# EXCLUSIVE $\omega \pi^{0}$ PRODUCTION WITH MUONS 

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#### Abstract

Using 160 GeV muon scattering data collected with the COMPASS Experiment at CERN, the exclusive production of $\omega \pi^{0}$ via virtual photons was studied. Selective population of a peak around 1250 MeV is observed. Possible contributions from spin-parity $1^{-}$are searched for, inspecting decay angular correlations. In particular, the orientation of the $\omega$ decay plane may allow a distinction from the $1^{+} b_{1}(1235)$ state. Our observation is compared with indications of a $\rho^{\prime}(1250)$ in annihilation and in $\gamma p$.


## 1 Motivation

Identification of the radially excited $\rho$ meson is debated since a long time 1) 2). An early photoproduction experiment ${ }^{3)}$, using photons with energy between 20 and 70 GeV , observed an enhancement in the $\omega \pi^{0}$ channel with mass around

1250 MeV and width of about 200 MeV . For spin-parity analysis it was assumed that the produced meson retains the helicity of the incoming photon (s-channel helicity conservation, SCHC ). A dominant $1^{-}$contribution was deduced.

However, subsequent investigations at the CERN SPS ${ }^{4)}$ and at SLAC ${ }^{5)}$ employing linearly polarized photons, revealed a dominance of the well known $J^{P C}=1^{+-}$state $b_{1}(1235)$, leaving only about $20 \%$ for a $\rho^{\prime}\left(1^{--}\right)$contribution at the same mass. Angular distributions were found inconsistent with SCHC in these experiments where the mean photon energy was $20-30 \mathrm{GeV}$.

Supportive evidence for a $\rho^{\prime}$ state at this mass came from a Crystal Barrel study ${ }^{6)}$ of the annihilation reaction $\bar{p} n \rightarrow \omega \pi^{-} \pi^{0}$, suggesting $\rho$ excitations at 1200,1400 and 1700 MeV . The lowest lying state stands out by dominant $\omega \pi$ decay, in contrast to other non- $\omega$ related $4 \pi$ decays.

The experimental situation has been reviewed by Donnachie and Kalashnikova ${ }^{2)}$, including results from $e^{+} e^{-}$annihilation and $\tau$ decay. In their interpretation, two $1^{--}$states with mixed configurations are present between the ground state $\rho(770)$ and the first orbital excitation $\left(1^{3} D_{1}\right) \rho^{\prime}(1700)$ : the one at 1250 MeV with dominant $q \bar{q}$ configuration $2^{3} S_{1}$ (the radial $\rho^{\prime}$ excitation), decaying preferably via $\omega \pi$, and the heavier one at $\sim 1450 \mathrm{MeV}$, with dominant hybrid or quartet configuration, preferring alternative decay channels like e.g. $a_{1} \pi$.

Concerning $b_{1}$ and $\rho^{\prime}$ competition in photoproduction, it was suggested 7) that helicity-flip Regge exchange, resulting in $b_{1}$, prevails at the mean photon energies of Ref. ${ }^{4)}$ 5), while helicity conserving Pomeron exchange, resulting in $\rho^{\prime}$, wins at higher energy.

We report on the first study of $\omega \pi^{0}$ production with virtual, quasi-real photons in inelastic muon scattering. According to the suggested systematics ${ }^{7)}, b_{1}$ and $\rho^{\prime}$ production should be of comparable size at the available $\gamma^{*} p$ c.m. energy $\mathrm{W} \approx 13 \mathrm{GeV}$.

## 2 Experimental setup

COMPASS ${ }^{8)}$ is a two stage magnetic spectrometer installed at the end of the M2 beam extraction line at the CERN SPS machine. The extracted $\mu^{+}$beam of an intensity of about $2 \cdot 10^{8}$ per spill, with 5 s spill length and 16 s repetition, had an energy of 160 GeV and a polarisation of about $80 \%$. It was directed on a two-cells polarized ${ }^{6} \mathrm{LiD}$ target, where the (longitudinal) polarisation was +
and $-56 \%$.
Charged particle tracking involves silicon strip detectors, scintillation fibers, micromegas and GEMs at small angles and straw drift tubes and multiwire proportional chambers at large angles. In addition, muon-hadron separation is obtained with $\mu$-filters.

For neutral particle detection in 2004 a lead glass detector, covering angles up to $\pm 35 \mathrm{mrad}$ as viewed from the target, served as electromagnetic calorimeter (ECAL2).

