



Transverse Spin and Momentum effects in the COMPASS Experiment



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on behalf of the COMPASS Collaboration



International Committee for Spin Physics Symposia

XXIII WORKSHOP ON HIGH ENERGY PHYSICS

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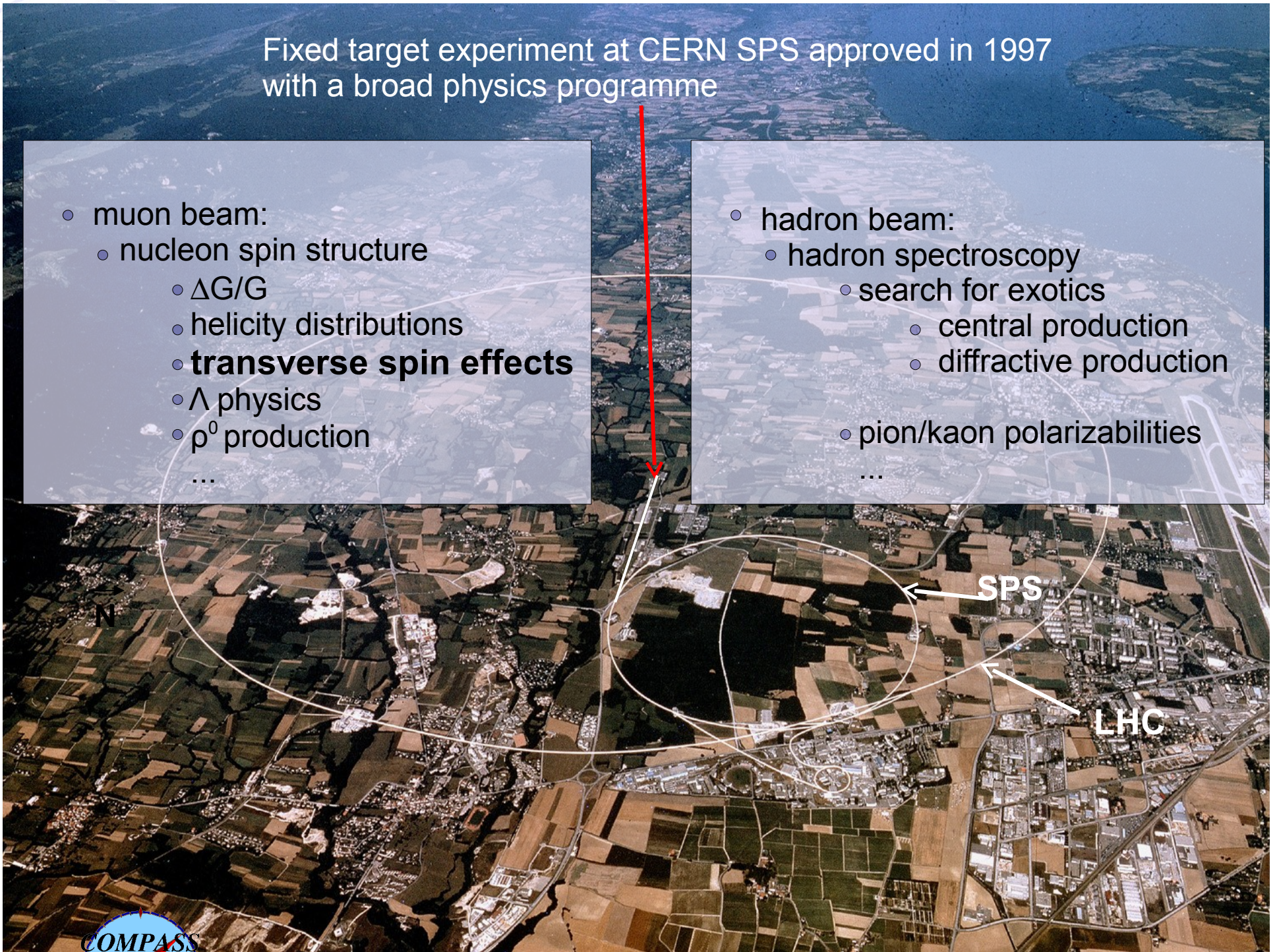
OUTLINE

- **The COMPASS experiment**
- **SIDIS**
- **Hadron azimuthal asymmetries**
 - unpolarized target
 - transversely polarized target
- **Two hadrons azimuthal asymmetries**
- **Conclusions and outlook**

Fixed target experiment at CERN SPS approved in 1997
with a broad physics programme

- muon beam:
 - nucleon spin structure
 - $\Delta G/G$
 - helicity distributions
 - **transverse spin effects**
 - Λ physics
 - ρ^0 production
 - ...

- hadron beam:
 - hadron spectroscopy
 - search for exotics
 - central production
 - diffractive production
 - pion/kaon polarizabilities
 - ...



Common Muon and Proton Apparatus for Structure and Spectroscopy



Data taking since 2002

muon beam	deuteron (${}^6\text{LiD}$) polarized target	2002	L/T target polarization
		2003	
		2004	
	proton (NH_3) polarized target	2006	L target polarization
		2007	L/T target polarization

hadron beam	LH target	2008
		2009

longitudinally polarised muon beam
beam intensity: $2 \cdot 10^8 \mu_+/\text{spill}$ (4.8s/16.2s)
beam momentum: 160 GeV/c
luminosity: $\sim 5 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$





The COMPASS Spectrometer – muon beam

built to cover a large kinematical range

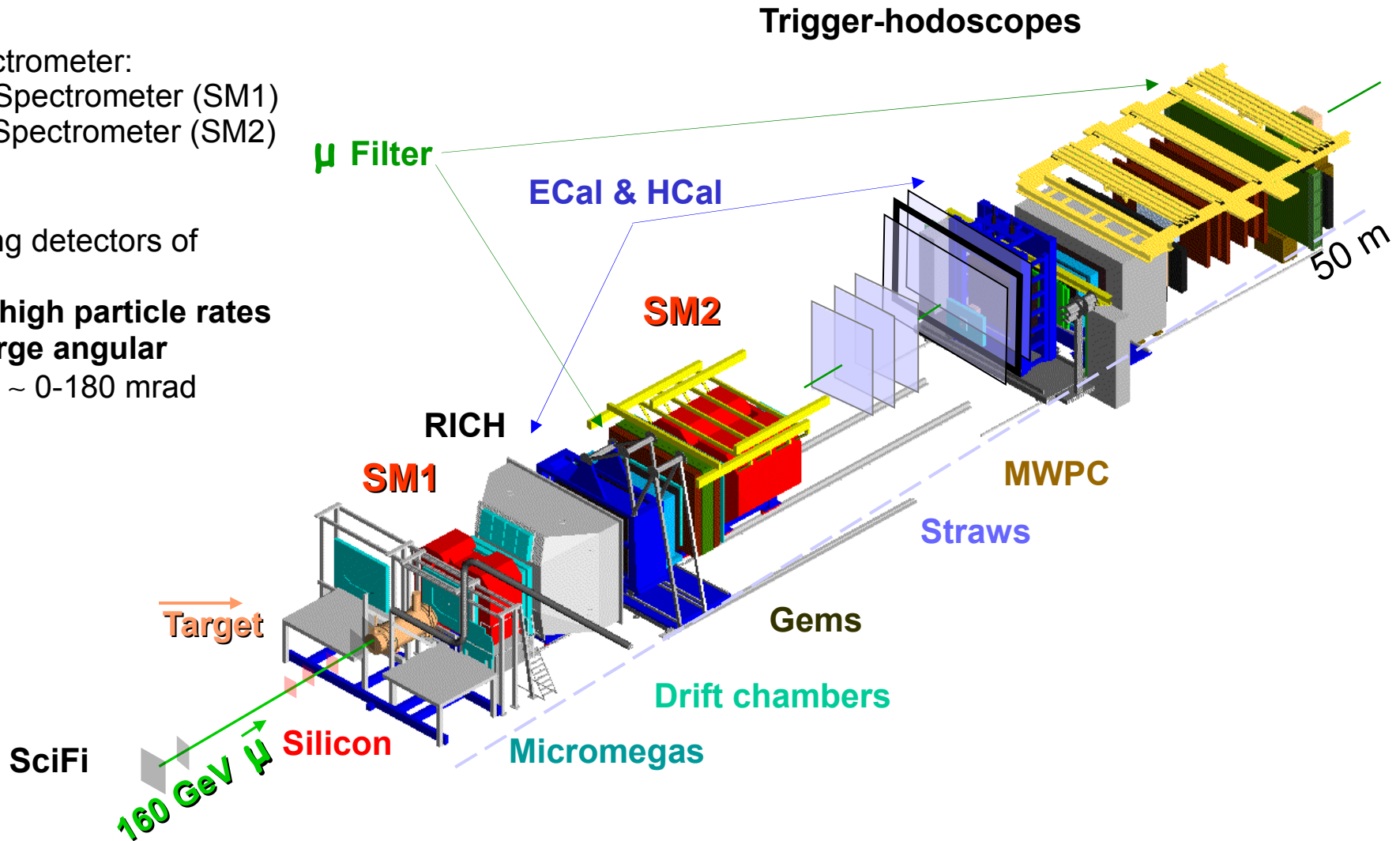
- two stage spectrometer:
- Large Angle Spectrometer (SM1)
 - Small Angle Spectrometer (SM2)

several tracking detectors of different type

- to cope with **high particle rates**
- to cover a **large angular acceptance** ~ 0-180 mrad

particle ID:

- calorimeters
- muon walls
- RICH



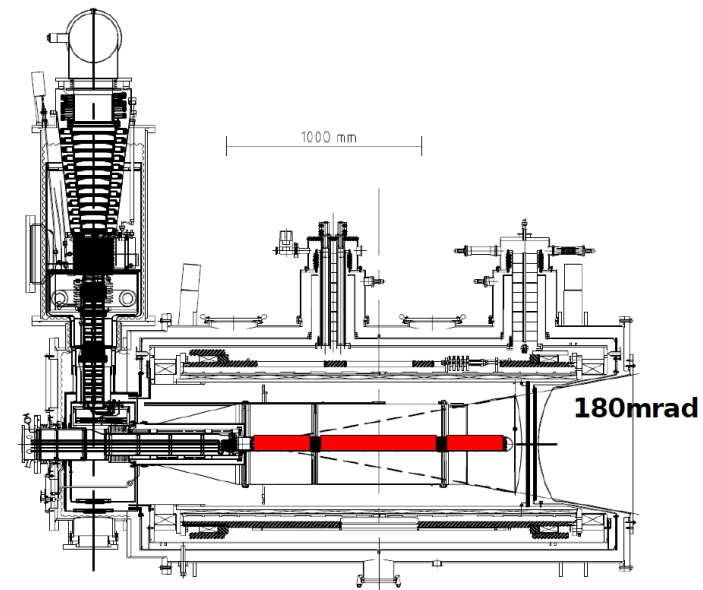
Common Muon and Proton Apparatus for Structure and Spectroscopy



The polarized target

Solid state target operating in frozen spin mode 120 cm long

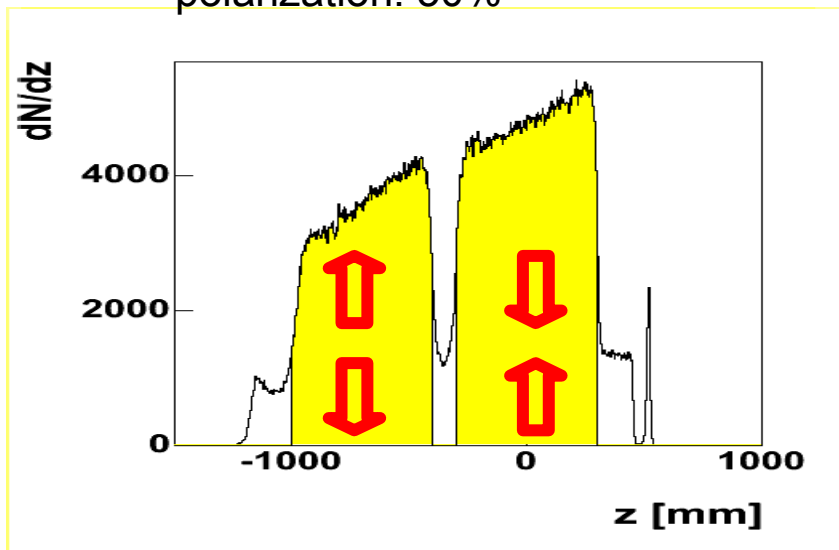
Composed by 2 or 3 cells with **opposite polarization** reversed every week during transverse data taking



2002-4 ${}^6\text{LiD}$

dilution factor: 0.38

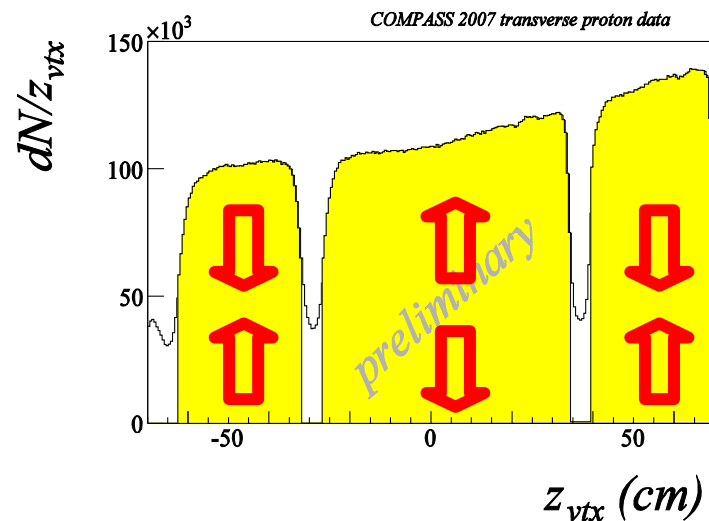
polarization: 50%



2007 NH_3

dilution factor: 0.14

polarization: 90%





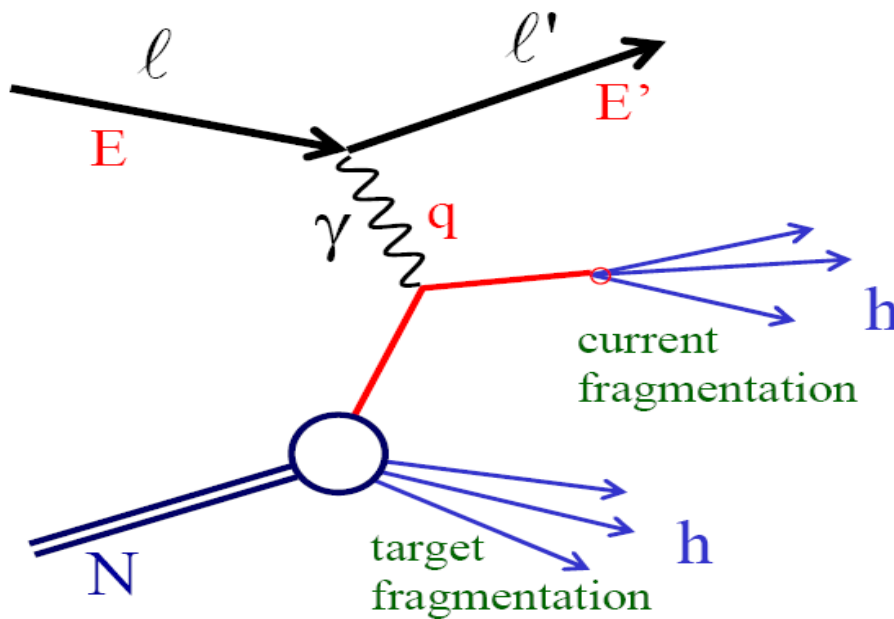
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SIDIS: a key process to investigate the spin structure of the nucleon

