

Gluon polarisation in the nucleon from COMPASS



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in High-Energy Physics
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COmmon Muon Proton Apparatus for Structure and Spectroscopy COMPASS

Bielefeld, Bochum, Bonn, Burdwan/Calcutta, CERN, Dubna, Erlangen, Freiburg, Lisbon, Mainz, Moscow, Munich, Nagoya, Prague, Protvino, Saclay, Tel Aviv, Turin, Trieste, Warsaw, **~240 physicists**

■ Muon beam program:
gluon polarisation,
polarised quark distributions,
polarised fragmentation functions,
transversity,
Lambda polarisation,
vector meson production,
DVCS (future)

■ Hadron beam program:
Primakoff reaction,
glueballs,
charmed baryons,
exotic charm states.

- longitudinally polarised muon beam
- longitudinally or transversely polarised deuteron (${}^6\text{LiD}$) target
- momentum and calorimetry measurement
- particle identification

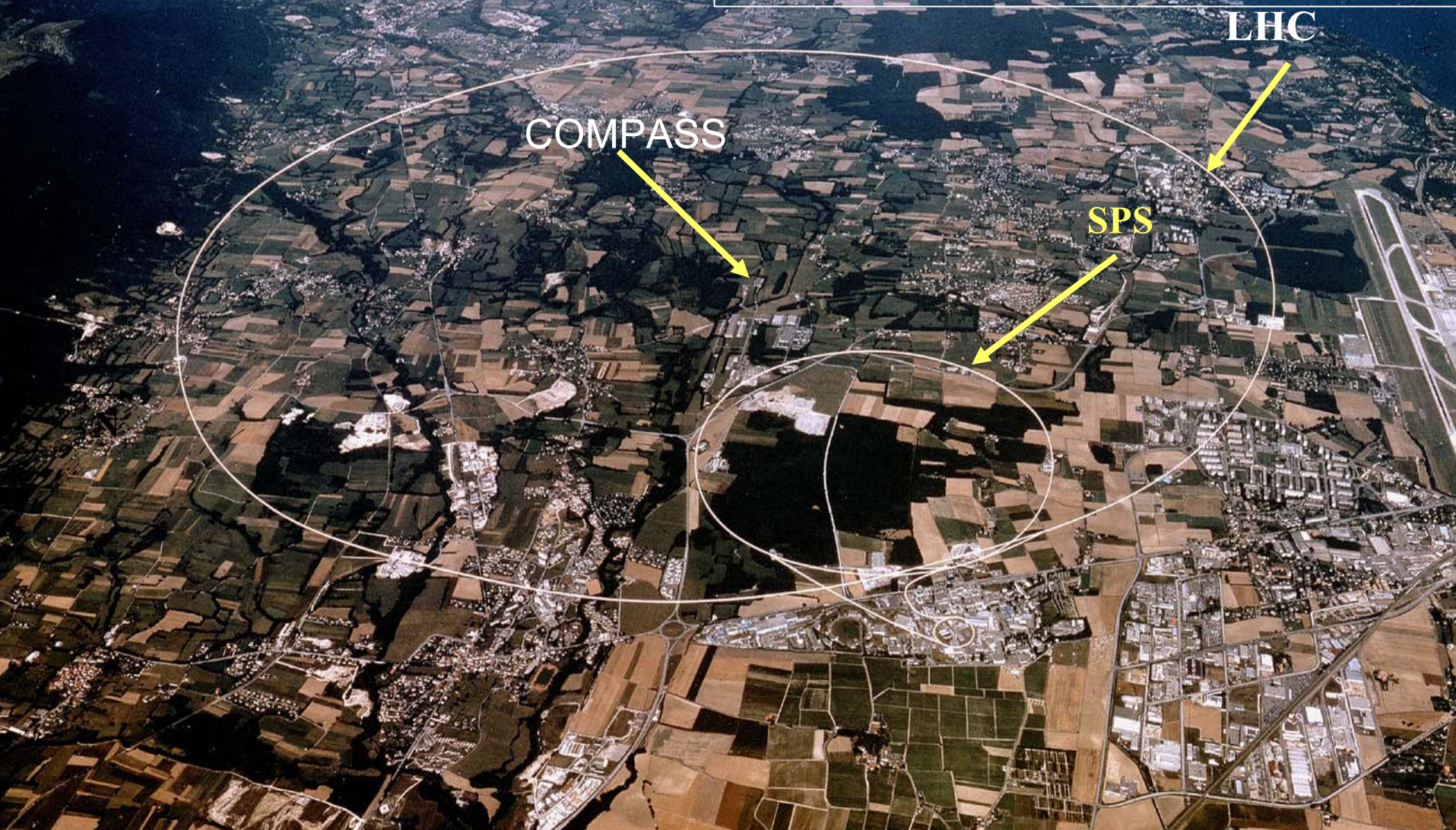
luminosity: $\sim 5 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$

beam intensity: $2 \cdot 10^8 \mu^+/\text{spill}$ (4.8s/16.2s)

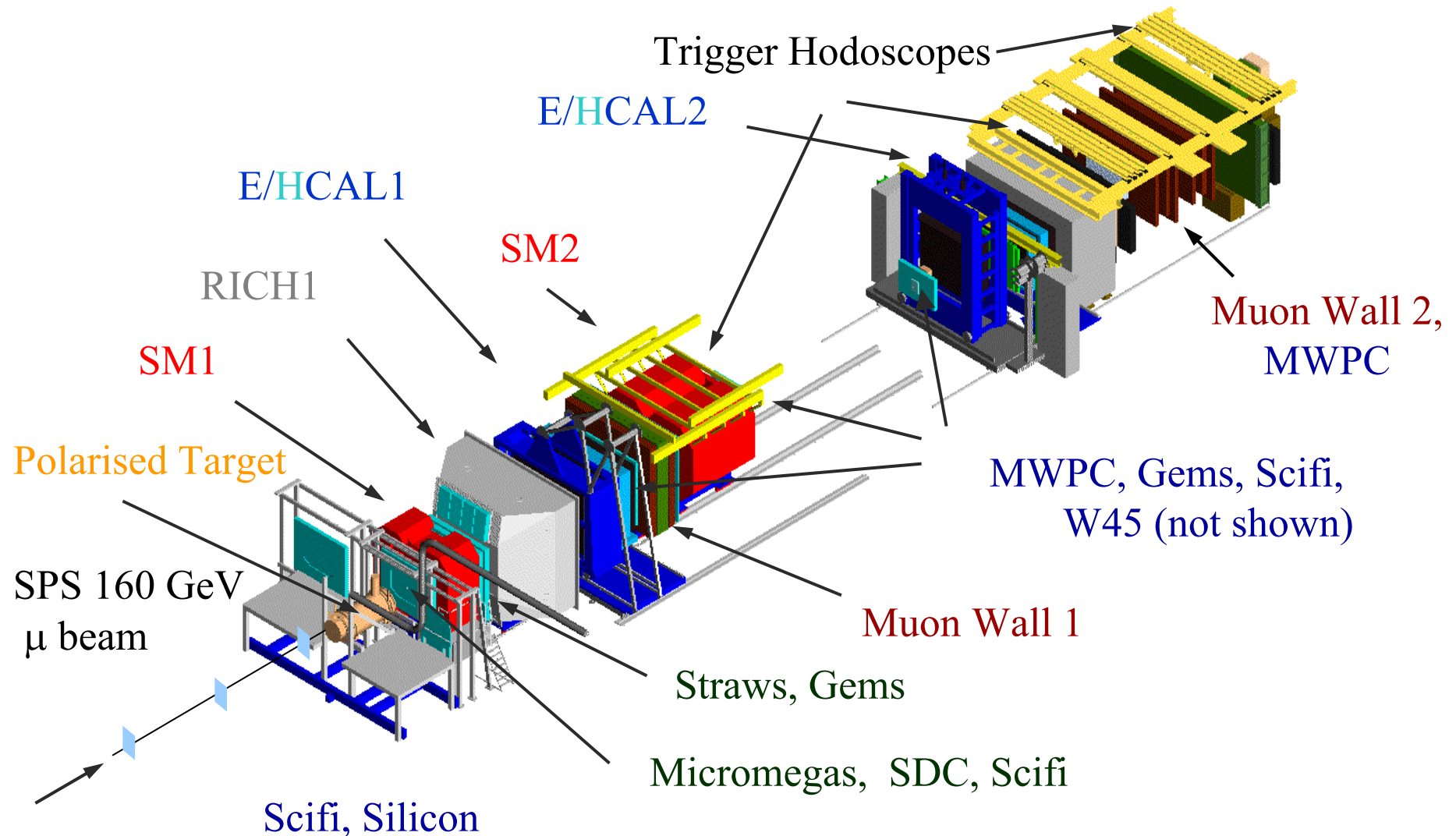
beam momentum: 160 GeV/c

beam polarization: $\sim 76 \%$

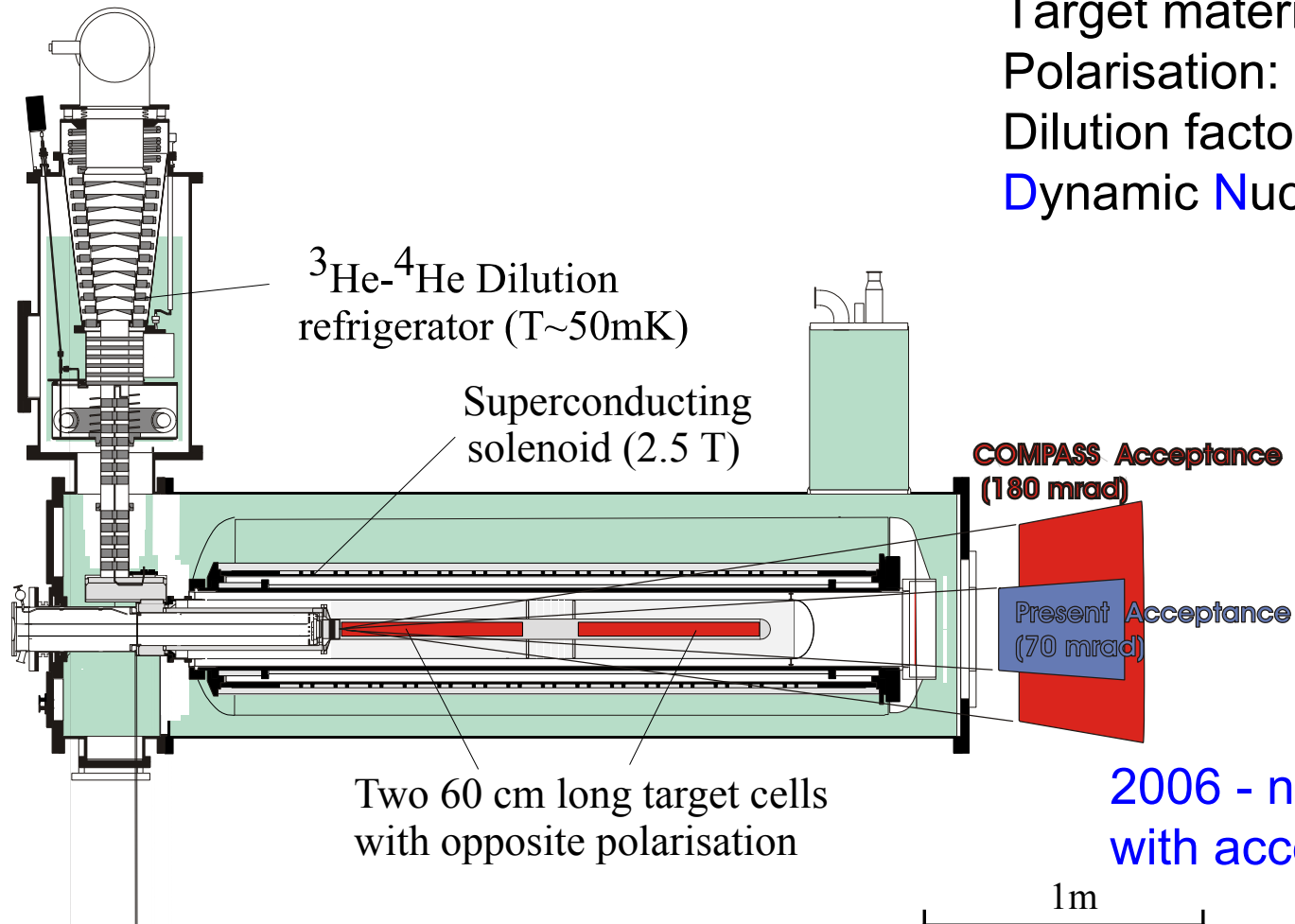
target polarization: $\sim 50 \%$



The COMPASS Spectrometer



The COMPASS polarised target



Target material: ^6LiD

Polarisation: $>50\%$

Dilution factor: ~ 0.4

Dynamic Nuclear Polarisation

2006 - new solenoid
with acceptance 180 mrad

Content

- Motivation and Nucleon spin decomposition.
- Inclusive asymmetry A_1^d , structure function g_1^d and QCD analysis for $Q^2 > 1 \text{ GeV}^2$ (fits).

Gluon polarisation $\frac{\Delta G}{G}$

Two methods of accessing directly $\frac{\Delta G}{G}$ in Compass.
Open charm channel method and results.

High p_T hadron pairs method:

- Results for events with low Q^2 .
- Results for events with $Q^2 > 1 \text{ GeV}^2$.

Conclusions.

