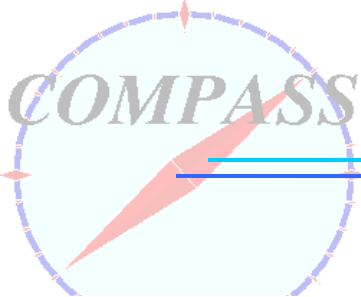


Transverse spin and transverse momentum structure of the nucleon from the COMPASS experiment

Federica Sozzi
INFN Sezione di Trieste




on behalf of the **COMPASS** Collaboration





Quark structure of the nucleon

At leading order, the inner structure of the nucleon can be described with three **Parton Distribution Function** (PDF):

| | | nucleon polarisation | | |
|-----------------------|---|--|---|---|
| | | U | L | T |
| quark polarisation | U | f_1  <i>number density</i> q | | |
| | L | | g_1  <i>helicity</i> Δq | |
| | T | | | h_1  <i>transversity</i> |

helicity and transversity PDF:

probability of finding a quark with a momentum fraction x and spin parallel to that of the parent nucleon...

...in a longitudinally polarised nucleon
















...in a transversely polarised nucleon

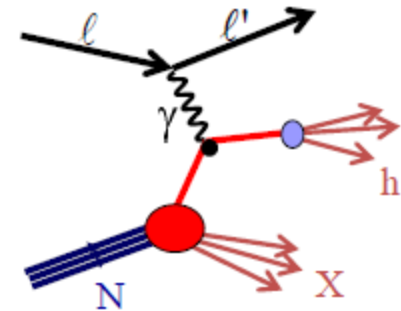
Taking into account the intrinsic parton transverse momentum, the nucleon structure description becomes more complex and needs 8 **“Transverse Momentum Dependent” PDF**.

TMDs describe the correlations between the spin and the momentum of quarks and of the parent nucleon

nucleon polarisation

quark polarisation

| | U | L | T |
|---|---|--|--|
| U | f_1  <i>number density</i> q | | f_{1T}^\perp  -  Sivers |
| L | | g_1  -  <i>helicity</i> Δq | g_{1T}  -  Worm-gear |
| T | h_1^\perp  -  Boer Mulders | h_{1L}^\perp  -  Worm-gear | h_1  -  <i>transversity</i> h_{1T}^\perp  -  pretzelosity |



One experimental way to access TMDs :
SIDIS off polarised targets

SIDIS cross section

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} =$$

$$\frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h \boxed{F_{UU}^{\cos\phi_h}} \right.$$

$$+ \varepsilon \cos(2\phi_h) \boxed{F_{UU}^{\cos 2\phi_h}} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h \boxed{F_{LU}^{\sin\phi_h}}$$

$$+ \boxed{S_{\parallel}} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h \boxed{F_{UL}^{\sin\phi_h}} + \varepsilon \sin(2\phi_h) \boxed{F_{UL}^{\sin 2\phi_h}} \right]$$

$$+ \boxed{S_{\parallel} \lambda_e} \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h \boxed{F_{LL}^{\cos\phi_h}} \right]$$

$$+ \boxed{S_{\perp}} \left[\sin(\phi_h - \phi_S) \left(\boxed{F_{UT,T}^{\sin(\phi_h - \phi_S)}} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right.$$

$$+ \varepsilon \sin(\phi_h + \phi_S) \boxed{F_{UT}^{\sin(\phi_h + \phi_S)}} + \varepsilon \sin(3\phi_h - \phi_S) \boxed{F_{UT}^{\sin(3\phi_h - \phi_S)}} \left. \right]$$

$$+ \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S \boxed{F_{UT}^{\sin\phi_S}} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) \boxed{F_{UT}^{\sin(2\phi_h - \phi_S)}} \left. \right]$$

$$+ \boxed{S_{\perp} |\lambda_e} \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) \boxed{F_{LT}^{\cos(\phi_h - \phi_S)}} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S \boxed{F_{LT}^{\cos\phi_S}} \right.$$

$$\left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) \boxed{F_{LT}^{\cos(2\phi_h - \phi_S)}} \right\},$$

several structure functions,
depending on different
combinations of azimuthal angles

Each structure function :
convolution of PDF and FF

COMPASS measures all of them

Common Muon and Proton Apparatus for Structure and Spectroscopy

Fixed target experiment
at CERN SPS
Data taking since 2002

Nucleon spin structure
with high energy muon beams
on longitudinally polarized targets:

- gluon polarization
- helicity PDF

Transversely polarized targets:
transversity PDF
TMDs

this talk



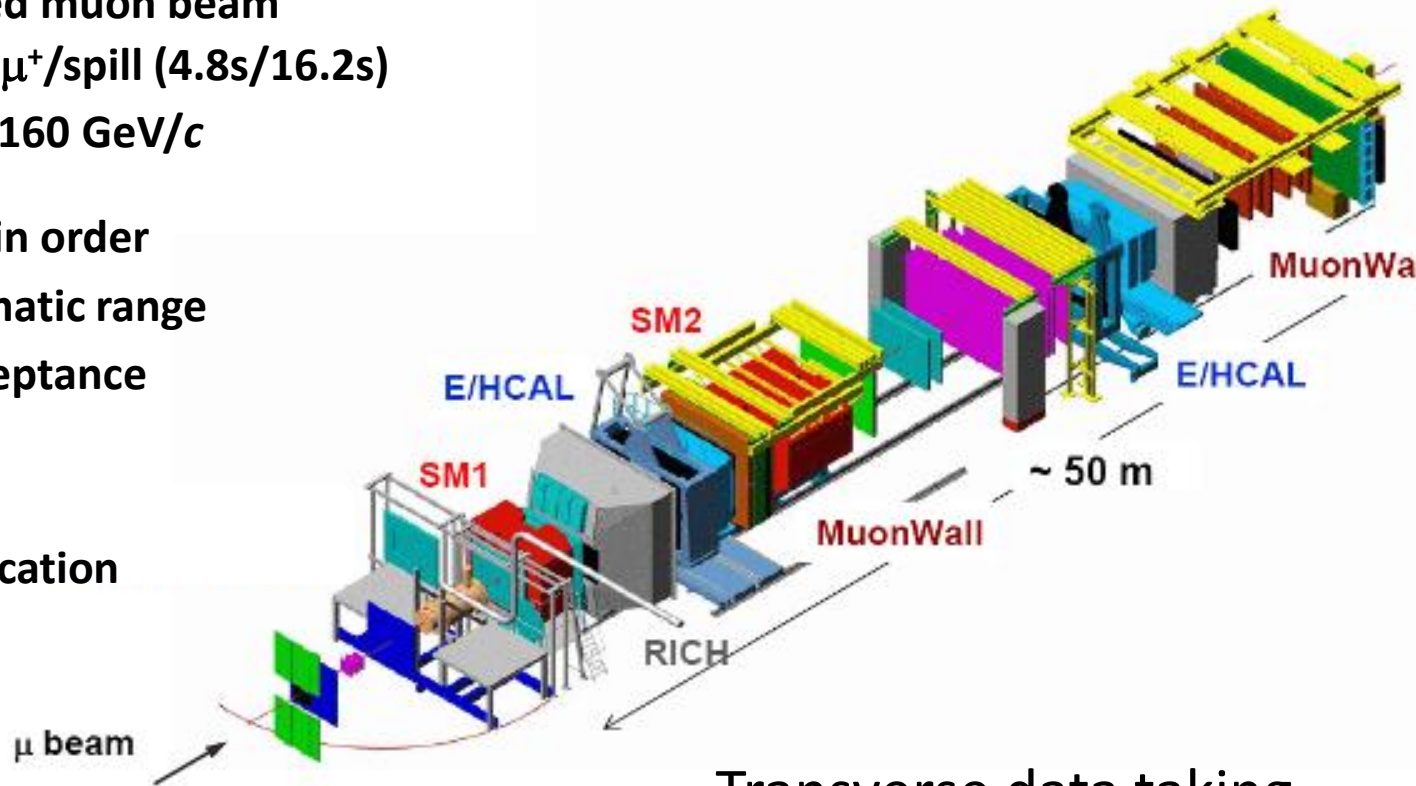
Meson and baryon spectroscopy
with high energy hadron beams

*See talks by C. Marchand, C. Quintans
Poster by J. Bernhard*

longitudinally polarised muon beam
 beam intensity: $2 \cdot 10^8 \mu^+/\text{spill}$ (4.8s/16.2s)
 beam momentum: 160 GeV/c

