

Photons Interacting with Pions at COMPASS

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COMPASS collaboration



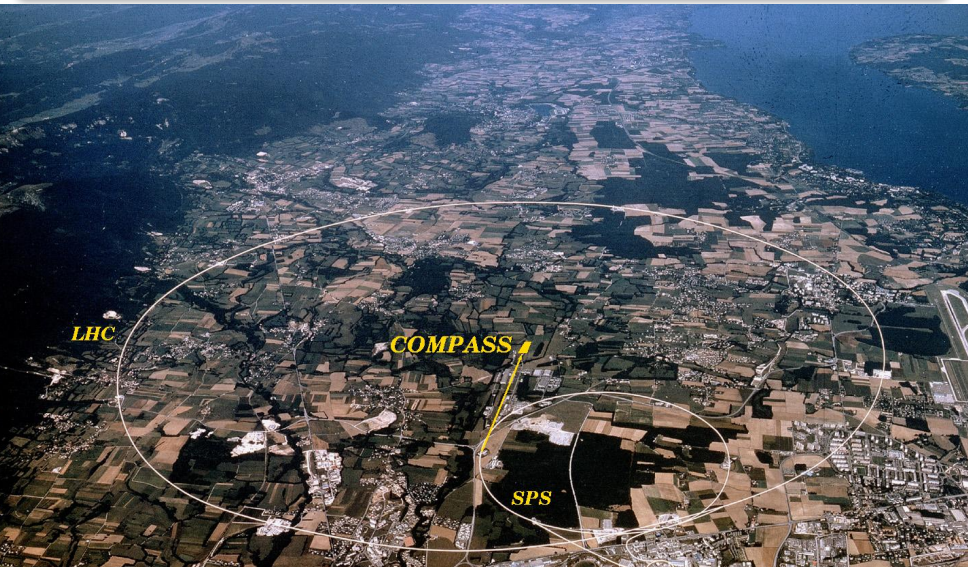
June 18, 2015



Bundesministerium
für Bildung
und Forschung



Common Muon and Proton Apparatus for Structure and Spectroscopy





Common Muon and Proton Apparatus for Structure and Spectroscopy

CERN SPS: protons ~ 400 GeV (5 – 10 sec spills)

- secondary $\pi, K, (\bar{p})$: up to $2 \cdot 10^7/s$ (typ. $5 \cdot 10^6/s$)
Nov. 2004, 2008-09, 2012:
hadron spec. & Primakoff reactions
- tertiary muons: $4 \cdot 10^7 / s$
2002-04, 2006-07, 2010-11: spin structure of the nucleon

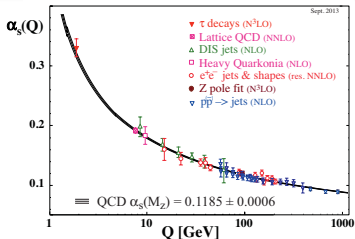
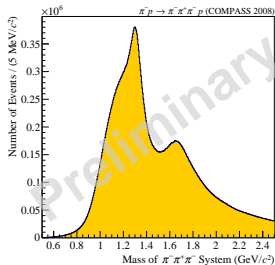
LHC

COMPASS

SPS



- lepton scattering
at high momentum transfer
→ partonic structure of the nucleons

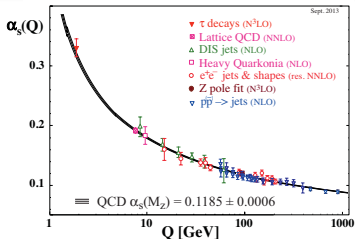
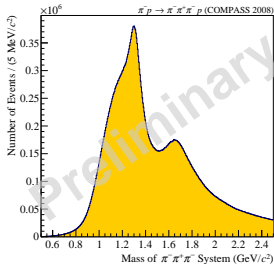


- diffractive dissociation of pions and kaons
→ meson spectrometry

- scattering of pions (and kaons) in nuclear Coulomb field
→ low-energetic meson-photon reactions
 $\pi\gamma \rightarrow \pi\gamma$ (pion polarisability), $\pi\gamma \rightarrow 3\pi$ (chiral dynamics, radiative couplings)



- lepton scattering
at high momentum transfer
→ partonic structure of the nucleons
s. talk tomorrow by A. Ferrero



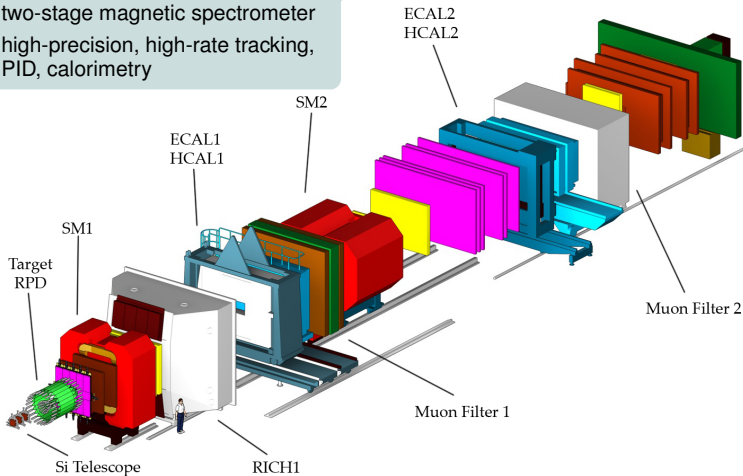
- diffractive dissociation of pions and kaons
→ meson spectrometry
s. following talk by B. Grube

- scattering of pions (and kaons) in nuclear Coulomb field
→ low-energetic meson-photon reactions *(this talk)*
 $\pi\gamma \rightarrow \pi\gamma$ (pion polarisability), $\pi\gamma \rightarrow 3\pi$ (chiral dynamics, radiative couplings)



Fixed-target experiment

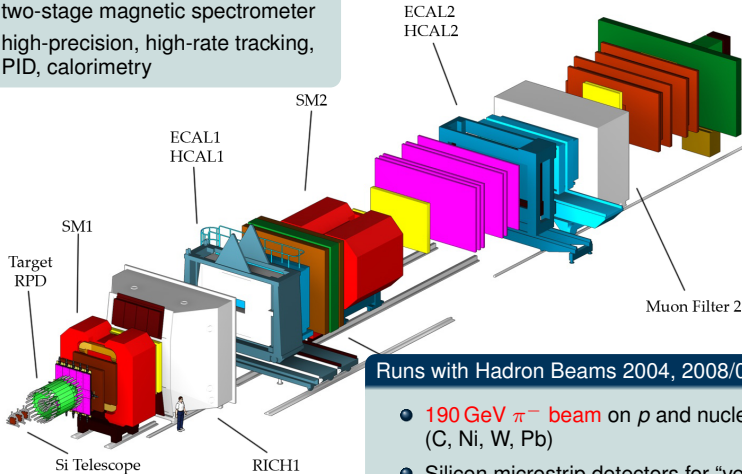
- two-stage magnetic spectrometer
- high-precision, high-rate tracking, PID, calorimetry





Fixed-target experiment

- two-stage magnetic spectrometer
- high-precision, high-rate tracking, PID, calorimetry

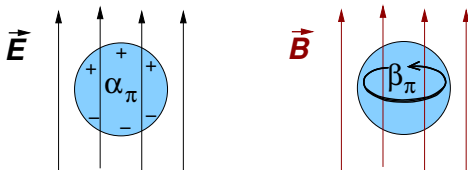


Runs with Hadron Beams 2004, 2008/09, 2012

- **190 GeV π^- beam** on p and nuclear targets (C, Ni, W, Pb)
- Silicon microstrip detectors for “vertexing”
- recoil and (digital) ECAL triggers



Pion polarisability and ChPT



pion polarisabilities α_π, β_π in units of 10^{-4} fm^3

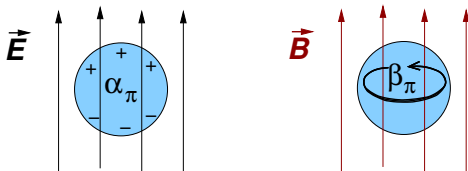
size of the pion $\sim 1 \text{ fm}^3$ [cf. atoms: polarisability \approx size $\approx 1 \text{ \AA}^3$]

Theory: ChPT (2-loop) prediction:

$$\begin{aligned} \alpha_\pi - \beta_\pi &= 5.7 \pm 1.0 \\ \alpha_\pi + \beta_\pi &= 0.16 \pm 0.1 \end{aligned}$$

experiments for $\alpha_\pi - \beta_\pi$ lie in the range $4 \dots 14$

($\alpha_\pi + \beta_\pi = 0$ assumed)



pion polarisabilities α_π, β_π in units of 10^{-4} fm^3

size of the pion $\sim 1 \text{ fm}^3$ [cf. atoms: polarisability \approx size $\approx 1 \text{ \AA}^3$]

