

Measurement of the charged-pion polarisability and more on chiral dynamics from COMPASS

Jan M. Friedrich

on behalf of the COMPASS collaboration



Symmetries and Spin

Prague 2015, June 28



Bundesministerium
für Bildung
und Forschung



Common Muon and Proton Apparatus for Structure and Spectroscopy





Common Muon and Proton Apparatus for Structure and Spectroscopy

CERN SPS: protons ~ 400 GeV (5 – 10 sec spills)

- secondary $\pi, K, (\bar{p})$: up to $2 \cdot 10^7/s$ (typ. $5 \cdot 10^6/s$)
Nov. 2004, 2008-09, 2012:
hadron spec. & Primakoff reactions
- tertiary muons: $4 \cdot 10^7 / s$
2002-04, 2006-07, 2010-11: spin structure of the nucleon

LHC

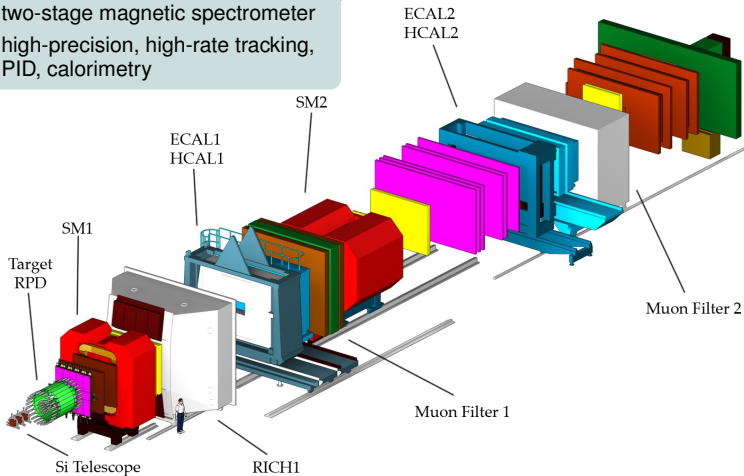
COMPASS

SPS



Fixed-target experiment

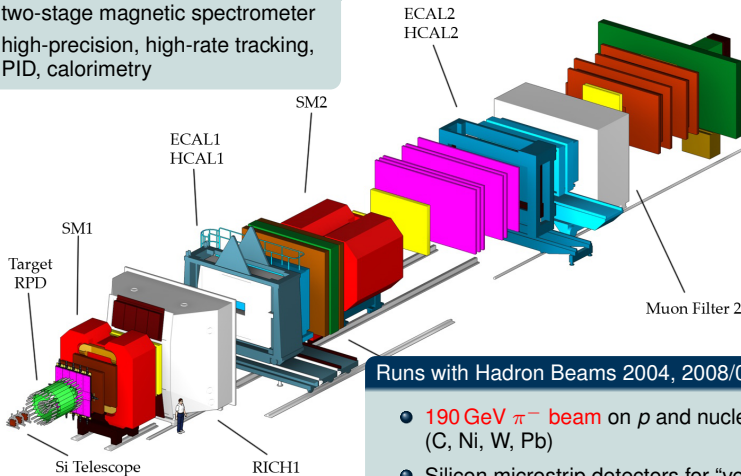
- two-stage magnetic spectrometer
- high-precision, high-rate tracking, PID, calorimetry





Fixed-target experiment

- two-stage magnetic spectrometer
- high-precision, high-rate tracking, PID, calorimetry

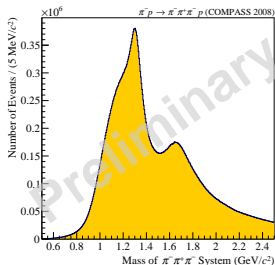


Runs with Hadron Beams 2004, 2008/09, 2012

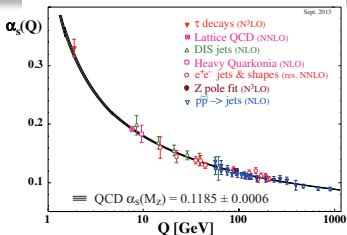
- **190 GeV π^- beam** on p and nuclear targets (C, Ni, W, Pb)
- Silicon microstrip detectors for “vertexing”
- recoil and (digital) ECAL triggers



- lepton scattering at high Q^2
→ partonic structure of the nucleons
- since 2012: exclusive hard processes (DVCS, HEMP) for GPD, pol. Drell-Yan



- scattering of pions (and kaons) in nuclear Coulomb field
→ low-energetic meson-photon reactions
 $\pi\gamma \rightarrow \pi\gamma$ (pion polarisability), $\pi\gamma \rightarrow 3\pi$ (chiral dynamics, radiative couplings)



from: PDG, S. Bethke et al.

- diffractive dissociation of pions and kaons into multi-meson final states
→ meson spectrometry

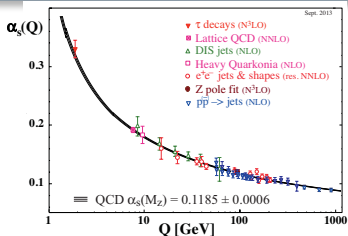
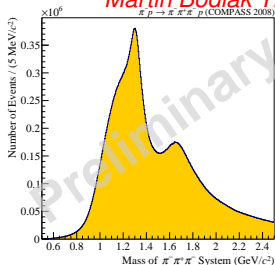


GPDP's: *Oleg Kouznetsov tomorrow*
 transverse spin: *Anna Martin tomorrow,*
Jan Matousek Thursday

pol. DY: *Riccardo Longo yesterday*

pol. DY target: *Michael Pesek tomorrow*

new DAQ: *Josef Novy tomorrow,*
Martin Bodlak Thursday



from: PDG, S. Bethke et al.

- diffractive dissociation of pions and kaons into multi-meson final states
 → meson spectrometry

- scattering of pions (and kaons) in nuclear Coulomb field
 → low-energetic meson-photon reactions
 $\pi\gamma \rightarrow \pi\gamma$ (pion polarisability), $\pi\gamma \rightarrow 3\pi$ (chiral dynamics, radiative couplings)

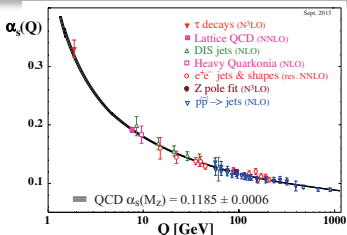
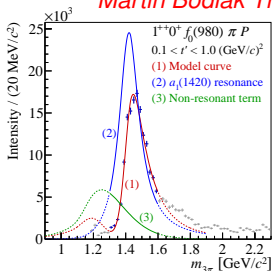


GPDP's: *Oleg Kouznetsov tomorrow*
 transverse spin: *Anna Martin tomorrow,*
Jan Matousek Thursday

pol. DY: *Riccardo Longo yesterday*

pol. DY target: *Michael Pesek tomorrow*

new DAQ: *Josef Novy tomorrow,*
Martin Bodlak Thursday



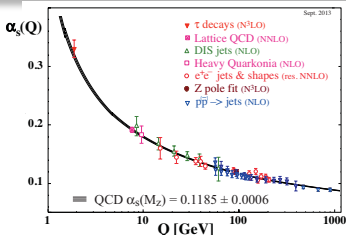
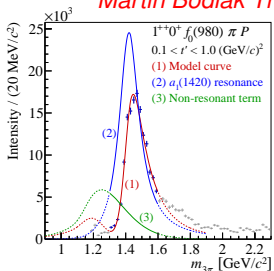
from: PDG, S. Bethke et al.

- diffractive dissociation of pions and kaons into multi-meson final states
 → meson spectrometry
PRL on discovery of $a_1(1420)$ acc'd last week
candidate for $s\bar{s}u\bar{d}$

- scattering of pions (and kaons) in nuclear Coulomb field
 → low-energetic meson-photon reactions
 $\pi\gamma \rightarrow \pi\gamma$ (pion polarisability), $\pi\gamma \rightarrow 3\pi$ (chiral dynamics, radiative couplings)



GPD's: *Oleg Kouznetsov tomorrow*
 transverse spin: *Anna Martin tomorrow,*
 Jan Matousek Thursday
 pol. DY: *Riccardo Longo yesterday*
 pol. DY target: *Michael Pesek tomorrow*
 new DAQ: *Josef Novy tomorrow,*
 Martin Bodlak Thursday



- diffractive dissociation of pions and kaons into multi-meson final states
 → meson spectrometry
PRL on discovery of $a_1(1420)$ acc'd last week
candidate for $s\bar{s}u\bar{d}$

- scattering of pions (and kaons) in nuclear Coulomb field
 → low-energetic meson-photon reactions *(focus of this talk)*
 $\pi\gamma \rightarrow \pi\gamma$ (pion polarisability), $\pi\gamma \rightarrow 3\pi$ (chiral dynamics, radiative couplings)



- decoupling of left- and right-handed massless quarks into $SU(N_f)_R \times SU(N_f)_L$
- hidden by spontaneous breaking
- explicit breaking by $m_q \neq 0$
- series expansion of QCD in chiral-symmetric terms and symmetry-breaking “perturbation”
- pion identified with the “Goldstone mode”



Quark-gluon dynamics when the interaction *is* strong

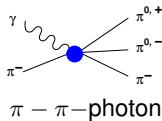
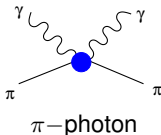
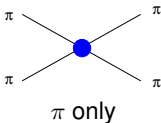


complicated system of interacting quarks and gluons

ChPT \rightarrow



effective degrees of freedom
at low energy: mass, charge, spin, effective (self-)coupling

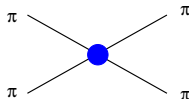




Chiral Perturbation Theory vs. Experiment

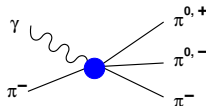
- pion scattering lengths: 2-loop predictions

- $a_0^0 m_\pi = 0.220 \pm 0.005$ confirmed by E865 in $K^+ \rightarrow \pi^+ \pi^- e^+ \nu_e$
- $(a_0^0 - a_0^2) m_\pi = 0.264 \pm 0.006$ confirmed by NA48 in 0.268 ± 0.010 $K^+ \rightarrow \pi^+ \pi^0 \pi^0$



- pion scattering including a real photon

- leading-order prediction from ChPT
 - \leftrightarrow pion scattering lengths combined with photon coupling
- **chiral loop contribution**
theory prediction available, no measurement

