

# Spin structure of the proton at low $x$ and low $Q^2$ in two-dimensional bins from COMPASS

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Acknowledgements: Ana Sofia Nunes (main analysis investigator)

# Motivation

- low  $x \leftrightarrow$  high parton densities
- low  $x$  and low  $Q^2 \leftrightarrow$  **transition from the regime of photoproduction to the regime of DIS (described by pQCD)**
- theoretical predictions for  $g_1^p$  as function of two kinematic variables:
  - ▶ Badełek et al., Eur.Phys.J. C26 (2002) 45  
*"Spin structure function  $g_1(x, Q^2)$  and the DHGHY integral  $I(Q^2)$  at low  $Q^2$ : Predictions from the GVMD model"*
  - ▶ Ermolaev et al., Eur.Phys.J. C58 (2008) 29  
*"Comment on the recent COMPASS data on the spin structure function  $g_1$ "*
  - ▶ Ermolaev et al., Riv.Nuovo Cim. 33 (2010) 57  
*"Overview of the spin structure function  $g_1$  at arbitrary  $x$  and  $Q^2$ "*  
"one can parameterize  $g_1$  by the set of variables  $x$ ,  $Q^2$  or, alternatively,  $\omega[\equiv 2pq = 2M(E - E')]$ ,  $Q^2$ , or  $\nu$ ,  $Q^2$ "
- COMPASS'  $\sim 7 \times 10^8$  events allow a 2D extraction
  - ▶ extraction, for the first time, in 4 2D grids:  $(x, Q^2)$ ,  $(\nu, Q^2)$ ,  $(x, \nu)$ ,  $(Q^2, x)$

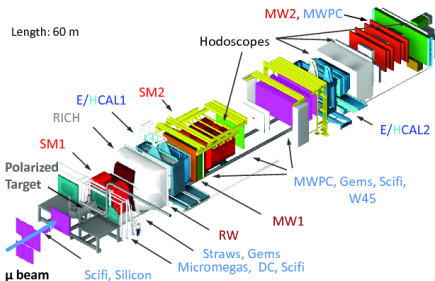
# The COMPASS experiment at CERN

## COMPASS @ CERN

**CO**mmun **M**uon **P**roton  
**A**pparatus for **S**tructure  
and **S**pectroscopy

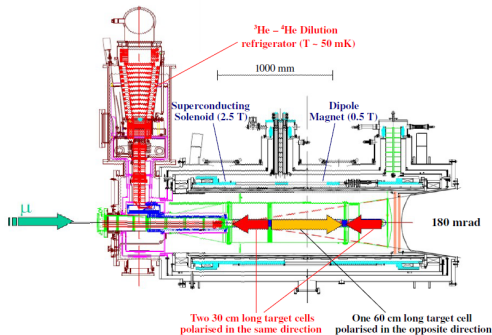


- **Fixed target experiment** at the SPS using a tertiary **muon beam**
- Collaboration of about 200 members from 11 countries and 23 institutions



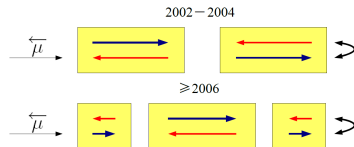
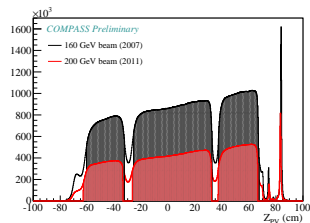
- 160/200 GeV  $\mu^+$   
**polarised beam**,  $P_b \sim -80\%$
- ${}^6\text{LiD}$  or  $\text{NH}_3$ , 1.2 m long, **polarised target**  
@ 2.5 T and 60 mK,  $P_{\text{target}} \sim 50/85\%$
- large acceptance, two staged spectrometer
- tracking, calorimetry, PID

# Polarised target

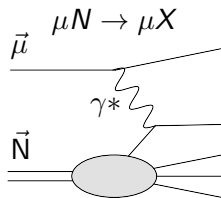


Material	Dilution factor (f)	Polarisation ( $P_{\text{target}}$ )
<sup>6</sup> LiD	0.40	50%
NH <sub>3</sub>	0.16	85%