## 3 Event selection

A data sample collected in 8 weeks of the 2004 COMPASS run was analyzed. To select the exclusive process

$$
\begin{equation*}
\mu+N \rightarrow \mu^{\prime}+\omega\left(\pi^{+} \pi^{-} \pi^{0}\right) \pi^{0}+N \tag{1}
\end{equation*}
$$

with $\pi^{0} \rightarrow \gamma \gamma$, the following criteria were applied:

- a primary reaction vertex with an identified incoming and scattered $\mu$ and (only) two additional particles of opposite charge is fully reconstructed;
- 4 and only 4 clusters not associated with a reconstructed charged track are found in ECAL2. To reduce background, only clusters with energy above 1 GeV are accepted.
- $\pi^{0}$ 's are selected cutting on the 2 photon invariant mass, $120 \mathrm{MeV}<$ $m(\gamma \gamma)<150 \mathrm{MeV}$, and on the decay opening angle, $\theta_{\gamma \gamma}<0.025 \mathrm{rad} ;$
- a $\omega$ candidate is selected imposing the cut $750 \mathrm{MeV}<m\left(\pi^{+} \pi^{-} \pi^{0}\right)<815$ MeV;
- exclusivity is defined by means of the missing energy

$$
\begin{equation*}
E_{m i s s}=\frac{M_{m i s s}^{2}-M_{P}^{2}}{2 M_{P}} \tag{2}
\end{equation*}
$$

where $M_{P}$ is the proton mass and $M_{\text {miss }}$ is the missing mass. The exclusive $\omega \pi^{0}$ final sample is selected with the cut $-6<E_{\text {miss }}<4 \mathrm{GeV}$.

Figure 1, left, shows the missing energy versus the $4 \pi$ invariant mass for events with a uniquely identified $\omega \pi^{0}$ without the exclusivity cut: evident is the presence of an exclusive sample around $E_{\text {miss }}=0$. The $E_{\text {miss }}$ window


Figure 1: Missing energy $E_{\text {miss }}$ vs. $\pi^{+} \pi^{-} \pi^{0} \pi^{0}$ invariant mass for events with a single reconstructed $\omega\left(\pi^{+} \pi^{-} \pi^{0}\right) \pi^{0}$ (left) and projection on the $E_{\text {miss }}$ axis (right).
used for selection was adapted to the exclusivity peak visible in the projection (right).

## 4 Results

Figure 2 shows the $\omega \pi^{0}$ invariant mass spectrum. A peak with a mean value of about 1250 MeV and a width of about 300 MeV is observed. The acceptance variation over the peak range is estimated to be less than $20 \%$. Our observation is consistent with the results of the quoted photoproduction experiments.

To access non- $\omega$ background, the $\pi^{+} \pi^{-} \pi^{0}$ invariant mass cut was somewhat relaxed. Figure 3 (left) shows the $3 \pi$ versus the $4 \pi$ invariant mass: events in the $\omega$ mass region correspond to the $4 \pi$ invariant mass interval around 1250 MeV . The projection on the $3 \pi$ mass axis (right), puts in evidence the $\omega$ contribution; the width is due to the experimental resolution.

For a quantitative determination of the non- $\omega$ background, we have considered the $\lambda$ distribution, defined by

$$
\begin{equation*}
\lambda=\frac{\left|\overrightarrow{p_{1}} \times \overrightarrow{p_{2}}\right|^{2}}{\left|\overrightarrow{p_{1}} \times \overrightarrow{p_{2}}\right|_{\max }^{2}} \tag{3}
\end{equation*}
$$

where $\overrightarrow{p_{1}}$ and $\overrightarrow{p_{2}}$ are the momenta of any two of the three pions. In this analysis, the two charged ones were chosen. The observed linear increase of the intensity with $\lambda$ is a unique signature of $J^{P}=1^{-}$, already applied in the original $J^{P}$


Figure 2: Invariant mass spectrum of exclusively produced $\omega \pi^{0}$.
assignment for the $\omega^{9)}$. In contrast, the $\lambda$ distribution for events outside the exclusivity window is flat. From the linear fit in figure 4, we deduce a background contribution of $12 \%$ in the final sample.

Figure 5 shows some important kinematic distributions for the final sample: the virtual photon mass squared $Q^{2}=-q^{2}$, the Bjorken scale variable $x_{B}$, the $\gamma^{*} p$ center of mass energy $W$, and the $\omega \pi^{0}$ momentum in the laboratory system. The mean value of the latter corresponds to $E\left(\gamma^{*}\right) \approx 90 \mathrm{GeV}$. The


Figure 3: $3 \pi$ vs. $4 \pi$ invariant mass for events in the exclusivity region (left) and corresponding $3 \pi$ mass projection (right).