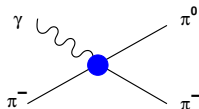




More pion-photon reactions

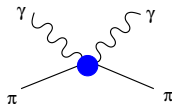
- **chiral anomaly** $F_{3\pi}$

- established on 10% level
- further development: inclusion of the ρ resonance
theoretical work by Kubis, Hoferichter, Sakkas
PRD86(2012)116009



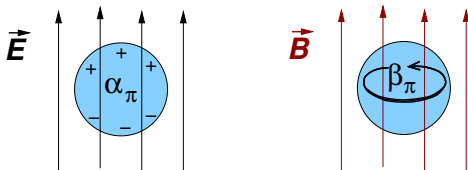
- **pion polarisability: electric** α_π , **magnetic** β_π

- contribution to Compton scattering
- **ChPT prediction** obtained by the relation to
 $\pi^+ \rightarrow e^+ \nu_e \gamma$ [Gasser, Ivanov, Sainio, Nucl. Phys. B745, 2006]
[PIBETA, M. Bychkov et al., PRL 103, 051802, 2009]
- ChPT prediction **contradicting** the experimental findings
(prior to our analysis)





Pion polarisability and ChPT



pion polarisabilities α_π, β_π in units of 10^{-4} fm^3

size of the pion $\sim 1 \text{ fm}^3$ [cf. atoms: polarisability \approx size $\approx 1 \text{ \AA}^3$]

Theory: ChPT (2-loop) prediction:

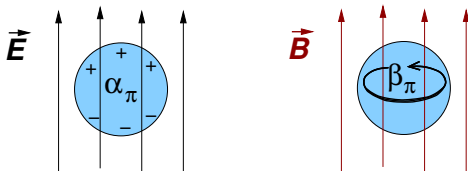
$$\begin{aligned} \alpha_\pi - \beta_\pi &= 5.7 \pm 1.0 \\ \alpha_\pi + \beta_\pi &= 0.16 \pm 0.1 \end{aligned}$$

experiments for $\alpha_\pi - \beta_\pi$ lie in the range $4 \dots 14$

($\alpha_\pi + \beta_\pi = 0$ assumed)



Pion polarisability and ChPT



pion polarisabilities α_π, β_π in units of 10^{-4} fm^3

size of the pion $\sim 1 \text{ fm}^3$ [cf. atoms: polarisability \approx size $\approx 1 \text{ \AA}^3$]

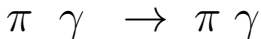
Theory: ChPT (2-loop) prediction:

$$\alpha_\pi = 2.93 \pm 0.5$$

$$\beta_\pi = -2.77 \pm 0.5$$

experiments for α_π lie in the range $2 \dots 7$

($\alpha_\pi + \beta_\pi = 0$ assumed)



- Two kinematic variables, in CM: total energy \sqrt{s} , scattering angle θ_{cm}

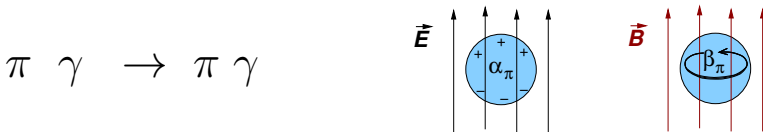
$$\frac{d\sigma_{\pi\gamma}}{d\Omega_{cm}} = \frac{\alpha^2 (s^2 z_+^2 + m_\pi^4 z_-^2)}{s (s z_+ + m_\pi^2 z_-)^2} - \frac{\alpha m_\pi^3 (s - m_\pi^2)^2}{4s^2 (s z_+ + m_\pi^2 z_-)} \cdot \mathcal{P}$$

$$\mathcal{P} = z_-^2 (\alpha_\pi - \beta_\pi) + \frac{s^2}{m_\pi^4} z_+^2 (\alpha_\pi + \beta_\pi) - \frac{(s - m_\pi^2)^2}{24s} z_-^3 (\alpha_2 - \beta_2)$$

$$z_\pm = 1 \pm \cos \theta_{cm}$$



Pion Compton Scattering



- Two kinematic variables, in CM: total energy \sqrt{s} , scattering angle θ_{cm}

$$\frac{d\sigma_{\pi\gamma}}{d\Omega_{cm}} = \frac{\alpha^2 (s^2 z_+^2 + m_\pi^4 z_-^2)}{s (s z_+ + m_\pi^2 z_-)^2} - \frac{\alpha m_\pi^3 (s - m_\pi^2)^2}{4s^2 (s z_+ + m_\pi^2 z_-)} \cdot \mathcal{P}$$

$$\mathcal{P} = z_-^2 (\alpha_\pi - \beta_\pi) + \frac{s^2}{m_\pi^4} z_+^2 (\alpha_\pi + \beta_\pi) - \frac{(s - m_\pi^2)^2}{24s} z_-^3 (\alpha_2 - \beta_2)$$

$$z_\pm = 1 \pm \cos \theta_{cm}$$



latest publication on the pion polarisability



Technische Universität München

PRL **114**, 062002 (2015)

PHYSICAL REVIEW LETTERS

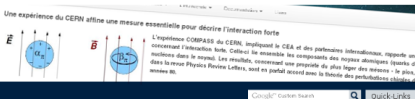
week ending
13 FEBRUARY 2015

Measurement of the Charged-Pion Polarizability

C. Adolph,⁸ R. Akhunzyanov,⁷ M. G. Alexeev,²⁷ G. D. Alexeev,⁷ A. Amoroso,^{27,29} V. Andrieux,²² V. Anosov,⁷
... [213 authors]
(COMPASS Collaboration)

(Received 2 June 2014; revised manuscript received 24 December 2014; published 10 February 2015)

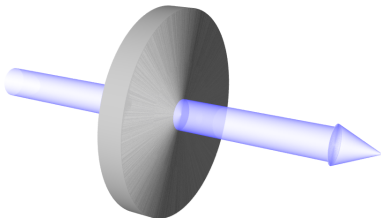
The COMPASS collaboration at CERN has investigated pion Compton scattering, $\pi^- \gamma \rightarrow \pi^- \gamma$, at center-of-mass energy below 3.5 pion masses. The process is embedded in the reaction $\pi^- \text{Ni} \rightarrow \pi^- \gamma \text{Ni}$, which is initiated by 190 GeV pions impinging on a nickel target. The exchange of quasireal photons is selected by isolating the sharp Coulomb peak observed at smallest momentum transfers, $Q^2 < 0.0015 \text{ (GeV}/c)^2$. From a sample of 63 000 events, the pion electric polarizability is determined to be $\alpha_\pi = (2.0 \pm 0.6_{\text{stat}} \pm 0.7_{\text{sys}}) \times 10^{-4} \text{ fm}^3$ under the assumption $\alpha_\pi = -\beta_\pi$, which relates the electric and magnetic dipole polarizabilities. It is the most precise measurement of this fundamental low-energy parameter of strong





Principle of the COMPASS measurement

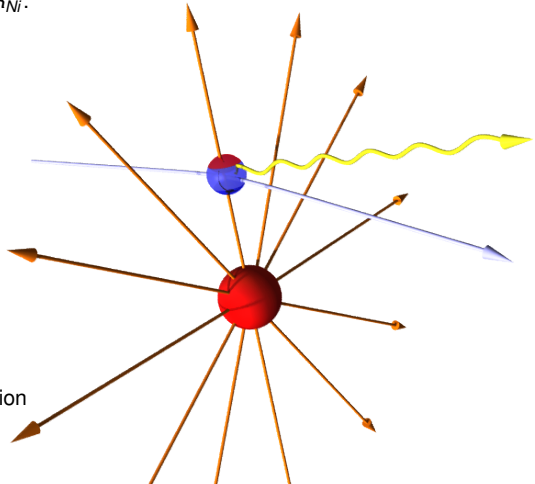
- high-energetic pion beam on 4mm nickel disk
- observe scattered pions in coincidence with produced hard photons
- study of cross-section shape





Polarisability effect in Primakoff technique

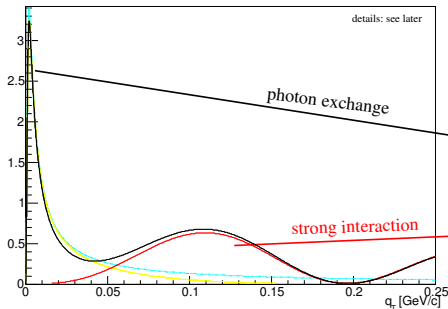
- Charged pions traverse the nuclear **electric field**
 - typical field strength at $d = 5R_N$:
 $E \approx 300 \text{ kV/fm}$
- Bremsstrahlung process:
 - particles scatter off **equivalent photons**
 - tiny momentum transfer
 $Q^2 \approx 10^{-5} \text{ GeV}^2/c^2$
 - pion/muon (quasi)-real Compton scattering
- Polarisability contribution
 - Compton cross-section typically diminished
 - corresponding charge separation
 $\approx 10^{-5} \text{ fm} \cdot e$





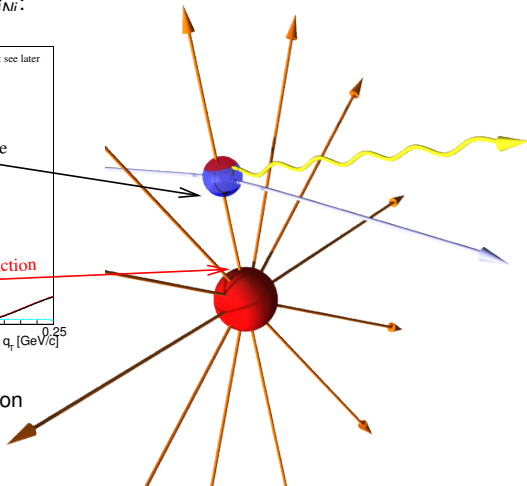
Polarisability effect in Primakoff technique

- Charged pions traverse the nuclear **electric field**
 - typical field strength at $d = 5R_{MI}$:



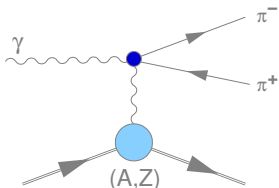
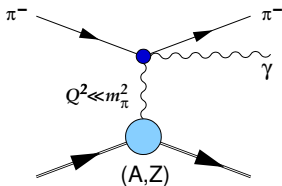
typically diminished

