

Review of longitudinal spin physics

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IWHSS 2013

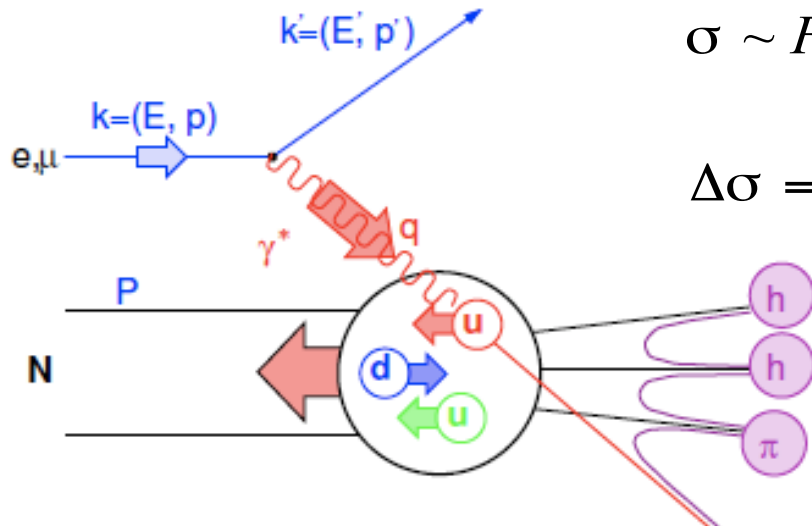
International Workshop on Hadron Structure and Spectroscopy 2013
Erlangen, Germany, 22-24 July

Contents

- Introduction: Nucleon spin structure
- Inclusive longitudinal spin asymmetries and g_1 structure function
- Semi inclusive asymmetries and flavour separations
- Strange quark sea polarisation
- Gluon polarisation
- Summary

- Spin is a fundamental degree of freedom originated from space-time symmetry quantum and relativistic object - survives high-energy limit
- Spin plays a critical role in determining the basic structure of fundamental interactions
- Spin provides a unique opportunity to probe the inner structure of a composite system such as the nucleon

After many years still the driving question for QCD spin physics is where the nucleon spin comes from?



$$\sigma \sim F_1(x) = \frac{1}{2} \sum_i e_q^2 q_i(x) \quad F_2(x) \approx 2xF_1$$

$$\Delta\sigma = \overset{\uparrow\uparrow}{\sigma} - \overset{\uparrow\downarrow}{\sigma} \sim g_1(x) = \frac{1}{2} \sum_i e_q^2 \Delta q_i(x) \quad g_2$$

$$\Delta q(x) = q(x)^+ - q(x)^-$$

$$q(x) = q(x)^+ + q(x)^-$$

+ quark $\uparrow\uparrow$ nucleon

- quark $\uparrow\downarrow$ nucleon

$$A_{meas} = \frac{1}{fP_T P_B} \left(\frac{N^{\uparrow\uparrow} - N^{\uparrow\downarrow}}{N^{\uparrow\uparrow} + N^{\uparrow\downarrow}} \right) \approx DA_1$$

$$A_1(x, Q^2) = \frac{\sigma_{\uparrow\downarrow} - \sigma_{\uparrow\uparrow}}{\sigma_{\uparrow\downarrow} + \sigma_{\uparrow\uparrow}} \approx \frac{\sum_q e_q^2 \Delta q(x, Q^2)}{\sum_q e_q^2 q(x, Q^2)} = \frac{g_1(x, Q^2)}{F_1(x, Q^2)}$$

Sum rules

first moment $\Gamma_1 = \int g_1(x) dx$

Bjorken s.r. $\Gamma_1^p - \Gamma_1^n = \frac{g_A}{6g_V} C_1^{NS}$

Ellis-Jaffe s.r. $\Gamma_1^{p,n} = \left(\pm a_3 + \frac{a_8}{3} \right) \frac{C_1^{NS}}{12} + a_0 \frac{C_1^S}{9}$

$a_3, a_8, g_{A,V}$ measured in weak β decays (+Su(3)_f)

$C_1^{S,NS}$ calculable in pQCD

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$C_1^{S,NS}$

calculable in pQCD

$a_0 = \Delta\Sigma = \Delta u + \Delta d + \Delta s$

Quark contribution to nucleon helicity

$$\Gamma_1^N(Q^2) = \frac{1}{9} \left(1 - \frac{\alpha_s(Q^2)}{\pi} + O(\alpha_s^2) \right) \left(a_0(Q^2) + \frac{1}{4} a_8 \right)$$

$$a_{0|Q_0^2=3(GeV/c)^2} = 0.35 \pm 0.03(stat) \pm 0.05(syst)$$

from Y. Goto *et al.*, PRD62 (2000)
 034017:
 (SU(3)_f assumed for weak decays)
 $a_8 = 0.585 \pm 0.025$

QCD NLO

$$\Gamma_1^N(Q^2) = \frac{1}{9} C_1^S(Q^2) \hat{a}_0 + \frac{1}{36} C_1^{NS}(Q^2) a_8$$

beyond NLO

$$\hat{a}_{0|Q^2 \rightarrow \infty} = 0.33 \pm 0.03(stat) \pm 0.05(syst)$$

C_1 calculated behind 3 loops app.
 S.A.Larin *et al.*, Phys.Lett.B404(1997)153

quark contribution to the proton spin: $\sim 1/3$

$$(\Delta s + \Delta \bar{s}) = \frac{1}{3} (\hat{a}_0 - a_8) = -0.08 \pm 0.01(stat) \pm 0.02(syst)$$

comment: a_8 (under SU(3) symmetry) maybe reduced (model)

“Switching on” spin leads us to two complications:

$$q^- \sim \psi \frac{1}{2} (1 - \gamma_5) \gamma_\mu \psi \quad q^+ \sim \psi \frac{1}{2} (1 + \gamma_5) \gamma_\mu \psi$$

$$\Delta q = u^+ - u^- \sim \psi \frac{1}{2} \gamma_\mu \gamma_5 \psi$$

- axial current is not conserved due to Adler-Bell-Jackiw triangle anomaly
- there is no local, gauge invariant dimension-3 axial operator for gluons

$\overline{\text{MS}}$ scheme - a_0 depends on the scale

AB scheme - a_0 does not depend on the scale but now

$$a_0 = \Delta \Sigma - \frac{3}{2} \frac{\alpha_s}{\pi} \Delta G$$

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anomaly gives a possible interpretation of the measured value of a_0

- if gluon polarization is large enough then E-J sum rule can be restored and quark contribution to the nucleon spin is significant as expected in simple QPM

Semi-inclusive asymmetry

Inclusive asymmetry:
$$A_1(x, Q^2) = \frac{\sigma_{\uparrow\downarrow} - \sigma_{\uparrow\uparrow}}{\sigma_{\uparrow\downarrow} + \sigma_{\uparrow\uparrow}} \approx \frac{\sum_q e_q^2 \Delta q(x, Q^2)}{\sum_q e_q^2 q(x, Q^2)} = \frac{g_1(x, Q^2)}{F_1(x, Q^2)}$$

Semi-inclusive asymmetry:
$$A_1^h(x, z, Q^2) = \frac{\sigma_{\uparrow\downarrow}^h - \sigma_{\uparrow\uparrow}^h}{\sigma_{\uparrow\downarrow}^h + \sigma_{\uparrow\uparrow}^h} \approx \frac{\sum_q e_q^2 \Delta q(x, Q^2) D_q^h(z, Q^2)}{\sum_q e_q^2 q(x, Q^2) D_q^h(z, Q^2)}$$

inclusive and semi-inclusive asymmetry measured on proton and neutron or deuteron allows to full flavour separation

2002-2006 deuteron data taking in COMPASS (160 GeV): Phys. Lett. B 647 (2007) 8

2007: proton data; DIS events - $Q^2 > 1$ (GeV/c)² (160 GeV): Phys. Lett. B 690 (2010) 466–472

New: 2011 proton data 200 GeV

Difference asymmetry

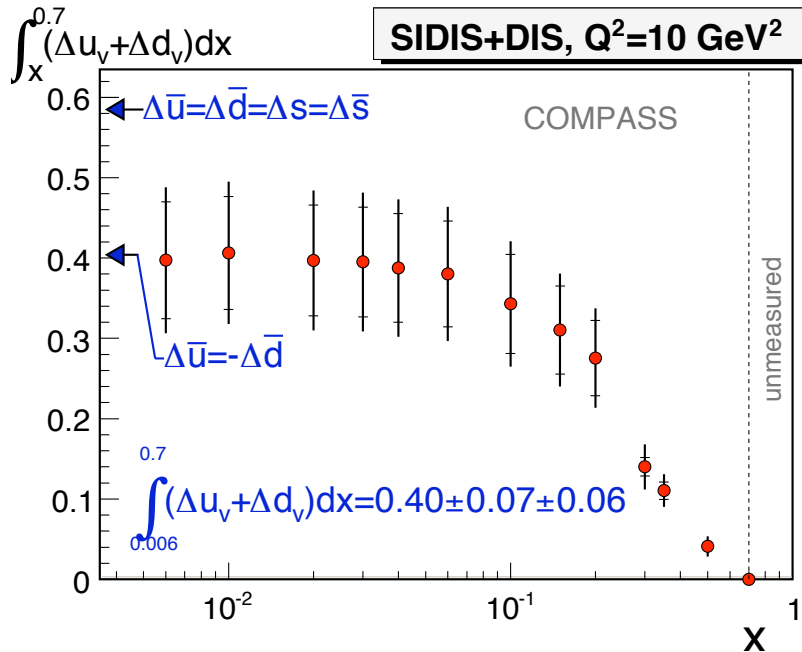
Idea: Phys.Lett.B230(1989)141,
 SMC:Phys.Lett.B369(1996)93,
 COMPASS: Phys.Lett.B660(2008)458

$$A_d^{\pi^+ - \pi^-}(x) = A_d^{K^+ - K^-}(x) = \frac{\Delta u_v(x) + \Delta d_v(x)}{u_v(x) + d_v(x)}$$

$$A^{+-} = \frac{(\sigma_{\uparrow\downarrow}^{h^+} - \sigma_{\uparrow\downarrow}^{h^-}) - (\sigma_{\uparrow\uparrow}^{h^+} - \sigma_{\uparrow\uparrow}^{h^-})}{(\sigma_{\uparrow\downarrow}^{h^+} - \sigma_{\uparrow\downarrow}^{h^-}) + (\sigma_{\uparrow\uparrow}^{h^+} - \sigma_{\uparrow\uparrow}^{h^-})}$$

Only valence quarks!

Fragmentation functions cancel out in LO and under the assumption of independent fragmentation.



$$\Delta \bar{u} = \Delta \bar{d} = \Delta s = \Delta \bar{s}$$

symmetric

$$\Delta \bar{u} = -\Delta \bar{d}$$

asymmetric

scenario

$$\Gamma_v = \int_0^1 (\Delta u_v(x) + \Delta d_v(x)) dx$$

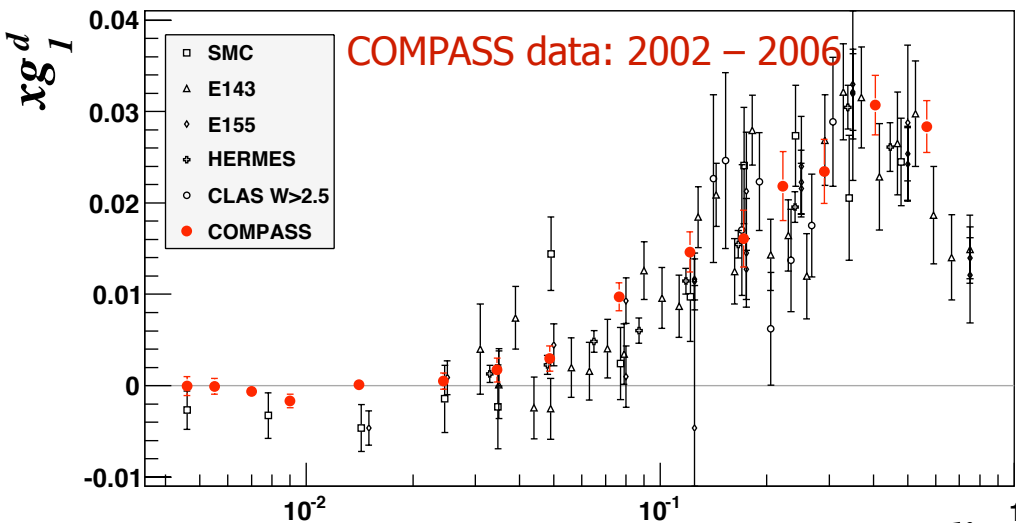
$$\Delta \bar{u} + \Delta \bar{d} = 3\Gamma_1^N - \frac{1}{2}\Gamma_v + \frac{1}{12}a_8 = (\Delta s + \Delta \bar{s}) + \frac{1}{2}(a_8 - \Gamma_v)$$

Experiments:

- **Inclusive spin-dependent DIS**
EMC, SMC, COMPASS & COMPASS II (CERN),
E142,E143,E154,E156,
HERMES (DESY) ,JLAB-Hall,A,B(CLAS)
- **Semi-inclusive DIS:** SMC, COMPASS,HERMES
- **Polarized pp collision** RHIC-PHENIX & STAR, BRAHMS (Brookhaven)
- **ee:** BELLE (KEK) (**Fragmentation functions**)

