

Study of hadron structure using Drell-Yan scattering

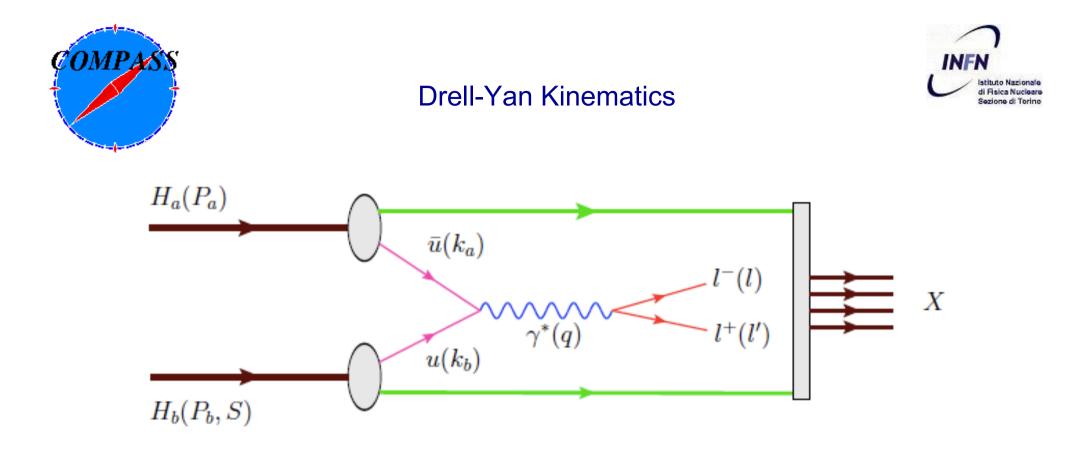
Oleg Denisov INFN section of Turin 11.07.2013



Outline



- Drell-Yan process, status of the theory •
- Drell-Yan, polarised case •
- Transversity & TMD PDFs:
 - Proton description at LO
 - TMDs factorisation and universality crucial test of modern QCD
- Unpolarised pion Drell-Yan:
 - Pion PDFs
 - EMC effect in DY
 - Sea guark distributions
- TMDs study choice of kinematic domain
- Upcoming/started DY experiments: ٠
 - E906 @ FermiLab
 - Unpolarised proton induced DY •
 - Future polarised DY experiment
 - Polarised DY@COMPASS _
 - Set-up
 - Kinematics & Projections
 - RHIC?
- Some conclusions 15-07-2013



 $P_{a(b)}$ $s = (P_a + P_b)^2,$ $x_{a(b)} = q^2 / (2P_{a(b)} \cdot q),$ $x_F = x_a - x_b,$ $M_{\mu\mu}^2 = Q^2 = q^2 = s \ x_a \ x_b,$ $k_{Ta(b)}$ $q_T = P_T = k_{Ta} + k_{Tb}$

the momentum of the beam (target) hadron, the total centre-of-mass energy squared, the momentum fraction carried by a parton from $H_{a(b)}$, the Feynman variable, the invariant mass squared of the dimuon, the transverse component of the quark momentum, the transverse component of the momentum of the virtual photon.

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Drell-Yan cross-section – general (full) angular distribution



Direct extraction of transversity and its accompanying T-odd distribution from the unpolarized and single-polarized Drell-Yan processes $\frac{d\sigma^{(1)}(z)}{dQd\phi}$

> A.N. Sissakian,^{*} O.Yu. Shevchenko,[†] A.P. Nagaytsev,[‡] and O.N. Ivanov[§] Joint Institute for Nuclear Research, 141980 Dubna, Russia

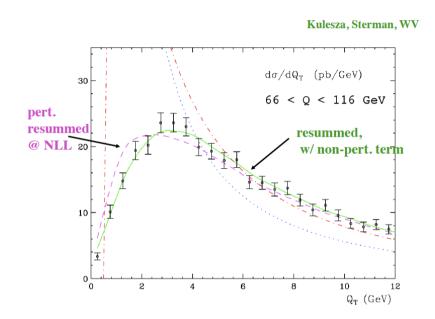
> > Phys. Rev. D72 (2005) 054027

 $\frac{d\sigma^{(1)}(H_1H_2^{\uparrow} \to l\bar{l}X)}{d\Omega d\phi_{S^2} dx_1 dx_2 d^2 \mathbf{q}_T} = \frac{\alpha^2}{12Q^2} \sum_q e_q^2$ $\times \left\{ (1 + \cos^2\theta) \mathcal{F}[\bar{f}_{1q}f_{1q}] + \sin^2\theta \cos(2\phi) \right.$ $\times \mathcal{F} \left[(2\hat{\mathbf{h}} \cdot \mathbf{k}_{1T} \, \hat{\mathbf{h}} \cdot \mathbf{k}_{2T} - \mathbf{k}_{1T} \cdot \mathbf{k}_{2T}) \frac{\bar{h}_{1q}^{\perp} h_{1q}^{\perp}}{M_1 M_2} \right]$ + $(1 + \cos^2\theta)\sin(\phi - \phi_{S_2})\mathcal{F}\left[\hat{\mathbf{h}} \cdot \mathbf{k}_{2T} \frac{\bar{f}_1^q f_{1T}^{\perp q}}{M_2}\right]$ $-\sin^2\theta\sin(\phi+\phi_{S_2})\mathcal{F}\left[\hat{\mathbf{h}}\cdot\mathbf{k}_{1T}\frac{\bar{h}_{1q}^{\perp}h_{1q}}{M_1}\right]$

 $\frac{d\sigma}{d^4q \, d\Omega} = \frac{\alpha_{em}^2}{F \, q^2} \times$ $\frac{2008: S. Arnold, (Ruhr U., Bochum), A. Metz, (Temple U.), M. Schlegel, (Jefferson Lab) Phys.Rev.D79:034005,2009, e-Print: arXiv:0809.2262$ $<math display="block"> \left\{ \left((1 + \cos^2 \theta) F_{UU}^1 + (1 - \cos^2 \theta) F_{UU}^2 + \sin 2\theta \cos \phi F_{UU}^{\cos \phi} + \sin^2 \theta \cos 2\phi F_{UU}^{\cos 2\phi} \right) + S_{aL} \left(\sin 2\theta \sin \phi F_{LU}^{\sin \phi} + \sin^2 \theta \sin 2\phi F_{LU}^{\sin 2\phi} \right) + S_{bL} \left(\sin 2\theta \sin \phi F_{UL}^{\sin \phi} + \sin^2 \theta \sin 2\phi F_{UL}^{\sin 2\phi} \right) \right\} + \left| \vec{S}_{aT} \right| \left[\sin \phi_a \left((1 + \cos^2 \theta) F_{1U}^1 + (1 - \cos^2 \theta) F_{2U}^2 + \sin 2\theta \cos \phi F_{UU}^{\cos \phi} + \sin^2 \theta \cos 2\phi F_{UU}^{\cos 2\phi} \right) + \cos \phi_a \left(\sin 2\theta \sin \phi F_{UT}^{\sin \phi} + \sin^2 \theta \sin 2\phi F_{UT}^{\sin 2\phi} \right) \right]$ $+ \left| \vec{S}_{bT} \right| \left[\sin \phi_b \left((1 + \cos^2 \theta) F_{UT}^1 + (1 - \cos^2 \theta) F_{UT}^2 + \sin 2\theta \cos \phi F_{UT}^{\cos \phi} + \sin^2 \theta \cos 2\phi F_{UT}^{\cos 2\phi} \right) + \cos \phi_b \left(\sin 2\theta \sin \phi F_{UT}^{\sin \phi} + \sin^2 \theta \sin 2\phi F_{UT}^{\sin 2\phi} \right) \right]$ $+ S_{aL} S_{bL} \left((1 + \cos^2 \theta) F_{LL}^1 + (1 - \cos^2 \theta) F_{LL}^2 + \sin 2\theta \cos \phi F_{LL}^{\cos \phi} + \sin^2 \theta \cos 2\phi F_{LL}^{\cos 2\phi} \right)$ $+ S_{aL} S_{bL} \left((1 + \cos^2 \theta) F_{LL}^1 + (1 - \cos^2 \theta) F_{LL}^2 + \sin 2\theta \cos \phi F_{LL}^{\cos \phi} + \sin^2 \theta \cos 2\phi F_{LL}^{\cos 2\phi} \right)$ $+ S_{aL} S_{bL} \left((1 + \cos^2 \theta) F_{LL}^1 + (1 - \cos^2 \theta) F_{LL}^2 + \sin 2\theta \cos \phi F_{LL}^{\cos \phi} + \sin^2 \theta \cos 2\phi F_{LL}^{\cos 2\phi} \right)$ $+ S_{aL} S_{bL} \left((1 + \cos^2 \theta) F_{LL}^1 + (1 - \cos^2 \theta) F_{LL}^2 + \sin 2\theta \cos \phi F_{LL}^{\cos \phi} + \sin^2 \theta \cos 2\phi F_{LL}^{\cos 2\phi} \right)$ $+ S_{aL} S_{bL} \left((1 + \cos^2 \theta) F_{LL}^1 + (1 - \cos^2 \theta) F_{LL}^2 + \sin 2\theta \cos \phi F_{LL}^{\cos \phi} + \sin^2 \theta \cos 2\phi F_{LL}^{\cos 2\phi} \right)$

