

COMPASS results on longitudinal spin physics

IWHSS11 - Paris



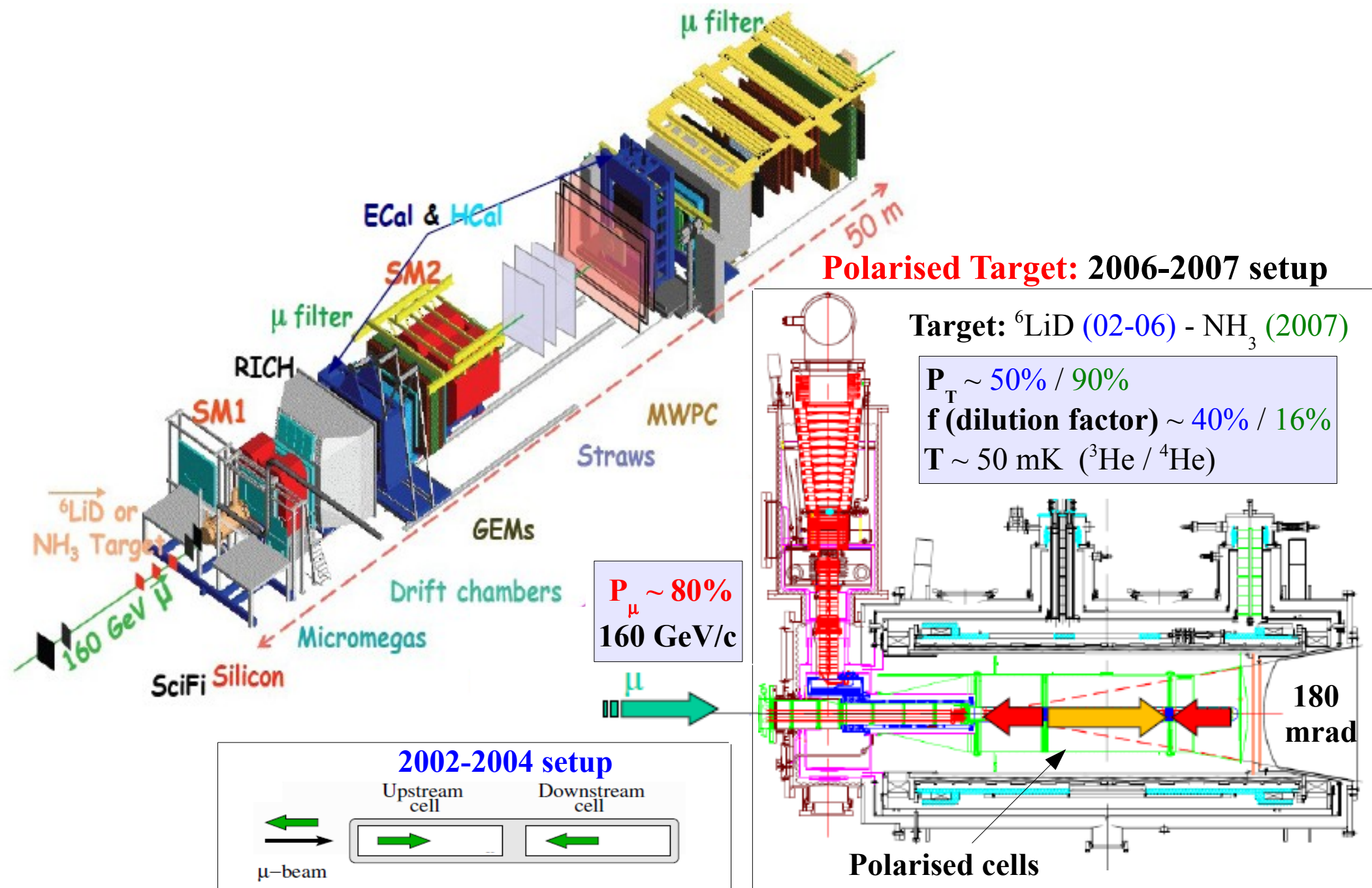
Celso Franco (*LIP – Lisboa*)
on behalf of the COMPASS collaboration

Outline

Longitudinally polarised DIS:

- $A_1^{d/p}$, $g_1^{d/p}$, first moments of g_1^d and the Bjorken sum rule
- Semi-inclusive asymmetries and flavour separation
- Gluon polarisation in LO:
 - Open Charm
 - High- p_T hadron pairs
- **Gluon polarisation in NLO:** → NEW
 - Open Charm

The COMPASS spectrometer and target



Inclusive asymmetries and spin structure functions

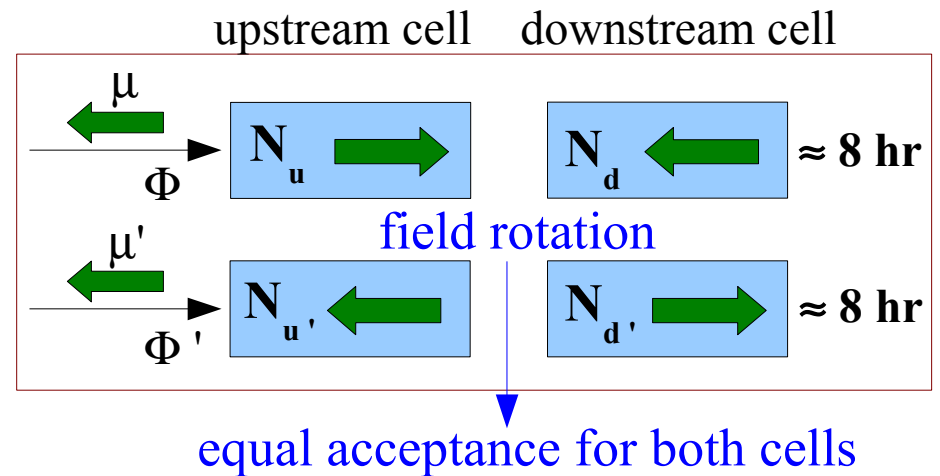
Asymmetry measurement: $A_1^N := \frac{\Delta \sigma_{\gamma^*N}}{\sigma_{\gamma^*N}} = \frac{(\sigma_{\gamma^*N}^{\leftarrow\rightarrow} - \sigma_{\gamma^*N}^{\leftarrow\leftarrow})}{\sigma_{\gamma^*N}^{\text{unpol}}}$

- The number of reconstructed events inside each spin configuration, N_t ($t = u, d, u', d'$), can be used to extract the γ^* -deuteron / proton (A_1^d / A_1^p) asymmetries:

$$A^{\text{exp}} = \frac{1}{2} \left(\frac{N_u - N_d}{N_u + N_d} + \frac{N_{d'} - N_{u'}}{N_{d'} + N_{u'}} \right)$$

$$= f \cdot P_\mu \cdot P_T \cdot \mathbf{D} \cdot \mathbf{A}_1 \rightarrow A^{\mu N}$$

$D = \underline{\text{Depolarisation factor}}$

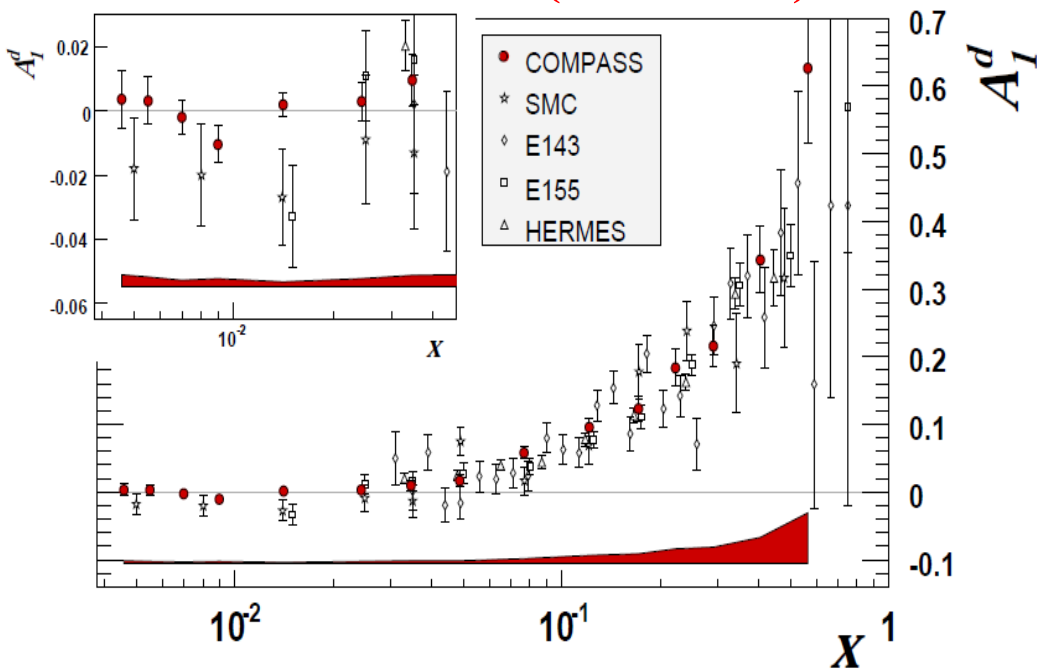


- Weighting each event with $\omega = (f P_\mu D)$:

$$A_1 = \frac{1}{P_T} \times \frac{1}{2} \left(\frac{\sum_u \omega - \sum_d \omega}{\sum_u \omega + \sum_d \omega} + \frac{\sum_{d'} \omega - \sum_{u'} \omega}{\sum_{d'} \omega + \sum_{u'} \omega} \right) \text{ with statistical gain: } \frac{\langle \omega^2 \rangle}{\langle \omega \rangle^2}$$

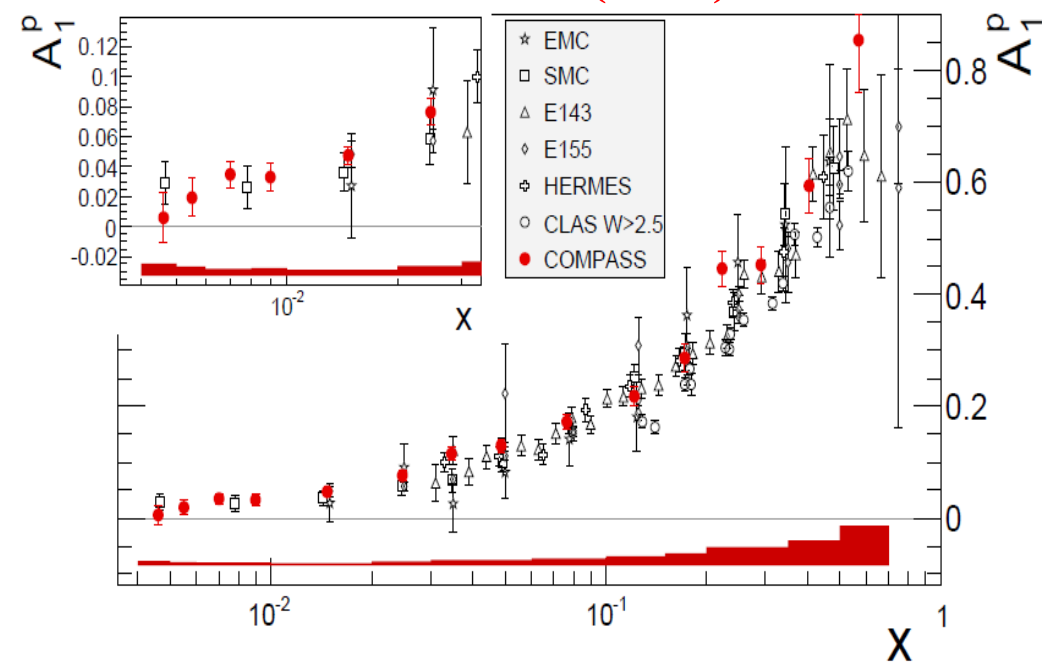
Inclusive asymmetries $A_1^{d/p}$: $Q^2 > 1$ (GeV/c)²

Deuteron data (2002-2006)



COMPASS, PLB 647 (2007) 330-340

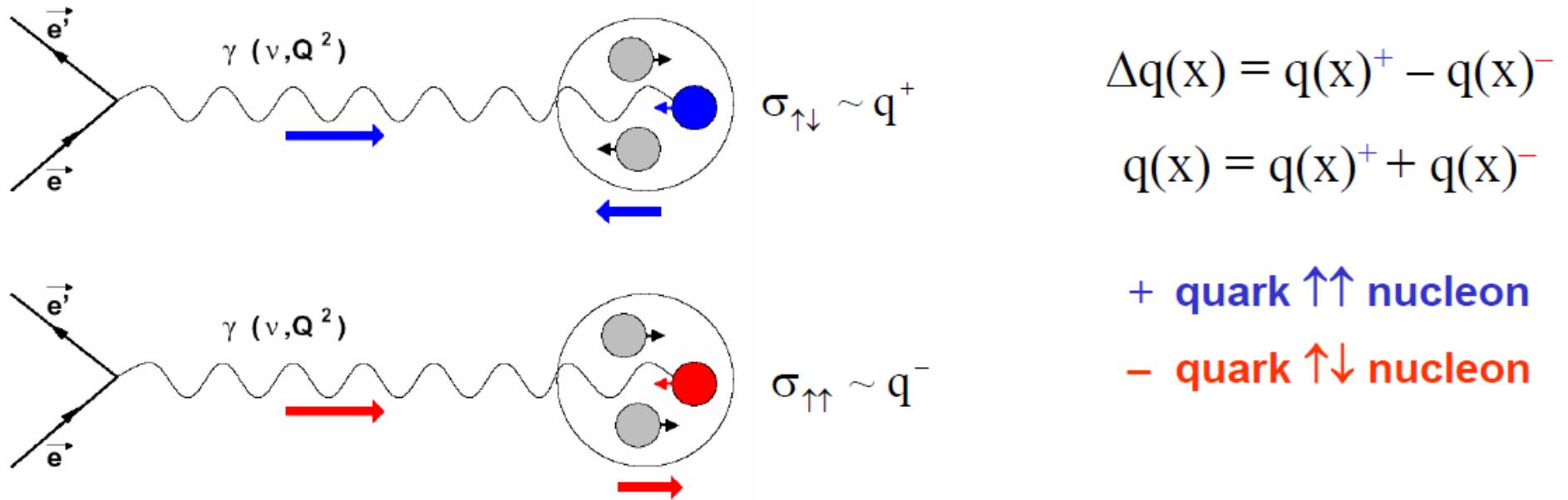
Proton data (2007)



COMPASS, PLB 690 (2010) 466-472

- Good agreement between all experimental points
- Significant improvement of precision in the low x region: compatible with zero for $x < 0.01$
- No negative trend for A_1^d

