

# Di-hadron asymmetries and interplay between transversity induced asymmetries in hadron leptonproduction at COMPASS

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**COMmmon  
Muon and  
Proton  
Apparatus for  
Structure and  
Spectroscopy**

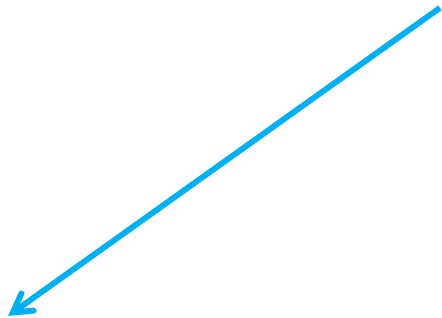
Collaboration  
~ 250 physicists  
28 institutions  
12 countries





# Common Muon and Proton Apparatus for Structure and Spectroscopy

wide physics program carried on using both **muon** and hadron beam



luminosity:  $\sim 5 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$   
 beam intensity:  $2 \cdot 10^8 \mu^+/\text{spill}$  (4.8s/16.2s)  
 beam momentum: 160 GeV/c

|   |   |                      |       |             |                 |              |
|---|---|----------------------|-------|-------------|-----------------|--------------|
| <b>longitudinally polarized muon beam</b> | deuteron ( <sup>6</sup> LID) polarized target | 2002<br>2003<br>2004 | } L/T | hadron beam | nuclear targets | 2004         |
|   | proton (NH <sub>3</sub> ) polarized target    | 2006                 | L     |             | LH target       | 2008<br>2009 |
|   | H <sub>2</sub> target                         | 2007                 | L/T   |             | 2012            |              |
|   |   | 2010                 | T     |             | 2014            |              |
|   |   | 2011                 | L     |             | 2015            |              |
|   |   | 2012                 |       |             | polarised DY    |              |
|   |   |                      |       |             |                 |              |
|   |   |                      |       |             |                 |              |



# the COMPASS spectrometer



- high energy beams
- large angular acceptance
- broad kinematical range

two stages spectrometer

Large Angle Spectrometer (SM1)

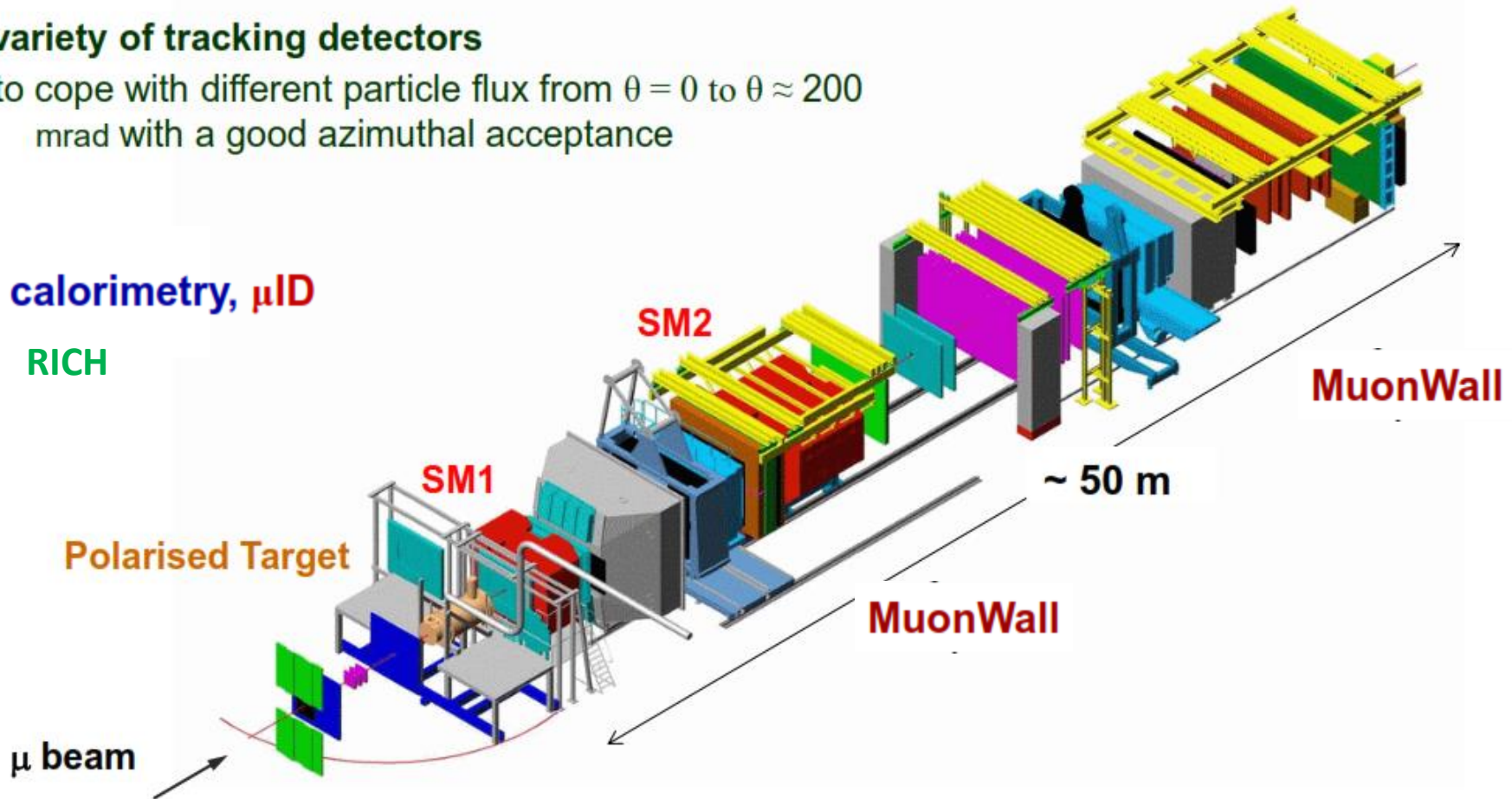
Small Angle Spectrometer (SM2)

variety of tracking detectors

to cope with different particle flux from  $\theta = 0$  to  $\theta \approx 200$  mrad with a good azimuthal acceptance

calorimetry,  $\mu$ ID

RICH



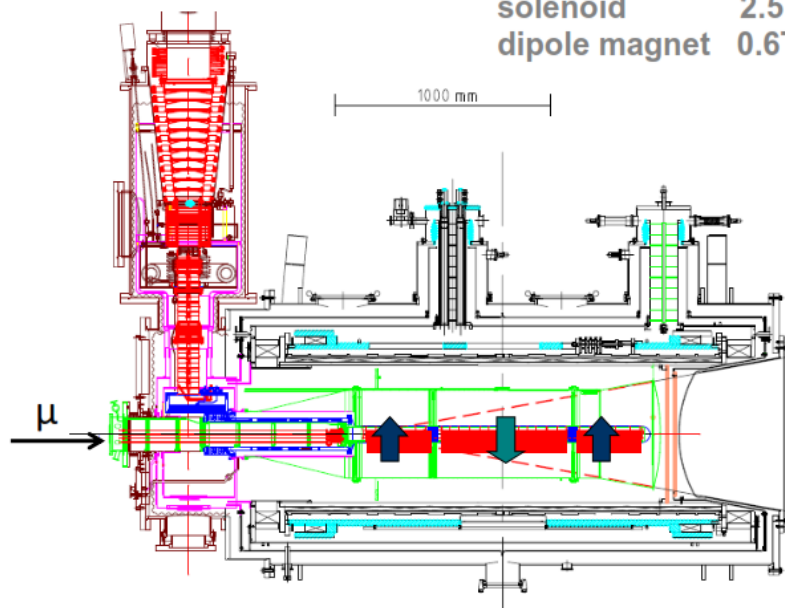
# polarized target system (>2005)

*solid state target operating in frozen spin mode*

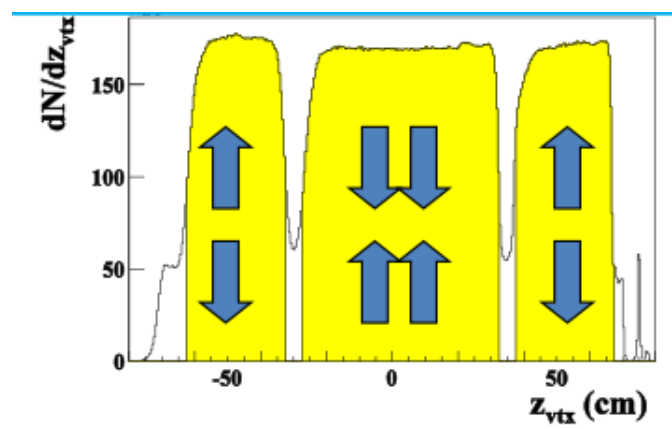
|                        |                            |                           |
|------------------------|----------------------------|---------------------------|
|                        | <b>d (<sup>6</sup>LiD)</b> | <b>p (NH<sub>3</sub>)</b> |
| <b>polarization</b>    | <b>50%</b>                 | <b>90%</b>                |
| <b>dilution factor</b> | <b>40%</b>                 | <b>16%</b>                |