Theory: ChPT (2-loop) prediction:

$$\alpha_\pi = 2.93 \pm 0.5$$

$$\beta_\pi = -2.77 \pm 0.5$$

experiments for α_π lie in the range $2 \dots 7$

($\alpha_\pi + \beta_\pi = 0$ assumed)



latest publication on the pion polarisability



Technische Universität München

PRL **114**, 062002 (2015)

PHYSICAL REVIEW LETTERS

week ending
13 FEBRUARY 2015

Measurement of the Charged-Pion Polarizability

C. Adolph,⁸ R. Akhunzyanov,⁷ M. G. Alexeev,²⁷ G. D. Alexeev,⁷ A. Amoroso,^{27,29} V. Andrieux,²² V. Anosov,⁷

... [213 authors]

(COMPASS Collaboration)

(Received 2 June 2014; revised manuscript received 24 December 2014; published 10 February 2015)

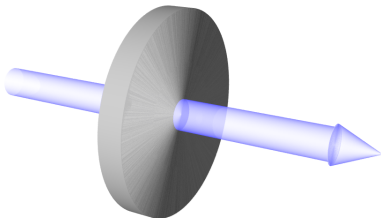
The COMPASS collaboration at CERN has investigated pion Compton scattering, $\pi^- \gamma \rightarrow \pi^- \gamma$, at center-of-mass energy below 3.5 pion masses. The process is embedded in the reaction $\pi^- \text{Ni} \rightarrow \pi^- \gamma \text{Ni}$, which is initiated by 190 GeV pions impinging on a nickel target. The exchange of quasireal photons is selected by isolating the sharp Coulomb peak observed at smallest momentum transfers, $Q^2 < 0.0015 \text{ (GeV}/c^2)^2$. From a sample of 63 000 events, the pion electric polarizability is determined to be $\alpha_\pi = (2.0 \pm 0.6_{\text{stat}} \pm 0.7_{\text{sys}}) \times 10^{-4} \text{ fm}^3$ under the assumption $\alpha_\pi = -\beta_\pi$, which relates the electric and magnetic dipole polarizabilities. It is the most precise measurement of this fundamental low-energy parameter of strong





Principle of the COMPASS measurement

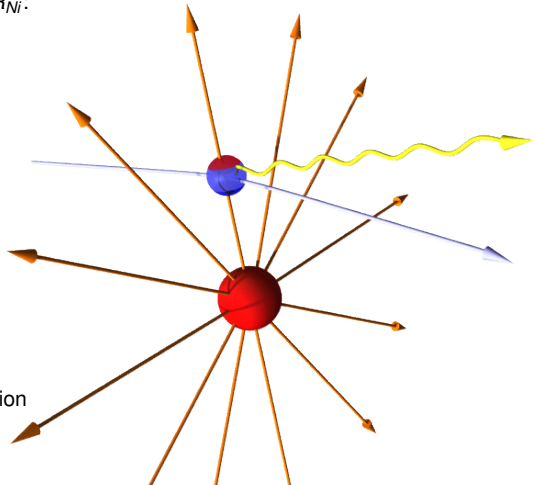
- high-energetic pion beam on 4mm nickel disk
- observe scattered pions in coincidence with produced hard photons
- study of cross-section shape





Polarisability effect in Primakoff technique

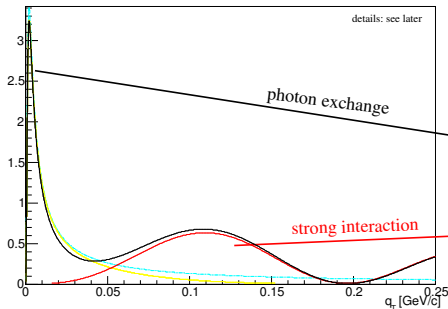
- Charged pions traverse the nuclear **electric field**
 - typical field strength at $d = 5R_N$:
 $E \approx 300 \text{ kV/fm}$
- Bremsstrahlung process:
 - particles scatter off **equivalent photons**
 - tiny momentum transfer
 $Q^2 \approx 10^{-5} \text{ GeV}^2/c^2$
 - pion/muon (quasi)-real Compton scattering
- Polarisability contribution
 - Compton cross-section typically diminished
 - corresponding charge separation
 $\approx 10^{-5} \text{ fm} \cdot e$





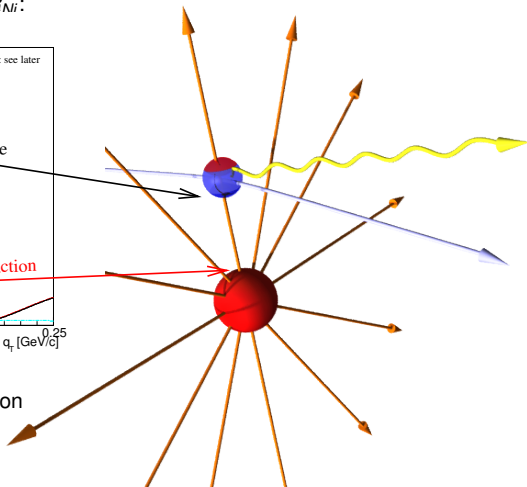
Polarisability effect in Primakoff technique

- Charged pions traverse the nuclear **electric field**
 - typical field strength at $d = 5R_M$:



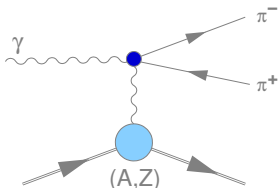
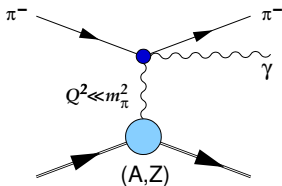
typically diminished

- corresponding charge separation $\approx 10^{-5} \text{ fm} \cdot e$

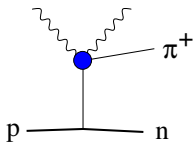




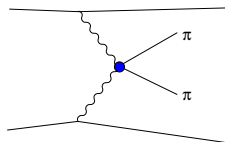
Pion Compton scattering: embedding the process



Primakoff processes



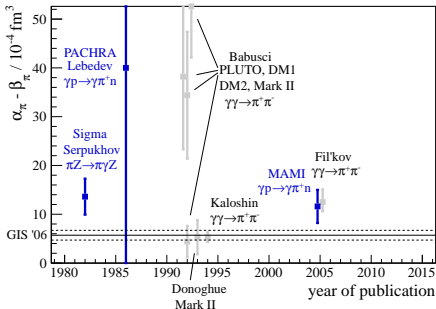
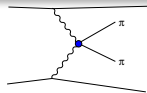
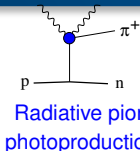
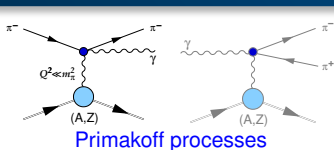
Radiative pion photoproduction



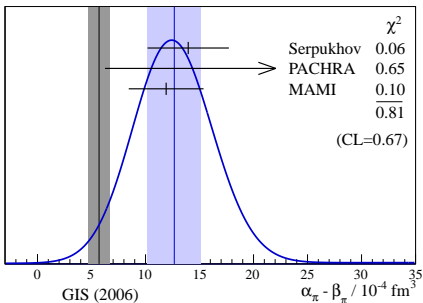
Photon-Photon fusion



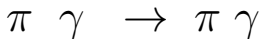
Pion polarisability: world data before COMPASS



world avg.: 12.7 ± 2.5



GIS'06: ChPT prediction, Gasser, Ivanov, Sainio, NPB745 (2006), plots: T. Nagel, PhD
 Fil'kov analysis objected by Pasquini, Drechsel, Scherer PRC81, 029802 (2010)

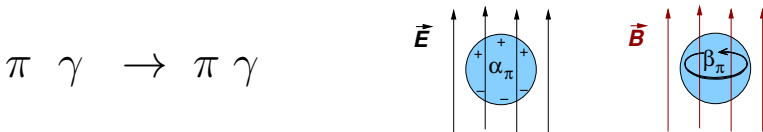


- Two kinematic variables, in CM: total energy \sqrt{s} , scattering angle θ_{cm}

$$\frac{d\sigma_{\pi\gamma}}{d\Omega_{cm}} = \frac{\alpha^2 (s^2 z_+^2 + m_\pi^4 z_-^2)}{s (s z_+ + m_\pi^2 z_-)^2} - \frac{\alpha m_\pi^3 (s - m_\pi^2)^2}{4s^2 (s z_+ + m_\pi^2 z_-)} \cdot \mathcal{P}$$

$$\mathcal{P} = z_-^2 (\alpha_\pi - \beta_\pi) + \frac{s^2}{m_\pi^4} z_+^2 (\alpha_\pi + \beta_\pi) - \frac{(s - m_\pi^2)^2}{24s} z_-^3 (\alpha_2 - \beta_2)$$

$$z_\pm = 1 \pm \cos \theta_{cm}$$



- Two kinematic variables, in CM: total energy \sqrt{s} , scattering angle θ_{cm}

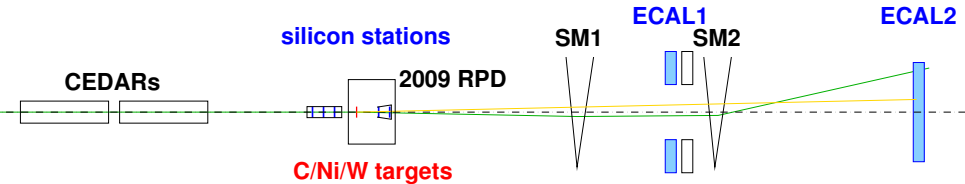
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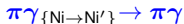
Principle of the measurement





Extraction of the pion polarisability

- Identify exclusive reactions



at smallest momentum transfer $< 0.001 \text{ GeV}^2/c^2$

- Assuming $\alpha_\pi + \beta_\pi = 0$, from the cross-section

$$R = \frac{\sigma(x_\gamma)}{\sigma_{\alpha_\pi=0}(x_\gamma)} = \frac{N_{meas}(x_\gamma)}{N_{sim}(x_\gamma)} = 1 - \frac{3}{2} \cdot \frac{m_\pi^3}{\alpha} \cdot \frac{x_\gamma^2}{1-x_\gamma} \alpha_\pi$$

is derived, depending on $x_\gamma = E_{\gamma(lab)}/E_{Beam}$.

