

Light Exotics

-

Have we observed mesons beyond ?

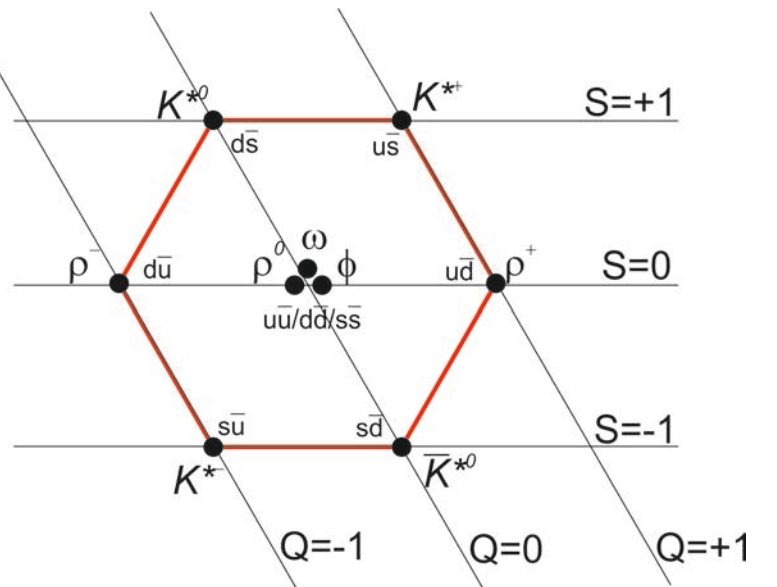
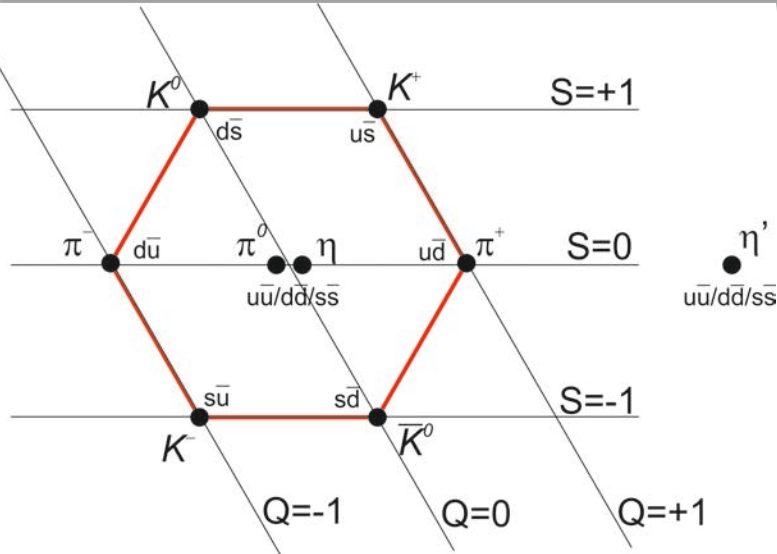
Bernhard Ketzer

Rheinische Friedrich-Wilhelms-Universität
Bonn

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Bormio, Italy

28 January 2015

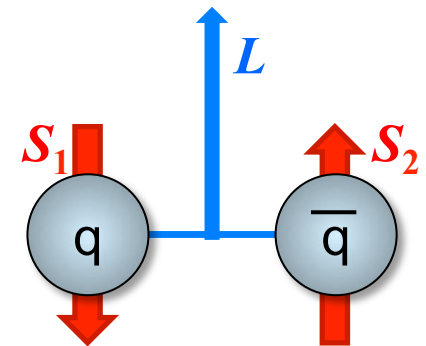


Quark model:

- bound state of $q\bar{q}$
- $SU(3)_{\text{flavor}}$: $q \otimes \bar{q}' = 3 \otimes \bar{3} = 8 \oplus 1$
- color singlets

Quantum numbers:

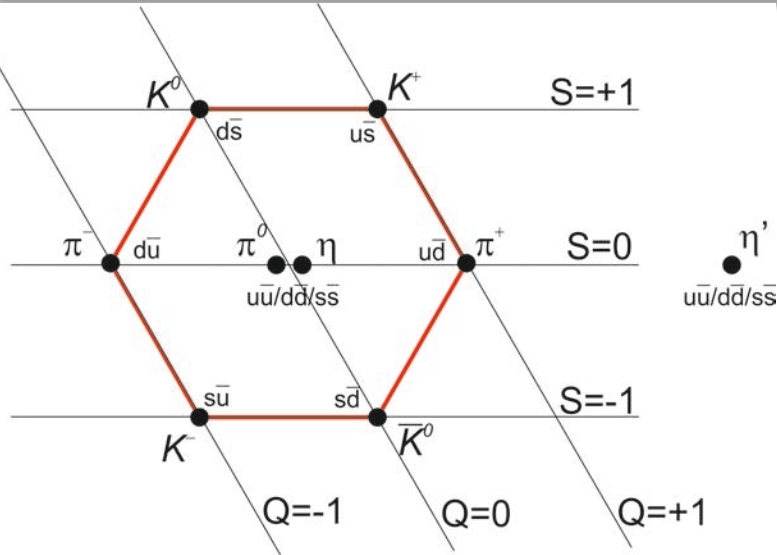
- measured: $I^G (J^{PC})$
- quark model: ${}^{2S+1}L_J$
 $S = S_1 + S_2$, $J = L + S$



$$P = (-1)^{L+1}$$

$$C = (-1)^{L+S}$$

$$G = (-1)^{L+S+I}$$



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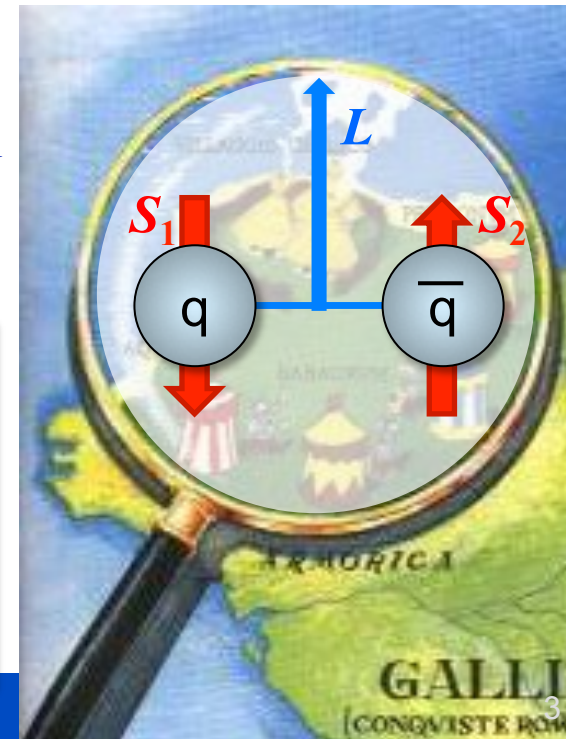
Allowed combinations

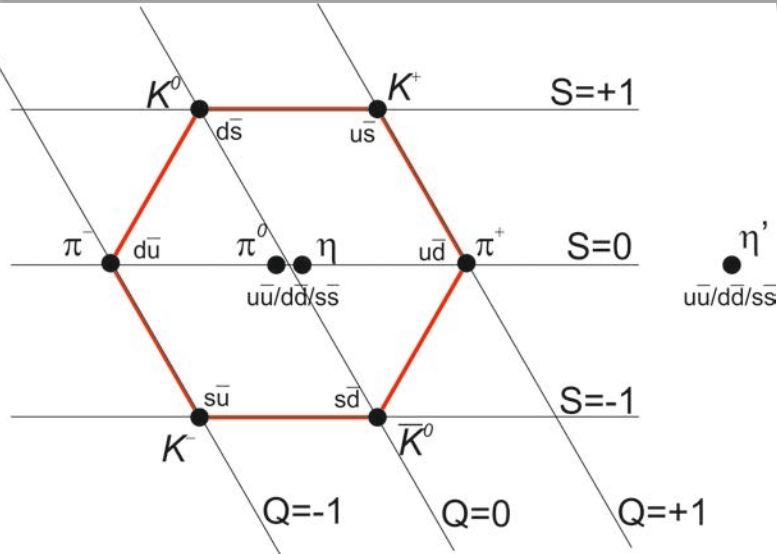
$$J^{PC} = 0^{-+}, 0^{++}, 1^{--}, 1^{+-}, 2^{++}, \dots$$

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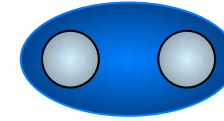


Allowed combinations

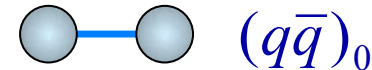
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“Forbidden” combinations

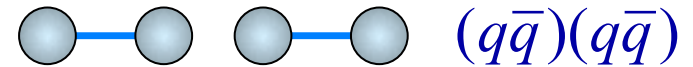
$$J^{PC} = 0^{--}, 0^{+-}, 1^{-+}, 2^{+-}, \dots$$



=



+



+ Molecule / 4 quarks



+

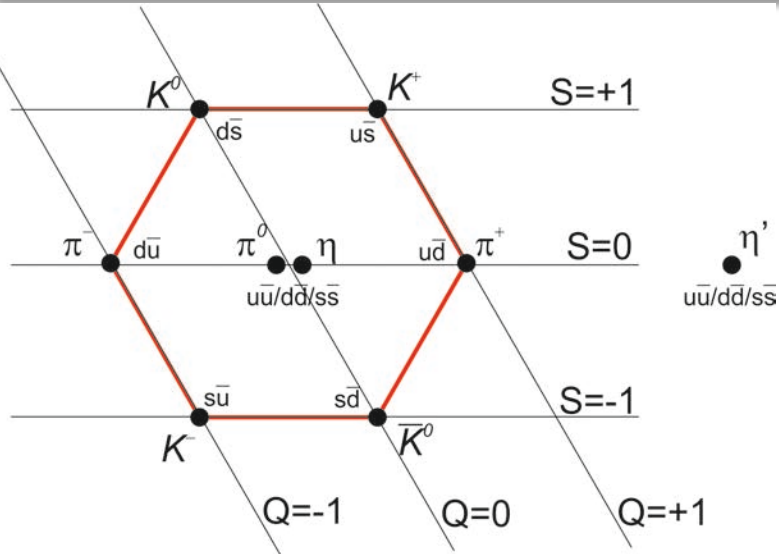
Hybrids



gg

Glueballs

+

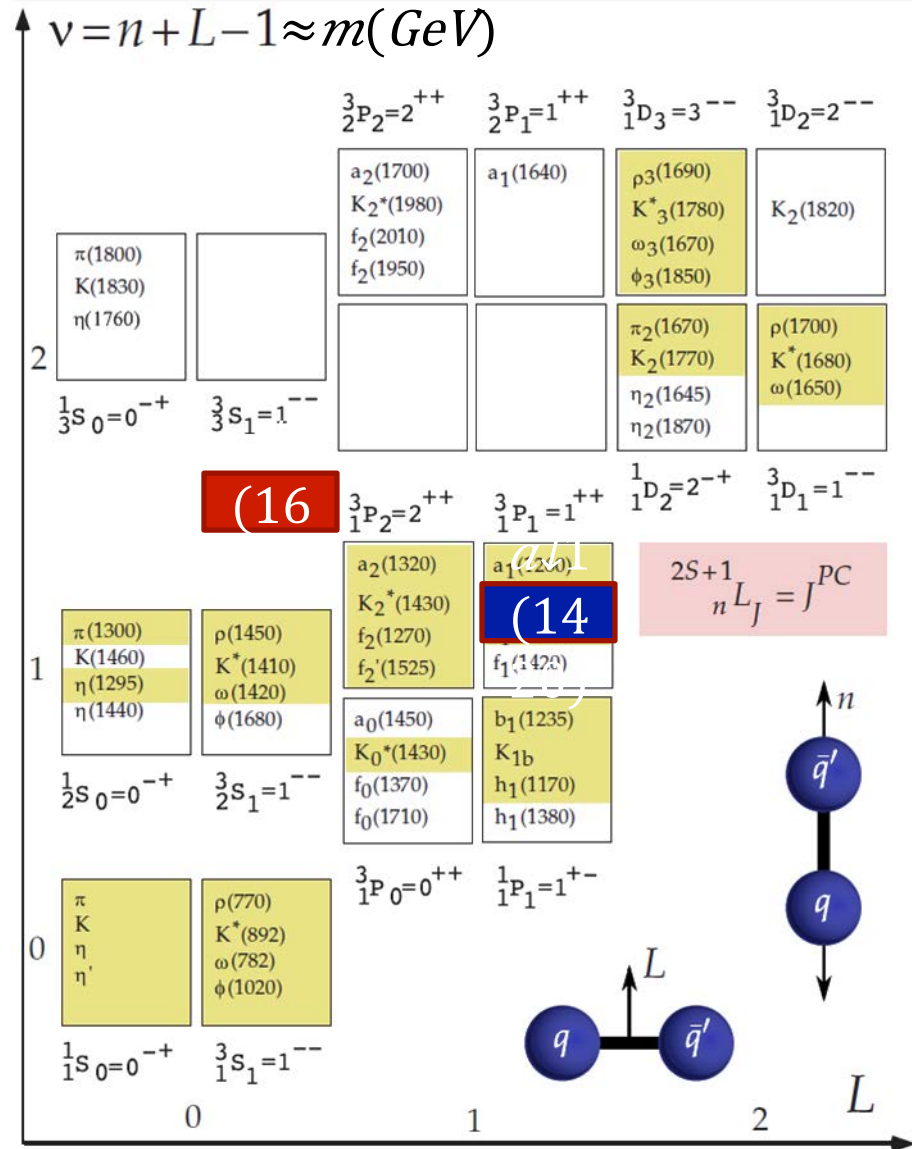


Allowed combinations

$$J^{PC} = 0^{-+}, 0^{++}, 1^{--}, 1^{+-}, 2^{++}, \dots$$

“Forbidden” combinations

$$J^{PC} = 0^{--}, 0^{+-}, 1^{-+}, 2^{+-}, \dots$$



- Ground state 0^{-+} , 1^{--} nonets ok
- Many predicted radial and orbital excitations missing / unclear
- Identification of exotics
 - overpopulation of meson spectrum
 - spin-exotic quantum numbers