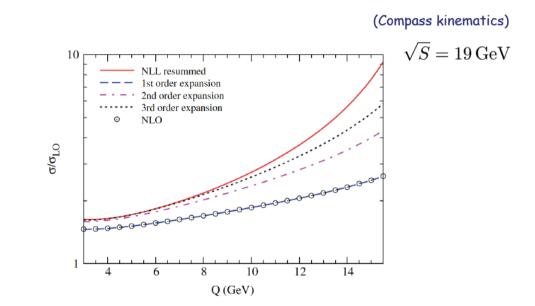
Pion-induced DY, DY Cross section calculations theoretical calculations - COMPASS





Aicher et al. (PRL 105, 252003 (2010))

Very big progress recently achieved: NLO threshold re-summation mechanism with non - perturbative term (Vogelsang and collaborators) - good experimental data description – K-factor issue is not there anymore



Aicher, Schäfer, WV (earlier studies: Shimizu, Sterman, WV, Yokoya)

15-07-2013

OMPAS

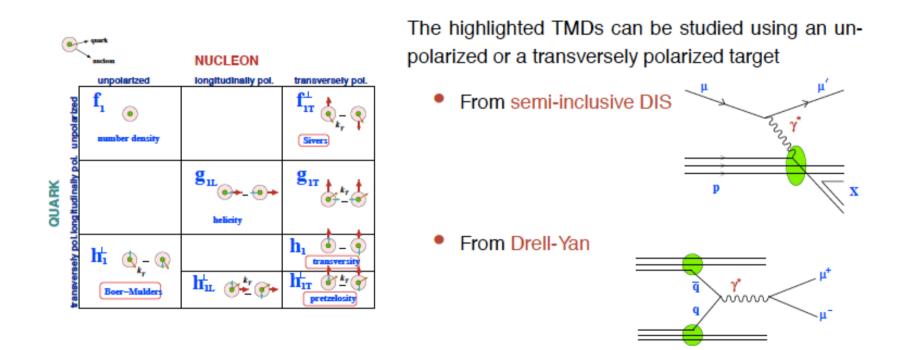
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Leading Order PDFs



In LO, 8 Transverse Momentum Dependent PDFs (TMDs) are needed to describe the nucleon structure when the intrinsic transverse momentum is taken into account:



By measuring the Single Transverse Spin Asymmetries (SSA) in these processes one can access the correlations between the partons k_T and the nucleon spin.



Were we studying TMD PDFs?



So far we were doing SIDIS thinking that we are studying the structure of the nucleon, no proof so far (artefact of our analysis tools or some different physics: higher twist or so)....

In order to obtain an important and convincing proof we have to change the "probe" \rightarrow the physics process we are using to access the nucleon structure. We have to go to Drell-Yan and see whether we see the similar effects and we can describe them in the framework of the TMD formalism and if the TMDs as we see them in DY are the same as in SIDIS.

Actually it was explicitly written in a very first version of the COMPASS DY proposal which was written by O.Teryaev and myself in the very beginning of 2005.



The three TMD PDFs below describe important properties of spin dynamics of nucleon

- $f_{1T}^{\perp}(x, k_T^2)$: the Sivers effect is related to an azimuthal asymmetry in the parton intrinsic transverse momentum distribution induced by the nucleon spin
- ▶ h[⊥]₁ (x, k²_T): the Boer-Mulders function describes the correlation between the transverse spin and the transverse momentum of a quark inside the unpolarised hadron
- *h*[⊥]_{1T} (*x*, *k*²_T): the Pretzelosity function describes the polarisation of a quark along its intrisic *k*_T direction making accessible the orbital angular momentum information





Single-polarised DY cross-section: Leading order QCD parton model



At LO the general expression of the DY cross-section simplifies to (Aram Kotzinian) :

$$\frac{d\sigma^{LO}}{d^4q \, d\Omega} = \frac{\alpha_{em}^2}{F \, q^2} \hat{\sigma}_U^{LO} \left\{ \left(1 + D_{[\sin^2 \theta]}^{LO} A_U^{\cos 2\phi} \cos 2\phi \right) \right. \\
\left. + S_L D_{[\sin^2 \theta]}^{LO} A_L^{\sin 2\phi} \sin 2\phi \right. \\
\left. + \left. \left. |\vec{S}_T| \left[A_T^{\sin \phi_S} \sin \phi_S + D_{[\sin^2 \theta]}^{LO} \left(A_T^{\sin(2\phi + \phi_S)} \sin(2\phi + \phi_S) \right. \right. \\
\left. + A_T^{\sin(2\phi - \phi_S)} \sin(2\phi - \phi_S) \right) \right] \right\},$$

Thus the measurement of 4 asymmetries (modulations in the DY cross-section):

- $-A_U^{\cos 2\phi}$ gives access to the Boer-Mulders functions of the incoming hadrons, $-A_T^{\sin \phi_S}$ to the Sivers function of the target nucleon, $-A_T^{\sin(2\phi+\phi_S)}$ to the Boer-Mulders functions of the beam hadron and to h_{1T}^{\perp} , the
- pretzelosity function of the target nucleon,
- $-A_T^{\sin(2\phi-\phi_S)}$ to the Boer-Mulders functions of the beam hadron and h_1 , the transversity function of the target nucleon.

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The time-reversal odd character of the Sivers and Boer-Mulders PDFs lead to the prediction of a sign change when accessed from SIDIS or from Drell-Yan processes:

 \hookrightarrow Check the predictions:

 $f_{1T}^{\perp}(DY) = -f_{1T}^{\perp}(SIDIS)$

 $h_1^{\perp}(DY) = -h_1^{\perp}(SIDIS)$

Its experimental confirmation is considered a crucial test of non-perturbative QCD.

Universality test includes not only the sing-reversal character of the TMDs but also the comparison of the amplitude as well as the shape of the corresponding TMDs COMPASS

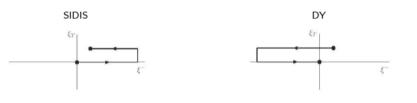
Sivers, Boer-Mulders functions SIDIS $\leftarrow \rightarrow$ DY

QCD $\sigma_{h} \cong \sigma_{p} \times PDF$

QCD factorization, valid for hard processes only (Q, $q_{\rm T}\,are\,large)$

Cross-sections are gauge-invariant objects, to provide the gauge invariance of the PDFs the gauge-link was introduced (intrinsic feature of PDF). The presence of gauge-link provides the possibility of existence of non-zero T-odd TMD PDFs

Direction of the gauge-link of the k_T dependent PDF is process-dependent (gauge-link is resummation of all collinear soft gluons) and it changes to the opposite in SIDIS wrt DY



i Fisica Nucleare

Sivers and Boer-Mulders functions are T-odd, and to provide the time-invariance they change the sign in SIDIS wrt DY due to the opposite direction of the gauge-link

$$f_{1T}^{\perp}(x, \mathbf{k}_T) \Big|_{SIDIS} = -f_{1T}^{\perp}(x, \mathbf{k}_T) \Big|_{DY}$$
$$h_1^{\perp}(x, \mathbf{k}_T) \Big|_{SIDIS} = -h_1^{\perp}(x, \mathbf{k}_T) \Big|_{DY}$$

J.C. Collins, Phys. Lett. B536 (2002) 43

J. Collins, talk at LIGHT CONE 2008



SIDIS ← → DY – QCD test

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Andreas Metz (Trento-TMD'2010):

Sign reversal of the Sivers function

• Prediction based on operator definition (Collins, 2002)

 $f_{1T}^{\perp}\big|_{DY} = - \left.f_{1T}^{\perp}\right|_{DIS}$

- What if sign reversal of f_{1T}^{\perp} is not confirmed by experiment?
 - Would not imply that QCD is wrong
 - Would imply that SSAs not understood in QCD
 - Problem with TMD-factorization
 - Problem with resummation of large logarithms
 - \rightarrow Resummation relevant if more than one scale present
 - \rightarrow CSS resummation in Drell-Yan (Collins, Soper, Sterman, 1985); resum logarithms of the type

