

Azimuthal asymmetries in the unpolarized SIDIS mu-D cross section at COMPASS

Giulio Sbrizzai Trieste University and INFN

*on behalf of the COMPASS
collaboration*

COmmon

Muon and

Proton

Apparatus for

Structure and

Spectroscopy

NA58

**Czech Republic, Finland, France, Germany, India, Israel, Italy, Japan,
Poland, Portugal, Russia**

Bielefeld, Bochum, Bonn, Calcutta, CERN,

Dubna, Erlangen, Freiburg, Lisbon,

Mainz, Moscow, Munich, Prague, Protvino,

Saclay, Tel Aviv, Torino, Trieste, Warsaw, Yamagata

29 Institutes, ~230 physicists

muon beam

deuteron (${}^6\text{LiD}$)
polarized target

2002
2003
2004
2006
2007
2010

hadron beam

LH target

2008
2009

luminosity:

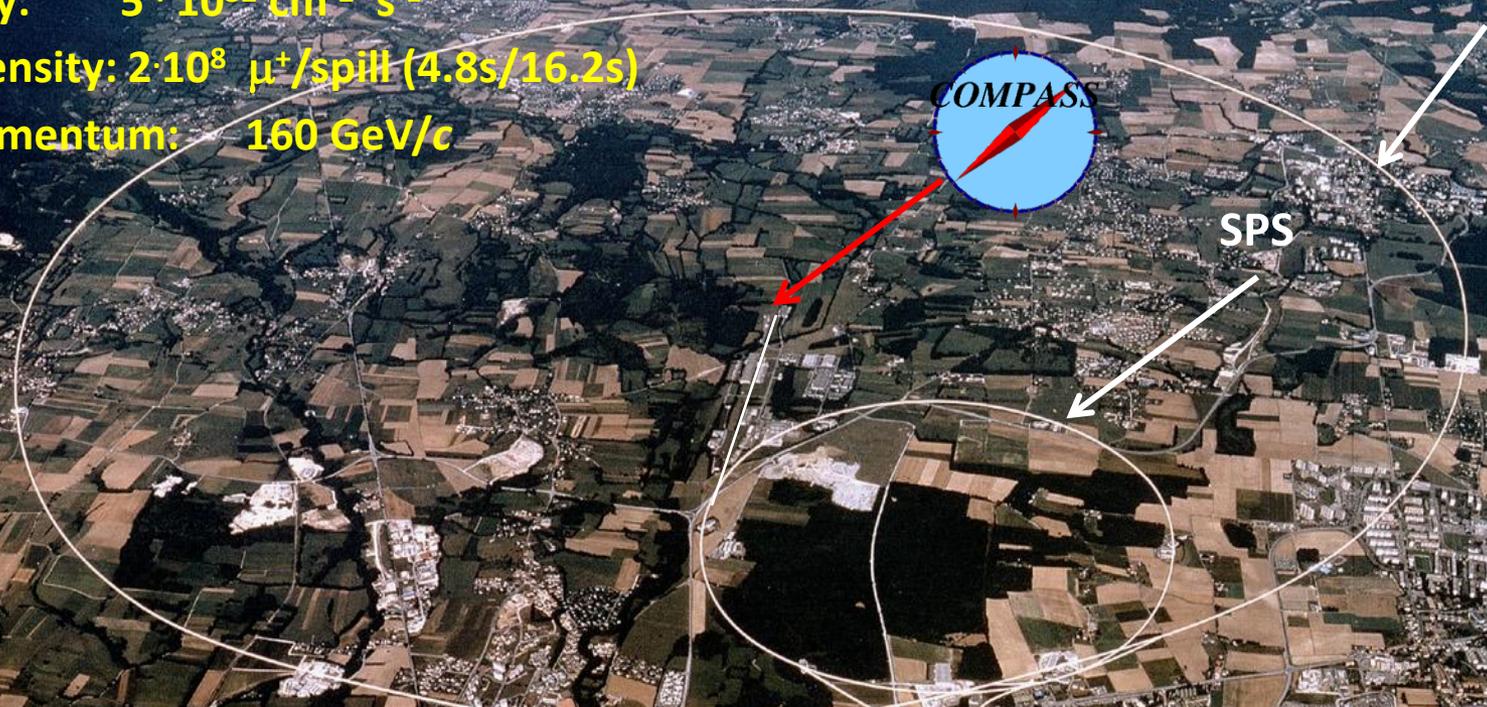
$\sim 5 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$

beam intensity:

$2 \cdot 10^8 \mu^+/\text{spill} (4.8\text{s}/16.2\text{s})$

beam momentum:

160 GeV/c



LHC

SPS

COMPASS

This talk will focus on the results extracted from the 2004 data (longitudinally pol. muon beam, transversely pol. deuteron target)

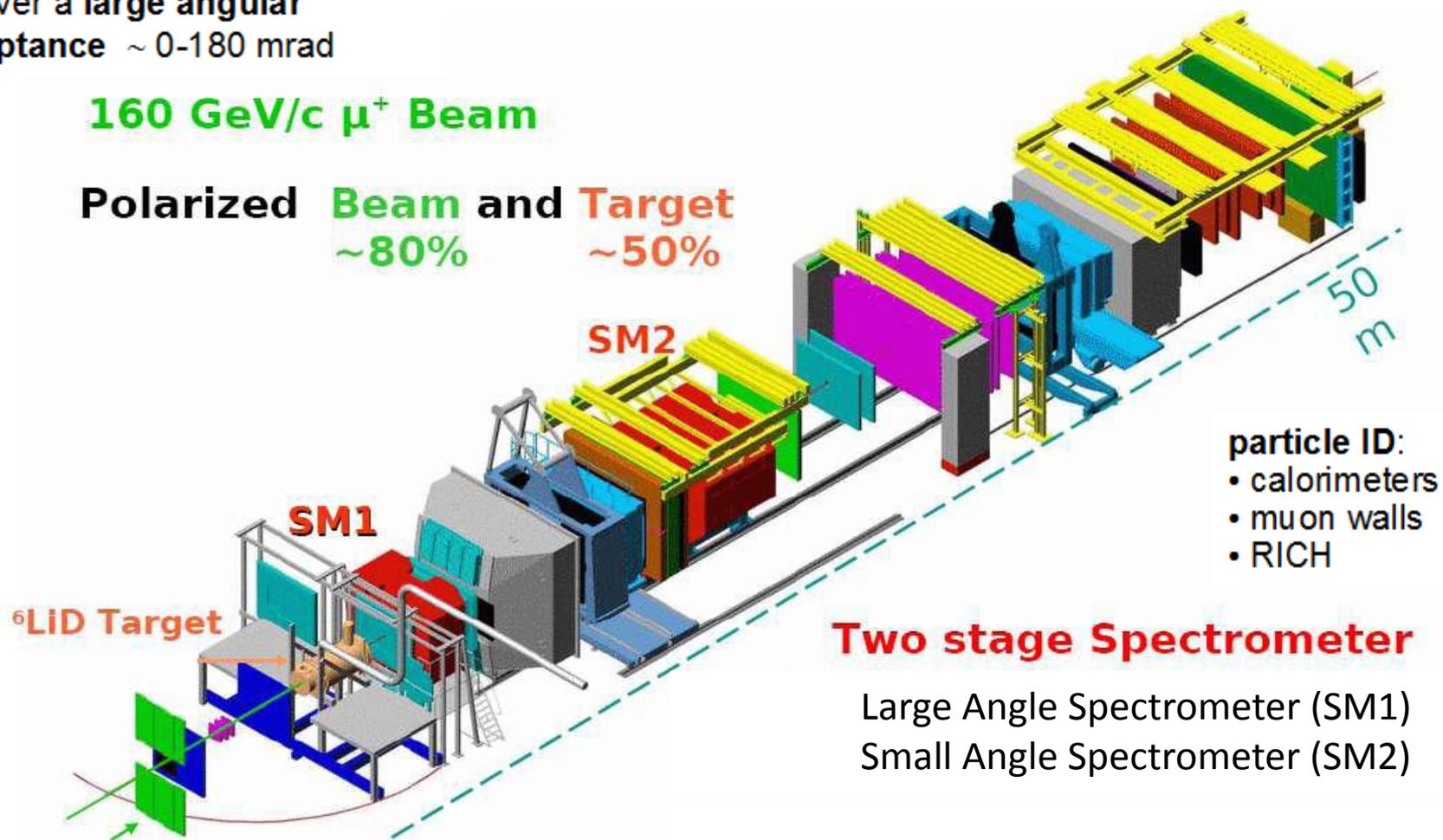
The COMPASS experiment (2004 setup)

several tracking detectors of different type

- to cope with **high particle rates**
- to cover a **large angular acceptance** $\sim 0-180$ mrad

160 GeV/c μ^+ Beam

Polarized Beam and Target
~80% **~50%**



OUTLINE

Introduction

The Analysis

- asymmetries extraction
- MC description
- acceptance studies
- systematic errors

Final results

Conclusions

SIDIS: a key process to investigate the structure of the nucleon

lepton interacts with a **single constituent** of the nucleon ($Q^2 > 1 \text{ GeV}^2/c^2$)

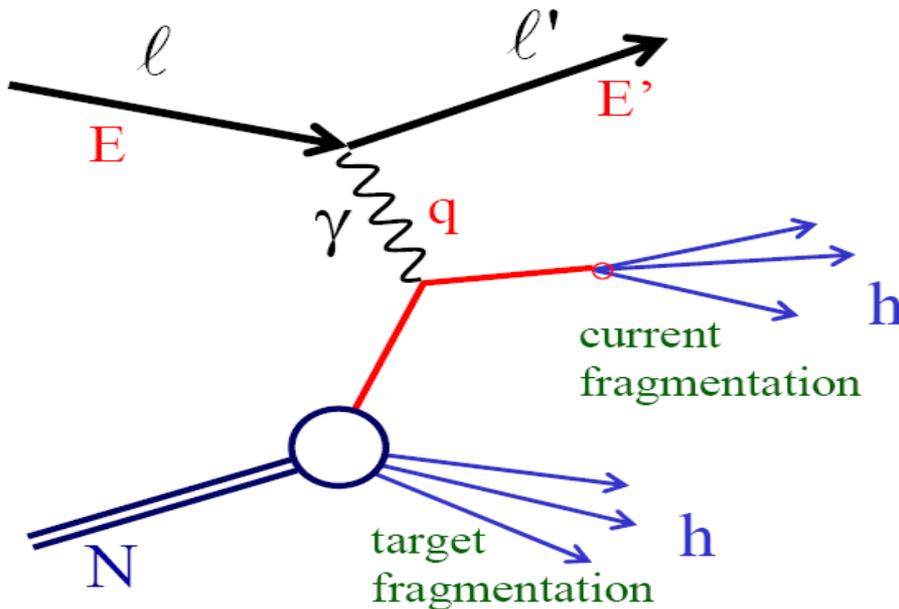
$$q = \ell - \ell'$$

$$Q^2 = -q^2 \quad W^2 = (P + q)^2$$

$$x = \frac{Q^2}{2P \cdot q} \quad \text{Bjorken scaling variable}$$

$$y = \frac{P \cdot q}{P \cdot \ell} =_{LAB} \frac{E - E'}{E}$$

$$z = \frac{P \cdot P_h}{P \cdot q} =_{LAB} \frac{E_h}{E - E'}$$



at least one hadron is detected
in the final state
(information on the **struck quark**)

The SIDIS cross section (one photon exchange approximation)

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dI_{h\perp}^2} =$$

unpolarized {

independent modulations in the azimuthal angle of the spin ϕ_S and of the hadron ϕ_h (around the virtual photon direction) coupled with different Structure Functions related to the:

- nucleon and quark polarization
- quark transverse momentum
- ...

nucleon transv. pol.

$$\begin{aligned}
 & \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\
 & + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \\
 & + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \\
 & + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \\
 & + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
 & + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\
 & + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \\
 & + |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right. \\
 & \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \left. \right\},
 \end{aligned}$$