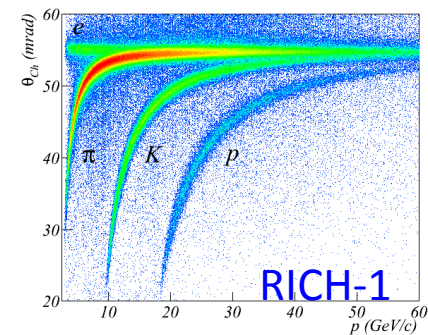
2 stage spectrometer in order
 to cover a large kinematic range
 180 mrad angular acceptance

Muon identification
 Detection and identification
 of hadrons for SIDIS
 measurements



Transverse data taking
 2002-4: ${}^6\text{LiD}$ target
 $p_T \sim 50\%$; $f \sim 0.38$
 2007/2010: NH_3 target
 $p_T \sim 90\%$; $f \sim 0.15$

π , K, p separation
 (from 2, 9, 17 GeV/c
 to ~ 50 GeV/c)



$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} =$$

$$\frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right.$$

$$+ \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h}$$

$$\left. + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \right.$$

18 structure functions,
8 transverse target
dependent spin asymmetries
with different azimuthal dependences

$$+ S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right]$$

$$+ |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right.$$

$$+ \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)}$$

$$\left. + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right]$$

$$+ |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right.$$

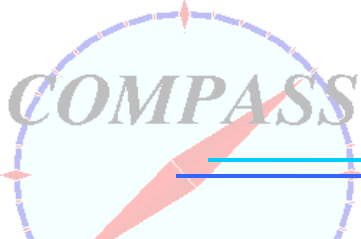
$$\left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \left. \right\},$$

Sivers

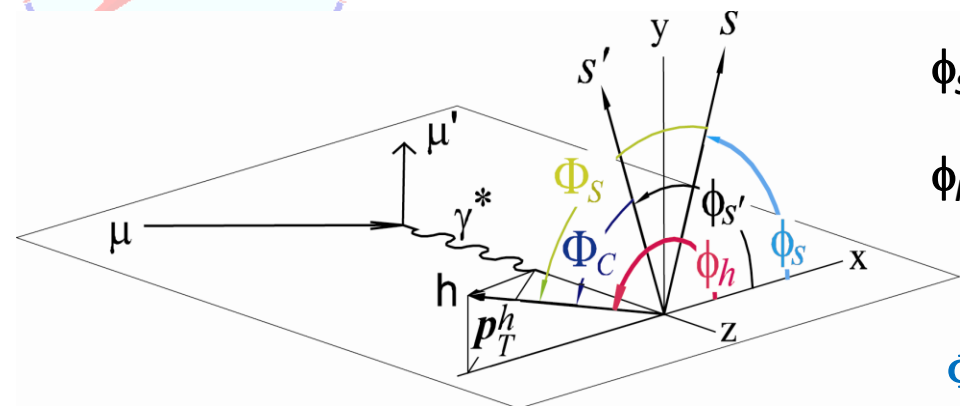
Collins

From [A. Bacchetta et al.](#),

[JHEP 0702:093,2007](#). e-Print: [hep-ph/0611265](#)



Azimuthal modulations: Sivers SSA



ϕ_s azimuthal angle of spin vector of initial quark

ϕ_h azimuthal angle of hadron momentum

$$\Phi_s = \phi_h - \phi_s \quad \text{Sivers angle}$$

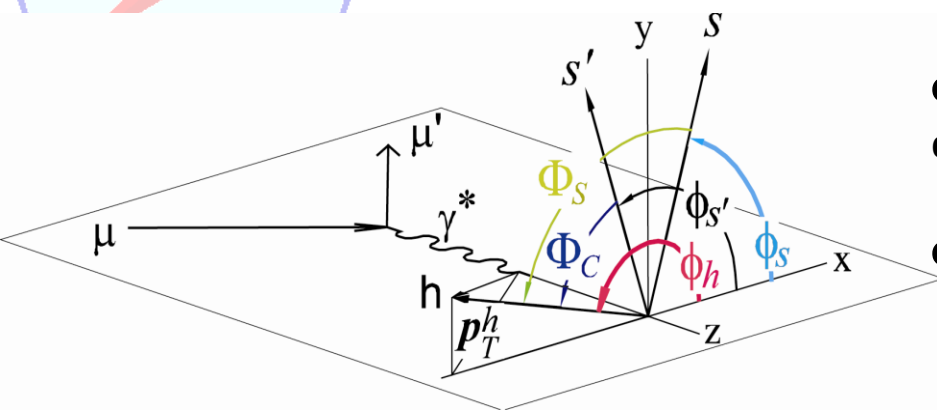
$$\mathbf{A}_{\text{Siv}} = \frac{\mathbf{A}_S^h}{\mathbf{f} \cdot \mathbf{P}_T} = \frac{\sum_q e_q^2 \cdot \mathbf{f}_{1T}^{\perp q} \otimes \mathbf{D}_{1q}^h}{\sum_q e_q^2 \cdot \mathbf{f}_1^q \cdot \mathbf{D}_{1q}^h}$$

Sivers PDF \otimes FF



Sivers PDF : correlation between the intrinsic transverse momentum of unpolarized quarks in a transversely polarized nucleon

Sensitive to orbital angular momentum



$\phi_{S'}$, azimuthal angle of spin vector of fragmenting quark ($\phi_{S'} = \pi - \phi_S$)

ϕ_h azimuthal angle of hadron momentum

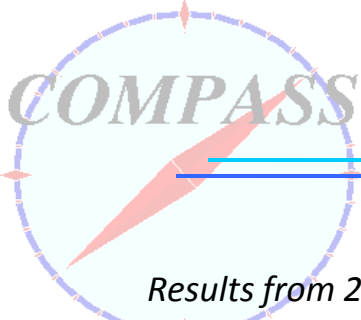
$$\Phi_C = \phi_h - \phi_{S'} = \phi_h + \phi_S - \pi \quad \text{Collins angle}$$

$$A_{\text{Coll}} = \frac{A_C^h}{\mathbf{f} \cdot \mathbf{P}_T \cdot D_{nn}} = \frac{\sum_q e_q^2 \cdot \mathbf{h}_1^q \otimes \mathbf{H}_{1q}^{\perp h}}{\sum_q e_q^2 \cdot \mathbf{f}_1^q \cdot \mathbf{D}_{1q}^h}$$

← “transversity” PDF \otimes Collins FF

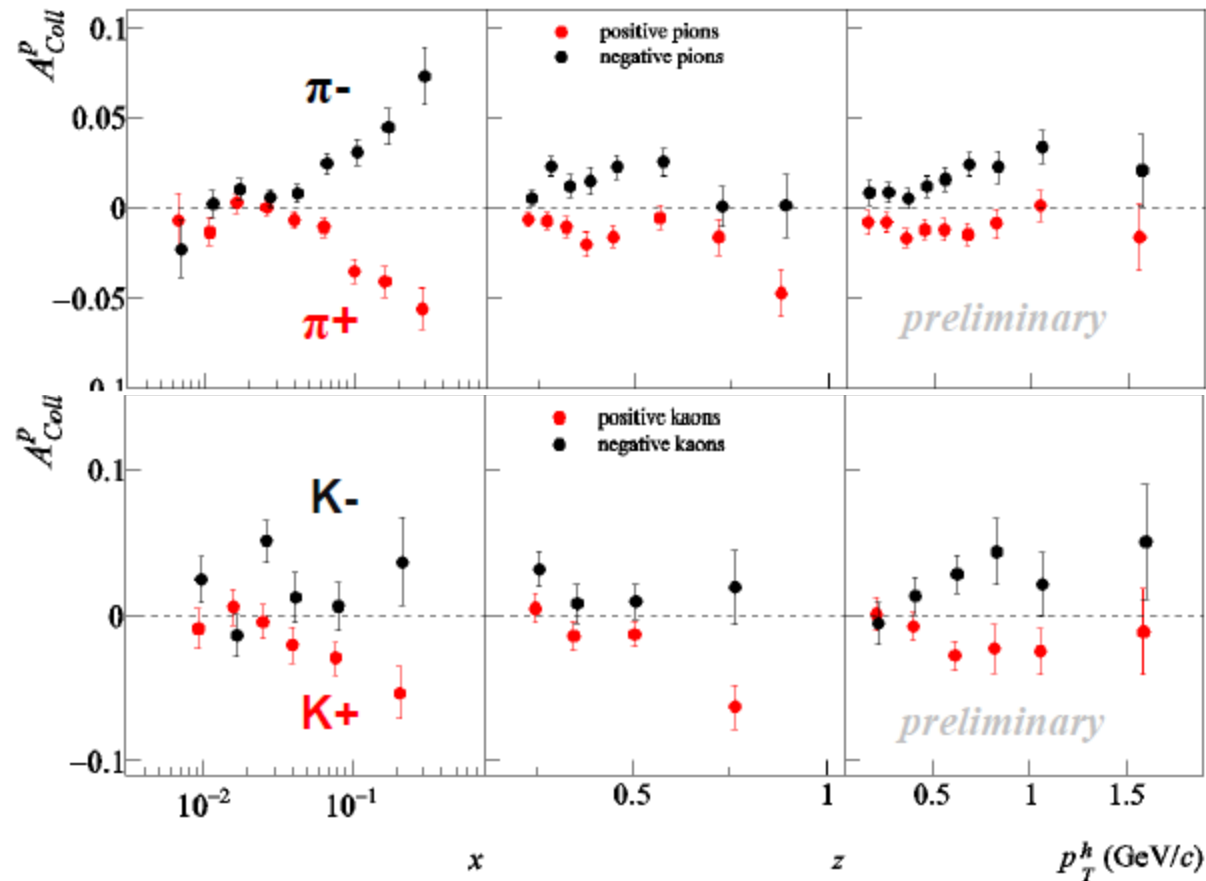


Transversity coupled to another chiral-odd function: [Collins fragmentation function](#) describing the correlation between the fragmenting quark spin and the hadron momentum



