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# EXCLUSIVE $\omega \pi^0$ PRODUCTION WITH MUONS

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### Abstract

Using 160 GeV muon scattering data collected with the COMPASS Experiment at CERN, the exclusive production of  $\omega \pi^0$  via virtual photons was studied. Selective population of a peak around 1250 MeV is observed. Possible contributions from spin-parity 1<sup>-</sup> are searched for, inspecting decay angular correlations. In particular, the orientation of the  $\omega$  decay plane may allow a distinction from the 1<sup>+</sup>  $b_1(1235)$  state. Our observation is compared with indications of a  $\rho'(1250)$  in annihilation and in  $\gamma p$ .

# 1 Motivation

Identification of the radially excited  $\rho$  meson is debated since a long time <sup>1) 2)</sup>. An early photoproduction experiment <sup>3)</sup>, using photons with energy between 20 and 70 GeV, observed an enhancement in the  $\omega \pi^0$  channel with mass around 1250 MeV and width of about 200 MeV. For spin-parity analysis it was assumed that the produced meson retains the helicity of the incoming photon (s-channel helicity conservation, SCHC). A dominant  $1^-$  contribution was deduced.

However, subsequent investigations at the CERN SPS <sup>4</sup>) and at SLAC <sup>5</sup>) employing linearly polarized photons, revealed a dominance of the well known  $J^{PC} = 1^{+-}$  state  $b_1(1235)$ , leaving only about 20% for a  $\rho'(1^{--})$  contribution at the same mass. Angular distributions were found inconsistent with SCHC in these experiments where the mean photon energy was 20–30 GeV.

Supportive evidence for a  $\rho'$  state at this mass came from a Crystal Barrel study <sup>6</sup>) of the annihilation reaction  $\bar{p}n \to \omega \pi^- \pi^0$ , suggesting  $\rho$  excitations at 1200, 1400 and 1700 MeV. The lowest lying state stands out by dominant  $\omega \pi$  decay, in contrast to other non- $\omega$  related  $4\pi$  decays.

The experimental situation has been reviewed by Donnachie and Kalashnikova<sup>2)</sup>, including results from  $e^+e^-$  annihilation and  $\tau$  decay. In their interpretation, two 1<sup>--</sup> states with mixed configurations are present between the ground state  $\rho(770)$  and the first orbital excitation  $(1^3D_1) \rho'(1700)$ : the one at 1250 MeV with dominant  $q\bar{q}$  configuration  $2^3S_1$  (the radial  $\rho'$  excitation), decaying preferably via  $\omega \pi$ , and the heavier one at ~1450 MeV, with dominant hybrid or quartet configuration, preferring alternative decay channels like e.g.  $a_1\pi$ .

Concerning  $b_1$  and  $\rho'$  competition in photoproduction, it was suggested <sup>7</sup>) that helicity-flip Regge exchange, resulting in  $b_1$ , prevails at the mean photon energies of Ref. <sup>4</sup>) <sup>5</sup>), while helicity conserving Pomeron exchange, resulting in  $\rho'$ , wins at higher energy.

We report on the first study of  $\omega \pi^0$  production with virtual, quasi-real photons in inelastic muon scattering. According to the suggested systematics <sup>7</sup>),  $b_1$  and  $\rho'$  production should be of comparable size at the available  $\gamma^* p$  c.m. energy W $\approx$ 13 GeV.

### 2 Experimental setup

COMPASS <sup>8</sup>) is a two stage magnetic spectrometer installed at the end of the M2 beam extraction line at the CERN SPS machine. The extracted  $\mu^+$  beam of an intensity of about  $2 \cdot 10^8$  per spill, with 5 s spill length and 16 s repetition, had an energy of 160 GeV and a polarisation of about 80%. It was directed on a two-cells polarized <sup>6</sup>LiD target, where the (longitudinal) polarisation was +

and -56%.

Charged particle tracking involves silicon strip detectors, scintillation fibers, micromegas and GEMs at small angles and straw drift tubes and multiwire proportional chambers at large angles. In addition, muon-hadron separation is obtained with  $\mu$ -filters.

For neutral particle detection in 2004 a lead glass detector, covering angles up to  $\pm 35$  mrad as viewed from the target, served as electromagnetic calorimeter (ECAL2).

#### **3** Event selection

A data sample collected in 8 weeks of the 2004 COMPASS run was analyzed. To select the exclusive process

$$\mu + N \to \mu' + \omega (\pi^+ \pi^- \pi^0) \pi^0 + N, \tag{1}$$

with  $\pi^0 \to \gamma \gamma$ , the following criteria were applied:

— a primary reaction vertex with an identified incoming and scattered  $\mu$  and (only) two additional particles of opposite charge is fully reconstructed;

- 4 and only 4 clusters not associated with a reconstructed charged track are found in ECAL2. To reduce background, only clusters with energy above 1 GeV are accepted.