Nucleon spin decomposition



Nucleon spin decomposition

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + \Delta G + L$$

- A very small fraction of the proton spin is carried by the spin of the quarks - put the naive but well-accepted quark model into serious questioning!
(EMC (1988): $a_0 = \Delta\Sigma = 12 \pm 9 \pm 14\%$ while $\approx 60\%$ expected, confirmed by SMC, SLAC and Hermes : $\Delta\Sigma = 20 - 30\%$)
- The possible role of axial anomaly:
measured quantity $a_0 = \Delta\Sigma - (3\alpha_s/2\pi) \Delta G$


- EMC – 1988 – 18 years ago: An impressive follow-ups:
SLAC E142,E143,E155,E156, SMC, HERMES, JLAB spin physics, COMPASS, RHIC Spin, JLAB12 GeV Upgrade

Nucleon spin decomposition

$$g_1(x) = \frac{1}{2} \sum e_q^2 \Delta q(x) \quad \text{and} \quad \Delta q(x) = q^+(x) - q^-(x);$$

Well defined in terms of quark helicity densities but:

$$q^+ \sim \Psi (1 + \gamma_5) \gamma_\mu \Psi, \quad q^- \sim \Psi (1 - \gamma_5) \gamma_\mu \Psi \Rightarrow \Delta q(x) \sim \Psi \gamma_5 \gamma_\mu \Psi$$

Axial vector current is not conserved due to
 Adler-Bell-Jackiw anomaly ;

In a consequence: measured quantity $a_0 = \Delta\Sigma - (3\alpha_S/2\pi) \Delta G$.
where $\Delta\Sigma = \Delta u + \Delta d + \Delta s$, $\Delta q = \int \Delta q(x) dx$ and
 $\Delta G = \int \Delta G(x) dx$ is a gluon polarization contribution.

Nucleon spin decomposition

$$\Gamma_1 = \int g_1(x) dx$$

$$\Gamma_1^p - \Gamma_1^n = \frac{g_A}{6g_V} C_1^{NS} \quad (\text{Bjorken sum rule})$$

$$\Gamma_1^{p,n} = \left(\pm a_3 + \frac{a_8}{\sqrt{3}} \right) \frac{C_1^{NS}}{12} + a_0 \frac{C_1^S}{9} \quad (\text{Ellis-Jaffe sum rule})$$

$a_3, a_8, g_{A,V}$ - hyperon β decay + $SU_f(3)$;

$C_1^{S,NS}$ - calculable in QCD

But - due to Δ anomaly - $a_0 = \Delta\Sigma - (3\alpha_S/2\pi) \Delta G$ and
if $\Delta G \approx 2.5 \rightarrow \Delta\Sigma \approx 0.6 \rightarrow$ can “solve the spin crisis”



Need direct measurement of ΔG



Inclusive asymmetry A_1^d and
structure function g_1^d

A_1^d and structure function g_1^d

$$A^{\mu d} = A_{||} = \frac{\sigma^{\uparrow\downarrow} - \sigma^{\uparrow\uparrow}}{\sigma^{\uparrow\downarrow} + \sigma^{\uparrow\uparrow}} = D (A_1^d + \eta A_2^d)$$

$$|\eta A_2^{d,p,n}| \ll |A_1^{d,p,n}|,$$

$$A_1^{p,n} = A^{\gamma N} = \frac{\sigma^{1/2} - \sigma^{3/2}}{\sigma^{1/2} + \sigma^{3/2}} \quad \text{for nucleon}$$

$$A_1^d = A^{\gamma d} = \frac{\sigma^0 - \sigma^2}{\sigma^0 + \sigma^2} \quad \text{for deuteron}$$

$$A_{\text{meas}} \sim A^{\mu d} \sim A_1^d;$$

Measurement of A_1 gives access to g_1 structure function

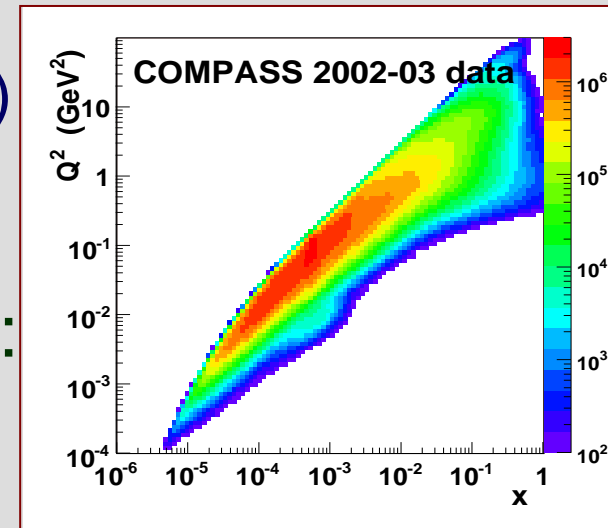
$$g_1^d = \frac{1}{2} (g_1^p + g_1^n) \left(1 - \frac{3}{2} \omega_d\right) \simeq A_1^d F_1^d = A_1^d \frac{F_2^d}{2x(1+R)}$$

A_1^d and structure function g_1^d

- A_1^d and g_1^d for small Q^2 ($Q^2 < 1 \text{ GeV}^2$):
physics at small x , parton saturation,
non-perturbative models (Regge, VDM)
poorly known (only SMC data)
- A_1^d and g_1^d for high Q^2 ($Q^2 > 1 \text{ GeV}^2$):
QCD analysis possible: ΔG estimation

A_1^d and structure function g_1^d

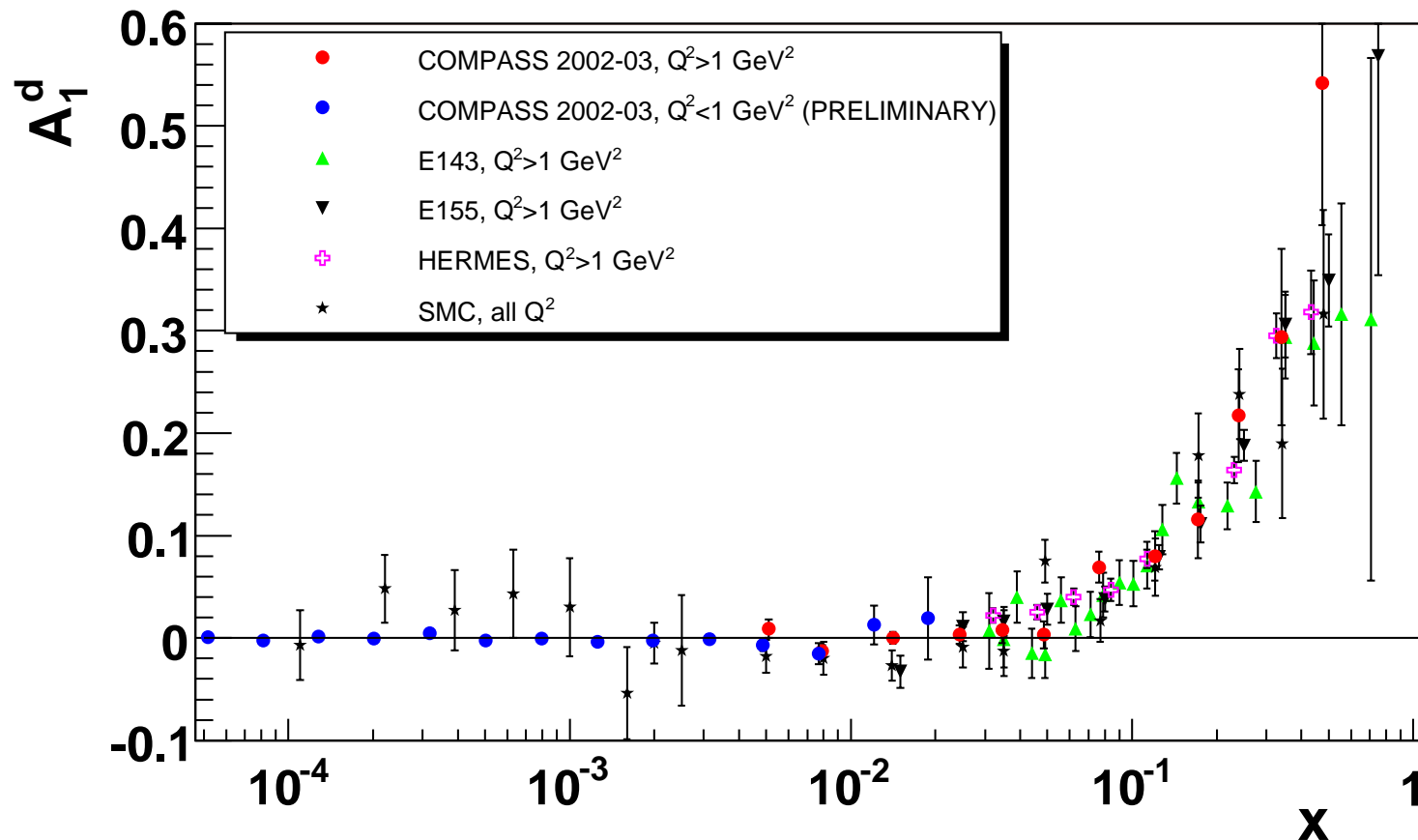
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- A_1^d and g_1^d for high Q^2 ($Q^2 > 1 \text{ GeV}^2$):
QCD analysis possible: ΔG estimation



Unfortunately - Q^2 and x are strongly correlated for
small Q^2 in COMPASS (fixed target)

A_1^d and structure function g_1^d

blue points – Compass 2002-2003 data for $Q^2 < 1 \text{ GeV}^2$
10-20 times lower statistical errors compared to SMC



A_1^d and structure function g_1^d

g_1^d - Compass 2002-2003
data for $Q^2 < 1 \text{ GeV}^2$

F_2^d taken from SMC param.

R depends on x: $x > 0.12$

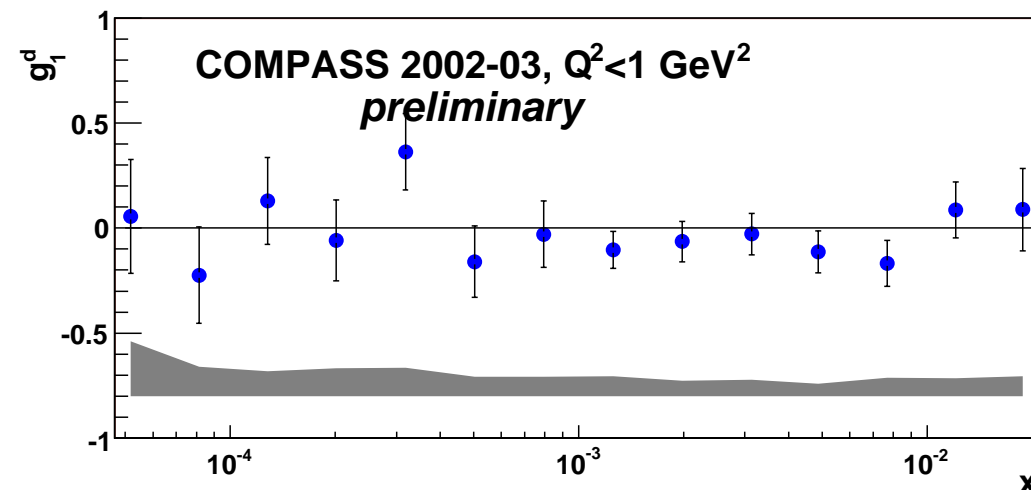
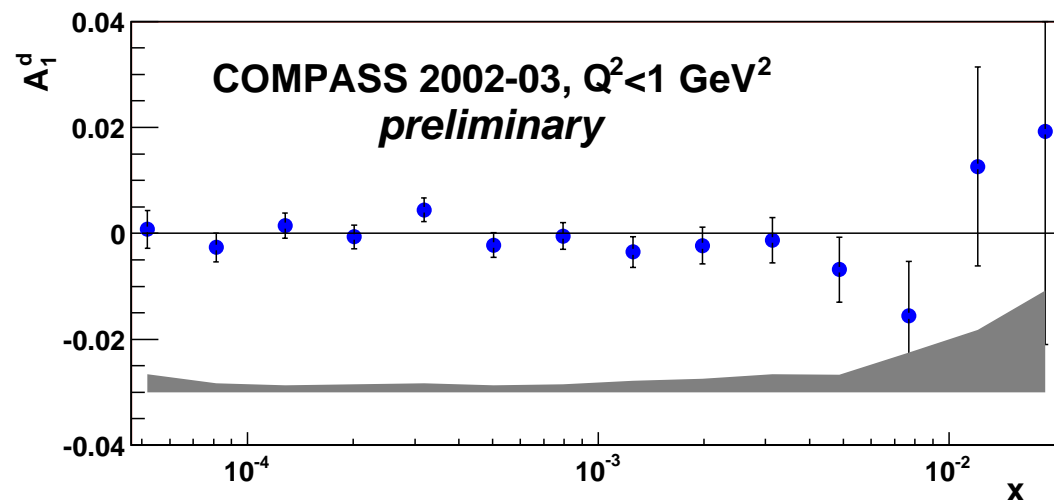
SLAC (Phys.Lett.B250(1990)193,
B52(1999)194)

$0.003 < x < 0.12$ NMC

(unpublished)