$$l N \longrightarrow l' h X$$



lepton interacts with a single parton of the nucleon (via photon exchange)

q virtual photon four-momentum

$$Q^2 = -q^2 > 0 \quad v = E - E'$$

$$x = Q^2 / 2Mv \quad y = v/E$$

$$W^2 = (P_N + q)^2$$

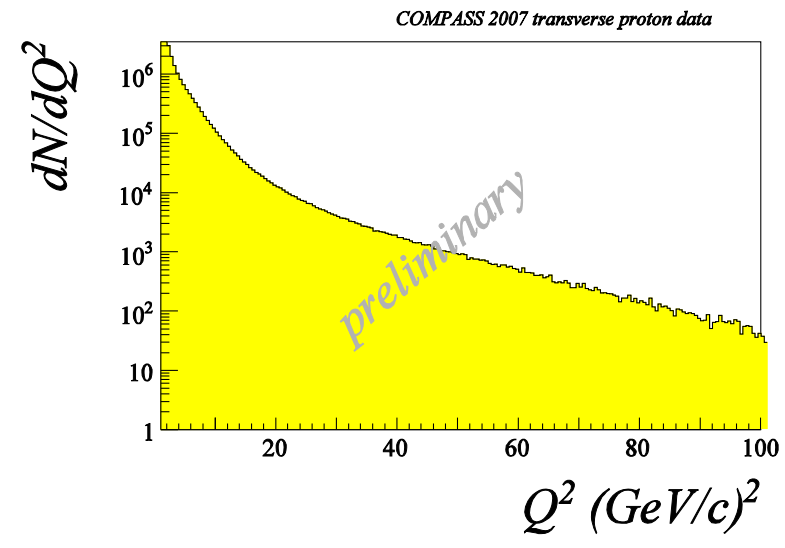
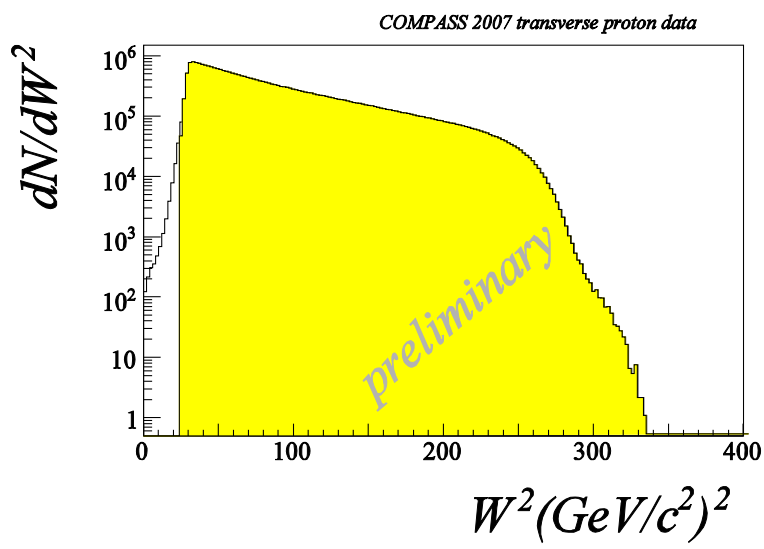
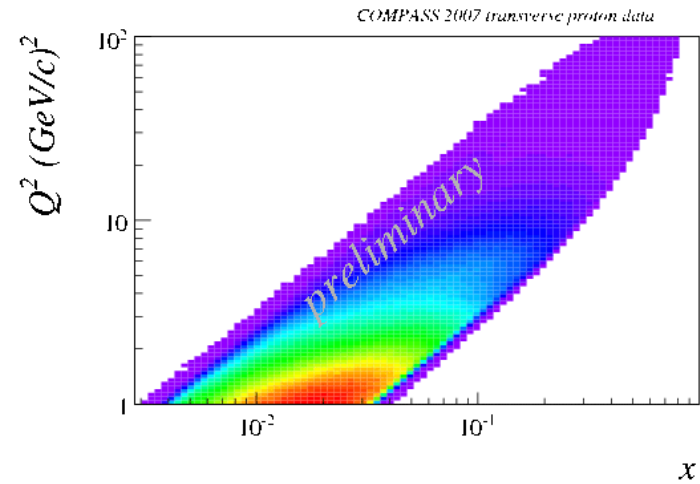
at least one hadron is detected in the final state

$$z = E_h / v$$



Data selection and kinematics

DIS event selection:
 $Q^2 > 1 \text{ (GeV/c)}^2$
 $0.1 < y < 0.9$
 $W > 5 \text{ GeV/c}^2$





Data selection and kinematics

DIS event selection:

$$Q^2 > 1 \text{ (GeV/c)}^2$$

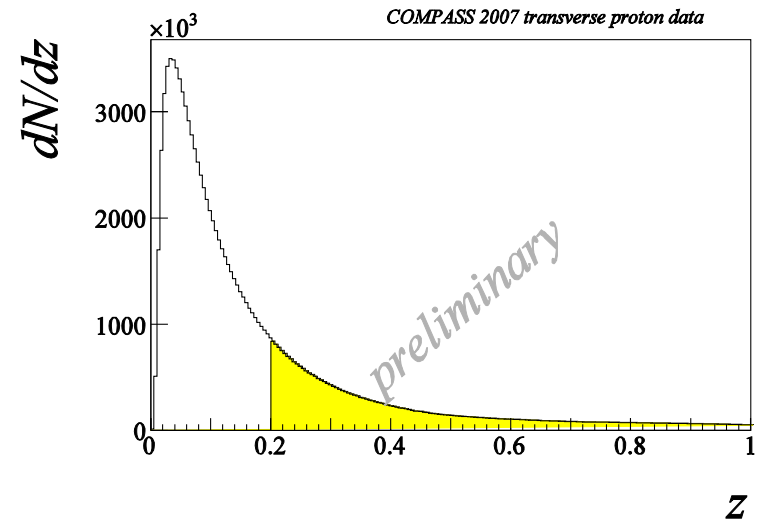
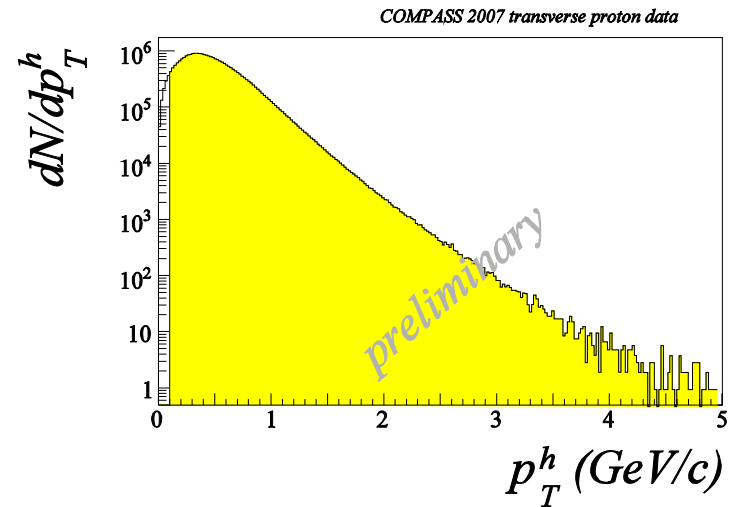
$$0.1 < y < 0.9$$

$$W > 5 \text{ GeV/c}^2$$

Hadron selection

$$p_T > 0.1 \text{ GeV/c}$$

$$z > 0.2$$





SIDIS cross section

From [A. Bacchetta et al.](#),
[JHEP 0702:093,2007.](#)

unpolarized

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} =$$

$$\left\{ \frac{\alpha^2}{xy Q^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \right.$$

$$+ \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h}$$

$$+ S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right]$$

$$+ S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right]$$

these modulations
will be addressed in this talk

transverse
polarization

$$+ |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right.$$

$$+ \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)}$$

$$\left. + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right]$$

$$+ |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right.$$

$$\left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \Big\},$$



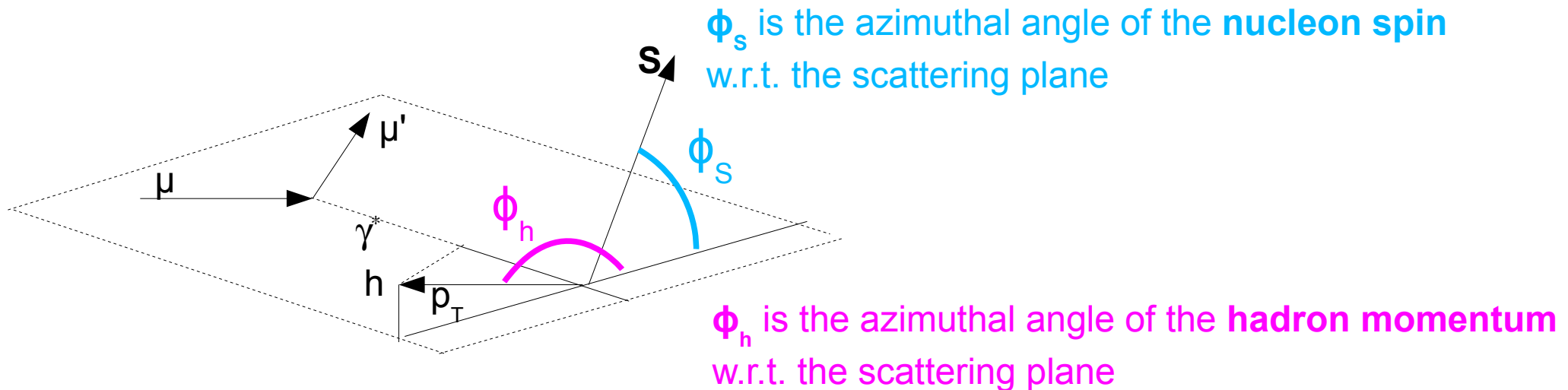
SIDIS cross section

$$\sigma^{lN \rightarrow l' h X} \propto \sum_q q(x) \otimes \sigma^{lq \rightarrow l' q'} \otimes D(z)$$

↑ PDF
 ↑ FF

the complete **SIDIS cross section**, including the quarks transverse momentum, has **18 structure functions** (PDF \otimes FF), 8 leading order

most of them can be accessed by measuring the corresponding azimuthal modulation (**independent** azimuthal modulations in $\phi_s \phi_h$)



the relevant **azimuthal angles** $\phi_s \phi_h$ are evaluated in the **Gamma Nucleon System**



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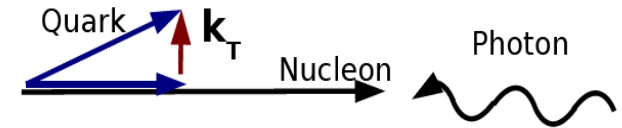


three azimuthal modulations in the unpolarized cross section:

$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1+\epsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\ \left. + \epsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\epsilon(1-\epsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \right.$$

$$F_{UU}^{\cos\phi}$$

mainly **Cahn** effect
kinematical effect due to the **quark transverse momentum**



$$F_{UU}^{\cos 2\phi}$$

given by **Boer-Mulders** function, one of the most famous **TMD PDF**, convoluted with the **Collins FF**

recently increasing interest !

the **Boer-Mulders** function correlates the **quark transverse momentum** and the **quark spin** in an **unpolarized nucleon**



$$F_{UU}^{\sin\phi}$$

kinematical effect proportional to the beam polarization
 (no clear interpretation in term of PM)

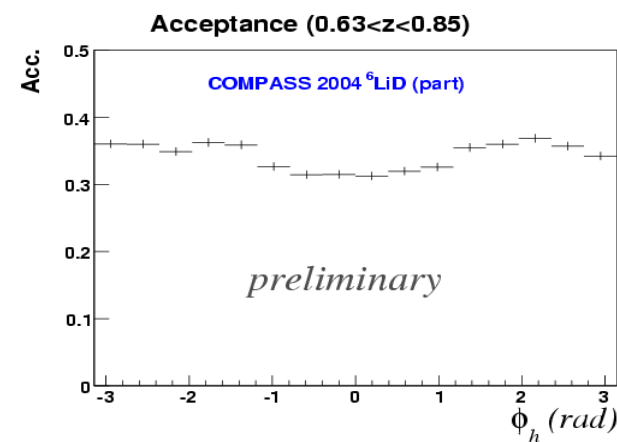
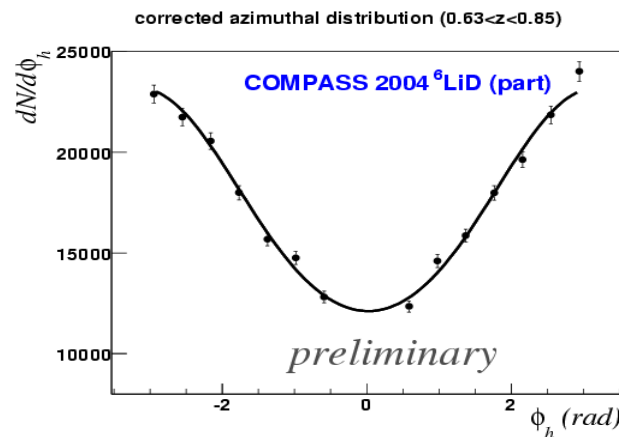
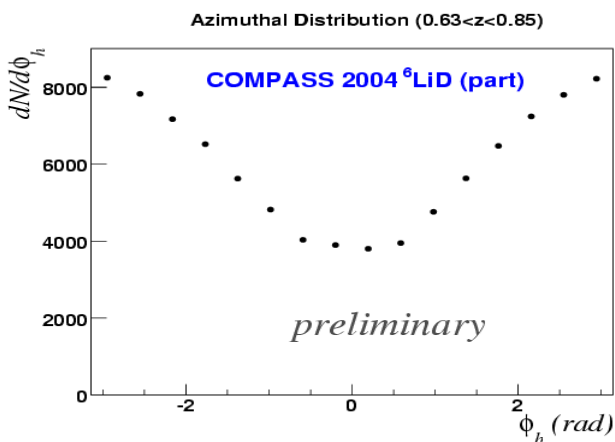
measurement done on deuteron target

data combined to cancel possible polarization dependent terms

To extract the asymmetries:

the azimuthal distributions have to be **corrected by the apparatus acceptance**

→ **dedicated MC simulations** for L and T target polarization data



fit function

$$N_{corr}(\phi) = N_0 \cdot \left\{ 1 + A_{\cos\phi} \cdot \cos(\phi) + A_{\cos 2\phi} \cdot \cos(2\phi) + A_{\sin\phi} \cdot \sin(\phi) \right\}$$

these asymmetries have been measured for the **first time** separately for **positive and negative hadrons**