<sup>3</sup>He – <sup>4</sup>He dilution refrigerator (T~50mK)

solenoid 2.5T  
dipole magnet 0.6T



*3 cells target with opposite polarizations*



**2 configurations:  
polarisation reversed each week  
to minimize possible systematic errors**

**the results shown in these slides are obtained from:**

- **analysis performed on charged hadrons (unidentified)**
- **data taken during 2010**
- **on a transversely polarised proton target (NH<sub>3</sub>)  
and using a 160 GeV positive muon beam**

# SIDIS: a key process to investigate the structure of the nucleon

lepton interacts with a **single constituent** of the nucleon ( $Q^2 > 1 \text{ GeV}^2/c^2$ )

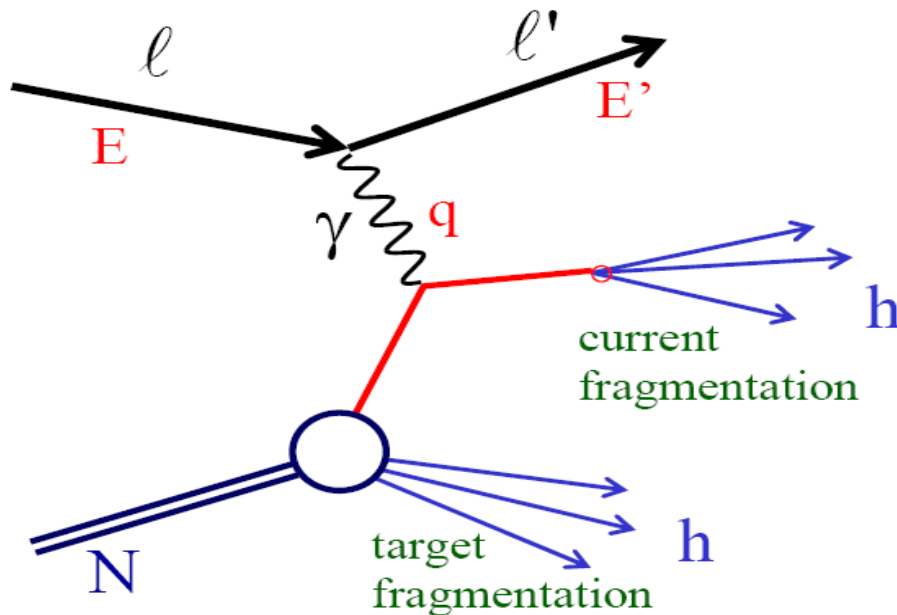
$$q = \ell - \ell'$$

$$Q^2 = -q^2 \quad W^2 = (P + q)^2$$

$$x = \frac{Q^2}{2P \cdot q} \quad \text{Bjorken scaling variable}$$

$$y = \frac{P \cdot q}{P \cdot \ell} = \frac{E - E'}{E}$$

$$z = \frac{P \cdot P_h}{P \cdot q} = \frac{E_h}{E - E'}$$



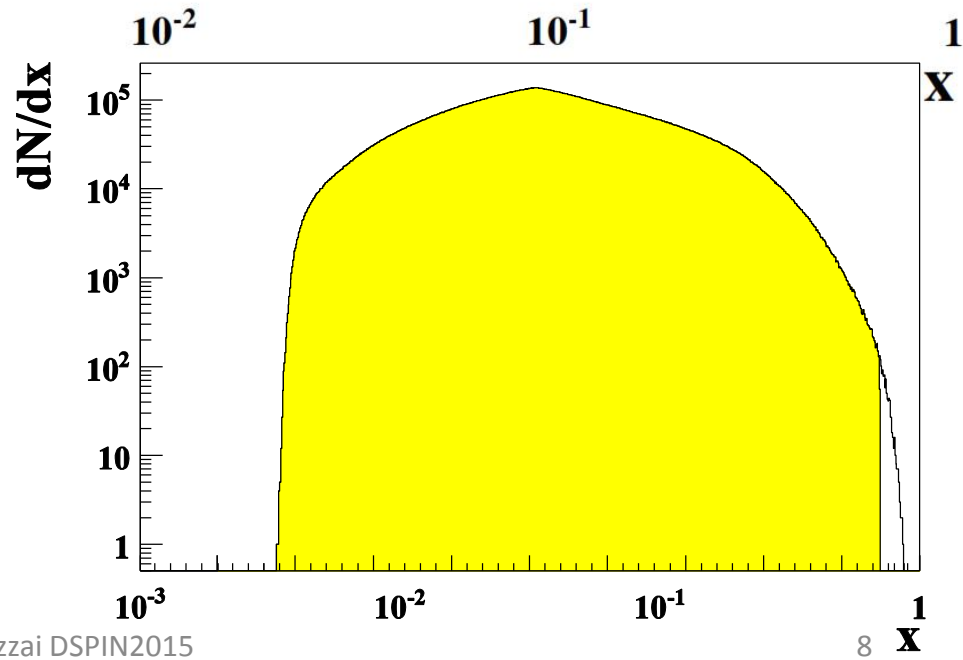
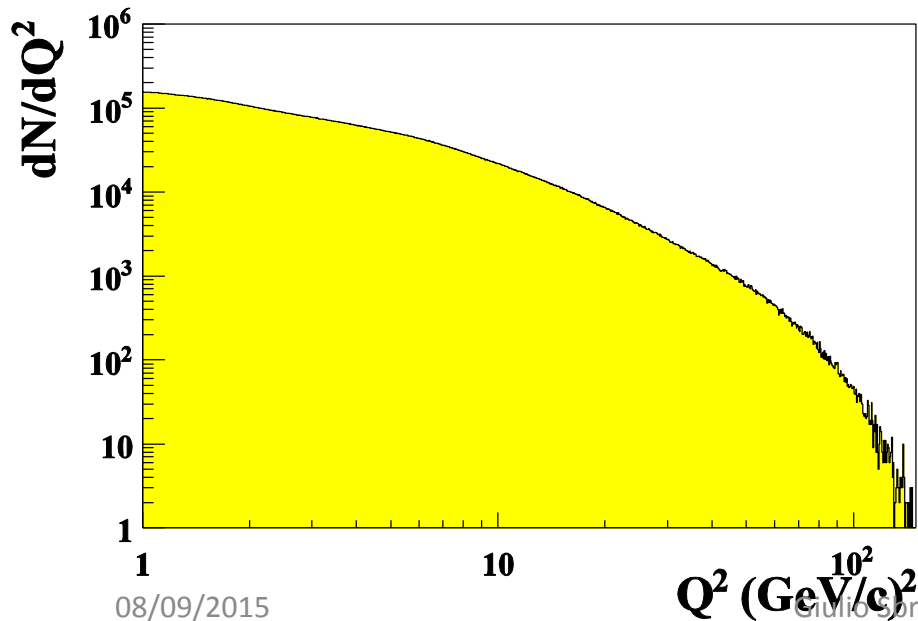
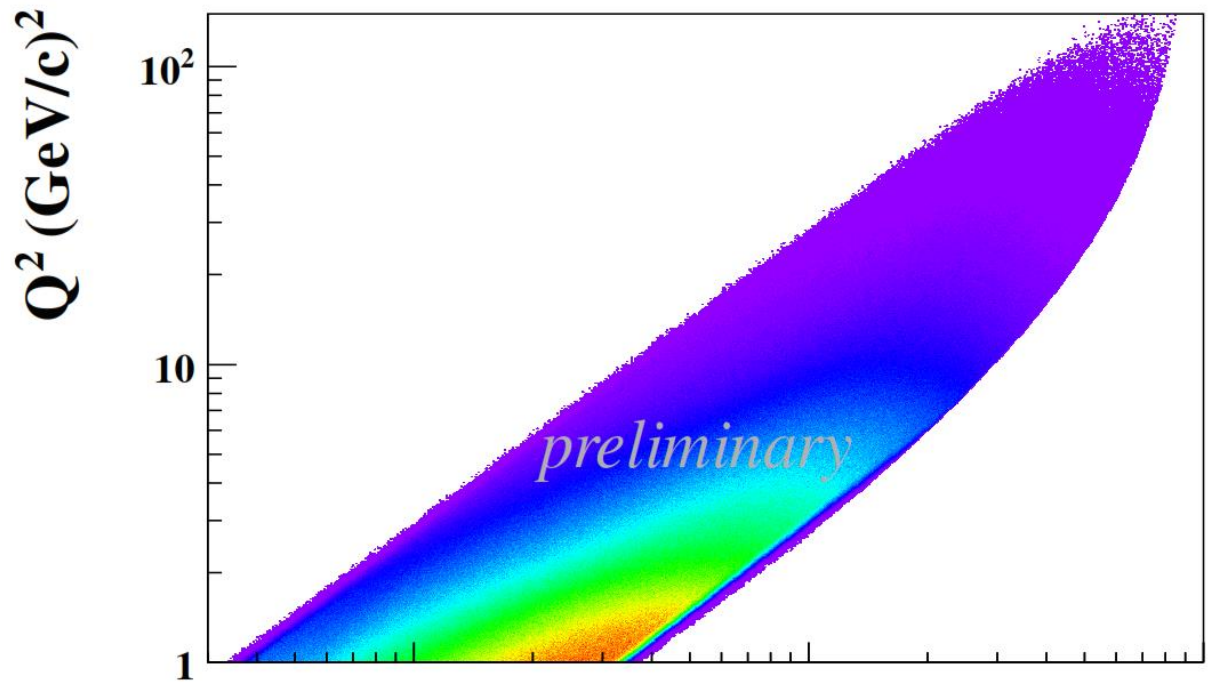
at least one hadron is detected in the final state  
(information on the **struck quark**)