$x < 0.003$ ZEUS

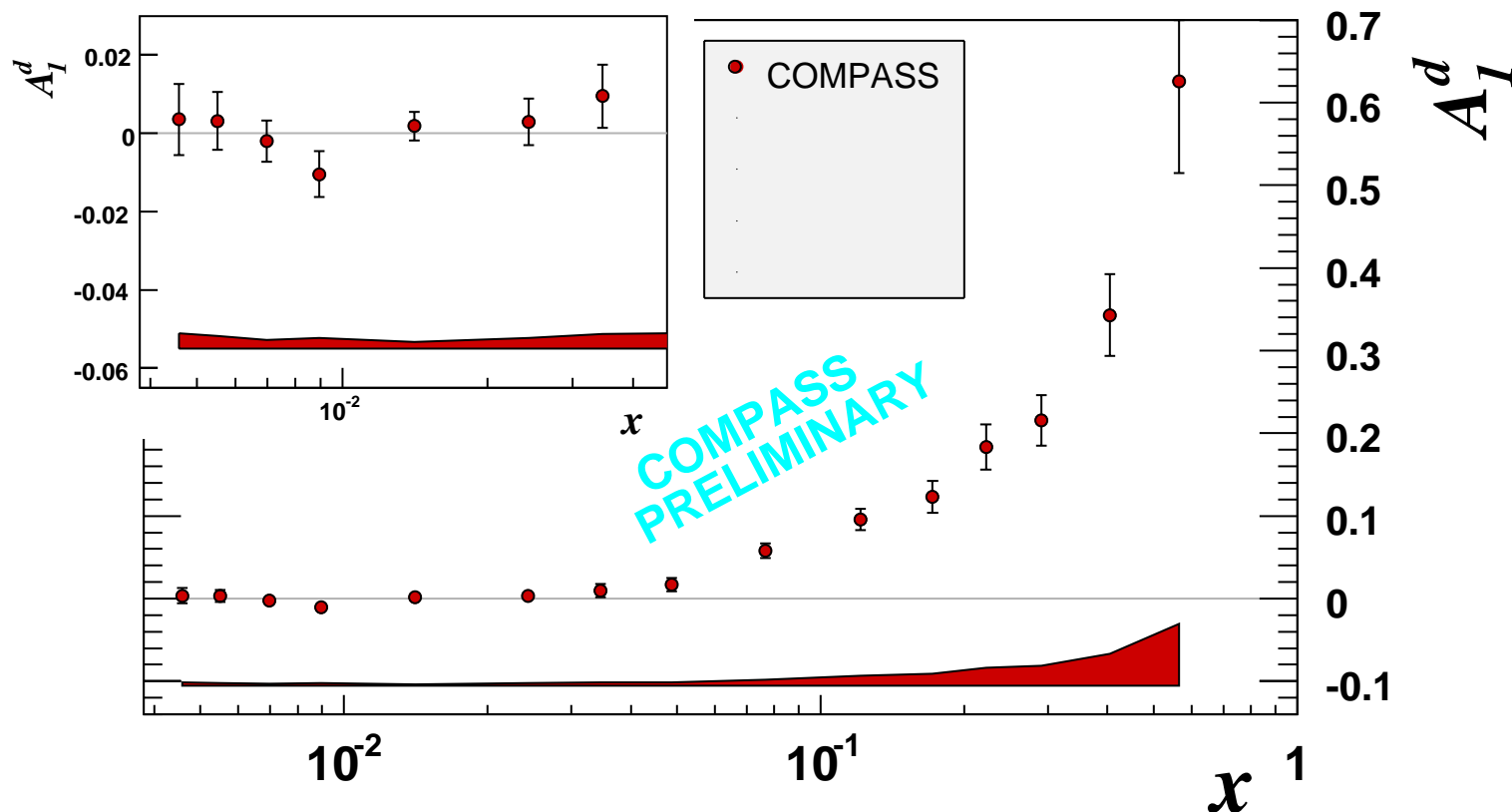
(Eur.Phys.JC7(1999)609, σ_L , σ_T
cross sections param.)



A_1^d and structure function g_1^d

A_1^d - Compass 2002-2004 data for $Q^2 > 1 \text{ GeV}^2$

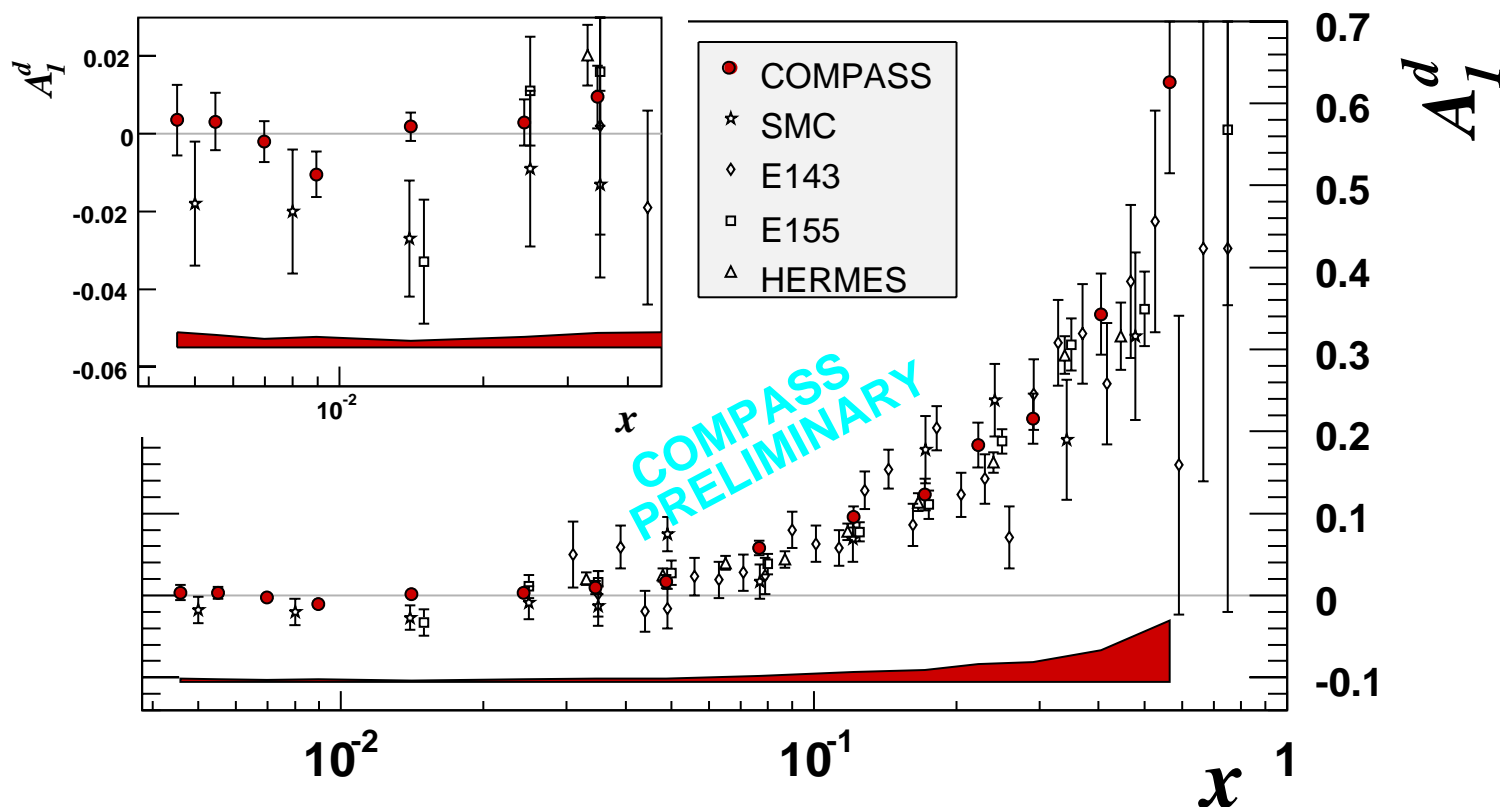
~90 mln events (60 inclusive ones), Q^2 depends on x bins



A_1^d and structure function g_1^d

A_1^d - Compass 2002-2004 data for $Q^2 > 1 \text{ GeV}^2$

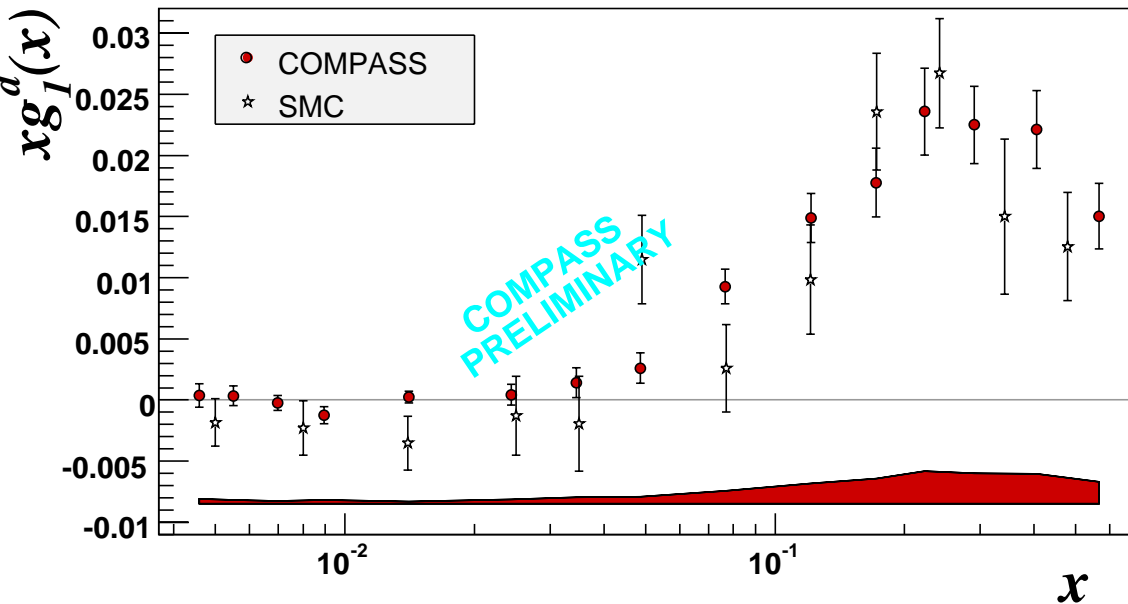
For $x < 0.03$: statistical errors reduced by factor 4;
no tendency toward negative values



A_1^d and structure function g_1^d

g_1^d - Compass 2002-2004 data for $Q^2 > 1 \text{ GeV}^2$

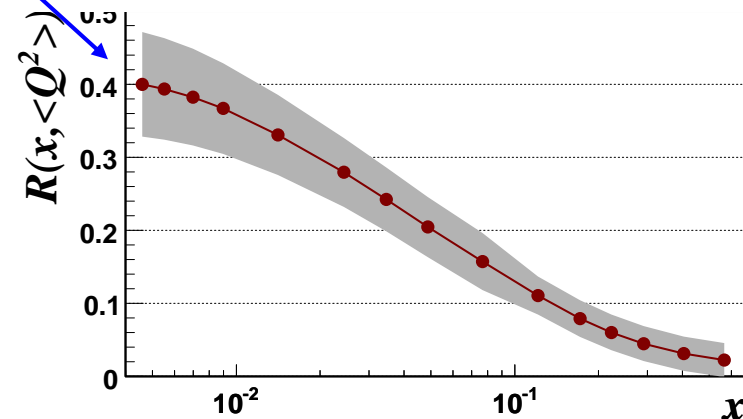
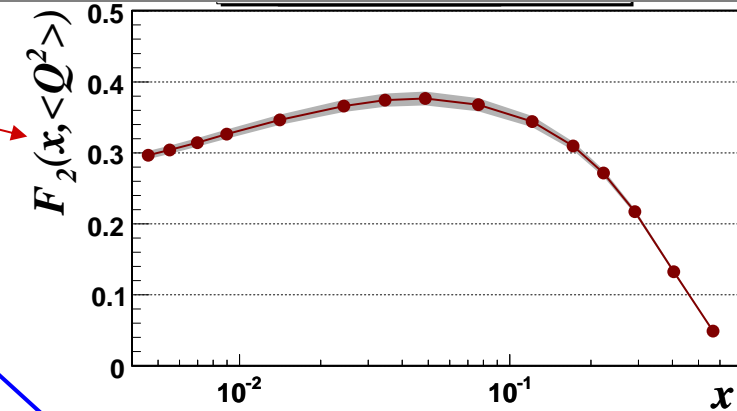
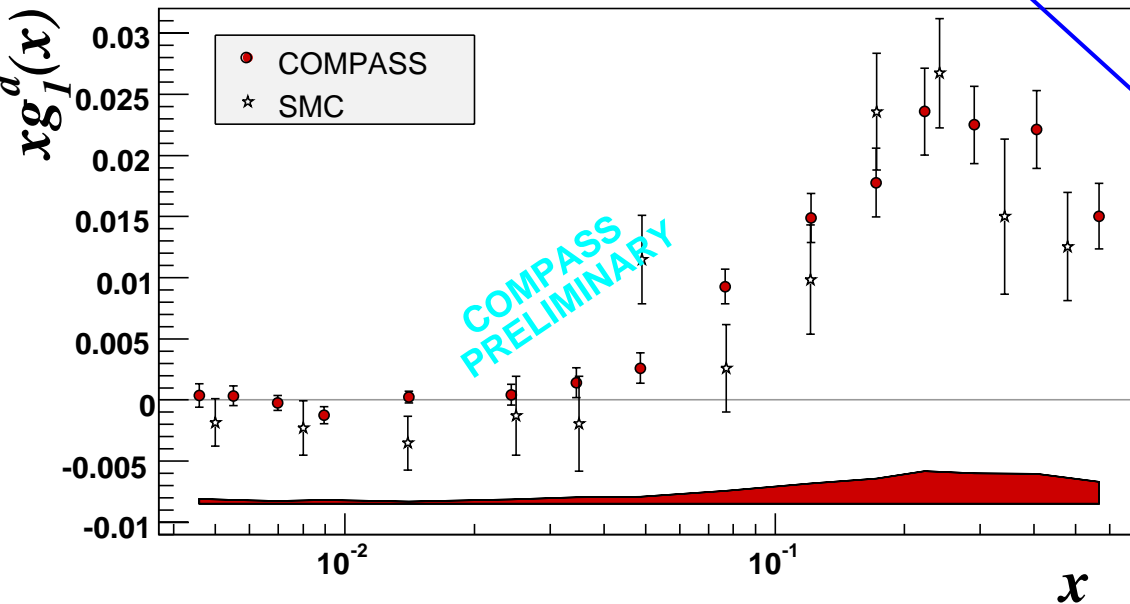
$$g_1^d = g_1^N \left(1 - \frac{3}{2} \omega_d\right) = \frac{F_2^d}{2x(1+R)} A_1^d$$



A_1^d and structure function g_1^d

g_1^d - Compass 2002-2004 data for $Q^2 > 1 \text{ GeV}^2$

$$g_1^d = g_1^N \left(1 - \frac{3}{2} \omega_d\right) = \frac{F_2^d}{2x(1-R)} A_1^d$$



QCD fits

- Measured structure function $g_1^{p,n,d}$ (different x and Q^2)

$$g_1(x, Q^2) = \frac{1}{2} \langle e^2 \rangle \left[C_q^S \otimes \Delta\Sigma^{NS} + C_q^{NS} \otimes \Delta q^{NS} + 2n_f C_G \otimes \Delta G \right]$$