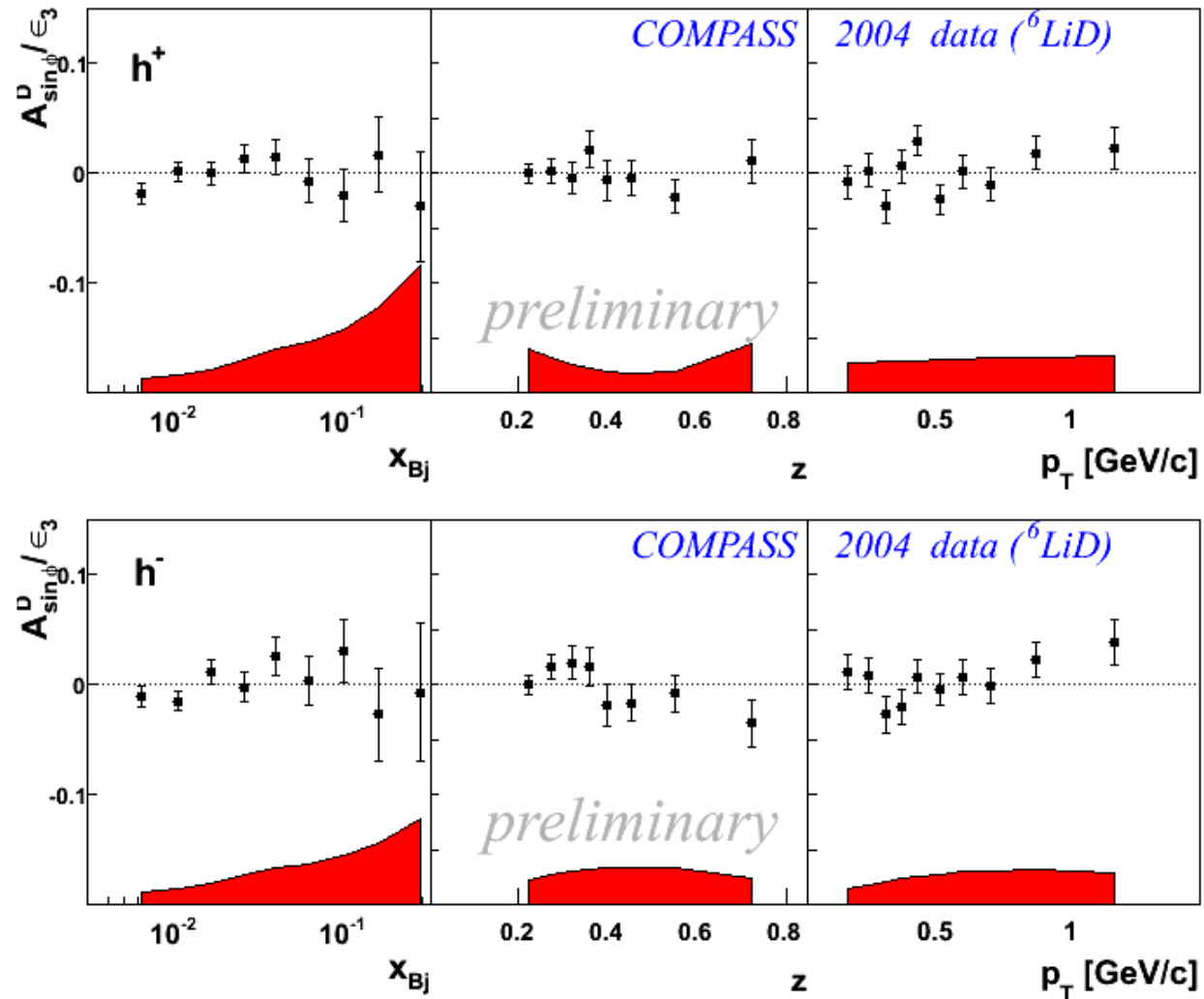


Results

$\sin\phi$: small amplitudes, compatible with zero

$$A_{\sin\phi} / \varepsilon_s$$

$$\varepsilon_s = \frac{2y\sqrt{1-y}}{1+(1-y)^2}$$



red bands are systematic errors

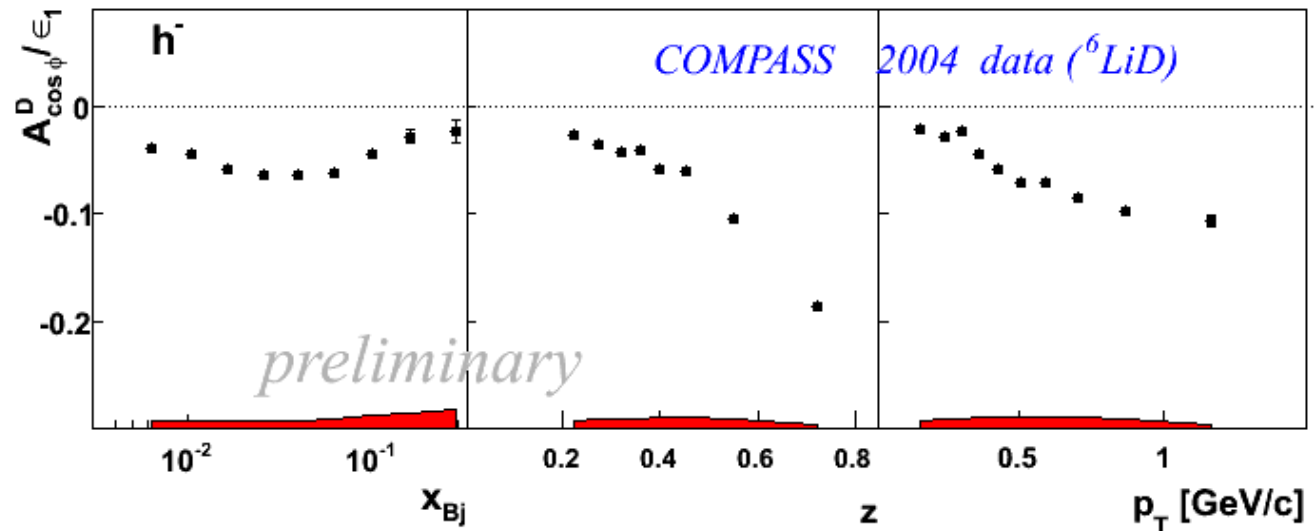
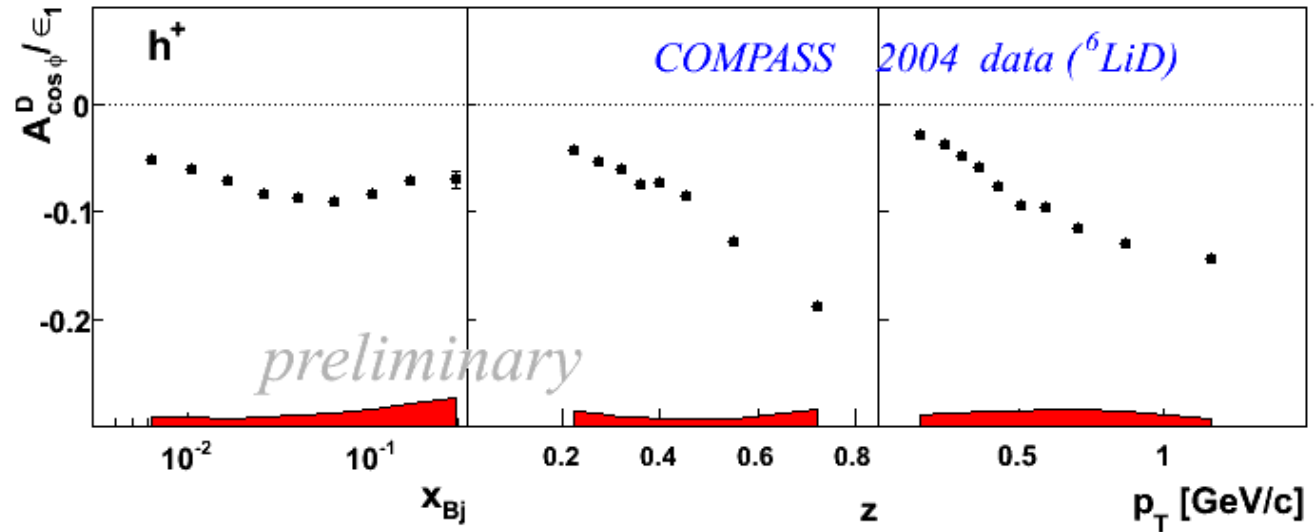


Results

$\cos\phi$: strong signal both for positive and negative hadrons

$$A_{\cos\phi} / \varepsilon_c$$

$$\varepsilon_c = \frac{2(2-y)\sqrt{1-y}}{1+(1-y)^2}$$



red bands are systematic errors

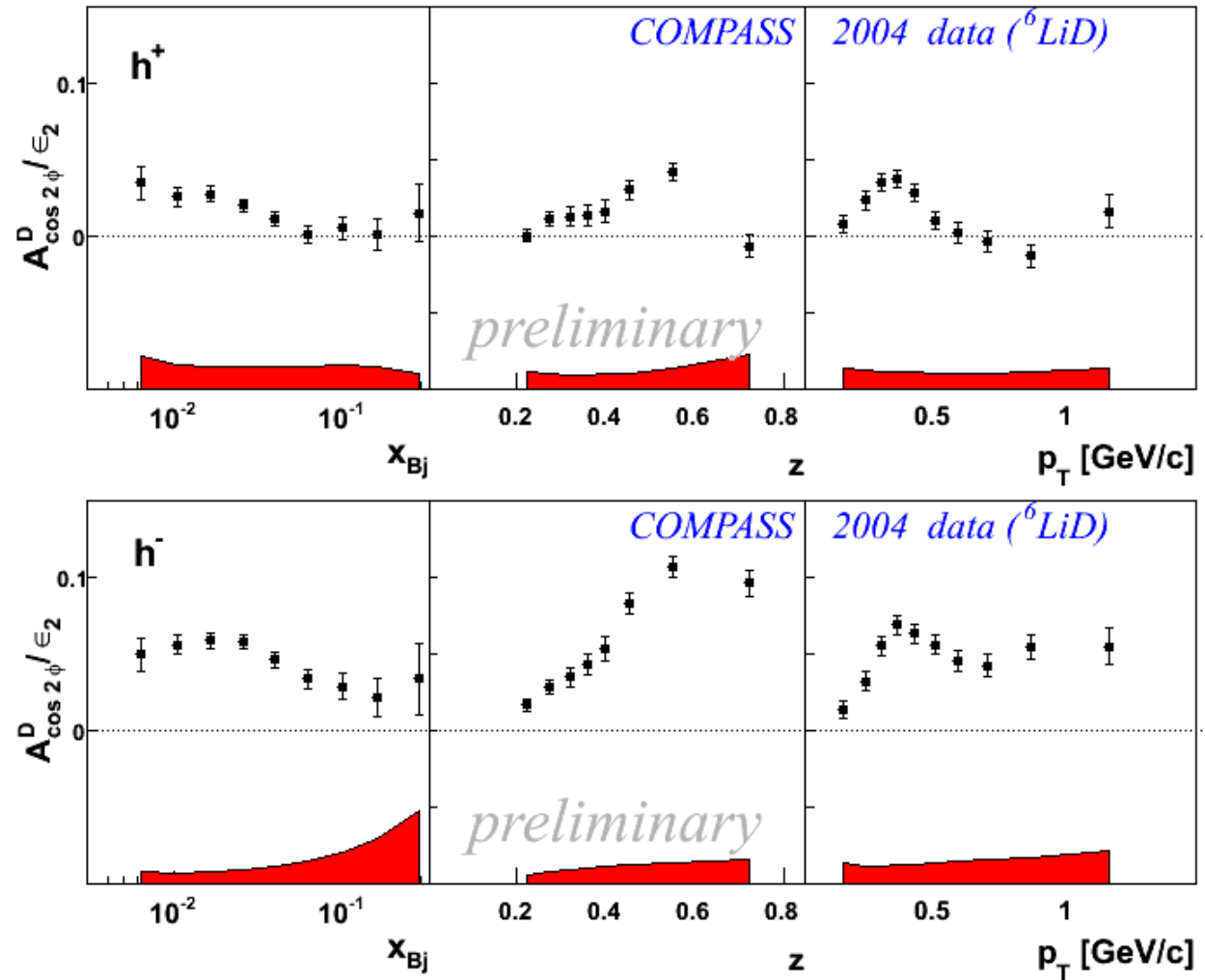


Results

$\cos 2\phi$: different from zero and different
for positive and negative hadrons

