Lambda asymmetries

Andrea Ferrero - for the COMPASS collaboration

Transversity2005 Workshop Como, September 7-11 2005



Outline:

- **9** Λ polarization & transversity
- ${}_{ \mbox{\scriptsize \emph{D}}}$ Method of extraction of P^{Λ}
- ${}_{m{D}}$ P^{Λ} vs. x_{Bj} in 2002+2003 data
- Conclusions





Self-analyzing weak decay: $\Lambda ~\rightarrow ~p~\pi^-$, B.R. $\simeq 64\%$

The Λ polarization P_S^{Λ} along a certain direction \vec{S} is measured from the angular distribution of the decay proton:

 $W(heta^*) \propto 1 + \alpha P_S^{\Lambda} \cos(heta^*)$,

where θ^* is the proton emission angle wrt. \vec{S} in the Λ rest frame





Self-analyzing weak decay: $\Lambda ~\rightarrow ~p~\pi^-$, B.R. $\simeq 64\%$

The Λ polarization P_S^{Λ} along a certain direction $\overrightarrow{\mathbf{S}}$ is measured from the angular distribution of the decay proton:

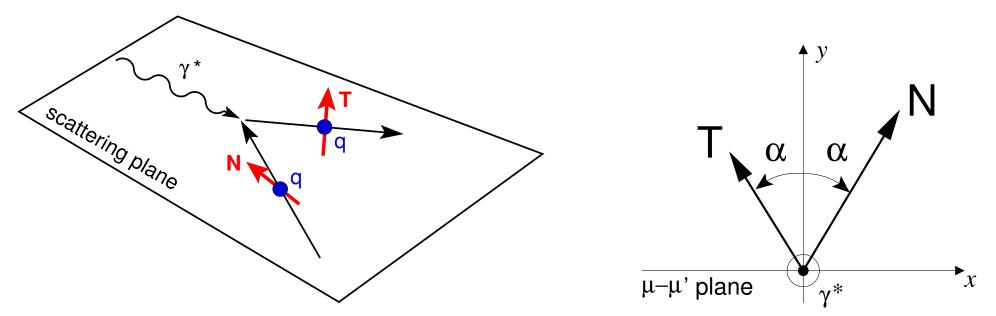
 $W(heta^*) \propto 1 + \alpha P_S^{\Lambda} \cos(heta^*)$,

where θ^* is the proton emission angle wrt. \vec{S} in the Λ rest frame

In general, the proton angular distribution is distorted by the non-ideal experimental acceptance:

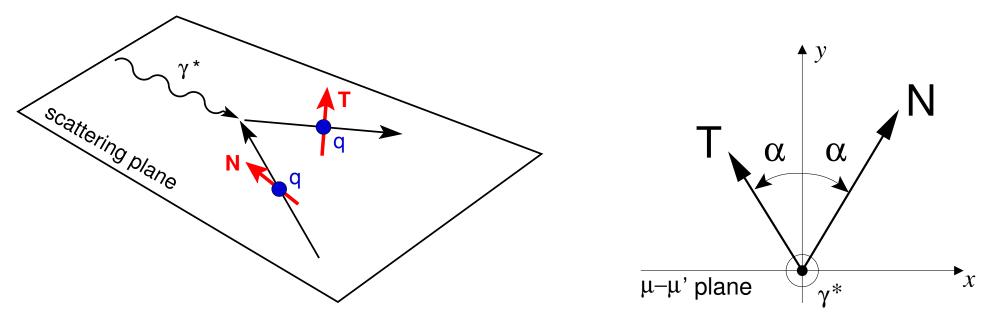
$$W_{exp}(\theta^*) \propto (1 + \alpha P_S^{\Lambda} \cos(\theta^*)) \cdot Acc(\theta^*)$$

Definition of the Λ polarization axis



- N: component of target spin perpendicular to p_{γ^*}
- T: symmetric of N wrt. the normal to the scattering plane, scaled by the virtual photon depolarization factor $D(y) = \frac{2(1-y)}{1+(1-y)^2}$

Definition of the Λ polarization axis



N: component of target spin perpendicular to p_{γ^*}

T: symmetric of N wrt. the normal to the scattering plane, scaled by the virtual photon depolarization factor $D(y) = \frac{2(1-y)}{1+(1-y)^2}$

If q fragments into a Lambda hyperon, then the measurement of P_T^{Λ} gives information about the initial quark polarization in the nucleon



