

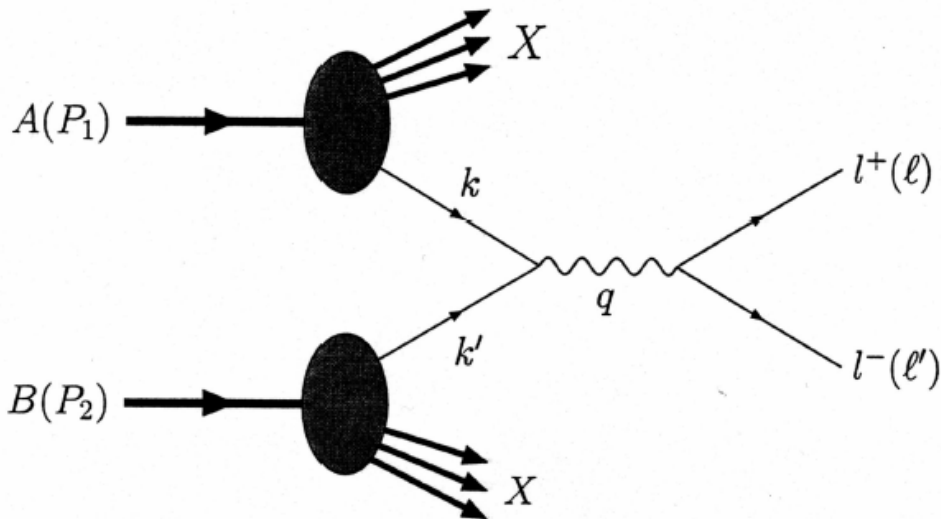
Drell-Yan Di-Lepton Production $\bar{p}p \rightarrow \mu^+ \mu^- X$

Why Drell-Yan?

Asymmetries depend on PD only
(SIDIS \rightarrow convolution with QFF)

Why ?

Each valence quark can contribute to the diagram



Kinematics

$$x_1 = \frac{M^2}{2P_1 \cdot q}$$

$$x_2 = \frac{M^2}{2P_2 \cdot q}$$

$$x_F = x_1 - x_2$$

$$\tau = x_1 x_2 = \frac{M^2}{s}$$

Drell-Yan Di-Lepton Production $\bar{p}p \rightarrow \mu^+ \mu^- X$

$$\frac{d^2\sigma}{dM^2 dx_F} = \frac{4\alpha^2\pi}{9M^2s} \frac{1}{x_1 + x_2} \sum_a e_a^2 \left[f^a(x_1) f^{\bar{a}}(x_2) + f^{\bar{a}}(x_1) f^a(x_2) \right]$$

Scaling:

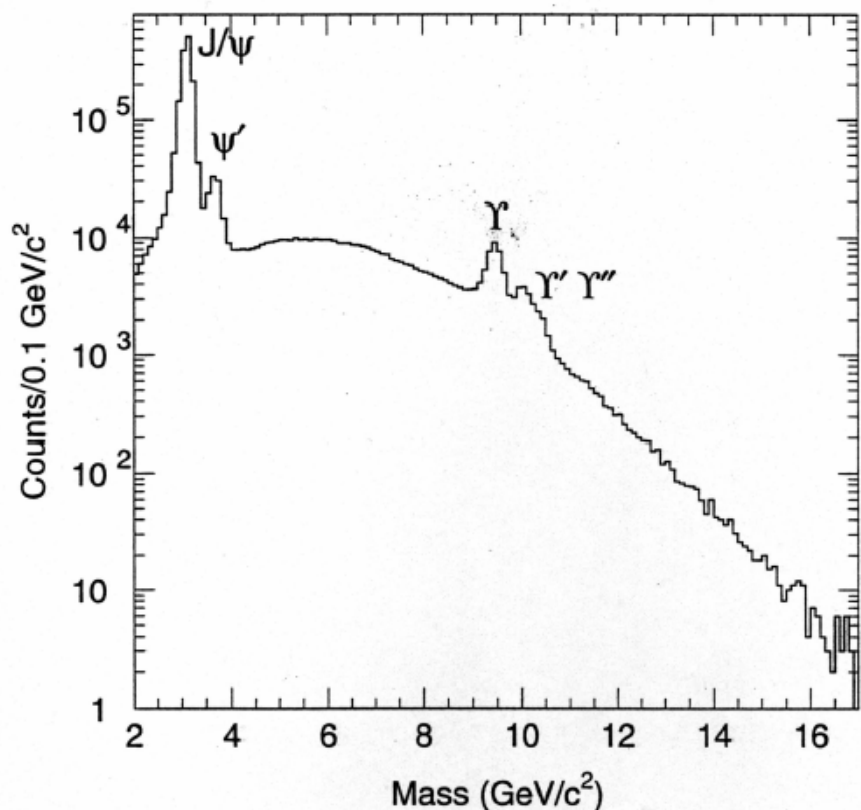
$$\frac{d^2\sigma}{d\sqrt{\tau} dx_F} \propto \frac{1}{s}$$

