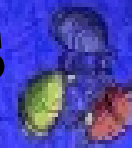


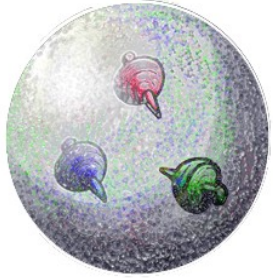
Hadron production from μ - d scattering at $\sqrt{s} = 17\text{GeV}$ at Compass



$$\mu^+d \rightarrow \mu^+h^\pm X$$

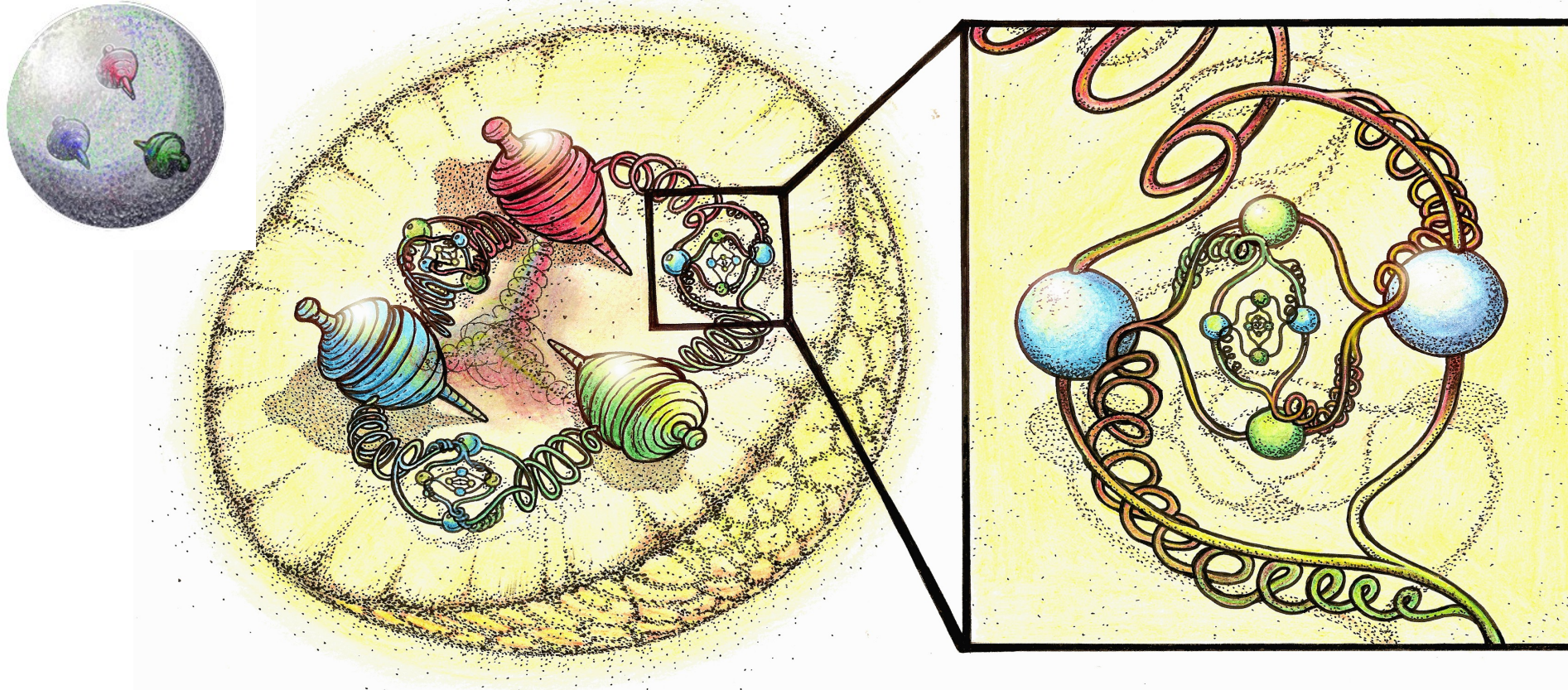


The QCD Proton Picture



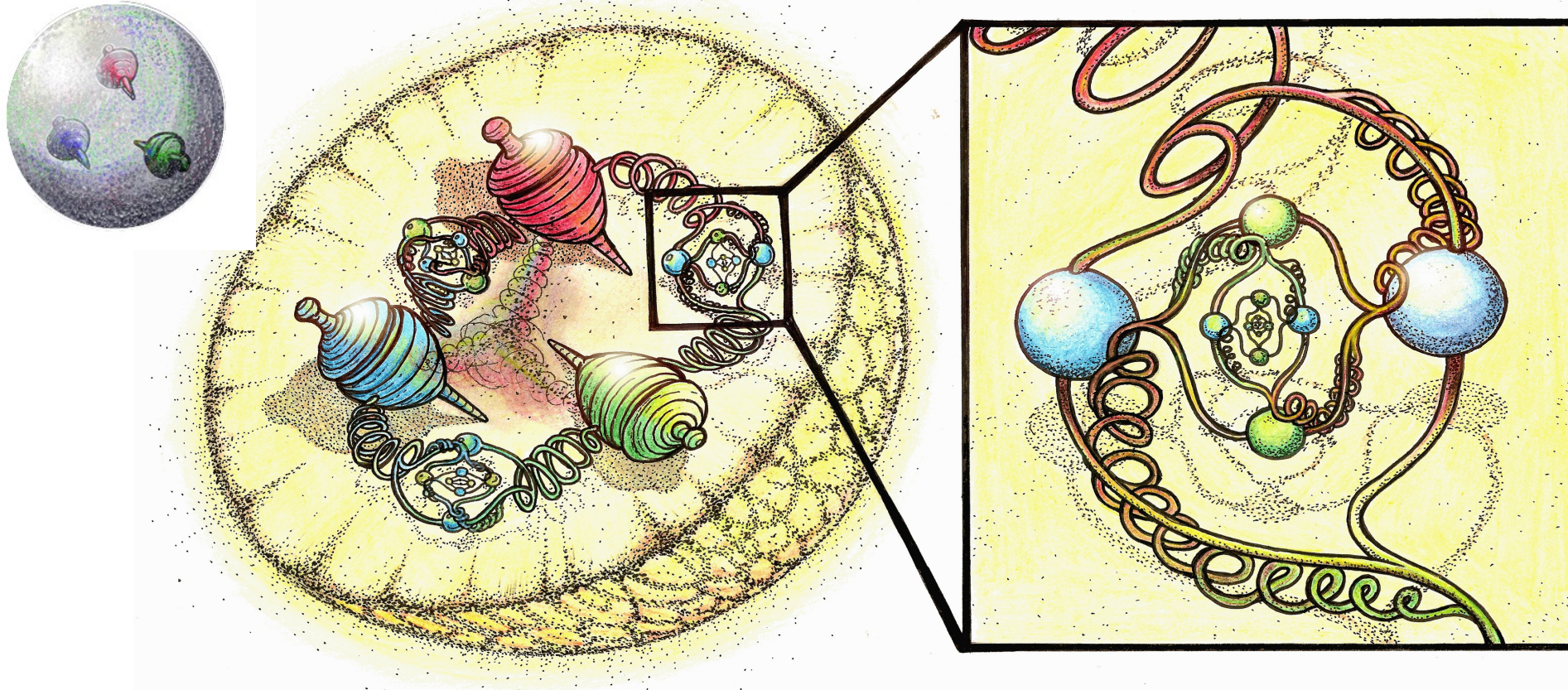
That nucleon has a large anomalous magnetic moment proves that this is not *fundamental* spin1/2 Dirac particle.

The QCD Proton Picture



That nucleon has a large anomalous magnetic moment proves that this is not **fundamental** spin $1/2$ Dirac particle.
Nucleon Spin is Subtle: Quarks, gluons and their angular momentum caused by their high speed motion within the nucleon are contributors to the Nucleon's spin.

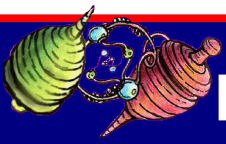
The QCD Proton Picture



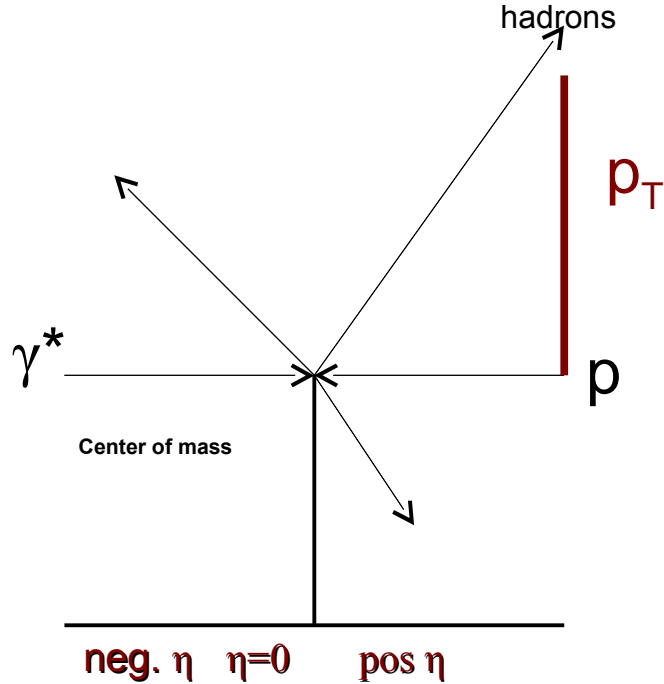
That nucleon has a large anomalous magnetic moment proves that this is not **fundamental** spin $1/2$ Dirac particle.

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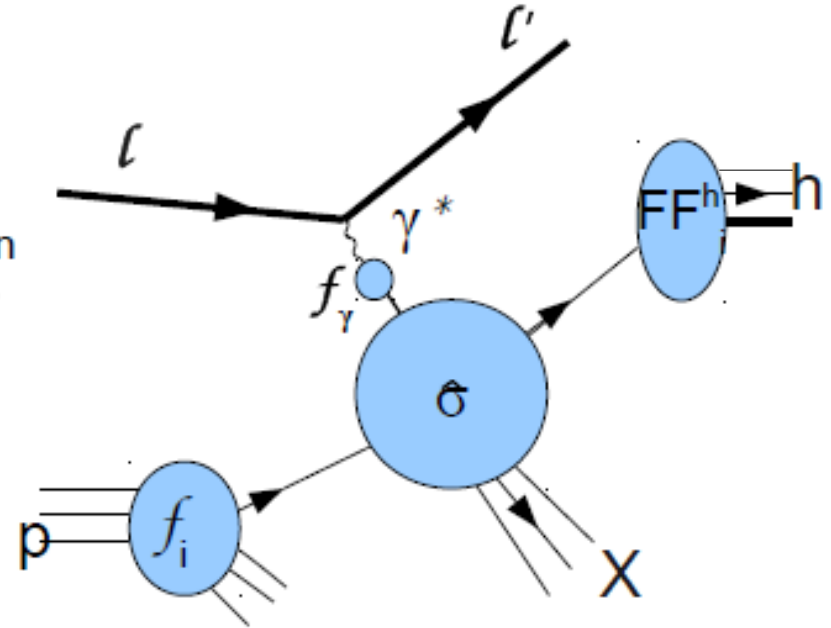
Spin physics has opened a box full of questions about matter, and it has also laid the groundwork to a plethora of scientific advancements: from the medical field, to astronomy research.



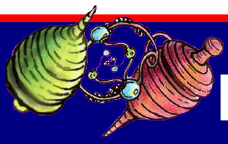
How can we access polarized parton information?



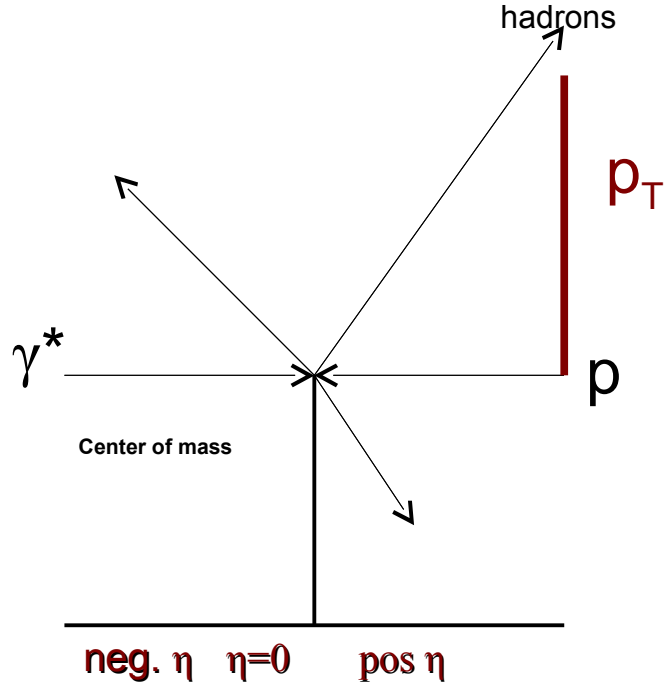
Factorization
Theorem



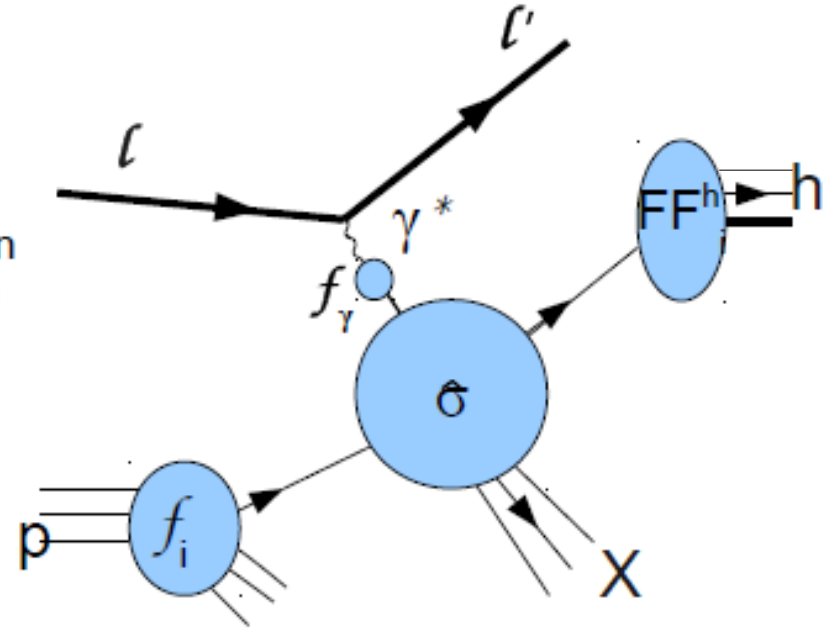
$$A_{LL} = \frac{\sigma_{++} - \sigma_{+-}}{\sigma_{++} + \sigma_{+-}} = \frac{\sum_{i=q,\bar{q},g} \Delta f_\gamma \otimes \Delta f_i \otimes \Delta \hat{\sigma} \otimes FF_{h/i}}{\sum_{i=q,\bar{q},g} f_\gamma \otimes f_i \otimes \hat{\sigma} \otimes FF_{h/i}}$$



How can we access polarized parton information?



Factorization
Theorem



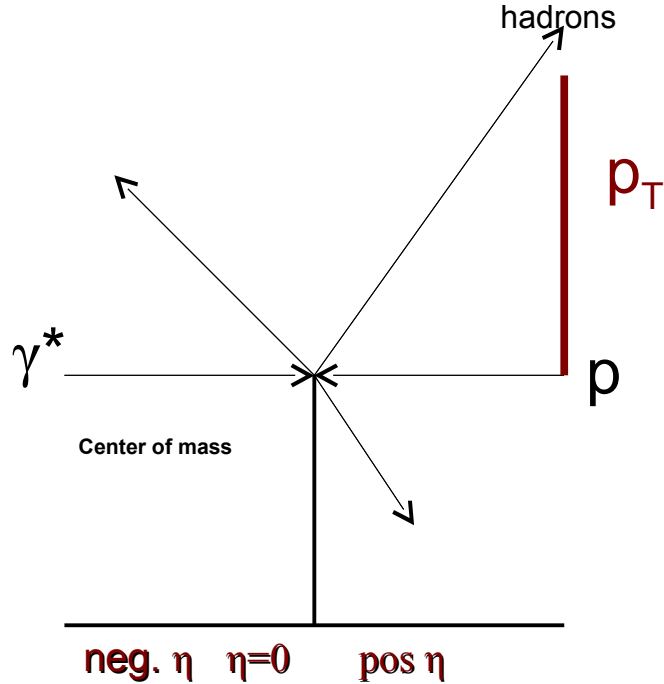
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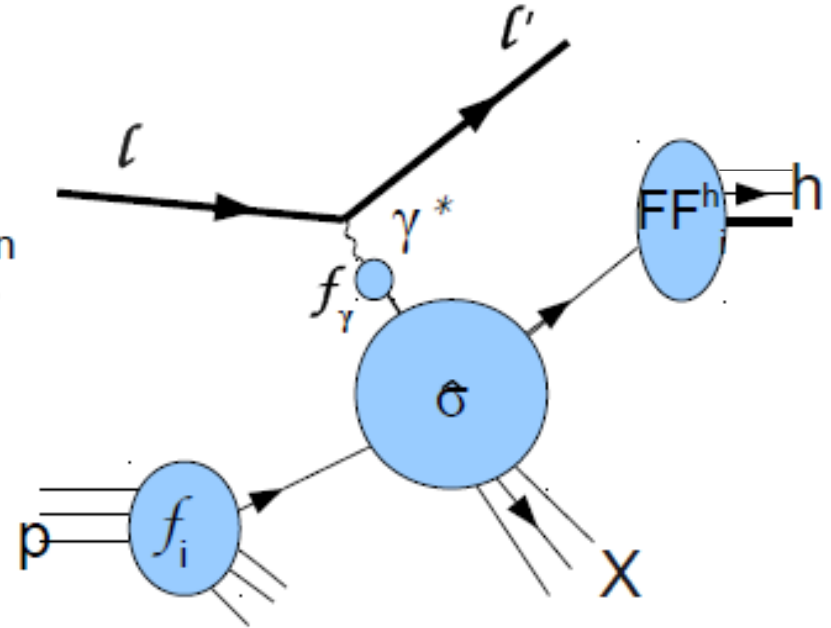
- (N) Number of events
- (F) Dilution Factor
- (P) Polarization in beam or target



How can we access polarized parton information?