- DGLAP equations:

$$t = \log\left(\frac{Q^2}{\Lambda^2}\right)$$

$$\left\{ \begin{array}{l} \frac{d}{dt} \Delta q^{NS} = \frac{\alpha_s(t)}{2\pi} P_{qq}^{NS} \otimes \Delta q^{NS} \\ \frac{d}{dt} \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix} = \frac{\alpha_s(t)}{2\pi} \begin{pmatrix} P_{qq}^S & 2n_f P_{qG}^S \\ P_{Gq}^S & P_{GG}^S \end{pmatrix} \otimes \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix} \end{array} \right.$$

- Initial parametrization:
x dependence at fixed Q^2

$$(\Delta\Sigma, \Delta q_s, \Delta q_8, \Delta G) = \eta \frac{x^\alpha (1-x)^\beta (1+\gamma x)}{\int_0^1 x^\alpha (1-x)^\beta (1+\gamma x) dx}$$

- Minimization routine

$$\chi^2 = \sum_{i=1}^N \frac{[g_1^{calc}(x, Q^2) - g_1^{exp}(x, Q^2)]^2}{[\sigma_{stat}^{exp}(x, Q^2)]^2}$$

QCD fits

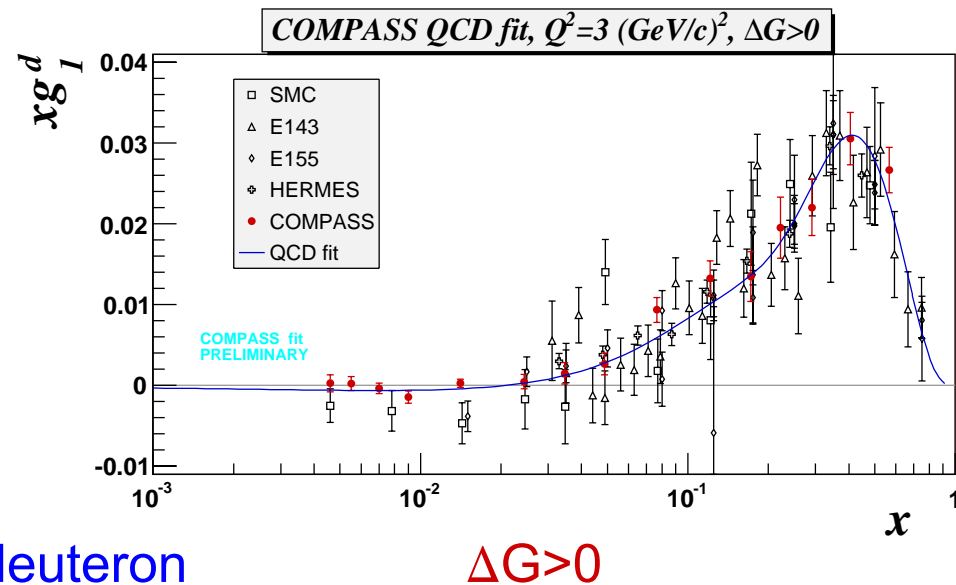
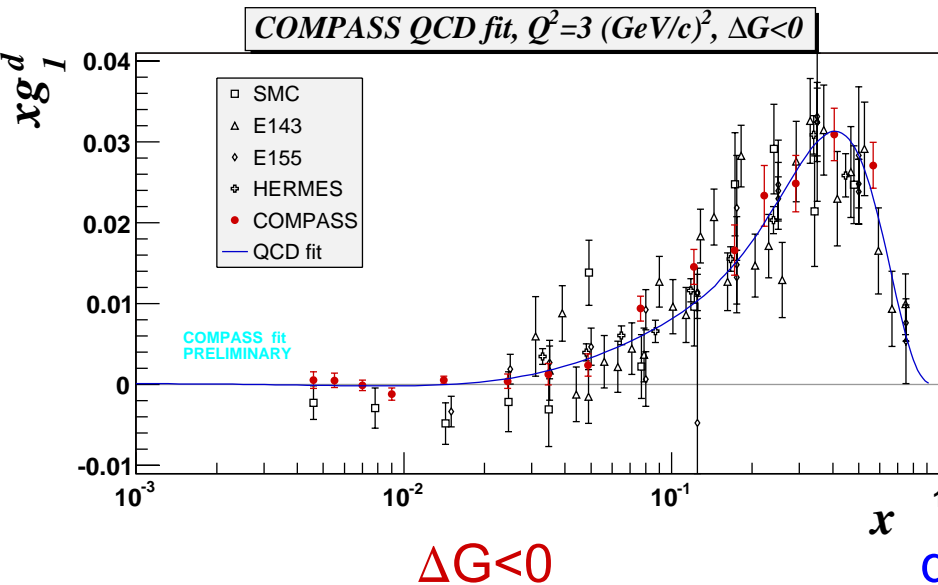
Two different codes in NLO $\overline{\text{MS}}$ scheme used:

- grid in (Q^2, x) space (Phys.Rev.D58(1998)112002)
- Mellin transform + moments space (Phys.Rev.D70(2004)074032)

World data fit: 9 experiments, 230 points

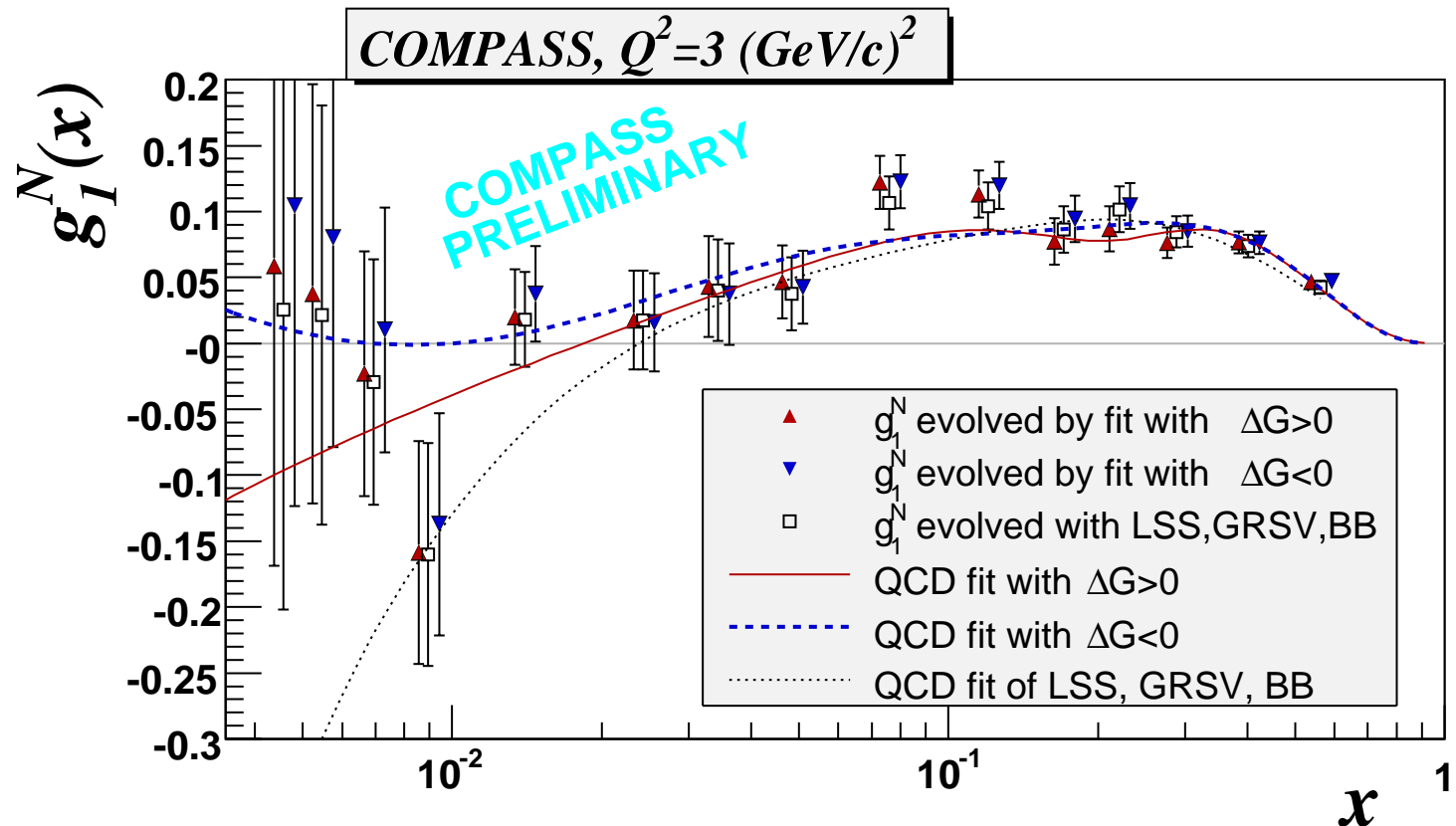
Two solutions describe data equally well: $\Delta G > 0$ and $\Delta G < 0$, $Q^2 = 3 \text{ GeV}^2$,

$$g_1(x, Q_0^2) = g_1(x, Q_i^2) + \left[g_1^{fit}(x, Q_0^2) - g_1^{fit}(x, Q_i^2) \right]$$

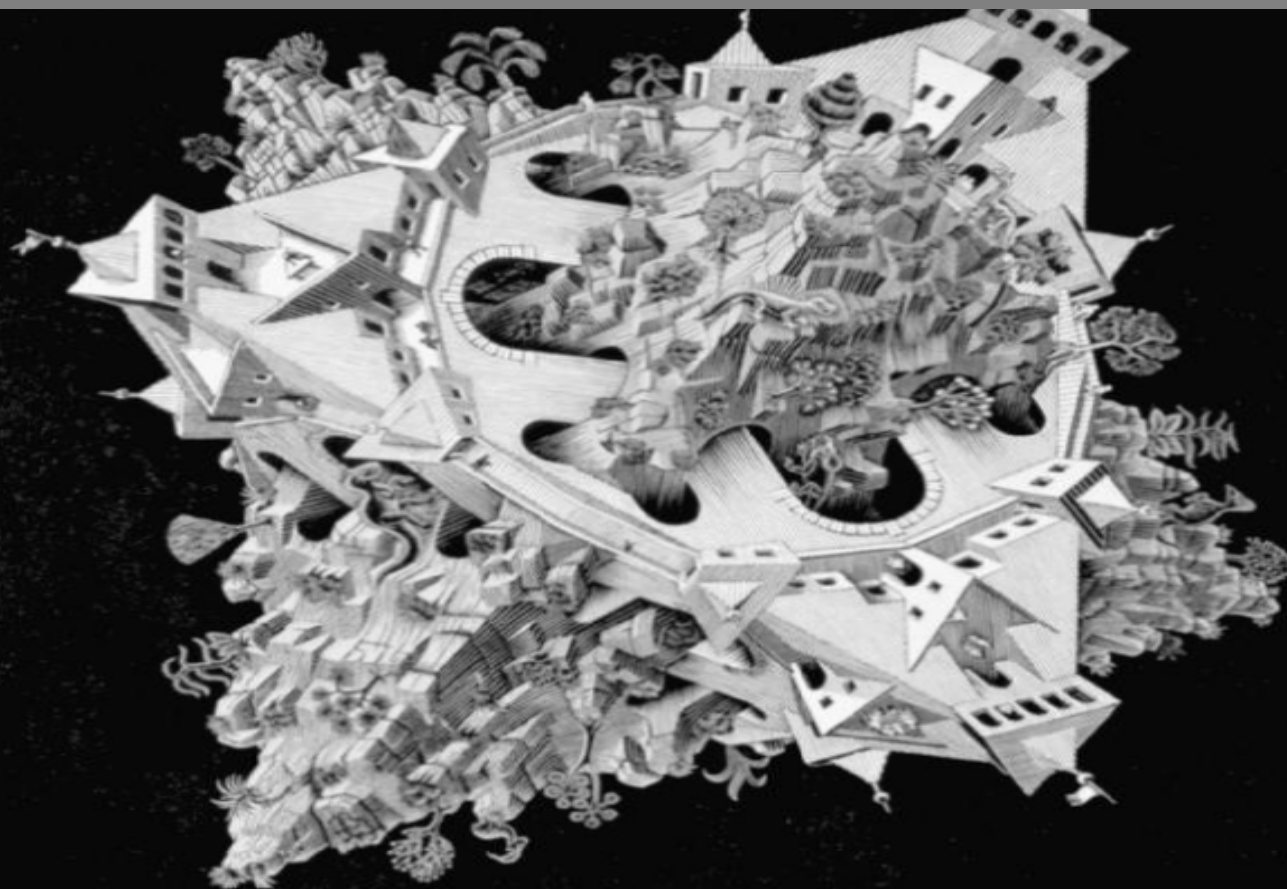


QCD fits

Comparison of fits - disagreement of data with previous QCD fits (LSS05, BB, GRSV)



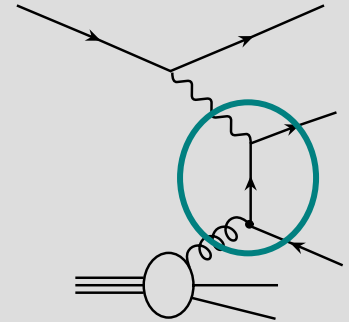
Direct measurements of $\Delta G/G$



Direct measurements of $\Delta G/G$

Photon Gluon Fusion (PGF) probes gluons

$$A_{||} = R_{PGF} a_{PGF}^{LL} \frac{\Delta G}{G} + A_{Bkg}$$



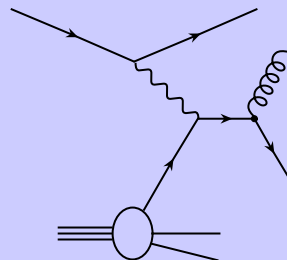
- Open charm „golden channel”
 - ☺ no background asymmetry, less MC dependent.
 - ☹ small statistics, NLO corrections can be important
- 2 high p_T hadrons ($p_T > 0.7\text{GeV}$, then selection on Σp_T^2)
 - ☺ Large statistics
 - ☹ physical background: „model” (MC) dependent, requires very good description of data by MC.