JLAB-12, EIC, eLHC

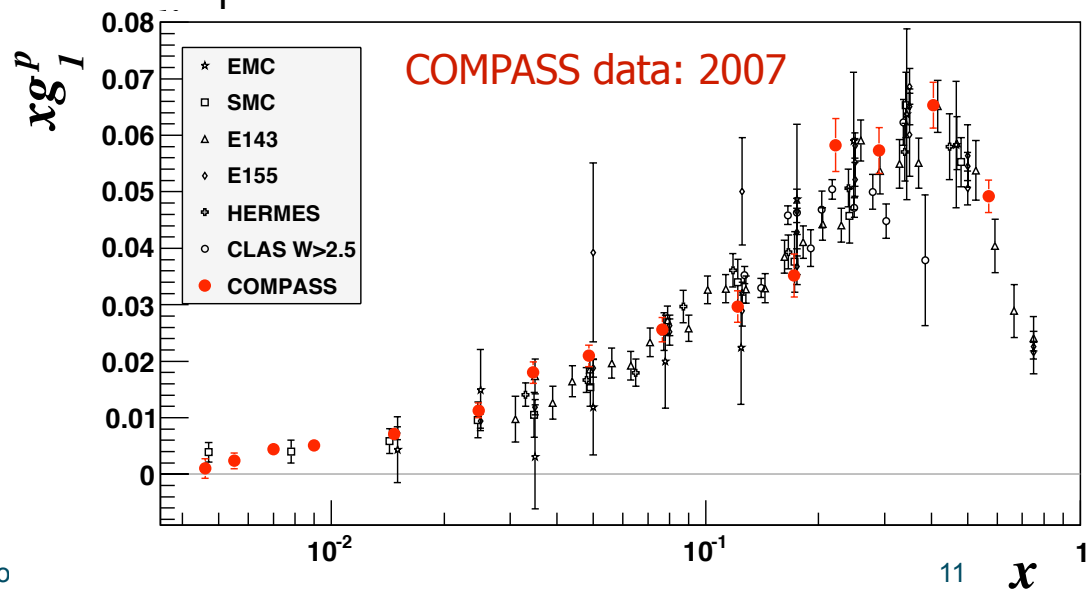
COMPASS data



Phys. Lett. B 690 (2010) 466–472

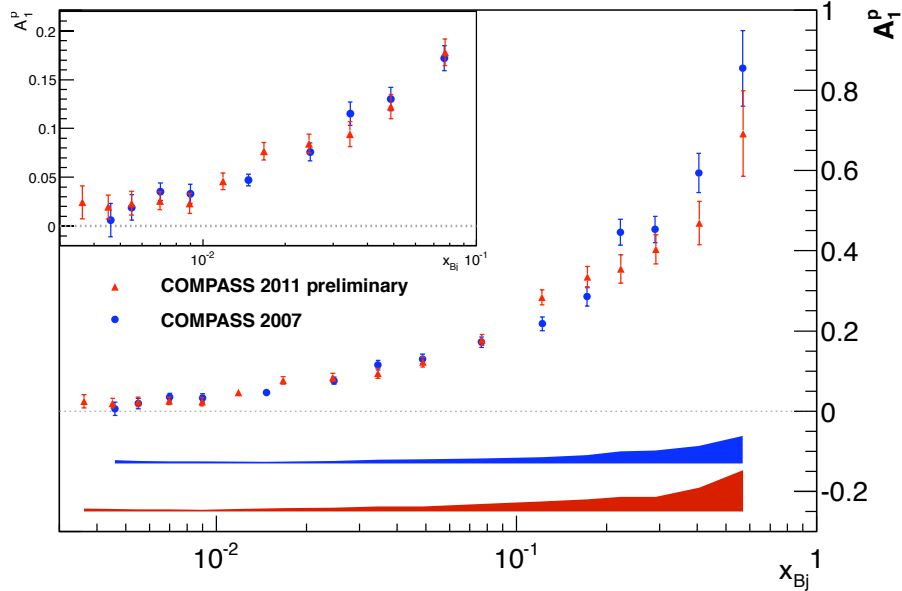
Phys. Lett. B 647 (2007) 8

Very good agreement between experiments. Points evolved to $Q^2 = 3 \text{ (GeV/c)}^2$



Proton g_1 structure function

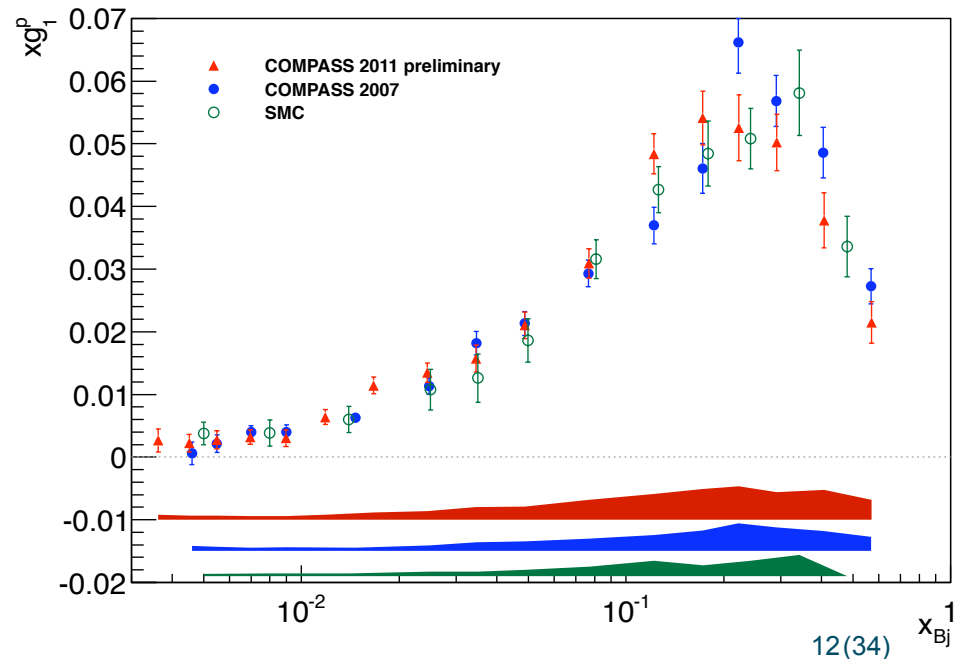
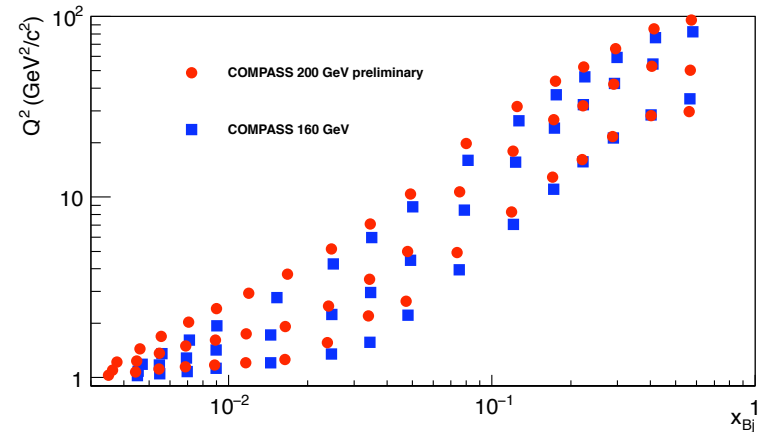
COMPASS data 2011, 200 GeV



SMC parameterization of F_2 -
 Phys.Rev.D 55 (1998) 112001
 R - Phys. Lett. B 647 (2007) 330

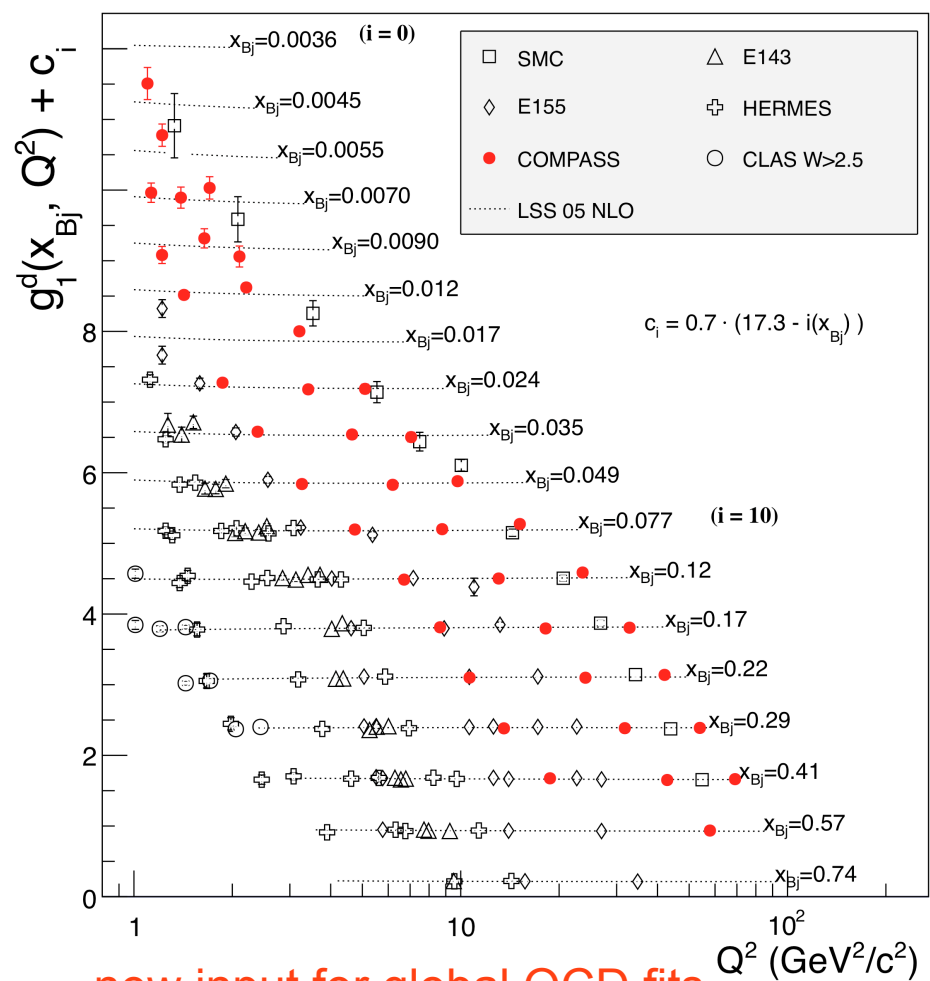
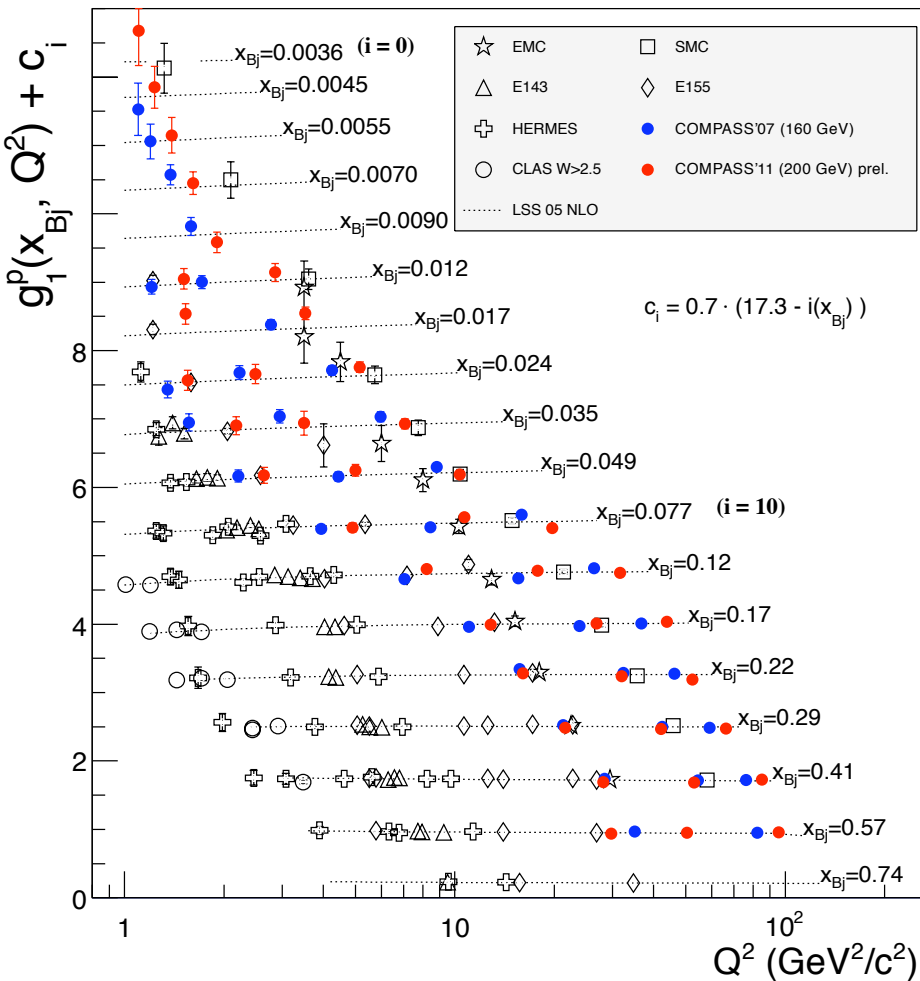
Lower x_B

data taken to balance the amount of deuteron data:
 to increase precision of the Bjorken s.r.determination,
 to better constrain the strange quark polarisation



World data for proton and deuteron g_1 structure function

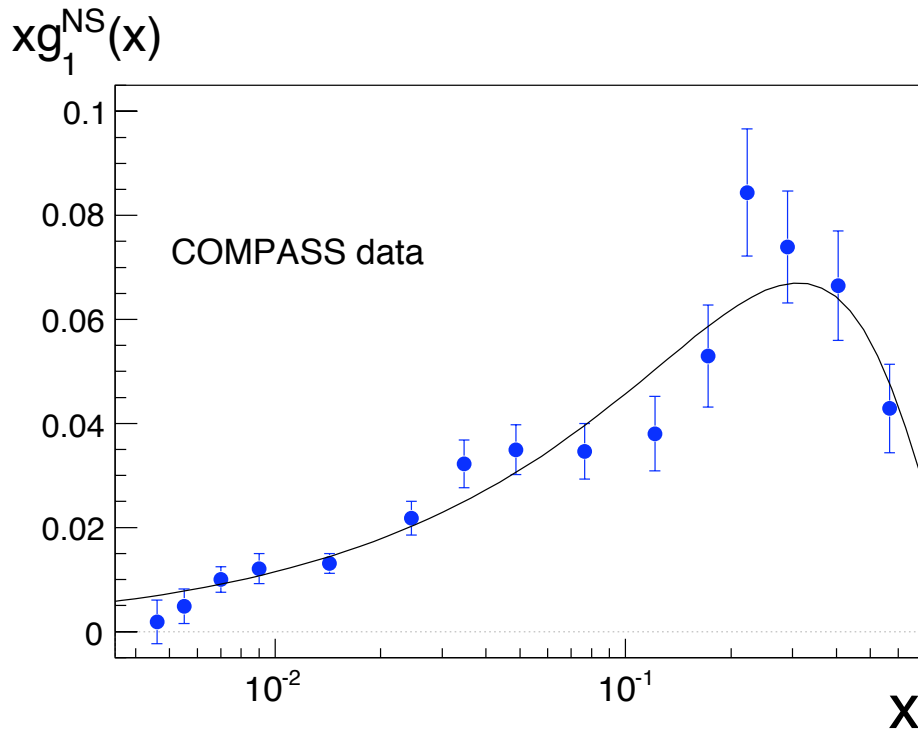
COMPASS proton data 2011@200 GeV included