## DIS event selection

$$Q^2 > 1 \text{ (GeV/c)}^2$$

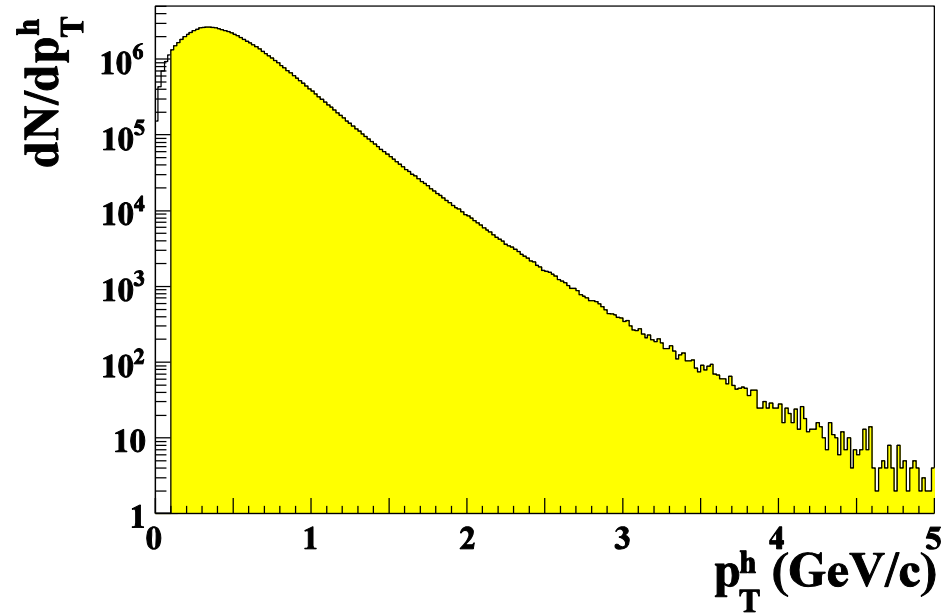
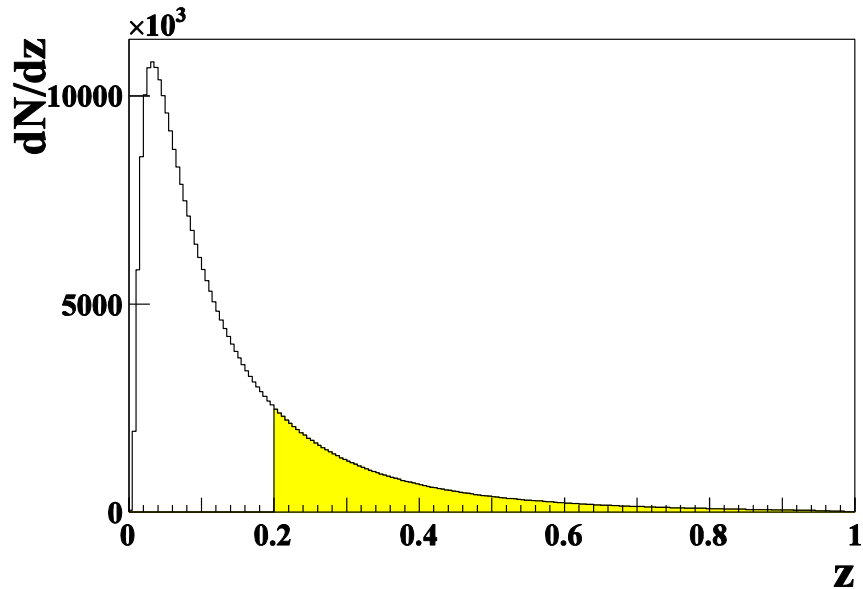
$$0.1 < y < 0.9$$

$$W > 5 \text{ GeV/c}^2$$





## SIDIS event selection



*charged hadron selection*

**$z > 0.2$**

**$p_T^h > 0.1$  GeV/c**



transverse momentum  
w.r.t. the virtual photon direction

# Collins Asymmetries

# Collins asymmetries

$$l N \rightarrow l' h X$$

direction of the polarization

kinematical factors

$$D_{NN} = \frac{1-y}{1-y-\frac{y^2}{2}}$$

$$N_h^\pm = N_h^0 \left[ 1 \pm f P_T D_{NN} A_{Coll} \sin(\phi_h + \phi_S + \pi) \right]$$

*Collins angle*

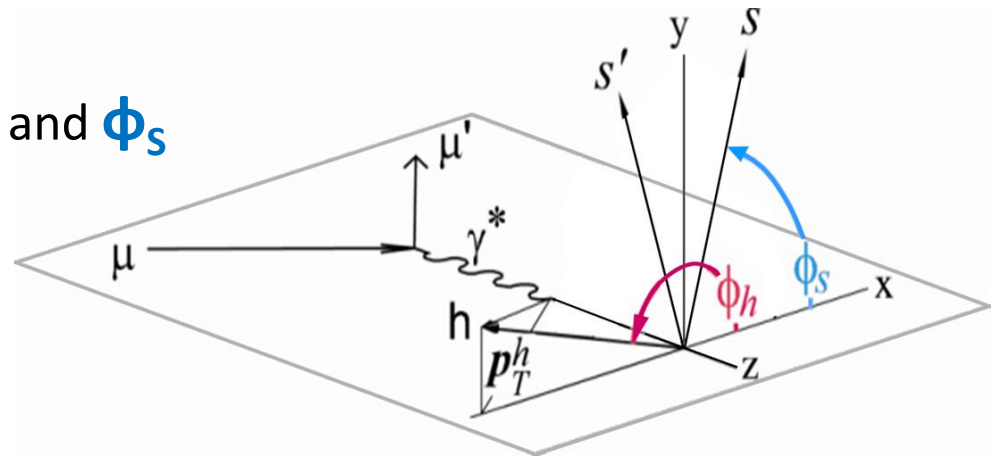
target polarisation  $P_T$   
dilution factor  $f$

convolution over the transverse momenta of the quarks in the nucleon  $k_\perp$  and in the fragmentation  $p_\perp$

transversity PDF      Collins FF

$$A_{Coll} \approx \frac{\sum_q e_q^2 \cdot h_1^q(k_\perp^2, x) \otimes H_{1q}^{\perp h}(p_\perp^2, z)}{\sum_q e_q^2 \cdot f_1^q(k_\perp^2, x) \otimes D_{1q}^{2h}(p_\perp^2, z)}$$

definition of the produced **hadron** and target polarisation azimuthal angles  $\phi_h$  and  $\phi_S$



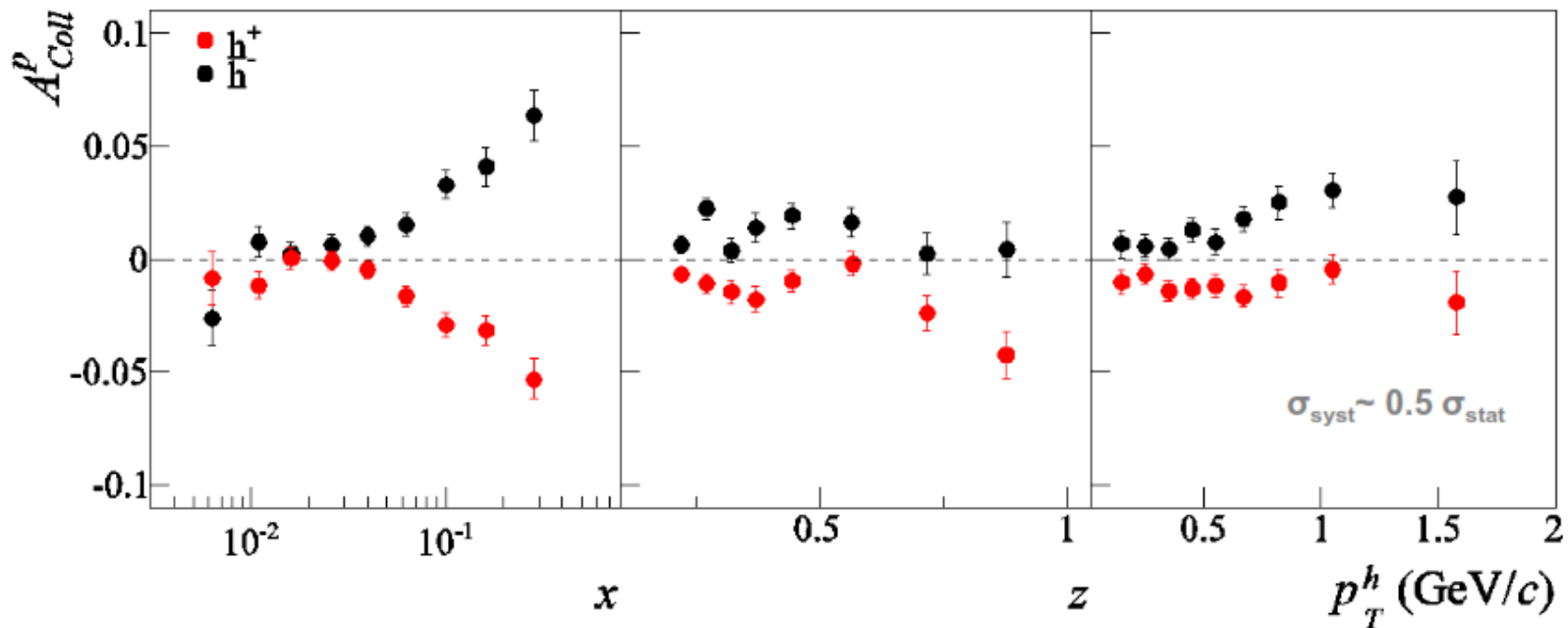
# The Collins asymmetry *charged hadrons*

azimuthal distributions calculated in the different bins of  $x$ ,  $z$  and  $p_T^h$