- corresponding charge separation $\approx 10^{-5} \text{ fm} \cdot e$

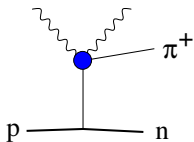




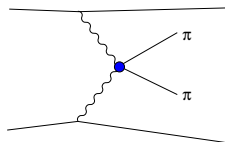
Pion Compton scattering: embedding the process



Primakoff processes



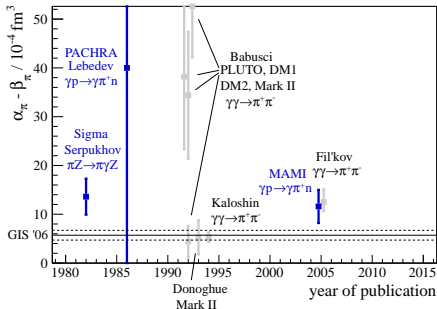
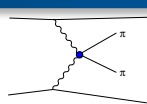
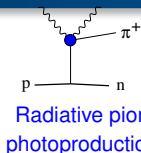
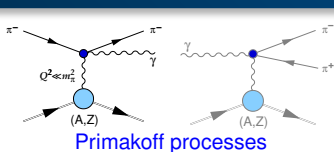
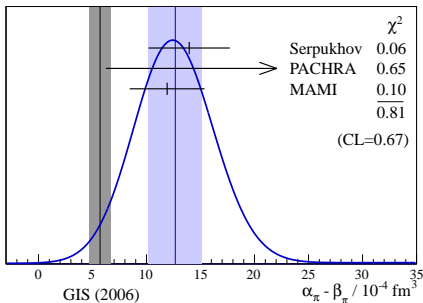
Radiative pion photoproduction



Photon-Photon fusion



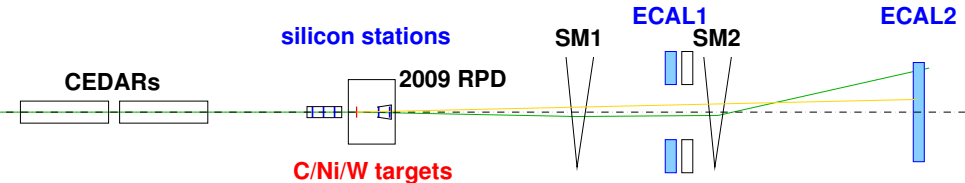
Pion polarisability: world data before COMPASS

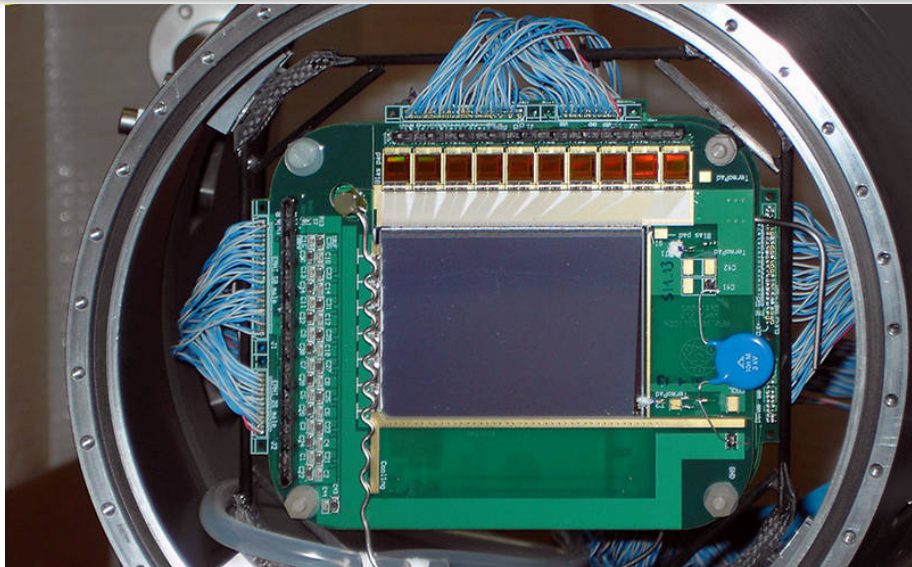
world avg.: 12.7 ± 2.5 

GIS'06: ChPT prediction, Gasser, Ivanov, Sainio, NPB745 (2006), plots: T. Nagel, PhD
 Fil'kov analysis objected by Pasquini, Drechsel, Scherer PRC81, 029802 (2010)



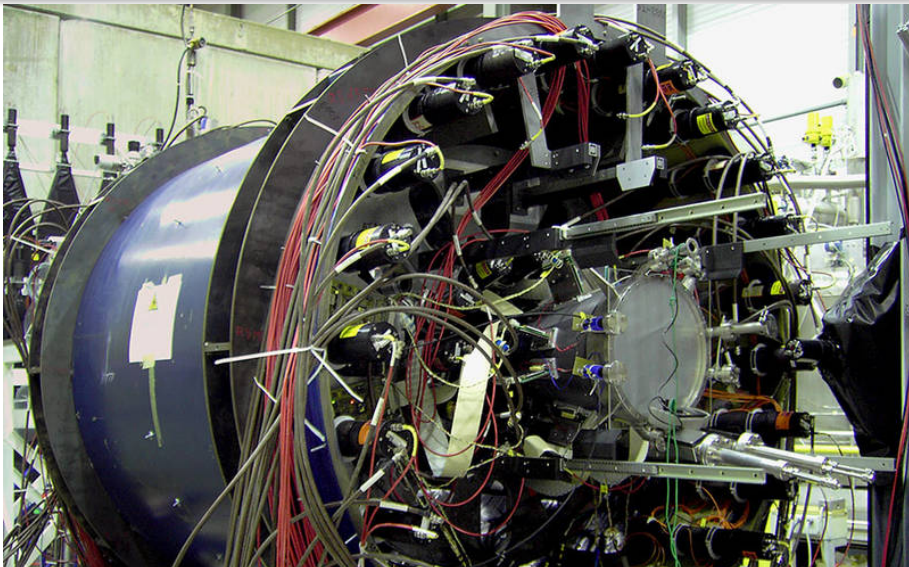
Principle of the measurement



Silicon detector module, two-sided $\sigma_{x,y} \sim 5\mu\text{m}$ 



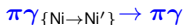
Silicon cryostat in the recoil detector





Extraction of the pion polarisability

- Identify exclusive reactions



at smallest momentum transfer $< 0.001 \text{ GeV}^2/c^2$

- Assuming $\alpha_\pi + \beta_\pi = 0$, from the cross-section

$$R = \frac{\sigma(x_\gamma)}{\sigma_{\alpha_\pi=0}(x_\gamma)} = \frac{N_{meas}(x_\gamma)}{N_{sim}(x_\gamma)} = 1 - \frac{3}{2} \cdot \frac{m_\pi^3}{\alpha} \cdot \frac{x_\gamma^2}{1-x_\gamma} \alpha_\pi$$

is derived, depending on $x_\gamma = E_{\gamma(lab)}/E_{Beam}$.

Measuring R the polarisability α_π can be concluded.

- Control systematics by



and





Extraction of the pion polarisability

- Identify **exclusive reactions**

$$\pi\gamma_{\{Ni \rightarrow Ni'\}} \rightarrow \pi\gamma$$

at smallest momentum transfer $< 0.001 \text{ GeV}^2/c^2$

- Assuming $\alpha_\pi + \beta_\pi = 0$, from the cross-section

$$R = \frac{\sigma(x_\gamma)}{\sigma_{\alpha_\pi=0}(x_\gamma)} = \frac{N_{meas}(x_\gamma)}{N_{sim}(x_\gamma)} = 1 - \frac{3}{2} \cdot \frac{m_\pi^3}{\alpha} \cdot \frac{x_\gamma^2}{1-x_\gamma} \alpha_\pi$$

is derived, depending on $x_\gamma = E_{\gamma(lab)}/E_{Beam}$.

Measuring R the polarisability α_π can be concluded.

- Control systematics by

$$\mu\gamma_{\{Ni \rightarrow Ni'\}} \rightarrow \mu\gamma$$

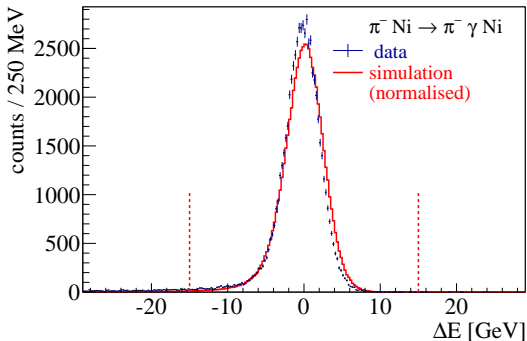
and

$$K^- \rightarrow \pi^- \pi^0 \rightarrow \pi\gamma\gamma$$



Identifying the $\pi\gamma \rightarrow \pi\gamma$ reaction

Phys. Rev. Lett. 114, 062002 (2015)

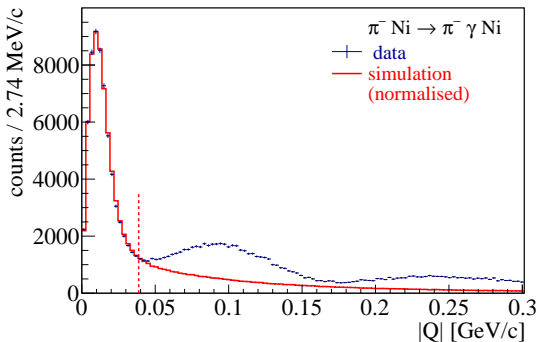


- Energy balance $\Delta E = E_\pi + E_\gamma - E_{\text{Beam}}$
- Exclusivity peak $\sigma \approx 2.6 \text{ GeV}$ (1.4%)
- ~ 63.000 exclusive events ($x_\gamma > 0.4$) (Serpukhov ~ 7000 for $x_\gamma > 0.5$)



Primakoff peak

Phys. Rev. Lett. 114, 062002 (2015)

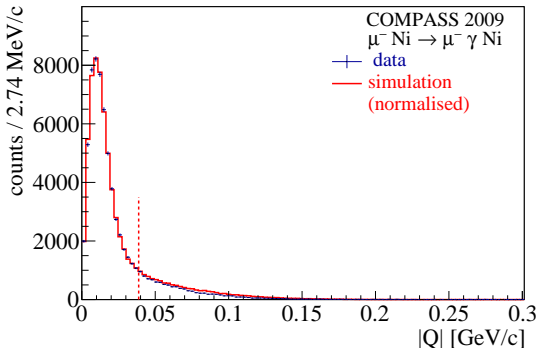


- $\Delta Q_T \approx 12 \text{ MeV}/c$ (190 GeV/c beam \rightarrow requires few- μrad angular resolution)
- first diffractive minimum on Ni nucleus at $Q \approx 190 \text{ MeV}/c$
- data a little more narrow than simulation \rightarrow negative interference?



