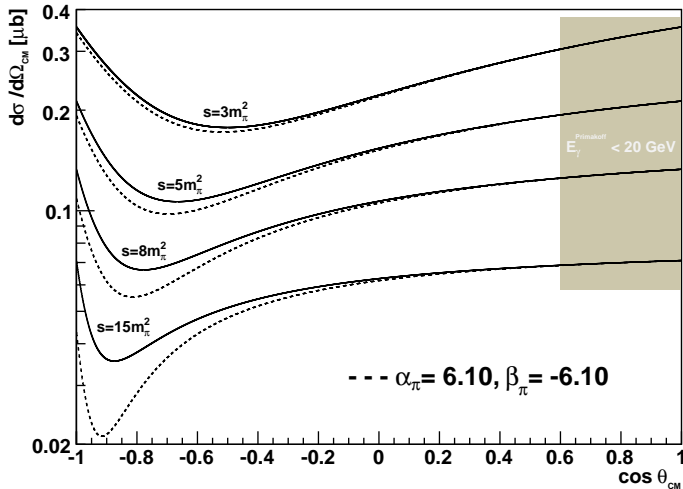


loop effects not shown





loop effects not shown





- Radiative π^+ production on the proton:



Mainz (2005) measurement: $\alpha_\pi - \beta_\pi = 11.6 \pm 1.5 \pm 3.0 \pm 0.5$

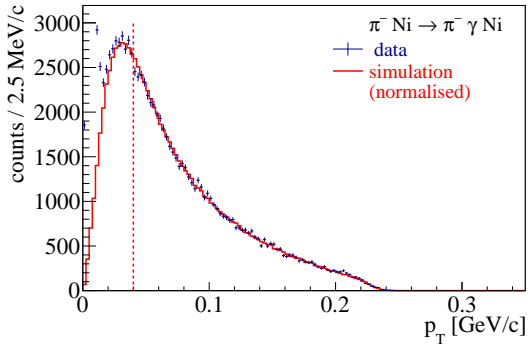
“ ± 0.5 ”: model error *only within the used ansatz*,
full systematics not under control

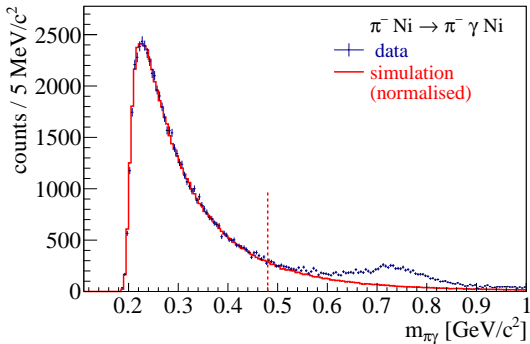
- Primakoff Compton reaction:



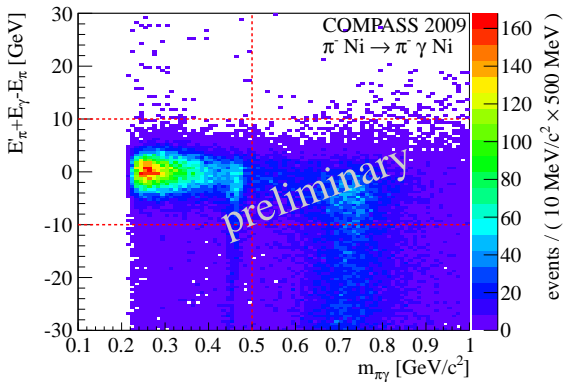
tiny extrapolation $\gamma^* \rightarrow \gamma \mathcal{O}(10^{-3} m_\pi^2)$
fully under theoretical control

[N. Kaiser, J.F., Nucl. Phys. A 812 (2008) 186]