**amplitudes** of the azimuthal modulations **extracted**

**fitting** the azimuthal distribution in  $\phi_h$  and  $\phi_S$  *using extended maximum likelihood estimator*

combined **2007** - PLB 692 (2010) 240 - and **2010** - PLB 717 (2012) 376 -



- **precise measurements**
- **clear signal at  $x > 0.03$ , with opposite sign for  $h^+$  and  $h^-$**

# dihadron asymmetries



# dihadron asymmetries

$$lN \rightarrow l' h^+ h^- X$$

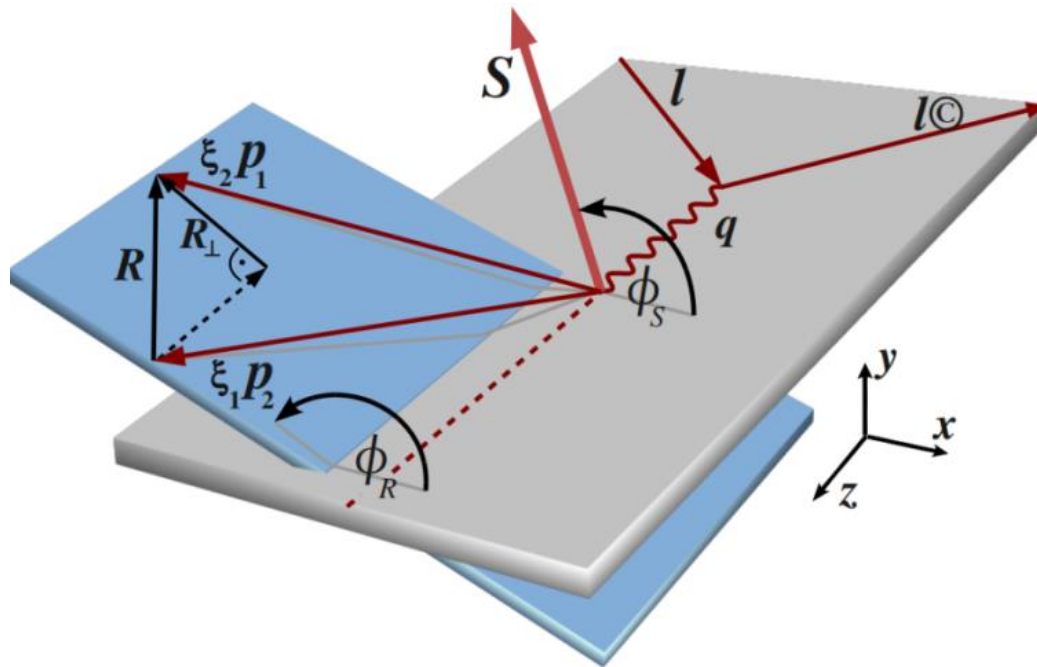
we measure

$$A_{UT}^{\sin \phi_{RS}}$$

from

$$N_{h^+h^-} = N_h^0 \left[ 1 \pm f P_T D_{NN} A_{UT}^{\sin \phi_{RS}} \sin(\phi_{RS}) \right]$$

all possible **oppositely charged hadrons pairs**



the azimuthal distribution of the hadrons pairs shows a modulation in the **azimuthal angle**:

$$\phi_{RS} = \phi_R + \phi_S - \pi$$

$$R = \xi_1 p_1 - \xi_2 p_2$$

$$\xi_1 = \frac{z_2}{z_1 + z_2}$$

$$\xi_2 = \frac{z_1}{z_1 + z_2}$$

Artru definition arXiv:0207309v2  
equivalent to Radici et al arXiv:1212.3568v1  
as seen from data

# dihadron asymmetries

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$$N_{h^+h^-} = N_h^0 \left[ 1 \pm f P_T D_{NN} A_{UT}^{\sin \phi_{RS}} \sin(\phi_{RS}) \right]$$

*all possible oppositely charged hadrons pairs*

transversity PDF

dihadron interference FF

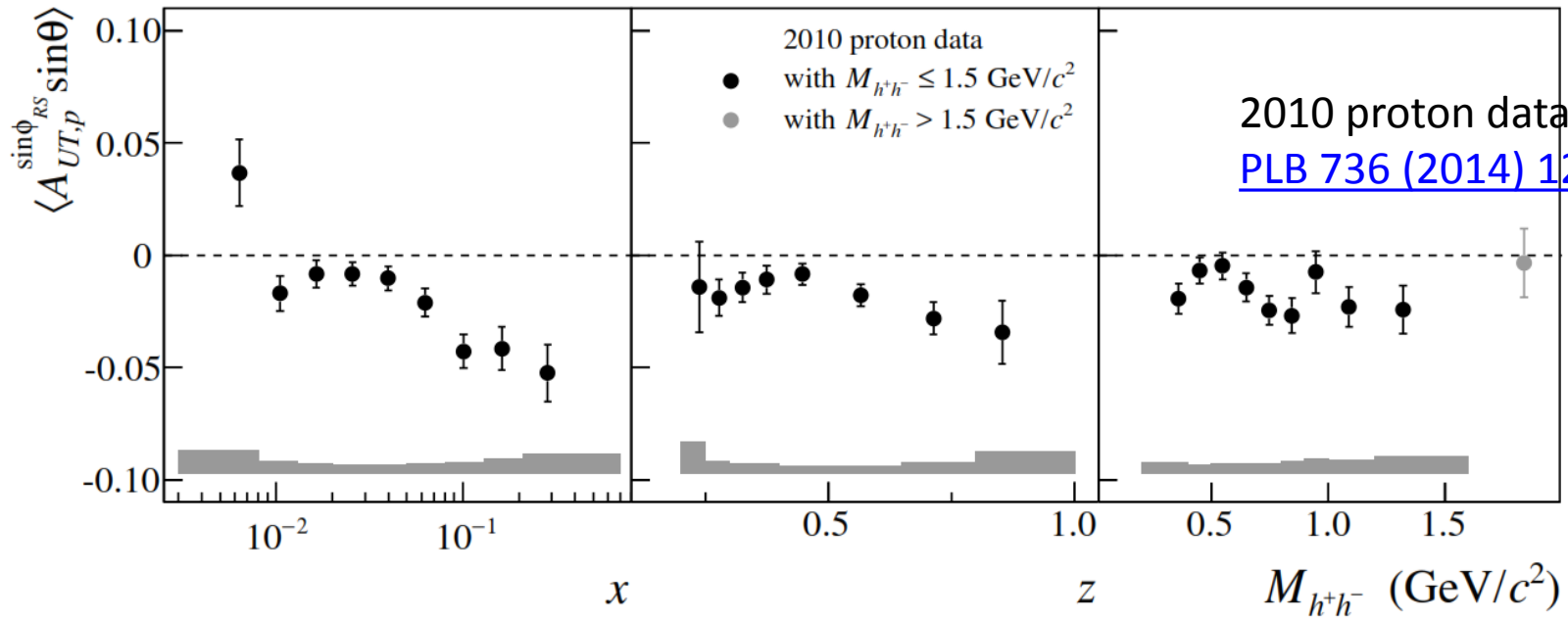
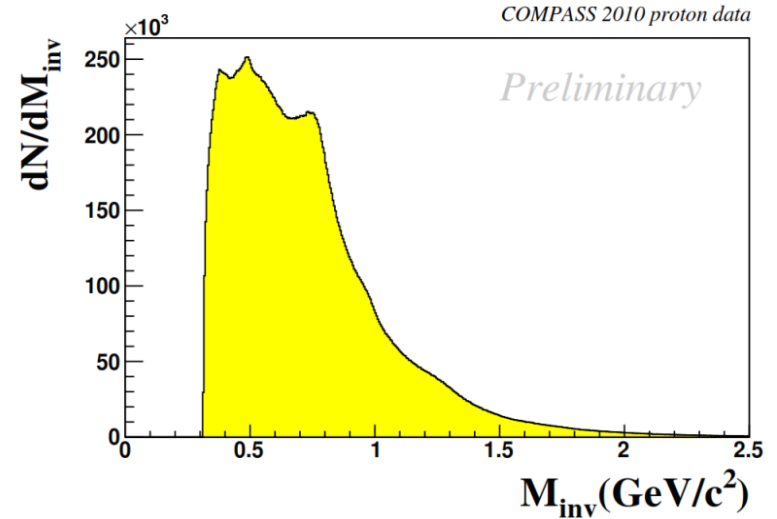
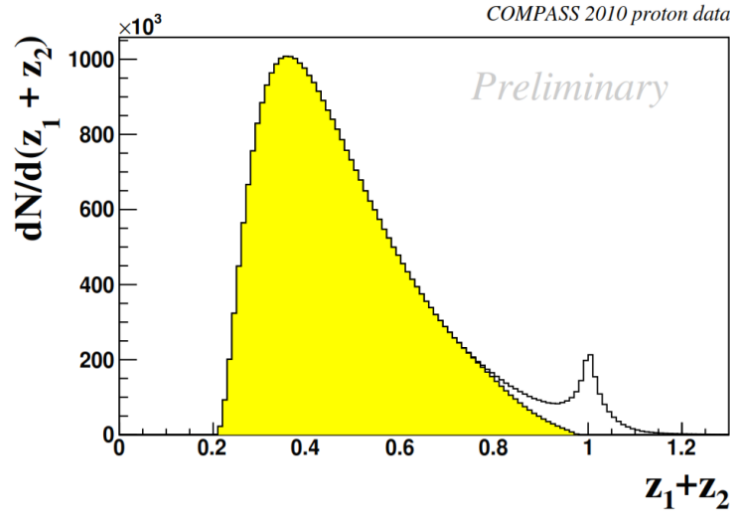
$$A_{UT}^{\sin \phi_{RS}} \approx \frac{\sum_q e_q^2 \cdot h_1^q(x) \cdot H_{1q}^{\angle}(z, M_{hh}^2)}{\sum_q e_q^2 \cdot f_1^q(x) \cdot D_{1q}^h(z, M_{hh}^2)}$$