Primakoff peak: muon data

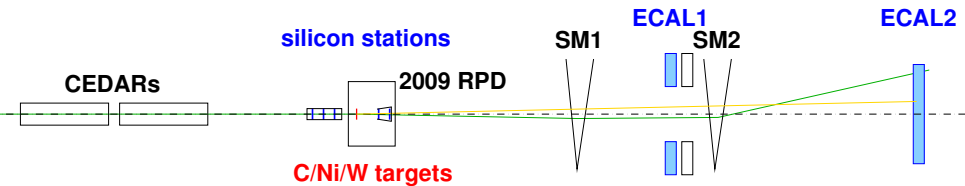
Phys. Rev. Lett. 114, 062002 (2015)



- **muon control measurement:** pure electromagnetic interaction
- e.m. nuclear effects well understood



Principle of the measurement





ECAL2: 3000 cells of different types



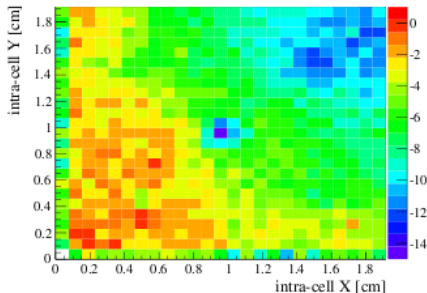


Figure 3.5: Profile of energy deviations shown for 1/4 of a shashlik block and for muon data photons within the range $133 \text{ GeV} < E_\gamma < 152 \text{ GeV}$.

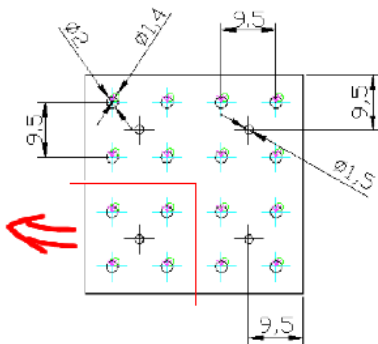
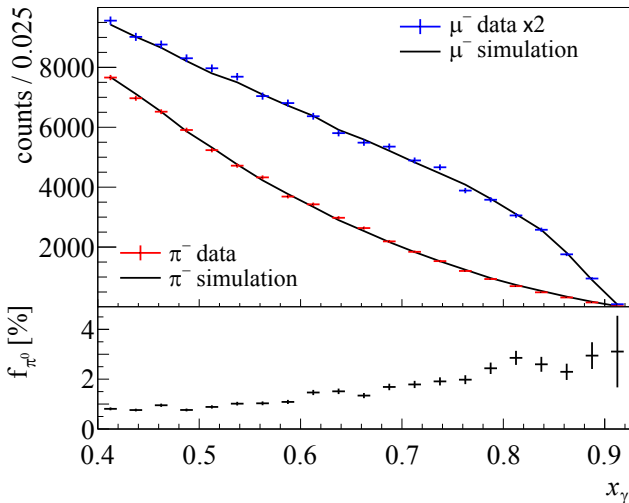


Figure 3.6: Technical drawing of a full shashlik cell to be compared with the figure to the left.

from: Th. Nagel, PhD thesis TUM 2012



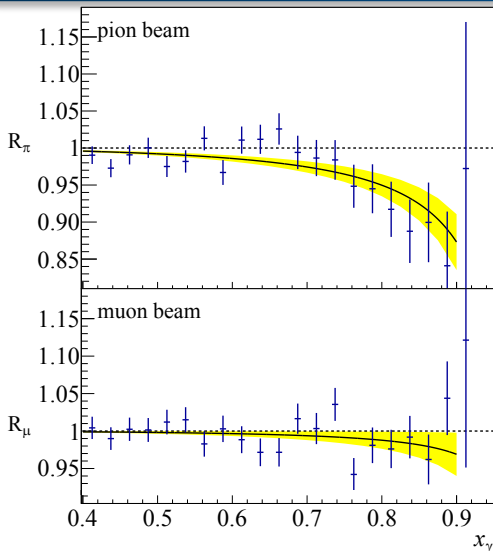
Photon energy spectra for muon and pion beam



Phys. Rev. Lett. 114, 062002 (2015)



Pion polarisability: COMPASS result



$$\alpha_\pi = (2.0 \pm 0.6_{\text{stat}}) \times 10^{-4} \text{ fm}^3$$

(assuming $\alpha_\pi = -\beta_\pi$)

“false polarisability” from muon data:

$$(0.5 \pm 0.5_{\text{stat}}) \times 10^{-4} \text{ fm}^3$$

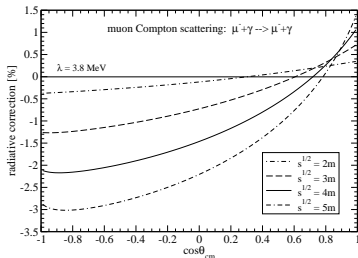
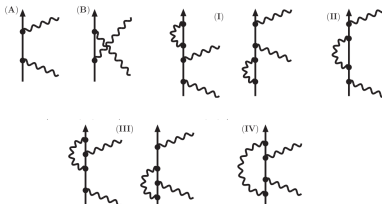
Phys. Rev. Lett. 114, 062002 (2015)



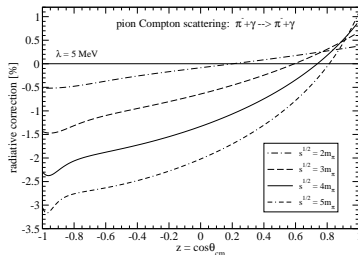
Radiative corrections (Compton scattering part)



Technische Universität München



Nucl.Phys. A837 (2010)



Eur.Phys.J. A39 (2009) 71



source of systematic uncertainty	estimated magnitude	
	CL = 68 %	$[10^{-4} \text{ fm}^3]$
determination of tracking-detector efficiencies		0.5
treatment of radiative corrections		0.3
subtraction of π^0 background		0.2
strong interaction background		0.2
pion-electron elastic scattering		0.2
contribution of muons in the beam		0.05
quadratic sum		0.7



source of systematic uncertainty	estimated magnitude CL = 68 % [10 ⁻⁴ fm ³]
determination of tracking-detector efficiencies	0.5
treatment of radiative corrections	0.3
subtraction of π^0 background	0.2
strong interaction background	0.2
pion-electron elastic scattering	0.2
contribution of muons in the beam	0.05
quadratic sum	0.7

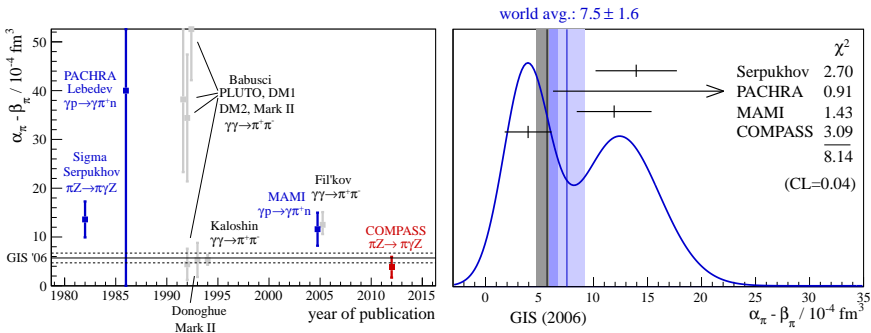
COMPASS result for the pion polarisability:

$$\alpha_\pi = (2.0 \pm 0.6_{\text{stat}} \pm 0.7_{\text{syst}}) \times 10^{-4} \text{ fm}^3$$

with $\alpha_\pi = -\beta_\pi$ assumed



Pion polarisability: world data including COMPASS



- The new COMPASS result is in significant tension with the earlier measurements of the pion polarisability
- The expectation from ChPT is confirmed within the uncertainties

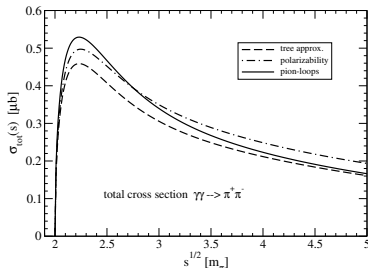


Photon-photon fusion process $\gamma\gamma \rightarrow \pi^+\pi^-$

- Planned measurements at ALICE and JLab

$$\sigma_{tot}(s) = \frac{2\pi\alpha^2}{\hat{s}^3 m_\pi^2} \left\{ [4 + \hat{s} + \hat{s} |C(\hat{s})|^2] \sqrt{\hat{s}(\hat{s} - 4)} + 8[2 - \hat{s} + \hat{s} \operatorname{Re} C(\hat{s})] \ln \frac{\sqrt{\hat{s}} + \sqrt{\hat{s} - 4}}{2} \right\},$$

$$C(\hat{s}) = -\beta_\pi \frac{m_\pi^3}{2\alpha} \hat{s} - \frac{m_\pi^2}{(4\pi f_\pi)^2} \left\{ \frac{\hat{s}}{2} + 2 \left[\ln \frac{\sqrt{\hat{s}} + \sqrt{\hat{s} - 4}}{2} - \frac{i\pi}{2} \right]^2 \right\}$$