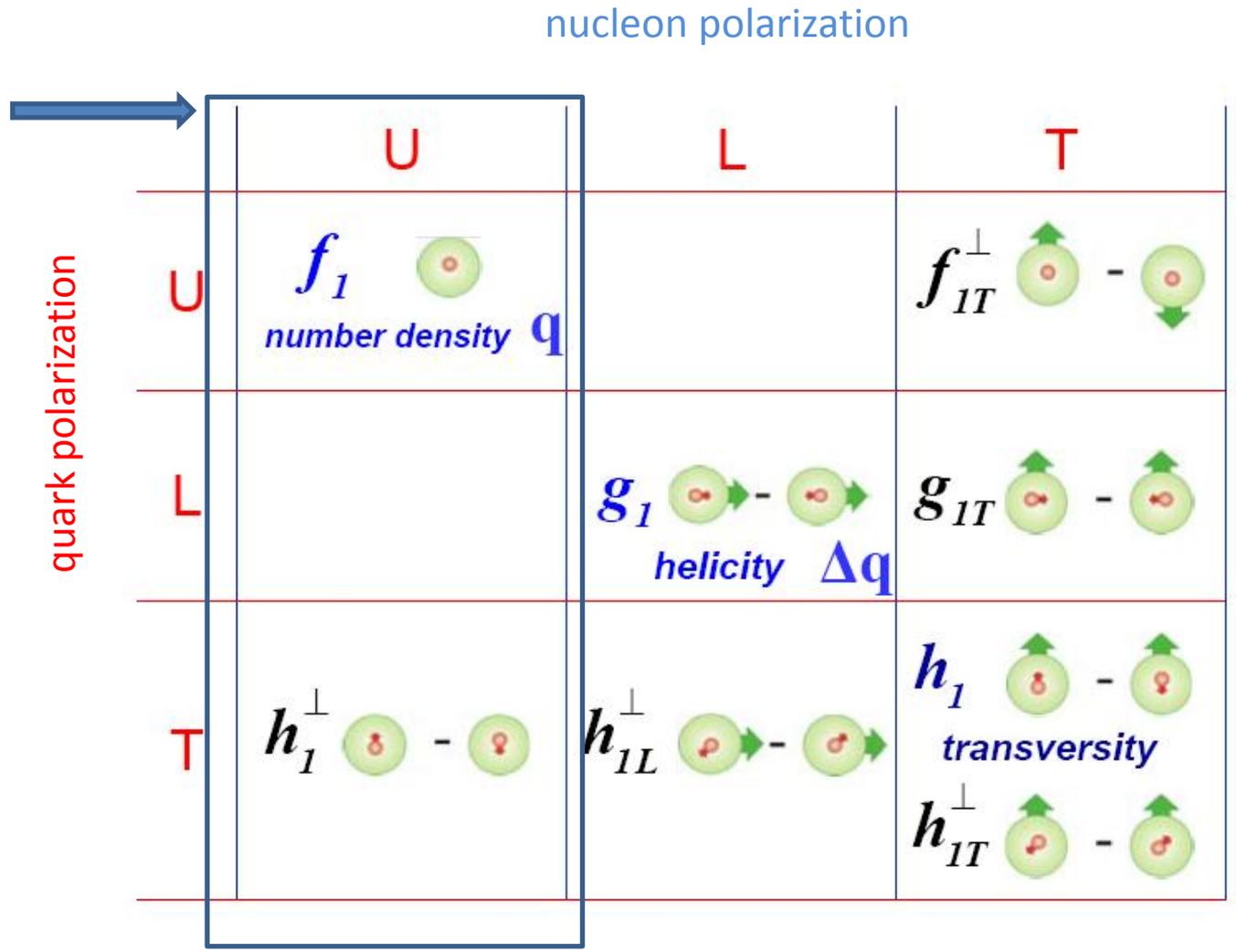
nucleon long. pol.

nucleon polarization

		U	L	T
quark polarization	U	f_1  number density q		f_{1T}^\perp  - 
	L		g_1  -  helicity Δq	g_{1T}  - 
	T	h_1^\perp  - 	h_{1L}^\perp  - 	h_1  -  transversity h_{1T}^\perp  - 

... allow to measure the **Parton Distribution Functions (PDF)** which describe the nucleon at leading order

PDF which appear in an unpolarized nucleon
 (balancing data taken with different target polarization)



3 independent azimuthal modulations *expected in the*
in the azimuthal hadrons distribution
if the Transverse Momentum (TM) of the quarks $\neq 0$

$$N(\phi_h) \propto N_0 \cdot (1 + \varepsilon_1 \boxed{A_{\cos \phi_h}^{UU}} \cos \phi_h + \varepsilon_2 \boxed{A_{\cos 2\phi_h}^{UU}} \cos 2\phi_h + \lambda_l \varepsilon_3 \boxed{A_{\sin \phi_h}^{LU}} \sin \phi_h)$$



$$\sum_q PDF_q(x) \otimes FF_q^h(z)$$

convolution on the TM of the quark

interesting terms: TMD PDFs , FFs

kinematical
factors

$$\varepsilon_1 = \frac{2(2-y)\sqrt{1-y}}{1+(1-y)^2}$$

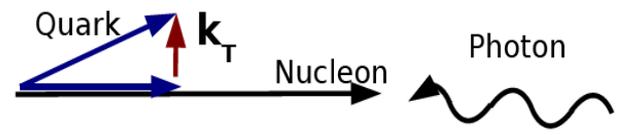
$$\varepsilon_2 = \frac{2(1-y)}{1+(1-y)^2}$$

$$\varepsilon_3 = \frac{2y\sqrt{1-y}}{1+(1-y)^2}$$

λ_l beam
polarization

$$A_{\cos \phi_h}^{UU}$$

mainly **Cahn** effect: **kinematical effect** proportional to the **quark transverse momentum**



$$d\sigma^{lq \rightarrow lq} \propto \hat{s}^2 + \hat{u}^2 \propto \left(1 + \varepsilon_1 \frac{k_{\perp}}{Q} \cos \varphi \right)$$

pQCD negligible for $P_{T}^h < 1 \text{ GeV}/c$

$$A_{\cos 2\phi_h}^{UU}$$

Boer-Mulders (*T-odd* !) function, one of the most famous **TMD PDF**, convoluted with the **Collins FF**



also Cahn effect contribution $\propto \left(\frac{k_{\perp}}{Q} \right)^2$

pQCD negligible for $P_{T}^h < 1 \text{ GeV}/c$

the **Boer-Mulders function** correlates the **quark transverse momentum** and the **quark spin** in an **unpolarized nucleon**

$$A_{\sin \phi_h}^{LU}$$

twist 3 effect due to beam polarization

The **amplitudes of the 3 azimuthal modulations** have been **measured at COMPASS** separately for **positive and negative hadrons**, as functions of the kinematical variables **x, z** and **P_T^h** (transv. mom. of the hadron w.r.t. the virtual photon)

Basic idea of the method

- To measure those amplitudes from experimental data the **apparatus acceptance is needed** (unlike transverse spin asymmetries measurement)

- **for each bin** (k) in **x, z** and **P_T^h**

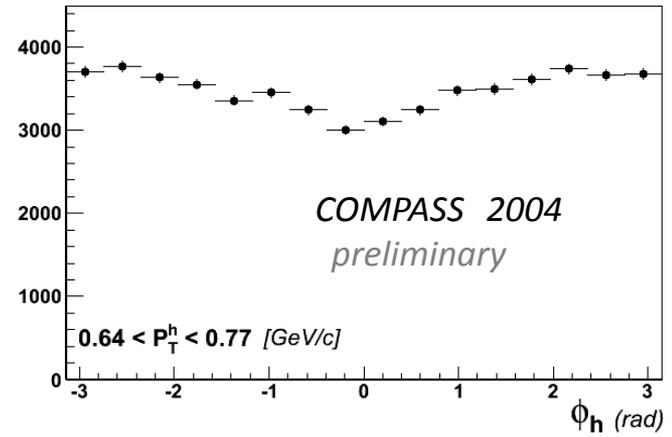
- The **measured azimuthal distributions** need to be **corrected** for the **apparatus acceptance** which may depend on ϕ_h

$$N_k^{corr}(\phi_h) = \frac{N_k(\phi_h)}{Acc_k(\phi_h)}$$

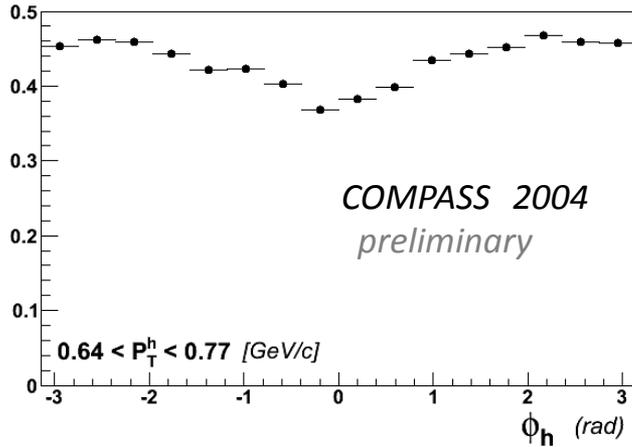
- Azimuthal acceptance calculated from **dedicated MC simulations**

$$Acc_k(\phi_h) = \frac{R_k^{mc}(\phi_h)}{G_k^{mc}(\phi_h)}$$

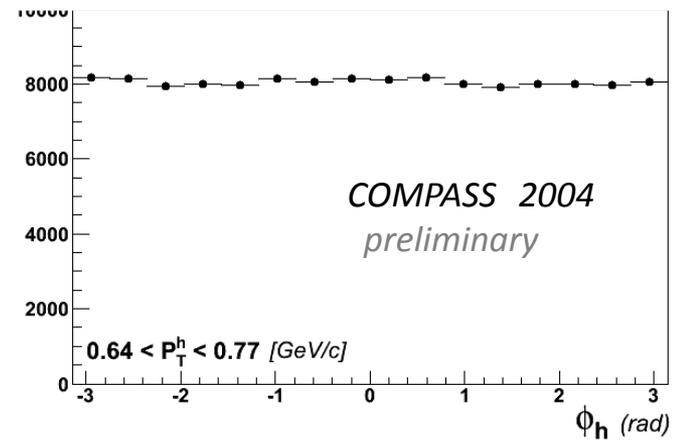
MC rec. azimuthal distribution



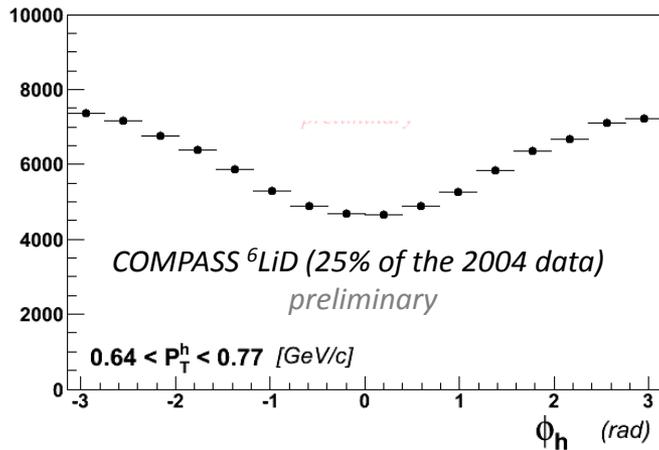
azimuthal acceptance



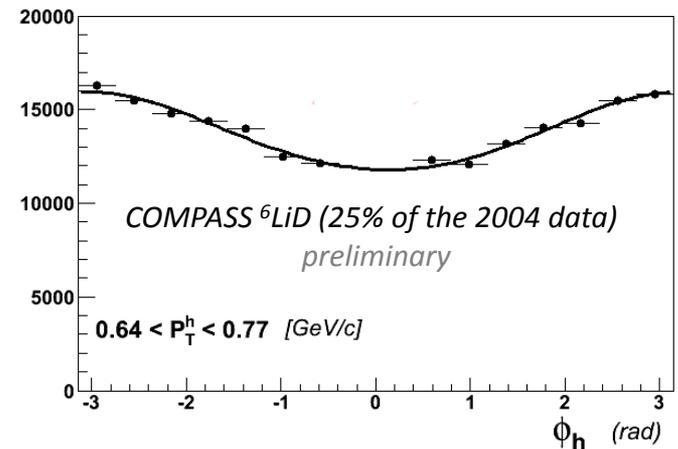
MC gen. azimuthal distribution



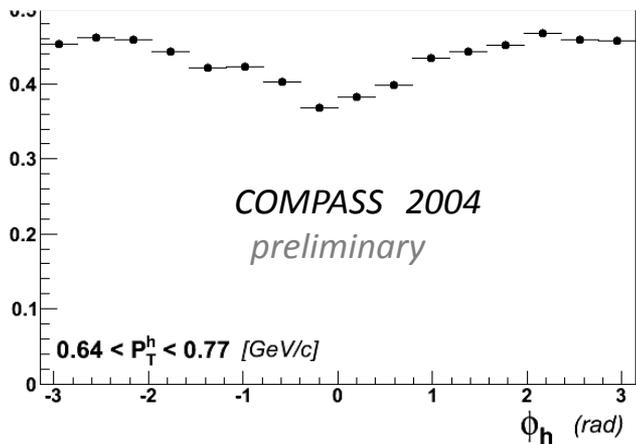
measured azimuthal distribution



measured azimuthal distributions corrected by the acceptance



azimuthal acceptance



amplitudes extracted with a fit:

$$p_0 \cdot (1 + p_1 \cdot \cos \phi_h + p_2 \cdot \cos 2\phi_h + p_3 \cdot \sin \phi_h)$$

ACCEPTANCE STUDIES

Full MC chain to reproduce apparatus acceptance:

- LEPTO (generator, CTEQ5L PDF)
- COMGEANT (dedicated software to simulate the spectrometer)
- CORAL (same software used for the event reconstruction)