**no convolution  
on the intrinsic  
transverse momenta**

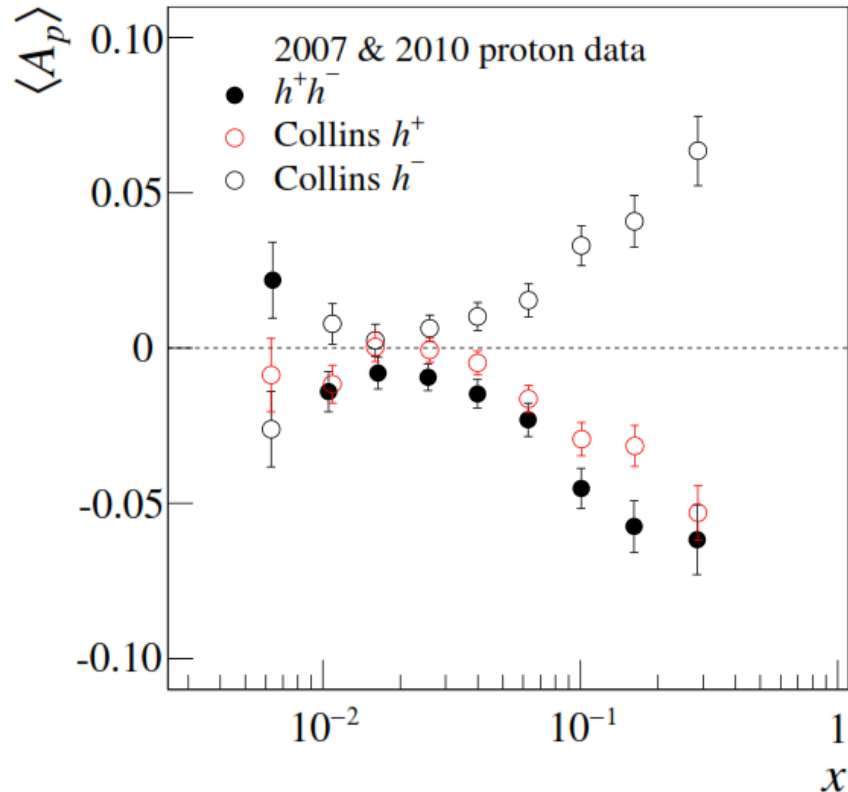
*collinear !*

$$A_{UT}^{\sin \phi_{RS}}$$

asymmetries measured as function of  $x$ ,  $z=z_1 + z_2$ ,  $M_{inv}$



interplay between **dihadron**  
and **Collins** asymmetries



[PLB 736 \(2014\) 124](#)

- Collins asymmetry for  $h^+$  and for  $h^-$  :  
“mirror symmetry”
- dihadron asymmetry vs Collins asymmetry:  
*only somewhat larger*

*motivated further studies on the interplay between the transversity induced asymmetries*



analysis of the single hadron and dihadron  
**transversity induced asymmetries** performed  
*on a common data sample* (2010 transversely polarised proton)



standard COMPASS **SIDIS sample**  
 but with  **$h^+ h^-$  at least detected** (each hadron with  $z > 0.1$ )

*essentially the same as the sample used to measure the dihadron asymmetries*

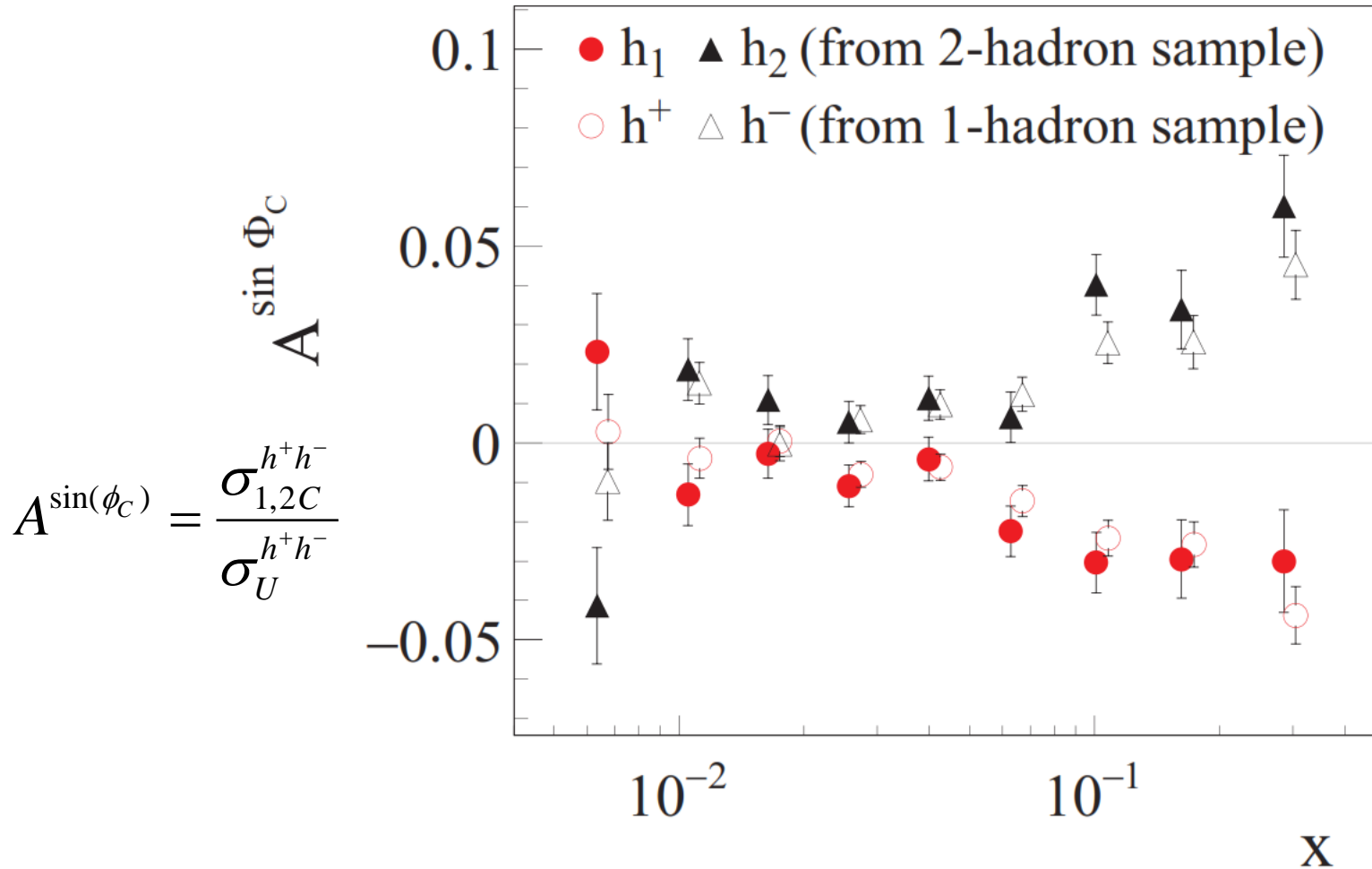
the general expression for the  $lN \rightarrow l' h^+ h^- X$  cross section

$$\frac{d\sigma^{h^+h^-}}{d\phi_{h^+} d\phi_{h^-} d\phi_S} = \sigma_U^{h^+h^-} + S_T \cdot \left[ \sigma_{1C}^{h^+h^-} \sin(\phi_{h^+} + \phi_S - \pi) + \sigma_{2C}^{h^+h^-} \sin(\phi_{h^-} + \phi_S - \pi) \right]$$