Measuring R the polarisability α_π can be concluded.

- Control systematics by



and





Extraction of the pion polarisability

- Identify **exclusive reactions**

$$\pi\gamma_{\{Ni \rightarrow Ni'\}} \rightarrow \pi\gamma$$

at smallest momentum transfer $< 0.001 \text{ GeV}^2/c^2$

- Assuming $\alpha_\pi + \beta_\pi = 0$, from the cross-section

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- Control systematics by

$$\mu\gamma_{\{Ni \rightarrow Ni'\}} \rightarrow \mu\gamma$$

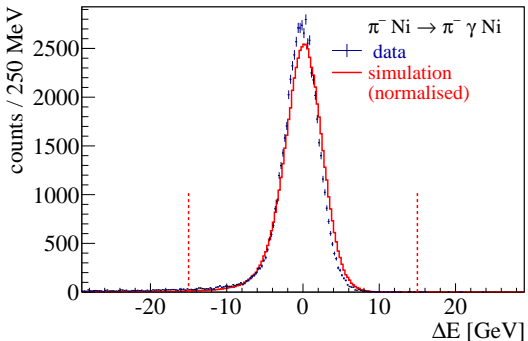
and

$$K^- \rightarrow \pi^- \pi^0 \rightarrow \pi\gamma\gamma$$



Identifying the $\pi\gamma \rightarrow \pi\gamma$ reaction

Phys. Rev. Lett. 114, 062002 (2015)

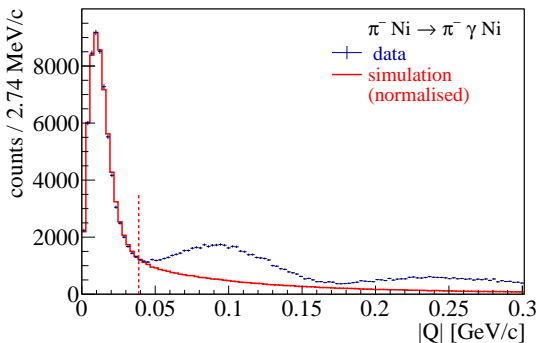


- Energy balance $\Delta E = E_\pi + E_\gamma - E_{\text{Beam}}$
- Exclusivity peak $\sigma \approx 2.6 \text{ GeV}$ (1.4%)
- ~ 63.000 exclusive events ($x_\gamma > 0.4$) (Serpukhov ~ 7000 for $x_\gamma > 0.5$)



Primakoff peak

Phys. Rev. Lett. 114, 062002 (2015)

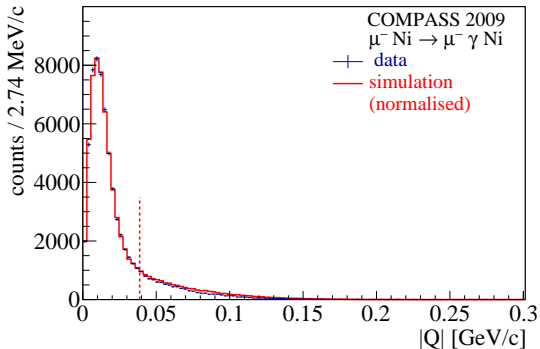


- $\Delta Q_T \approx 12 \text{ MeV}/c$ (190 GeV/c beam \rightarrow requires few- μrad angular resolution)
- first diffractive minimum on Ni nucleus at $Q \approx 190 \text{ MeV}/c$
- data a little more narrow than simulation \rightarrow negative interference?



Primakoff peak: muon data

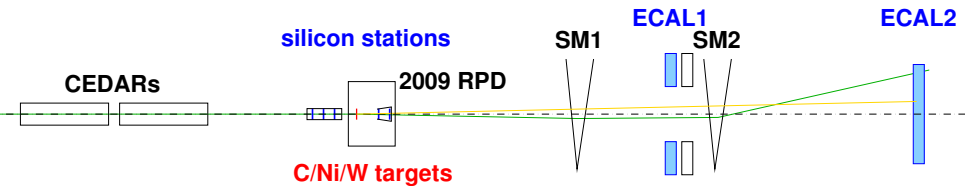
Phys. Rev. Lett. 114, 062002 (2015)



- **muon control measurement:** pure electromagnetic interaction
- e.m. nuclear effects well understood



Principle of the measurement





ECAL2: 3000 cells of different types



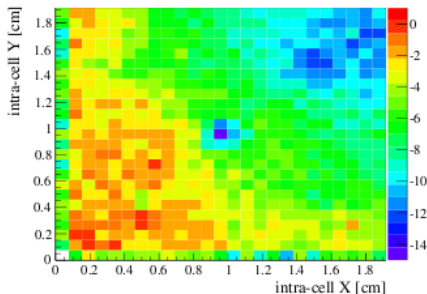


Figure 3.5: Profile of energy deviations shown for 1/4 of a shashlik block and for muon data photons within the range $133 \text{ GeV} < E_\gamma < 152 \text{ GeV}$.

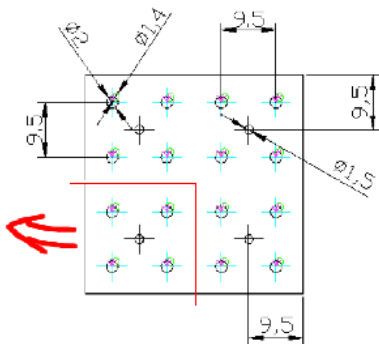
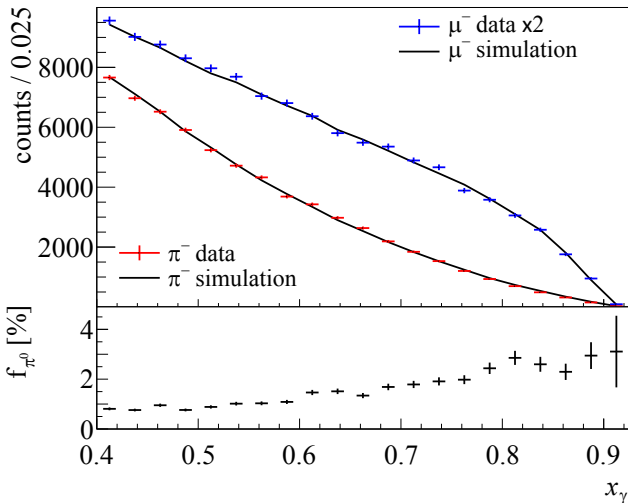


Figure 3.6: Technical drawing of a full shashlik cell to be compared with the figure to the left.

from: Th. Nagel, PhD thesis TUM 2012



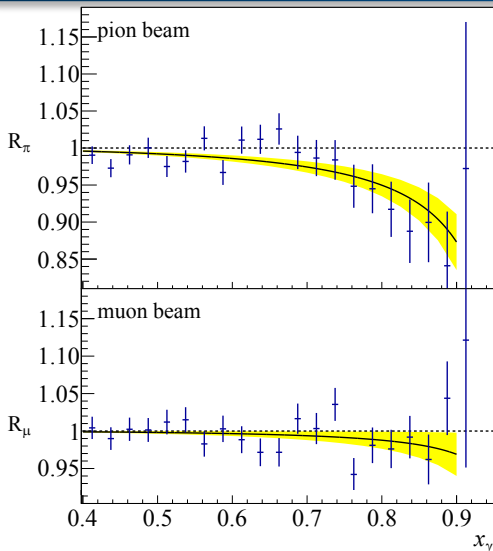
Photon energy spectra for muon and pion beam



Phys. Rev. Lett. 114, 062002 (2015)



Pion polarisability: COMPASS result



$$\alpha_\pi = (2.0 \pm 0.6_{\text{stat}}) \times 10^{-4} \text{ fm}^3$$

(assuming $\alpha_\pi = -\beta_\pi$)