## Vertex coordinate $z_{PV}$



# Reaction of interest



$Q^2$ : photon virtuality

$x$ : Bjorken scaling variable

(fraction of nucleon momentum carried by the struck quark)

$\nu$ : Virtual photon energy

Observables:  $A_1^P$  and  $g_1^P$

Related to  $\Delta q$  (high  $Q^2$ ) + non perturbative mechanisms (low  $Q^2$ )

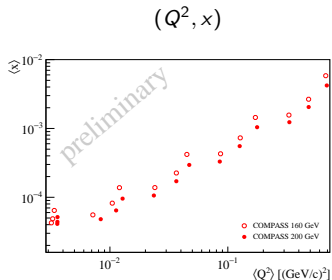
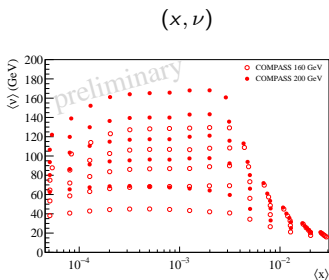
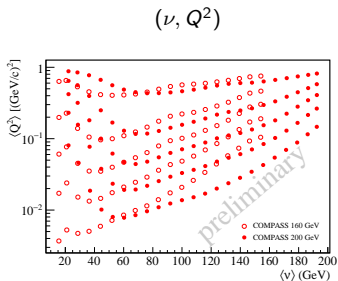
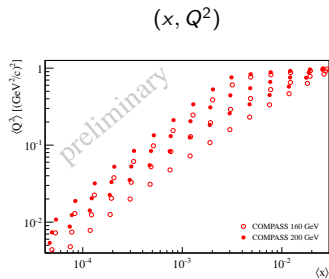
# Data samples for the extraction of $A_1^p$ and $g_1^p$

- Longitudinally polarised target ( $\text{NH}_3$ ):  **$676 \times 10^6$  events**  
( $447 \times 10^6$  with 160 GeV beam in 2007,  $229 \times 10^6$  with 200 GeV beam in 2011)
- Before, SMC low  $x$ , low  $Q^2$  proton data:  $4.5 \times 10^6$  events  
 $\Rightarrow$  The COMPASS data set has **150** $\times$  more events than SMC

## Main selection criteria:

- at least one additional track (besides the scattered muon) in the interaction point ("hadron method") - SMC proved there is no bias to the inclusive asymmetries at low  $x$
- not a  $\mu e$  elastic scattering event
- $Q^2 < 1 \text{ (GeV}/c)^2$
- $x \geq 4 \times 10^{-5}$
- $0.1 < y < 0.9$

# Phase-space coverage of the 2D analysis



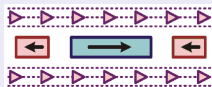
# Double longitudinal spin asymmetry measurement

## Asymmetry extraction at COMPASS

$$A_1 = \frac{1}{|P_B P_T| f D} \left( \frac{N_{\rightarrow\rightarrow}^{\leftarrow} - N_{\leftarrow\leftarrow}^{\leftarrow}}{N_{\rightarrow\leftarrow}^{\leftarrow} + N_{\leftarrow\rightarrow}^{\leftarrow}} \right)$$

Simultaneous recording of the two spin states in oppositely polarised target cells

## COMPASS target





# Double longitudinal spin asymmetry measurement

## Asymmetry extraction at COMPASS

$$A_1 = \frac{1}{|P_B P_T| f D} \left( \frac{N_{\rightarrow\rightarrow}^{\leftarrow} - N_{\leftarrow\leftarrow}^{\leftarrow}}{N_{\leftarrow\rightarrow}^{\leftarrow} + N_{\rightarrow\leftarrow}^{\leftarrow}} \right)$$

Simultaneous recording of the two spin states in oppositely polarised target cells

## COMPASS target



→ Reversal by field rotation every 24h to cancel out acceptance difference

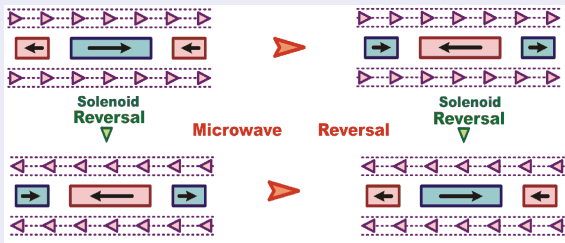
# Double longitudinal spin asymmetry measurement