Interpretation of A_1 in terms of structure functions

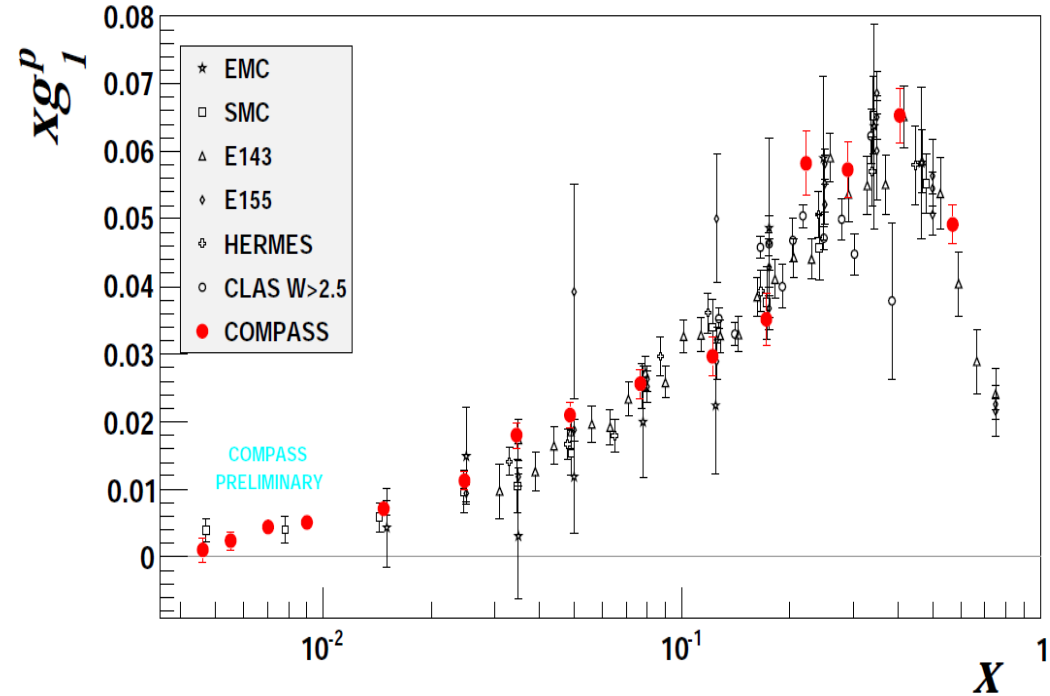
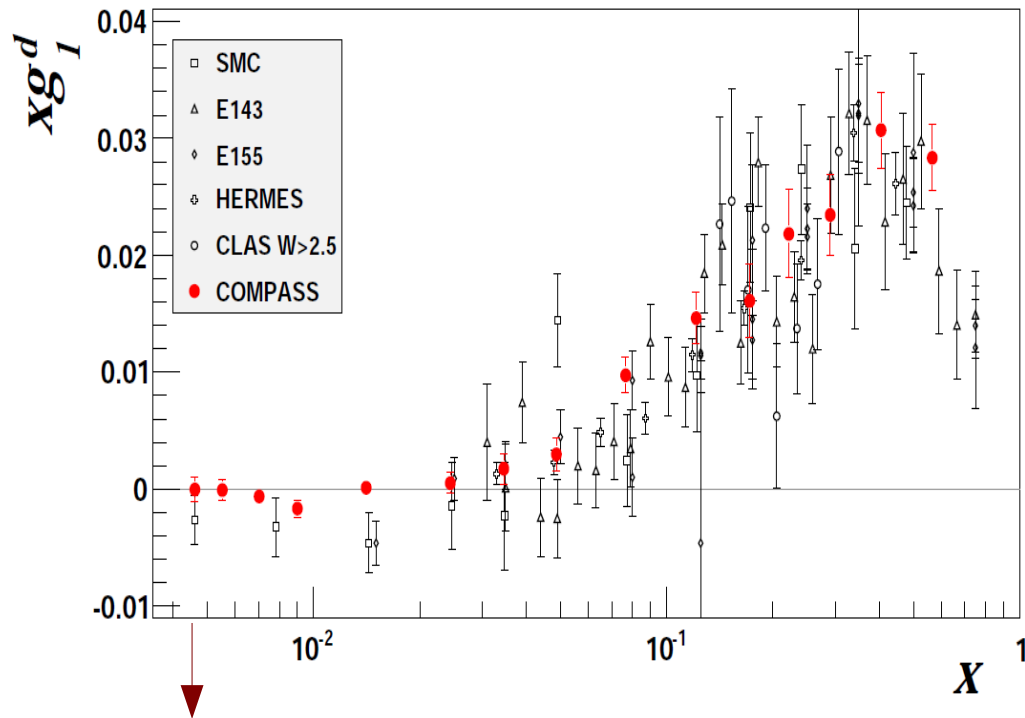


$$A_1(x, Q^2) = \frac{\sigma_{\uparrow\downarrow} - \sigma_{\uparrow\uparrow}}{\sigma_{\uparrow\downarrow} + \sigma_{\uparrow\uparrow}} \approx \frac{\sum_q e_q^2 \Delta q(x, Q^2)}{\sum_q e_q^2 q(x, Q^2)} = \frac{g_1(x, Q^2)}{F_1(x, Q^2)} = \frac{g_1(x, Q^2) 2x(1+R)}{F_2(x, Q^2)}$$

- g_1 (polarised structure function) is obtained from the asymmetry A_1 using:

$F_2 \rightarrow$ SMC parameterisation and $R = \sigma^L/\sigma^T \rightarrow$ SLAC parameterisation

COMPASS results for $g_1^{d/p}$ and first moments of g_1^d



$$\Gamma_1^N(Q_0^2=3(\text{GeV}/c)^2) = \int_0^1 g_1(x) dx = 0.0502 \pm 0.0028(\text{stat}) \pm 0.0020(\text{evol}) \pm 0.0051(\text{syst})$$

$$= \frac{1}{9} \left(1 - \frac{\alpha_s(Q^2)}{\pi} + \mathcal{O}(\alpha_s^2) \right) \left(\mathbf{a}_0(Q^2) + \frac{1}{4} \mathbf{a}_8 \right) \Rightarrow \mathbf{a}_0 = 0.35 \pm 0.03(\text{stat}) \pm 0.05(\text{syst})$$

$$\Delta \Sigma^{\overline{\text{MS}}} = 0.33 \pm 0.03(\text{stat}) \pm 0.05(\text{syst}) \quad (\Delta \Sigma^{\overline{\text{MS}}} = \mathbf{a}_0 \text{ @ } Q^2 \rightarrow \infty)$$

$$(\Delta \mathbf{s} + \Delta \bar{\mathbf{s}}) = \frac{1}{3} (\Delta \Sigma^{\overline{\text{MS}}} - \mathbf{a}_8) = -0.08 \pm 0.01(\text{stat}) \pm 0.02(\text{syst})$$

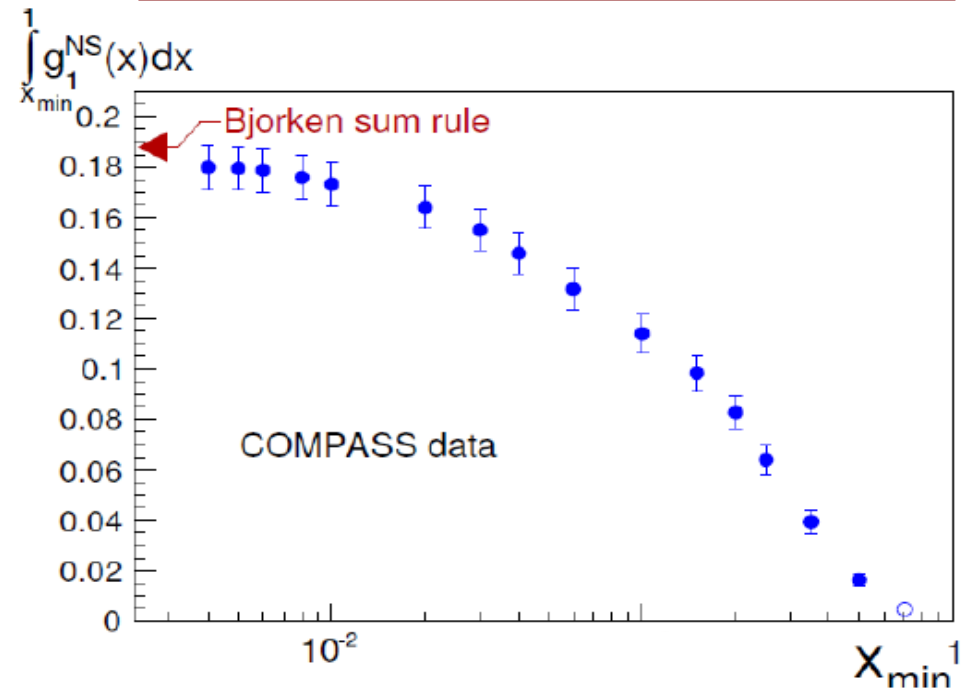
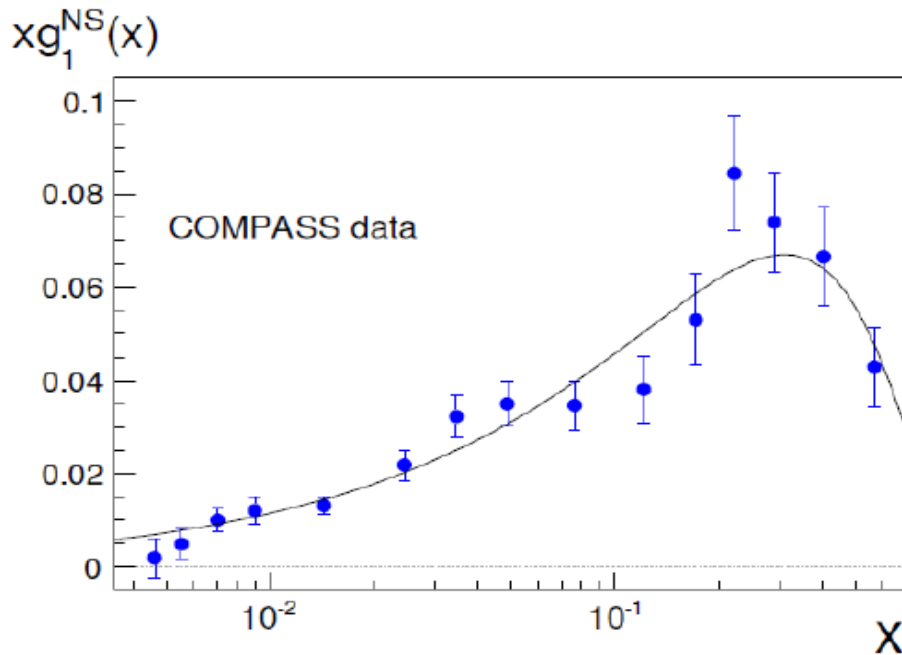
Bjorken sum rule

- According to the Bjorken sum rule the first moment of the non-singlet spin structure function, g_1^{NS} , is proportional to the ratio of axial and vector coupling constants g_A/g_V :

$$\int_0^1 g_1^{NS}(x, Q^2) dx = \frac{1}{6} \left| \frac{g_A}{g_B} \right| C_1^{NS}(Q^2)$$

using

$$g_1^{NS}(x, Q^2) = g_1^p(x, Q^2) - g_1^n(x, Q^2) \\ = 2g_1^p - 2g_1^d / (1 - 1.5\omega_D)$$

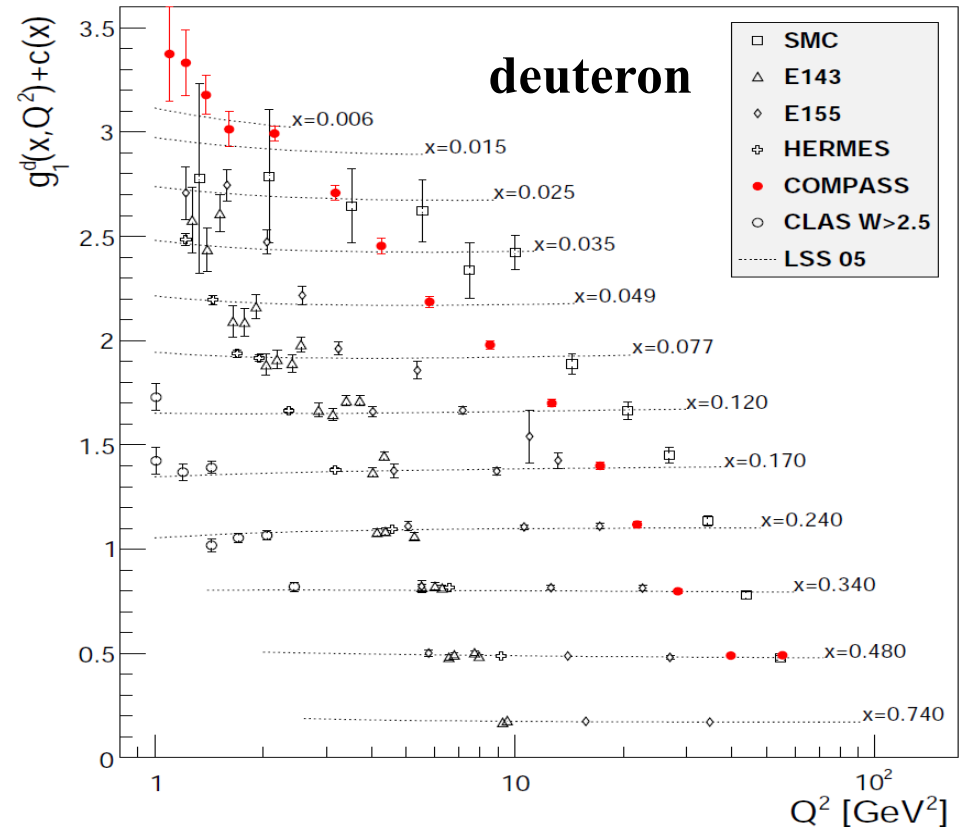
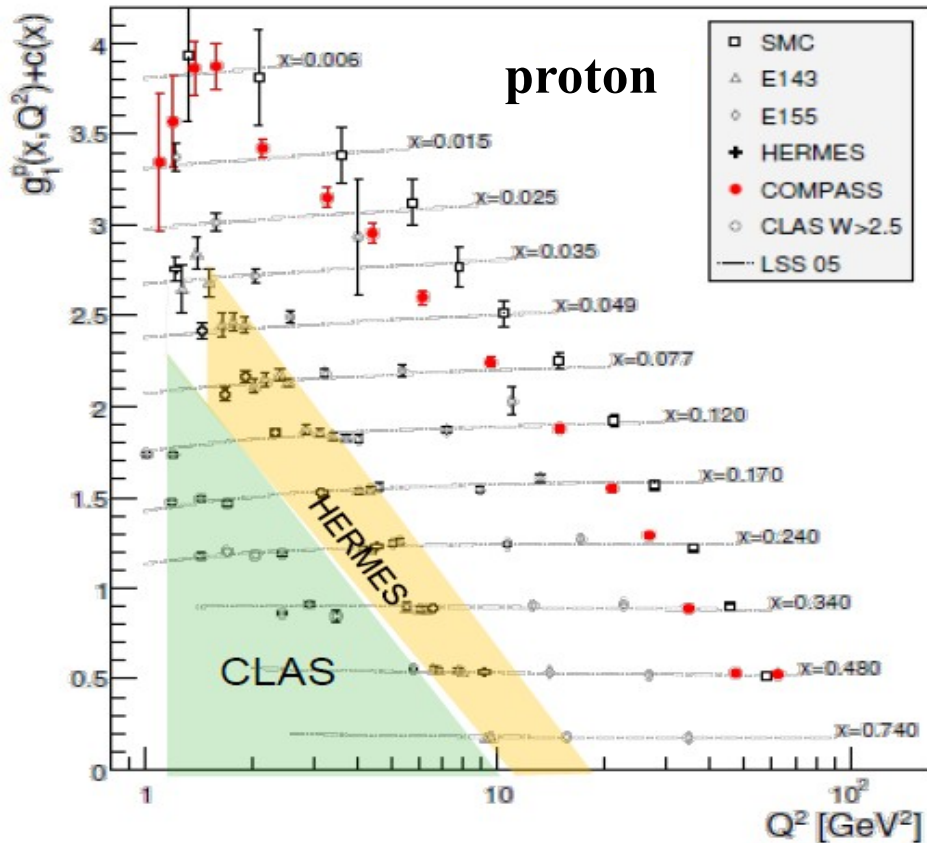


- QCD fit of COMPASS data using $\Delta q^{NS} = |g_A/g_V| x^\alpha (1-x)^\beta$:

$$\left| \frac{g_A}{g_V} \right| = 1.28 \pm 0.07(\text{stat}) \pm 0.10(\text{sys})$$

(PDG value: $|g_A/g_V| = 1.269 \pm 0.003$)