$$A_{\cos 2\phi} / \epsilon_2$$

$$\epsilon_2 = \frac{2(2-y)}{1+(1-y)^2}$$

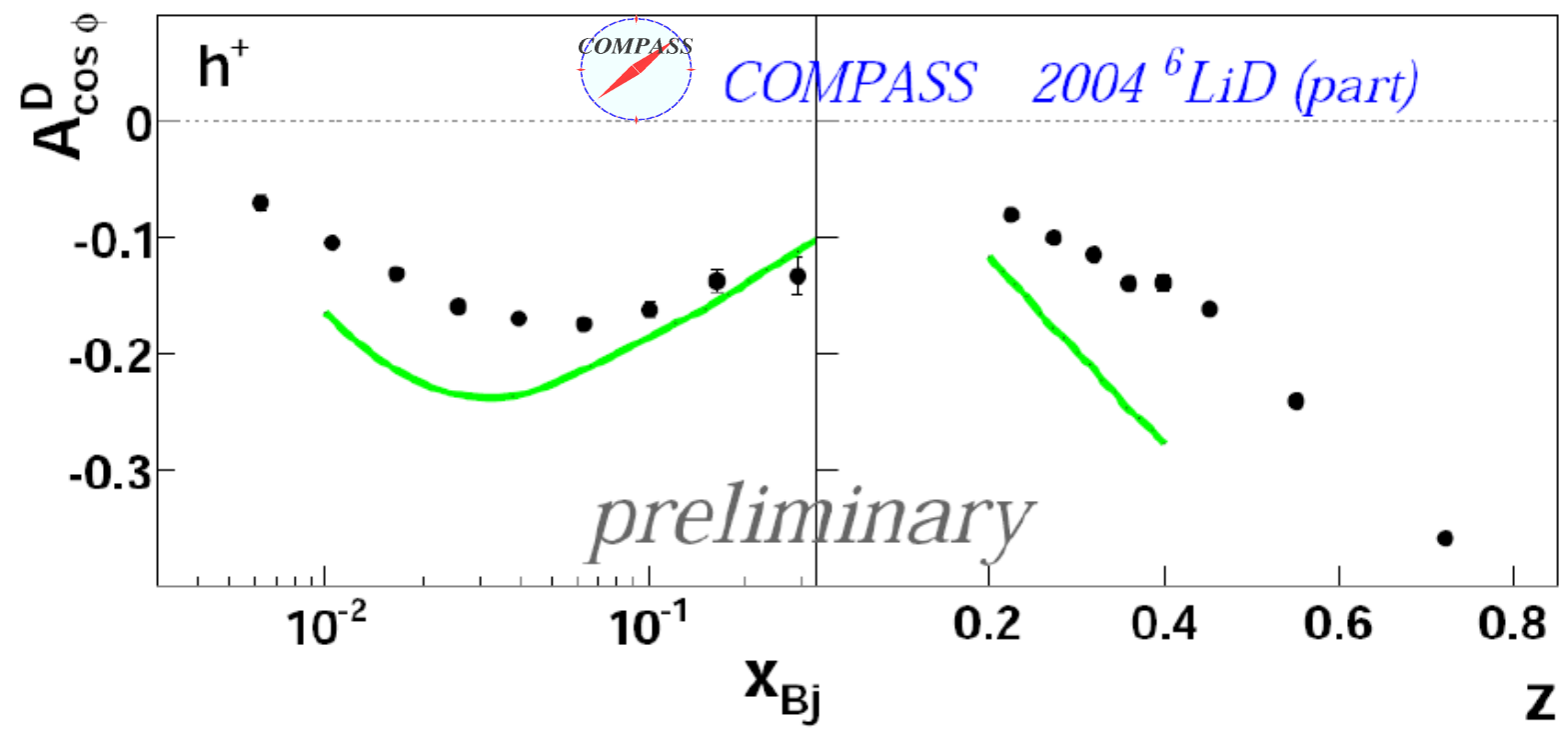


red bands are systematic errors



Comparison with predictions

$\cos\phi$

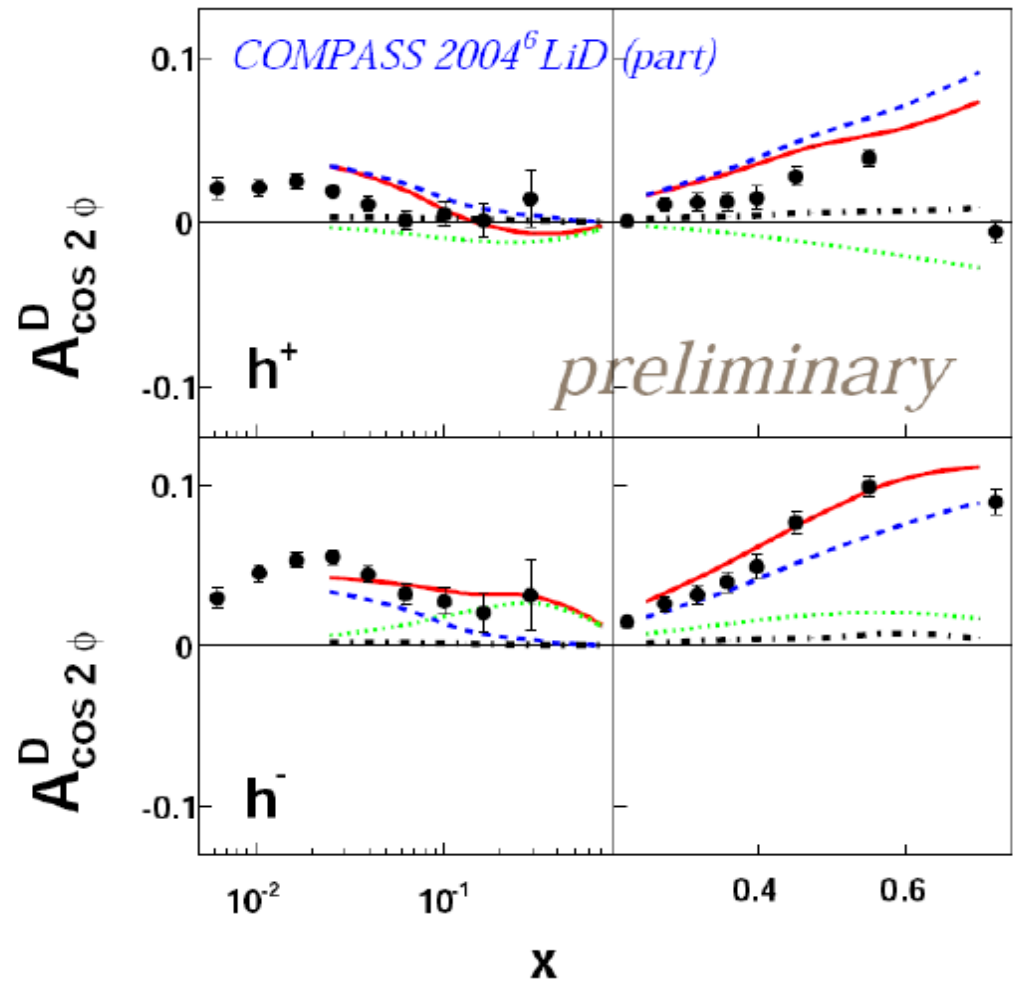


— M. Anselmino, M. Boglione, A. Prokudin, C. Türk
Eur. Phys. J. A 31, 373-381 (2007)
does not include Boer – Mulders contribution



Comparison with predictions

$\cos 2\phi$



V. Barone, A. Prokudin, B.Q. Ma
 Phys.Rev. D78: 045022, 2008

- sum of all contributions
- - - Cahn effect
- ⋯ Boer-Mulders
- · - · QCD (first order)



8 azimuthal modulations can be measured in SIDIS on transversely polarized target

$$\begin{aligned}
 \frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = & \\
 & \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\
 & + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \\
 & + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \\
 & + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \\
 & + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
 & + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\
 & \left. + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right] \\
 & + |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right. \\
 & \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \left. \right\},
 \end{aligned}$$

The most known are:

Sivers asymmetry

Collins asymmetry



Collins asymmetry is a very important measurement
it gives access to the:

Transversity PDF

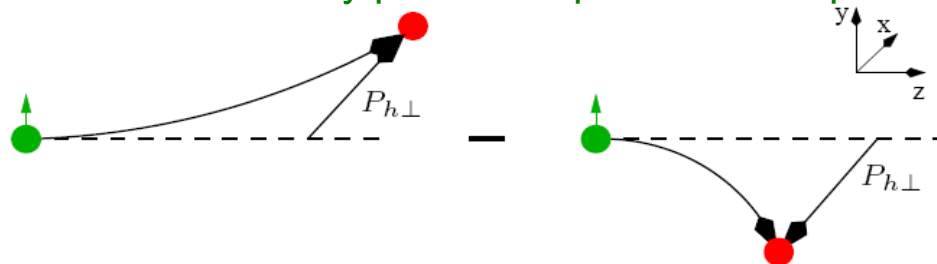
- which is one of the three PDF needed to **describe** the **spin structure of the nucleon** at **leading order** (with helicity PDF and unpolarized PDF)
- it gives the probability to find a **quark** with **spin parallel or anti-parallel to the nucleon's spin** in a **transversely polarized** nucleon



- it is **chiral odd** → can be measured only in **SIDIS** on transversely polarized target (convoluted with another chiral odd function)

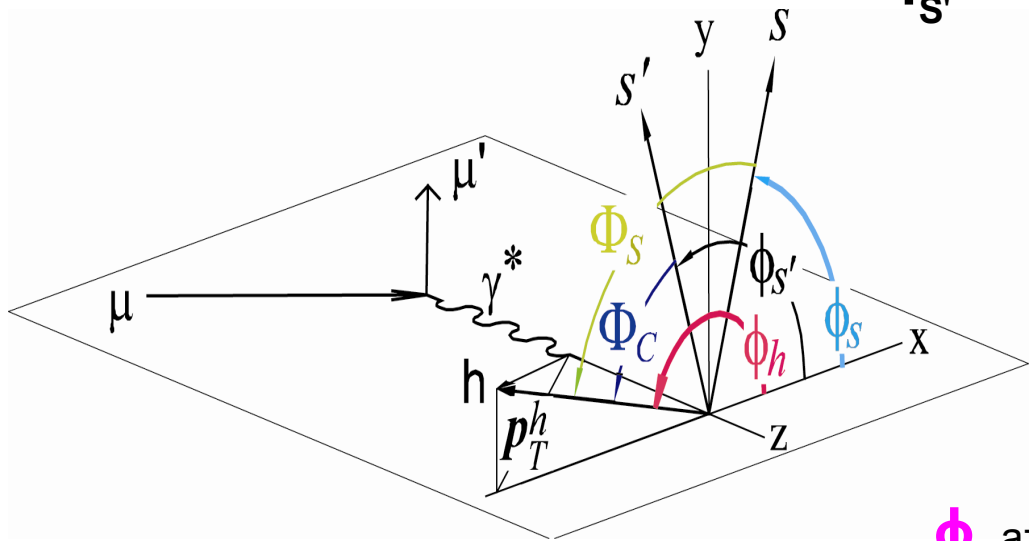
Transversity \otimes Collins FF

left right asymmetry in the fragmentation
of the transversely polarized quark in an unpol. hadron





Collins asymmetry appears as a modulation in Φ_C



$\phi_{S'}$ azimuthal angle of the **spin** vector of the **fragmenting quark**

$$\phi_{S'} = \pi - \phi_S$$

$$\Phi_C = \phi_h - \phi_{S'} = \phi_h + \phi_S - \pi$$

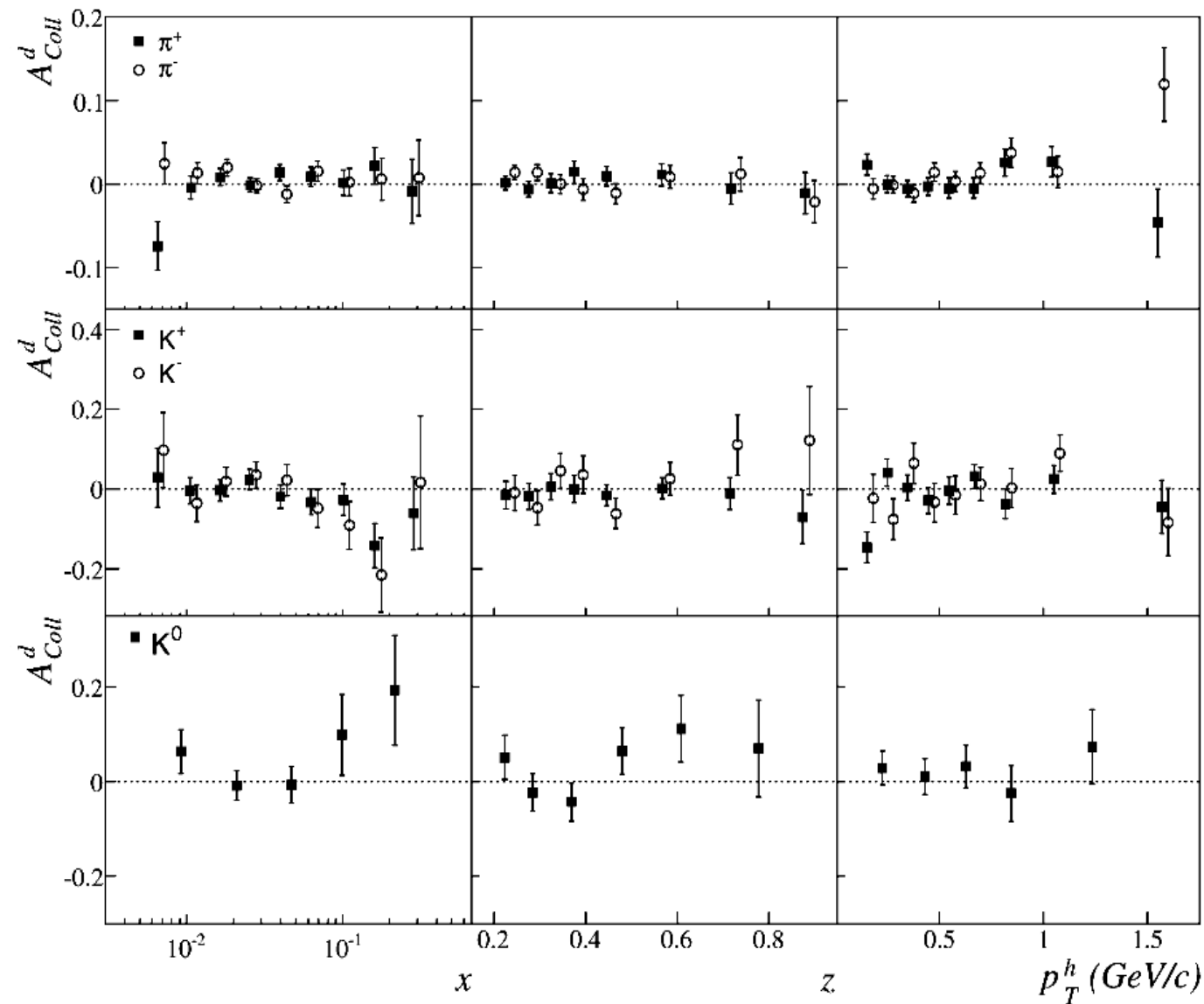
Collins angle

ϕ_h azimuthal angle of the **hadron momentum**

$$\mathbf{A}_{\text{Coll}} = \frac{\mathbf{A}_C^h}{\mathbf{f} \cdot \mathbf{P}_T \cdot D_{nn}} = \frac{\sum_q e_q^2 \Delta_T^0 \mathbf{q} \cdot \Delta_T^0 \mathbf{D}_q^h}{\sum_q e_q^2 \cdot \mathbf{q} \cdot D_q^h} \leftarrow \text{“transversity” PDF} \otimes \text{Collins FF}$$



COMPASS measurement on transversely polarized deuteron (2002-2004 data)



COMPASS Collaboration
Physics Letters B 673 (2009) 127–135

*Values corrected for the purity;
systematic error below
30% of the statistical one*



Collins asymmetries on deuteron are compatible with zero

at variance with **HERMES** which measured a **signal** for both positive and negative hadrons on **transversely polarized proton**

- Naïve interpretation of the results (parton model, valence region)

$$A_{Coll}^{d,\pi^+} \simeq \frac{\Delta_T u_v + \Delta_T d_v}{u_v + d_v} \frac{4\Delta_T^0 D_1 + \Delta_T^0 D_2}{4D_1 + D_2} \quad A_{Coll}^{d,\pi^-} \simeq \frac{\Delta_T u_v + \Delta_T d_v}{u_v + d_v} \frac{\Delta_T^0 D_1 + 4\Delta_T^0 D_2}{D_1 + 4D_2}$$