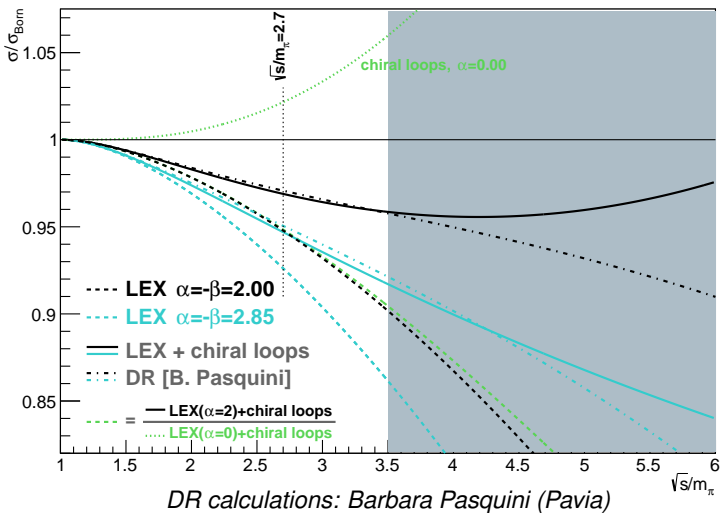


courtesy Norbert Kaiser (TUM)

limited sensitivity to the polarisability contribution



Polarisability and Loop Contributions $z=-1.0$





Pion polarisability on the lattice

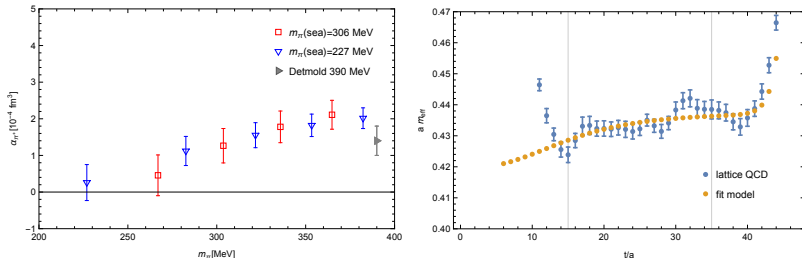
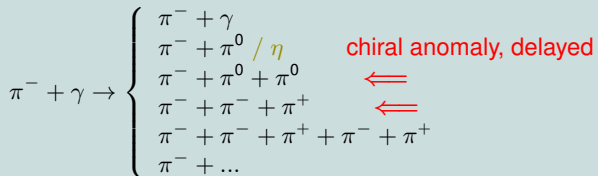


FIGURE 3. Left: electric polarizability for the charged pions as a function of the valence quark mass. The data for $m_{\pi} = 390 \text{ MeV}$ is taken from [5]. Right: effective mass for a charged pion correlator together with the scalar particle correlator determined from the fit. The fitting range is indicated by the vertical bars.

Alexandru *et al.*, Pion electric polarizability from lattice QCD, arXiv:1501.06516



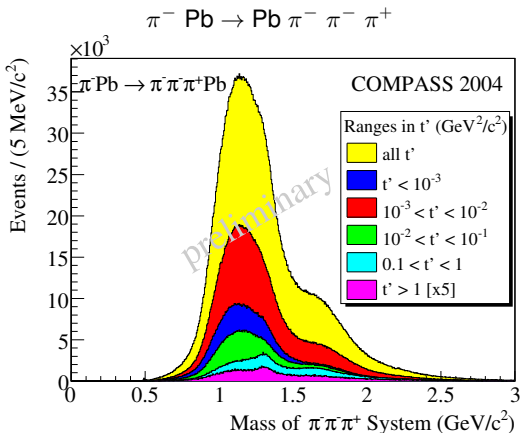
Primakoff reactions accessible at COMPASS



analogous: Kaon-induced reactions $K^- + \gamma \rightarrow \dots$



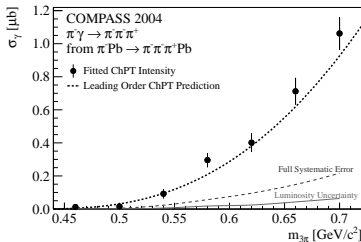
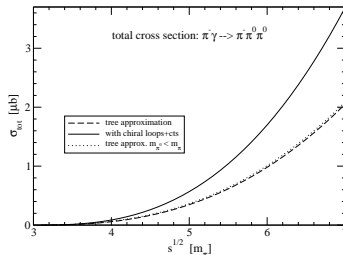
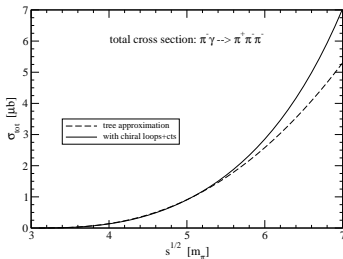
2004 Primakoff results



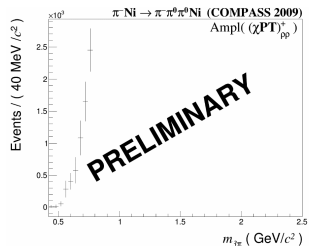
- "Low t' ": $10^{-3} \text{ (GeV/c)}^2 < t' < 10^{-2} \text{ (GeV/c)}^2 \quad \sim 2\,000\,000 \text{ events}$
- "Primakoff region": $t' < 10^{-3} \text{ (GeV/c)}^2 \quad \sim 1\,000\,000 \text{ events}$



Chiral dynamics in $\pi\gamma \rightarrow 3\pi$



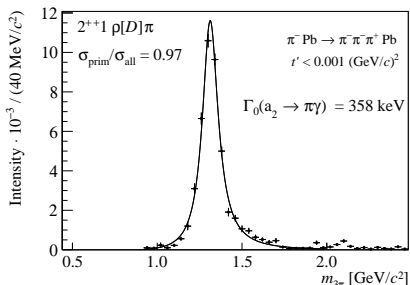
published in PRL 108 (2012) 192001



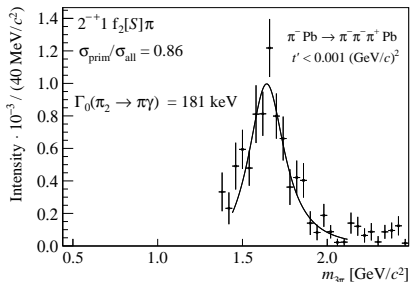
normalization: analysis ongoing



Radiative Coupling of $a_2(1320)$ and $\pi_2(1670)$



$\Gamma_0(a_2(1320) \rightarrow \pi\gamma)$ **M2**



$\Gamma_0(\pi_2(1670) \rightarrow \pi\gamma)$ **E2**

- radiative width for $a_2(1320)$: $358 \pm 6 \pm 42$ keV
larger than PDG value (287 ± 30 keV)
- first measurement for $\pi_2(1670)$

published in EPJ A50 (2014) 79



- COMPASS pursues a broad program in the field of strong interaction: partonic structure, hadron spectroscopy, precision measurements to test EFT
- Measurement of the **pion polarisability** at COMPASS
 - Via the Primakoff reaction, COMPASS has determined

$$\alpha_\pi = (2.0 \pm 0.6_{\text{stat}} \pm 0.7_{\text{syst}}) \times 10^{-4} \text{ fm}^3 \quad \text{assuming } \alpha_\pi + \beta_\pi = 0$$

- most direct access to the $\pi\gamma \rightarrow \pi\gamma$ process
 - Most precise experimental determination
 - Systematic control: $\mu\gamma \rightarrow \mu\gamma$, $K^- \rightarrow \pi^- \pi^0$
- COMPASS measures more aspects of chiral dynamics in $\pi^- \gamma \rightarrow \pi^- \pi^0$ and $\pi\gamma \rightarrow \pi\pi\pi$ reactions
- High-statistics run 2012
 - separate determination of α_π and β_π
 - s -dependent quadrupole polarisabilities
 - First measurement of the kaon polarisability



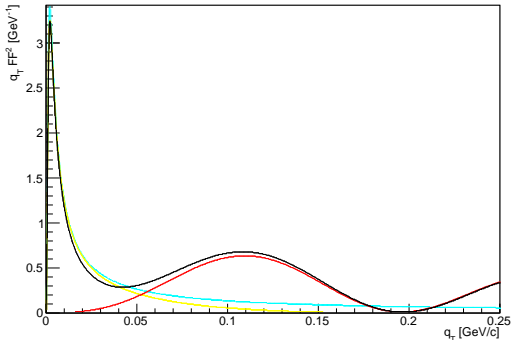
Thank you for your attention!



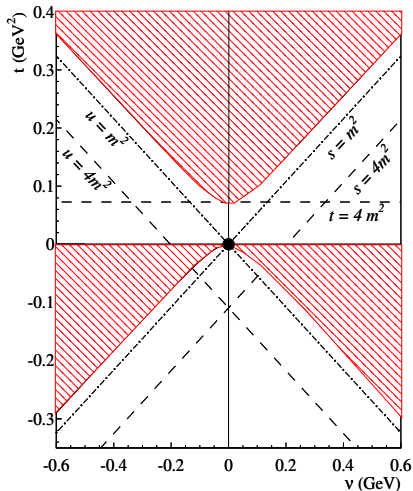




Photon density squared form factor

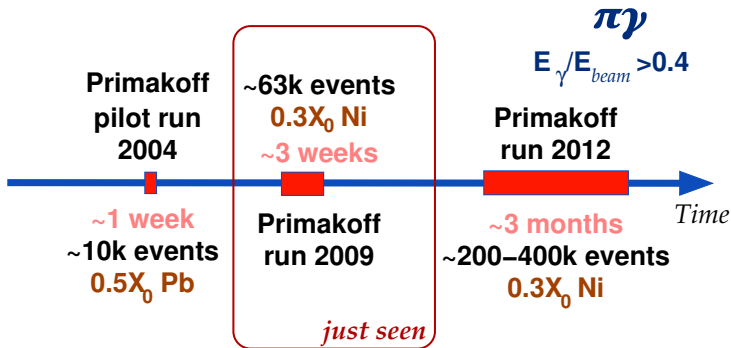


- calculation following G. Fäldt (Phys. Rev. C79, 014607)
- eikonal approximation: pions traverse Coulomb and strong-interaction potentials



- ▶ **red hatched:**
physical regions
 $\gamma + \gamma \rightarrow \pi + \pi$
 $\gamma + \pi \rightarrow \gamma + \pi$
- ▶ two-pion thresholds
at $s = 4m_\pi^2$, $u = 4m_\pi^2$,
 $t = 4m_\pi^2$
- ▶ DR integration paths
 $t = 0$ (forward),
 $\theta = 180^\circ$ (backward)
 $u = m_\pi^2$, $s = m_\pi^2$, ...