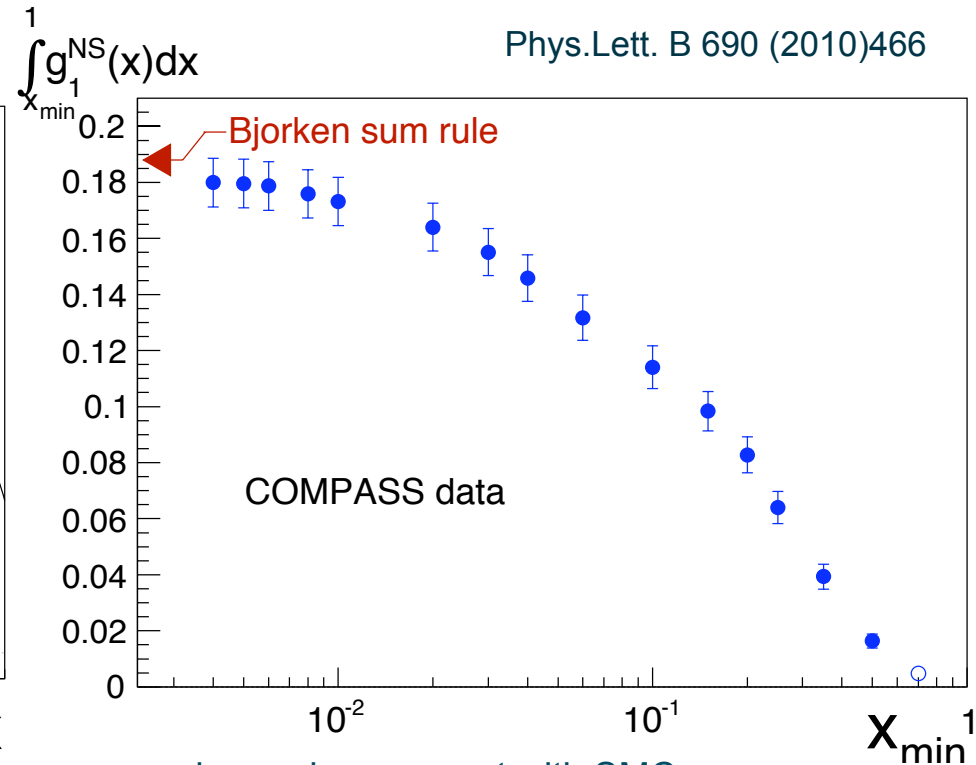
new input for global QCD fits

Bjorken sum rule

COMPASS data, SMC, HERMES



x range	Γ_1^{NS}
0 – 0.004	0.0098
0.004 – 0.7	$0.175 \pm 0.009 \pm 0.015$
0.7 – 1.0	0.0048
0 – 1	$0.190 \pm 0.009 \pm 0.015$



in good agreement with SMC:
 0.198 ± 0.023 ($Q^2 = 10$).

truncated moment from $x=0.021$:

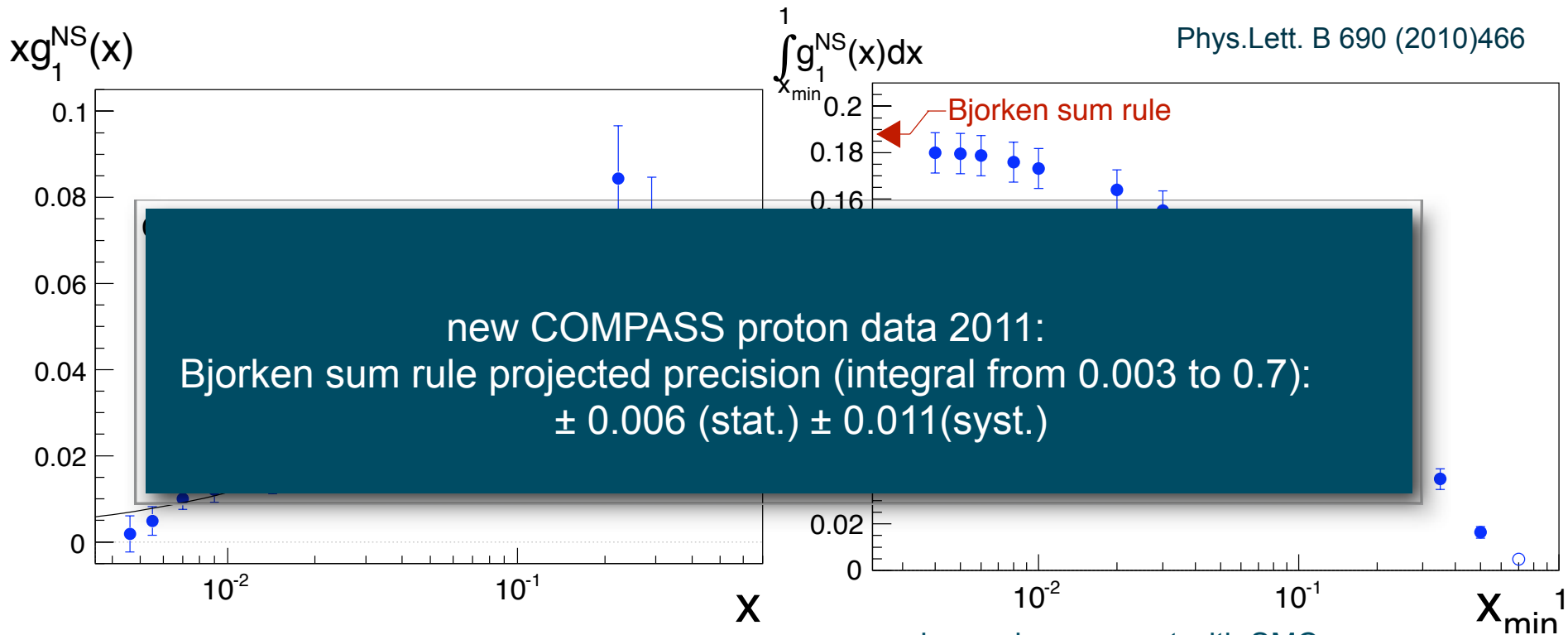
$0.1583 \pm 0.0085 \pm 0.014$

in good agreement with HERMES:

$0.1484 \pm 0.0055 \pm 0.016$

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COMPASS & HERMES data

- SIDIS

$$A_1^h = \frac{\sum e_q^2 [\Delta q(x) \int D_q^h(z) dz + \Delta \bar{q}(x) \int D_{\bar{q}}^h(z) dz]}{\sum e_q^2 [q(x) \int D_q^h(z) dz + \bar{q}(x) \int D_{\bar{q}}^h(z) dz]}$$

- $D_q^h \neq D_{\bar{q}}^h$
 yields quark and antiquark separation

- measured:

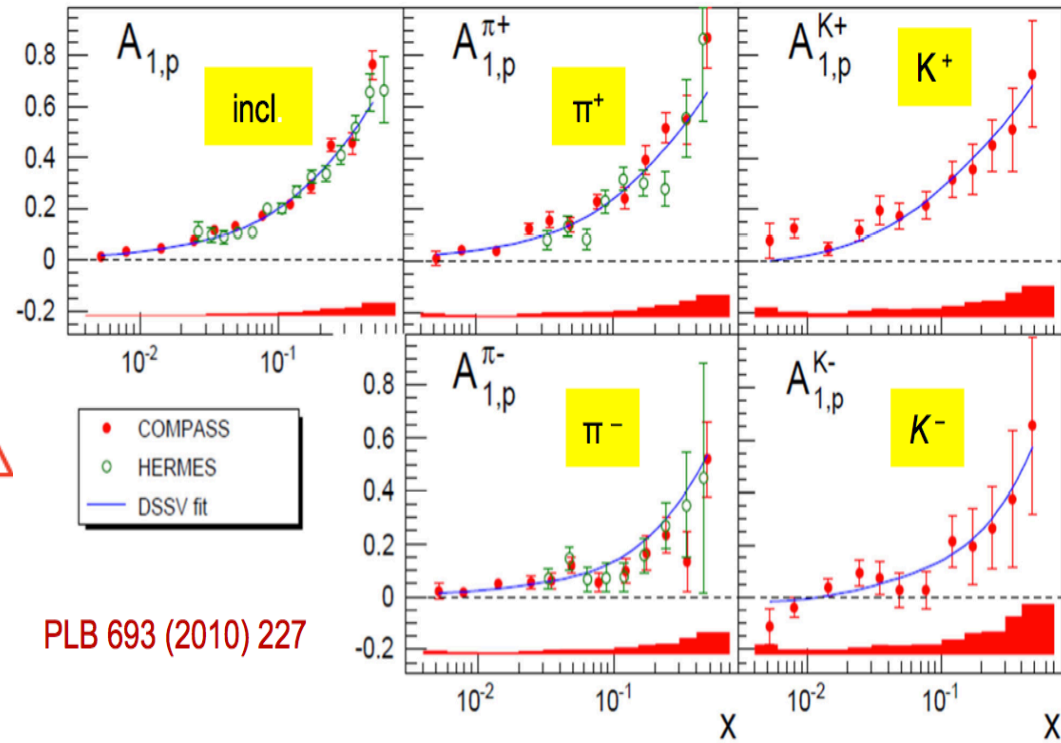
$$A_1^d, A_{1d}^{K^\pm}, A_{1d}^{\pi^\pm}, A_1^p, A_{1p}^{K^\pm}, A_{1p}^{\pi^\pm}$$

- determined: $\Delta u, \Delta \bar{u}, \Delta d, \Delta \bar{d}, \Delta s \equiv \Delta$

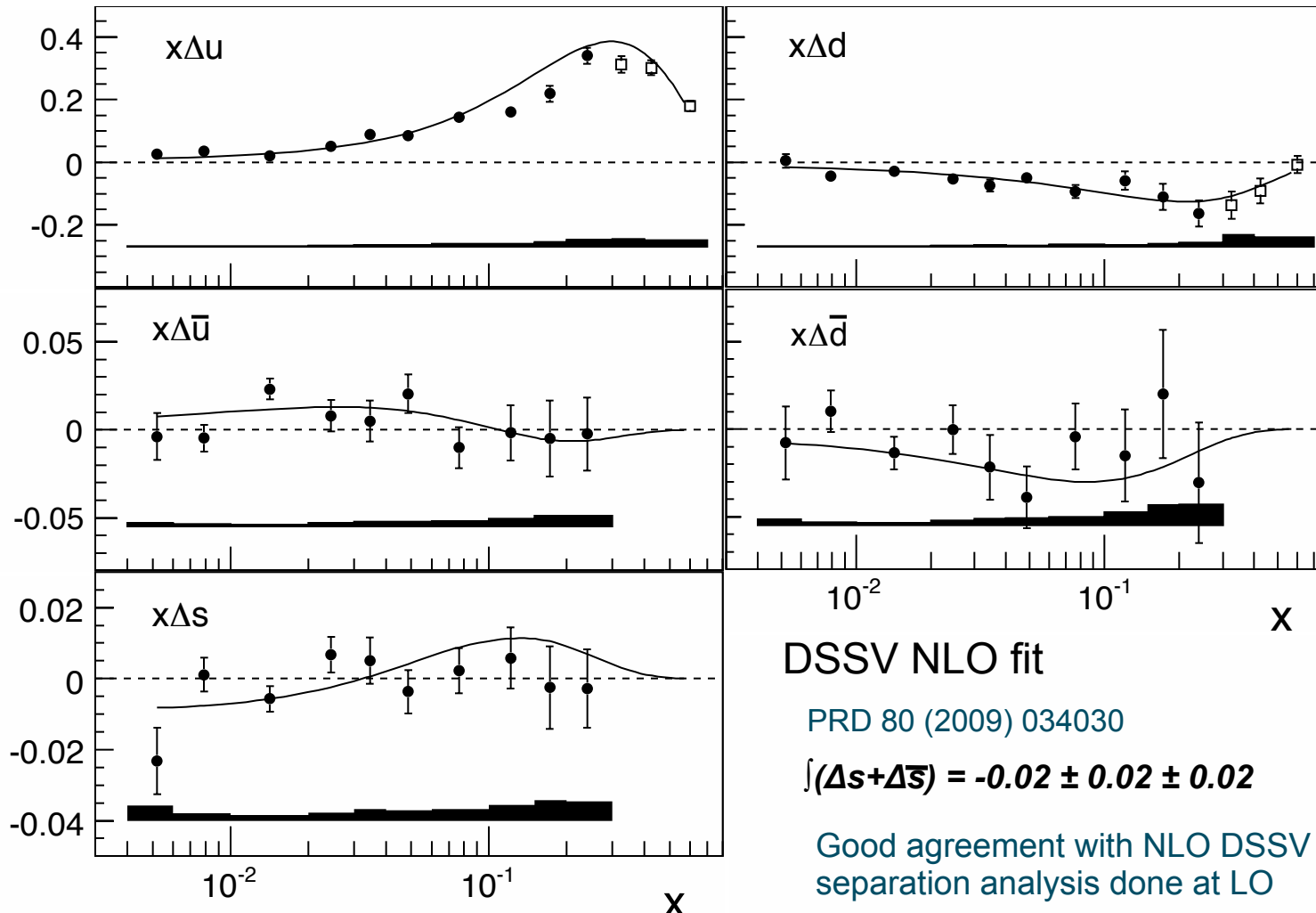
- system of linear equations in LO

- input: MRST04 unpolarised PDFs,
 DSS parametrisation of FFs
 (e^+e^- , DIS, hadron-hadron)

PRL 101 (2008) 072001; PRD 80 (2009) 034030



COMPASS data



DSSV NLO fit

PRD 80 (2009) 034030

$$\int (\Delta s + \Delta \bar{s}) = -0.02 \pm 0.02 \pm 0.02$$

Good agreement with NLO DSSV - qualitative,
 separation analysis done at LO

COMPASS data

Phys.Lett.B 693 (2010) 227

$\Delta\bar{u} \geq \Delta\bar{d}$?