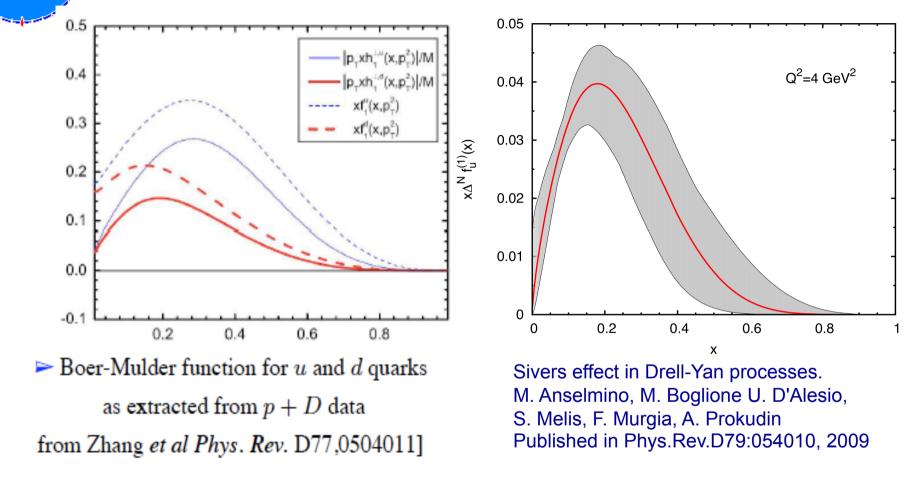
$$\alpha_s^k \ln^{2k} \frac{\vec{Q}_T^2}{Q^2}$$

 \rightarrow Has also implications for Fermilab and LHC physics

Some indications for the future Drell-Yan experiments



1. TMD PDFs – ALL are sizable in the valence quark region



2. $\Lambda_{QCD} < p_T < Q$: - P_T should be small (~ 1 GeV), can be generated by intrinsic motion of quarks and/or by soft gluon emission. This is the region where TMD formalism applies. **Oleg Denisov**





Pion-induced unpolarised DY, Pion structure functions I

The only way to access pion PDFs

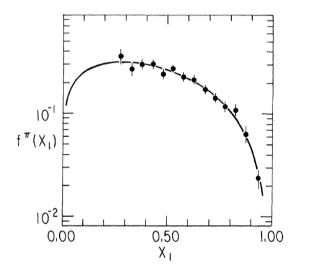
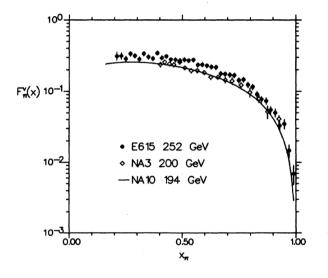


FIG. 3. The pion structure function $f^{\pi}(x_1) = x_1 \overline{u}^{\pi^-}(x_1)$.

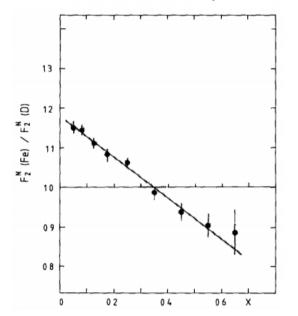
CIP (PRL 42, 951, (1979))



E615 (PRD 39, 92 (1989)):



EMC effect - experiment



Cloet et. al (PRL 102, 252301, 2009): Flavor dependence of the EMC effects ?

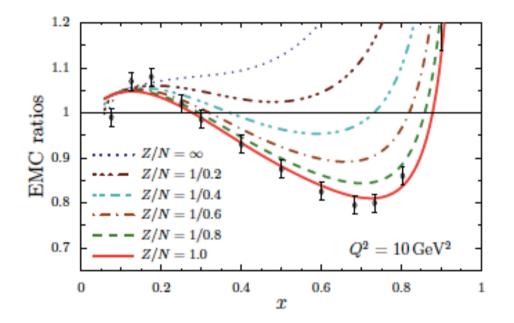
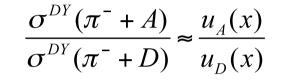


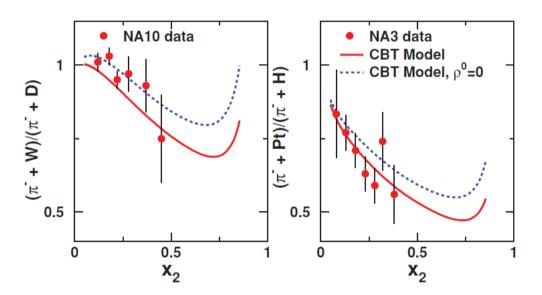
FIG. 1: Isospin dependence of the EMC effect for protonneutron ratios greater than one. The data is from Ref. [24] and corresponds to N = Z nuclear matter.

The isovector ρ^0 mean-field generated in Z \neq N nuclei can modify nucleon's *u* and *d* PDFs in nuclei.

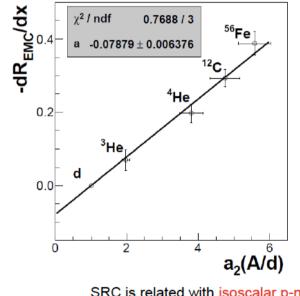
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Dutta et al. (PRC 83, 042201, 2011): Pion-induced Drell-Yan and the flavor-dependent EMC effect



SRC is related with isoscalar p-n interaction. \rightarrow NO flavor-dependence of EMC effect

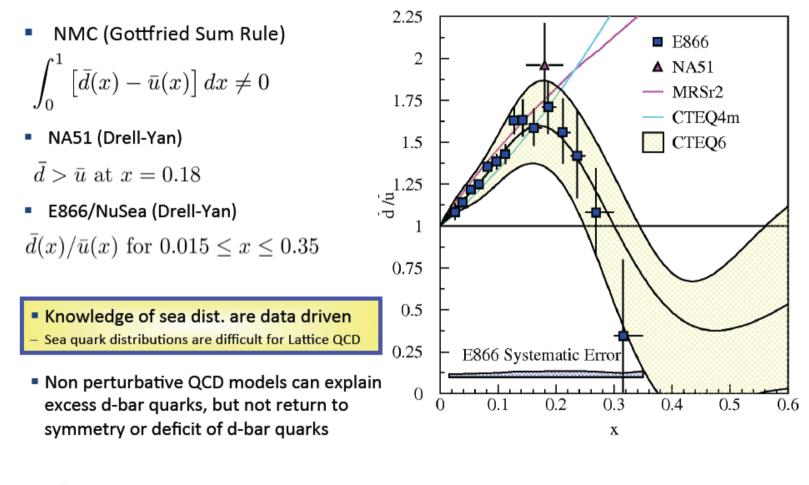
Weinstein et al. (PRL 106 , 052301 (2011)): EMC & Short Range Correlation (SRC)



Sea quarks contribution study



Light Antiquark Flavor Asymmetry: Brief History





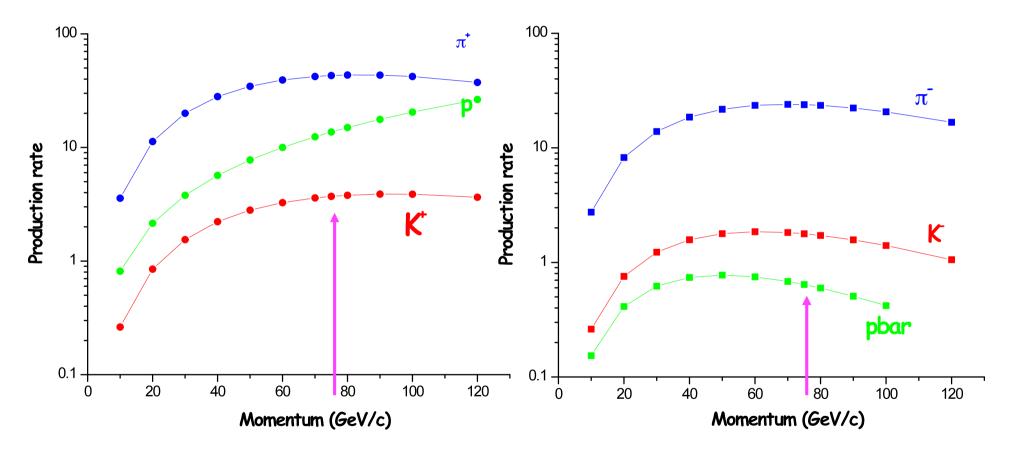
15



Kaon and Antiproton induced Drell-Yan I, Case of secondary SPS hadron beams



Particle production at 0 mrad



Preliminary rate estimates for RF separated antiproton beams



E906 Drell-Yan experiment at FermiLab (materials by Paul E. Reimer)



Advantages of 120 GeV Main Injector

The (very successful) past: Fermilab E866/NuSea

 $4\pi\alpha^2$ 1

 $9x_{t}x_{t}$

Data in 1996-1997

 $dx_{\rm t} dx_{\rm b}$

- ¹H, ²H, and nuclear targets
- 800 GeV proton beam

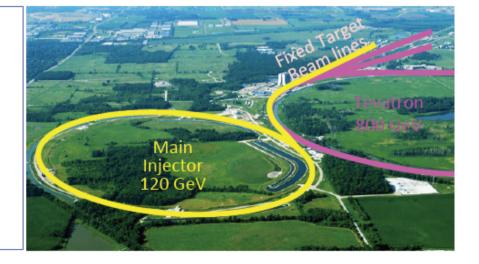
- The future: Fermilab E906
- First test run in 2011
- ¹H, ²H, and nuclear targets

 $e^{2} \left[\bar{q}_{t} \left(x_{t} \right) q_{b} \left(x_{b} \right) + q_{t} \left(x_{t} \right) \bar{q}_{b} \left(x_{b} \right) \right]$

120 GeV proton Beam

- Cross section scales as 1/s
 - 7× that of 800 GeV beam
- Backgrounds, primarily from J/ψ decays scale as s
 - 7× Luminosity for same detector rate as 800 GeV beam

50× statistics!!