—  $\pi^0$ 's are selected cutting on the 2 photon invariant mass, 120 MeV <  $m(\gamma\gamma) < 150$  MeV, and on the decay opening angle,  $\theta_{\gamma\gamma} < 0.025$  rad;

— a  $\omega$  candidate is selected imposing the cut 750 MeV  $< m(\pi^+\pi^-\pi^0) < 815$  MeV;

— exclusivity is defined by means of the missing energy

$$E_{miss} = \frac{M_{miss}^2 - M_P^2}{2M_P},\tag{2}$$

where  $M_P$  is the proton mass and  $M_{miss}$  is the missing mass. The exclusive  $\omega \pi^0$  final sample is selected with the cut  $-6 < E_{miss} < 4$  GeV.

Figure 1, left, shows the missing energy versus the  $4\pi$  invariant mass for events with a uniquely identified  $\omega\pi^0$  without the exclusivity cut: evident is the presence of an exclusive sample around  $E_{miss} = 0$ . The  $E_{miss}$  window

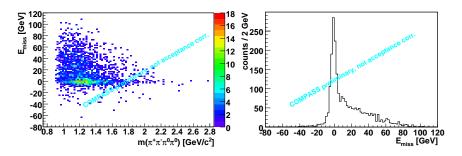


Figure 1: Missing energy  $E_{miss}$  vs.  $\pi^+\pi^-\pi^0\pi^0$  invariant mass for events with a single reconstructed  $\omega(\pi^+\pi^-\pi^0)\pi^0$  (left) and projection on the  $E_{miss}$  axis (right).

used for selection was adapted to the exclusivity peak visible in the projection (right).

#### 4 Results

Figure 2 shows the  $\omega \pi^0$  invariant mass spectrum. A peak with a mean value of about 1250 MeV and a width of about 300 MeV is observed. The acceptance variation over the peak range is estimated to be less than 20%. Our observation is consistent with the results of the quoted photoproduction experiments.

To access non- $\omega$  background, the  $\pi^+\pi^-\pi^0$  invariant mass cut was somewhat relaxed. Figure 3 (left) shows the  $3\pi$  versus the  $4\pi$  invariant mass: events in the  $\omega$  mass region correspond to the  $4\pi$  invariant mass interval around 1250 MeV. The projection on the  $3\pi$  mass axis (right), puts in evidence the  $\omega$  contribution; the width is due to the experimental resolution.

For a quantitative determination of the non- $\omega$  background, we have considered the  $\lambda$  distribution, defined by

$$\lambda = \frac{|\vec{p_1} \times \vec{p_2}|^2}{|\vec{p_1} \times \vec{p_2}|^2_{max}},\tag{3}$$

where  $\vec{p_1}$  and  $\vec{p_2}$  are the momenta of any two of the three pions. In this analysis, the two charged ones were chosen. The observed linear increase of the intensity with  $\lambda$  is a unique signature of  $J^P = 1^-$ , already applied in the original  $J^P$ 

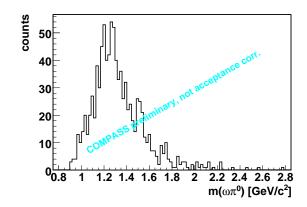


Figure 2: Invariant mass spectrum of exclusively produced  $\omega \pi^0$ .

assignment for the  $\omega$ <sup>9)</sup>. In contrast, the  $\lambda$  distribution for events outside the exclusivity window is flat. From the linear fit in figure 4, we deduce a background contribution of 12% in the final sample.

Figure 5 shows some important kinematic distributions for the final sample: the virtual photon mass squared  $Q^2 = -q^2$ , the Bjorken scale variable  $x_B$ , the  $\gamma^* p$  center of mass energy W, and the  $\omega \pi^0$  momentum in the laboratory system. The mean value of the latter corresponds to  $E(\gamma^*) \approx 90$  GeV. The

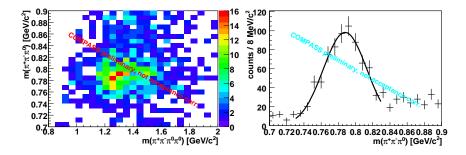


Figure 3:  $3\pi$  vs.  $4\pi$  invariant mass for events in the exclusivity region (left) and corresponding  $3\pi$  mass projection (right).

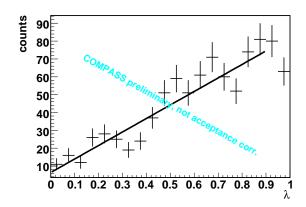


Figure 4:  $\lambda$  distribution, eq. (3).