A.Kotzinian PRD91 (2015) 054001

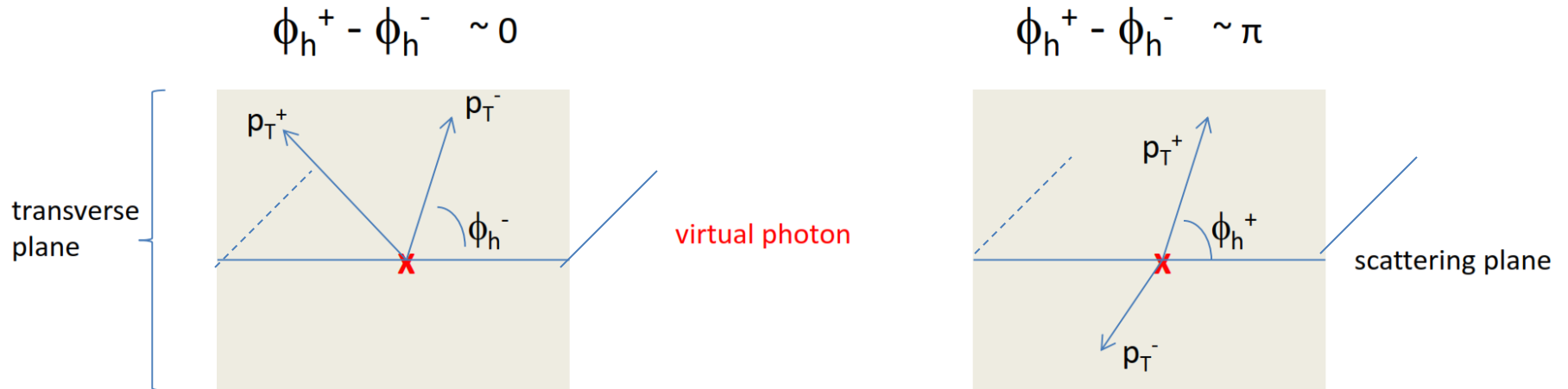
## single hadron asymmetries from the common sample

are **very similar** to the Collins asymmetries calculated from the **1 hadron sample** (which are the usual Collins asymmetries)



we studied the **azimuthal correlation** among the hadrons forming the pair

$$\Delta\phi = \phi_{h^+} - \phi_{h^-}$$



we measured **single hadron** transversity induced as function of  $\Delta\Phi$

the two hadron cross section can be written as function of  $\Delta\phi$  and  $\phi_{+,-}$

*dependencies  $\sigma(\Delta\phi)$  omitted ...*

$$\phi_{h^+}, \phi_{h^-} \rightarrow \phi_{h^+}, \Delta\phi$$

$$\frac{d\sigma^{h^+h^-}}{d\phi_{h^+} d\Delta\phi d\phi_S} = \sigma_U^{h^+h^-} + S_T \cdot \left[ \left( \sigma_{1C}^{h^+h^-} + \sigma_{2C}^{h^+h^-} \cos \Delta\phi \right) \sin(\phi_{h^+} + \phi_S - \pi) + \dots \right]$$

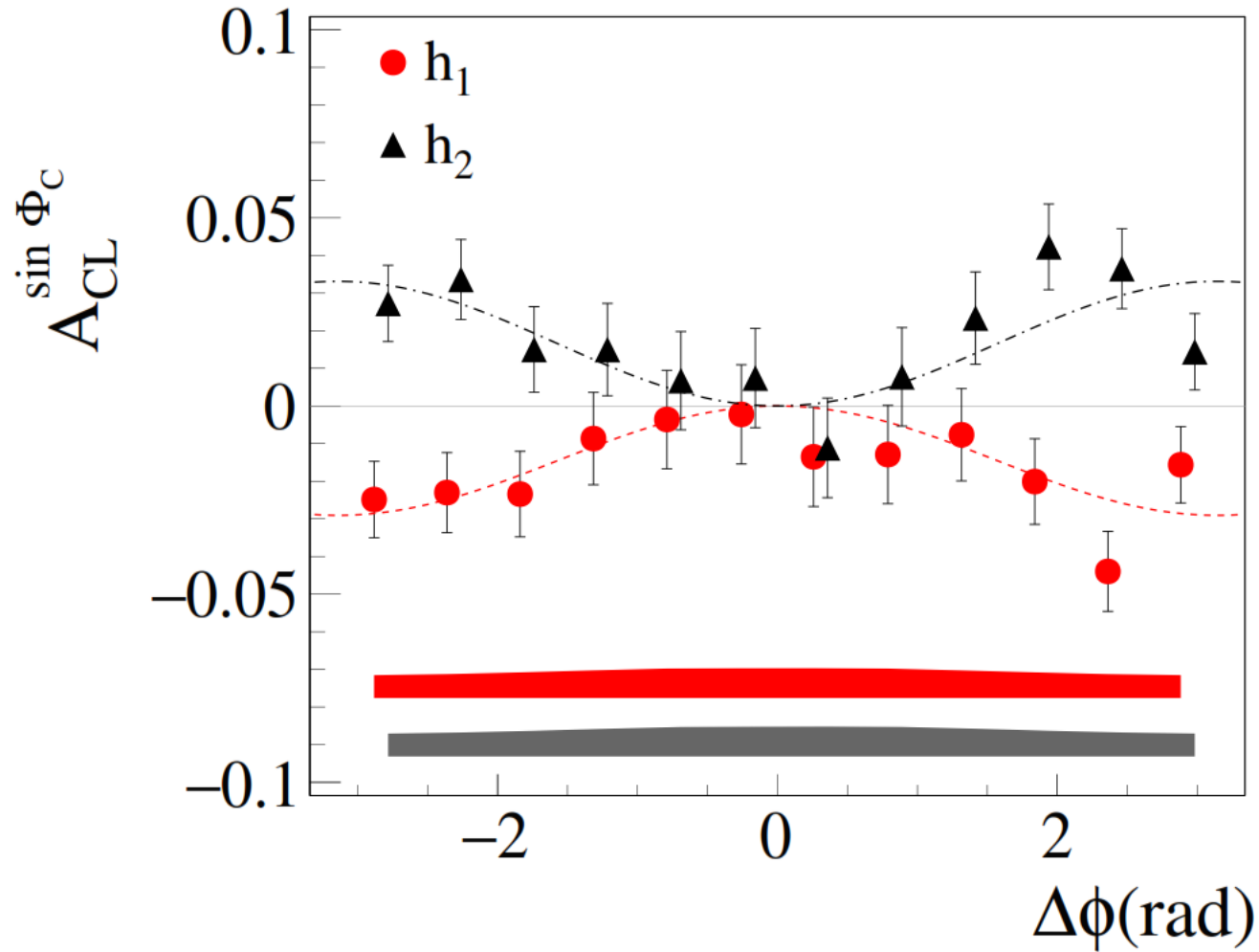
$$\phi_{h^+}, \phi_{h^-} \rightarrow \Delta\phi, \phi_{h^-}$$

$$\frac{d\sigma^{h^+h^-}}{d\phi_{h^-} d\Delta\phi d\phi_S} = \sigma_U^{h^+h^-} + S_T \cdot \left[ \left( \sigma_{2C}^{h^+h^-} + \sigma_{1C}^{h^+h^-} \cos \Delta\phi \right) \sin(\phi_{h^-} + \phi_S - \pi) + \dots \right]$$

we measured **single hadron** transversity induced asymmetries as function of  $\Delta\phi$

$$A_{1CL}^{\sin(\phi_{h^+} + \phi_{h^-} - \pi)} = \frac{\sigma_{1C}^{h^+h^-}(\Delta\phi) + \sigma_{2C}^{h^+h^-}(\Delta\phi) \cdot \cos \Delta\phi}{\sigma_U^{h^+h^-}(\Delta\phi)}$$

$$A_{2CL}^{\sin(\phi_{h^+} + \phi_{h^-} - \pi)} = \frac{\sigma_{2C}^{h^+h^-}(\Delta\phi) + \sigma_{1C}^{h^+h^-}(\Delta\phi) \cdot \cos \Delta\phi}{\sigma_U^{h^+h^-}(\Delta\phi)}$$



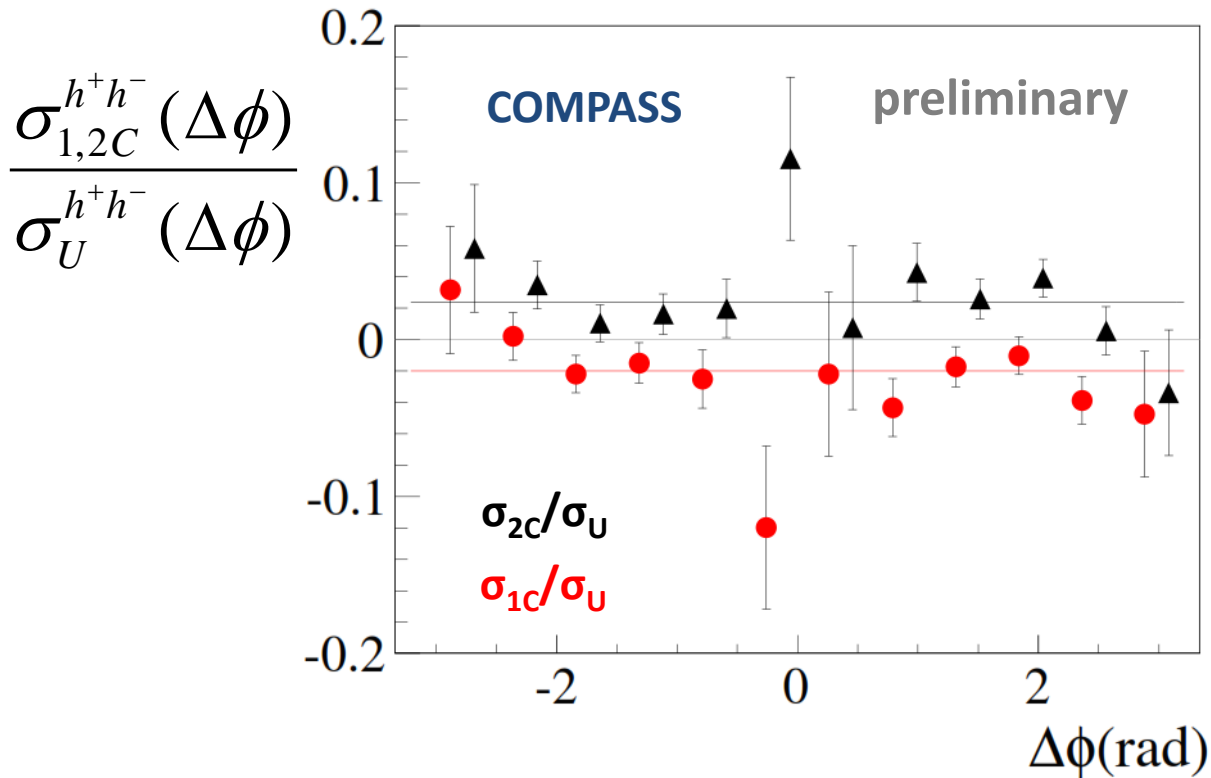
*signal bigger when  $h^+$  and  $h^-$  go in opposite direction*