Direct measurements of $\Delta G/G$

2 high p_T hadrons:

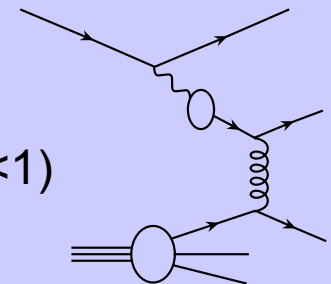
- Low Q^2 analysis ($Q^2 < 1 \text{ GeV}^2$): perturbative scale fixed by p_T , complicated physical background e.g. resolved γ , low p_T
- High Q^2 analysis ($Q^2 > 1 \text{ GeV}^2$): scale Q^2 , physical background better controlled in the frame of pQCD.

Physical background:

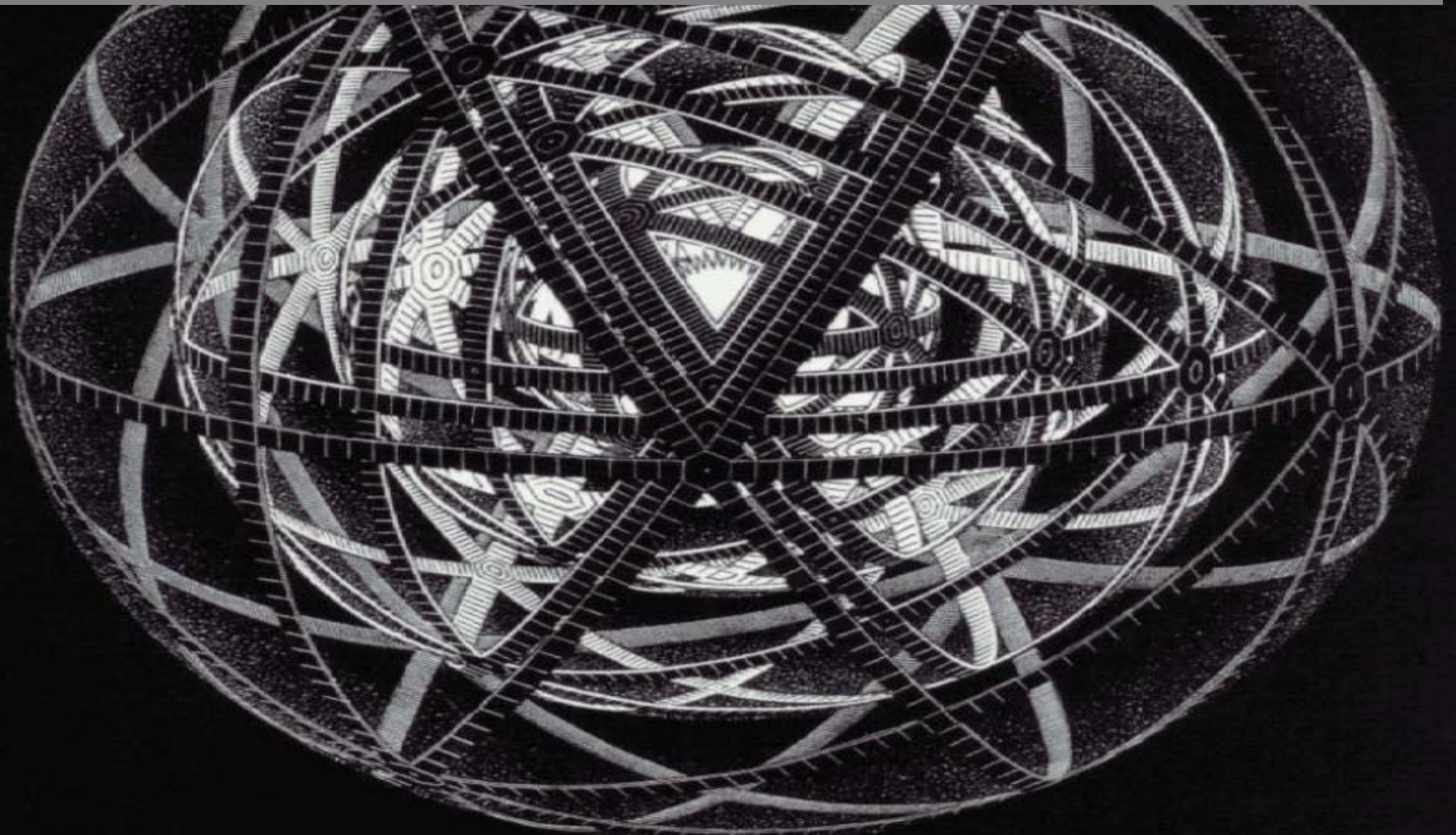


QCD-Compton

resolved γ ($Q^2 < 1$)



$\Delta G/G$ from open charm channel



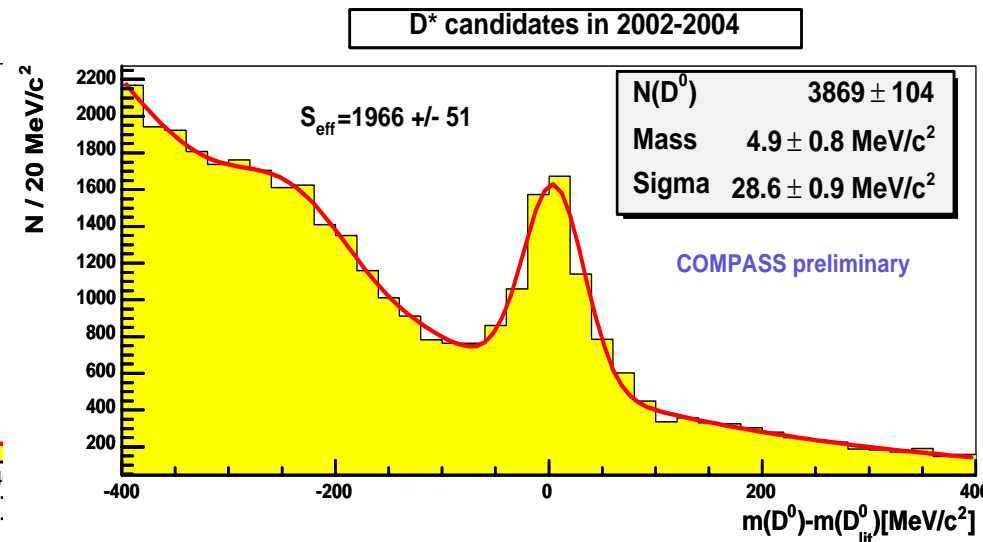
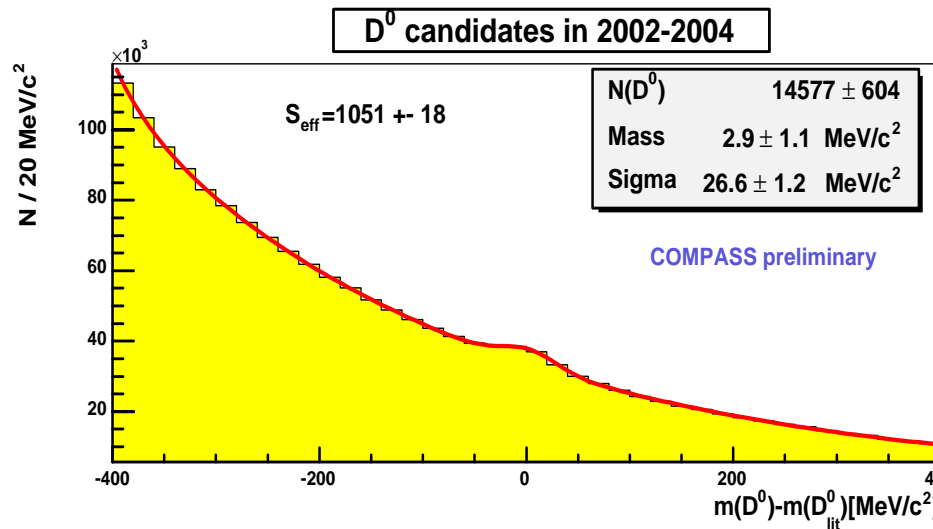
$\Delta G/G$ from open charm channel

$$A_{LL} = \frac{S}{S+B} a_{LL} \frac{\Delta G}{G}(x_g)$$

Scale: $\langle Q^2 \rangle \approx 13 \text{ GeV}^2$
 $\sim 4^* m_c^2$

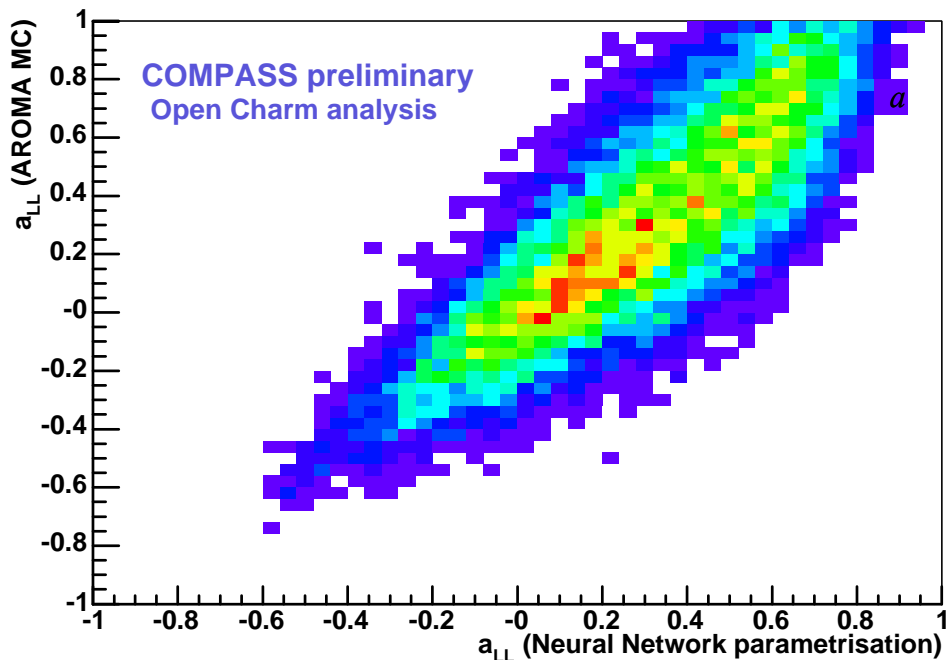
$D^0 \rightarrow K^+ \pi$ (untagged)

$D^* \rightarrow D^0 + \pi \rightarrow K^+ \pi^+ \pi$



$\Delta G/G$ from open charm channel

$$A_{LL} = \frac{S}{S+B} a_{LL} \frac{\Delta G}{G}(x_g)$$



a_{LL} – calculated with help of MC and parametrized by measured quantities (Neural Network used)

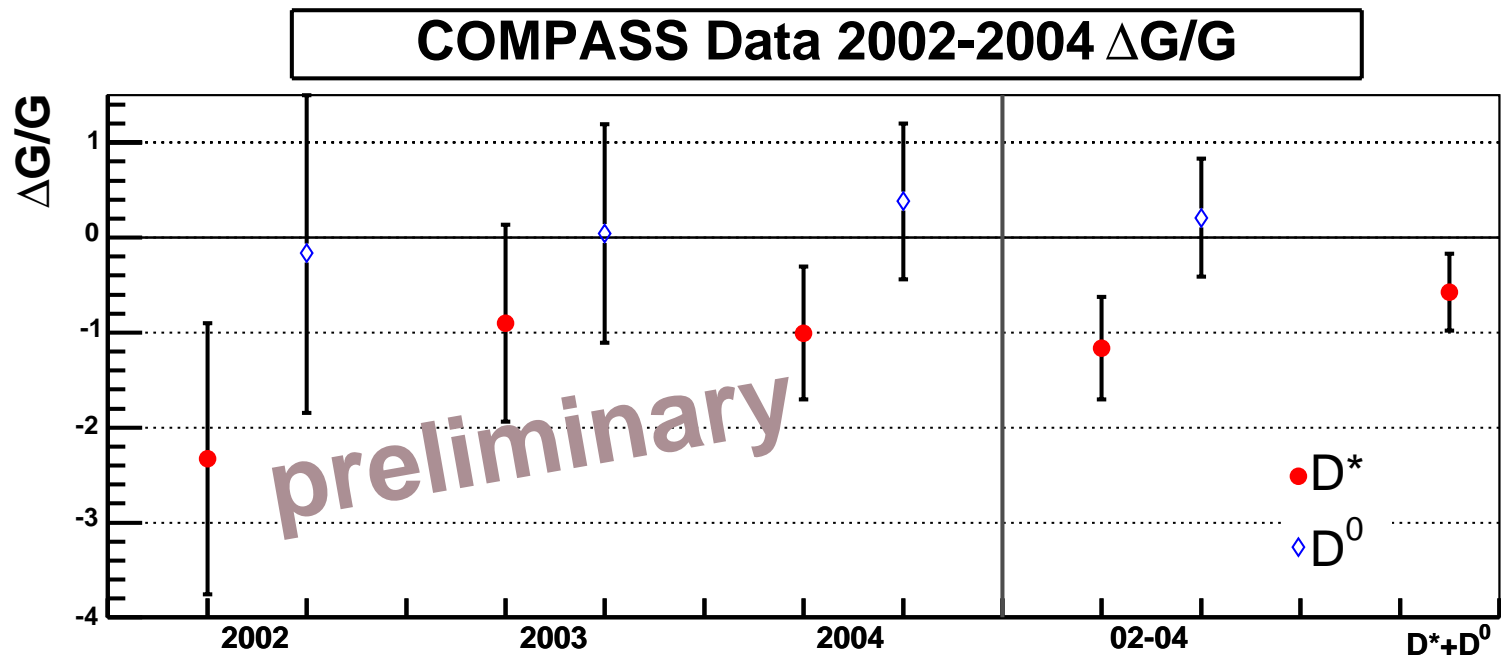
$\Delta G/G$ from open charm channel

$D^0 + D^*$ result 2002 – 2004:

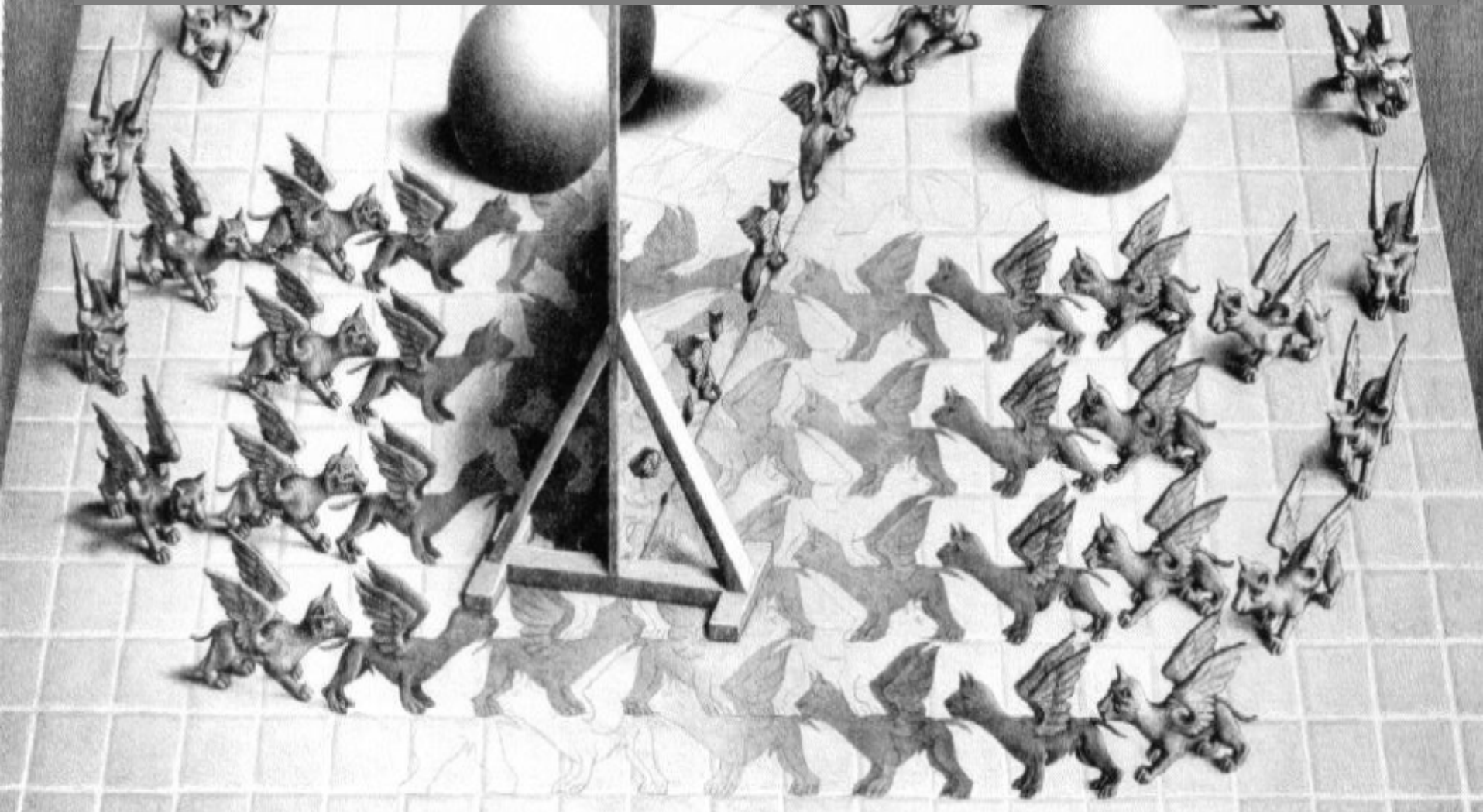
$$\frac{\Delta G}{G} = -0.57 \pm 0.41(\text{stat.})$$

The studies on the systematical uncertainty are ongoing

@ $x_g \approx 0.15$,
scale $\approx 13 \text{ GeV}^2$



$\Delta G/G$ from 2 high p_T hadrons



$\Delta G/G$ from 2 high p_T hadrons (low Q^2)

Low Q^2 : $Q^2 < 1 \text{ GeV}^2$

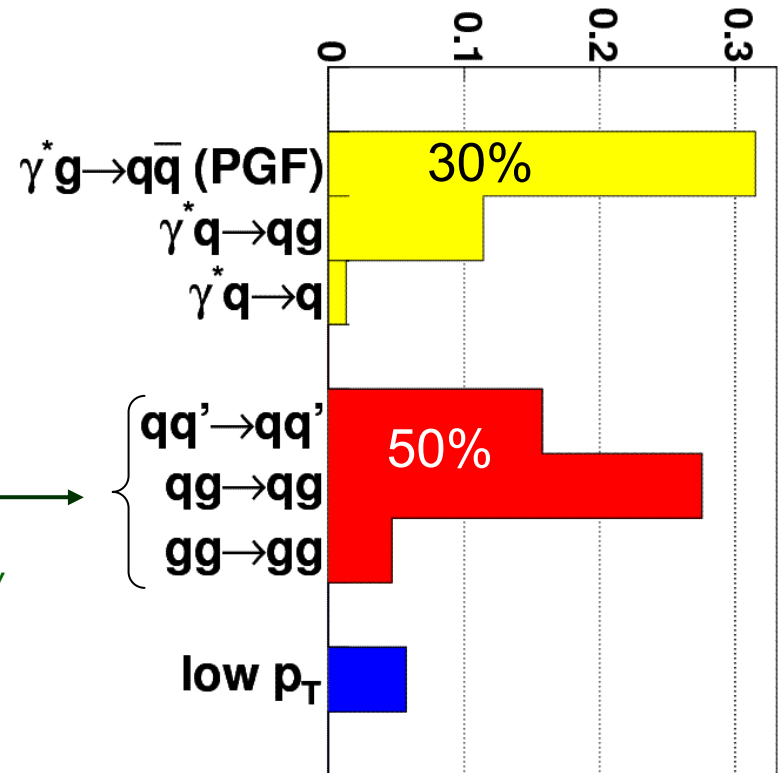
$$A_{LL}/D = R_{pgf} \Delta G/G a_{LL}^{pgf} + R_{qcdc} \Delta q/q a_{LL}^{qcdc}$$

Resolved γ {

$$+ R_{qq} \Delta q/q a_{LL}^{gq} (\Delta G/G)^\gamma$$

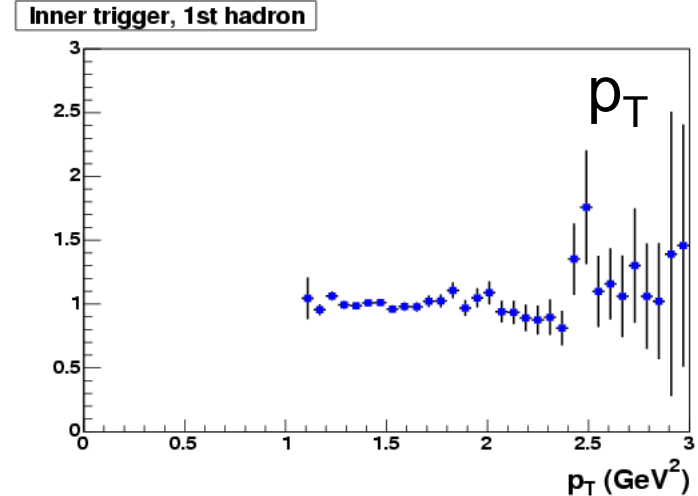
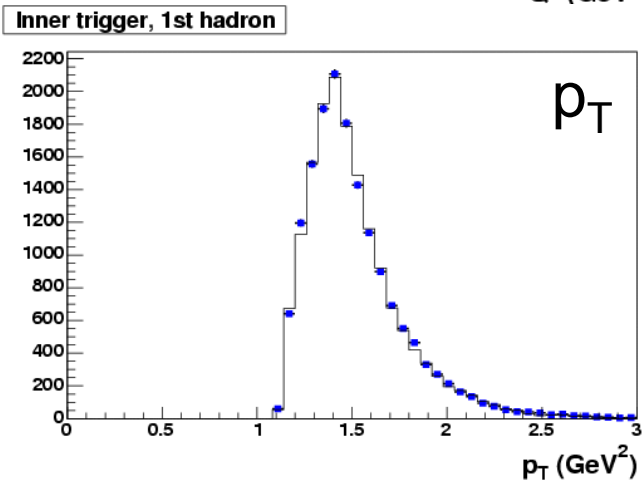
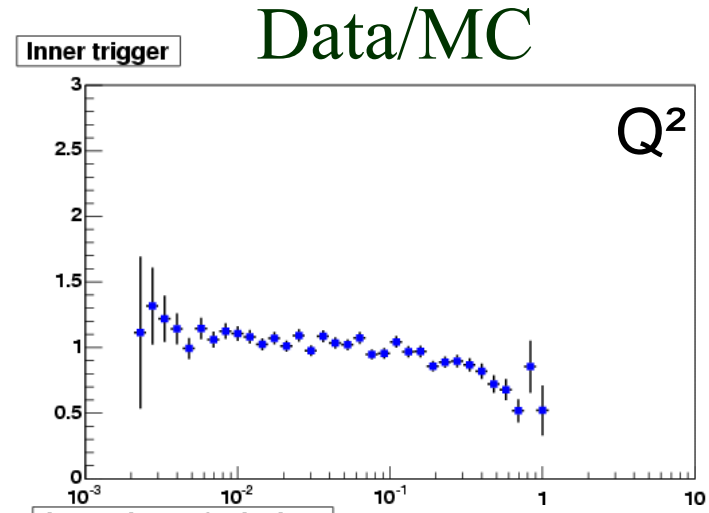
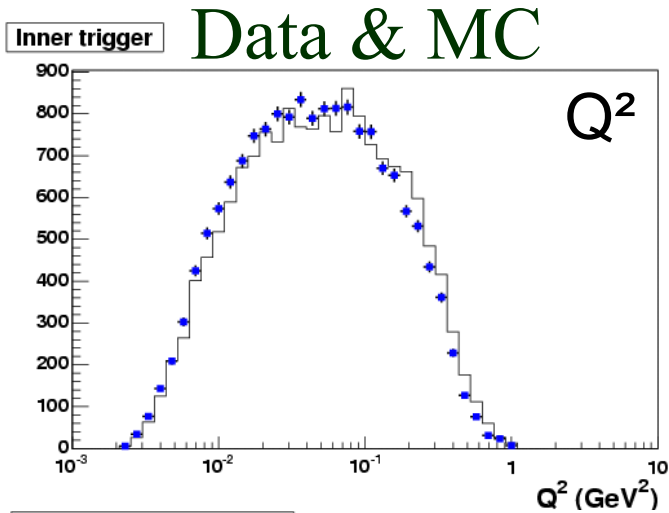
$$+ R_{qg} \Delta G/G a_{LL}^{qg} (\Delta q/q)^\gamma$$

+ ...



MC event generator *PYTHIA* is used for low Q^2 analysis

$\Delta G/G$ from 2 high p_T hadrons (low Q^2)



Example of good description of the data by the simulation

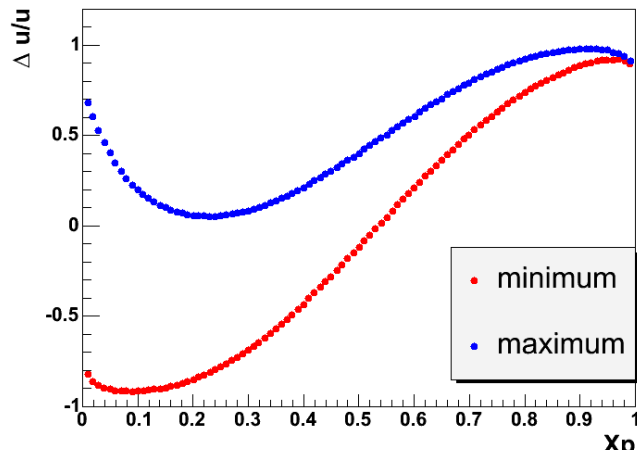
$\Delta G/G$ from 2 high p_T hadrons (low Q^2)

Problem: polarized PDF's of virtual photons not measured!