from: D. Drechsel, talk at IWHSS 2011 Paris

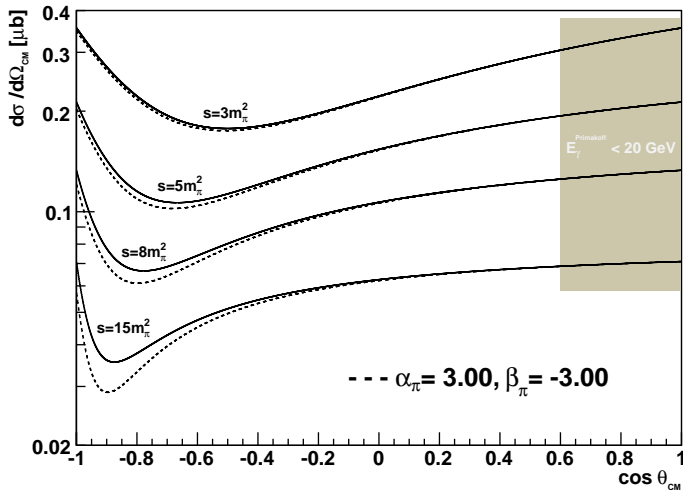


$$\mathcal{P} = z_-^2 (\alpha_{\pi} - \beta_{\pi}) + \frac{s^2}{m_{\pi}^4} z_+^2 (\alpha_{\pi} + \beta_{\pi}) - \frac{(s - m_{\pi}^2)^2}{24s} z_-^3 (\alpha_2 - \beta_2)$$

$$z_{\pm} = 1 \pm \cos \theta_{cm}$$

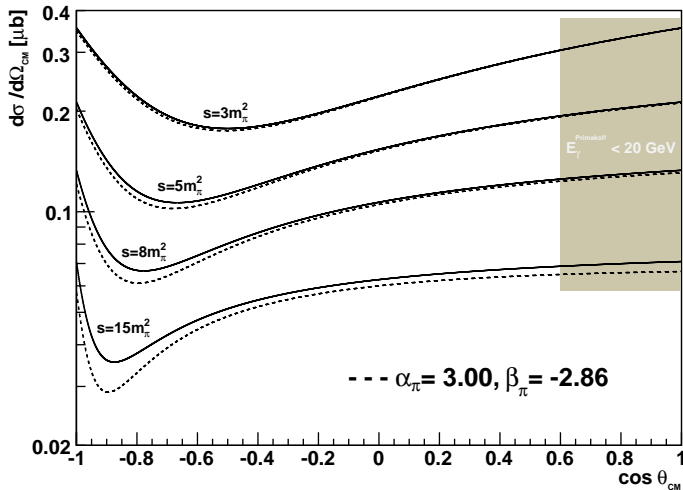


loop effects not shown



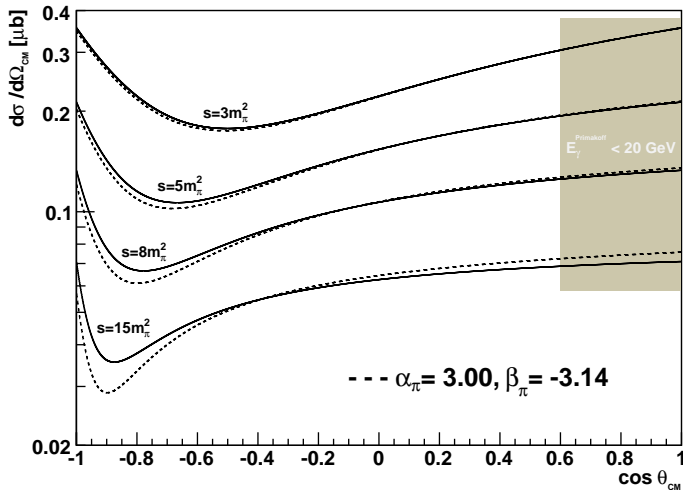


loop effects not shown



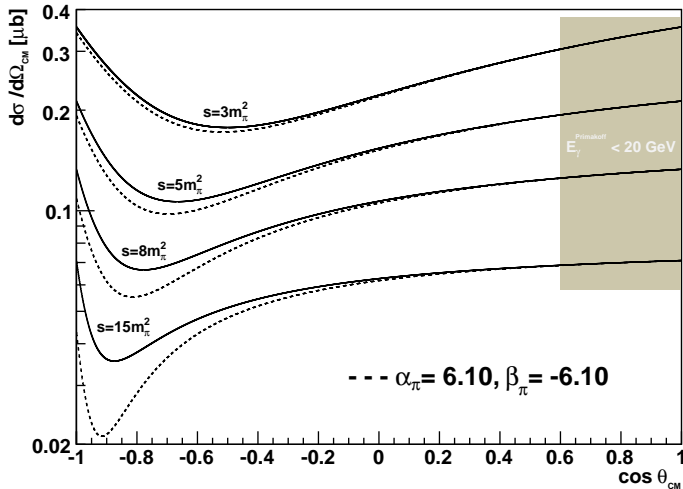


loop effects not shown





loop effects not shown





- Radiative π^+ production on the proton:

$$\gamma \pi^* \longrightarrow \pi \gamma \quad [\text{via } \gamma p \rightarrow n \pi^+ \gamma]$$

Mainz (2005) measurement: $\alpha_\pi - \beta_\pi = 11.6 \pm 1.5 \pm 3.0 \pm 0.5$

“ ± 0.5 ”: model error *only within the used ansatz*,

full systematics not under control

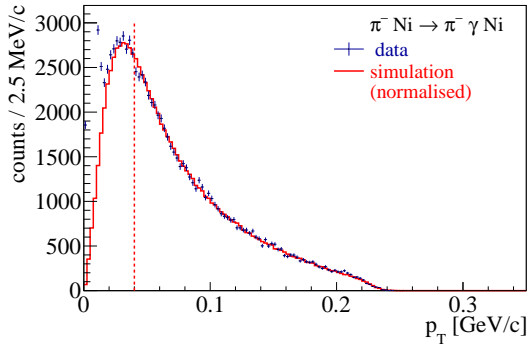
- Primakoff Compton reaction:

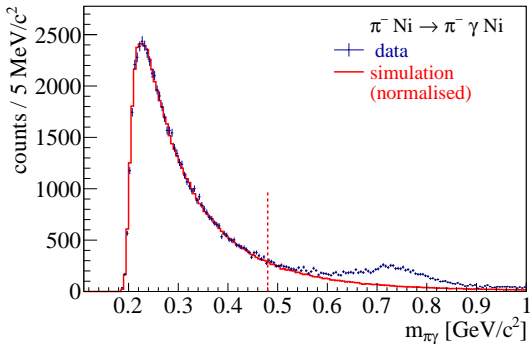
$$\gamma^* \pi \longrightarrow \pi \gamma \quad [\text{via } \pi Z \rightarrow Z \pi \gamma]$$

tiny extrapolation $\gamma^* \rightarrow \gamma$ $\mathcal{O}(10^{-3} m_\pi^2)$

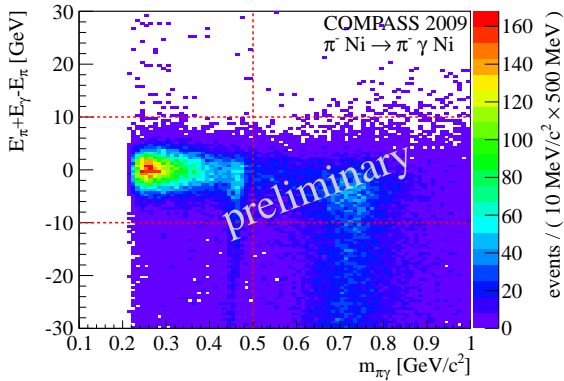
fully under theoretical control

[N. Kaiser, J.F., Nucl. Phys. A 812 (2008) 186]





- ρ contribution from $\pi\gamma \rightarrow \pi\pi^0$

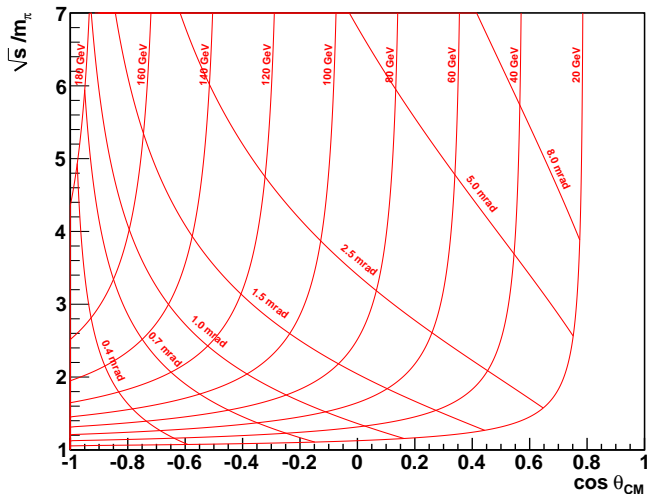


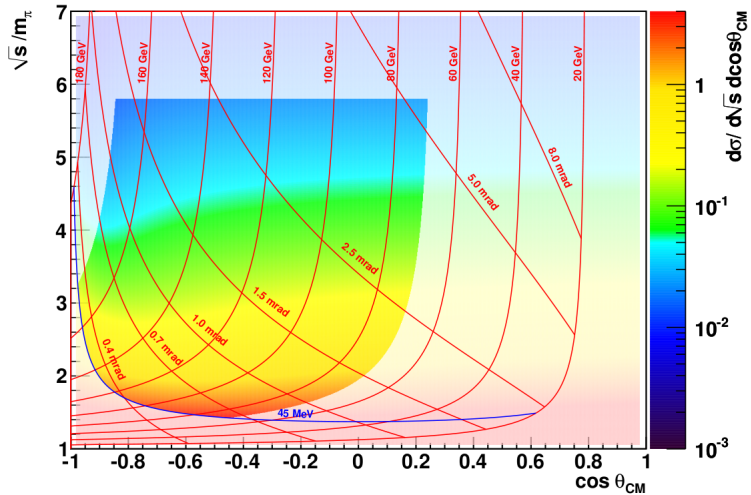
- ρ contribution from $\pi\gamma \rightarrow \pi\pi^0$



Mandelstam $\{s, t\} \leftrightarrow$ Laboratory $\{E_\gamma, \theta_\gamma\}$

for $\pi\gamma \rightarrow \pi\gamma$

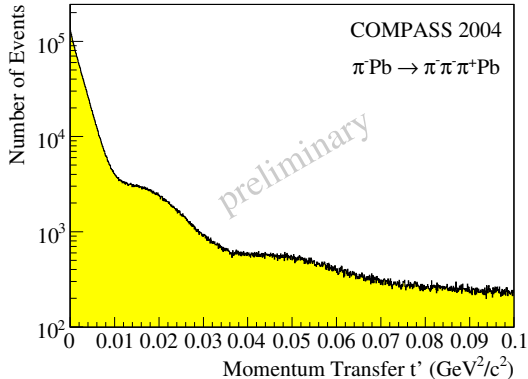




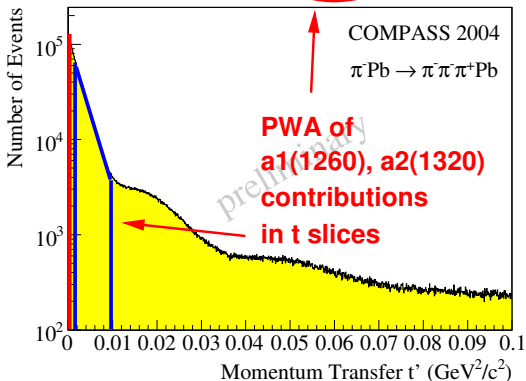


M.R. Pennington in the 2nd DAΦNE Physics Handbook,
“What we learn by measuring $\gamma\gamma \rightarrow \pi\pi$ at DAΦNE”:

All this means that the only way to measure the pion polarisabilities is in the Compton scattering process near threshold and not in $\gamma\gamma \rightarrow \pi\pi$. Though the low energy $\gamma\gamma \rightarrow \pi\pi$ scattering is seemingly close to the Compton threshold (...) and so the *extrapolation* not very far, the dominance of the pion pole (...) means that the energy scale for this continuation is m_π . Thus the polarisabilities cannot be determined accurately from $\gamma\gamma$ experiments in a model-independent way and must be measured in the Compton scattering region.