## Asymmetry extraction at COMPASS

$$A_1 = \frac{1}{|P_B P_T| f D} \left( \frac{N_{\rightarrow\leftarrow} - N_{\leftarrow\rightarrow}}{N_{\rightarrow\leftarrow} + N_{\leftarrow\rightarrow}} \right)$$

Simultaneous recording of the two spin states in oppositely polarised target cells

## COMPASS target



- Reversal by field rotation every 24h to cancel out acceptance difference
- Reversal by micro-wave once in a while to cancel out acceptance/field correlation

# Double longitudinal spin asymmetry $A_1^p$

- Each event is given a weight  $\omega = \mathbf{f} \cdot \mathbf{D} |\mathbf{P}_{\text{beam}}|$  to optimize the statistical errors
- Unpolarised radiative corrections (RC), included in the dilution factor, from TERAD  
[A.A. Akhundov, *et al.*, Fortschr. Phys. 44 (1996) 373]
- Polarised radiative corrections ( $A^{\text{RC}} \leq 0.25 \delta A_1^{\text{stat}}$ ) from POLRAD  
[I. Akushevich *et al.*, Comput.Phys.Commun. 104 (1997) 201]
- Corrected for polarisable  $^{14}\text{N}$  ( $A^{^{14}\text{N}} \leq 0.01 \delta A_1^{\text{stat}}$ )
- Thorough checks on possible sources of false asymmetries  $\Rightarrow$  systematic uncertainties similar to the statistical errors

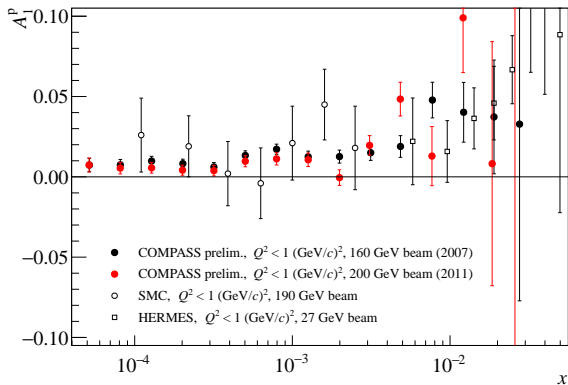
# Spin dependent structure function $g_1^p$

- The structure function is obtained in bins of  $x$  or  $\nu$  according to:

$$g_1^p(\langle x \rangle, \langle Q^2 \rangle) = \frac{F_2^p(\langle x \rangle, \langle Q^2 \rangle)}{2x[1 + R(\langle x \rangle, \langle Q^2 \rangle)]} A_1^p(\langle x \rangle, \langle Q^2 \rangle)$$

- $F_2^p(\langle x \rangle, \langle Q^2 \rangle)$  from the SMC fit on data or from a model (for low  $x$  and  $Q^2$ )  
[SMC, Phys.Rev. D58 (1998), 112001; B. Badełek & J. Kwieciński, Phys.Lett. B295 (1992) 263]
- $R(\langle x \rangle, \langle Q^2 \rangle)$  based on SLAC parameterization, extended to low  $Q^2$   
[COMPASS, PLB 647 (2007) 330]

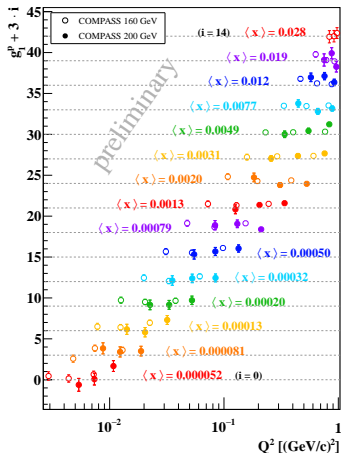
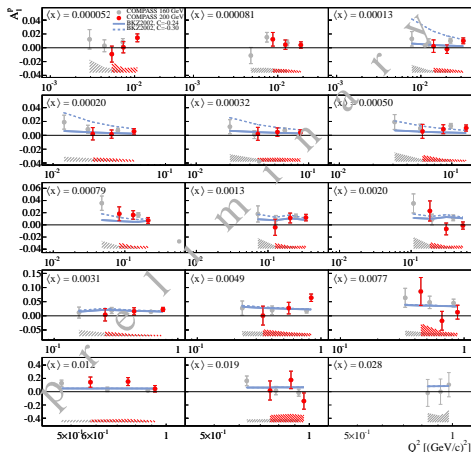
# $A_1^P(x)$ & comparison with previous experiments