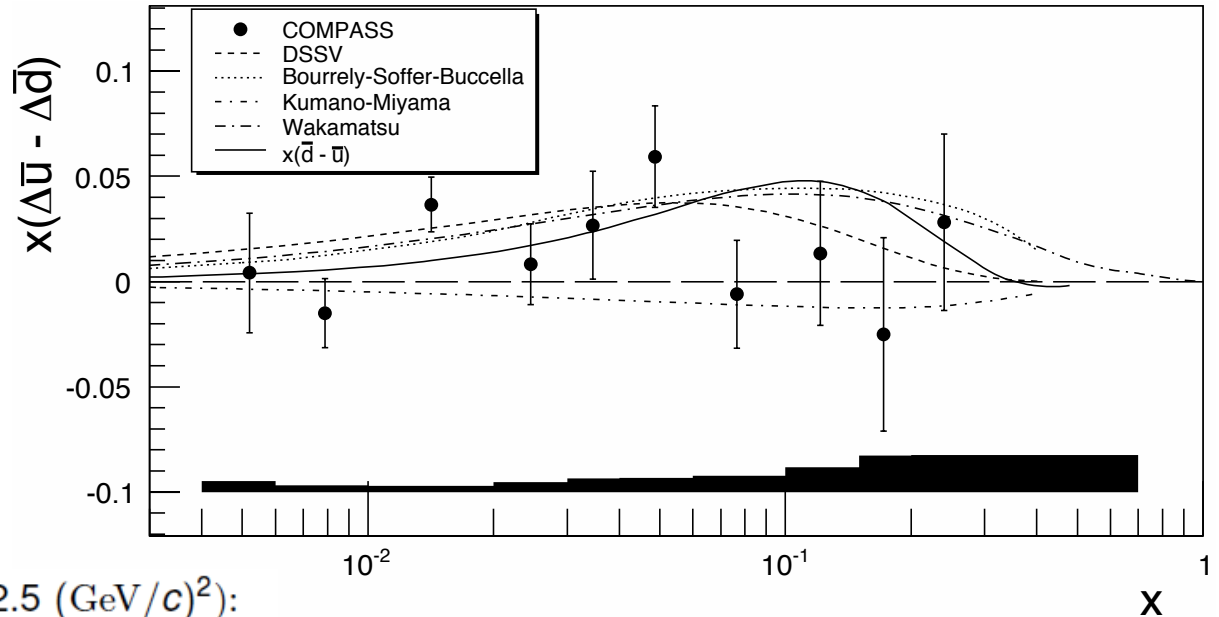
Phys. Rev. Lett. 18, (1967) 1174

Pauli exclusion principle, Δ resonance etc.

NMC, E866

unpolarized asymmetry:

$$\int_0^1 (\bar{d} - \bar{u}) dx = 0.118 \pm 0.012$$



HERMES ($Q^2 = 2.5 \text{ (GeV/c)}^2$):

$$\int_{0.023}^{0.3} (\Delta\bar{u} - \Delta\bar{d}) dx = 0.048 \pm 0.057(\text{stat.}) \pm 0.028(\text{syst.})$$

COMPASS ($Q^2 = 3 \text{ (GeV/c)}^2$):

$$\int_{0.004}^{0.3} (\Delta\bar{u} - \Delta\bar{d}) dx = 0.06 \pm 0.04 (\text{stat.}) \pm 0.02 (\text{syst.})$$

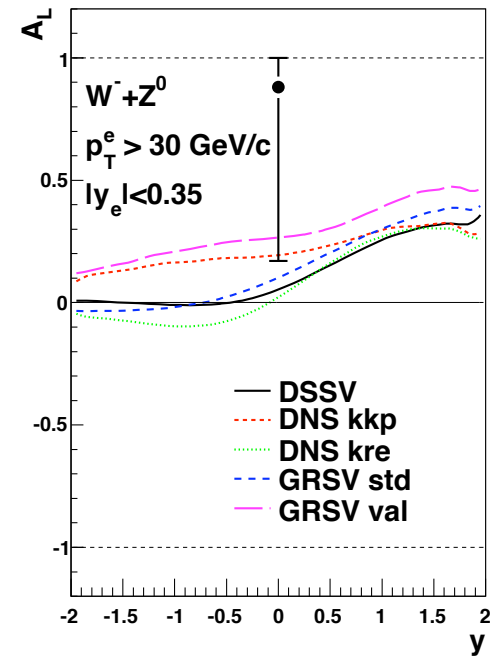
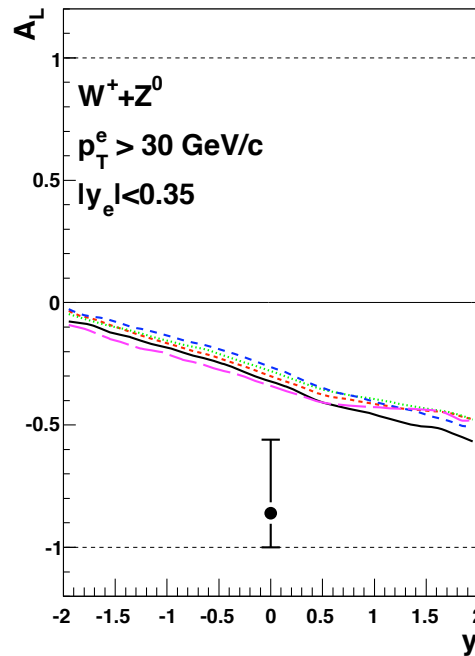
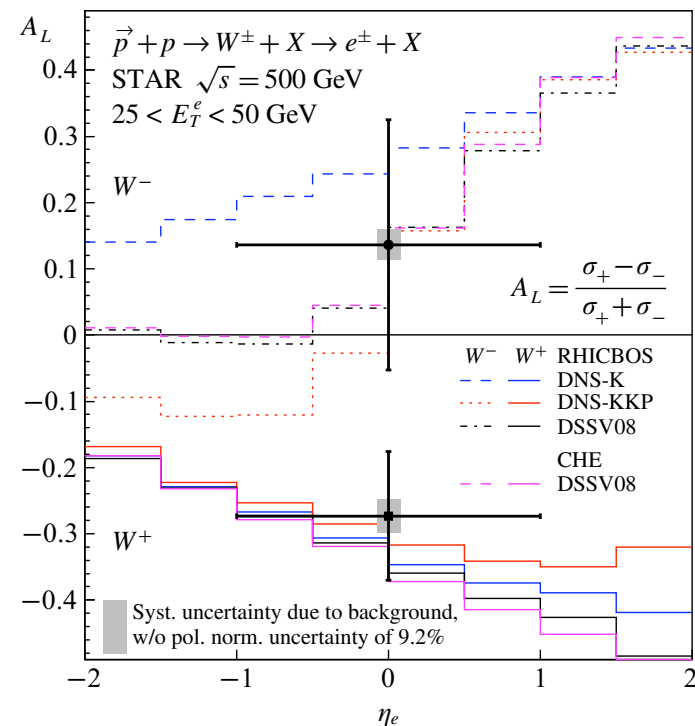
The data disfavour models predicting $\Delta\bar{u} - \Delta\bar{d} \gg \bar{d} - \bar{u}$

RHIC W data

$$A_L^{W^+} = \frac{\Delta u(x_1)\bar{d}(x_2) - \Delta\bar{d}(x_1)u(x_2)}{u(x_1)\bar{d}(x_2) + \bar{d}(x_1)u(x_2)},$$

$$A_L^{W^+} \rightarrow \Delta u/u$$

$$A_L^{W^-} \rightarrow \Delta\bar{d}/\bar{d}$$

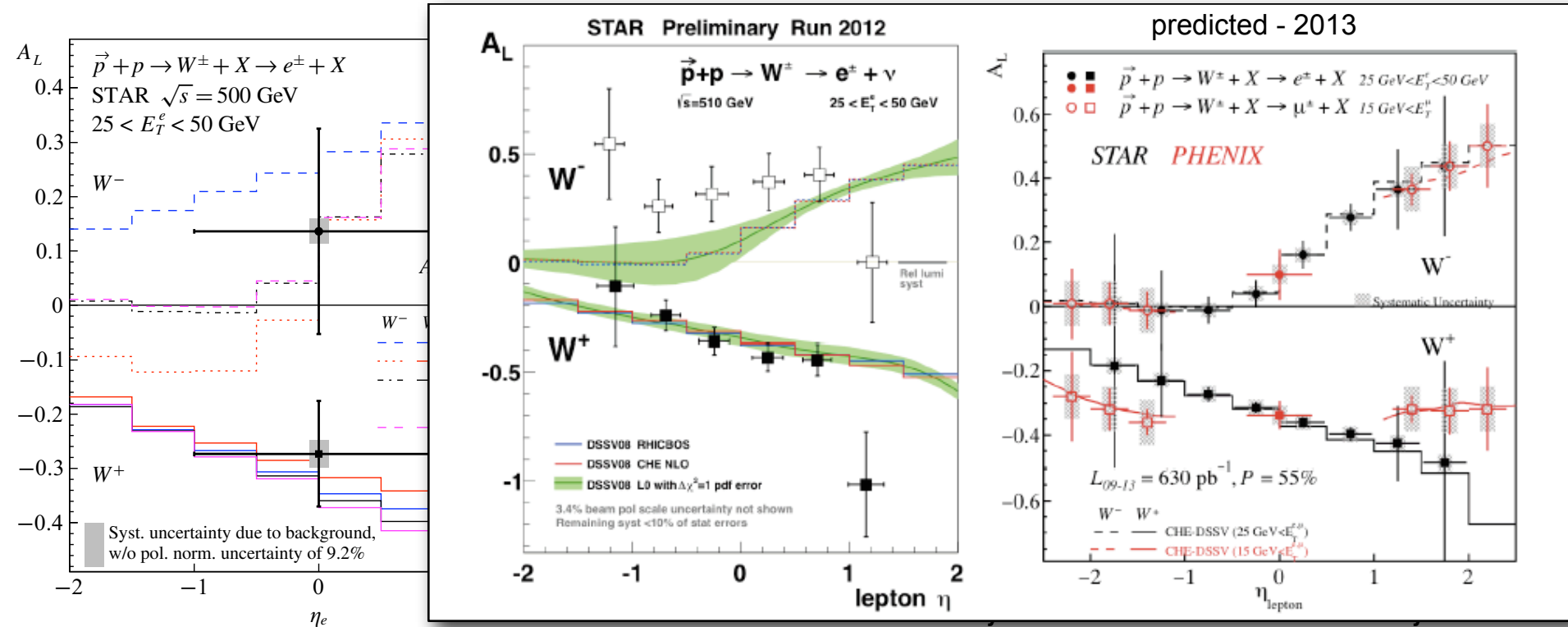


RHIC W data

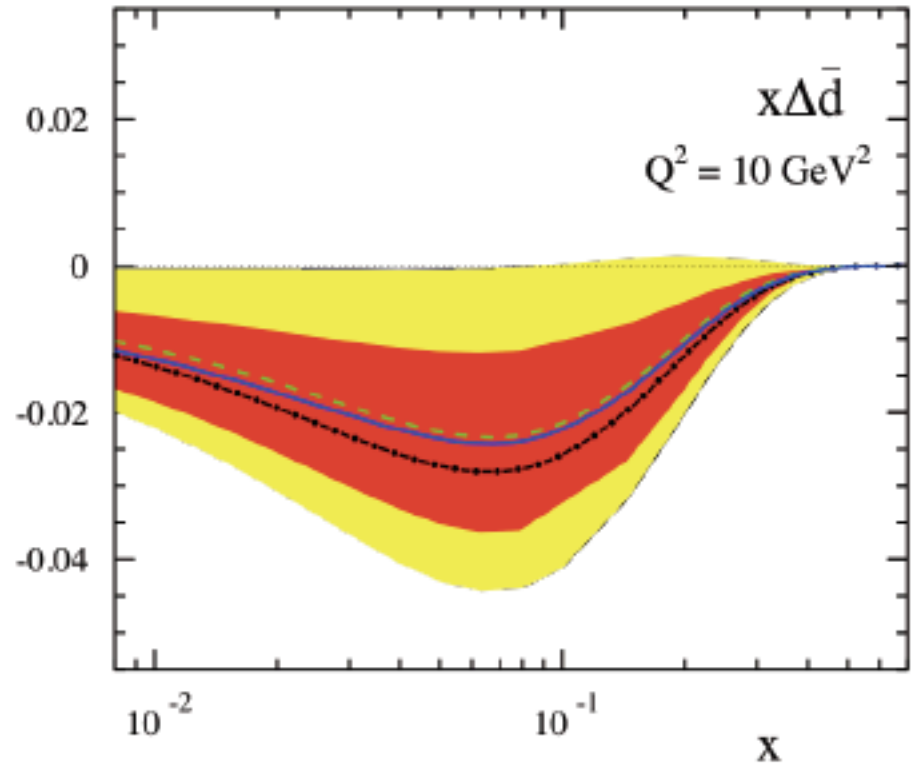
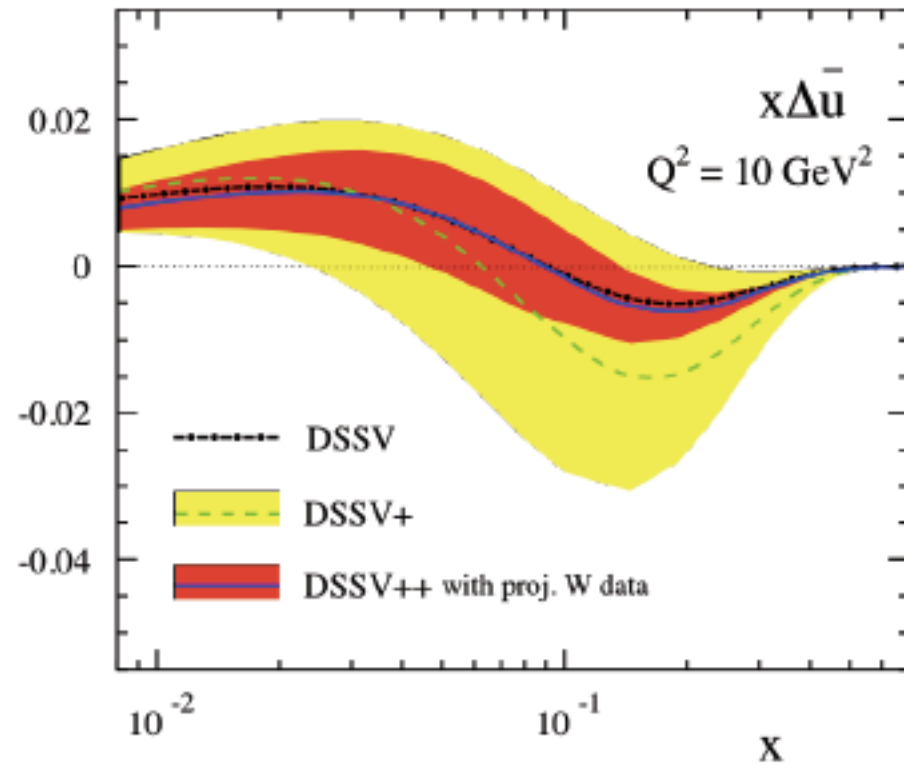
$$A_L^{W^\pm} = \frac{\Delta u(x_1)\bar{d}(x_2) - \Delta\bar{d}(x_1)u(x_2)}{u(x_1)\bar{d}(x_2) + \bar{d}(x_1)u(x_2)},$$