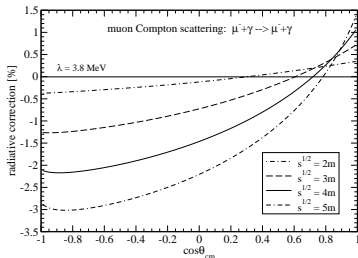
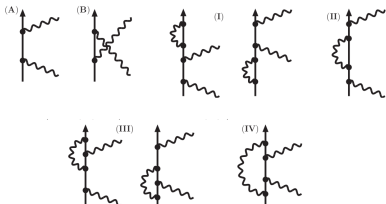
“false polarisability” from muon data:

$$(0.5 \pm 0.5_{\text{stat}}) \times 10^{-4} \text{ fm}^3$$

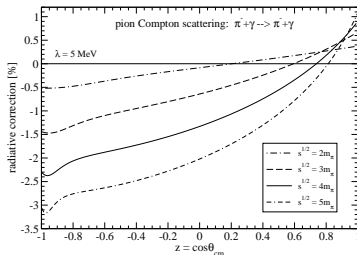
Phys. Rev. Lett. 114, 062002 (2015)



Radiative corrections (Compton scattering part)



Nucl.Phys. A837 (2010)



Eur.Phys.J. A39 (2009) 71



| source of systematic uncertainty | estimated magnitude | |
|---|---------------------|--------------------------|
| | CL = 68 % | $[10^{-4} \text{ fm}^3]$ |
| determination of tracking-detector efficiencies | | 0.5 |
| treatment of radiative corrections | | 0.3 |
| subtraction of π^0 background | | 0.2 |
| strong interaction background | | 0.2 |
| pion-electron elastic scattering | | 0.2 |
| contribution of muons in the beam | | 0.05 |
| quadratic sum | | 0.7 |



| source of systematic uncertainty | estimated magnitude CL = 68 % [10 ⁻⁴ fm ³] |
|---|--|
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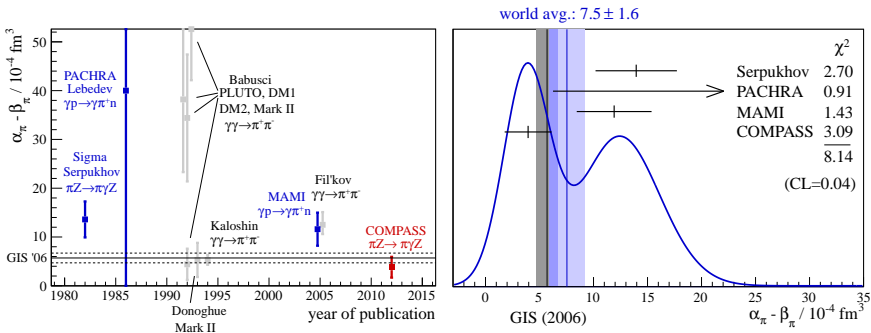
COMPASS result for the pion polarisability:

$$\alpha_\pi = (2.0 \pm 0.6_{\text{stat}} \pm 0.7_{\text{syst}}) \times 10^{-4} \text{ fm}^3$$

with $\alpha_\pi = -\beta_\pi$ assumed



Pion polarisability: world data including COMPASS



- The new COMPASS result is in significant tension with the earlier measurements of the pion polarisability
- The expectation from ChPT is confirmed within the uncertainties

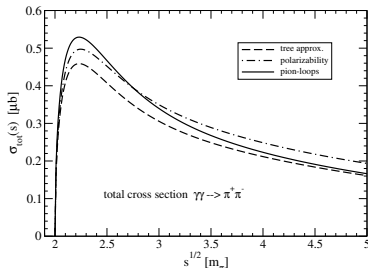


Photon-photon fusion process $\gamma\gamma \rightarrow \pi^+\pi^-$

Planned measurements at ALICE and JLab

$$\sigma_{tot}(s) = \frac{2\pi\alpha^2}{\hat{s}^3 m_\pi^2} \left\{ [4 + \hat{s} + \hat{s} |C(\hat{s})|^2] \sqrt{\hat{s}(\hat{s} - 4)} + 8[2 - \hat{s} + \hat{s} \operatorname{Re} C(\hat{s})] \ln \frac{\sqrt{\hat{s}} + \sqrt{\hat{s} - 4}}{2} \right\},$$

$$C(\hat{s}) = -\beta_\pi \frac{m_\pi^3}{2\alpha} \hat{s} - \frac{m_\pi^2}{(4\pi f_\pi)^2} \left\{ \frac{\hat{s}}{2} + 2 \left[\ln \frac{\sqrt{\hat{s}} + \sqrt{\hat{s} - 4}}{2} - \frac{i\pi}{2} \right]^2 \right\}$$

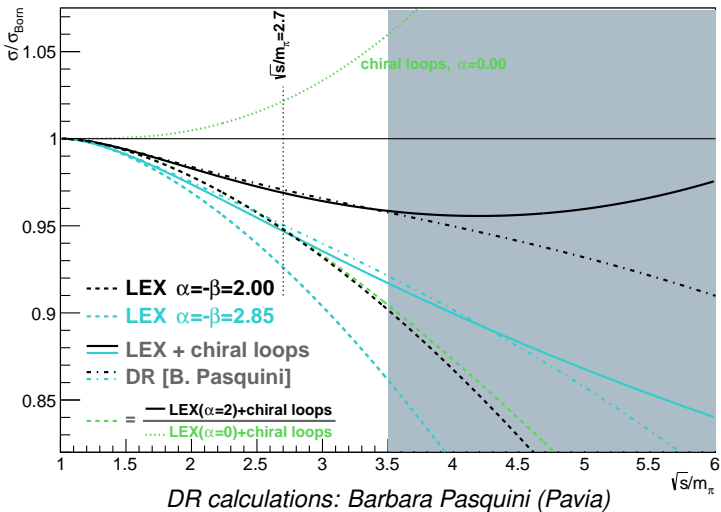


courtesy Norbert Kaiser (TUM)

limited sensitivity to the polarisability contribution



Polarisability and Loop Contributions $z=-1.0$





Pion polarisability on the lattice

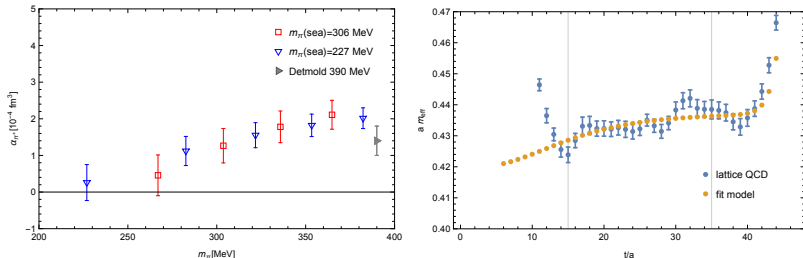


FIGURE 3. Left: electric polarisability for the charged pions as a function of the valence quark mass. The data for $m_{\pi} = 390 \text{ MeV}$ is taken from [5]. Right: effective mass for a charged pion correlator together with the scalar particle correlator determined from the fit. The fitting range is indicated by the vertical bars.

Alexandru *et al.*, Pion electric polarisability from lattice QCD, arXiv:1501.06516



Primakoff reactions accessible at COMPASS

Access to $\pi + \gamma$ reactions via the **Primakoff effect**:

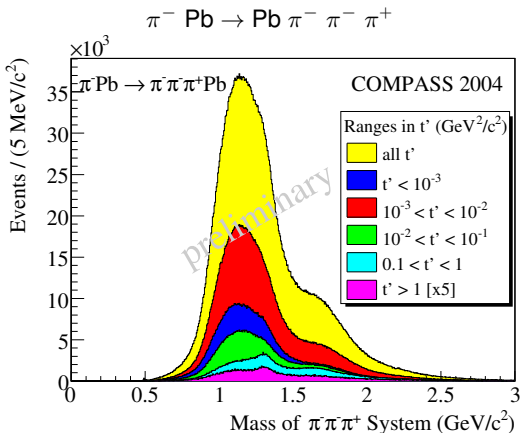
At smallest momentum transfers to the nucleus, high-energetic particles scatter predominantly off the **electromagnetic field** quanta ($\sim Z^2$)

$$\pi^- + \gamma \rightarrow \left\{ \begin{array}{l} \pi^- + \gamma \\ \pi^- + \pi^0 / \eta \\ \pi^- + \pi^0 + \pi^0 \\ \pi^- + \pi^- + \pi^+ \quad \leftarrow \leftarrow \\ \pi^- + \pi^- + \pi^+ + \pi^- + \pi^+ \\ \pi^- + \dots \end{array} \right.$$

analogously: Kaon-induced reactions $K^- + \gamma \rightarrow \dots$



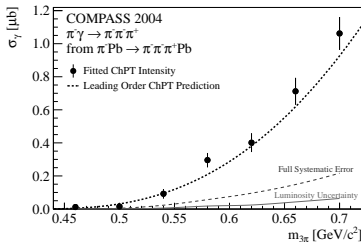
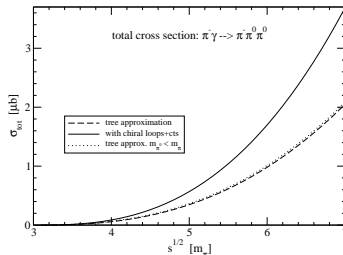
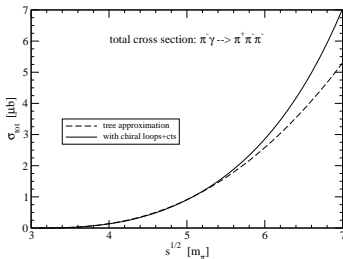
2004 Primakoff results