In this talk: 2 examples

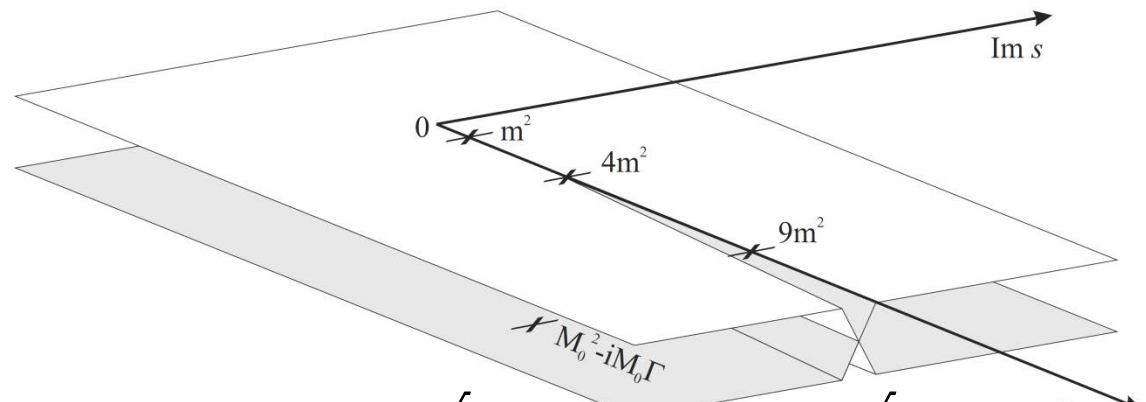
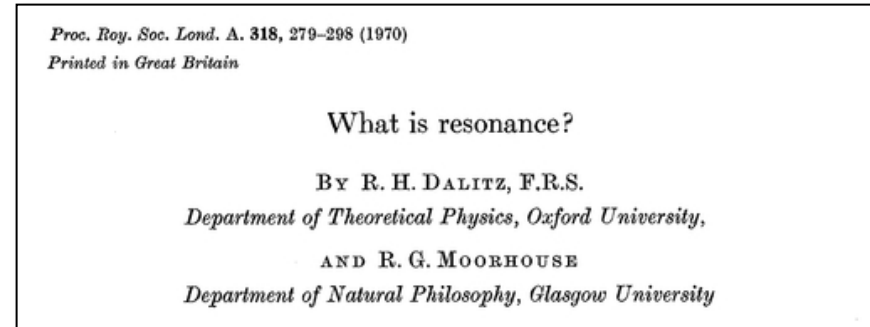
- $a_1(1420)$: $J^{PC} = 1 \uparrow ++$
- $\pi_1(1600)$: $J^{PC} = 1 \uparrow - +$

[Amsler et al., Phys. Rept. 389, 61 (2004)]

- “Not every bump is a resonance, and not every resonance is a bump”
- Resonances have complex properties like mass and width, which do not depend on the experiment or the specific model
- Resonances correspond to **poles in the S-matrix** on unphysical Riemann sheets
- S-matrix: $S = I + iT$
 - unitary
 - analytic

Transition (reaction) matrix:

$$T_{ab} = (2\pi)^4 \delta^4(P_b - P_a) \prod_{i \in a} 1/\sqrt{2E_i} \prod_{j \in b} 1/\sqrt{2E_j} M_{ab}$$



For a 2-body reaction: expand scattering amplitude in partial waves

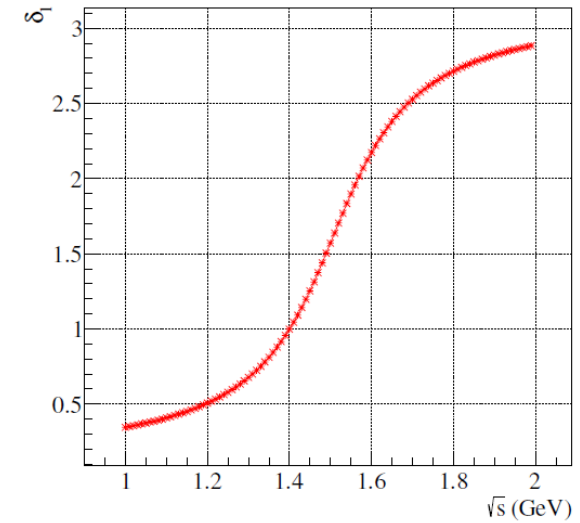
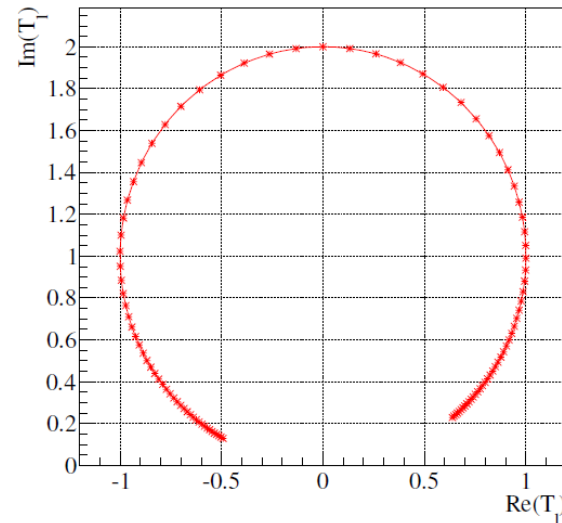
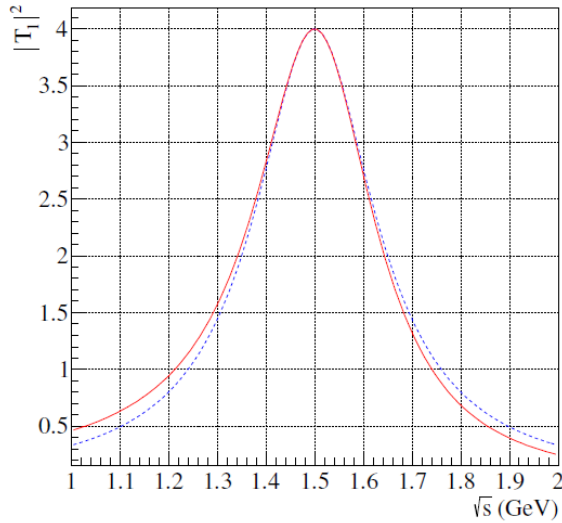
$$M_{ab} \equiv A(s, t) = \sum_{\ell=0}^{\infty} (2\ell+1) A_{\ell}(s) P_{\ell}(\cos\theta)$$

- P_{ℓ} Legendre polynomials \Rightarrow angular distribution
- A_{ℓ} transition amplitudes \Rightarrow dynamics
- General parameterization for scattering through resonance

$$A_{\ell}(s) = 8\pi\sqrt{s}/k \cdot \eta_{\ell}(s) e^{i\delta_{\ell}(s)} - 1/2i = 8\pi\sqrt{s}/k f_{\ell}(s)$$

- For isolated, narrow resonance: Breit-Wigner parameterization

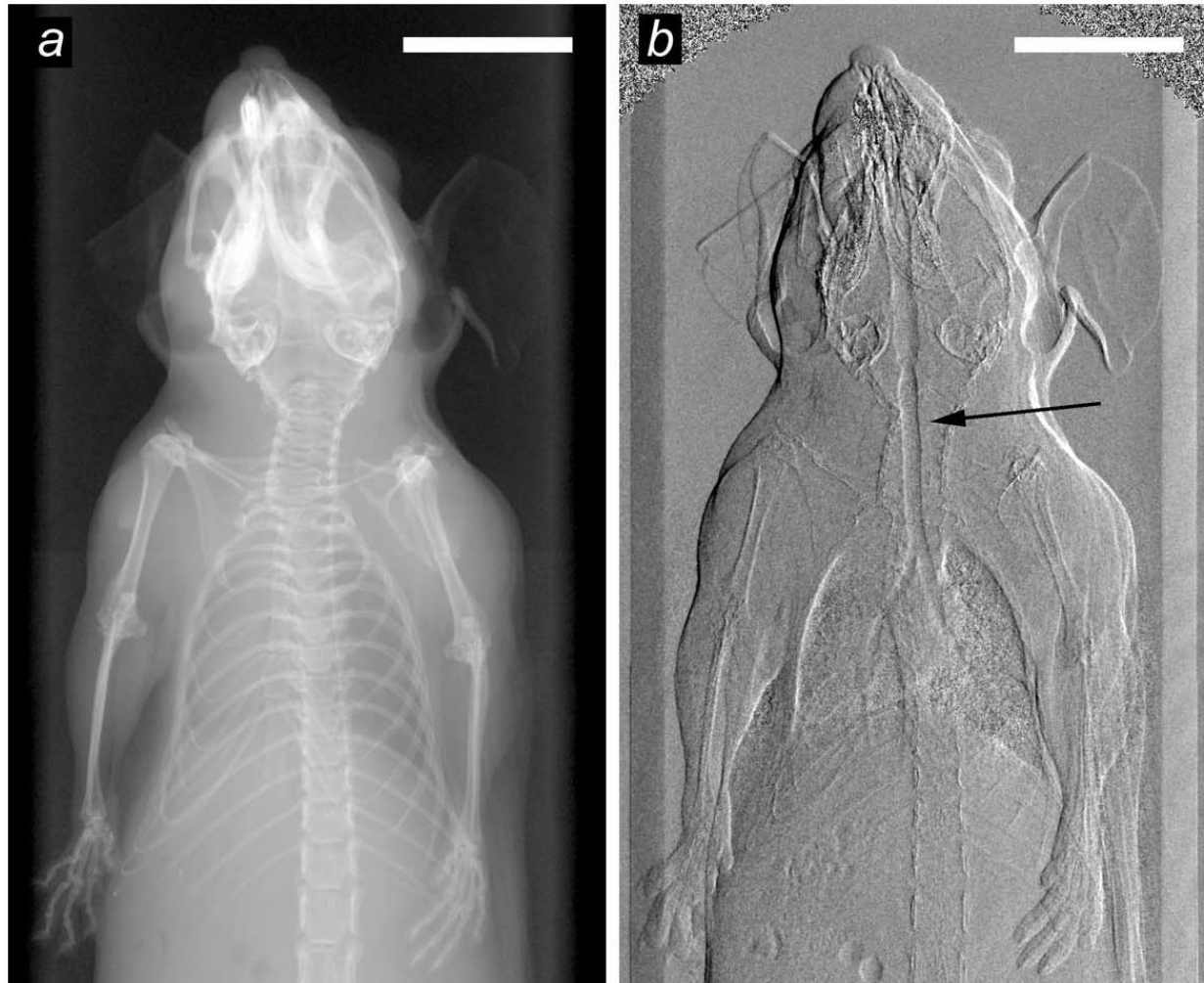
$$f_{\ell}(s) \cong m_0 \Gamma / m_0^2 - s - im_0 \Gamma$$