Collins asymmetries, results on proton

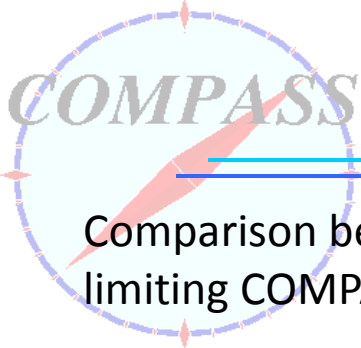
Results from 2007 and 2010 data taking



syst error \sim 0.6 stat error

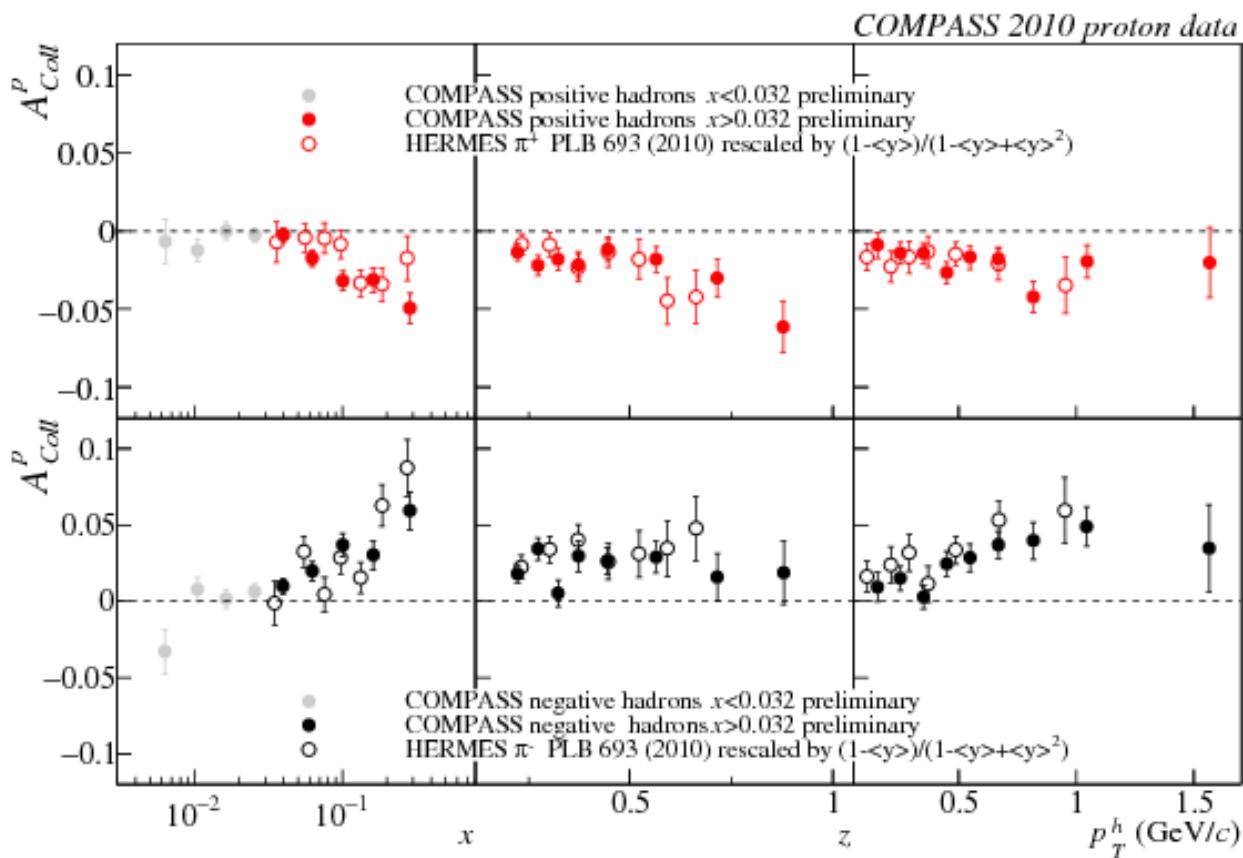
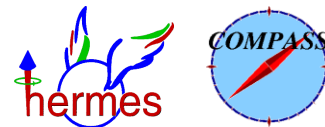
- at small x asymmetries compatible with zero
- Strong signal in the valence region of opposite sign for π^+ and π^-
- opposite sign
 - Dunf \sim Dfav

K^+ negative trend in the valence region
 K^- positive in average



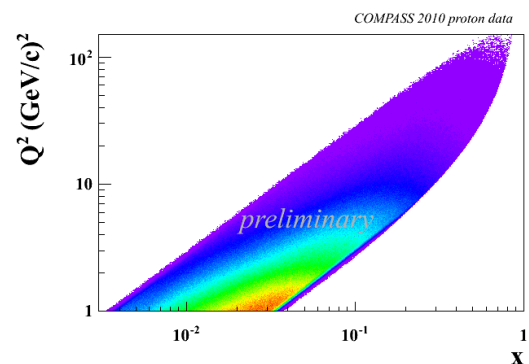
Collins asymmetries, results on proton

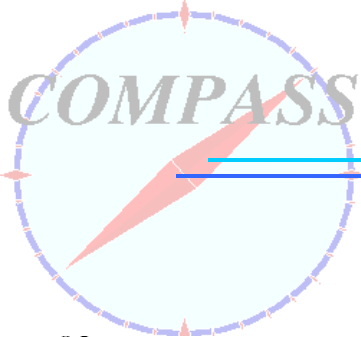
Comparison between HERMES and COMPASS,
limiting COMPASS range to the $x > 0.032$ region, overlap with HERMES



Good agreement :

- Non trivial result:
 Q^2 COMPASS larger of HERMES's of a factor 2-3 in the last x bins
→ weak Q^2 dependence of the Collins effect