- ρ contribution from $\pi\gamma \rightarrow \pi\pi^0$

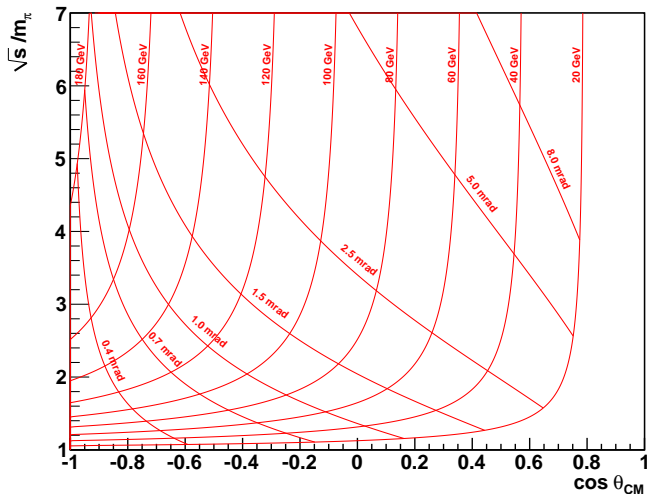


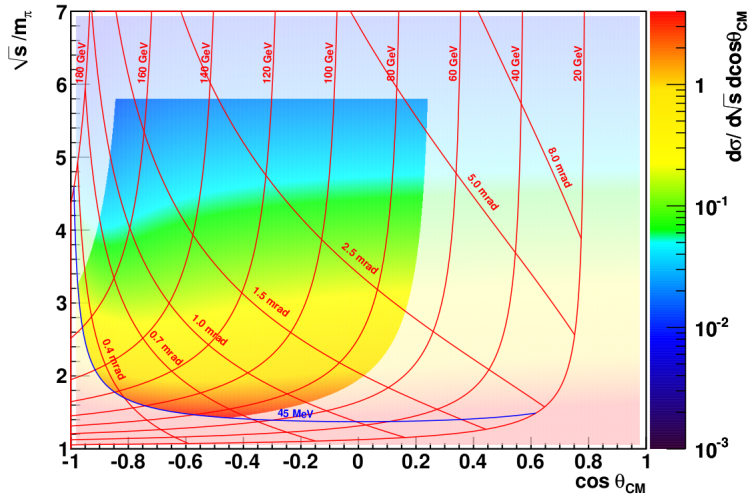
- ρ contribution from $\pi\gamma \rightarrow \pi\pi^0$



Mandelstam $\{s, t\} \leftrightarrow$ Laboratory $\{E_\gamma, \theta_\gamma\}$

for $\pi\gamma \rightarrow \pi\gamma$

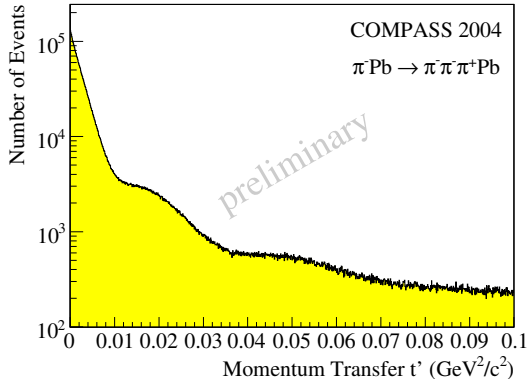




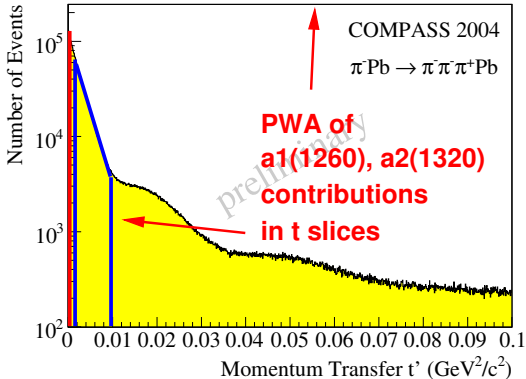


M.R. Pennington in the 2nd DAΦNE Physics Handbook,
“What we learn by measuring $\gamma\gamma \rightarrow \pi\pi$ at DAΦNE”:

All this means that the only way to measure the pion polarisabilities is in the Compton scattering process near threshold and not in $\gamma\gamma \rightarrow \pi\pi$. Though the low energy $\gamma\gamma \rightarrow \pi\pi$ scattering is seemingly close to the Compton threshold (...) and so the *extrapolation* not very far, the dominance of the pion pole (...) means that the energy scale for this continuation is m_π . Thus the polarisabilities cannot be determined accurately from $\gamma\gamma$ experiments in a model-independent way and must be measured in the Compton scattering region.