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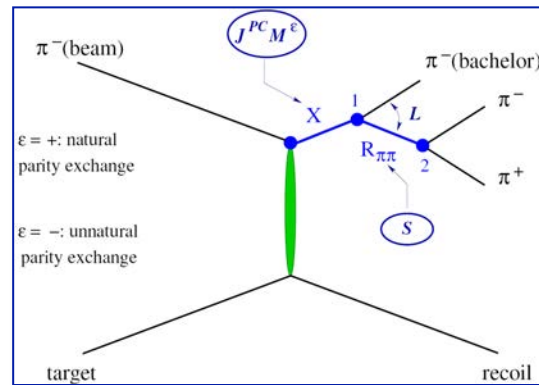
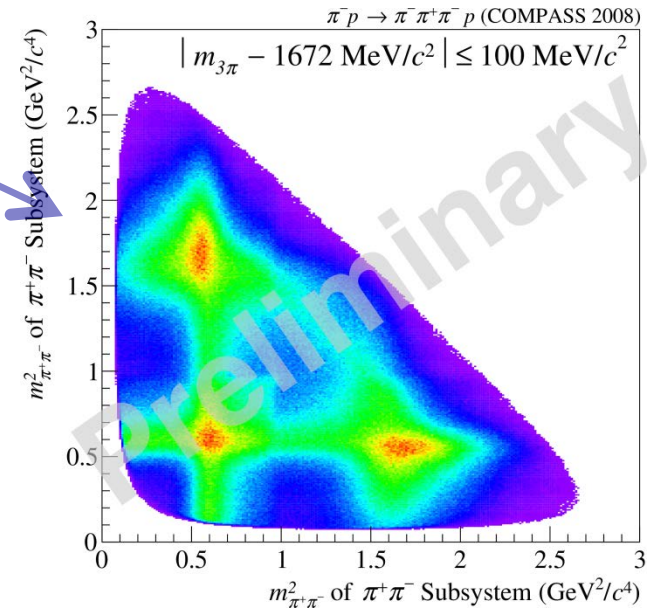
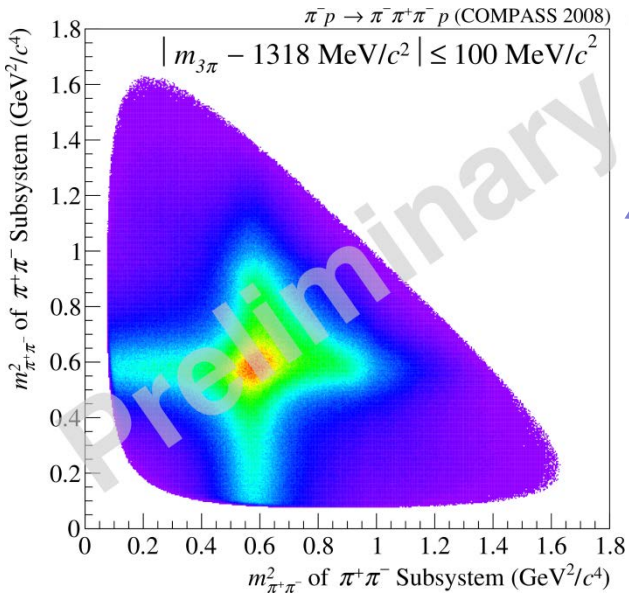
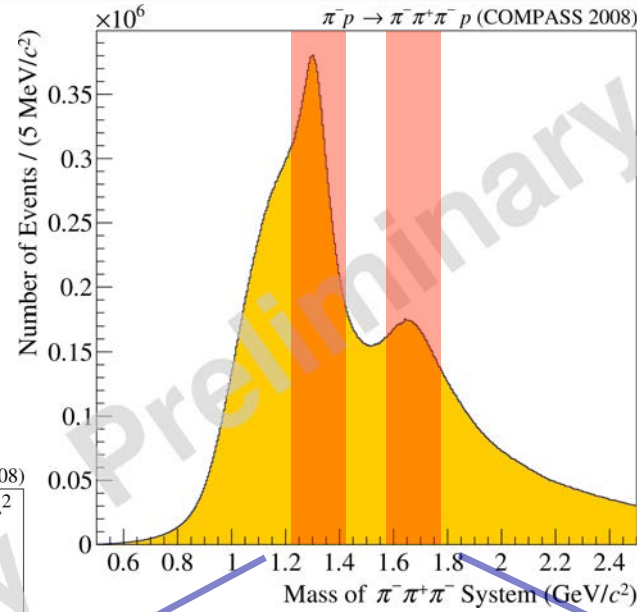


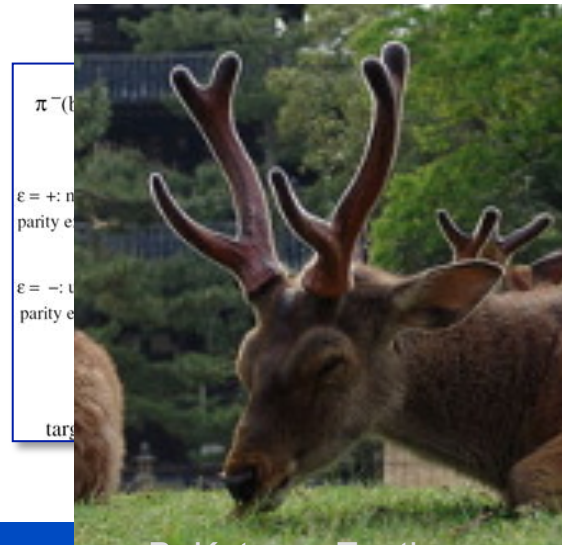
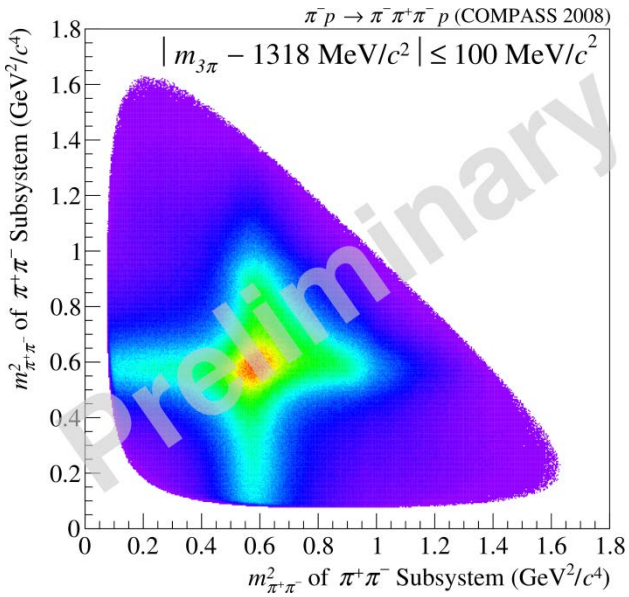
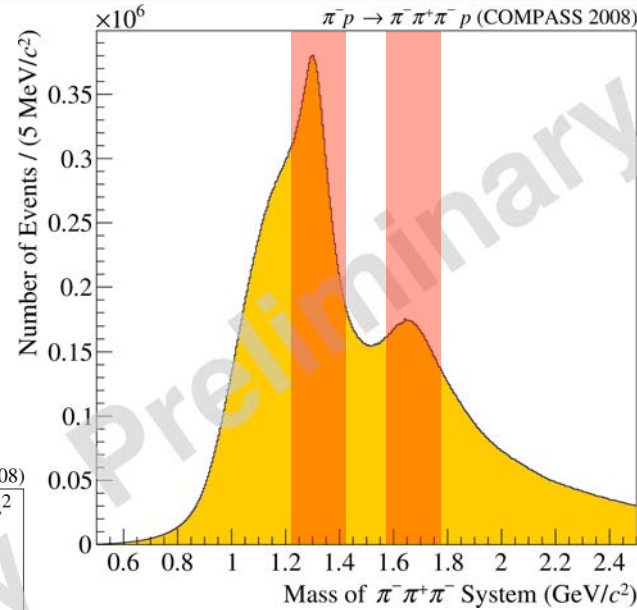
[M. Beck et al., Sci. Rep. 3, 3209; DOI:10.1038/srep03209 (2013).]

- Many different resonances are produced which decay into same final state
- Goal:
 - find and disentangle (all) contributing resonances
 - determine **mass, width** and **quantum numbers** J^P of resonances
 - ⇒ **angular distributions** of decay products
- Interference effects ⇒ **small resonances** may be enhanced
- Take into account experimental **acceptance**

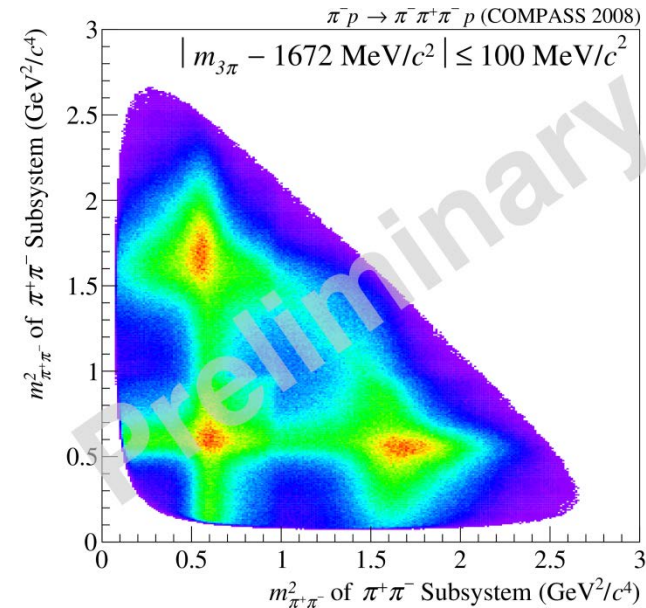
Assumptions:

- Production and decay of a state factorize
- Decay into multi-particle final state can be described by a sequence of 2-body decays





$\pi^- (\dots)$
 $\epsilon = +$: n
 parity e
 $\epsilon = -$: u
 parity e
 targ



1. PWA of angular distributions in **mass bins** and **t' bins**

$$I(\tau) = \left| \sum_{\xi} T_{\xi} A_{\xi}(\tau) \right|^2$$

- $T_{\downarrow\xi} =$ production amplitude for state with $\chi = I^{\uparrow}G (J^{\uparrow}PC)M$ decaying to $\downarrow\xi$
- $A_{\downarrow\xi}(\tau) =$ decay amplitude (calculable without free parameters)
- Result: spin-density matrix $\rho_{\downarrow\xi\xi'}^{\uparrow} = T_{\downarrow\xi} T_{\downarrow\xi'}^{\uparrow*}$

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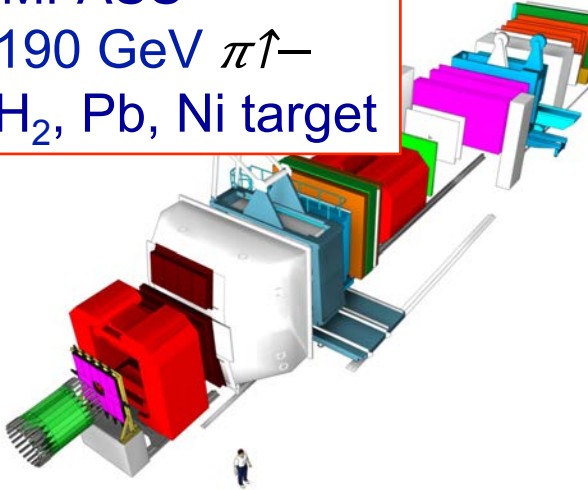
2. χ^2 -Fit of mass- and t-dependence of spin-density matrix

- Resonant contributions: Breit-Wigner functions
- Non-resonant contributions: empirical functions
- Only subset of spin-density matrix is considered for computational reasons

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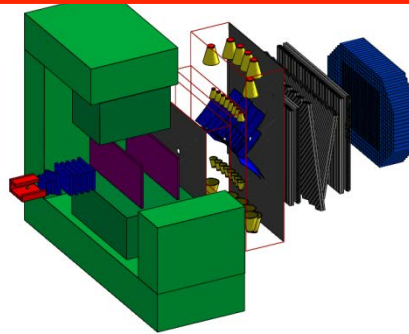
COMPASS