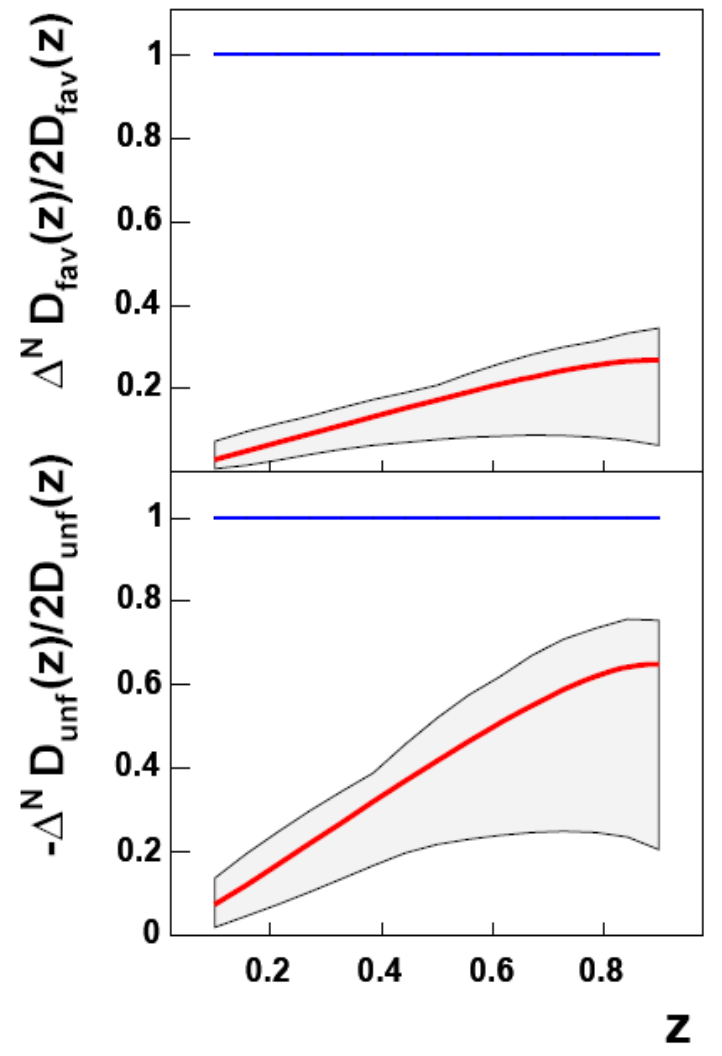
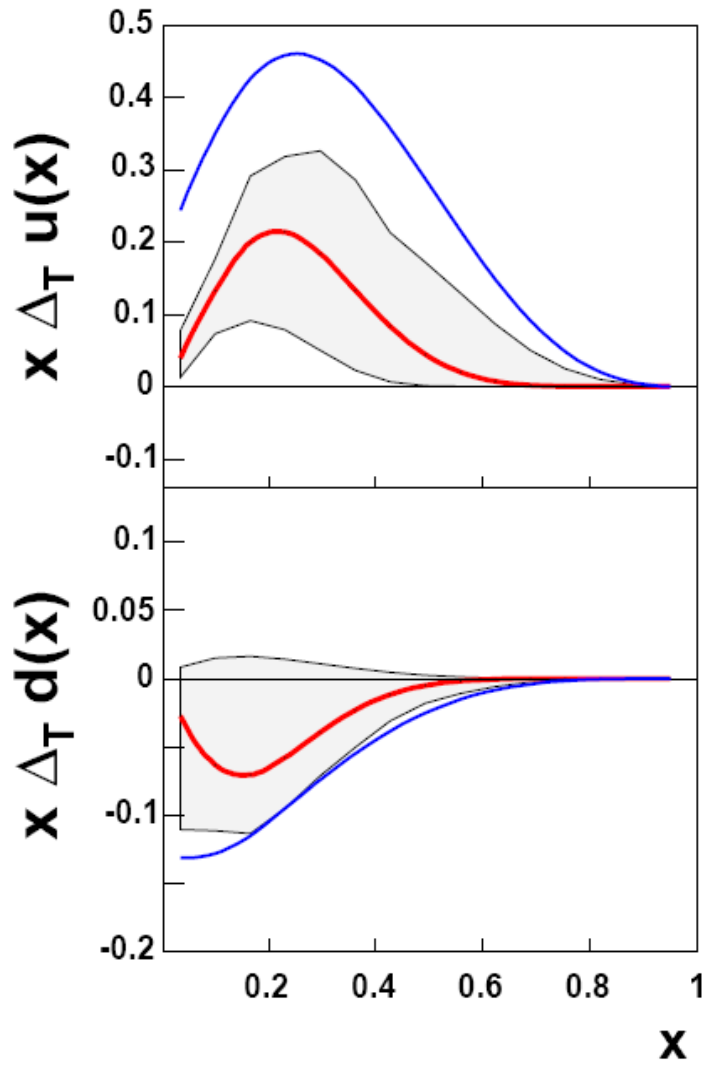
Small asymmetries $\rightarrow \Delta_T u(\mathbf{x}) + \Delta_T d(\mathbf{x}) \sim 0$

measurements on deuteron (isoscalar target) important to access $\Delta_T d$

- data taken with deuteron (COMPASS), proton (HERMES) and $e^+ e^- \rightarrow$ hadrons (BELLE)
global fit \rightarrow consistent description



first extraction of the:
transversity PDF and the **Collins FF** for **u** and **d** quarks

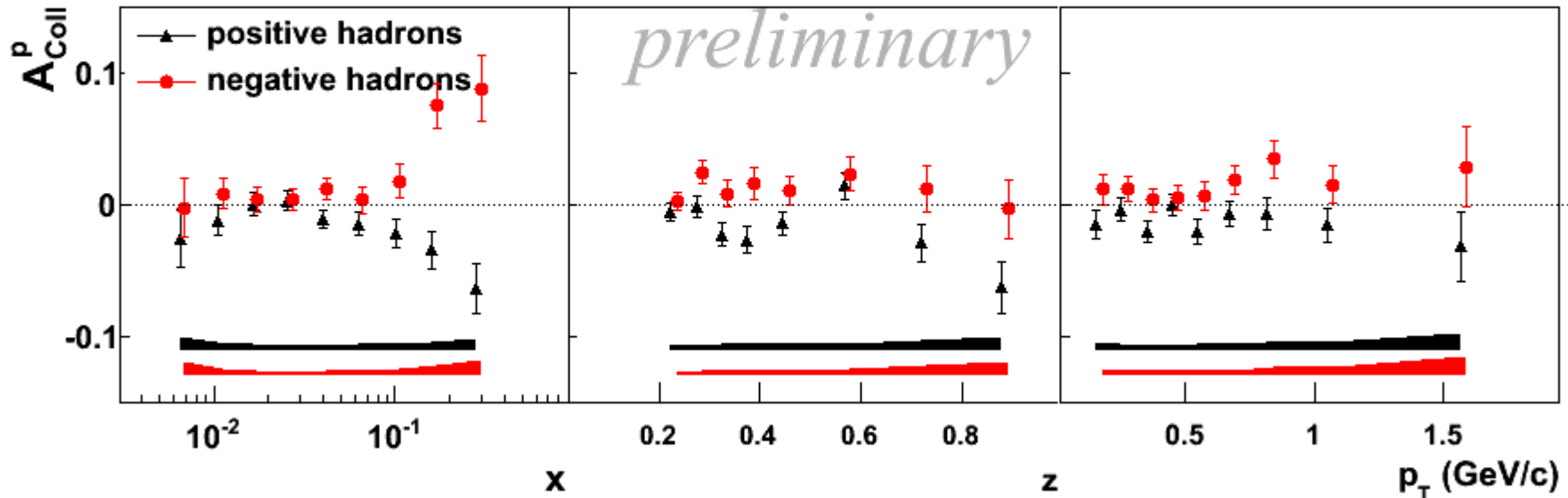


Anselmino et al. Phys.Rev. D75 (2007) 054032



COMPASS Collins asymmetry on transversely polarized proton (DIS09)

COMPASS 2007 proton data



bands are systematic errors

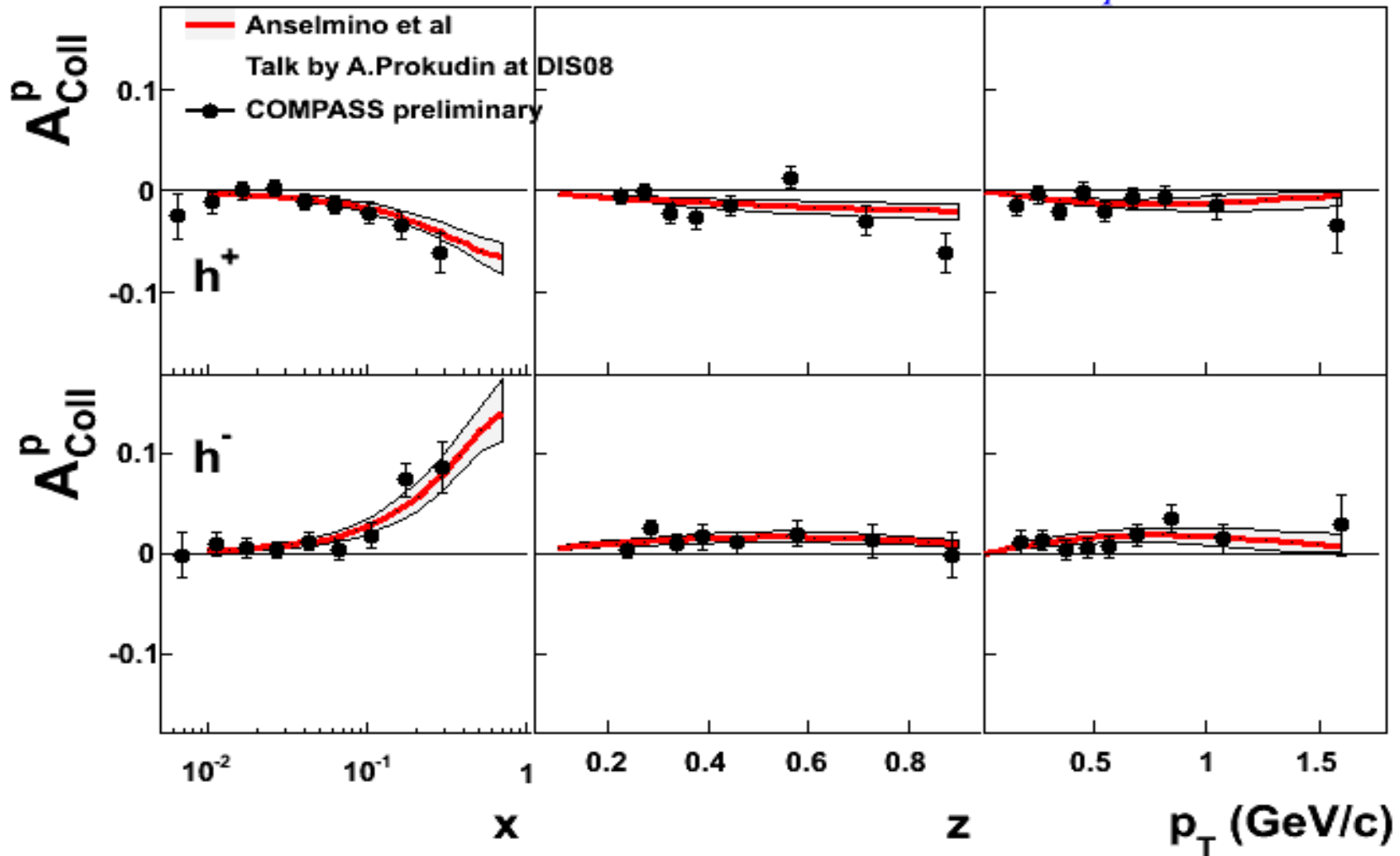
The effect is still there at higher Q^2 (w.r.t. HERMES)

at small x , the asymmetries are compatible with zero
in the valence region the asymmetries are different from zero,
of opposite sign for positive and negative hadrons,
and have the same strength and sign as HERMES



COMPASS measurements on proton confirm present picture and are **consistent with prediction** (from global fit)

COMPASS 2007 proton data



last prediction from Anselmino group



$$\begin{aligned}
 \frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = & \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\
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 & + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
 & \quad \left. + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \right. \\
 & \quad \left. + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right] \\
 & + |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right. \\
 & \quad \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \left. \right\},
 \end{aligned}$$

Sivers asymmetry

Collins asymmetry



- The **Sivers** function is one of the most studied **TMD** parton distribution functions

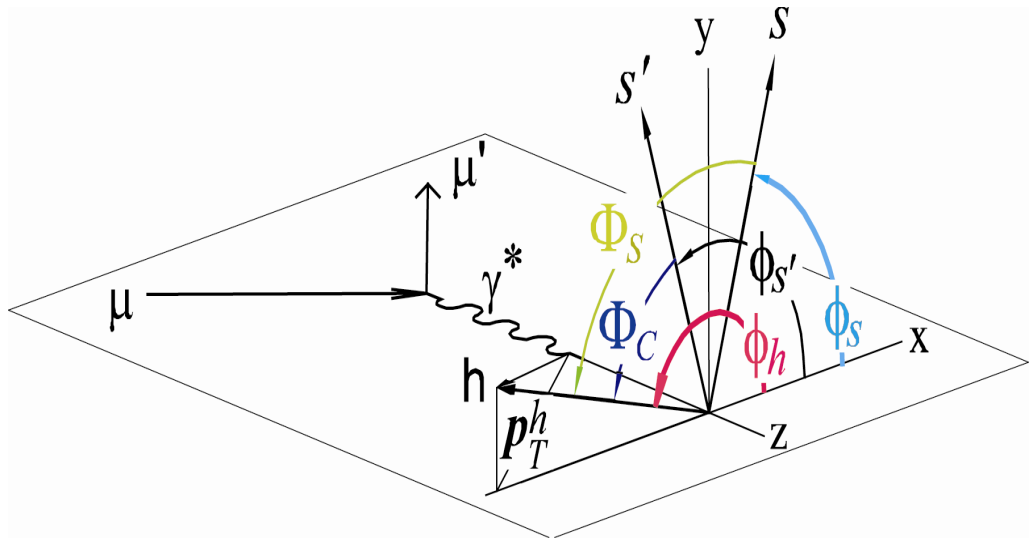
- correlates the **quark transverse momentum** and the **nucleon spin** in a transversely polarized nucleon



- it is related to **angular orbital momentum** of the quark in a transversely polarized nucleon



Sivers asymmetry appears as a modulation in Φ_S



ϕ_h azimuthal angle of the **hadron momentum**

ϕ_S azimuthal angle of the **nucleon spin**

$$\Phi_S = \phi_h - \phi_S$$

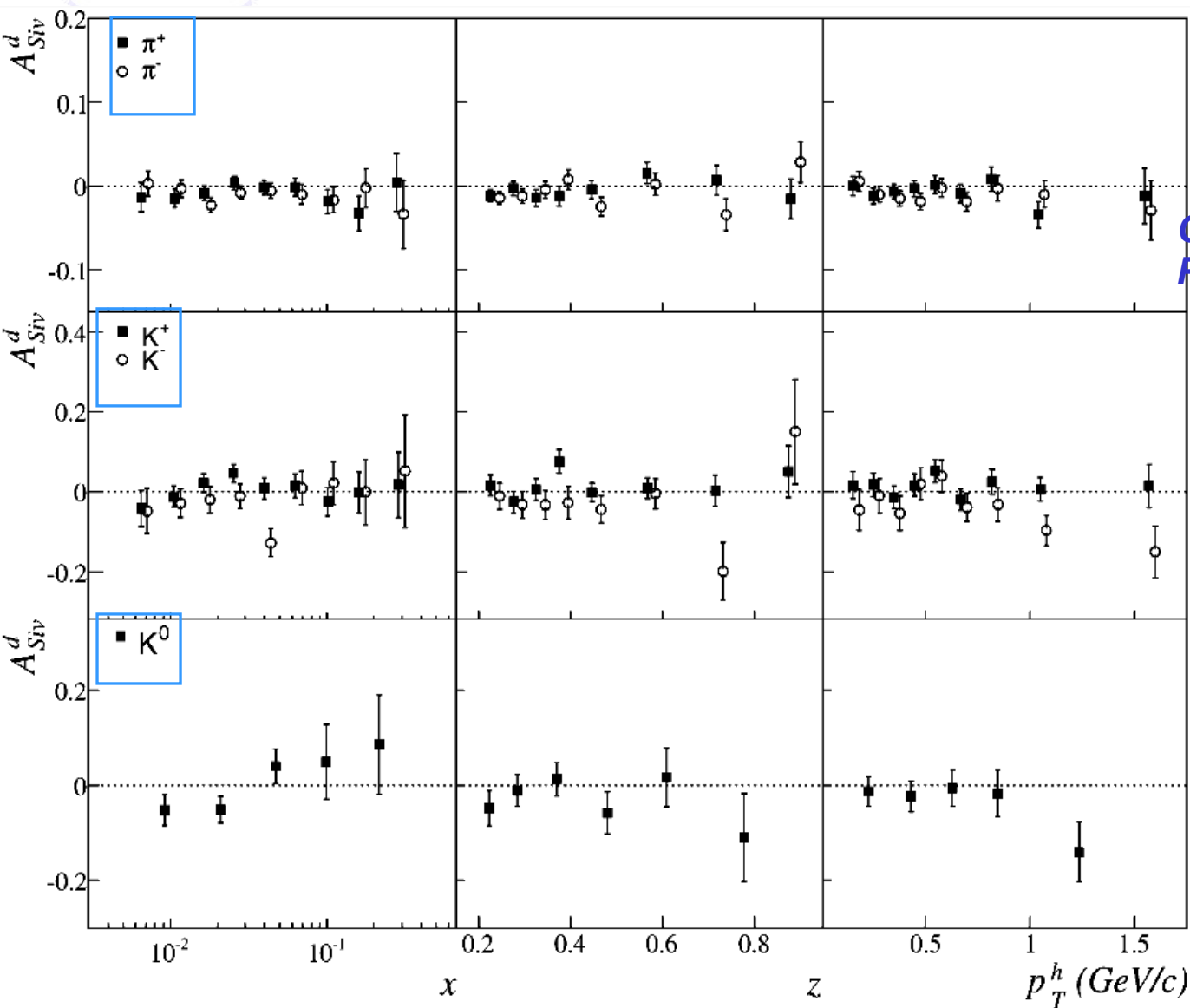
Sivers angle \blacktriangleright measured azimuthal modulation

$$\mathbf{A}_{\text{Siv}} = \frac{\mathbf{A}_S^h}{\mathbf{f} \cdot \mathbf{P}_T} = \frac{\sum_q \mathbf{e}_q^2 \cdot \Delta_0^T \mathbf{q} \cdot \mathbf{D}_q^h}{\sum_q \mathbf{e}_q^2 \cdot \mathbf{q} \cdot \mathbf{D}_q^h}$$