SIDIS cuts applied

$$Q^2 > 1 \text{ (GeV/c)}^2$$

$$0.1 < y < 0.9$$

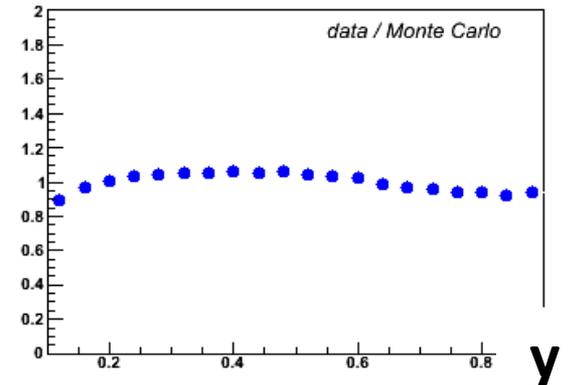
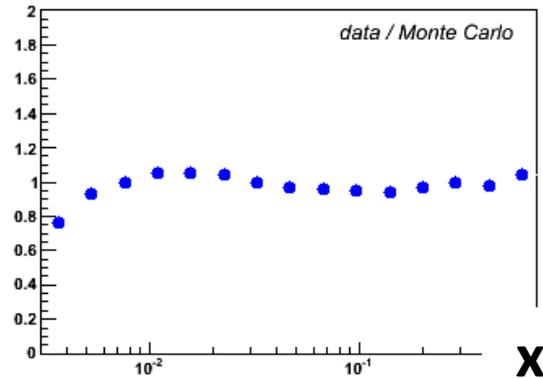
$$W > 5 \text{ GeV/c}^2$$

$$0.2 < z < 0.85$$

$$P_h^T > 0.1 \text{ GeV/c}$$

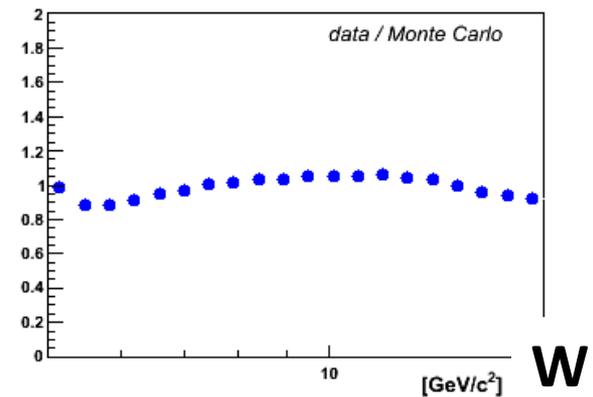
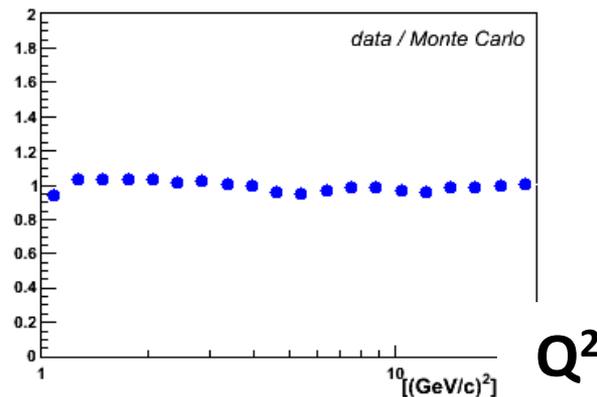
DATA / MONTE CARLO

DIS variables



COMPASS 2004

preliminary



overall good agreement between real data and MC is a check on the reliability of the MC description (no absolute normalization needed)

SIDIS cuts applied

$$Q^2 > 1 \text{ (GeV/c)}^2$$

$$0.1 < y < 0.9$$

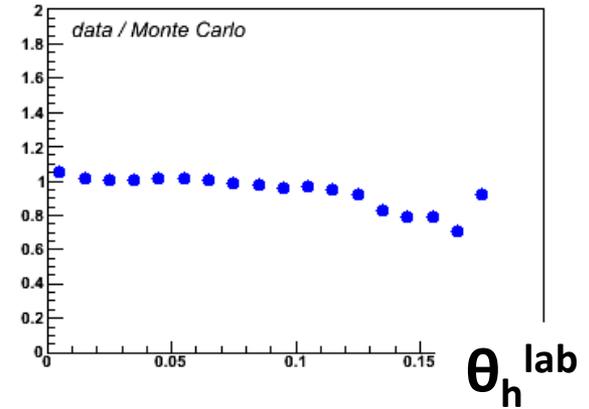
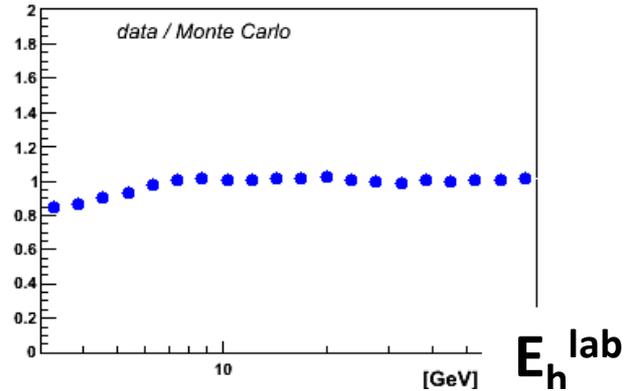
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$$0.2 < z < 0.85$$

$$P_h^T > 0.1 \text{ GeV/c}$$

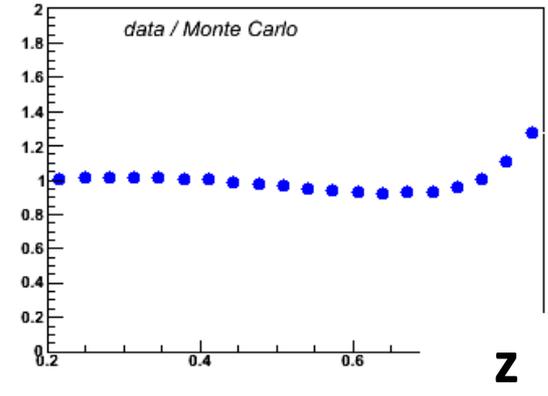
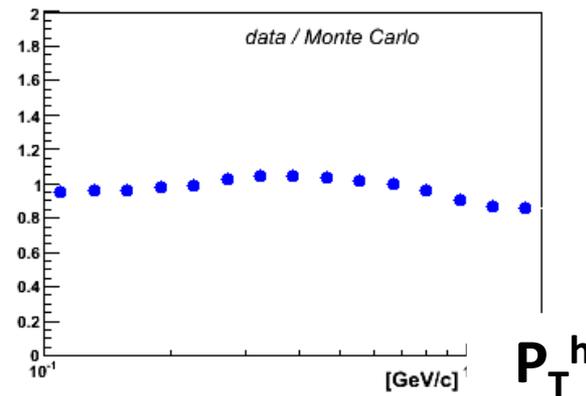
DATA / MONTE CARLO

Hadron variables



COMPASS 2004

preliminary



The apparatus azimuthal acceptance as a function of the event kinematics has been **studied** at length **in order to exclude** the regions giving rise to **large azimuthal modulations in the acceptance** (above 50% for some bin)

→ **new cuts have been tuned:**

$$Q^2 > 1 \text{ (GeV/c)}^2$$

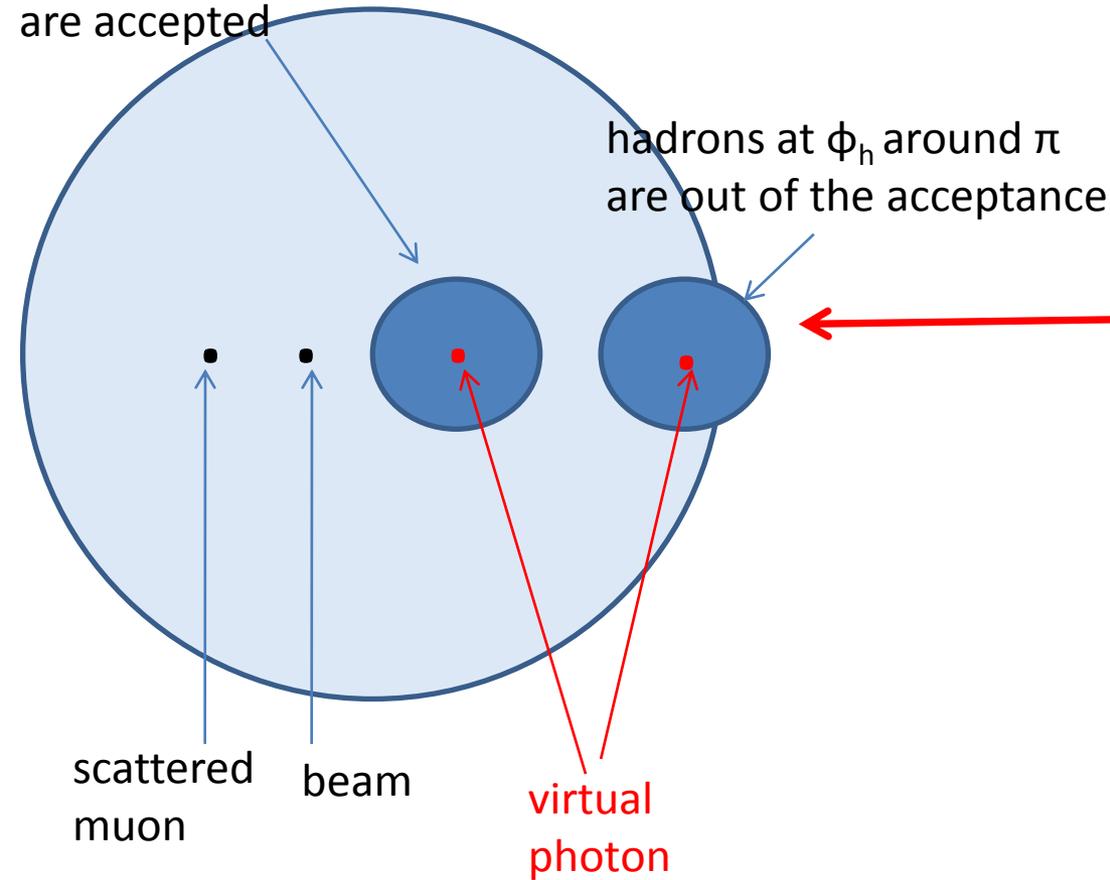
$$-0.9 < y < 0.9$$

$$W > 5 \text{ GeV/c}^2$$

$$0.2 < z < 0.85$$

$$0.1 < P_h^T \text{ GeV/c}$$

most of the produced hadrons
are accepted



$$Q^2 > 1 \text{ (GeV/c)}^2$$

$$\theta_v^{\text{lab}} < 0.06$$

$$0.003 < x < 0.13$$

$$0.2 < y < 0.9$$

$$W > 5 \text{ GeV/c}^2$$

$$0.2 < z < 0.85$$

$$0.1 < P_h^T < 1 \text{ GeV/c}$$

the **large corrections** given by the acceptance **disappeared**

- the dependence of the **acceptance** from all the ϕ_h dependent terms turned out to be **flat** over the kinematical range



checked by extracting the acceptance in one more dimension (x in the example)

$$Acc_k(\phi_h, x) = \frac{R_k^{mc}(\phi_h, x)}{G_k^{mc}(\phi_h, x)}$$



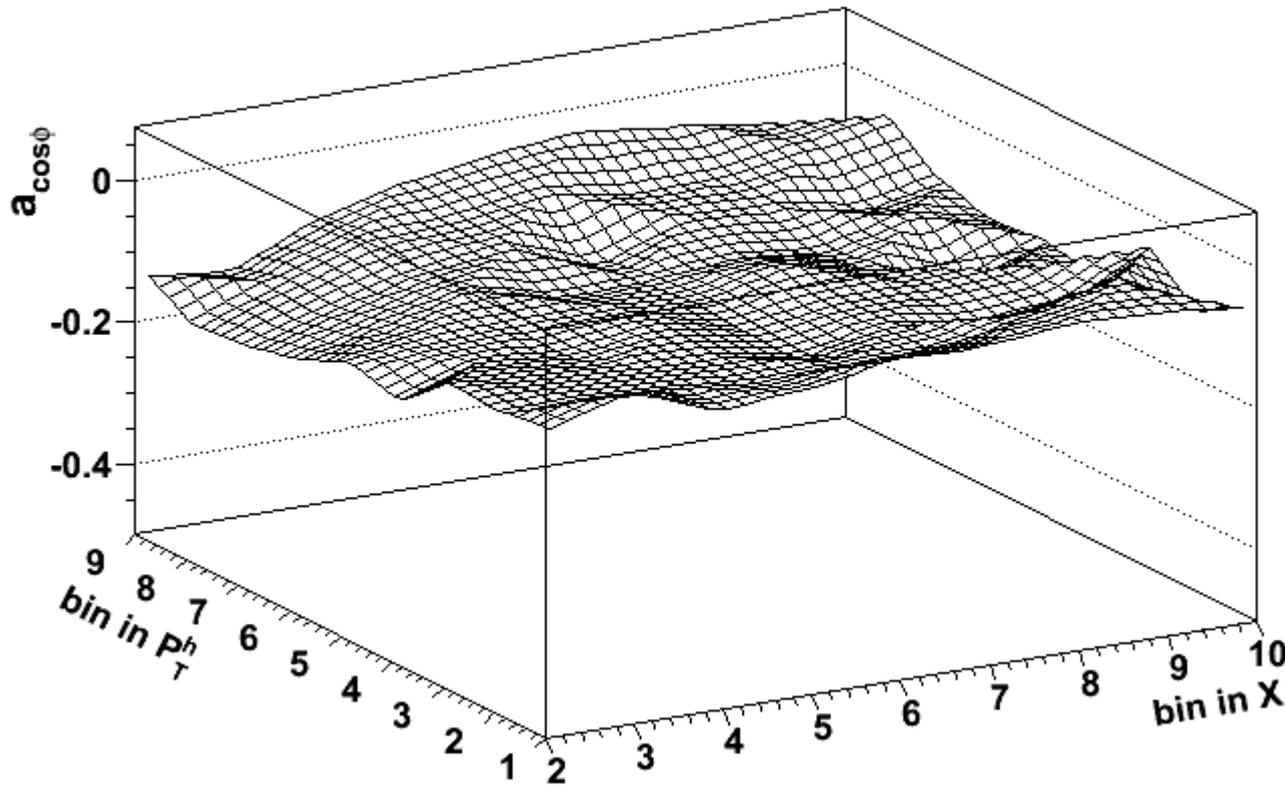
$$Acc_k(\phi_h) = \frac{R_k^{mc}(\phi_h)}{G_k^{mc}(\phi_h)}$$

$$Acc_k(\phi, x) = a_0^k(x) \cdot (1 + a_1^k(x) \cos \phi + a_2^k(x) \cos 2\phi + a_3^k(x) \sin \phi + \dots)$$