*mirror asymmetries*

$$\sigma_{1C}^{h^+h^-} = -\sigma_{2C}^{h^+h^-}$$



the dependence from  $\Delta\Phi$  in the *polarised* part and *unpolarised* part of the cross section has been found to be roughly the same



*the  $\Delta\Phi$  dependence looks similar*

and also mirror symmetry

## comparison between **dihadron** and **Collins** asymmetries

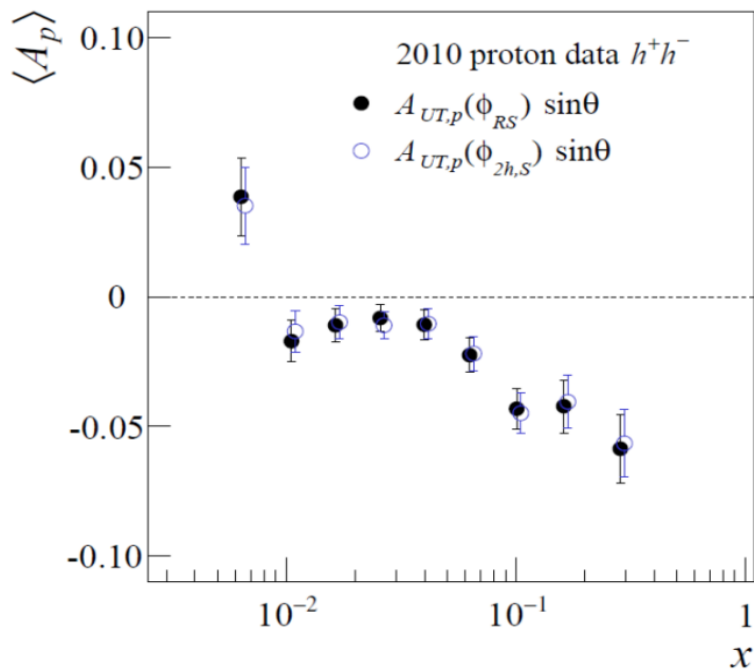
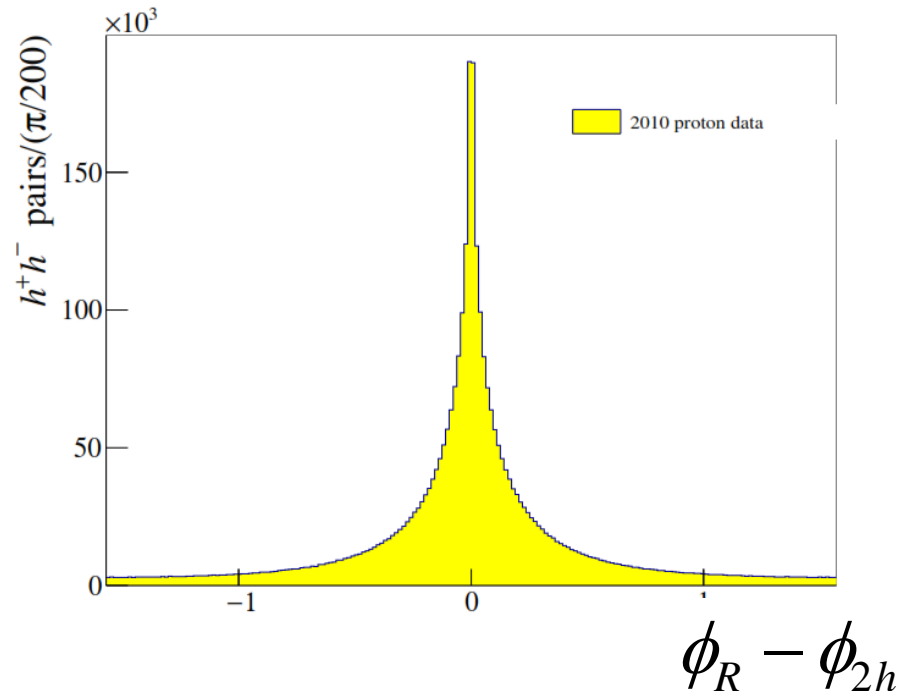
# we defined another angle for the 2 hadrons pair *(more simple to use...)*

$\phi_{2h}$  which is the azimuthal angle of the vector  $\hat{\mathbf{P}}_T^{h^+} - \hat{\mathbf{P}}_T^{h^-}$

strongly correlated with

$$\phi_R = \frac{z^{h^-} \cdot \hat{\mathbf{P}}_T^{h^+} - z^{h^+} \cdot \hat{\mathbf{P}}_T^{h^-}}{z^{h^-} + z^{h^+}}$$

used to extract the “standard” dihadron asymmetries



azimuthal transverse spin asymmetries calculated using the 2 angles are the same

**we defined another angle for the 2 hadrons pair** (*more simple to use...*)

and by using the relation implied by the mirror asymmetries  $\sigma_{1C}^{h^+h^-} = -\sigma_{2C}^{h^+h^-}$   
the **cross section** becomes

$$\frac{d\sigma^{h^+h^-}}{d\phi_{2h} d\Delta\phi d\phi_S} = \sigma_U^{h^+h^-} + S_T \cdot \sigma_{1C}^{h^+h^-} \cdot \sqrt{2(1 - \cos \Delta\phi)} \sin(\phi_{2h} + \phi_S - \pi)$$



measure **di-hadron asymmetries** (*in  $\phi_{2h}$* )

as function of  **$\Delta\Phi$**

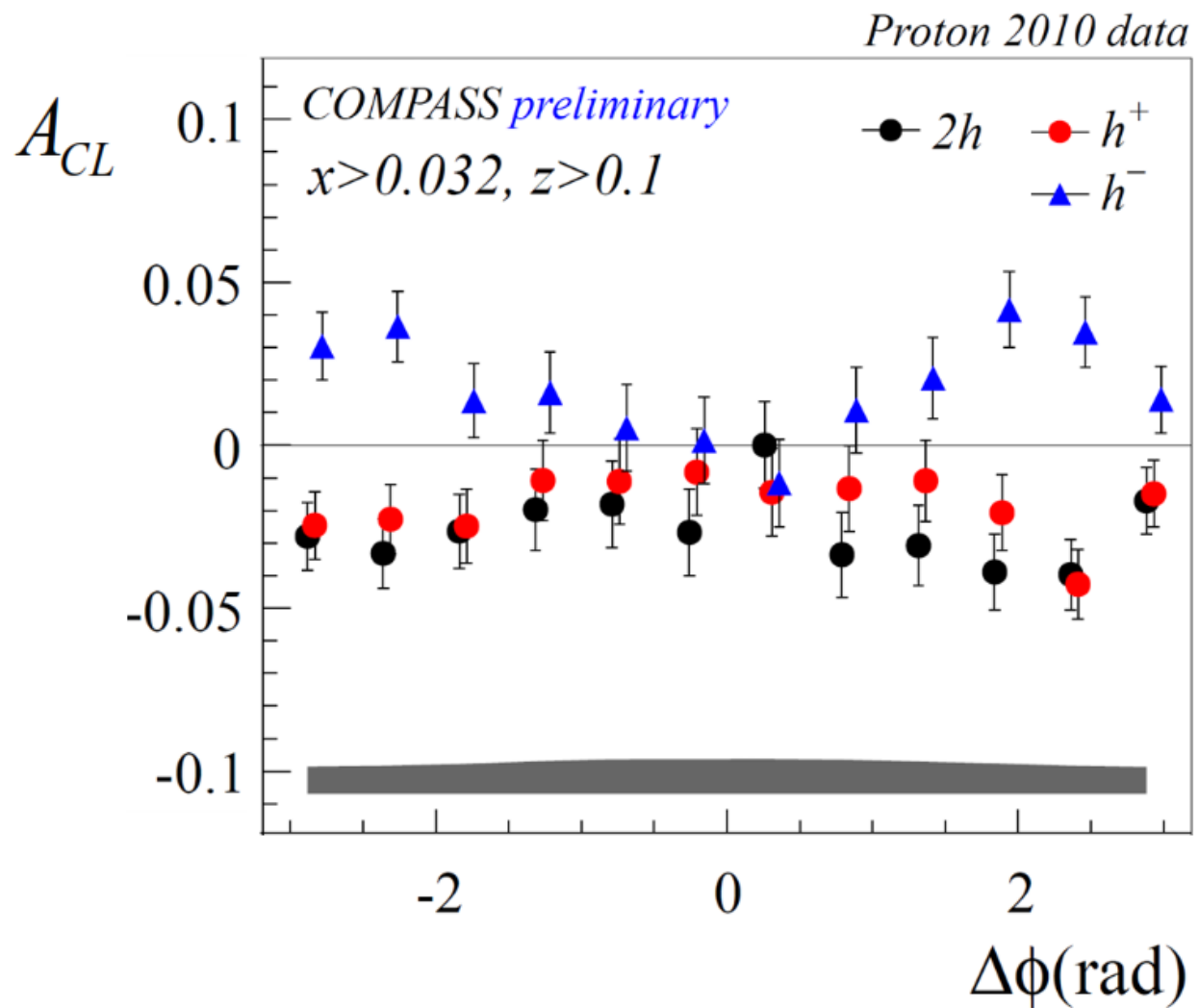
to compare with single hadron collins like asymmetries