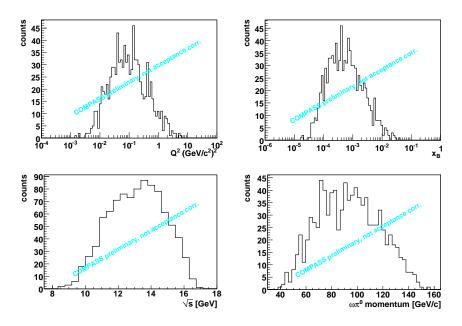


Figure 5: Kinematic distributions for the exclusive  $\omega \pi^0$  final sample. Top-left: Virtual photon mass squared  $Q^2$ ; Top-right: Bjorken scale variable; Bottomleft:  $\gamma^* p$  c.m. energy W; Bottom-right:  $\omega \pi^0$  momentum in laboratory frame.

4-momentum transfer squared  $t = (q - v)^2$  (not shown) is characterized by an exponential shape, as is typical of diffractive processes.

## 5 Angular distributions

Three types of angular correlations are suited for spin-parity studies. The first two characterize the decay of the  $\omega \pi^0$  resonance:

(i) the angle  $\psi$  of the  $\omega$  momentum  $\vec{p}_{\omega}$  relative to the  $\omega \pi^0$  direction (reference axis z) in the overall  $\gamma^* p$  c.m. system;

(ii) the angle  $\theta$  between the vector  $\vec{n}_{\omega}$  perpendicular to the  $\omega$  decay plane (in the  $\omega$  rest frame) and the z axis.

For electroproduction via quasi-real photons, one can assume linear polarization of the  $\gamma^*$  in the primary scattering plane and adopt the corresponding angular correlation formalism <sup>10</sup>). Following Ballam *et al.* <sup>11</sup>), we define appropriate "spin analyzers"  $\vec{a} = \vec{n}_{\omega} \times \vec{p}_{\omega}$  and  $\vec{a} = \vec{n}_{\omega}$  for  $J^P = 1^-$  and  $1^+$  states, respectively. Their direction with respect to the  $\gamma^*$  polarization is described by:

(iii) the azimuthal angle  $\Psi$  between  $\mu$  scattering plane and  $\vec{a}$ .

Assuming SCHC, the two sets of angular distributions in table 1 are predicted <sup>3</sup>) for the two different  $J^P$  assignments to  $\omega \pi^0$ . The quantity  $x \approx 0.07$  is the known D/S-wave amplitude ratio squared of  $b_1$ .

Monte Carlo simulations for pure  $1^+$  and  $1^-$  states reveal a strong acceptance dependance of the distribution (i), whereas (ii) is only weakly affected. As shown in figure 6, the characteristic shapes of  $I(\cos\theta)$  are roughly maintained after taking into account detector and selection acceptance. Our preliminary experimental results (not shown) are in favour of the  $1^-$  case. However the de-

Table 1: Decay angular distributions for  $J^P = 1^{\pm}$  assignments to  $\omega \pi^0$ .

| $J^P$                         | $I(cos\psi)$             | $I(cos\theta)$          |
|-------------------------------|--------------------------|-------------------------|
| $1^+$ (b <sub>1</sub> (1235)) | $\sim 1 + x \cos^2 \psi$ | $\sim sin^2 \theta$     |
| $1^{-}(\rho')$                | $\sim 1 + cos^2 \psi$    | $\sim 1 + cos^2 \theta$ |

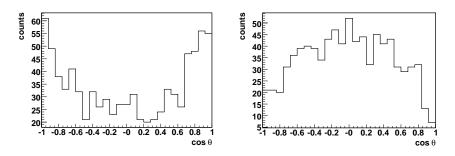


Figure 6: Estimate of the  $\cos\theta$  distributions, based on Monte Carlo simulations of the detector and selection acceptance, for  $J^P(\omega\pi^0) = 1^-$  (left) and  $1^+$ (right).

pendence on the SCHC assumption should be kept in mind. This holds as well for the distribution (iii), which shows an indication of a  $cos2\Psi$  contribution, characteristic of  $J^P = 1^-$ . Interference between S- and P-wave, corresponding to  $1^+$  and  $1^-$  decay in  $\omega \pi^0$ , would give rise to a forward-backward anisotropy in the distribution (i), irrespective of the SCHC assumption. Detailed acceptance studies are required for this analysis.

# 6 Conclusion

We have observed the exclusive production of  $\omega \pi^0$  in muon scattering via virtual photons in the energy range around 90 GeV lab. energy. The mass spectrum is dominated by a peak at 1250 MeV and width 300 MeV, which is consistent with previous photoproduction experiments. Preliminary results on angular correlations are consistent with the presence of a 1<sup>-</sup> contribution, if SCHC holds. An appreciable increase in statistics is expected with the 2006 and 2007 COMPASS data.

### 7 Acknowledgements

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