- the dependence of the **acceptance** from all the ϕ_h dependent terms turned out to be **flat** over the kinematical range



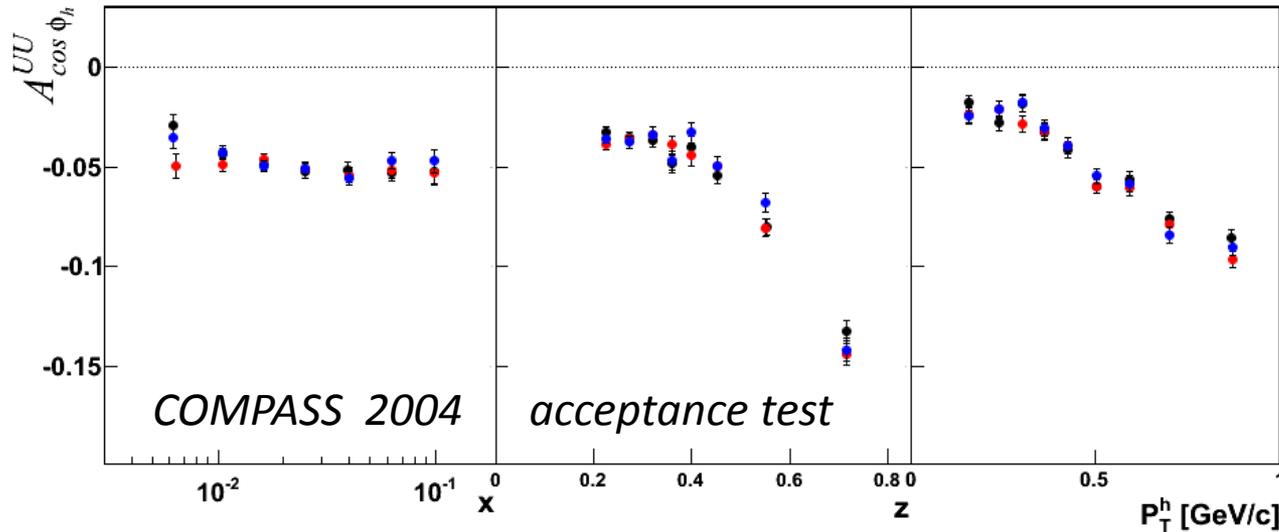
amplitude of the $\cos\phi$ acceptance modulation ($a_{\cos\phi}$)



SYSTEMATIC ERROR

the sources of the systematic error which have been checked and included in the evaluation are the following:

- differences between amplitudes of the azimuthal modulations extracted using 3 different MC samples (obtained by changing PDFs and FF parameters)



the differences are small as expected if

$$Acc_k(\phi_h) = \frac{R_k^{mc}(\phi_h)}{G_k^{mc}(\phi_h)} = \frac{\int G_k^{mc}(\phi_h, x) \cdot Acc_k(\phi_h) dx}{G_k^{mc}(\phi_h)}$$

SYSTEMATIC ERROR

the sources of the systematic error which have been checked and included in the evaluation are the following:

- differences between amplitudes of the azimuthal modulations extracted using the 3 different MC samples
- differences between amplitudes of the azimuthal modulations extracted from transverse and longitudinal data (different apparatus setup → 2 different MC descriptions)
- estimation of possible azimuthal effects from unknown inefficiencies of detectors in low redundancies region

the different contributions calculated in each kinematical bin have been summed in quadrature :

$$\sigma_{sys} = 2 \cdot \sigma_{stat}$$

other studies ...

- a first calculation of the possible effects of radiative corrections have been performed using RADGEN and have been found to be relatively small
- more amplitudes have been added in the fitting function to the ones expected from the unpolarized cross section  compatible with zero

... not included in the systematic errors

$$A_{\cos \phi_h}^{UU}$$

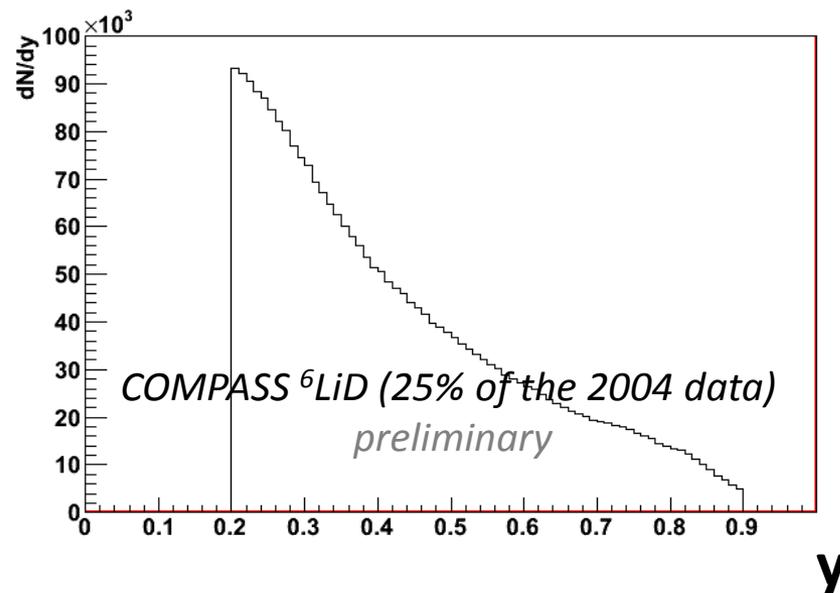
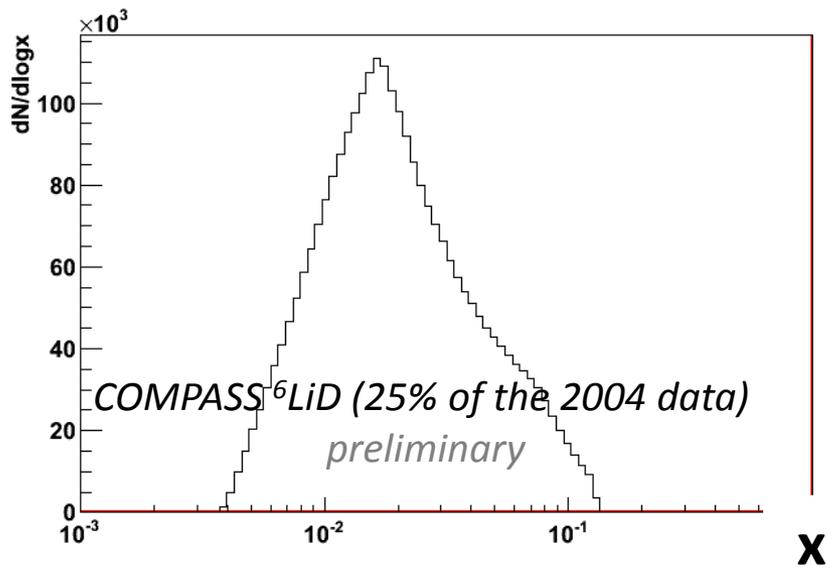
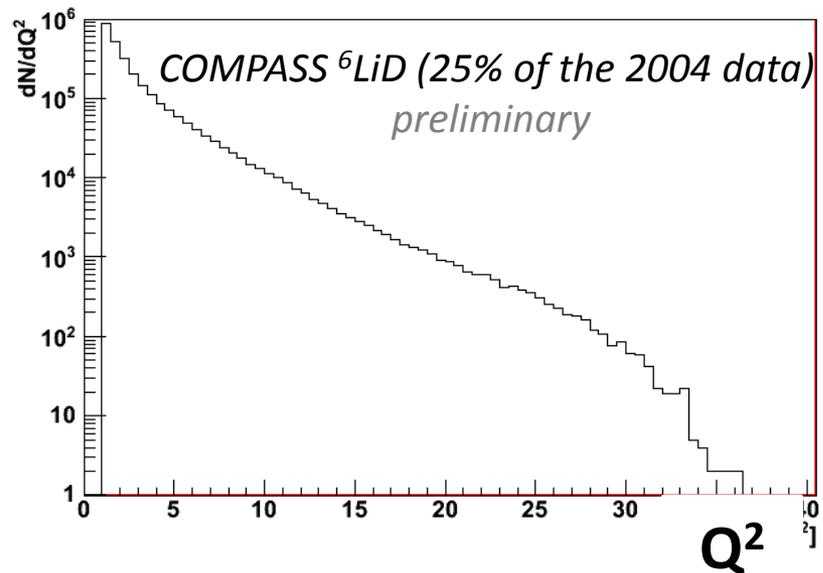
$$A_{\cos 2\phi_h}^{UU}$$

$$A_{\sin \phi_h}^{UU}$$

have been extracted from the 2004 COMPASS data on the ${}^6\text{LiD}$ transversely polarized target

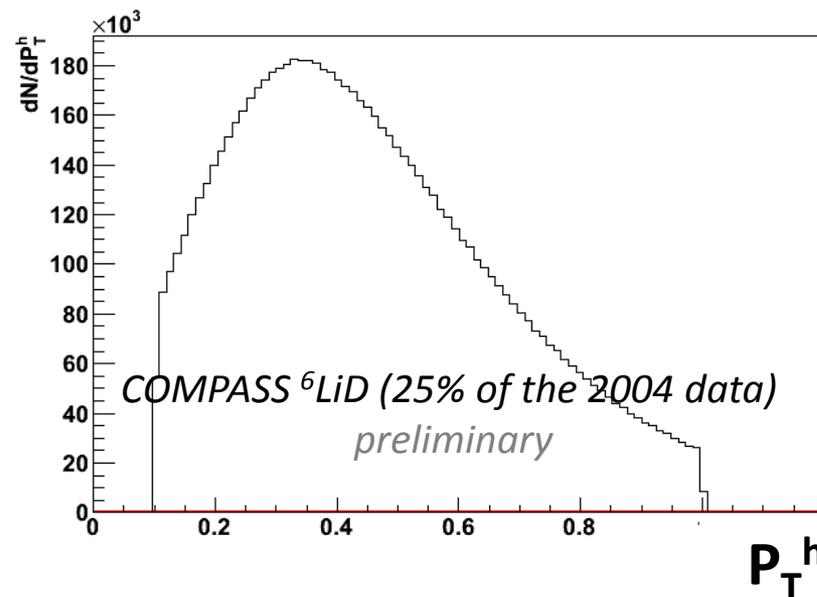
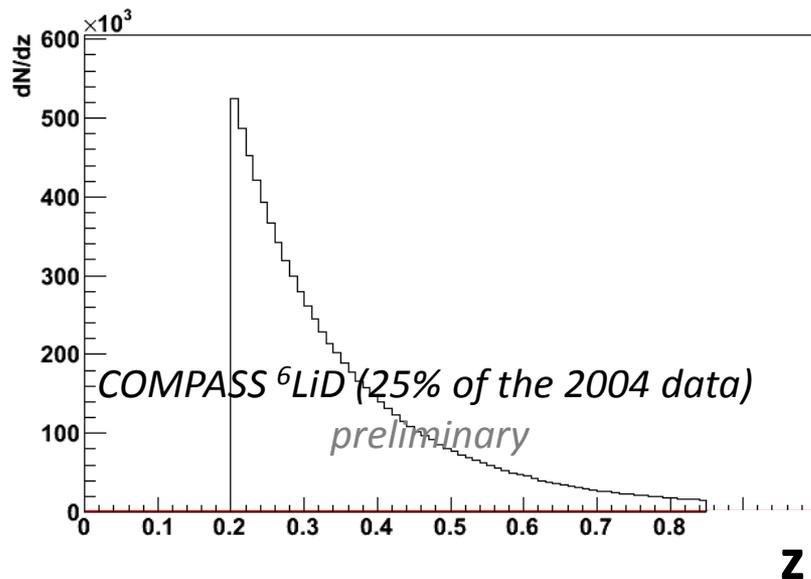
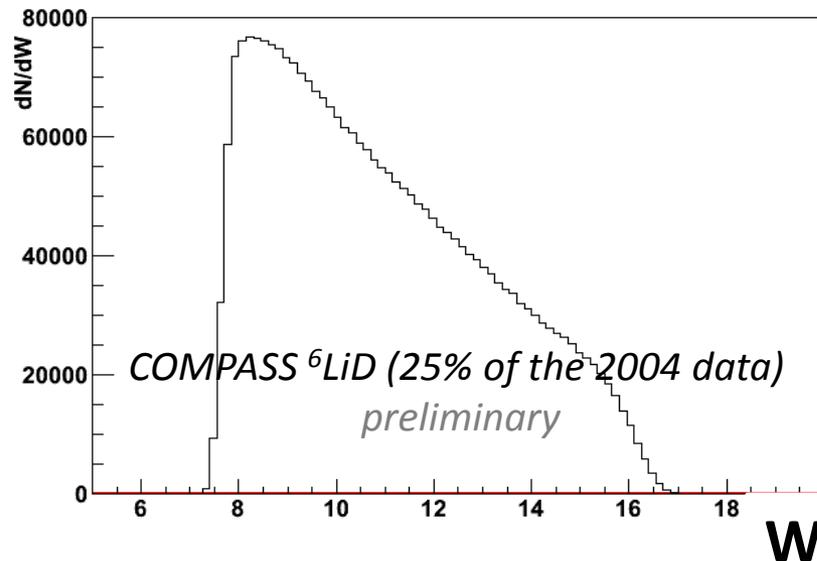
$Q^2 > 1 \text{ (GeV/c)}^2$
 $0.003 < x < 0.13$
 $0.2 < y < 0.9$
 $W > 5 \text{ GeV/c}^2$
 $0.2 < z < 0.85$
 $0.1 < P_h^T < 1 \text{ GeV/c}$