Informations on $\Delta_T q(x)$ (or $h_1(x)$) can be accessed in the process

 $\mu \ N^{\uparrow} \
ightarrow \ \mu' \Lambda^{\uparrow} X$



Informations on $\Delta_T q(x)$ (or $h_1(x)$) can be accessed in the process

$$\mu \ N^{\uparrow} \rightarrow \mu' \Lambda^{\uparrow} X$$

$$P_{T,exp}^{\Lambda} = \frac{d\sigma^{\mu N^{\uparrow} \to \mu' \Lambda^{\uparrow} X} - d\sigma^{\mu N^{\downarrow} \to \mu' \Lambda^{\uparrow} X}}{d\sigma^{\mu N^{\uparrow} \to \mu' \Lambda^{\uparrow} X} + d\sigma^{\mu N^{\downarrow} \to \mu' \Lambda^{\uparrow} X}} = fP_N D(y) \frac{\sum_q e_q^2 \Delta_T q(x) \Delta_T D_{\Lambda/q}(z)}{\sum_q e_q^2 q(x) D_{\Lambda/q}(z)}$$

f = target dilution factor, P_N = target polarization,

D(y) = virtual photon depolarization factor



 $\sum 2 \cdot \frac{1}{2} \cdot \frac{1}{2}$

Informations on $\Delta_T q(x)$ (or $h_1(x)$) can be accessed in the process

$$\mu \ N^{\uparrow} \rightarrow \mu' \Lambda^{\uparrow} X$$

$$\begin{split} P_{T,exp}^{\Lambda} &= \frac{d\sigma^{\mu N^{\uparrow} \to \mu' \Lambda^{\uparrow} X} - d\sigma^{\mu N^{\downarrow} \to \mu' \Lambda^{\uparrow} X}}{d\sigma^{\mu N^{\uparrow} \to \mu' \Lambda^{\uparrow} X} + d\sigma^{\mu N^{\downarrow} \to \mu' \Lambda^{\uparrow} X}} = fP_N D(y) \frac{\sum_q e_q^2 \Delta_T q(x) \Delta_T D_{\Lambda/q}(z)}{\sum_q e_q^2 q(x) D_{\Lambda/q}(z)} \\ f &= \text{target dilution factor, } P_N = \text{target polarization,} \end{split}$$

D(y) = virtual photon depolarization factor

Alternative approach to Collins or Drell-Yan processes

 $\sum 2 \cdot \frac{1}{2} \cdot \frac{1}{2}$

Informations on $\Delta_T q(x)$ (or $h_1(x)$) can be accessed in the process

$$\mu N^{\uparrow} \rightarrow \mu' \Lambda^{\uparrow} X$$

$$\begin{split} P_{T,exp}^{\Lambda} &= \frac{d\sigma^{\mu N^{\uparrow} \to \mu' \Lambda^{\uparrow} X}}{d\sigma^{\mu N^{\uparrow} \to \mu' \Lambda^{\uparrow} X}} + \frac{d\sigma^{\mu N^{\downarrow} \to \mu' \Lambda^{\uparrow} X}}{d\sigma^{\mu N^{\downarrow} \to \mu' \Lambda^{\uparrow} X}} = fP_N D(y) \frac{\sum_q e_q^2 \Delta_T q(x) \Delta_T D_{\Lambda/q}(z)}{\sum_q e_q^2 q(x) D_{\Lambda/q}(z)} \\ f &= \text{target dilution factor, } P_N = \text{target polarization,} \end{split}$$

D(y) = virtual photon depolarization factor

- Alternative approach to Collins or Drell-Yan processes
- **(**) The chiral-odd partner of $\Delta_T q(x)$ is the fragmentation function

$$\Delta_T D_{\Lambda/q}(z) = D_{\Lambda^{\uparrow}/q^{\uparrow}}(z) - D_{\Lambda^{\downarrow}/q^{\uparrow}}(z)$$