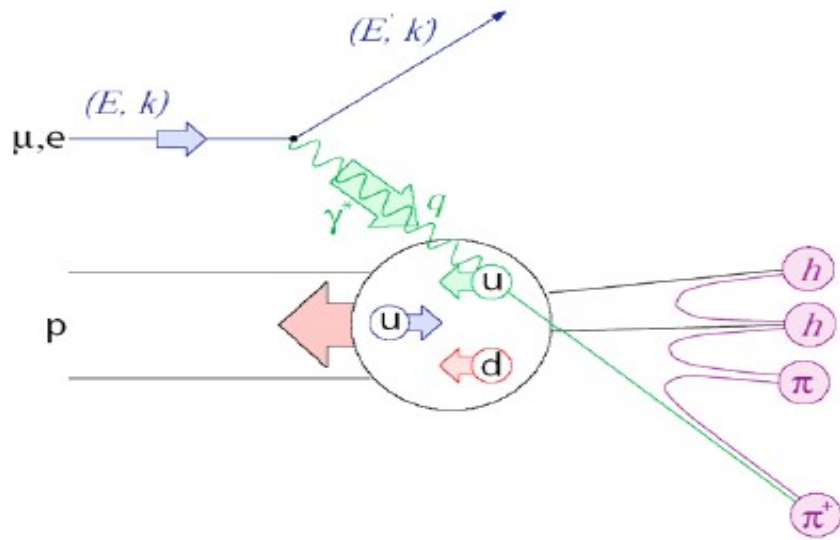
Q^2 dependence of $g_1(x, Q^2)$ for DGLAP evolution



- The kinematic range is still limited (compared to the unpolarised F_2)
 - ➔ additional data from colliders is required!
- $(\Delta u + \Delta \bar{u})$ and $(\Delta d + \Delta \bar{d})$ are well constrained by the data (*LSS PRD 80 2009*)
- Δs comes out negative and Δg is small (< 0.5) ➔ Still with large uncertainties

Semi-inclusive asymmetries and flavour separation

Extraction of the quark helicity distributions from SIDIS



- The outgoing hadron tags the quark flavour
- Required: fragmentation function of a quark q to a hadron h : $D_q^h(z, Q^2)$

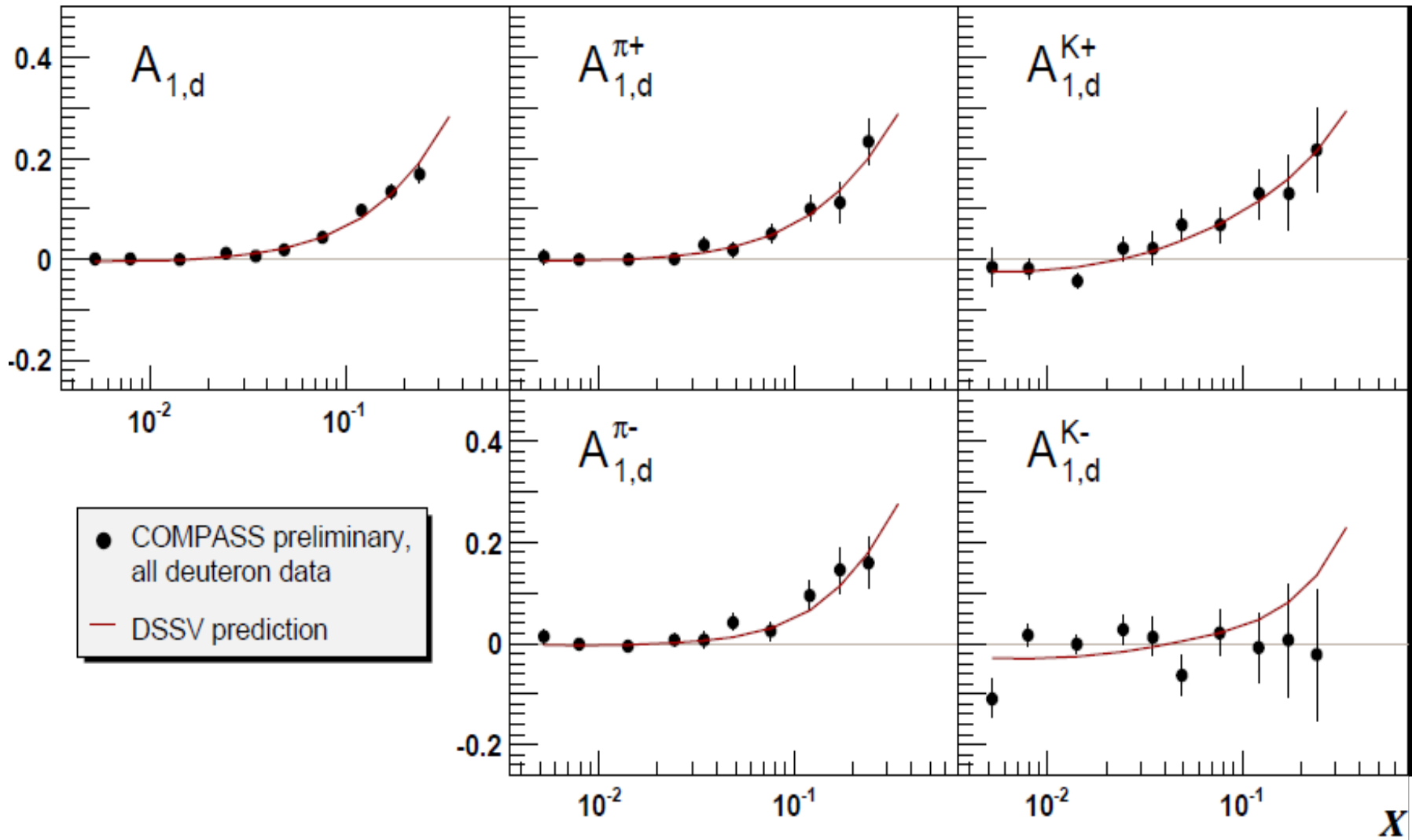
$$z = E_h / (E_\mu - E'_\mu)$$

- The semi-inclusive asymmetries have the following interpretation (in LO):

$$A_1^{h(p/d)}(x, z, Q^2) \approx \frac{\sum_q e_q^2 \Delta q(x, Q^2) D_q^h(z, Q^2)}{\sum_q e_q^2 q(x, Q^2) D_q^h(z, Q^2)}$$

- Inputs needed for the extraction of $\Delta q(x, Q^2)$:
 - Unpolarised PDFs ($q(x, Q^2)$) → [MRST04](#)
 - $D_q^h(z, Q^2)$ → [DSS parameterisation](#)

Inclusive and semi-inclusive spin asymmetries: Deuteron data

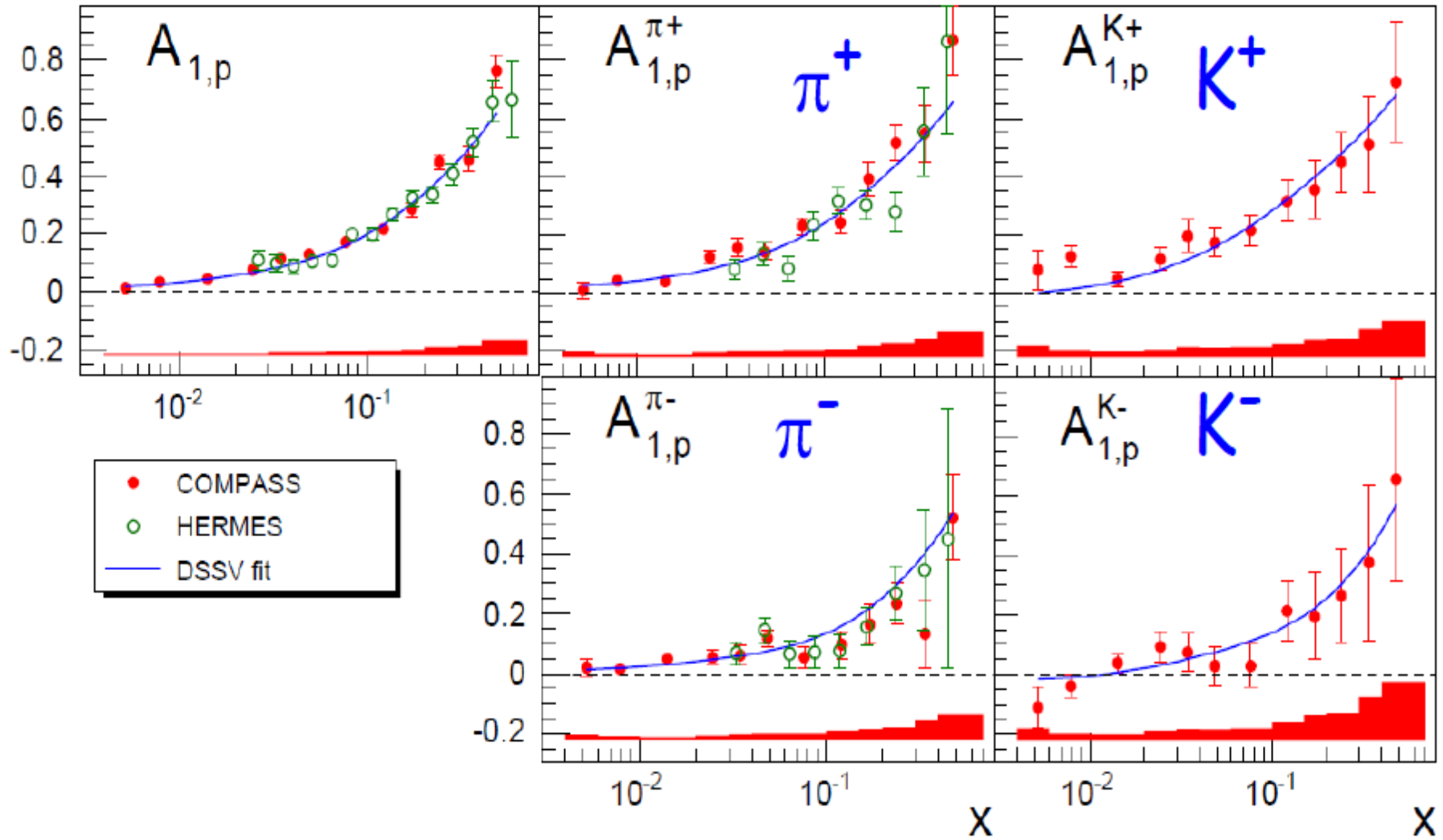


- From these asymmetries one can extract:

$$\Delta u + \Delta d, \quad \Delta \bar{u} + \Delta \bar{d} \quad \text{and} \quad \Delta s = \Delta \bar{s}$$

Inclusive and semi-inclusive spin asymmetries: Proton data

First measurement ever of $A_{1,p}^K$

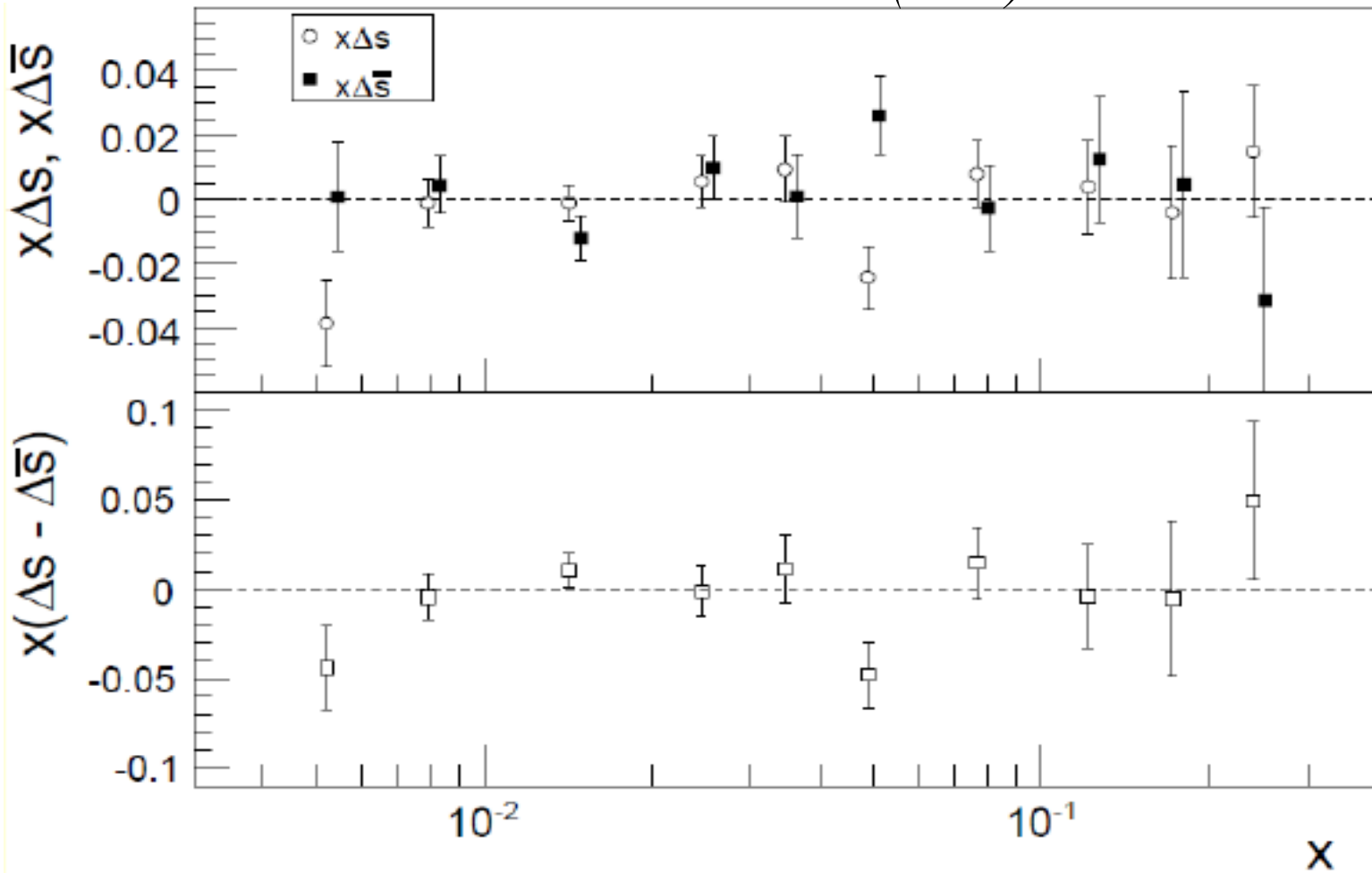


- Using $A_{1,p}^h$ and $A_{1,d}^h$, we can separately extract:

Δu , Δd , $\Delta \bar{u}$, $\Delta \bar{d}$, Δs and $\Delta \bar{s}$

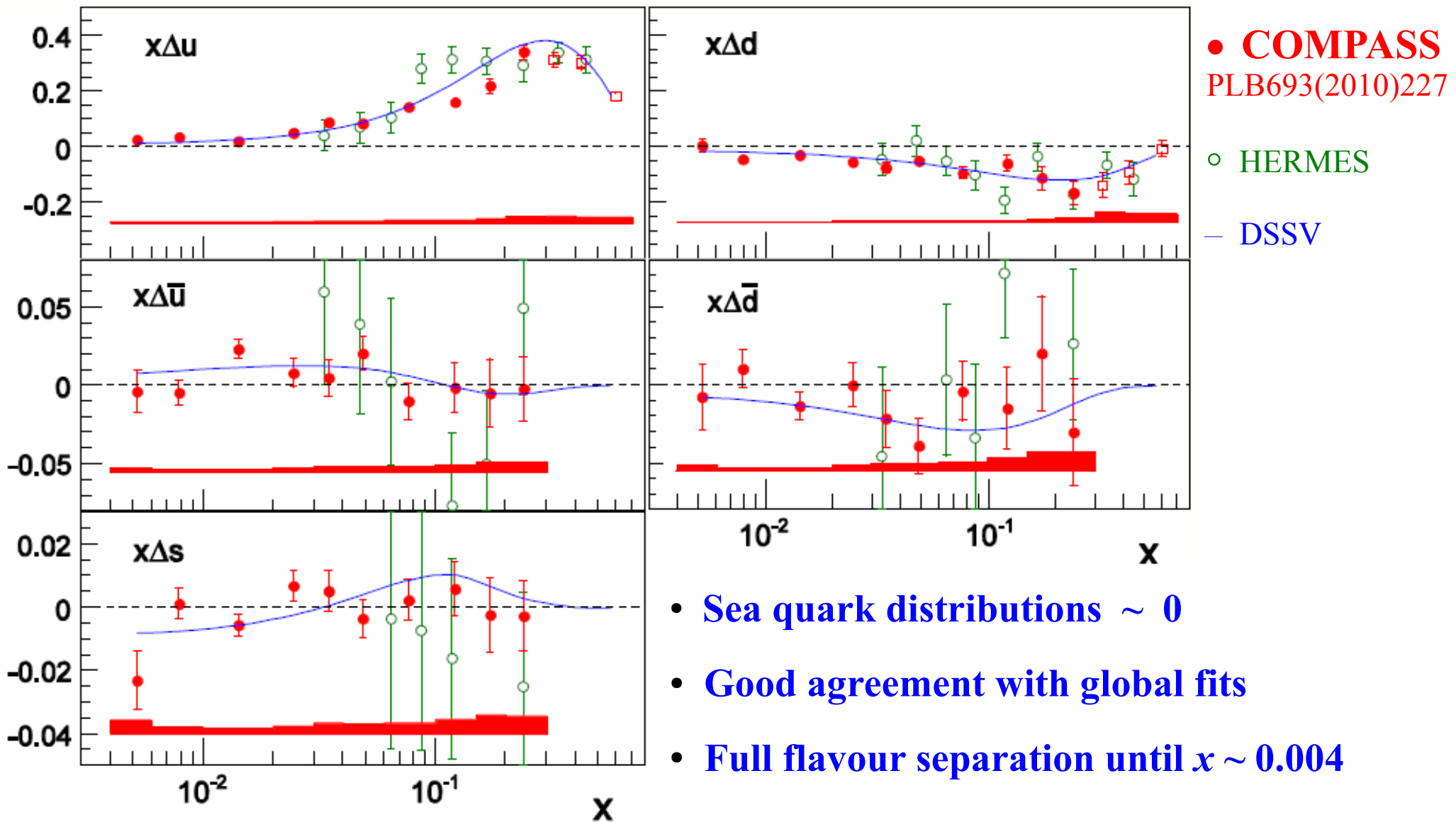
Comparison of Δs with $\Delta \bar{s}$

COMPASS PLB 693 (2010) 227



$\Delta s - \Delta \bar{s}$ is compatible with 0 \rightarrow $\Delta s = \Delta \bar{s}$ is assumed in the subsequent analysis

Quark helicities from SIDIS ($Q^2 = 3 \text{ (GeV/c)}^2$ and $x < 0.3$)

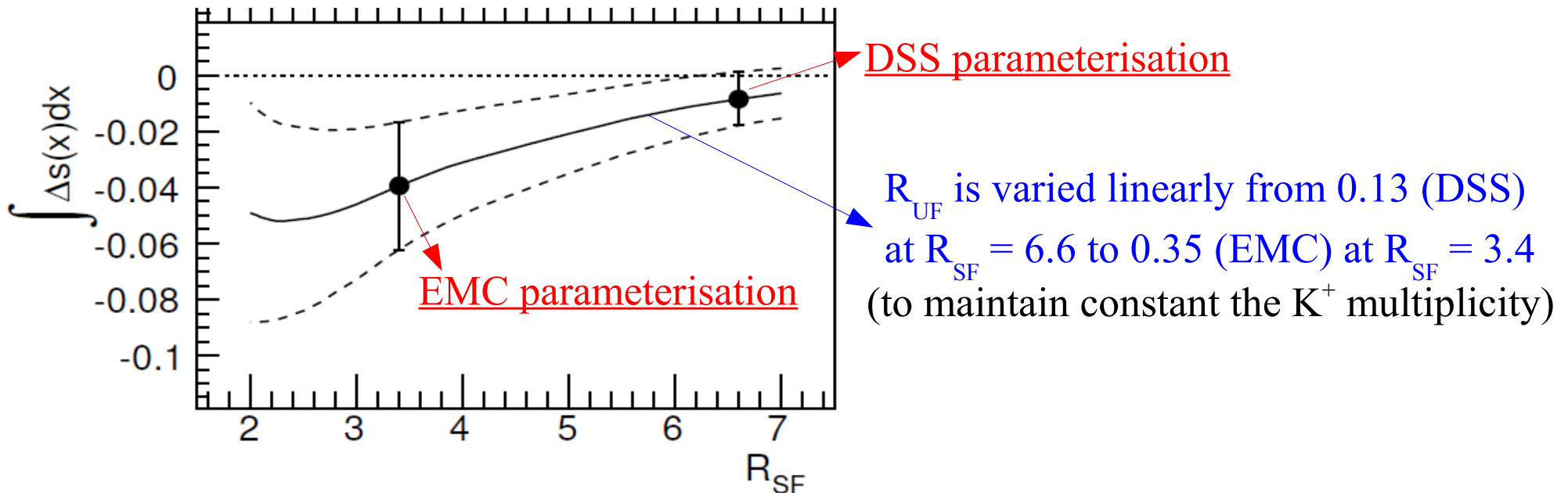


$$\Delta s(\text{SIDIS}) = -0.01 \pm 0.01(\text{stat}) \pm 0.01(\text{syst}) \quad @ \quad 0.003 < x < 0.3$$

Δs dependence on FFs

- The relation between the semi-inclusive asymmetries and Δs depends only on the following ratios:

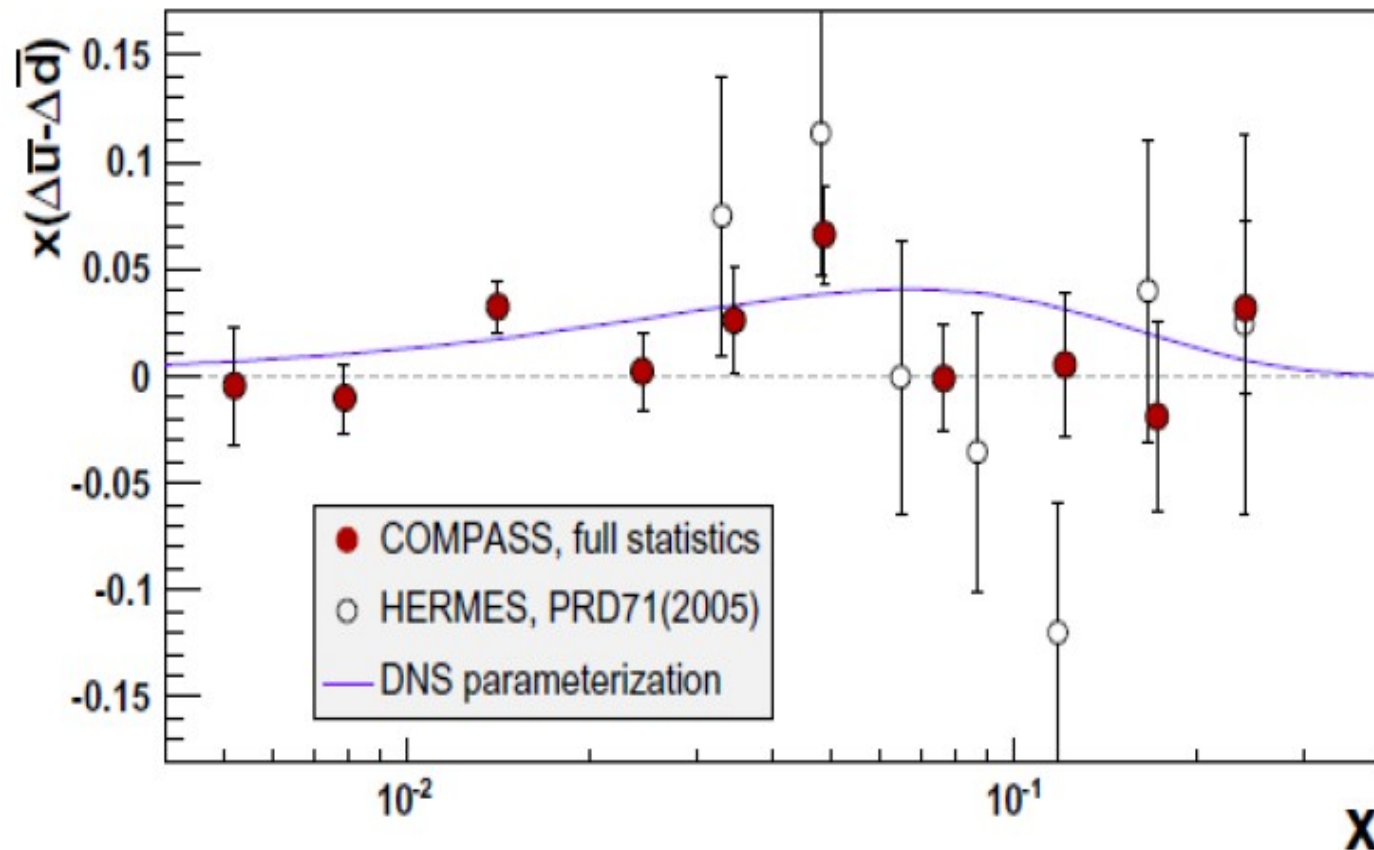
$$\mathbf{R}_{\text{UF}} = \frac{\int_{0.2}^{0.85} \mathbf{D}_d^{\text{K}^+}(z) dz}{\int_{0.2}^{0.85} \mathbf{D}_u^{\text{K}^+}(z) dz}, \quad \mathbf{R}_{\text{SF}} = \frac{\int_{0.2}^{0.85} \mathbf{D}_{\bar{s}}^{\text{K}^+}(z) dz}{\int_{0.2}^{0.85} \mathbf{D}_u^{\text{K}^+}(z) dz}$$



- Determination of R_{SF} from hadron multiplicities on the way

$\Delta\bar{u} - \Delta\bar{d}$: Flavour asymmetry?

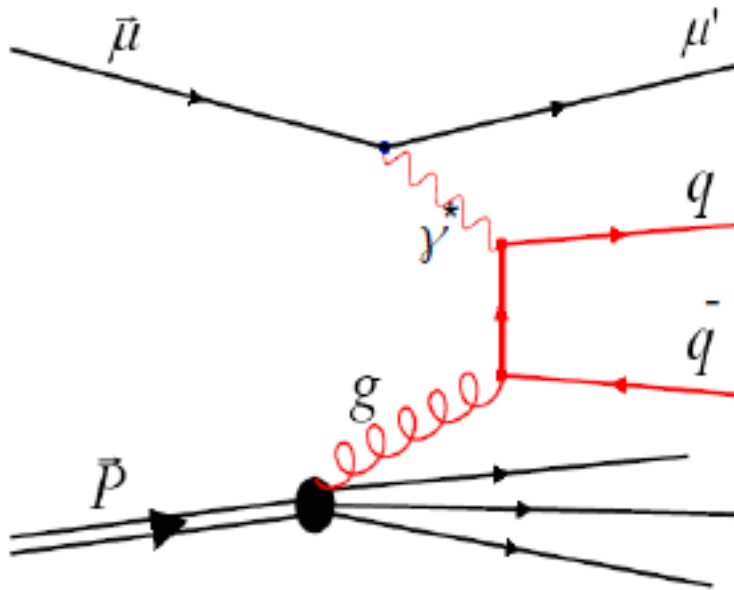
- The considerable asymmetry observed for $(\bar{u} - \bar{d})$ is not verified in the polarised case :
- $\Delta\bar{u} - \Delta\bar{d}$ is slightly positive but compatible with zero!



Gluon Polarisation

Direct measurement of $\Delta G/G$ in LO

photon-gluon fusion process (PGF)



$$A_{\gamma N}^{\text{PGF}} = \frac{\int d\hat{s} \Delta \sigma^{\text{PGF}} \Delta G(\mathbf{x}_G, \hat{s})}{\int d\hat{s} \sigma^{\text{PGF}} G(\mathbf{x}_G, \hat{s})}$$

$$\approx \langle \mathbf{a}_{\text{LL}}^{\text{PGF}} \rangle \frac{\Delta G}{G}$$

analysing power

There are two methods to tag this process:

- **Open Charm production**

- $\gamma^* g \rightarrow c\bar{c} \Rightarrow$ reconstruct D^0 mesons
- **Hard scale:** M_c^2
- **No intrinsic charm in COMPASS kinematics**
- **No physical background**
- **Weakly Monte Carlo dependent**
- **Low statistics**

- **High- p_T hadron pairs**

- $\gamma^* g \rightarrow q\bar{q} \Rightarrow$ reconstruct 2 jets or h^+h^-
- **Hard scale:** Q^2 or Σp_T^2 [$Q^2 > 1$ or $Q^2 < 1$ (GeV/c) 2]
- **High statistics**
- **Physical background**
- **Strongly Monte Carlo dependent**

Open Charm

Open Charm analysis: Simultaneous extraction of $\Delta G/G$ and A^{bg}

- The relation between the number of reconstructed D^0 (for each target cell configuration) and $\Delta G/G$ is given by:

$$N_t = a \phi n (S+B) \left(1 + f P_T P_\mu \left[a_{LL} \frac{S}{S+B} \frac{\Delta G}{G} + D \frac{B}{S+B} A^{bg} \right] \right), \quad t=(u, d, u', d')$$