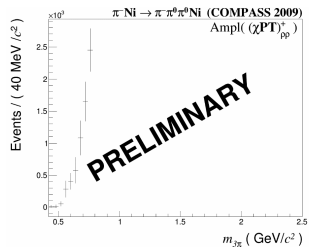
- "Low t' ": $10^{-3} \text{ (GeV/c)}^2 < t' < 10^{-2} \text{ (GeV/c)}^2 \sim 2\,000\,000$ events
- "Primakoff region": $t' < 10^{-3} \text{ (GeV/c)}^2 \sim 1\,000\,000$ events



Chiral dynamics in $\pi\gamma \rightarrow 3\pi$



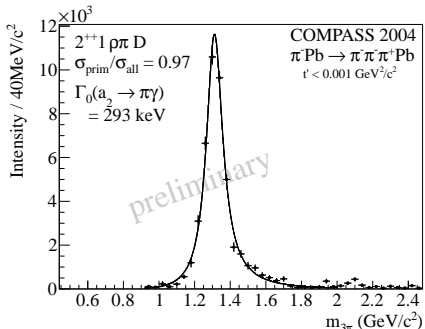
published in PRL 108 (2012) 192001



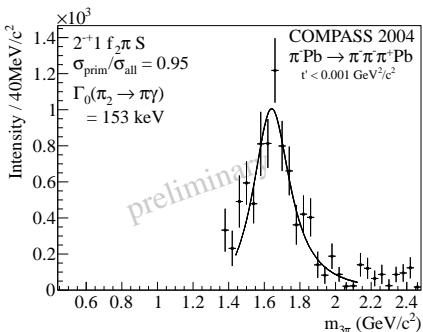
normalization: analysis ongoing



Radiative Coupling of $a_2(1320)$ and $\pi_2(1670)$



$\Gamma_0(a_2(1320) \rightarrow \pi\gamma)$ **M2**



$\Gamma_0(\pi_2(1670) \rightarrow \pi\gamma)$ **E2**

\Leftrightarrow meson w.f.'s: $\Gamma_{i \rightarrow f} \propto |\langle \Psi_f | e^{-i\vec{q} \cdot \vec{r}} \hat{\epsilon} \cdot \vec{p} | \Psi_i \rangle|^2$, VMD

- normalization via beam kaon decays
- large Coulomb correction

published in EPJ A50 (2014) 79



- Measurement of the **pion polarisability** at COMPASS
 - Via the Primakoff reaction, COMPASS has determined

$$\alpha_\pi = (2.0 \pm 0.6_{\text{stat}} \pm 0.7_{\text{syst}}) \times 10^{-4} \text{ fm}^3 \quad \text{assuming } \alpha_\pi + \beta_\pi = 0$$

- most direct access to the $\pi\gamma \rightarrow \pi\gamma$ process
 - Most precise experimental determination
 - Systematic control: $\mu\gamma \rightarrow \mu\gamma$, $K^- \rightarrow \pi^- \pi^0$
- COMPASS measures more aspects of chiral dynamics in $\pi^- \gamma \rightarrow \pi^- \pi^0$ and $\pi\gamma \rightarrow \pi\pi\pi$ reactions
- High-statistics run 2012
 - separate determination of α_π and β_π
 - s -dependent quadrupole polarisabilities
 - First measurement of the kaon polarisability



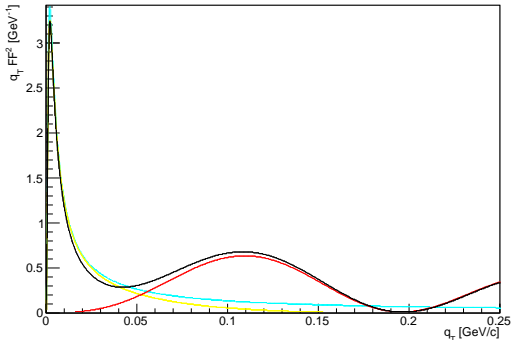
Thank you for your attention!



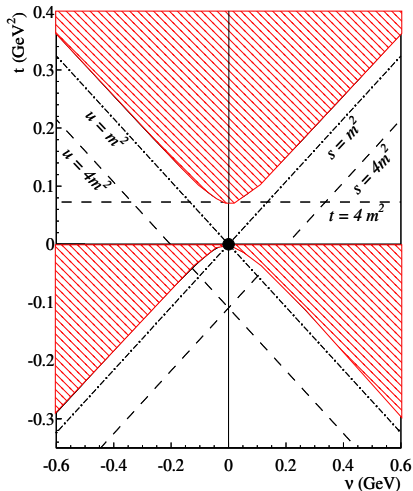




Photon density squared form factor

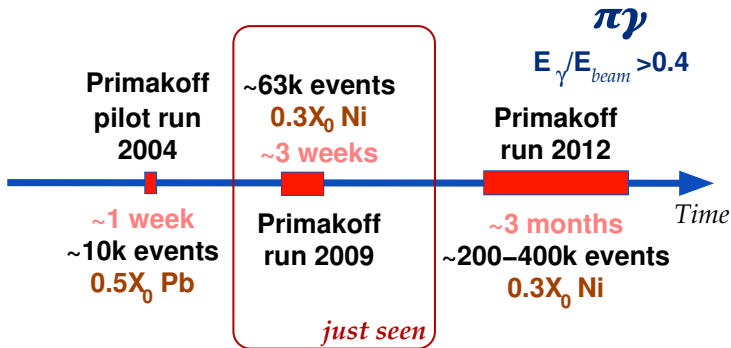


- calculation following G. Fäldt (Phys. Rev. C79, 014607)
- eikonal approximation: pions traverse Coulomb and strong-interaction potentials



- ▶ **red hatched:**
physical regions
 $\gamma + \gamma \rightarrow \pi + \pi$
 $\gamma + \pi \rightarrow \gamma + \pi$
- ▶ two-pion thresholds
at $s = 4m_\pi^2$, $u = 4m_\pi^2$,
 $t = 4m_\pi^2$
- ▶ DR integration paths
 $t = 0$ (forward),
 $\theta = 180^\circ$ (backward)
 $u = m_\pi^2$, $s = m_\pi^2$, ...

from: D. Drechsel, talk at IWHSS 2011 Paris

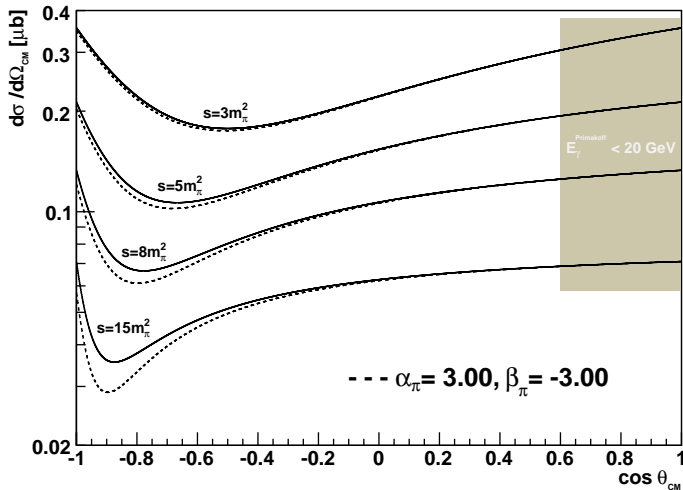


$$\mathcal{P} = z_-^2 (\alpha_\pi - \beta_\pi) + \frac{s^2}{m_\pi^4} z_+^2 (\alpha_\pi + \beta_\pi) - \frac{(s - m_\pi^2)^2}{24s} z_-^3 (\alpha_2 - \beta_2)$$