Full x_1, x_2 range $\Rightarrow \tau \in [0, 1]$

$\vec{p}_{\text{BEAM}} \geq 40 \text{ GeV}/c$ needed

$$\sigma_{\bar{p}p \rightarrow \mu^+ \mu^- X} \approx 0.3 \text{ nb}^{[1]}$$

^[1] Anassontzis et al., Phys. Rev. D38 (1988) 1377



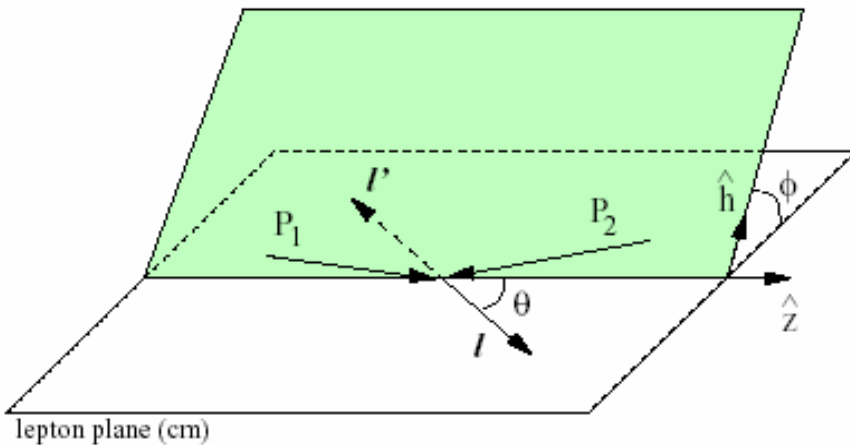
Drell-Yan Asymmetries

Polarised beam and target

$$A_{LL} = \frac{\sum_a e_a^2 g_1^a(\mathbf{x}_1) g_1^{\bar{a}}(\mathbf{x}_2)}{\sum_a e_a^2 f_1^a(\mathbf{x}_1) f_1^{\bar{a}}(\mathbf{x}_2)}$$

$$A_{TT} = \frac{\sin^2\theta \cos 2\phi}{1 + \cos^2\theta} \frac{\sum_a e_a^2 h_1^a(\mathbf{x}_1) h_1^{\bar{a}}(\mathbf{x}_2)}{\sum_a e_a^2 f_1^a(\mathbf{x}_1) f_1^{\bar{a}}(\mathbf{x}_2)}$$

$$A_{LT} = \frac{2 \sin 2\theta \cos \phi}{1 + \cos^2\theta} \frac{M}{Q} \frac{\sum_a e_a^2 \left(g_1^a(\mathbf{x}_1) x_2 g_T^{\bar{a}}(\mathbf{x}_2) - x_1 h_L^a(\mathbf{x}_1) h_1^{\bar{a}}(\mathbf{x}_2) \right)}{\sum_a e_a^2 f_1^a(\mathbf{x}_1) f_1^{\bar{a}}(\mathbf{x}_2)}$$