15-07





Drell-Yan Spectrometer Guiding Principles

- Follow basic design of MEast spectrometer (don't reinvent the wheel):
 - Two magnet spectrometer
 - Beam dump within first magnet
- Hadron absorber within first magnet
- Muon-ID wall before final elements
- Where possible and practical, reuse elements of the E866 spectrometer.
 - Tracking chamber electronics
 - Hadron absorber, beam dump, muon ID walls
 - Station 2 tracking chambers
 - Hodoscope array PMT's
 - SM3 Magnet
- New Elements
 - 1st magnet (different boost)
 Experiment shrinks from 60m to 26m
 - Sta. 1 tracking (rates)
 - Scintillator (age)
 - Trigger (flexibility)
 - St. 3 tracking (acceptance)
 - Paul E. Reimer ECT* Workshop on the Flavor Structure of the Nucleon Sea





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Event: 18 Coda Event

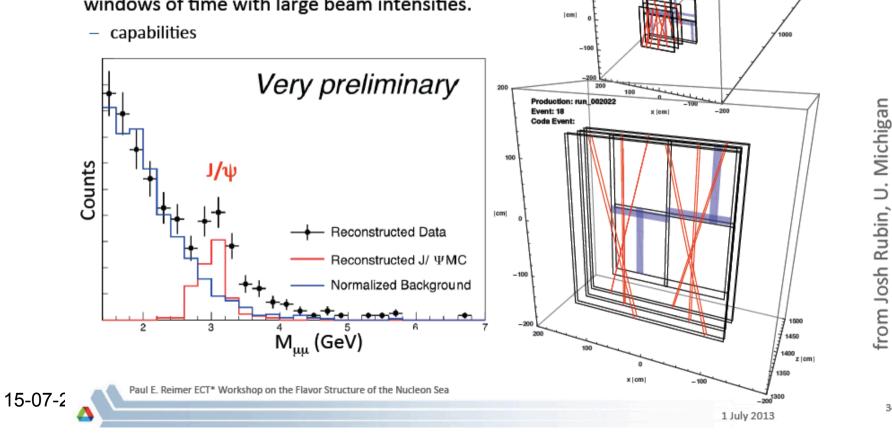


z [cm]

34

The Splat-Block Card

- A card was developed to keep a running average of the multiplicity over a 160 ns window (8 RF buckets).
- If average multiplicity above threshold, raises a trigger veto
- Luminosity greatly reduced, but trigger suppresses windows of time with large beam intensities.

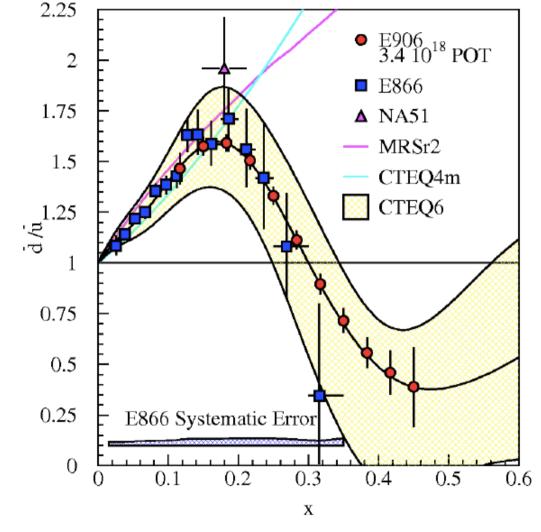






Extracting d-bar/-ubar From Drell-Yan Scattering

- E906/Drell-Yan will extend these measurements and reduce statistical uncertainty.
- E906 expects systematic uncertainty to remain at approx. 1% in cross section ratio.







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Polarized Drell-Yan at Fermilab

Test Fundamental Prediction of QCD Gauge formalism and factorization

$$\left.f_{1T}^{\perp}\left(x,k_{T}
ight)
ight|_{\mathrm{SIDIS}}=-\left.f_{1T}^{\perp}\left(x,k_{T}
ight)
ight|_{\mathrm{DY}}$$

Does the Sivers' function change sign?

Polarized beam

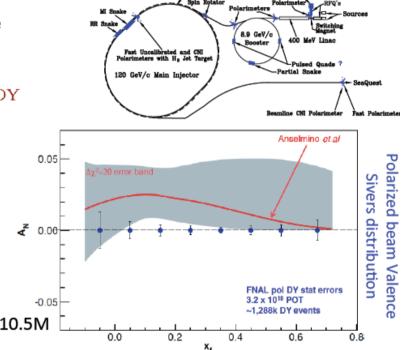
- Measure valence-quark distributions
 - Directly comparable to SIDIS measurements
 - Fits expect measureable asymmetry
- Stage-I Fermilab Approval
- Fermilab intensity frontier: Luminosity
- Cost: \$6.5M + 65%(cont. and manage.) = \$10.5M
- Argonne Spokesperson (PER)

Polarized Target

15-07-2013

- Measure Single Spin Asymmetry (SSA) in Sea Quarks
- Double Spin Asymmetry w/polarized beam
- LOI submitted Spring 2013

Paul E. Reimer ECT* Workshop on the Flavor Structure of the Nucleon Sea



1 July 2013

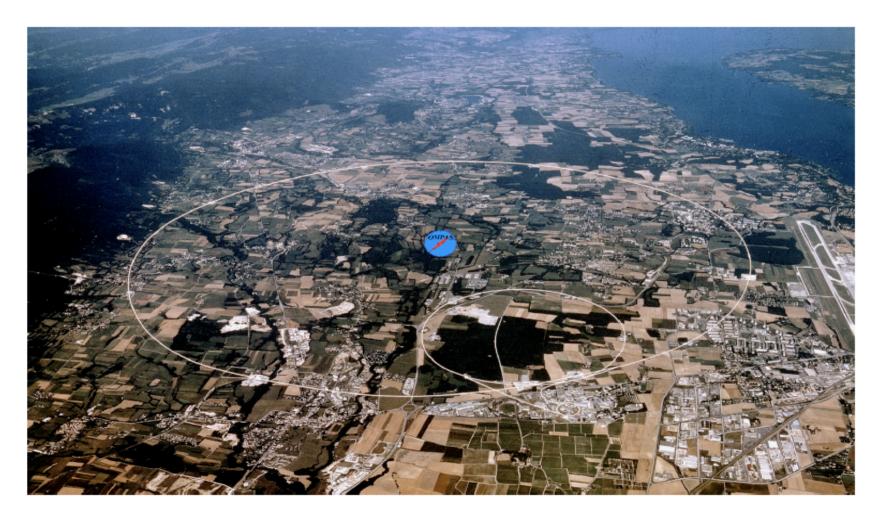
46



COMPASS facility at CERN (SPS)



COmmon Muon Proton Apparatus for Structure and Spectroscopy





Two stages spectrometer

Large Angle Spectrometer (LAS)

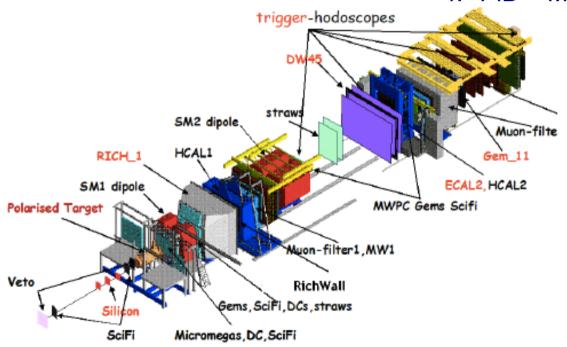
Small Angle Spectrometer (SAS)

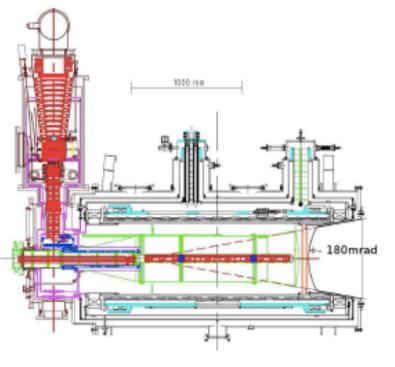
COMPASS facility at CERN



Most important features:

- 1. Muon, electron or hadron secondary beams
- 2. Solid state polarised targets (NH₃ or ⁶LiD) as well as liquid hydrogen target and nuclear
 - targets
- 3. Powerful tracking system 350 planes
- 4. PiD Muon Walls, Calorimeters, RICH





15-07-2013

COMPASS

COMPASS-II (New Physics) a piece of history



- COMPASS is very sophisticated, universal and flexible facility → Physics beyong SIDIS and hadron spectroscopy is possible:
 - Unique COMPASS Polarised Target
 - Both hadron and lepton beams
 - Flexible spectrometer structure \rightarrow easy-accessible spectrometer components
- All together that has generated new physics proposals with COMPASS DVCS(GPDs) and polarised DY (here I will stick naturally to DY):
 - For the first time the idea of the polarised π-p DY were reported at the Villars SPSC meeting in September 2004, DY was basically considered as the alternative method (wrt SiDIS) to access transversity and TMD PDFs
 - Beginning of 2005 A.V.Efremov brought an attention to the issue of the sign-change of the Sivers function from SIDIS to DY (Brodsky, Collins), since then it became the flagship measurement
 - The first version of the polarised π⁻p DY Letter of Intent was written in 2005 (Oleg Teryaev has contributed as well)
 - Since then 5 International Workshops (Torino, Dubna, CERN, Trento), > 70 COMPASS DY subgroup meetings, 4 Beam Tests, > 50 presentations at the international Conferences.... In the end we succeed to convince ourselves that we can do it.
- The COMPASS-II proposal was submitted to the CERN SPSC on May 17th 2010
- Approved by the CERN research board on December 1st 2010, 1 year for Drell-Yan and 2 years for GPDs in the time interval between two LHC shutdowns.





- 1. Large angular acceptance spectrometer
- 2. SPS M2 secondary beams with the intensity up to 10⁸ particles per second
- 3. Transversely polarized solid state proton target with a large relaxation time and high polarization, when going to spin frozen mode;
- 4. a detection system designed to stand relatively high particle fluxes;
- 5. a Data Acquisition System (DAQ) that can handle large amounts of data at large trigger rates;
- 6. The dedicated muon trigger system

For the moment we consider two step DY program:

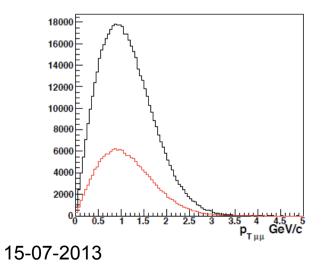
- •The program with high intensity pion beam
- •The program with Radio Frequency separated antiproton beam

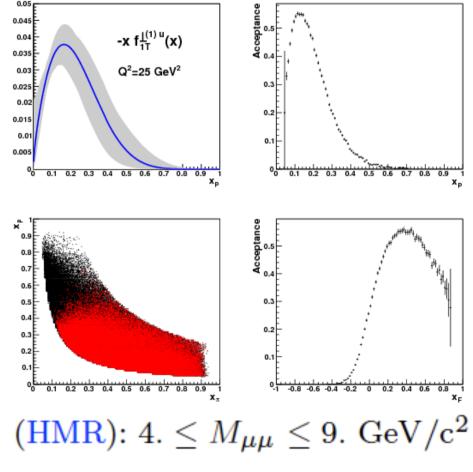


DY@COMPASS – kinematics - valence quark range $\pi^{-}p \rightarrow \mu^{-} \mu X$ (190 GeV pion beam)



- In our case (π⁻ p → μ⁻ μ X) contribution from valence quarks is dominant
- In COMPASS kinematics uubar dominance
- <P_T> ~ 1GeV TMDs induced effects expected to be dominant with respect to the higher QCD corrections





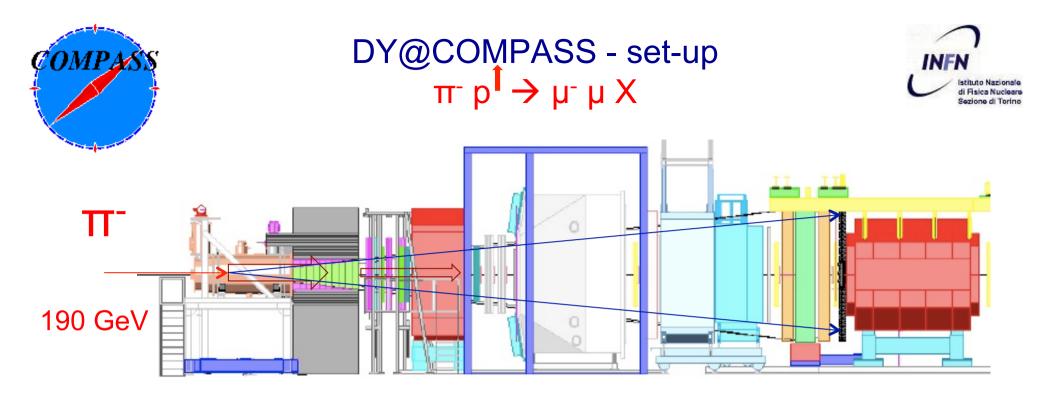


DY@COMPASS - set-up $\pi^{-} p \xrightarrow{I} \mu^{-} \mu X (190 \text{ GeV})$



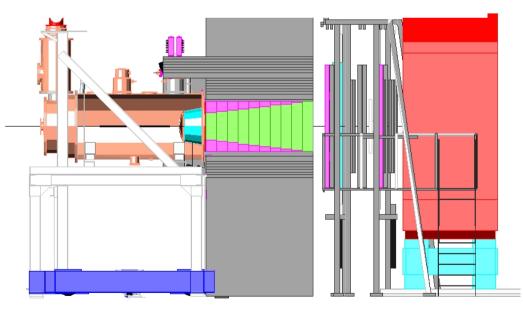
The main characteristics of the future fixed-target Drell-Yan experiment:

- Small cross section → High intensity hadron beam (up to 10⁹ pions per spill) on the COMPASS PT
- 2. High intensity hadron beam on thick target \rightarrow
 - 1. Hadron absorber to stop secondary particles flux
 - 2. Beam plug to stop the non interacted beam
 - 3. Radioprotection shielding around to protect things and people
 - 4. High-rate-capable radiation hard beam telescope
- 3. Hadron absorber + shielding \rightarrow PT has to be moved by 2.2 meters upstream
- LAS dominates in the acceptance → The performance of the LAS tracking system must be improved and muon trigger in LAS has to be well tuned.
- 5. Hadron absorber -> vertex detector is very welcome to improve cell-to-cell separation



Key elements:

- 1. COMPASS PT
- 2. Tracking system (both LAS abs SAS) and beam telescope in front of PT
- 3. Muon trigger (in LAS is of particular importance 60% of the DY acceptance)
- RICH1, Calorimetry also important to reduce the background (the hadron flux downstream of the hadron absorber ~ 10 higher then muon flux)



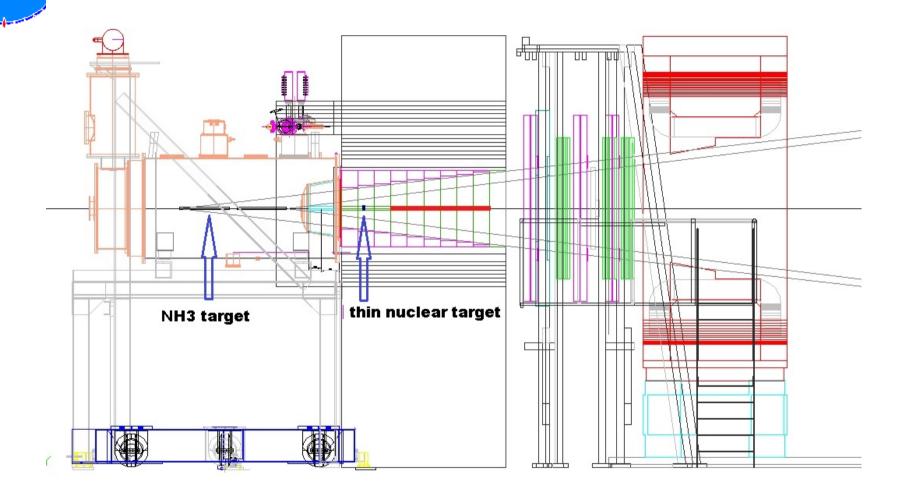
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Flavour-dependent EMC effect study, possible set-up