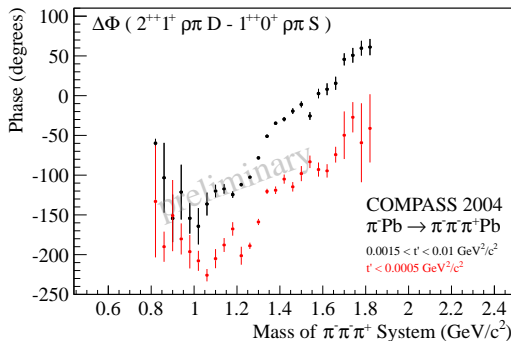
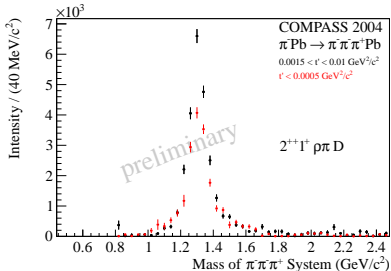
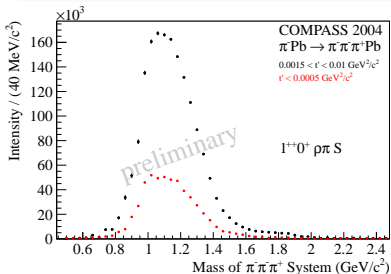
- "Low t' ": $10^{-3} (\text{GeV}/c)^2 < t' < 10^{-2} (\text{GeV}/c)^2$ $\sim 2\,000\,000$ events
- "Primakoff region": $t' < 10^{-3} (\text{GeV}/c)^2$ $\sim 1\,000\,000$ events

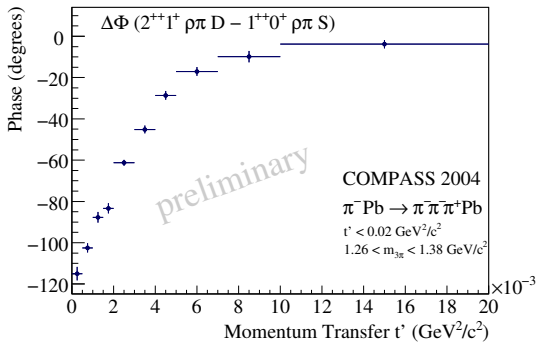
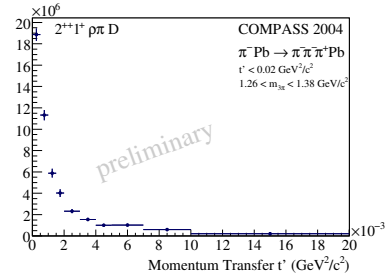
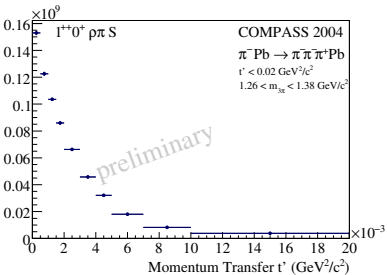


- "Low t' ": $10^{-3} (\text{GeV}/c)^2 < t' < 10^{-2} (\text{GeV}/c)^2$ $\sim 2\,000\,000$ events
- "Primakoff region": $t' < 10^{-3} (\text{GeV}/c)^2$ $\sim 1\,000\,000$ events