Factorization
Theorem



$$A_{LL} = \frac{\sigma_{++} - \sigma_{+-}}{\sigma_{++} + \sigma_{+-}} = \frac{\sum_{i=q,\bar{q},g} \Delta f_\gamma \otimes \Delta f_i \otimes \Delta \hat{\sigma} \otimes FF_{h/i}}{\sum_{i=q,\bar{q},g} f_\gamma \otimes f_i \otimes \hat{\sigma} \otimes FF_{h/i}}$$

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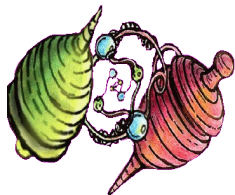
- (N) Number of events
- (F) Dilution Factor
- (P) Polarization in beam or target

Spin Asymmetries give us access to **polarized parton distribution functions**

What motivates the Measurement of an Unpolarized Cross-section?

$$A_{LL} = \frac{\sigma_{+++} - \sigma_{+-}}{\sigma_{+++} + \sigma_{+-}} = \frac{\sum_{i=q,\bar{q},g} \Delta f_{\gamma} \otimes \Delta f_i \otimes \Delta \hat{\sigma} \otimes FF_{h/i}}{\sum_{i=q,\bar{q},g} f_{\gamma} \otimes f_i \otimes \hat{\sigma} \otimes FF_{h/i}}$$

1. Measure the unpolarized differential Cross Section to confirm that the pQCD framework is applicable to data.



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What motivates the Measurement of an Unpolarized Cross-section?

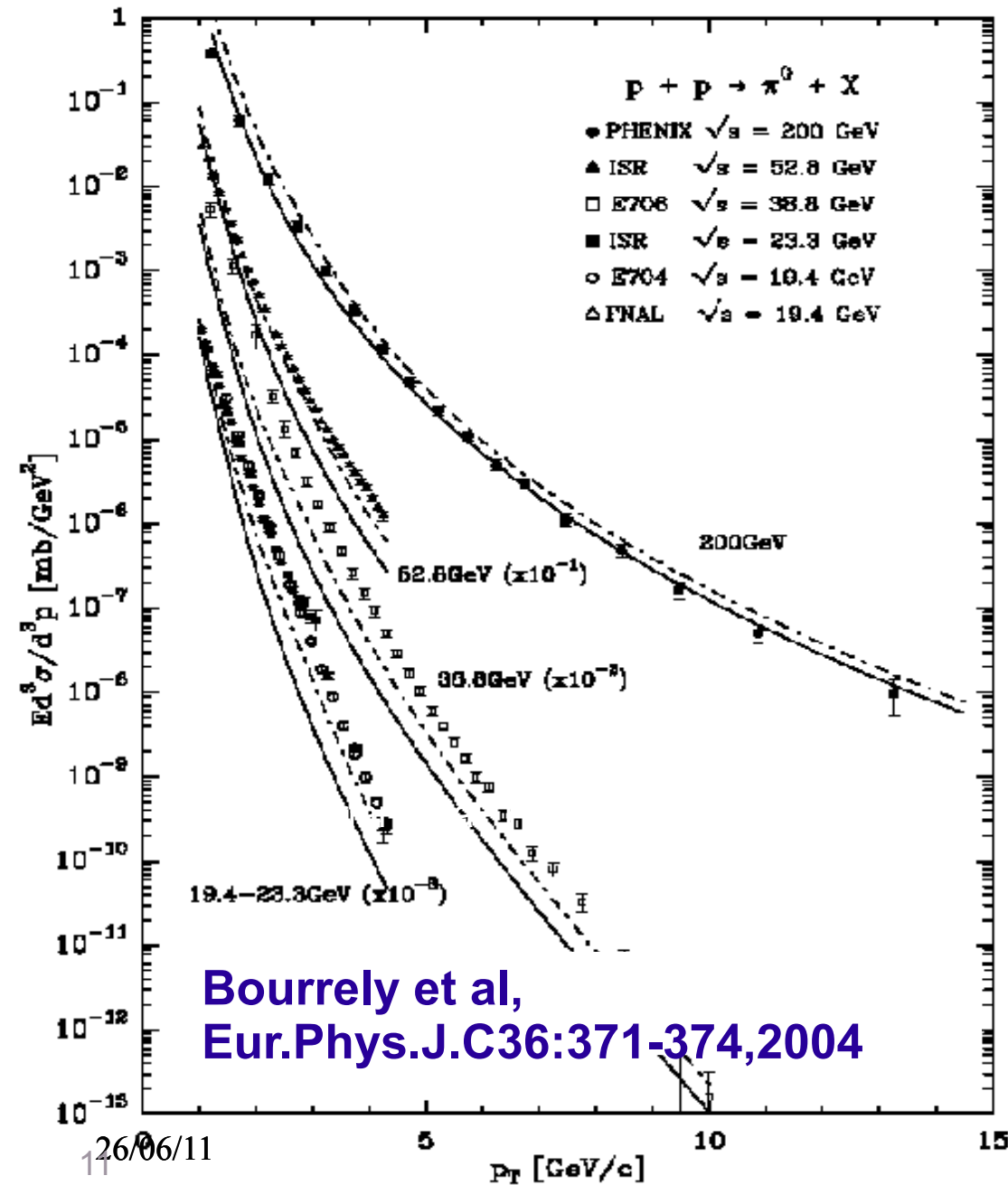
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→ The pQCD unpolarized cross section

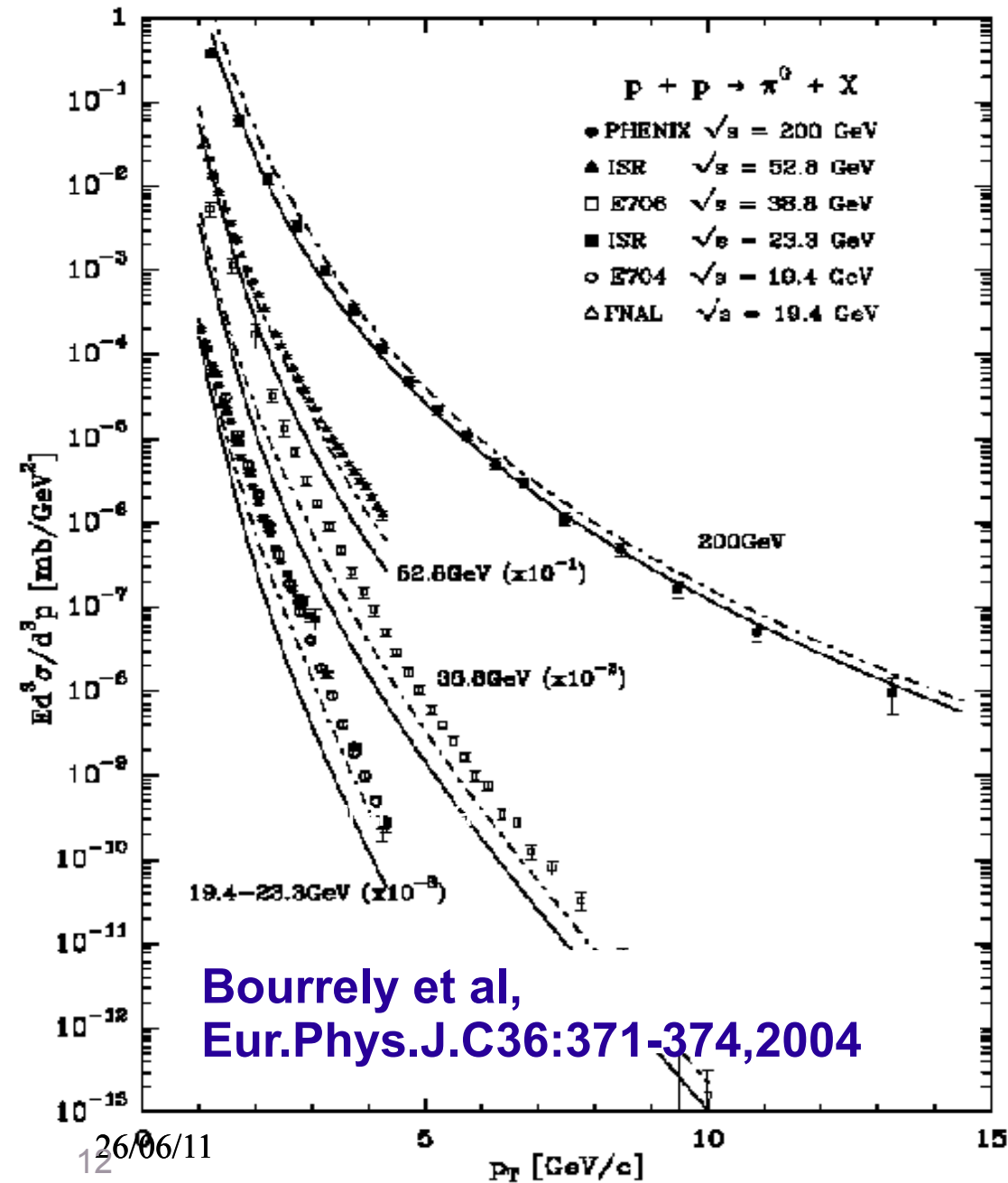
Pion Cross-Sections compared to pQCD: The p+p case



π^0 Cross-Sections at $\eta_{\text{cm}} = 0$
Experiment vs Theory in p+p

→ $\sqrt{s}=200$ GeV (RHIC) Agreement?

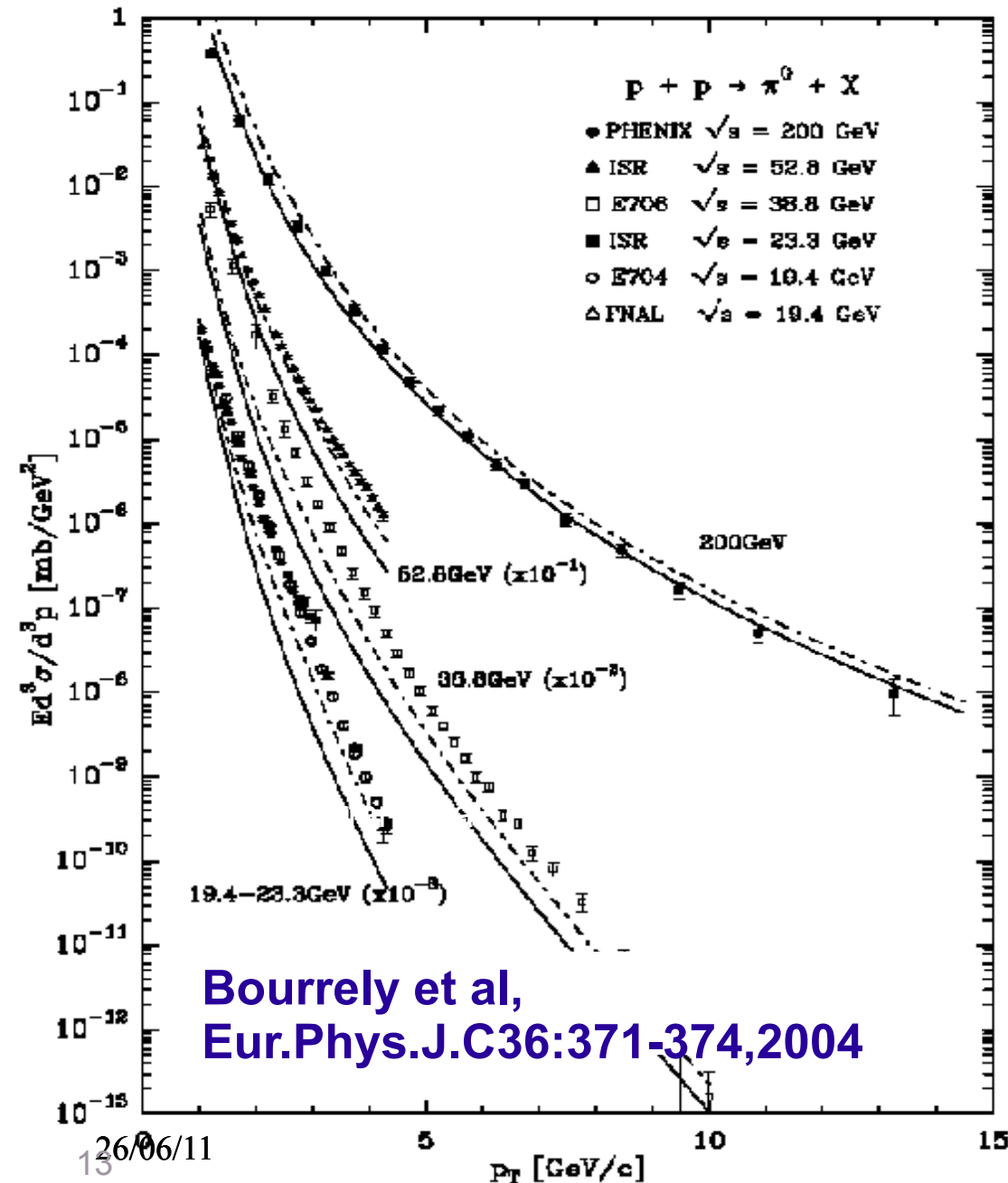
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- $\sqrt{s}=52.8$ GeV (ISR) Factor 2 disagreement

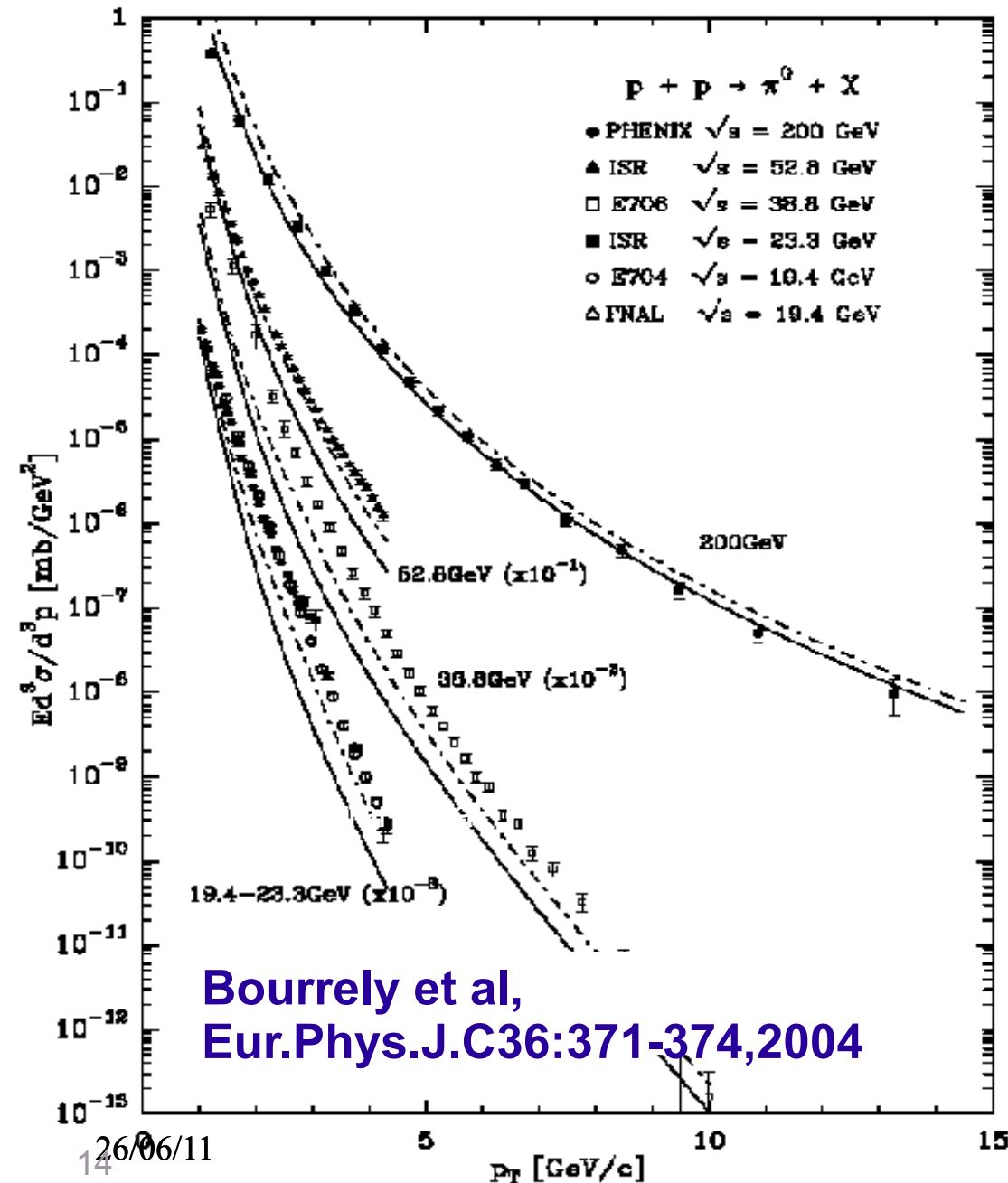
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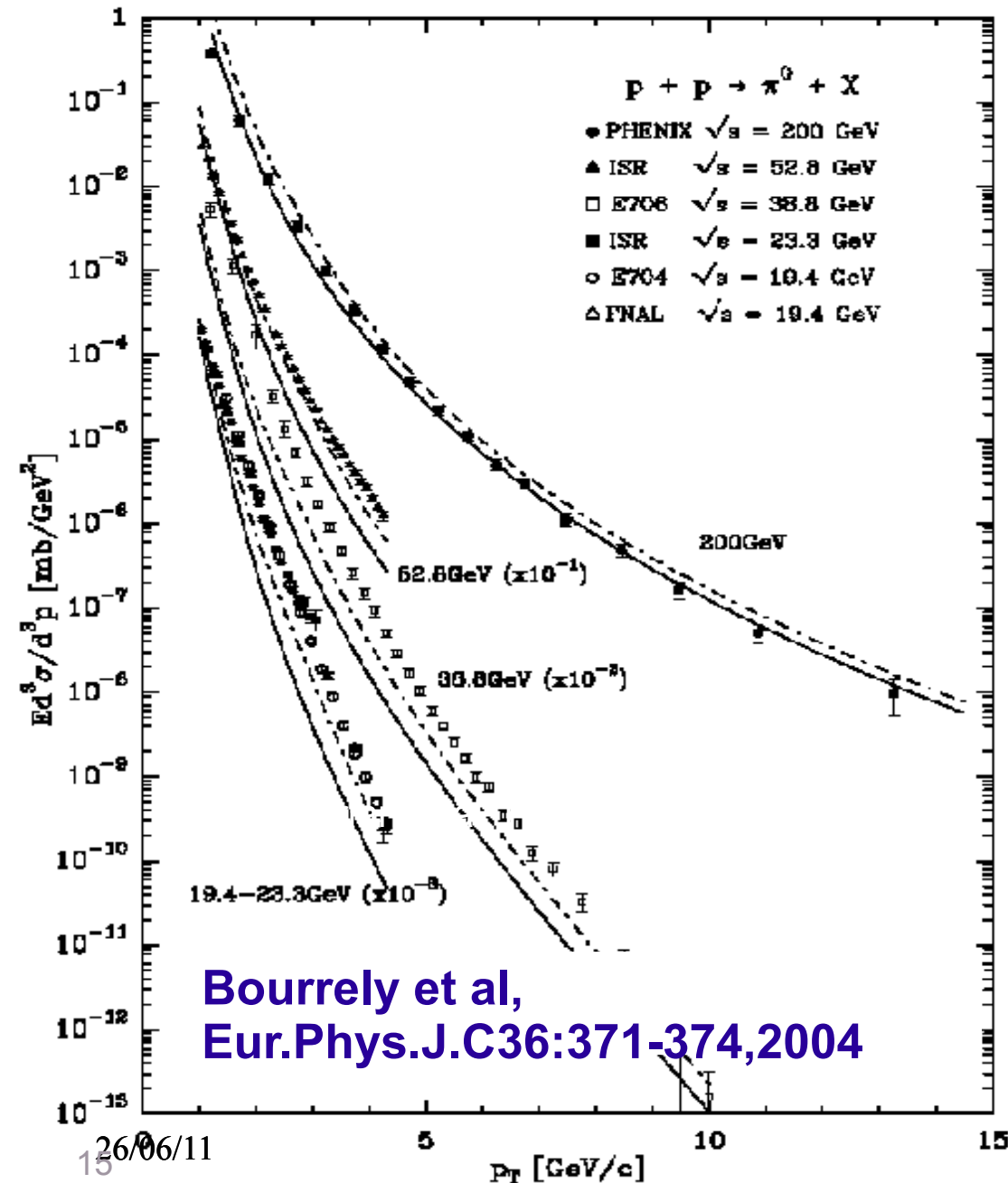
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Disagreement at lower center of mass energies observed in p+p data