- 190 GeV $\pi^+\pi^-$
- H₂, Pb, Ni target

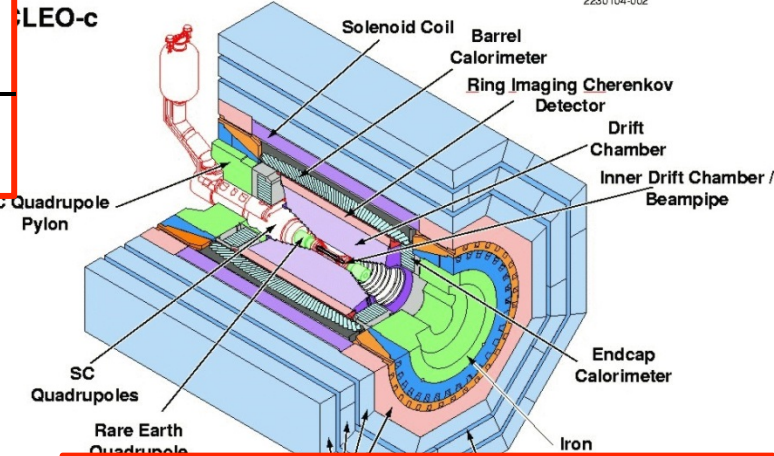


VES:

- 37/29 GeV $\pi^+\pi^-$
- Be target

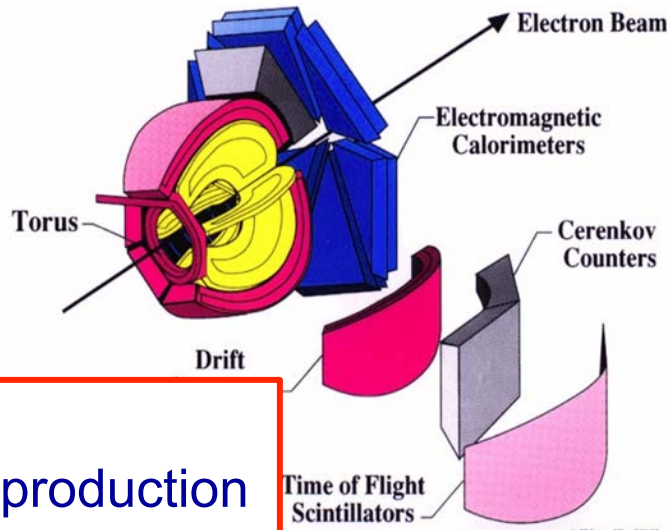


CLEO-c



CLEO-c:

- e^+e^- , $\sqrt{s} = 3.1 - 4.1$ GeV



CLAS:

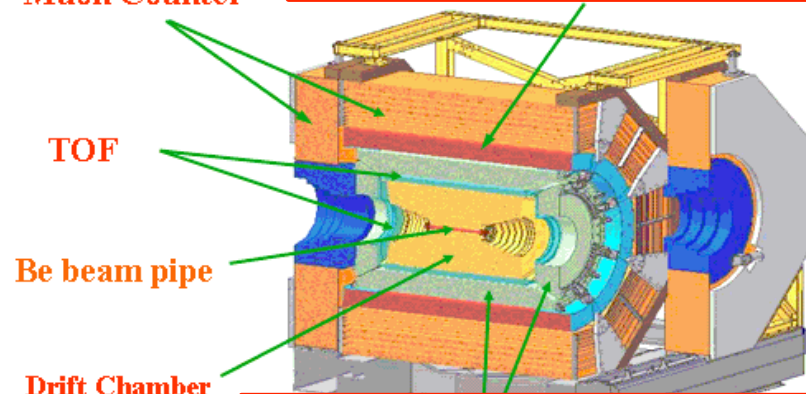
- γ -, e-production
- p target

Muon Counter

TOF

Be beam pipe

Drift Chamber



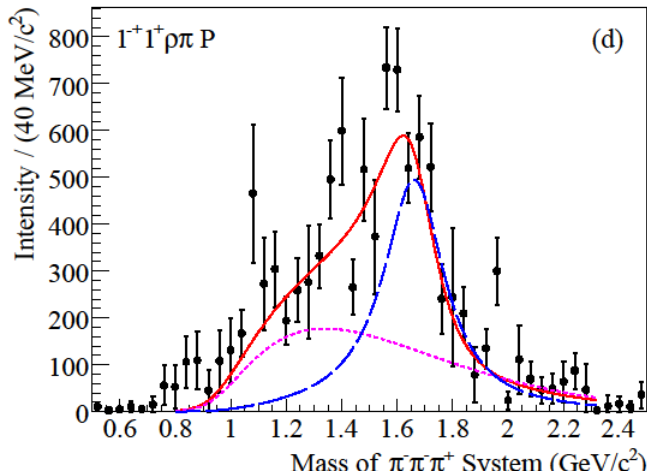
BES III

- e^+e^- , $\sqrt{s} = 3.1 - 4.6$ GeV

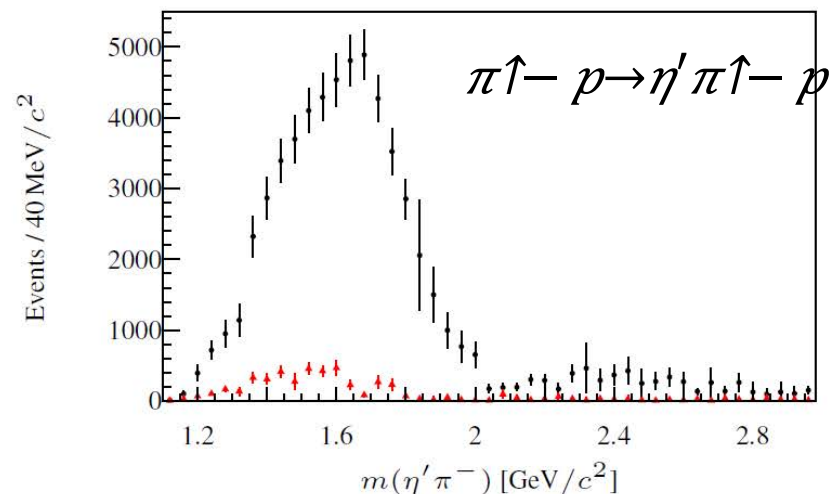
The π_1 (1600)



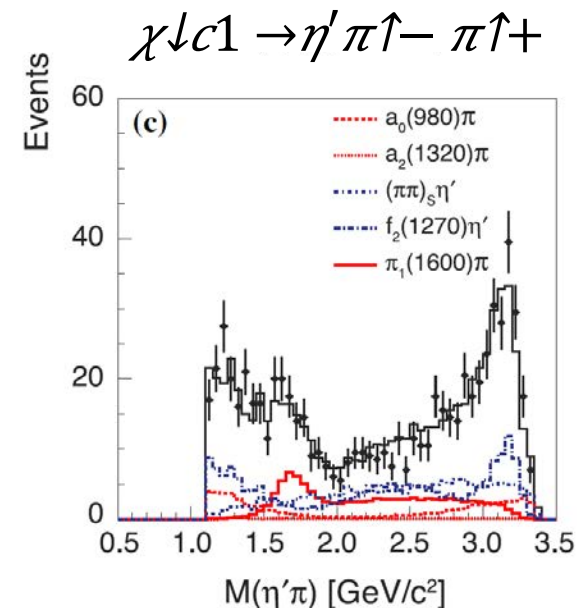
- Resonance-like signal observed
- Large contribution of non-resonant background
- Need to understand origin for a reliable fit of spin-density matrix of high-statistics H_2 data



[Alekseev et al., Phys. Rev. Lett. 104, 241803 (2010)]

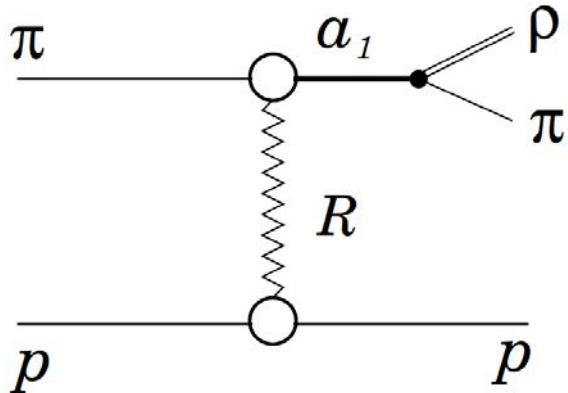


[C. Adolph et al., subm. PLB, arXiv:1408.4286]

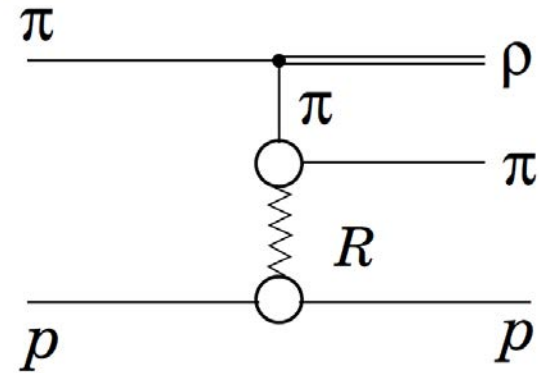


[G.S. Adams et al., Phys. Rev. D 84, 112009 (2011)]

Resonant production



Non-resonant production

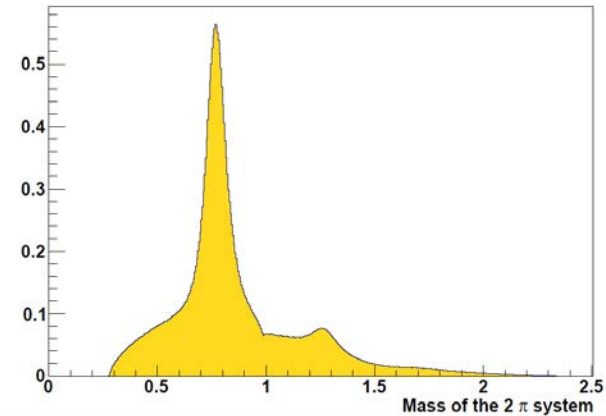


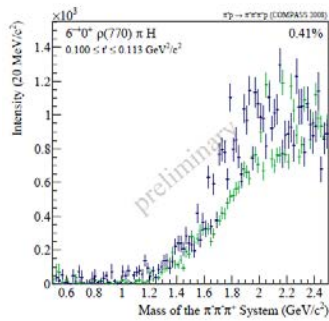
- Generate pure Deck-like events

$$\psi(M_{\pi\pi}, t_{\pi}, t) = \frac{A_{\pi\pi}(M_{\pi\pi}, t_{\pi}) A_{\pi p}(s_{\pi p}, t)}{m_{\pi}^2 - t_{\pi}}$$

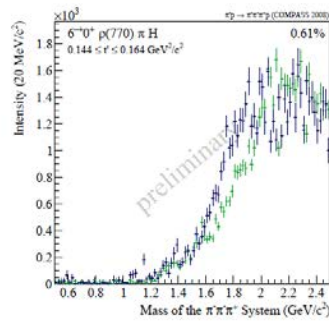
[G. Ascoli et al., Phys. Rev. D 8, 3894 (1973)]

- Pass through Monte Carlo & PWA
- Normalize intensity to data for each wave and sum over t'
- Benchmark on waves w/o resonances, test on exotic wave





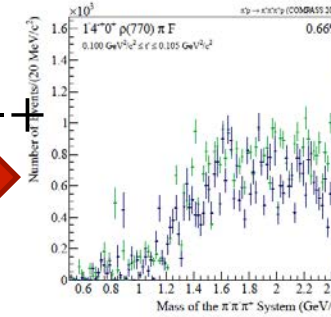
(a)



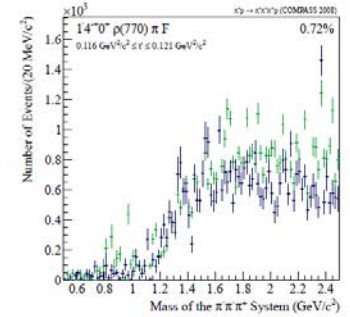
(b)

$6\uparrow-+$

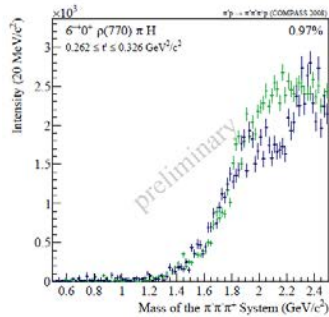
$4\uparrow-+$



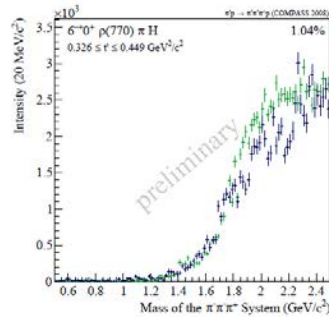
(a)



(b)