OMPAS



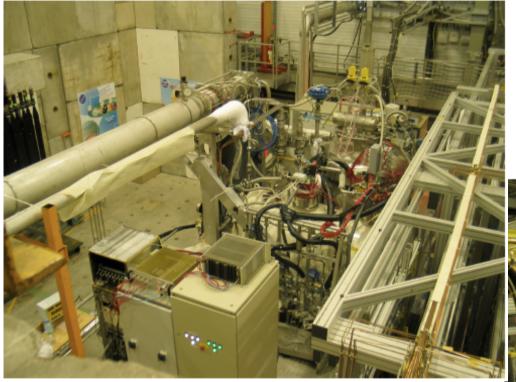


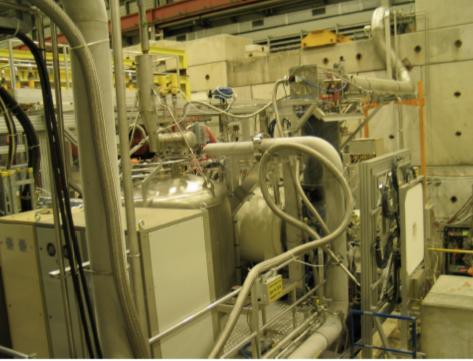
Feasibility and set-up optimization still has to be completed



PT movement

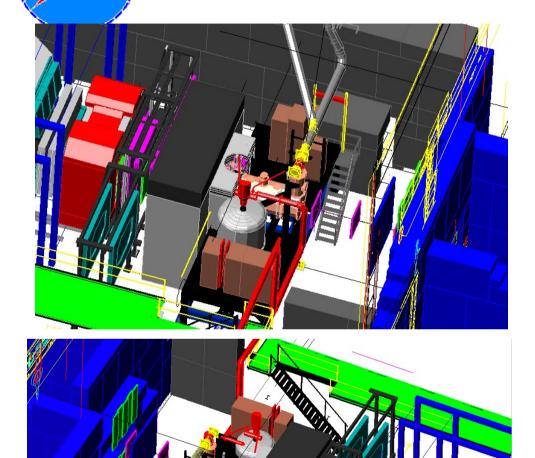


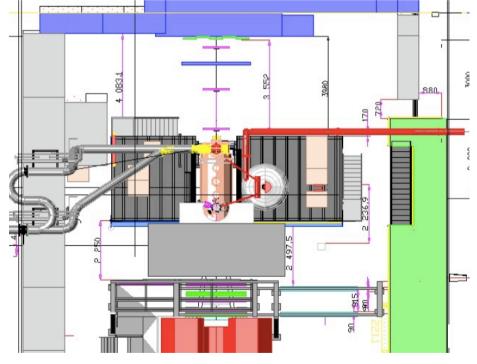




Drell-Yan experiment lay-out Hadron absorber&R.P.shieldings and Polarised Target







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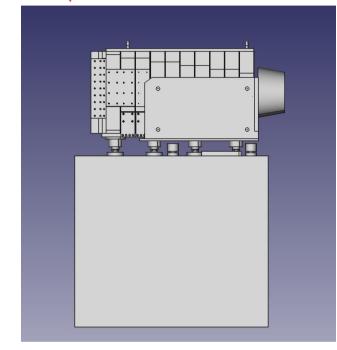
OMP A

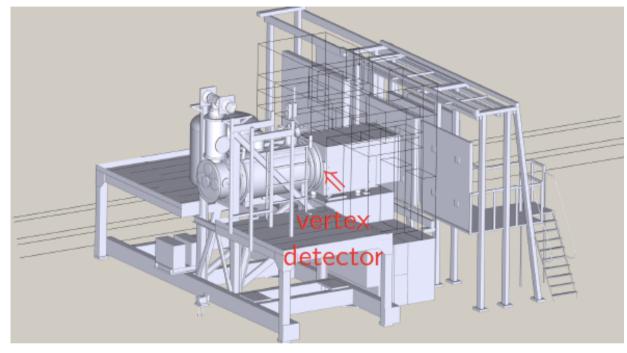
Oleg Denisov

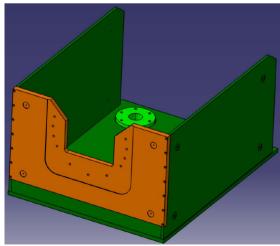


Drell-Yan experiment lay-out Sci-Fi based Vertex detector

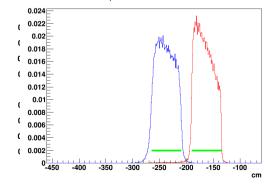




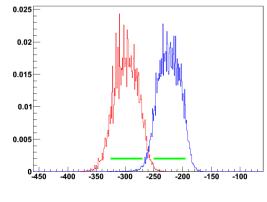




cells separation with vertex detector



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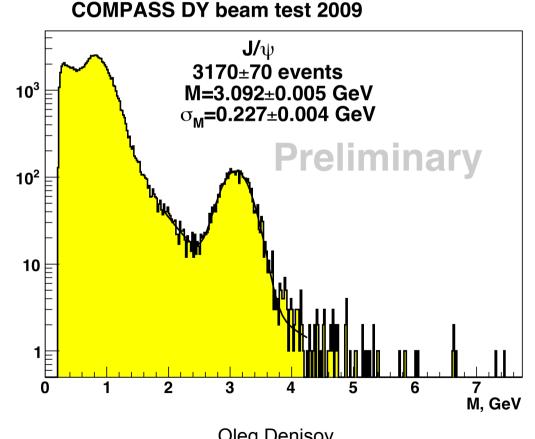




DY@COMPASS - feasibility - Signal



- Expected according to the proposal J/Psi and Drell-Yan yields: 3600±600 and 110±22 (normalized to 2009 beam flux ~3.7 x 10¹¹)
- Measured in 2009 beam test J/Psi yield is 3170±70, and DY yield • is 84±10



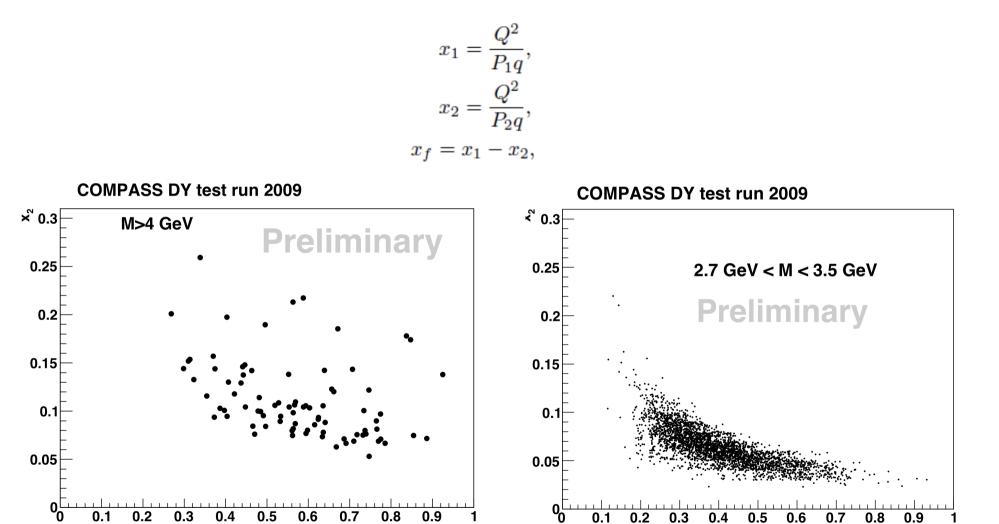


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DY@COMPASS - feasibility - Kinematics I



• Valence quark range for both J/Psi and DY



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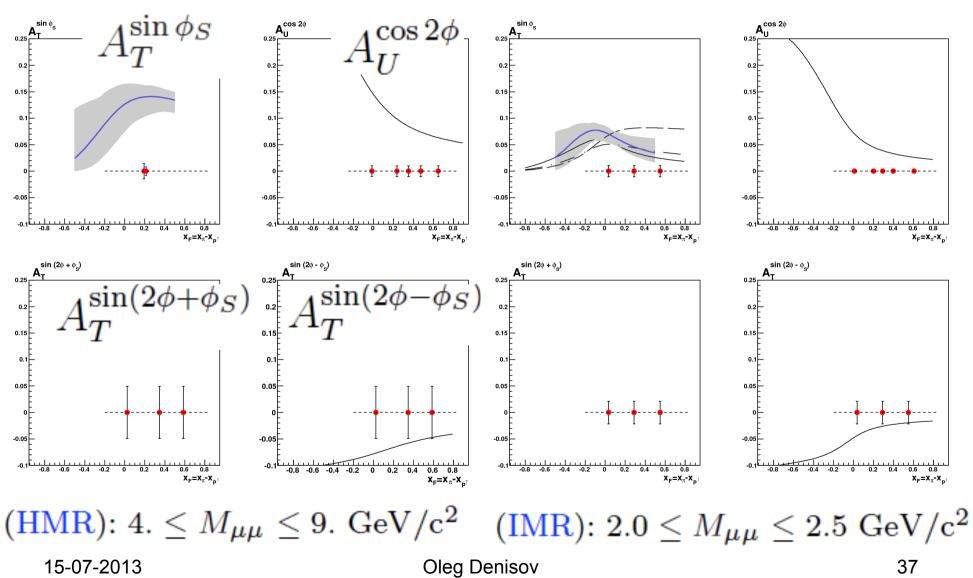
X₁