To be corrected for:

$$\frac{1}{P_B f P_T}$$

NH₃ polarised target:

$$f = \frac{3}{17} = 0.176$$

$$P_T \approx 0.85$$

Drell-Yan Asymmetries

Unpolarised beam and target

Di-Lepton Rest Frame

$$\frac{1}{\sigma} \frac{d\sigma}{d\Omega} = \frac{3}{4\pi} \frac{1}{\lambda+3} \left(1 + \lambda \cos^2\theta + \mu \sin^2\theta \cos\varphi + \frac{\nu}{2} \sin^2\theta \cos 2\varphi \right)$$

NLO pQCD: $\lambda \sim 1$, $\mu \sim 0$, $\nu \sim 0$

Experimental data ^[1]: $\nu \sim 30\%$

^[1] J.S.Conway et al., Phys. Rev. D39(1989)92.

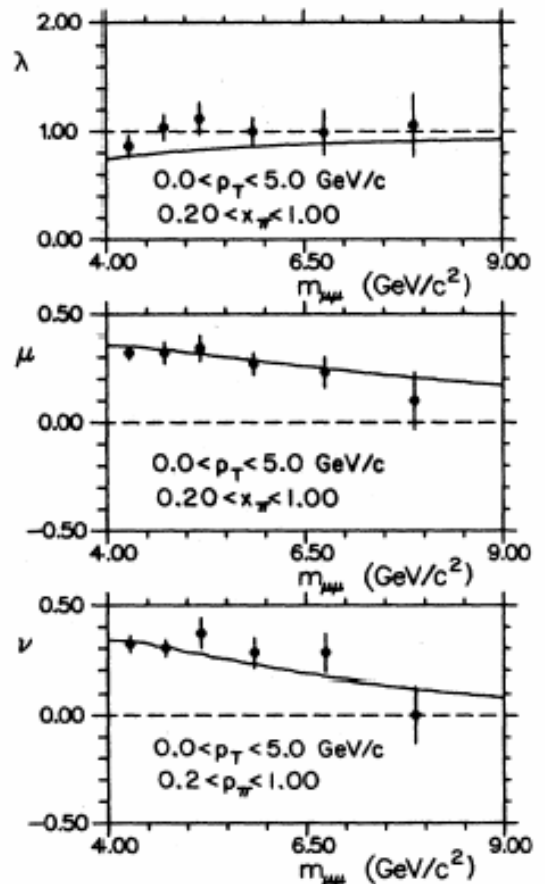
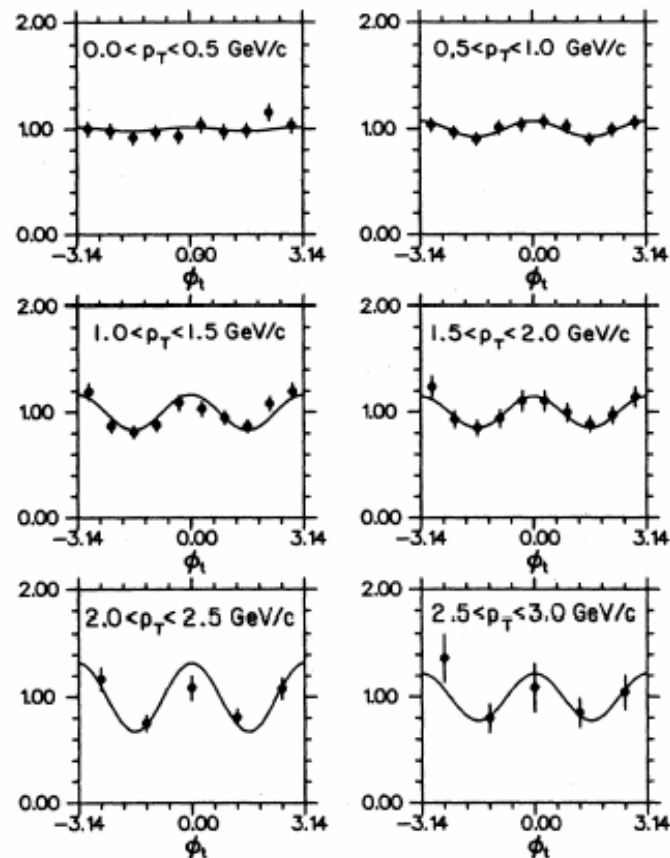
ν involves transverse spin effects at leading twist
^[2]:

$\cos 2\varphi$ contribution to angular distribution
provide:

$$\mathbf{h}_1^\perp(\mathbf{x}_2, \boldsymbol{\kappa}_\perp^2) \bar{\mathbf{h}}_1^\perp(\mathbf{x}_1, \boldsymbol{\kappa}'_\perp{}^2)$$

^[2] D. Boer et al., Phys. Rev. D60(1999)014012.

Angular distribution in CM frame



$$-0.6 < \cos\vartheta < 0.6$$

$$4 < M < 8.5 \text{ GeV}/c^2$$

- cut on \mathbf{P}_T selects asymmetry
- 30% asymmetry observed

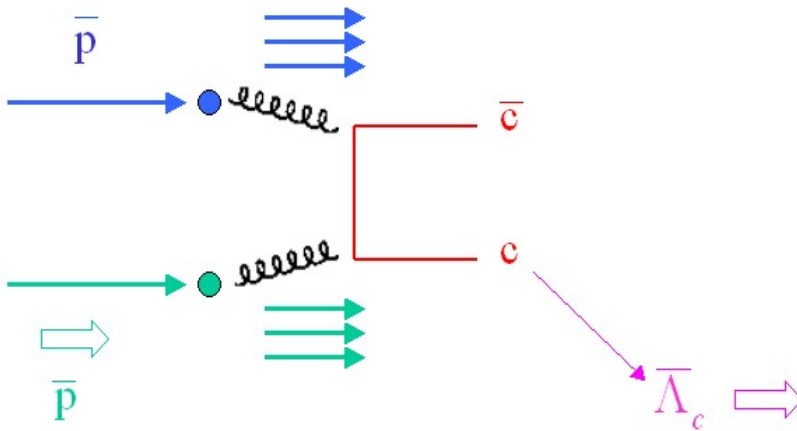
Drell-Yan Asymmetries

Unpolarised beam, polarised target

$$A_T = |S_{1T}| \frac{2 \sin 2\theta \sin(\varphi - \varphi_{S_1})}{1 + \cos^2\theta} \frac{M}{Q} \frac{\sum_a e_a^2 [x_1 f_1^{a\perp}(x_1) f_1^{\bar{a}}(x_2) + x_2 h_1^a(x_1) h_1^{\bar{a}\perp}(x_2)]}{\sum_a e_a^2 f_1^a(x_1) f_1^{\bar{a}}(x_2)}$$

Even unpolarised beam is
a powerful tool to investigate
 κ_T dependence of QDF

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$\vec{p}, \vec{\Lambda}_c$ longitudinally polarized

Weak decay

$$\Lambda_c^+ \rightarrow \Lambda \pi^+$$

B.R. = 0.9 %

$$\alpha = -0.98$$

$\sigma_{\bar{p}p \rightarrow \Lambda_c}$ at 40 GeV ?

Assume $1 \mu\text{b}$ [1] \Rightarrow # ev. $5 \cdot 10^{-3}$ /s $L = 2.25 \cdot 10^{31} \text{ cm}^{-2} \cdot \text{s}^{-1}$



36000 events / 100 days