Figure 4: $\lambda$ distribution, eq. (3).


Figure 5: Kinematic distributions for the exclusive $\omega \pi^{0}$ final sample. Top-left: Virtual photon mass squared $Q^{2}$; Top-right: Bjorken scale variable; Bottomleft: $\gamma^{*} p$ c.m. energy $W$; Bottom-right: $\omega \pi^{0}$ momentum in laboratory frame.

4-momentum transfer squared $t=(q-v)^{2}$ (not shown) is characterized by an exponential shape, as is typical of diffractive processes.

## 5 Angular distributions

Three types of angular correlations are suited for spin-parity studies. The first two characterize the decay of the $\omega \pi^{0}$ resonance:
(i) the angle $\psi$ of the $\omega$ momentum $\vec{p}_{\omega}$ relative to the $\omega \pi^{0}$ direction (reference axis $z$ ) in the overall $\gamma^{*} p$ c.m. system;
(ii) the angle $\theta$ between the vector $\vec{n}_{\omega}$ perpendicular to the $\omega$ decay plane (in the $\omega$ rest frame) and the $z$ axis.

For electroproduction via quasi-real photons, one can assume linear polarization of the $\gamma^{*}$ in the primary scattering plane and adopt the corresponding angular correlation formalism ${ }^{10)}$. Following Ballam et al. ${ }^{11)}$, we define appropriate "spin analyzers" $\vec{a}=\vec{n}_{\omega} \times \vec{p}_{\omega}$ and $\vec{a}=\vec{n}_{\omega}$ for $J^{P}=1^{-}$and $1^{+}$states, respectively. Their direction with respect to the $\gamma^{*}$ polarization is described by:
(iii) the azimuthal angle $\Psi$ between $\mu$ scattering plane and $\vec{a}$.

Assuming SCHC, the two sets of angular distributions in table 1 are predicted ${ }^{3)}$ for the two different $J^{P}$ assignments to $\omega \pi^{0}$. The quantity $x \approx 0.07$ is the known $D / S$-wave amplitude ratio squared of $b_{1}$.

Monte Carlo simulations for pure $1^{+}$and $1^{-}$states reveal a strong acceptance dependance of the distribution (i), whereas (ii) is only weakly affected. As shown in figure 6 , the characteristic shapes of $I(\cos \theta)$ are roughly maintained after taking into account detector and selection acceptance. Our preliminary experimental results (not shown) are in favour of the $1^{-}$case. However the de-

Table 1: Decay angular distributions for $J^{P}=1^{ \pm}$assignments to $\omega \pi^{0}$.

| $J^{P}$ | $I(\cos \psi)$ | $I(\cos \theta)$ |
| :---: | :---: | :---: |
| $1^{+}\left(b_{1}(1235)\right)$ | $\sim 1+x \cos ^{2} \psi$ | $\sim \sin ^{2} \theta$ |
| $1^{-}\left(\rho^{\prime}\right)$ | $\sim 1+\cos ^{2} \psi$ | $\sim 1+\cos ^{2} \theta$ |



Figure 6: Estimate of the cos $\theta$ distributions, based on Monte Carlo simulations of the detector and selection acceptance, for $J^{P}\left(\omega \pi^{0}\right)=1^{-}$(left) and $1^{+}$ (right).
pendence on the SCHC assumption should be kept in mind. This holds as well for the distribution (iii), which shows an indication of a $\cos 2 \Psi$ contribution, characteristic of $J^{P}=1^{-}$. Interference between S - and P -wave, corresponding to $1^{+}$and $1^{-}$decay in $\omega \pi^{0}$, would give rise to a forward-backward anisotropy in the distribution (i), irrespective of the SCHC assumption. Detailed acceptance studies are required for this analysis.

## 6 Conclusion

We have observed the exclusive production of $\omega \pi^{0}$ in muon scattering via virtual photons in the energy range around 90 GeV lab. energy. The mass spectrum is dominated by a peak at 1250 MeV and width 300 MeV , which is consistent with previous photoproduction experiments. Preliminary results on angular correlations are consistent with the presence of a $1^{-}$contribution, if SCHC holds. An appreciable increase in statistics is expected with the 2006 and 2007 COMPASS data.

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