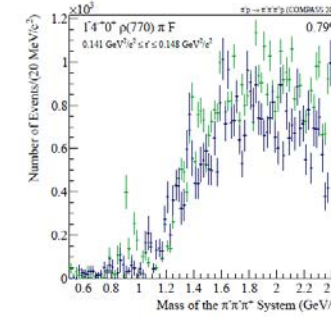


(c)

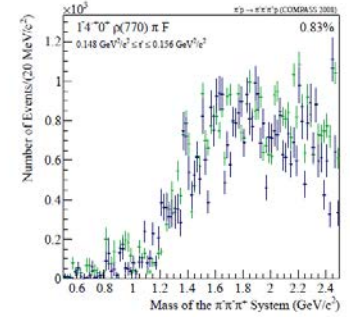


(d)

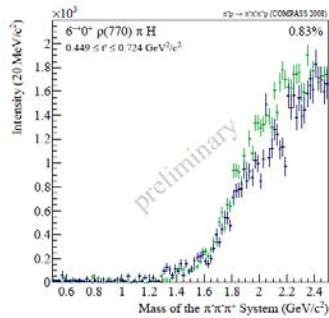
t'



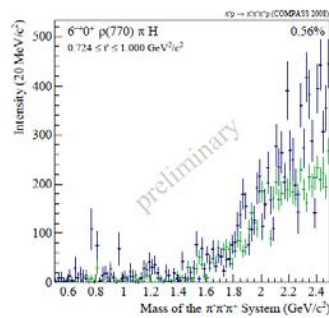
(c)



(d)

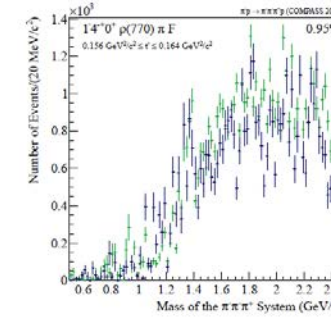


(e)

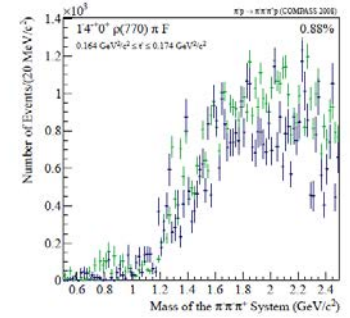


(f)

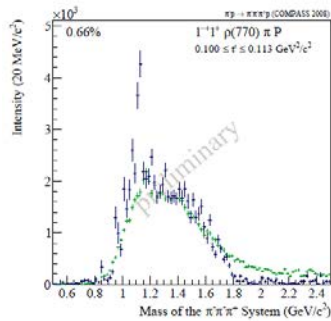
— Data
— Deck



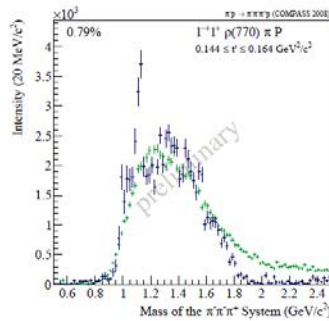
(e)



(f)

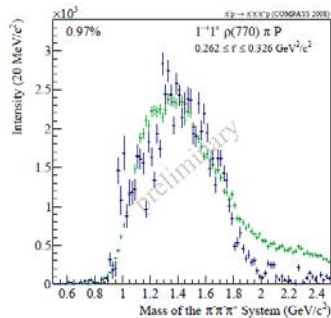


(a)

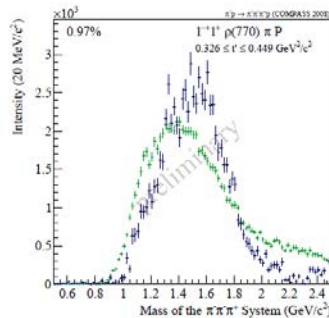


(b)

$1 \uparrow - +$
←

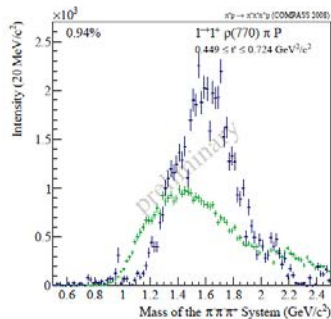


(c)

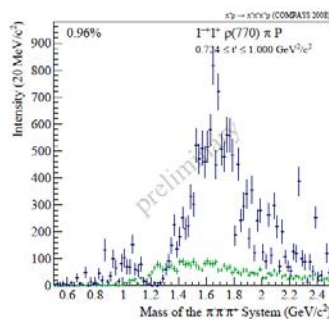


(d)

↓ t'



(e)



(f)

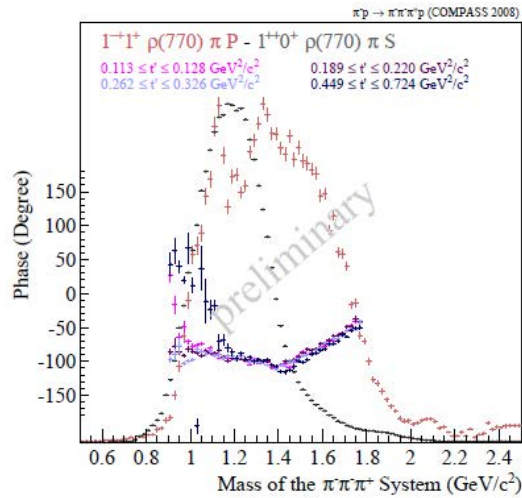
— Data
— Deck

Low values of t' :

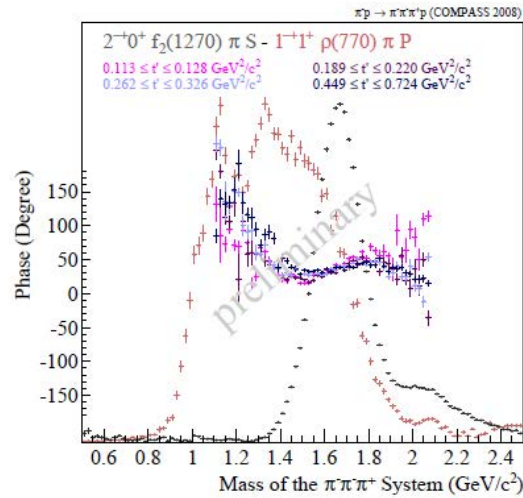
- Mostly non-resonant production
- Good description by Deck model

High values of t' :

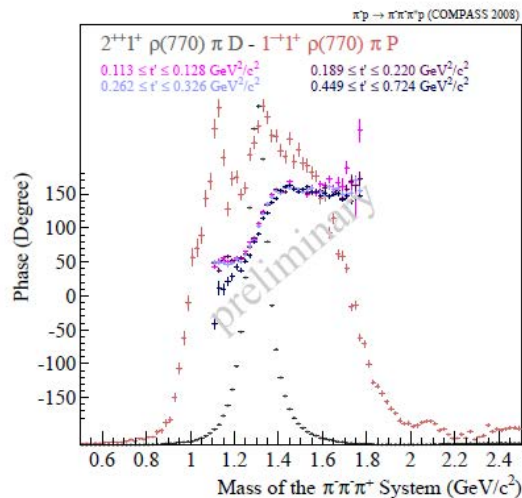
- Resonance appears
- Dominates highest t' - bin



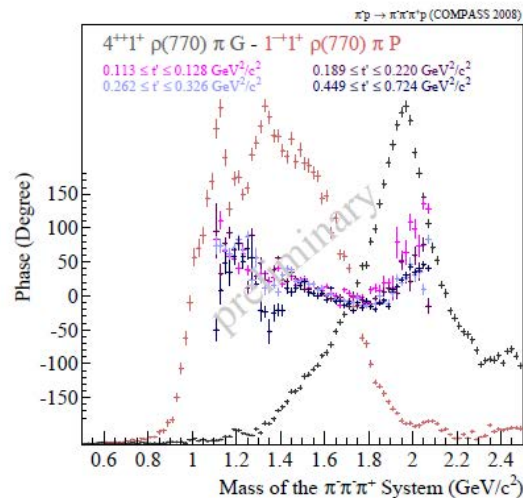
(a)



(b)



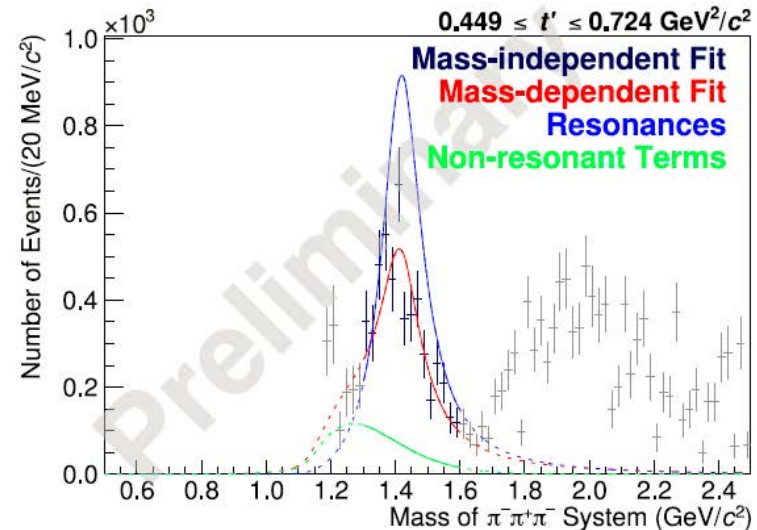
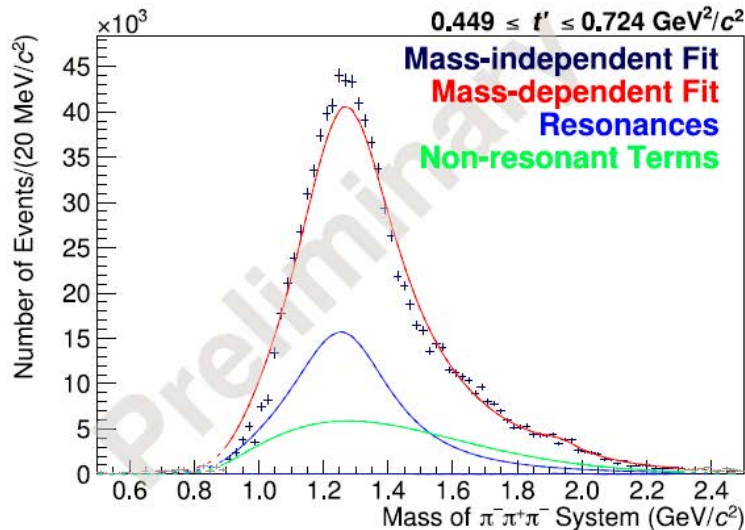
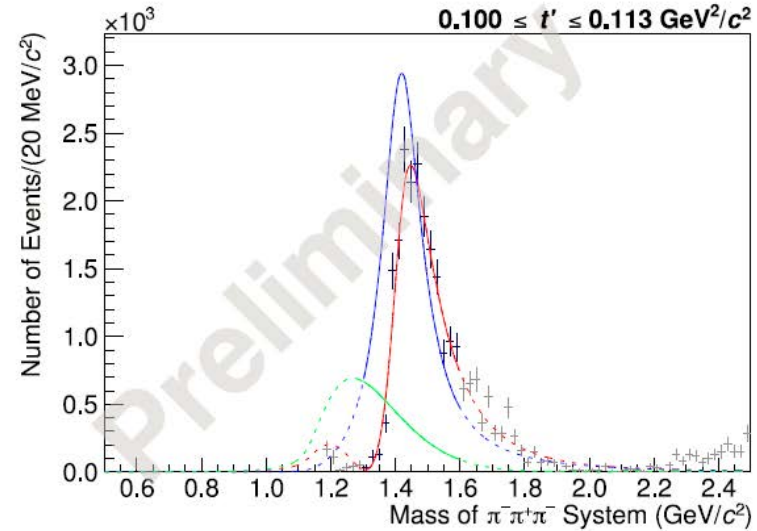
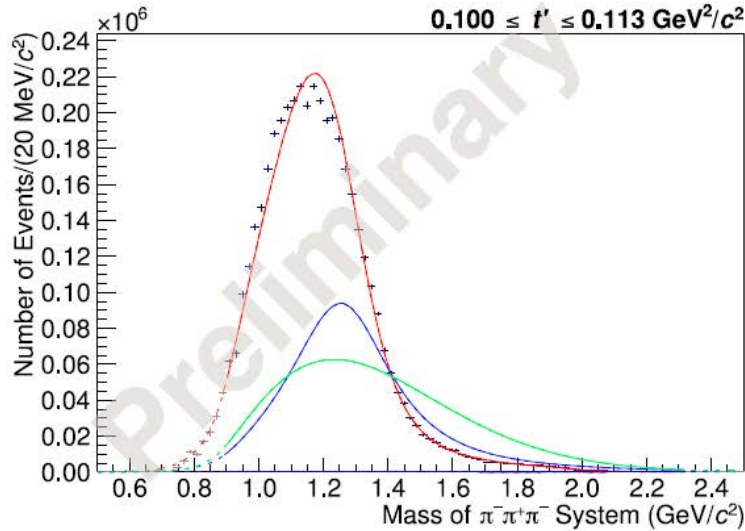
(c)