$$A_L^{W^+} \rightarrow \Delta u/u$$

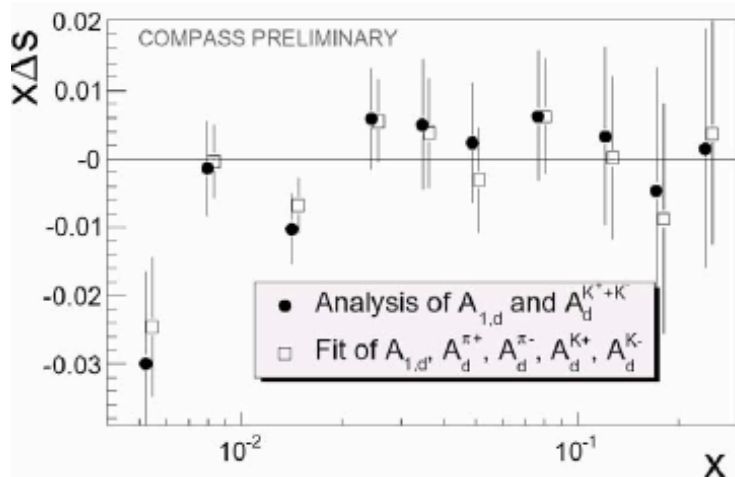
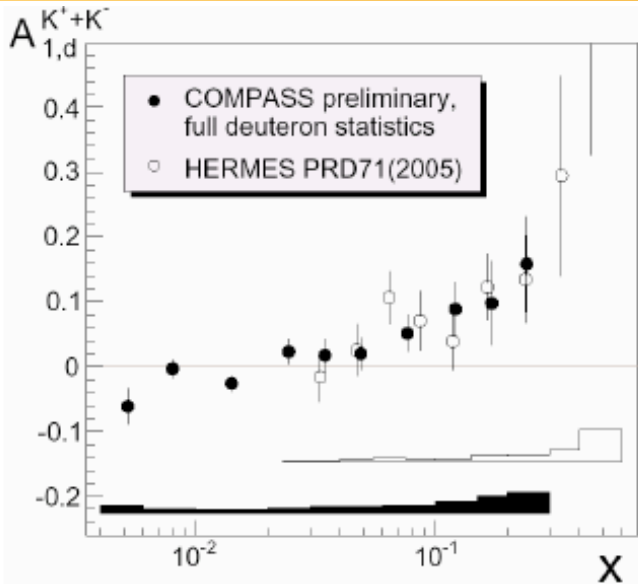
$$A_L^{W^-} \rightarrow \Delta\bar{d}/\bar{d}$$



RHIC W data



COMPASS data



$$\frac{\Delta s}{s} = A_1^d + \left(A_1^{K^+K^-} - A_1^d \right) \frac{Q/s + \alpha}{\alpha - 0.8}$$

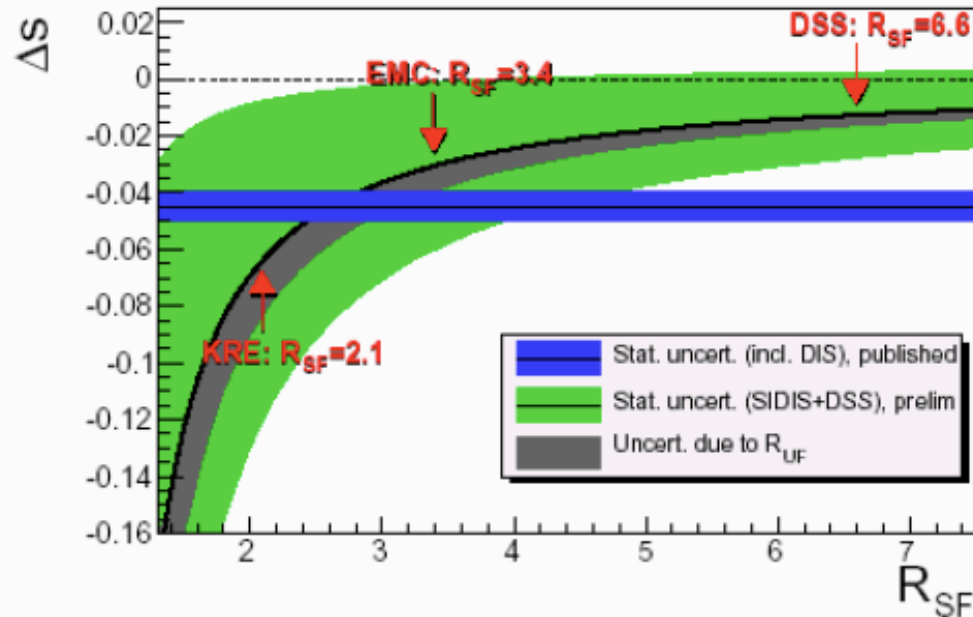
$$\alpha = \frac{2R_{UF} + 2R_{SF}}{3R_{UF} + 2} \quad Q = u + \bar{u} + d + \bar{d}$$

$$R_{UF} = \frac{\int D_d^{K^+}(z) dz}{\int D_u^{K^+}(z) dz} \quad R_{SF} = \frac{\int D_{\bar{s}}^{K^+}(z) dz}{\int D_u^{K^+}(z) dz}$$

if $A_1^d = A^{K^+K^-} \Rightarrow \Delta s \geq 0$, insensitive to FFs

if $A^{K^+K^-} < 0$ (at low x) $\Rightarrow \Delta s < 0$

COMPASS data



Phys.Lett. B 680 (2009)217-114

$$R_{SF} = \frac{\int D_s^{K^+}(z) dz}{\int D_u^{K^+}(z) dz}$$

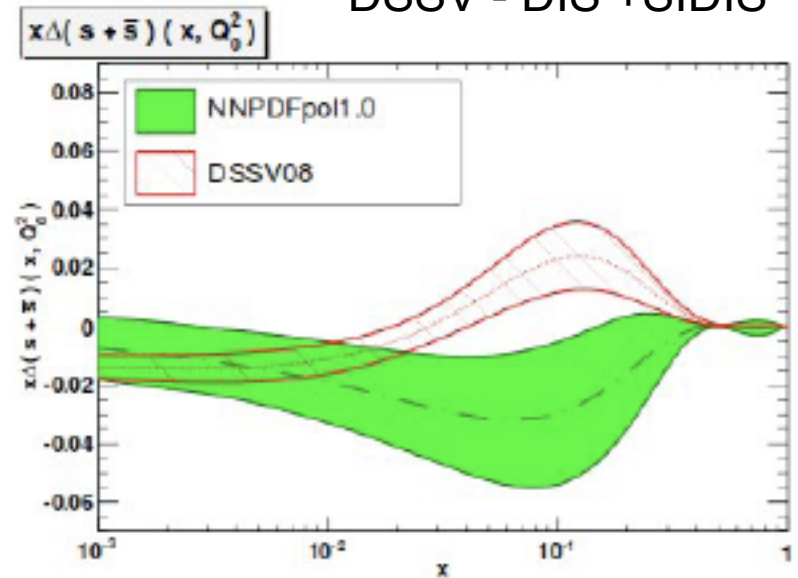
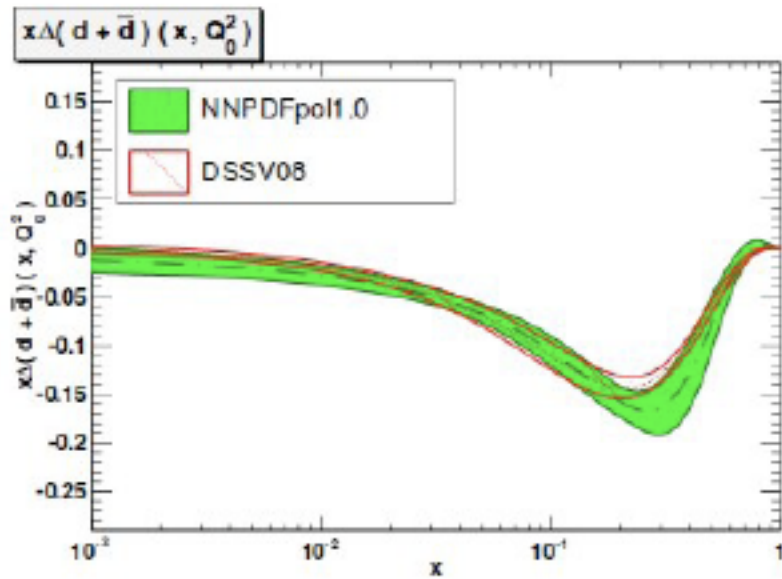
$$R_{UF} = \frac{\int D_d^{K^+}(z) dz}{\int D_u^{K^+}(z) dz}$$

- R_{UF} fixed at 0.14 from the DSS fragmentation functions
- Large statistical uncertainty due to R_{SF} , slight dependence on R_{UF}
- If $R_{SF} > 5$ $\Delta s(\text{SIDIS}) > \Delta s(\text{DIS})$ and $\Delta s < 0$ for $x < 0.004$
- If $R_{SF} < 4$: A^K becomes insensitive to Δs

$$\Delta s \text{ (inclusive)} = -0.045 \pm 0.005 \pm 0.010$$

in LO pQCD

$$\Delta s \text{ (SIDIS)} = -0.01 \pm 0.01 \pm 0.01$$

Δs puzzle

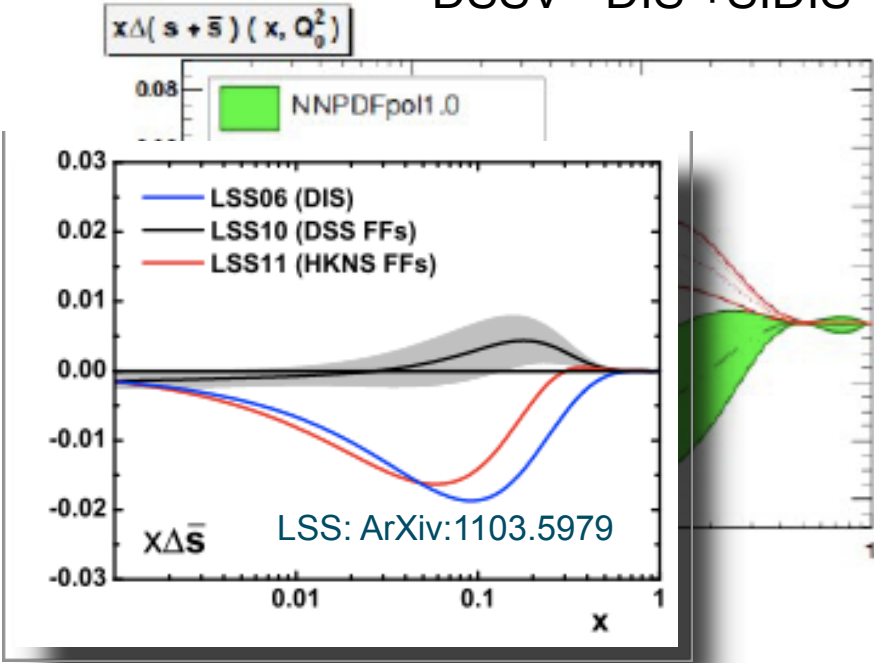
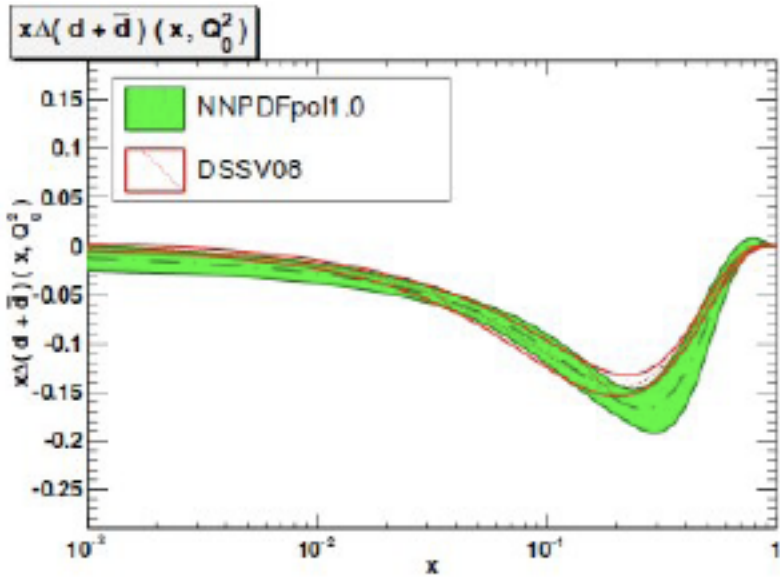
NNPDF - only DIS
DSSV - DIS +SIDIS