kinematical distributions

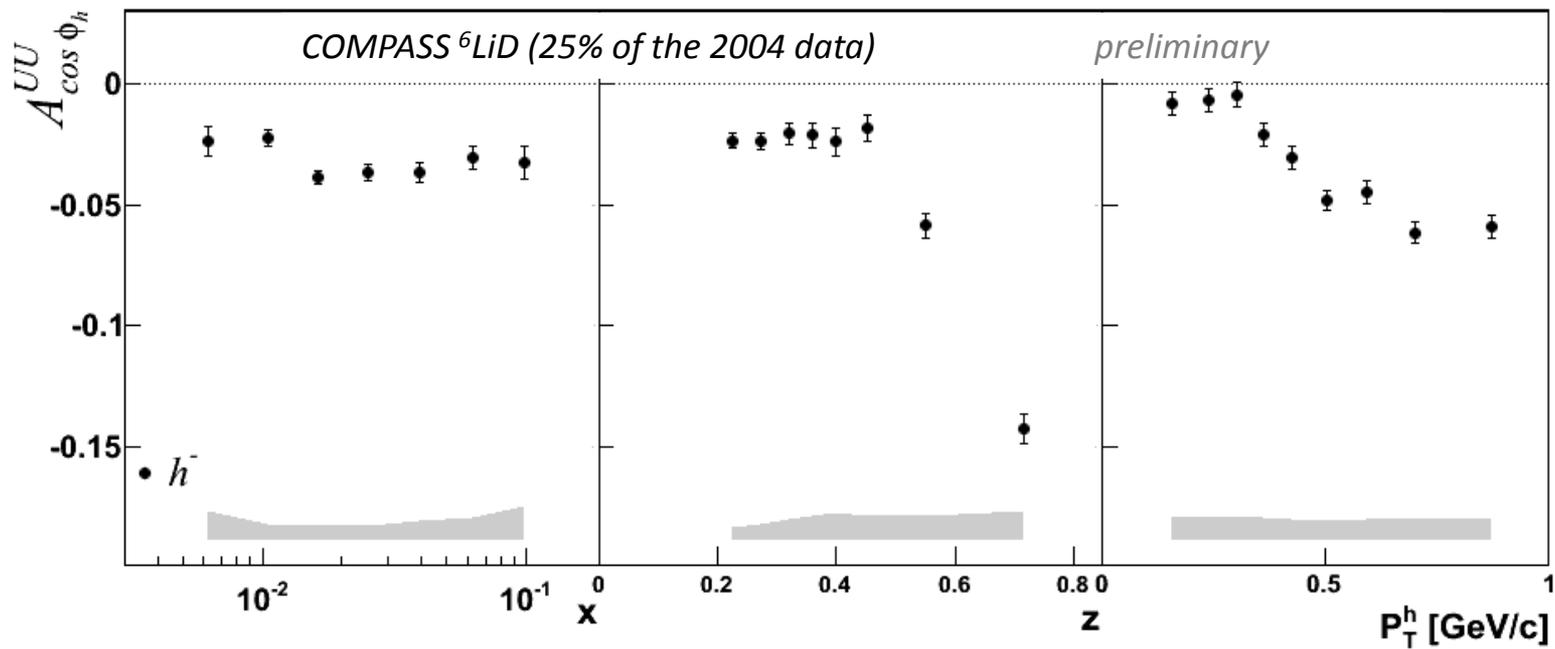
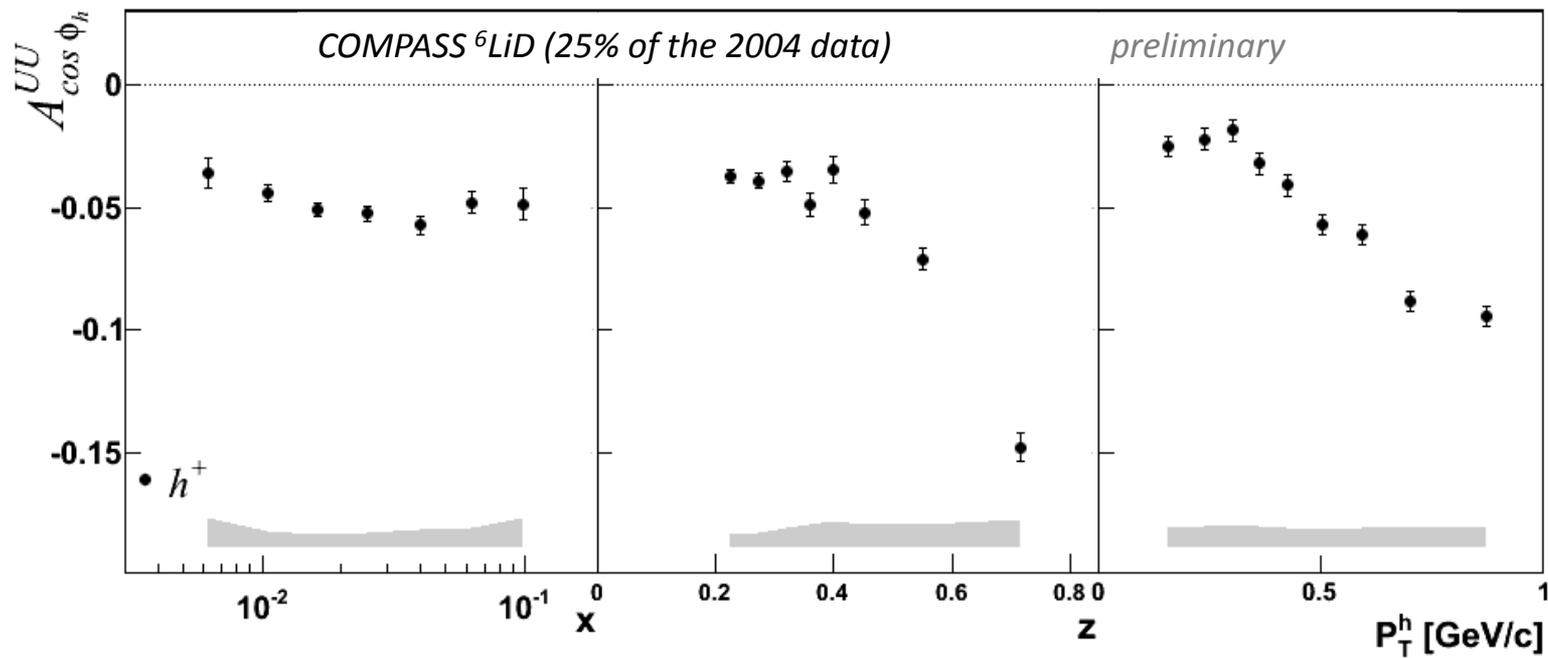


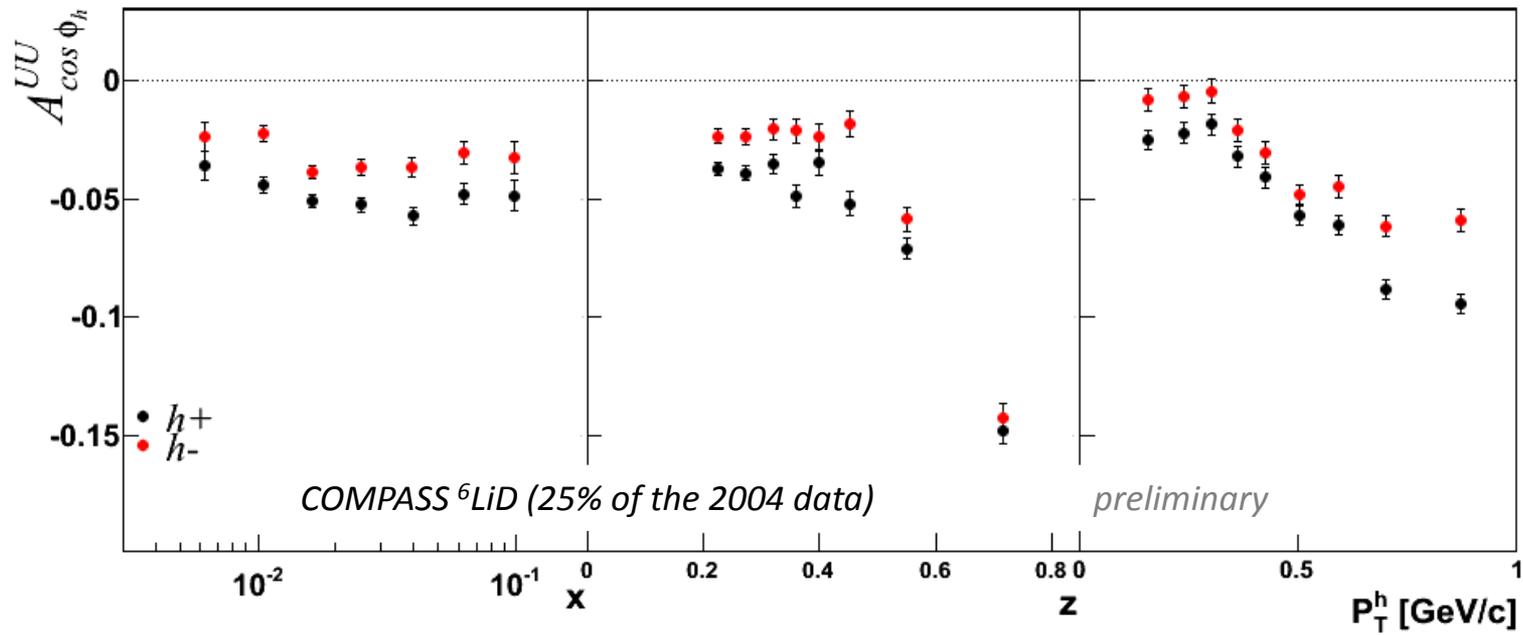
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 $0.1 < P_h^T < 1 \text{ GeV/c}$

kinematical distributions



$$A^{UU} \cos \phi_h$$





different results for positive and negative hadrons

- different transverse momentum for the u and d quarks
- Boer-Mulders TMD PDF (higher twist)

considering only the Cahn effect one expects:

Schweitzer et al.

$$A_{\cos \phi_h}^{UU}(z) = \frac{z \langle k_{\perp}^2 \rangle \sqrt{\pi}}{2 \langle Q \rangle \sqrt{z^2 \langle k_{\perp}^2 \rangle + \langle p_{\perp}^2 \rangle}}$$

k_{\perp} quark transverse momentum w.r.t. the direction of the virtual photon

p_{\perp} hadron transverse momentum w.r.t. the direction of the fragmenting quark

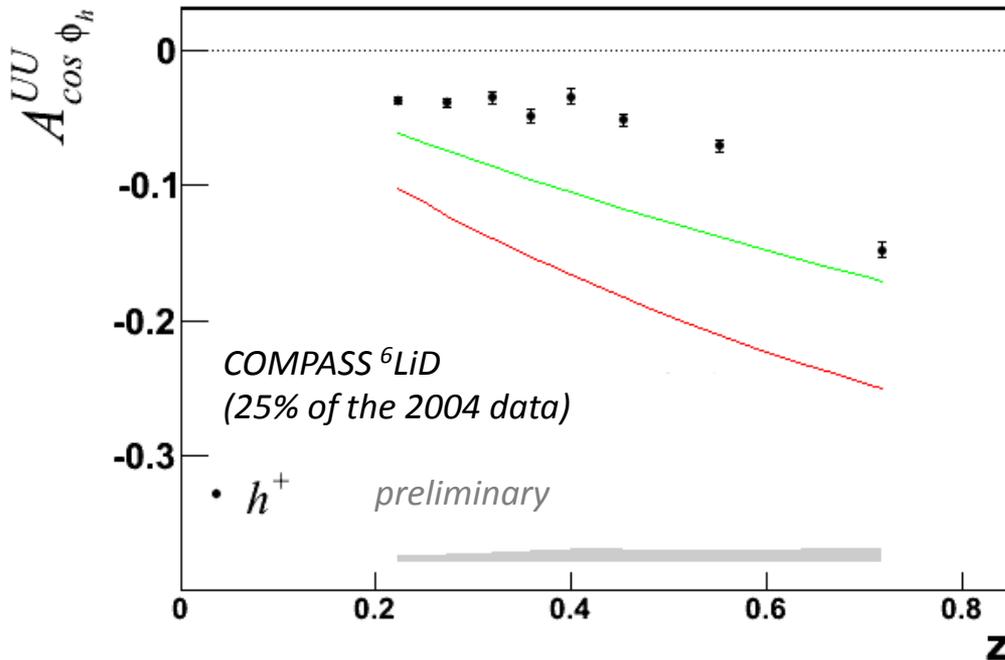


expected curve with:

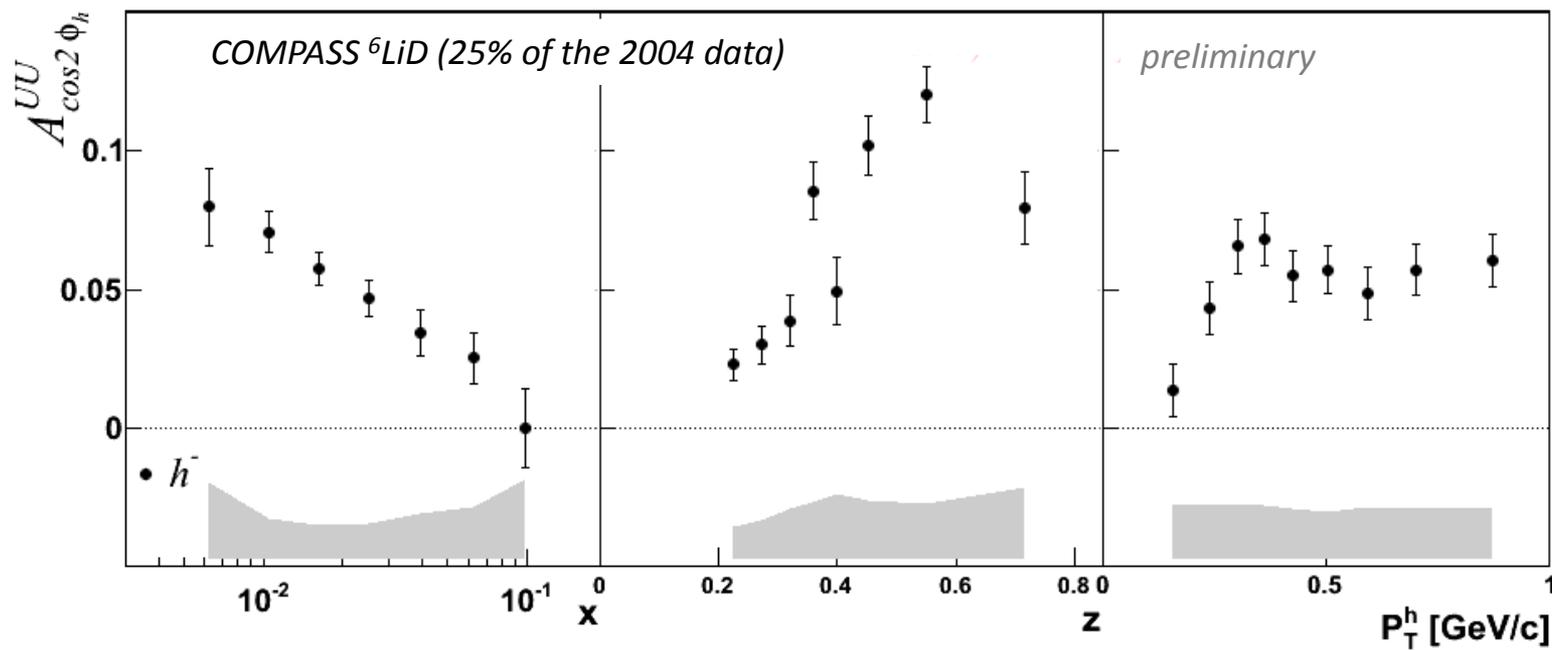
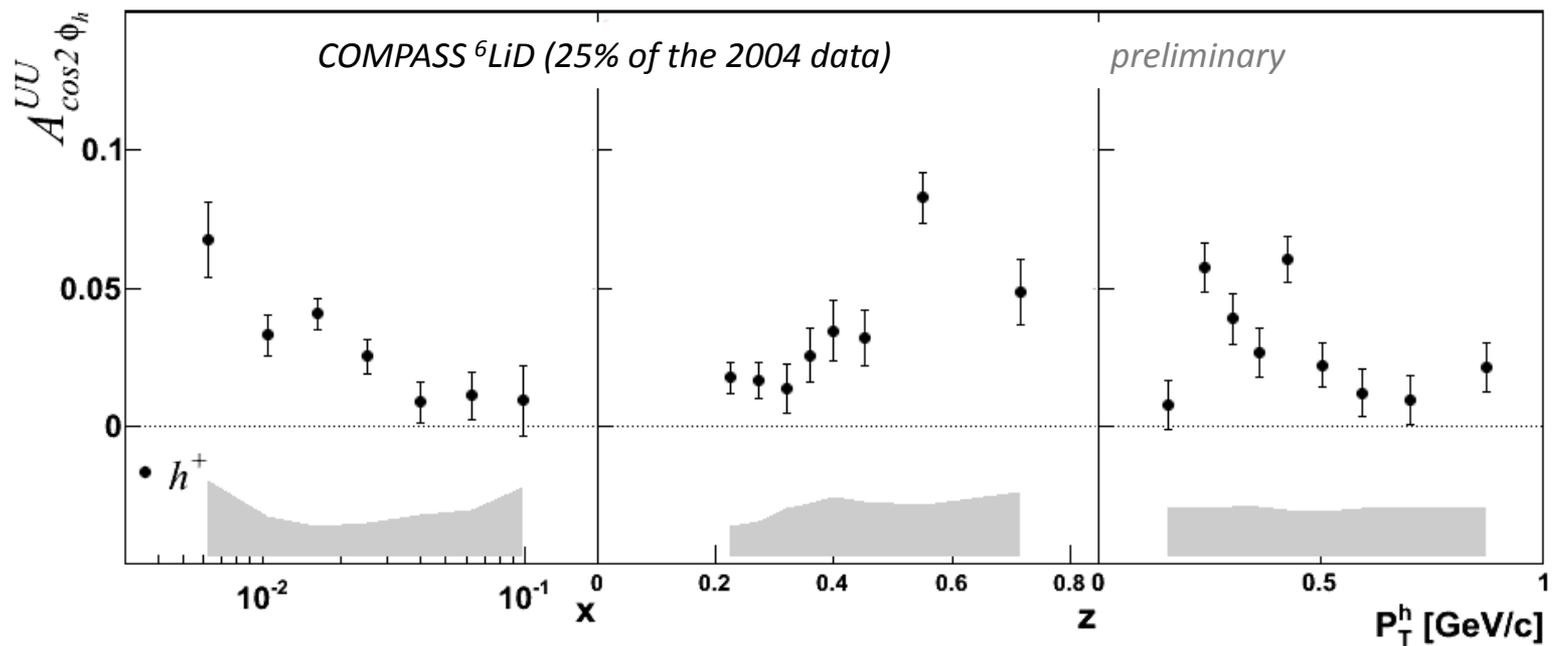
$\langle k_{\perp}^2 \rangle, \langle p_{\perp}^2 \rangle$ (GeV/c)²

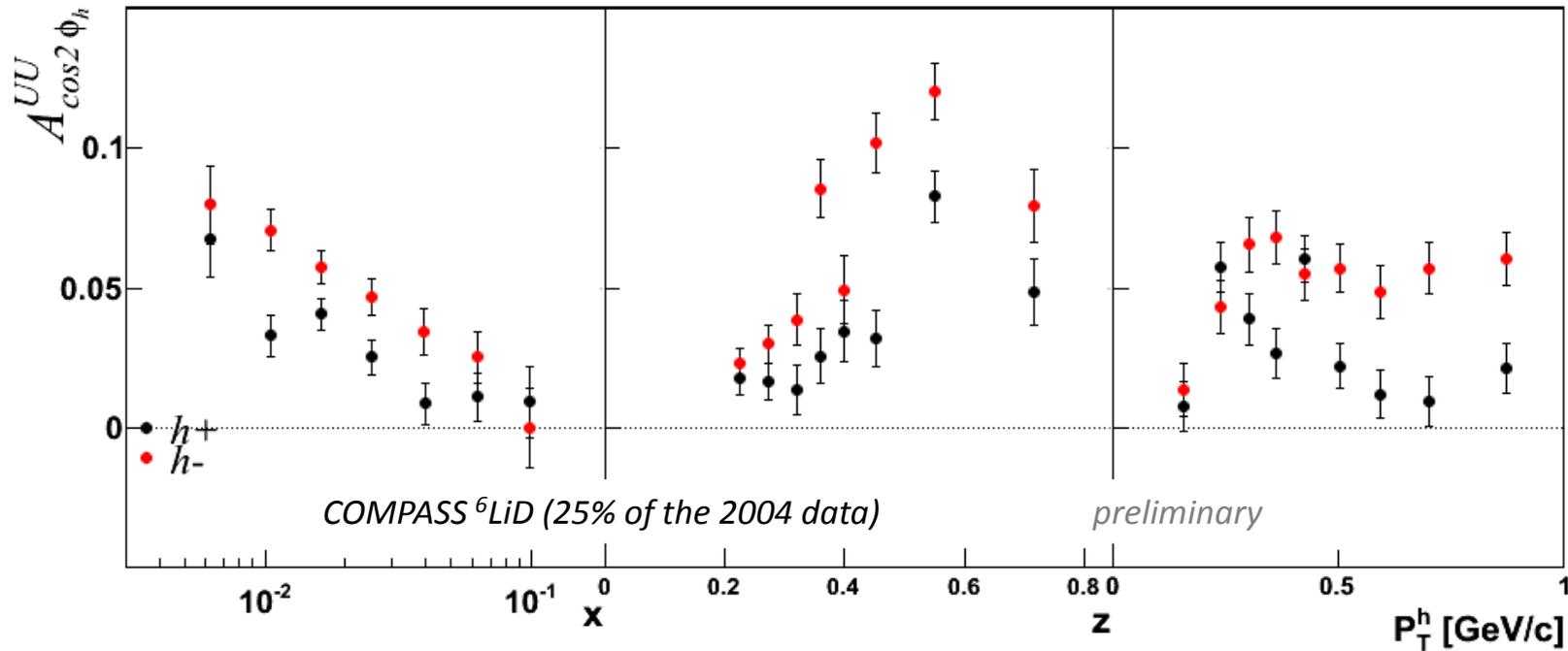
0.25, 0.20 *Anselmino et al. (EMC data)*

0.38, 0.16 *Schweitzer et al. (HERMES data)*



$$A_{\cos 2\phi_h}^{UU}$$





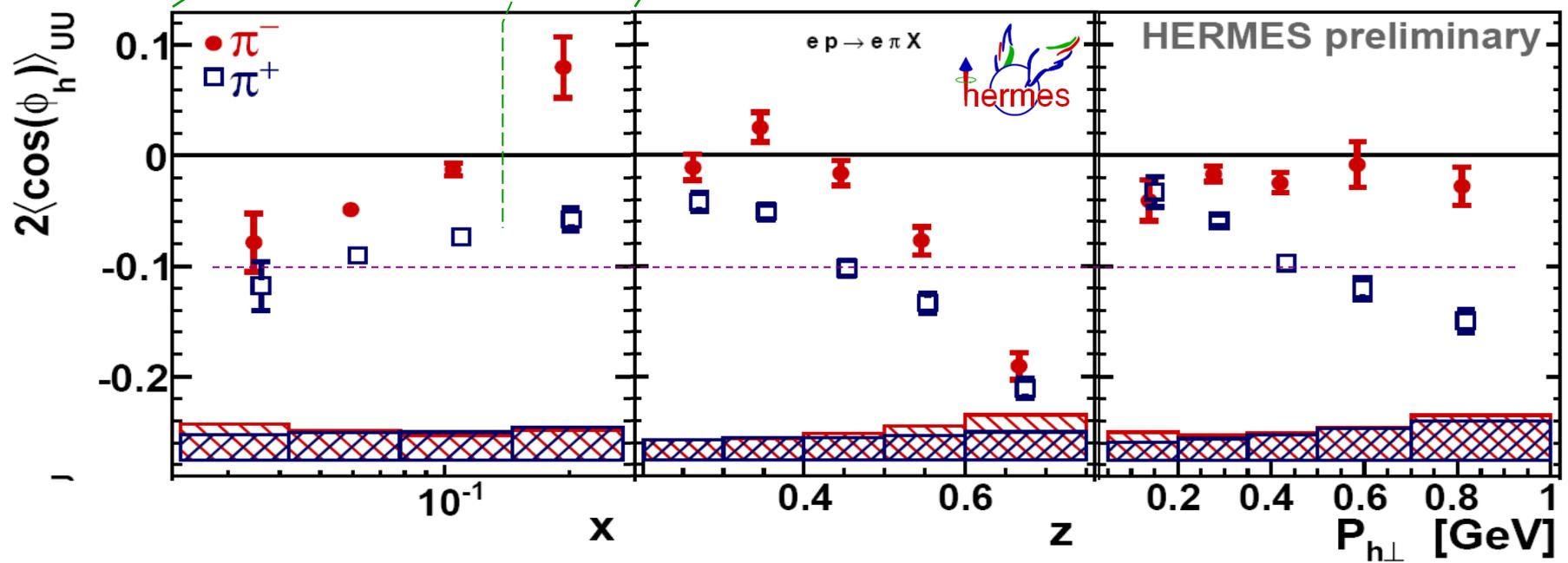
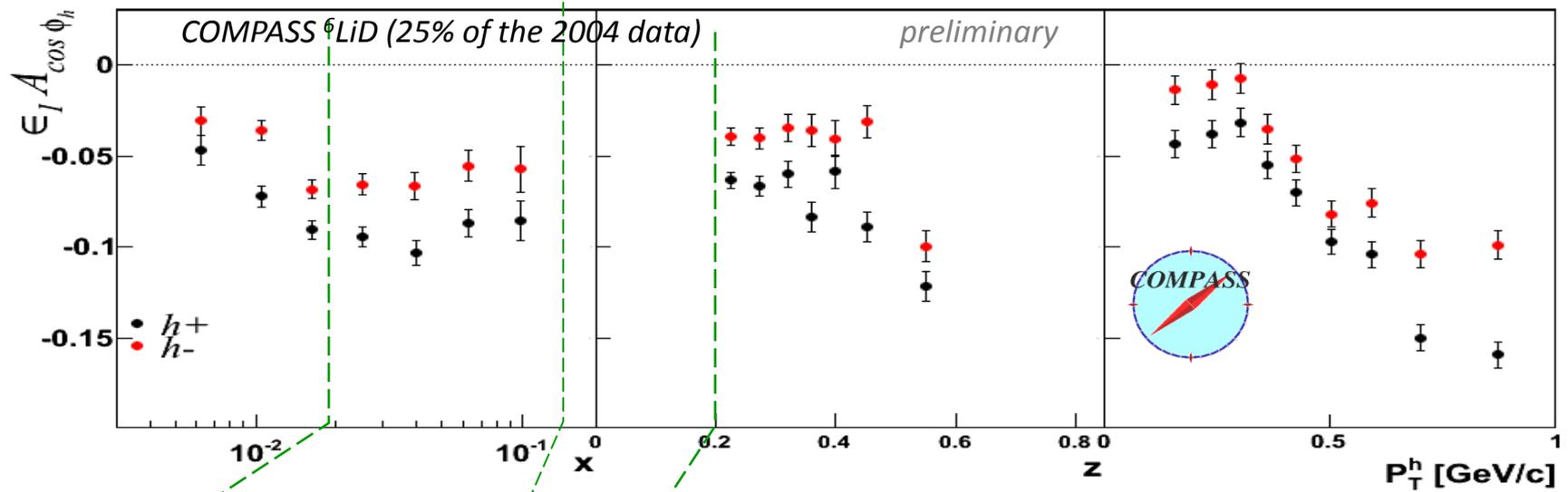
Different results for positive and negative hadrons
→ Boer-Mulders TMD PDF (flavour dependent)