Work needed to understand lower energies?

Pion Cross-Sections compared to pQCD: The p+p case



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Disagreement at lower center of mass energies observed in p+p data

Work needed to understand lower energies?

What about $\mu+d \rightarrow \mu'+hX$?

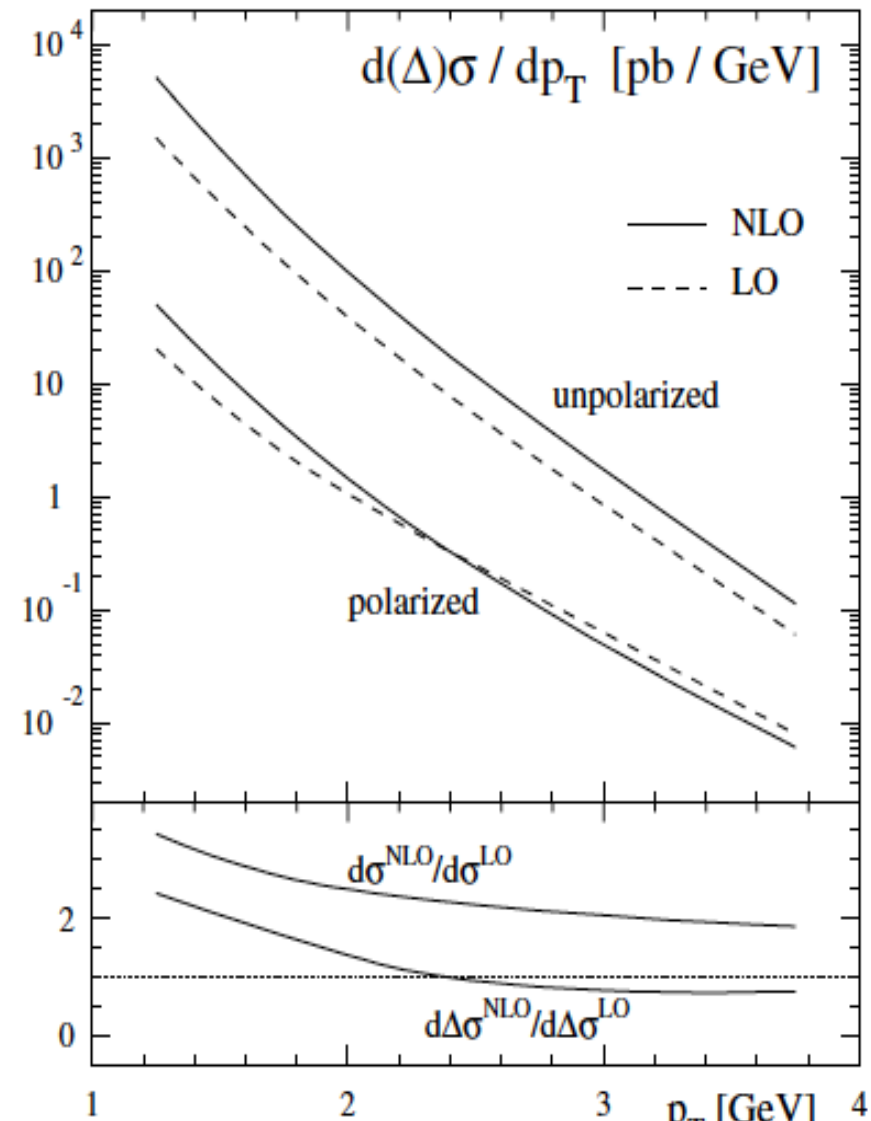
→ Quasi-real Photo Production

The quasi-real photo production case: NLO pQCD.

Eur. Phys. J. C 36, 371–374 (2004)

→ π^0 production prediction:
difficult to compare at the time.

→ The updated calculations for
charged hadron production at
Compass' full kinematics
have been performed by **V.
Wogelsang**



Unpolarized and polarized cross-sections
at LO and NLO

B. Jaeger, M. Stratmann, V. Wogelsang

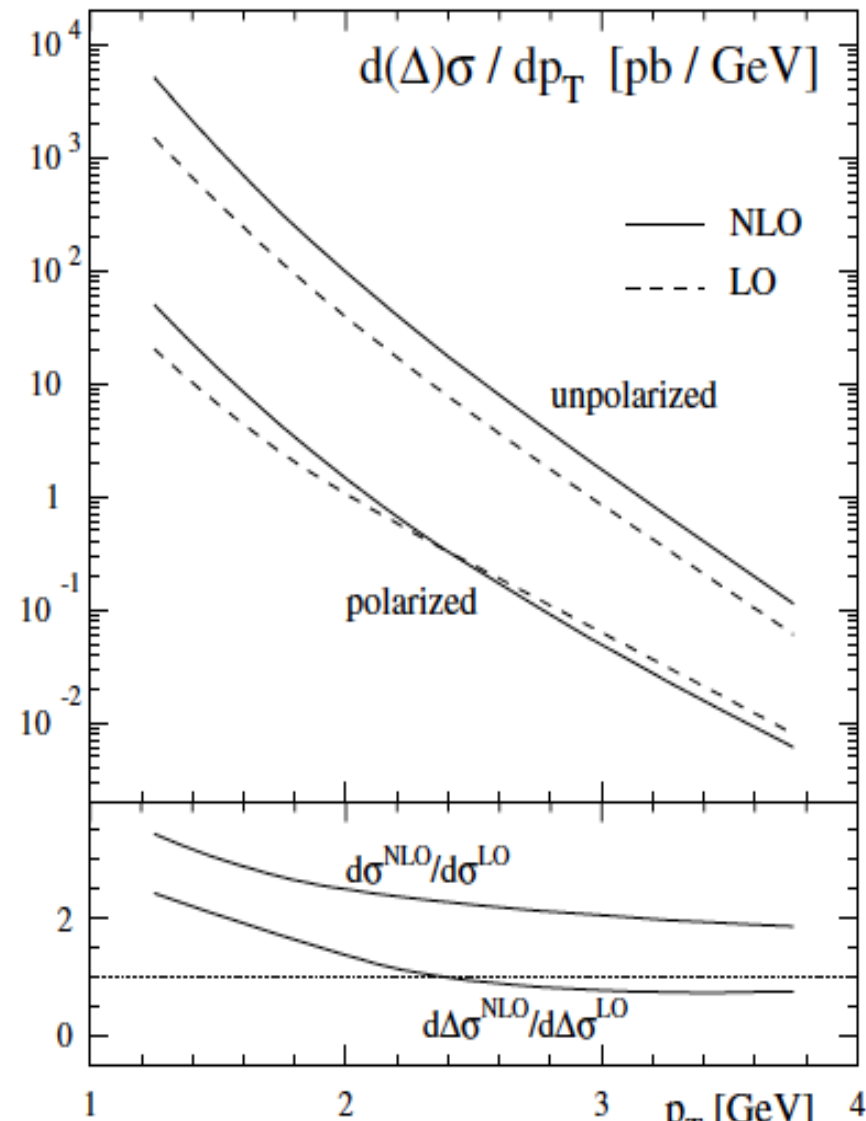
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→ Radiative contributions calculatic
however small, has been calculatec
A. Afanasiev.



Unpolarized and polarized cross-sections
at LO and NLO

B. Jaeger, M. Stratmann, V. Wogelsang

Unpolarized Cross sections: The Measurement

$$\frac{1}{2\pi p_T} \cdot \frac{d^2\sigma}{dp_T dy} \Rightarrow \frac{1}{L^*} \cdot \frac{1}{2\pi p_T} \cdot \frac{d^2N}{dp_T dy} \Rightarrow \frac{1}{2\pi p_T} \frac{N_{h^\pm}(p_T)}{\int L dt \epsilon_{Acc} \Delta p_T \Delta \eta}$$

$\int L dt$ → Integrated luminosity: 142.4 pb⁻¹

$N_{h^\pm}(p_T)$ → Number of hadrons in a p_T bin.

ϵ_{Acc} → Detector's geometrical acceptance, reconstruction algorithm, detection efficiencies.

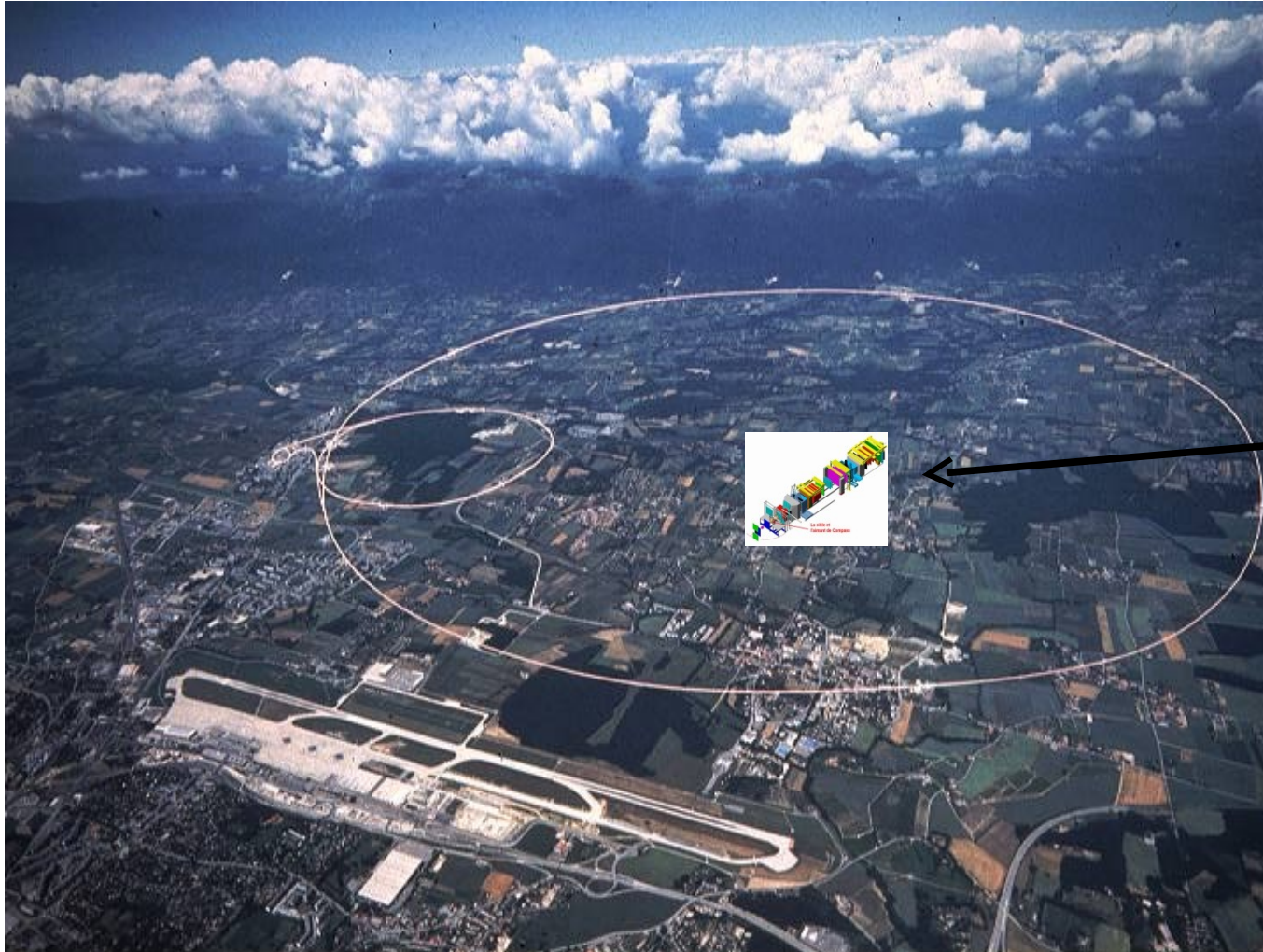
Δp_T → Bin width (250 MeV)

$\Delta \eta$ → Rapidity width (w.r.t. to theory calculation)

The Experiment:

COMPASS @CERN (Prevessin):

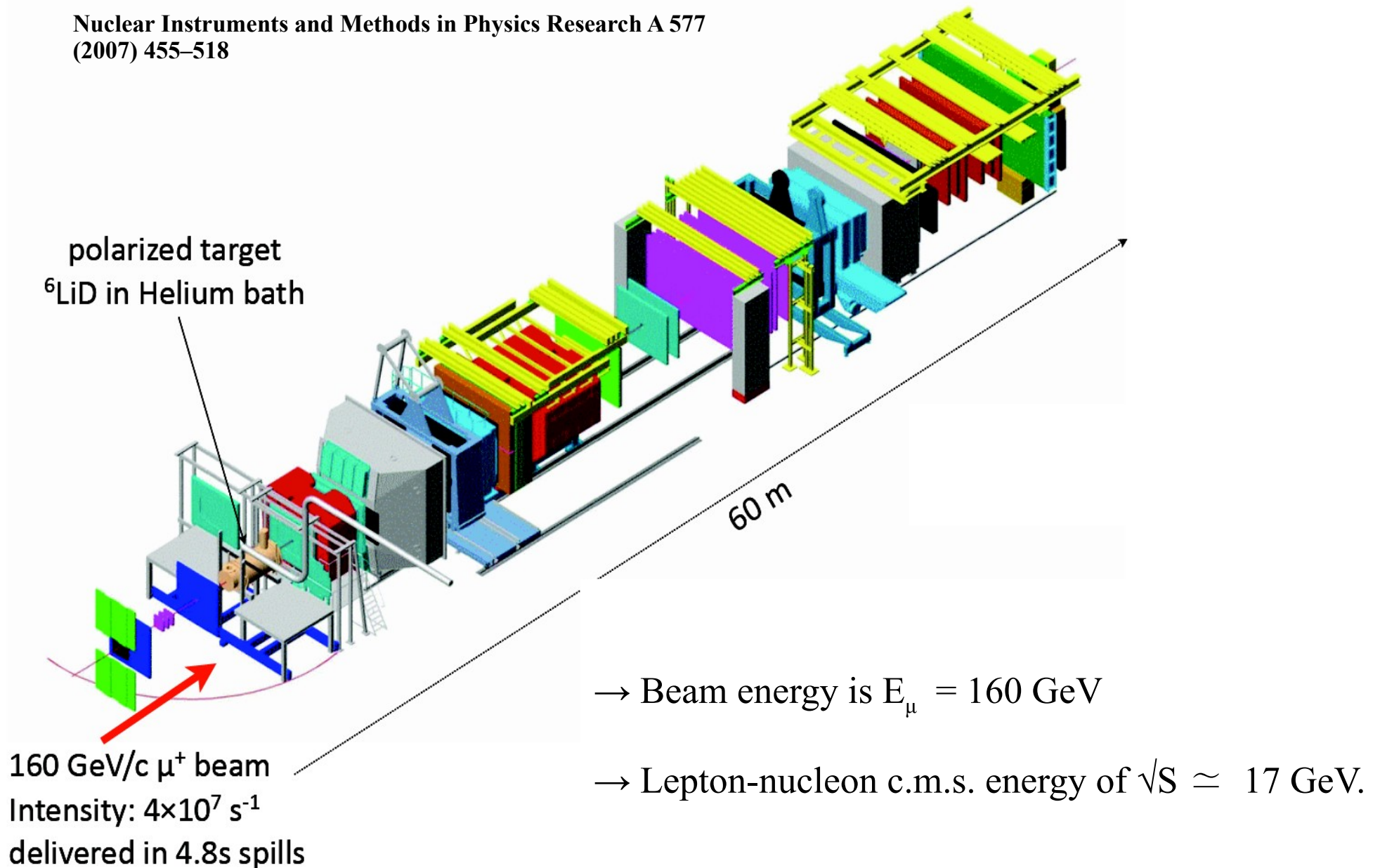
COMMON APPARATUS FOR MUON SPECTROSCOPY:



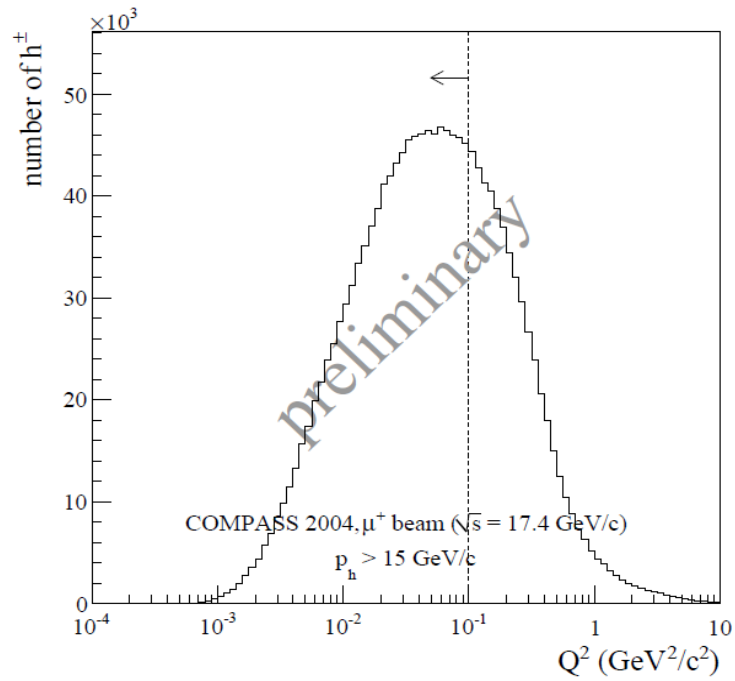
19 Tertiary beam of positive muons produced in the M2 beamline at the CERN SPS.

The COMPASS detector at CERN's SPS

Nuclear Instruments and Methods in Physics Research A 577
(2007) 455–518



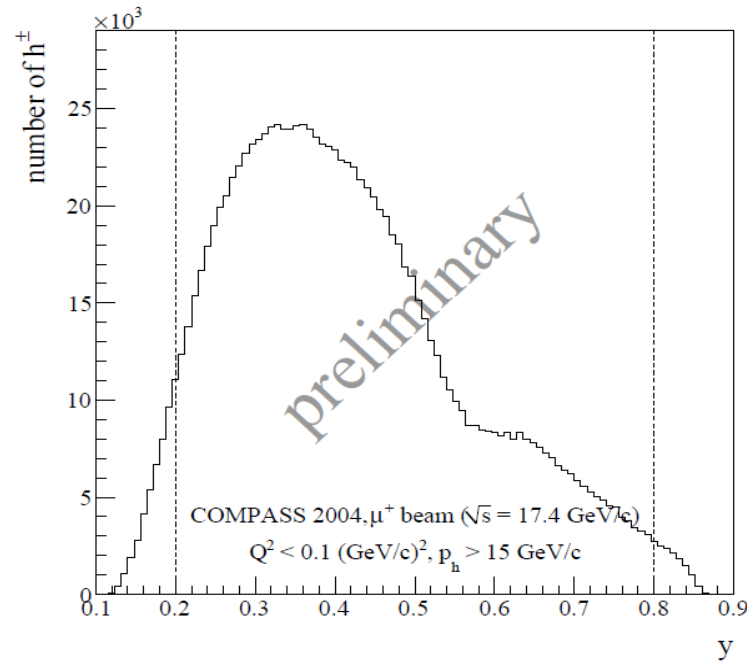
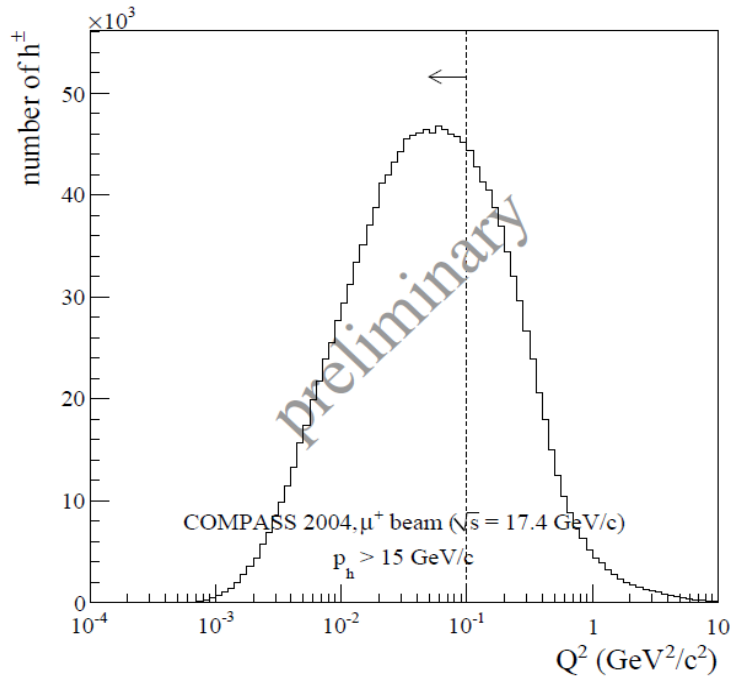
Kinematic Distributions



The kinematic cuts:

- $Q^2 < 0.1 \text{ GeV}^2/c^2$

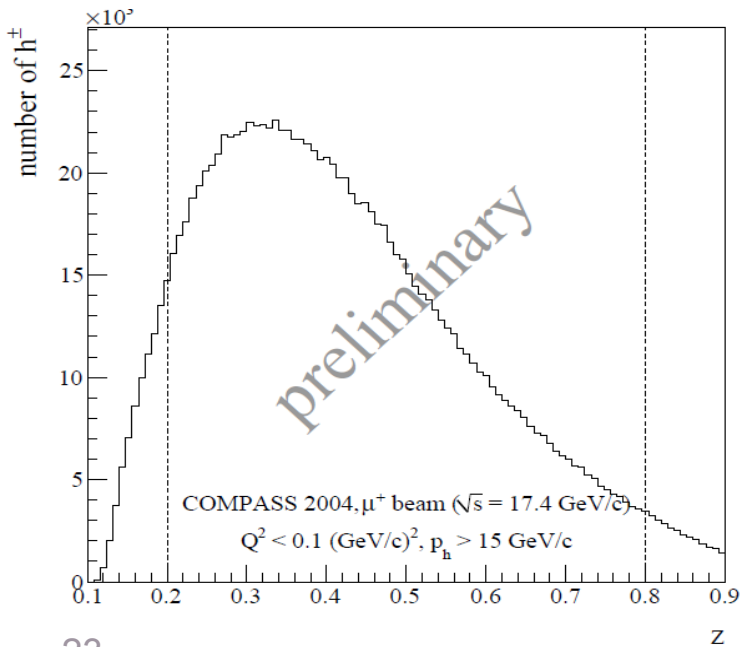
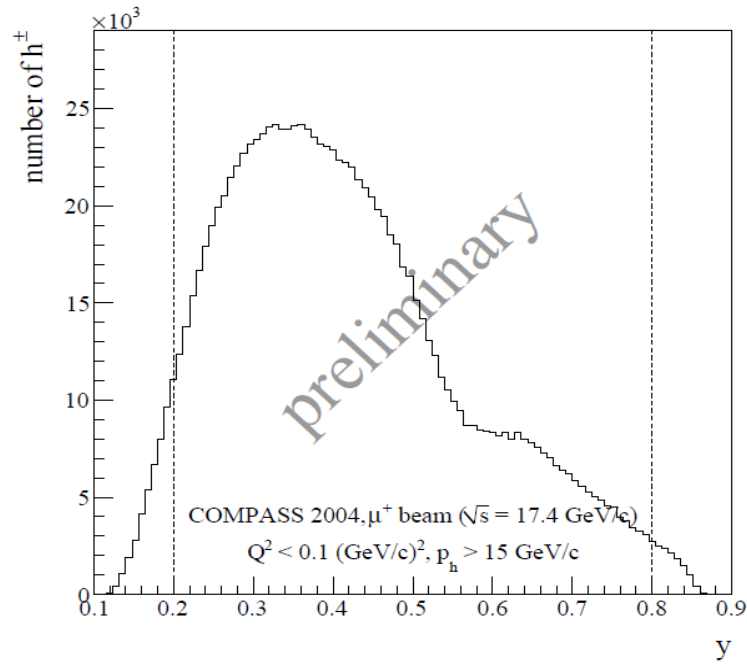
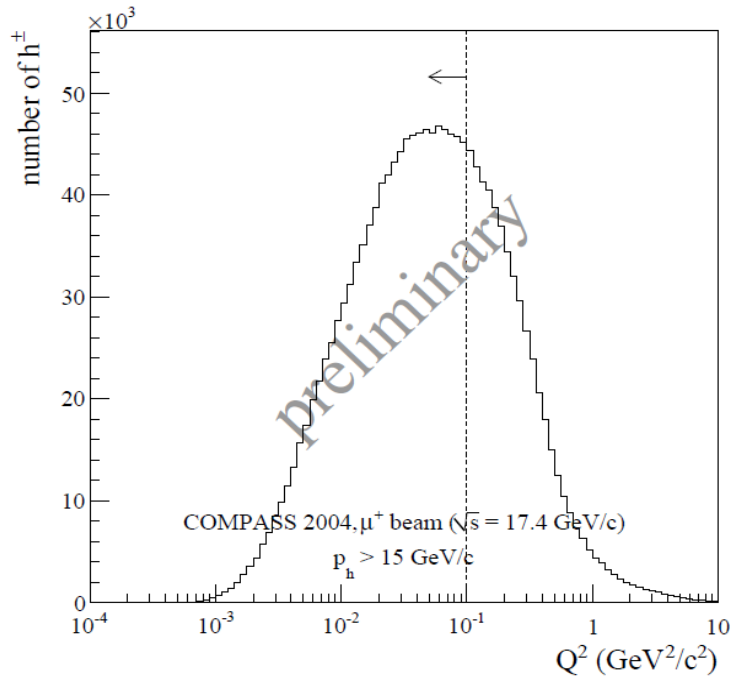
Kinematic Distributions



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- Y: From 0.2 to 0.8

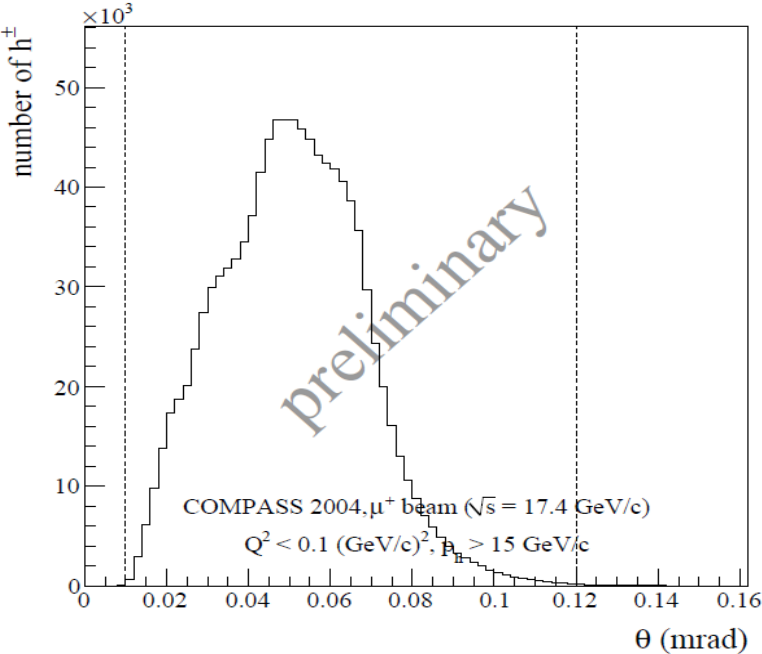
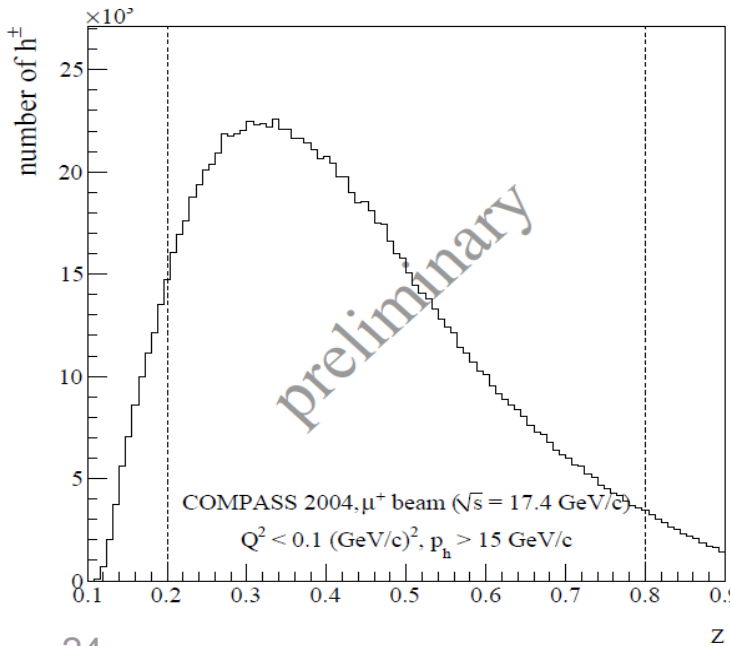
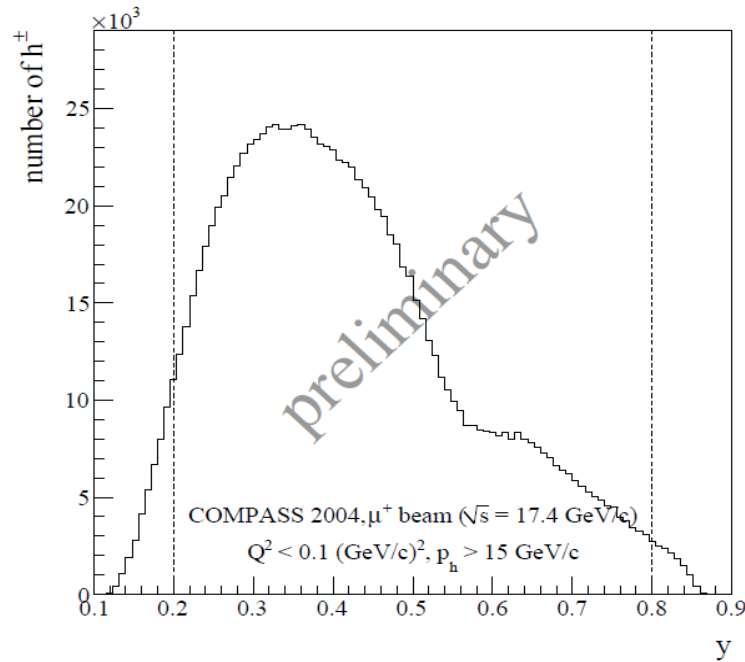
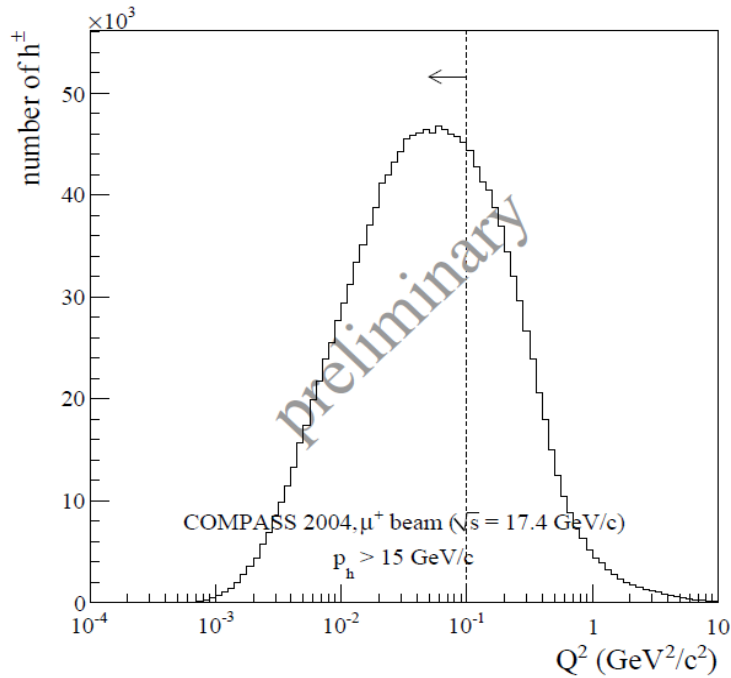
Kinematic Distributions



The kinematic cuts:

- $Q^2 < 0.1$ GeV $^2/c^2$
- Y: From 0.2 to 0.8
- Z: 0.2~0.8

Kinematic Distributions



The kinematic cuts:

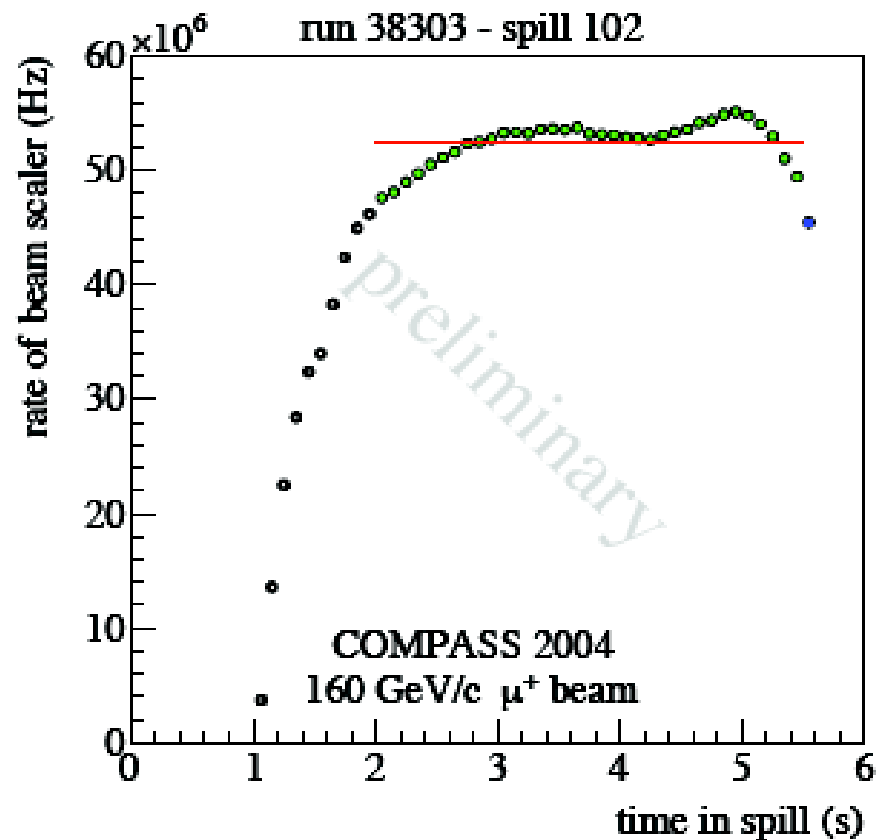
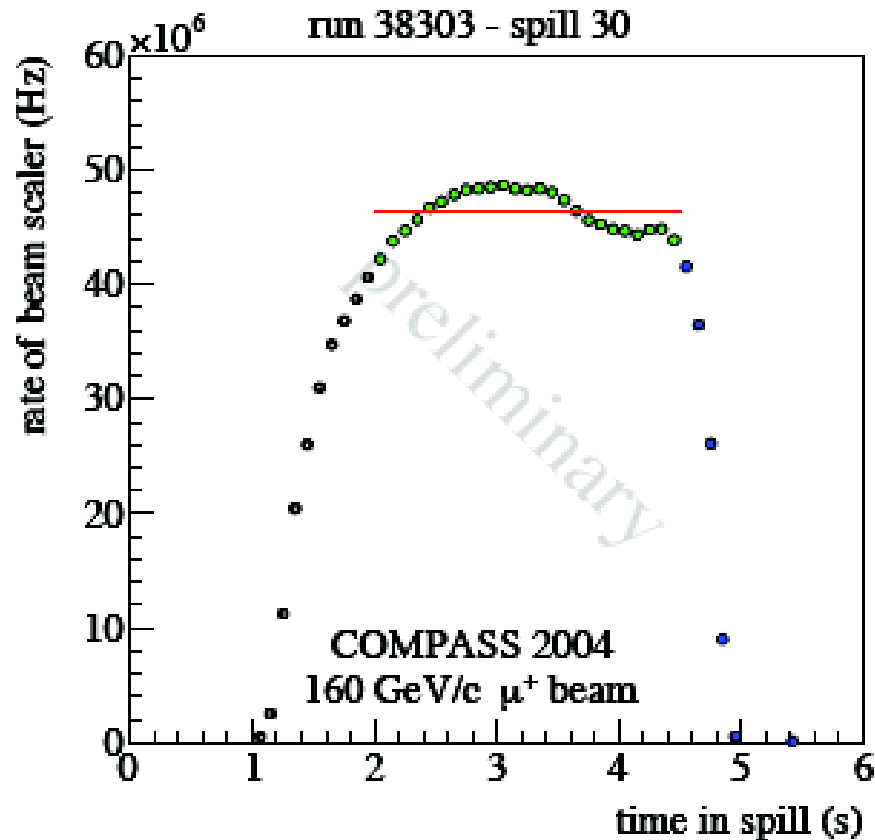
- $Q^2 < 0.1 \text{ GeV}^2/c^2$
- Y : From 0.2 to 0.8
- Z : 0.2~0.8
- Θ :
 20 - 120mrad
 → translates to
 $-0.1 < \eta_{\text{cm}} < 2.4$
- $p_T > 1 \text{ GeV}/c$

Luminosity

- Luminosity is measured using the beam scalers spill by spill.
- Each spill corresponds to the beam delivered by SPS.
 - Every 16.8 seconds (Flatop duration of 4.8 seconds).
- Detector effects which affect the total beam rate are accounted for: Acquisition, veto (beam halo) deadtimes.

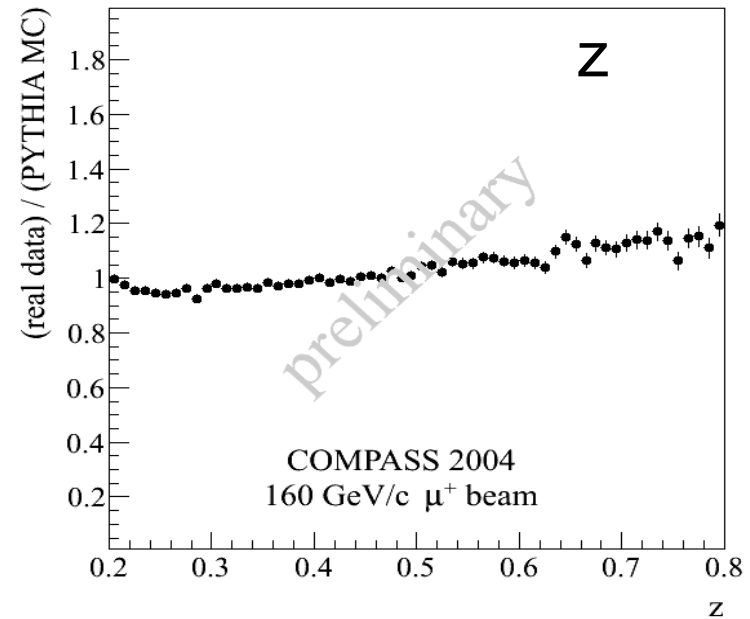
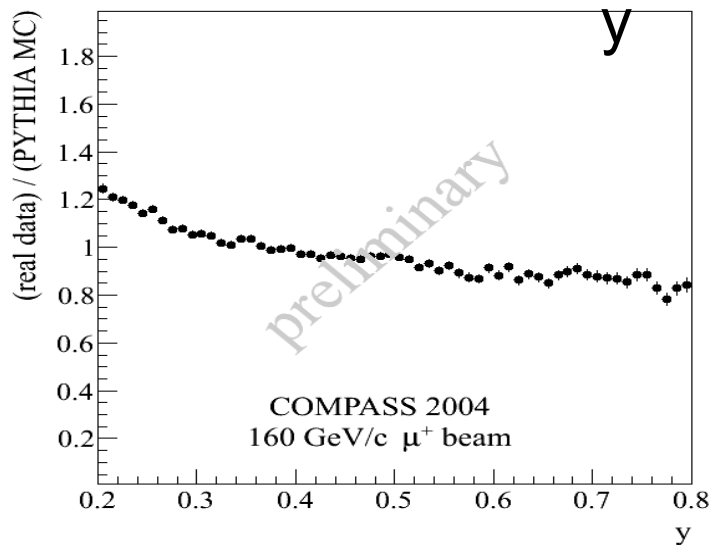
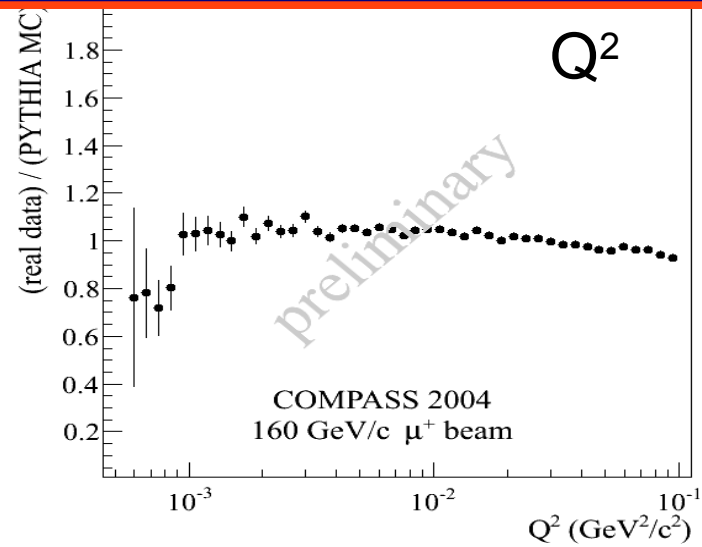
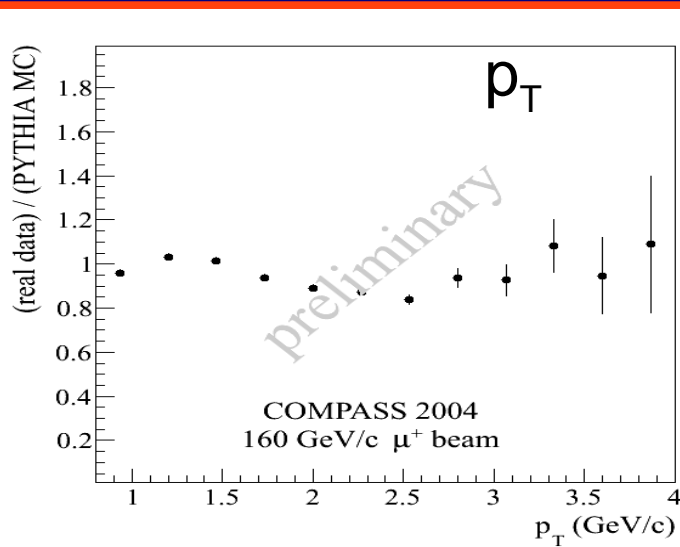
For the data sample of interest the luminosity has been determined within 10% accuracy.

Luminosity: Flattop selection



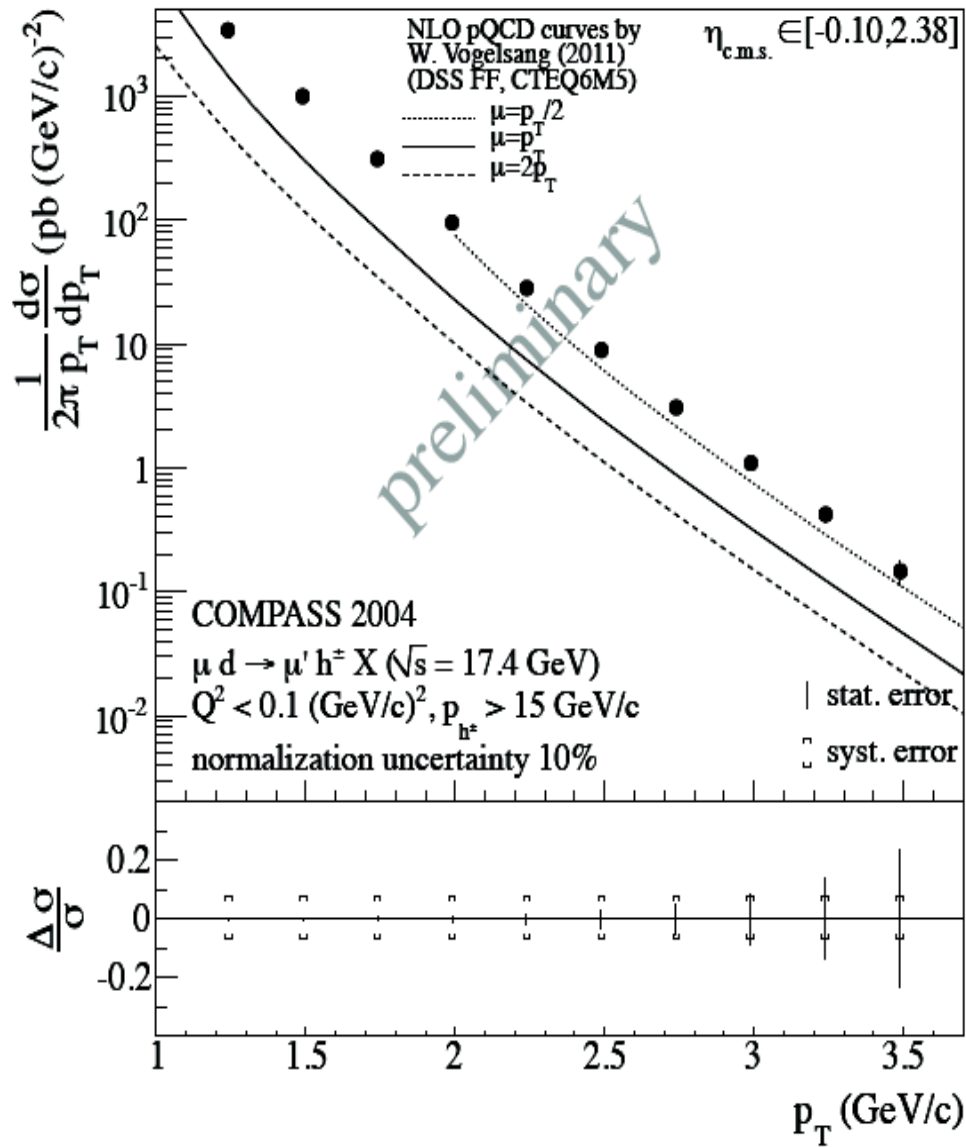
- Evaluate the spill rates as seen by the beam scaler as a function of time.
- Only events within the flattop are considered (red line)

Data/MonteCarlo



- MC used only for extraction of acceptance
(PYTHIA6, GEANT3, COMPASS' data reconstruction)
- 7% Systematic error

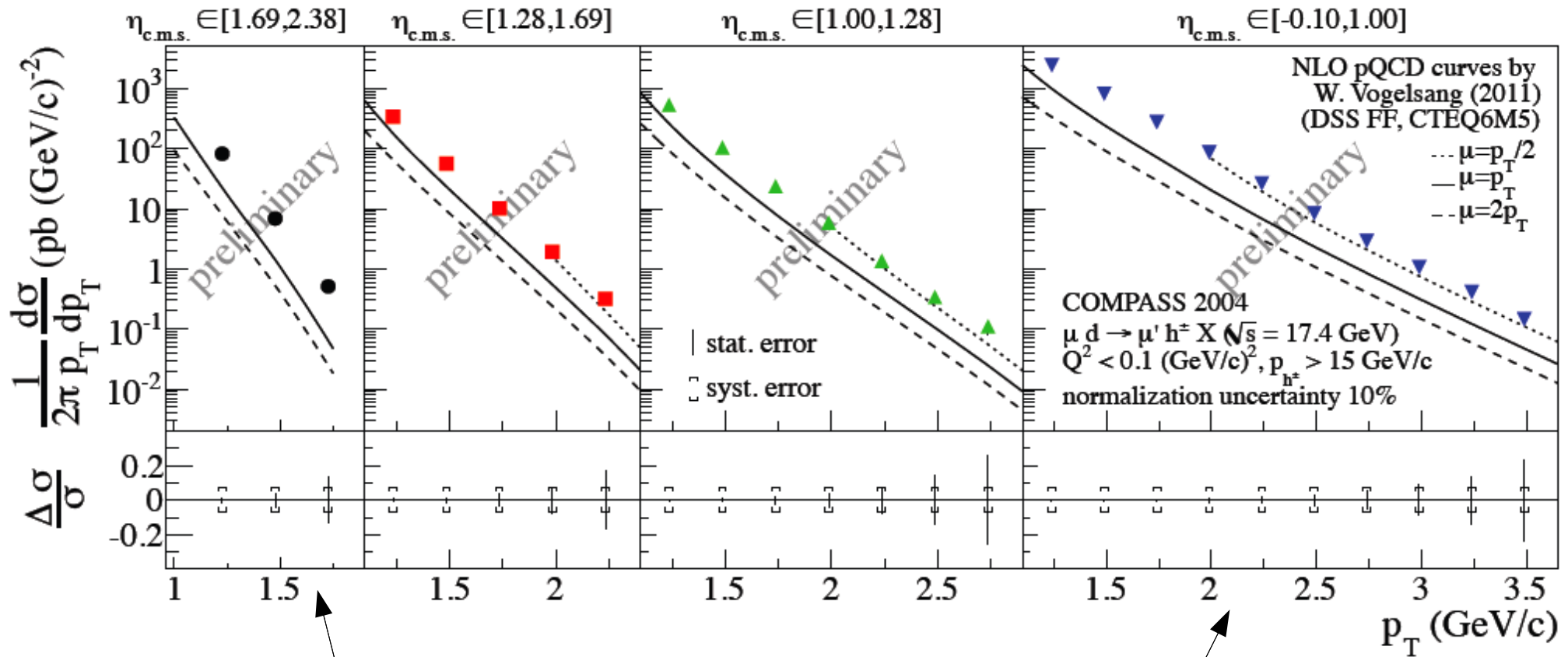
Cross-sections of h^\pm



$\mu^+d \rightarrow \mu^+h^\pm X$: Cross-sections.

- 2004 μ^+d : 4 Weeks of Data.
- Systematic error contributions:
- Background and multi-dimensional acceptance extraction.