NNPDF, R.D.Ball et al. arXiv: 1303.7236

Δs puzzle

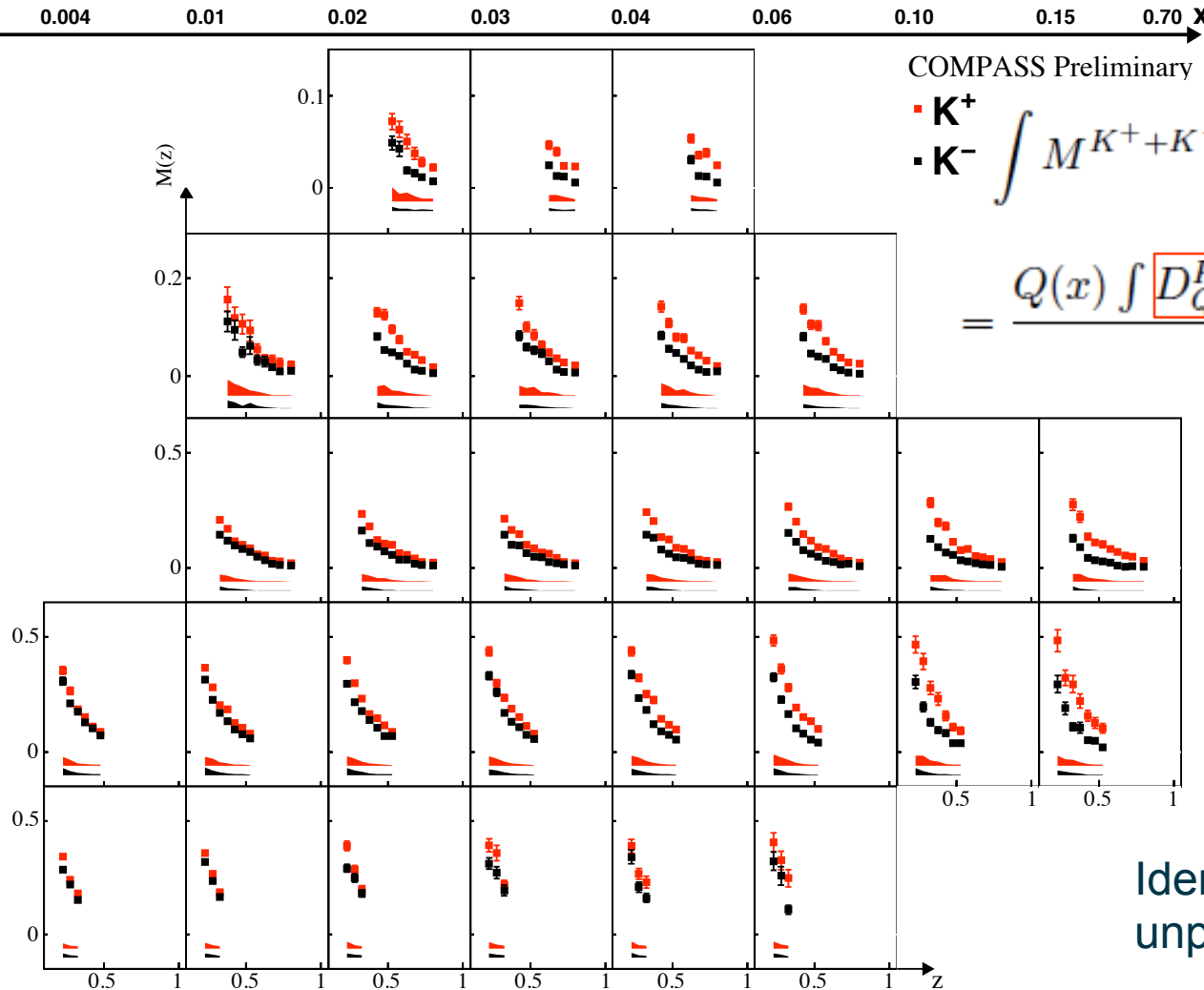
The lesson:
 Important piece of information:
 precise measurement of FFs

NNPDF - only DIS
 DSSV - DIS +SIDIS



NNPDF, R.D.Ball et al. arXiv: 1303.7236

Hadron multiplicities

 z dependence in x & y bins

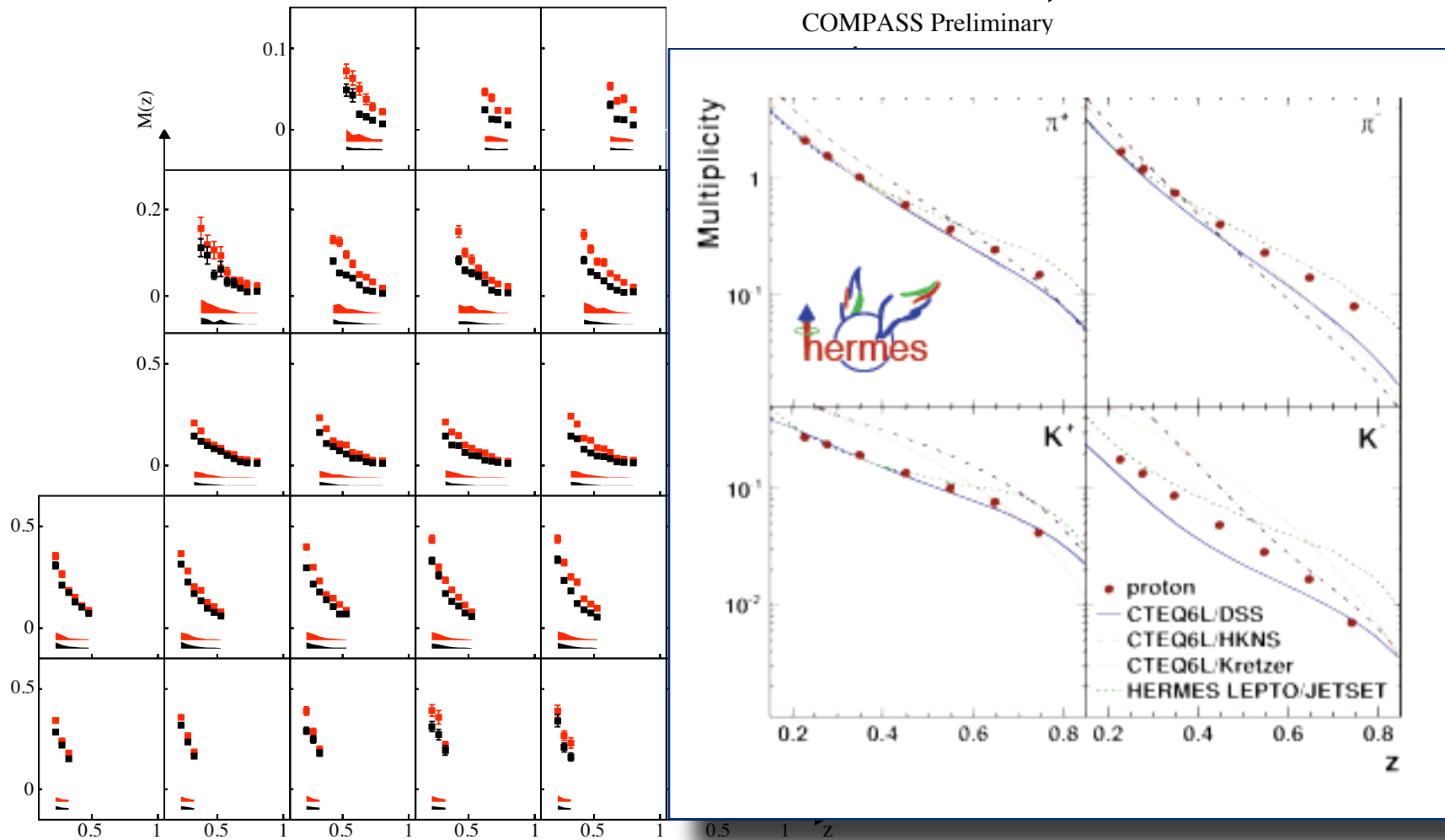
Identified K 's and Π 's;
unpolarised - well known

Hadron multiplicities

z dependence in x & y bins

0.004 0.01 0.02 0.03 0.04 0.06 0.10 0.15 0.70 x

COMPASS Preliminary



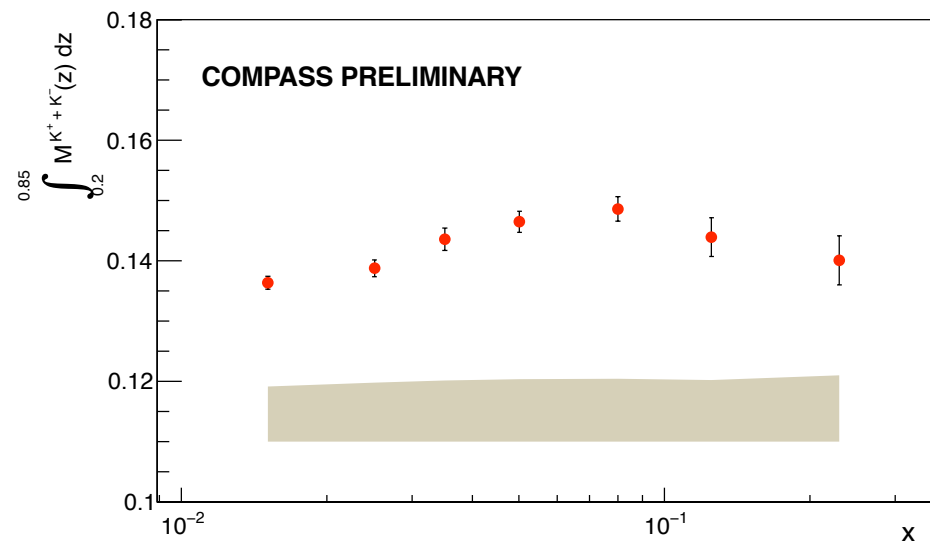
$$\frac{dK}{dx} = \dots$$

Hadron multiplicities

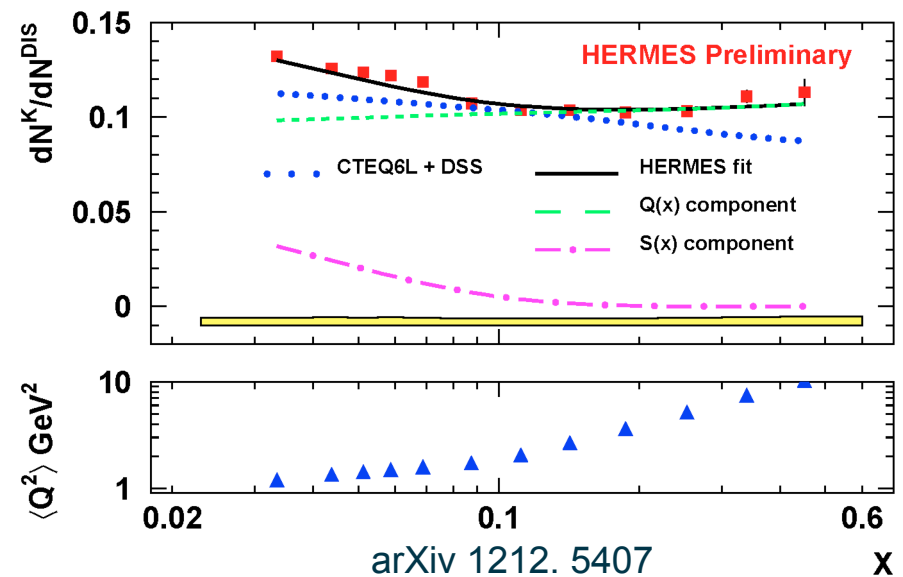
$$2S(x) \ll 5Q(x)$$

$$\int_{0.2}^{0.85} M^{K^+ + K^-}(x, z) dz = \frac{1}{5} \left(\int D_Q^K(z) dz + \frac{S(x)}{Q(x)} \int D_S^K(z) dz \right)$$

Directly related to strange PDF and FF of strange quark into K,
should help in understanding strange sea



small x dependence



Gluon polarization

QCD evolution

► QCD fits - Idea:

$$g_1(x, Q^2) = \frac{1}{2} \langle e^2 \rangle \left[C_q^S \otimes \Delta\Sigma + C_q^{NS} \otimes \Delta q^{NS} + 2n_f C_G \otimes \Delta G \right]$$

DGLAP eqs

$$\frac{d}{dt} \Delta q^{NS} = \frac{\alpha_s(t)}{2\pi} P_{qq}^{NS} \otimes \Delta q^{NS} \quad t = \log\left(\frac{Q^2}{\Lambda^2}\right)$$

$$\frac{d}{dt} \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix} = \frac{\alpha_s(t)}{2\pi} \begin{pmatrix} P_{qq}^S & 2n_f P_{qG}^S \\ P_{Gq}^S & P_{GG}^S \end{pmatrix} \otimes \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix}$$

Initial parameterization in x at fixed Q^2

$$(\Delta\Sigma, \Delta q_s, \Delta q_8, \Delta G) = \eta \frac{x^\alpha (1-x)^\beta (1+\gamma x)}{\int_0^1 x^\alpha (1-x)^\beta (1+\gamma x) dx}$$

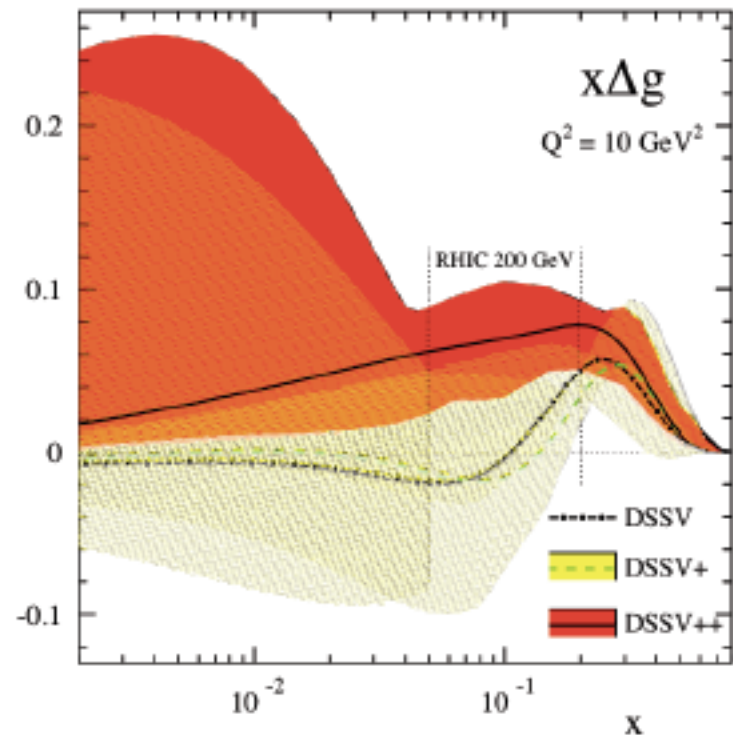
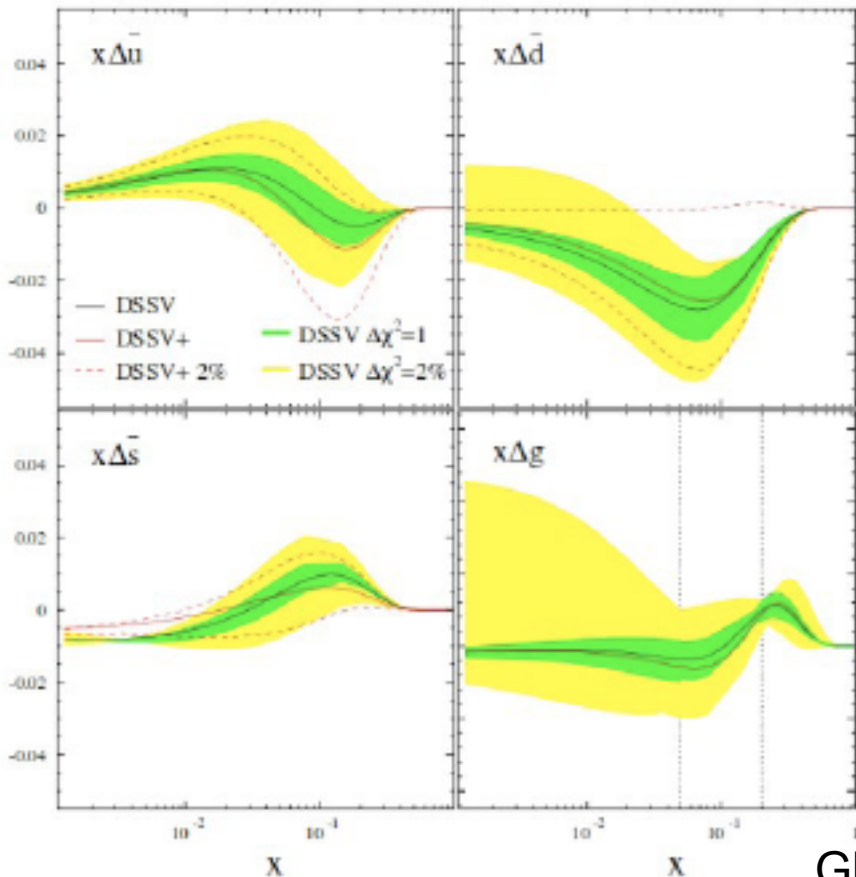
Minimalization procedure

$$\chi^2 = \sum_{i=1}^N \frac{[g_1^{calc}(x, Q^2) - g_1^{exp}(x, Q^2)]^2}{[\sigma_{stat}^{exp}(x, Q^2)]^2}$$