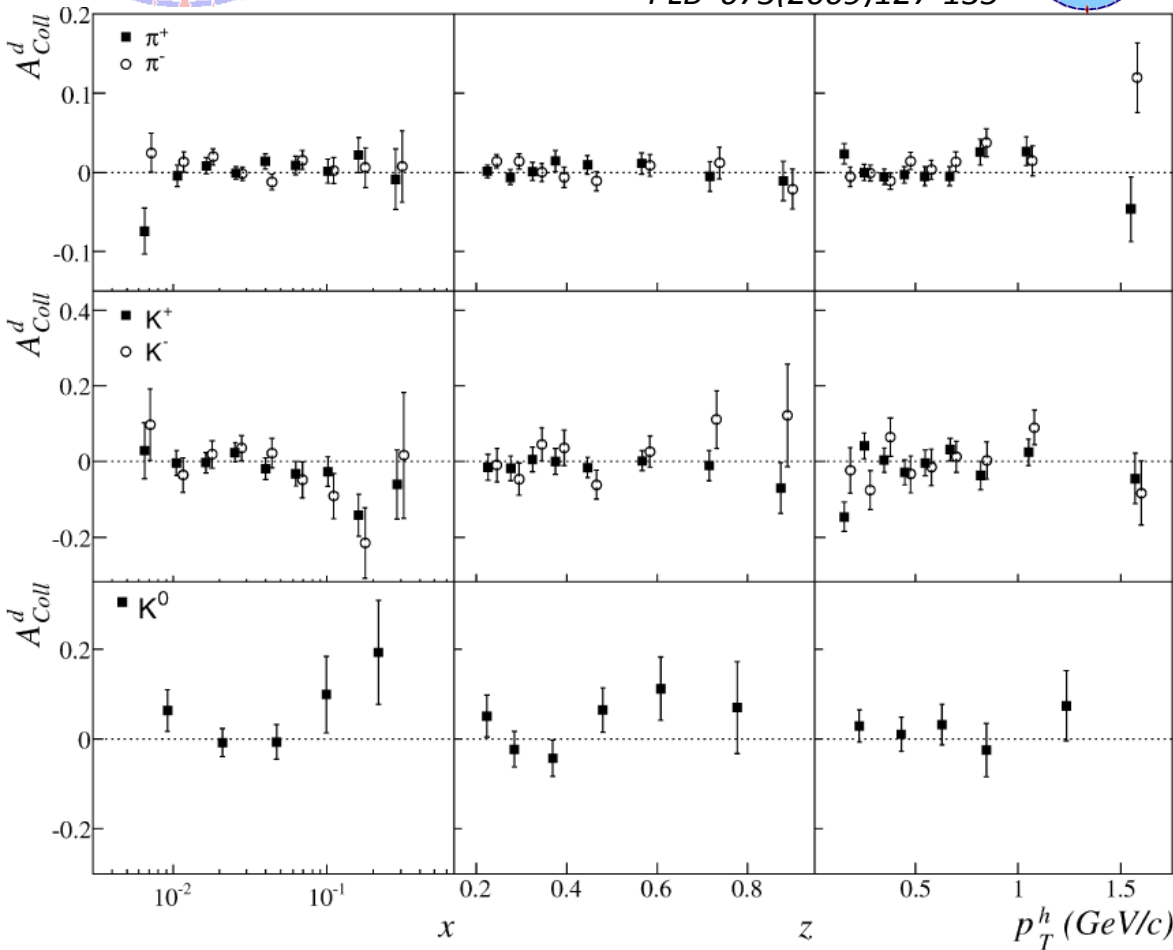




Collins asymmetries, results on deuterium



PLB 673(2009)127-135

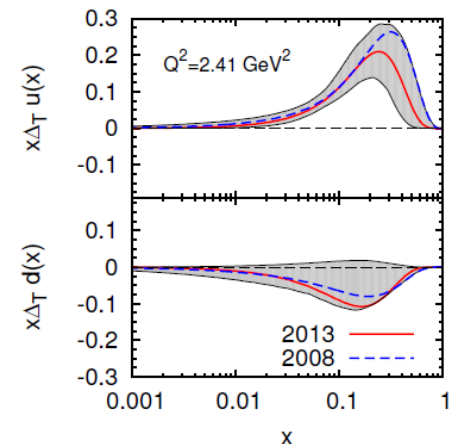


systematic error below 30% of the statistical one

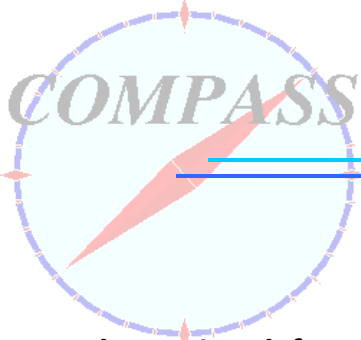
Asymmetries on deuteron target compatible with zero

Some small effects expected even if $H_{unf}^1 \sim -H_{fav}^1$
 \rightarrow cancellation between $\Delta_T u(x)$ and $\Delta_T d(x)$

handle on $\Delta_T d(x)$



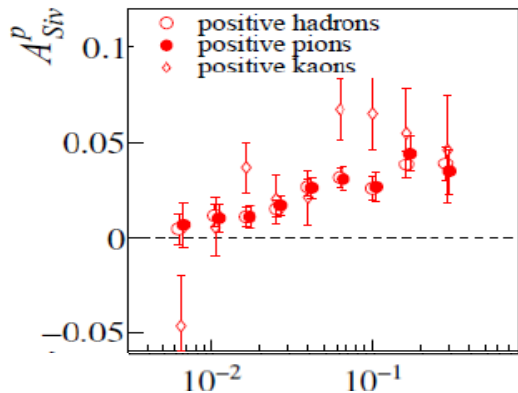
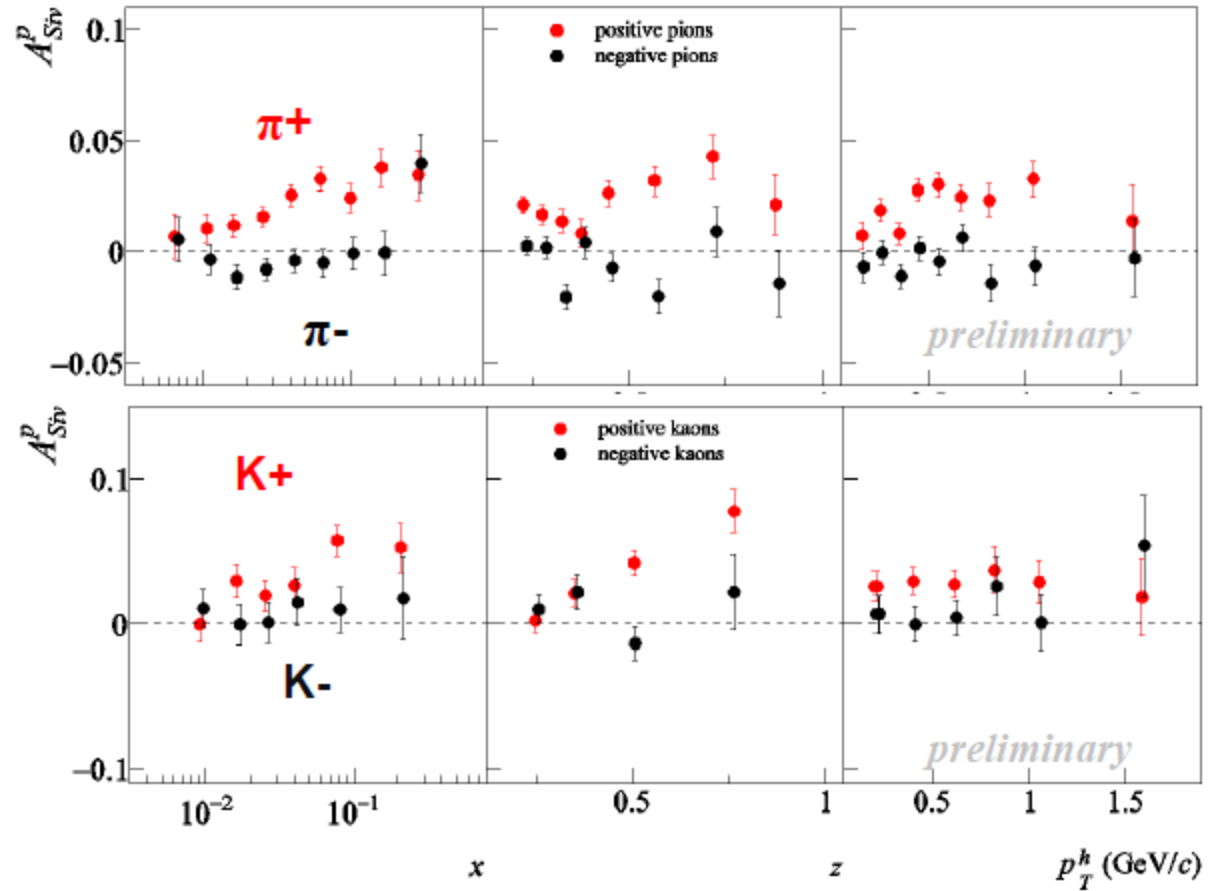
fit by Torino group
(arxiv 1303.3822)



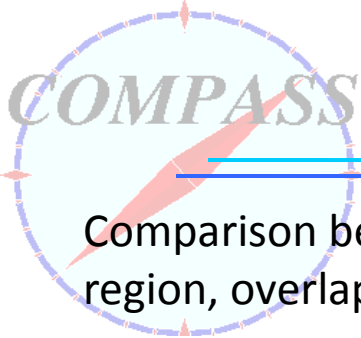
Sivers asymmetries, results on proton

Results from 2007 and 2010 data taking

- large signal for π^+ and K^+ over all the measured x range
- increasing with z
- linear behavior at small p_T , saturation for $P_T^h > 0.4 \text{ GeV}/c$
- K^+ positive in average
 K^- compatible with 0
- difference between K^+ and π^+ :
important role of sea quarks?