$$z_\pm = 1 \pm \cos \theta_{cm}$$

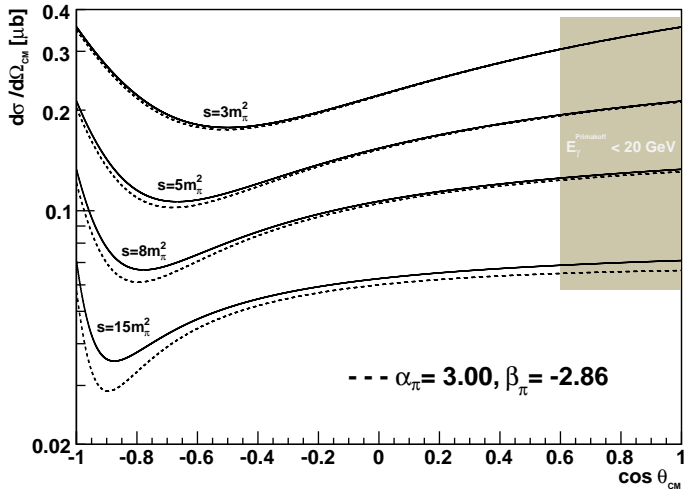


loop effects not shown



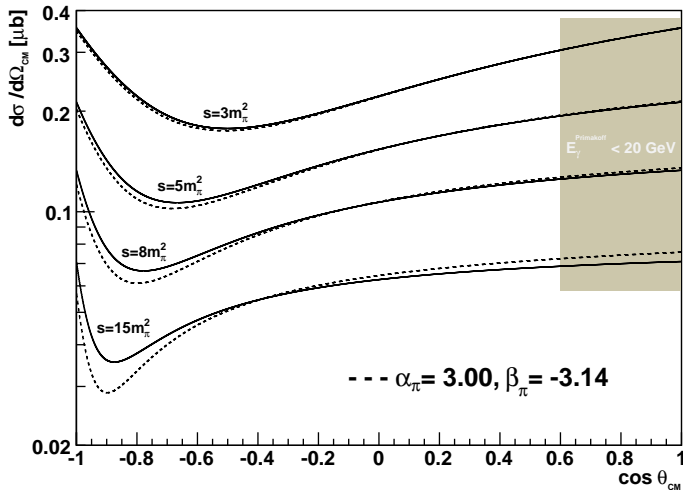


loop effects not shown



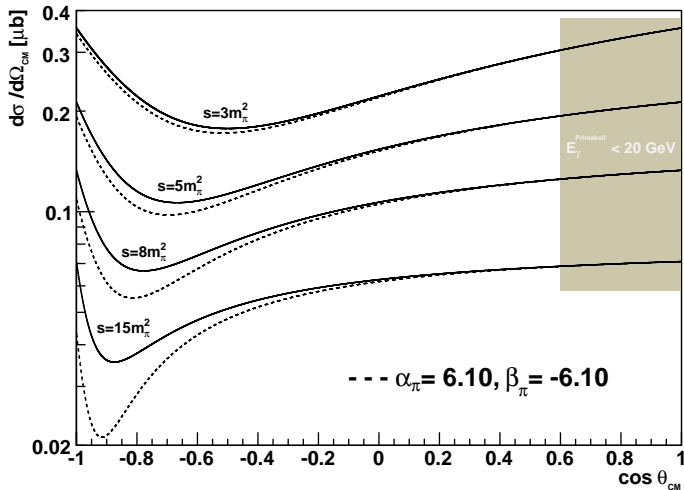


loop effects not shown





loop effects not shown





- Radiative π^+ production on the proton:



Mainz (2005) measurement: $\alpha_\pi - \beta_\pi = 11.6 \pm 1.5 \pm 3.0 \pm 0.5$

“ ± 0.5 ”: model error *only within the used ansatz*,

full systematics not under control

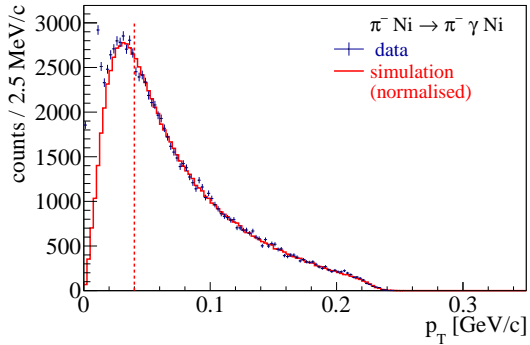
- Primakoff Compton reaction:

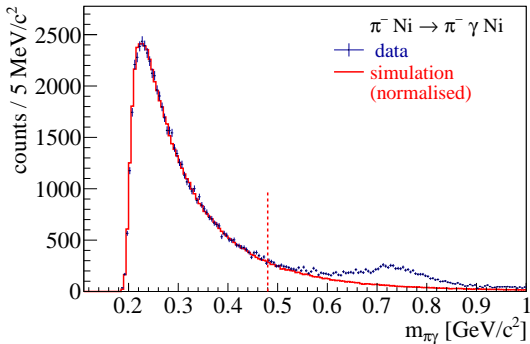


tiny extrapolation $\gamma^* \rightarrow \gamma \mathcal{O}(10^{-3} m_\pi^2)$

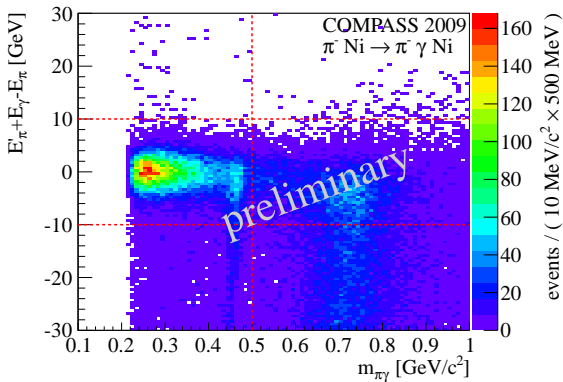
fully under theoretical control

[N. Kaiser, J.F., Nucl. Phys. A 812 (2008) 186]





- ρ contribution from $\pi\gamma \rightarrow \pi\pi^0$

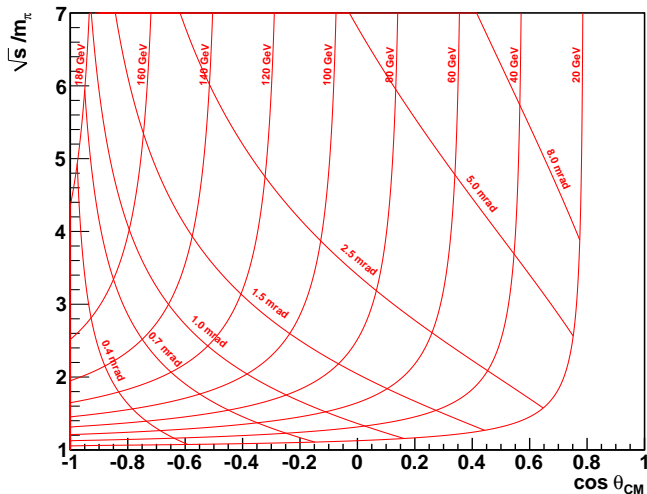


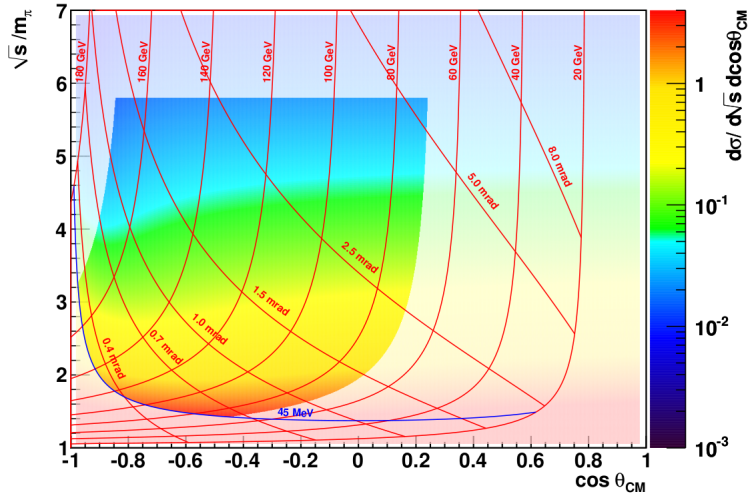
- ρ contribution from $\pi\gamma \rightarrow \pi\pi^0$



Mandelstam $\{s, t\} \leftrightarrow$ Laboratory $\{E_\gamma, \theta_\gamma\}$

for $\pi\gamma \rightarrow \pi\gamma$

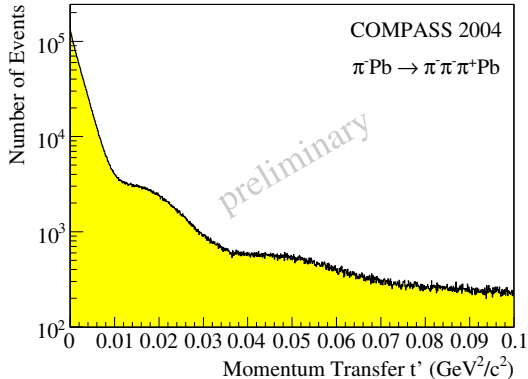




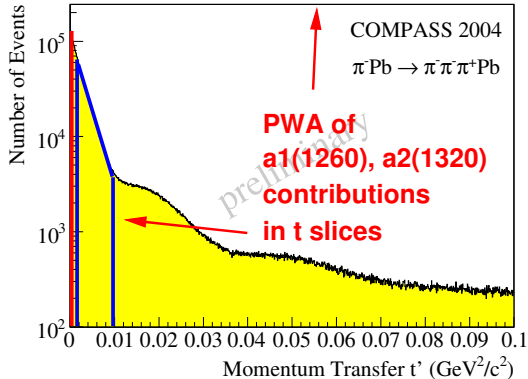


M.R. Pennington in the 2nd DAΦNE Physics Handbook,
“What we learn by measuring $\gamma\gamma \rightarrow \pi\pi$ at DAΦNE”:

All this means that the only way to measure the pion polarisabilities is in the Compton scattering process near threshold and not in $\gamma\gamma \rightarrow \pi\pi$. Though the low energy $\gamma\gamma \rightarrow \pi\pi$ scattering is seemingly close to the Compton threshold (...) and so the *extrapolation* not very far, the dominance of the pion pole (...) means that the energy scale for this continuation is m_π . Thus the polarisabilities cannot be determined accurately from $\gamma\gamma$ experiments in a model-independent way and must be measured in the Compton scattering region.