Gluon polarization from QCD fits

QCD fits - DSSV

Global fits (DSSV/DSSV+/DSSV++) include: spin-dependent DIS data, SIDIS data with identified π and K , and proton-proton data



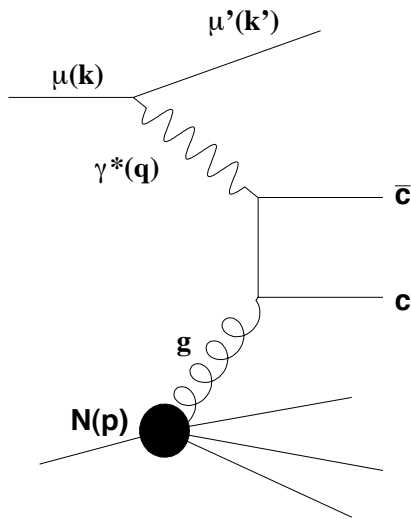
Latest PHENIX and STAR data included in DSSV++

Gluons are not constrained by inclusive DIS data

arXiv: 1108.1713

arXiv: 1304.0079

Direct gluon polarisation determination via tagging PGF events



$$\sigma^{PGF} = G \otimes \hat{\sigma}^{PGF} \otimes H$$

$$\Delta\sigma^{PGF} = \Delta G \otimes \Delta\hat{\sigma}^{PGF} \otimes H$$

To select PGF process two methods are used:

- **Open-charm D meson production:**
charm quark pairs produced in PGF, “clean” channel however with huge combinatorial background, low statistics but analysis less MC dependent
- **High transverse momentum hadron pairs production:**
light quark pairs produced, high statistics but physical background; strongly model and MC dependent analysis, requires a very good agreement between data and MC

$$A \approx \frac{\Delta G}{G}(\bar{x}_G) < \hat{a}_{LL}^{PGF} >_G$$

from MC

signal asymmetry from data

COMPASS@Low Q^2 - Phys. Lett. B 633 (2006) 25-32COMPASS@high Q^2 - Phys. Lett. B 718 (2013) 922

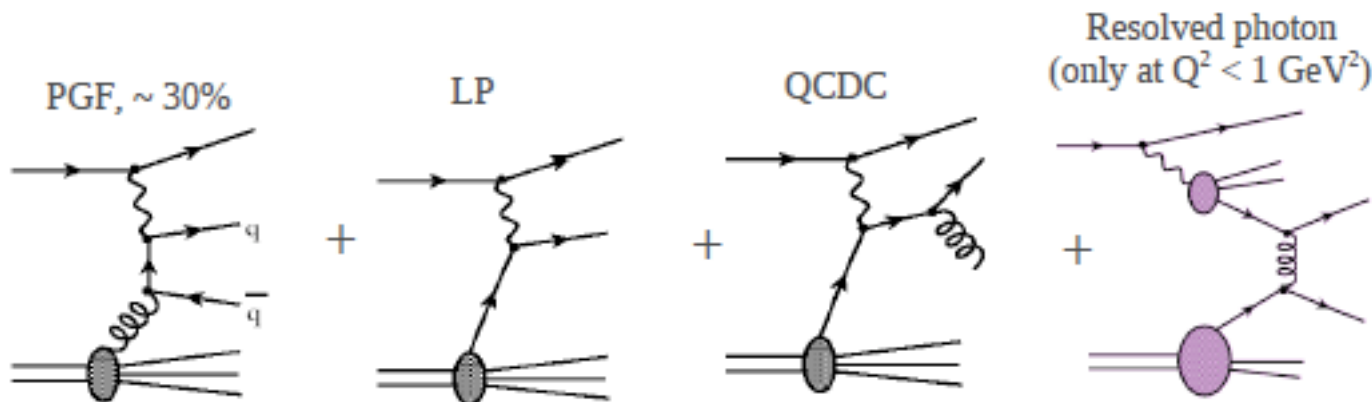
R.D.Carlitz, J.C.Collins and A.H.Mueller, Phys.Lett.B 214, 229 (1988)

Revisited by A.Bravar,D.von Harrach and A.Kotzinian, Phys.Lett.B 421, 349 (1998)

Applied by SMC, HERMES and COMPASS

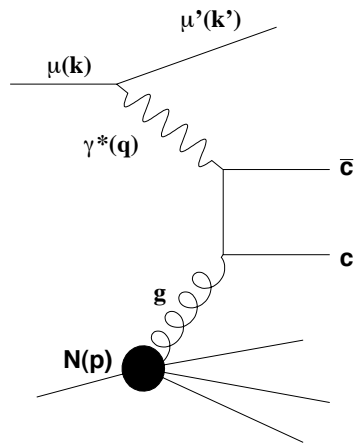
Two kinematical regions: low and large Q^2 :

- low Q^2 - $Q^2 < 1$ (GeV/c)² - here p_T is a perturbative scale, also **resolved photon contribution** important (~50%) - COMPASS 2002-2003 data published
- large Q^2 - $Q^2 > 1$ (GeV/c)² - scale Q^2 - 2002-2006 COMPASS data, method based on Neural Network approach used



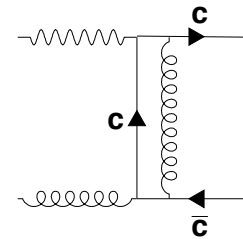
gluon polarisation from charm meson production

COMPASS: Phys. Rev. D(2013) 052018

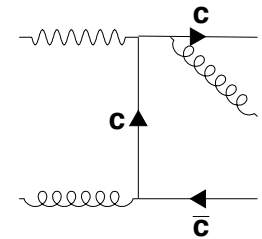


LO

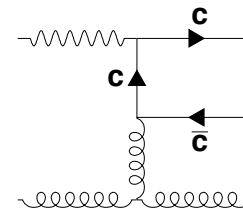
ex. of NLO



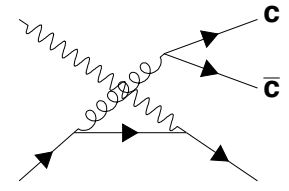
(a)



(b)



(c)

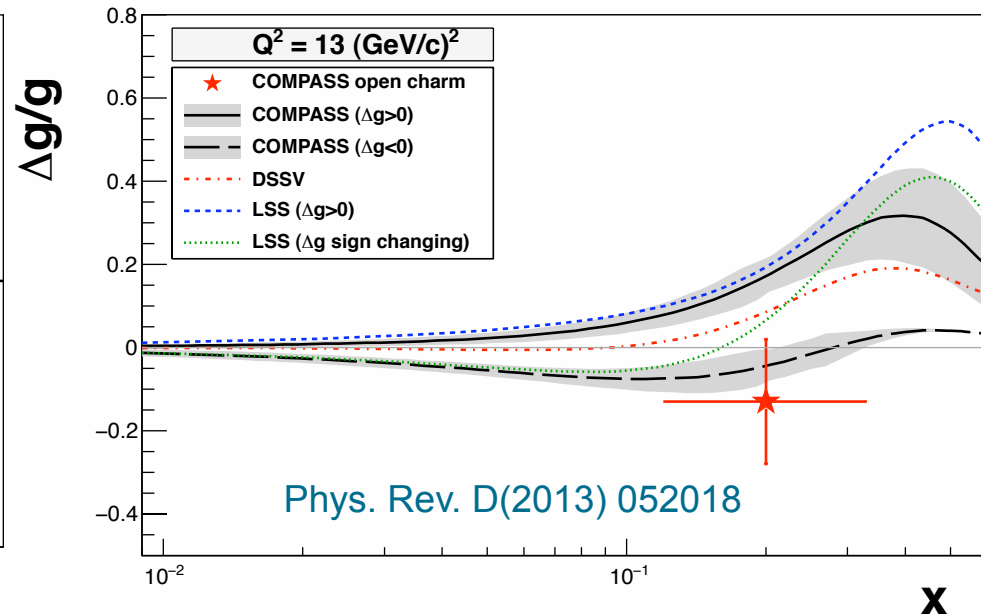
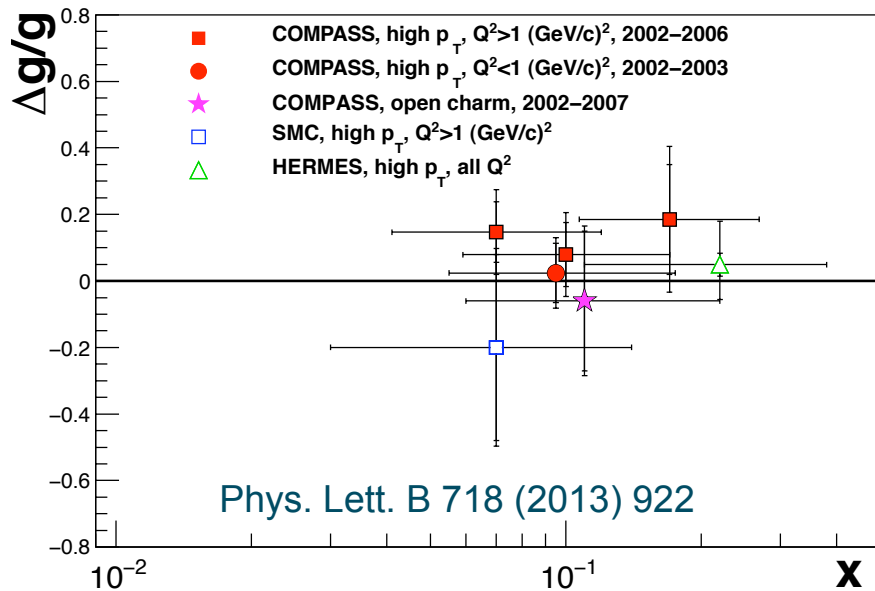


(d)

$$A_{\gamma N} = \frac{a_{LL}}{D} \frac{\Delta g}{g} + A_{\text{corr}}$$

Aroma MC generator used for simulating Phase Space (kinematics) (NLO - Parton Shower) calculations
 NLO Calculations from I. Bojak and M. Stratmann: Phys. Lett. B 433 (1998) 411; Nucl. Phys. B 540 (1999) 345

gluon polarisation from charm meson production



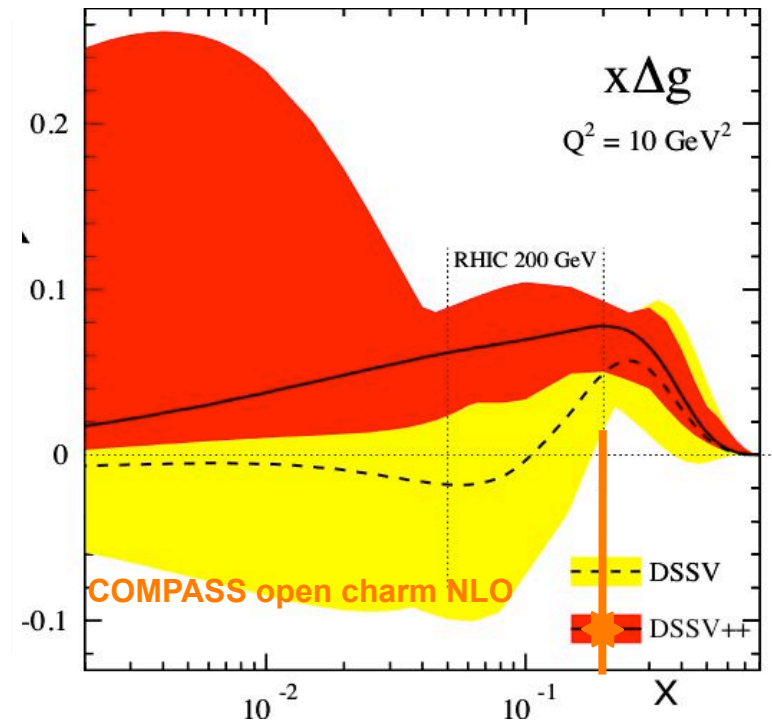
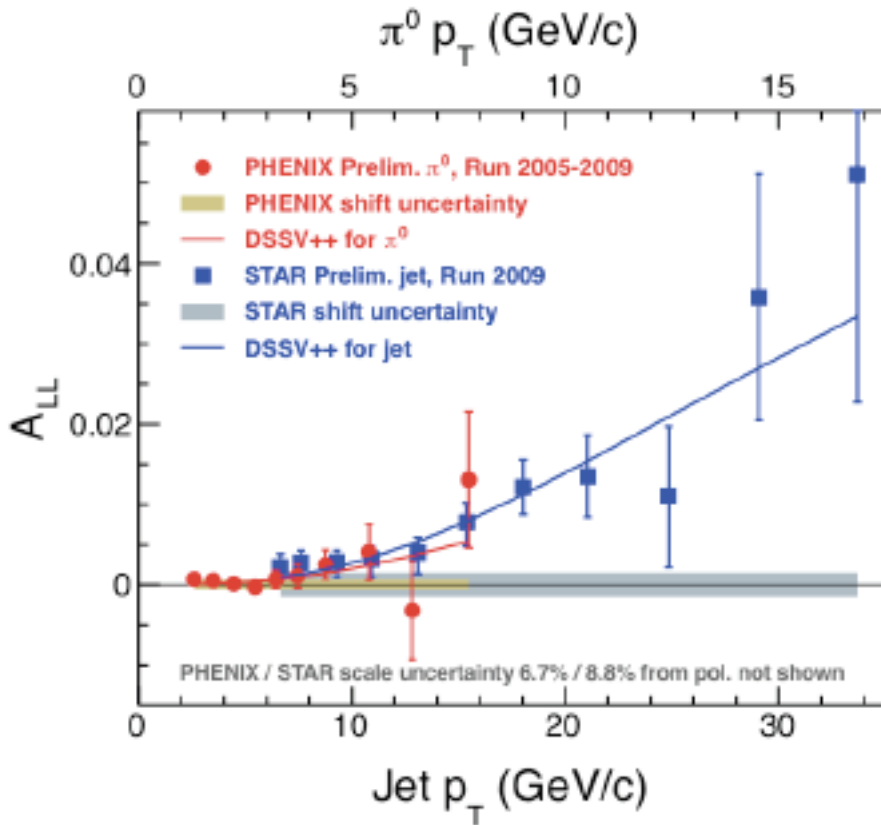
LO: all data consistent and point toward small gluon polarisation

NLO QCD result of COMPASS, at $\langle x \rangle \approx 0.20$, influences a $\Delta g(x) > 0$ fit, reducing $\Delta G = 0.39 \pm 0.07$ (stat.) to 0.24 ± 0.09 (stat.) at $Q^2 = 3$ (GeV/c)².

$$\left\langle \frac{\Delta g}{g} \right\rangle^{\text{NLO}} = -0.13 \pm 0.15 \text{ (stat.)} \pm 0.15 \text{ (syst.)}$$

RHIC results.