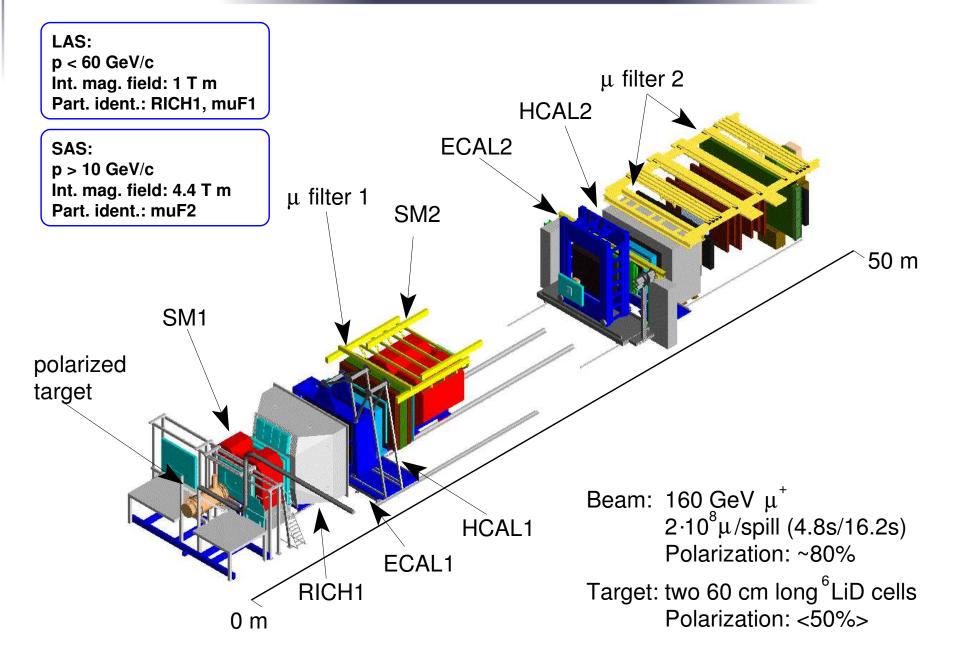


Literature on P_T^{Λ} measurement

- An early proposal for this measurement can be found in Baldracchini $et \ al.$, Fortsch. Phys. 30 (1981) 505
- Also discussed in Artru and Mekhfi, Nucl. Phys. A532 (1991) 351c, and by Artru in the proceedings of the SPIN-93 conference, LYCEN-93-53, without giving estimates of the expected Lambda polarization
- An estimate of $P_T^{\Lambda} \simeq 5\% 6\%$ at $x_{Bj} \simeq 0.2$ is given in Kunne *et al.*, LNS-Ph-93-01, assuming a target polarization of 80% and f = 1
- A more recent proposal by M. Anselmino can be found in the proceedings of the ''Future Physics at COMPASS'' workshop
- A large uncertainty on the expected P_T^{Λ} value comes from the almost unknown properties of the fragmentation functions $\Delta_T D_{\Lambda/q}(z)$

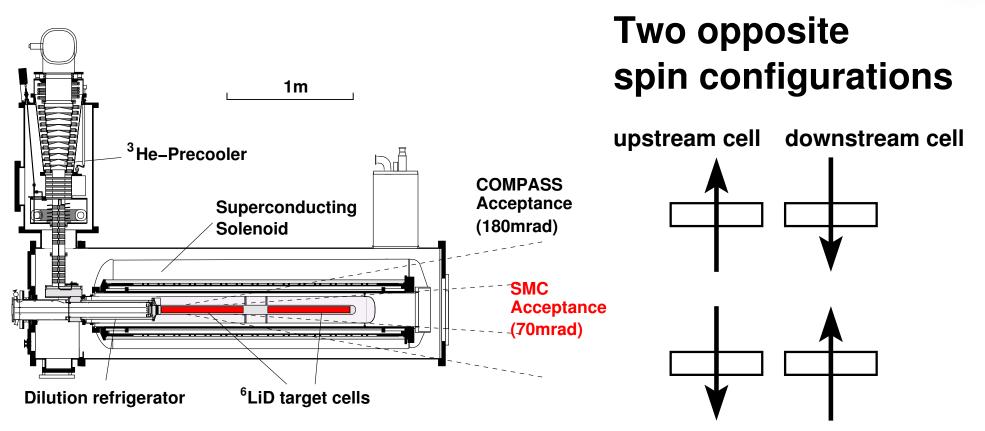
COMPASS spectrometer (2002 setup)