acceptance, muon flux, number of target nucleons

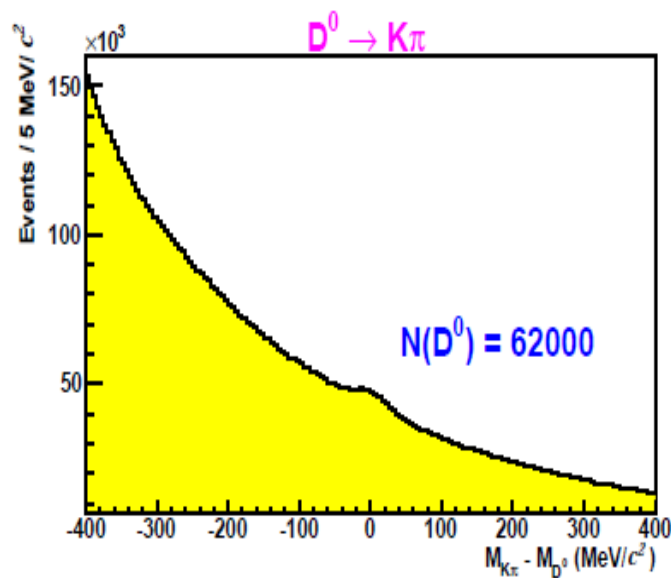
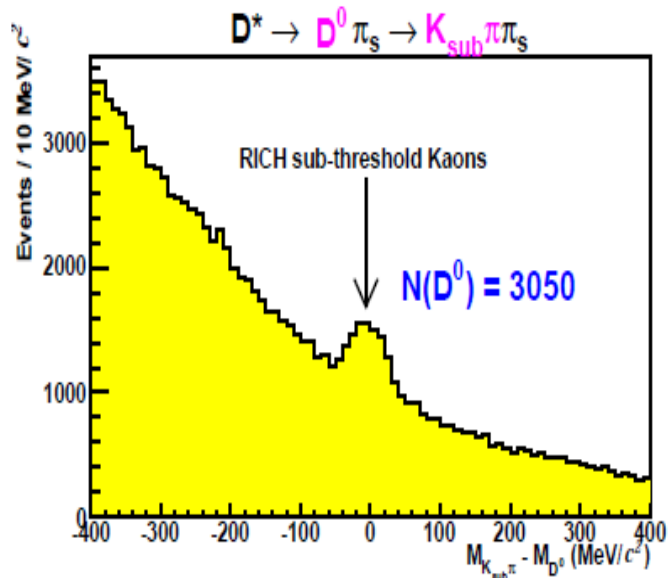
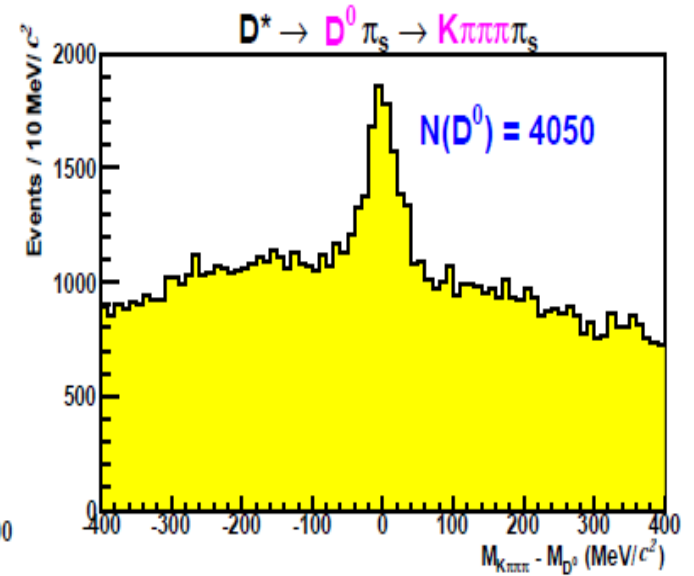
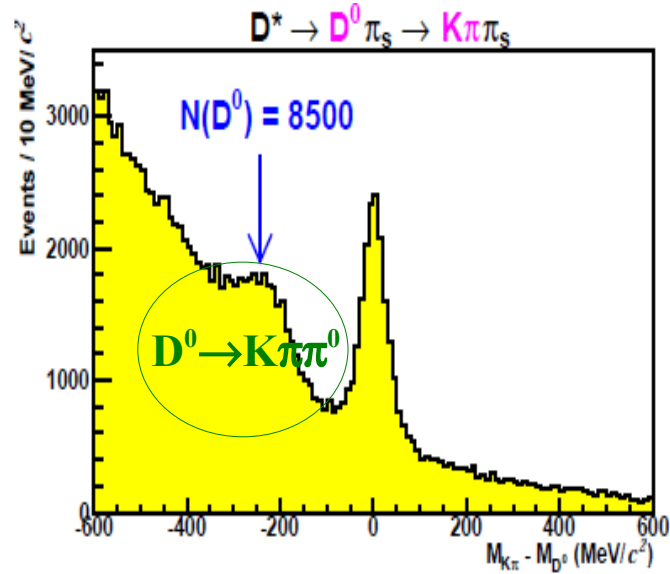
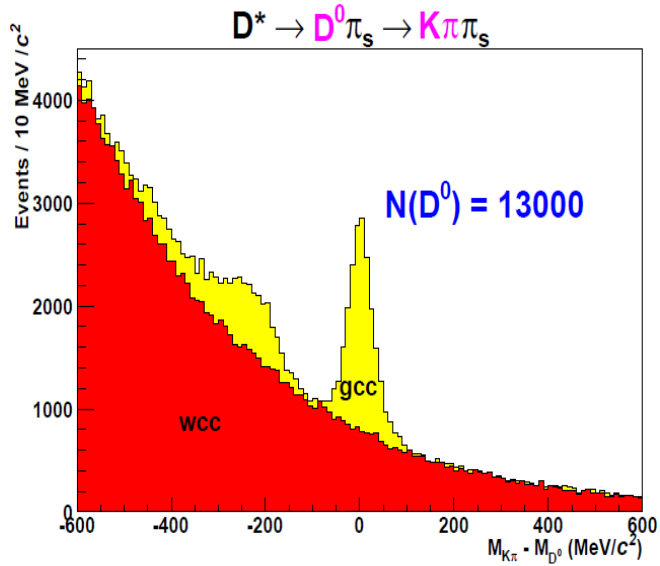
Open Charm event probability

- Each equation is weighted with a signal weight $\omega_S = f P_\mu a_{LL} S/(S+B)$ and also with a background weight $\omega_B = f P_\mu D B/(S+B)$:

8 equations with 7 unknowns: $\Delta G/G$, A^{bg} + 5 independent $\alpha = (a\phi n)$ factors

The system is solved by a χ^2 minimisation

D^0 invariant mass spectra: All samples (2002-2007 data)



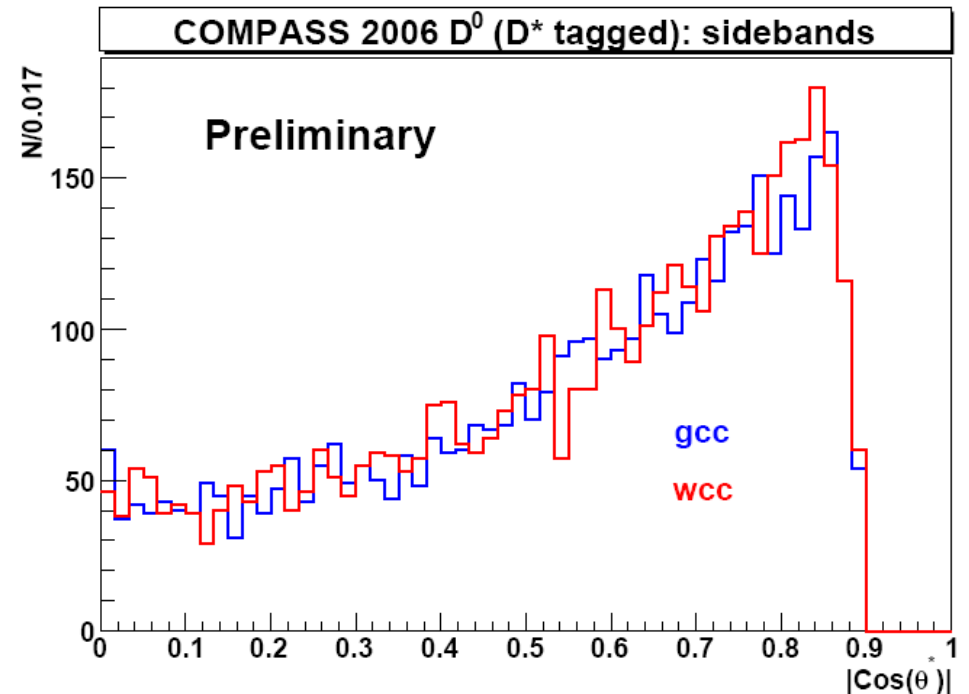
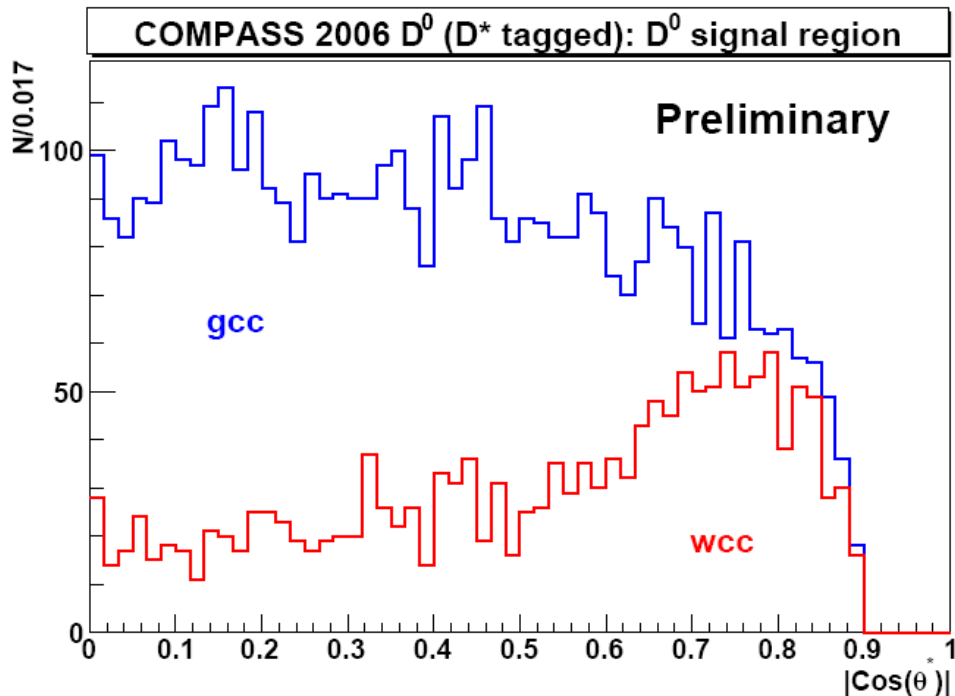
Number of D^0 :

- Total \rightarrow 90600
- ${}^6\text{LiD}$ \rightarrow 65600
- NH_3 \rightarrow 25000

Neural Network qualification of events

- **Two real data samples** (with the same cuts applied) **are compared by a Neural Network** (using some kinematic variables as a learning vector):
 - **Signal model** \rightarrow **gcc** = $\mathbf{K}^+ \pi^- \pi_s^- + \mathbf{K}^- \pi^+ \pi_s^+$ (D^0 spectrum: signal + background)
 - **Background model** \rightarrow **wcc** = $\mathbf{K}^+ \pi^+ \pi_s^- + \mathbf{K}^- \pi^- \pi_s^+$ (no D^0 is allowed)
- **If the background model is good enough:** The Neural Network is able to distinguish the signal from the combinatorial background on a event by event basis (inside gcc)

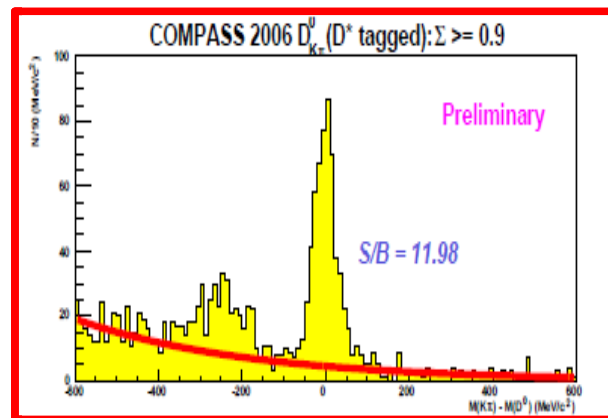
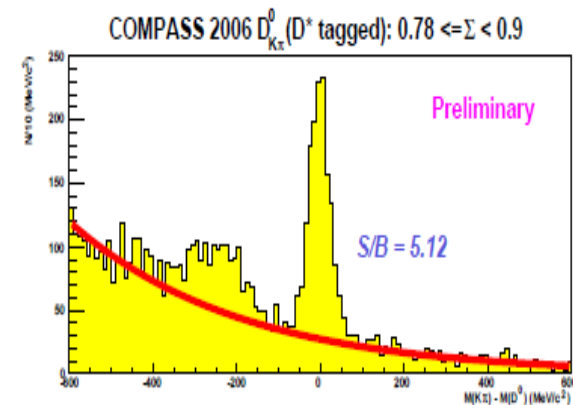
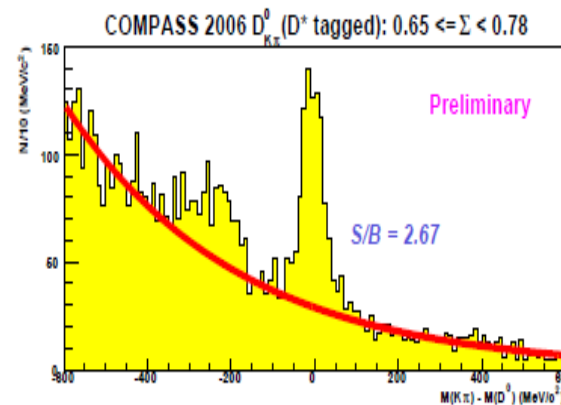
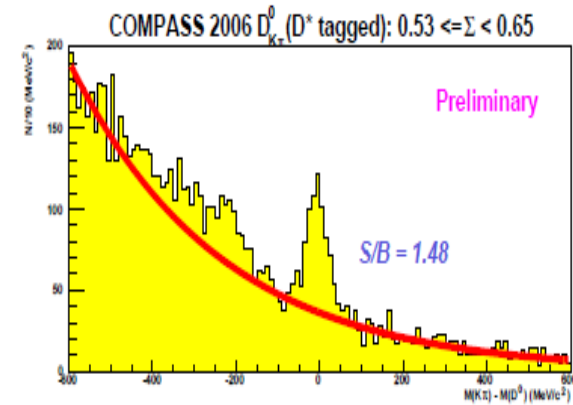
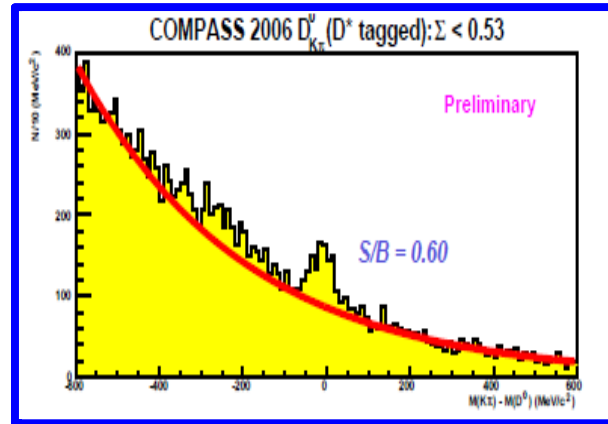
Example of a good learning variable



S/(S+B): Obtaining final probabilities for a D⁰ candidate

D⁰_{Kπ} tagged spectrum in bins of $\Sigma = S/(S+B)_{NN}$

- Events with small $S/(S+B)_{NN}$
 - Mostly combinatorial background is selected



S/(S+B) is obtained from a fit inside this bins (correcting with the NN parameterisation)

- Events with large $S/(S+B)_{NN}$
 - Mostly Open Charm events are selected

$$\delta \left(\frac{\Delta G}{G} \right) \propto \frac{1}{\text{FOM}}$$

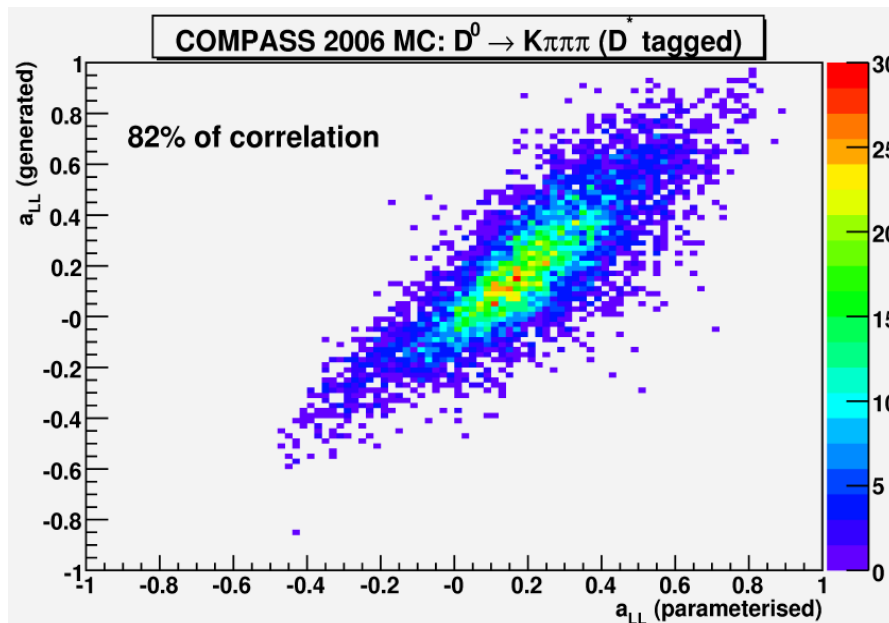
Analysing power (muon-gluon asymmetry a_{LL})

- a_{LL} is dependent on the full knowledge of the partonic kinematics:

$$a_{LL} = \frac{\Delta \sigma^{\text{PGF}}}{\sigma_{\text{PGF}}}(y, Q^2, x_g, z_C, \phi)$$