Gluon polarisation summary

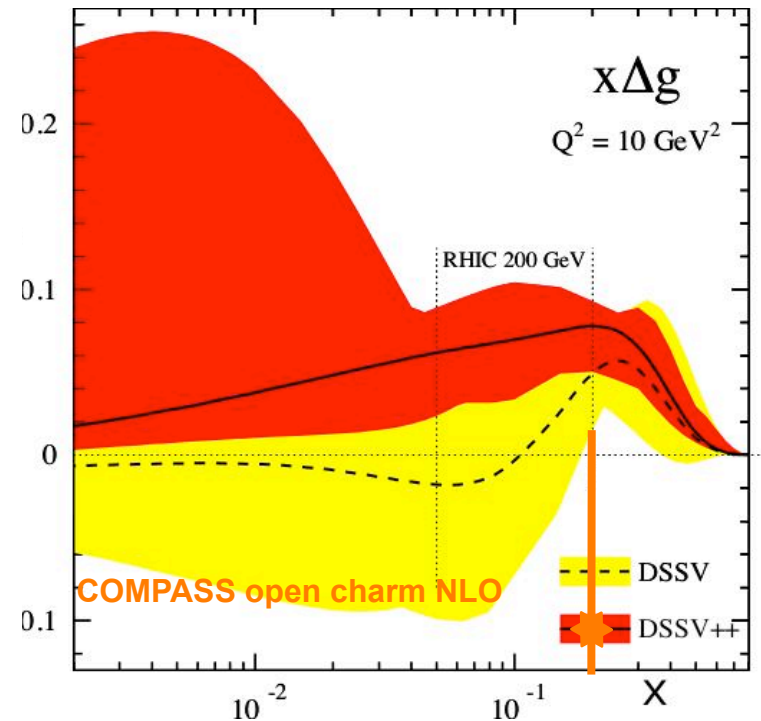
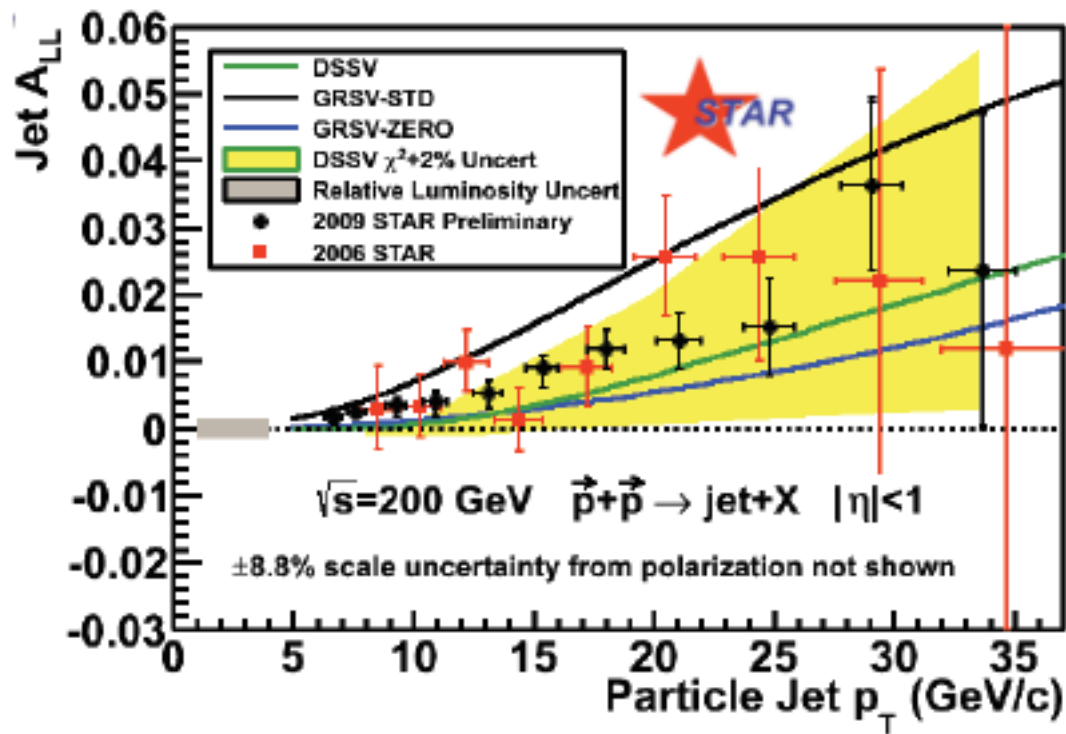


$$\text{At } Q^2 = 10 \text{ GeV}^2, \text{ DSSV : } \int_{0.05}^{0.2} \Delta g(x) dx = 0.005_{-0.164}^{+0.129}; \text{ DSSV++ : } \int_{0.05}^{0.2} \Delta g(x) dx = 0.10_{-0.07}^{+0.06}$$

$$\text{But DSSV : } \int_0^1 \Delta g(x) dx = 0.013_{-0.314}^{+0.702}; \text{ DSSV++ : } \int_0^1 \Delta g(x) dx = ?$$

RHIC results.

Gluon polarisation summary

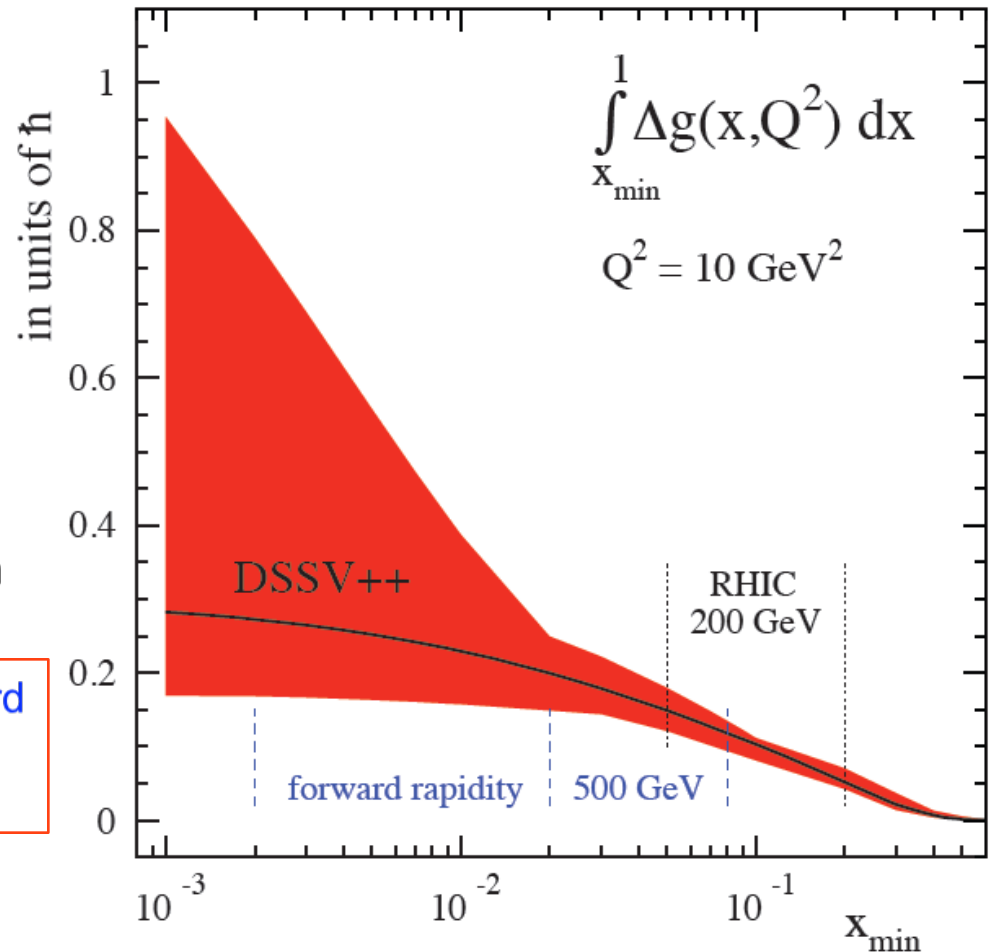


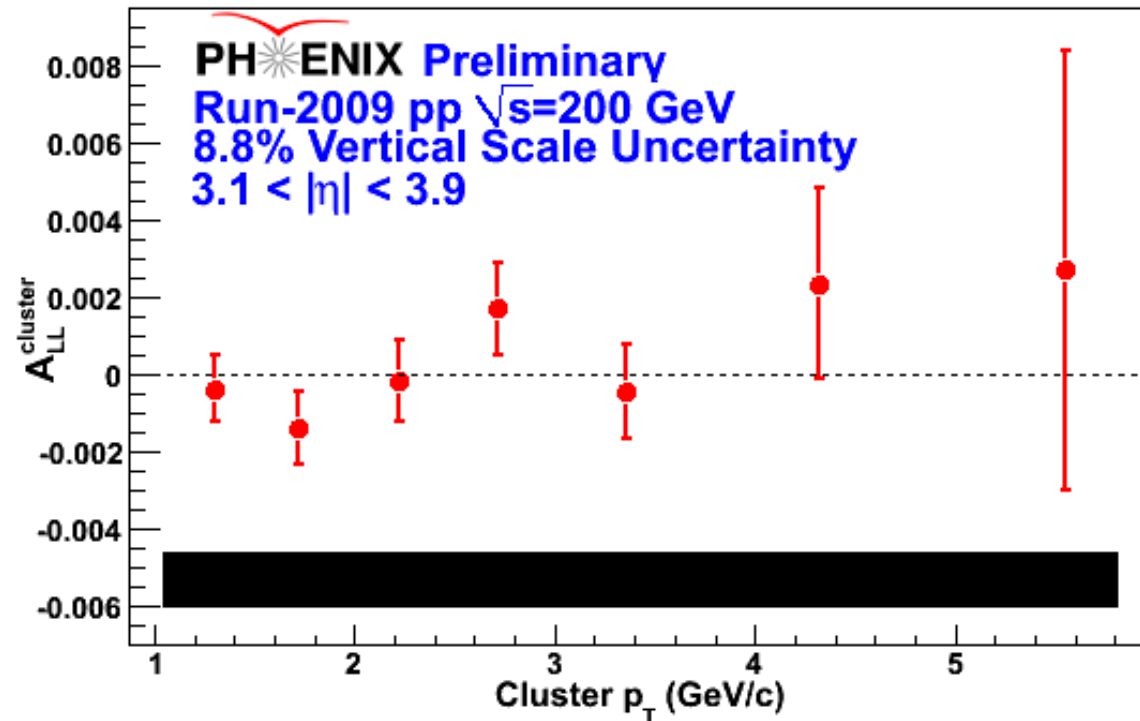
$$\text{At } Q^2 = 10 \text{ GeV}^2, \text{ DSSV: } \int_{0.05}^{0.2} \Delta g(x) dx = 0.005^{+0.129}_{-0.164}; \text{ DSSV++: } \int_{0.05}^{0.2} \Delta g(x) dx = 0.10^{+0.06}_{-0.07}$$

$$\text{But DSSV: } \int_0^1 \Delta g(x) dx = 0.013^{+0.702}_{-0.314}; \text{ DSSV++: } \int_0^1 \Delta g(x) dx = ?$$

RHIC results. Current uncertainty on gluon polarisation at RHIC

- **Large uncertainty still remains in both the shape of $\Delta g(x)$ and its integral, particularly for small x region not yet accessed by current RHIC data**
- **RHIC will continue to improve constraints through the following:**
 - Reduced statistical and systematic uncertainties for all A_{LL} channels
 - Better handle on measured x through correlation measurements
 - Access lower x region through **forward measurements** and $\sqrt{s} = 500 \text{ GeV}$





- High p_T EM Cluster Asymmetry, forward pseudo-rapidity $3.1 < |\eta| < 3.9$
 - >80% Merged π^0
- 510 GeV Datasets: Run09, Run11, Run12, Run13
 - Run12 and Run13 use new MPC electronics with $\sim 4x$ higher purity

Summary

- New precise data have been presented - new input to global QCD analysis
- New multiplicities for K and Pi can constrain FFs, especially D_s^K . It can explain the difference between DIS and SIDIS strange quark sea polarisation
- New RHIC W boson data can significantly improve precision on $\overline{\Delta u}$ and $\overline{\Delta d}$ (DSSV++ prediction)
- New RHIC data increase the value of the gluon polarisation in global QCD analysis (DSSV++). Still open question what the value of integral is - large uncertainties for small x region. New measurements for small x needed to constrain ΔG
- After more than 30 years - still the spin structure of the nucleon is understood only qualitatively. What we learnt is that even fast moving nucleon is a 3-dimensional complicated object with probably non-negligible orbital momentum contribution to its spin.

Thank you