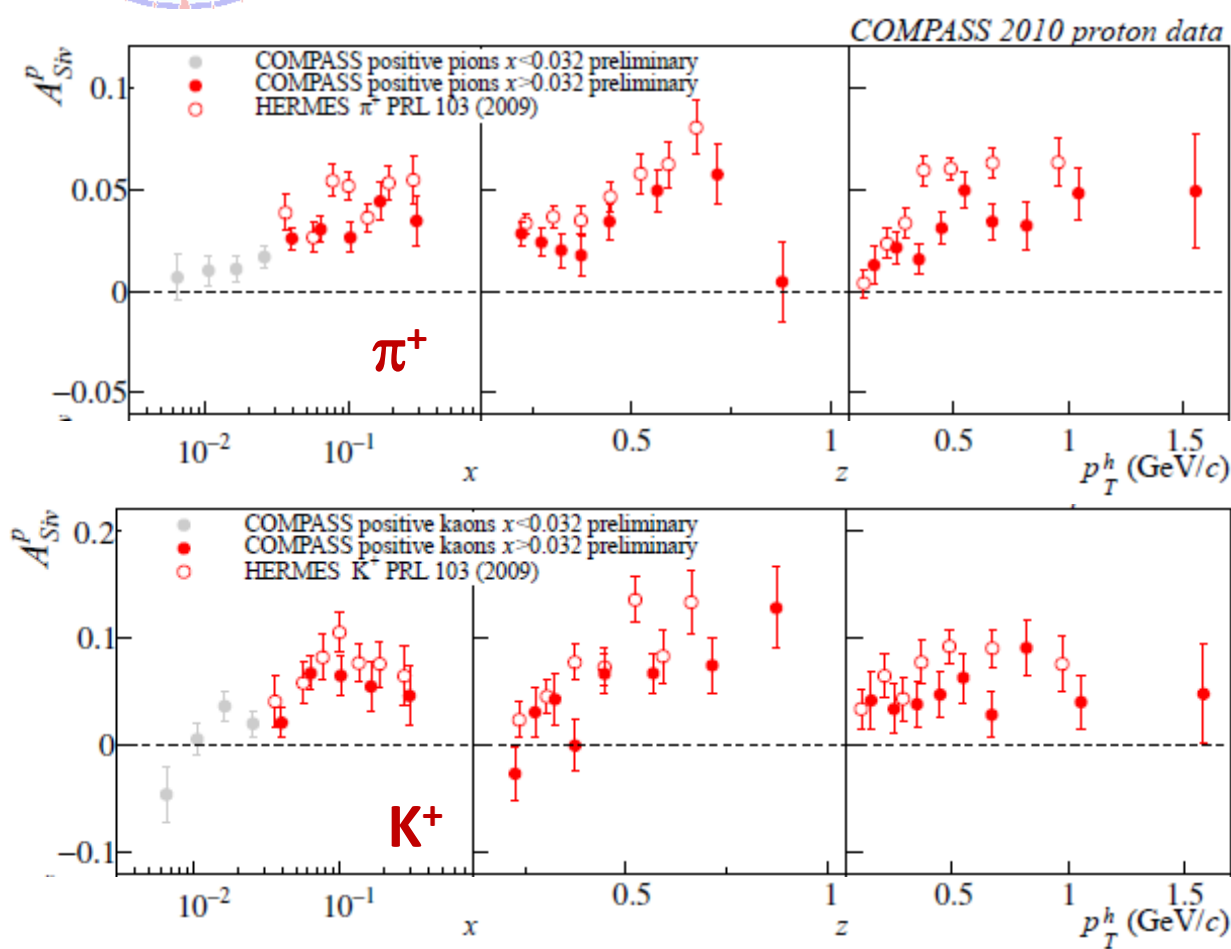


syst error ~ 0.6 stat error



Sivers asymmetries, results on proton

Comparison between HERMES and COMPASS, limiting COMPASS range to the $x > 0.032$ region, overlap with HERMES



HERMES π^+ and K^+ asymmetries larger than COMPASS

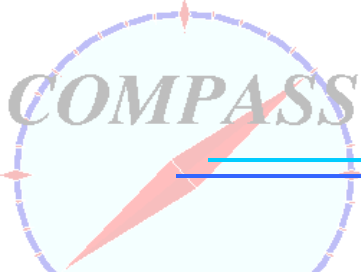
Q^2 COMPASS larger of HERMES's of a factor 2-3 in the last x bins
→ Q^2 dependence of the Sivers effect plays a role



TMD Q^2 evolution has been worked out and added in global fits very recently

S. M. Aybat, A. Prokudin, T. C. Rogers
PRL 108 (2012) 242003

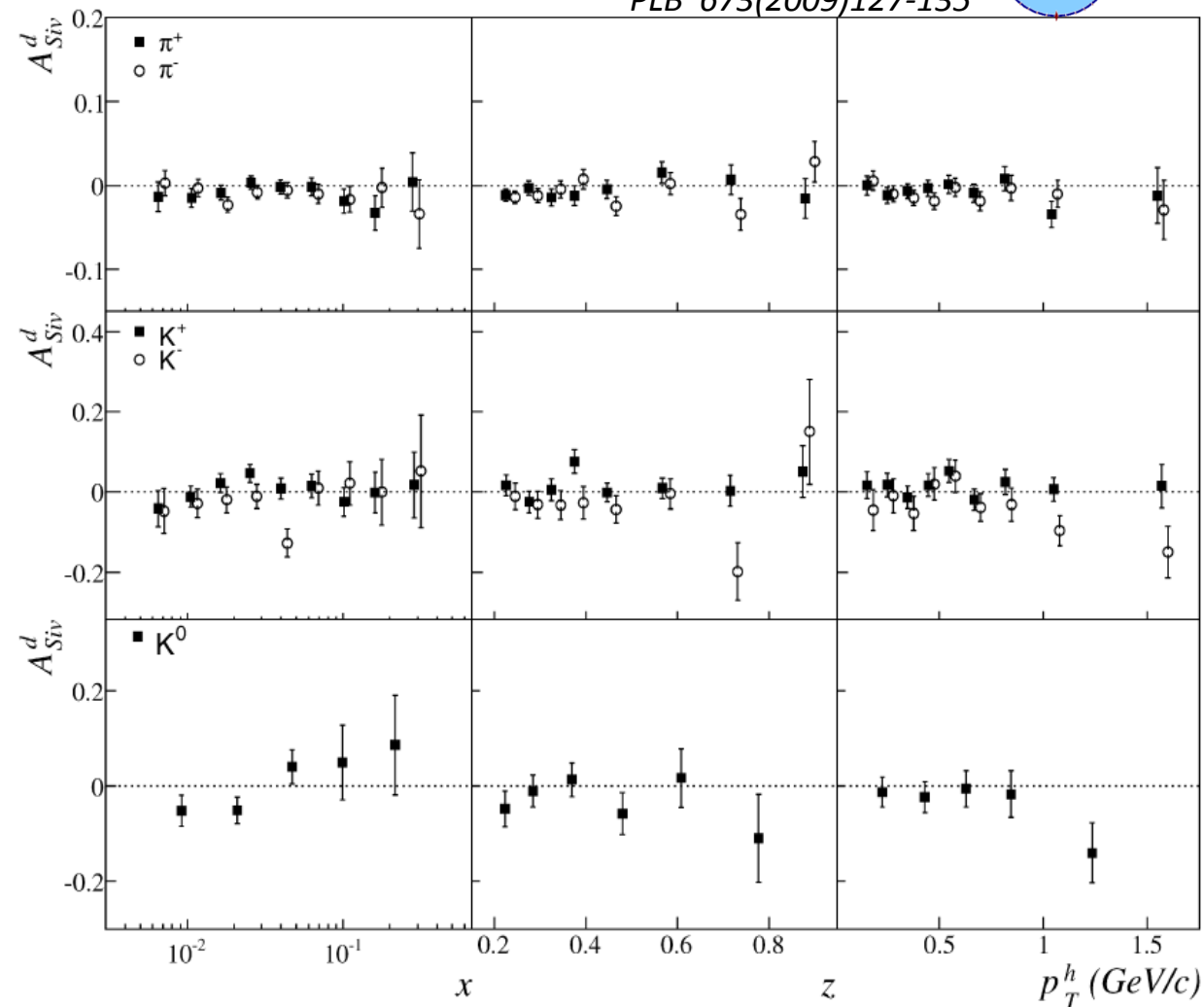
M. Anselmino, M. Boglione, S. Melis
PRD 86 (2012) 014028