Sivers TMD PDF \otimes Unpolarized FF



COMPASS measurement on transversely polarized deuteron (2002-2004 data)



COMPASS Collaboration
Physics Letters B 673 (2009) 127–135

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systematic error below
30% of the statistical one



COMPASS measurements on polarized deuteron **small Sivers asymmetries compatible with zero**

global fit together with *HERMES* data on proton
extraction of Sivers function

asymmetries $\neq 0$ for π^+

naïve interpretation (parton model, valence region)

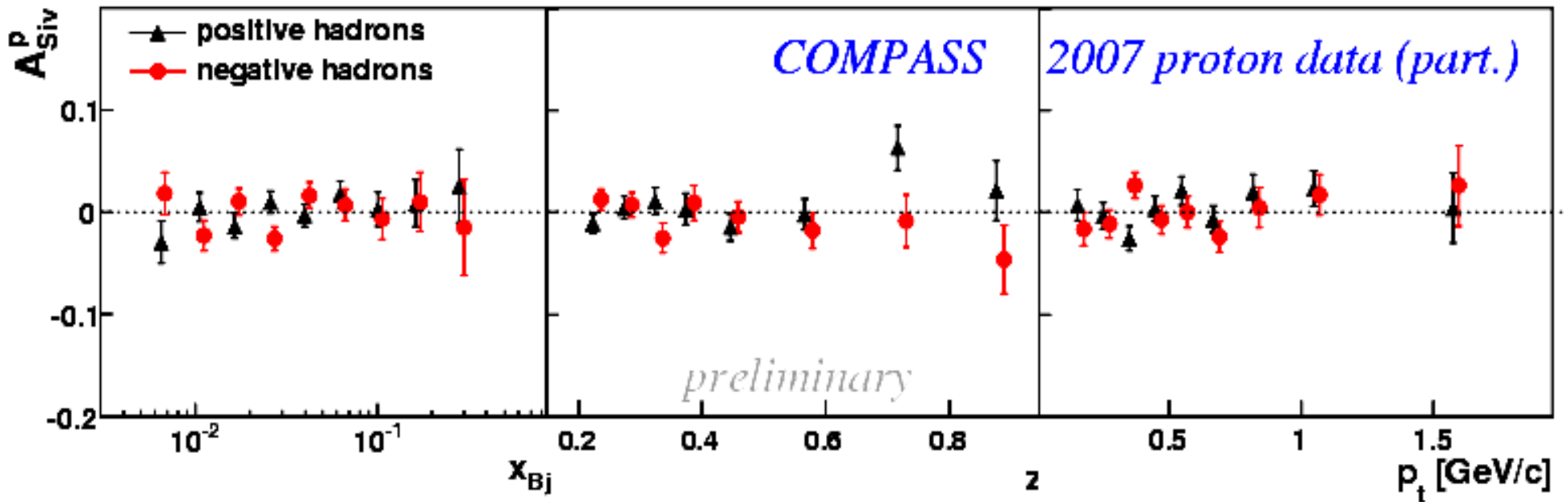
$$A_{Siv}^{d,\pi^+} \simeq A_{Siv}^{d,\pi^-} \simeq \frac{\Delta_0^T u_v + \Delta_0^T d_v}{u_v + d_v}$$

small asymmetries suggest

$$\Delta_0^T d_v \simeq -\Delta_0^T u_v$$



COMPASS Siverts asymmetry: measurement on transversely polarized proton (DIS09)



systematic errors $\sim 0.5 \sigma$ stat

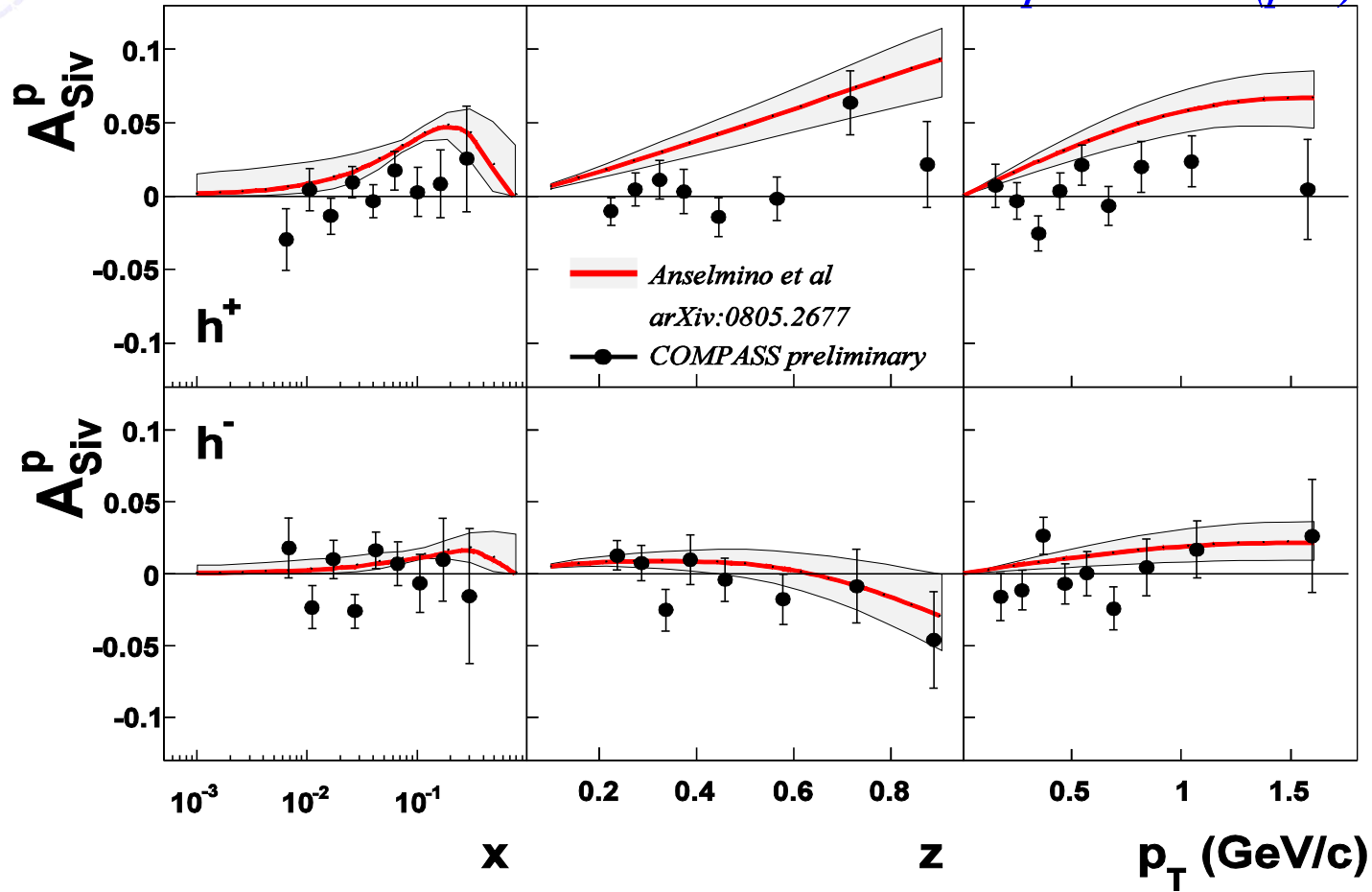
the measured asymmetries are small, compatible with zero

intriguing results , not easy to explain in the present theoretical picture



COMPASS Siverts asymmetries on proton: comparison with prediction from Anselmino et al.

COMPASS 2007 transverse proton data (part)



marginal agreement between data and predictions

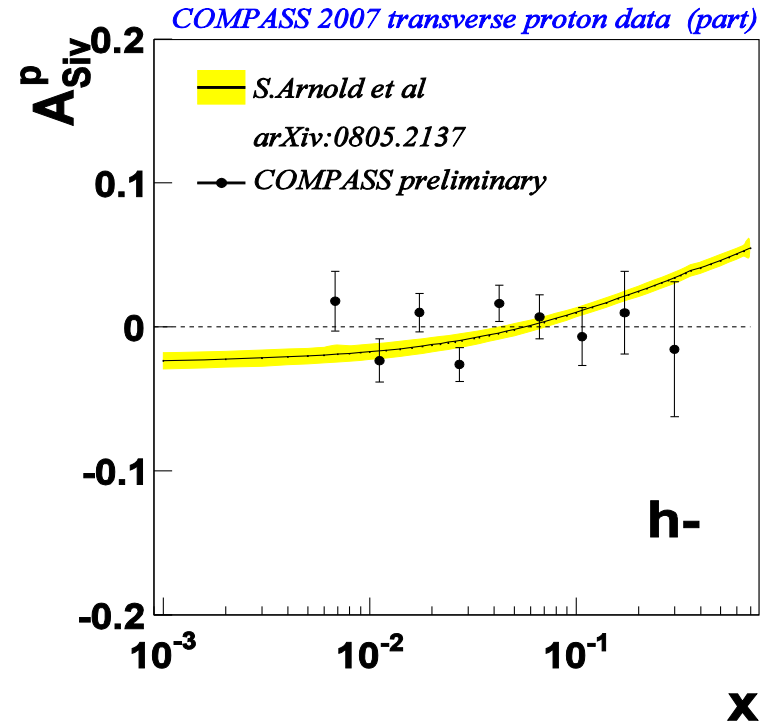
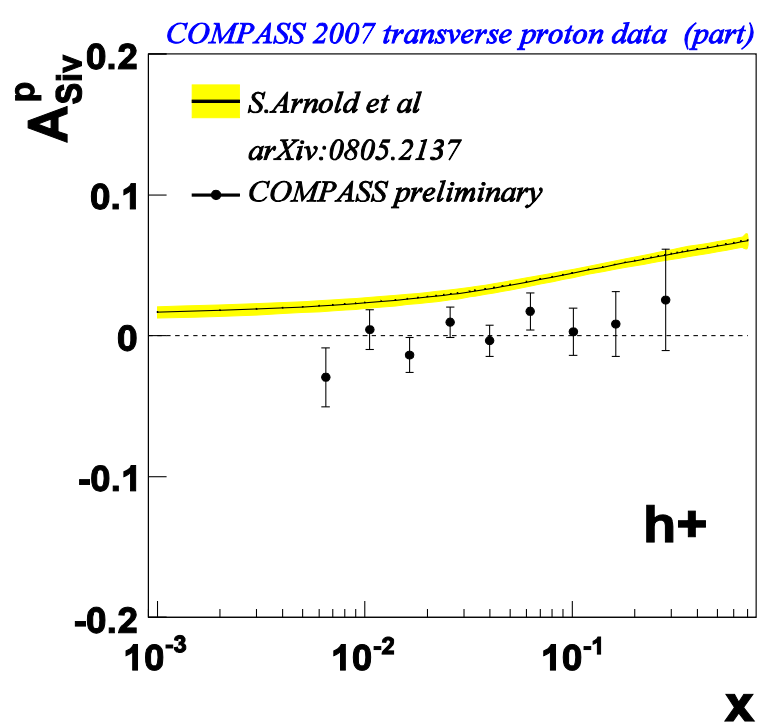
more precise high energy data urgently needed



COMPASS Siverts asymmetries on proton: comparison with prediction from

based on Quark Soliton Model

S.Arnold, A.V.Efremov, K.Goeke, M.Schlegel and P.Schweitzer, arXiv:0805.2137



marginal agreement between data and predictions



From [A. Bacchetta et al., JHEP 0702:093,2007.](#)

$$\begin{aligned}
 \frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = & \frac{\alpha^2}{xy Q^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\
 & + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \\
 & + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \\
 & + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \\
 & + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
 & + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\
 & + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \left. \right] \\
 & + |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right. \\
 & \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \left. \right\},
 \end{aligned}$$

other **6 single spin asymmetries**
 on a transversely polarized target,
 all measured by **COMPASS** on **deuteron**
 and **found to be compatible with zero**



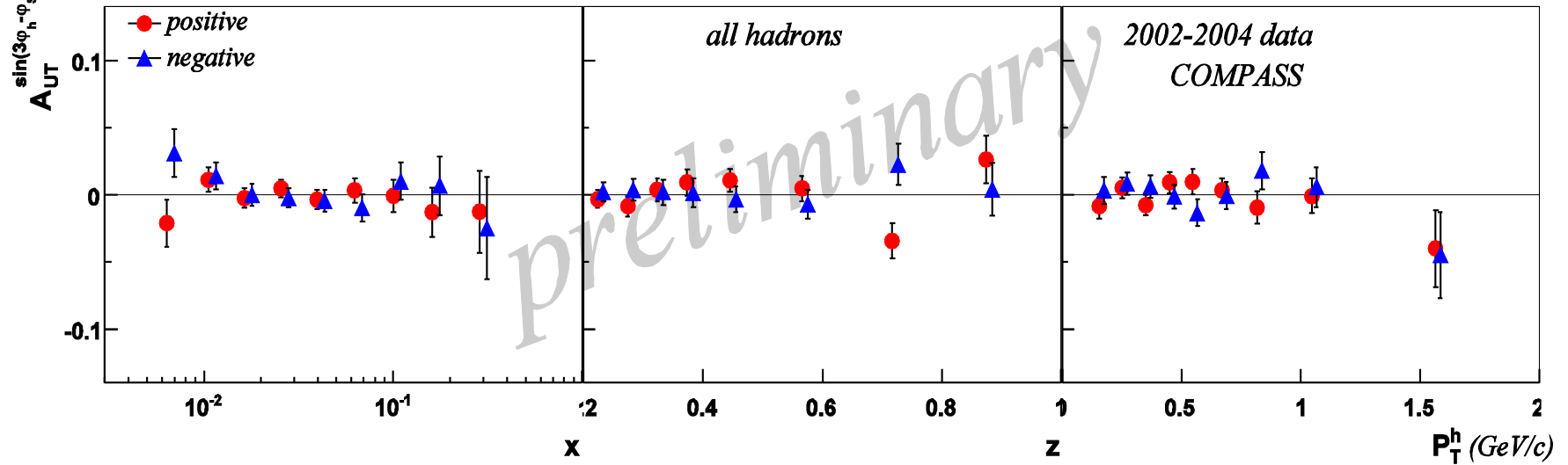
$$F_{UT}^{\sin(3\phi_h - \phi_S)}$$

given by the so called “pretzelosity” PDF

correlates the **quark transverse momentum** and the **quark spin** in a transversely polarized nucleon



$$A_{UT}^{\sin(3\phi_h - \phi_S)}$$





OUTLINE

- The COMPASS experiment
- SIDIS
- Hadron azimuthal asymmetries
 - unpolarized target
 - transversely polarized target
- **Two hadrons azimuthal asymmetries**
- Conclusions and outlook