PWA: a_1 , a_2 and $\Delta\Phi$ in separated t' regions

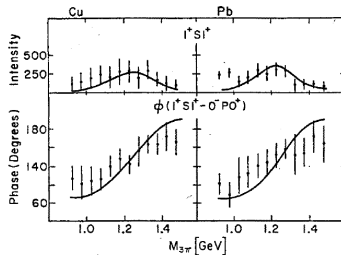
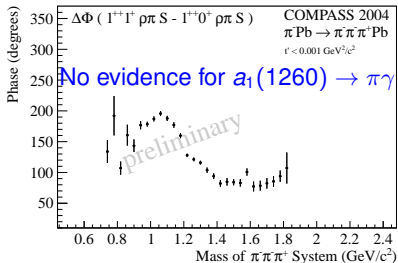
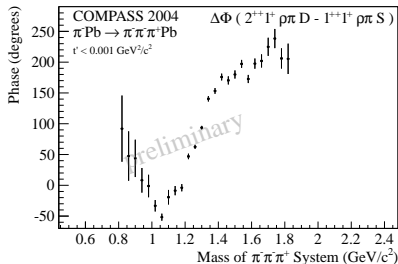
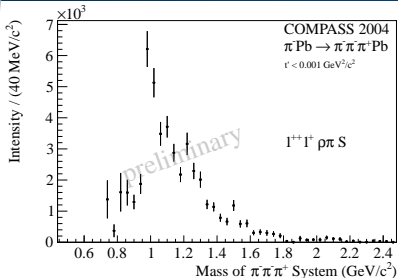




- transition of $\pi\gamma$ to $\pi IP \rightarrow a_2$ production
- work in progress
- interference can be used to map details of resonances and production mechanisms



Primakoff production of $a_1(1260)$ vs. E272 result



M. Zielinski et al, Phys. Rev. Lett 52 (1984) 1195



- **Mass-independent PWA** (narrow mass bins):

$$\sigma_{\text{indep}}(\tau, m, t') = \sum_{\epsilon=\pm 1} \sum_{r=1}^{N_r} \left| \sum_i T_{ir}^\epsilon f_i^\epsilon(t') \psi_i^\epsilon(\tau, m) / \sqrt{\int |f_i^\epsilon(t')|^2 dt'} \sqrt{\int |\psi_i^\epsilon(\tau', m)|^2 d\tau'} \right|^2$$

- Production strength assumed constant in single bins
- Decay amplitudes $\psi_i^\epsilon(\tau, m)$, with t' dependence $f_i^\epsilon(t')$
- Production amplitudes $T_{ir}^\epsilon \rightarrow$ Extended log-likelihood fit
- Acceptance corrections included
- **Spin-density matrix:** $\rho_{ij}^\epsilon = \sum_r T_{ir}^\epsilon T_{jr}^{\epsilon*}$

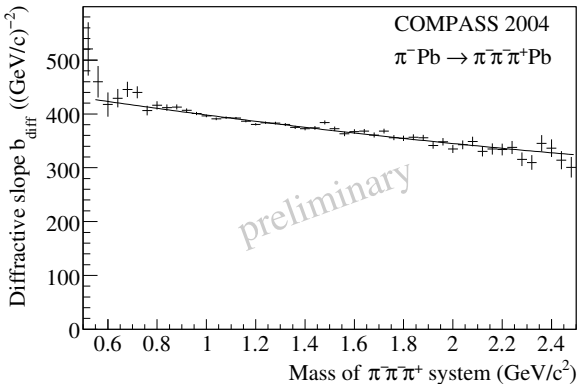
\rightarrow Physical parameters:

$$\text{Intens}_i^\epsilon = \rho_{ii}^\epsilon,$$

relative phase Φ_{ij}^e

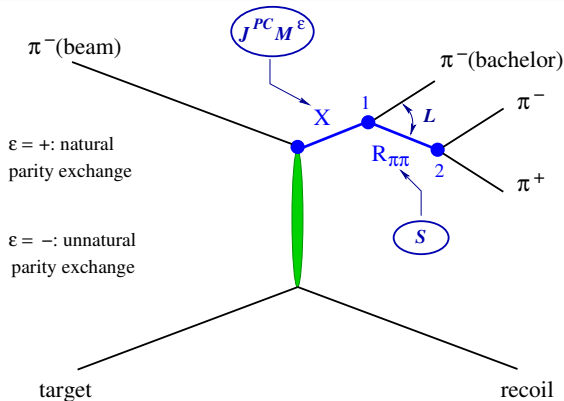
$$\text{Coh}_{i,j}^\epsilon = \sqrt{(\text{Re } \rho_{ij}^\epsilon)^2 + (\text{Im } \rho_{ij}^\epsilon)^2} / \sqrt{\rho_{ii}^\epsilon \rho_{jj}^\epsilon}$$