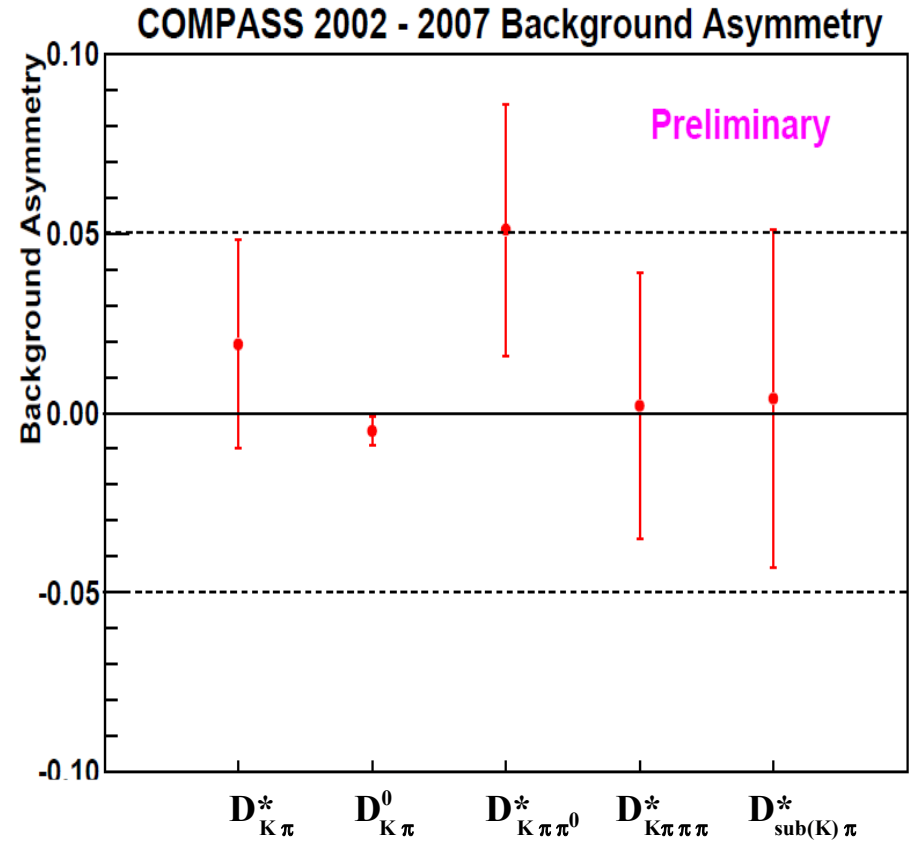
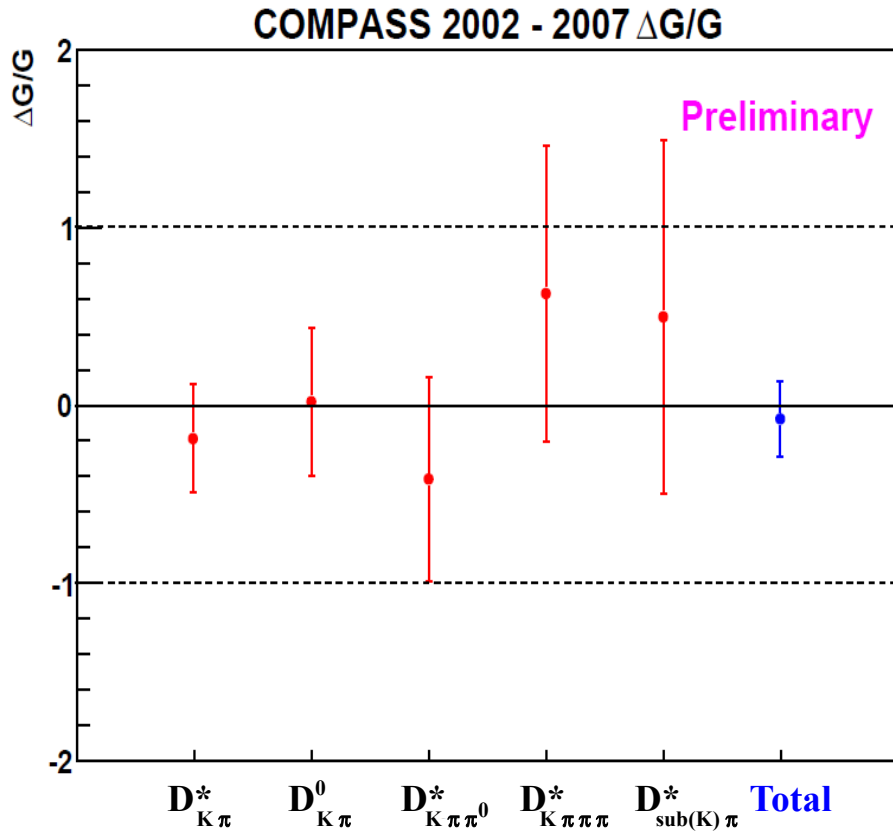
↓
Can't be experimentally obtained: only one charmed meson is reconstructed

- a_{LL} is obtained from Monte-Carlo (in LO), to serve as input for a Neural Network parameterisation on some reconstructed kinematical variables: y , x_{Bj} , Q^2 , z_D and p_T



↘
Parameterised a_{LL} , shows a strong correlation with the generated one (using AROMA)

Open Charm results in LO



$$\frac{\Delta G}{G} = -0.08 \pm 0.21(\text{stat}) \pm 0.08(\text{syst}) \quad @ \langle \mathbf{x}_g \rangle = 0.11_{-0.05}^{+0.11}, \quad \langle \mu^2 \rangle = 13 \text{ (GeV/c)}^2$$

High- p_T hadron pairs

High- p_T asymmetries (2002-2006): $Q^2 > 1$ (GeV/c)²

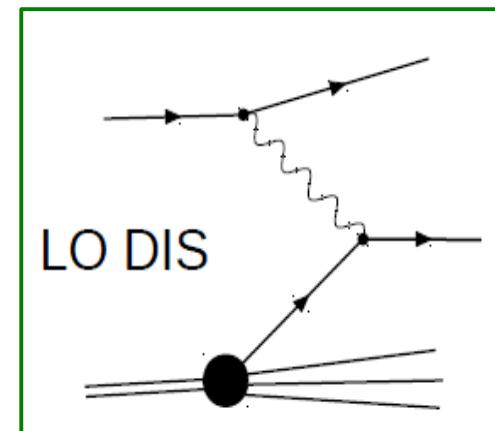
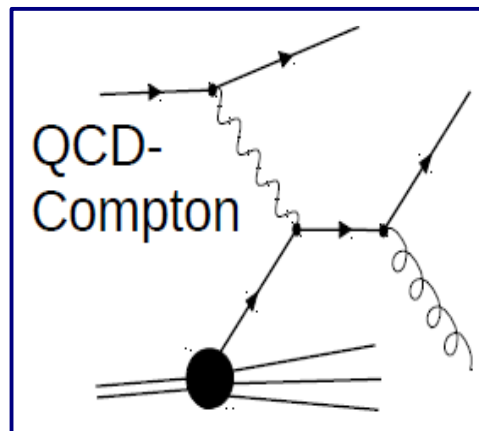
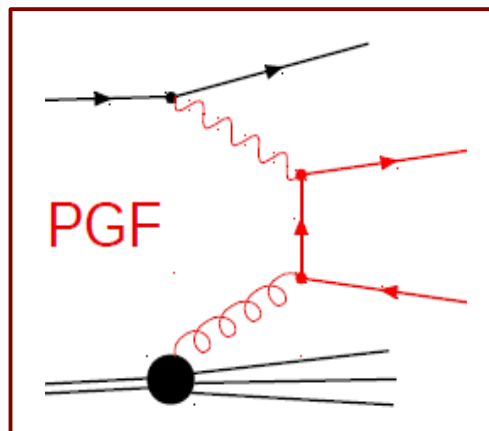
- Two samples are considered:

Inclusive asymmetry

$$A_1^d(\mathbf{x}) = \frac{\Delta G}{G}(\mathbf{x}_g) \left(a_{LL}^{\text{PGF,inc}} \frac{\sigma^{\text{PGF,inc}}}{\sigma^{\text{Tot,inc}}} \right) + A_1^{\text{LO}}(\mathbf{x}_C) \left(a_{LL}^{\text{C,inc}} \frac{\sigma^{\text{C,inc}}}{\sigma^{\text{Tot,inc}}} \right) + A_1^{\text{LO}}(\mathbf{x}_{\text{Bj}}) \left(D \frac{\sigma^{\text{LO,inc}}}{\sigma^{\text{Tot,inc}}} \right)$$

$$A_{LL}^{2h}(\mathbf{x}) = \left(\frac{A^{\text{exp}}}{f P_\mu P_T} \right) = \frac{\Delta G}{G}(\mathbf{x}_g) \left(a_{LL}^{\text{PGF}} \frac{\sigma^{\text{PGF}}}{\sigma^{\text{Tot}}} \right) + A_1^{\text{LO}}(\mathbf{x}_C) \left(a_{LL}^{\text{C}} \frac{\sigma^{\text{C}}}{\sigma^{\text{Tot}}} \right) + A_1^{\text{LO}}(\mathbf{x}_{\text{Bj}}) \left(D \frac{\sigma^{\text{LO}}}{\sigma^{\text{Tot}}} \right)$$

high- p_T hadron pairs ($p_{T1} / p_{T2} > 0.7 / 0.4$ GeV/c) \Rightarrow enhancement of the PGF contribution



Extraction of $\Delta G/G$ from high- p_T : $Q^2 > 1$ (GeV/c)²

- The gluon polarisation is determined from two asymmetry samples: the **two high- p_T hadrons** and the **inclusive** data samples. **The final formula is:**

$$\frac{\Delta g}{g}(x_g) = \frac{1}{\beta} \left[A_{LL}^{2h}(x) + A_{\text{corr}} \right] \quad A_{\text{corr}} = - \left(A_1(x_{Bj}) D \frac{R_{LO}}{R_{LO}^{\text{inc}}} - A_1(x_C) \beta_1 + A_1(x_{C'}) \beta_2 \right)$$

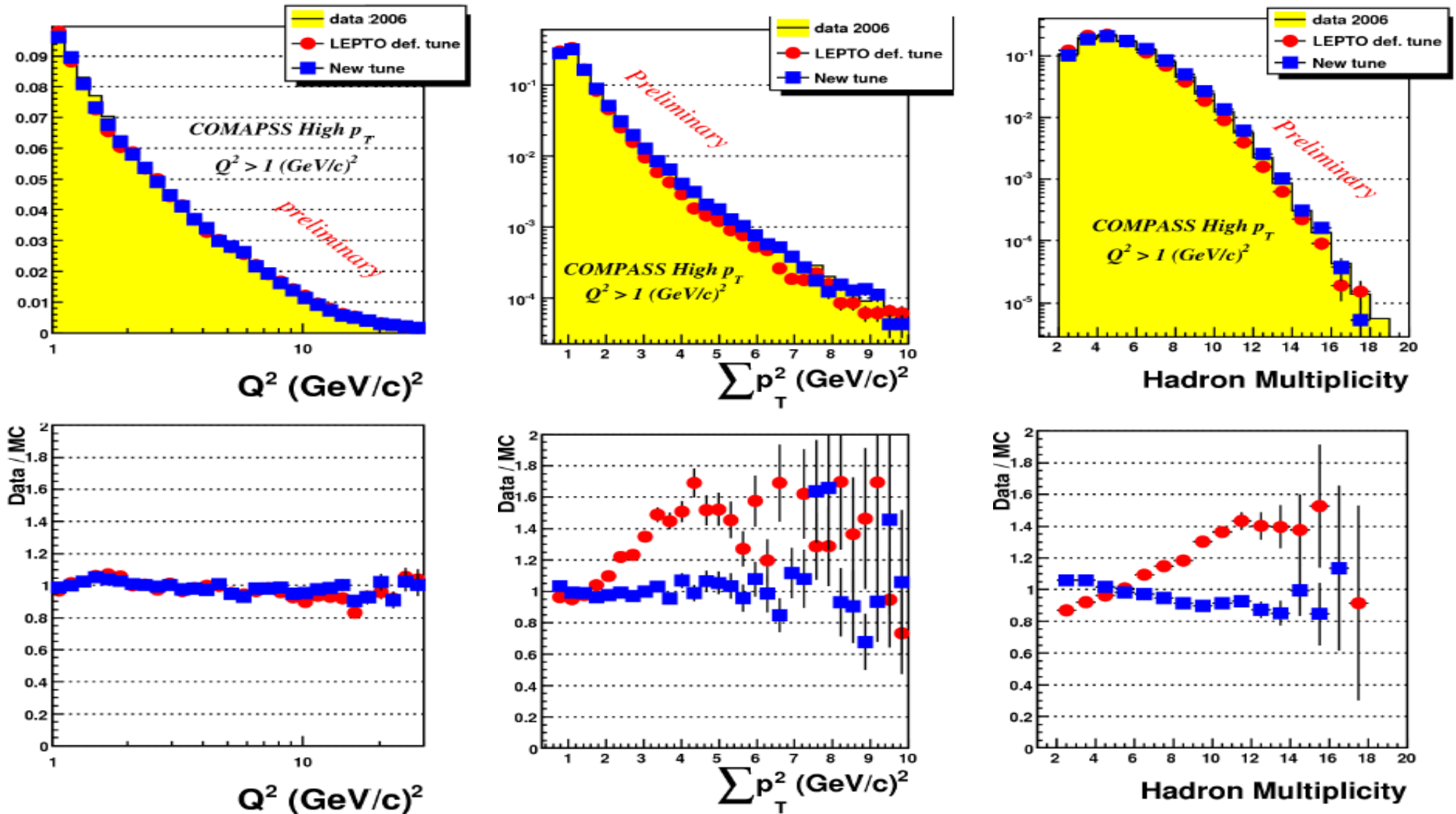
$$\beta = a_{LL}^{\text{PGF}} R_{\text{PGF}} - a_{LL}^{\text{PGF,inc}} R_{\text{PGF}}^{\text{incl}} \frac{R_{LO}}{R_{LO}^{\text{inc}}} - a_{LL}^{\text{PGF,inc}} \frac{R_C R_{\text{PGF}}^{\text{inc}} a_{LL}^C}{R_{LO}^{\text{inc}} D}$$

- β_1 and β_2 are factors depending on a_{LL}^i and R_i
- Each event is weighted with $\omega = f D P_\mu \beta \rightarrow$ statistical improvement!
- The following parameters are obtained from Monte Carlo, and then they are parameterised event-by-event by a Neural Network (to allow for their use in data):

$$R_{\text{PGF}}, R_C, R_{LO}, R_{\text{PGF}}^{\text{inc}}, R_C^{\text{inc}}, R_{LO}^{\text{inc}}, a_{LL}^{\text{PGF}}, a_{LL}^C, a_{LL}^{LO}, a_{LL}^{\text{PGF,inc}}, a_{LL}^{C,inc} \text{ and } a_{LL}^{LO,inc}$$