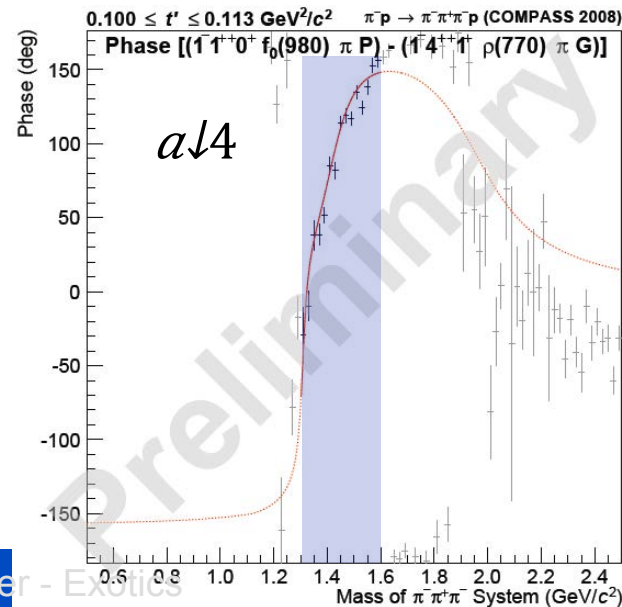
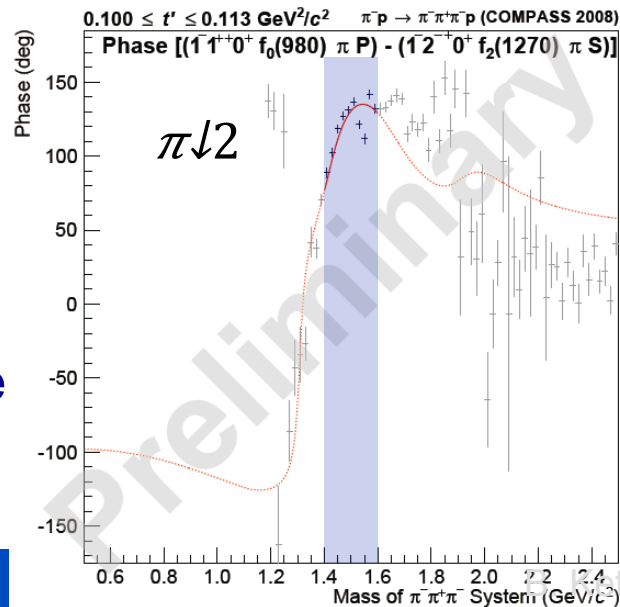
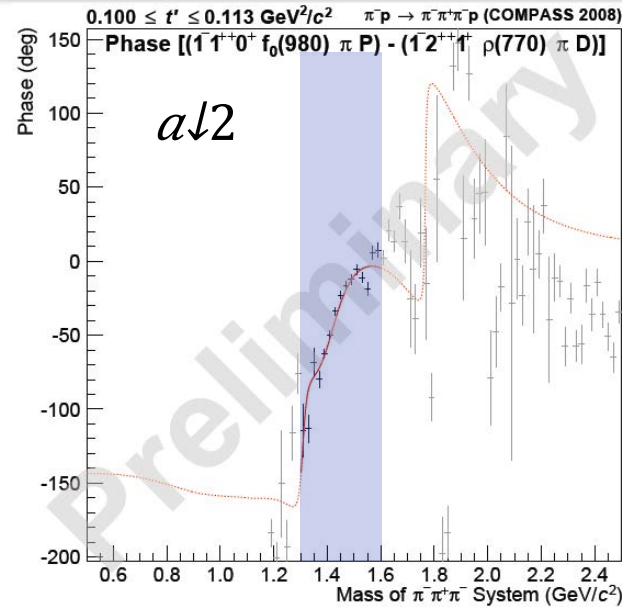
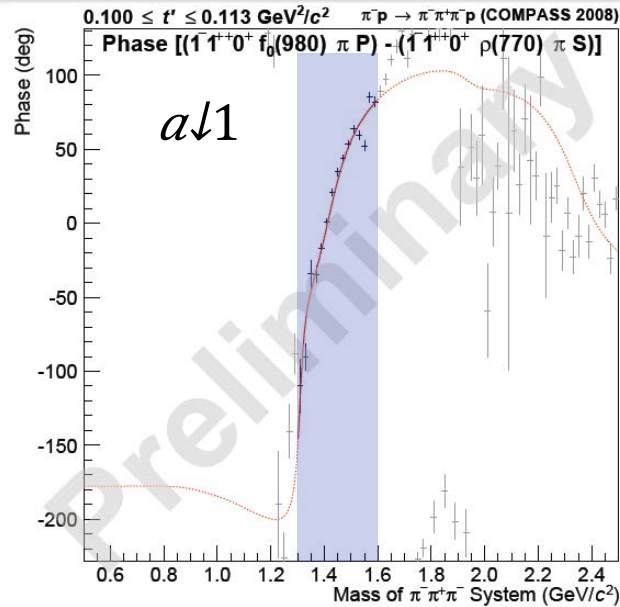
(d)

- Clear phase variation
- Variation $\mathcal{O}(50^\circ)$
- Mild variation with t'

The $a\downarrow 1$ (1420)

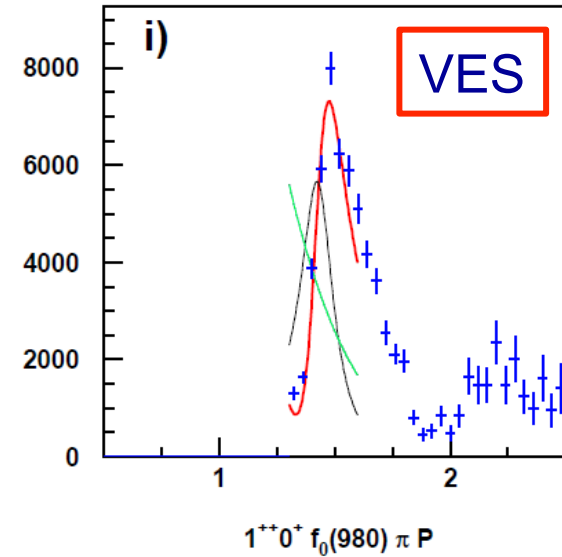
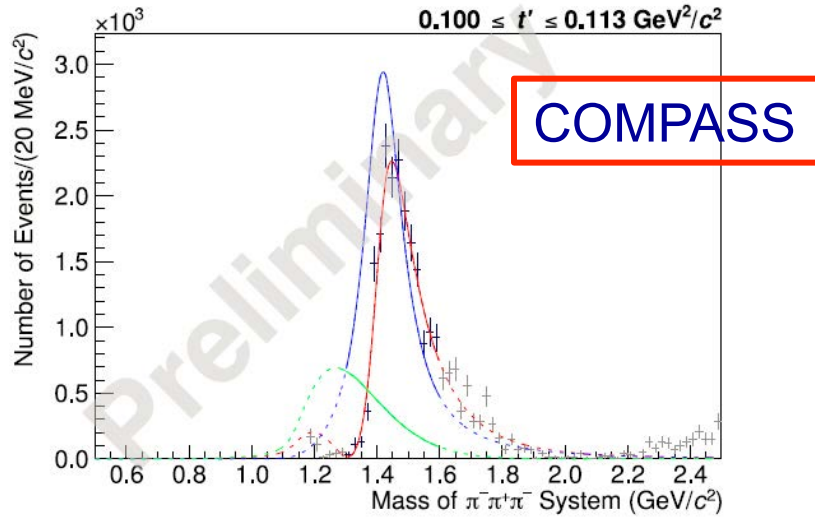


[C. Adolph et al., COMPASS, subm. PRL, arXiv:1501.05732]

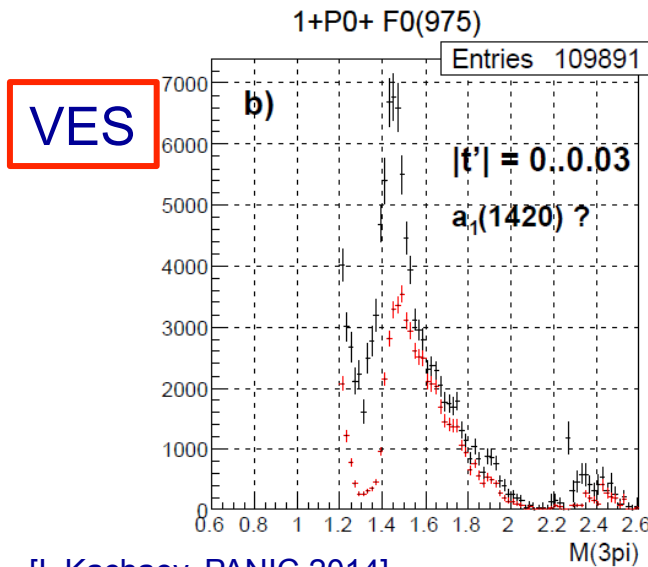


fit range

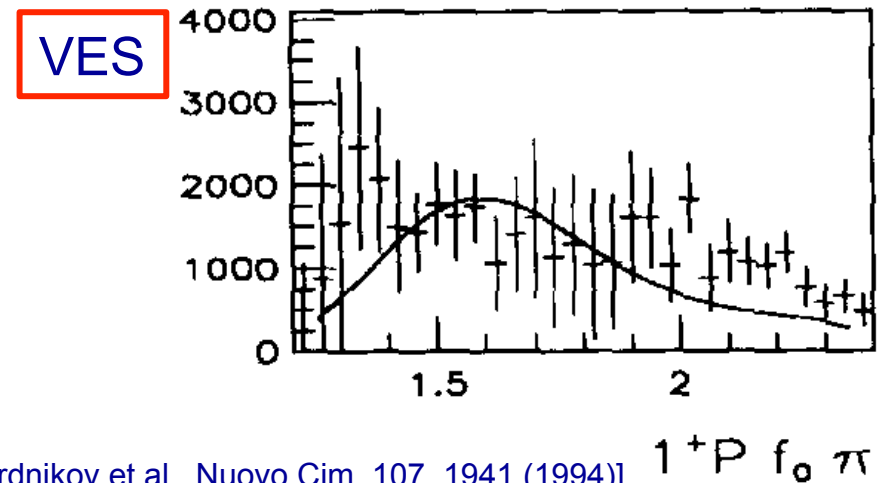
Low t'



[Yu. Khokhlov, PoS Hadron 088 (2014)]



[I. Kachaev, PANIC 2014]



[E.B. Berdnikov et al., Nuovo Cim. 107, 1941 (1994)]

- Has all features of resonance!
 - Narrow peak in intensity
 - Sharp phase motion
 - Description with Breit-Wigner gives

$$M_{a_1} = 1412 - 1422 \text{ MeV}/c^2$$

$$\Gamma_{a_1} = 130 - 150 \text{ MeV}/c^2$$
- Does not fit into quark model
 - Mass difference to $a_1(1260)$ only 160 MeV
 - Much narrower than ground state