- **Mass-dependent χ^2 -fit** (not presented here):
 - X parameterized by Breit-Wigner (BW) functions
 - Background can be added





Isobar Model

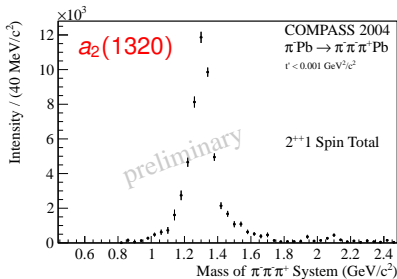
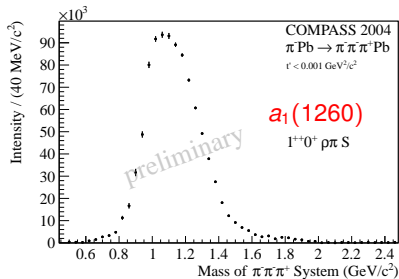
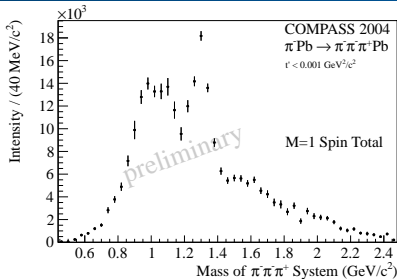
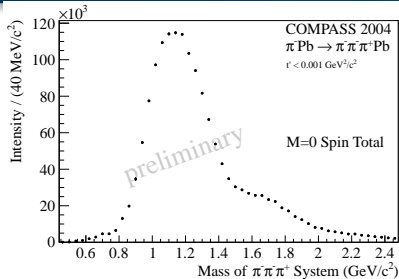


- Isobar model: Intermediate 2-particle decays
- Partial wave in reflectivity basis: $J^{PC} M^\epsilon [isobar] L$

- **Mass-independent** PWA ($40 \text{ MeV}/c^2$ mass bins): **38 waves**
Fit of angular dependence of partial waves, interferences
- **Mass-dependent** χ^2 -fit (Not presented here)



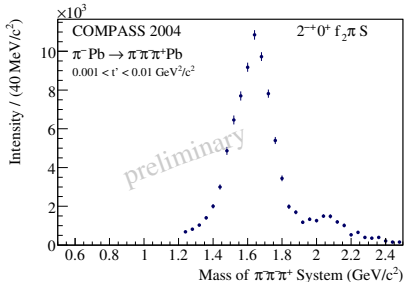
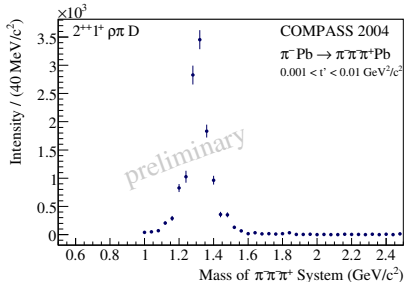
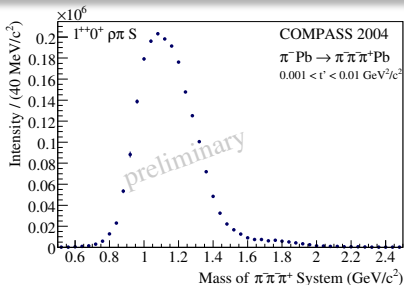
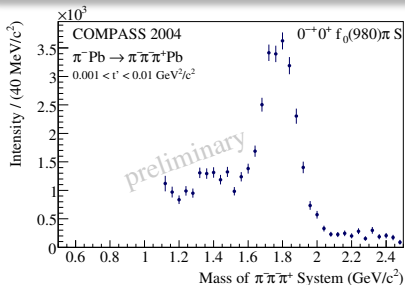
Major intensities in $m(3\pi)$ -bins (acceptance corrected)