Cross Sections Separated by Rapidity Bins

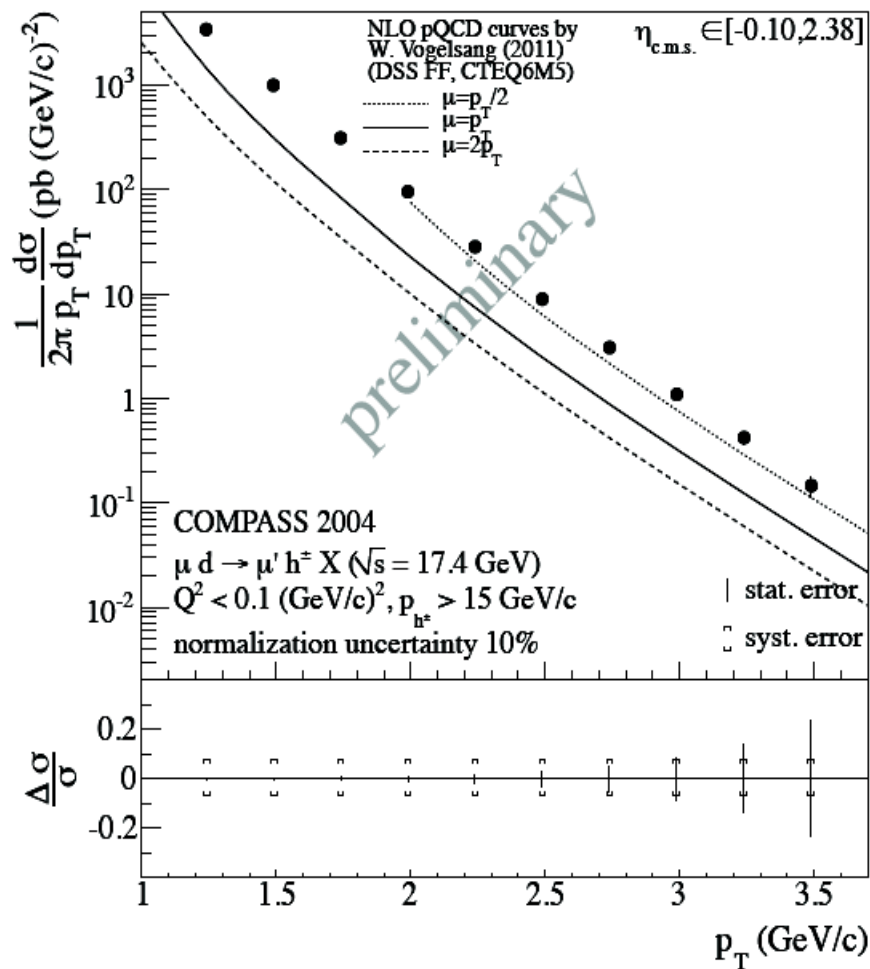


Most Forward Rapidity
Accessible in the Data

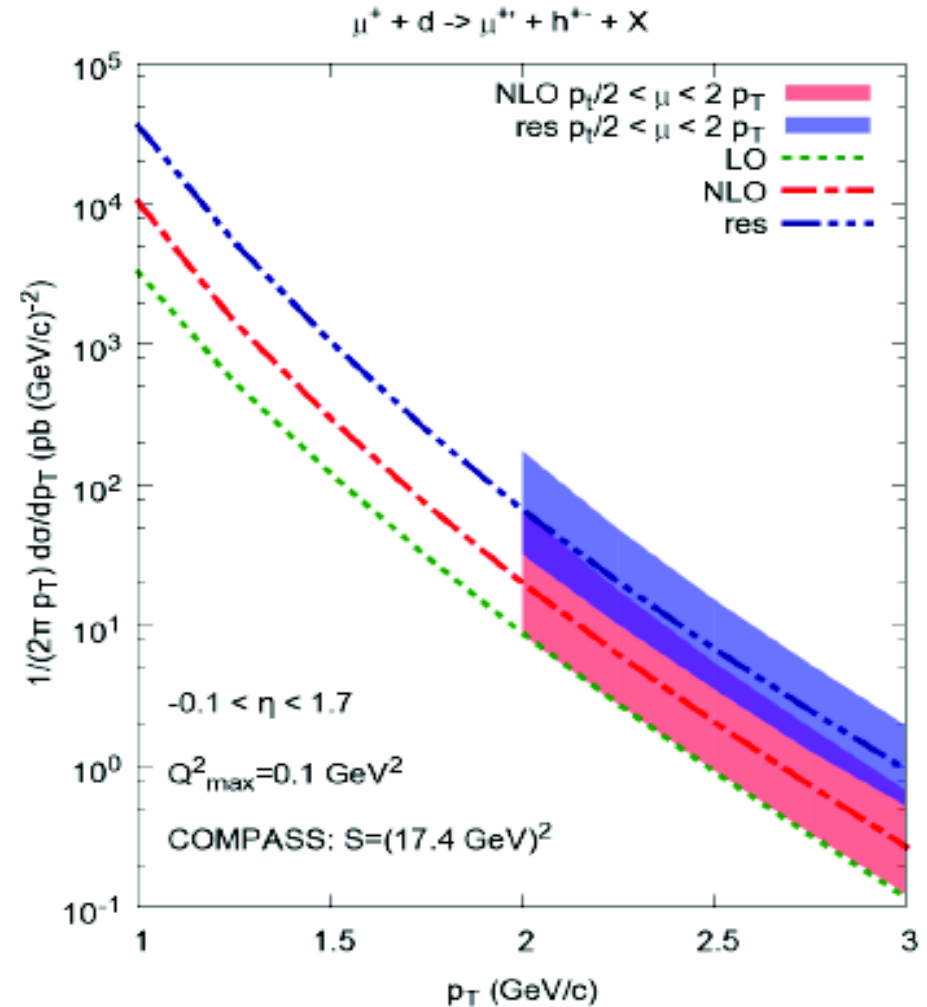
Most Central rapidity

Slope variations as a function of rapidity

Preliminary updated calculations with resummations. (de Florian, Pfeuffer, Schaeffer, Vogelsang)



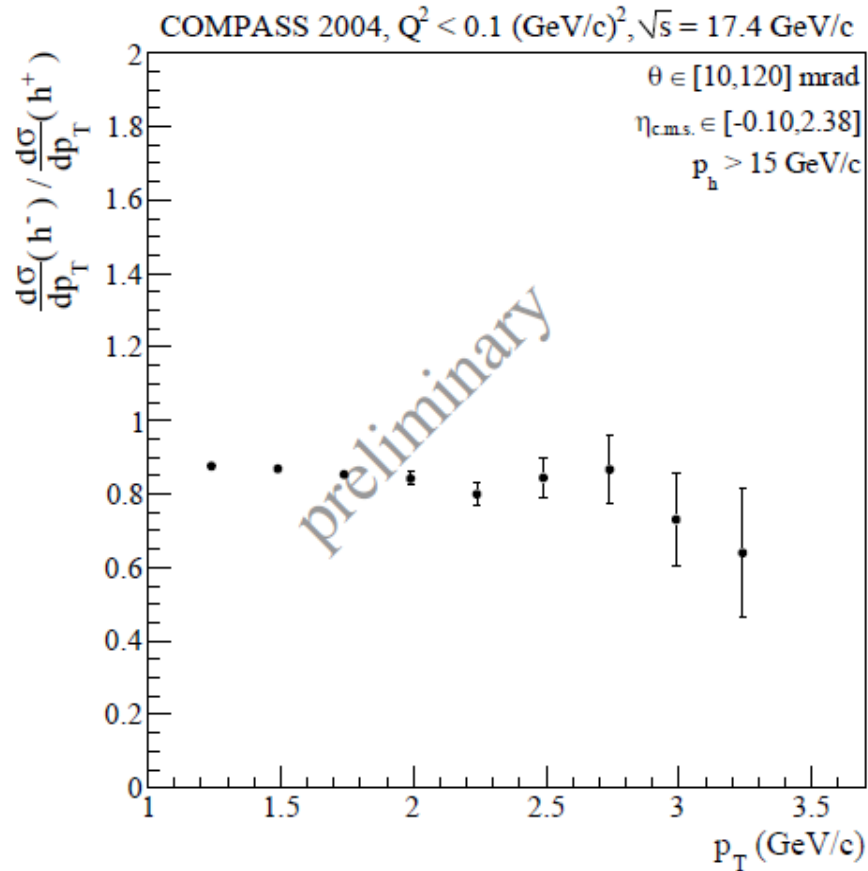
COMPASS



de Florian, Pfeuffer,
Schäfer, WV (prel.)

Ratio h^-/h^+

More up quarks in the final state: Conservation of charge.

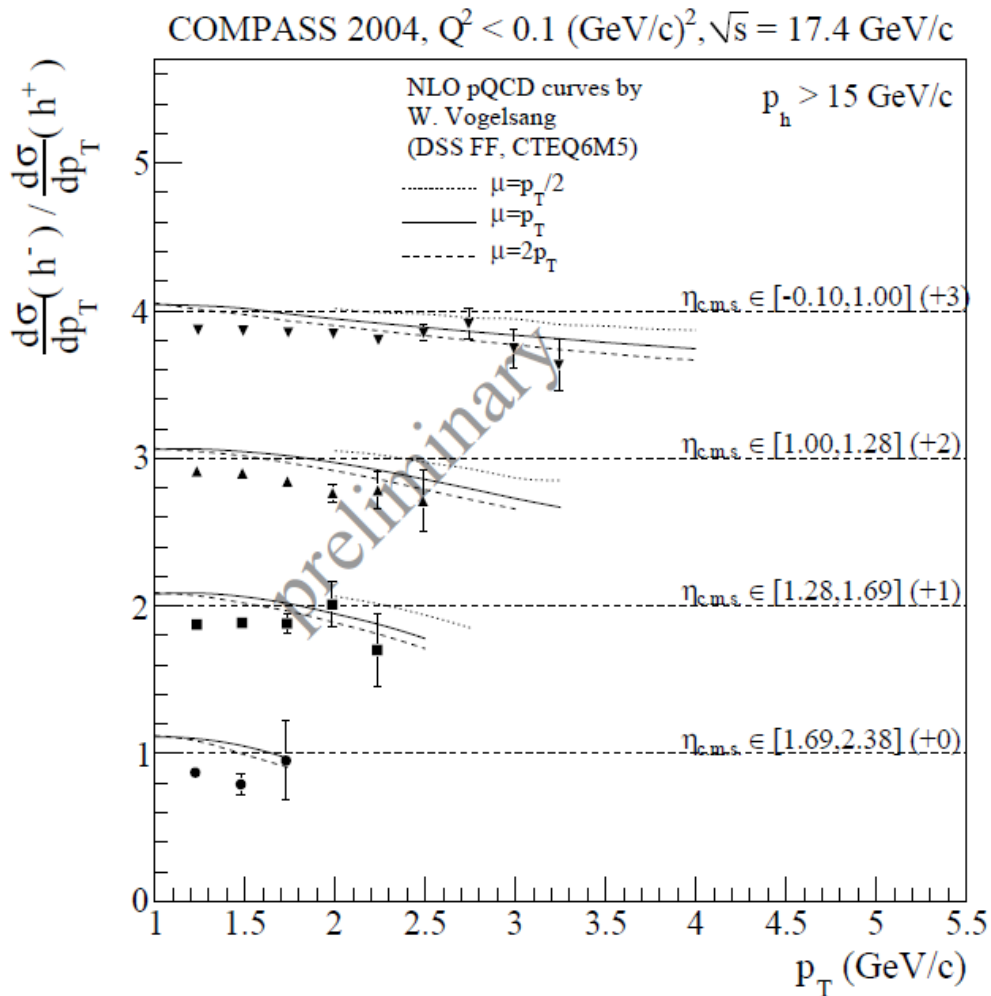


→ Flat as a function of p_T

→ Sensitive to fragmentation functions.

Integrated over all rapidity accessible in the dataset

Ratio h^-/h^+



→ Flat as a function of p_T

→ Sensitive to Fragmentation functions.

→ Disagreement observed with calculation.

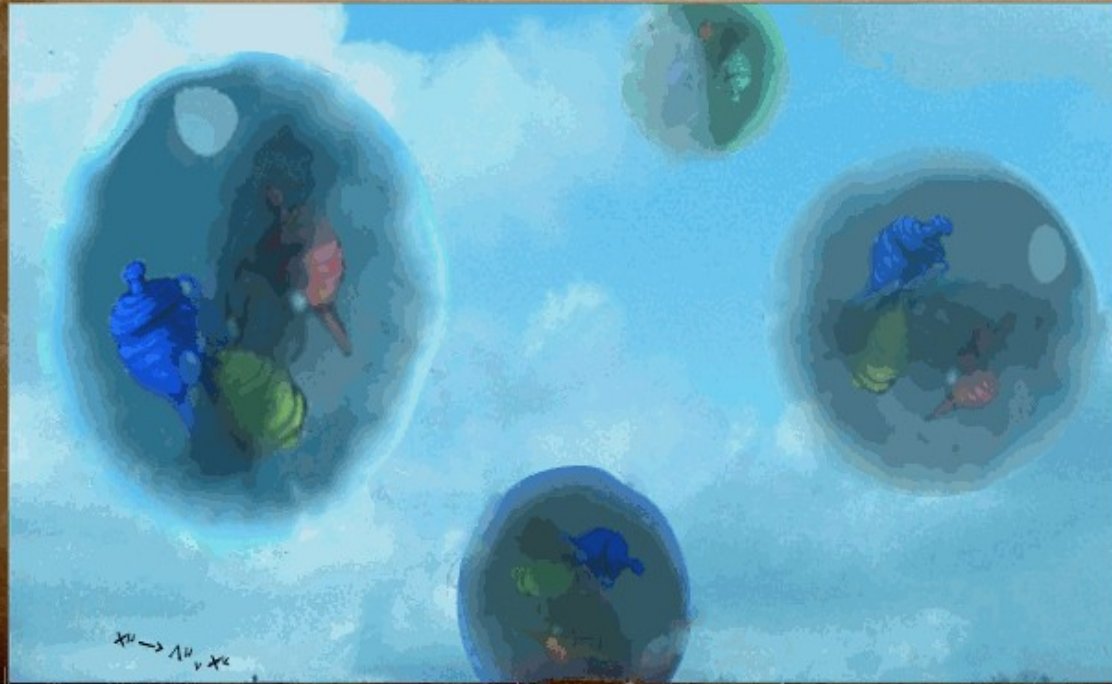
← Most Forward rapidity
accessible to the data.

Theory **does not agree** with experimental results. → Fragmentation Functions?

Summary

- COMPASS has measured the cross sections of the SIDIS process $\mu^+d \rightarrow \mu^+h^{\pm}X$ at high p_T and low Q^2
- Results show that within the theory's errors, the data is well described. *settles the theory framework (ΔG) through that process.*
- **Ratio** of negative to positive cross-sections show a discrepancy at low p_T
- Updated **theory calculations** with resummations are under paper preparation.
- **Results** of these work is finished → under paper preparation.
- **Asymmetries and ΔG . Ongoing work**

Thank You



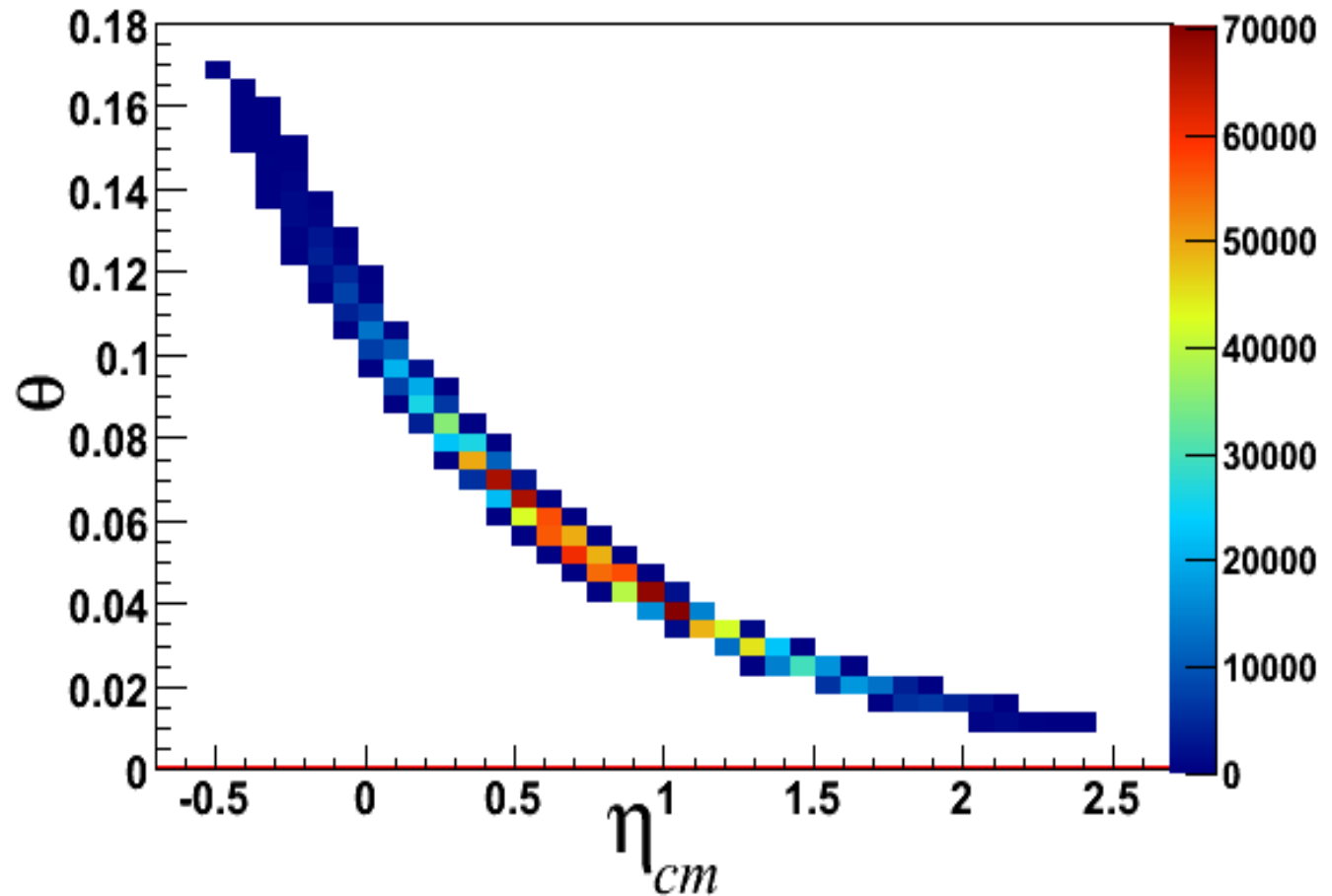
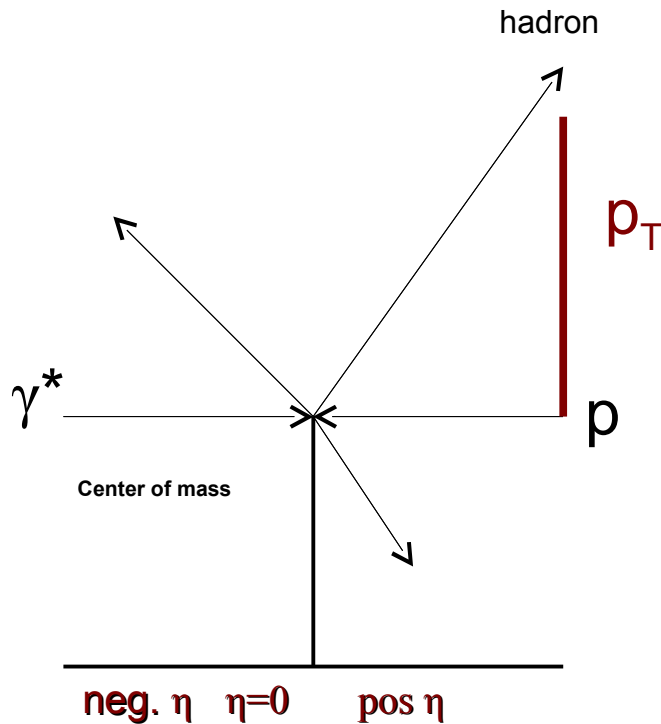
26/06