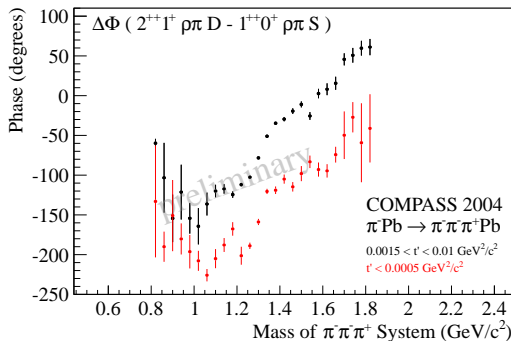
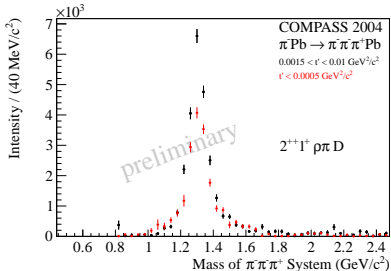
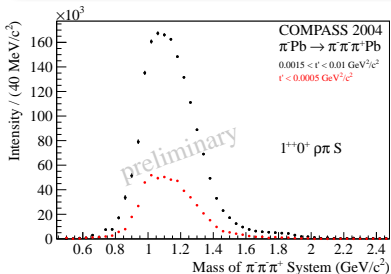
- "Low t' ": $10^{-3} (\text{GeV}/c)^2 < t' < 10^{-2} (\text{GeV}/c)^2$ $\sim 2\,000\,000$ events
- "Primakoff region": $t' < 10^{-3} (\text{GeV}/c)^2$ $\sim 1\,000\,000$ events

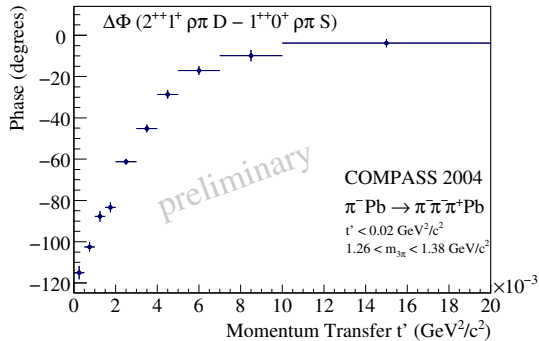
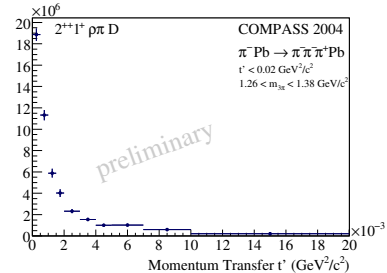
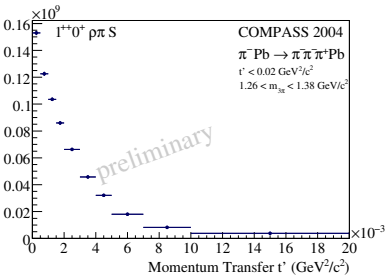


- "Low t' ": $10^{-3} (\text{GeV}/c)^2 < t' < 10^{-2} (\text{GeV}/c)^2$ $\sim 2\,000\,000$ events
- "Primakoff region": $t' < 10^{-3} (\text{GeV}/c)^2$ $\sim 1\,000\,000$ events



PWA: a_1 , a_2 and $\Delta\Phi$ in separated t' regions

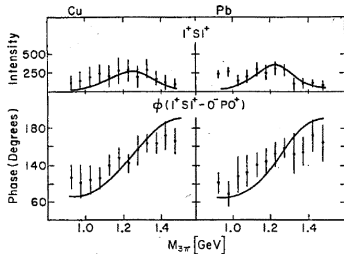
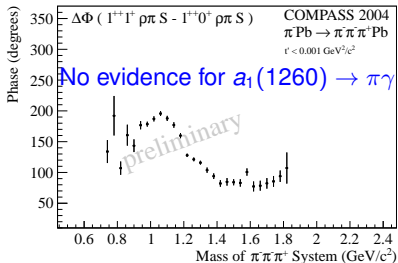
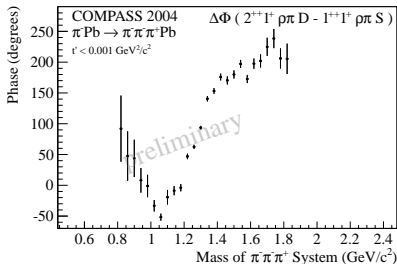
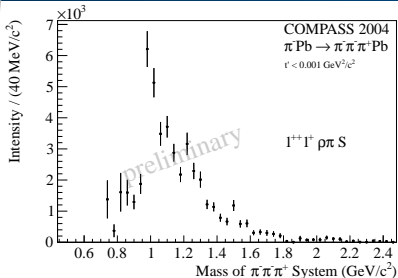




- transition of $\pi\gamma$ to $\pi IP \rightarrow a_2$ production
- work in progress
- interference can be used to map details of resonances and production mechanisms



Primakoff production of $a_1(1260)$ vs. E272 result



M. Zielinski et al, Phys. Rev. Lett 52 (1984) 1195



- **Mass-independent PWA** (narrow mass bins):

$$\sigma_{\text{indep}}(\tau, m, t') = \sum_{\epsilon=\pm 1} \sum_{r=1}^{N_r} \left| \sum_i T_{ir}^\epsilon f_i^\epsilon(t') \psi_i^\epsilon(\tau, m) / \sqrt{\int |f_i^\epsilon(t')|^2 dt'} \sqrt{\int |\psi_i^\epsilon(\tau', m)|^2 d\tau'} \right|^2$$

- Production strength assumed constant in single bins
- Decay amplitudes $\psi_i^\epsilon(\tau, m)$, with t' dependence $f_i^\epsilon(t')$
- Production amplitudes $T_{ir}^\epsilon \rightarrow$ Extended log-likelihood fit
- Acceptance corrections included
- **Spin-density matrix:** $\rho_{ij}^\epsilon = \sum_r T_{ir}^\epsilon T_{jr}^{\epsilon*}$

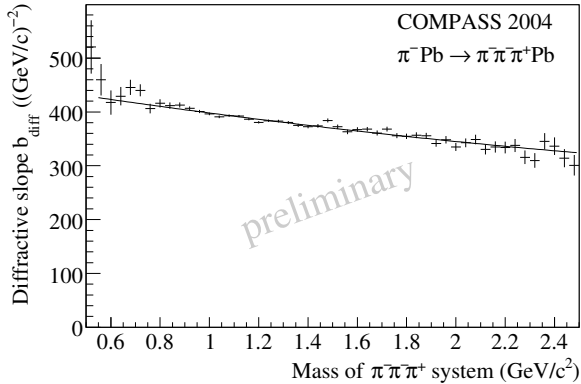
\rightarrow Physical parameters:

$$\text{Intens}_i^\epsilon = \rho_{ii}^\epsilon,$$

relative phase Φ_{ij}^e

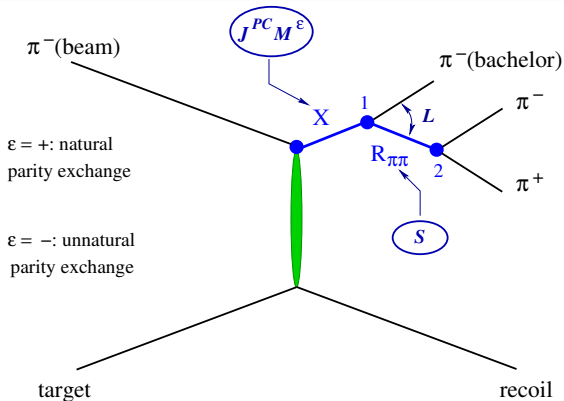
$$\text{Coh}_{i,j}^\epsilon = \sqrt{(\text{Re } \rho_{ij}^\epsilon)^2 + (\text{Im } \rho_{ij}^\epsilon)^2} / \sqrt{\rho_{ii}^\epsilon \rho_{jj}^\epsilon}$$

- **Mass-dependent χ^2 -fit** (not presented here):
 - X parameterized by Breit-Wigner (BW) functions
 - Background can be added





Isobar Model

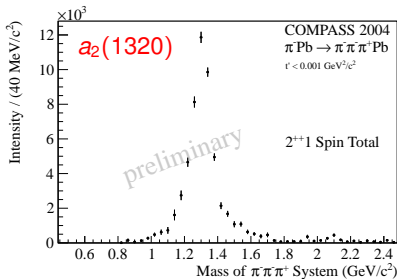
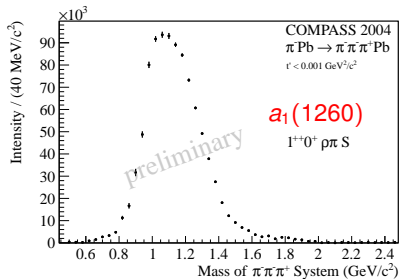
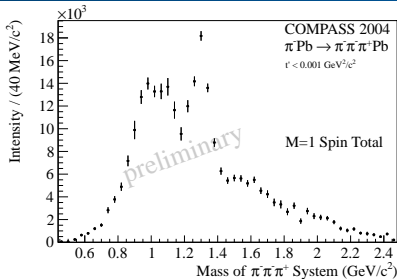
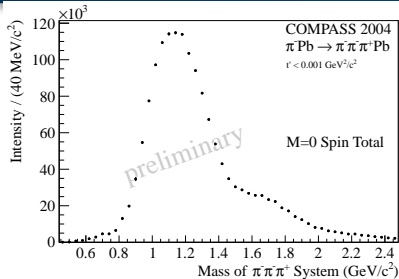