Sivers asymmetries, results on deuterium



PLB 673(2009)127-135



Asymmetries on deuteron target compatible with zero

→ cancellation between $\Delta_T u(x)$ and $\Delta_T d(x)$

handle on $\Delta_T d(x)$

systematic error below 30% of the statistical one

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} =$$

$$\frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right.$$

$$+ \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h}$$

$$\left. + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \right.$$

Sivers

other twist-2 contributions

$$+ S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right]$$

Collins

$$+ |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right.$$

pretzelosity $h_{IT}^{\perp} \otimes H_1^{\perp}$

All asymmetries on p and d,
compatible with zero

$$+ \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)}$$

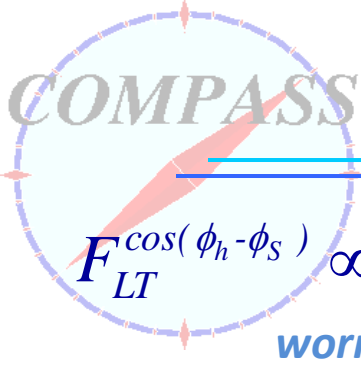
wormgear $g_{IT} \otimes D_1$

$$\left. + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right]$$

$$+ |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right.$$

Remaining four can be
interpreted as twist-3
contributions

$$\left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right\},$$



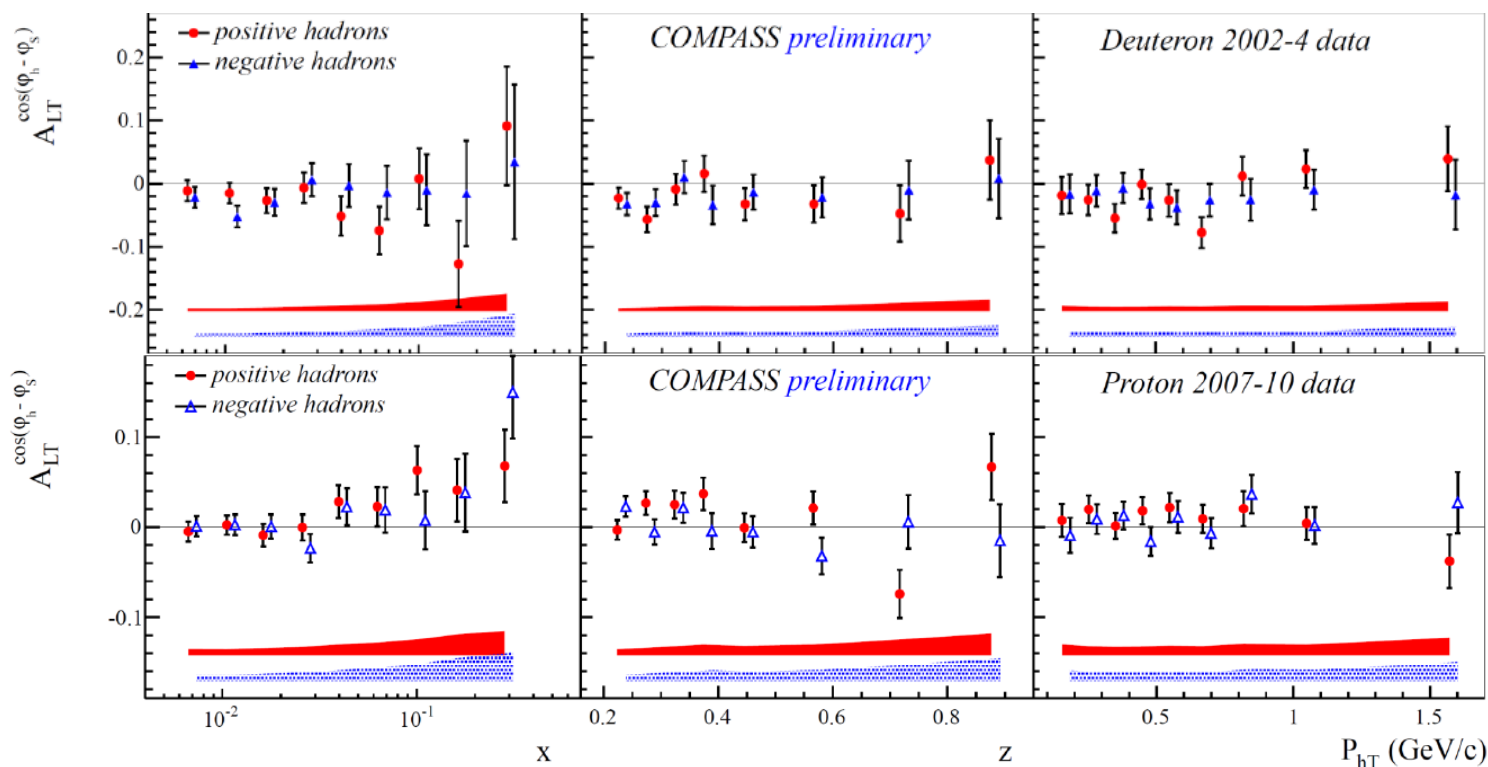
$$F_{LT}^{\cos(\phi_h - \phi_S)} \propto g_{1T} \otimes D_1$$

wormgear

Other transverse spin asymmetries

Probability of finding a longitudinally polarized quark inside a transversely polarized nucleon.

Double spin asymmetries, requiring both longitudinally polarized beam and transversely polarized target



Similar trend is present in HERMES preliminary results

Positive signal for π^- also seen by JLab E06-00, on neutron.

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} =$$

$$\frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right.$$

$$+ \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h}$$

$$+ S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right]$$

**3 independent modulations
in the hadron
azimuthal distribution**

$$+ S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right]$$

$$+ |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right.$$

$$+ \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)}$$

$$+ \left. \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right]$$

$$+ |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right.$$

$$\left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right\},$$

Sivers

Collins

twist-2 contribution

Remaining four can be
interpreted as twist-3
contributions

From [A. Bacchetta et al.](#),

[JHEP 0702:093,2007](#). e-Print: [hep-ph/0611265](#)



SIDIS cross section: unpolarized part

Also the azimuthal asymmetries in the unpolarized cross section give information on TMD effects.

$A_{\sin\phi_h}$ higher twist effect proportional to beam polarization
no clear interpretation in terms of PM

$$A_{\cos\phi_h}^{UU} = \frac{1}{Q} Cahn + \frac{1}{Q} BM \quad \text{Cahn effect} + \text{Boer-Mulders DF}$$

$$A_{\cos 2\phi_h}^{UU} = BM + \frac{1}{Q^2} Cahn \quad \text{Boer- Mulders} \times \text{Collins FF} + \text{Cahn effect}$$

Cahn effect

kinematical effect due to quark transverse momentum

$$\frac{d\sigma}{d\phi_h} \propto 1 - 4 \frac{\langle k_t^2 \rangle z P_t}{Q \langle P_t^2 \rangle} D_{\cos\phi_h}(y) \cos\phi_h + \dots$$