$\Delta q^{\gamma, \text{pert}}(x_p, \mu^2)$ calculable

$$-q^{\gamma, \text{npert}}(x_p, \mu_0^2) < \Delta q^{\gamma, \text{npert}}(x_p, \mu_0^2) < q^{\gamma, \text{npert}}(x_p, \mu_0^2)$$

measured

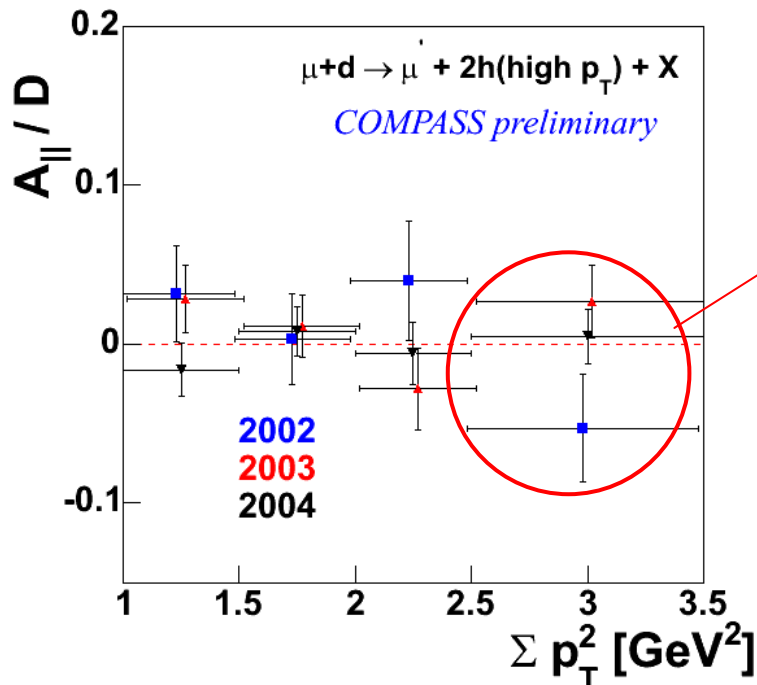


⇒ Allows us to obtain a range for $(\Delta q/q)^\gamma$ and $(\Delta G/G)^\gamma$

→ Adds a **limited** uncertainty to the estimation of $(\Delta G/G)(x_g)$

Resolved photon asymmetry

$\Delta G/G$ from 2 high p_T hadrons (low Q^2)



Values used for extraction
of $\Delta G/G$

Scale: 3 GeV^2

data	$(\Delta G/G)(x_g)$	stat	exp.syst	MC.syst	γ
02-03	0.024	0.089	0.014	0.052	0.018
02-04	0.016	0.058	0.014	0.052	0.013

$$@ x_g = 0.085^{+0.071}_{-0.035}$$

2002-2003 published: PLB **633** (2006) 25-32

$\Delta G/G$ from 2 high p_T hadrons ($Q^2 > 1$)

- Statistics smaller than in low Q^2 analysis (10%)
- Background better controlled – pQCD (QCD-C, LP)
- *LEPTO* MC generator has been used for data description (tunning similar to SMC)
- $\Sigma p_T^2 > 2.5 \text{ GeV}^2$ used.

Preliminary 2002-2003 data result:

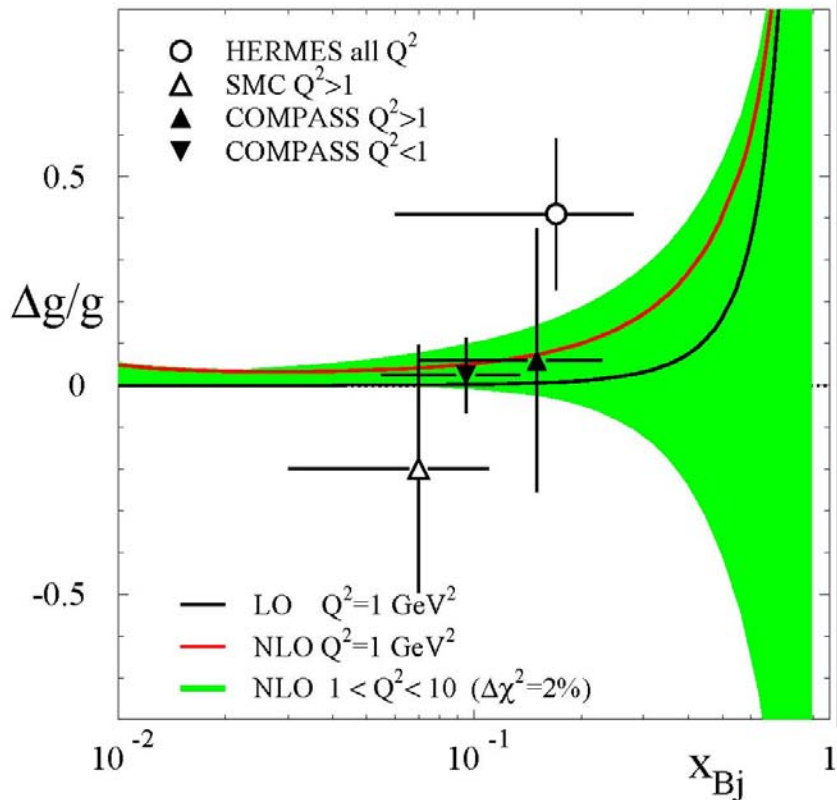
$$\frac{\Delta G}{G} = 0.06 \pm 0.31(\text{stat.}) \pm 0.06(\text{syst.}) @ x_g = 0.13 \pm 0.08$$

$\Delta G/G$ from 2 high p_T hadrons ($Q^2 > 1$)

- Analysis is ongoing; 2002-2004 results expected soon
- Scale is determined by Q^2 and – in contrast to low Q^2 analysis – the cut $\Sigma p_T^2 > 2.5 \text{ GeV}^2$ can be released to smaller value to optimize „working point”
(question: higher fraction R_{PGF} and small statistics or lower fraction and higher statistics?)

- Neural Network is tested to improve selection of PGF subprocess and optimize „working point”.
The significant improvement is expected.

Comment

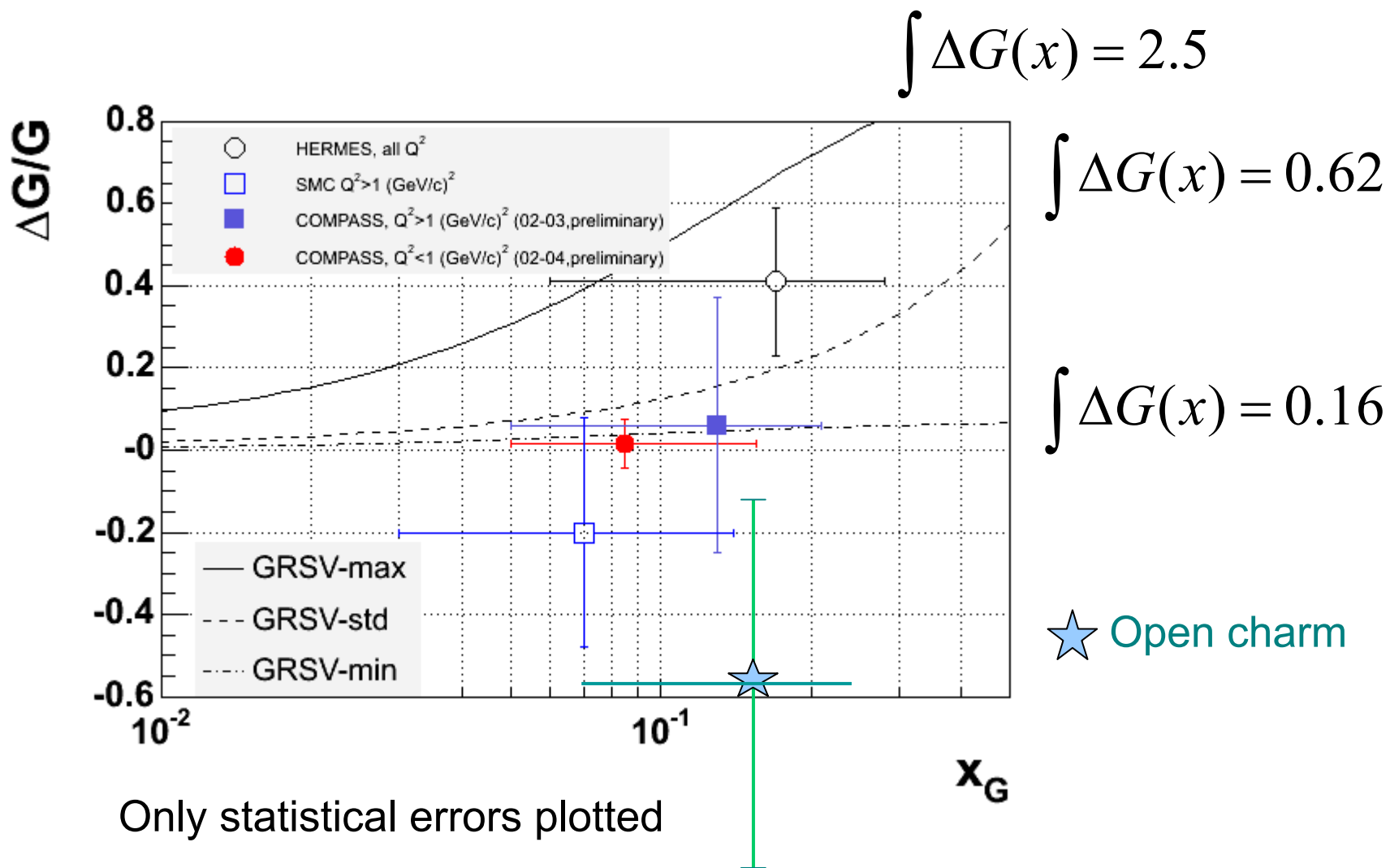


Phys.Rev.D71, 094018 (2005),
hep-ph/0602236

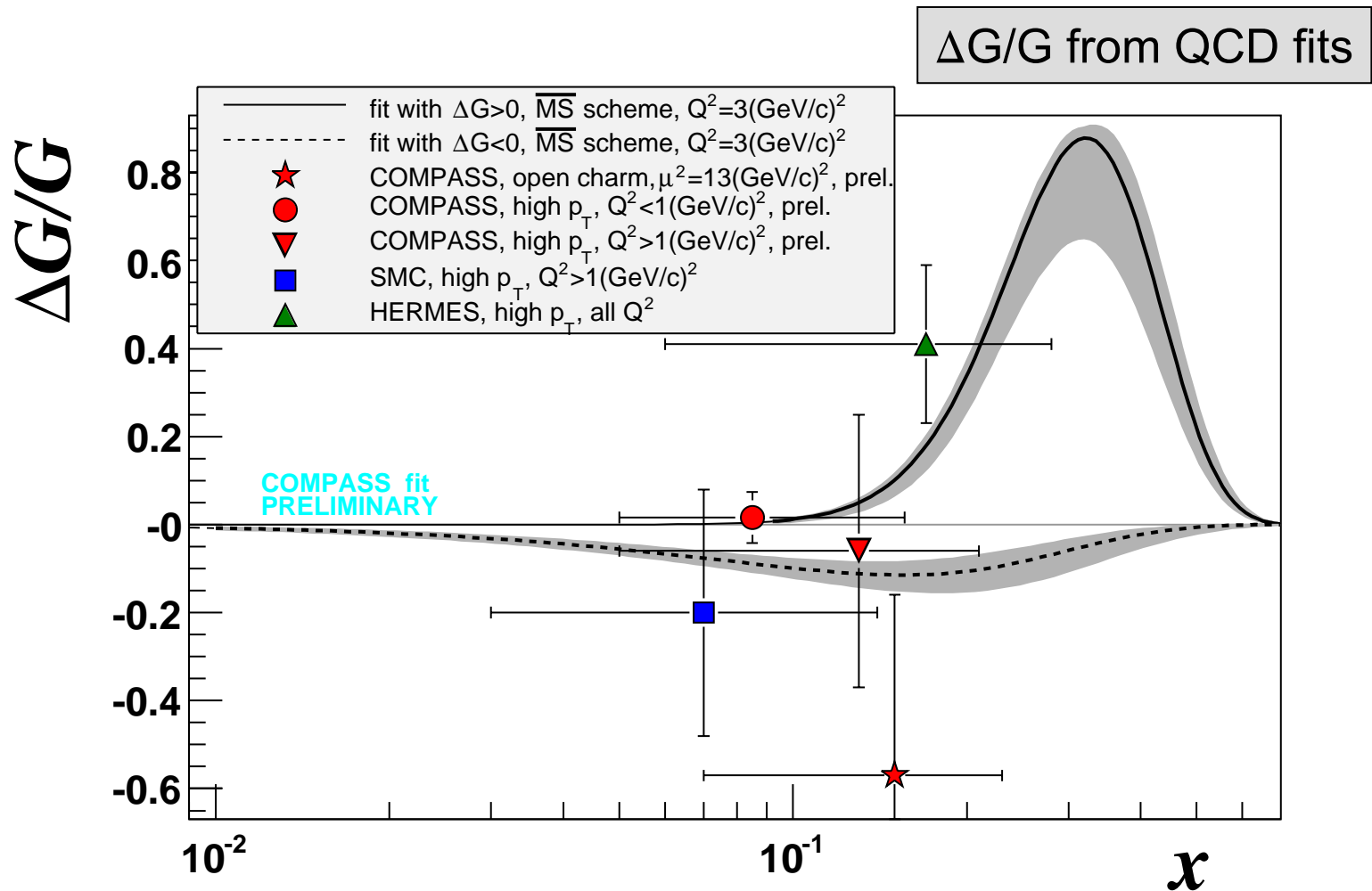
MC based high- p_T LO analysis often criticized **but**:

- NLO corrections partially simulated via so-called parton shower
- Part of the NLO effect taken into account via modification of fragmentation function (internal k_T)
- NLO fit:
de Florian, Navarro, Sassot, Jiang;
Compass results not taken in the fit!

Results for $\Delta G/G$ - summary



Results for $\Delta G/G$ - summary



Outlook

- More results soon available with 2004 data for $Q^2 > 1 \text{ GeV}^2$ high p_T events.
- For the future:
Optimization of event selection with a **neural network**,
Bins in x_g (requires improvement of x_g reconstruction),
NLO + resolved γ in open charm analysis.

- 2006 data with **new COMPASS magnet** (larger x_g)
- Expected precision with 2006 data (stat.error):
open charm - **0.28**,
high p_T $Q^2 > 1 \text{ GeV}^2$: **0.14**,
high p_T $Q^2 < 1 \text{ GeV}^2$: **0.045**.

Summary

- New measurements of A_1^d , g_1^d have been presented.
- Good agreement with results from previous experiments in the region of middle and high x .
- Improvement in statistical precision factor 4 for $x < 0.03$.
- No tendency toward negative values at $x < 0.03$.

- Disagreement of data with previous QCD fits (small x)
- Existing QCD parametrization need to be revised.