PWA of data with low t'

Intensity of selected waves: $0^{-+}0^{+}f_0(980)\pi S$, $1^{++}0^{+}\rho\pi S$, $2^{++}1^{+}\rho\pi D$, $2^{-+}0^{+}f_2(1270)\pi S$





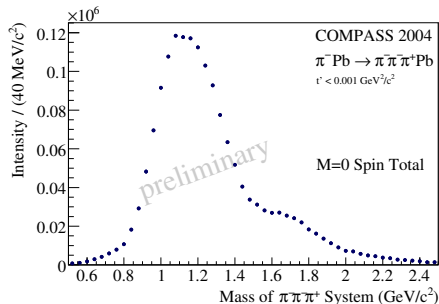
“Spin Totals”: Sum of all contributions for given M (i.e. z-projection of J)

t' -dependent amplitudes:

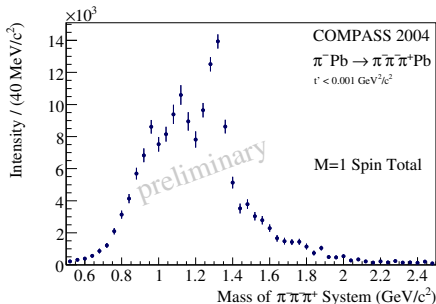
Primakoff production: $M=1: \sigma(t') \propto e^{-b_{\text{Prim}}t'} \rightarrow$ arises at $t' \approx 0$ (resolved shape!)

Diffractive production: $M=0: \sigma(t') \propto e^{-b_{\text{diff}}(m)t'}$

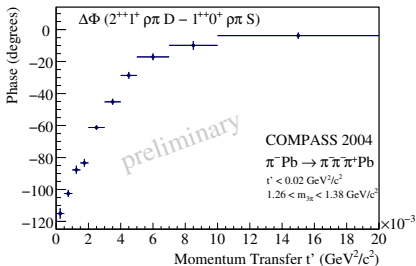
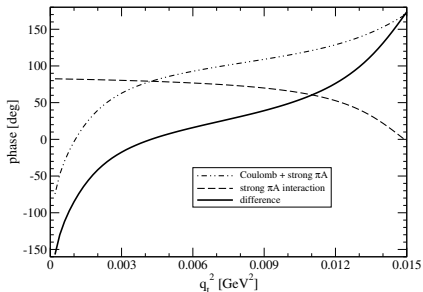
$M=1: \sigma(t') \propto t' e^{-b_{\text{diff}}(m)t'} \rightarrow$ vanishes for $t' \approx 0$



M=0



M=1



Glauber modell

G. Fäldt and U. Tengblad, Phys. Rev. C79, 014607 (2009)

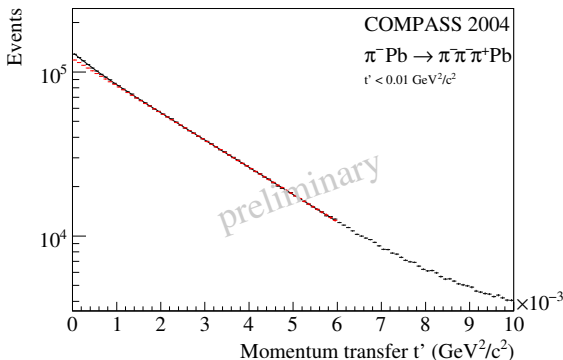
Plot: N. Kaiser (TU München)

- ⇒ indicates confirmation of interference Coulomb-interaction - strong interaction
- ⇒ detailed studies of the nature of resonances