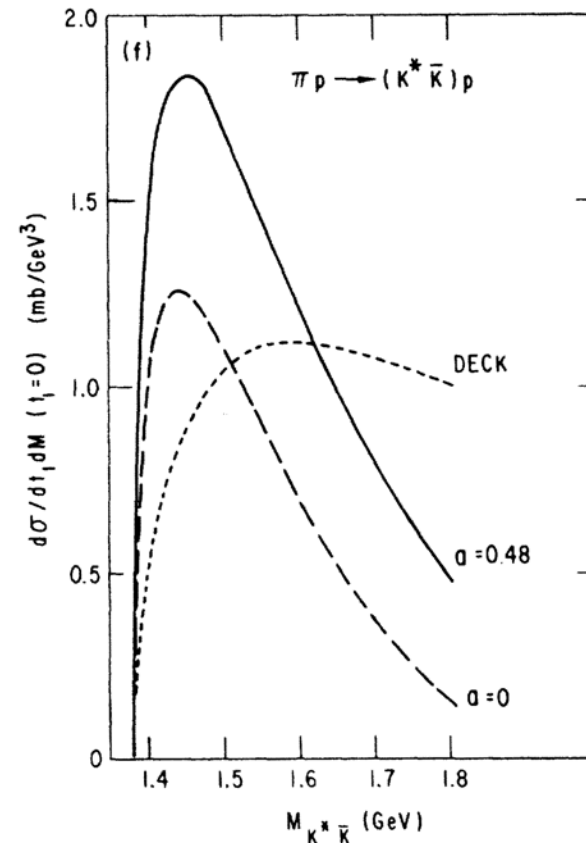
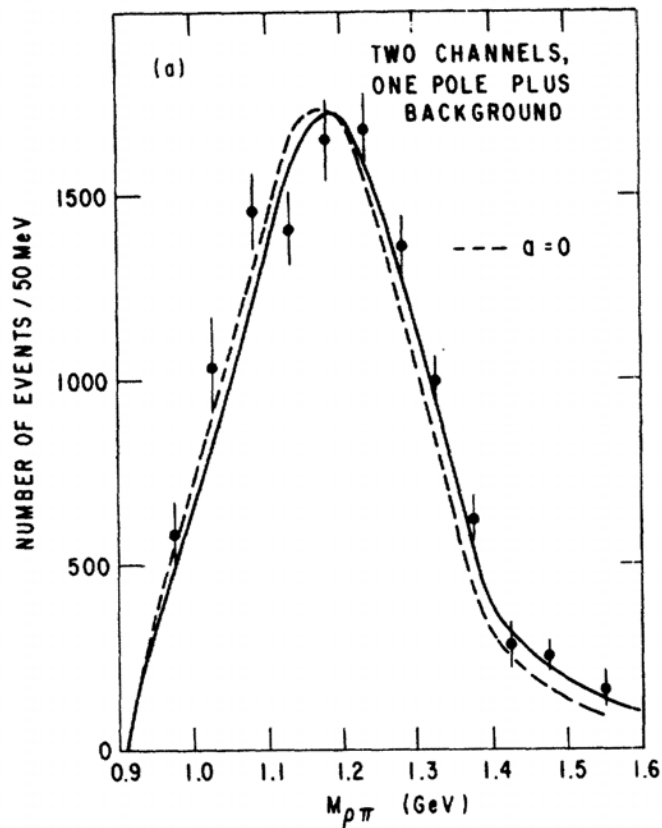
Interpretation:

- 4-quark / molecular state, isospin partner of $f_1(1420)$?
- Dynamic interpretation?

$a_1(1260) \rightarrow \rho\pi, K^* K$

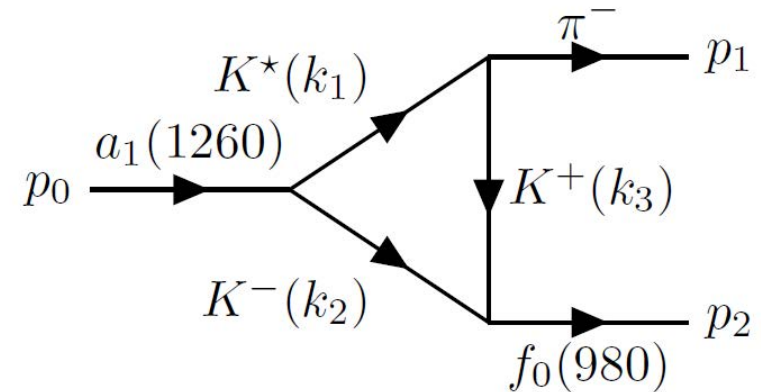
[J.L. Basdevant et al., PRD 16, 657 (1977)]

- Unitary coupled-channel analysis
- Gives sharp rise at $K^* K$ threshold
- Phase motion inherited from $a_1(1260)$?



$$a_1(1260) \rightarrow K^* K + c.c. \rightarrow f_0(980) \pi$$

- Decay of $a_1(1260) \rightarrow K^* K$
- Rescattering of K's to $f_0(980)$
- Triangle diagram
- Logarithmic singularity in amplitude



Two isospin combinations contribute

- $a_1^-(1260) \rightarrow K^{*0} K^- \rightarrow \pi^- K^+ K^- \rightarrow \pi^- f_0,$
- $a_1^-(1260) \rightarrow K^{*-} K^0 \rightarrow \pi^- \bar{K}^0 K^0 \rightarrow \pi^- f_0.$

Landau, 1959:

- Positions of singularity in scalar theory given by

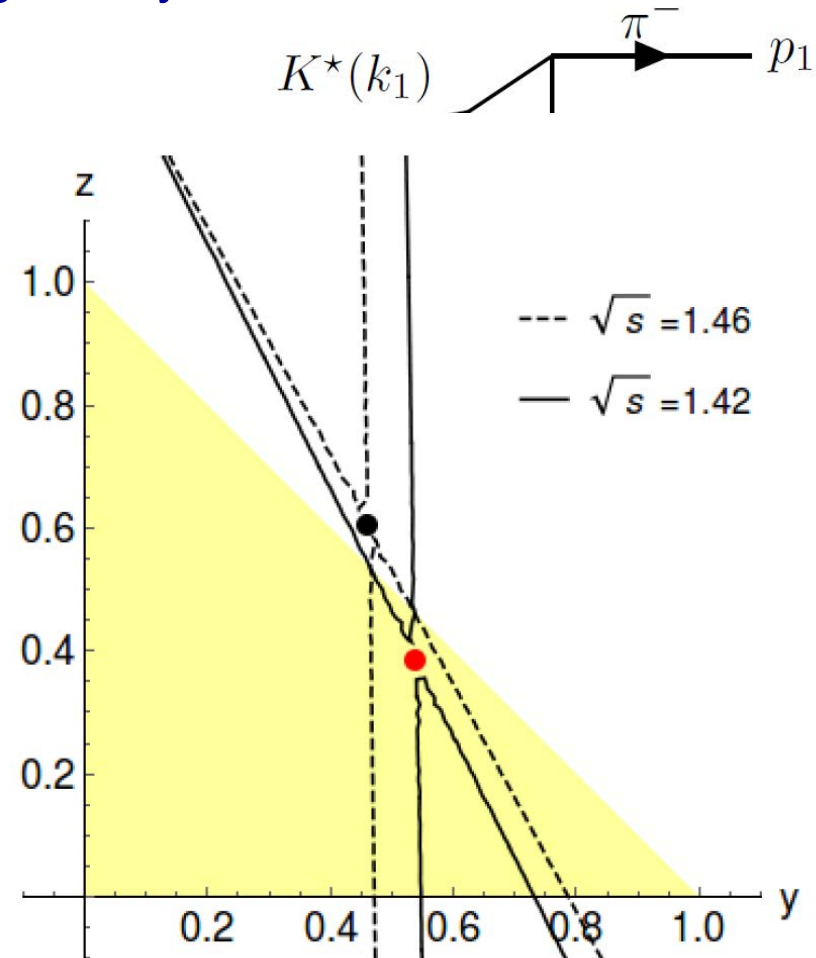
$$\begin{cases} k_i^2 = m_i^2, & i = 1 \dots 3, \\ x k_{1\mu} - y k_{2\mu} + z k_{3\mu} = 0, & x, y, z \in [0, 1], \\ x + y + z = 1, \end{cases}$$

- For case of $a\downarrow 1$ (1260):

- $E\downarrow 1 = 1.42 \text{ GeV}$
- $E\downarrow 2 = 1.46 \text{ GeV}$

Interpretation:

- All intermediate particles on mass shell
- Collinear to each other
- K and K have same velocity only for $E\downarrow$



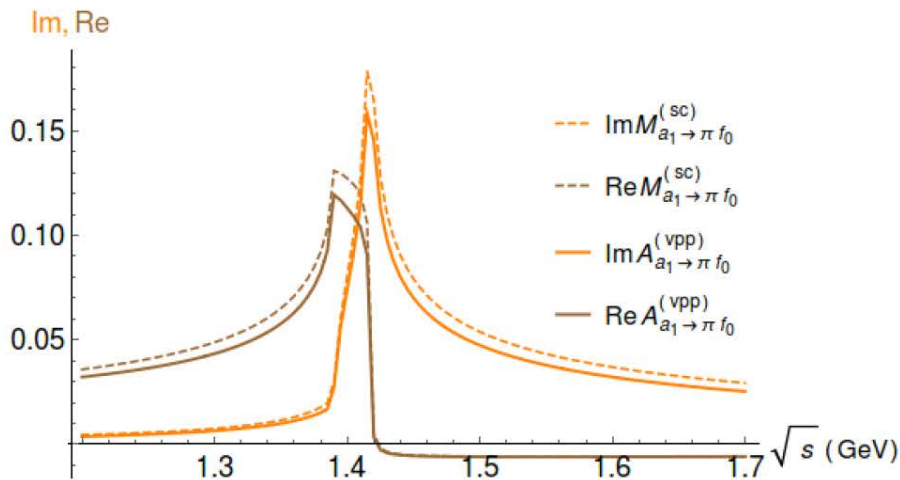
Feynman rules for hadronic processes:

- Scalar case

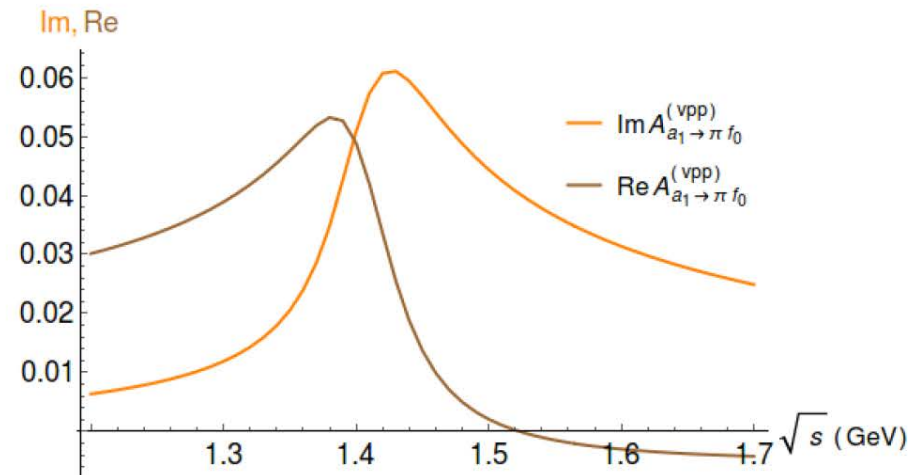
$$M_{a_1 \rightarrow f_0 \pi}^{(sc)} = g^3 \int \frac{d^4 k_1}{(2\pi)^4 i} \frac{1}{(m_1^2 - k_1^2 - i\epsilon)(m_2^2 - (p_0 - k_1)^2 - i\epsilon)(m_3^2 - (k_1 - p_1)^2 - i\epsilon)}$$

- VPP case: denominator carries spin structure

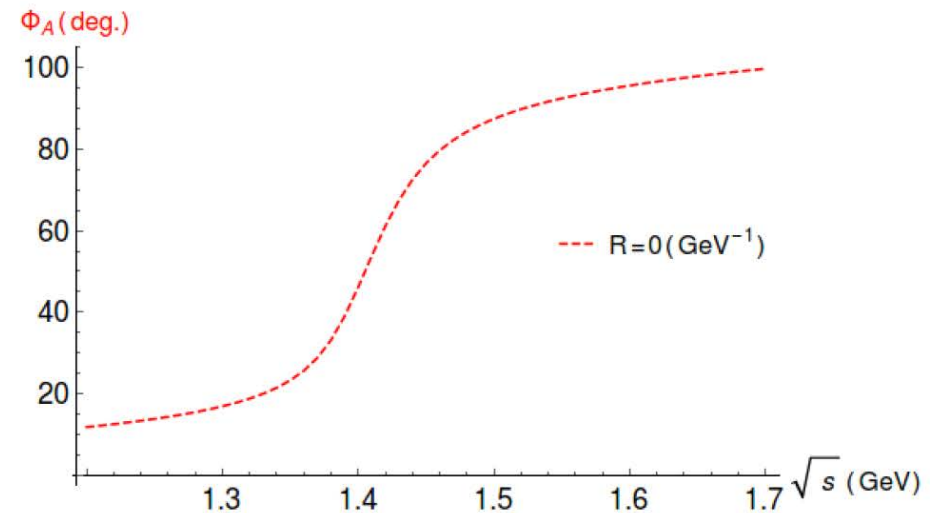
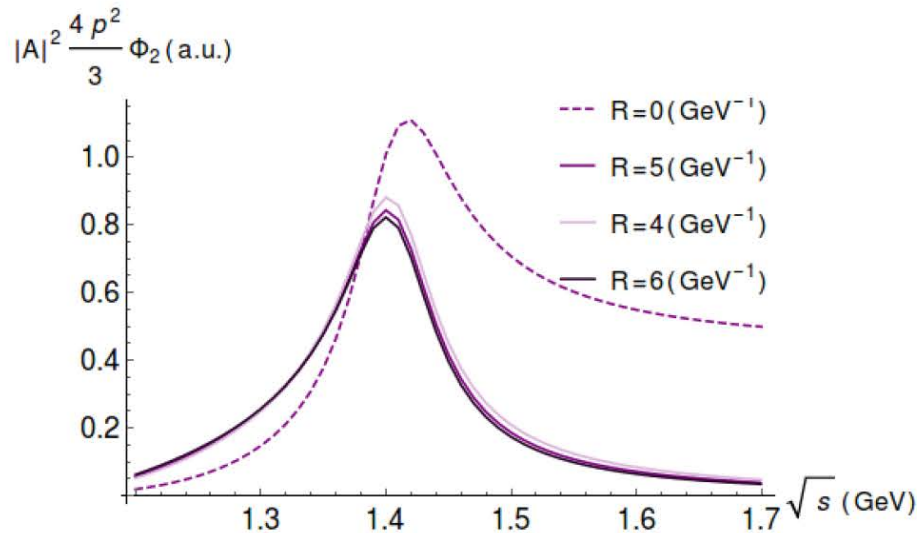
Scalar vs VPP