X₁



DY@COMPASS projections II



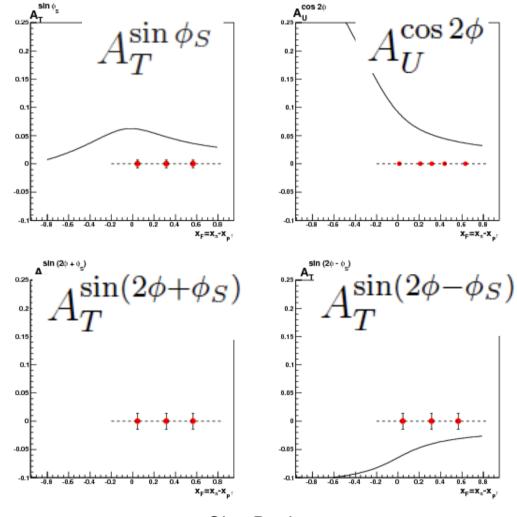




DY@COMPASS projections III



 J/ψ region: $2.9 \le M_{\mu\mu} \le 3.2 \text{ GeV/c}^2$

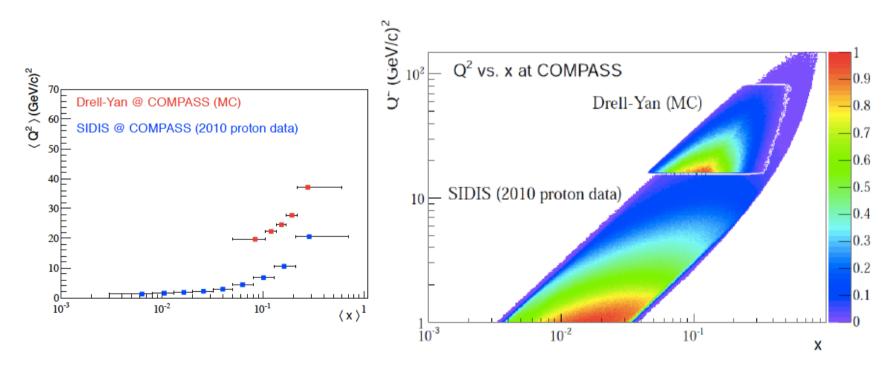


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In COMPASS we have the opportunity to access these TMD PDFs from both DY and SIDIS processes.



There is a phase space overlap between the two measurements.

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Drell-Yan Experiments 5 years Running plan



- 1. Fermilab E-906: data taking will resume in September 2013 and will last for at least 1 year (experiment is approved for 2 years running period)
- 2. Polarised DY at Fermilab hopefully in a few years from now
- 3. COMPASS polarized Drell-Yan measurement will be started in the mid of October 2014, with a short beam test. Physics data taking will take place over the whole 2015.
- 4. A second year of DY data-taking is planned, in case of LS2 delay, in 2018.
- 5. RHIC ?





- 1. Two (almost) running experiments for Drell-Yan physics after long no data period
- 2. 100% complementary experiments
- 3. What is missing collider experiment, DY@NICA is very welcome:
 - 4π geometry access to the exclusive channels (no absorber)
 - Collider mode no dilution factor as in case of fixed target solid state PT experiments (factor ~5 in statistical error)
 - We are entered the era of the global fits and global data analysis → NICA data will certainly contribute a lot



DY Experiments: Summary



- Proton and pion induced DY data are coming up in the next 2 years
 - Some statistics with kaons and antiprotons will be available
- Statistical error on single spin asymmetries on the level of 1÷2% will be achieved by the 2016
 - A lots of new unique data is just behind the corner



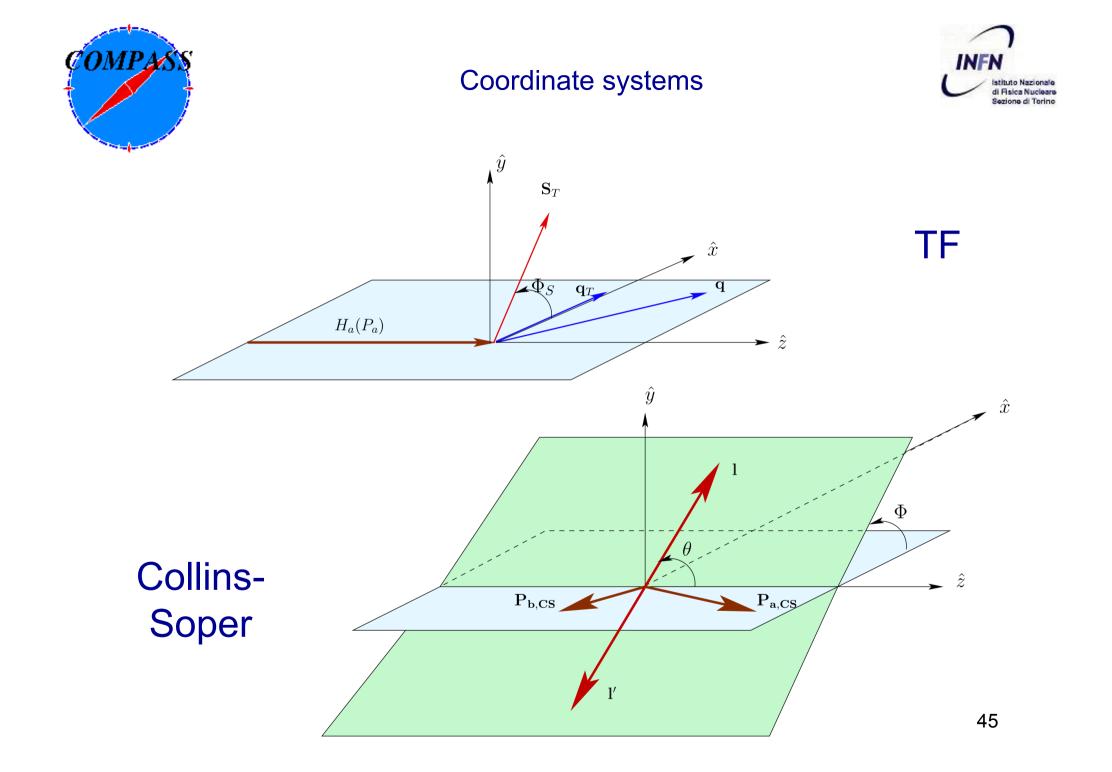


Spares





- 2010 COMPASS polarised SIDIS data (Sivers, transversity via global data fit)
- 2010 2013? E906 (SeaQuest) pp Drell-Yan Boer-Mulders of the proton
- 2013 2016 COMPASS polarised Drell-Yan pi-p data TMDs universality and T-odd TMDs sign change SIDIS ← → DY (for Boer-Mulders function study the input from E906 as well as new transversity fit from the global data analysis is very welcome)
- 2015 → …… RHIC, NICA pp (un)polarised DY data very welcome – complimentary to COMPASS
- 2017 → more COMPASS data, antiprotons?.....
- MANY NEW data just behind the corner





Drell-Yan Workshop at CERN, April 26-27





Studying the hadron structure in Drell-Yan reactions

26-27 April 2010 CERN

Overview	Since a long time the Drell-Yan (DY) process is considered to be a powerful tool to study hadron structure. In the past, several experiments were successfully carried out using unpolarised beams and targets. Nowadays, taking into account
Programme	the much advanced understanding of the spin structure of the nucleon, we are
Registration	discussing a new generation of DY measurements using polarised beams and/or targets.
Registration Form	
List of registrants	The COMPASS collaboration is currently preparing a proposal for future studies of nucleon structure beyond 2011. One of the main aims is a first measurement of
Laptop and Wireless access	transverse-momentum-dependent parton distributions (TMDs) using the Drell-Yan process on a transversely polarised proton target hit by a pion beam. Among
Access Cards	the distributions to be studied are Sivers, Boer-Mulders and pretzelosity TMDs as well as transversely polarised quark distributions.
Accomodation	The workshop will review ongoing theoretical and experimental efforts related
How to get to CERN	to the Drell-Yan process. Detailed presentations and discussions of the theoretical aspects will be complemented by descriptions of planned fixed-target and collider

Support

Organizers: Paula Bordalo (LIP-Lisbon and IST/UTL) Oleg Denisov (CERN/INFN-Torino) Eva-Maria Kabuss (Mainz) Fabienne Kunne (CEA Saclay) Alain Magnon (CEA Saclay) Gerhard Mallot (CERN) Anna Martin (Univ. Trieste and INFN-Trieste) Wolf-Dieter Nowak (CERN) Daniele Panzieri (Univ. Alessandria and INFN-Torino)

Dates: from 26 April 2010 09:00 to 27 April 2010 18:00

Location: CERN

experiments.