Astrid Morreale,
CEA Saclay, Irfu/SPhN, NSF

34

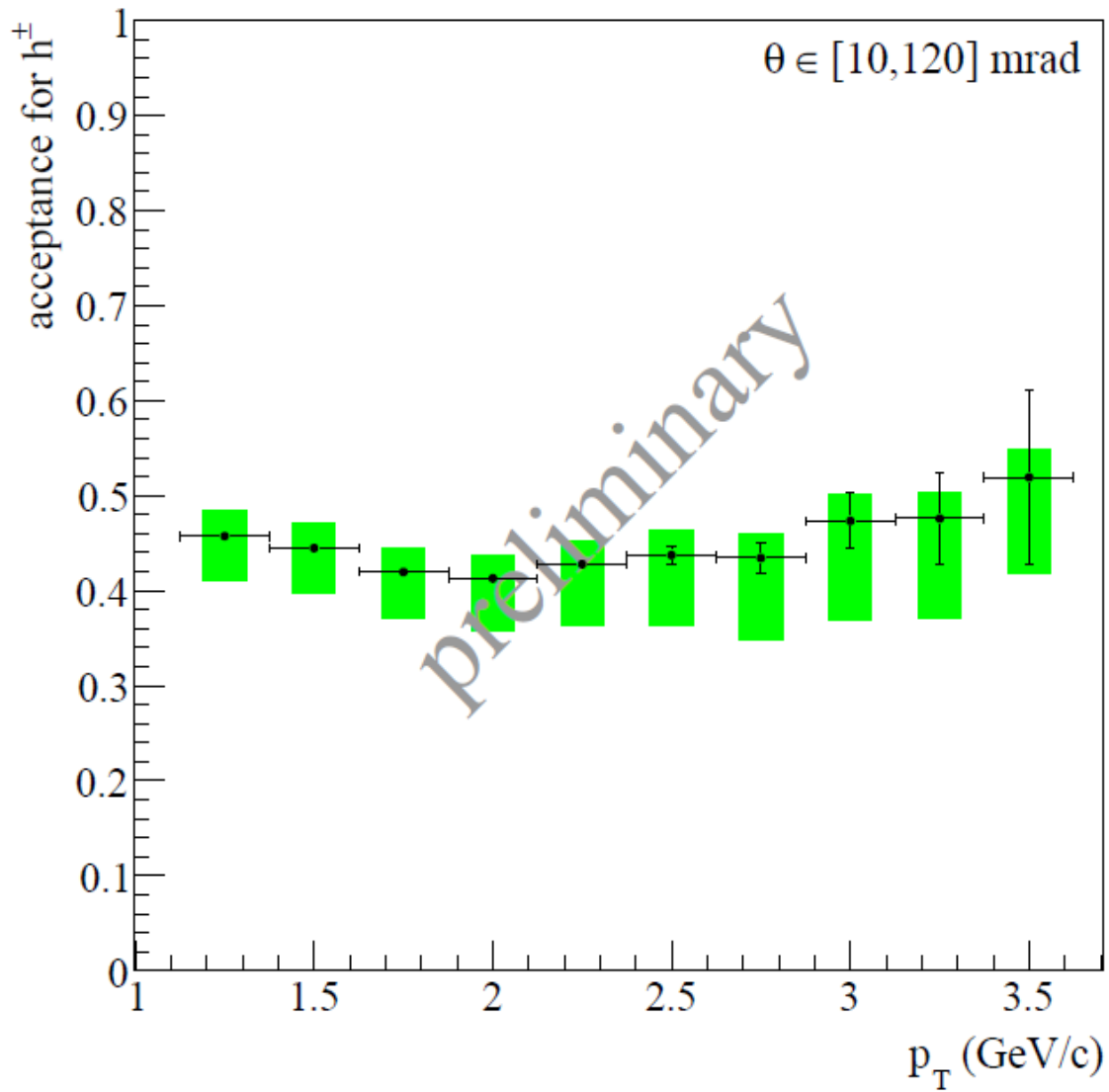
η_{cms}

$$\eta_{\text{cms}} = -\log(\tan(\theta/2)) - 0.5 \cdot \log(2P_{\text{beam}}/M_p)$$



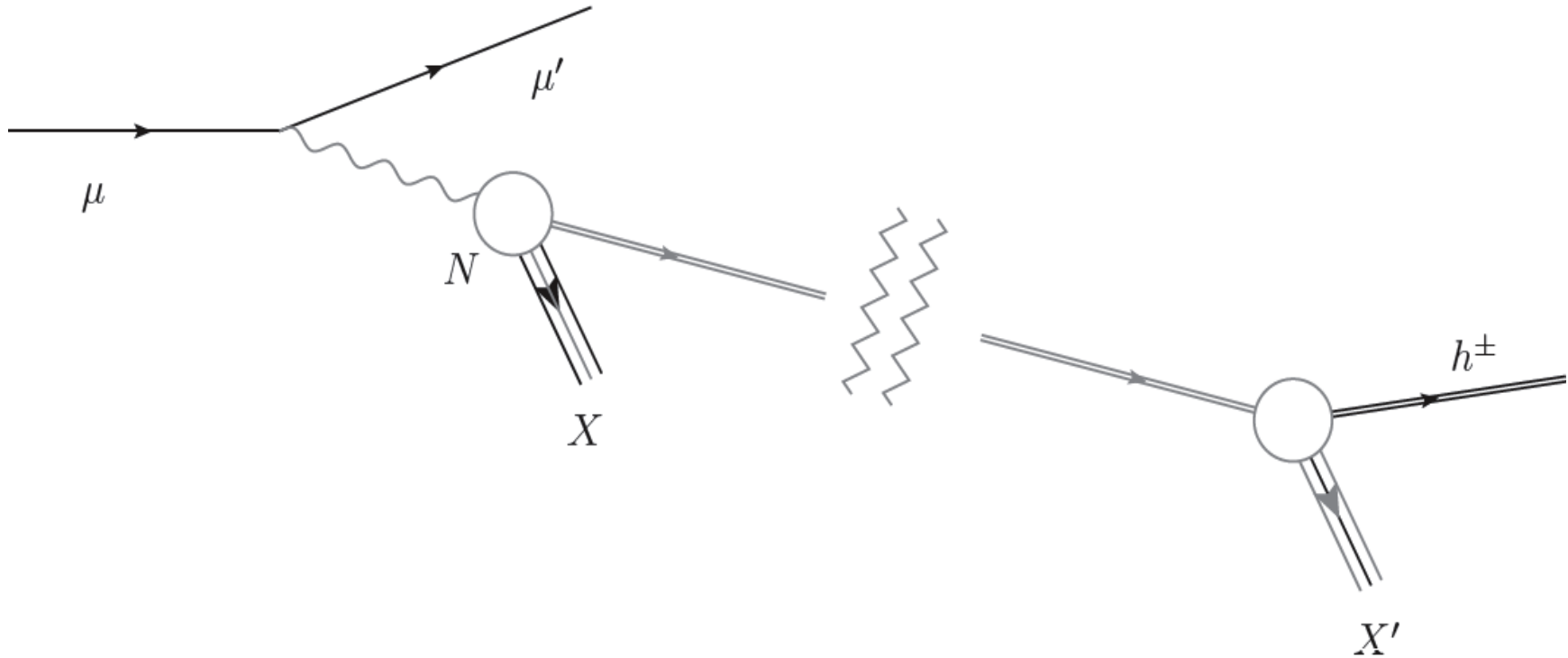
Low $\theta \rightarrow$ Forward η_{cm} . High $\theta \rightarrow$ Central η_{cm} . (Hard Scattering)

Acceptance



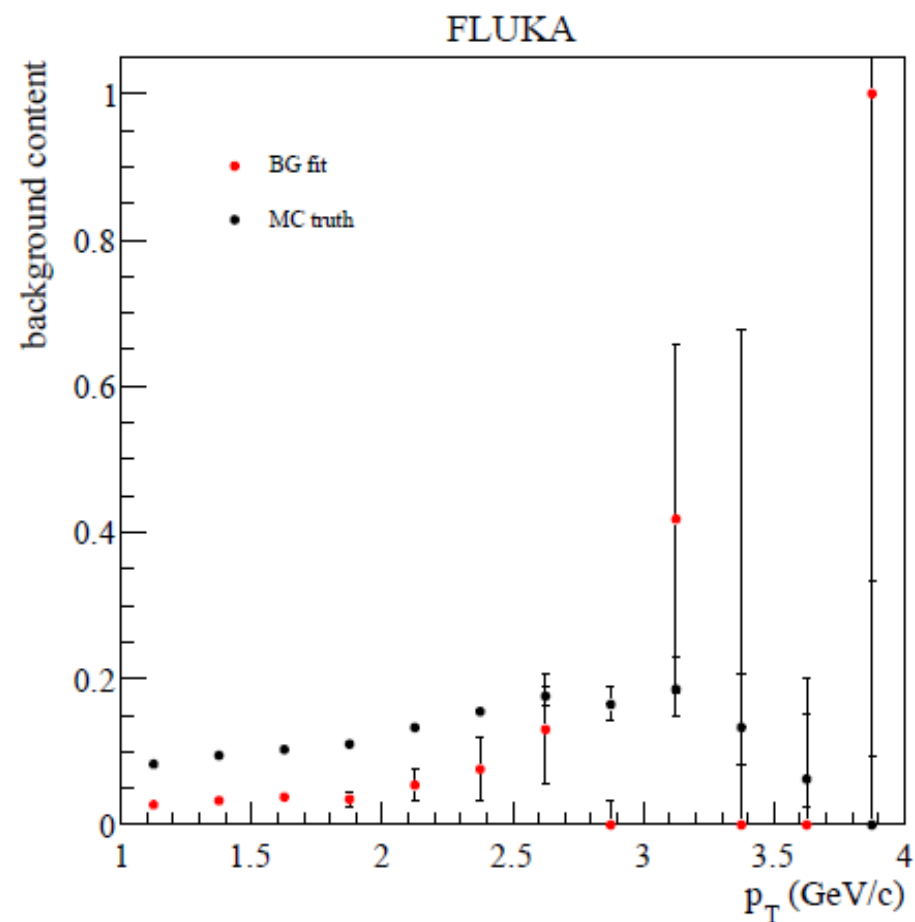
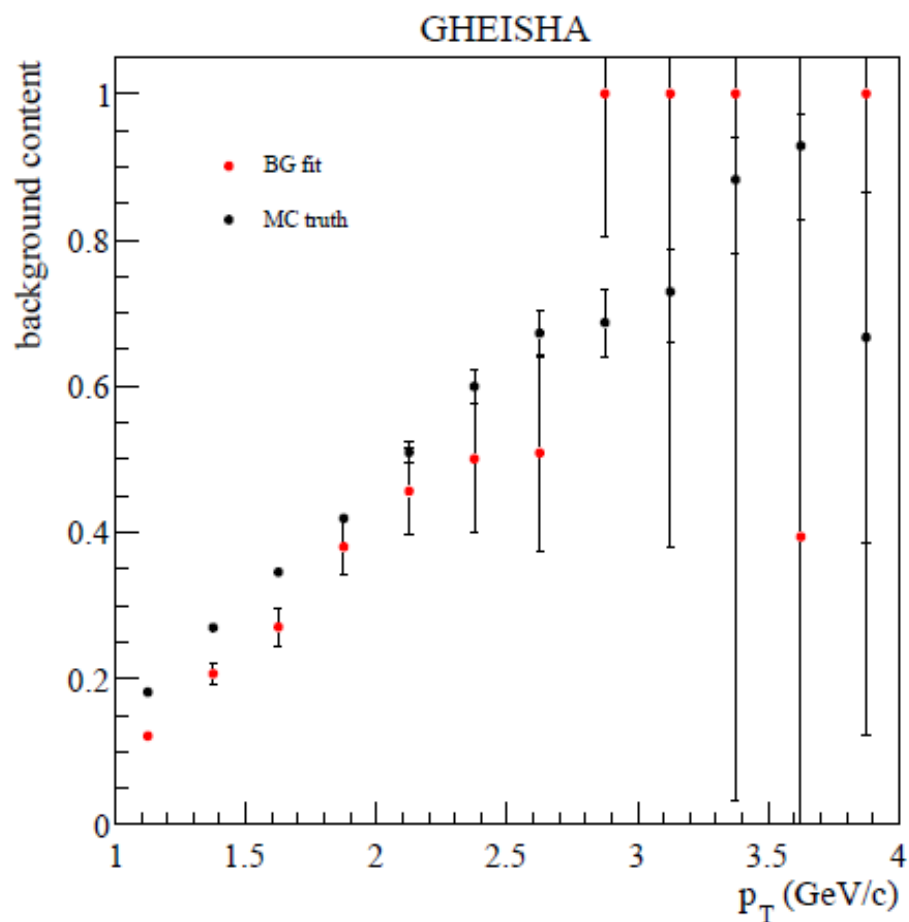
Flat as a function of p_T

Secondary interactions in the target



- A hadron produced in the primary vertex
- Sometimes it will re-interact with nuclear material in the target
- Kinematics will be modified introducing a smeared spectrum

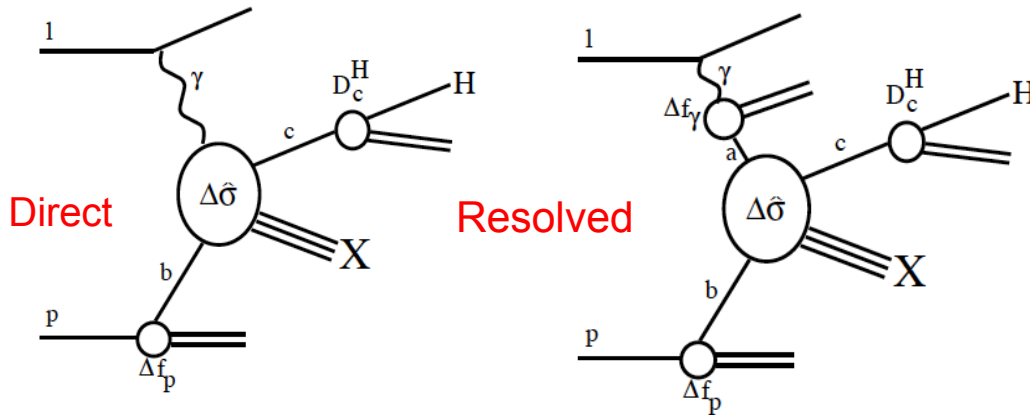
Background Estimate



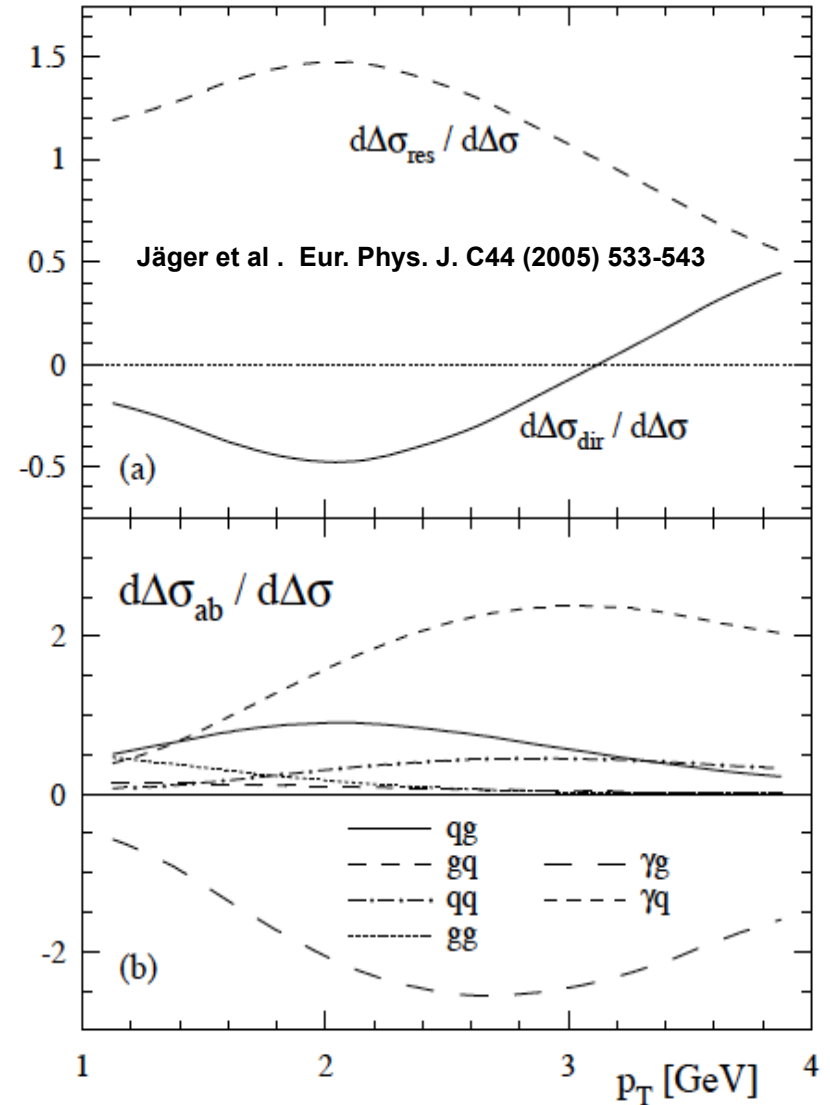
Both models were parametrized and compared to data
Fluka agreed while GHEISHA failed to reproduce
the data → **GHEISHA discarded**