- results for the two beam energies are compatible within errors
- systematic errors are similar to the statistical errors (not shown here)
- $A_1^P$  is **significantly positive**
- no dependence on  $x$  is seen (nor on  $\nu$ , not shown here)
- the COMPASS results **improve the precision** of the measurement

# $A_1^P$ and $g_1^P$ at low $x$ and low $Q^2$ : results for the grid ( $x, Q^2$ )

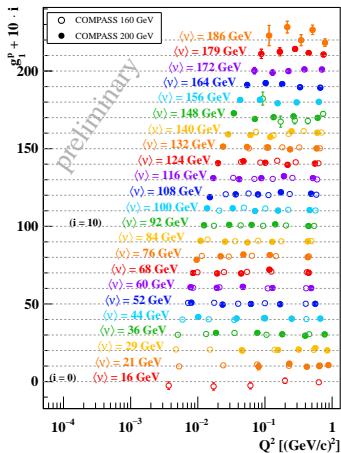
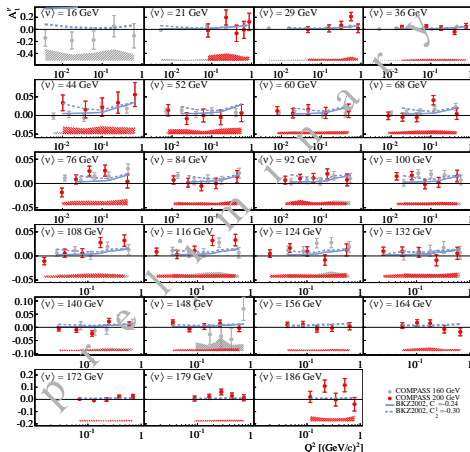
Data: 2007&2011,  $\mu^+p \rightarrow \mu^+X$



- **no strong dependence** on  $x$  or  $Q^2$
- results **compatible with theoretical model (GVMD)** [Eur.Phys.J. C26 (2002) 45]

# $A_1^P$ and $g_1^P$ at low $x$ and low $Q^2$ : results for the grid ( $\nu, Q^2$ )

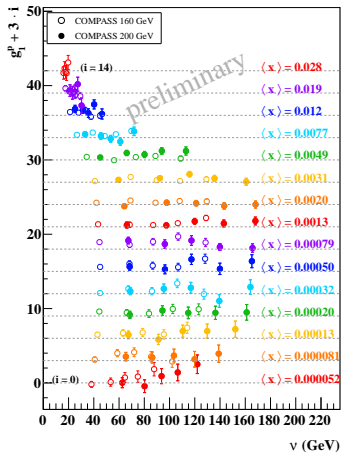
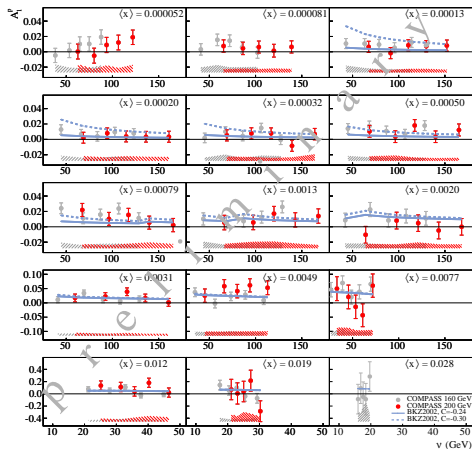
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Data: 2007&2011,  $\mu^+p \rightarrow \mu^+X$