Boer-Mulders PDF

Correlation between the **quark transverse momentum**

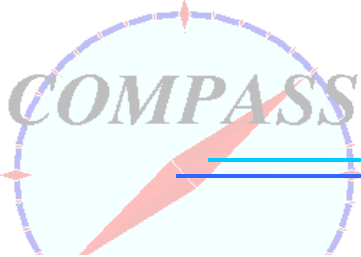
and the **quark spin**

in an **unpolarized nucleon**

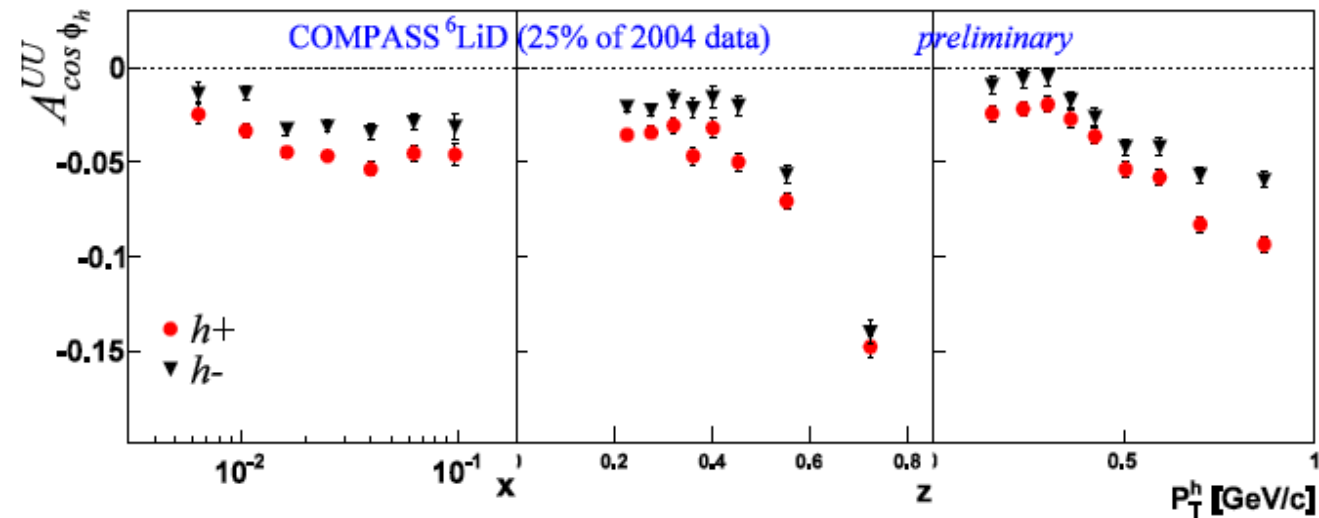
quark



pQCD contributions expected to be important for $p_T > 1 \text{ GeV}/c$

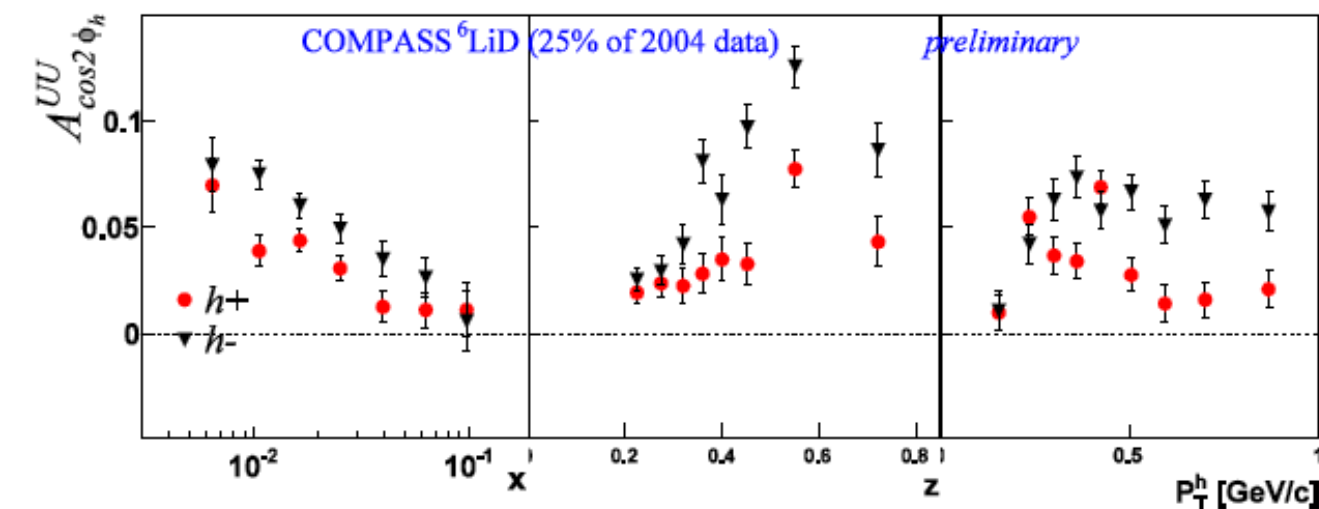


Unpolarized asymmetries, results on deuterium



$\cos\Phi$

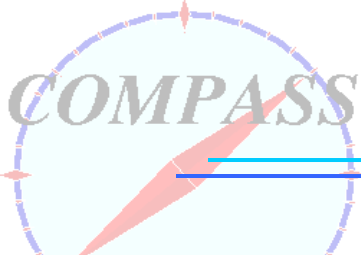
- Large signal over all the x range
- Strong z dependence, for $z > 0.5$



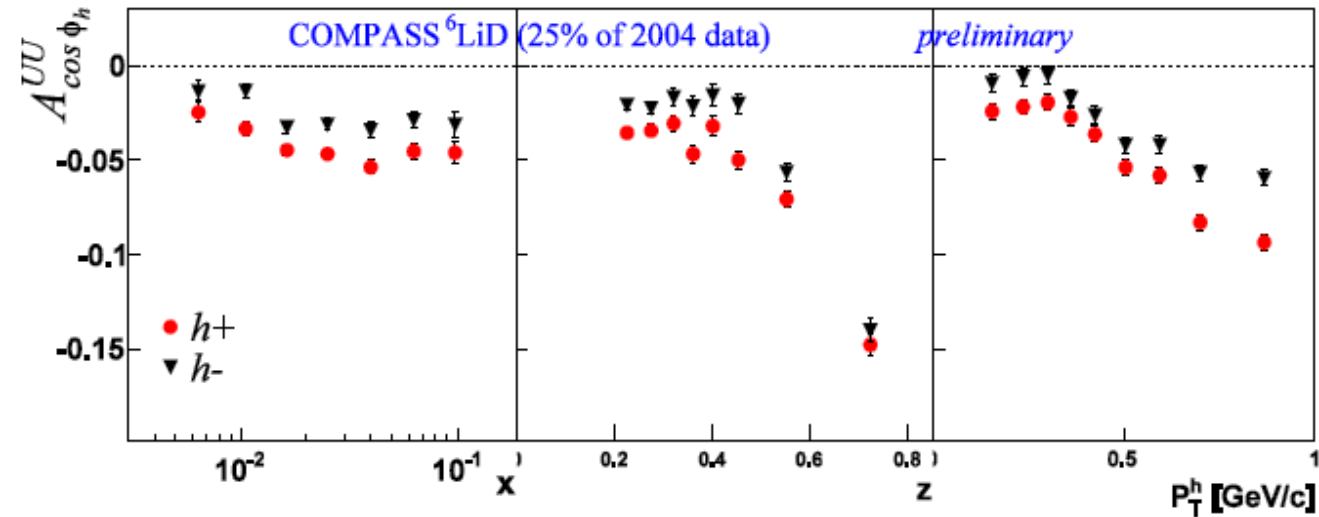
$\cos 2\Phi$

- Different for positive and negative hadrons
- Large signal at small x
- Strong dependence on x , z , and p_T , difficult to describe

$\text{sys} \approx 2 \cdot \text{stat}$

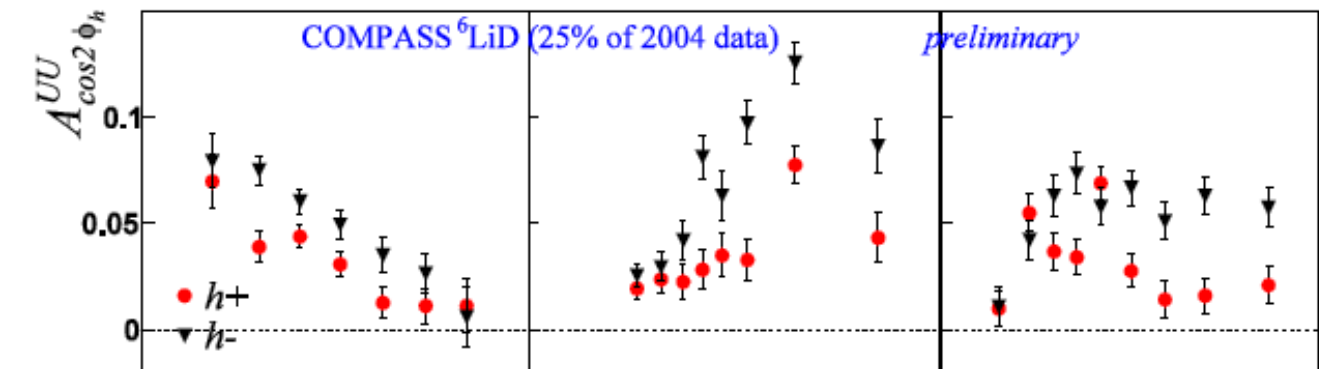


Unpolarized asymmetries, results on deuterium



$\cos\Phi$

- Large signal over all the x range
- Strong z dependence, for $z > 0.5$



$\cos 2\Phi$

- Different for positive and negative hadrons
- Large signal at small x
- Strong dependence on x , z , and p_T , difficult to describe