+ finite width of K^*

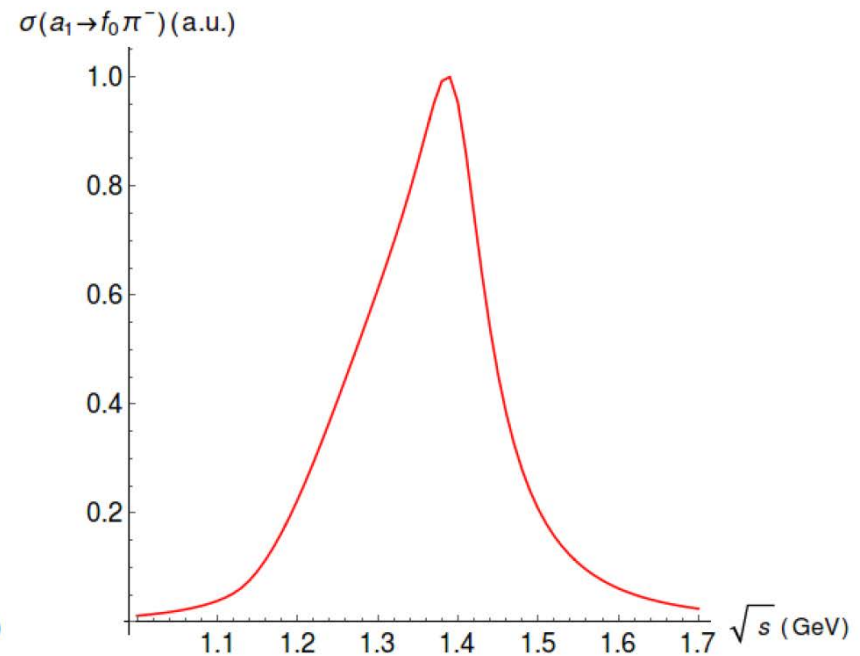
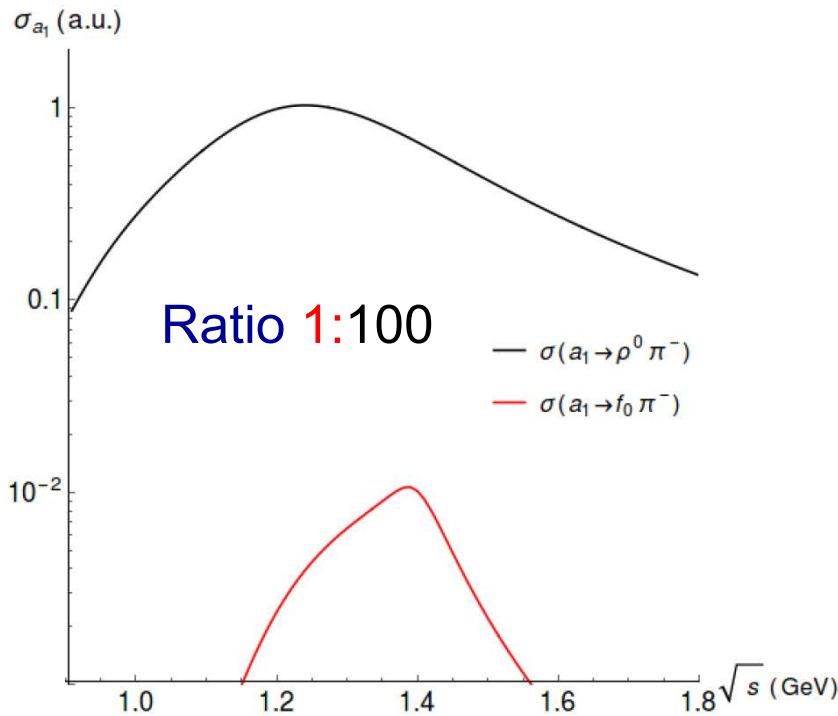
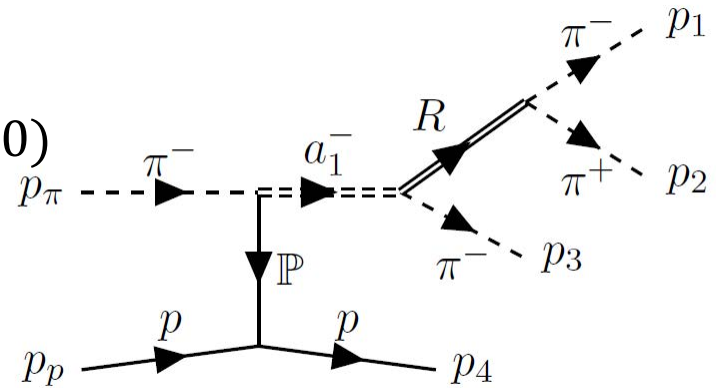


- Finite width of $K1^*$
- Suppression of P-wave tail due to $K1^* \rightarrow K\pi$ decay
 - Blatt-Weisskopf barrier factors
 - Exponential correction factors for finite meson-size
 - Introduce left-hand singularity in the amplitude



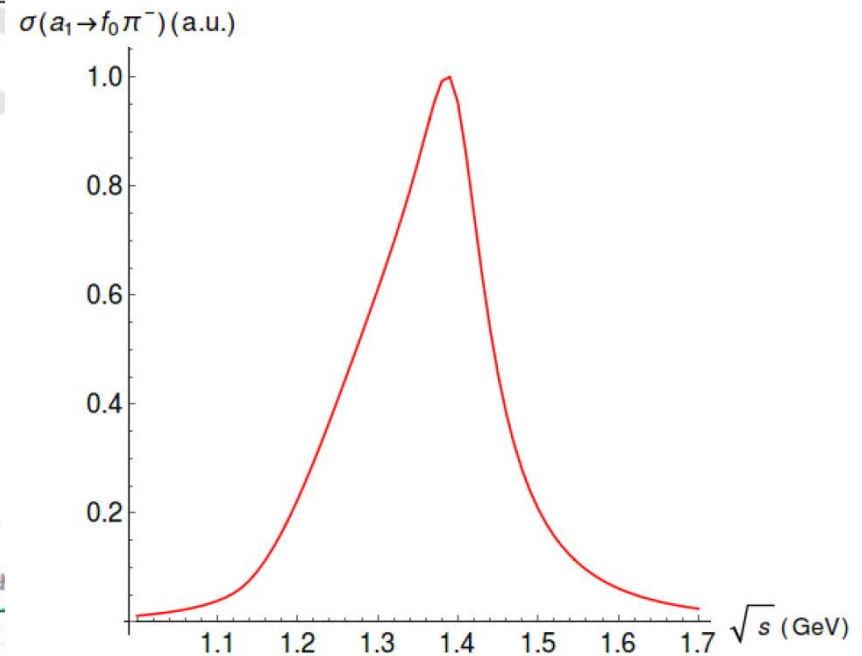
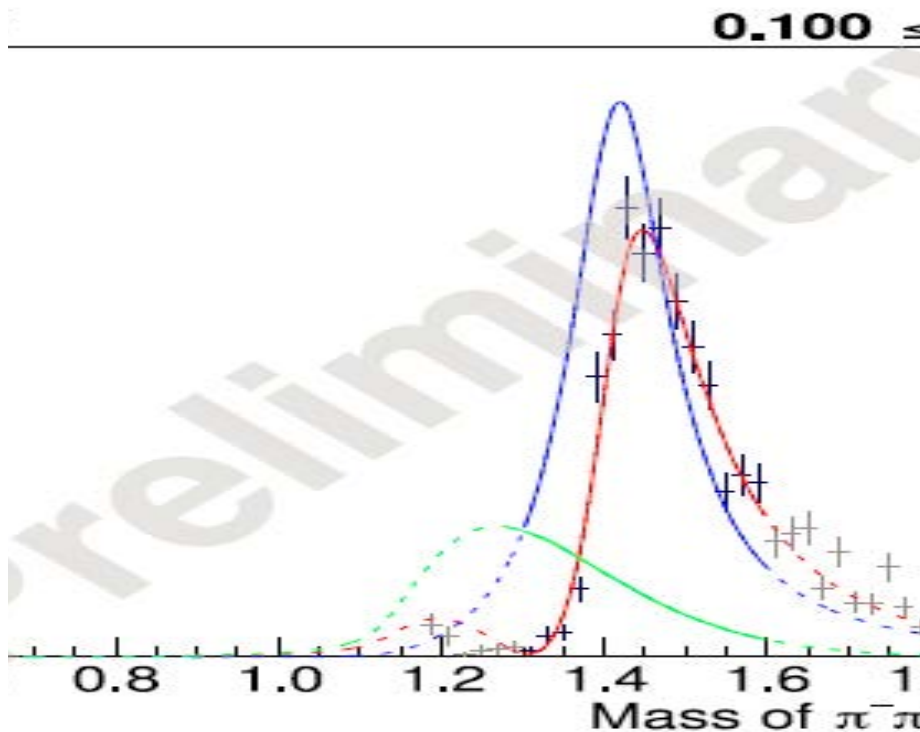
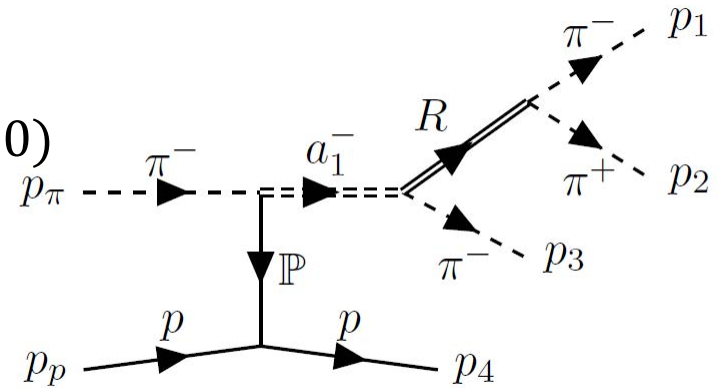
Consider full process:

- Diffractive production of resonant $a_1(1260)$
- Direct decay to $\rho\pi$
- Decay to $f_0\pi$ only via triangle diagram



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- Diffractive production of resonant $a_1(1260)$
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- Hadron spectroscopy is entering **precision era**
- Statistical uncertainties very small
- Systematic **model uncertainties** become dominant
- Spin-exotic $\pi_1(1600)$: (re-) observed by COMPASS, VES, CLEO-c
 - Large statistics of COMPASS \Rightarrow **2D-PWA in bins of $m_{\downarrow X}$ and t_{\uparrow}**
 - \Rightarrow Strong non-resonant contribution to $\rho\pi, \eta'\pi$
 - \Rightarrow Can be well described by Deck effect
 - \Rightarrow Resonant part dominates at high t_{\uparrow}
- New axial vector meson observed in $a_{\downarrow 1}(1420) \rightarrow f_{\downarrow 0}(980)\pi$?
 - Has all features of a genuine resonance
 - Possible explanation: pseudo-resonance due to triangle singularity
- Future: identify exotic multiplets and measure decay pattern
- General: **amplitudes need to satisfy analyticity and unitarity!**



Spare Slides