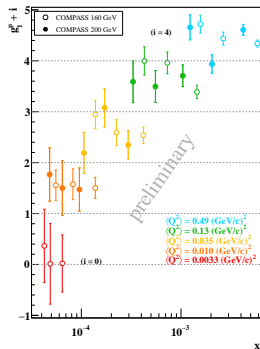
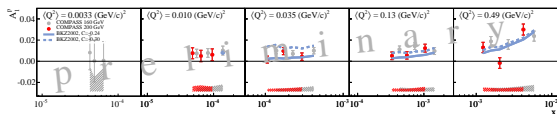


- **no strong dependence** on  $\nu$  or  $x$
- results **compatible with theoretical model (GVMD)** [Eur.Phys.J. C26 (2002) 45]



# $A_1^p$ and $g_1^p$ at low $x$ and low $Q^2$ : results for the grid ( $Q^2, x$ )

Data: 2007&2011,  $\mu^+ p \rightarrow \mu^+ X$



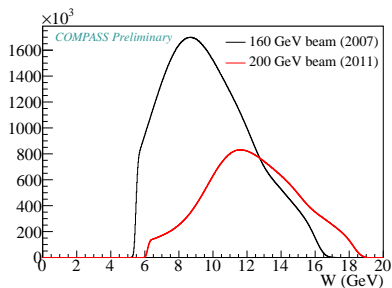
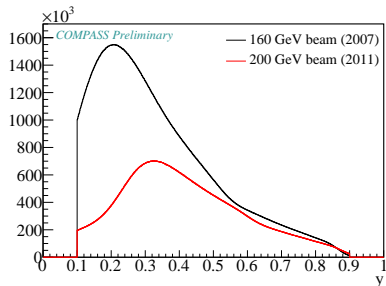
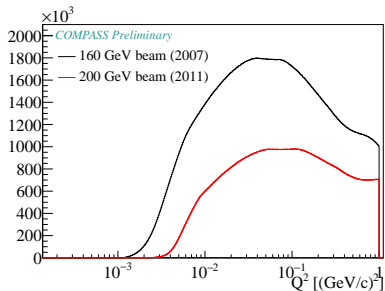
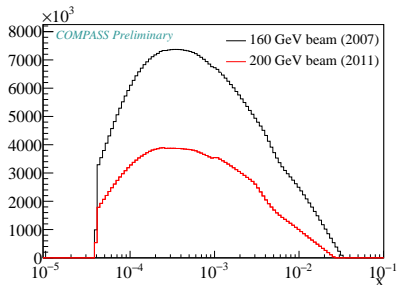
- no strong dependence on  $x$  or  $Q^2$
- results compatible with theoretical model (GVMD) [Eur.Phys.J. C26 (2002) 45]

# Summary and outlook

- Longitudinal double spin asymmetries  $A_1^P$  and the spin dependent structure function  $g_1^P$  extracted in 4 two-dimensional grids:
  - ▶  $(x, Q^2)$
  - ▶  $(\nu, Q^2)$
  - ▶  $(\nu, x)$
  - ▶  $(Q^2, x)$
- **Positive** spin asymmetries at very low  $x$
- **No significant dependence** on studied kinematic variables
- **Compatibility with GVMD model** predictions [Eur.Phys.J. C26 (2002) 45]

# BACKUP

# Characteristics of the final sample



# GVMD model [Eur.Phys.J. C26 (2002) 45]

[Badelek et al., Eur.Phys.J. C26 (2002) 45]

$$\begin{aligned}g_1(x, Q^2) &= g_1^\perp(x, Q^2) + g_1^{\text{AS}}(\bar{x}, Q^2 + Q_0^2) \\&= C \left[ \frac{4}{9} (\Delta u_{\text{val}}^{(0)}(x) + \Delta \bar{u}^{(0)}(x)) \right. \\&\quad \left. + \frac{1}{9} (\Delta d_{\text{val}}^{(0)}(x) + \Delta \bar{d}^{(0)}(x)) \right] \frac{M_\rho^4}{(Q^2 + M_\rho^2)^2} \\&\quad + C \left[ \frac{1}{9} (2\Delta \bar{s}^{(0)}(x)) \right] \frac{M_\phi^4}{(Q^2 + M_\phi^2)^2} \\&\quad + g_1^{\text{AS}}(\bar{x}, Q^2 + Q_0^2).\end{aligned}\tag{5}$$