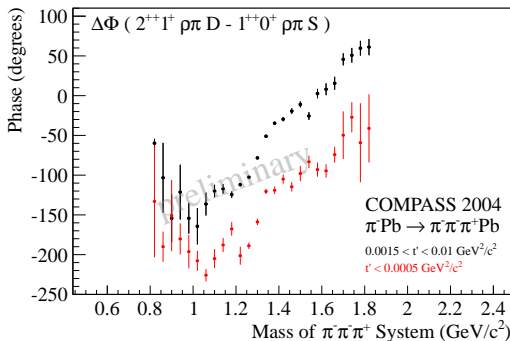
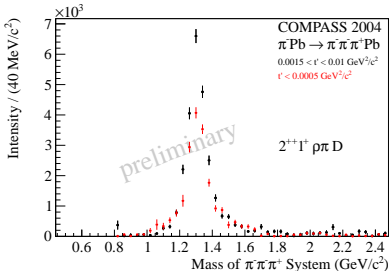
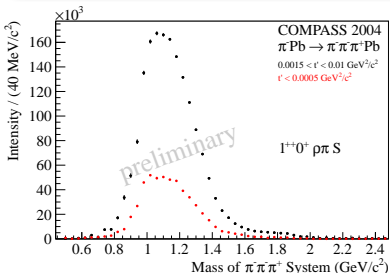
- "Low t' ": $10^{-3} (\text{GeV}/c)^2 < t' < 10^{-2} (\text{GeV}/c)^2$ $\sim 2\,000\,000$ events
- "Primakoff region": $t' < 10^{-3} (\text{GeV}/c)^2$ $\sim 1\,000\,000$ events

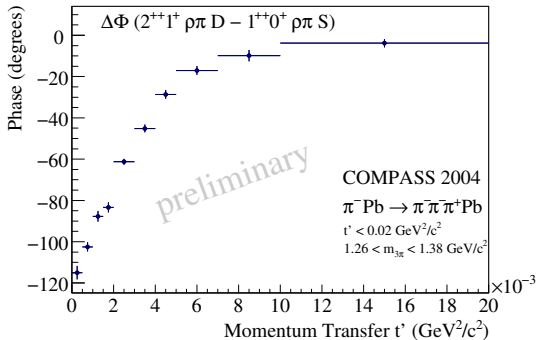
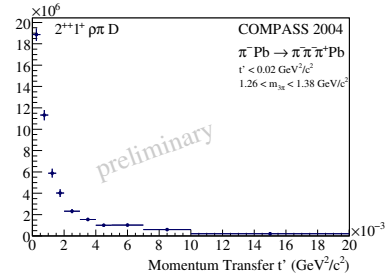
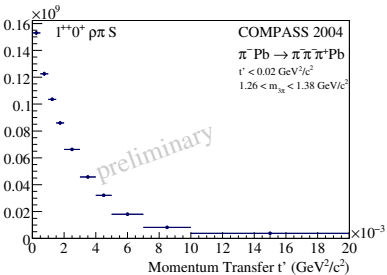


- "Low t' ": $10^{-3} (\text{GeV}/c)^2 < t' < 10^{-2} (\text{GeV}/c)^2$ $\sim 2\,000\,000$ events
- "Primakoff region": $t' < 10^{-3} (\text{GeV}/c)^2$ $\sim 1\,000\,000$ events



PWA: a_1 , a_2 and $\Delta\Phi$ in separated t' regions

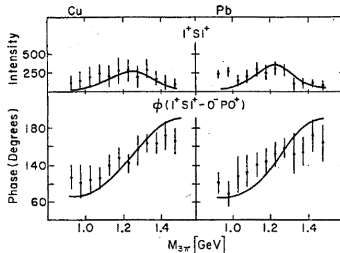
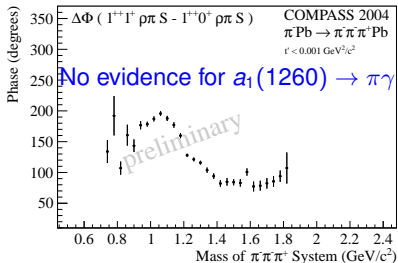
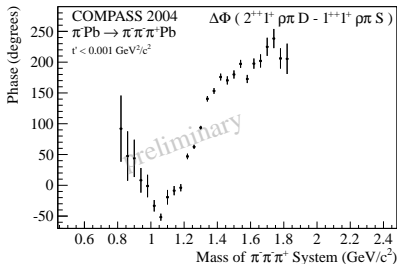
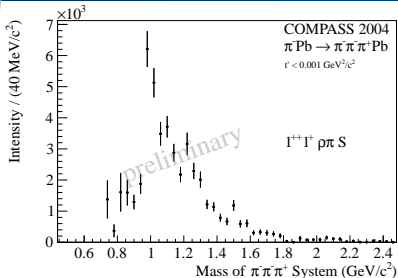




- transition of $\pi\gamma$ to $\pi IP \rightarrow a_2$ production
- work in progress
- interference can be used to map details of resonances and production mechanisms



Primakoff production of $a_1(1260)$ vs. E272 result



M. Zielinski et al, Phys. Rev. Lett 52 (1984) 1195



- **Mass-independent PWA** (narrow mass bins):

$$\sigma_{\text{indep}}(\tau, m, t') = \sum_{\epsilon=\pm 1} \sum_{r=1}^{N_r} \left| \sum_i T_{ir}^\epsilon f_i^\epsilon(t') \psi_i^\epsilon(\tau, m) / \sqrt{\int |f_i^\epsilon(t')|^2 dt'} \sqrt{\int |\psi_i^\epsilon(\tau', m)|^2 d\tau'} \right|^2$$

- Production strength assumed constant in single bins
 - Decay amplitudes $\psi_i^\epsilon(\tau, m)$, with t' dependence $f_i^\epsilon(t')$
 - Production amplitudes $T_{ir}^\epsilon \rightarrow$ Extended log-likelihood fit
 - Acceptance corrections included
- **Spin-density matrix:** $\rho_{ij}^\epsilon = \sum_r T_{ir}^\epsilon T_{jr}^{\epsilon*}$

\rightarrow Physical parameters:

$$\text{Intens}_i^\epsilon = \rho_{ii}^\epsilon,$$

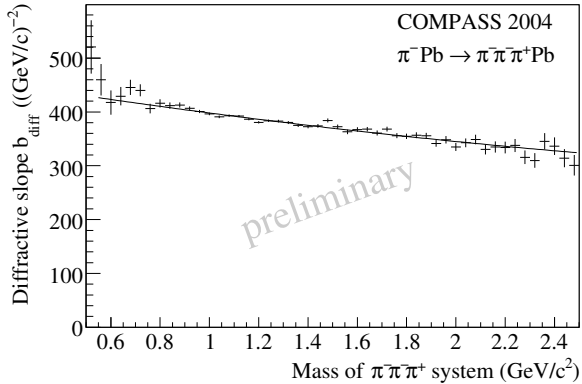
relative phase Φ_{ij}^e

$$\text{Coh}_{i,j}^\epsilon = \sqrt{(\text{Re } \rho_{ij}^\epsilon)^2 + (\text{Im } \rho_{ij}^\epsilon)^2} / \sqrt{\rho_{ii}^\epsilon \rho_{jj}^\epsilon}$$

- **Mass-dependent χ^2 -fit** (not presented here):
 - X parameterized by Breit-Wigner (BW) functions
 - Background can be added

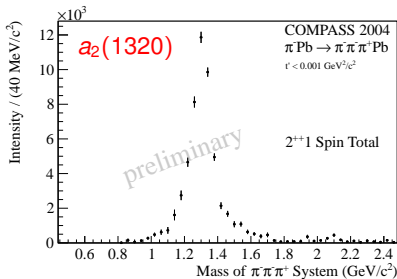
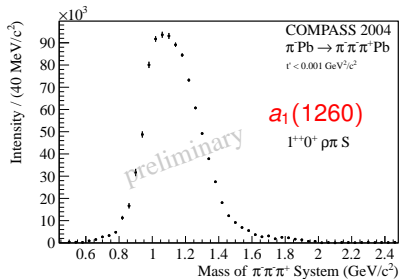
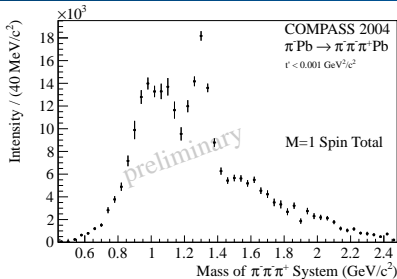
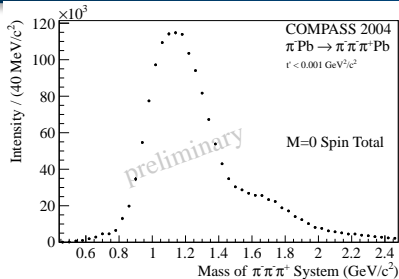