amplitudes  
of the azimuthal  
modulations in the  
Collins angle

$h^+$

$h^-$

dihadron



same mechanism seems to give collins asymmetries and di-hadron asymmetries



there is a **quantitative relation** between **dihadron** and the **single hadron Collins-like asymmetries** that is derived using

$$\sigma_{1C}^{h^+h^-} = -\sigma_{2C}^{h^+h^-}$$

which comes from the **mirror asymmetries** seen from data

writing explicitly the measured asymmetries we have

**h<sup>+</sup>**

$$A_{1CL}^{\sin(\phi_{h^+} + \phi_S - \pi)} = \frac{\sigma_{1C}^{h^+h^-}}{\sigma_U^{h^+h^-}} \cdot (1 - \cos \Delta\phi)$$

**h<sup>-</sup>**

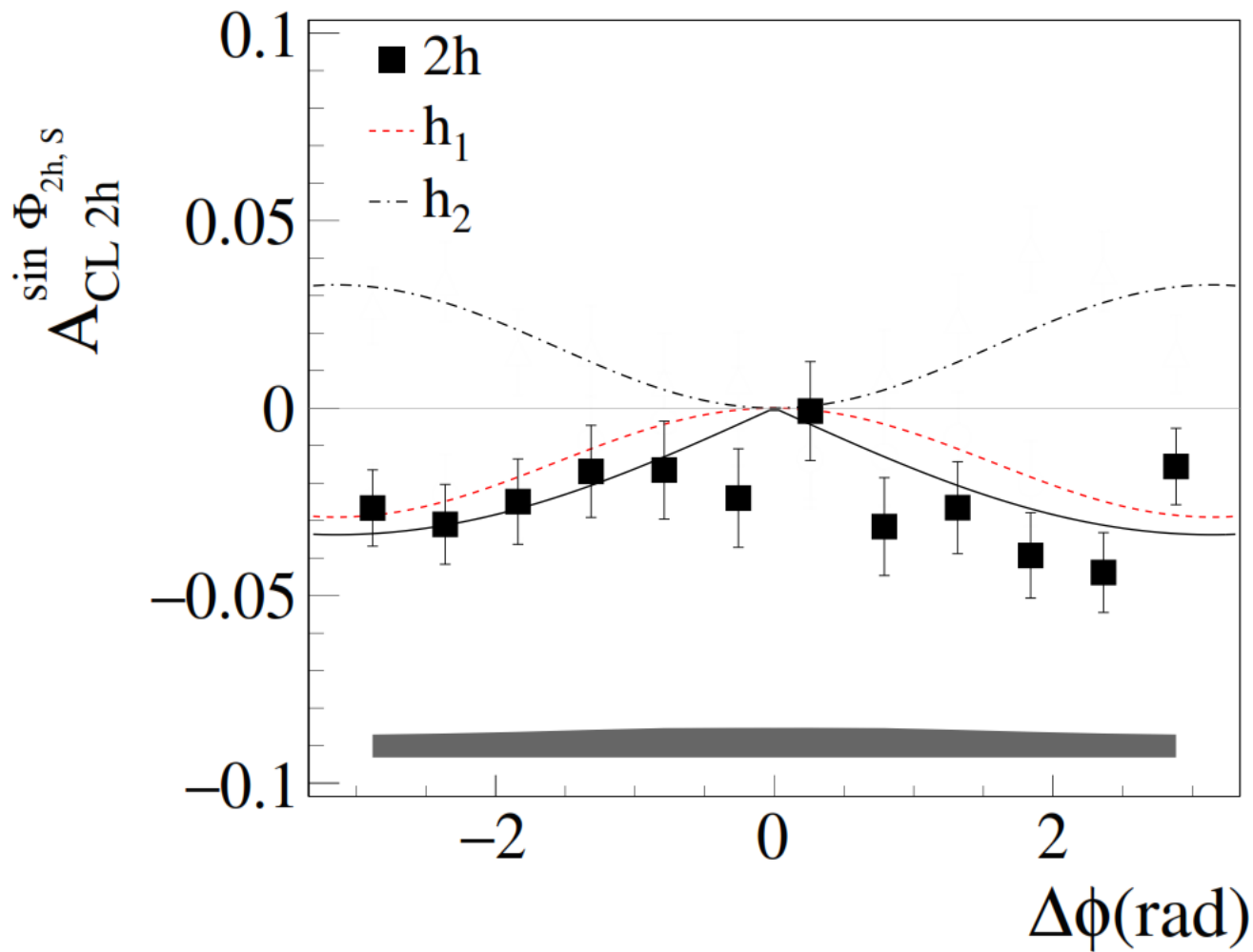
$$A_{2CL}^{\sin(\phi_{h^-} + \phi_S - \pi)} = -\frac{\sigma_{1C}^{h^+h^-}}{\sigma_U^{h^+h^-}} \cdot (1 - \cos \Delta\phi)$$

**dihadron**

$$A_{2h,CL}^{\sin(\phi_{2h} + \phi_S - \pi)} = \frac{\sigma_{1C}^{h^+h^-}}{\sigma_U^{h^+h^-}} \cdot \sqrt{2 \cdot (1 - \cos \Delta\phi)}$$



**same quantity !**

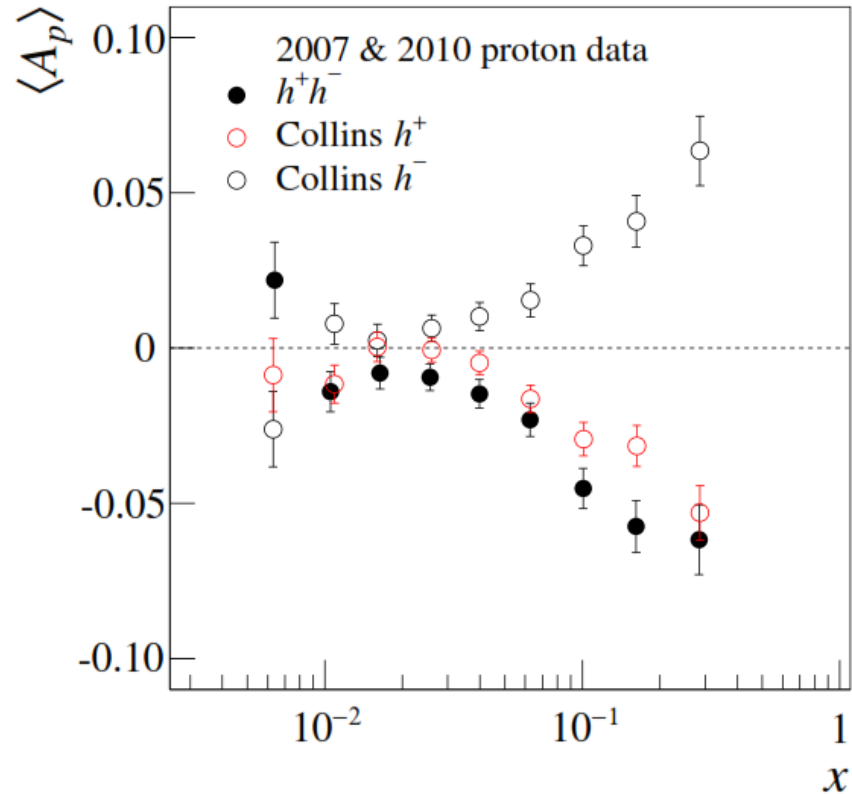


$$c \cdot (1 - \cos \Delta\phi)$$

$$-c \cdot (1 - \cos \Delta\phi)$$

$$c \cdot \sqrt{2} \cdot (1 - \cos \Delta\phi)$$

*integrating over  $\Delta\phi$  the ratio between dihadron and positive hadron asymmetries is  $4/\pi$*



[PLB 736 \(2014\) 124](#)

- Collins asymmetry for  $h^+$  and for  $h^-$  :  
“mirror symmetry”
- dihadron asymmetry vs Collins asymmetry:

*only somewhat larger*



*integrating over  $\Delta\phi$  the ratio between dihadron and positive hadron asymmetries is  $4/\pi$*

**interesting results**  
**the paper has been written**

interplay among transversity induced asymmetries  
in hadron lepton production

arXiv:1507.07593v1

CERN-PH-EP-2015–199

*submitted to Phys.Rev.Lett.*