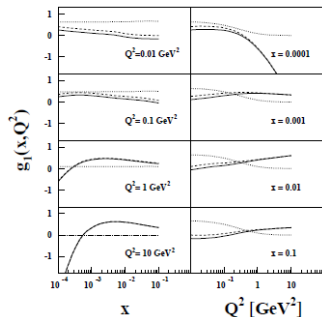
To obtain the value of  $C$  from (12), the contribution of resonances was evaluated using the preliminary data taken at ELSA/MAMI by the GDH Collaboration [16] at the photoproduction, for  $W_t = 1.8$  GeV. The asymptotic part of  $g_1$  was parametrized using the GRSV2000 fit for the “standard scenario” of polarized parton distributions with a flavor symmetric light sea,  $\Delta \bar{u} = \Delta \bar{d} = \Delta s = \Delta \bar{s}$ , at the NLO accuracy [9]. The non-perturbative parton distributions,  $\Delta p_j^{(0)}(x)$ , in the light vector meson component of  $g_1$ , (3), were evaluated at fixed  $Q^2 = Q_0^2$ , using, either

- (i) the GRSV2000 fit, or
- (ii) a simple, “flat” input:

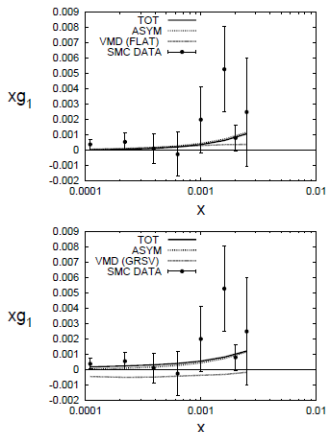
$$\Delta p_i^{(0)}(x) = N_i(1-x)^{\eta_i},\tag{13}$$

with  $\eta_{u_v} = \eta_{d_v} = 3$ ,  $\eta_u = \eta_s = 7$  and  $\eta_g = 5$ . The normalization constants  $N_i$  were determined by imposing the Bjorken sum rule for  $\Delta u_v^{(0)} - \Delta d_v^{(0)}$ , and requiring that the first moments of all other distributions are the same as those determined from the QCD analysis [18]. It was checked that the parametrization (13) combined with the unified equations gives a reasonable description of the SMC data on  $g_1^{\text{NS}}(x, Q^2)$  [19] and on  $g_1^p(x, Q^2)$  [5]. This fit was also used to investigate the magnitude of the double logarithmic corrections,  $\ln^2(1/x)$ , to the spin structure function of the proton at low  $x$  [20]. We have assumed  $Q_0^2 = 1.2$  GeV<sup>2</sup>, cf. (1) and (3), in accordance with the analysis of  $F_2$  [7, 8]. As a result the constant  $C$  was found to be  $-0.30$  in case (i) and  $-0.24$  in case (ii). These values change at most by 13% when  $Q_0^2$  changes in the interval  $1.0 < Q_0^2 < 1.6$  GeV<sup>2</sup>.

# GVMD model predictions [Eur.Phys.J. C26 (2002) 45]



**Fig. 1.** Values of  $g_1$  for the proton as a function of  $x$  and  $Q^2$ . The asymptotic contribution,  $g_1^{\text{AS}}$ , is marked with broken lines, the VMD part,  $g_1^{\text{V}}$ , with dotted lines and the continuous curves mark their sum, according to (5)



**Fig. 2.** Values of  $xg_1$  for the proton as a function of  $x$  at the measured values of  $Q^2$  in the non-resonant region,  $x < x_t = Q^2/2Mv_t(Q^2)$ . The upper plot corresponds to the VMD part parametrized using (13), the lower plot corresponds to the GRSV parametrization [9] of the VMD input. The  $g_1^{\text{AS}}$  in both plots has been calculated using the GRSV fit for standard scenario at the NLO accuracy. The contributions of the VMD and of the  $xg_1^{\text{AS}}$  are shown separately. Points are the SMC measurements at  $Q^2 < 1 \text{ GeV}^2$  [3]; errors are total. The curves have been calculated at the measured  $x$  and  $Q^2$  values