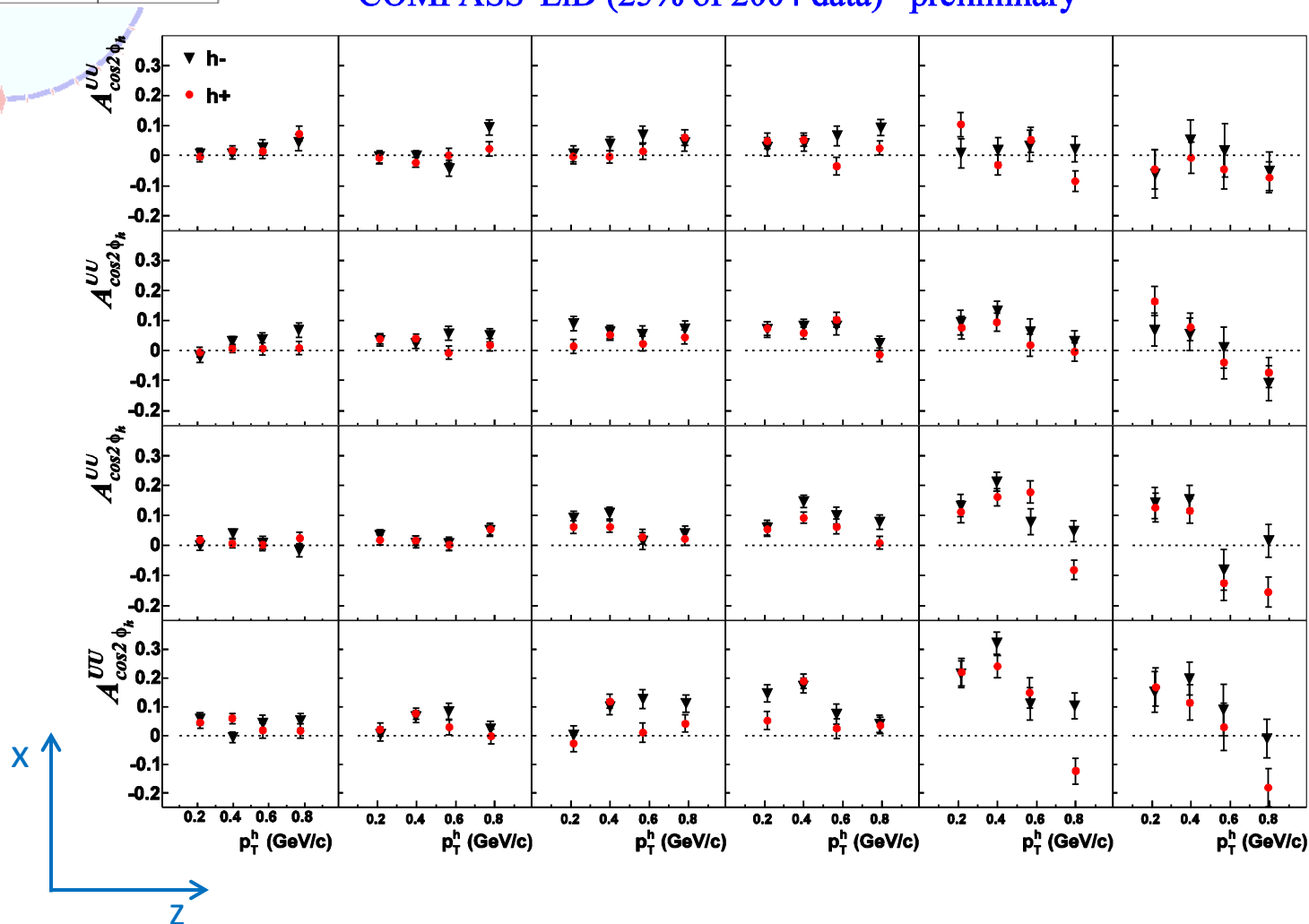
To investigate deeper the complicated and unexpected kinematical dependencies found, a multi-dimensional analysis has been done, binning simultaneously in x , z and p_T

→ interesting input for theory

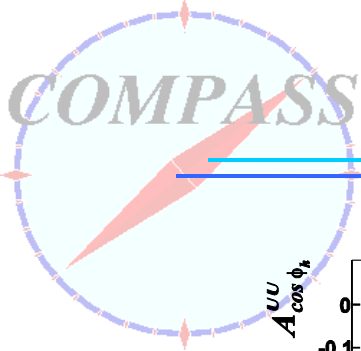
| x | P_T^h | z |
|---------------|------------|-------------|
| 0.003 - 0.012 | 0.1 - 0.3 | 0.2 - 0.25 |
| 0.012 - 0.02 | 0.3 - 0.5 | 0.25 - 0.32 |
| 0.02 - 0.038 | 0.5 - 0.64 | 0.32 - 0.40 |
| 0.038 - 0.13 | 0.64 - 1.0 | 0.40 - 0.55 |
| | | 0.55 - 0.70 |
| | | 0.70 - 0.85 |

cos2Φ asymmetries multi-dimensional analysis

COMPASS⁶LiD (25% of 2004 data) preliminary

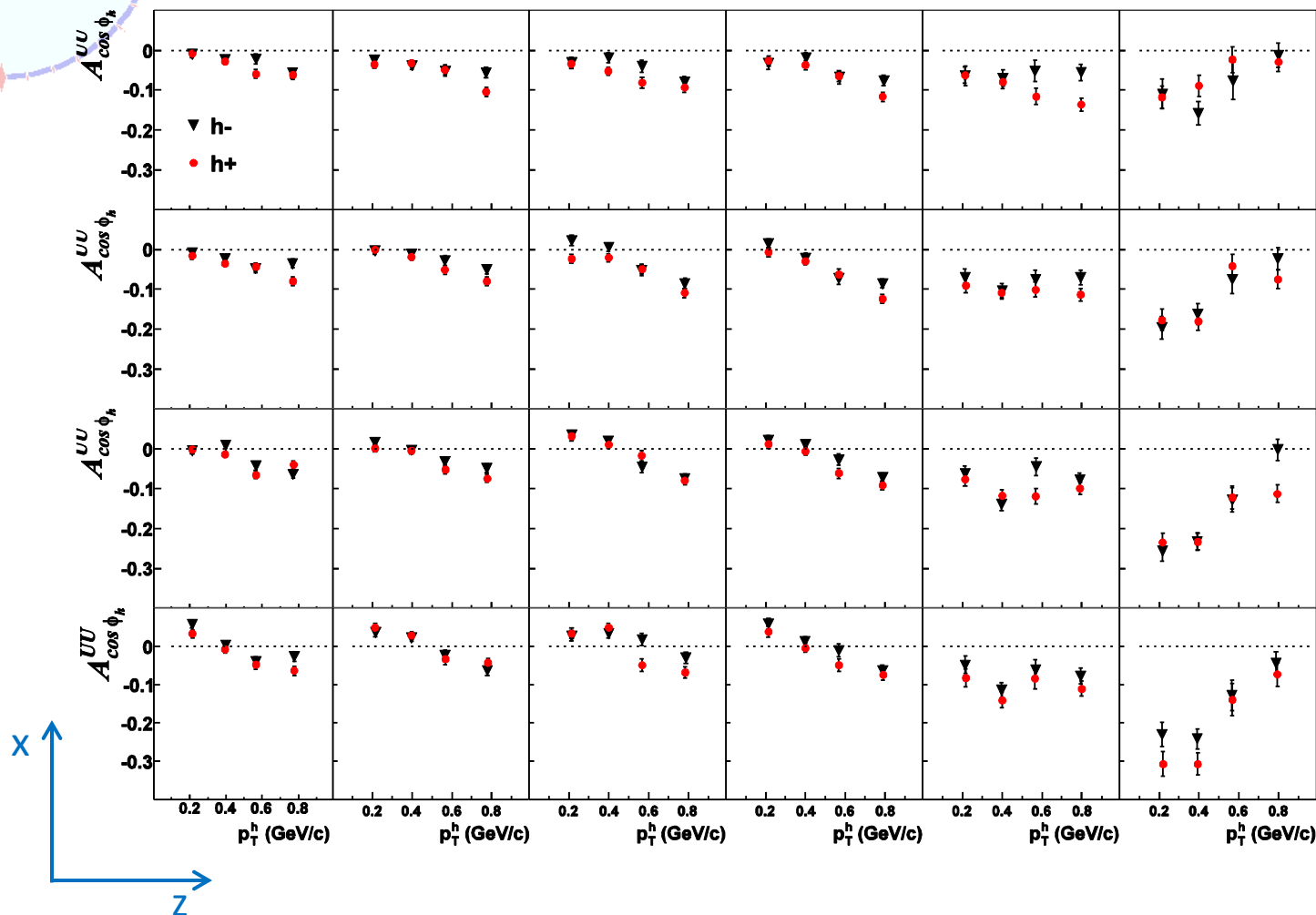


The p_T trend difficult to reproduce by models is there for large z and low x



cos Φ asymmetries multi-dimensional analysis

COMPASS ${}^6\text{LiD}$ (25% of 2004 data) preliminary



The p_T trend changes with z , and it is roughly the same over all the x range



COMPASS has investigated transverse spin and TMD effects using deuterium and proton targets.

Full set of results on **Collins and Sivers asymmetries**, on pions and kaons
interesting effects on p, to be investigated deeper with a multi-dimensional analysis

Results on many other channels available: other 6 transverse asymmetries, 2h asymmetries,
longitudinal azimuthal spin asymmetries

On a longer time scale, possible SIDIS measurements on p and d,
with different beam energies

Unpolarized asymmetries on d different from zero, showing complex
and interesting behavior in the kinematical variables.

New measurement of the unpolarized azimuthal asymmetries in parallel to DVCS
with a LH target measurement in 2012 and at COMPASS II starting from 2015