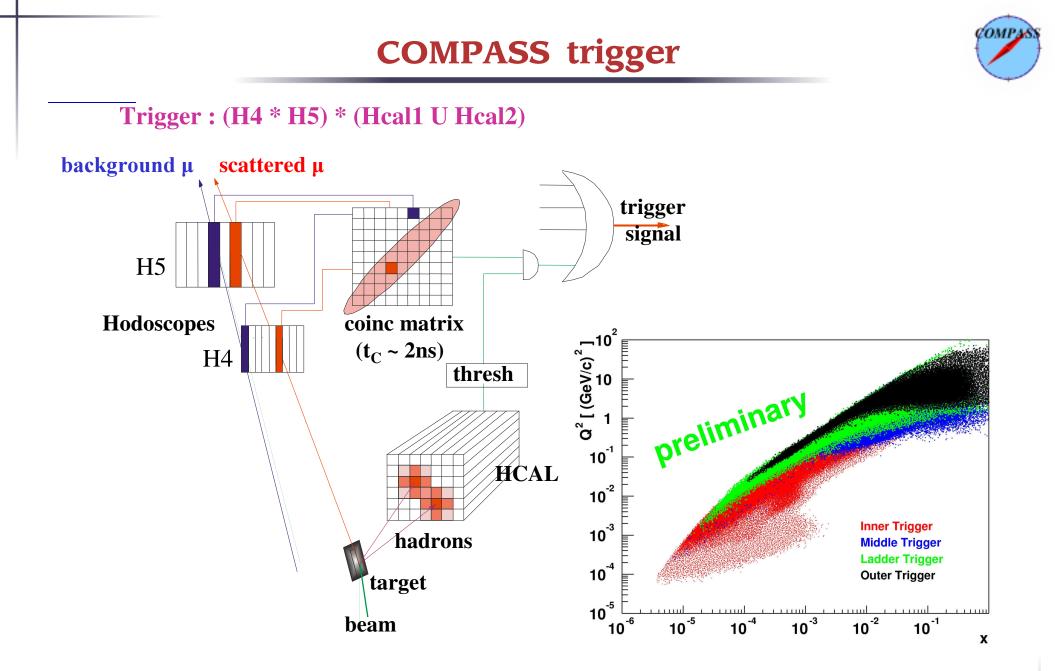
The polarized target





60 cm long ⁶LiD target cells, $\langle P_N \rangle \sim \pm 50\%$, dilution factor $f \sim 0.4$

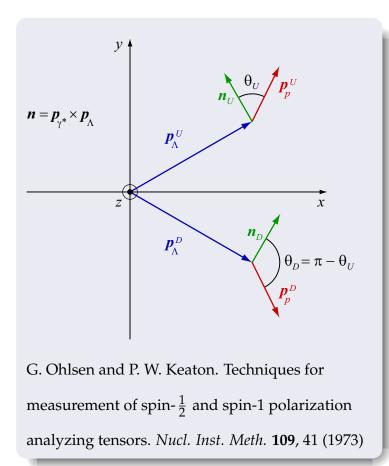
A 0.5 T dipole magnetic field is used to maintain the transverse spin configuration





The ''familiar'' case: polarization along $\overrightarrow{n} = \overrightarrow{p_{\Lambda}} \times \overrightarrow{p_{\gamma^*}}$

The distortion due to the apparatus acceptance is corrected using the up/down symmetry of the apparatus



$$\epsilon_n(\theta^*) = \frac{\sqrt{U_+ D_+} - \sqrt{U_- D_-}}{\sqrt{U_+ D_+} + \sqrt{U_- D_-}} = \alpha P_n \cos \theta^*$$

where

U: Lambda pointing upwards; D: Lambda pointing downwards

$$U_{\pm} = \frac{N_0^U}{2} A_U(\pm \cos \theta^*) (1 \pm \alpha P_n \cos \theta^*)$$
$$D_{\pm} = \frac{N_0^D}{2} A_D(\pm \cos \theta^*) (1 \pm \alpha P_n \cos \theta^*)$$

with the assumption

$$A_U(\cos\theta^*) = A_D(\cos(\pi - \theta^*)).$$

Extraction of the Λ **polarization**



In the transversity case we have additional symmetries to exploit:

two cells (+ & -) and two spin configurations (1 & 2)

Extraction of the Λ **polarization**



In the transversity case we have additional symmetries to exploit:

two cells (+ & -) and two spin configurations (1 & 2)

$$\epsilon_{T}(\theta^{*}) = \frac{\left[\sqrt{N_{1}^{+}(\theta^{*}) \cdot N_{2}^{+}(\theta^{*})} + \sqrt{N_{1}^{-}(\pi - \theta^{*}) \cdot N_{2}^{-}(\pi - \theta^{*})}\right] - \left[\sqrt{N_{1}^{+}(\pi - \theta^{*}) \cdot N_{2}^{+}(\pi - \theta^{*})} + \sqrt{N_{1}^{-}(\theta^{*}) \cdot N_{2}^{-}(\theta^{*})}\right]}{\left[\sqrt{N_{1}^{+}(\theta^{*}) \cdot N_{2}^{+}(\theta^{*})} + \sqrt{N_{1}^{-}(\pi - \theta^{*}) \cdot N_{2}^{-}(\pi - \theta^{*})}\right] + \left[\sqrt{N_{1}^{+}(\pi - \theta^{*}) \cdot N_{2}^{+}(\pi - \theta^{*})} + \sqrt{N_{1}^{-}(\theta^{*}) \cdot N_{2}^{-}(\theta^{*})}\right]}$$