$$\text{Cahn effect} \propto \left(\frac{k_{\perp}}{Q} \right)^2$$

comparison with HERMES results for $\cos\phi_h$ and $\cos 2\phi_h$
(shown at TMD 2010 – ECT ,Trento)

different kinematics (lower Q^2 and W)
→ needs to be taken into account for a comparison

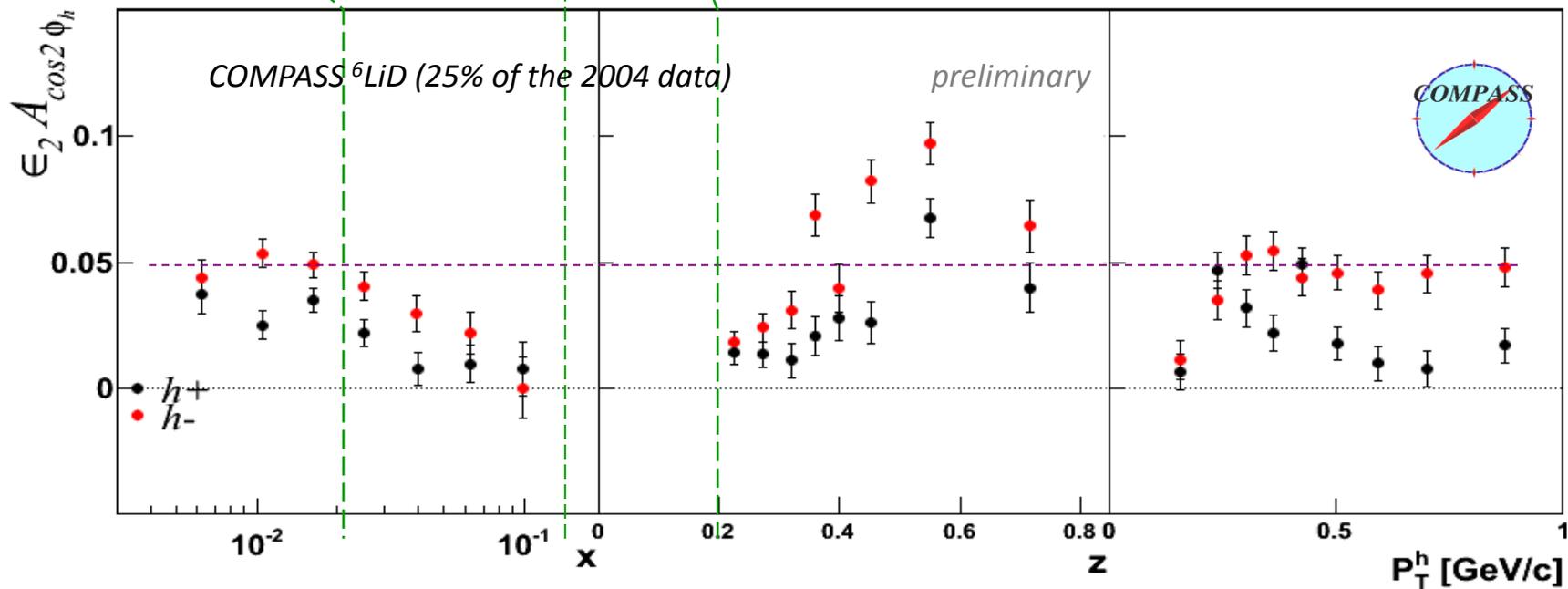
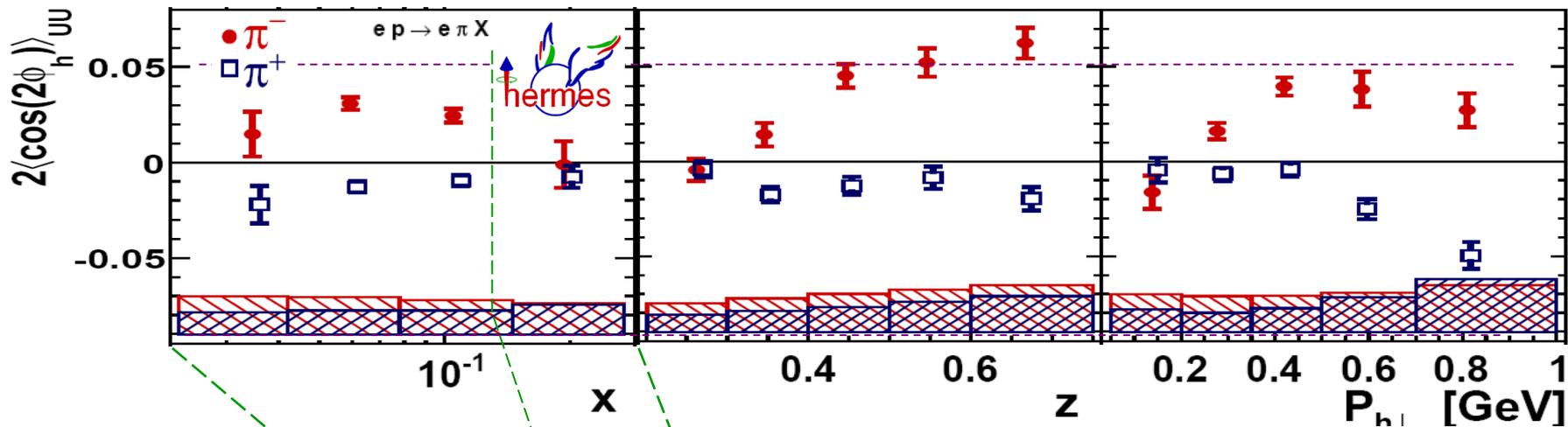
$$\varepsilon_1 A_{\cos \phi_h}^{UU}$$



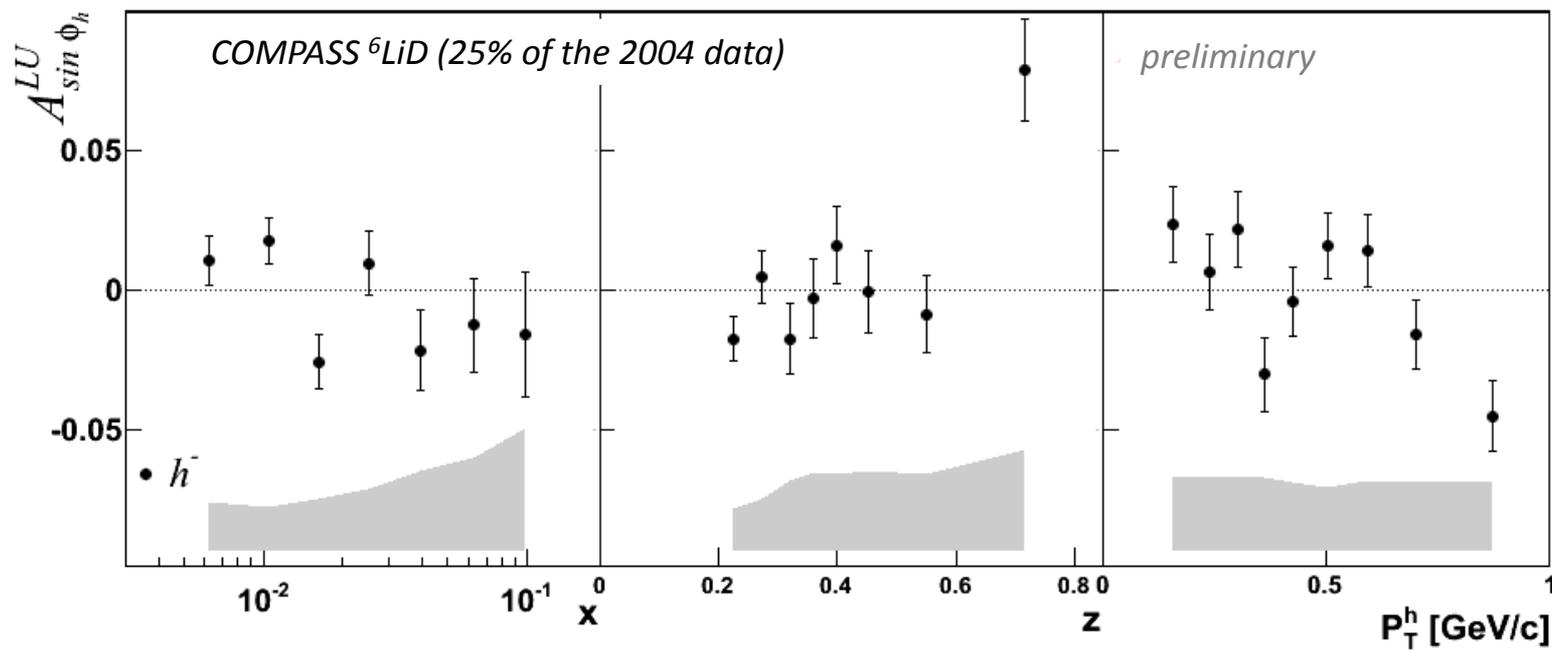
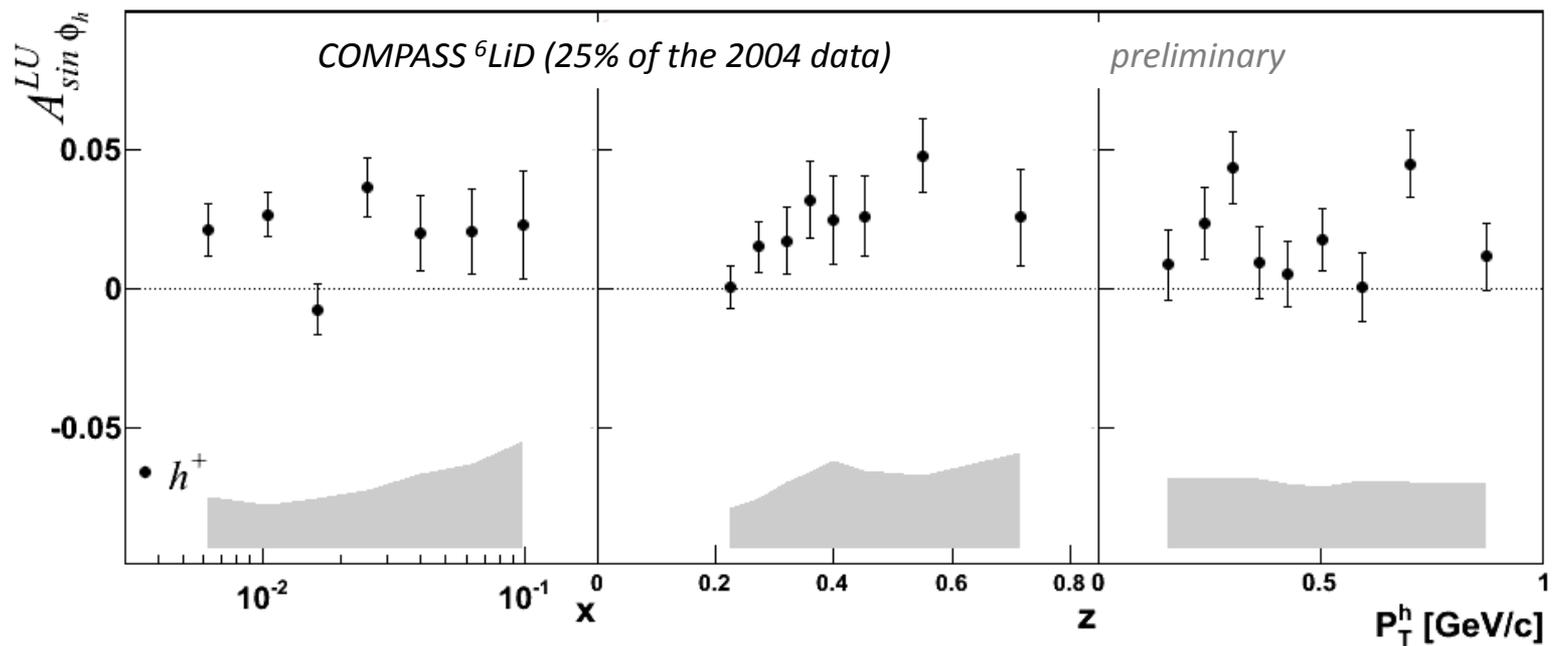
HERMES results shown by A.Rostomyan - TMD2010, Trento