- Isobar model: Intermediate 2-particle decays
- Partial wave in reflectivity basis: $J^{PC} M^\epsilon [isobar] L$

- **Mass-independent** PWA ($40 \text{ MeV}/c^2$ mass bins): **38 waves**
Fit of angular dependence of partial waves, interferences
- **Mass-dependent** χ^2 -fit (Not presented here)



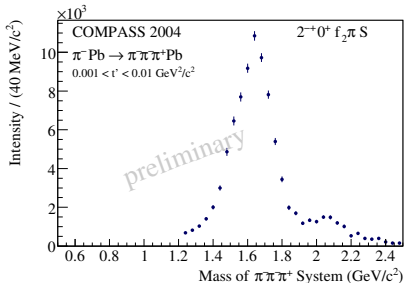
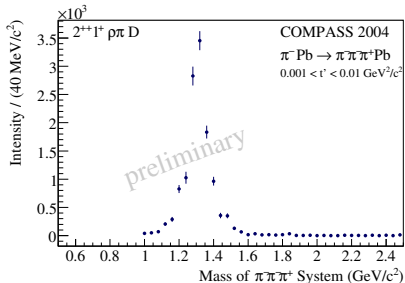
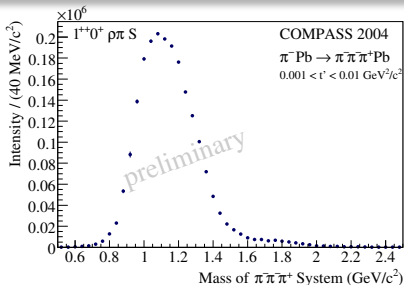
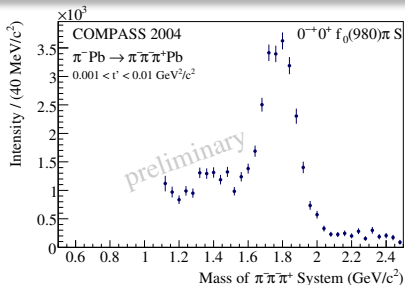
Major intensities in $m(3\pi)$ -bins (acceptance corrected)





PWA of data with low t'

Intensity of selected waves: $0^{-+}0^{+}f_0(980)\pi S$, $1^{++}0^{+}\rho\pi S$, $2^{++}1^{+}\rho\pi D$, $2^{-+}0^{+}f_2(1270)\pi S$





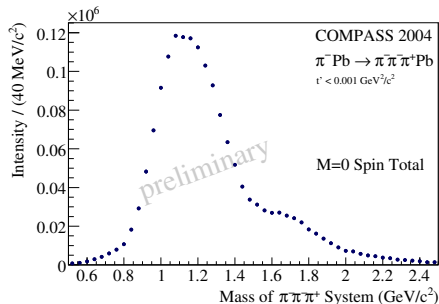
“Spin Totals”: Sum of all contributions for given M (i.e. z-projection of J)

t' -dependent amplitudes:

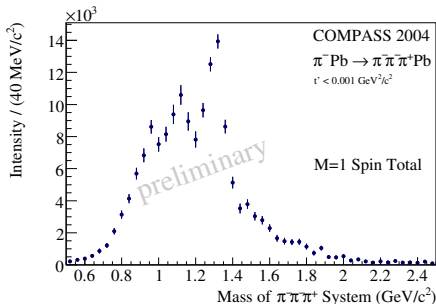
Primakoff production: **M=1**: $\sigma(t') \propto e^{-b_{\text{Prim}}t'}$ → arises at $t' \approx 0$ (resolved shape!)

Diffractive production: **M=0**: $\sigma(t') \propto e^{-b_{\text{diff}}(m)t'}$

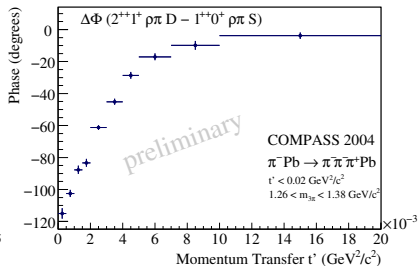
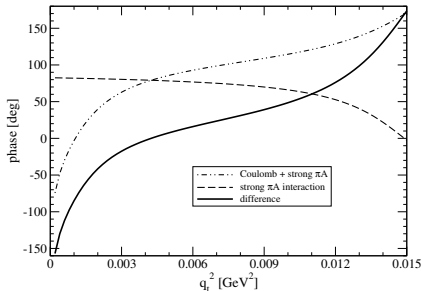
M=1: $\sigma(t') \propto t' e^{-b_{\text{diff}}(m)t'}$ → vanishes for $t' \approx 0$



M=0



M=1



Glauber modell

G. Fäldt and U. Tengblad, Phys. Rev. C79, 014607 (2009)

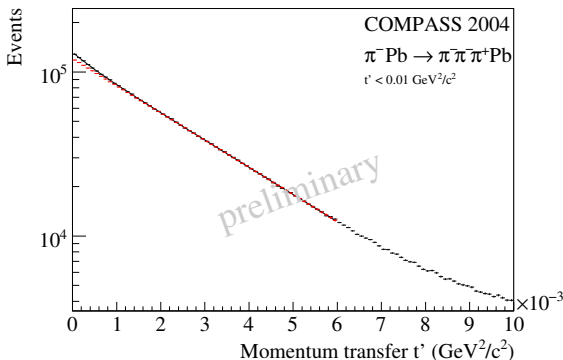
Plot: N. Kaiser (TU München)

- ⇒ indicates confirmation of interference Coulomb-interaction - strong interaction
- ⇒ detailed studies of the nature of resonances



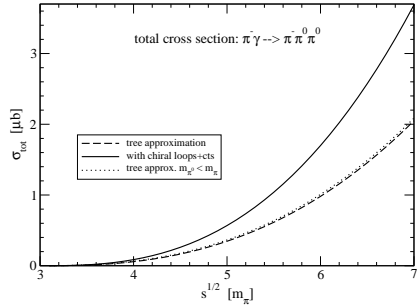
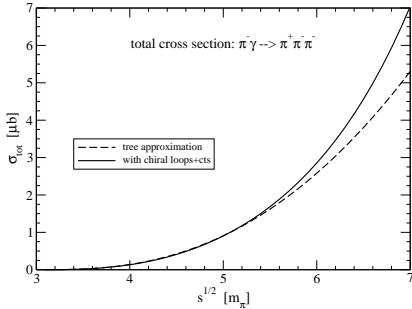
Primakoff: $\sigma(t') \propto e^{-b_{\text{Prim}} t'}$, $b_{\text{Prim}} \approx 2000 \text{ (GeV/c)}^{-2}$ (mainly resolution)

Diffractive: $\sigma(t') \propto e^{-b_{\text{diff}} t'}$, $b_{\text{diff}} \approx 400 \text{ (GeV/c)}^{-2}$ for lead target



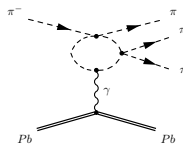
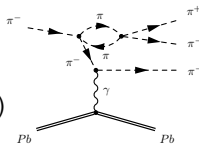
(Mass) spectrum of this Primakoff contribution?

⇒ Statistical subtraction of diffractive background (for bins of $m_{3\pi}$)

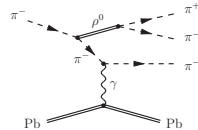


Chiral loops, e.g.

(N. Kaiser, NPA848 (2010) 198)

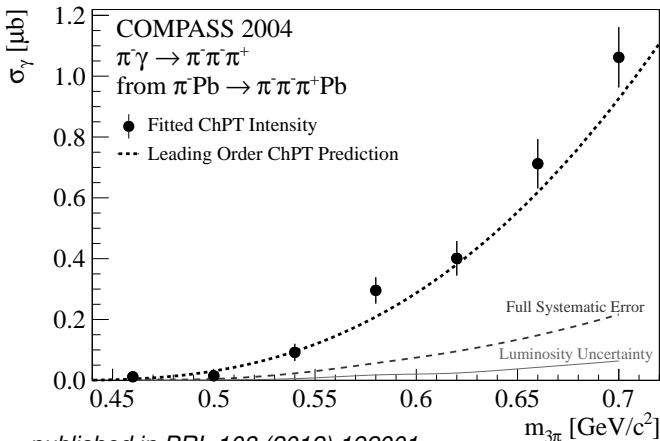
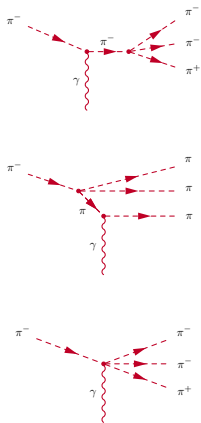


ρ terms:





Measured absolute cross-section of $\pi^-\gamma \rightarrow \pi^-\pi^-\pi^+$



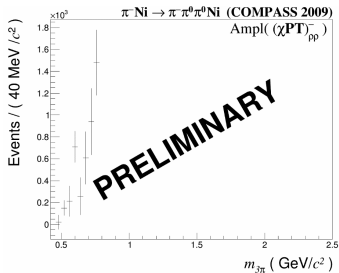
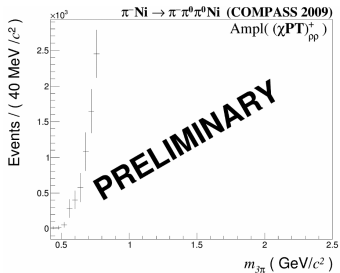
published in *PRL* 108 (2012) 192001



$$\pi^- \gamma \rightarrow \pi^- \pi^0 \pi^0$$

Partial Wave Analysis

Isobaric Model – Chiral Wave





$$\pi^- \gamma \rightarrow \pi^- \pi^0 \pi^0$$

Partial Wave Analysis

Chiral Model - Amplitudes