 $= \alpha P_T^{\Lambda} \cos \theta^*, \quad (= \alpha P_T^{\Lambda}/2 \text{ if only 2 } \theta^* \text{ bins are considered})$

where θ^* is the proton decay angle wrt. the T axis in the Λ rest frame, and $N_{1(2)}^{\pm}(\theta^*) = \Phi_{1(2)}^{\pm} \left(\frac{d\sigma}{d\Omega}\right)^0 \left(1 \pm \alpha P_T^{\Lambda} \cos \theta^*\right) \cdot Acc_{1(2)}^{\pm}(\cos \theta^*)$

Extraction of the Λ **polarization**



In the transversity case we have additional symmetries to exploit:

two cells (+ & -) and two spin configurations (1 & 2)

$$\epsilon_{T}(\theta^{*}) = \frac{\left[\sqrt{N_{1}^{+}(\theta^{*}) \cdot N_{2}^{+}(\theta^{*})} + \sqrt{N_{1}^{-}(\pi - \theta^{*}) \cdot N_{2}^{-}(\pi - \theta^{*})}\right] - \left[\sqrt{N_{1}^{+}(\pi - \theta^{*}) \cdot N_{2}^{+}(\pi - \theta^{*})} + \sqrt{N_{1}^{-}(\theta^{*}) \cdot N_{2}^{-}(\theta^{*})}\right]}{\left[\sqrt{N_{1}^{+}(\theta^{*}) \cdot N_{2}^{+}(\theta^{*})} + \sqrt{N_{1}^{-}(\pi - \theta^{*}) \cdot N_{2}^{-}(\pi - \theta^{*})}\right] + \left[\sqrt{N_{1}^{+}(\pi - \theta^{*}) \cdot N_{2}^{+}(\pi - \theta^{*})} + \sqrt{N_{1}^{-}(\theta^{*}) \cdot N_{2}^{-}(\theta^{*})}\right]}$$

 $= \alpha P_T^{\Lambda} \cos \theta^*, \quad (= \alpha P_T^{\Lambda}/2 \text{ if only 2 } \theta^* \text{ bins are considered})$

where θ^* is the proton decay angle wrt. the T axis in the Λ rest frame, and $N_{1(2)}^{\pm}(\theta^*) = \Phi_{1(2)}^{\pm} \left(\frac{d\sigma}{d\Omega}\right)^0 \left(1 \pm \alpha P_T^{\Lambda} \cos \theta^*\right) \cdot Acc_{1(2)}^{\pm}(\cos \theta^*)$

The only assumptions in the derivation are:

- 1. The target polarization is constant: $P_T^{(1)} = P_T^{(2)}$
- 2. The experimental acceptance does not change in time: $Acc_1^+(\theta^*) = Acc_2^-(\theta^*), \qquad Acc_1^-(\theta^*) = Acc_2^+(\theta^*)$



Selection of Λ events

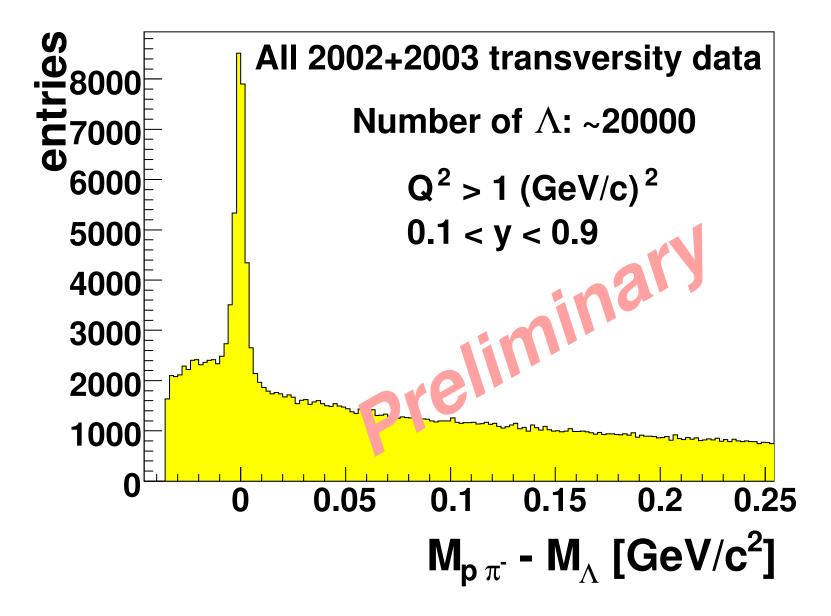
- Primary vertex in target cell material, beam crossing both cells
- **(**) μ' traverses at least 30 radiation lengths
- Tracks of p and π^- candidates traverse at least the SM1 magnet
- momentum of both decay particles > 1 GeV/c