Data vs Monte Carlo: Comparison of Q^2 and hadron variables

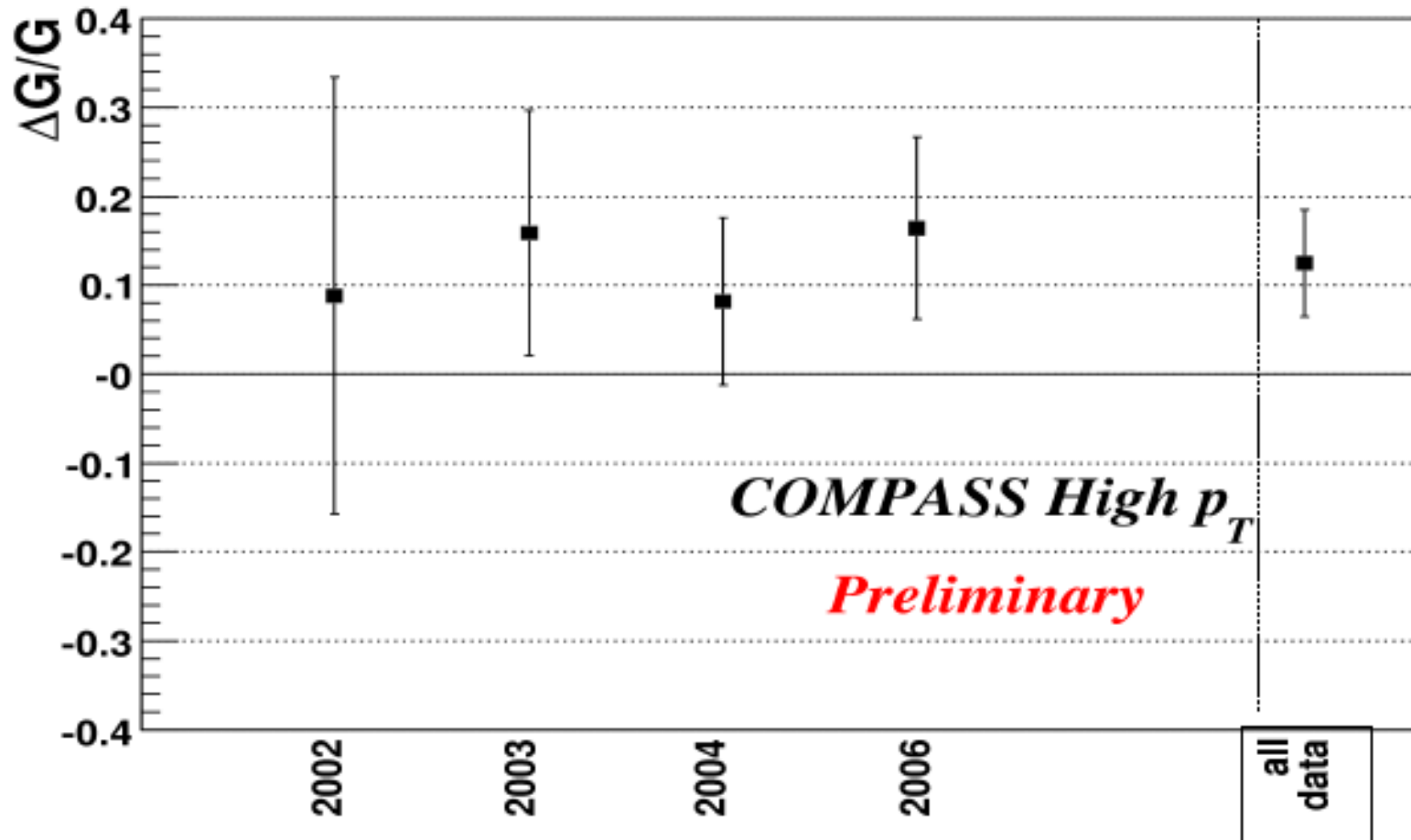
Monte Carlo (PS on): **LEPTO** generator with PDFs from **MSTW2008LO**



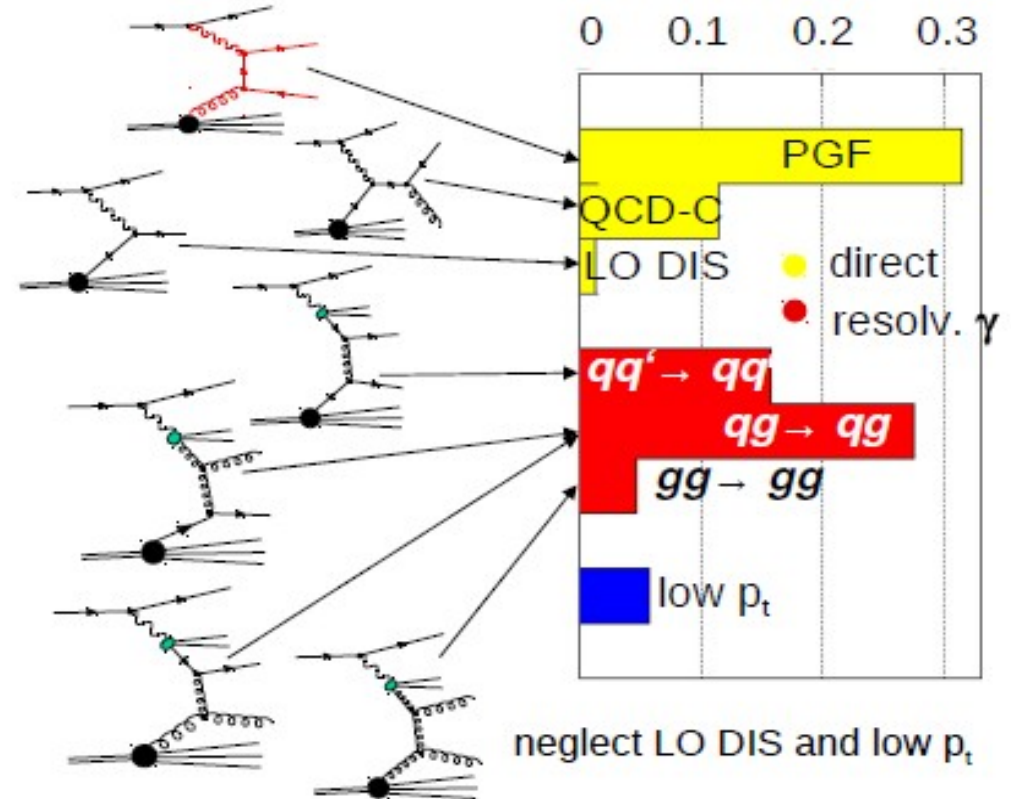
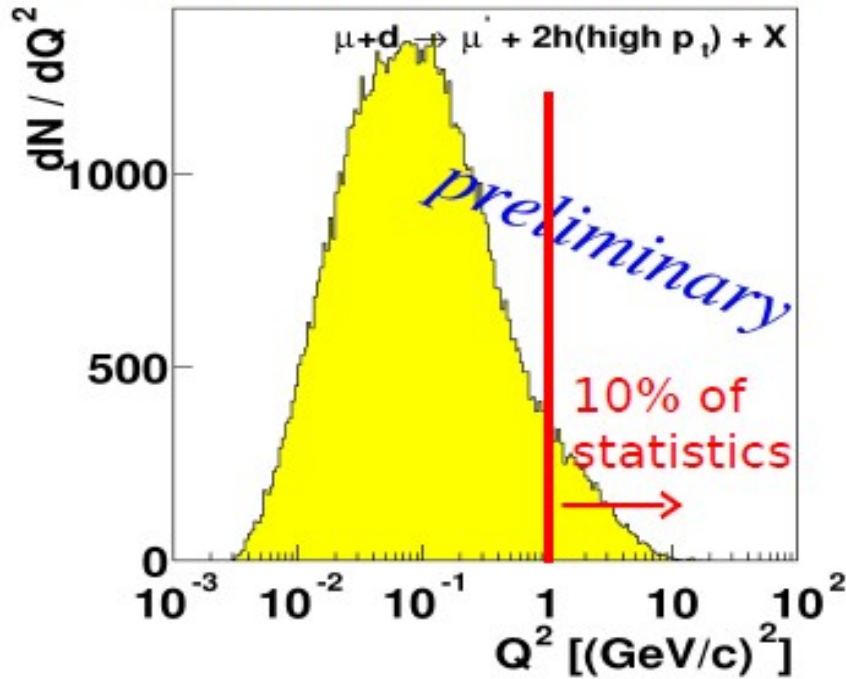
The impact of this tuning is included in the systematic error

High- p_T results: $Q^2 > 1 \text{ (GeV/c)}^2$

$$\frac{\Delta G}{G} = 0.125 \pm 0.060(\text{stat}) \pm 0.063(\text{syst}) \quad @ \langle x_g \rangle = 0.09_{-0.04}^{+0.08}, \langle \mu^2 \rangle = 3.4 \text{ (GeV/c)}^2$$



High- p_T analysis: $Q^2 < 1 \text{ (GeV/c)}^2$



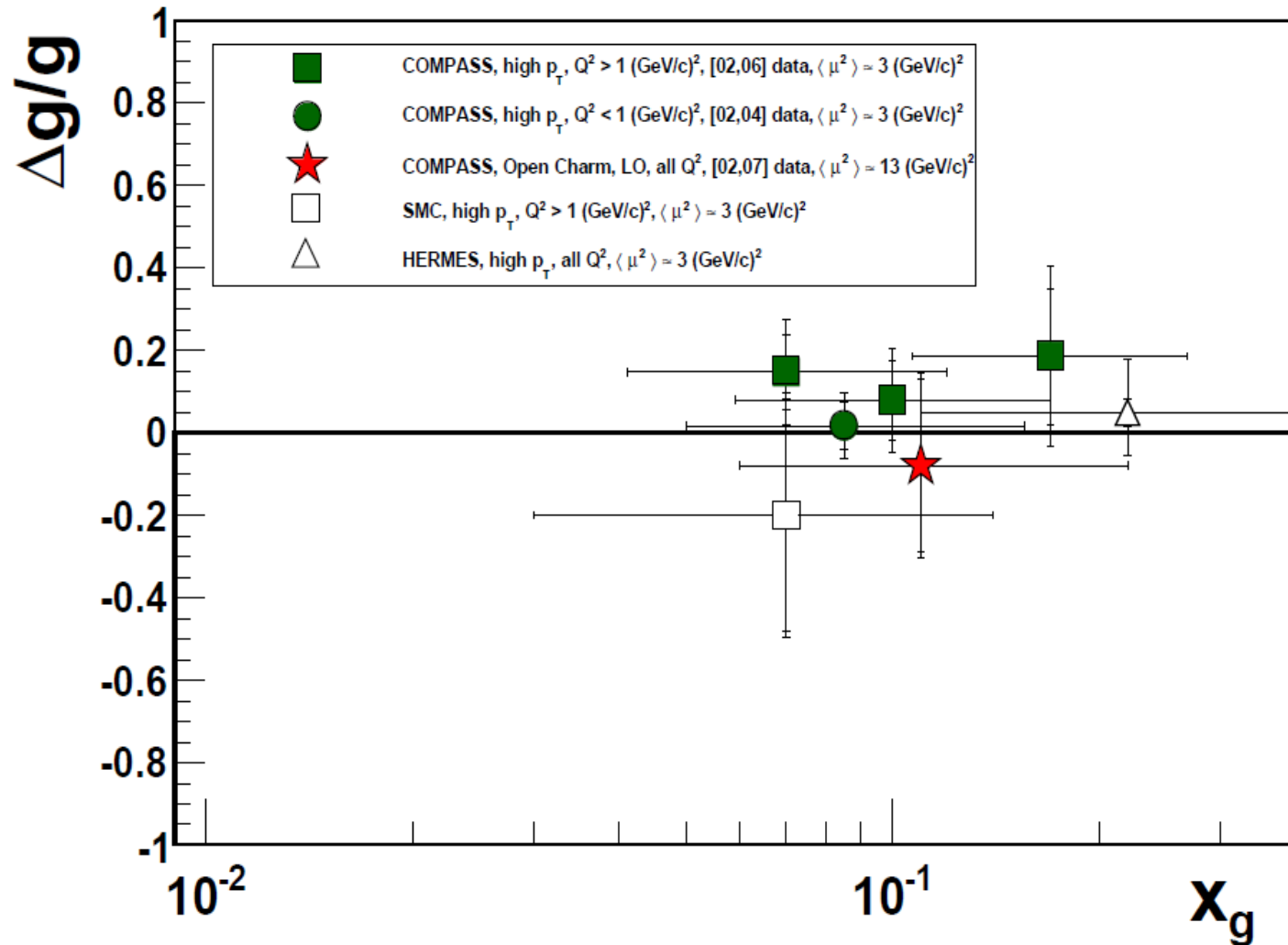
2002-2004 Preliminary:

$$\Delta G/G = 0.016 \pm 0.058 \text{ (stat)} \pm 0.055 \text{ (syst)}$$

2002-2003 Published:

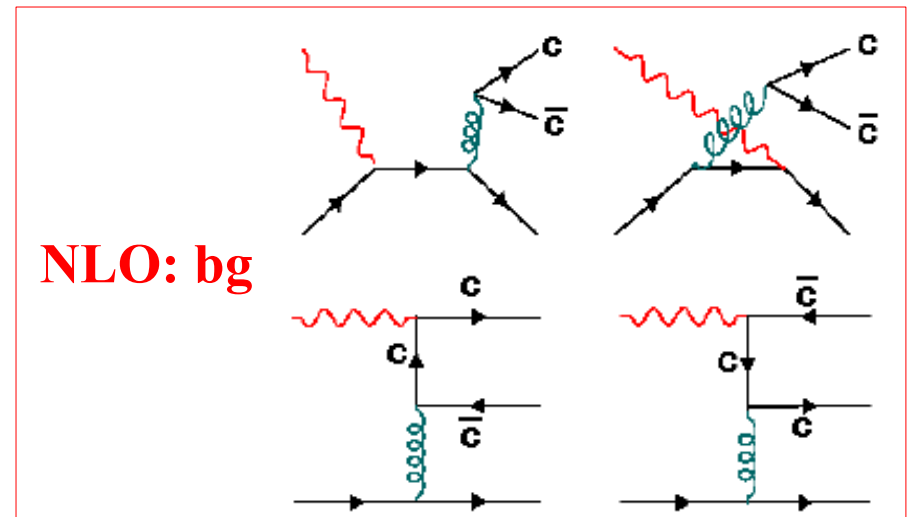
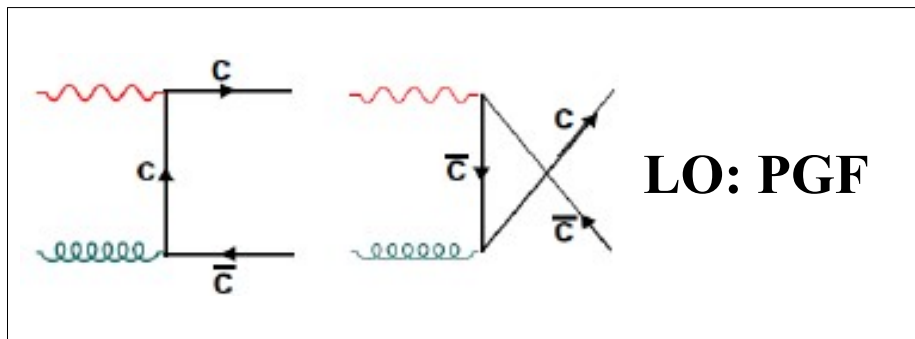
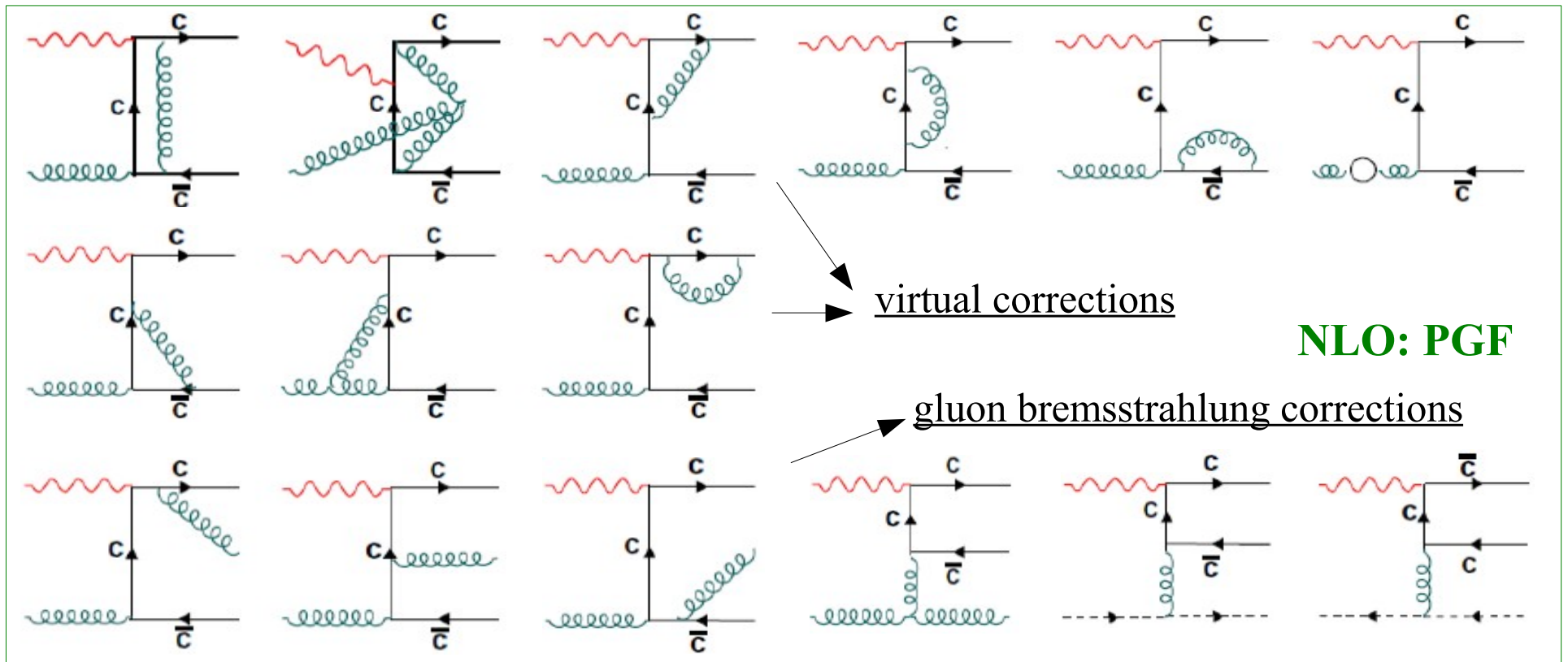
$$\Delta G/G = 0.024 \pm 0.089 \text{ (stat)} \pm 0.057 \text{ (syst)} \rightarrow \textit{Phys. Lett. B 633 (2006) 25 - 32}$$