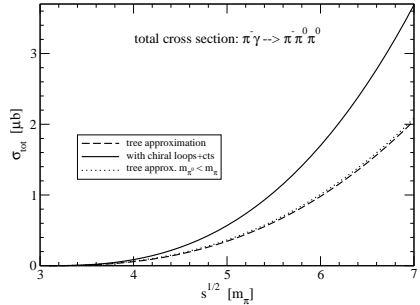
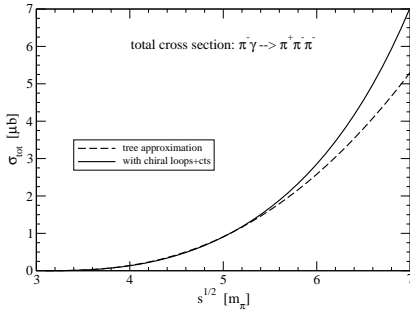
Primakoff: $\sigma(t') \propto e^{-b_{\text{Prim}} t'}$, $b_{\text{Prim}} \approx 2000 \text{ (GeV/c)}^{-2}$ (mainly resolution)

Diffractive: $\sigma(t') \propto e^{-b_{\text{diff}} t'}$, $b_{\text{diff}} \approx 400 \text{ (GeV/c)}^{-2}$ for lead target



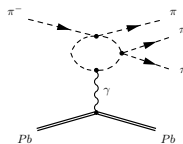
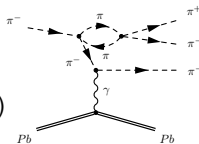
(Mass) spectrum of this Primakoff contribution?

⇒ Statistical subtraction of diffractive background (for bins of $m_{3\pi}$)

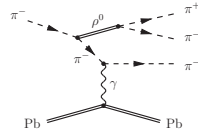


Chiral loops, e.g.

(N. Kaiser, NPA848 (2010) 198)

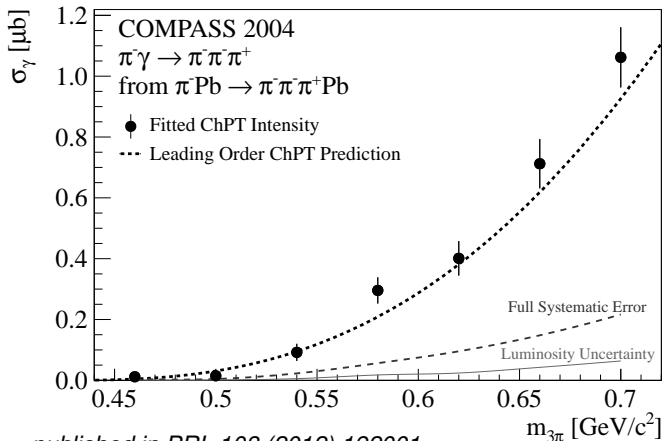
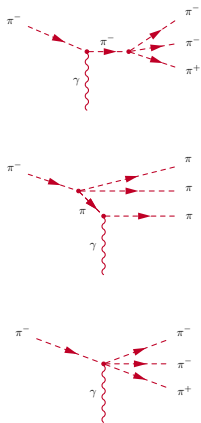


ρ terms:





Measured absolute cross-section of $\pi^-\gamma \rightarrow \pi^-\pi^-\pi^+$



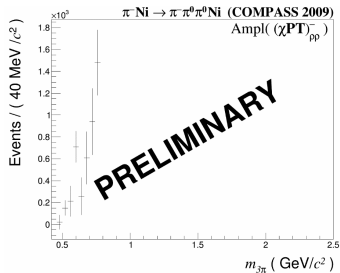
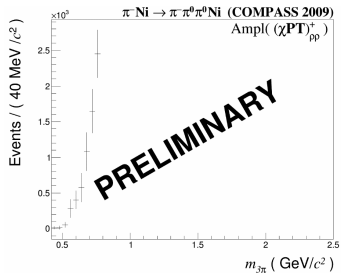
published in PRL 108 (2012) 192001



$$\pi^- \gamma \rightarrow \pi^- \pi^0 \pi^0$$

Partial Wave Analysis

Isobaric Model – Chiral Wave





$$\pi^- \gamma \rightarrow \pi^- \pi^0 \pi^0$$

Partial Wave Analysis

Chiral Model - Amplitudes