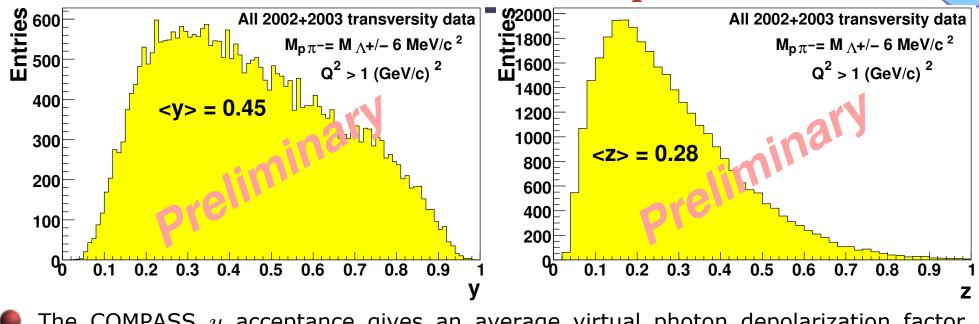
Armenteros $p_T > 23$ MeV/c

- ${}_{m{D}}$ The candidate Λ decay is downstream of the target and outside of it
- **(**) collinearity < 10 mrad
- 0.1 < y < 0.9

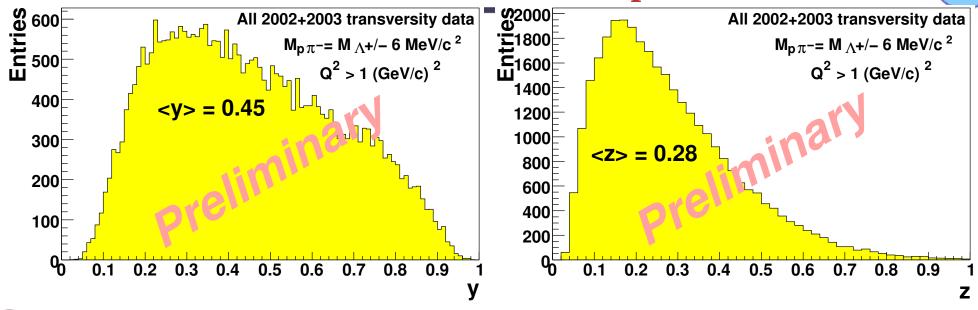
Overall available statistics





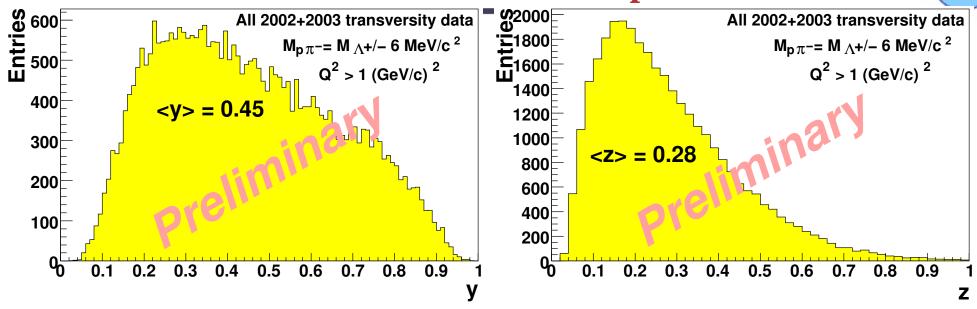


The COMPASS y acceptance gives an average virtual photon depolarization factor of $< D(y) > \simeq 0.8$