Salle Andersson Room: 40-S2-A01

15-07-2013

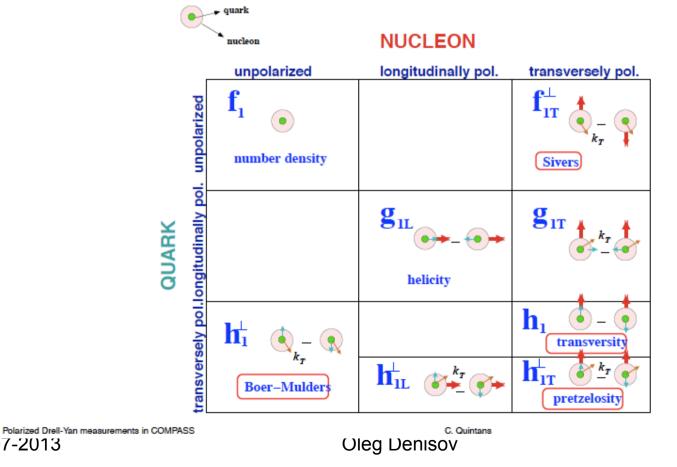


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Parton distribution functions

Taking into account the intrinsic transverse momentum k_T of quarks, at LO 8 PDFs are needed for a full description of the nucleon:



- p.2



WHAT ABOUT A RF SEPARATED BEAM ???



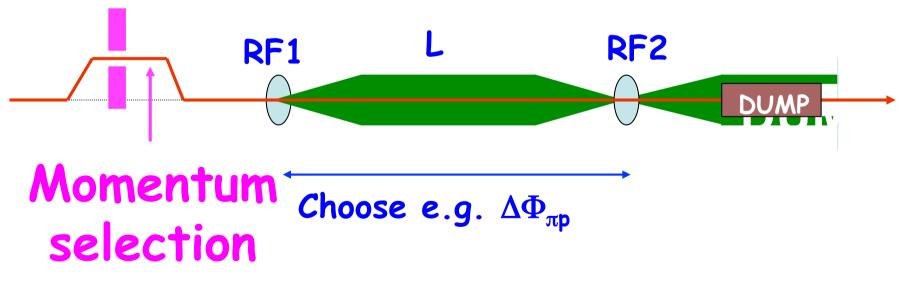
First and very preliminary thoughts, guided by

recent studies for P326

CKM studies by J.Doornbos/TRIUMF, e.g.

http://trshare.triumf.ca/~trjd/rfbeam.ps.gz

E.g. a system with two cavities:



 $\Delta \Phi = 2\pi (L f / c) (\beta_1^{-1} - \beta_2^{-1}) \text{ with } \beta_1^{-1} - \beta_2^{-1} = (m_1^2 - m_2^2)/2p^2$

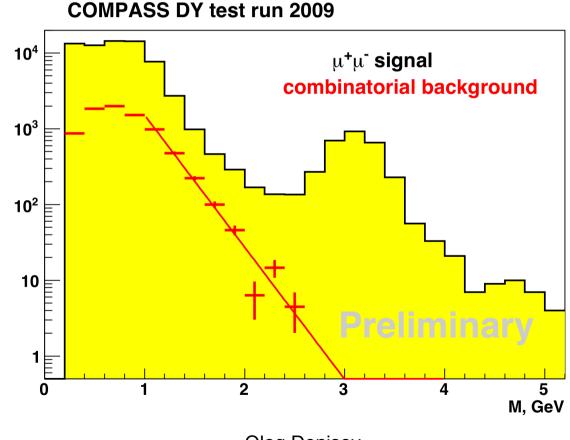
Preliminary rate estimates for RF separated antiproton beams





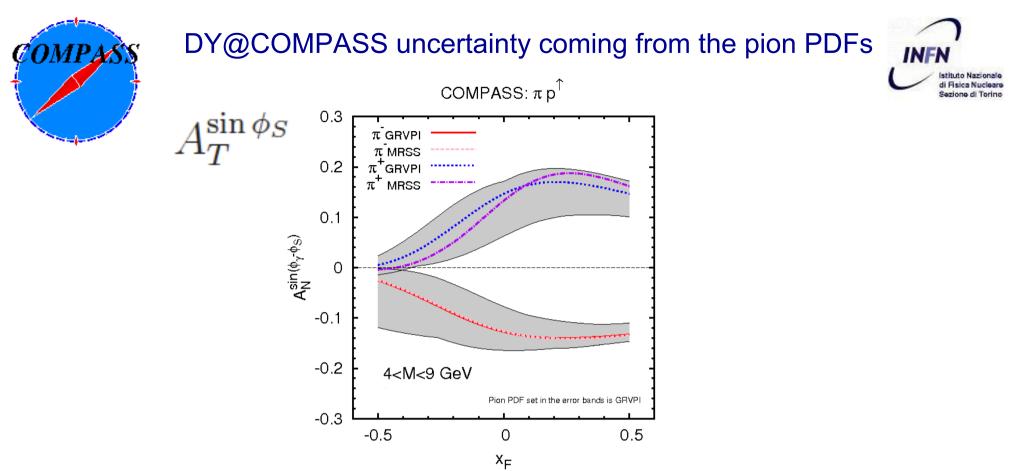


- 2009 beam test id very important
- Combinatorial background suppressed by ~10 at 2.0 GeV/c dimuon invariant mass (beam intensity ~8 times lower wrt Proposal)



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In case of $\pi^- p$ scattering the valence pion \bar{u} unpolarised PDF is well known and there is no difference between two pdf sets. In case of $\pi^+ p$ scattering there is a little contamination coming from sea \bar{u} of the pion, which annihilates with valence u quark of the proton, because the distribution functions are weighted in the cross section with e_q^2 , and the $\bar{u}u$ contribution is multiplied by factor 4/9 while the $\bar{d}d$ by factor 1/9. Thus, the contribution from the sea \bar{u} of the pion can not be neglected, it is less known with respect to valence PDFs and it explains the difference from one data set (GRVPI) to another (MRSS).

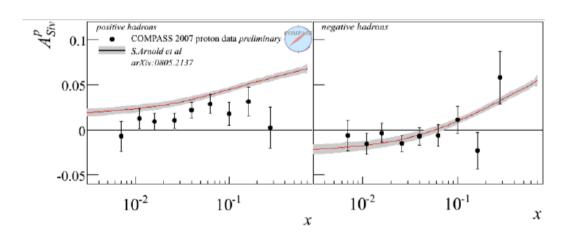
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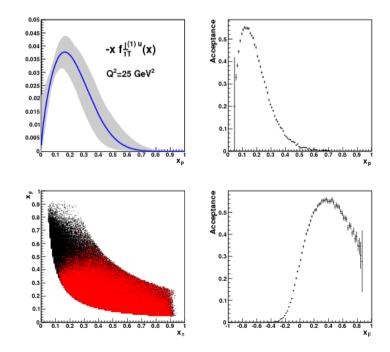




DY@COMPASS → SIDIS complemenatrity

- TMD PDFs study in SIDIS is an important part of the COMPASS-I program
- COMPASS-II, TMDs study in Drell-Yan processes:
 - We change the probe (elementary process)
 - We upgrade the spectrometer and we change its lay-out
 - We change the kinematic range









$$\delta A = \frac{1}{P_b f} \frac{1}{\sqrt{N_{sig}}} \sqrt{1 + \frac{N_{sig}}{N_{backg}}} \quad \tau = x_a x_b = M^2 / s$$

1. Drell-Yan experiments:

- High luminosity (DY Cross Section is a fractions of nanobarns) and large angular acceptance, better pion or antiproton beams (valence anti-quark)
- Sufficiently high energy to access 'safe' of background free M_{\parallel} range (4 GeV/c < M_{\parallel} < 9 GeV/c)
- Good acceptance in the valence quark range $x_B > 0.05$ and kinematic range: $\tau = x_A x_B = M^2/s > 0.1$
- 2. Polarised Drell-Yan:
 - Good factor of merit (F_m), which can be represented as a product of the luminosity and beam (target) polarisation (dilution factor) ($F_m \sim L \times P_{beam}$ (f))

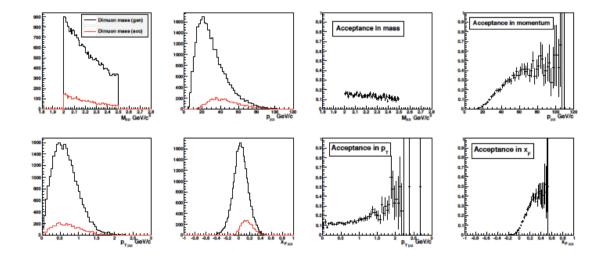






- Calculated by MC
- Negligible in both HM and IM ranges (~15% contribution in IM)

Acceptance for open-charm 2.0 - 2.5 GeV/c^2



As in the IMR the acceptances are 14% for open-charm and 43% for DY, the ratio of observable events in the dimuon mass spectra will be $N_{D\bar{D}}/N_{DY} = (5.47 \times 0.14)/(12.46 \times 0.43) = 0.14$.