$$\varepsilon_2 A_{\cos 2\phi_h}^{UU}$$

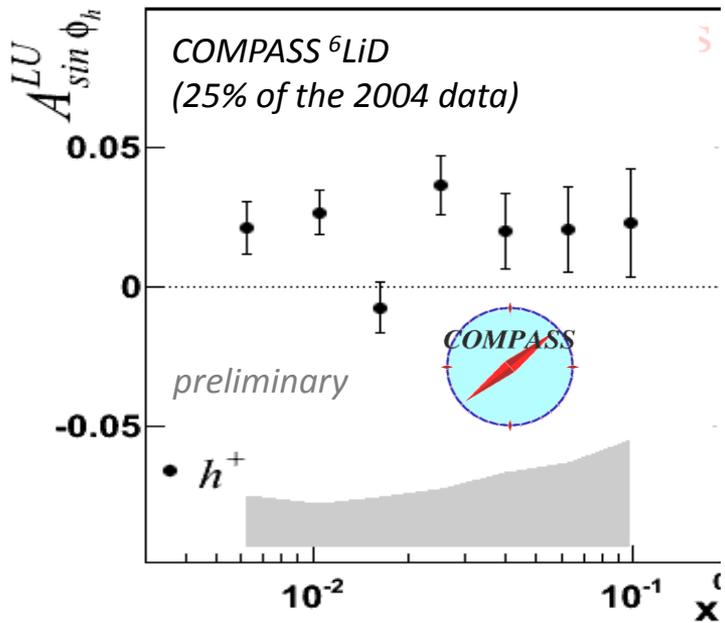
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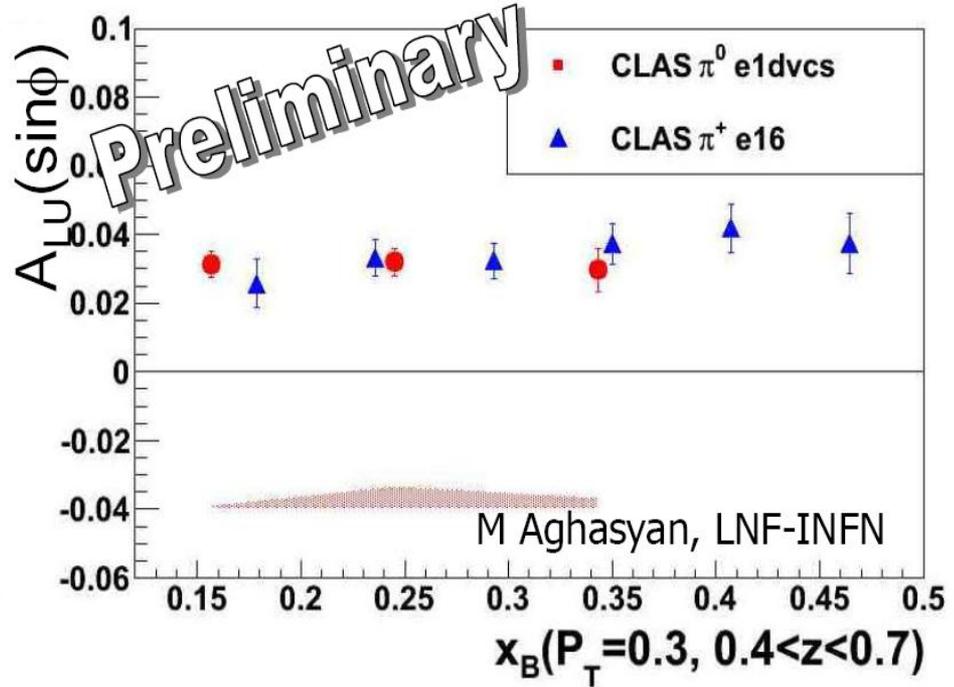
$$A_{\sin \phi_h}^{UU}$$



comparison with JLab results (beam energy 6 GeV)
 (shown at TMD 2010 – ECT ,Trento)



CLAS results shown by H.Avakian - TMD2010, Trento



positive signal for positive hadrons found in both experiments

CONCLUSIONS:

results on the unpolarized SIDIS azimuthal asymmetries from the COMPASS data (6LiD) have been produced as functions of the kinematical variables x , z and P_T^h for positive and negative hadrons

the measured amplitude of the $\cos\phi_h$ modulation shows a clear negative signal (up to 15%)

the measured amplitude of the $\cos 2\phi_h$ modulation shows a positive signal of the order of 5%

the difference between the results for positive and negative hadrons could be interpreted as
flavour dependent intrinsic TM
Boer-Mulders TMD PDF

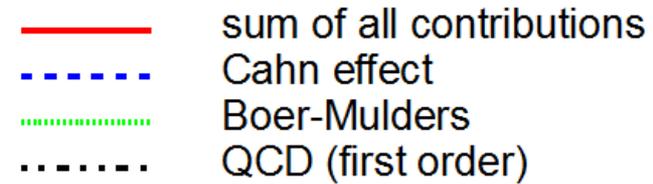
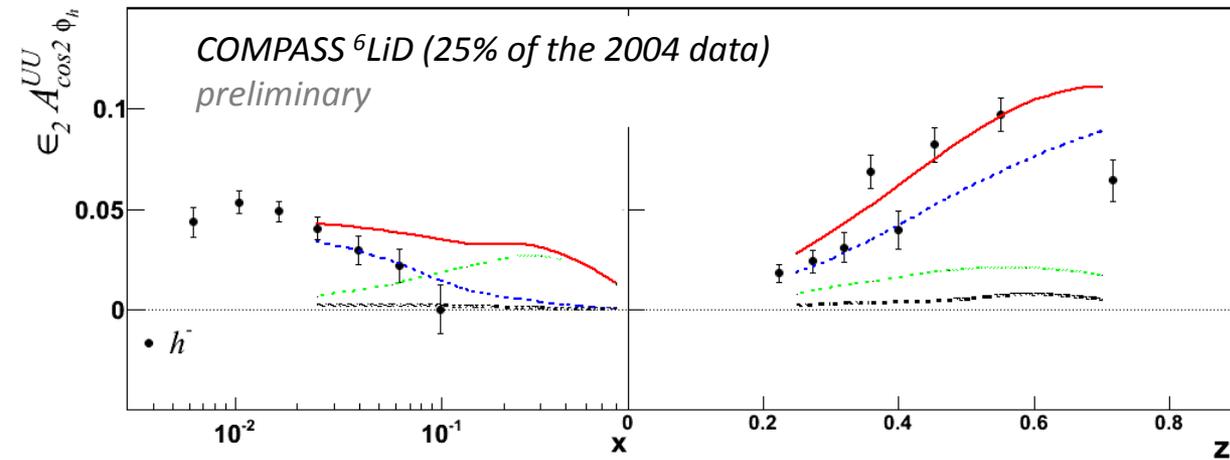
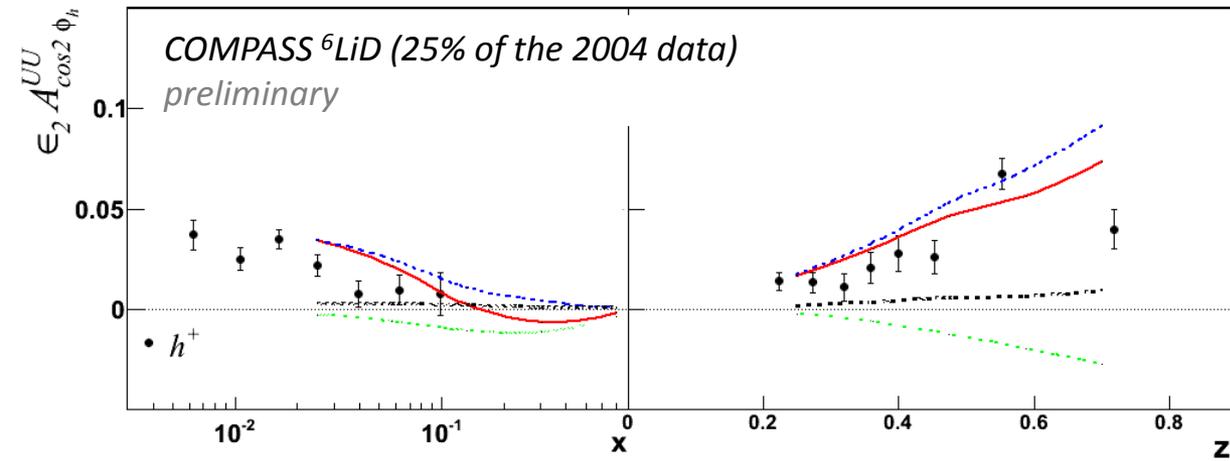
interesting input for theoretical work and global analysis

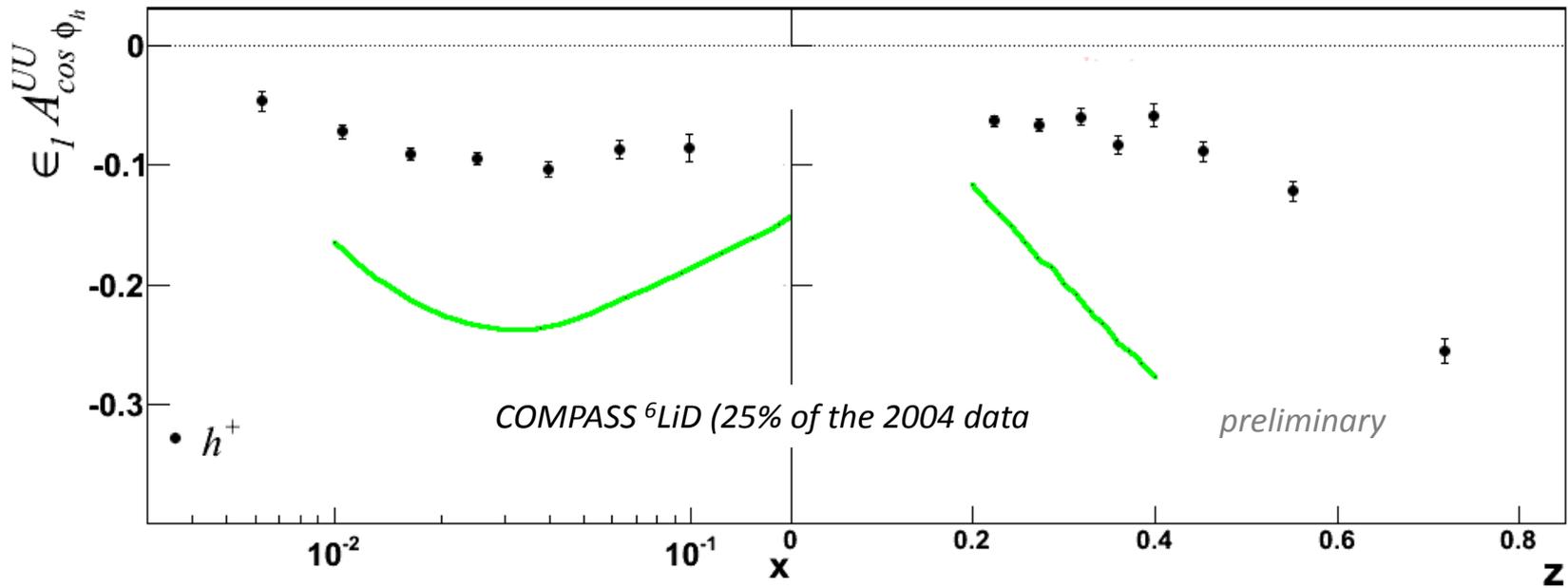
THANKS

backup

Predictions from
Barone, Prokudin, Ma

Phys.Rev.D78:045022





comparison with predictions from Anselmino et al. (**Eur.Phys.J.A31:373-381,2007**)
 only Cahn effect
 no dependence on the quark flavour

based on EMC, E665 data

Tuning of the LEPTO parameters

transverse momentum
of the hadron in the
fragmentation

PARJ(21): 0.34 GeV/c

PARJ(23): 0.04

PARJ(24): 2.8

PARJ(41): 0.025

PARJ(42): 0.075 GeV⁻²

LUND FF parameters