NLO pQCD: Quasi Real Photo-production of high p_T hadrons

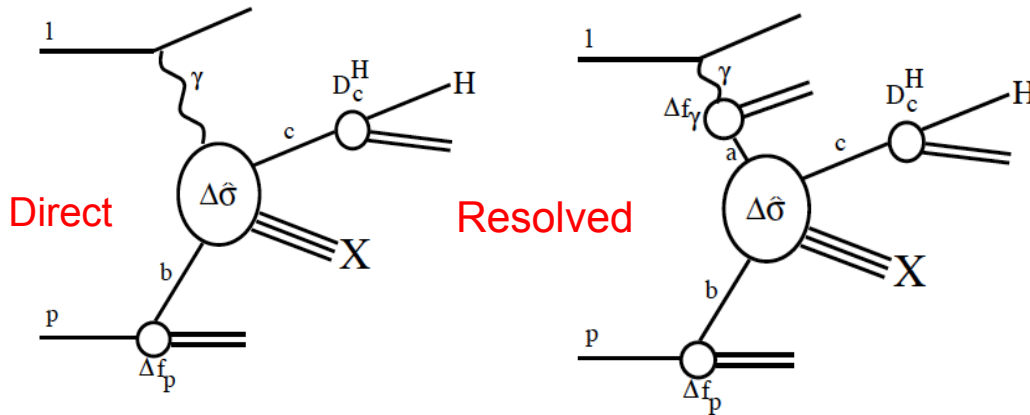


- Low Q^2 ($<0.1\text{GeV}/c$) and high p_T .
- NLO Calculation exists for Compass Kinematics.
- This process has an advantage of higher production rates of hadrons than in (DIS) electro-production.
- The selected hadron H is at high p_T ($>1\text{GeV}/c$) : large momentum transfer, p_T sets the scale

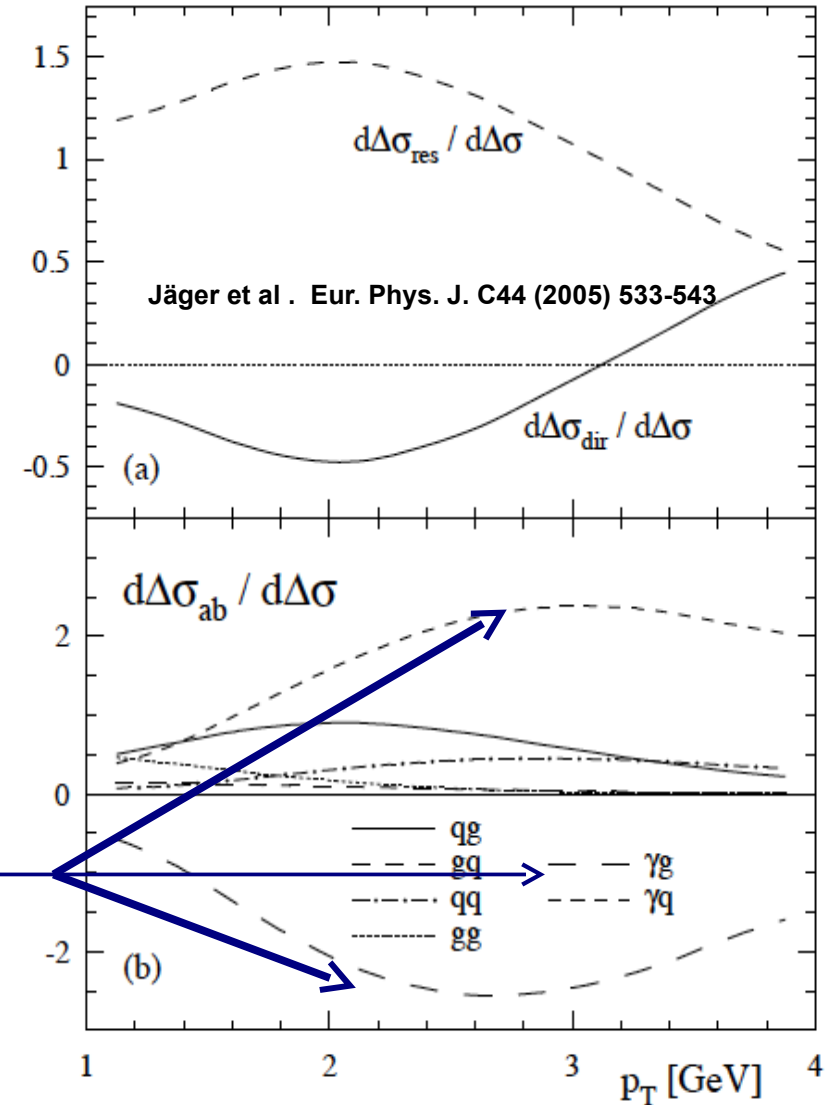




NLO pQCD: Quasi Real Photo-production of high p_T hadrons



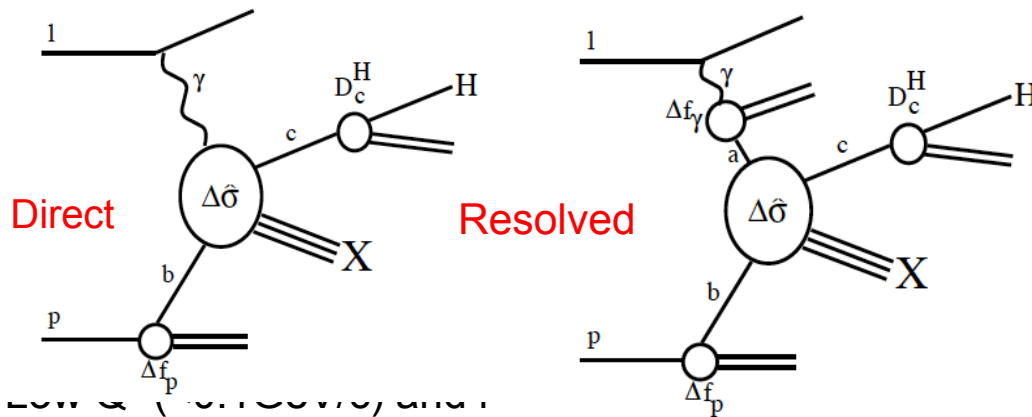
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- **Direct Processes** contribute with different sign ($\gamma g, \gamma q$)



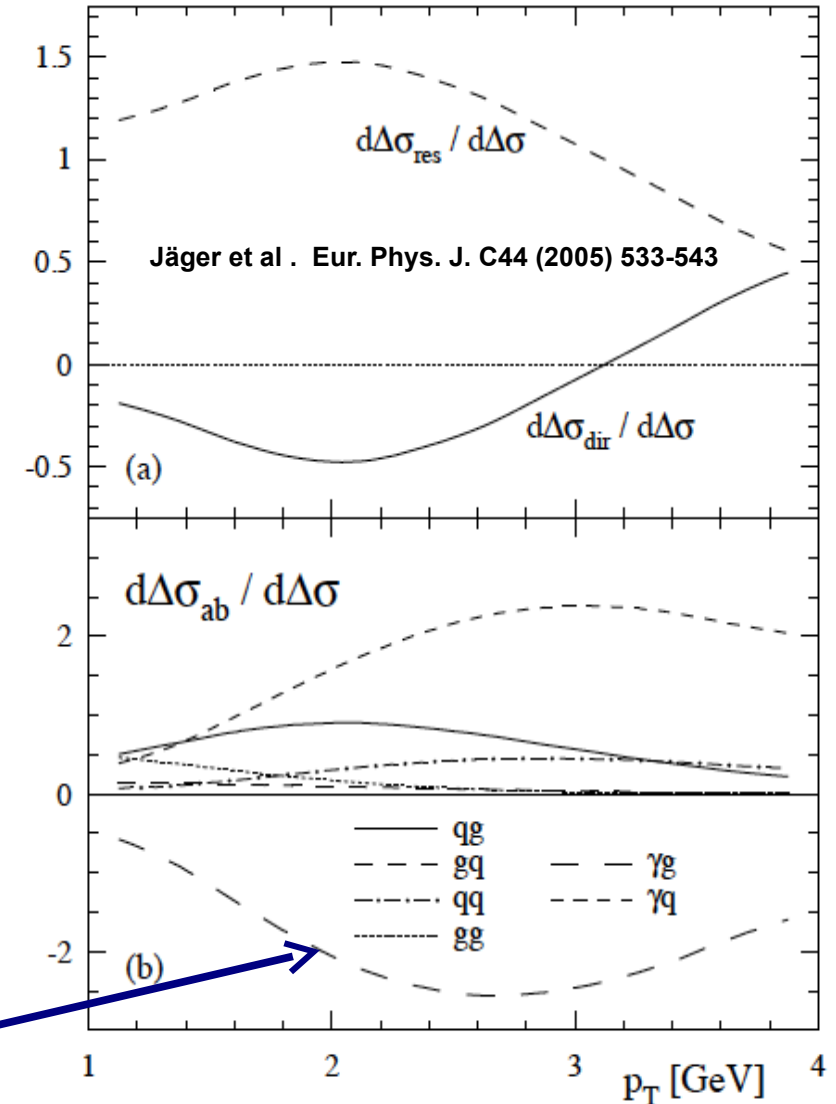
Partonic contributions to the production of high p_T hadrons at low Q^2 ($<0.5\text{GeV}^2$) in lepton nucleon scattering. $\text{c.m.} = 18\text{ GeV}$



NLO pQCD: Quasi Real Photo-production of high p_T hadrons



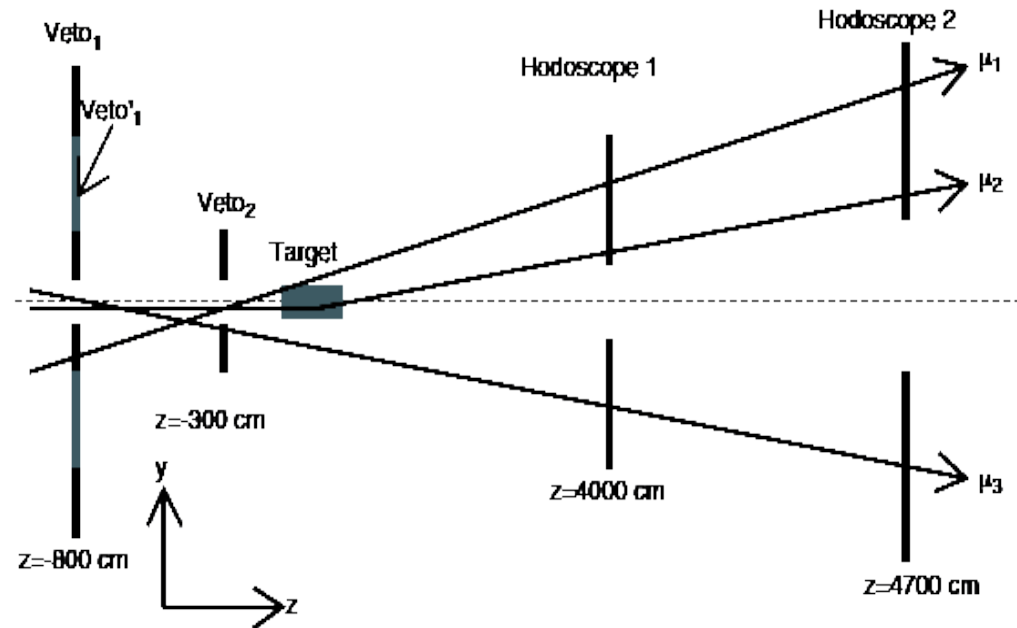
- NLO Calculation exists for Compass Kinematics.
- This process has an advantage of higher production rates of hadrons than in (DIS) electro-production.
- The selected hadron H is at high $p_T (> 1 \text{ GeV}/c)$: large momentum transfer, p_T sets the scale
- Direct Processes contribute with different sign ($\gamma g, \gamma q$)
- **Resolved processes** contribute with the same sign for a positive ΔG (qg, gq, qq, gg)



Partonic contributions to the production of high p_T hadrons at low $Q^2 (< 0.5 \text{ GeV}^2)$ in lepton nucleon scattering. c.m= 18 GeV

COMPASS Veto System

Bernet et al. NIMA 550(2005) 217-240.

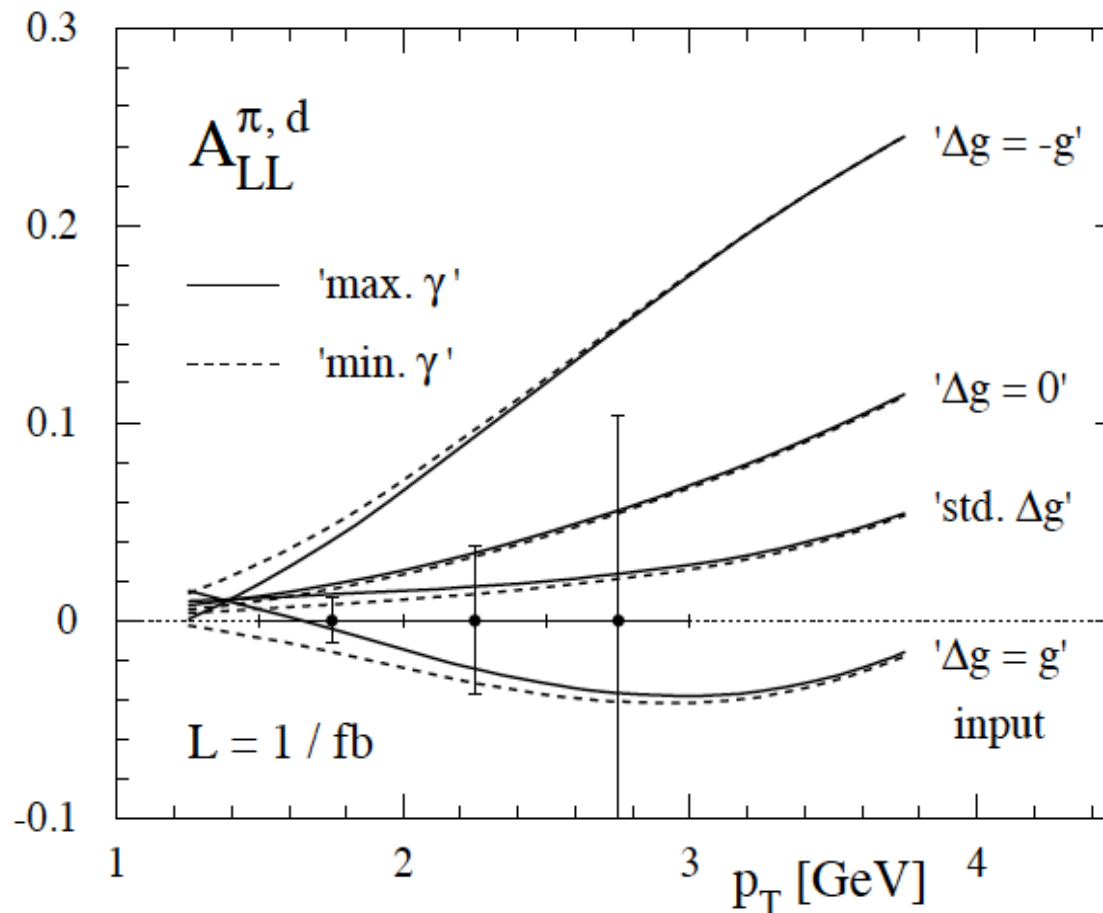


The compass veto system prevents a large fraction of beam halo tracks from contaminating the trigger sample

During high veto signal times no triggers are taken (this includes good events)
→ **The “dead time” effect is taken into account in the Luminosity calculation**



The NLO pQCD parametrizations: Asymmetries



Max γ and min γ correspond to minimal and maximal saturation of the polarized photon densities.

B. Jäger et al. Eur. Phys J. C 36, 371-374 (2004)

→ Additional ongoing effort by A. Afanasiev to produce a parametrization that takes into account higher twist effects.

→ We expect to make available the comparisons between these Models with compass measurements soon.

Sources of Systematic Errors

1. **Luminosity extraction** Dead times, acquisition. 10%

2. **Background Estimate:** MonteCarlo + Data studies.
Two different hadronic shower models studied
GHEISHA and Fluka, compared to real data. 5% .

3. **Acceptance calculation:**

Multi-dimensional acceptance studies in terms of y , w , charge, p_T 3%