Mass dependence of the diffractive slope





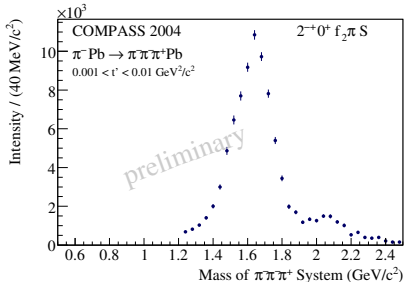
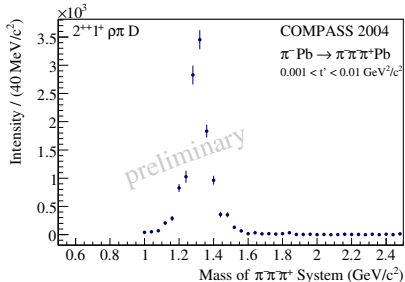
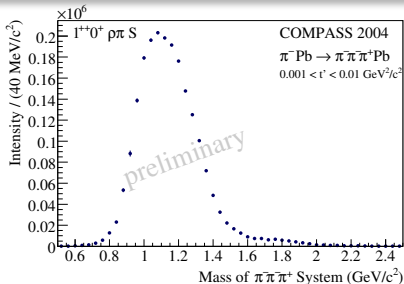
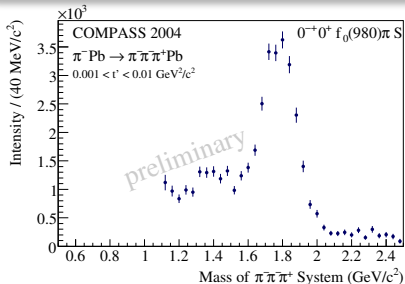
Major intensities in $m(3\pi)$ -bins (acceptance corrected)





PWA of data with low t'

Intensity of selected waves: $0^{-+}0^{+}f_0(980)\pi S$, $1^{++}0^{+}\rho\pi S$, $2^{++}1^{+}\rho\pi D$, $2^{-+}0^{+}f_2(1270)\pi S$





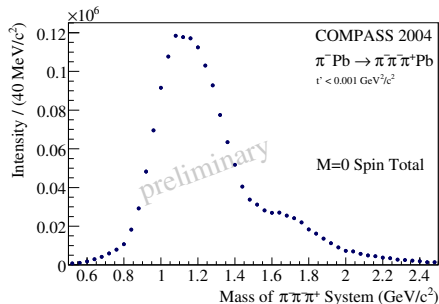
“Spin Totals”: Sum of all contributions for given M (i.e. z-projection of J)

t' -dependent amplitudes:

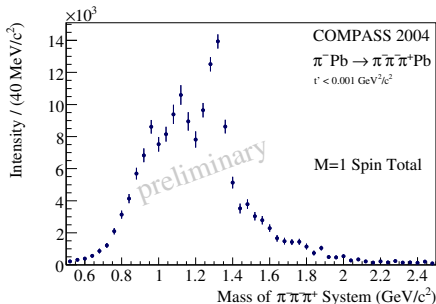
Primakoff production: **M=1**: $\sigma(t') \propto e^{-b_{\text{Prim}}t'}$ → arises at $t' \approx 0$ (resolved shape!)

Diffractive production: **M=0**: $\sigma(t') \propto e^{-b_{\text{diff}}(m)t'}$

M=1: $\sigma(t') \propto t' e^{-b_{\text{diff}}(m)t'}$ → vanishes for $t' \approx 0$



M=0

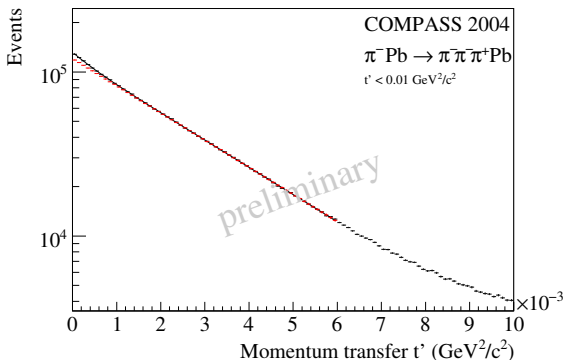


M=1



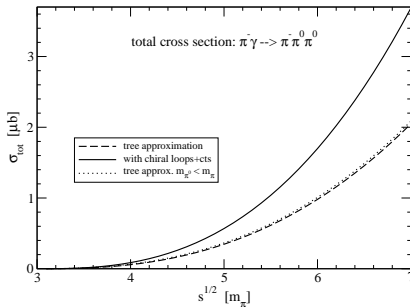
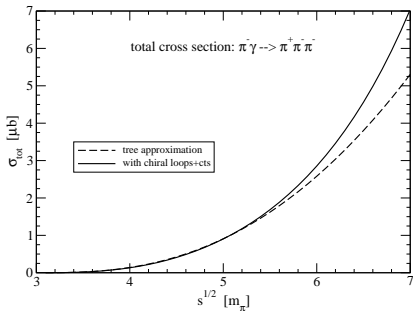
Primakoff: $\sigma(t') \propto e^{-b_{\text{Prim}} t'}$, $b_{\text{Prim}} \approx 2000 \text{ (GeV/c)}^{-2}$ (mainly resolution)

Diffractive: $\sigma(t') \propto e^{-b_{\text{diff}} t'}$, $b_{\text{diff}} \approx 400 \text{ (GeV/c)}^{-2}$ for lead target



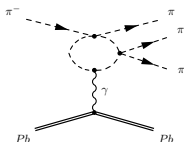
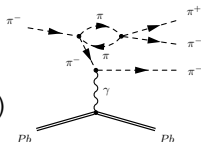
(Mass) spectrum of this Primakoff contribution?

⇒ Statistical subtraction of diffractive background (for bins of $m_{3\pi}$)

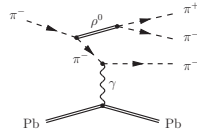


Chiral loops, e.g.

(N. Kaiser, NPA848 (2010) 198)

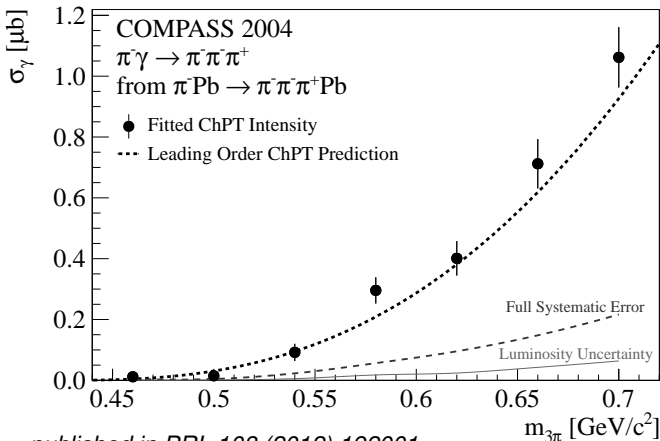
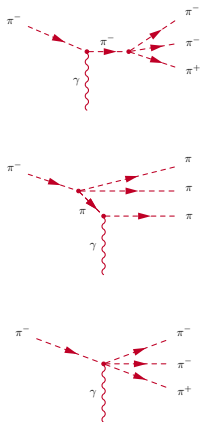


ρ terms:





Measured absolute cross-section of $\pi^-\gamma \rightarrow \pi^-\pi^-\pi^+$



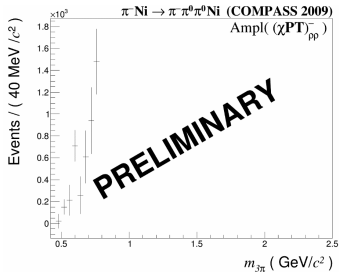
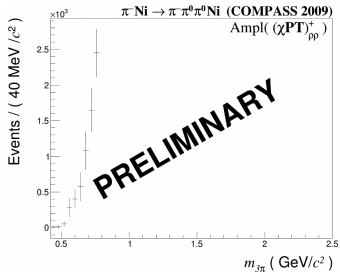
published in *PRL* 108 (2012) 192001



$$\pi^- \gamma \rightarrow \pi^- \pi^0 \pi^0$$

Partial Wave Analysis

Isobaric Model – Chiral Wave





$$\pi^- \gamma \rightarrow \pi^- \pi^0 \pi^0$$

Partial Wave Analysis

Chiral Model - Amplitudes