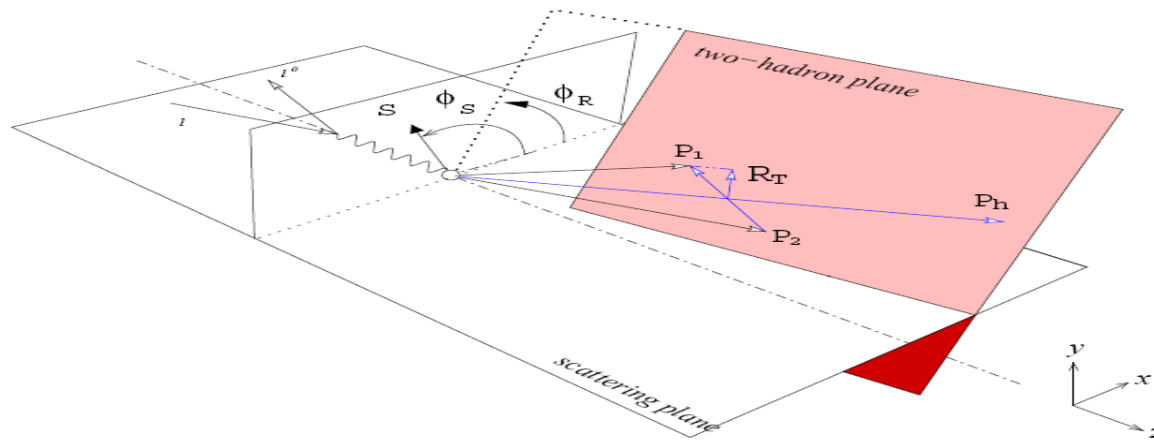
2 h asymmetries

Transversity can be also accessed in inclusive production of hadron pairs by measuring an **azimuthal asymmetries** in the angle:

$$\Phi_{RS} = \Phi_{RT} - \Phi_{S'}$$

is the angle of \mathbf{R}_T w.r.t the gamma

$$\mathbf{R}_T = \frac{z_2 \mathbf{P}_1 - z_1 \mathbf{P}_2}{z_1 + z_2}$$



$$A_{RS} = \frac{1}{\mathbf{f} \cdot \mathbf{P}_T \cdot \mathbf{D}} \cdot \mathbf{A} = \frac{\sum_q e_q^2 \Delta_T q(\mathbf{x}) H_q^\perp(\mathbf{z}, M_h^2)}{\sum_q e_q^2 \mathbf{q}(\mathbf{x}) \cdot \mathbf{D}_q^h(\mathbf{z}, M_h^2)}$$

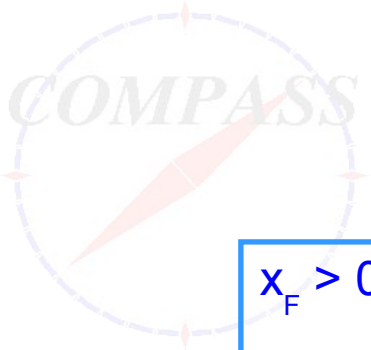
transversity distribution function

interference fragmentation function

A. Bacchetta, M. Radici, hep-ph/0407345

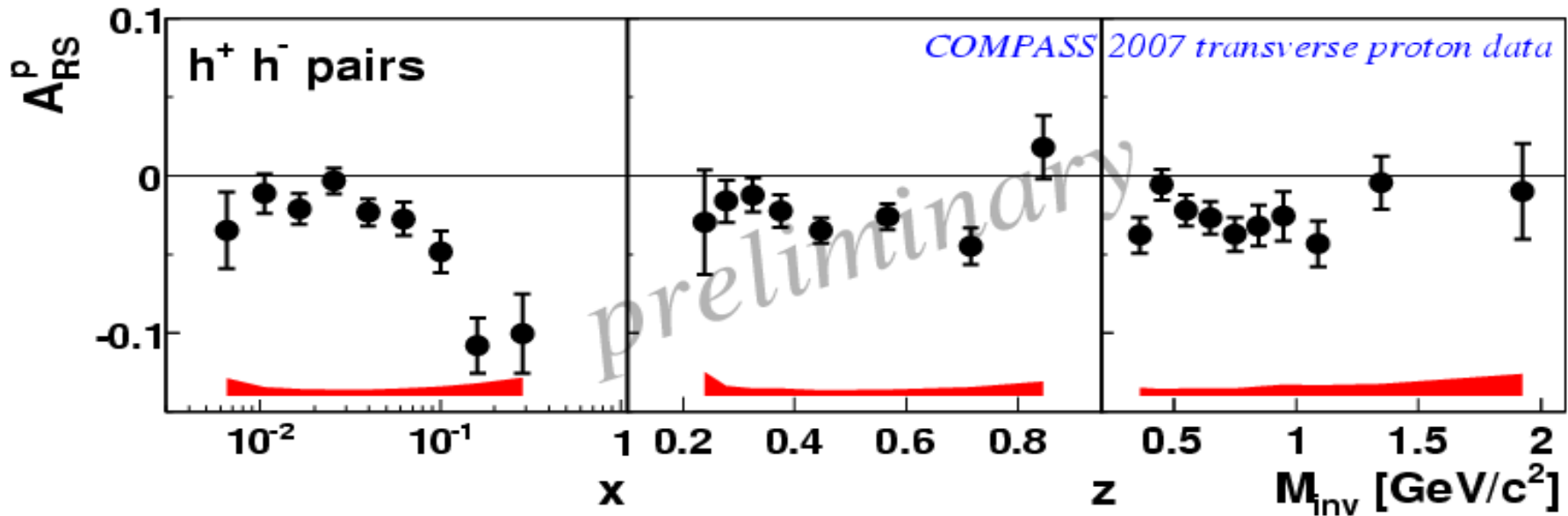
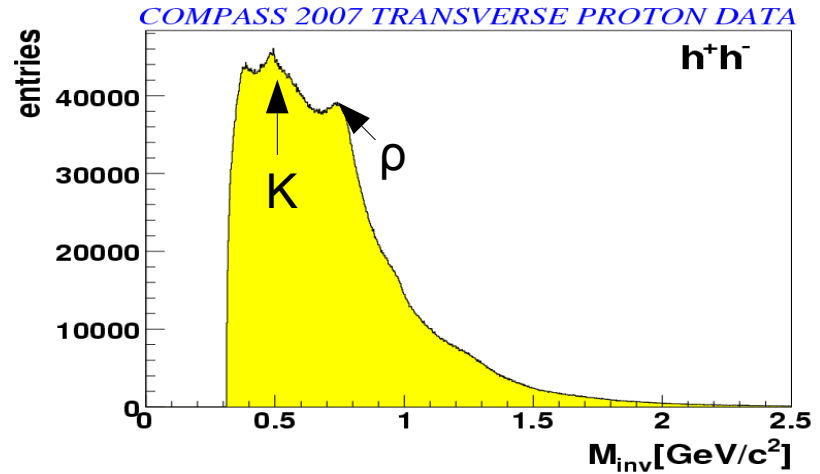
X. Artru, hep-ph/0207309

COMPASS measured small asymmetries (compatible with zero) on deuteron



2 hadrons COMPASS asymmetry on transversely polarized proton

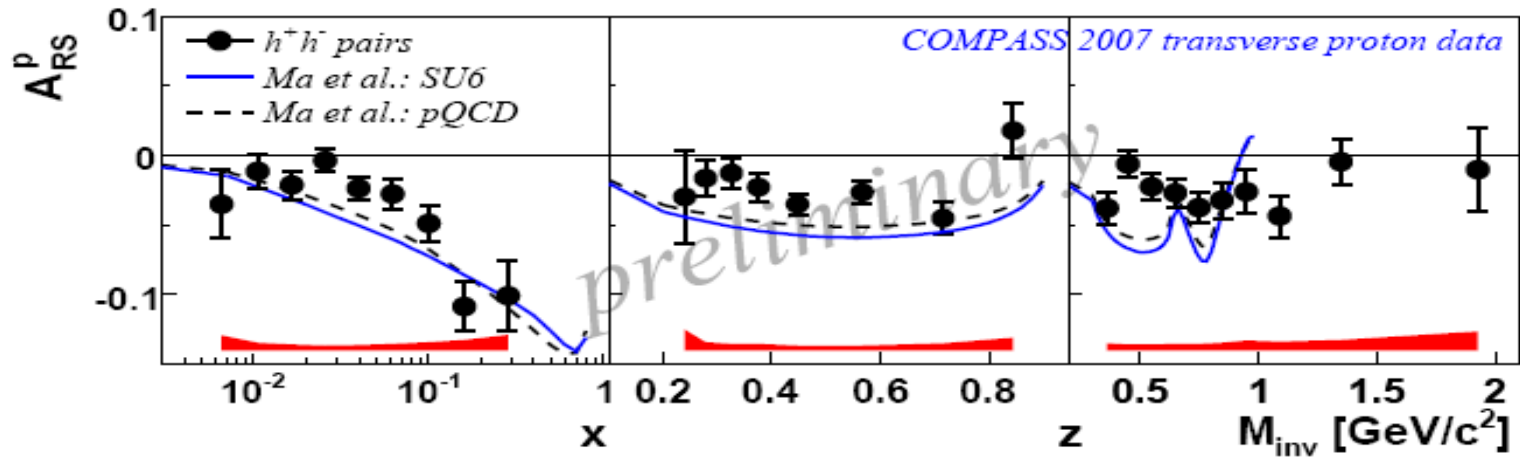
$x_F > 0.1$
 $z_1, z_2 > 0.1$
 $Z = z_1 + z_2 < 0.9$
 $R_T > 0.07 \text{ GeV}/c$



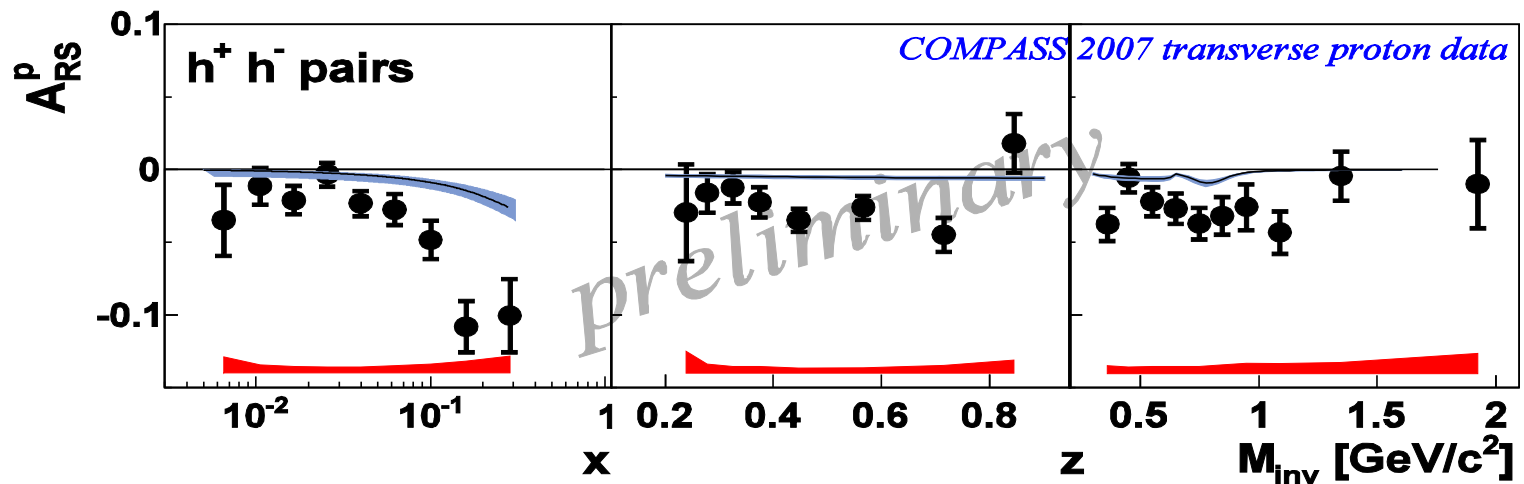
in the **valence region** the asymmetries are **different from zero**
and the signal is larger than for the Collins asymmetry

2 hadrons : comparison with predictions

Ma et al. private communication



Bacchetta et al. private communication





Conclusions and outlook

- many interesting results on **transverse momentum** and **transverse spin** from **COMPASS** data both on **deuteron** and **proton** polarized target
→ flavour separation

- **near future**

- new results on unpolarized azimuthal asymmetries on deuteron (final results)

on proton

- Collins and Sivers asymmetries on identified hadrons
- 2 hadrons asymmetries on identified hadrons



Conclusions and outlook

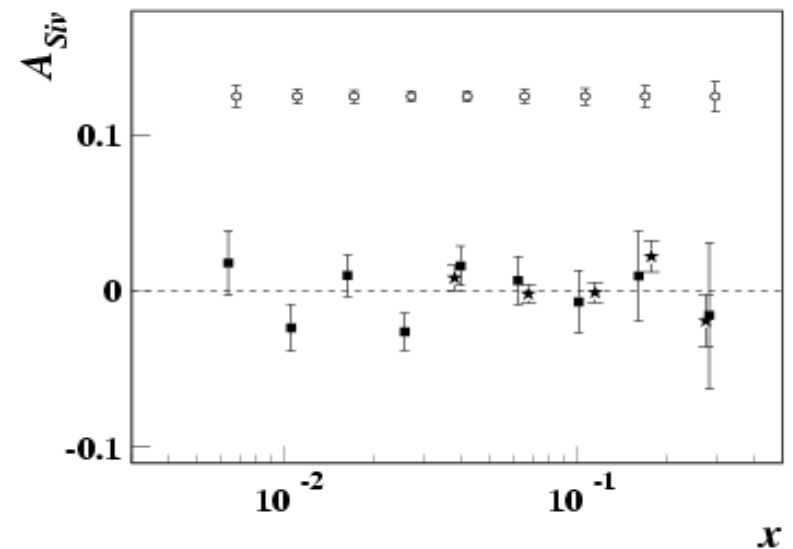
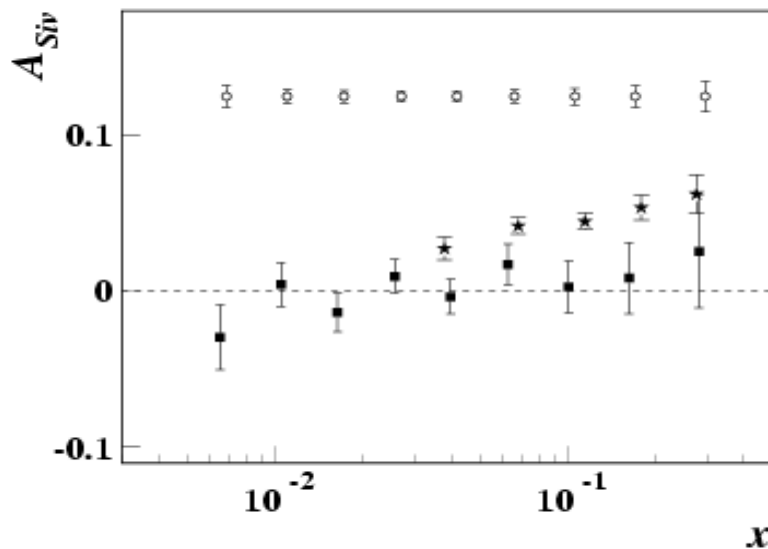
- Next year