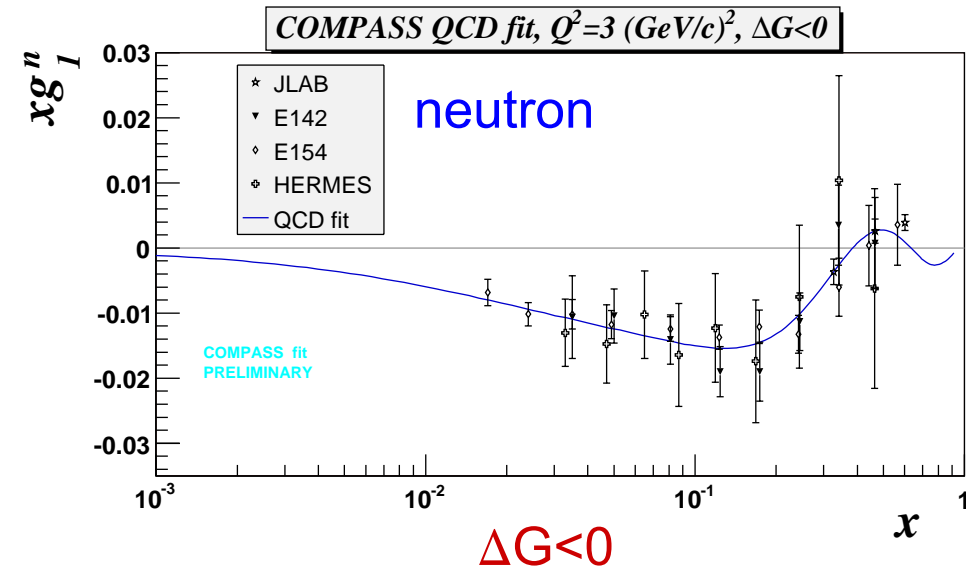
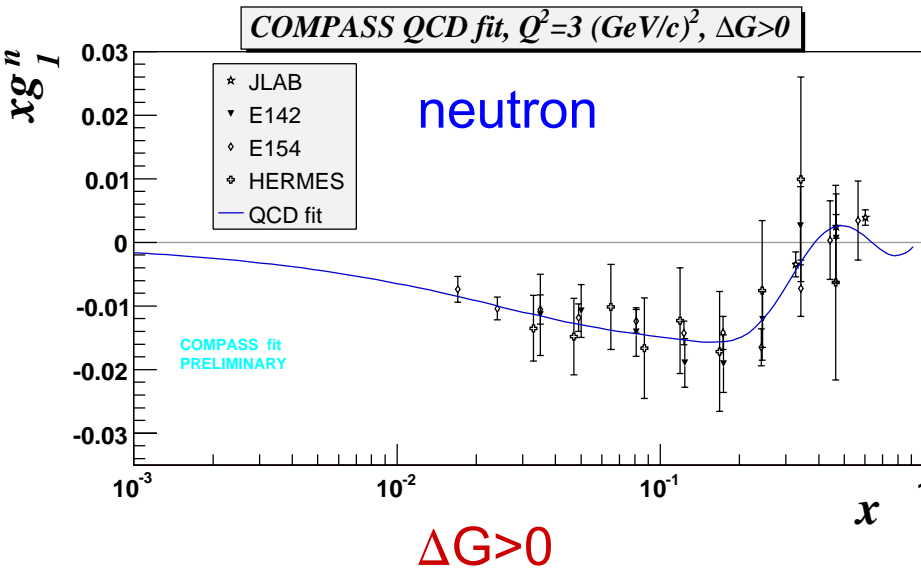
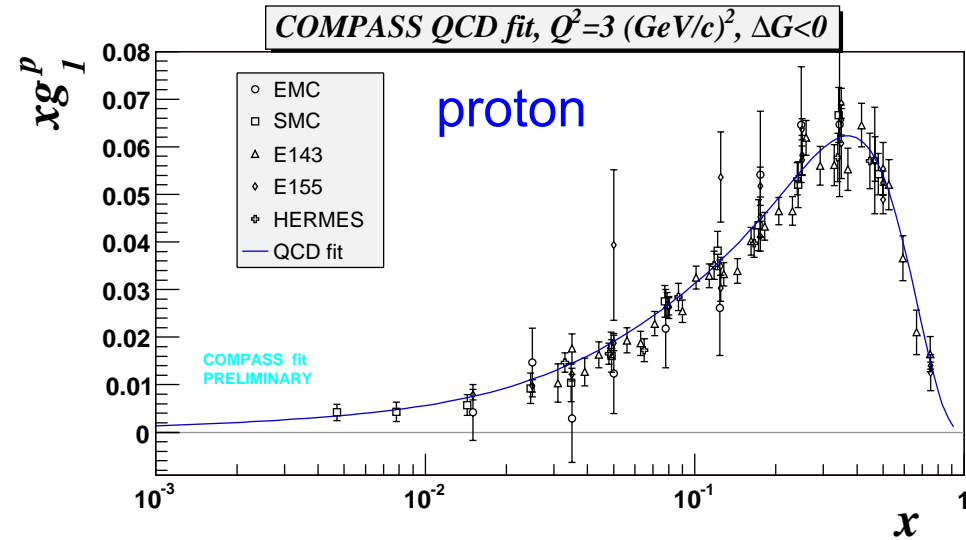
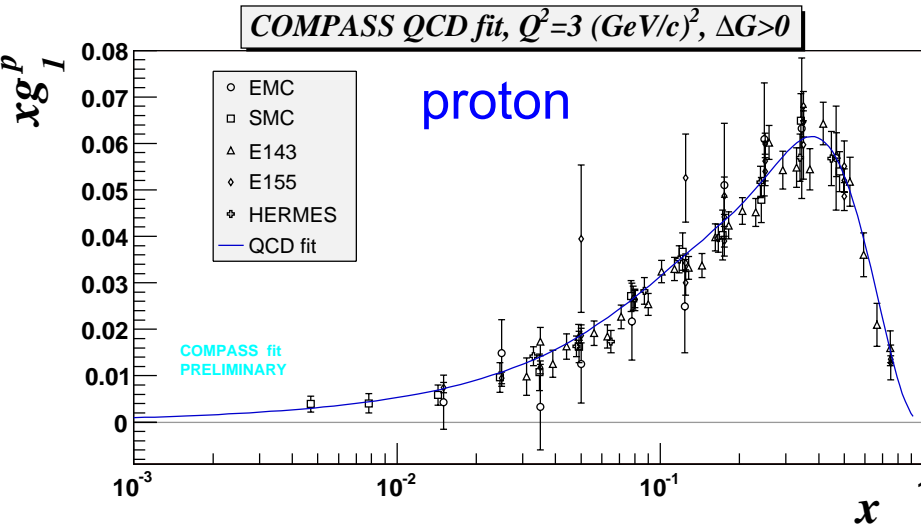
Summary

- New measurements of $\Delta G/G$ have been presented.
 - Small ΔG is preferred or $\Delta G(x_g)$ has a node around 0.1.
 - Ellis-Jaffe sum rule seems to be violated if large ΔG is excluded (axial anomaly).
-
- $\Delta G \approx 0.4$ not excluded and scenario when L is small still possible.
 - $\Delta G \approx 0$ indicates the important role of **angular orbital momentum** in nucleon spin decomposition described in the frame of parton model and pQCD.

Venus de Milo



Spare 1 – QCD fit



Spare 2 – parameters of QCD fit

Quark polarization η_Σ

- Well determined by data (proportional to the $\int_0^1 g_1^d(x, Q^2) dx$)
- No difference between results of two QCD-fit programs and the difference for two solutions ($\eta_G > 0$ and < 0) is also very small

	$\eta_G > 0$	$\eta_G < 0$
η_Σ	0.28 ± 0.01	0.32 ± 0.01

\Rightarrow

$\eta_\Sigma = 0.30 \pm 0.01(stat) \pm 0.02(evol)$
--

Gluon polarization η_G

- Indirect determination (via evolution questions)
- Solutions with $\eta_G > 0$: $\eta_G^{prog 1} = 0.26 \pm_{-0.06}^{+0.04}$, $\eta_G^{prog 2} = 0.19 \pm_{-0.10}^{+0.01}$
- Solutions with $\eta_G < 0$: $\eta_G^{prog 1} = -0.31 \pm_{-0.1}^{+0.1}$, $\eta_G^{prog 2} = -0.18 \pm_{-0.03}^{+0.04}$

$ \eta_G \simeq 0.2 - 0.3$

Spare 3 – parameters of QCD fit

Quark polarization with COMPASS data only

- The first moment of g_1^d at $Q^2=3 \text{ GeV}^2$:

$$\Gamma_1^N = \int_0^1 g_1^N(x, Q^2) dx = 0.0502 \pm 0.0028(\text{stat}) \pm 0.0020(\text{evol}) \pm 0.0051(\text{syst})$$

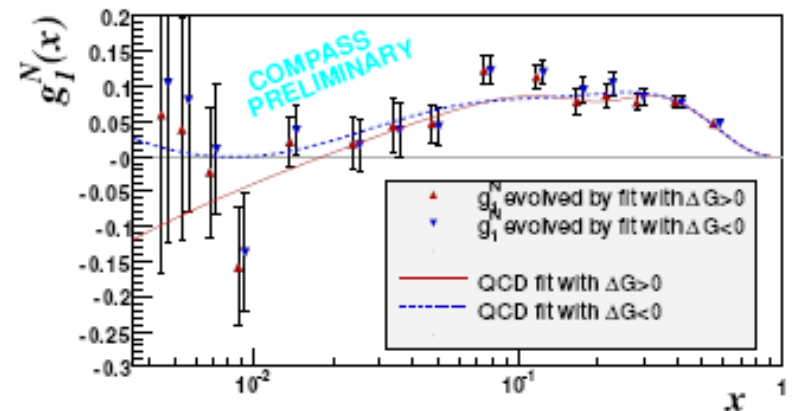
- a_0 can be extracted from the first moment of g_1^N

$$\Gamma_1^N(Q^2) \Big|_{NLO} = \frac{1}{9} \left(1 - \frac{\alpha_s(Q^2)}{\pi} + \mathcal{O}(\alpha_s^2) \right) \times \left(a_0(Q^2) + \frac{1}{4} a_8 \right)$$

- From hyperon β decays assuming $SU(3)_f$:

$$a_8 = 0.585 \pm 0.025$$

- Contribution from unmeasured x -range is $\approx 4\%$

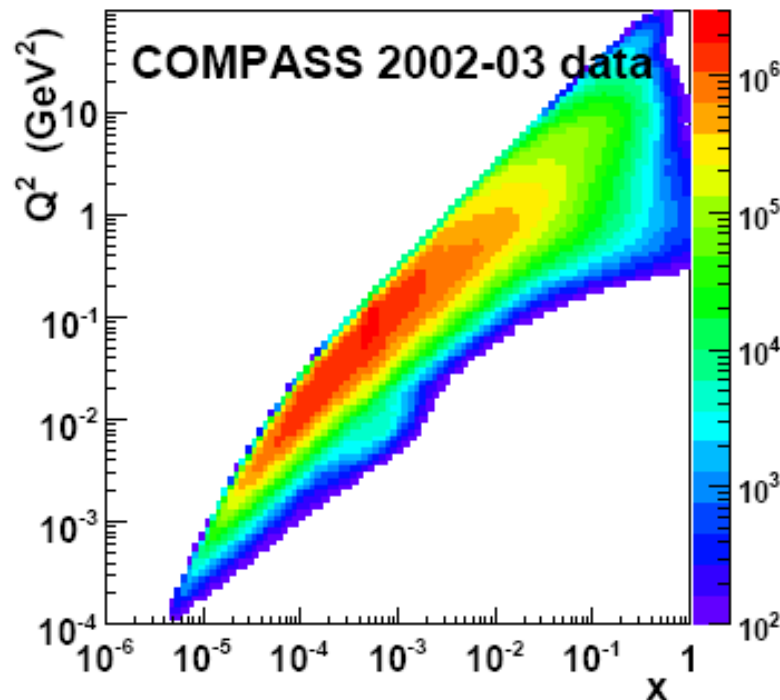
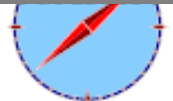


- Quark polarization at $Q^2=3 \text{ GeV}^2$:

$$a_0 = 0.35 \pm 0.03(\text{stat}) \pm 0.05(\text{syst})$$

$$\eta_\Sigma = 0.30 \pm 0.01(\text{stat}) \pm 0.02(\text{evol})$$

Spare 4 – data taking



Physics topics for longitudinal data

- inclusive asymmetries
- semi-inclusive asymmetries
- open charm production
- high p_T hadrons pairs
- Λ polarisation
- exclusive ρ production
- 20% of time for transverse data taking

	2002	2003	2004
Beam Time	106d	90d	110d
Preparation	30d	7d	3d
Integrated luminosity / fb ⁻¹	1	1.2	~ 2.4

Spare 5 – R for $Q^2 < 1 \text{ GeV}^2$

2.1 The R function

The R function which was previously used by the SMC, and it is commonly used by COMPASS [2] is composed of three different parameterizations in different regions of x (see [4] for references and explanations):

- SLAC, $x > 0.12$,
- NMC, $0.003 < x < 0.12$,
- ZEUS, $x < 0.003$.

Values of R have large discontinuities close to the validity limits of the parametrizations, Fig.4. To partially overcome the problem, a new SLAC parametrization was used for $Q^2 > 0.5 \text{ GeV}^2$, [5]. Below the $Q^2 = 0.5 \text{ GeV}^2$ the following formula was employed:

$$R(Q^2 < 0.5, x) = R_{SLAC}(0.5, x) \times \beta(1 - \exp(-Q^2/\alpha)) \quad (1)$$

where $\alpha = 0.2712$, $\beta = 1/(1 - \exp(-0.5/\alpha)) = 1.1880$. At $Q^2 = 0.5 \text{ GeV}^2$ the function and its first derivative are continuous. In the $Q^2=0$ limit: $R \sim Q^2$, which is expected from the current conservation. The new R parametrization is shown in the right plot of Fig.4. The error on R , δR , above $Q^2 = 0.5 \text{ GeV}^2$ was taken from [5] and below $Q^2 = 0.5 \text{ GeV}^2$ was set to $\delta R = 0.2$. For that value and for the simplest assumption about R for $Q^2 < 0.5 \text{ GeV}^2$ and any x , e.g. $R = 0.2$, there is an approximate agreement (within 1σ) with the value at the photo-production limit where $R=0$ and with measurements at higher Q^2 from HERA, where $R \approx 0.4$.