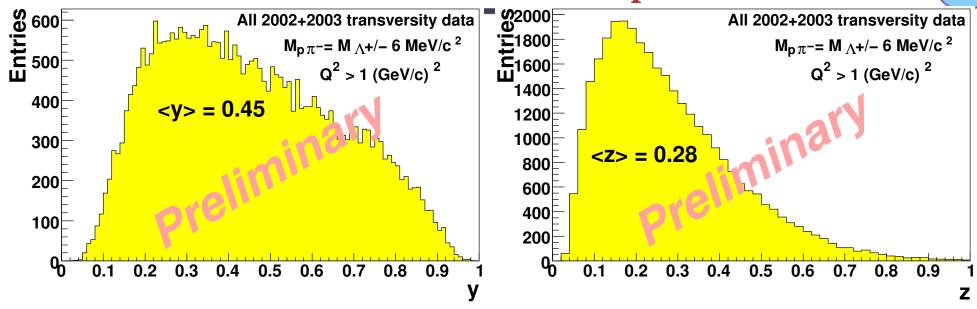
The COMPASS y acceptance gives an average virtual photon depolarization factor of $< D(y) > \simeq 0.8$

The majority of Lambda events are produced at $x_F > 0$ (current fragmentation region)



The COMPASS y acceptance gives an average virtual photon depolarization factor of $< D(y) > \simeq 0.8$

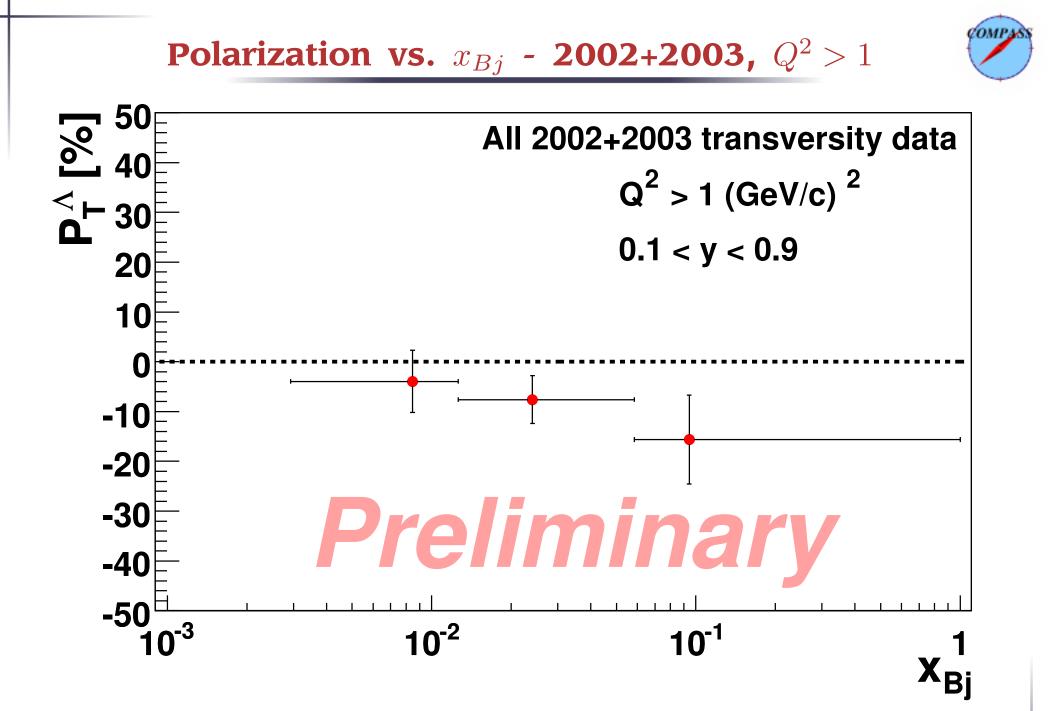
- The majority of Lambda events are produced at $x_F > 0$ (current fragmentation region)
- The accessible x_{Bj} ranges are:

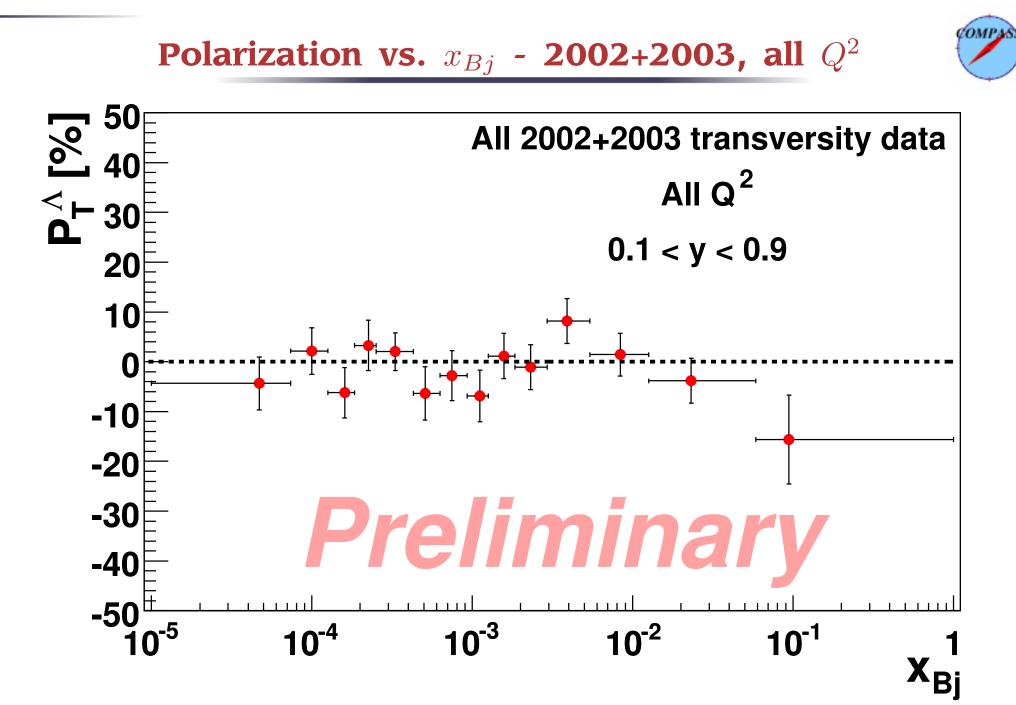


The COMPASS y acceptance gives an average virtual photon depolarization factor of $< D(y) > \simeq 0.8$

- The majority of Lambda events are produced at $x_F > 0$ (current fragmentation region)
- The accessible x_{Bj} ranges are:
 - $\sim 10^{-5} < x_{Bj} < 1$ for all Q^2
 - \bigcirc ~ 3 · 10⁻³ < x_{Bj} < 1 \bigcirc $@ Q^2 > 1$ (GeV/c)²

A binning on x_{Bj} has been applied to study the x_{Bj} -dependence of P_T^{Λ}







Possible systematic effects have been extensively studied, mainly by

() Checking the K^0 polarization



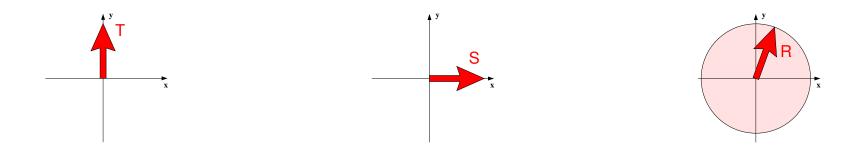
Possible systematic effects have been extensively studied, mainly by

- **(**) Checking the K^0 polarization
- **(**) Measuring P_T^{Λ} using two halves of the same target cell



Possible systematic effects have been extensively studied, mainly by

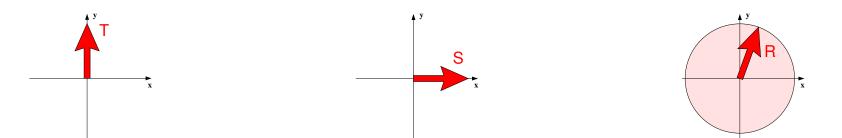
- **(**) Checking the K^0 polarization
- **9** Measuring P_T^{Λ} using two halves of the same target cell
- Assuming either an horizontal (S) or random (R) orientation of the target polarization, instead of the real vertical one (T)





Possible systematic effects have been extensively studied, mainly by

- **(**) Checking the K^0 polarization
- **9** Measuring P_T^{Λ} using two halves of the same target cell
- Assuming either an horizontal (S) or random (R) orientation of the target polarization, instead of the real vertical one (T)



The analysis shows that systematic effects are not larger than the statistical error

Conclusions



- **9** The P_T^{Λ} has been measured in the 2002+2003 transversity data sample
- The x_{Bj} dependence does not show a significant deviation from zero, but the statistics in the most interesting region ($x_{Bj} \sim 0.1$) is still poor
- The study of systematic effects shows that they are not larger than the statistical errors
- A statistics equivalent to the 2002+2003 sets is expected from the 2004 data, the analysis of which is now being started

Conclusions



- ${}$ The P_T^Λ has been measured in the 2002+2003 transversity data sample
- The x_{Bj} dependence does not show a significant deviation from zero, but the statistics in the most interesting region ($x_{Bj} \sim 0.1$) is still poor
- The study of systematic effects shows that they are not larger than the statistical errors
- A statistics equivalent to the 2002+2003 sets is expected from the 2004 data, the analysis of which is now being started

Tank you!