World measurements on $\Delta G/G$ in LO



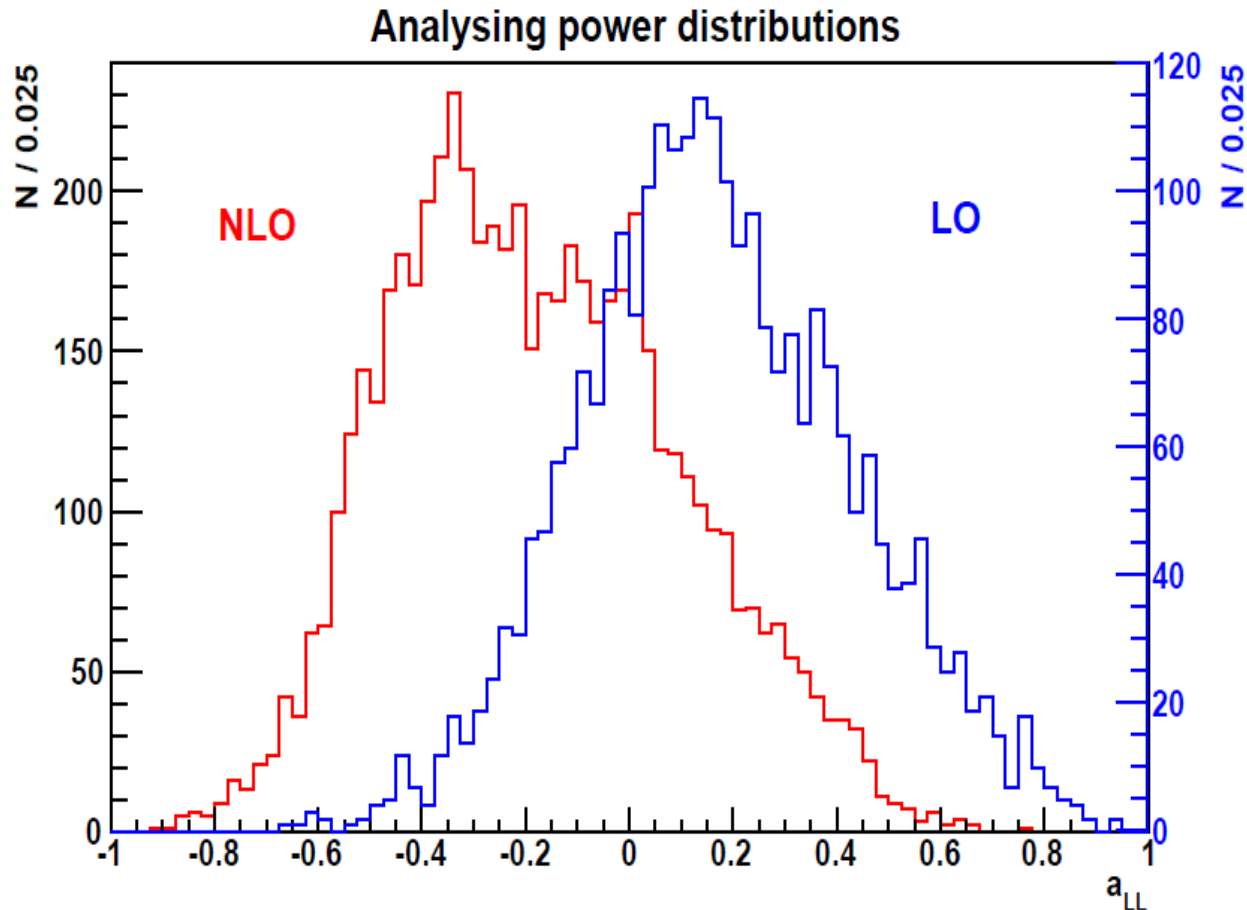
NLO results from Open Charm

NLO corrections to the analysing power $\langle a_{LL} \rangle$



Comparison of a_{LL} (LO) with a_{LL} (NLO)

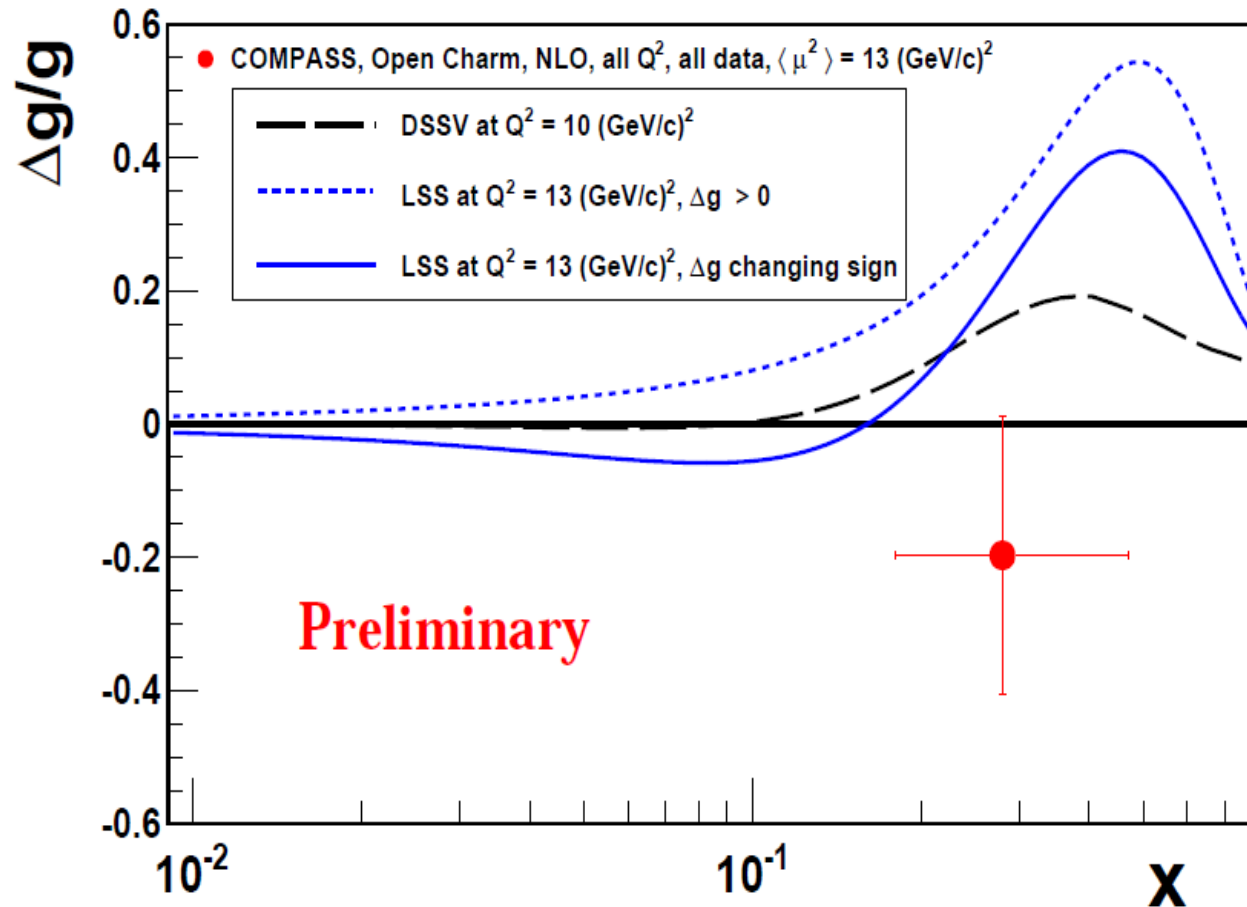
- The AROMA generator is used to simulate the phase space for the **NLO (PS on)** / **LO (PS off)** calculations of a_{LL} . The resulting D^0 mesons are reconstructed in the COMPASS spectrometer like real events. The respective a_{LL} distributions are:



NLO results for $A_{\gamma N}^{\text{PGF}}$: $A_{\gamma N} = \left(\frac{a_{\text{LL}}^{\text{PGF}}(\text{NLO})}{\mathbf{D}} \frac{\Delta \mathbf{G}}{\mathbf{G}} + \frac{a_{\text{LL}}^{\text{q}}(\text{NLO})}{\mathbf{D}} A_1 \right) \rightarrow A_{\text{corr}}$

Bins		$D^0 \rightarrow K\pi$ samples			$D^0 \rightarrow K\pi\pi^0$ sample			$D^0 \rightarrow K\pi\pi\pi$ sample		
$p_T(D^0)$ (GeV/c)	$E(D^0)$ (GeV)	$A_{\gamma N}$	$a_{\text{LL}}^{\text{PGF}}/\mathbf{D}$	A_{corr}	$A_{\gamma N}$	$a_{\text{LL}}^{\text{PGF}}/\mathbf{D}$	A_{corr}	$A_{\gamma N}$	$a_{\text{LL}}^{\text{PGF}}/\mathbf{D}$	A_{corr}
	[0, 30[-0.90±0.63	0.00	0.01	-0.63±1.29	-0.11	0.01	7.03±4.74	-0.09	0.01
[0, 0.3[[30, 50[-0.19±0.48	-0.06	0.01	0.27±1.17	-0.08	0.01	-2.05±1.10	-0.08	0.01
	> 50	0.07±0.68	-0.12	0.02	-2.55±2.00	-0.11	0.02	0.17±1.83	-0.09	0.01
	[0, 30[-0.18±0.37	-0.08	0.01	-0.24±0.80	-0.17	0.01	-0.59±1.74	-0.10	0.02
[0.3,0.7[[30, 50[0.10±0.26	-0.19	0.02	0.49±0.69	-0.23	0.02	1.00±0.54	-0.20	0.02
	> 50	-0.04±0.36	-0.22	0.02	-1.28±1.03	-0.18	0.02	-1.75±0.84	-0.21	0.02
	[0, 30[-0.42±0.44	-0.26	0.01	0.55±0.95	-0.29	0.02	2.91±2.61	-0.19	0.01
[0.7,1.0[[30, 50[-0.36±0.29	-0.29	0.01	-0.53±0.76	-0.32	0.02	1.42±0.57	-0.31	0.02
	> 50	1.49±0.42	-0.33	0.03	-0.17±1.00	-0.36	0.03	1.69±0.81	-0.32	0.03
	[0, 30[-0.30±0.35	-0.35	0.01	1.35±0.86	-0.40	0.02	-1.89±2.64	-0.36	0.02
[1.0,1.5[[30, 50[0.13±0.23	-0.40	0.02	-0.11±0.51	-0.44	0.03	-0.45±0.51	-0.41	0.02
	> 50	-0.20±0.33	-0.43	0.03	-0.05±0.78	-0.42	0.04	1.06±0.66	-0.45	0.03
	[0, 30[0.38±0.49	-0.49	0.02	-0.19±1.14	-0.52	0.02	1.64±3.52	-0.49	0.03
> 1.5	[30, 50[-0.00±0.25	-0.53	0.03	-0.23±0.51	-0.50	0.04	0.44±0.68	-0.54	0.03
	> 50	0.36±0.33	-0.53	0.04	0.26±0.90	-0.49	0.05	0.08±0.63	-0.54	0.05

$\Delta G/G$ result in NLO \rightarrow NEW



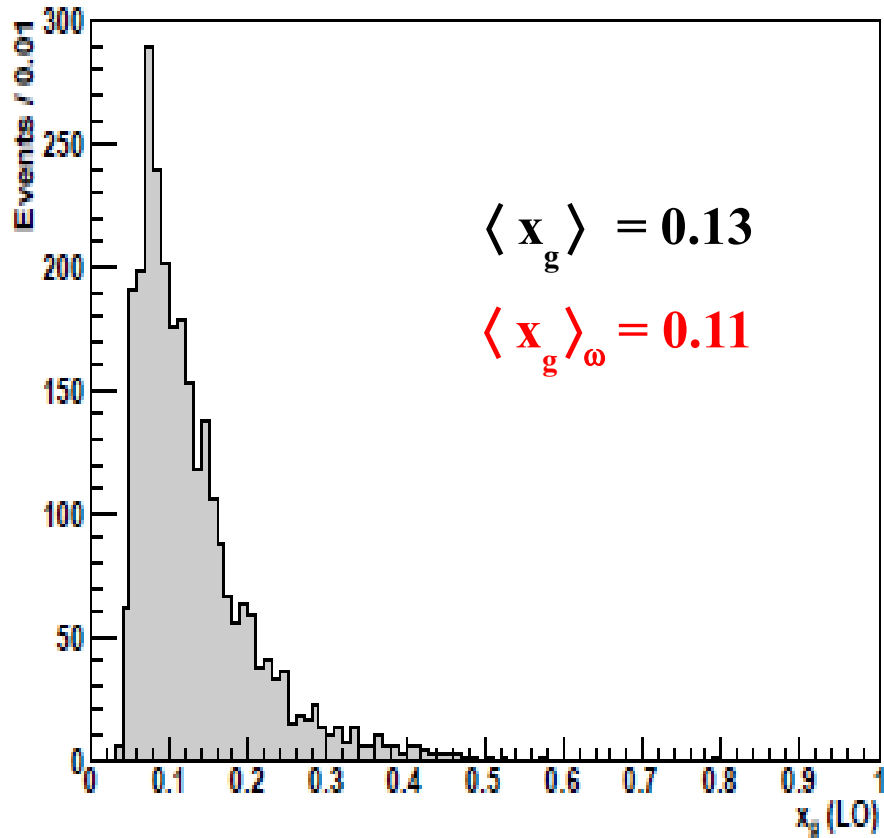
$$\frac{\Delta G}{G} = -0.20 \pm 0.21 \pm 0.08 \text{ (syst)} \quad @ \langle x_g \rangle = 0.28_{-0.10}^{+0.19}, \quad \langle \mu^2 \rangle = 13 \text{ (GeV/c)}^2$$

Preliminary: theoretical uncertainties still under study (a_{LL})

SPARES

Open Charm: Comparison of the x_g (LO) and x_g (NLO) distributions

LO



NLO