- Data taking with transversely polarized proton foreseen
→ new precise results

□ compass 2010
projected errors

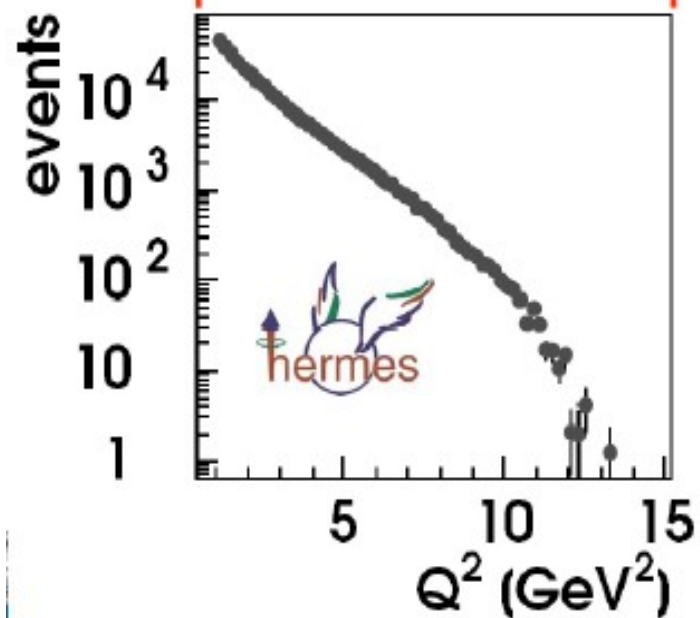
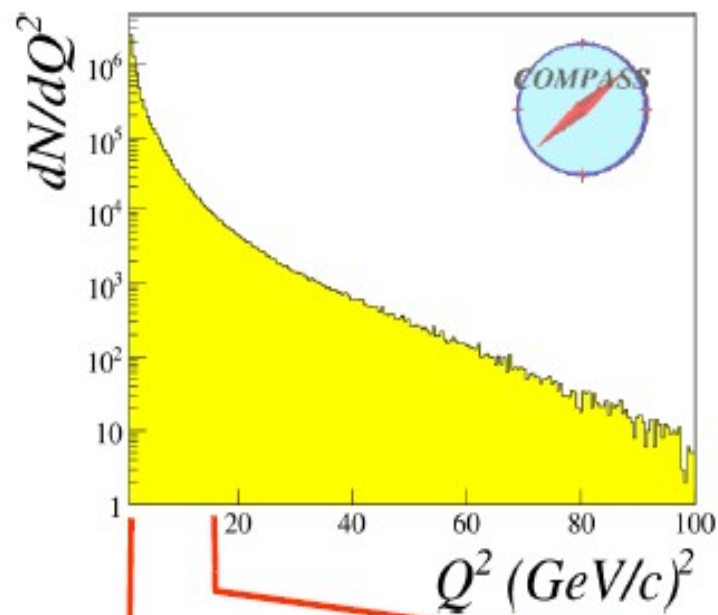
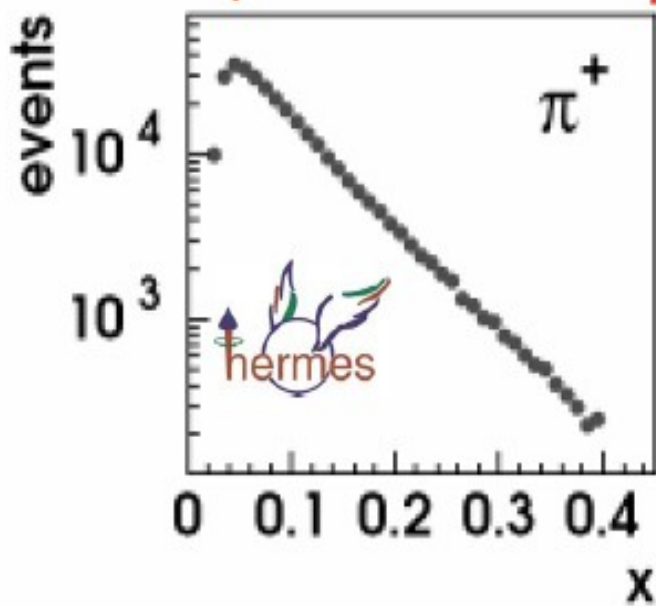
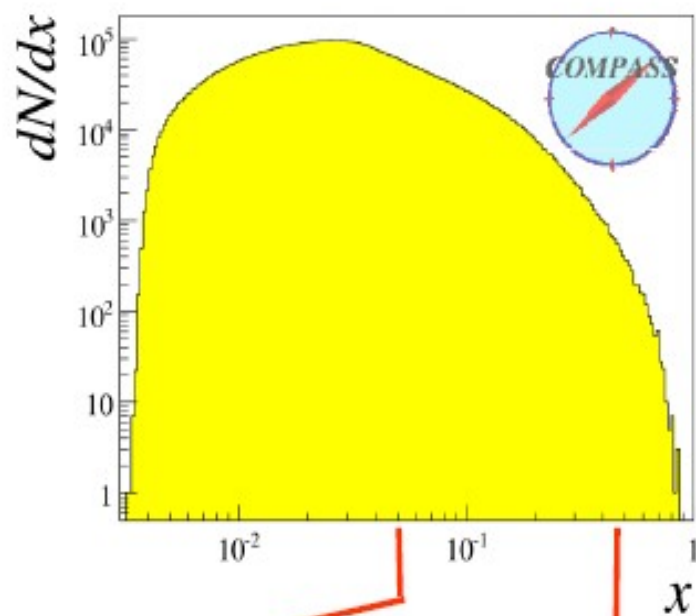
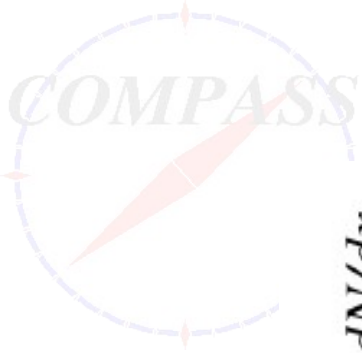
● compass 2007

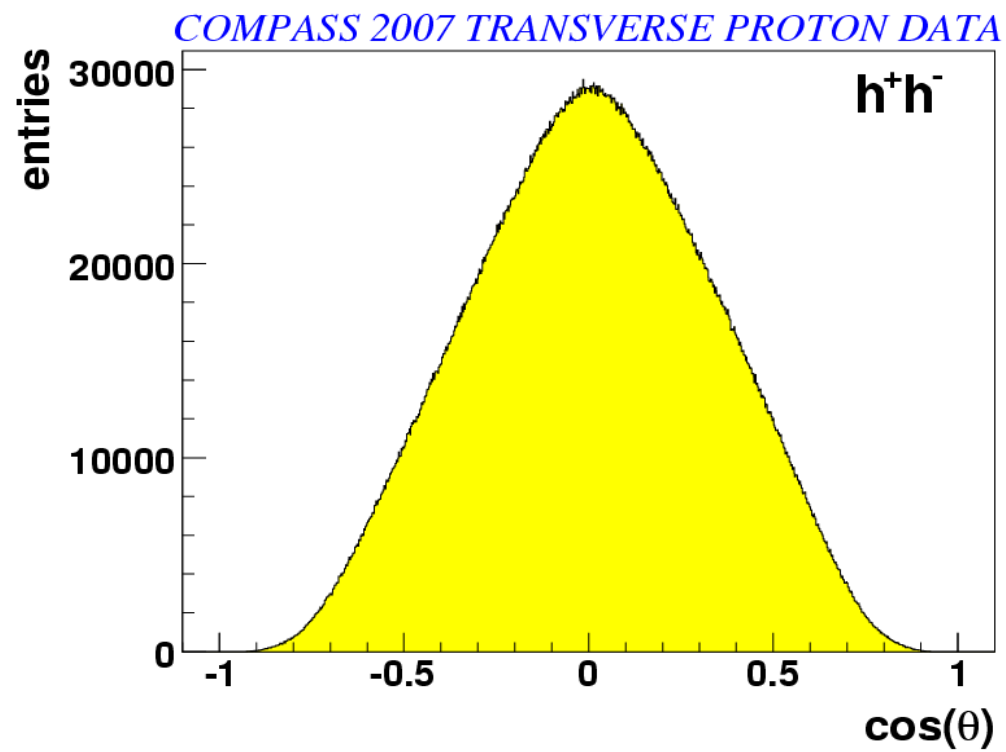
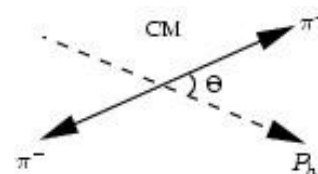
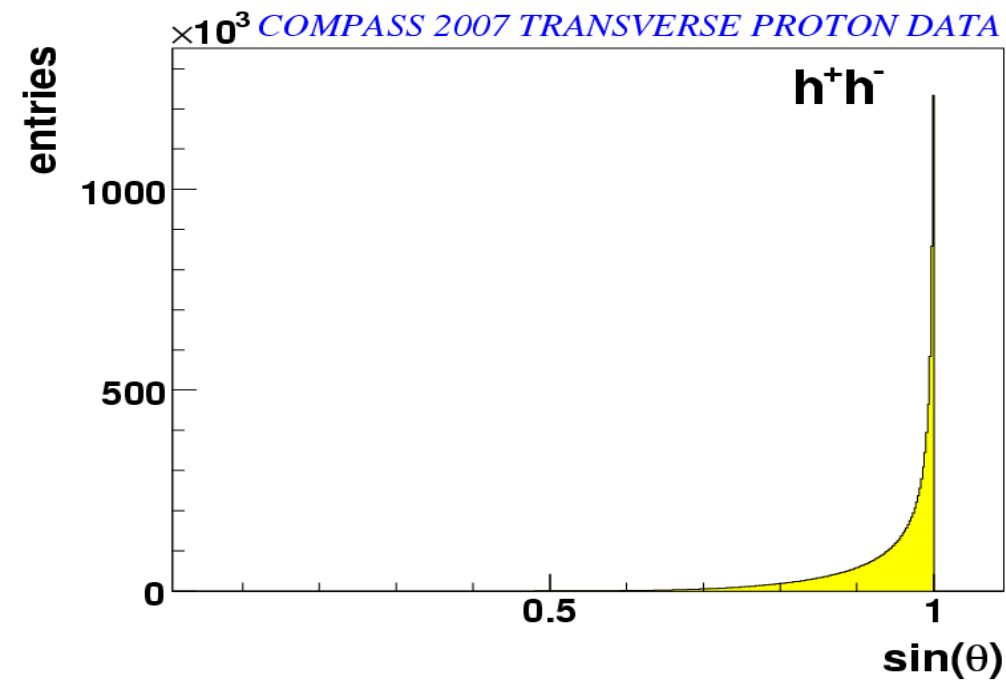
★ hermes





BACKUP

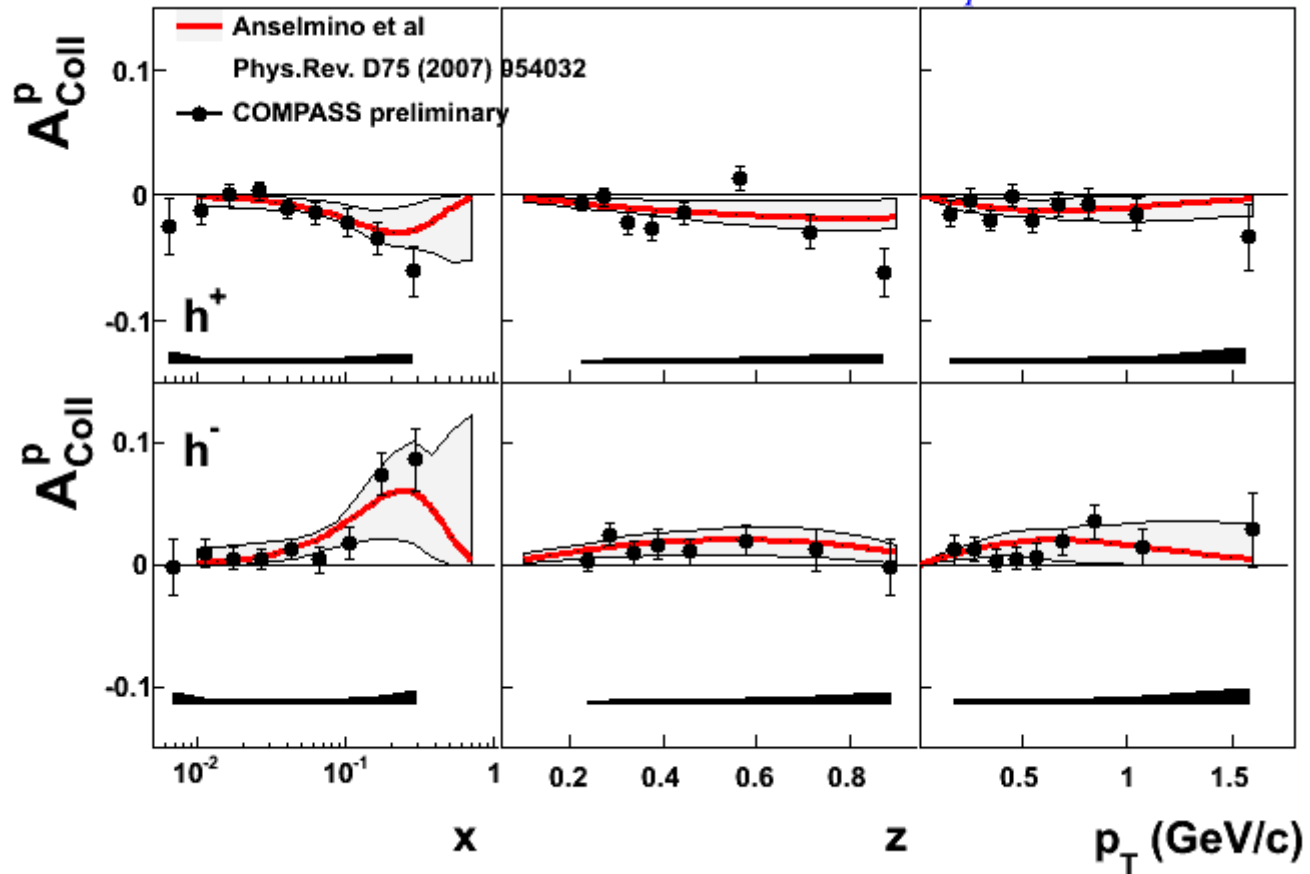






prediction by Anselmino group in Phys. Rev. D

COMPASS 2007 proton data



Anselmino et al. Phys.Rev. D75 (2007) 054032